

Where are the Icecube Neutrinos Beyond 2 PeV ?

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Neutrinos carry three types of information:

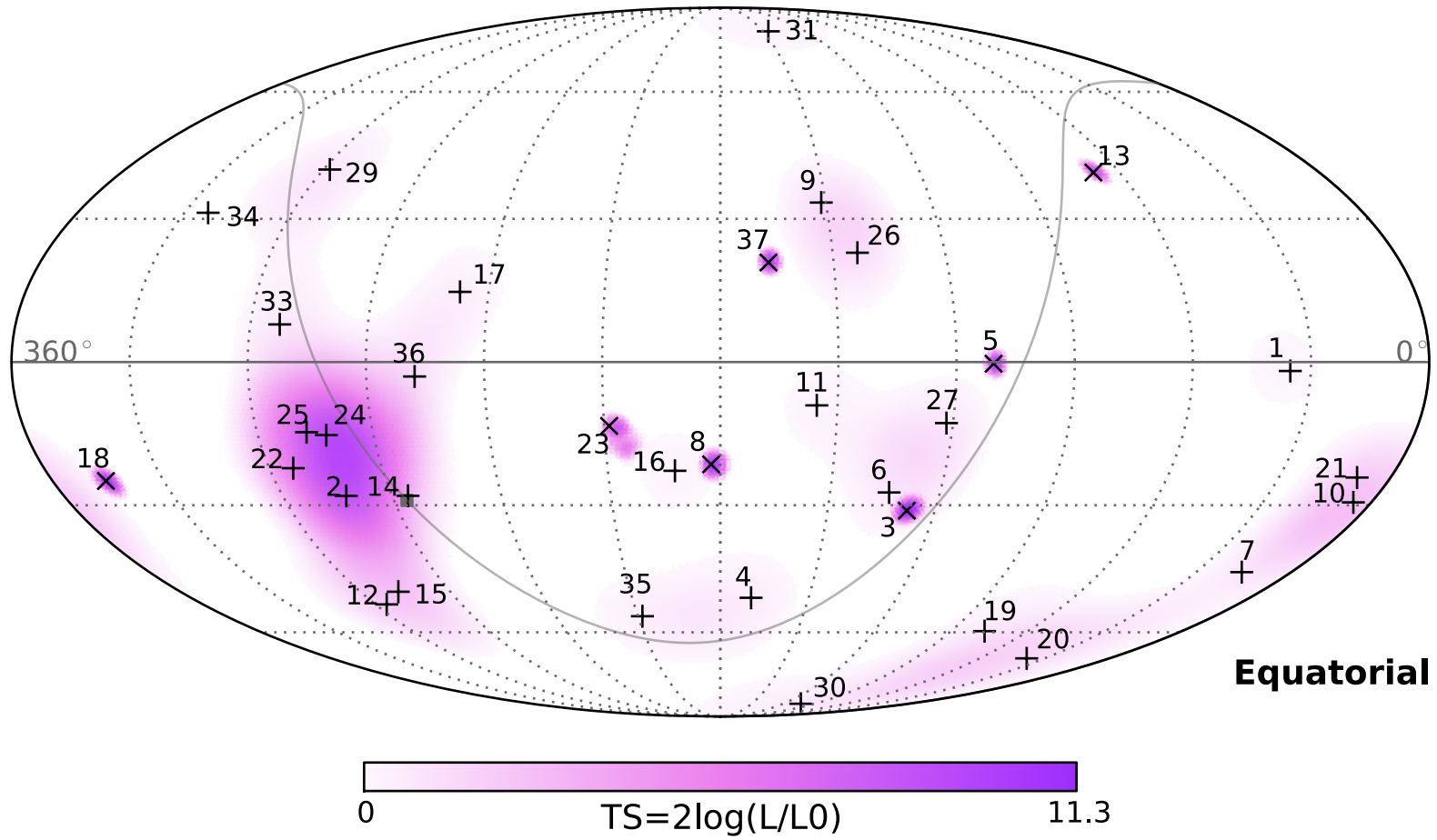
- (1) Direction
- (2) Energy
- (3) Flavor

