

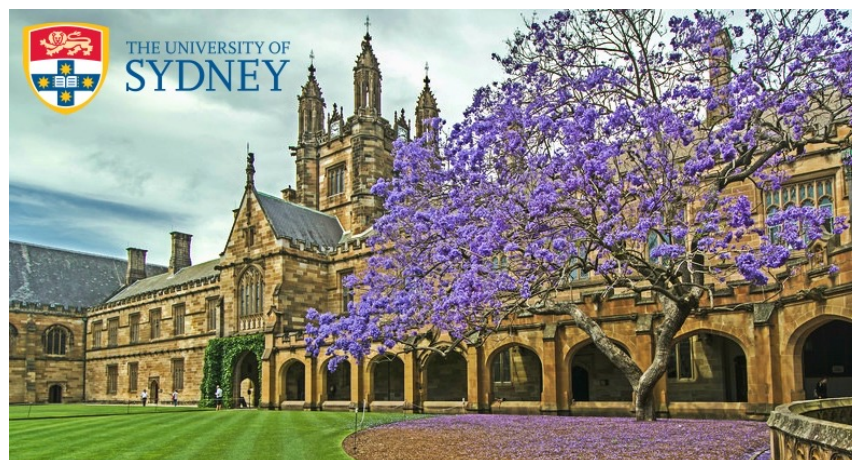
Ligand and Solvent Effects on the Colloidal Stability and Morphology of Apolar Nanoparticles

Asaph Widmer-Cooper

School of Chemistry

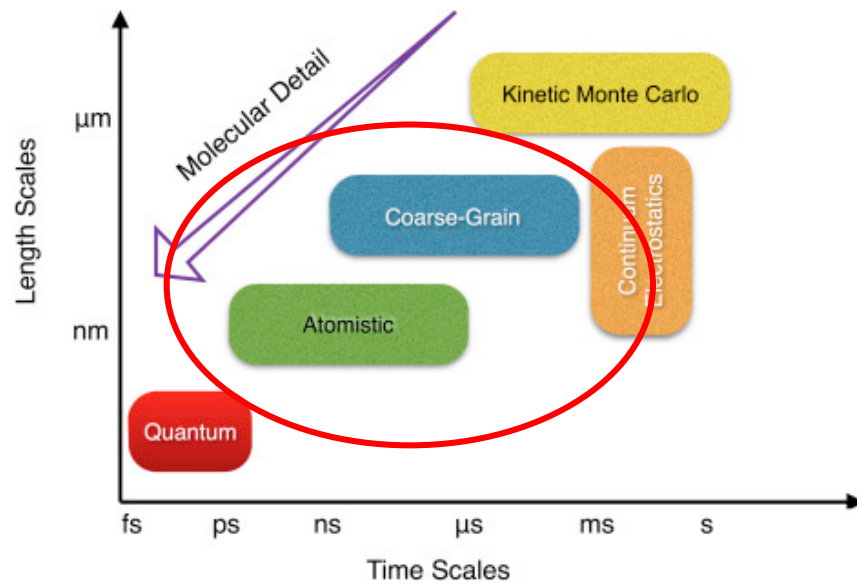


Australian Research Council Centre of Excellence in

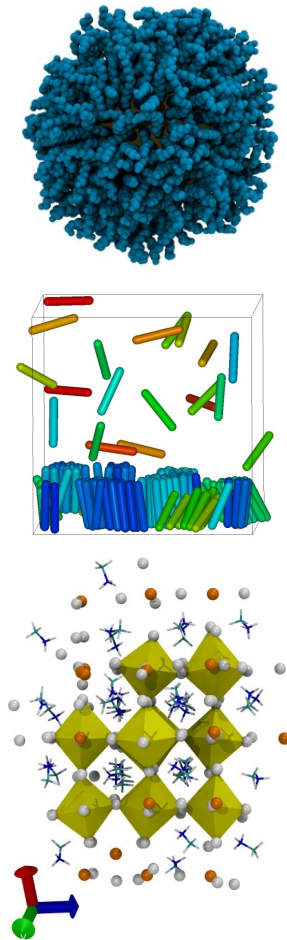


Soft Matter Theory and Simulation

We study how atoms, molecules and nanoparticles spontaneously organise themselves – driven by intermolecular and surface forces and external driving forces – and how that order affects materials properties



Molecular dynamics
Monte Carlo simulations
Free energy methods
Dissipative particle dynamics
Some quantum mechanics and continuum theory



ARC Centre of Excellence in Exciton Science

Australian Research Council Centre of Excellence in



>200 CIs, AIs, PIs,
ECRs, students

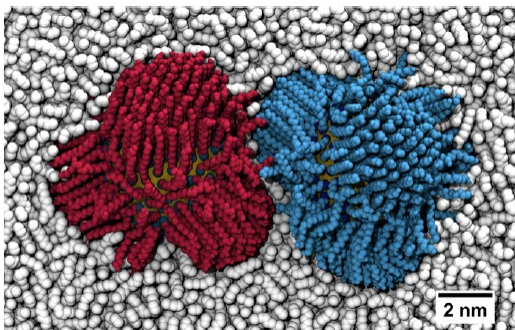


CSIRO
Defence Science Technology Group
Reserve Bank of Australia

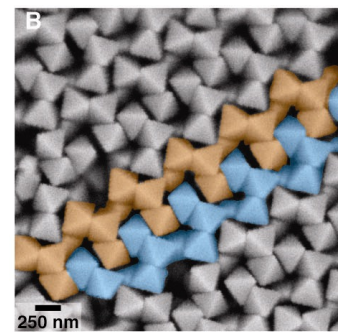
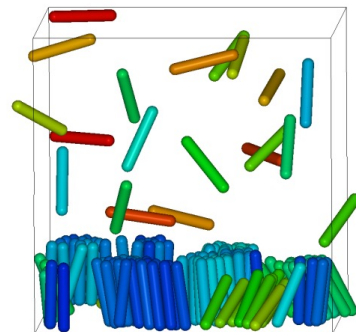
“Next-Generation Light Harvesting Materials for a Sustainable Energy Future”

Formation and Properties of Nanomaterials

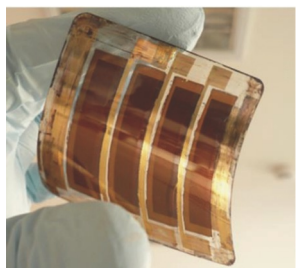
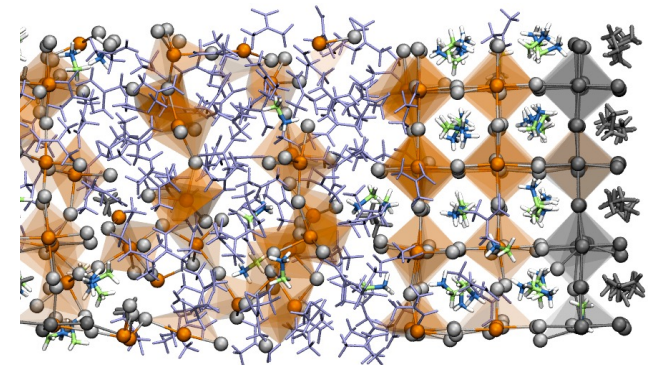
Colloidal properties of nanoparticles



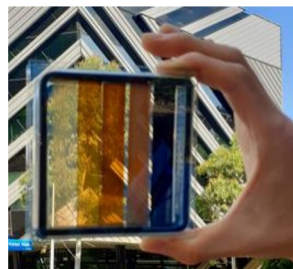
Nanoparticle assembly



Formation and stability of metal halide perovskites



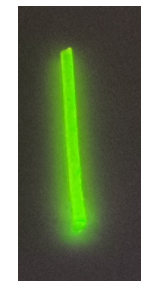
Printable solar



Solar windows



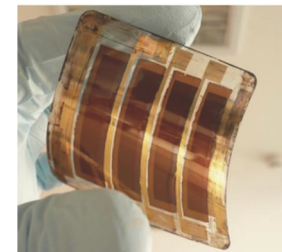
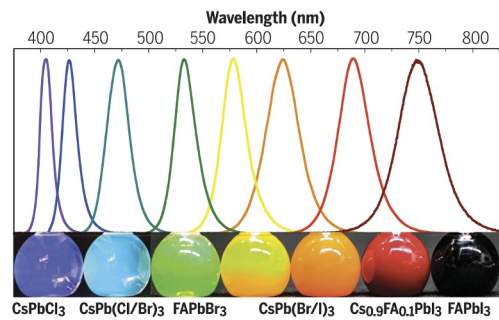
Luminescent solar concentrators



LEDs and lasers

Outline

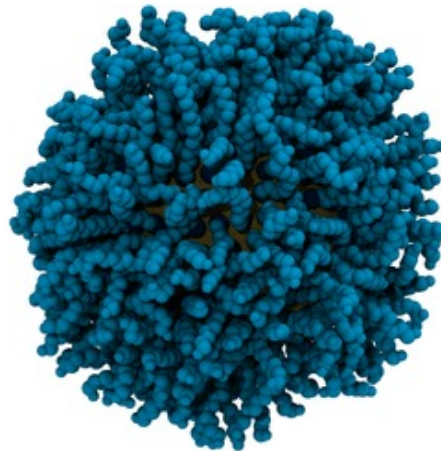
- 1) Colloidal stability of apolar nanoparticles
- 2) Morphology of nanoplatelets



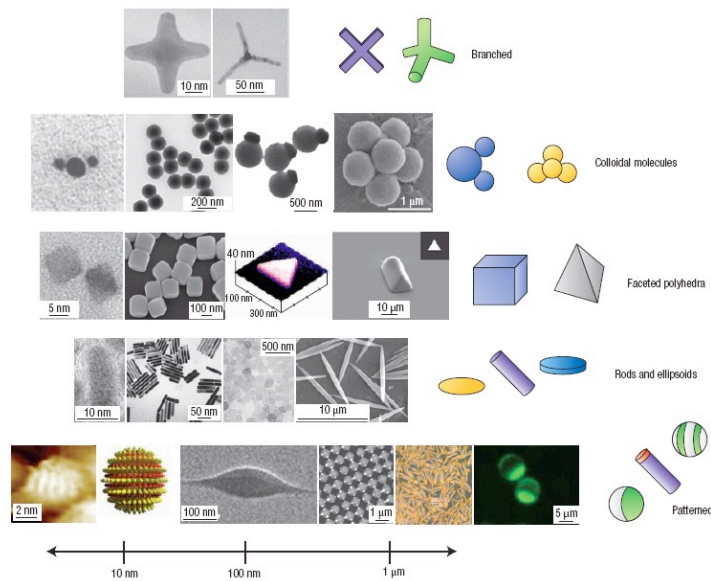
1) Colloidal Stability of Apolar Nanoparticles

Debora Monego, Rashed Hasan

Tobias Kraus, Paul Mulvaney



Colloidal stability is essential for processing and applications of nanoparticles



Glotzer & Solomon *Nature Mat.* 6, 2007



Long-term stability

Dense inks for printing

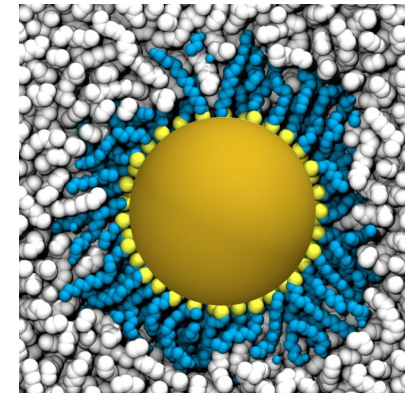
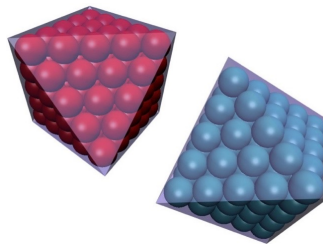
Ordered assembly

