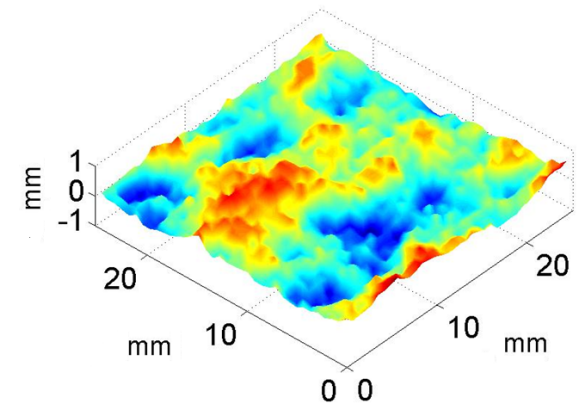
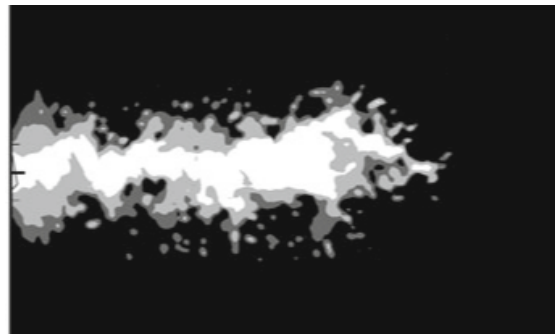
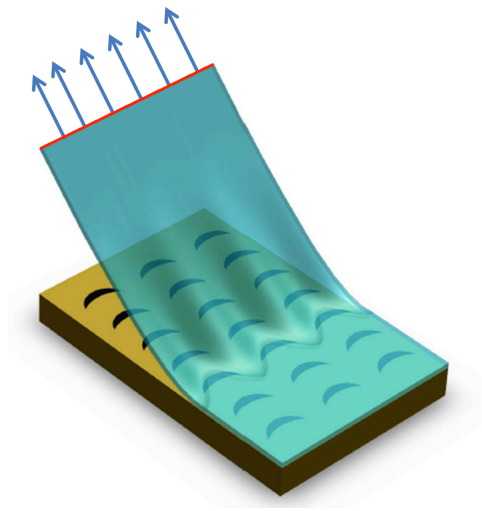




Fracture mechanics of homogeneous and heterogeneous brittle materials: An overview



The fracture of materials...

One generally wants



to avoid it...

The fracture of materials...

One generally wants



to avoid it...

But sometimes one
wants it badly

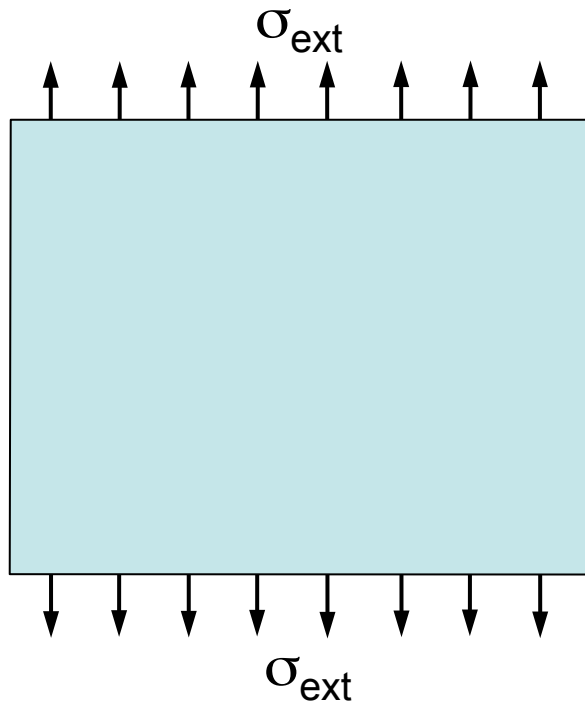


Good Question:

What in the world can be done about those
\$%^#&, customer-infuriating plastic clamshells?

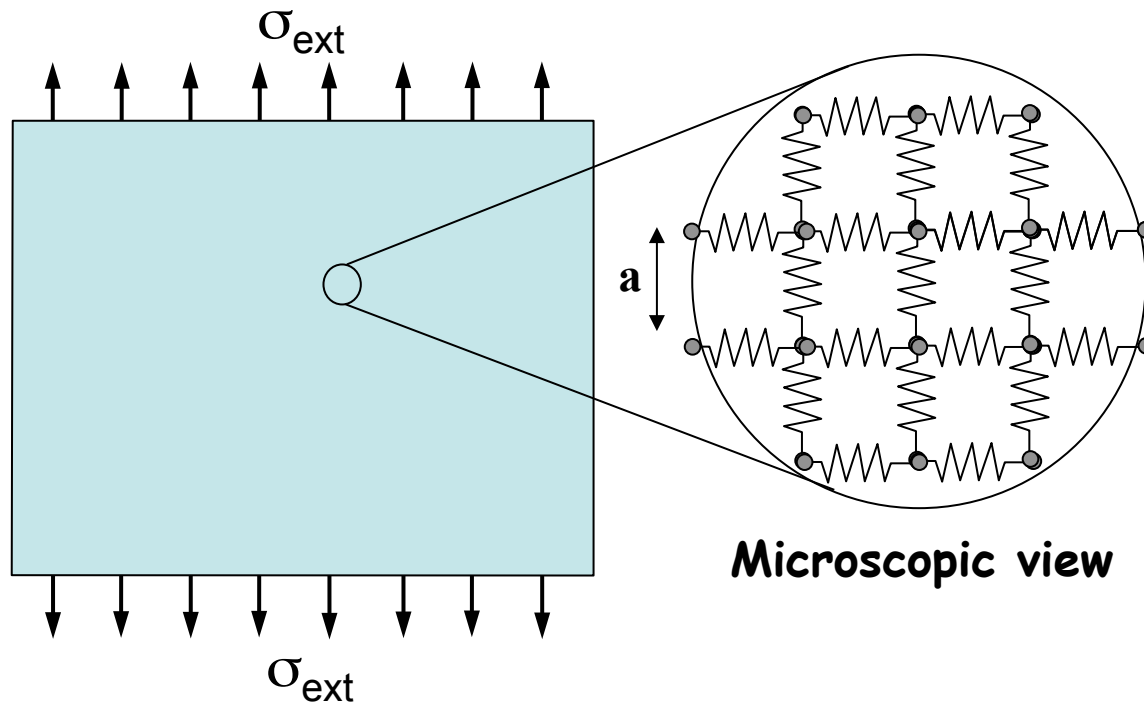
How to predict the strength of brittle solids

M. Marder and J. Fineberg 1996



How to predict the strength of brittle solids

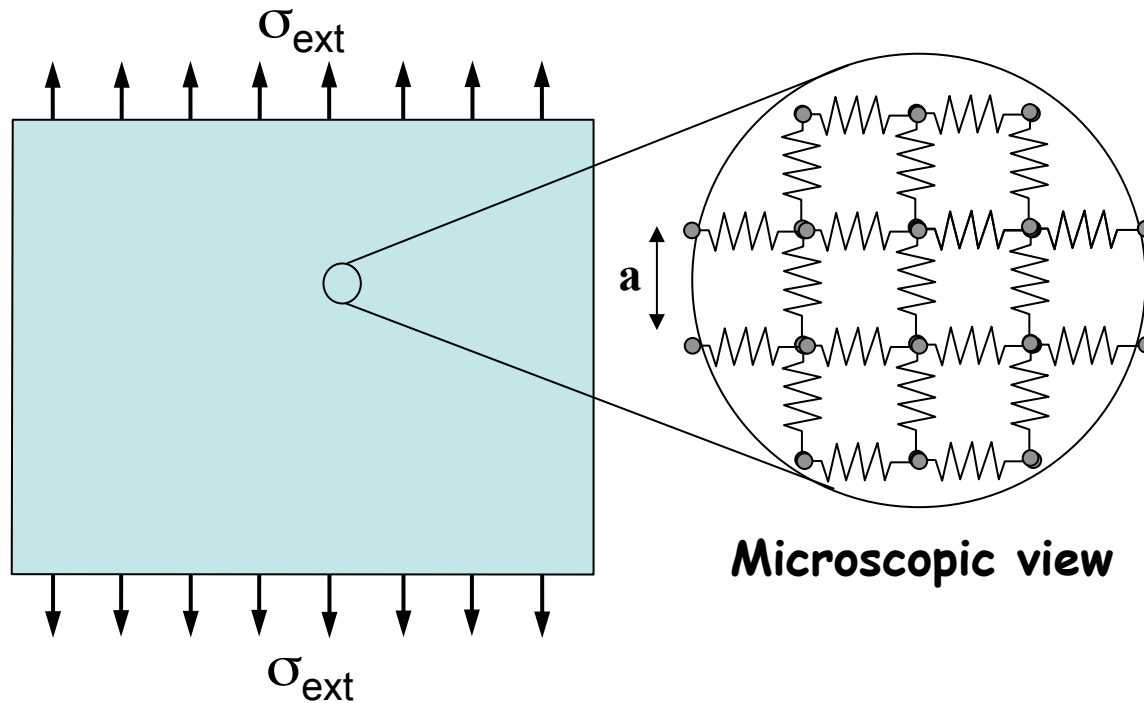
M. Marder and J. Fineberg 1996



E Young's modulus
a distance between atoms
 σ_F stress at failure

How to predict the strength of brittle solids

M. Marder and J. Fineberg 1996



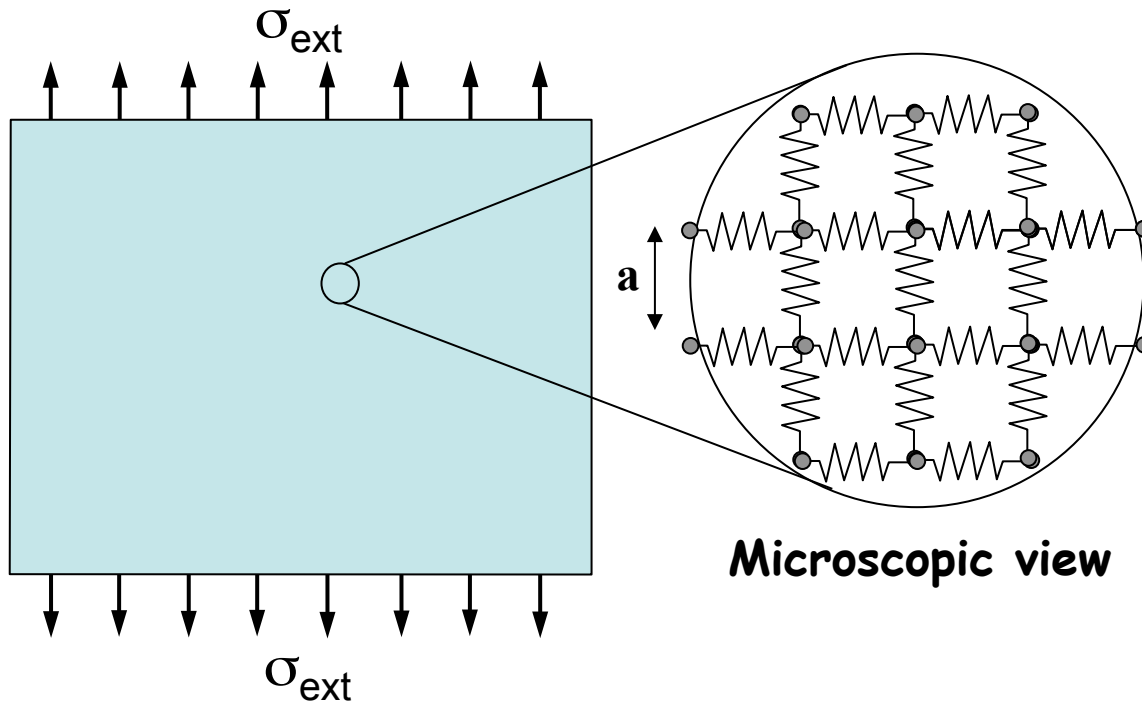
E Young's modulus
 a distance between atoms
 σ_F stress at failure

- ➔ Stiffness of elastic bonds $\sim E$
- ➔ Deformation at failure $\sim a$
- ➔ Force at failure $\sim E a$

➔ **Failure: $\sigma_F \sim E$**

How to predict the strength of brittle solids

M. Marder and J. Fineberg 1996



E Young's modulus
 a distance between atoms
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- ➔ Stiffness of elastic bonds $\sim E$
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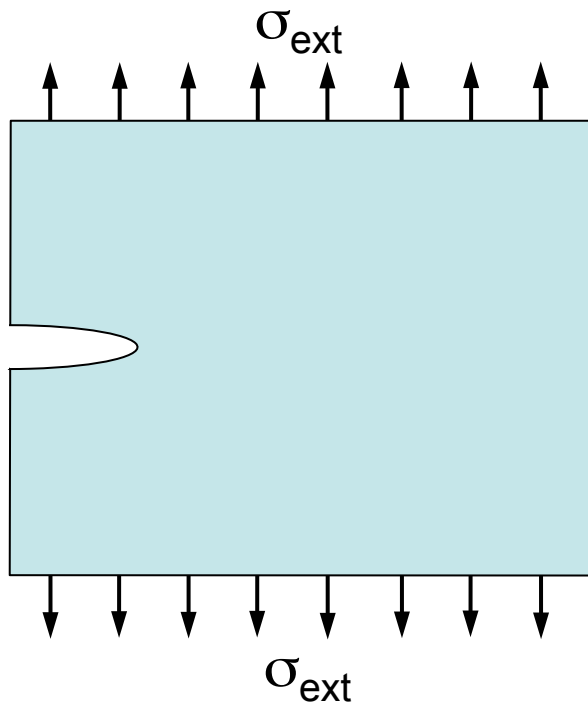
➔ Failure: $\sigma_F \sim E$

	E	σ_f
Steel	200 GPa	0.1-2 GPa
Glass	70 GPa	300 MPa
Al_2O_3	400 GPa	<100 Mpa

> Not like this!

