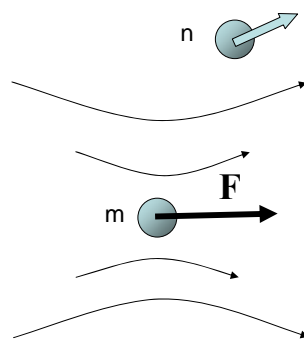


Hydrodynamic Interaction in Electrophoresis

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Toyota Physical and Chemical
Research Institute

Hydrodynamic Interaction



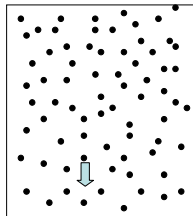
$$\mathbf{v}(\mathbf{r}) = \mathbf{H}(\mathbf{r}) \cdot \mathbf{F}$$

$$\mathbf{H}(\mathbf{r}) = \frac{\mathbf{I} + \hat{\mathbf{r}}\hat{\mathbf{r}}}{8\pi\eta r}$$

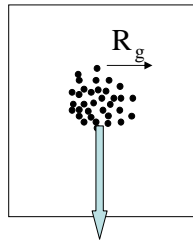
$$\mathbf{V}_n = \sum \mathbf{H}_{nm} \cdot \mathbf{F}_m$$

$$\mathbf{H}_{mn} = \mathbf{H}(\mathbf{r}_{mn})$$

Hydrodynamic interaction is long ranged



$$v_0 = \frac{mg}{6\pi\eta a}$$



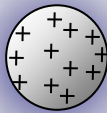
$$v = \frac{Nmg}{6\pi\eta R_g} = v_0 \left(\frac{R_g}{a} \right)^{D-1}$$

$$\mathbf{V}_n = \sum \mathbf{H}_{nm} \cdot \mathbf{F}_m$$

$$\mathbf{H}(\mathbf{r}) = \frac{\mathbf{I} + \hat{\mathbf{r}}\hat{\mathbf{r}}}{8\pi\eta r}$$

Electrophoresis

\mathbf{E}

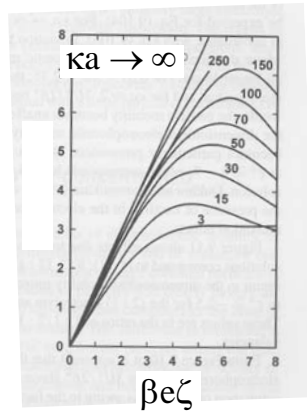
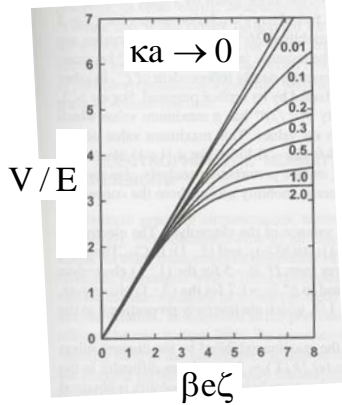


$$\mathbf{V} = \mu\mathbf{E}$$

Electrophoresis of Uniformly Charged Spherical Particle

$$V = \frac{\epsilon\zeta}{\eta} f(\kappa a, \beta e\zeta) E$$

Obrien White 1978



$\kappa a \rightarrow 0$

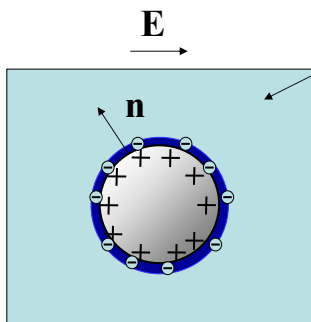
$$\frac{V}{E} = \frac{2\epsilon\zeta}{3\eta}$$

$\kappa a \rightarrow \infty$

$$\frac{V}{E} = \frac{\epsilon\zeta}{\eta}$$

Important limit :thin double layer

$$\kappa a \gg 1$$



Solution can be regarded as a conductive neutral fluid

Stokes equation $\eta \nabla^2 \mathbf{v} = -\nabla p$

$$\nabla \cdot \mathbf{v} = 0$$

Electric field

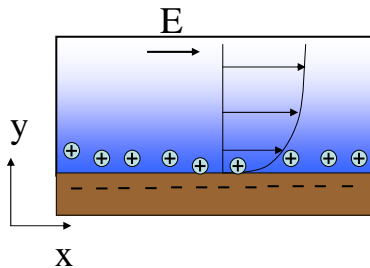
$$\mathbf{J} = -\sigma_b \nabla \psi \Rightarrow \nabla^2 \psi = 0$$