How particles interact is a key factor in stability and assembly

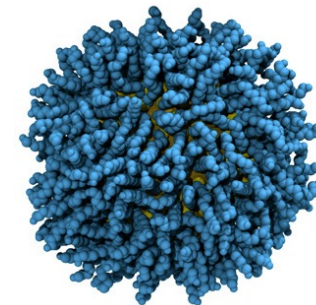
Ligand-coated (apolar) nanoparticles

Inorganic core + organic ligands
- e.g. metals, semiconductors, oxides

Bare cores would agglomerate

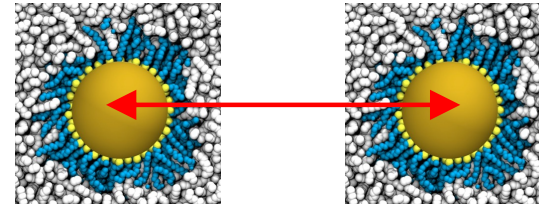


Add surface active molecules during synthesis or via exchange

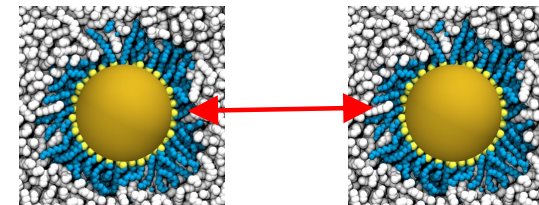


Classical colloid theory approach

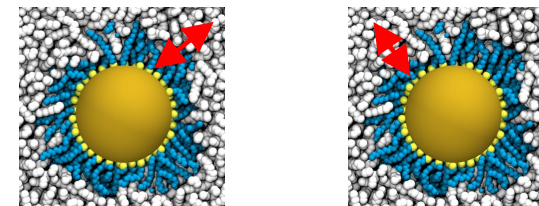
1) Core-core → Hamaker-Lifshitz theory



2) Ligand steric → Alexander-de Gennes theory



3) Ligand-solvent → Flory-Huggins theory

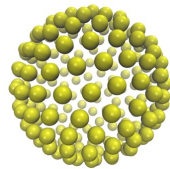


Langmuir 25 (2009) 13861

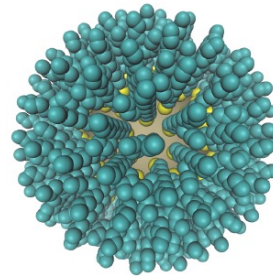
Fails in many cases

Our approach: SAXS + MD simulations

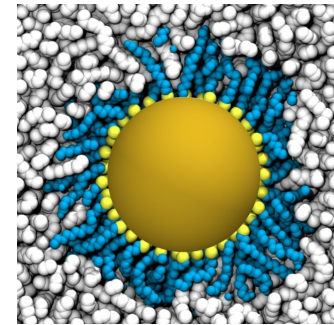
4-10 nm
spherical
core



3.6-5.5 ligands/nm²



add tails (e.g. SC₁₆)



add solvent

- United atom model (CH_x beads) for the ligand and solvent (TraPPE)

Martin and Siepmann *J. Phys. Chem. B* 102 (1998) 2569

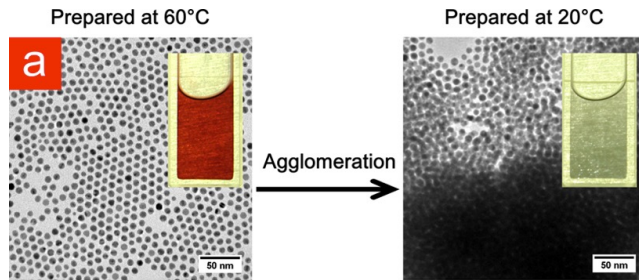
$$u(r_{ij}) = 4\epsilon_{ij} \left[\left(\frac{\sigma_{ij}}{r_{ij}} \right)^{12} - \left(\frac{\sigma_{ij}}{r_{ij}} \right)^6 \right] \quad u_{\text{bend}} = k_{\theta}(\theta - \theta_0)^2/2 \quad u_{\text{tors}} = c_1[1 + \cos \phi] + c_2[1 - \cos(2\phi)] + c_3[1 + \cos(3\phi)]$$

- Hamaker potential for core-core and core-CH_x interactions
- NPT or NVT molecular dynamics at the correct bulk solvent density

ACS Nano 12 (2018) 5969

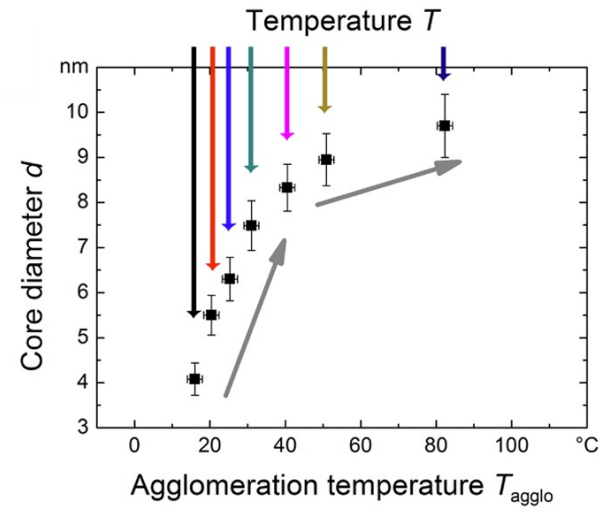
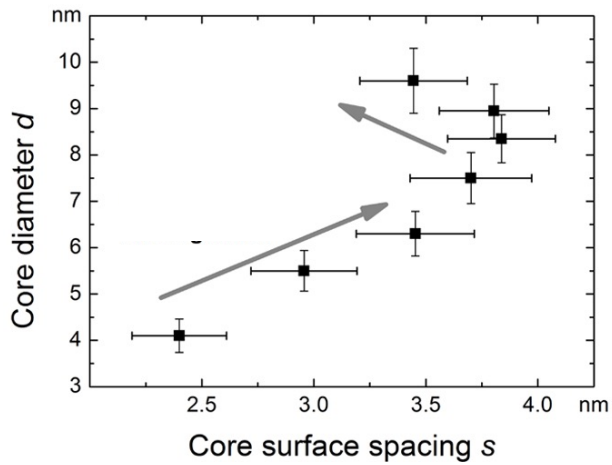
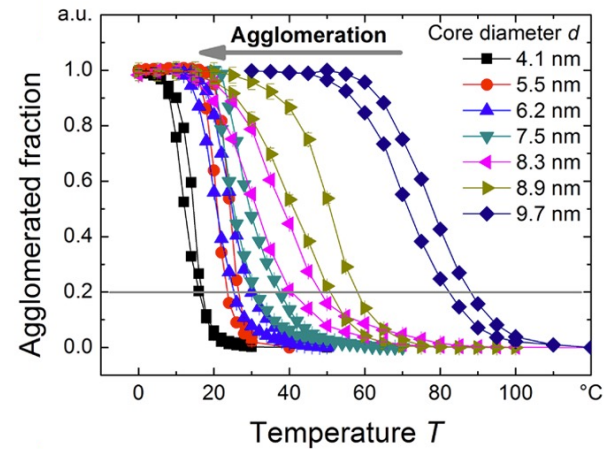
Changing core diameter: Au-SC₁₆ in decane

ACS Nano 12 (2018) 5969

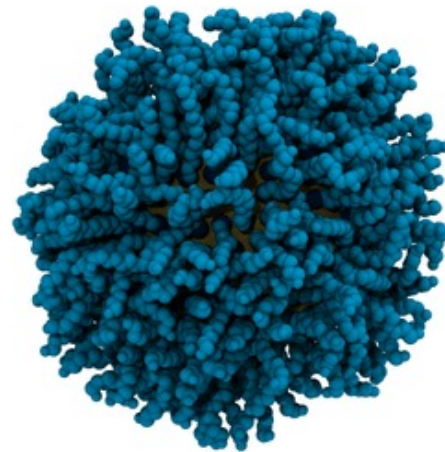


What is responsible for these trends?

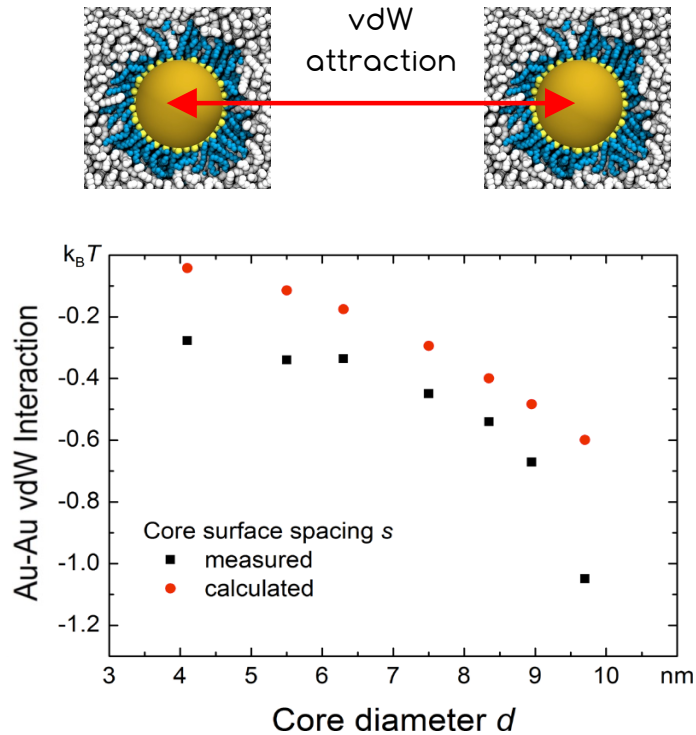
SAXS – Tobias Kraus' group



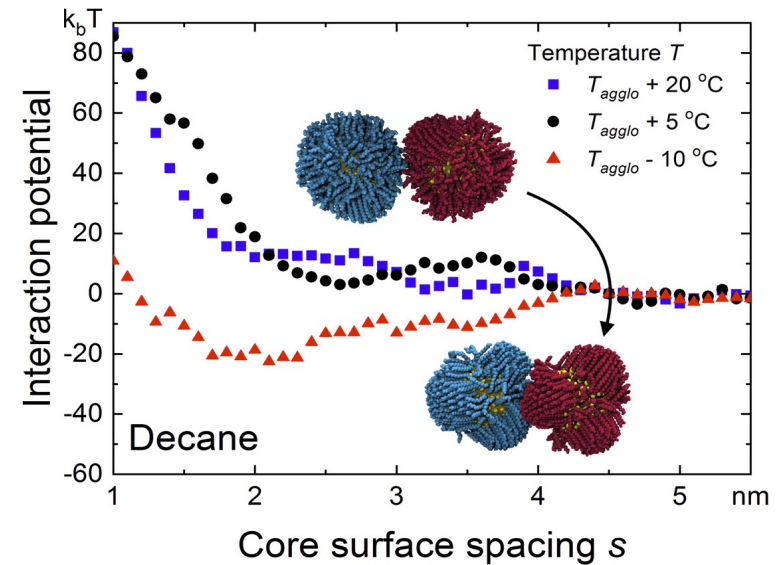
The conformational state of the ligand shell can change