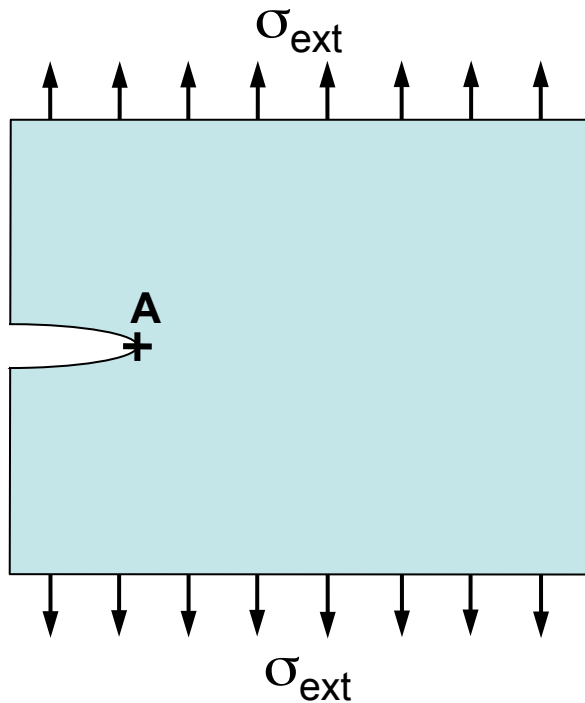
Stress concentration at defects

C. E. Inglis 1913



Stress concentration at defects

C. E. Inglis 1913



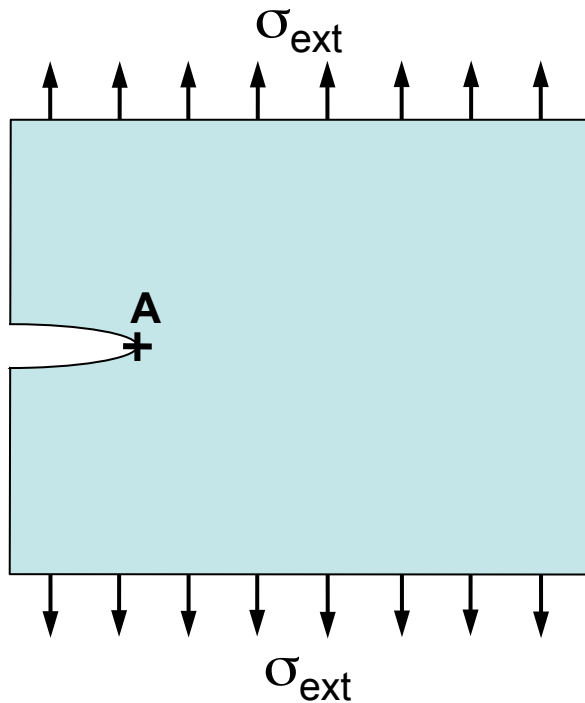
Stress amplification

$$\sigma_A \sim \frac{\sigma_{ext}}{\sqrt{\rho}}$$

- ➡ Applied remote stress $\sim \sigma_{ext}$
- ➡ Local stress in A $\sim \sigma_A$
- ➡ Radius of curvature of the defect in A $\sim \rho$

Stress concentration at defects

C. E. Inglis 1913



Stress amplification

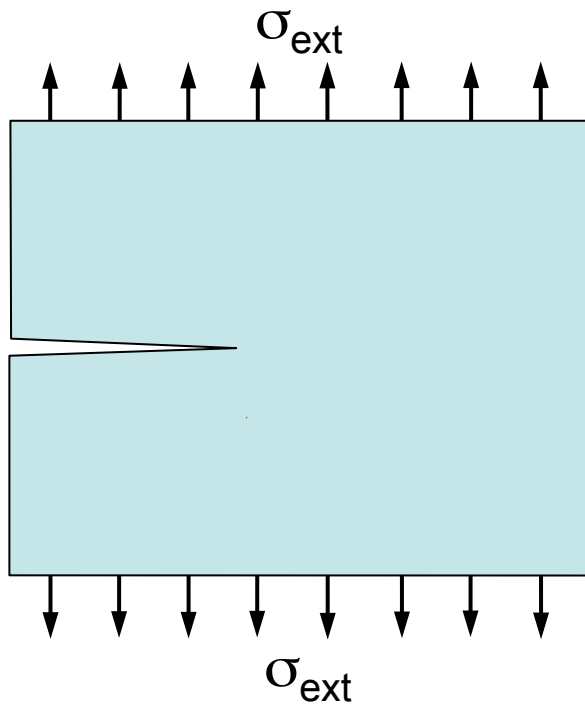
$$\sigma_A \sim \frac{\sigma_{ext}}{\sqrt{\rho}}$$

- ➔ Applied remote stress $\sim \sigma_{ext}$
- ➔ Local stress in A $\sim \sigma_A$
- ➔ Radius of curvature of the defect in A $\sim \rho$

➔ **Materials submitted to stress levels actually larger than the applied external stress**

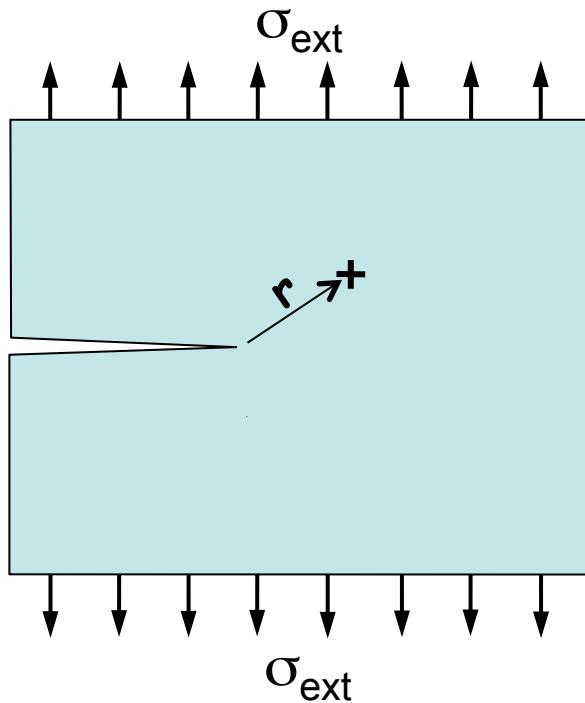
The slit crack model

G. Irwin 1957



The slit crack model

G. Irwin 1957



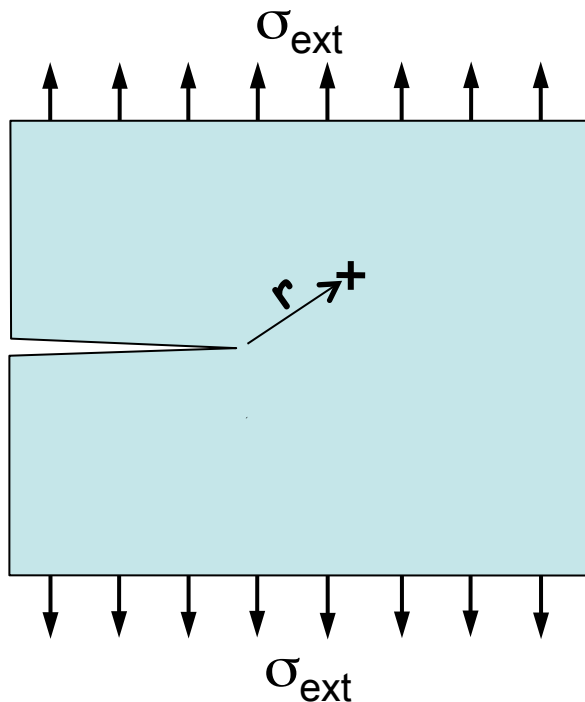
Stress field diverges
at the crack tip:

$$\sigma \sim \frac{K}{\sqrt{r}}$$

- ➔ Stress intensity factor $\sim K$
- ➔ Distance to crack tip $\sim r$

The slit crack model

G. Irwin 1957



Stress field diverges
at the crack tip:

$$\sigma \sim \frac{K}{\sqrt{r}}$$

- ➡ Stress intensity factor $\sim K$
- ➡ Distance to crack tip $\sim r$

Stress based
criterion for failure



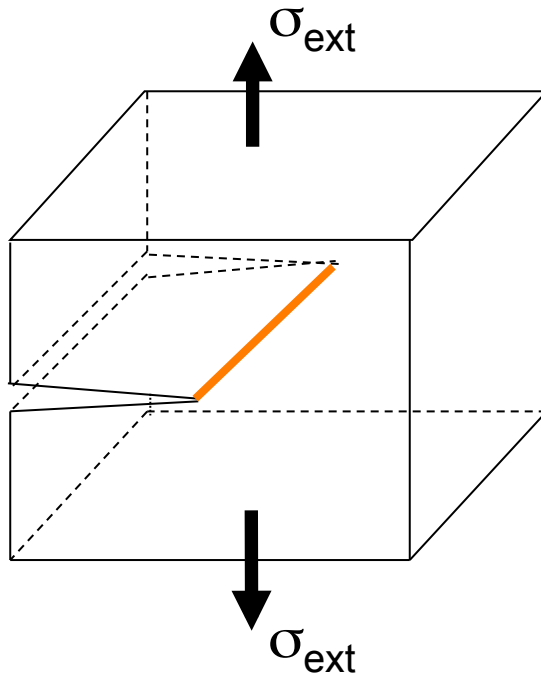
Materials would have no
resistance to failure

➤ Not like this neither!

A powerful predicting tool: Linear Elastic Fracture Mechanics

Predicting the stability of cracks in homogeneous media

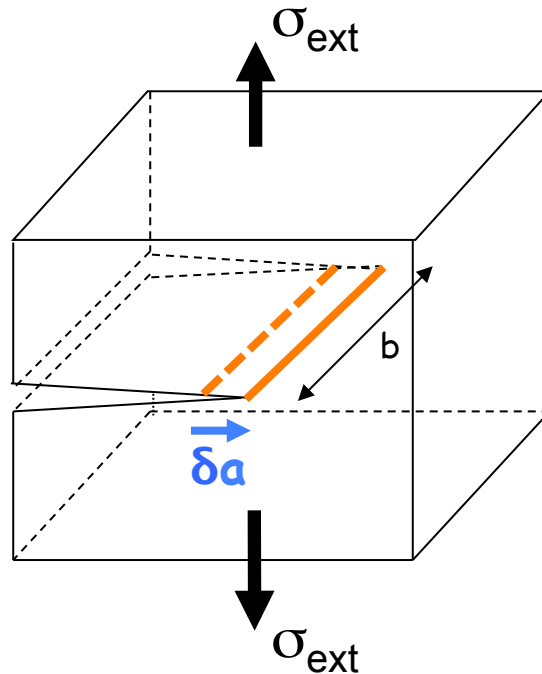
*A.A. Griffith 1920
J.R. Rice 1968*



A powerful predicting tool: Linear Elastic Fracture Mechanics

Predicting the stability of cracks in homogeneous media

A.A. Griffith 1920
J.R. Rice 1968



Energy balance:

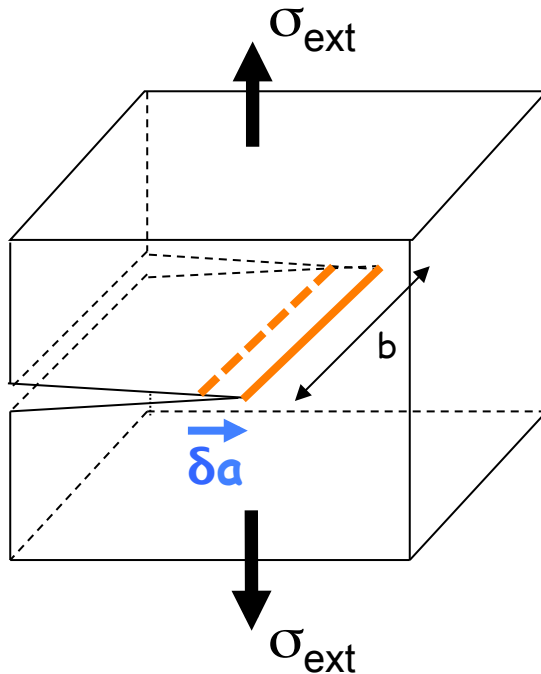
$$\delta W_{ext} = \delta E_{el} + \delta E_s$$

Work of the external force Variation of elastic energy Variation of surface energy

A powerful predicting tool: Linear Elastic Fracture Mechanics

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Work of the
external force

Variation of
elastic energy

Variation of
surface energy

Griffith's criterion:

Elastic energy
release rate

vs

Fracture energy

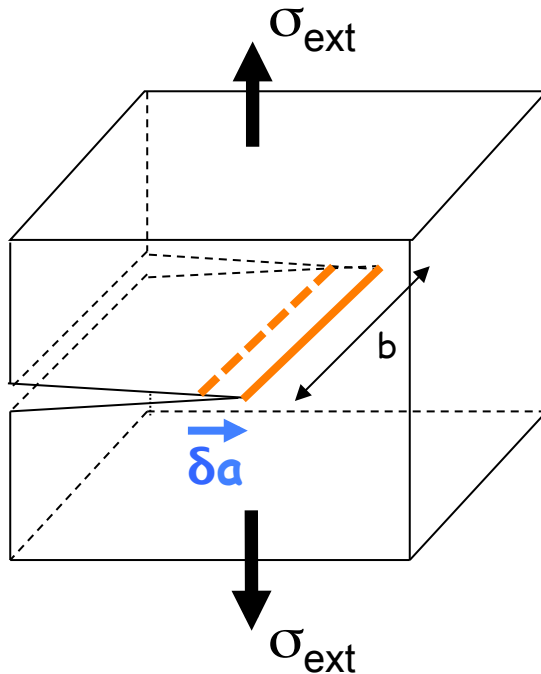
$$G = \delta(W_{ext} - \delta E_{el}) / (\delta a \cdot b)$$

$$G_c = \delta E_s / (\delta a \cdot b)$$

A powerful predicting tool: Linear Elastic Fracture Mechanics

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Griffith's criterion:

Elastic energy release rate vs Fracture energy

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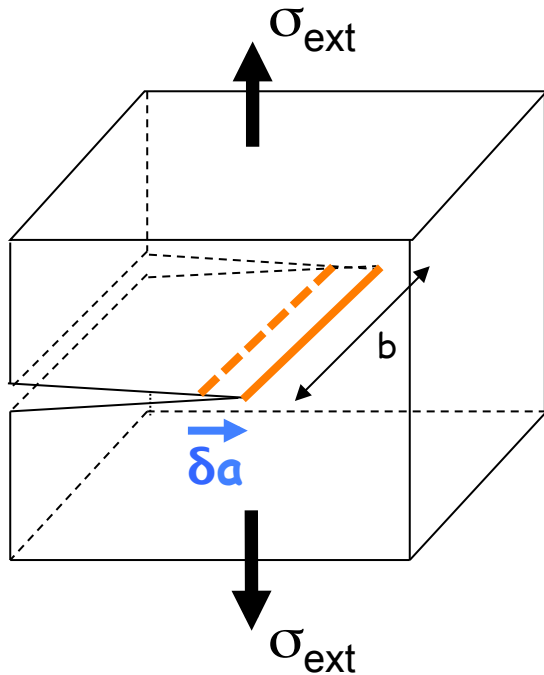
$G < G_c$ \longrightarrow Stable crack

$G = G_c$ \longrightarrow Propagating crack

A powerful predicting tool: Linear Elastic Fracture Mechanics

Predicting the stability of cracks in homogeneous media

A.A. Griffith 1920
J.R. Rice 1968



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$$\delta W_{\text{ext}} = \delta E_{\text{el}} + \delta E_{\text{s}}$$

