

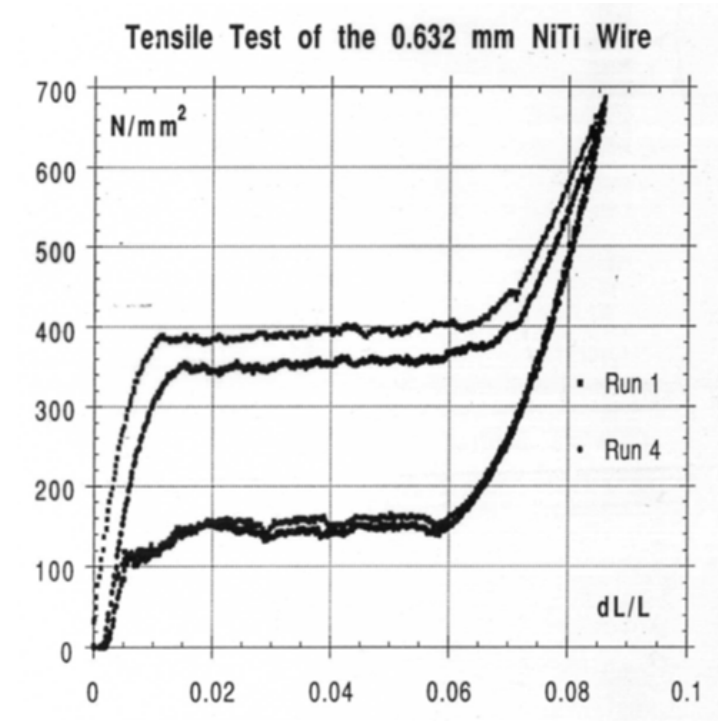
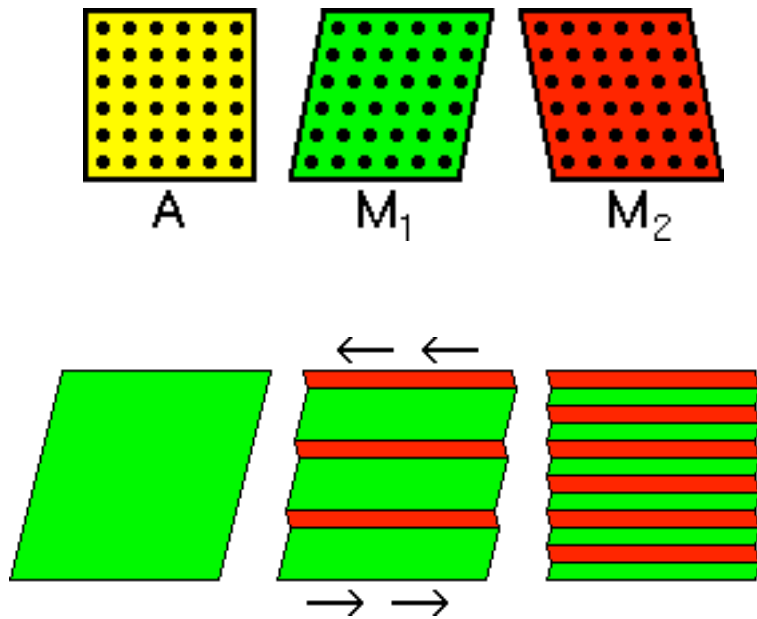
Criticality in martensites: plasticity or inertia?

Lev Truskinovsky

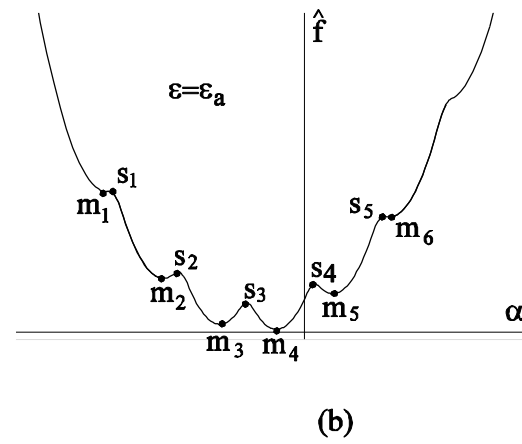
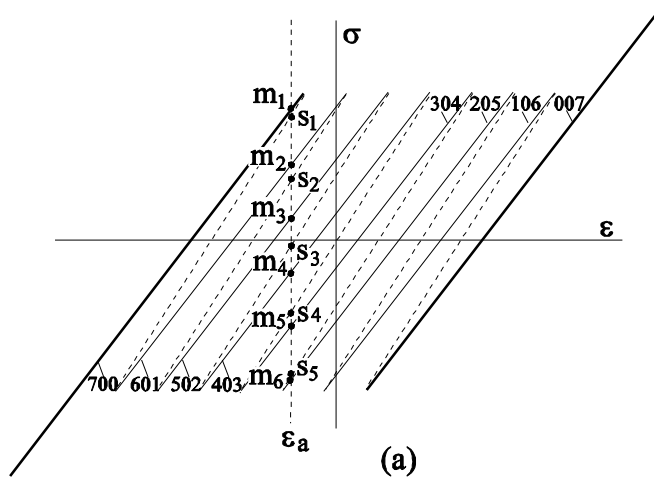
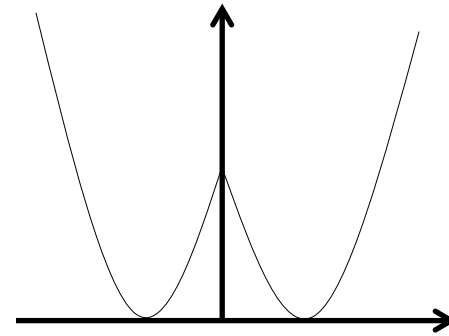
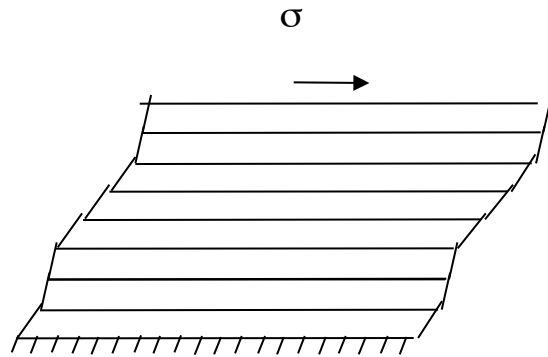
CNRS, Ecole Polytechnique

Joint work with F. Perez-Reche, G. Zanzotto, A. Finel and U. Salman

Martensitic phase transitions



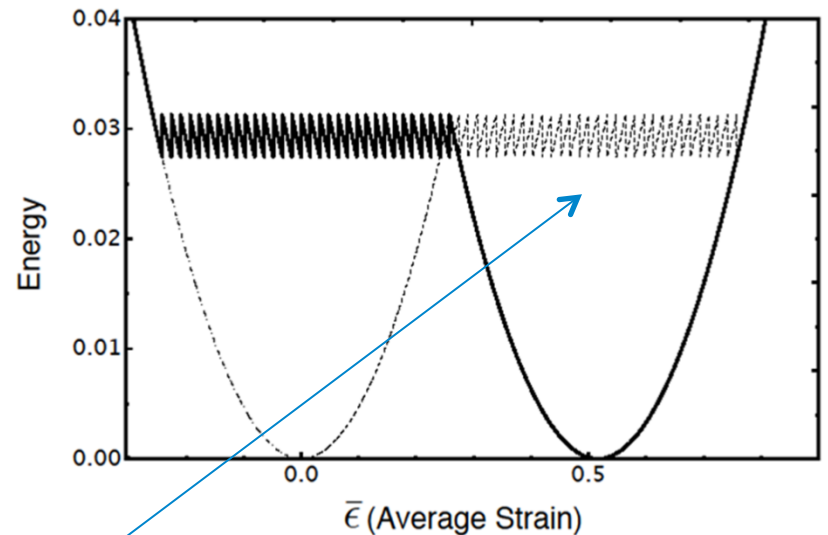
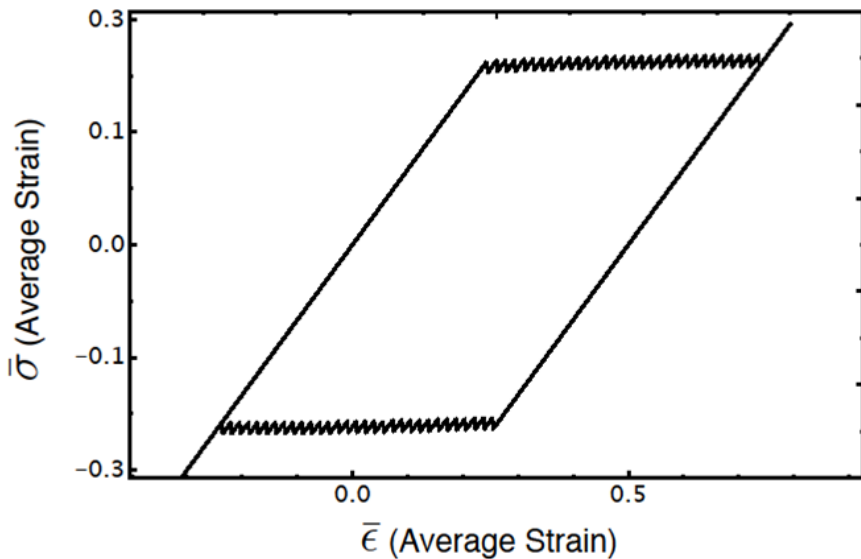
Double well energy



