

Microhydrodynamics of Liquid Crystals



- Hybrid LB approach
- Gallery of examples
- Nematic/cholesteric hydrodynamics
 - Phase diagram: the elusive BPIII
 - Particles in cholesterics: microrheology
 - Dimers in cholesterics: self-assembled rotors
 - Particles in active nematics: negative drag coefficient

With: *O. Henrich, J. Lintuvuori, K. Stratford,
G. Foffano, D. Marenduzzo*

Antecedents: *J. Yeomans Group, Oxford*



Lattice Boltzmann for Simple Fluids

= fast fluid mechanics on a lattice

- each site \mathbf{x} has velocity set $\{\mathbf{c}_i\}$: $\mathbf{c}_i \Delta t$ = lattice vector
- $f_i(\mathbf{x},t)$: population of fluid “particles” at \mathbf{x} with velocity \mathbf{c}_i

$$\rho(\mathbf{x},t) = \sum_i f_i \quad \text{fluid density}$$

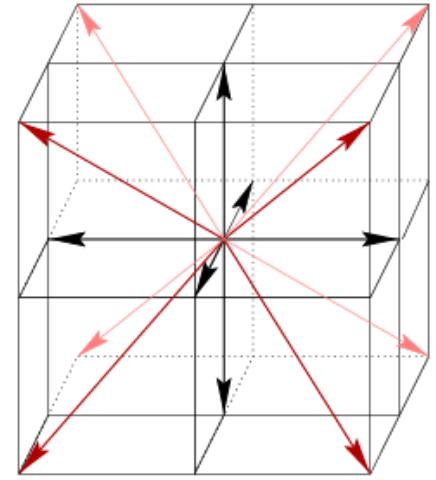
$$\rho \mathbf{v}(\mathbf{x},t) = \sum_i f_i \mathbf{c}_i + \mathcal{O}(\Delta t) \quad \text{fluid velocity}$$

- local streaming, and relaxation:

$$f_i(\mathbf{x} + \mathbf{c}_i, t+1) - f_i(\mathbf{x}, t) = \sum_j L_{ij} (f_j(\mathbf{x}, t) - f_j^0(\mathbf{x}, t)) - \mathbf{F} \cdot \nabla f_i + \mathcal{O}(\Delta t)$$

- continuum limit: Navier Stokes equation

$$\rho(\partial \mathbf{v} / \partial t + \mathbf{v} \cdot \nabla \mathbf{v}) = -\nabla P + \eta \nabla^2 \mathbf{v} + \mathbf{F}_{\text{ext}}$$



Lattice Boltzmann for Liquid Crystals

$$\mathbf{F}_{\text{ext}} = \nabla \cdot \Pi[\mathbf{Q}]$$

Stress tensor Π derives from a mesoscopic free energy functional $\mathcal{F}[\mathbf{Q}]$

Couple LB to finite difference code that updates \mathbf{Q}

$$\rho(\partial \mathbf{v}/\partial t + \mathbf{v} \cdot \nabla \mathbf{v}) = -\nabla P + \eta \nabla^2 \mathbf{v} + \nabla \cdot \Pi[\mathbf{Q}]$$

$$\mathcal{D}\mathbf{Q}/\mathcal{D}t = \Gamma \mathbf{H} + \boldsymbol{\zeta}$$

Total derivative knows about advection by \mathbf{v} and flow alignment by vorticity

$\boldsymbol{\zeta}$ = thermal noise

Molecular field: $\mathbf{H} = -\delta \mathcal{F}/\delta \mathbf{Q} + \frac{1}{3} \text{Tr} (\delta \mathcal{F}/\delta \mathbf{Q})$

A. N. Beris and B. J. Edwards, Thermodynamics of Flowing Systems, OUP (1994)

D. Marenduzzo, E. Orlandini, MEC, J. Yeomans, PRE 76, 031921 (2007)

Lattice Boltzmann

Discretizes momentum space

⇒ fast momentum transfer across lattice

Fully local dynamics

⇒ excellent for parallel supercomputers

Adding colloids (bounce-back of momentum):

K. Stratford et al., J Stat Phys 121, 163 (2005)

Adding thermal noise:

R. Adhikari et al., EPL 71, 473 (2005)

Adding subgrid stokeslet/stresslet particles:

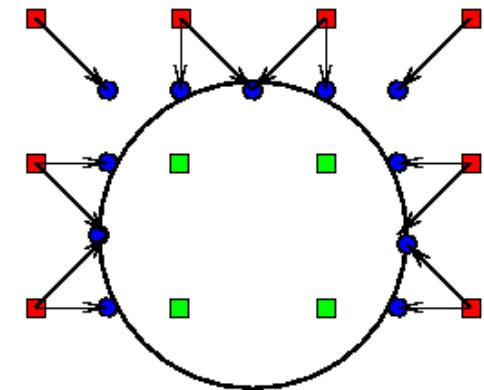
R. Nash et al., PRL 104, 258101 (2010)

Compressible flows e.g. liquid-vapour systems:

No hybrid route; noise etc now in place

M. Gross et al., PRE 82, 056714 (2010), JSTAT P03030 (2011), Phil Trans 369, 2274 (2011)

[Compare Dunweg Talk]

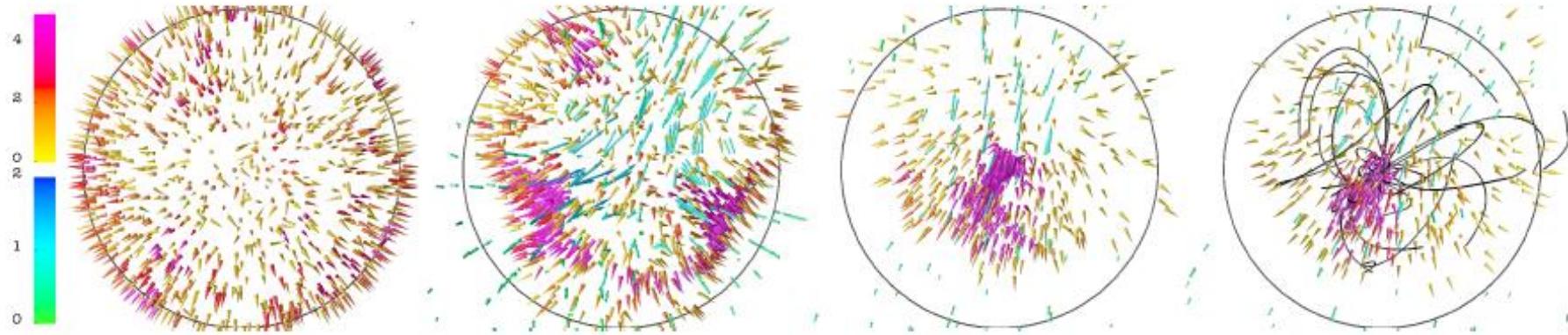


Lattice Boltzmann

Sub-grid swimming
stresslets (active
colloids/bacteria)
in harmonic trap

*R. Nash et al, PRL 104,
258101 (2010)*

MOVIE AVAILABLE AT PRL



Lattice Boltzmann

Binary fluid mixture

+ interfacial particles

+ gravity

*E. Kim et al PRE 85,
020403 (2012)*

MOVIE AVAILABLE AT PRE

Nematic Liquid Crystals

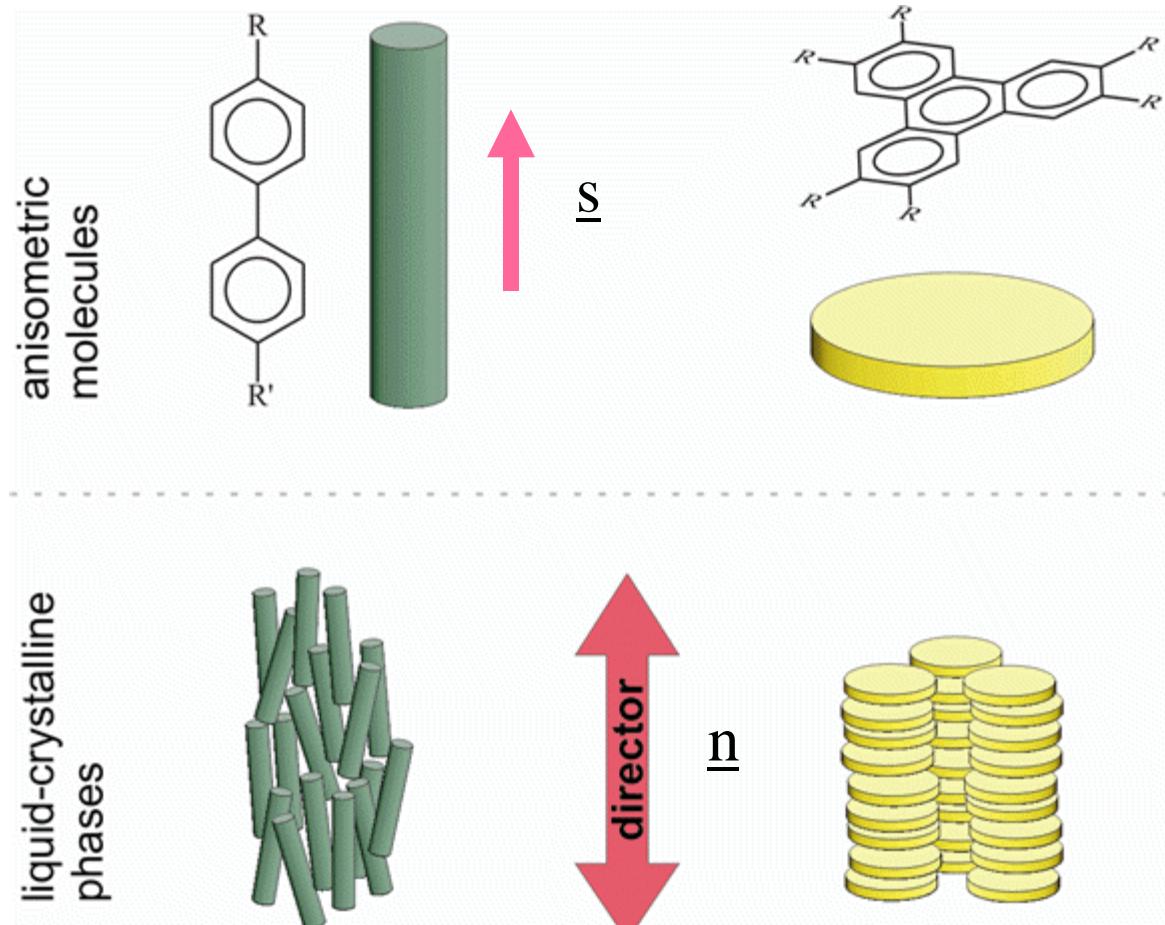
\mathbf{s} = molecular axis

$\langle s_\alpha \rangle = 0$ (apolar)