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Variation of
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Variation of
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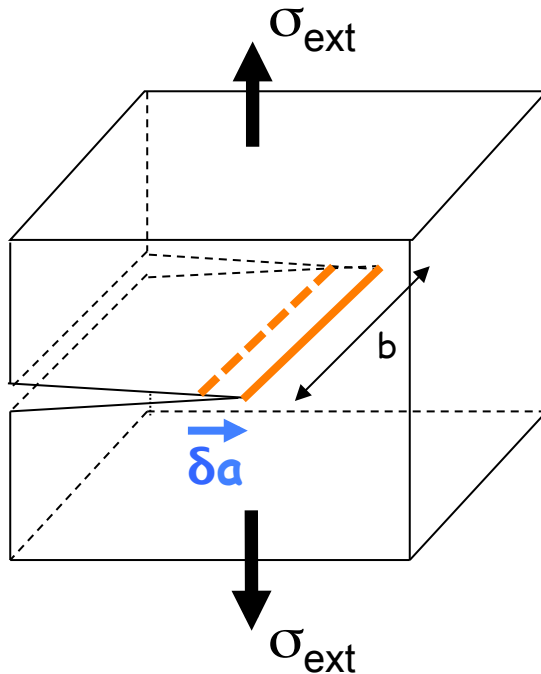
Relation between
stress intensity
factor and energy
release rate

$$G = \frac{K^2}{E} \sim \sigma_{\text{ext}}^2$$

A powerful predicting tool: Linear Elastic Fracture Mechanics

Predicting the stability of cracks in homogeneous media

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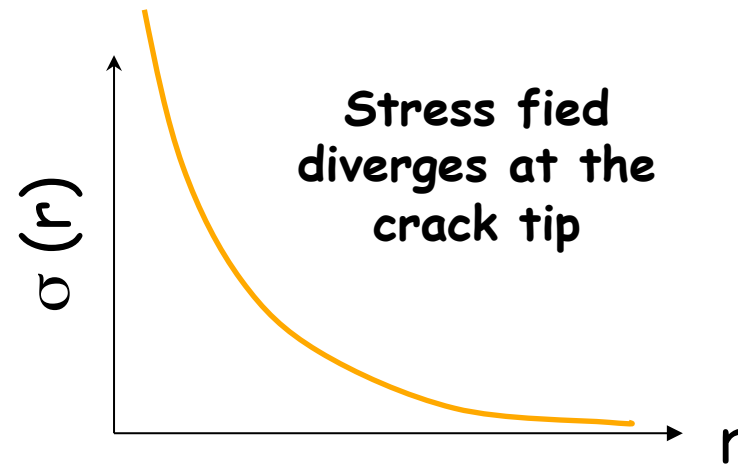
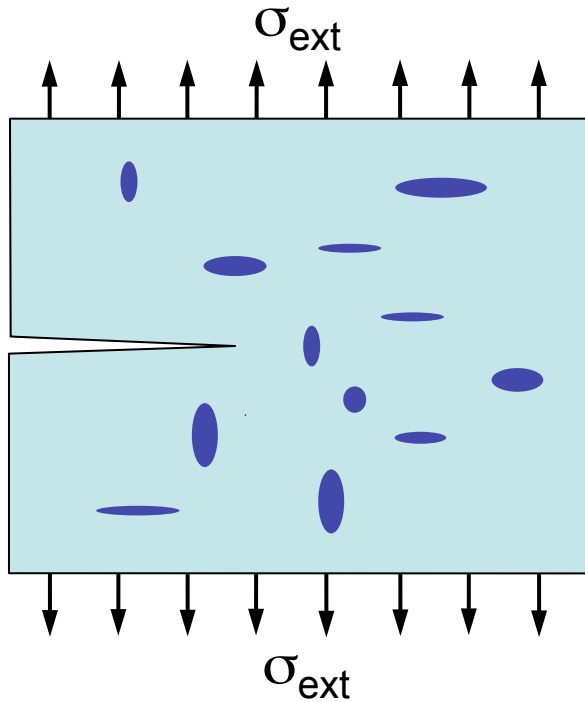
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Relation between stress intensity factor and energy release rate

$$G = \frac{K^2}{E} \sim \sigma_{\text{ext}}^2$$

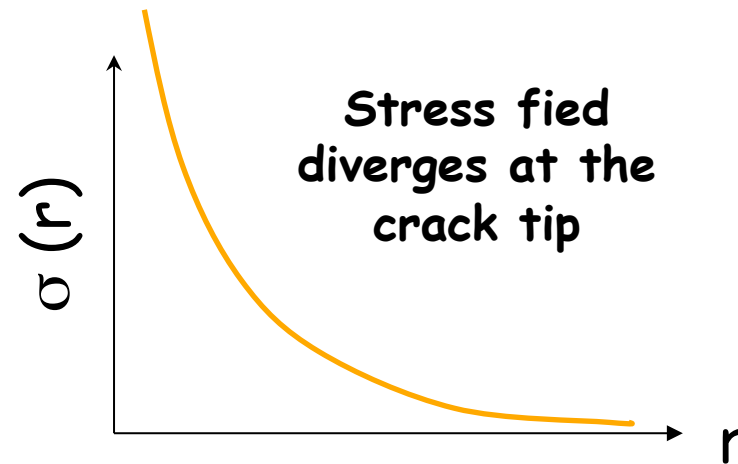
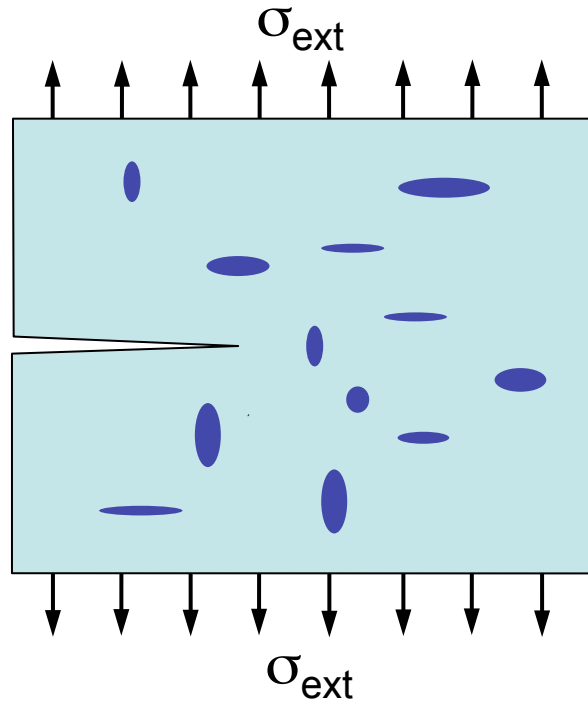
Powerful predictive approach, but limited to homogeneous media

The crack tip as a magnifying glass of the material heterogeneities



Macroscopic failure properties strongly dependent on microscopic material features

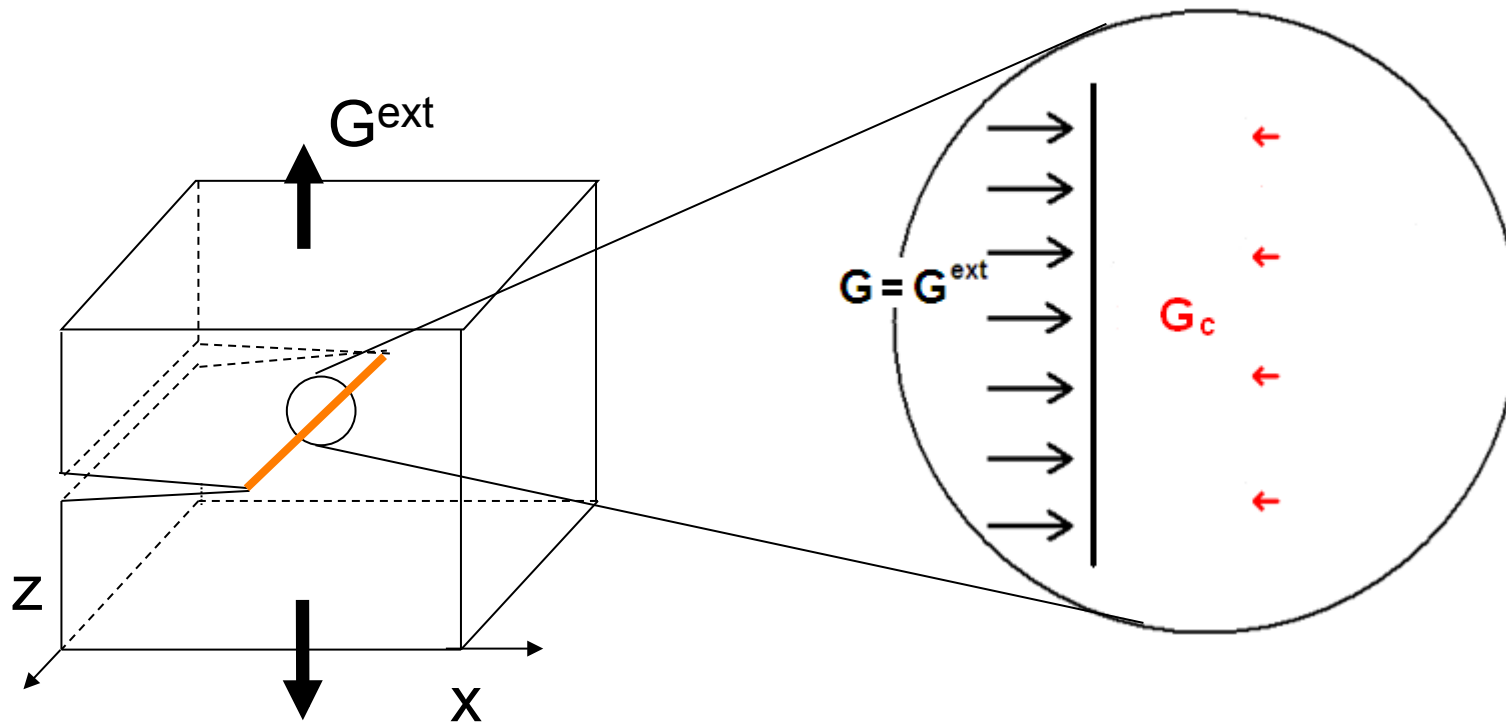
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Macroscopic failure properties strongly dependent on microscopic material features

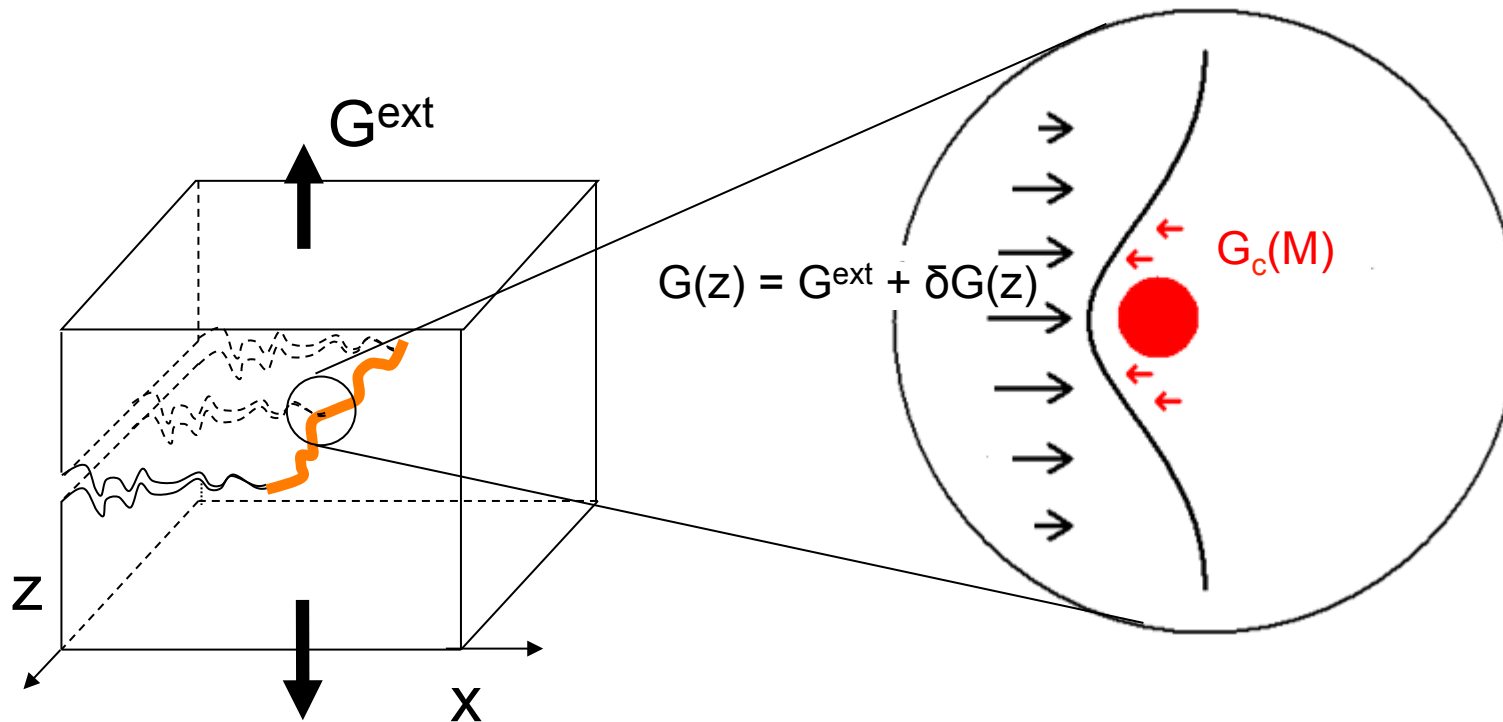
- ➔ Fracture as a complex and challenging multi-scale problem
- ➔ Opportunities for a rational design of materials with improved failure properties

What are the effects of heterogeneities on the propagation of a crack?



What are the effects of heterogeneities on the propagation of a crack?