There is a transition from core- to shell-driven agglomeration



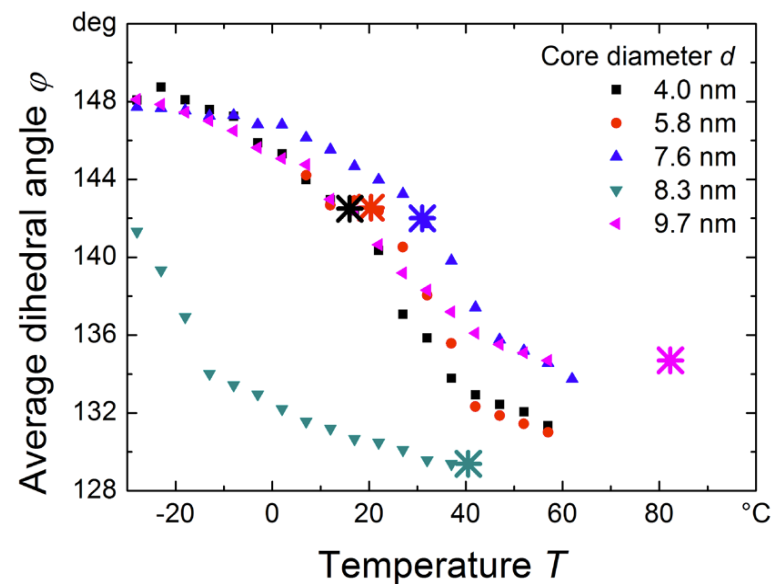
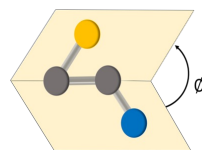
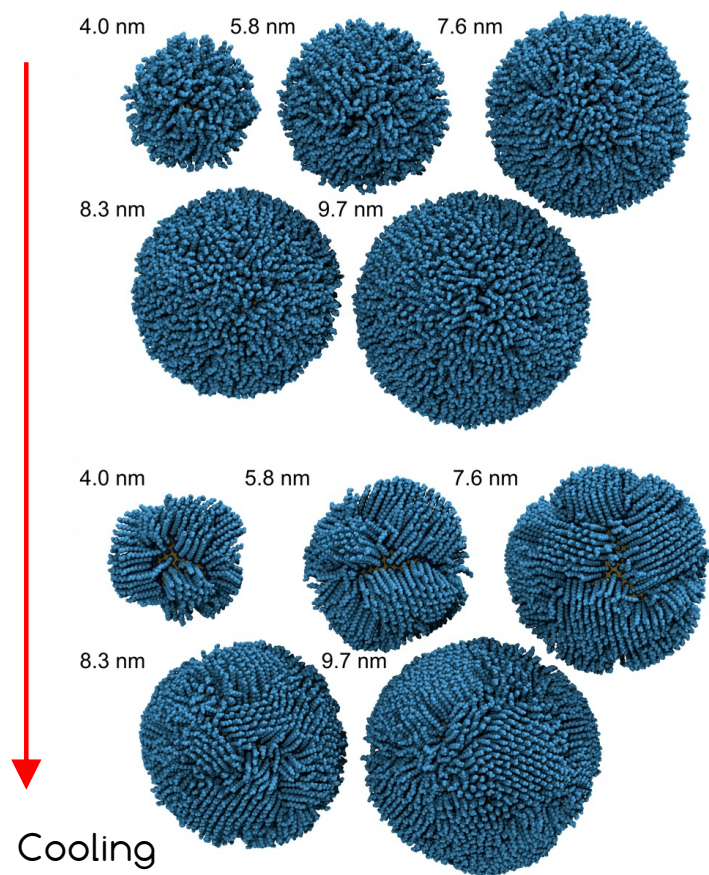
vdW attraction between cores is too weak to drive agglomeration except for the largest particles



In this case, ordering of the ligand shell can drive agglomeration

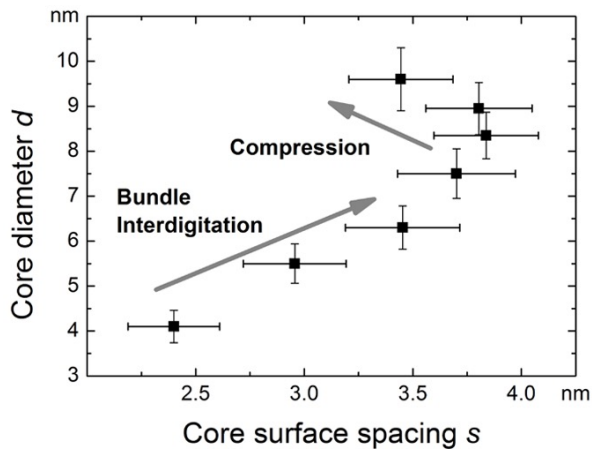
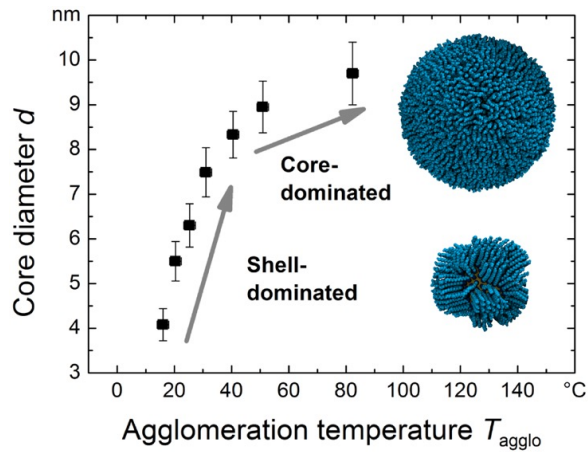
$$\phi_{MF}(r) = \int_r^\infty F_{mean}(s) ds$$

Smaller particles agglomerate when the ligands order



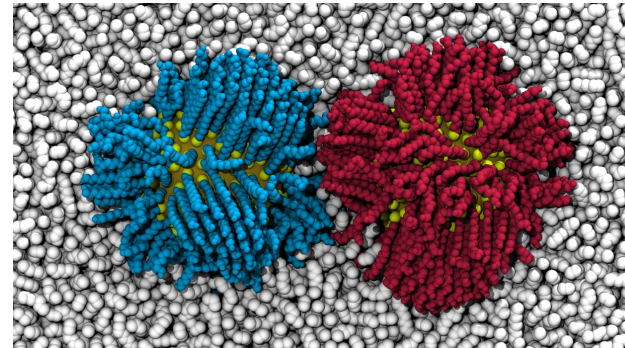
Ligand ordering precedes particle agglomeration (stars) only for particles smaller than ~9 nm

The ligand ordering transition ...



Explains:

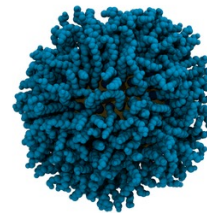
- non-linear dependence of T_{agglo} on d
- inversion of particle spacing



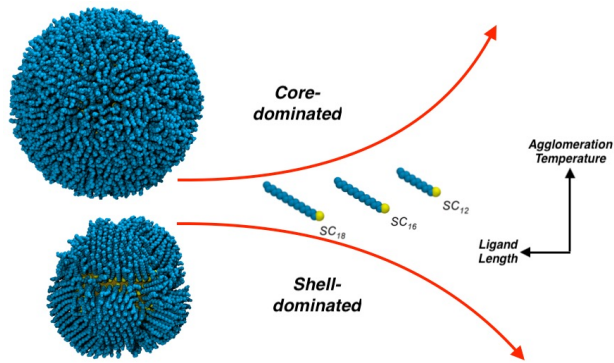
ACS Nano 12 (2018) 5969

Earlier work (with Phill Geissler):

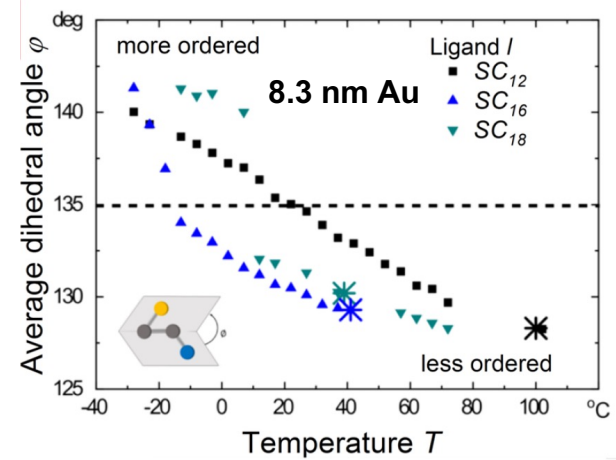
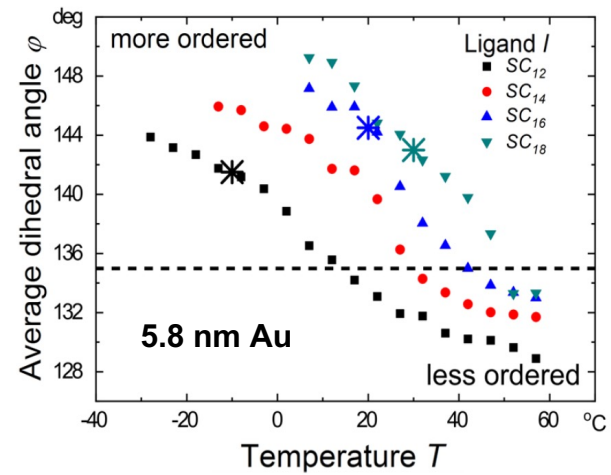
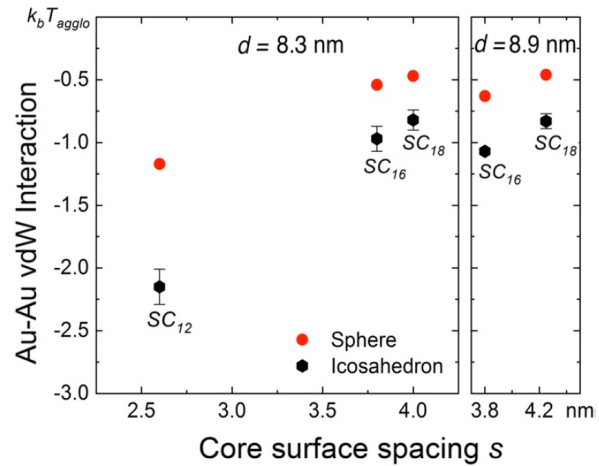
Nano Lett. 14 (2014) 57
ACS Nano 10 (2016) 1877



... inverts the effect of changing the ligand length

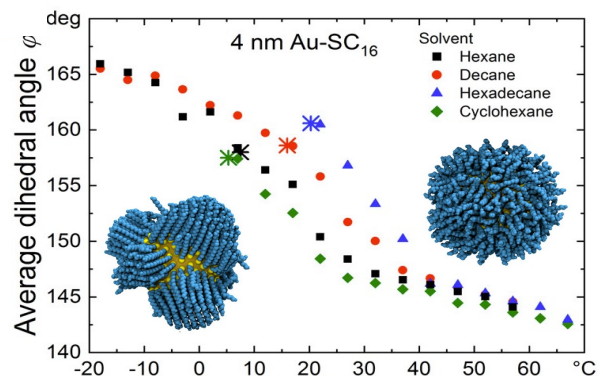
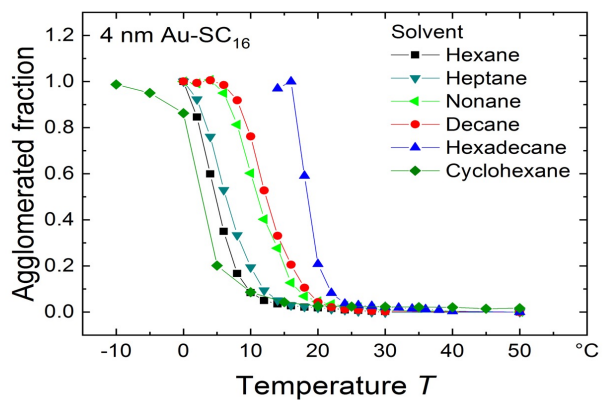


Langmuir 34 (2018) 12982



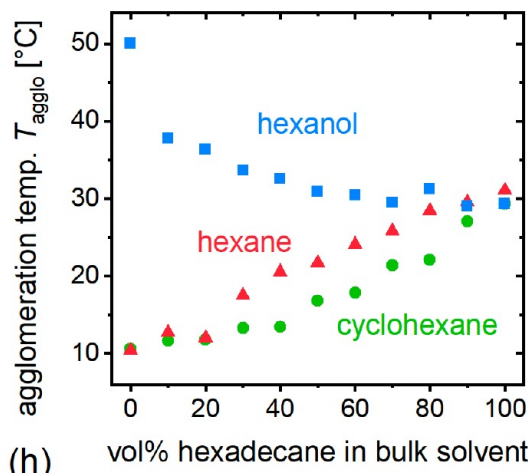
... and explains surprising solvent effects