Over-damped dynamics

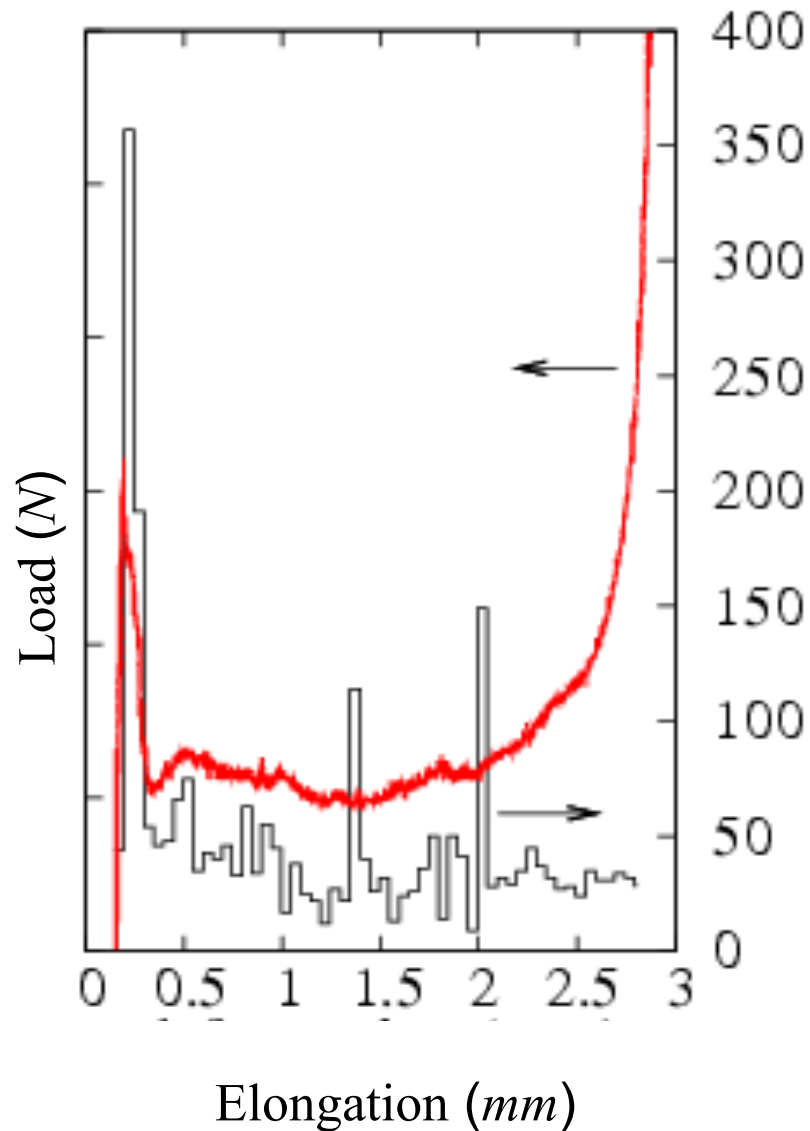
$$v \dot{\varepsilon}_i = -\frac{\partial \Phi}{\partial \varepsilon_i} + \sigma$$

$$\Phi = N^{-1} \sum_{i=1}^N (f_0(\varepsilon_i) - \sigma \varepsilon_i)$$



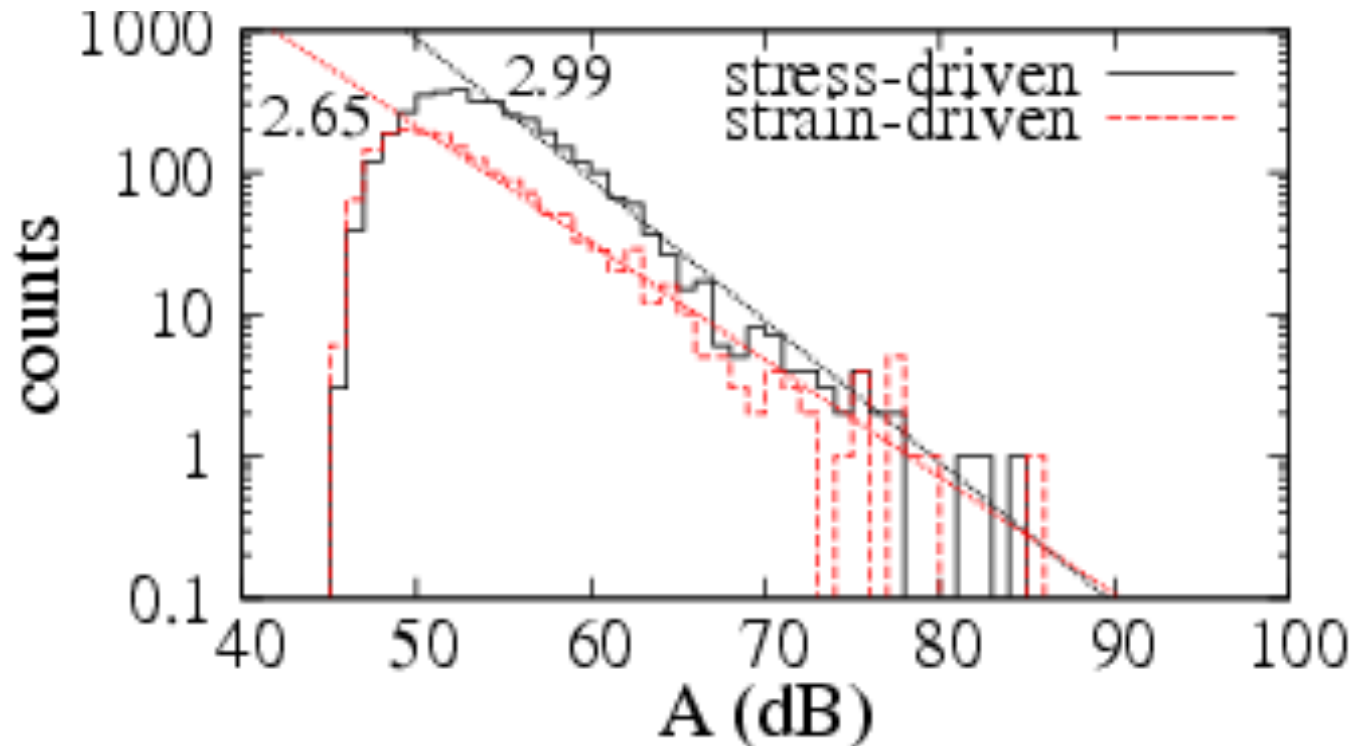
Marginally stable states

Martensitic phase transitions

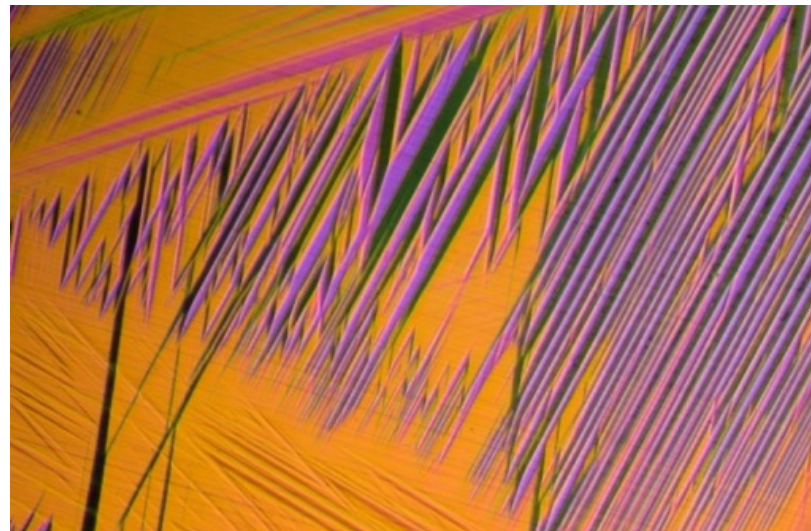
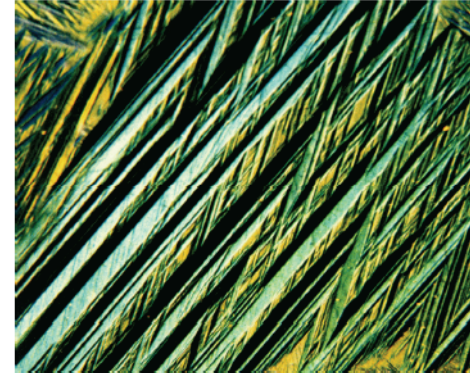
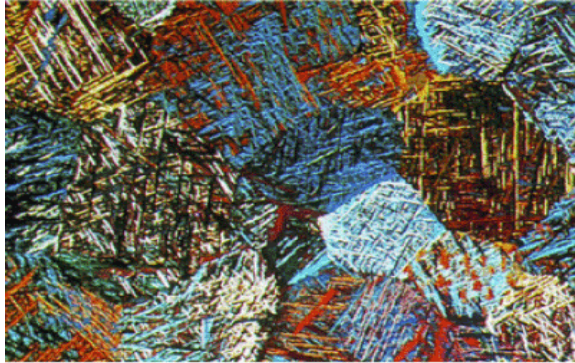


E.Vives et al,
PRB, 2009

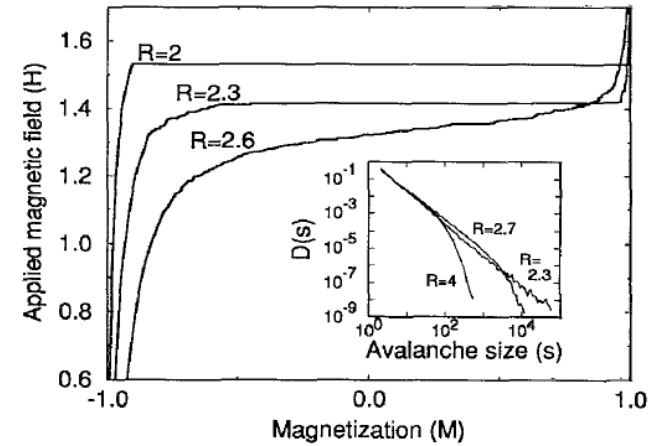
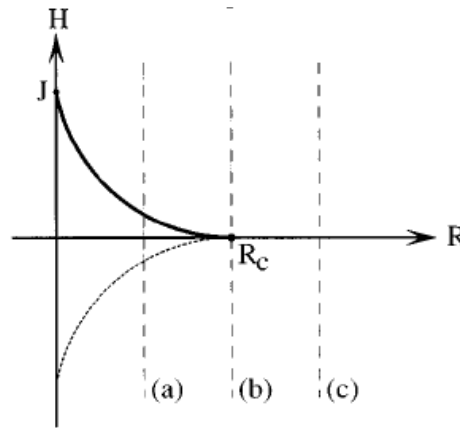
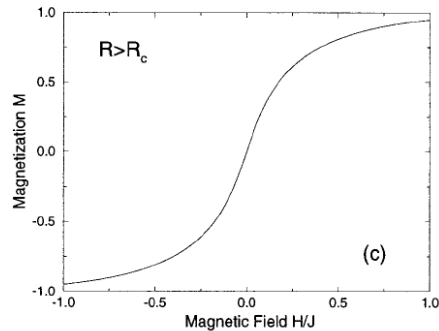
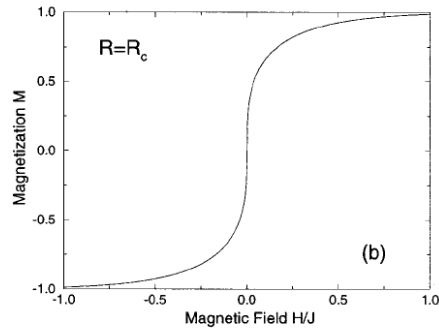
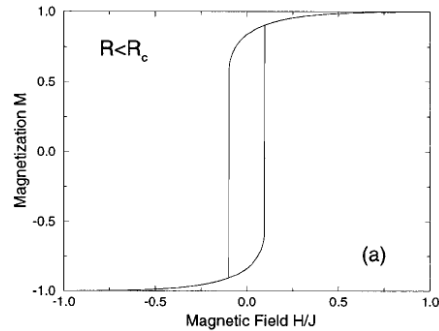
Statistics of avalanches