$$\nabla \cdot \mathbf{J} = 0$$

At the outer boundary of the sphere

$$\mathbf{n} \cdot \nabla \psi = 0$$

Smoluchowskii's slip velocity



$$v_x(y) \rightarrow U_s$$

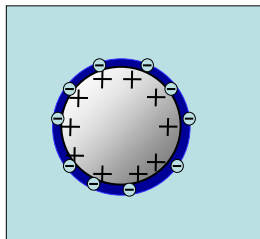
$$U_s = -\frac{\varepsilon\zeta}{\eta} E \quad \text{Slip velocity}$$

Smoluchowskii ~1910

$$\left. \begin{aligned} \eta \frac{d^2 v_x}{dy^2} &= -\sum c_i q_i E \\ \sum c_i q_i &= -\varepsilon \frac{d^2 \psi}{dy^2} \end{aligned} \right\} \eta \frac{d^2 v_x}{dy^2} = \varepsilon \frac{d^2 \psi}{dy^2} E$$

$$\eta(v_x(y) - v_x(0)) = \varepsilon(\psi(y) - \psi(0))E$$

Conventional equations for thin double layer limit



Bulk equation

$$\nabla^2 \phi = 0$$

$$\eta \nabla^2 \mathbf{v} = -\nabla p$$

$$\nabla \cdot \mathbf{v} = 0$$

Boundary condition

$$\mathbf{n} \cdot \nabla \phi = 0$$

$$\mathbf{v} = \mathbf{V} + \frac{\varepsilon\zeta}{\eta} \nabla_t \phi$$

$$\mathbf{V} = \frac{\varepsilon\zeta}{\eta} \mathbf{E}$$

$$\mathbf{E}(\mathbf{r}) = \left[1 - \frac{1}{2} \left(\frac{a}{r} \right)^3 \left(3 \frac{\mathbf{r}\mathbf{r}}{r^2} - \mathbf{I} \right) \right] \cdot \mathbf{E}^\infty$$

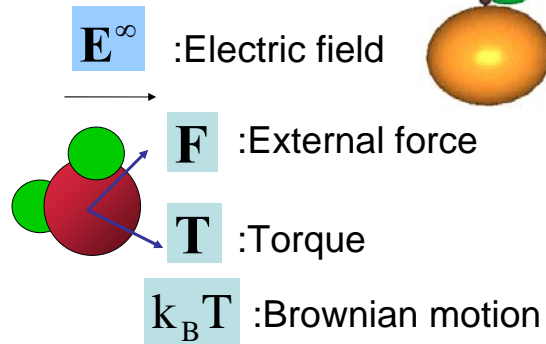
$$\mathbf{v} = \frac{1}{2} \left(\frac{a}{r} \right)^3 \left(3 \frac{\mathbf{r}\mathbf{r}}{r^2} - \mathbf{I} \right) \cdot \mathbf{V}$$

Particle Simulator MIKAN

-Microhydrodynamic Kinetic Analyzer-

Masato Makino

- Boundary element method
- Rigid particles
- Arbitrary shape and charge distribution



Velocity $\mathbf{V} = \mathbf{a} \cdot \mathbf{F} + \mathbf{b} \cdot \mathbf{T} + \mathbf{p} \cdot \mathbf{E}^\infty + \Delta \mathbf{V}$

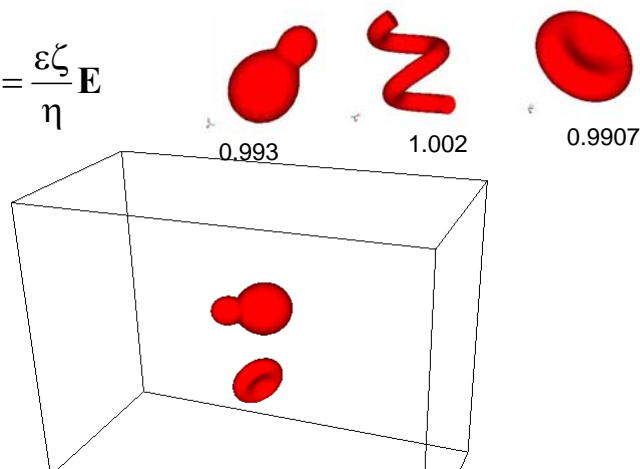
Angular velocity $\mathbf{\Omega} = \mathbf{b} \cdot \mathbf{F} + \mathbf{c} \cdot \mathbf{T} + \mathbf{q} \cdot \mathbf{E}^\infty + \Delta \mathbf{\Omega}$