Stability trends inconsistent with simple solvation models

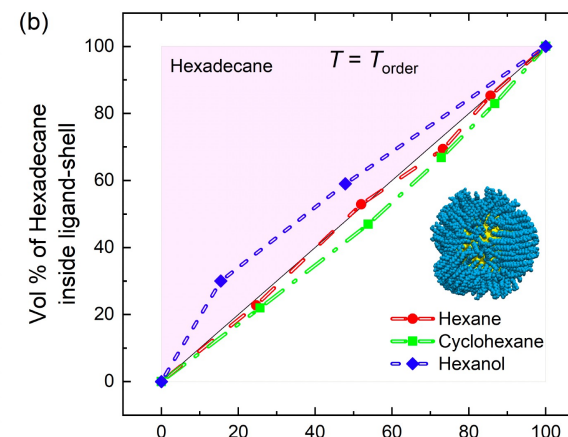


ACS Nano 14 (2020) 5278

Non-linear stability trends in solvent mixtures



(h)



ACS Nano (2023) to appear

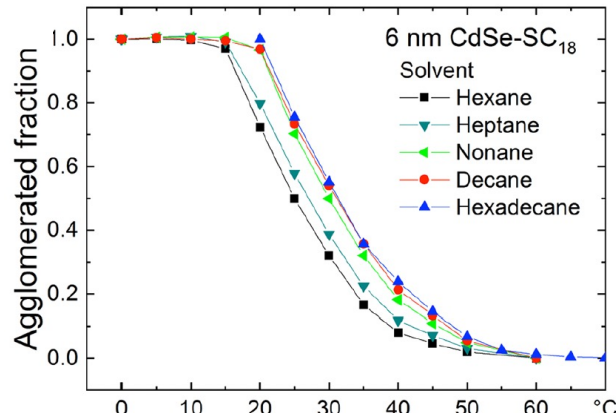
Rashed Hasan



Effect of changing solvent structure

SAXS
Tobias Kraus'
group

6 nm
CdSe-SC₁₈

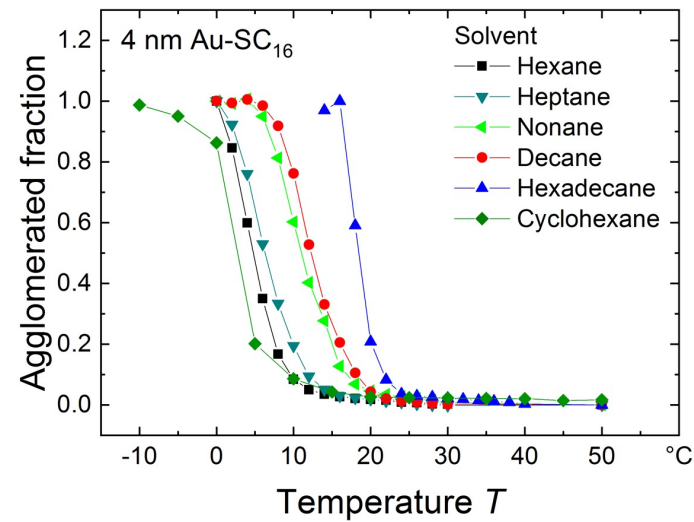
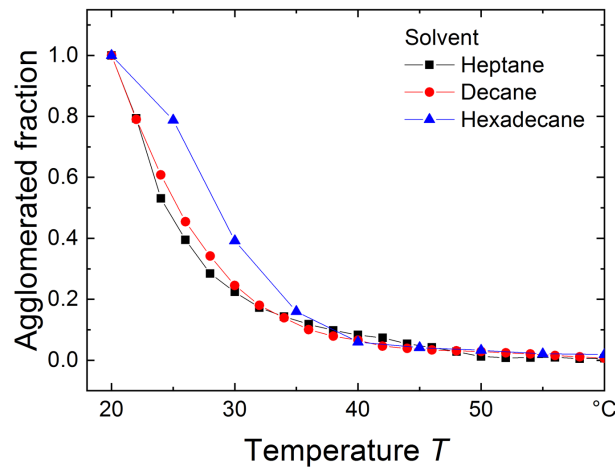


Opposite to 'like dissolves like' – why?

Solvent	Hildebrand Parameter ($Pa^{1/2}$)
Hexane	14900
Decane	15800
Hexadecane	16300
Cyclohexane	16732

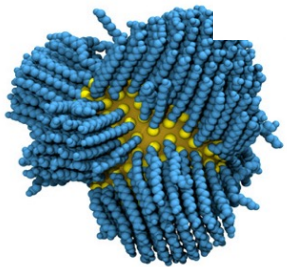
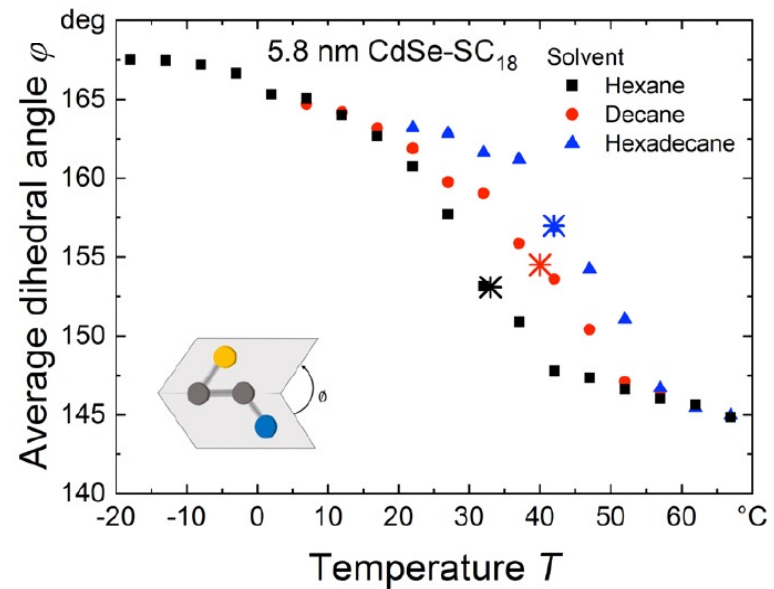
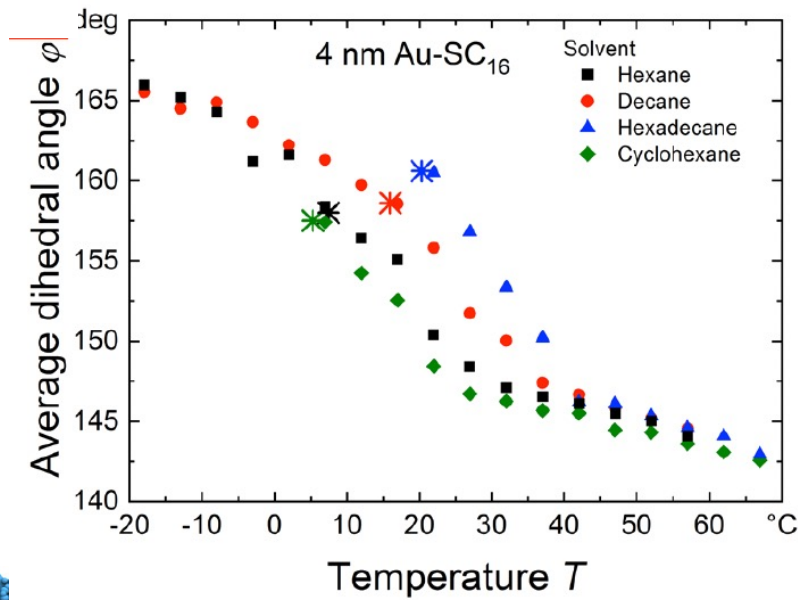
$$\chi = \frac{V_S}{RT}(\delta_L - \delta_S)^2 + 0.34$$

7.4 nm
Au-SC₁₆



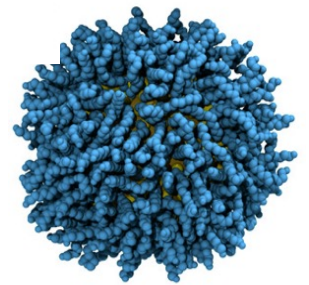
4 nm
Au-SC₁₆

Ligand ordering and particle agglomeration exhibit the same trend



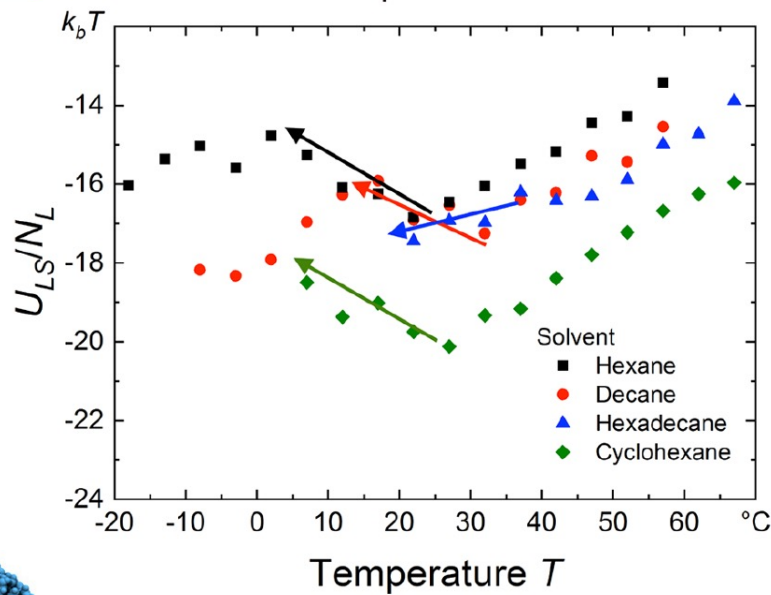
Why do the ligands order at different temperatures?

What is affecting the free energy difference?

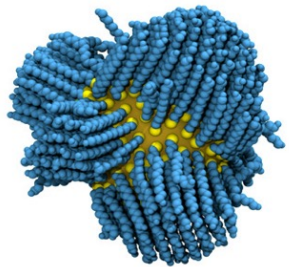
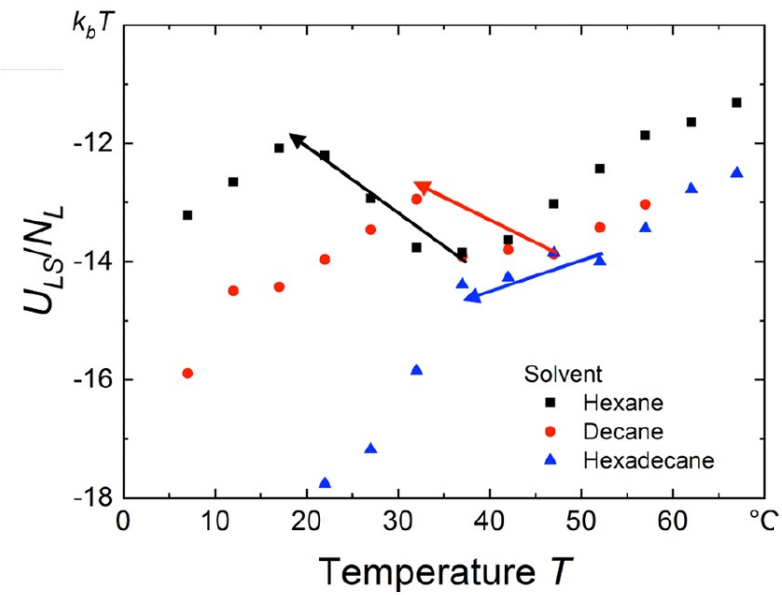