order parameter

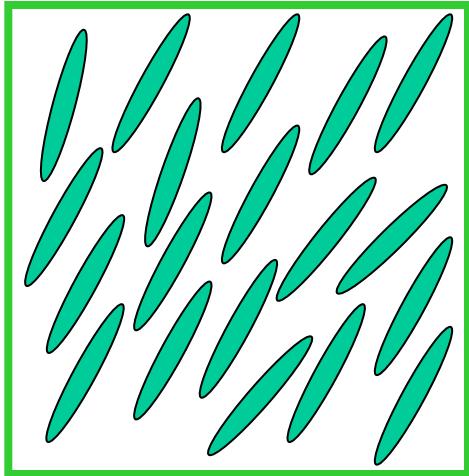
$$Q_{\alpha\beta} = \langle s_\alpha s_\beta \rangle - \delta_{\alpha\beta}/3$$

largest eigenvector = \mathbf{n}
eigenvalue q



Cholesterics: Nematics with a twist

Chiral nematic molecules



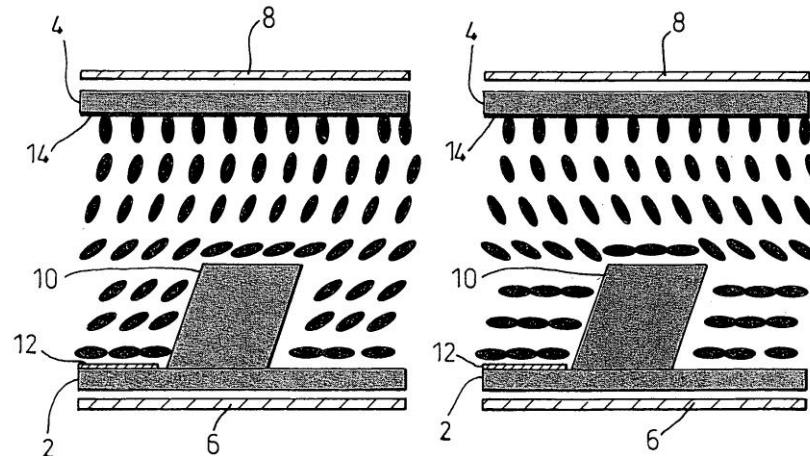
nematic: n uniform

$$Q_{\alpha\beta} = q(n_\alpha n_\beta - \delta_{\alpha\beta}/3)$$

in uniaxial case



cholesteric: n has helical pitch ($\sim 1\mu\text{m}$)



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Cholesterics: Bistable (powerless) LCDs



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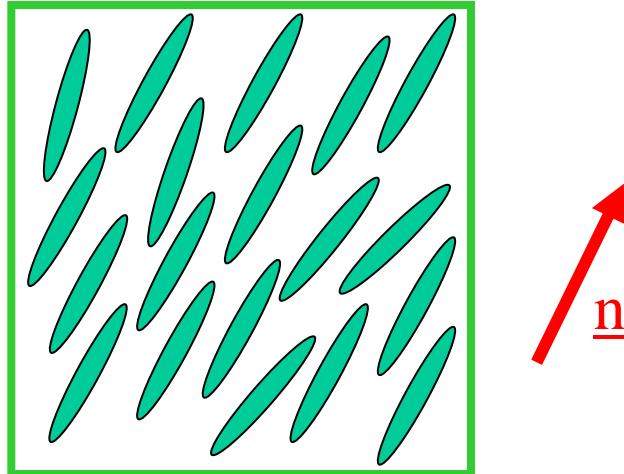
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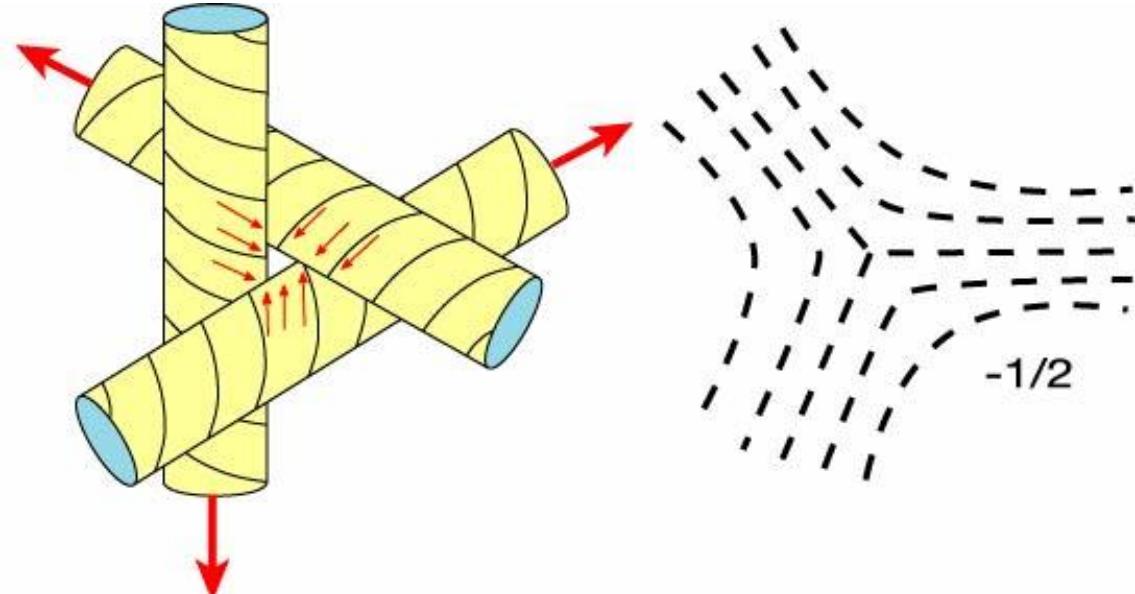
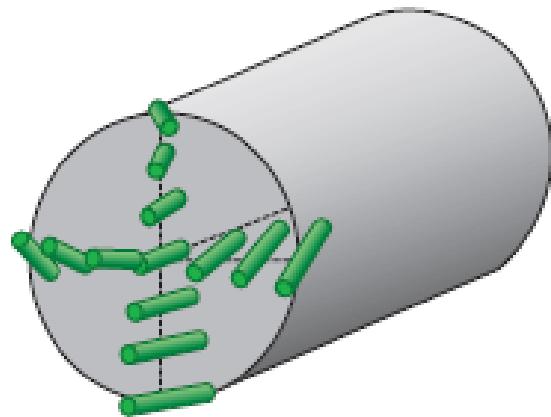
Cholesterics: Nematics with a twist

Chiral nematic molecules



Topological defects (disclination lines)

forced into system by too much chirality
and/or boundary conditions



Landau – de Gennes Free Energy Functional

$$\mathcal{F}[\mathbf{Q}] = \int (f_1 + f_2 + f_3) dV$$

f_1 = bulk nematic free energy density

$$f_1 = A_0 \left[\frac{1}{2} \left(1 - \frac{\gamma}{3} \right) Q_{\alpha\beta}^2 - \frac{\gamma}{3} Q_{\alpha\beta} Q_{\beta\gamma} Q_{\gamma\alpha} + \frac{\gamma}{4} Q_{\alpha\beta}^4 \right]$$

$A_0 \sim 10^6 \text{ Pa} \sim kT/\text{nm}^3$

I-N transition at $\gamma = 2.7$, spinodal at $\gamma = 3$

reduced temperature
 $\tau = 27/\gamma - 9$

$$f_2 = \frac{K}{2} \left[(\partial_\beta Q_{\alpha\beta})^2 + (\epsilon_{\alpha\gamma\delta} \partial_\gamma Q_{\delta\beta} + 2q_0 Q_{\alpha\beta})^2 \right]$$

K = elastic constant

$p = 2\pi/q_0$ = pitch of cholesteric helix

κ : reduced chirality
 $\kappa^2 = 108Kq_0^2/A_0\gamma$

LC dynamics: Beris-Edwards equations

Navier Stokes (η = viscosity)

$$\rho(\partial \mathbf{v}/\partial t + \mathbf{v} \cdot \nabla \mathbf{v}) = -\nabla P + \eta \nabla^2 \mathbf{v} + \nabla \cdot \boldsymbol{\Pi}$$

Stress tensor (ξ = flow alignment parameter)

$$\begin{aligned} \boldsymbol{\Pi}_{\alpha\beta} = & -P_0 \delta_{\alpha\beta} + 2\xi(Q_{\alpha\beta} + \frac{1}{3}\delta_{\alpha\beta})Q_{\gamma\epsilon}H_{\gamma\epsilon} - \xi H_{\alpha\gamma}(Q_{\gamma\beta} + \frac{1}{3}\delta_{\gamma\beta}) \\ & - \xi(Q_{\alpha\gamma} + \frac{1}{3}\delta_{\alpha\gamma})H_{\gamma\beta} - \partial_\beta Q_{\gamma\nu} \frac{\delta F}{\delta \partial_\alpha Q_{\gamma\nu}} + Q_{\alpha\gamma}H_{\gamma\beta} - H_{\alpha\gamma}Q_{\gamma\beta} \end{aligned}$$