All three have interesting features in IceCube data

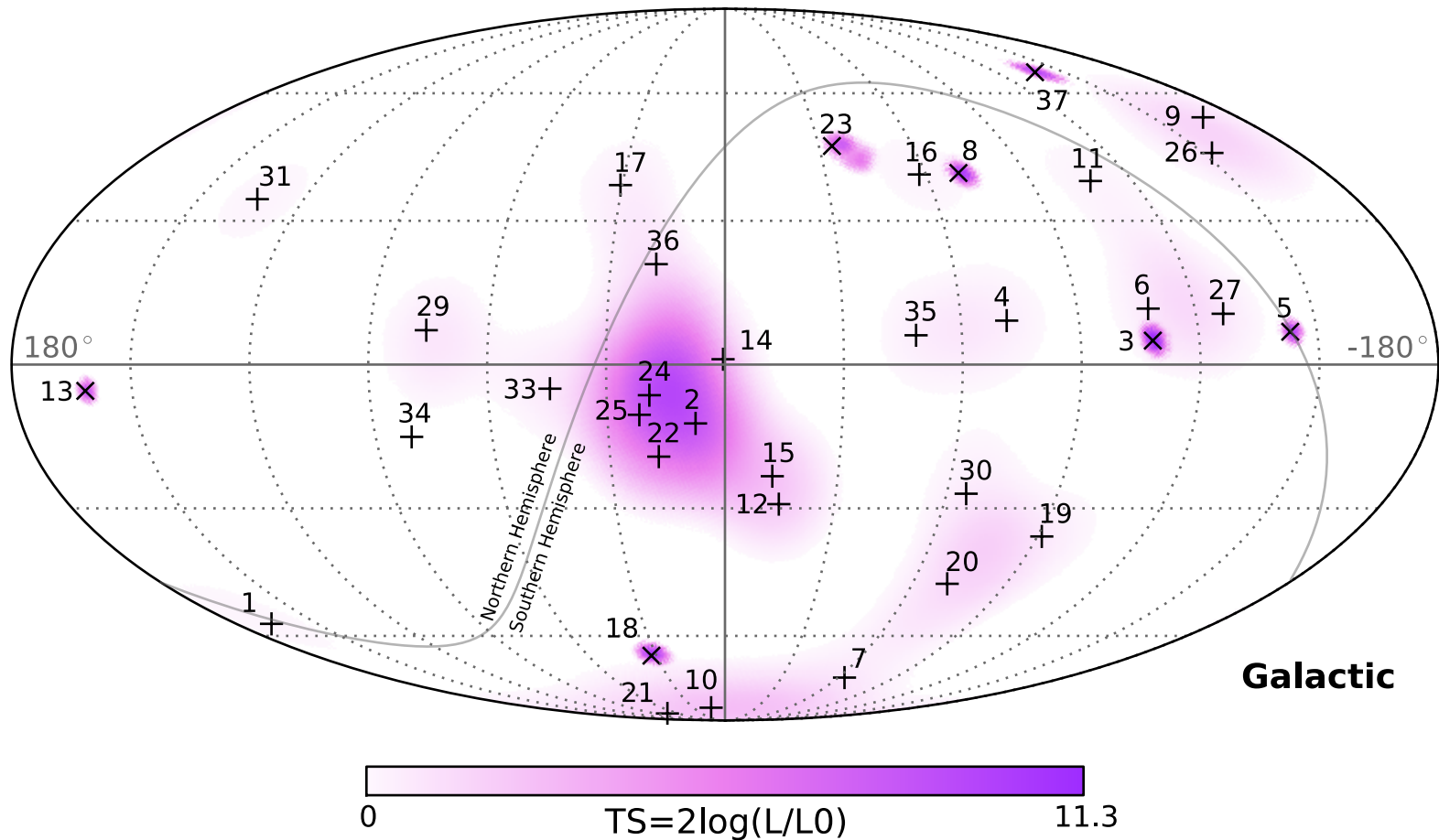
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- IceCube Data
- “missing” events above 2 PeV
 - continuum events
 - resonant at 6.3 PeV (“Glashow”)
- Statistics
- No Glashow is good Glashow?
 - End of the Neutrino spectrum
 - some consequences
 - some ideas

IceCube (Equatorial)



IceCube (Galactic)



Energy - the “gap”, and then three:

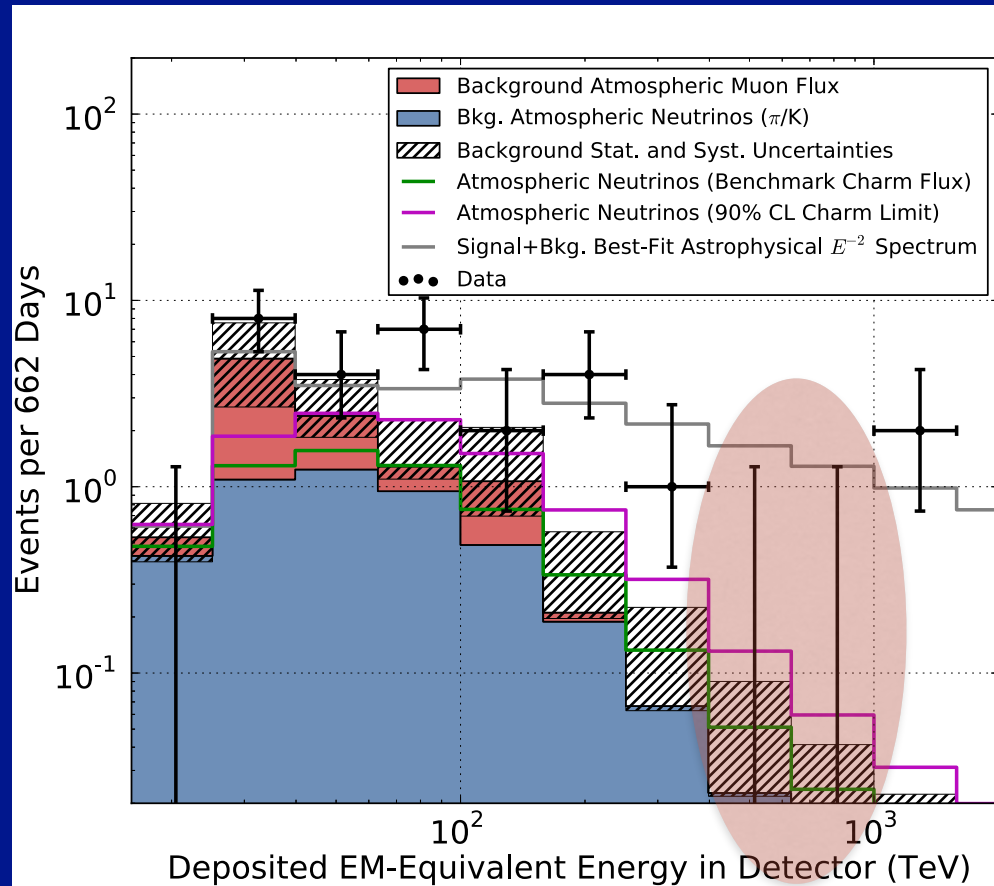


Figure 1: Distribution of the deposited energies of the observed events

Angle-Energy-Flavor Display:

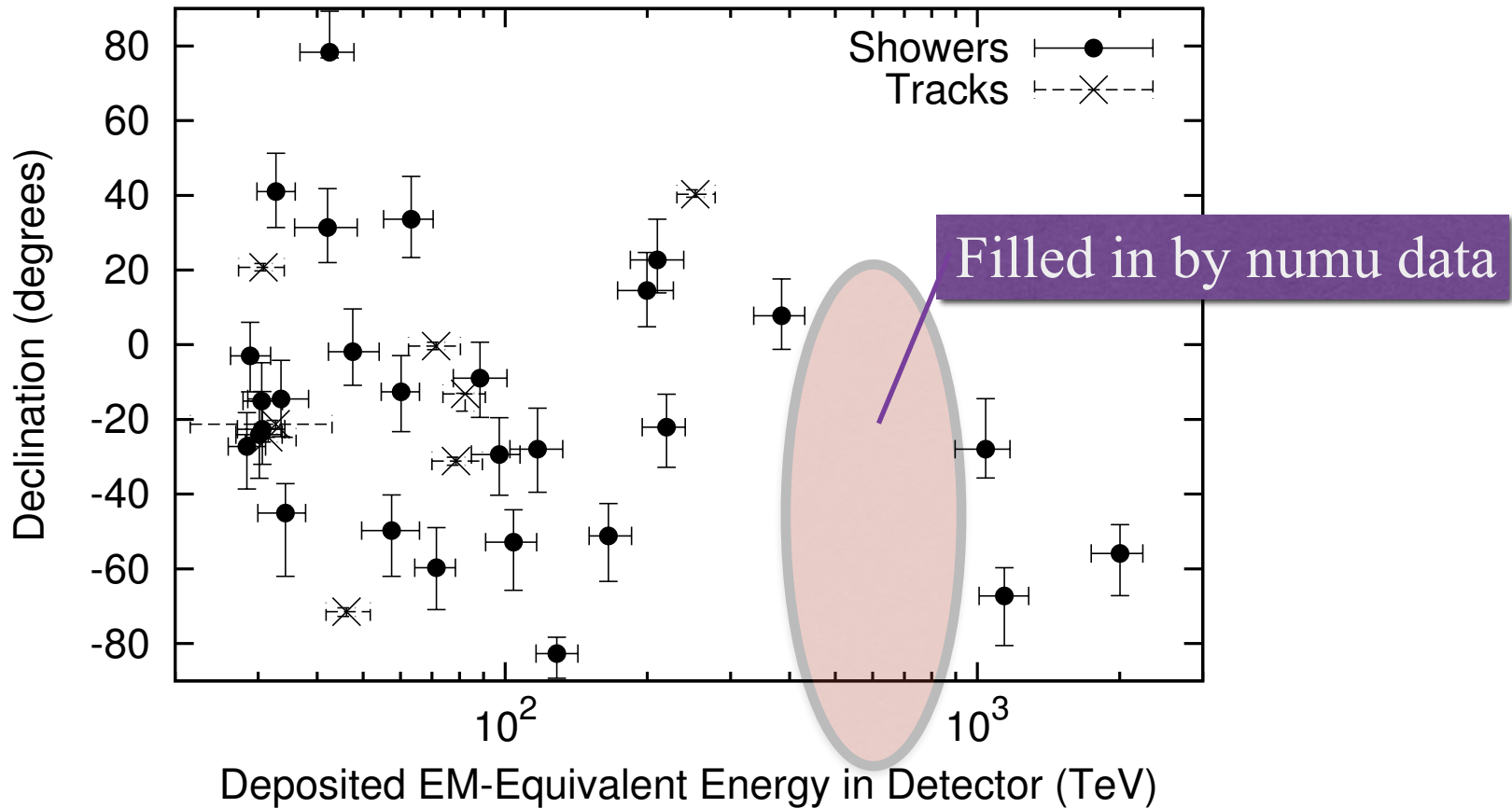
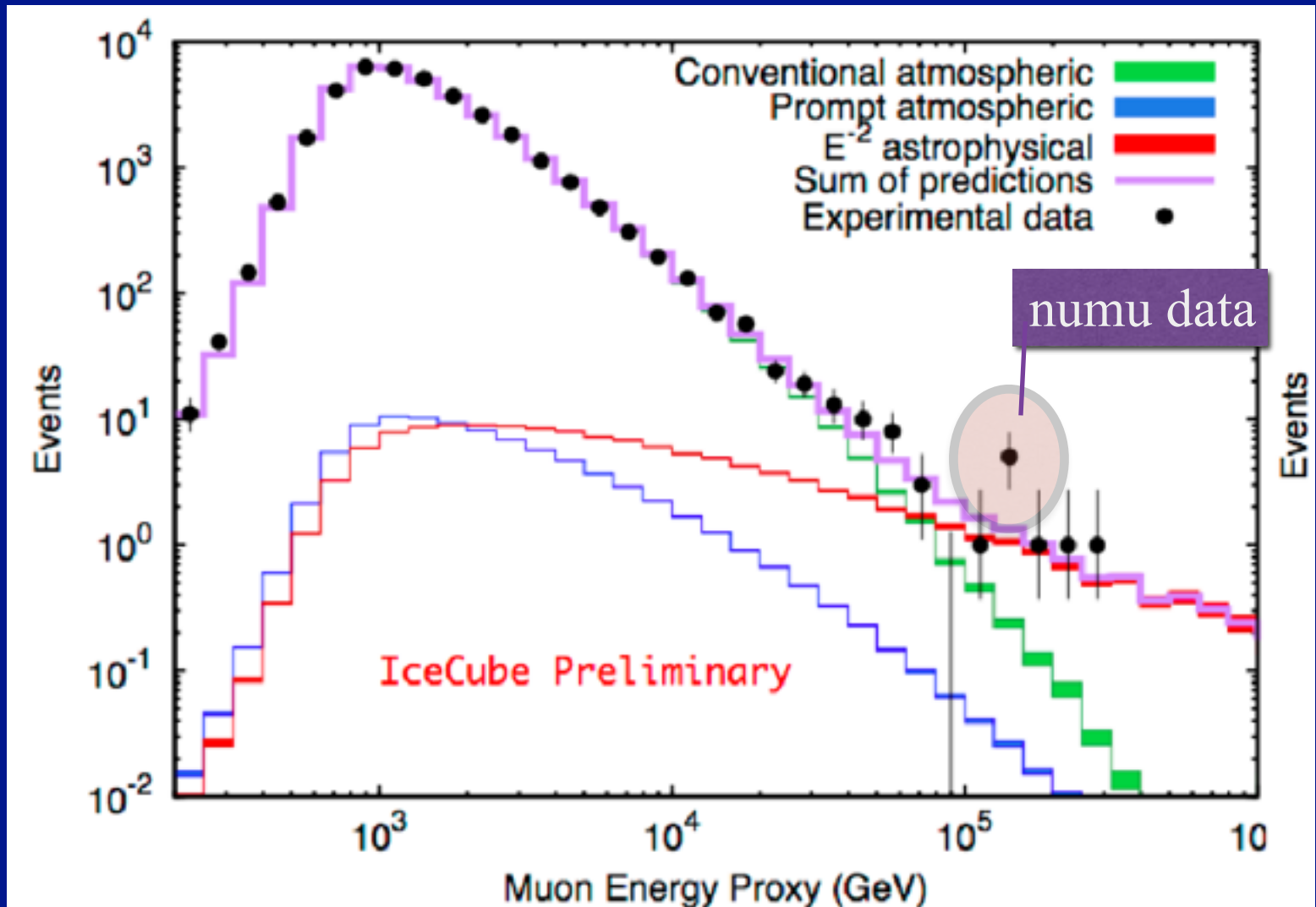