Spatial complexity



Random Field Ising model



$$\mathcal{H} = - \sum_{ij} J_{ij} s_i s_j - \sum_i (H + f_i) s_i$$

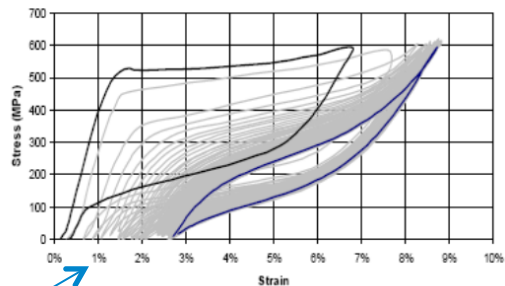
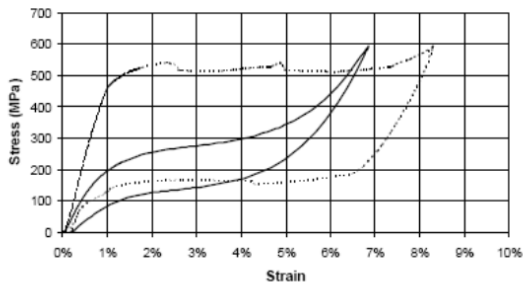
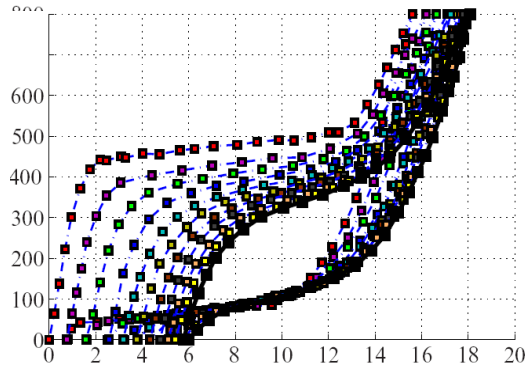
$$\rho(f_i) = \frac{1}{\sqrt{2\pi}R} \exp\left(-\frac{f_i^2}{2R^2}\right)$$

How the system finds the critical point?

- Critical domain is anomalously large
- Limited plasticity
- Inertia

Plasticity?

Training



residual strain

Local energy

$$\boxed{\phi = \phi(\mathbf{C}; T)} \quad C_{ab} = \mathbf{u}_a \cdot \mathbf{u}_b$$

Invariant with respect to the change between
lattice bases

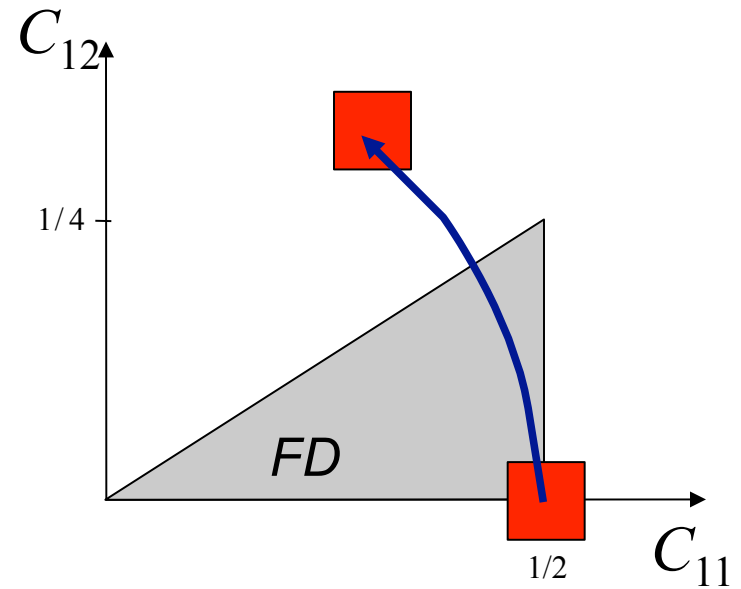
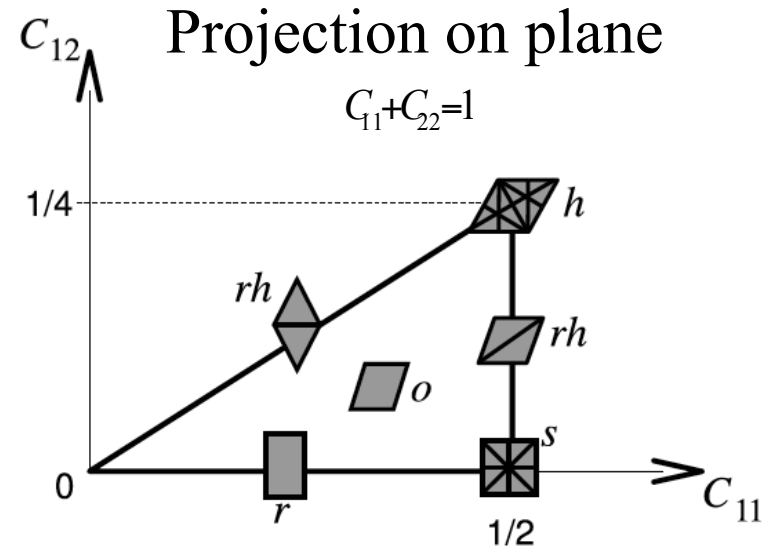
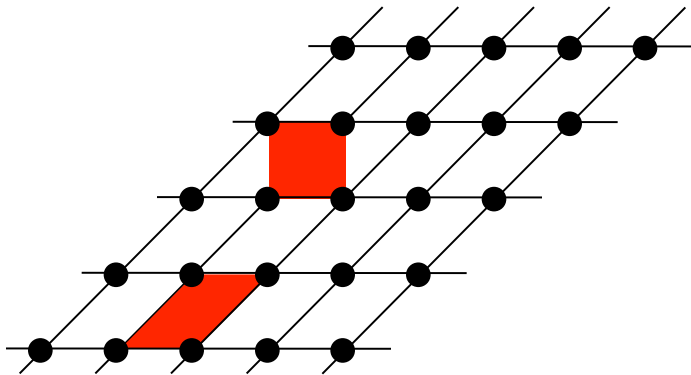
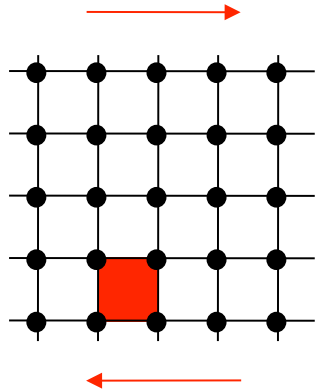


$GL(2, \mathbb{Z}) :$

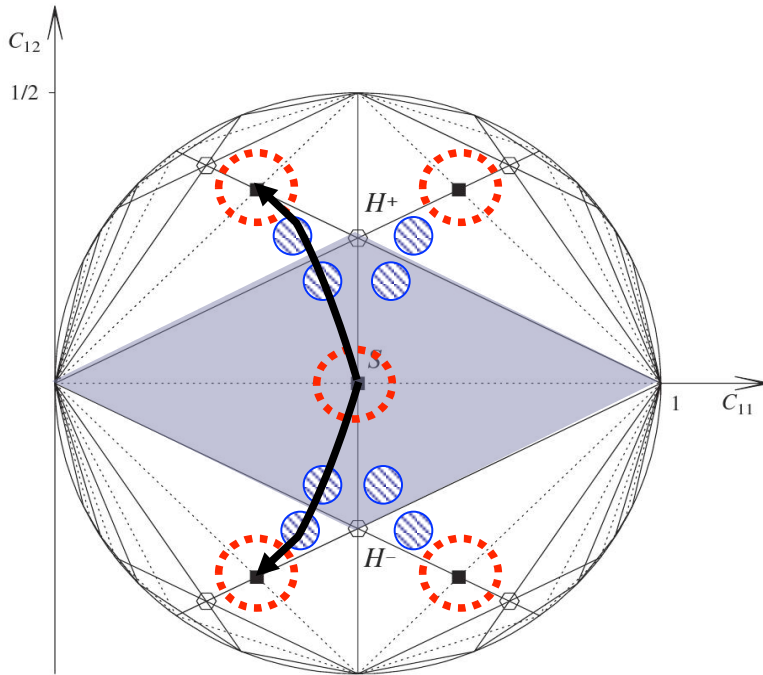
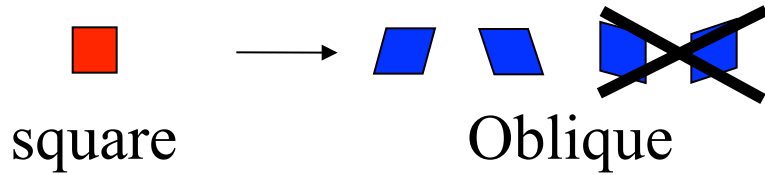
2×2 invertible matrices with integral entries

$$\phi(\mathbf{C}; T) = \phi(\mathbf{M}^t \mathbf{C} \mathbf{M}; T)$$

Plastic deformation



Energy structure



- Deformation along the path

$$\{e(\mathbf{C}), e'(\mathbf{C}), e''(\mathbf{C})\}$$

OP: Order parameter

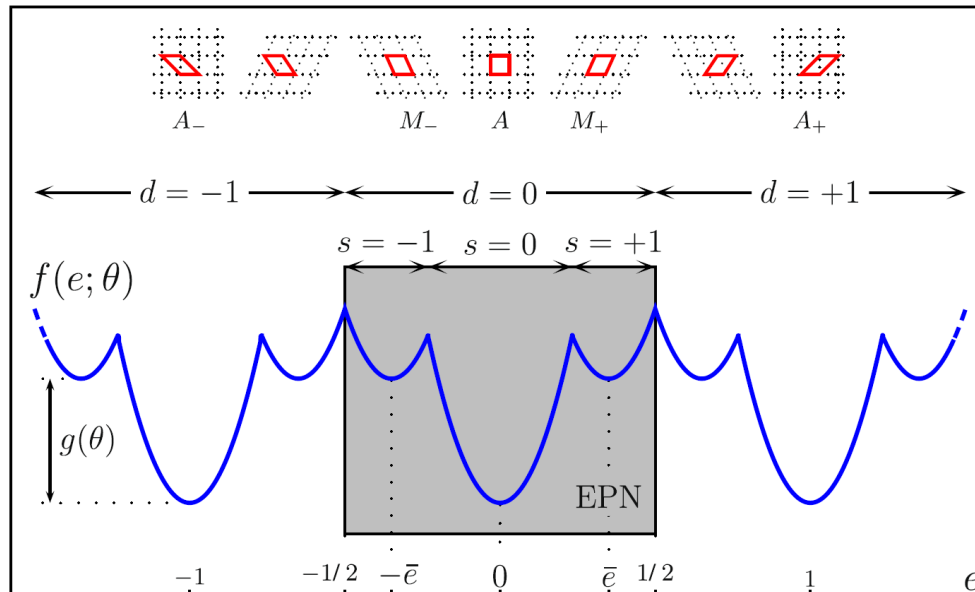
- Elimination of e' and e''
- Energy in terms of OP

$$\tilde{\Phi}(e; \theta) = \frac{1}{2} \sum_{i,j=1}^N K_{ij} e_i e_j + \sum_{i=1}^N f(e_i; \theta)$$

elastic
interaction

local energy
elastic units

Energy accounting of plastic deformation



$$\tilde{\Phi}(e; \theta) = \sum_{i=1}^N f(e_i; \theta) + \frac{1}{2} \sum_{i,j=1}^N K_{ij} e_i e_j$$

$$f(e; \theta) = \frac{1}{2}(e - w)^2 + g(\theta)s^2 \quad w_i (= d_i + \bar{e}s_i)$$