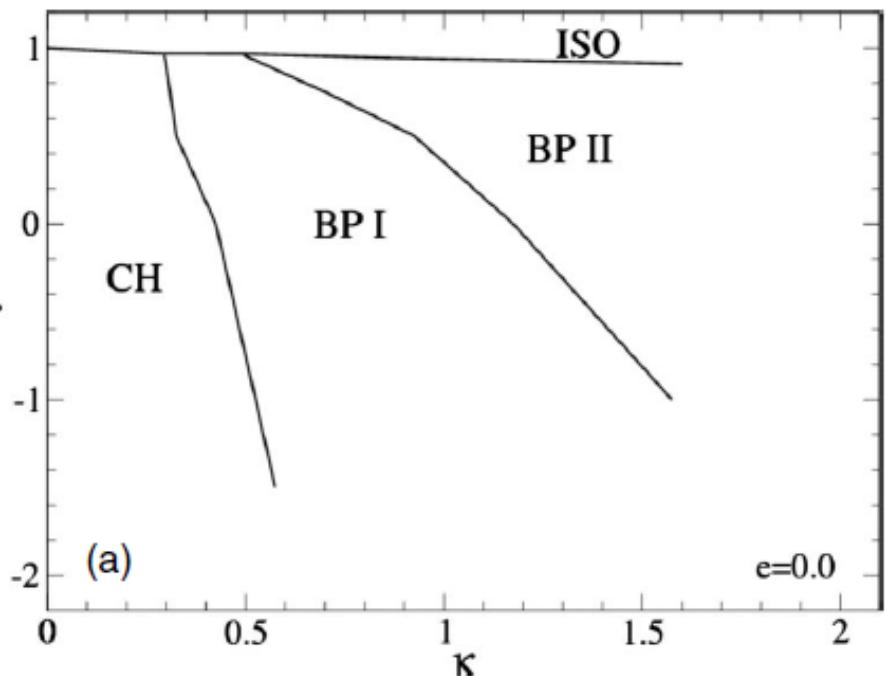
Order parameter relaxation (Γ = rotational mobility, \mathbf{S} = rotational advector)

$$(\partial_t + \mathbf{v} \cdot \nabla)Q_{\alpha\beta} + S_{\alpha\beta}(\xi, \nabla \mathbf{v}, \mathbf{Q}) = \Gamma H_{\alpha\beta}$$

Molecular field ($\mathbf{F} = \int \mathbf{f} dV$)

$$\mathbf{H} = -\frac{\delta F}{\delta \mathbf{Q}} + (\mathbf{I}/3) \text{Tr} \frac{\delta F}{\delta \mathbf{Q}}$$

Phase diagram of Landau–de Gennes model

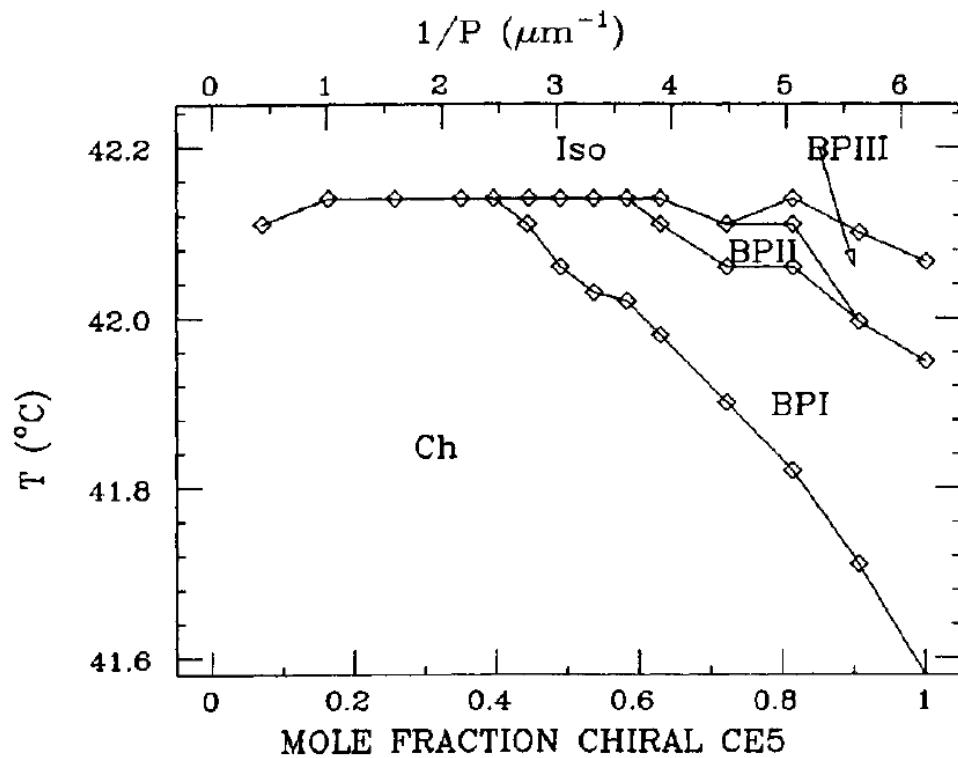


(a)

[Periodic phases only]

Numerical phase diagram

*G. Alexander + J. Yeomans,
PRE 74, 061706 (2006)*



[+Aperiodic BPIII]

Classic experiments

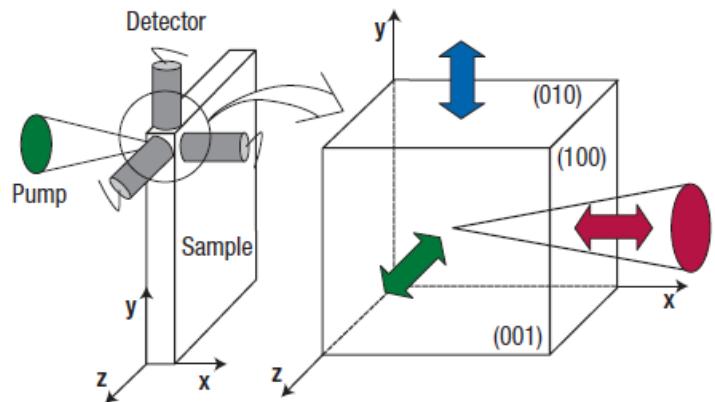
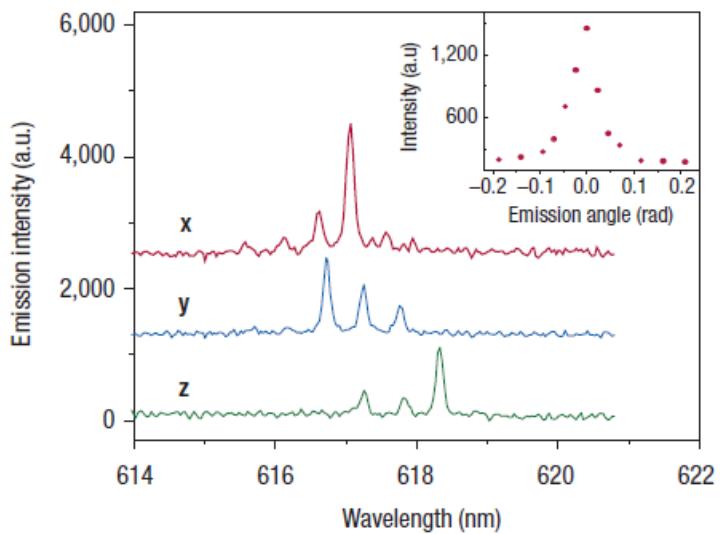
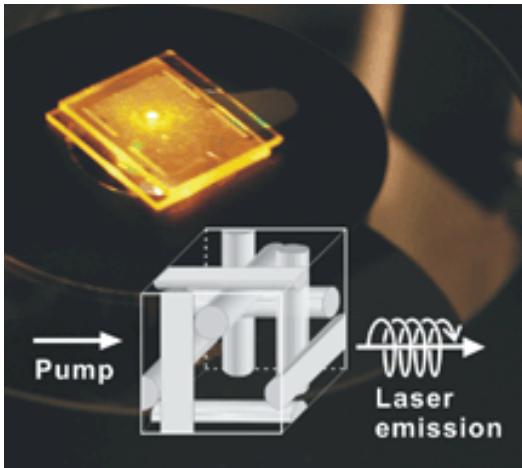
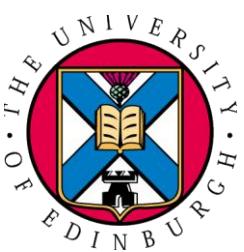
*Yang + Crooker,
PRA 35, 4419 (1987)*