FIG. 1. Arrival angles and deposited energies of the events.

muon neutrino energy:



Is there an Energy Cutoff at $\sim 2\text{PeV}$?

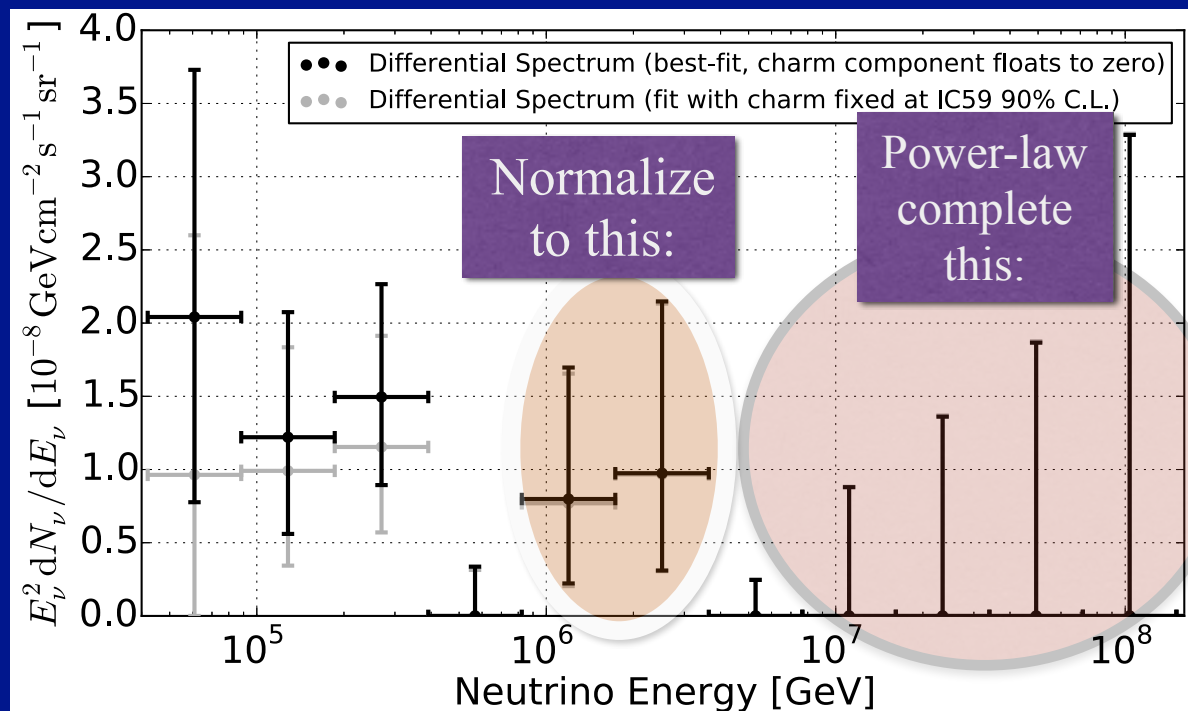


FIG. 4. Extraterrestrial neutrino flux ($\nu + \bar{\nu}$) as a function of energy. Vertical error bars indicate the $2\Delta\mathcal{L} = \pm 1$ contours of the flux in each energy bin, holding all other values, including background normalizations, fixed. These provide approximate 68% confidence ranges. An increase in the

Where are $E > 2$ PeV Neutrino Events?

“Missing” number depends on two parameters:

(1) spectral index, α

(2) number expected in 1-2 PeV interval,

$$N^{\text{exp}}(1-2 \text{ PeV})$$

μ

Integrated numbers:

$$N_{\text{cont}}^{\text{exp}}(E_1 < E < E_2) = \frac{N_{1-2\text{PeV}}^{\text{exp}}}{1 - 2^{-(\alpha-1.40)}} \left(E_1^{-(\alpha-1.40)} - E_2^{-(\alpha-1.40)} \right)$$

Taking $E_2 \rightarrow \infty$ as befits a continued power-law,

$$N_{\text{cont}}^{\text{exp}}(E > E_1) = \left(N_{1-2\text{PeV}}^{\text{exp}} \right) \frac{E_1^{-(\alpha-1.40)}}{1 - 2^{-(\alpha-1.40)}}.$$

Taking $\alpha = 2.0, 2.3,$ and $2.5,$ we get, respectively,

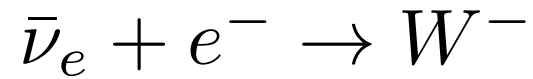
$$N_{\text{cont}}^{\text{exp}}(E > 2\text{PeV}) = \left(N_{1-2\text{PeV}}^{\text{exp}} \right) \times 1.94, 1.15, 0.87.$$

In addition, the number of Glashow resonance events is

$$\frac{N_{\text{Res}}^{\text{exp}}}{N_{\text{cont}}^{\text{exp}}(E > 2\text{PeV})} \approx 3.4 \times \mathcal{R},$$

where \mathcal{R} is the fraction of Earthly flux that is $\bar{\nu}_e$, ranging from 56% (β -beam) to zero (damped p - γ - no antineutrinos).

Where/what is Glashow?



$$s = M_W^2 = 2m_e E_\nu, \quad \text{so } E_R = \frac{M_W^2}{2m_e} = 6.3 \text{ PeV}$$