The ligand-solvent interactions differ

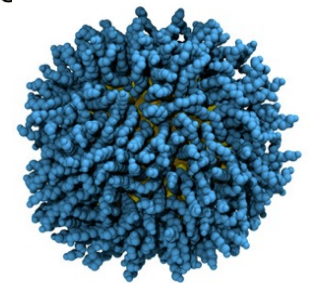
4 nm Au-SC₁₆



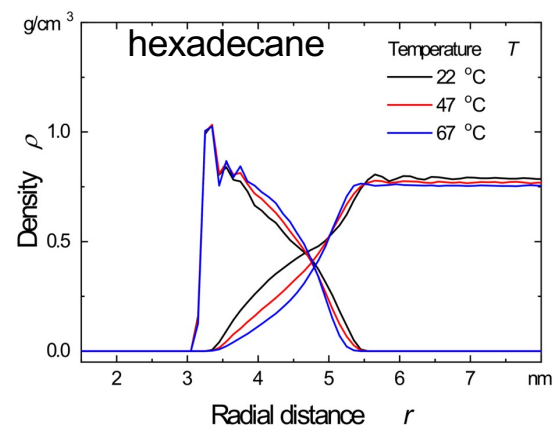
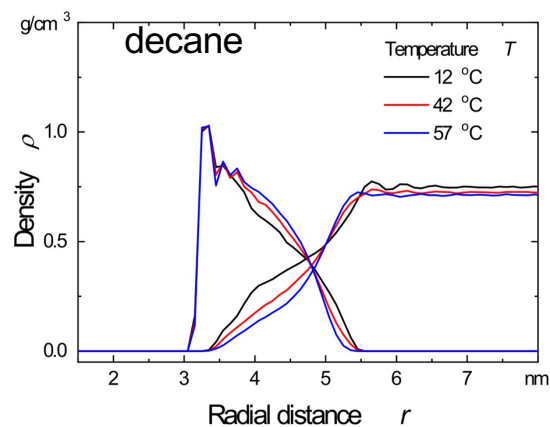
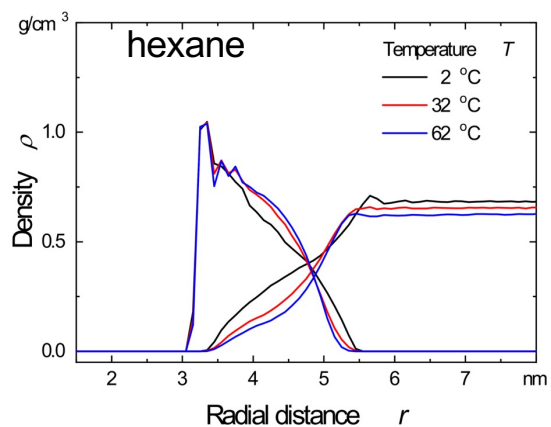
5.8 nm CdSe-SC₁₈



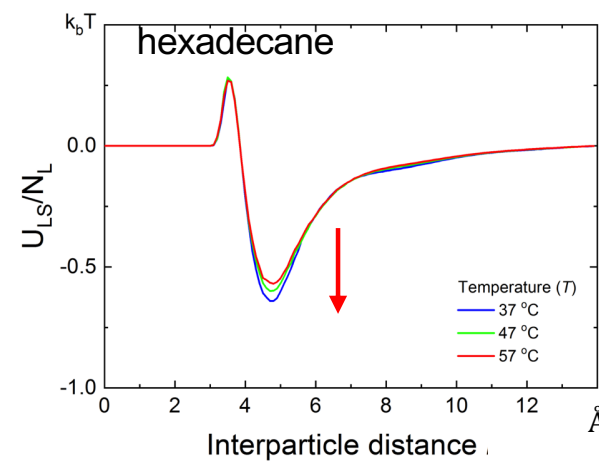
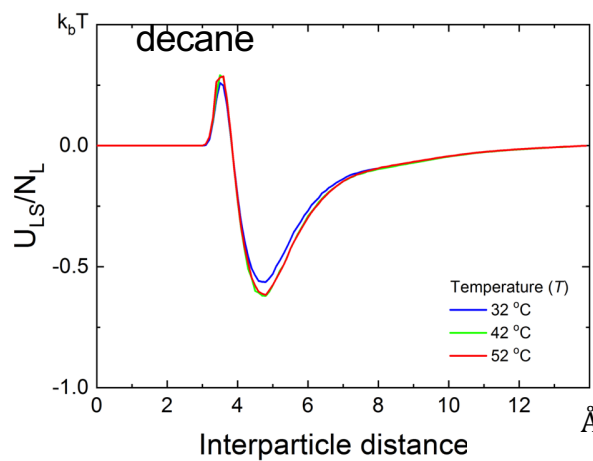
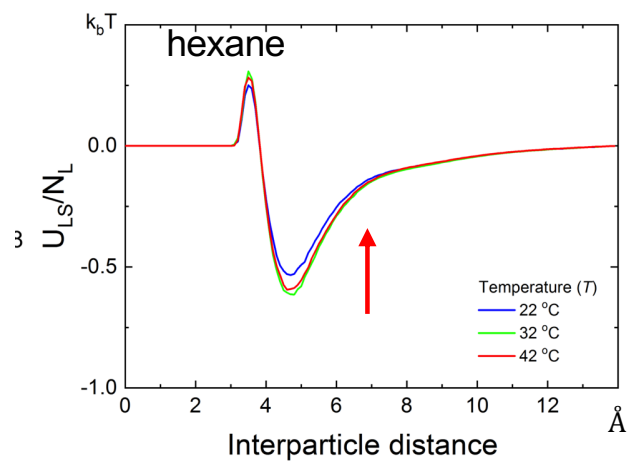
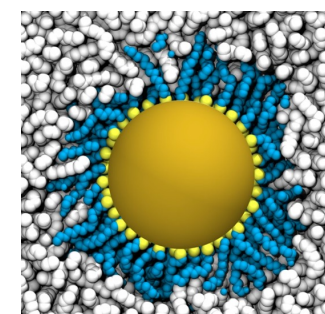
Hexadecane is energetically stabilizing the ligand ordered state
– why?



Due to details of ligand-solvent packing



5.8 nm CdSe-SC₁₈



The longer chain solvents pack better with the ligands in the ordered state

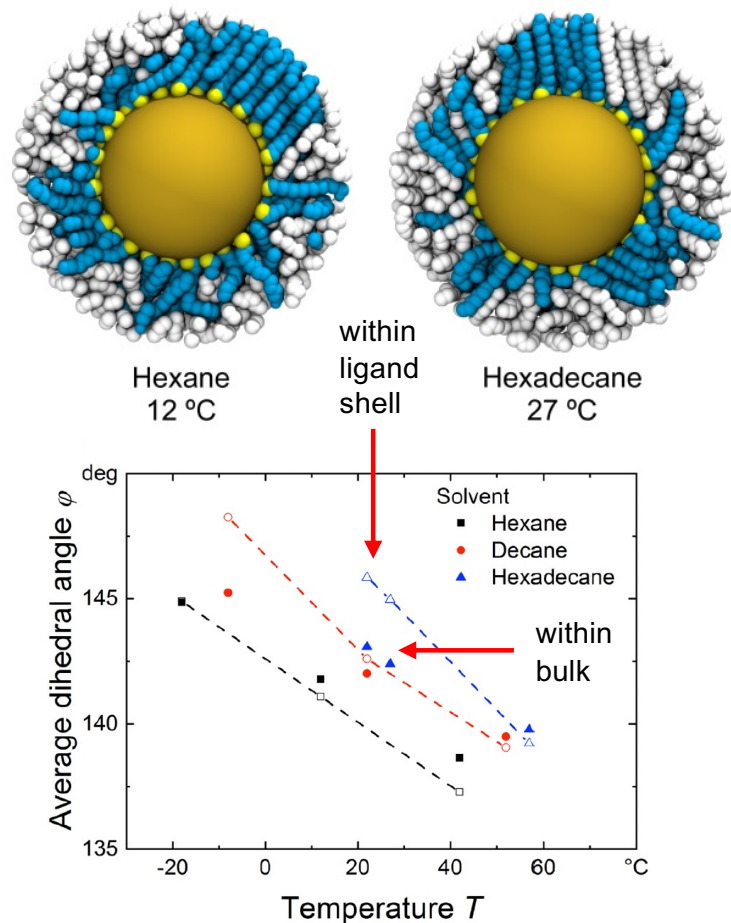
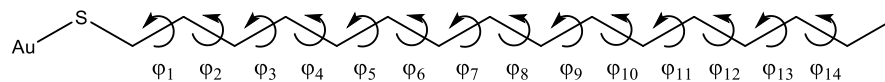


Table 1. Major Enthalpic and Entropic Contributions to the Difference in Free Energy between the Ordered and Disordered Ligand States^a $k_b T/\text{ligand}$

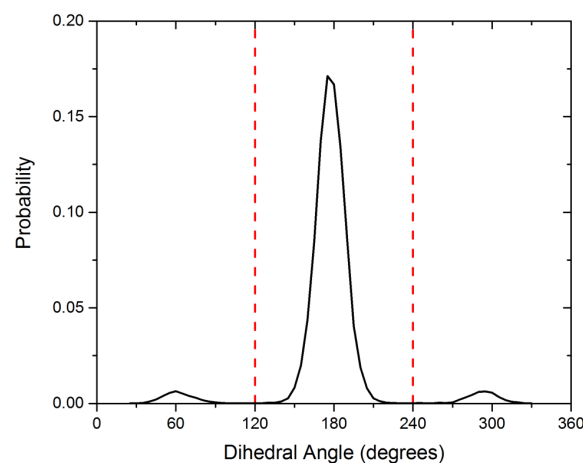
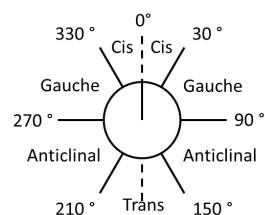
4 nm Au-SC ₁₆			
solvent	hexane	decane	hexadecane
T_{order}	290	295	300
ΔU_{LL}	-3.8	-3.7	-3.5
ΔU_{LS}	4.0	2.3	1.0^b
ΔU_{dih}^{lig}	-2.6	-2.3	-1.8
$-T\Delta S_{mix}$	0.32	0.22	0.14
$-T\Delta S_{conf}^{lig}$	2.5	3.0	2.5
$-T\Delta S_{conf}^{solv}$	0.00	0.00	0.30
5.8 nm CdSe-SC ₁₈			
solvent	hexane	decane	hexadecane
T_{order}	300	310	320
ΔU_{LL}	-4.4	-3.5	-3.1
ΔU_{LS}	4.3	2.6	1.3
ΔU_{dih}^{lig}	-1.9	-1.8	-1.3
$-T\Delta S_{mix}$	0.54	0.37	0.25
$-T\Delta S_{conf}^{lig}$	2.2	2.8	2.2
$-T\Delta S_{conf}^{solv}$	-0.17	0.12	-0.09

Estimating conformational entropy of ligands and solvent

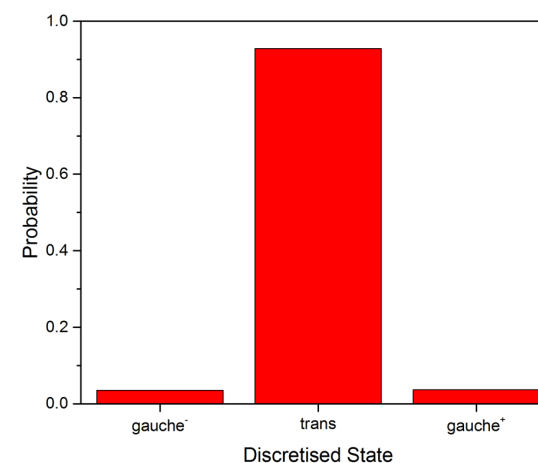
$$S = -k_B \sum_r p(r) \ln[p(r)]$$



- Treat each ligand independently (will overestimate more in ordered state)
- Discretise data



discretisation

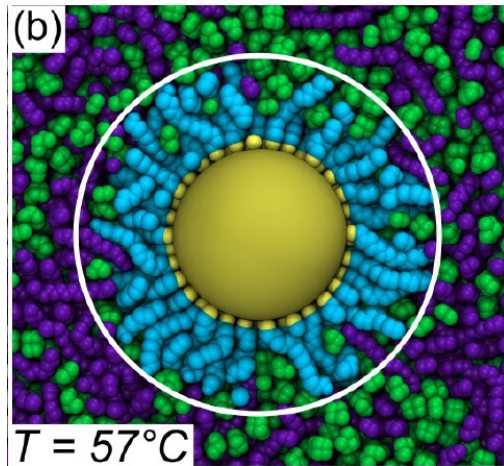


- # of ligand microstates = $3^{14} \approx 10^7$
- Estimate using the Cencalc CC-MLA method (converges rapidly)