Simple result for uniformly charged particle

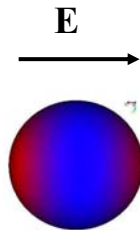
Morrison's theorem Morrison 1970

Uniformly charged particle translates uniformly

$$\mathbf{V} = \frac{\epsilon \zeta}{\eta} \mathbf{E}$$



Complex Result for Non-uniformly Charged Particle

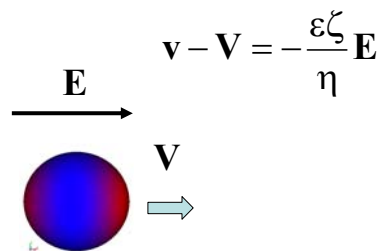
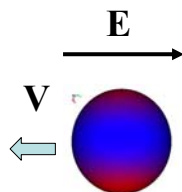
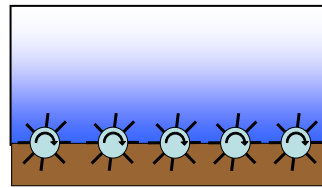
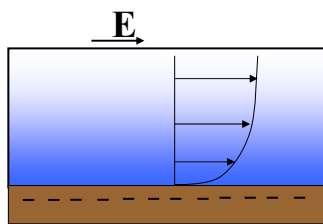


$$\zeta(\theta) = \zeta_0(3\cos^2\theta - 1)$$

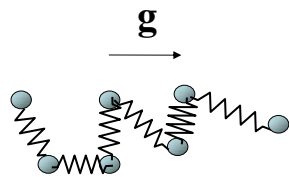
Place particle at origin in random



Surface Slippage Determines the Electrophoresis



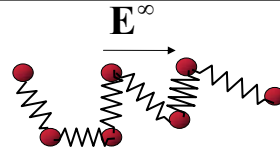
Hydrodynamic interaction in charged polymers



Hydrodynamic interaction for neutral polymers

$$\mathbf{V}_i = \sum_j \mathbf{H}_{ij} \cdot (\mathbf{F}_j + m\mathbf{g})$$

$$V \approx \frac{Nmg}{6\pi\eta R_h} \propto \sqrt{N}$$



Electrophoresis of charged polymers

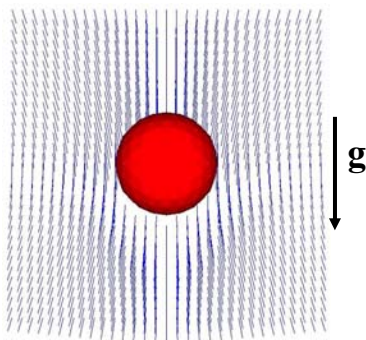
$$\mathbf{V}_i = \sum_j \mathbf{H}_{ij} \cdot (\mathbf{F}_j + q_j \mathbf{E}^\infty)$$

~~$$V \approx \frac{NqE}{6\pi\eta R_h} \propto \sqrt{N}$$~~

Correct answer is

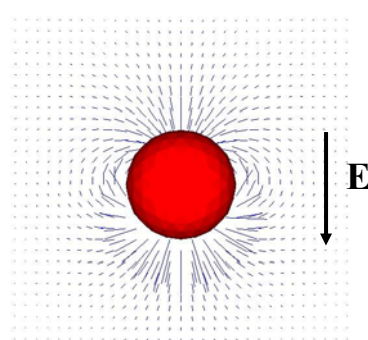
$$V = \frac{\epsilon\zeta}{\eta} E^\infty \propto \text{independent of } N$$