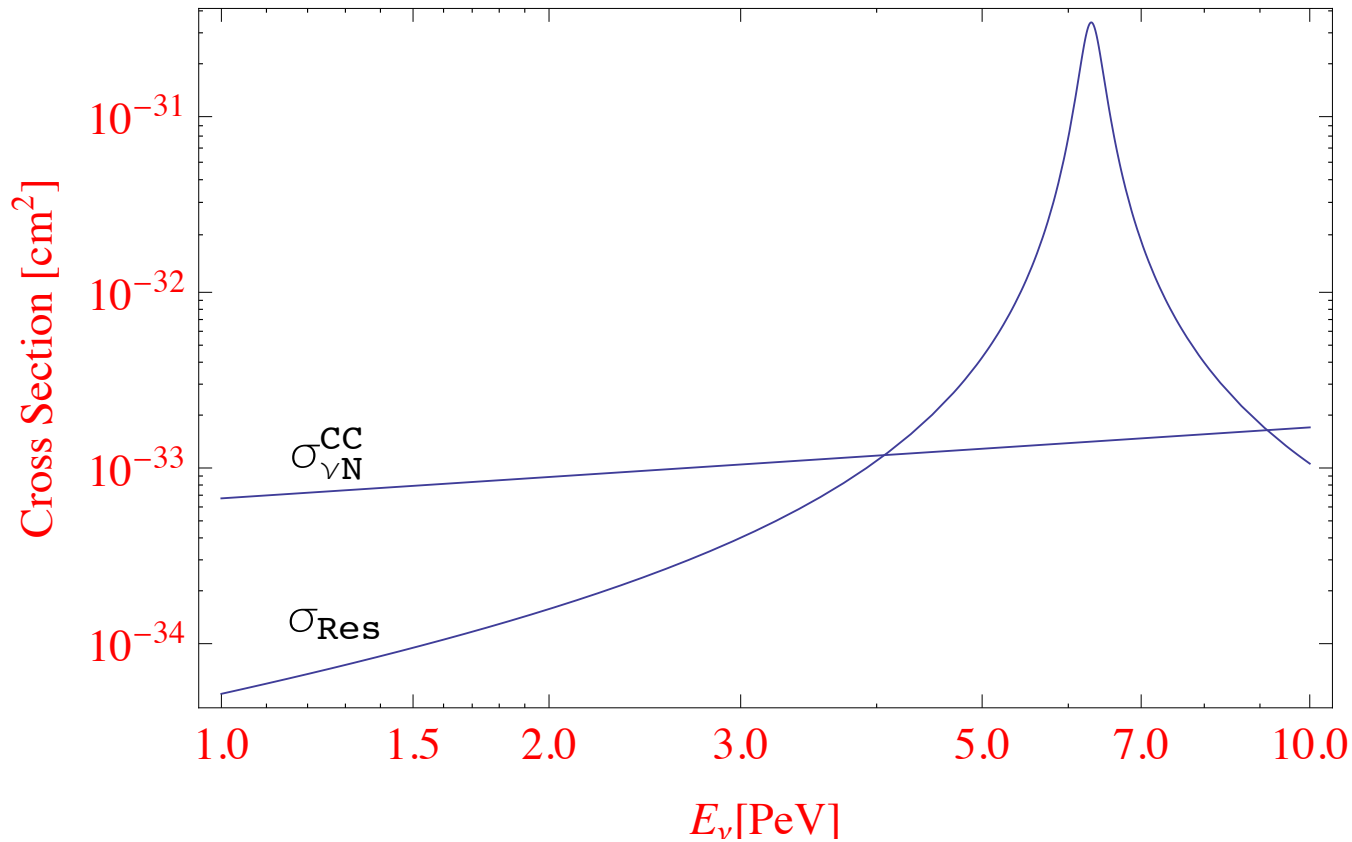


FIG. 1: Cross sections for the resonant process, $\bar{\nu}_e + e^- \rightarrow W^- \rightarrow \text{hadrons}$, and the non-resonant process, $\nu_e + N \rightarrow e^- + \text{hadrons}$, in the 1–10 PeV region.

Glashow Resonance - Formulas:

$$\left(\frac{N}{T\Omega} \right)_{\text{Res}} = \frac{N_p}{2m_e} (\pi M_W \Gamma_W) \sigma_{\text{Res}}^{\text{peak}} \left. \frac{dF_{\bar{\nu}_e}}{dE_{\bar{\nu}_e}} \right|_{E_{\bar{\nu}_e}=6.3\text{PeV}},$$

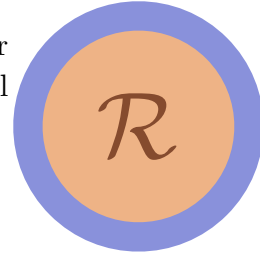
$$\sigma_{\text{Res}}^{\text{peak}} = \frac{24\pi \text{B}(W^- \rightarrow \bar{\nu}_e e^-) \text{B}(W^- \rightarrow \text{had})}{M_W^2} = 3.4 \times 10^{-31} \text{cm}^2.$$

Explored in: Anchordoqui, Goldberg, Halzen, TJW (hep-ph/0410003), ...
Bhattacharya, Gandhi, Rodejohann, Watanabe (1108.3163, 1209.2422),
Baerwald, Bustamente, Winter (1208.4600)

The “Resonometer” of Cosmic Nu Source Models:

Barger, Lingjun Fu, Learned, Marfatia, Pakvasa, TJW,
Phys Rev (Letters?) and arXiv 1407.3255

TABLE I: Neutrino flavor ratios at source, component of $\bar{\nu}_e$ in total neutrino flux at Earth after mixing and decohering, and consequent relative strength of Glashow resonance, for six astrophysical models. (Neutrinos and antineutrinos are shown separately, when they differ.)



| | Source flavor ratio | | Earthly flavor ratio | | $\bar{\nu}_e$ fraction in flux (\mathcal{R}) |