JCP 137 (2012) 084115

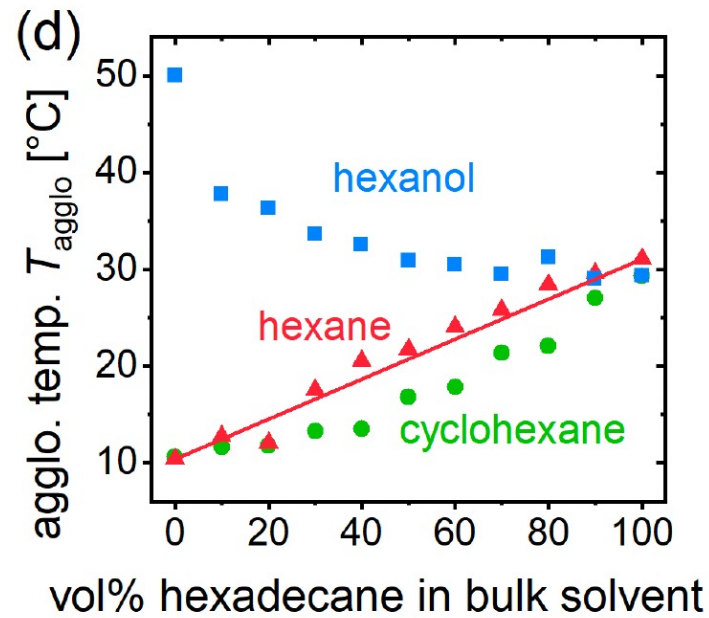
Effect of mixing solvents – solvent engineering

4 nm
Au-SC₁₆

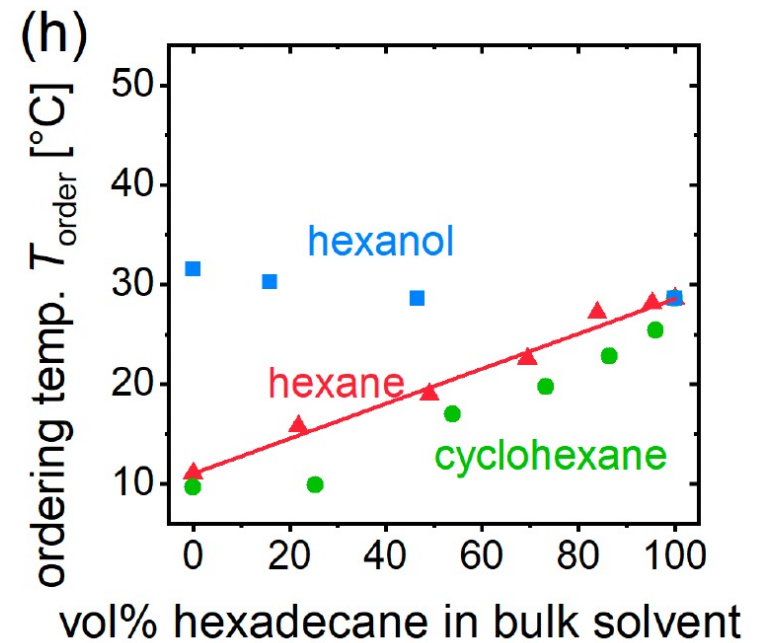


- 1) Hexadecane + hexane
- 2) Hexadecane + cyclohexane
- 3) Hexadecane + hexanol

SAXS

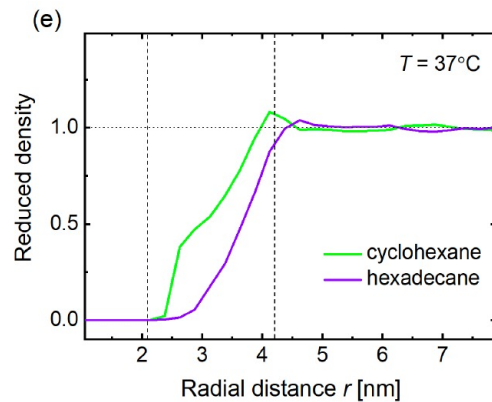
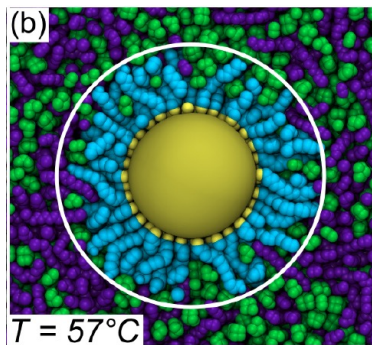
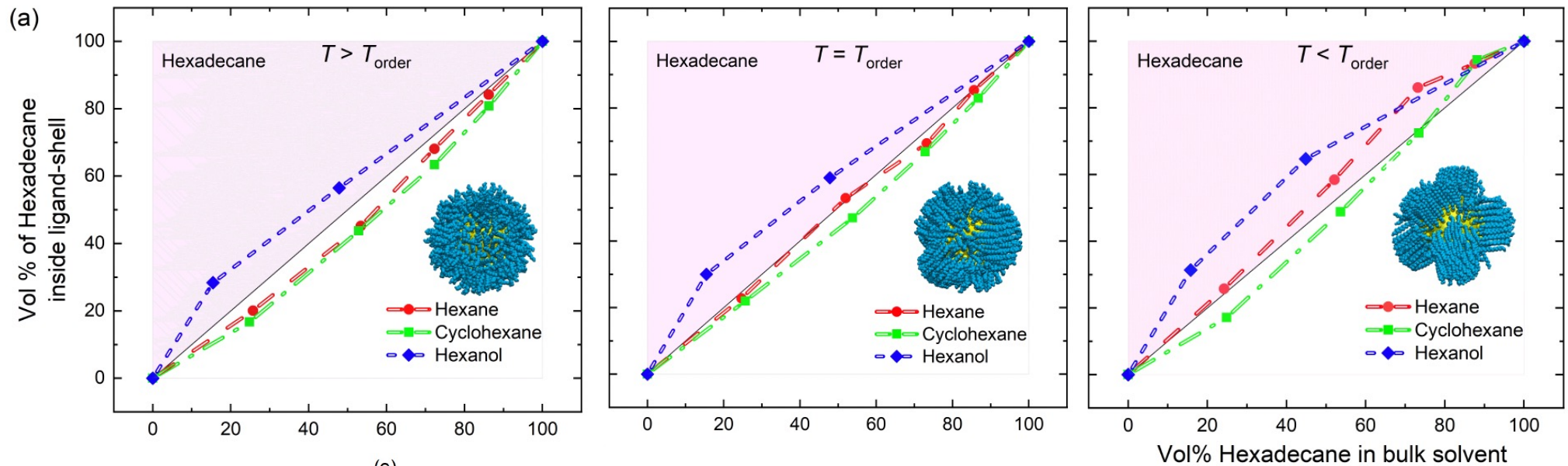


Simulation



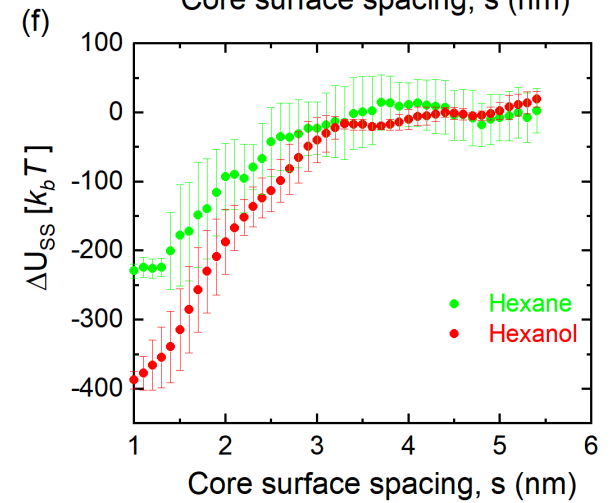
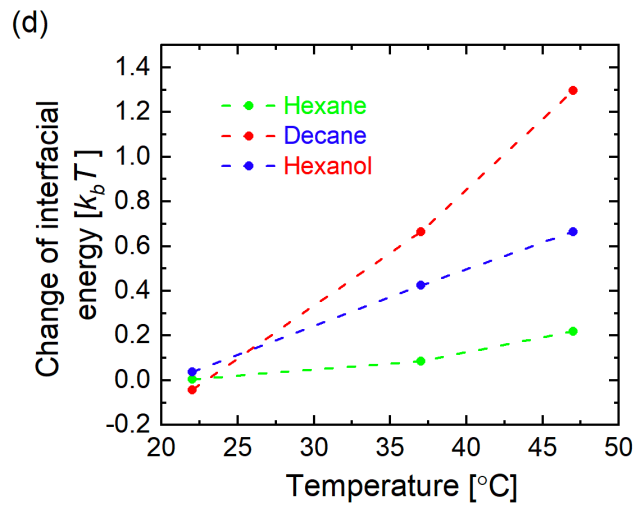
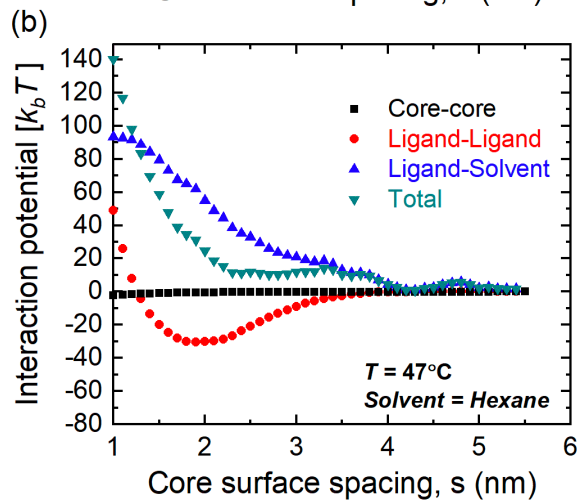
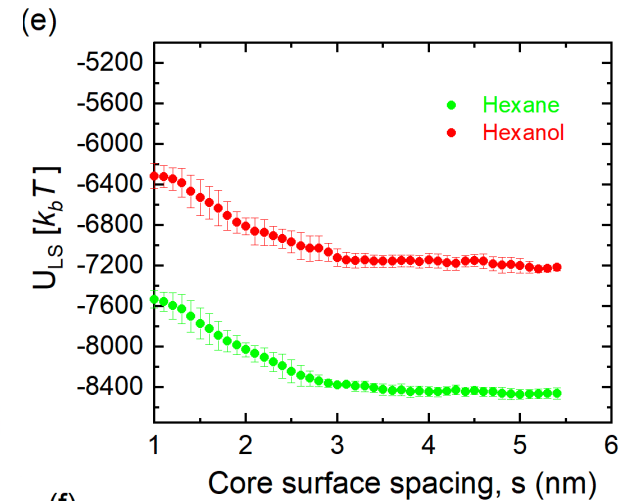
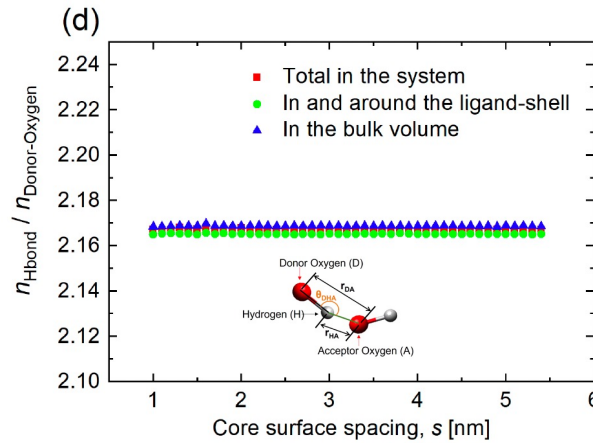
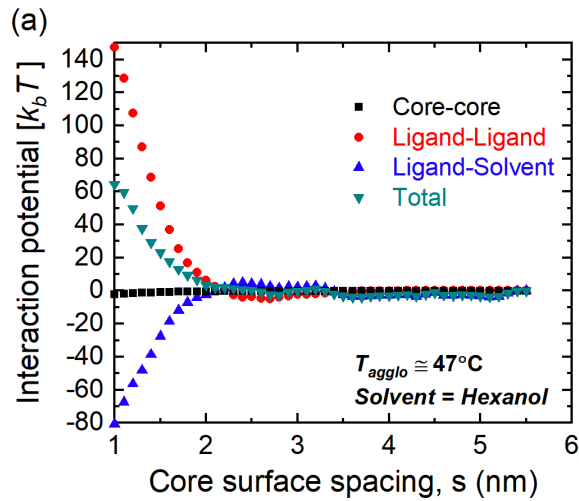
What is responsible for the non-linear trends?