Velocity field in sedimentation and electrophoresis



Sedimentation

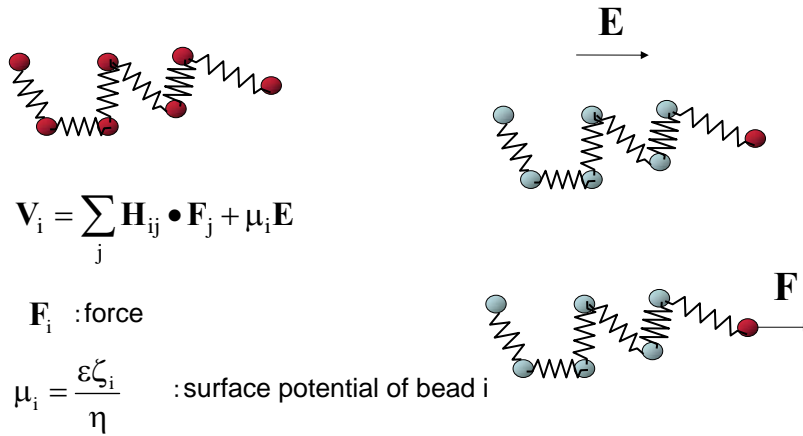
$$v(r) \propto \frac{1}{r}$$



Electrophoresis

$$v(r) \propto \frac{1}{r^3}$$

Rouse Zimm model for charged polymer



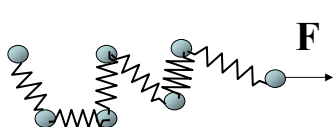
$$V_i = \sum_j \mathbf{H}_{ij} \cdot \mathbf{F}_j + \mu_i \mathbf{E}$$

\mathbf{F}_i : force

$\mu_i = \frac{\epsilon \zeta_i}{\eta}$: surface potential of bead i

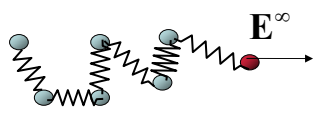
Approximate theory

$$V_i = \sum_j \mathbf{H}_{ij} \cdot \mathbf{F}_j + \mu_i \mathbf{E}^\infty$$



$$V = \frac{1}{\eta a} \sum_{j=0}^N \frac{1}{|i-j|^{1/2}} f_j \quad i = 0, 1, \dots, N$$

$$\sum_{j=0}^N f_j = F \quad V \approx \frac{F}{\eta a \sqrt{N}}$$

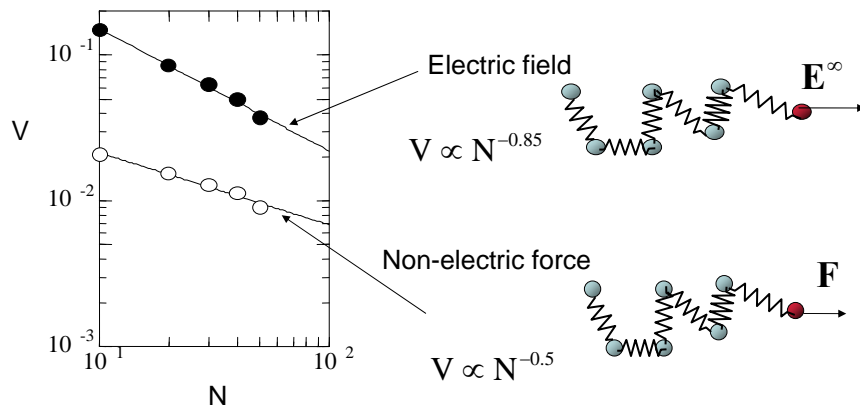


$$V = \frac{1}{\eta a} \sum_{j=0}^N \frac{1}{|i-j|^{1/2}} f_j + \delta_{i0} \mu E$$

$$\sum_{j=0}^N f_j = 0$$

Result of simulation

$$V_i = \sum_j H_{ij} \cdot F_j + \mu_i E^\infty$$



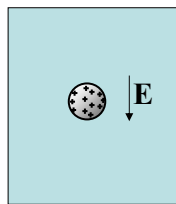
Reciprocal relations

Reciprocal relation

The conventional equation fails to explain the sedimentation potential

Electrophoresis

$$\mathbf{E} \Rightarrow \mathbf{V}$$

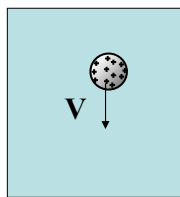


$$\mathbf{V} = \mu \mathbf{E}$$

Sedimentation potential

$$\mathbf{V} \Rightarrow \mathbf{E}$$

$$\mathbf{F} \Rightarrow \mathbf{J}$$



$$\mathbf{J} = \mu \mathbf{F}$$

Electro-hydrodynamic interaction

$$\mathbf{V}_i = \sum_j \mathbf{H}_{ij} \cdot \mathbf{F}_j + \mu_i \mathbf{E}$$

$$\mathbf{J} = \sigma \mathbf{E} + \sum_i \mu_i \mathbf{F}_i$$