BP Devices

Tri-directional pumped laser

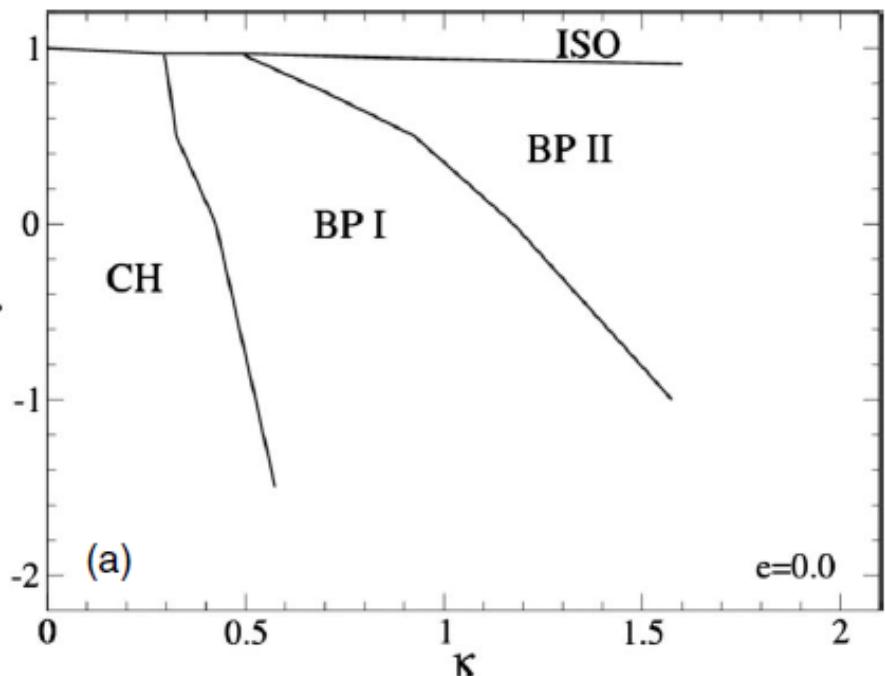
Mirrorless operation

Tunable colour



W. Cao et al, Nat Mat 1, 111 (2002)

Phase diagram of Landau–de Gennes model

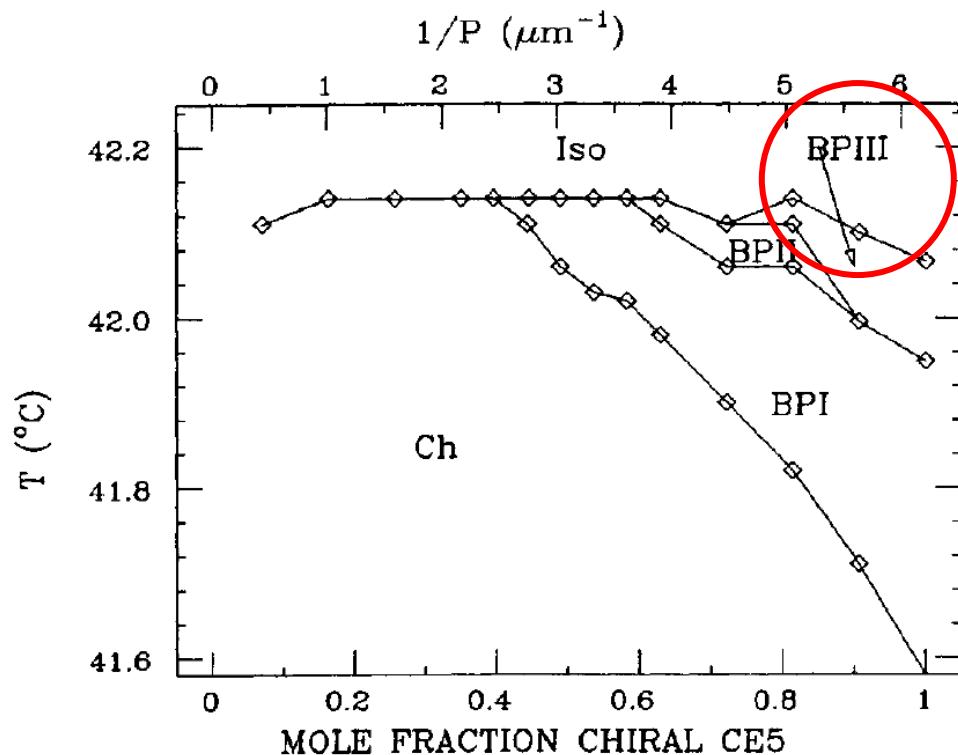


(a)

[Periodic phases only]

Numerical phase diagram

*G. Alexander + J. Yeomans,
PRE 74, 061706 (2006)*

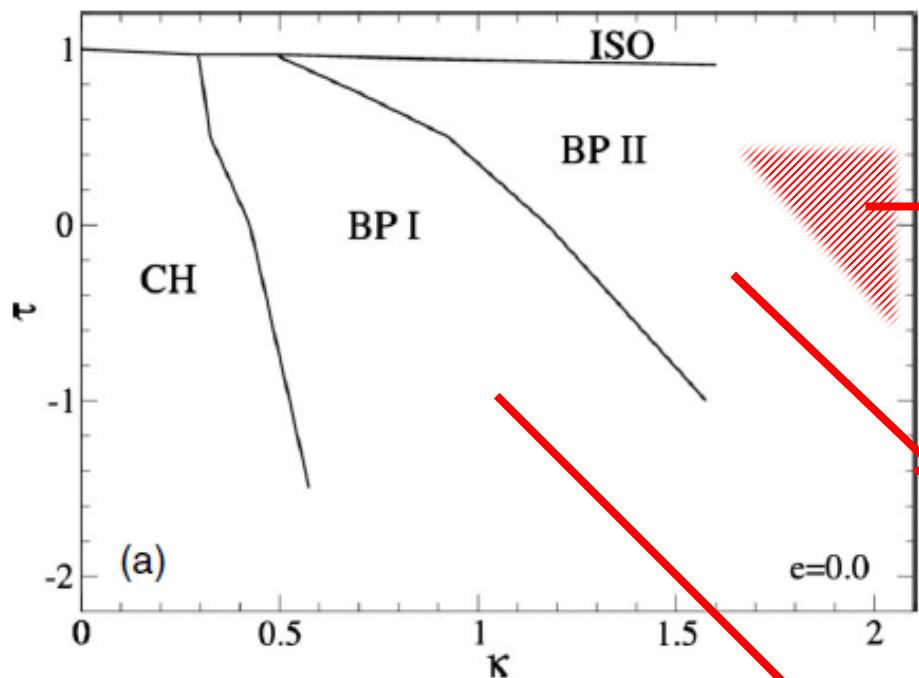


[+Aperiodic BPIII]

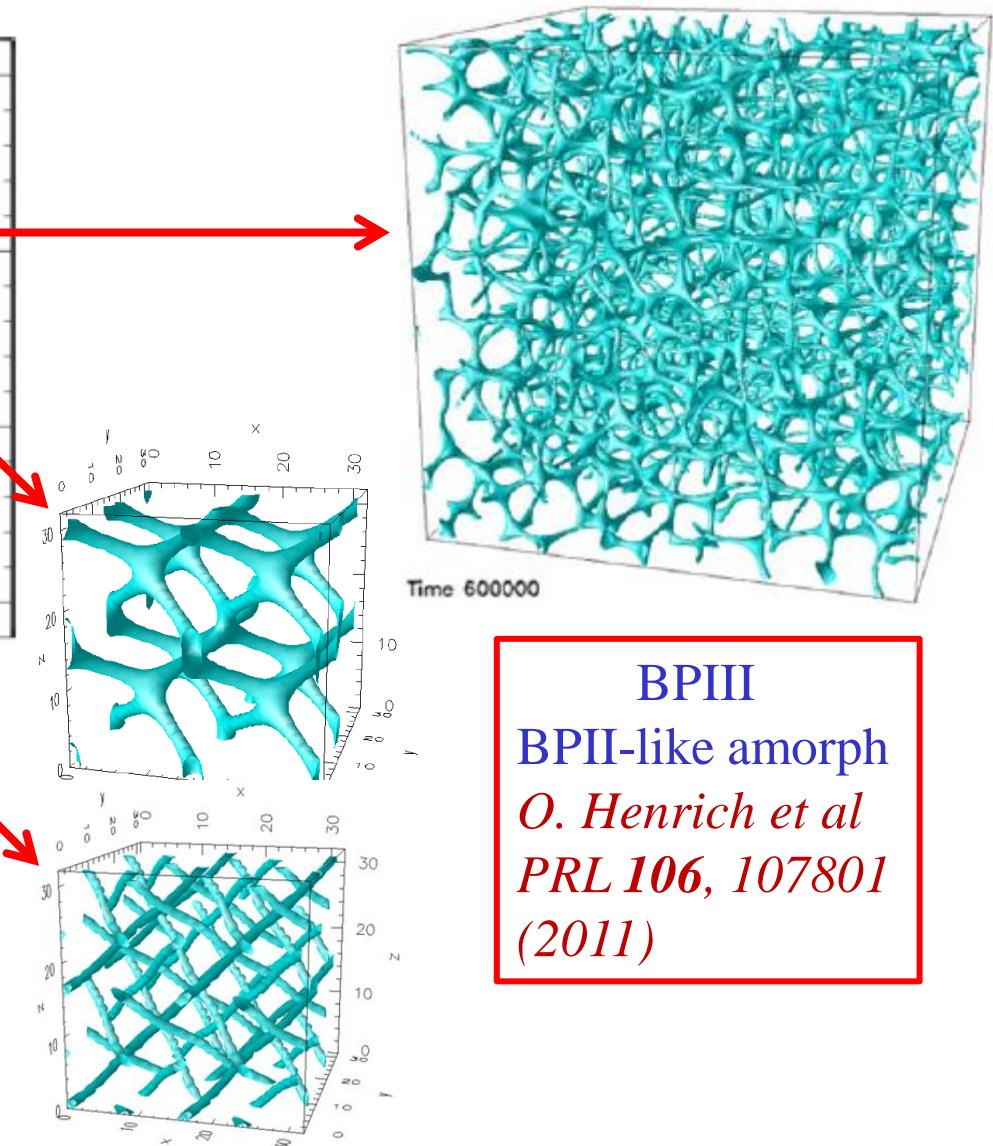
Classic experiments

*Yang + Crooker,
PRA 35, 4419 (1987)*

Phase diagram of Landau–de Gennes model



BPI,II
Ordered lattices of disclinations



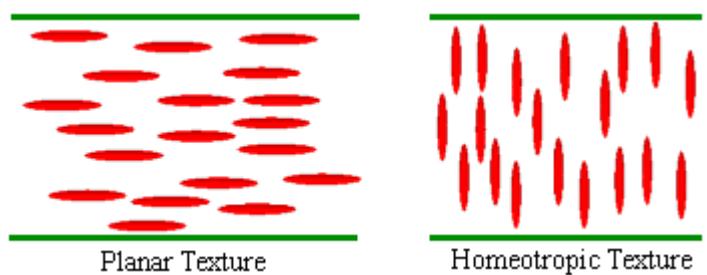
BPIII
BPII-like amorph
O. Henrich et al
PRL 106, 107801
(2011)

Colloids in Cholesterics

Statics

Equilibrium defect structure for planar or homeotropic anchoring

Dependence on colloid/pitch size ratio



Microrheology

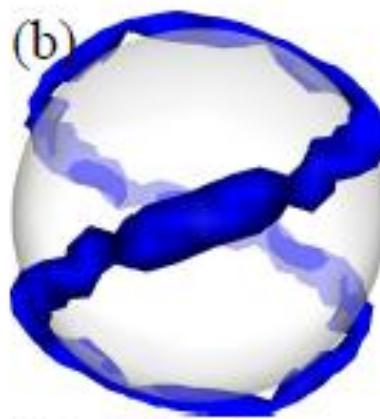
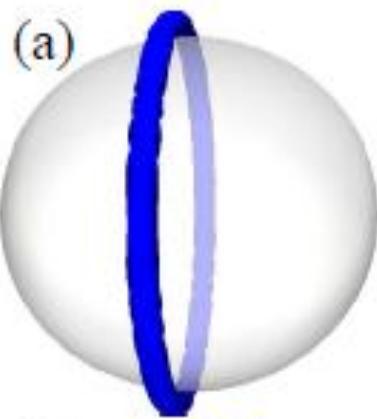
What is force \mathbf{f} to drag a colloid through the system at speed \mathbf{v}

(a) in the plane of cholesteric layers?