Solvent composition near ligand shell can differ from the bulk



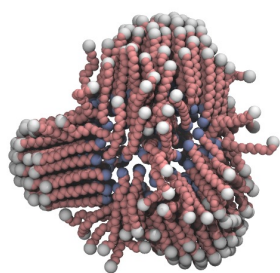
Ligands see solvent enhanced in one component, which can differ depending on the ligand state

Why is $T_{\text{aggl}} \gg T_{\text{order}}$ in pure hexanol?

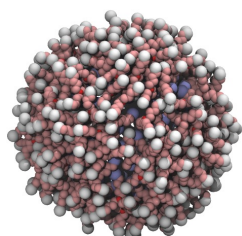


Ligands that inhibit ordering increase stability

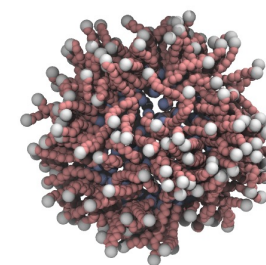
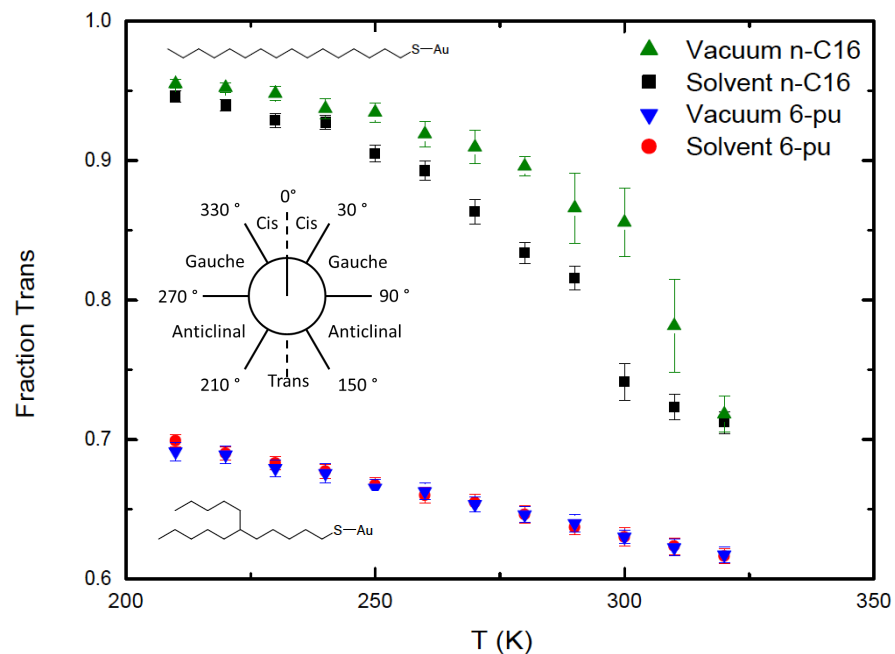
See e.g. Yang et al. *Nano Lett.* 16 (2016) 2127



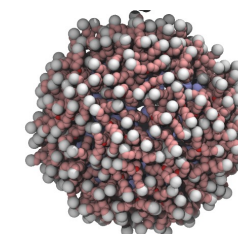
linear



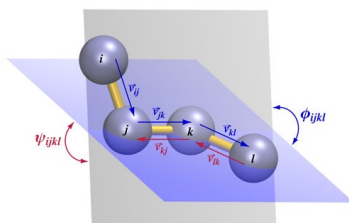
branched



linear



branched

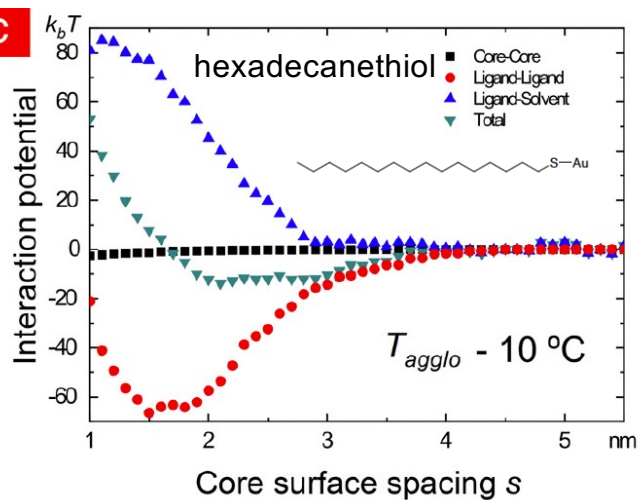


Changing the branch point or length yields similar results, as does adding a *cis* double bond

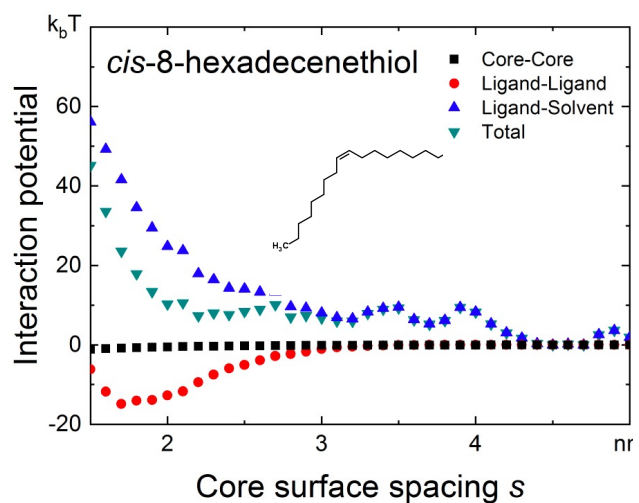
Ligands that inhibit ordering increase stability

4 nm AuNP in hexane

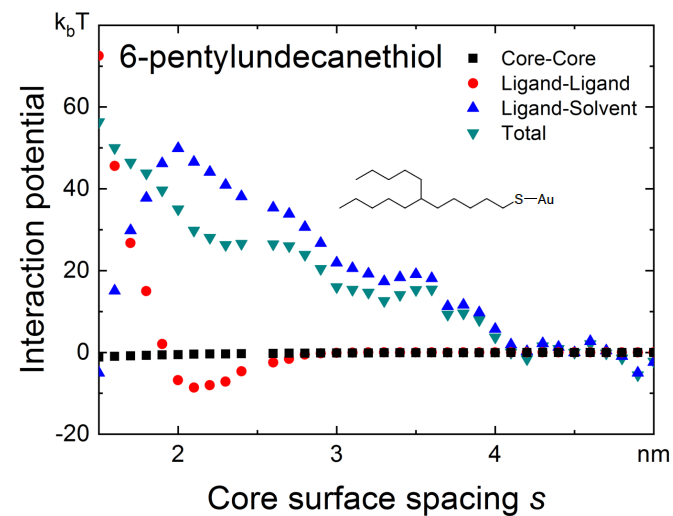
5 °C



-53 °C



-53 °C



Conclusions (1)

The ability of the ligand shell to order (or not) can explain colloidal stability in cases where classical colloid theory fails, including:

- Transition from core- to shell-dominated agglomeration
- Inversion of the effect of ligand length on stability
- Solvent trends inconsistent with simple solvation models
- Dramatic dependence on ligand structure



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Leibnitz Institute for New Materials

- Prof Tobias Kraus
- Thomas Kister
- David Doblas

University of Melbourne

- Prof Paul Mulvaney
- Nick Kirkwood

University of Sydney

- Debora Monego
- Rashed Hasan
- James Smith
- Leo Jiang

2) Morphology of nanoplatelets

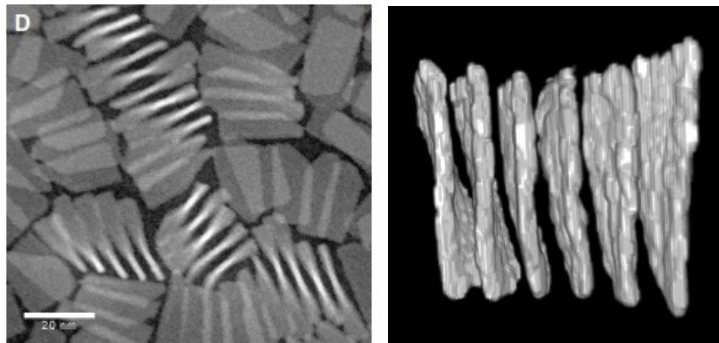
Asaph Widmer-Cooper

Marion Krapez, Debora Monego

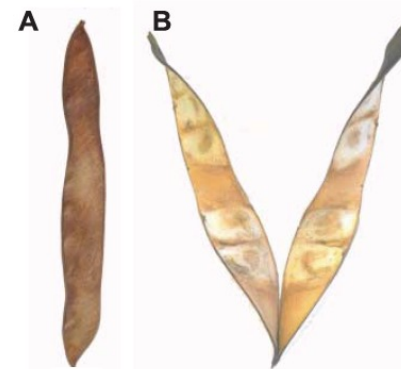
School of Chemistry, University of Sydney