|----------------------------------|---------------------|---------|----------------------|---------|--|
| $pp \rightarrow \pi^\pm$ pairs | (1:2:0) | | (1:1:1) | | 18/108 = 0.17 |
| w/ damped μ^\pm | (0:1:0) | | (4:7:7) | | 12/108 = 0.11 |
| $p\gamma \rightarrow \pi^+$ only | (1:1:0) | (0:1:0) | (14:11:11) | (4:7:7) | 8/108 = 0.074 |
| w/ damped μ^+ | (0:1:0) | (0:0:0) | (4:7:7) | (0:0:0) | 0 |
| charm decay | (1:1:0) | | (14:11:11) | | 21/108 = 0.19 |
| neutron decay | (0:0:0) | (1:0:0) | (0:0:0) | (5:2:2) | 60/108 = 0.56 |

(Kaons change
very little)

Poisson:

POISSON STATISTICS:

$$1 = e^{-\mu} e^{+\mu} = e^{-\mu} \sum_{n=0}^{\infty} \frac{\mu^n}{n!} .$$

from which we infer that

$$P(n|\mu) = e^{-\mu} \frac{\mu^n}{n!} .$$

is the probability of observing n events when the mean, or expected number, is μ .

Poisson:

Mean value of $P(n|\mu)$ is $\sum_n n \times P(n|\mu) = \mu$.

Moreover, variance $\sum_n (n^2 - \mu^2) \times P(n|\mu)$ is also μ .
(Source of $\frac{S}{N} \sim \frac{\mu}{\sqrt{\mu}} = \sqrt{\mu}$)

Moreover, $\int_0^\infty d\mu P(n|\mu) = \int_0^\infty d\mu \frac{e^{-\mu} \mu^n}{n!} = \frac{n!}{n!} = 1$.

FELDMAN-COUSINS

inverts the Poisson distribution:

instead of expected number μ as given, and observed number n as resulting parameter, it's observed number n as given, and range of expected number μ as the inferred parameter.