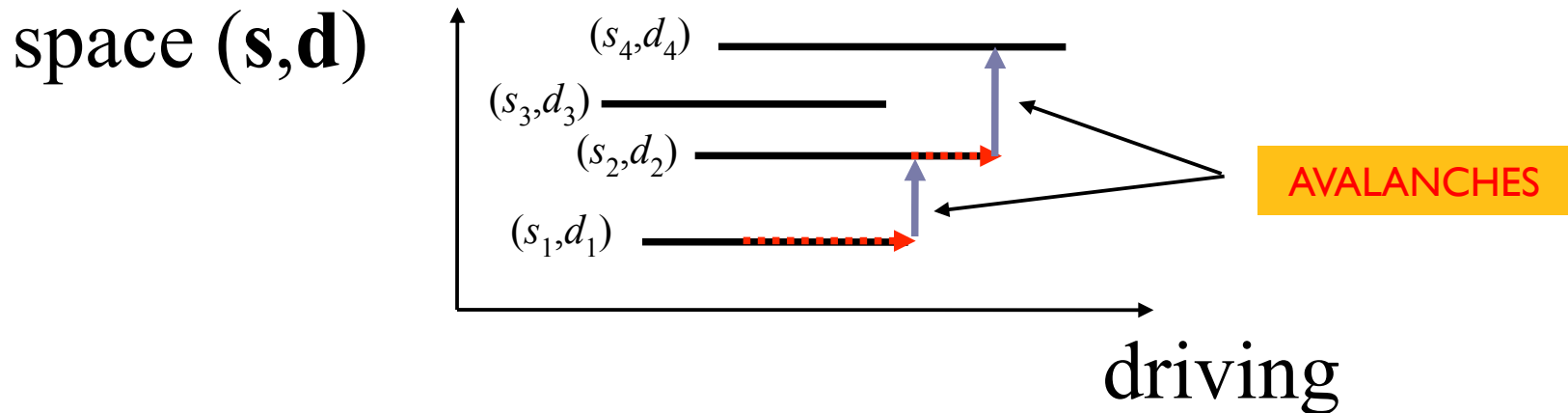
Dynamics of the order parameter

$$\frac{\partial \mathbf{e}}{\partial t} = -\gamma \frac{\partial \tilde{\Phi}(\mathbf{e}; \theta)}{\partial \mathbf{e}}$$

Slow driving: $\gamma \rightarrow \infty$

Dynamics projects as an automaton on the local minima of $\tilde{\Phi}$

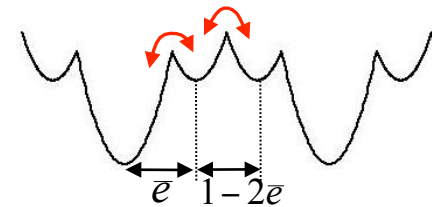
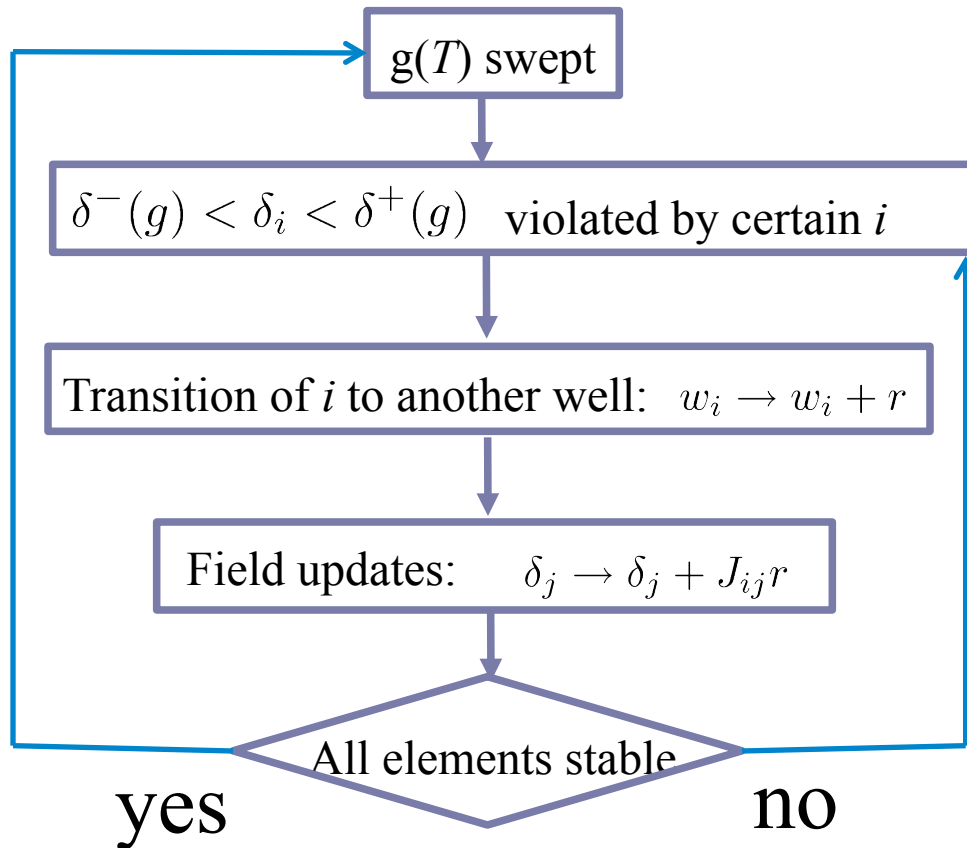
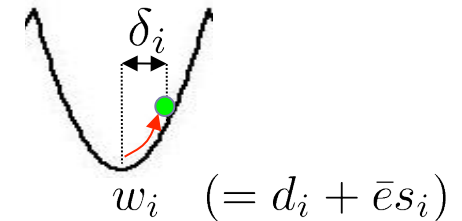
$$\frac{\partial \tilde{\Phi}}{\partial \mathbf{e}} = 0 \Rightarrow \mathbf{e}(\mathbf{s}, \mathbf{d}) = (\mathbf{1} + \mathbf{K})^{-1}(\mathbf{d} + \bar{\mathbf{e}}\mathbf{s})$$



Automaton

$$\delta_i = e_i - w_i \quad \text{elastic strain}$$

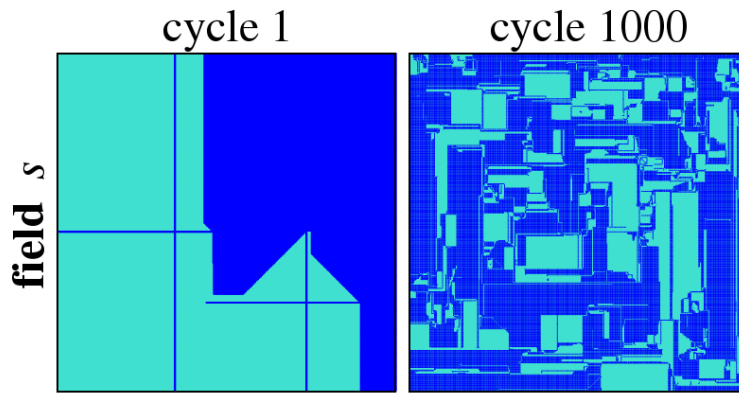
$$\delta^-(g) < \delta_i < \delta^+(g)$$



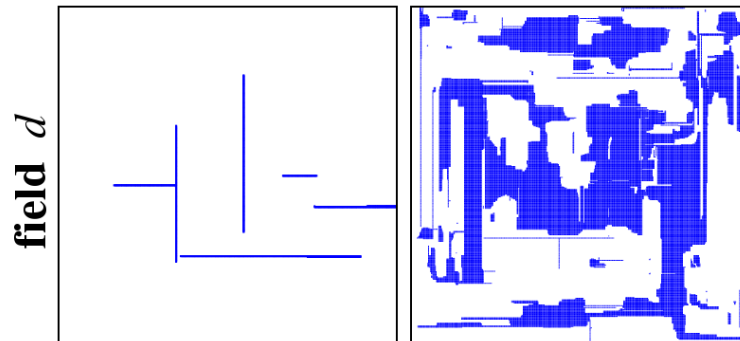
$$r = \begin{cases} \pm \bar{e}, & \text{phase} \\ \pm (1 - 2\bar{e}), & \text{slip} \end{cases}$$