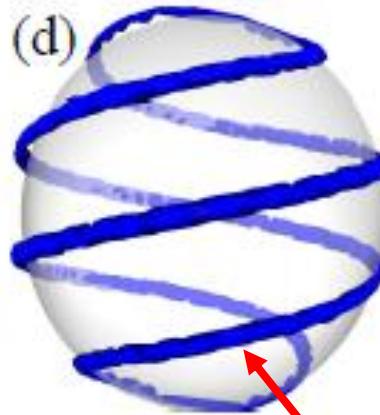
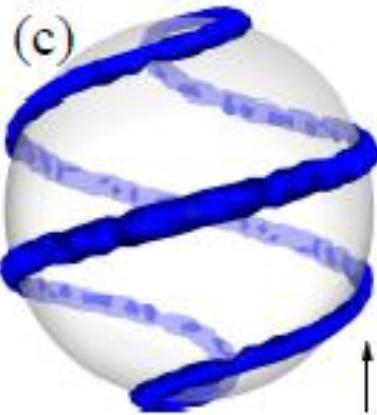
(b) along the helical axis?

Homeotropic (normal) anchoring

$R/p = 0$
(nematic)



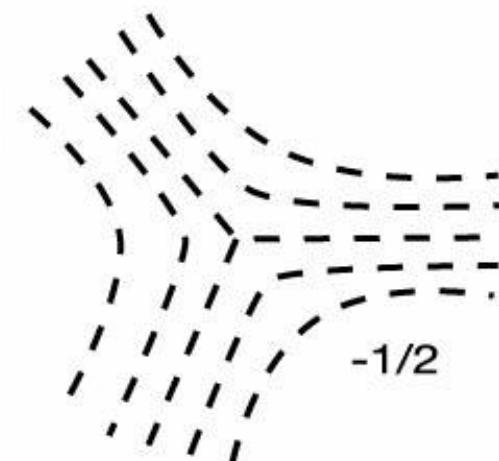
$R/p = 0.5$



$R/p = 0.25$

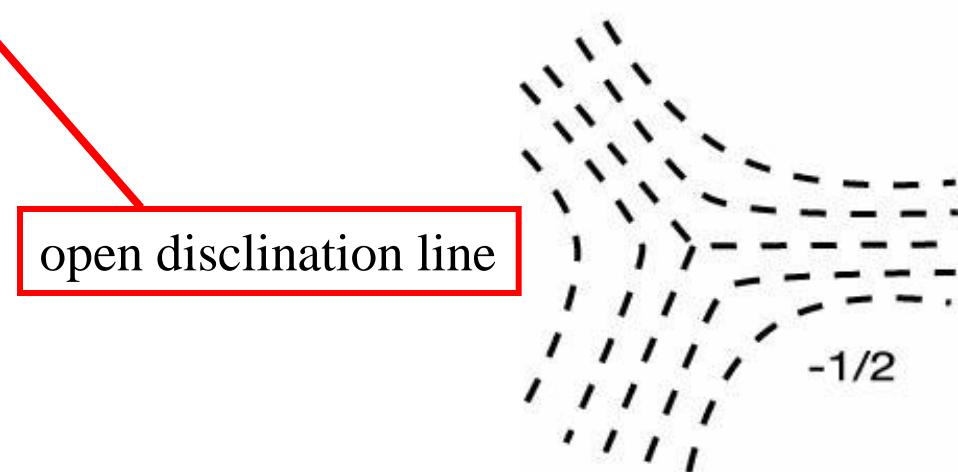
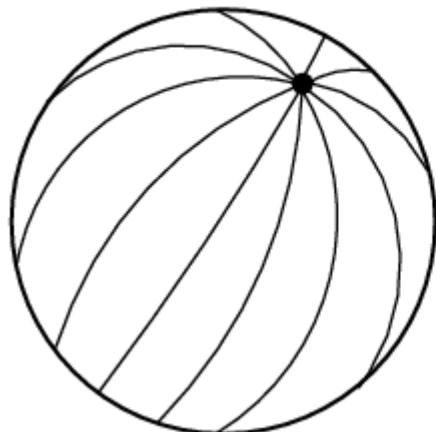
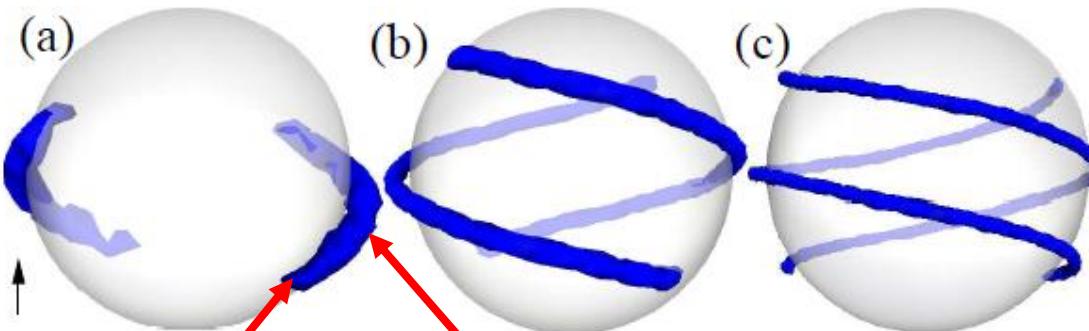
$R/p = 0.75$

closed disclination line



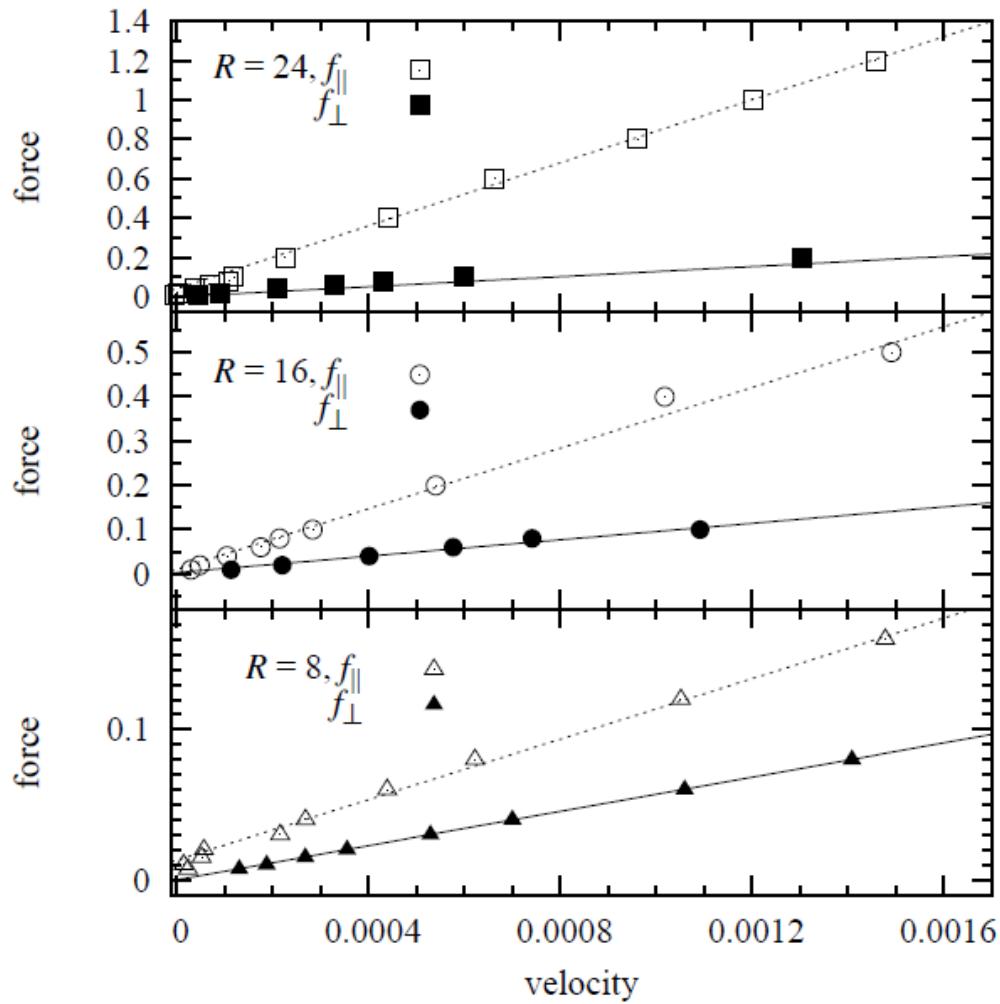
Planar anchoring

R/p: 0.25 0.50 0.75



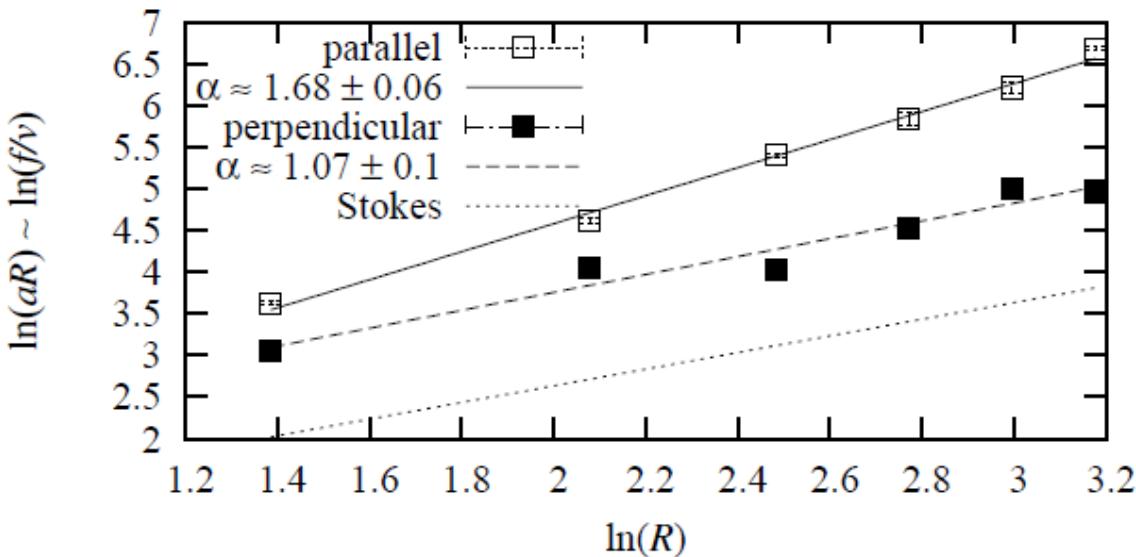
J. Lintuvuori et al, PRL, 105, 178302 (2010)

Microrheology: small speeds



force is linear in velocity
for both \parallel and \perp motion...