For a CL, have

$$CL = \int_{\mu_1}^{\mu_2} P(n|\mu),$$

with Feldman-Cousins providing a prescription for where to place the centroid of $[\mu_1, \mu_2]$. At 95% CL (roughly 2σ), F-C get $\mu \in [0.8, 8.0]$ when $n=3$ (and zero background).

But we have more (Bayesian) info:

there are no observed events above 2 PeV !

So take $\mu = 1, 2, 3$ as representative.

For the first two, three events viewed as an upward fluctuation.

Poisson:

SPECIAL CASES:

(1) using Stirling's approximation,

$$P(\mu|\mu) = \frac{1}{\sqrt{2\pi\mu}}, \text{ so}$$

$P(3|3) \approx \frac{1}{\sqrt{6\pi}} = 23\%$, while $P(3|2) = 18\%$, and $P(3|1) = 6\%$,
all roughly "equally" viable;

(2) when no events are observed,

$$P(0|\mu) = e^{-\mu},$$

so $P(0|1) = 1/e = 37\%$ and $P(0|3) = e^{-3} = 5.0\%$,
for the dearth of events above 2 PeV, also viable;
however, $P(0|5) = 0.7\%$, not so viable.

“missing” neutrinos at IceCube

Table : Neutrino events above 2 PeV, continuum and resonant.

| spectral index | $N_{\text{cont}}(E > 2 \text{ PeV})$ | $+N_{\text{res}} = 3.4\mathcal{R}N_{\text{cont}}(E > 2 \text{ PeV})$ |
|----------------|---|--|
| 2.0 | $1.94 N_{1-2 \text{ PeV}}^{\text{exp}}$ | $1.94 N_{1-2 \text{ PeV}}^{\text{exp}} (1 + 3.4\mathcal{R})$ |
| 2.3 | $1.15 N_{1-2 \text{ PeV}}^{\text{exp}}$ | $1.15 N_{1-2 \text{ PeV}}^{\text{exp}} (1 + 3.4\mathcal{R})$ |
| 2.5 | $0.87 N_{1-2 \text{ PeV}}^{\text{exp}}$ | $0.87 N_{1-2 \text{ PeV}}^{\text{exp}} (1 + 3.4\mathcal{R})$ |

“missing” neutrinos at IceCube

For example,

Table : Neutrino events above 2 PeV, continuum and resonant.

| $\alpha = 2.3, \mathcal{R} = 0.17$ | $N_{\text{cont}}(E > 2 \text{ PeV})$ | $+N_{\text{res}}$ |
|--|---|---|
| | $1.15 N_{1-2 \text{ PeV}}^{\text{exp}}$ | $1.81 N_{1-2 \text{ PeV}}^{\text{exp}}$ |
| $N_{1-2 \text{ PeV}}^{\text{exp}} = 1$ | 1.15 | 1.81 |
| $N_{1-2 \text{ PeV}}^{\text{exp}} = 2$ | 2.30 | 3.62 |
| $N_{1-2 \text{ PeV}}^{\text{exp}} = 3$ | 3.45 | 5.44 |

$$P(0|1.81) = 16 \%$$

$$P(0|3.62) = 2.7\%$$

$$P(0|5.44) = 0.43\%$$

No Glashow is good Glashow ?

“End of the Neutrino Energy Spectrum”

Anchordoqui, Barger, Goldberg, Learned, Marfatia, Pakvasa, Paul, TJW,
Phys Lett B and 1404.0622

(violates Learned's Theorem)

E_{maxed} leptons stabilize charged pion at hi E,
also stabilize neutron at hi E

(paper pending, with P. Denton, D. Marfatia)

Mass-Scales and Energy Cutoff in terms of Boost Factor

$$\Gamma_\nu = \left(\frac{E_\nu}{2 \text{ PeV}} \right) \left(\frac{0.05 \text{ eV}}{m_\nu} \right) \times 0.4 \cdot 10^{17}$$

whereas

$$\frac{M_{\text{Planck}}}{v_{\text{weak}}} = \frac{1.2 \times 10^{28} \text{ eV}}{247 \text{ GeV}} = 0.5 \times 10^{17} ;$$

Maybe suggests

$$\Gamma_\nu^{\text{max}} = M_{\text{Planck}} / M_{\text{weak}}$$

and a possible connection
to Gravity/spacetime foam.

Neutrino Energy Maximum:

$$\begin{aligned}
 E_{\nu}^{\max} &= \frac{m_{\nu} M_{\text{Planck}}}{M_{\text{weak}}} \\
 &= 2.5 \left(\frac{m_{\nu}}{0.05 \text{ eV}} \right) \left(\frac{M_{\text{Planck}}}{1.2 \times 10^{28} \text{ eV}} \right) \left(\frac{247 \text{ GeV}}{v_{\text{weak}}} \right) \text{ PeV} .
 \end{aligned}$$