Microstructure formation

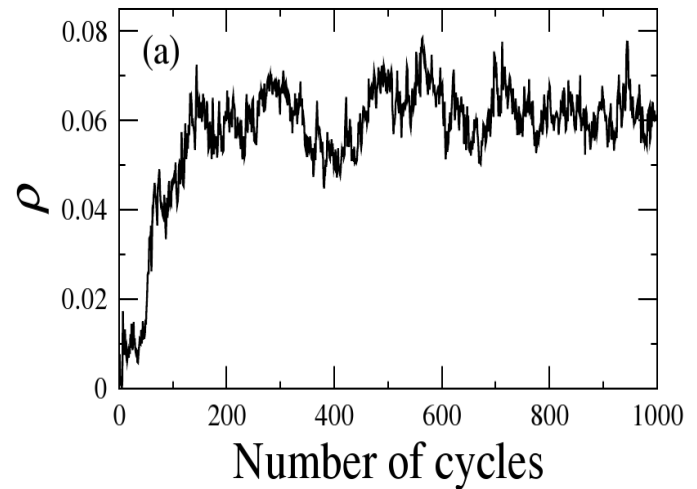
Transformation



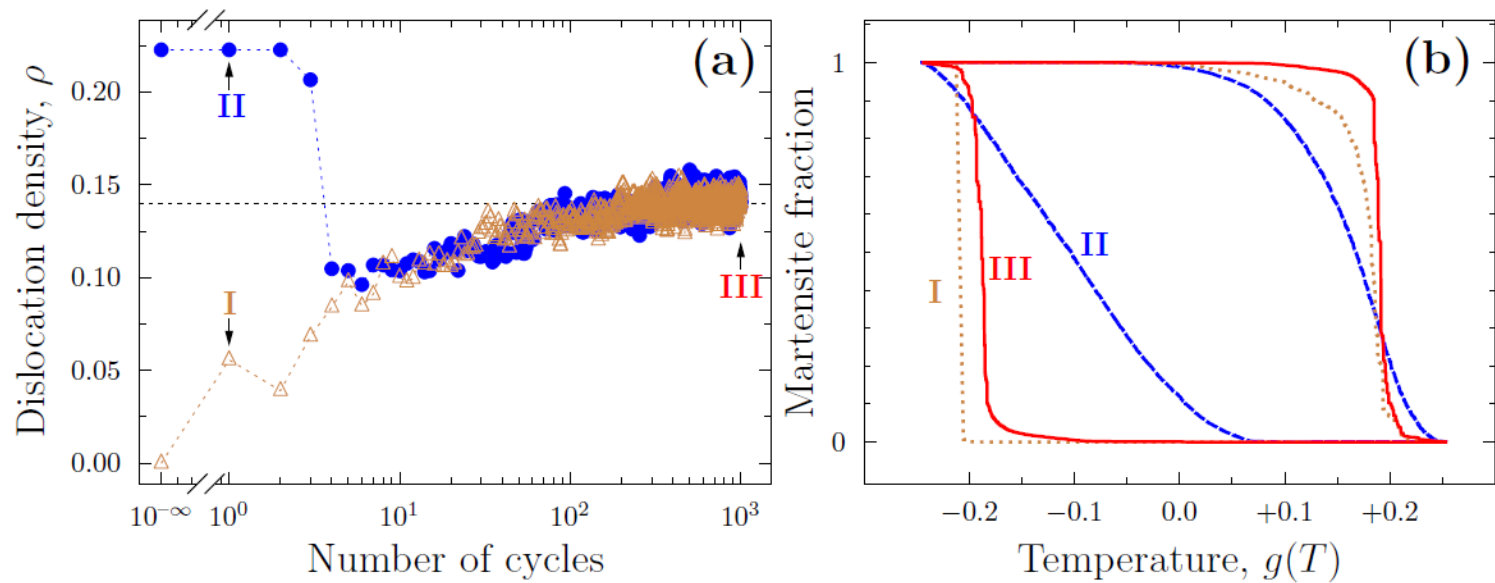
Plastic deformation



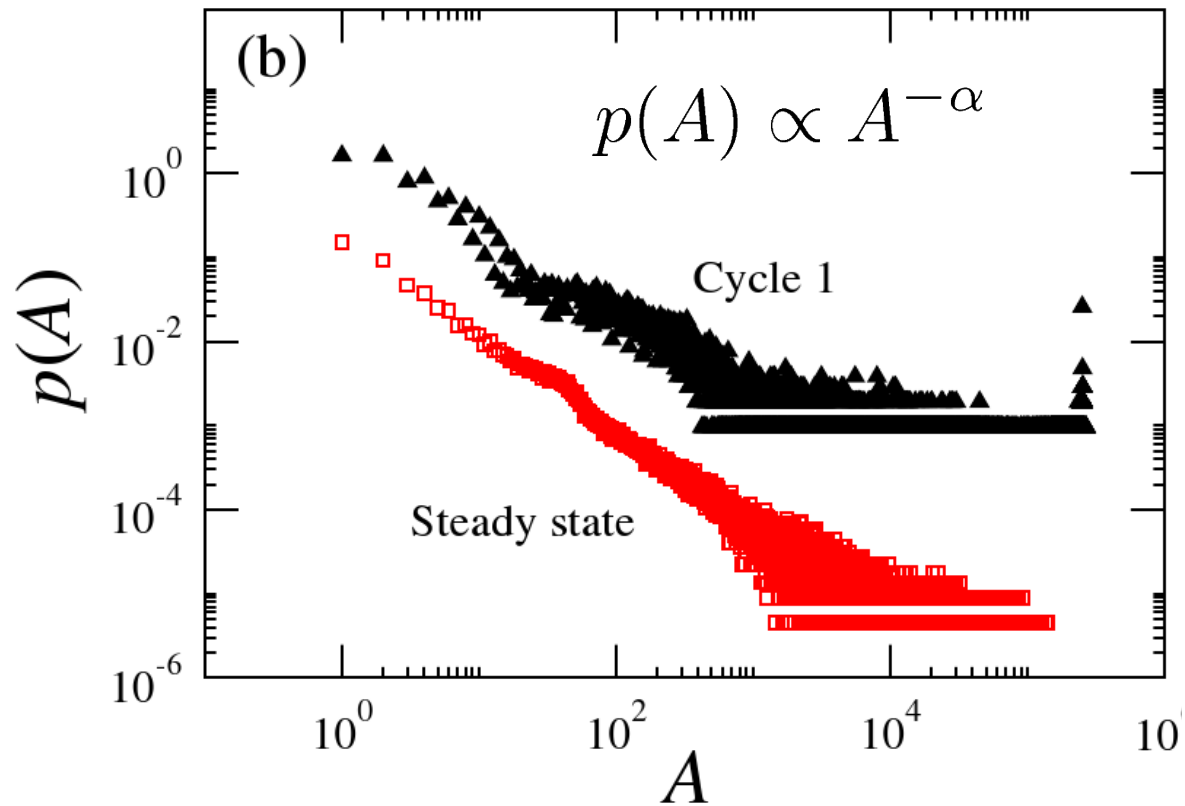
Density of dislocations



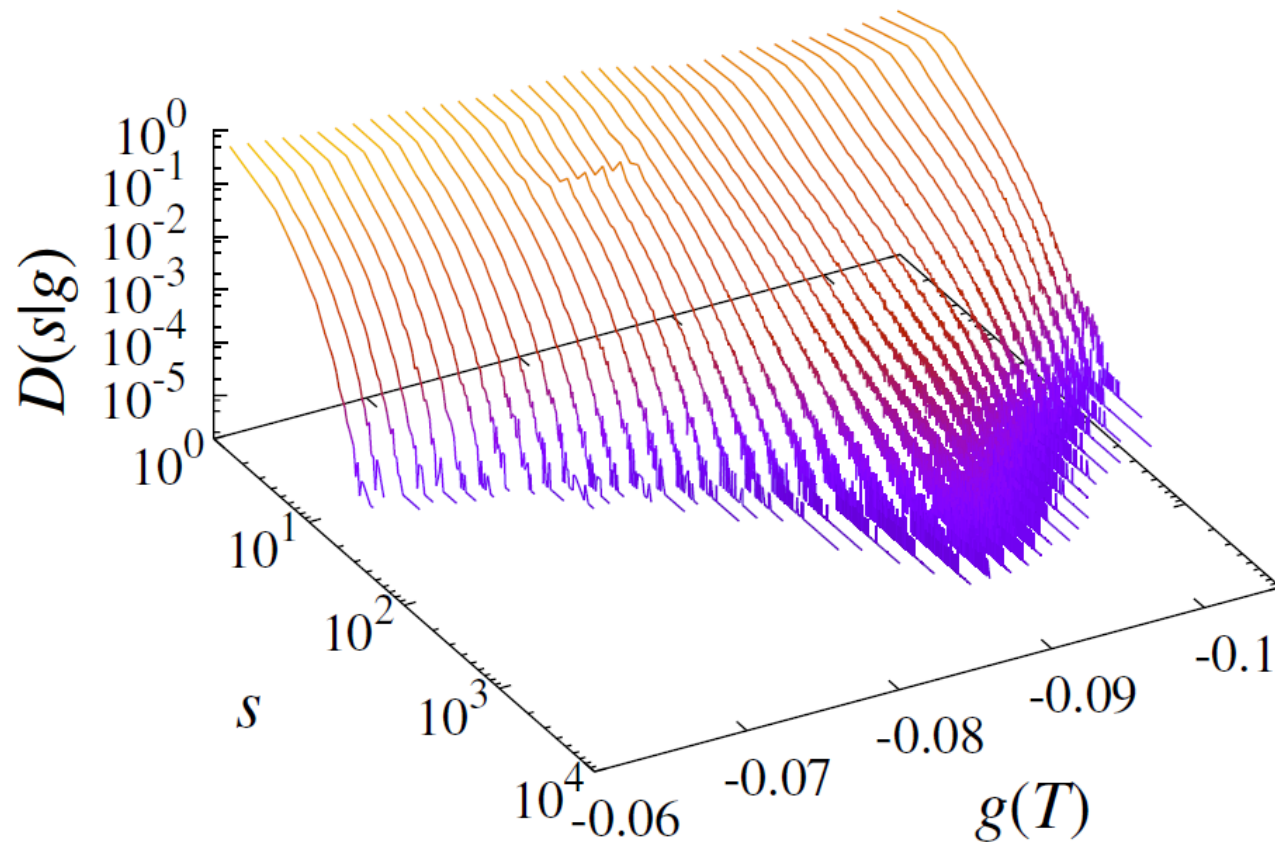
Shakedown state



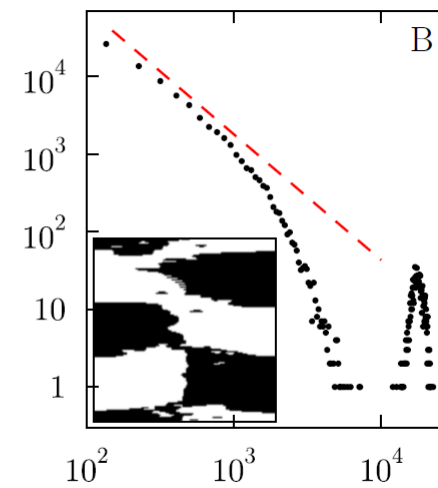
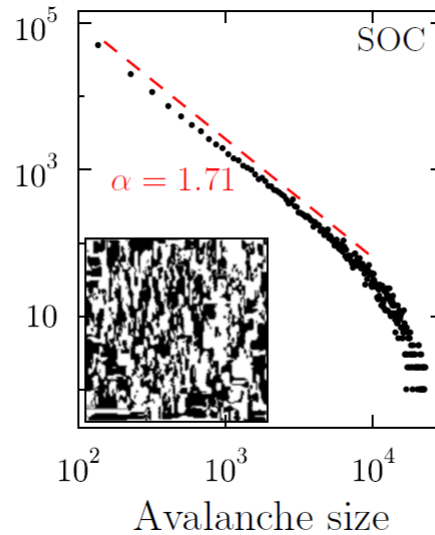
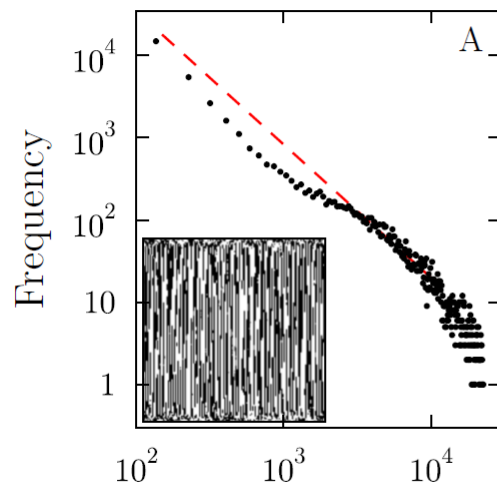
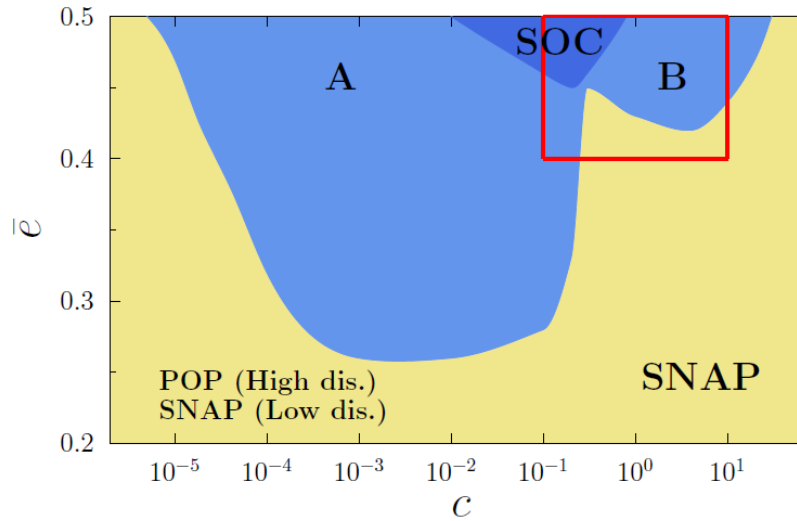
Power law behaviour



Fluctuation statistics varies along the cycle



Phase diagram

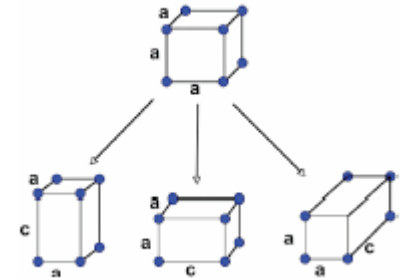


Inertia?