Sarit Dutta, Benjamin Abecassis, CNRS, ENS Lyon



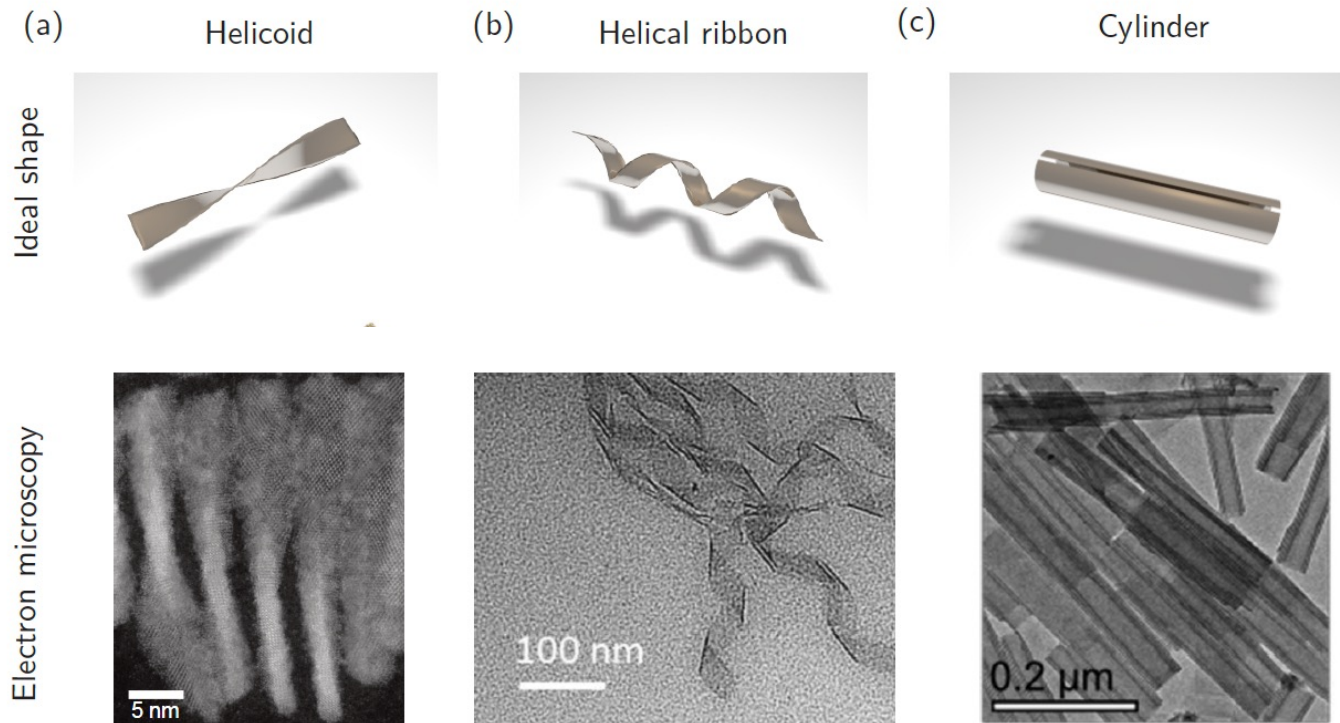
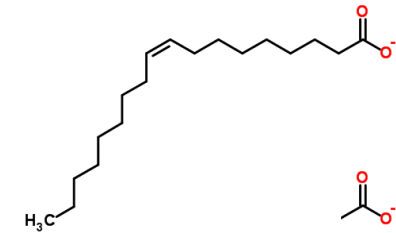
Science Advances 2017, 3, e1701483



Science 2011, 333, 1726

CdSe nanoplatelets

- 7-15 atomic layers thick (Cd on top and bottom)
- 10s-100s nm side length
- Zincblende crystal structure (dots and rods are wurtzite)



Simulation setup

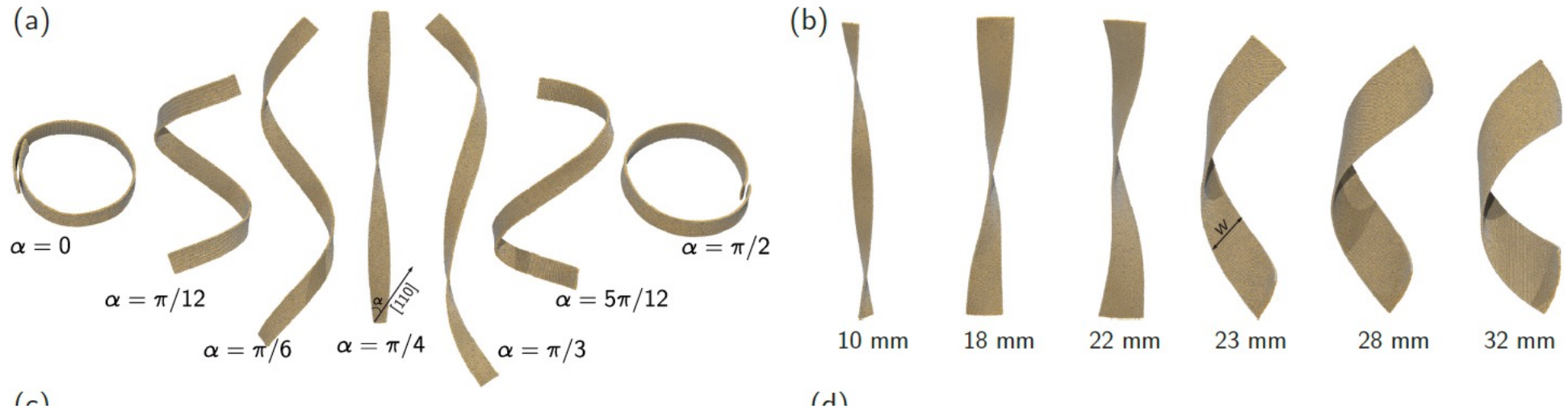
- Considered 3 different shapes (7 atomic layers thick, ~10 nm in side length)
- Used OPLS-AA force field for the ligands (also LJ/coulomb type)
- No explicit solvent



1. Add sufficient ligands to neutralise overall charge
2. Allow ligands to adsorb
3. Allow the full system to equilibrate

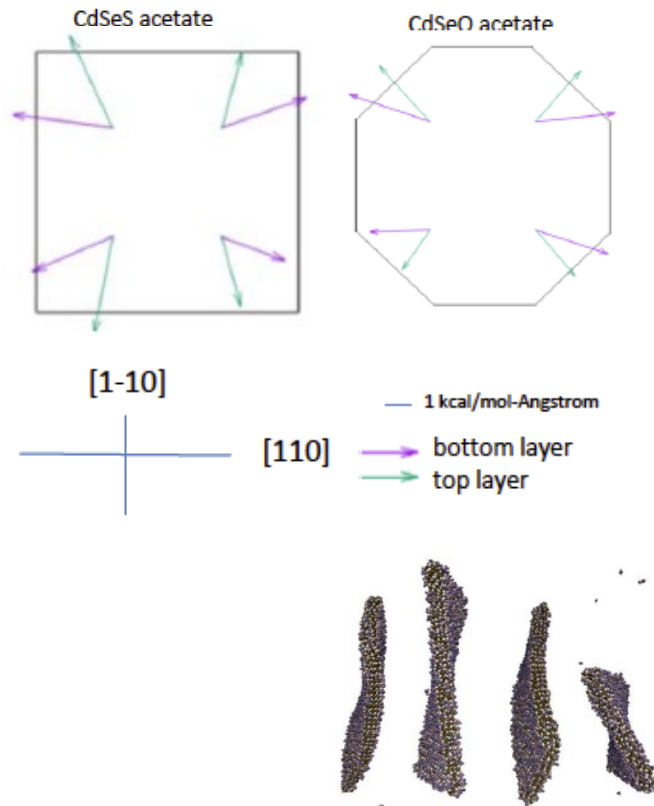


Results



Universal origin of twisting

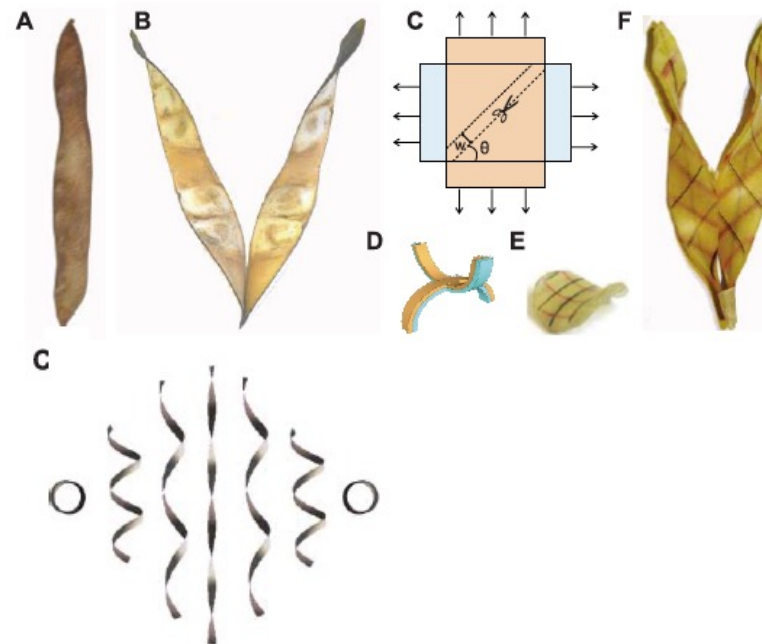
Stress anisotropy in NPLs



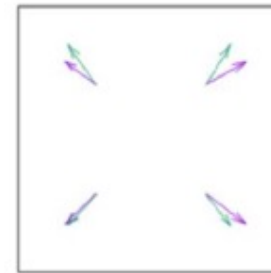
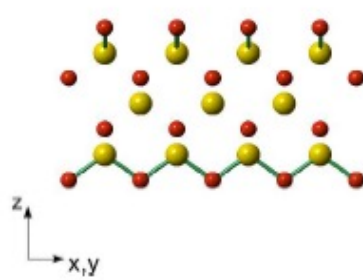
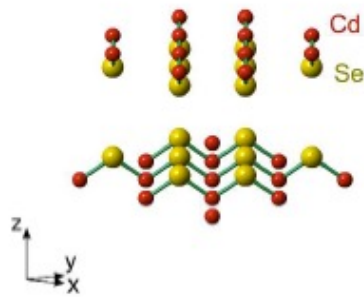
Geometry and Mechanics in the Opening of Chiral Seed Pods

Shahaf Armon,¹ Efi Efrati,¹ Raz Kupferman,² Eran Sharon^{1*}

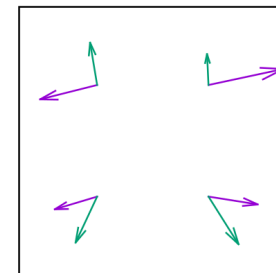
Science 2011, 333, 1726



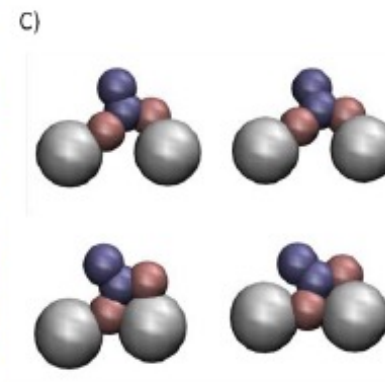
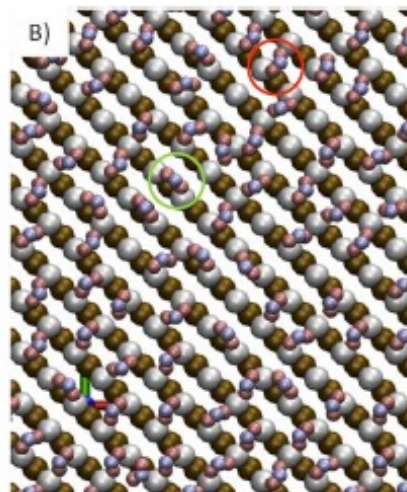
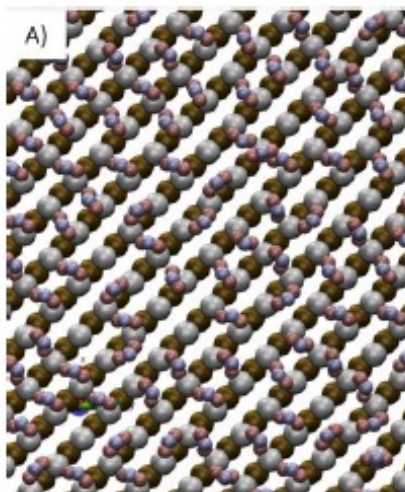
Chemical origin of twisting in CdSe NPLs



Bare crystal



With ligands



Ligands hold the rows together

Conclusions (2)

- Rolling and twisting of NPLs is due to the same physical mechanism that causes some seed pods to twist open, i.e. directional stress anisotropy in the top and bottom surfaces
- For CdSe NPLs, the anisotropy arises due to the interaction of a directionally anisotropic crystal surfaces (zincblende) and ligands preferentially binding in such a way that the top and bottom surfaces want to expand in different directions



Benjamin Abecassis



Marion Krapez



Debora Monego



Acknowledgements



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SYDNEY

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