Microrheology: small speeds



but force not linear in R
when $v \parallel$ to helix

violation of Stokes law: $f/v = 6\pi\eta R$

force law probes mesoscopic state, not local material properties

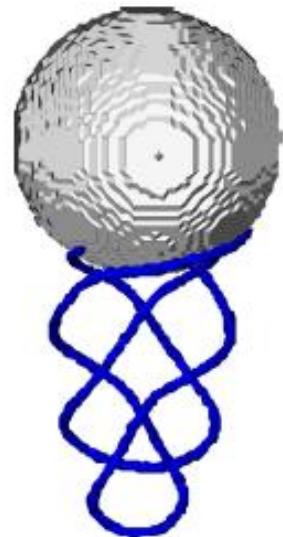
effective viscosity $\eta_{||}(R) \sim R^\alpha$, $\alpha \sim 0.7$

theory suggests $\alpha \sim 1$

Microrheology: larger speeds

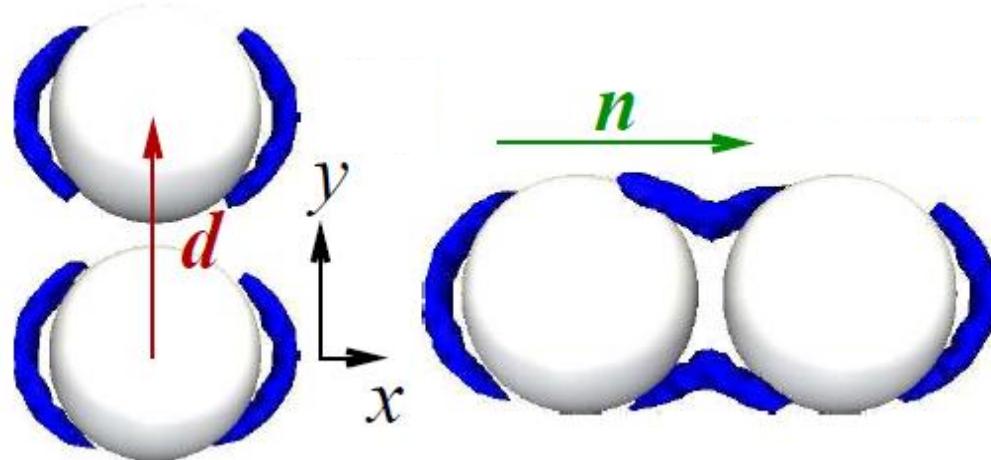
High Erickson number: $Er = 2q^2vR/\Gamma K = 0.7$

$R/p \sim 0.75$



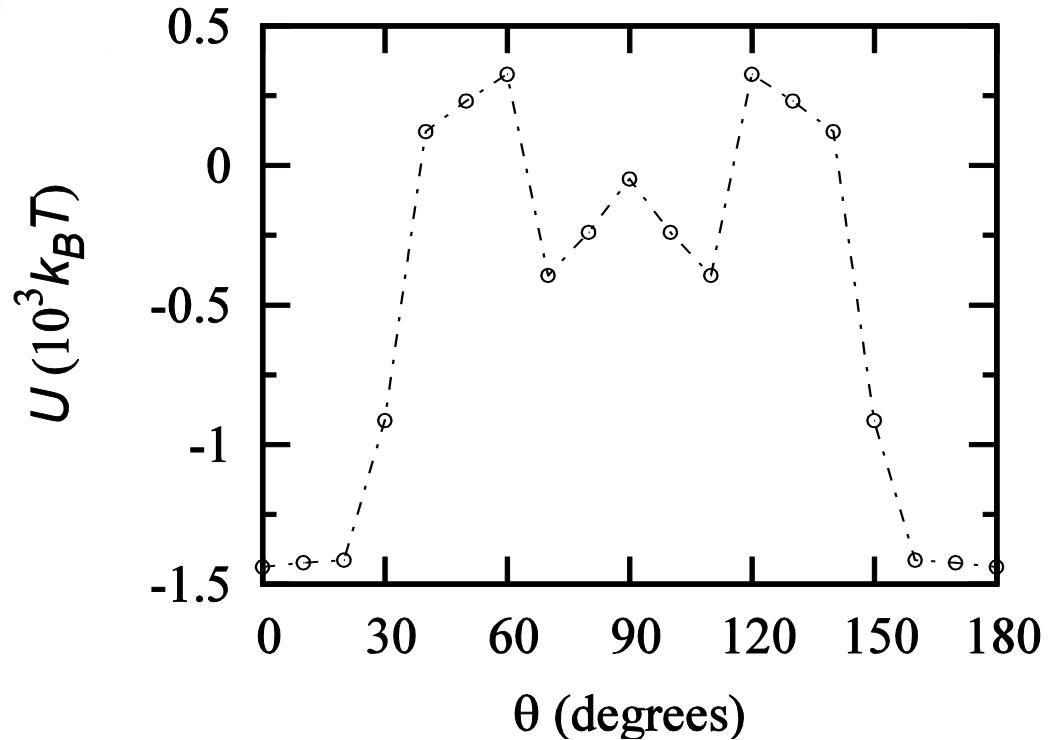
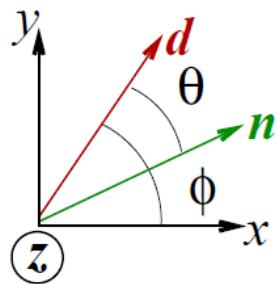
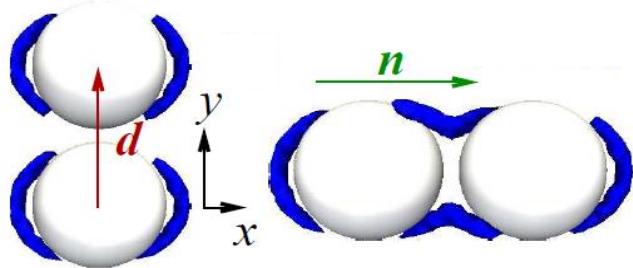
Dimers in cholesterics

- planar anchoring creates anisotropic short range attraction
- bonding by boojum exchange
- strong dependence on orientation of dimer axis to director



F. E. Mackay and C. Denniston, EPL, 94, 66003 (2011)

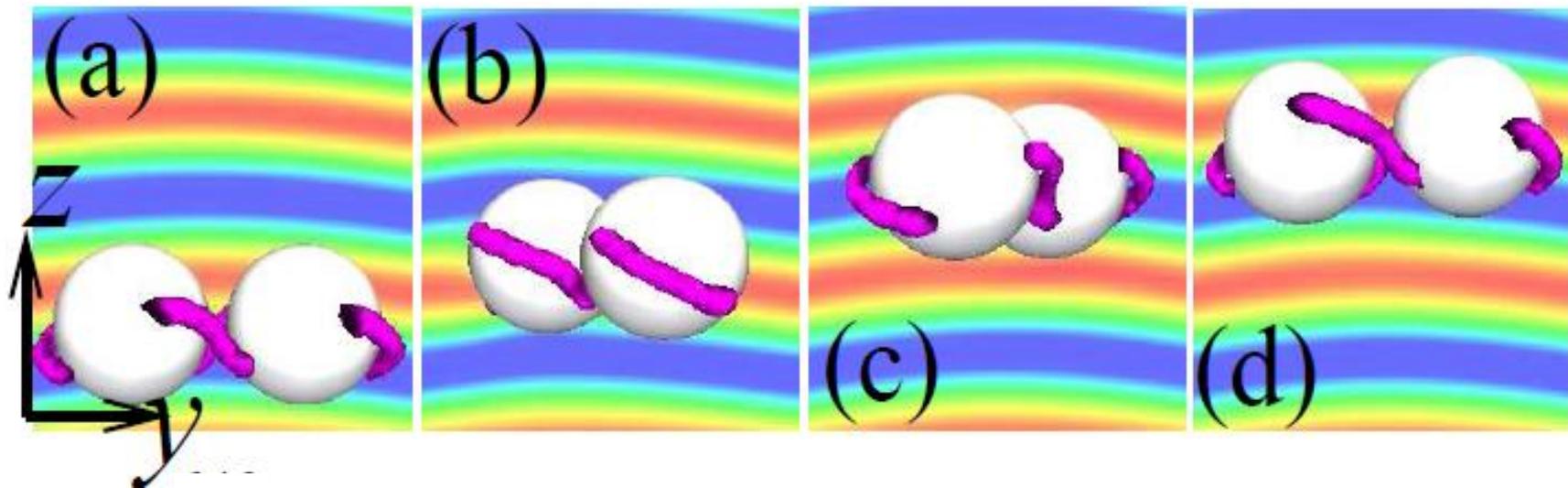
Dimers in x-y plane



*F. E. Mackay and C. Denniston, EPL, 94, 66003 (2011)
J. Lintuvuori et al, PRL 107, 268102 (2011)*