Direct continuum modeling

$$F_{\text{tetra}} = A_2(T) (\epsilon_2^2 + \epsilon_3^2) + A_3 \epsilon_3 (\epsilon_3^2 - 3\epsilon_2^2) + A_4 (\epsilon_2^2 + \epsilon_3^2)^2$$

$$F_{\text{grad}} = \frac{\beta}{2} (|\nabla e_2|^2 + |\nabla e_3|^2)$$



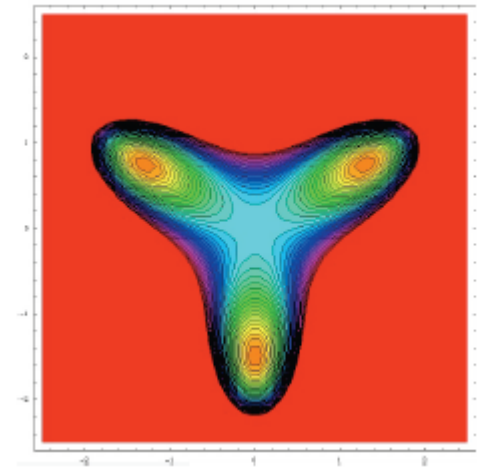
$$e_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

$$e_2 = (e_{11} - e_{22}) / \sqrt{2}$$

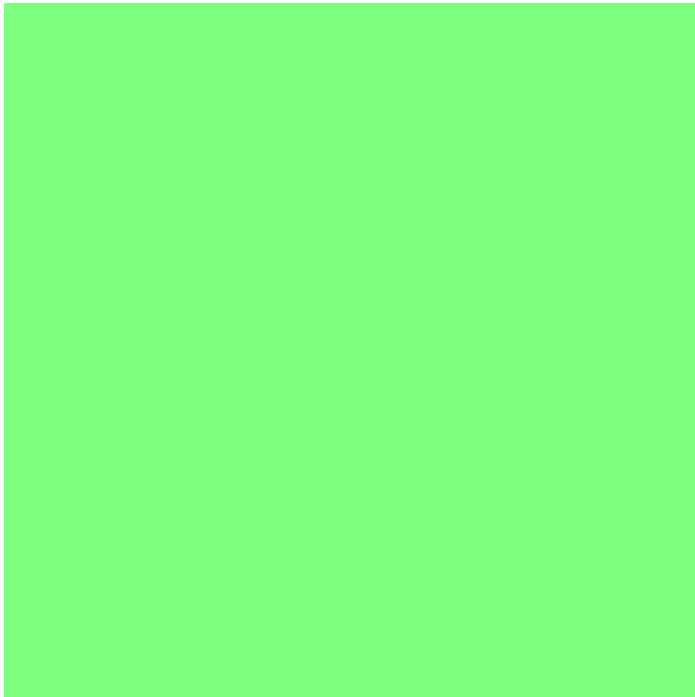
$$e_3 = (e_{11} + e_{22} - 2e_{33}) / \sqrt{6}$$

$$T = \frac{1}{2} \rho \dot{u}_i^2$$

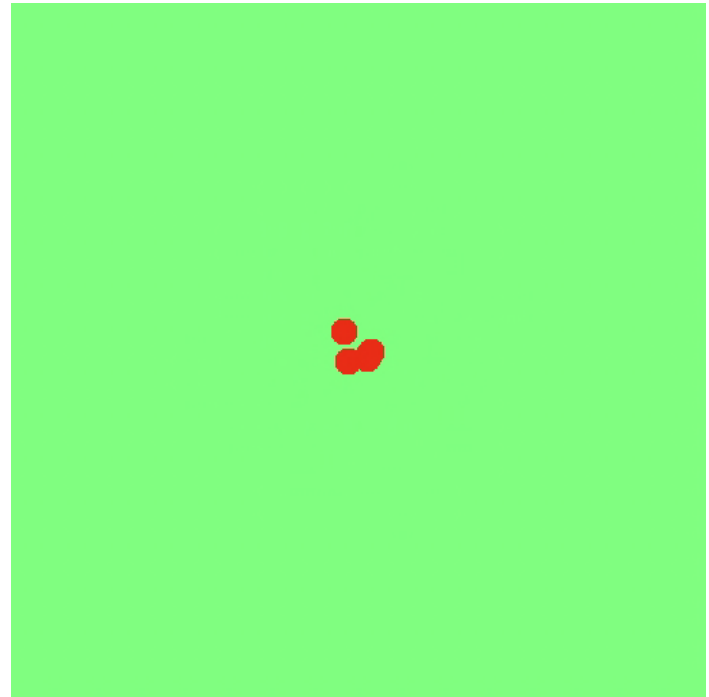
$$R = \frac{1}{2} \gamma (\dot{e}_2^2 + \dot{e}_3^2)$$



Numerical simulations

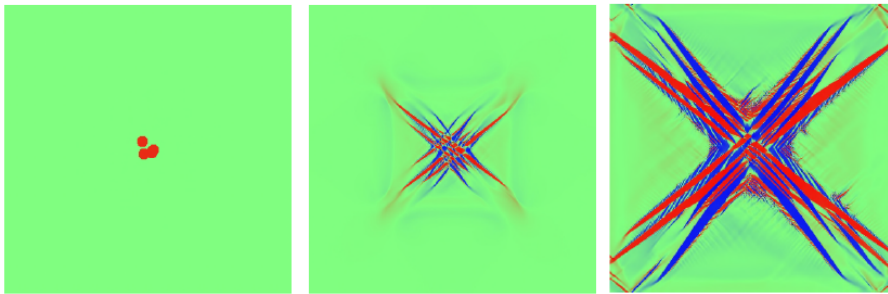


over-derdamped



underdamped

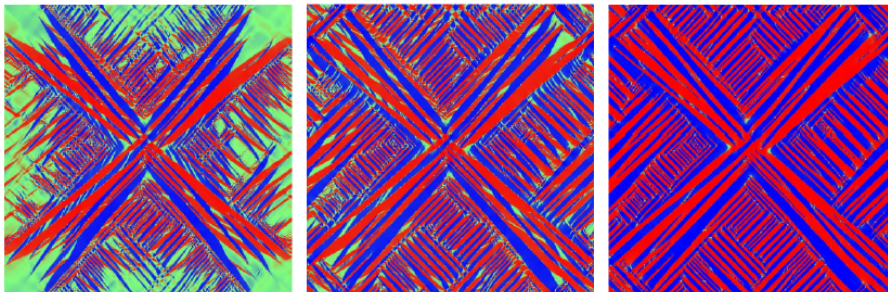
Inertia-induced complexity



(a)

(b)

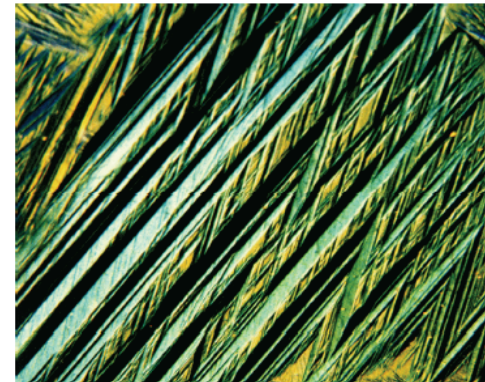
(c)



(d)

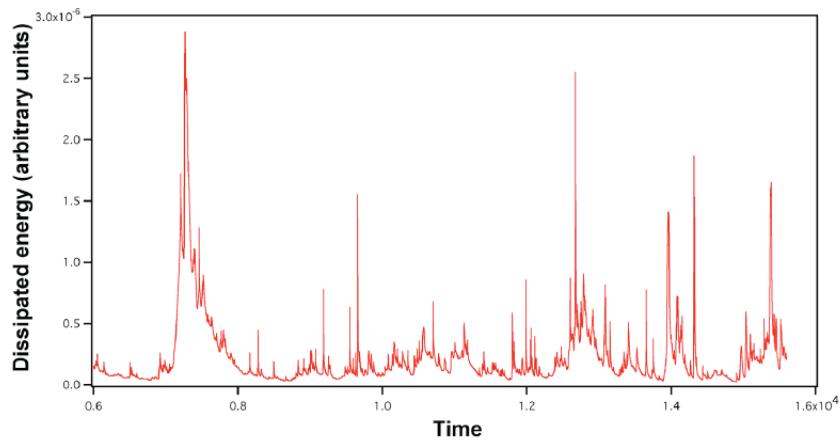
(e)

(f)



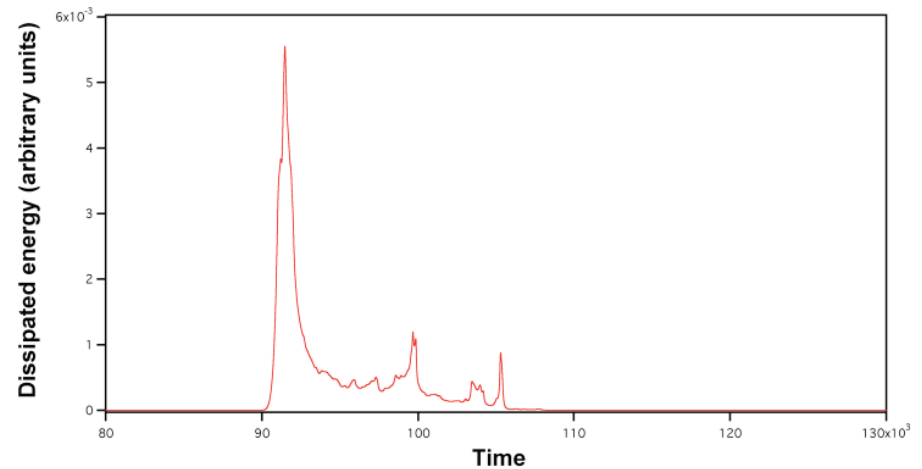
Cu-Zn-Al (3 mm x 2 mm)(Ortin - Lyon INSA)