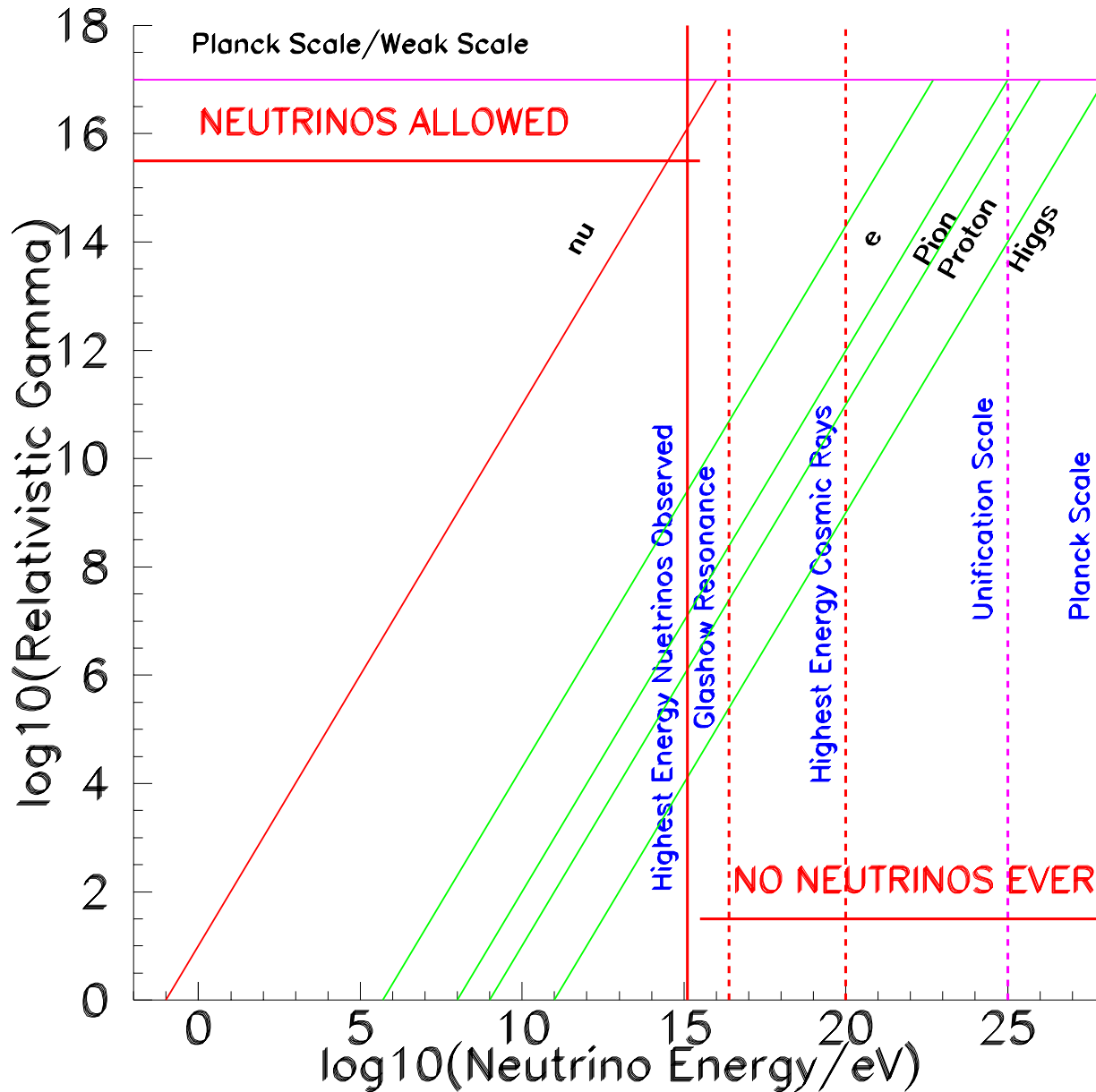
In what frame?

Nature provides THE preferred frame, the Cosmic Rest Frame.

So E_{ν}^{\max} can be written as $u_{\beta}^{\text{CRF}} (p_{\nu}^{\max})^{\beta}$, where $u_{\beta}^{\text{CRF}} = (1, \vec{0})$.

And $(p_{\nu}^{\max})^{\beta}$ transforms as usual four-vector.

The End of the Neutrino Spectrum



and Weinberg's mass operator:

Inserting Weinberg's neutrino-mass generating operator,

$$\frac{1}{\Lambda}(HL)(HL) \Rightarrow m_\nu = \frac{vev^2}{\Lambda},$$

find $\Lambda \sim M_{\text{GUT}}$, so

$$\sqrt{G} E_\nu^{\text{max}} \sim \frac{vev}{M_{\text{GUT}}}.$$

Reasons (excuses):

- 1- LI is “emergent” low-energy symmetry;
- 2- Weak int’n “size” is Higgs vev fluctuation,
contracted by Lorentz to Planck size:
 - a) spacetime foam
 - b) strong gravity/geometry
 - c) stringy relations, G and SM
 - d) BH entropy: area and G
 - c) extra dim’ns open up: geometry and M_P
 - d) Planck scale LIV
 - e) ...make your own model
(like Bj; like Illana, Masip, Meloni)

Outguessing Nature doesn’t matter, as it’s an Xpnl issue!

and yes,

Incredible claims require incredible evidence ...

data coming in, but not yet there

In Summary:

Multi-PeV continuum and Glashow resonance rates in borderline danger, worth watching.

Glashow resonance can reveal ν source dynamics on other side of Universe.

If events above a few PeV do not arrive, then $\%\$#@~* \&$, and either

(a) Nature cuts off the sources ($E_p \sim 20 E_\nu$ for pion chain);

e.g., Reno, Sarcevic, ... non-pion chain,
choked jets/charm prodct'n and decay

(b) the power-law is broken;

(c) new fundamental physics at scale $\Gamma_\nu \sim 10^{17}$.