Pinning and deformation of the crack front:

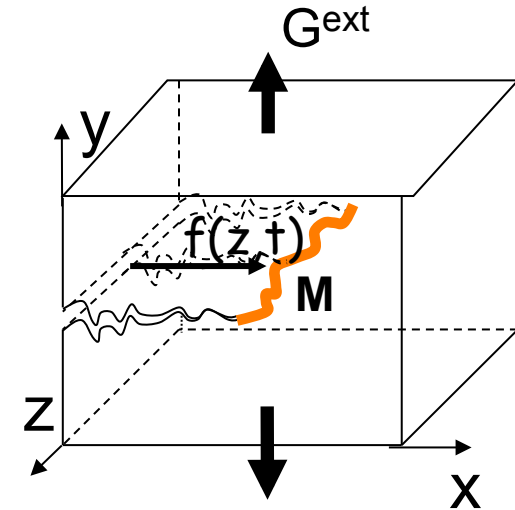


Fracture Mechanics of heterogeneous materials



Heterogeneous field of fracture energy $G_c(M)$

$$G_c(M) = \langle G_c \rangle + \delta G_c(M)$$



Fracture Mechanics of heterogeneous materials

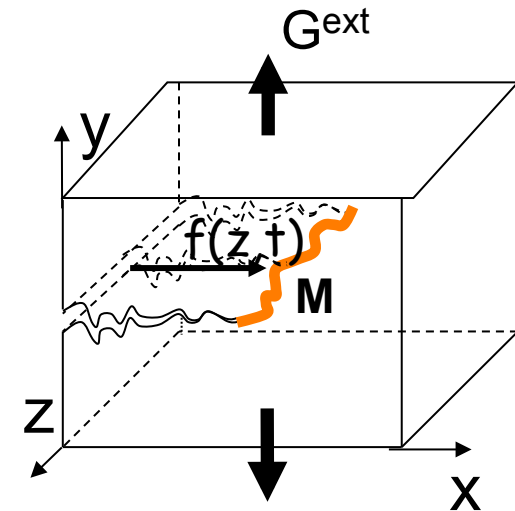
→ Heterogeneous field of fracture energy $G_c(M)$

$$G_c(M) = \langle G_c \rangle + \delta G_c(M)$$

→ Material elasticity: The crack front as an elastic interface

$$G(z) = G^{ext} + \frac{G^{ext}}{\pi} \int_{-\infty}^{\infty} \frac{f(z',t) - f(z,t)}{(z' - z)^2} dz'$$

J. R. Rice 1985



Hypothesis:

- Slow crack growth velocity
- Weakly heterogeneous material
- Very large sample

Fracture Mechanics of heterogeneous materials

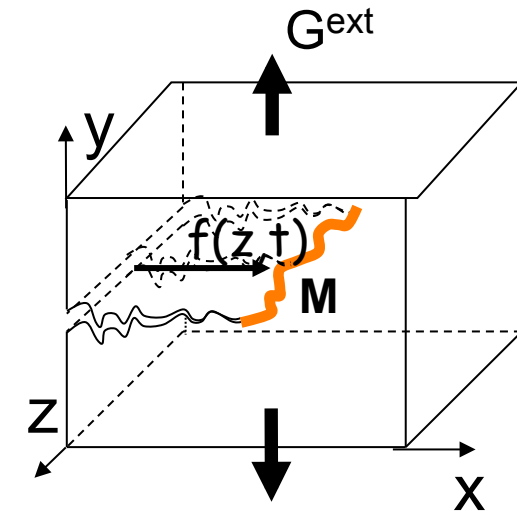
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Equation of motion for a crack

$$\mu \frac{\partial f(z, t)}{\partial t} \Big|_M = G(M) - G_c(M)$$

L. B. Freund 1990

Fracture Mechanics of heterogeneous materials

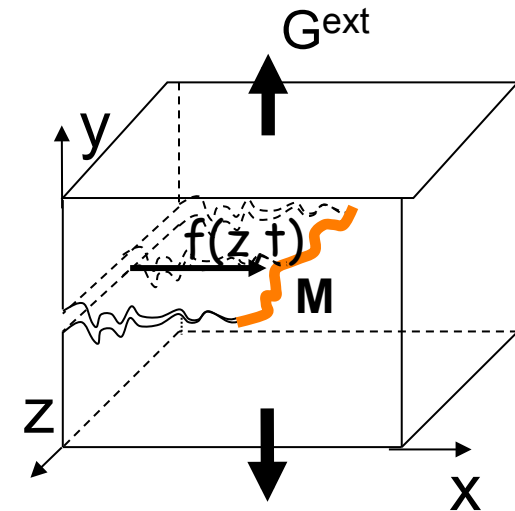
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J. R. Rice 1985



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Equation of motion for a crack

$$\mu \frac{\partial f(z,t)}{\partial t} = \left(G^{ext} - \langle G_c \rangle \right) + \frac{G^{ext}}{\pi} \int_{-\infty}^{\infty} \frac{f(z',t) - f(z,t)}{(z' - z)^2} dz' - \delta G_c(z, f(z,t))$$

J. Schmitbuhl et al. 1995, D. Bonamy et al. 2008, L. Ponson et al. 2010

Fracture Mechanics of heterogeneous materials

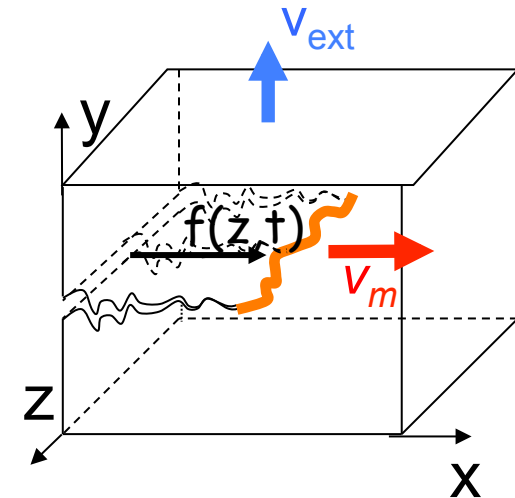
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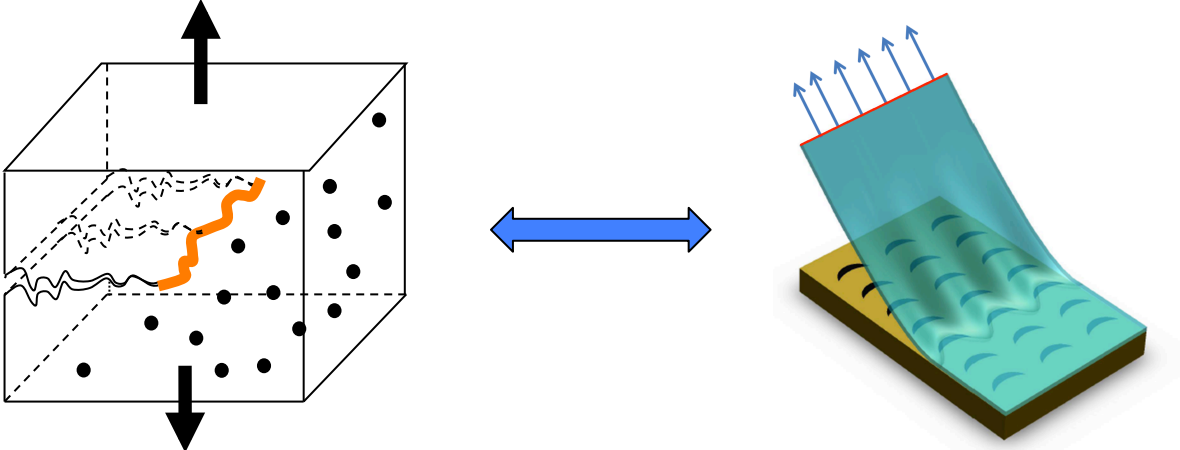
Equation of motion for a crack

$$\mu \frac{\partial f(z,t)}{\partial t} = k[v_m t - f(z,t)] + \frac{G^{ext}}{\pi} \int_{-\infty}^{\infty} \frac{f(z',t) - f(z,t)}{(z' - z)^2} dz' - \delta G_c(z, f(z,t))$$

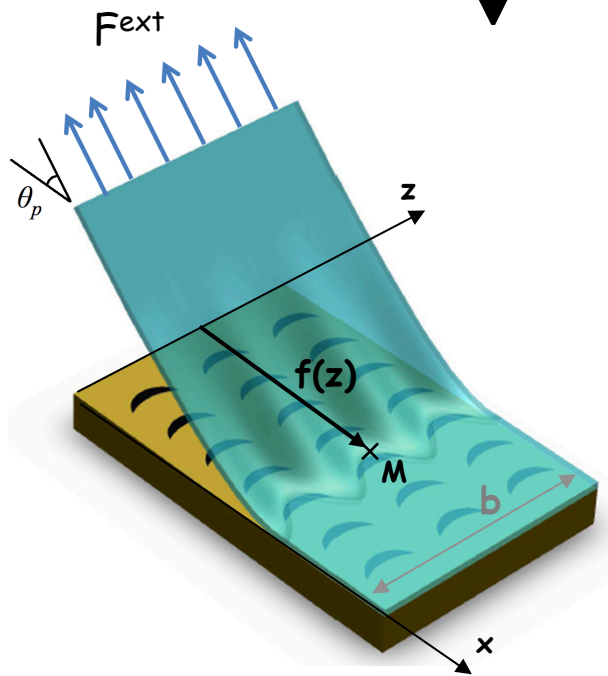
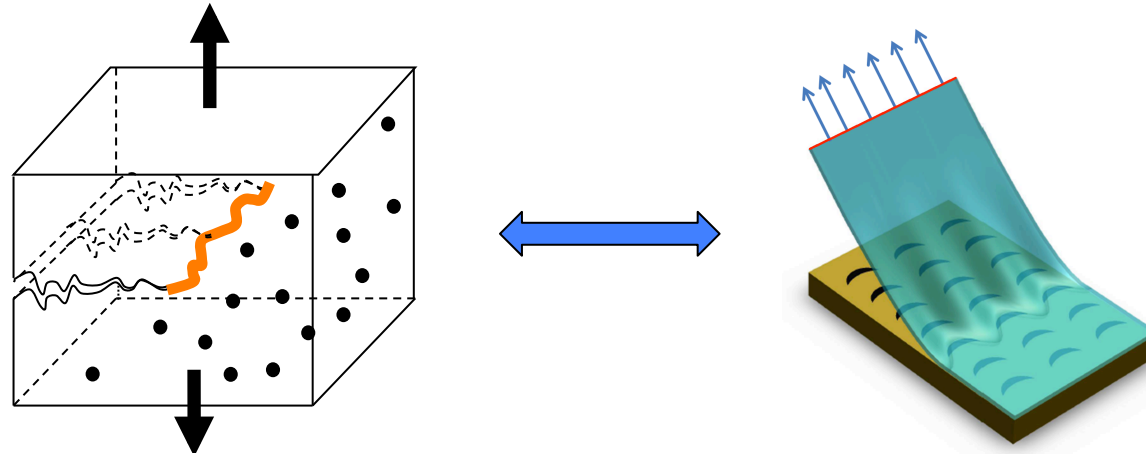
For displacement control experiments: $G^{ext} = \langle G_c \rangle + k[v_m t - f(z,t)]$

Critical driving given by $G_c^{eff} = \langle G_c \rangle + k[v_m t - \langle f(z,t) \rangle_z] \longleftrightarrow G_c^{eff} = \langle \delta G_c(z, f(z)) \rangle_z$

Thin film adhesive as a model system for exploring the failure of heterogeneous materials



Thin film adhesive as a model system for exploring the failure of heterogeneous materials



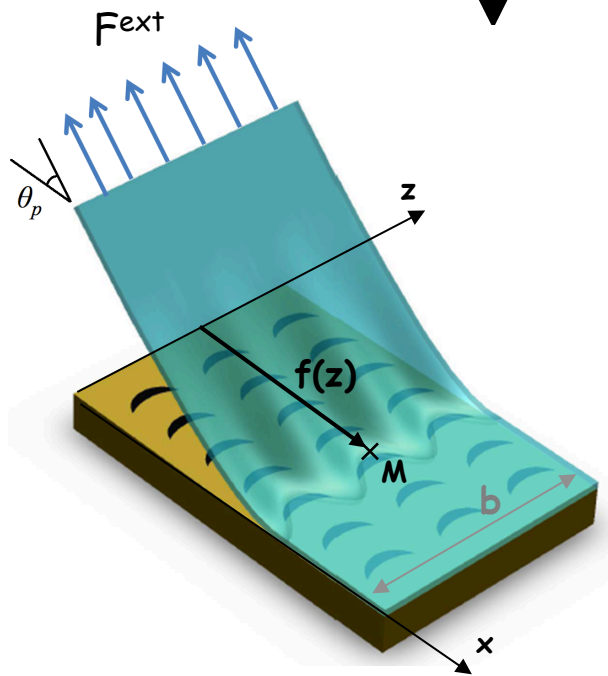
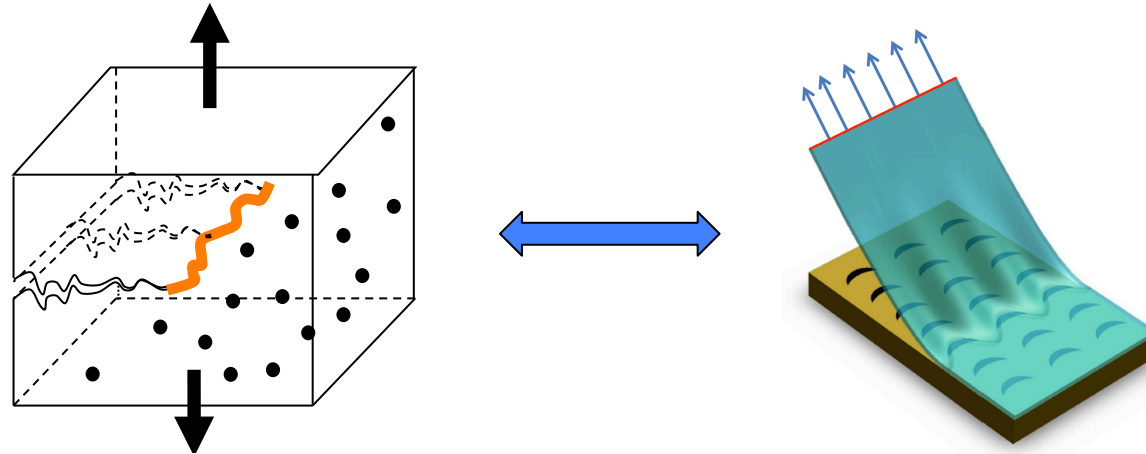
Equation of motion of the peeling front

$$\mu \left. \frac{\partial f(z,t)}{\partial t} \right|_M = G(z) - G_c(z, x = f(z,t))$$

Local driving force: $G(z) = G^{ext} + 4 \frac{G^{ext}}{\pi} \int \frac{f(z') - f(z)}{(z' - z)^2} dz'$

Applied driving force: $G^{ext} = \frac{F_p}{b} (1 - \cos \theta_p)$

Thin film adhesive as a model system for exploring the failure of heterogeneous materials



Equation of motion of the peeling front

$$\mu \left. \frac{\partial f(z,t)}{\partial t} \right|_M = G(z) - G_c(z, x = f(z,t))$$