Acoustic emission

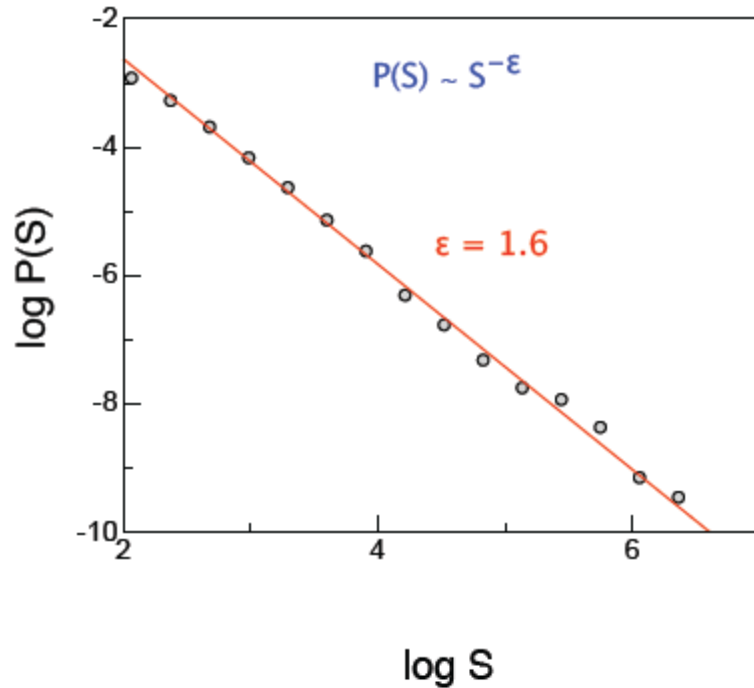


Inertial
dynamics

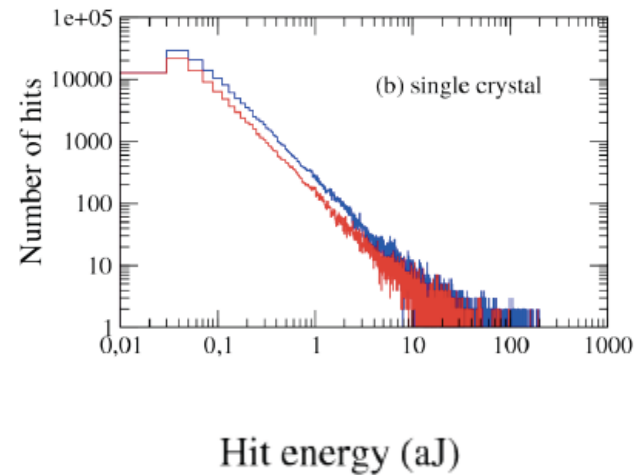
Overdamped
dynamics



Critical regime



comparison with AE experiment on $\text{Fe}_{68.8}\text{Pd}_{31.2}$
(E. Bonnot et al., PRB, 2009)



	ϵ
experiment	1.64 ± 0.1
simulation	1.60 ± 0.1

Mass-spring chain

$$\rho \ddot{u}_i = \frac{1}{\epsilon} \left[\phi' \left(\frac{u_{i+1} - u_i}{\epsilon} \right) - \phi' \left(\frac{u_i - u_{i-1}}{\epsilon} \right) + \frac{\gamma}{\epsilon} (\dot{u}_{i+1} + \dot{u}_{i-1} - 2\dot{u}_i) + \frac{\mu}{\epsilon} (u_{i+2} + u_{i-2} - 2u_i) \right]$$

inertia

NN interaction

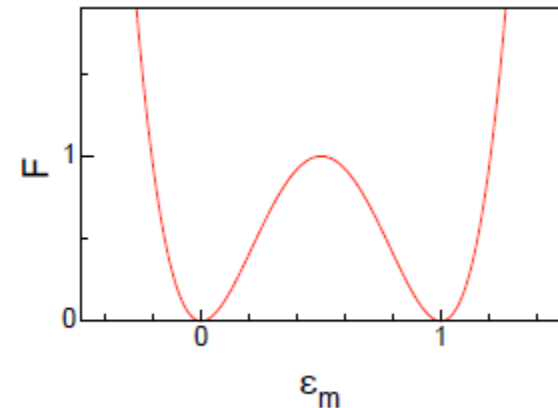
viscosity

NNN interaction

Quasi-static driving

$$u_0 = u_1 = 0$$

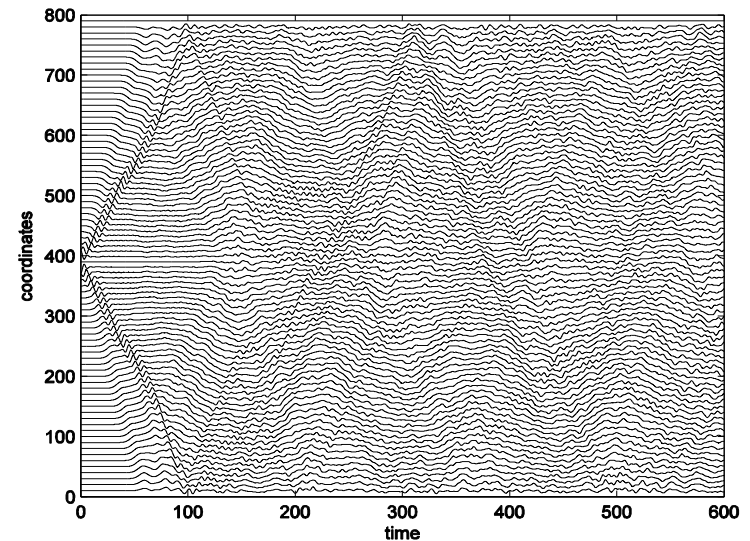
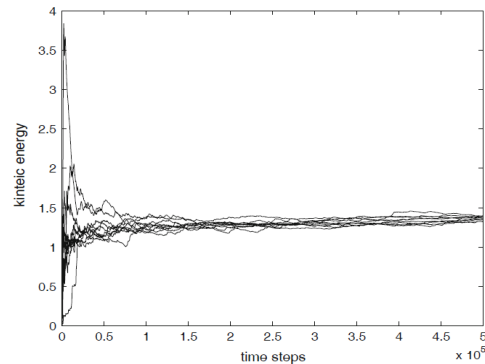
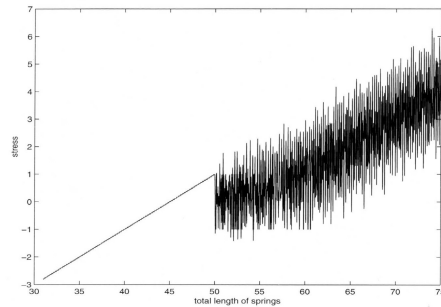
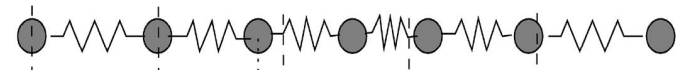
$$\dot{u}_{N-1} = \dot{u}_{N-2} = v$$



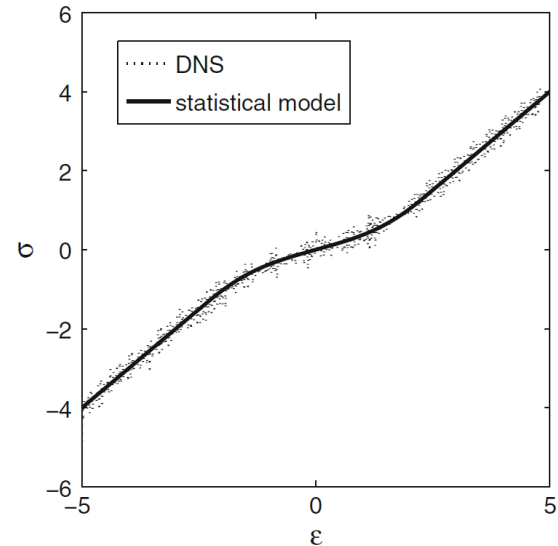
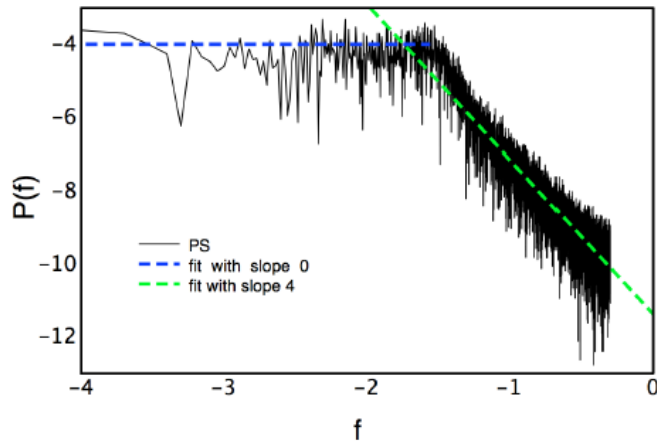
NN potential

Hamiltonian system (FPU)

$$\rho \ddot{u}_i = \frac{1}{\epsilon} \left[\phi' \left(\frac{u_{i+1} - u_i}{\epsilon} \right) - \phi' \left(\frac{u_i - u_{i-1}}{\epsilon} \right) + \frac{\gamma}{\epsilon} (\dot{u}_{i+1} + \dot{u}_{i-1} - 2\dot{u}_i) + \frac{\mu}{\epsilon} (u_{i+2} + u_{i-2} - 2u_i) \right]$$