Pulling dimers through a cholesteric

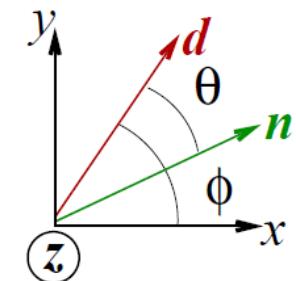
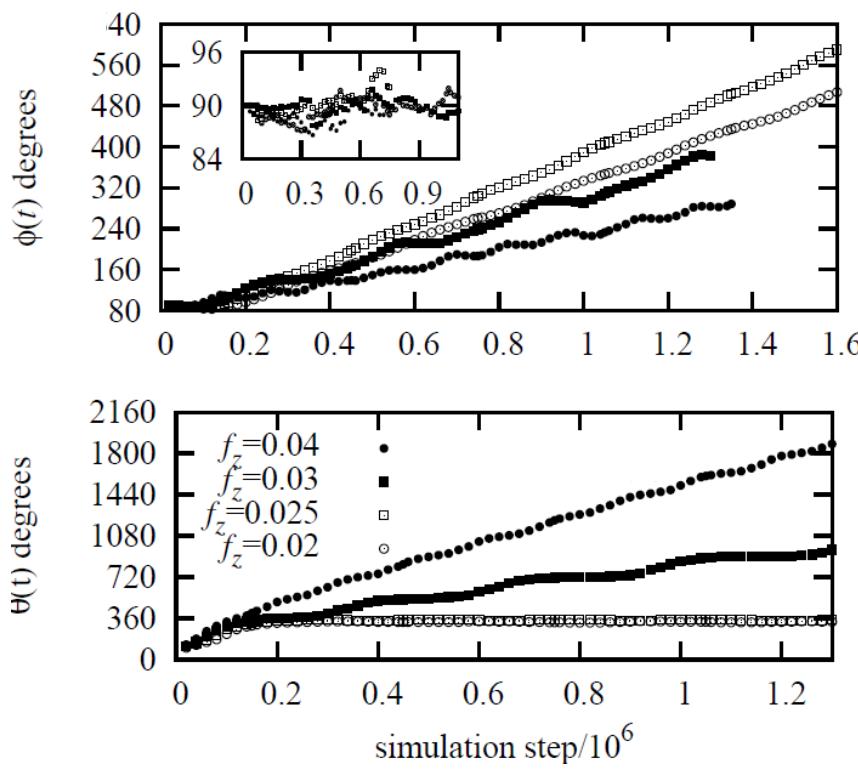
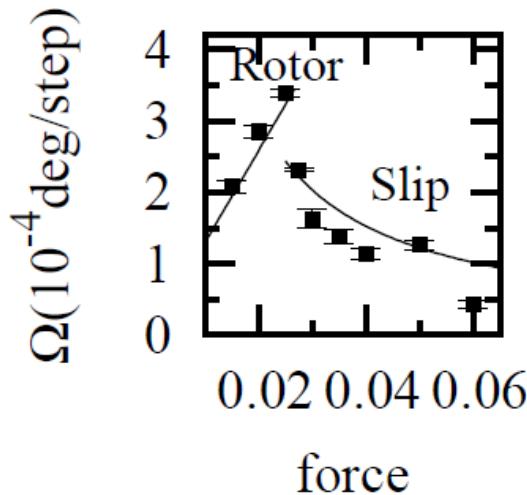
- rotation of **d** axis resulting from z translation
- small **f**: rotor phase, lock-in (with wobbles) between **d** and **n**
- larger **f**: phase slip



J. Lintuvuori et al, PRL 107, 268102 (2011)

Pulling dimers through a cholesteric

- rotation of **d** axis resulting from z translation
- small **f**: rotor phase, lock-in (with wobbles) between **d** and **n**
- larger **f**: phase slip



And finally...

Microhydrodynamics of Liquid Crystals



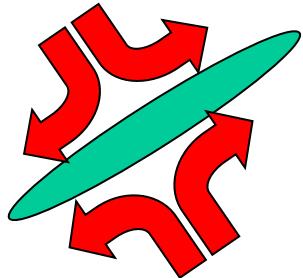
- Hybrid LB approach
- Gallery of examples
- Nematic/cholesteric hydrodynamics
 - Phase diagram: the elusive BPIII
 - Particles in cholesterics: microrheology
 - Dimers in cholesterics: self-assembled rotors
 - **Particles in active nematics: negative drag coefficient**

With: *O. Henrich, J. Lintuvuori, K. Stratford,
G. Foffano, D. Marenduzzo*

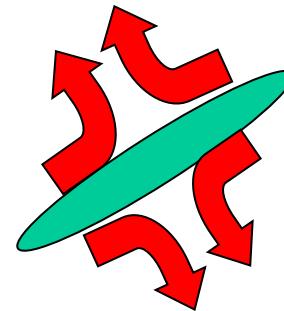
Antecedents: *J. Yeomans Group, Oxford*

Microrheology of active nematics

Active nematics:



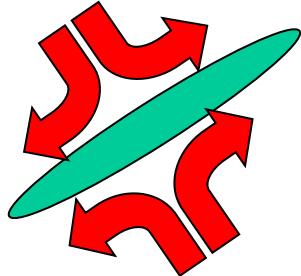
dense bacterial suspensions
(extensile, $\zeta > 0$)



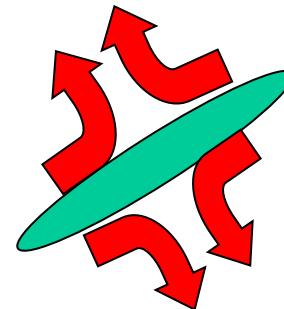
cytoskeletal actomyosin
(contractile, $\zeta < 0$)

Microrheology of active nematics

Active nematics:



dense bacterial suspensions
(extensile, $\zeta > 0$)



cytoskeletal actomyosin
(contractile, $\zeta < 0$)

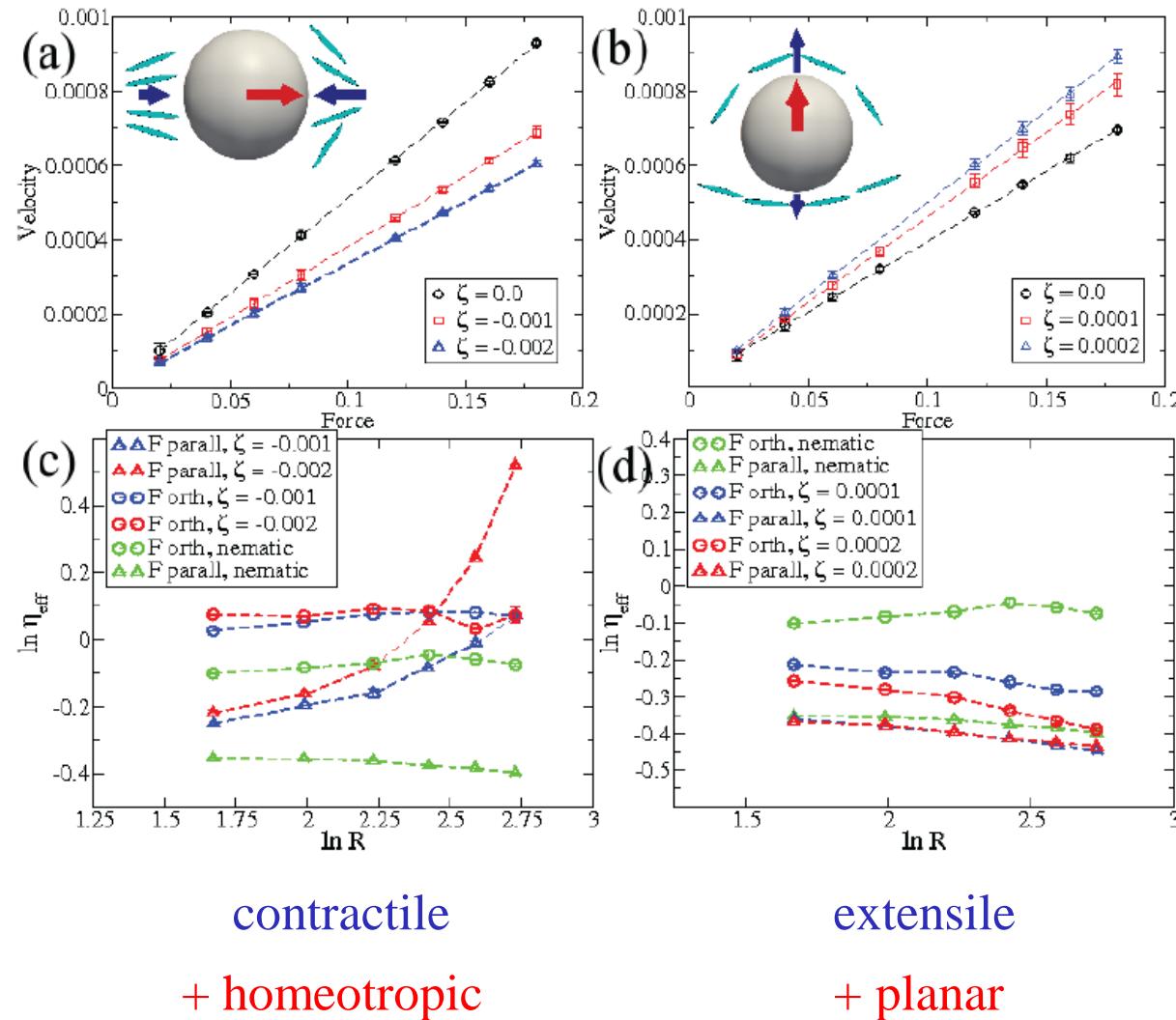
Equations of motion as previous ($p \rightarrow \infty$) with additional stress