Thermodynamic behavior

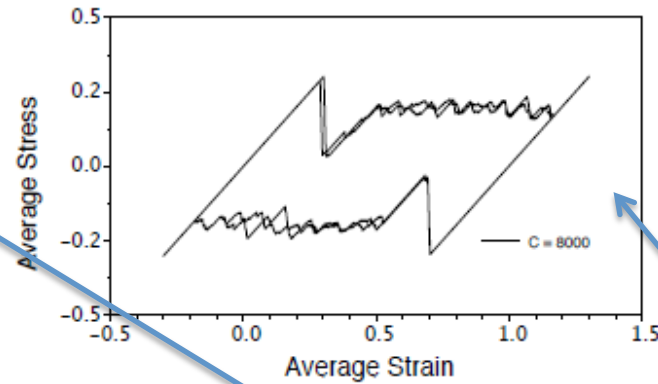
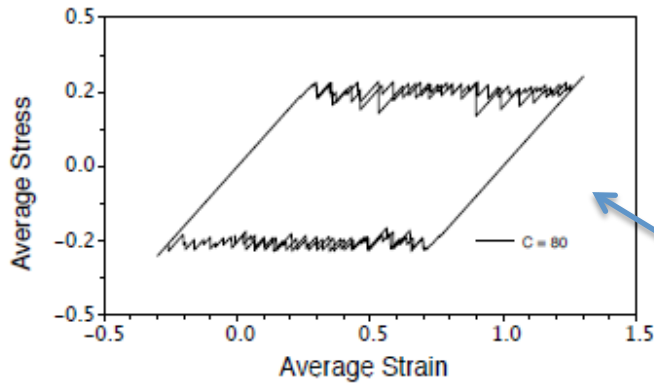


$$E_p = \sum_{i=1}^N \frac{\mu_i}{2} (\epsilon_i + \alpha_i)^2$$



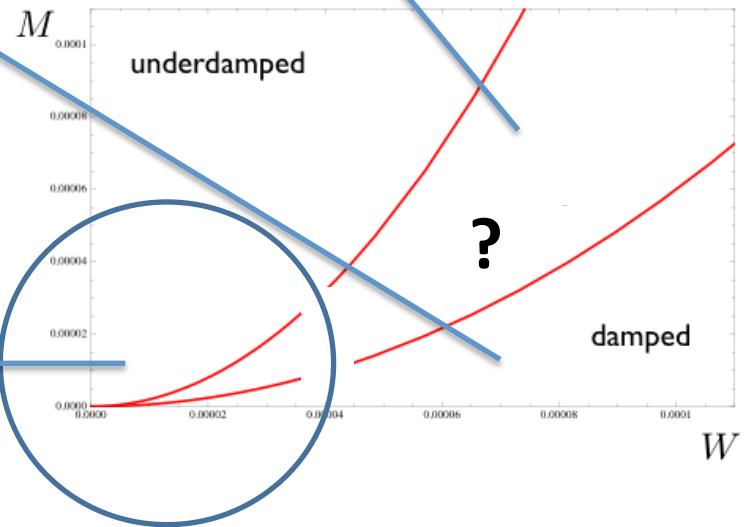
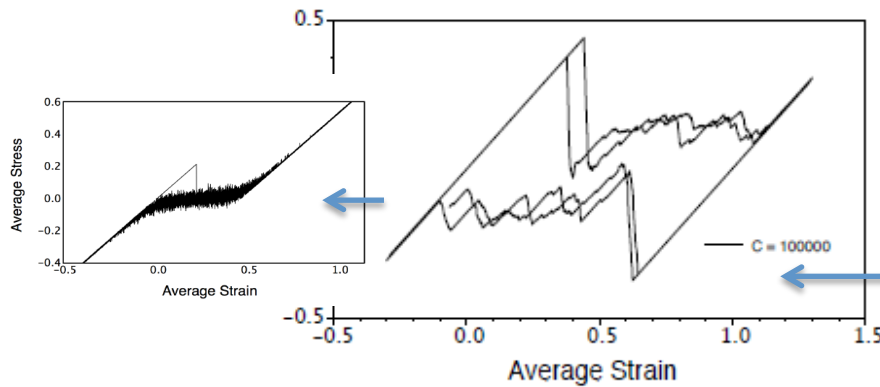
$$Z_p(\bar{\epsilon}, \beta) = N^{-1/2} \frac{(2\pi)^{\frac{N-1}{2}}}{\beta^{\frac{N-1}{2}}} \sum_{p=0}^N \frac{N!}{p!(N-p)!} (\mu_1^{p/N} \mu_2^{1-p/N})^{-N/2} \left(\frac{p/N}{\mu_1} + \frac{1-p/N}{\mu_2} \right)^{-1/2} \\ \times \exp\left(-\frac{\beta}{2} N \left(\frac{p/N}{\mu_1} + \frac{1-p/N}{\mu_2} \right)^{-1} \left(\bar{\epsilon} + \frac{p}{N} a_1 + \left(1 - \frac{p}{N} \right) a_2 \right)^2 \right).$$

From under to over damping



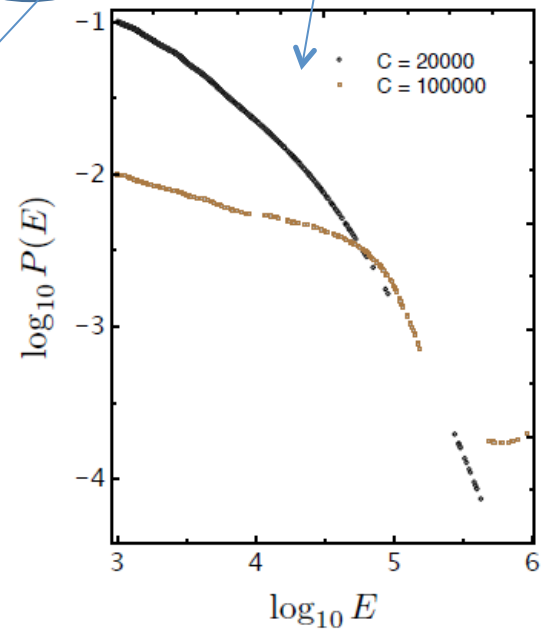
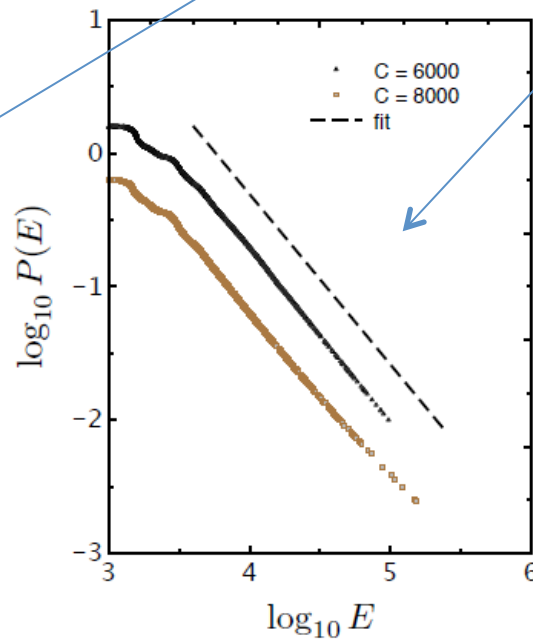
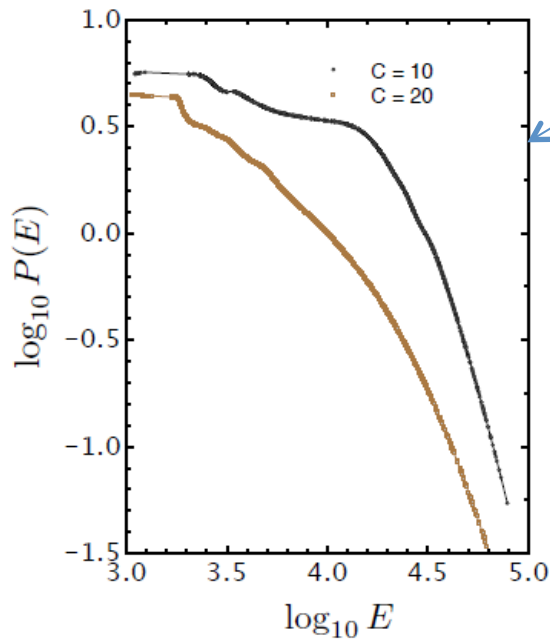
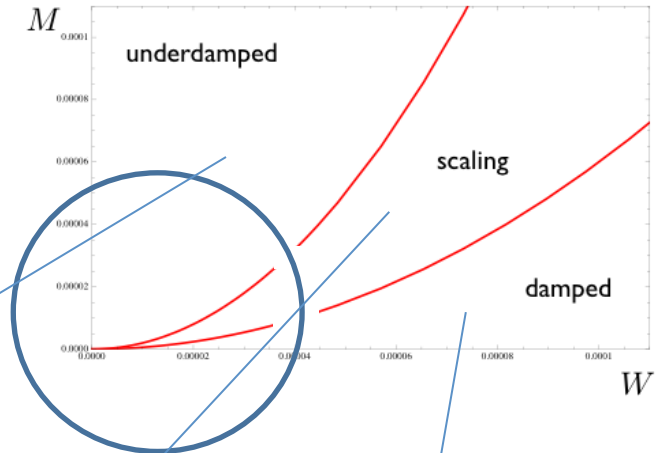
$$M = \frac{v^2 \rho}{\kappa}, W = \frac{\gamma v}{L \kappa}$$

$$C = M/W^2$$



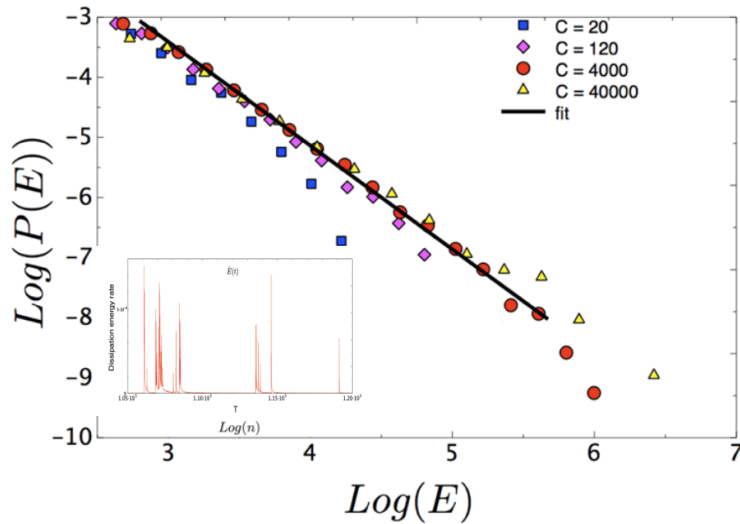
(c)

Scaling regime

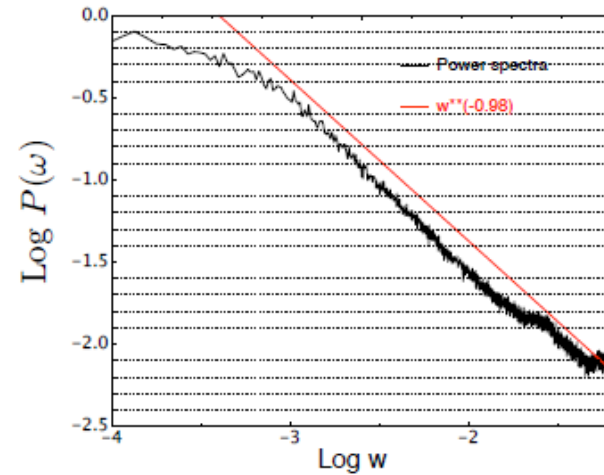


Scaling regime

$$P(E) \sim E^{-\epsilon} \quad \epsilon = 1.75 \pm 0.05$$



$$P(\omega) \sim \omega^{-\alpha} \quad \alpha = 0.98 \pm 0.1$$



Conclusions

- There are (at least) two paths to scaling in martensites. The domineering mechanism depends on micro-parameters: transformation strain, elastic moduli, barrier structure.
- Limited plasticity allows the system to generate (after training) a statistically stable configuration of defects which makes subsequent phase transformation scale free.
- Even small inertia can become crucial if the system is close to marginal stability because it can generate a dynamic disorder (ergodicity on a subset of the configurational space?)