$$\Delta \Pi = -\zeta [Q + I/3]$$

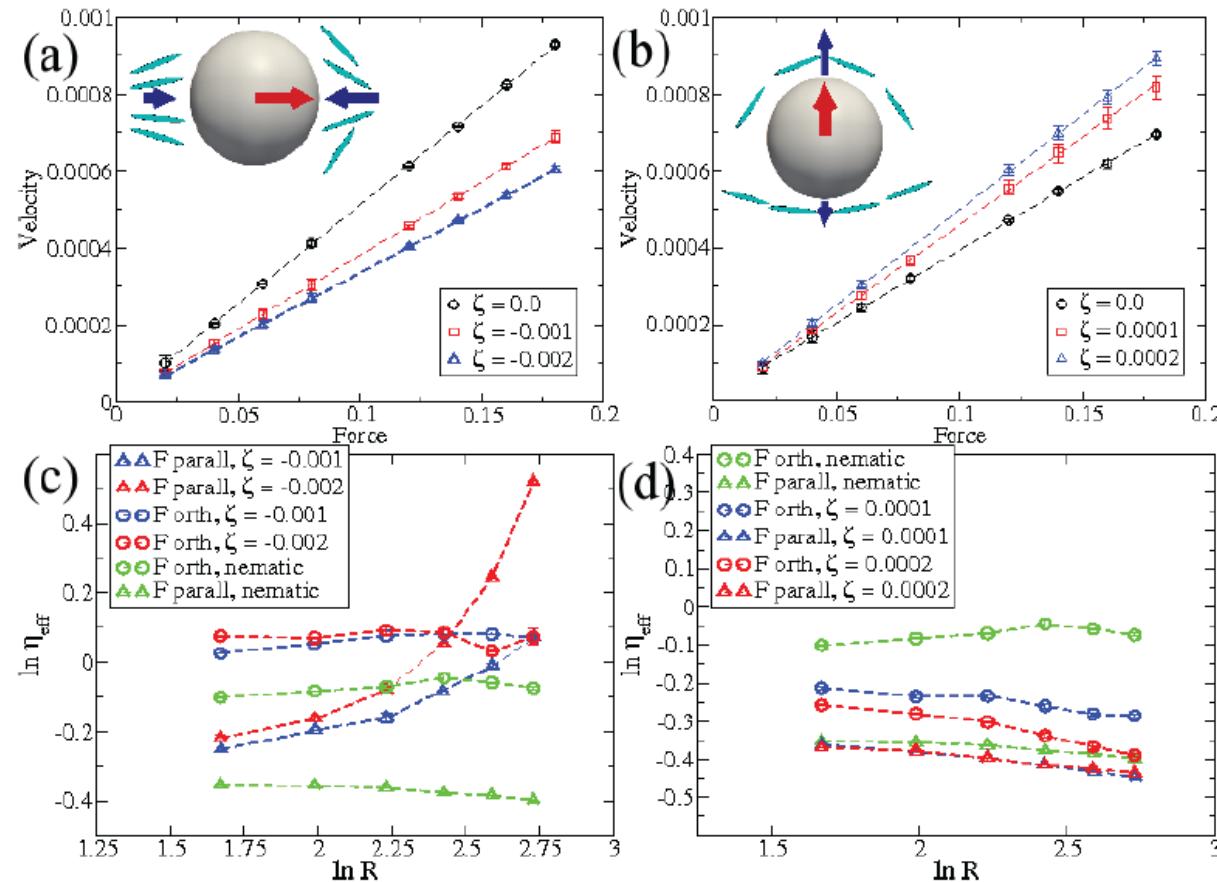
Flow instabilities in bulk samples (avoided here via PBCs)

*R. Simha and S. Ramaswamy, PRL 89, 058101 (2002), Y. Hatwalne et al
PRL 92, 118101 (2004), K. Kruse et al PRL 92, 078101 (2004)*

Microrheology of active nematics



Microrheology of active nematics



contractile
+ homeotropic

extensile
+ planar

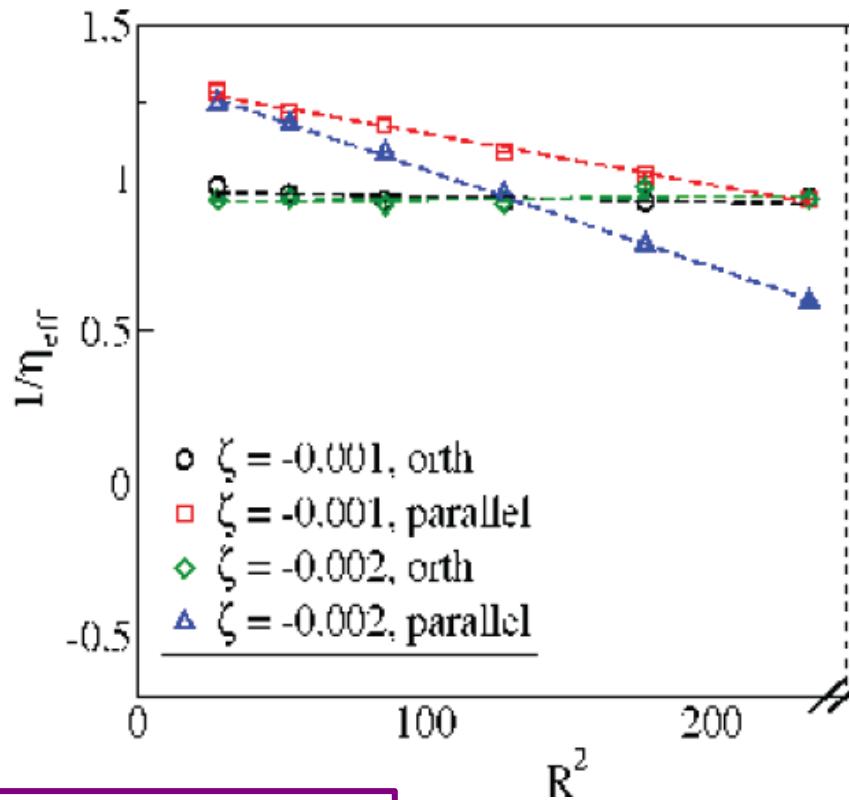
Microrheology of active nematics

Colloid in contractile fluid:

$$\text{net active force} \propto \zeta F R^2 / K$$

$$\frac{1}{\eta_{\text{eff}}} = \frac{1}{\eta_{\text{passive}}} \left(1 - \frac{\alpha \zeta R^2}{K} \right)$$

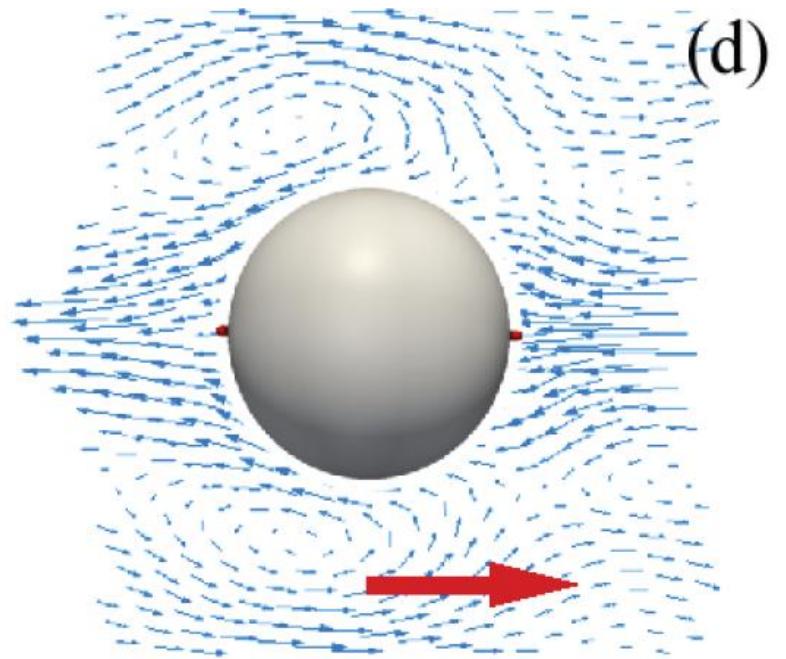
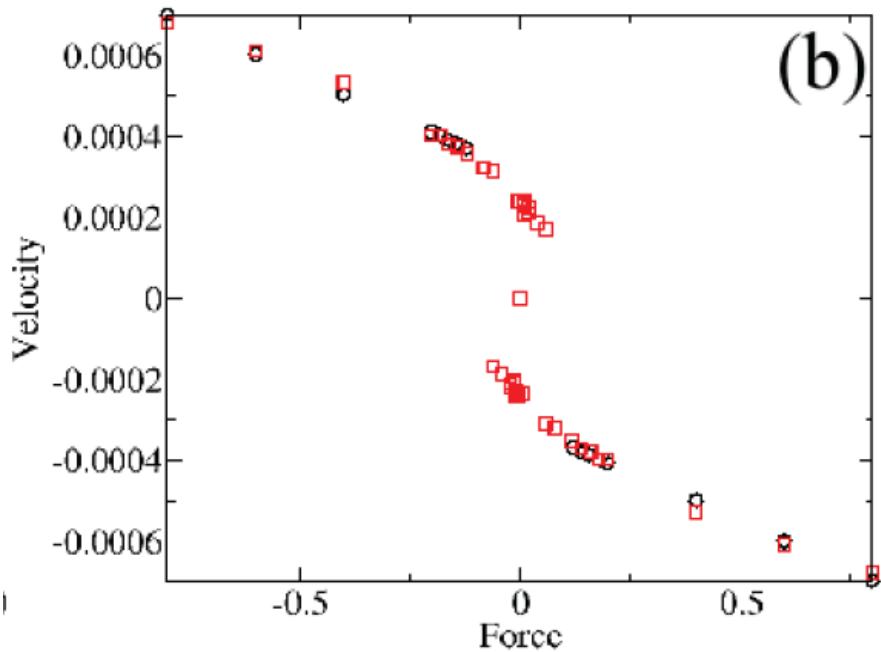
Good fit to data at modest R



Negative drag coefficient at large R ?

Microrheology of active nematics

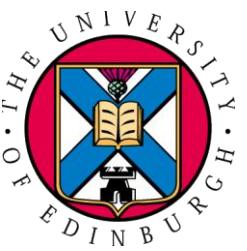
$$\alpha\zeta R^2/K = 3.4$$



Negative drag coefficient at large R

Activity + anchoring creates a packet of fluid moving in opposite direction to \mathbf{F}

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