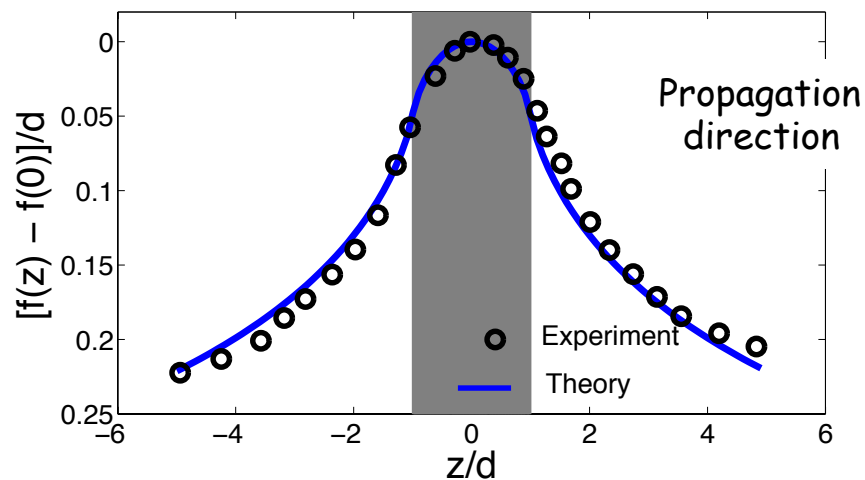
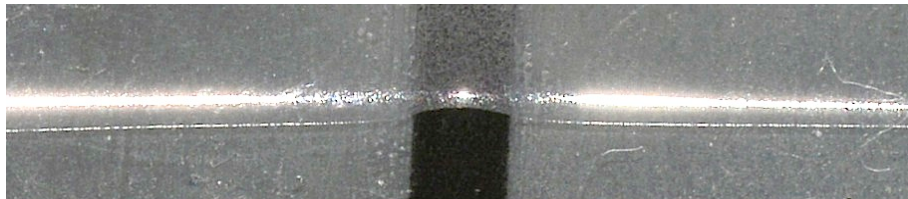
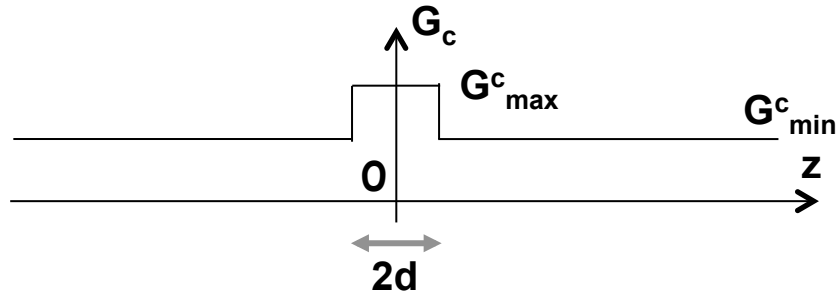
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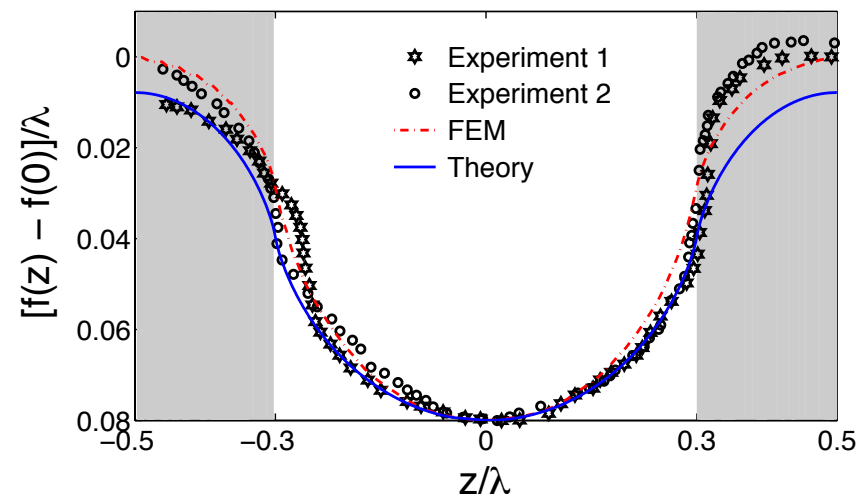
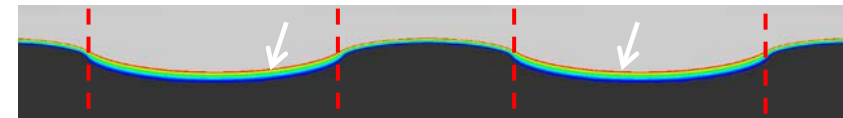
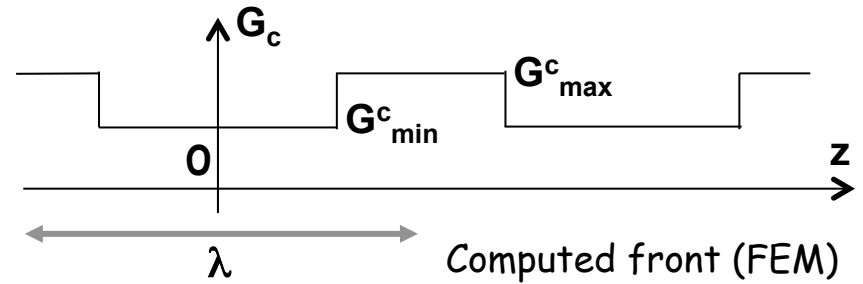
Application: Toughness field invariant in the propagation direction $\longrightarrow G[f(z)] \approx G_c(z)$

Application to crack pinning by...

...a single heterogeneity



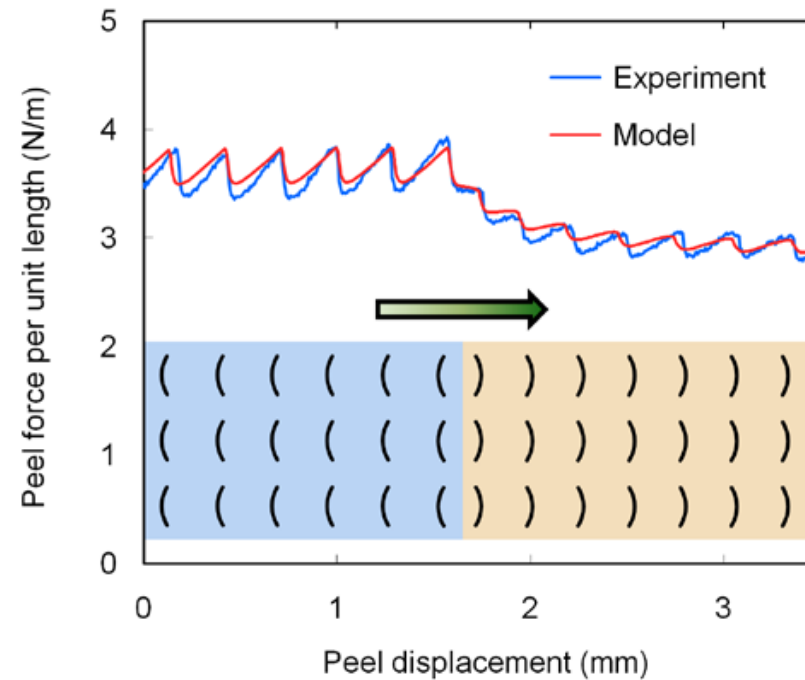
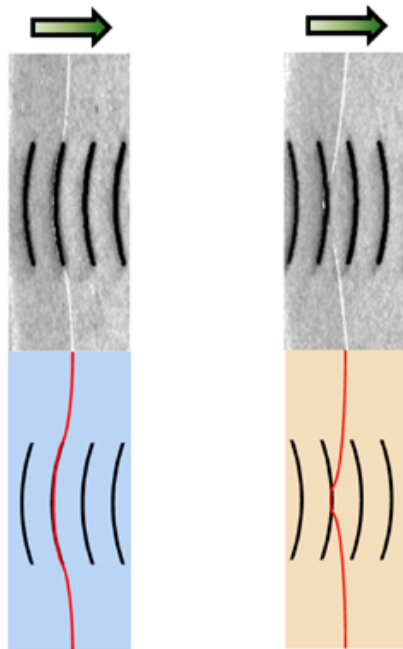
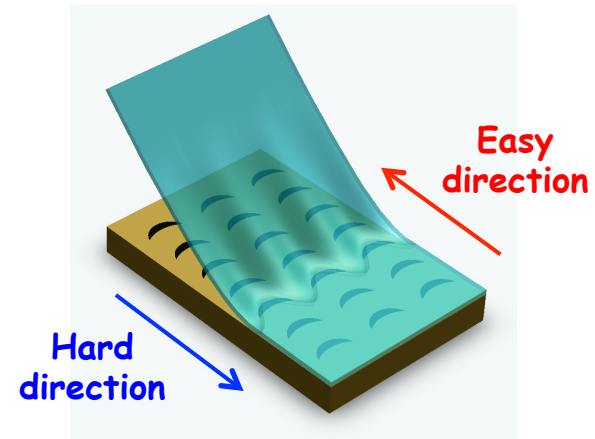
...a periodic array of heterogeneities



Application to material design

*S. Xia, L. Ponson, G. Ravichandran and K. Bhattacharya 2012
and international patent 2011*

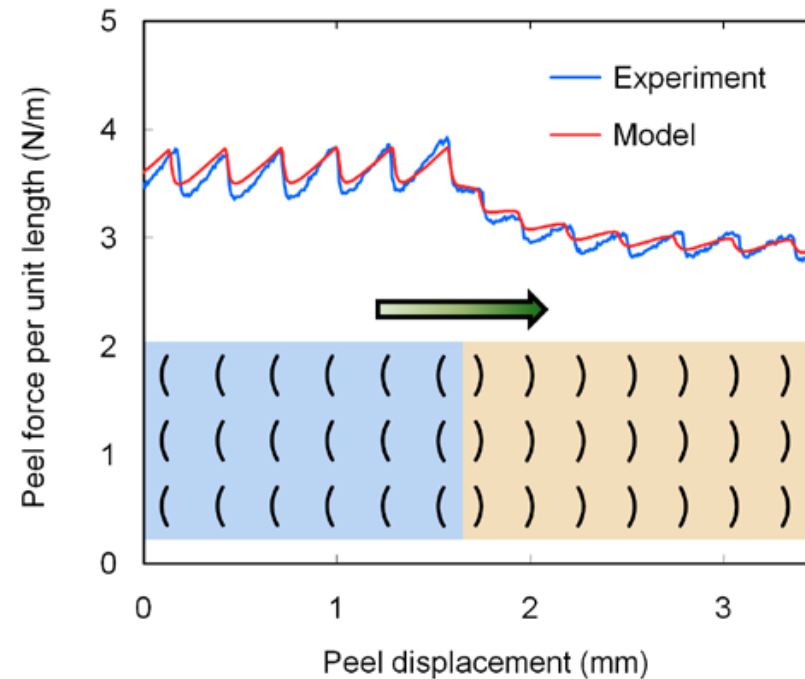
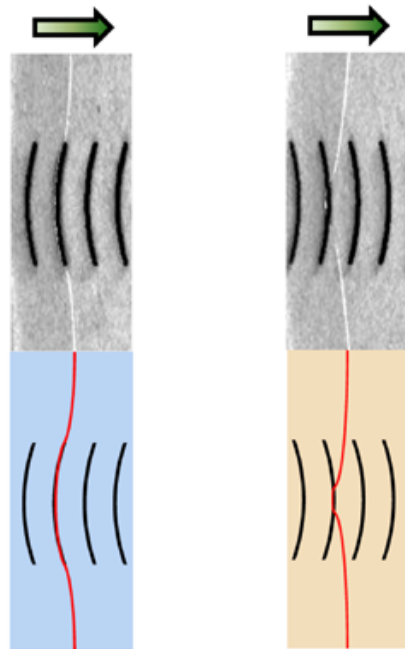
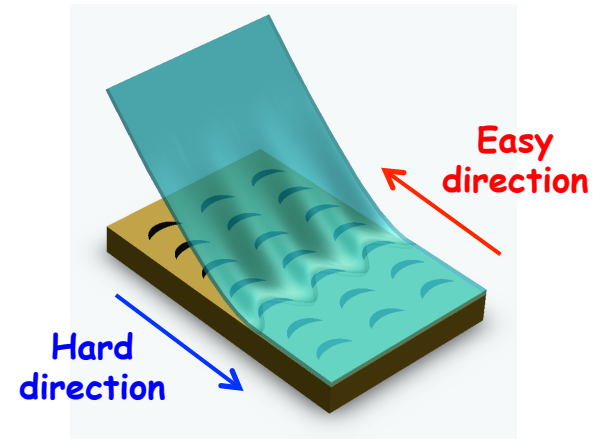
Asymmetric adhesives



Application to material design

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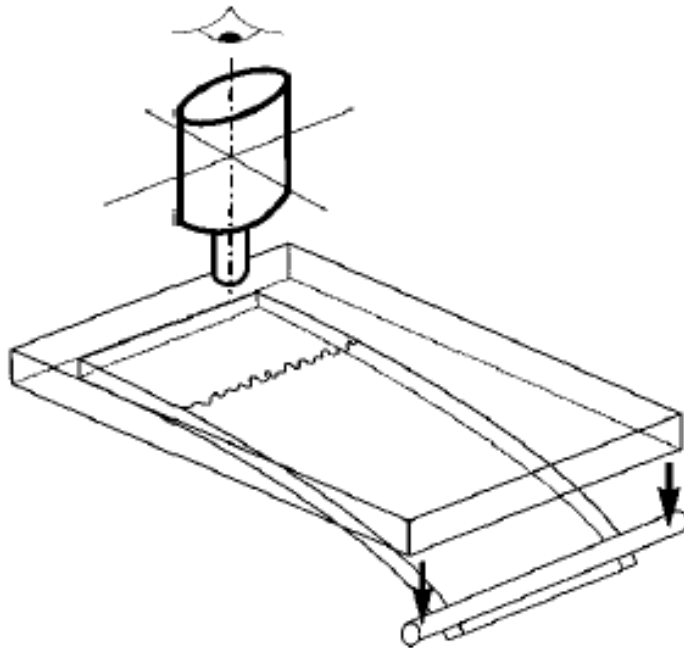
Asymmetric adhesives



➔ **Determination of the optimal shape and distribution of pinning sites to achieve targeted properties**

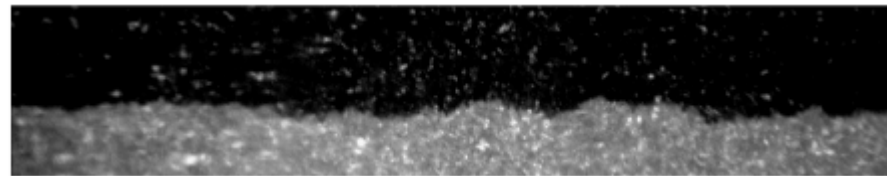
Crack propagation within disordered materials: Bridging the gap between experiments and theory

Experimental setup

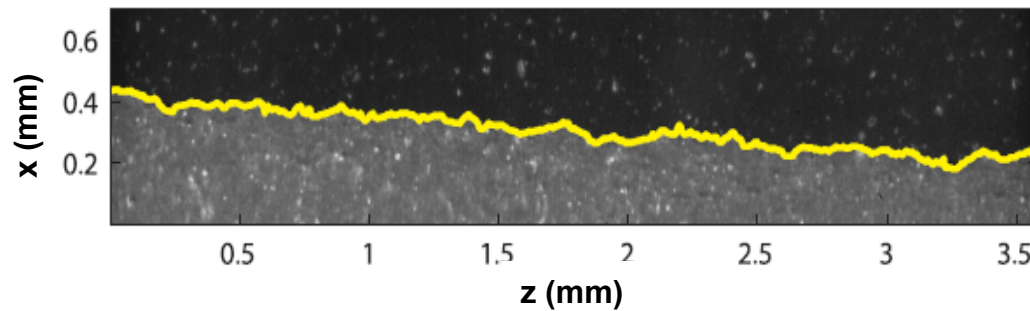


Movie: Courtesy of S. Santucci

Crack front position measured through fast camera:

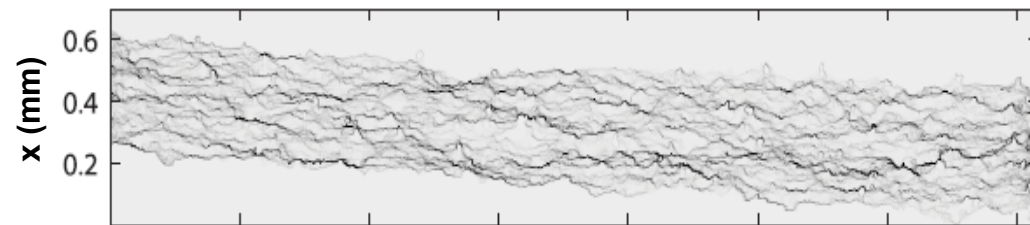


Characterizing the *local* crack dynamics



Position of the front at a given time

0	0	0	0	0	0	0	0	0	0
1	0	0	0	0	1	0	0	0	1
0	1	0	1	1	0	1	0	1	0
0	0	1	0	0	0	0	1	0	0



Matrix of waiting time T_w

1	1	3	4	3	3	3	1	1	1
3	3	3	1	1	1	3	4	1	1
1	4	3	1	2	5	5	5	3	1
1	1	4	3	5	3	1	1	4	4

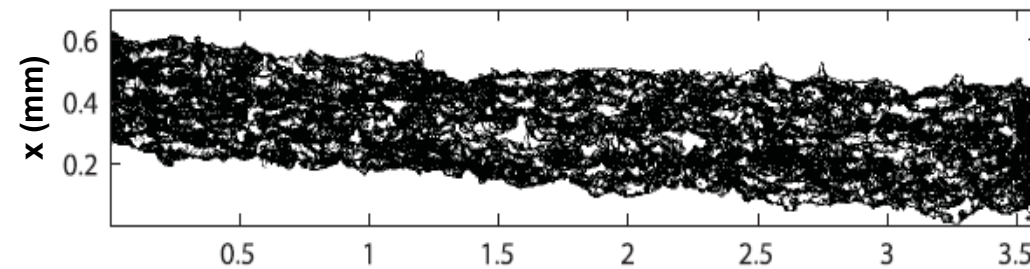
Definition of clusters such as $T_w < C < T_w >$



1	1	3	4	3	3	3	1	1	1
3	3	3	1	1	1	3	4	1	1
1	4	3	1	2	5	5	5	3	1
1	1	4	3	5	3	1	1	4	4

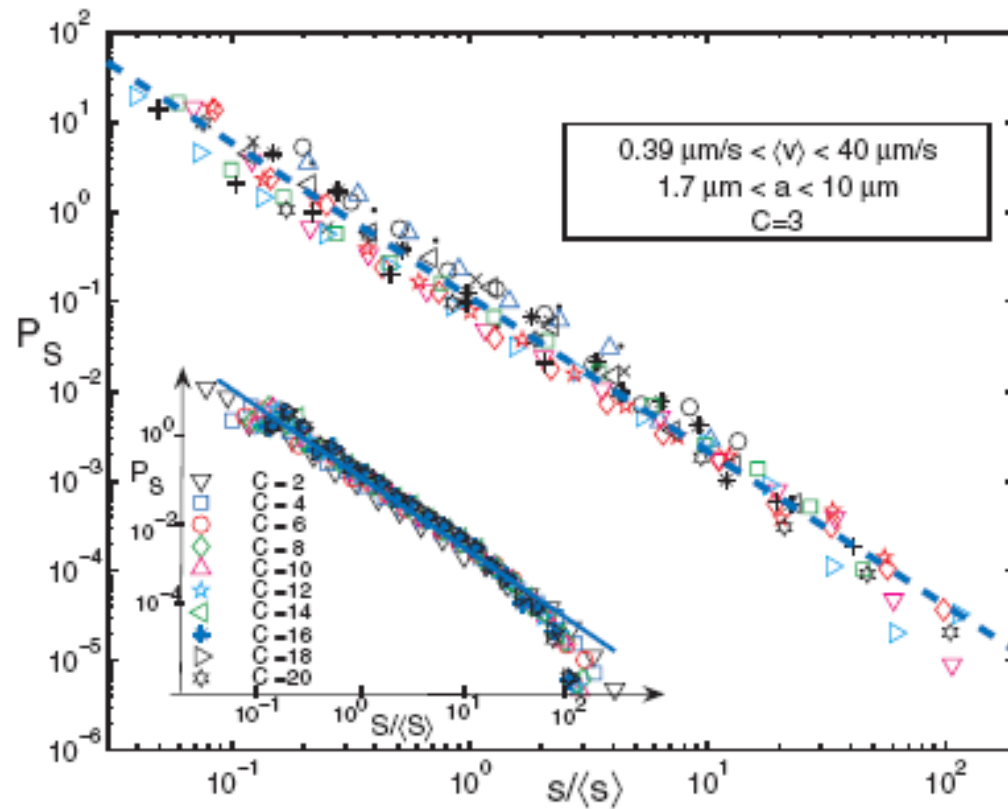


Avalanche of size $S=5$



Distribution of *local* avalanche sizes

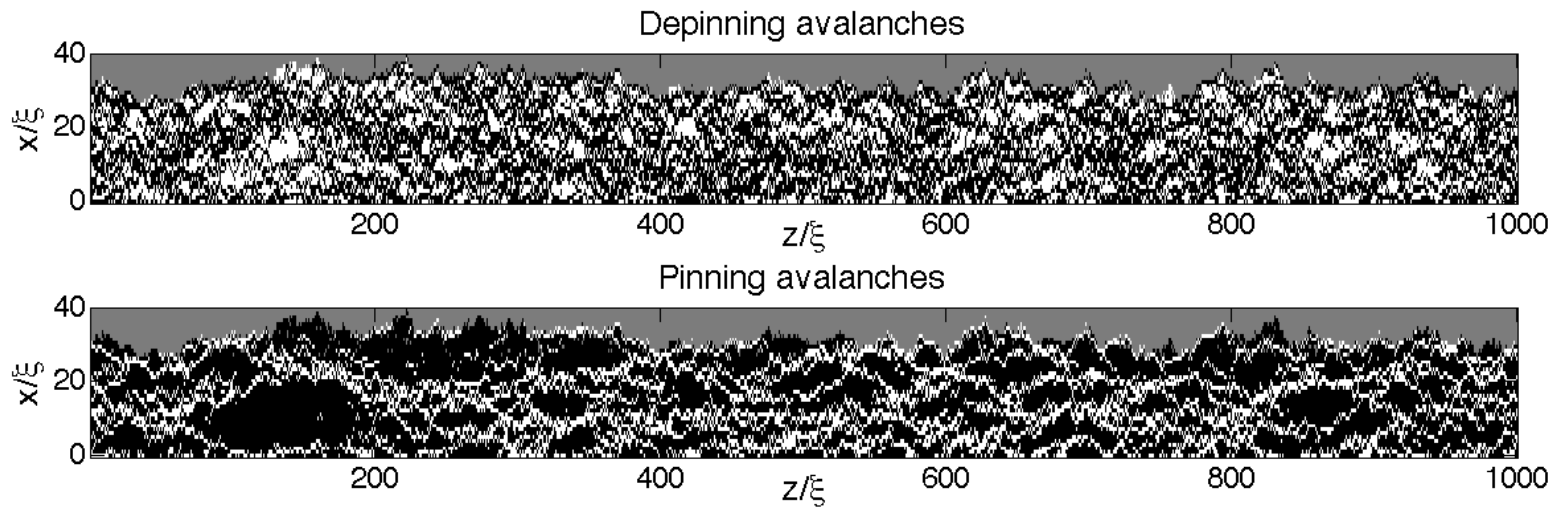
K. Maloy et al. 2006, K. Tallakstat et al. 2011



$$P(S) \sim S^{-\gamma}$$

with $\gamma \approx 1.55$

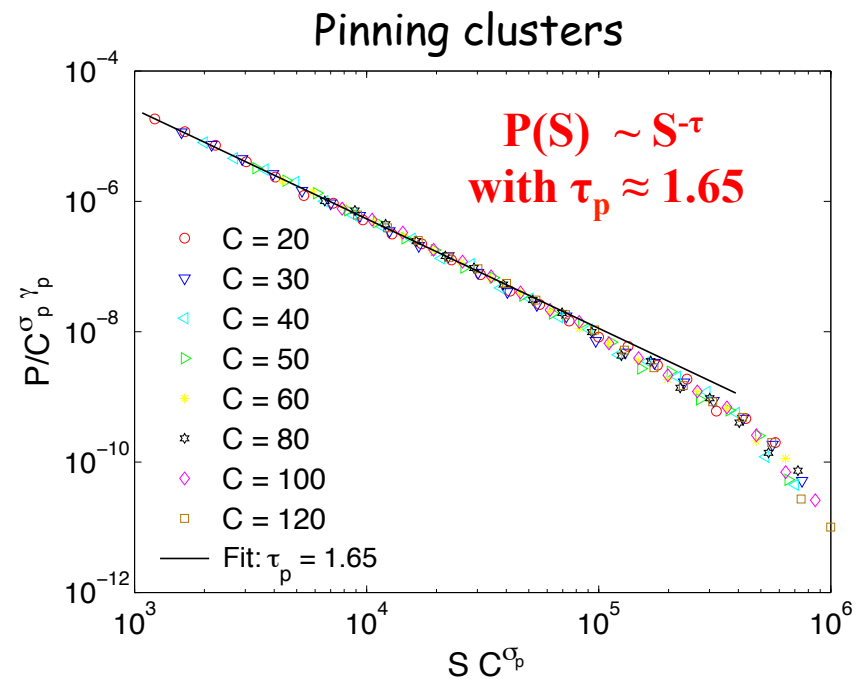
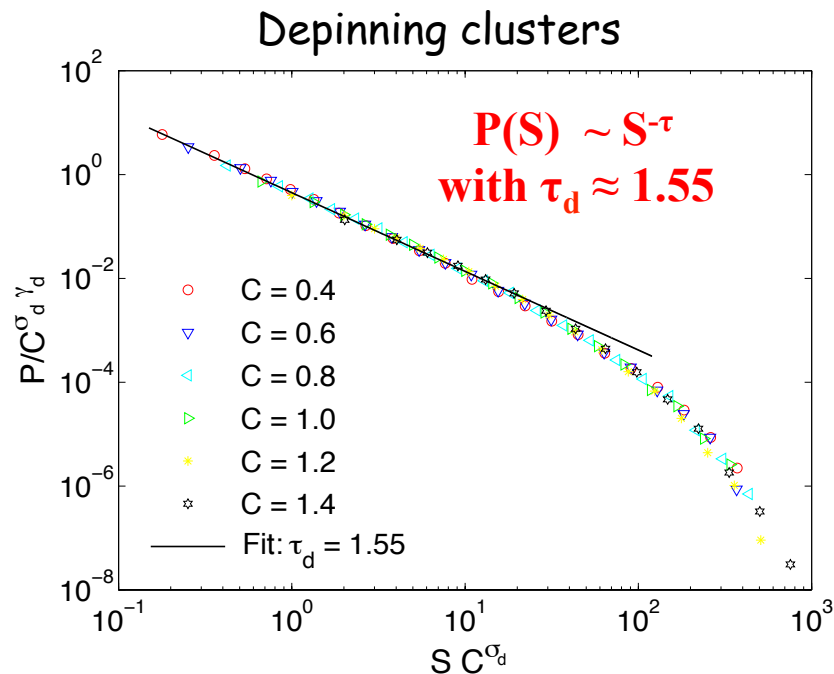
Comparison with the interface depinning model



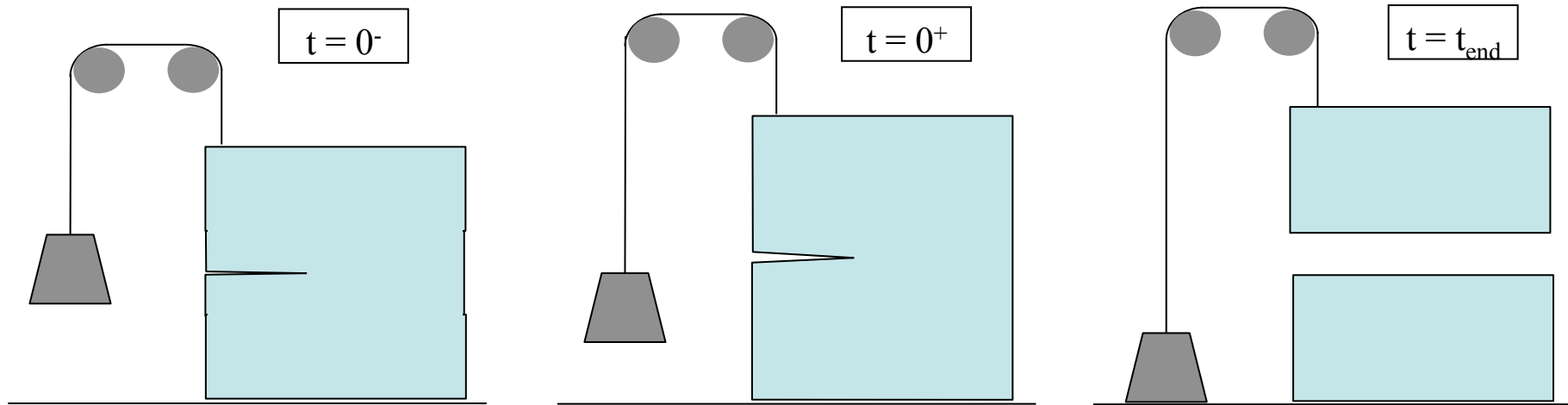
White pixels correspond to

$$T_w < C \langle T_w \rangle$$

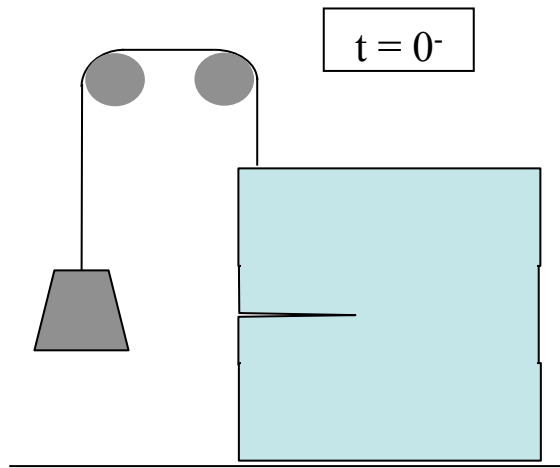
$$T_w > C \langle T_w \rangle$$



Energy transfer during brittle fracture

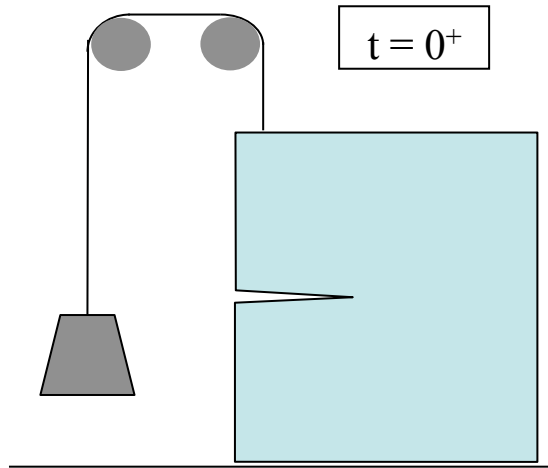


Energy transfer during brittle fracture



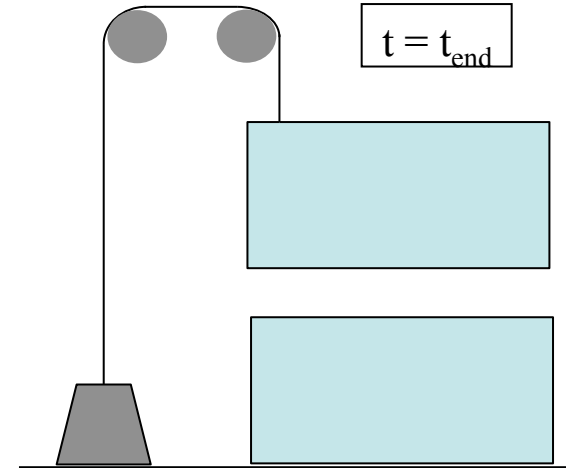
Total energy:

$$E_{tot} = E_0 = E_{pot} + \underbrace{E_{el} + E_s}_{E_m}$$



At time $t=0$:

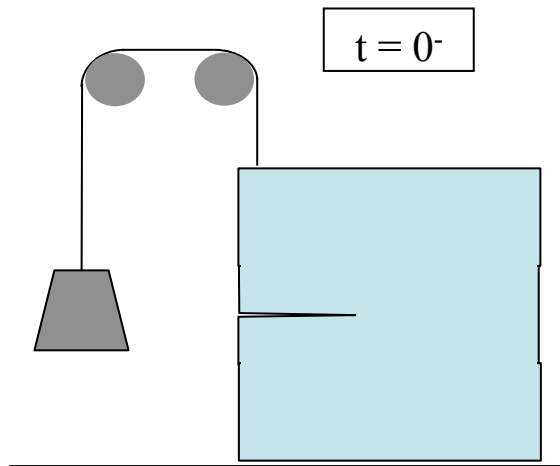
$$\begin{cases} E_m = E_0 \\ E_s = 0 \end{cases}$$



At time $t=t_{end}$:

$$\begin{cases} E_m = 0 \\ E_s = E_0 \end{cases}$$

Energy transfer during brittle fracture



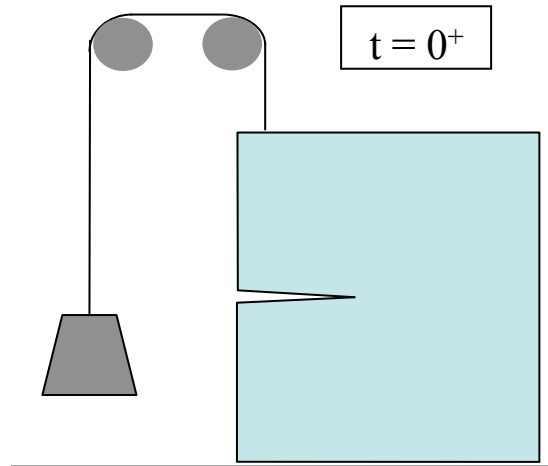
Total energy: $E_{tot} = E_0 = E_{pot} + \underbrace{E_{el} + E_s}_{E_m}$

Rate of energy transfer:

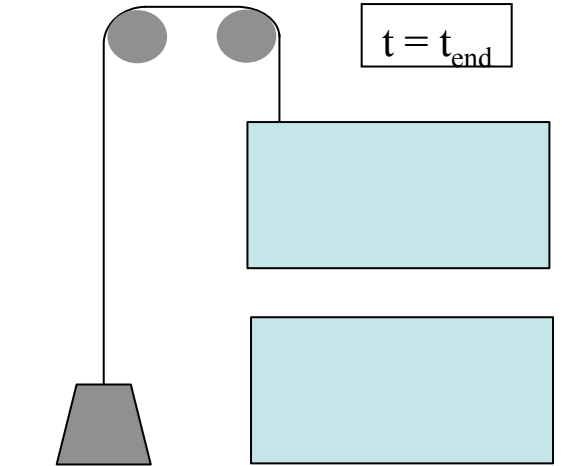
$$P(t) = \frac{\delta E_s}{\delta t} = -\frac{\delta E_m}{\delta t} = b \int G(z,t) \dot{f}(z,t) dz \approx b \langle G_c \rangle v_m$$

→ $P(t) \sim v_m(t)$

Rate of energy transfer
given by the average crack
growth velocity

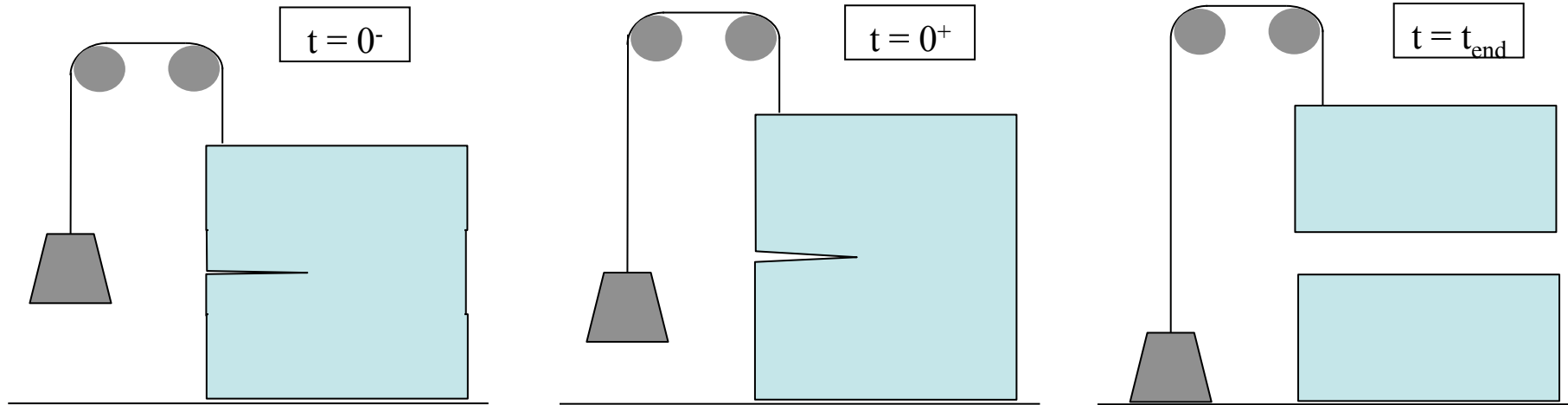


At time $t=0$: $\begin{cases} E_m = E_0 \\ E_s = 0 \end{cases}$



At time $t=t_{end}$: $\begin{cases} E_m = 0 \\ E_s = E_0 \end{cases}$

Energy transfer during brittle fracture



Total energy: $E_{tot} = E_0 = E_{pot} + E_{el} + E_s$

$\underbrace{E_{el} + E_s}_{E_m}$

At time $t=0$: $\begin{cases} E_m = E_0 \\ E_s = 0 \end{cases}$

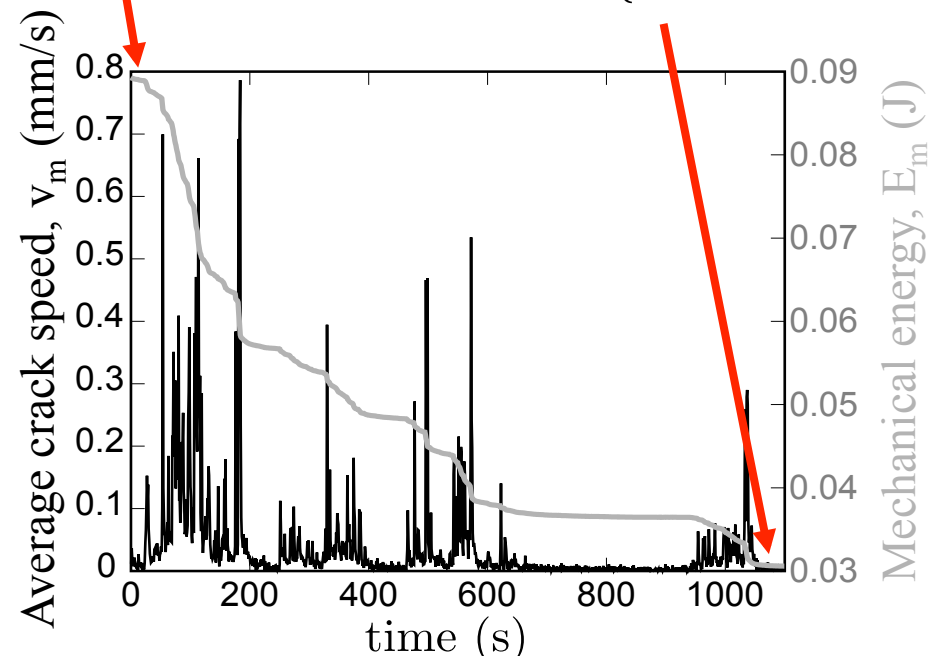
At time $t=t_{end}$: $\begin{cases} E_m = 0 \\ E_s = E_0 \end{cases}$

Rate of energy transfert:

$$P(t) = \frac{\delta E_s}{\delta t} = -\frac{\delta E_m}{\delta t} = b \int G(z,t) \dot{f}(z,t) dz \approx b \langle G_c \rangle v_m$$

$\longrightarrow P(t) \sim v_m(t)$

Rate of energy transfer given by the average crack growth velocity

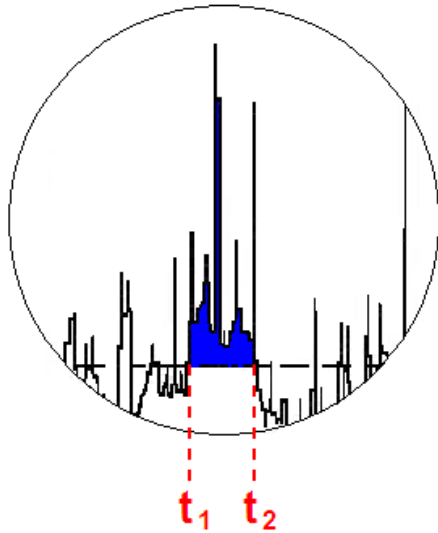


Statistics of *global* avalanches

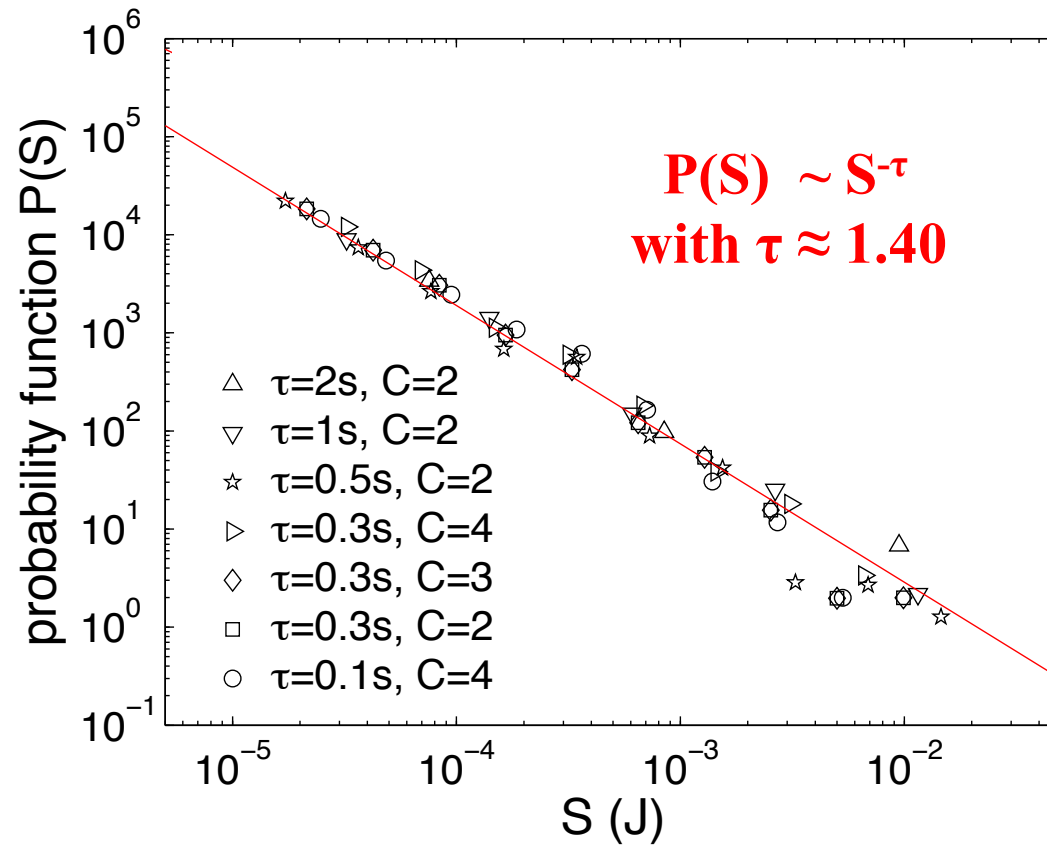
J. Barès, D. Bonamy et al. 2014

Avalanche size:

$$S \sim \int_{t_1}^{t_2} v_m(t) dt$$



Avalanche duration: $T = t_2 - t_1$

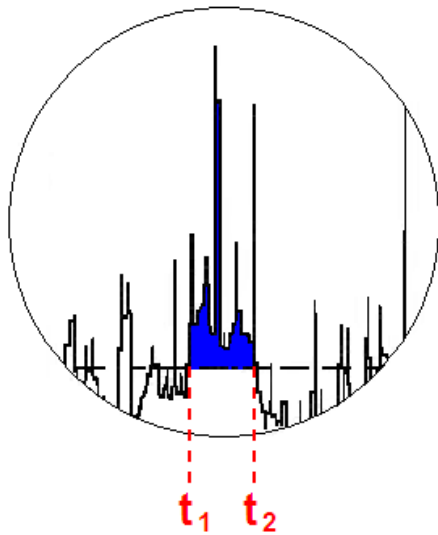


Statistics of *global avalanches*

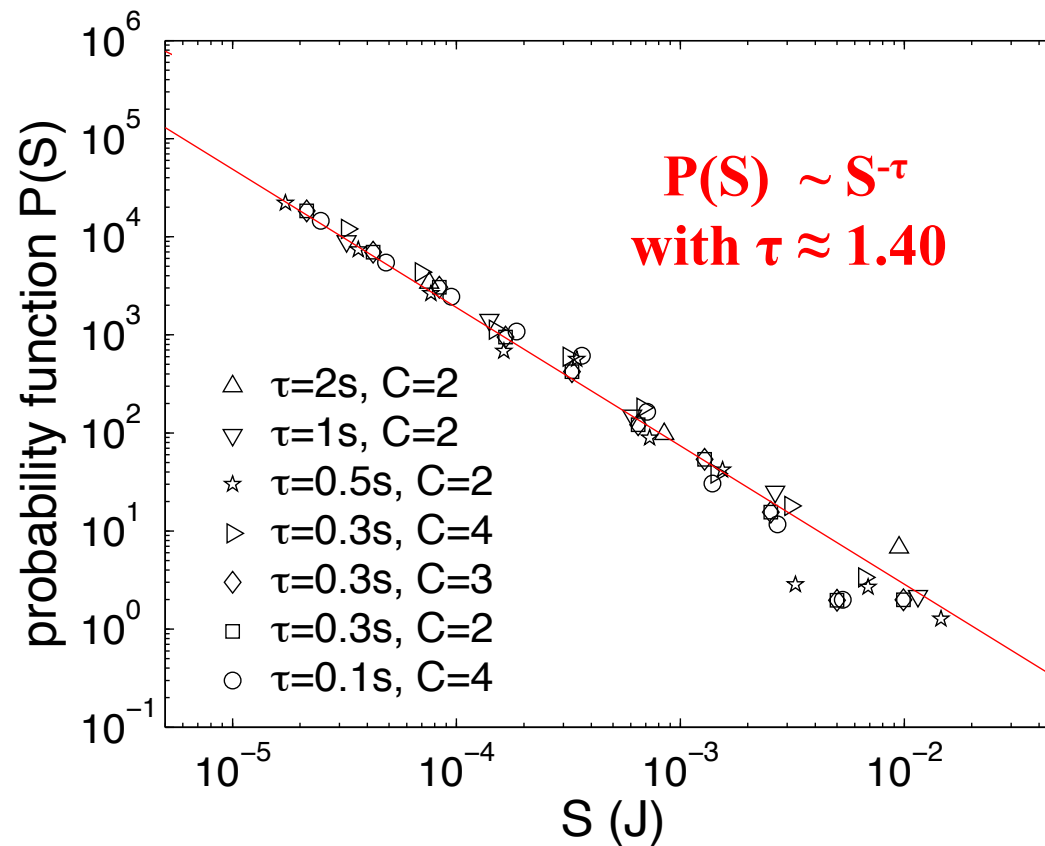
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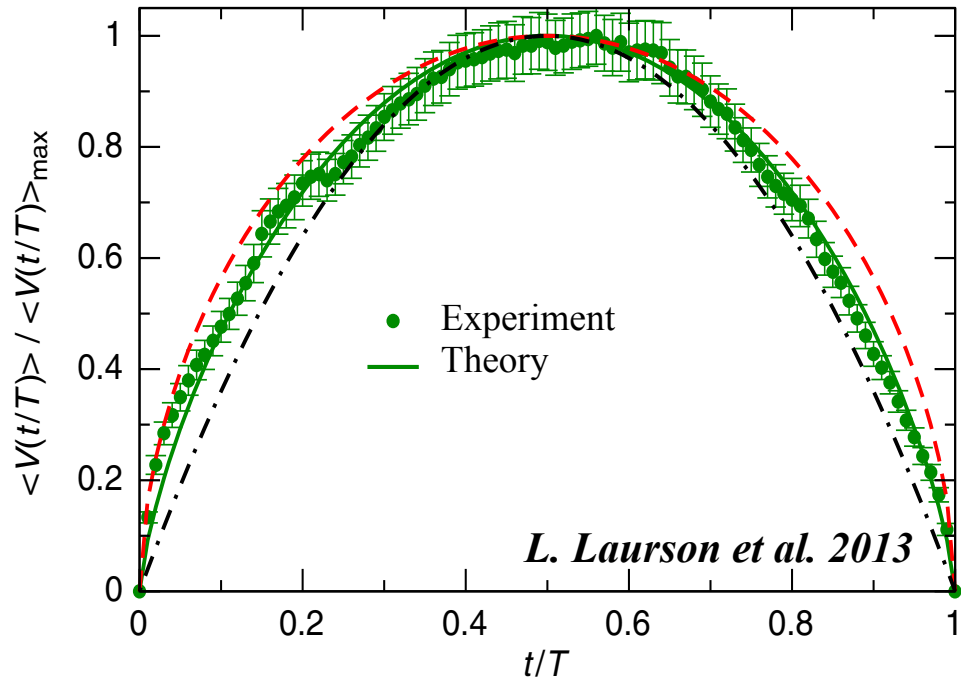
Prediction of the interface depinning model:

**$P(S) \sim S^{-\tau}$
with $\tau \approx 1.28$**

O. Narayan and D. Fisher 1992, A. Dobrinevski, K. Wiese, P. Le Doussal 2014

Avalanche shape defined...

...at the *global scale*



Size vs duration

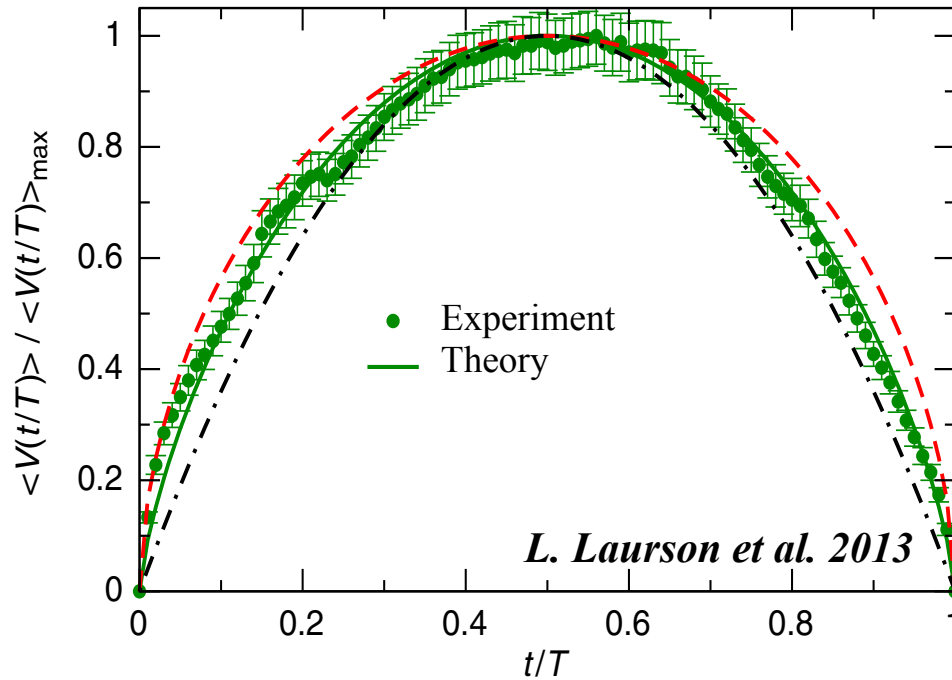
$$S \sim T^\gamma \text{ with } \gamma^{\text{exp}} \approx 1.75$$

→ Asymmetric shape and scaling exponent ($\gamma^{\text{th}} = 1.80$) consistent with the depinning model

See S. Santucci's talk at the conference

Avalanche shape defined...

...at the *global scale*

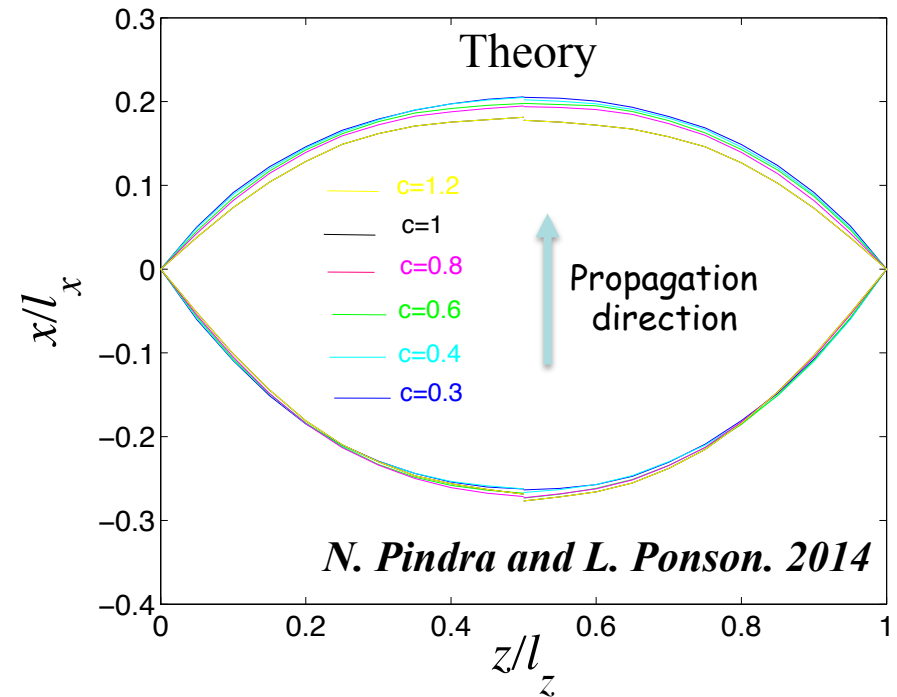


Size vs duration

$$S \sim T^\gamma \text{ with } \gamma^{\text{exp}} \approx 1.75$$

→ Asymmetric shape and scaling exponent ($\gamma^{\text{th}} = 1.80$) consistent with the depinning model

...at the *local scale*



Aspect ratio

$$l_x \sim l_z^\alpha \text{ with } \alpha^{\text{exp}} \approx 0.65$$

→ Exponent ($\gamma^{\text{th}} = 0.65$) consistent with the depinning model. Shape ??

See S. Santucci's talk at the conference

Interface depinning model: A relevant framework to describe the failure of brittle disordered materials?

Yes !

Depinning concepts capture

→ the avalanches dynamics at the *local* and *global* scale

→ the self-affine roughness of crack fronts

S. Santucci et al. 2010

→ the average crack dynamics

Depinning transition: $v_m \sim (G - G_c)^\theta$

L. Ponson 2009

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Depinning transition: $v_m \sim (G - G_c)^\theta$ *L. Ponson 2009*

But under some
specific conditions

→ Very large sample size

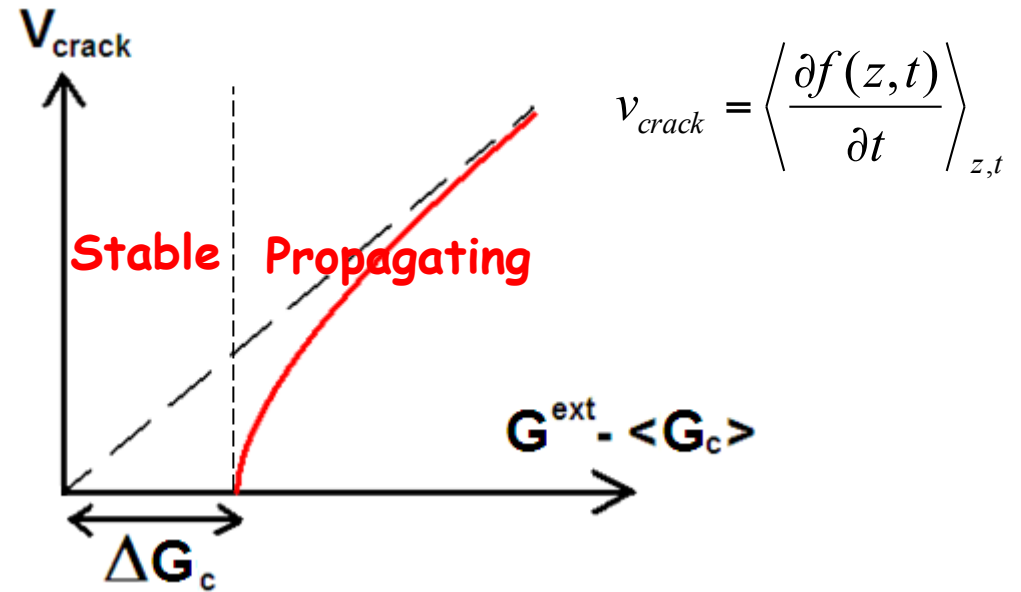
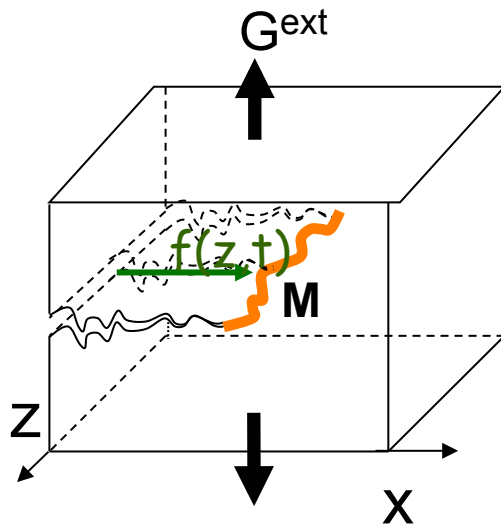
Specimen width \gg spatial extent of avalanches

→ Scale separation between the material heterogeneity and the process zone

Process zone size \ll material heterogeneity size

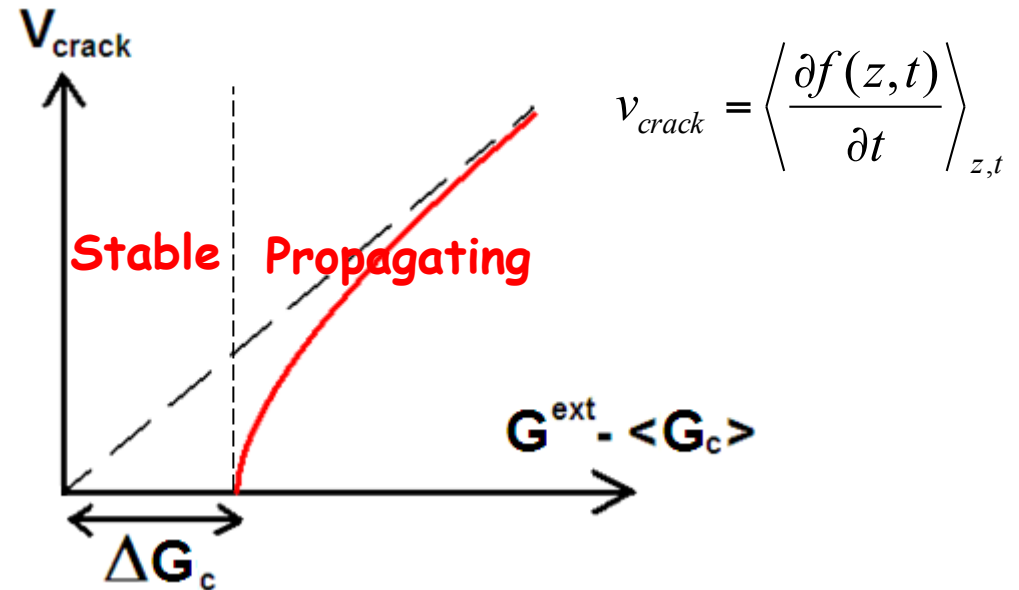
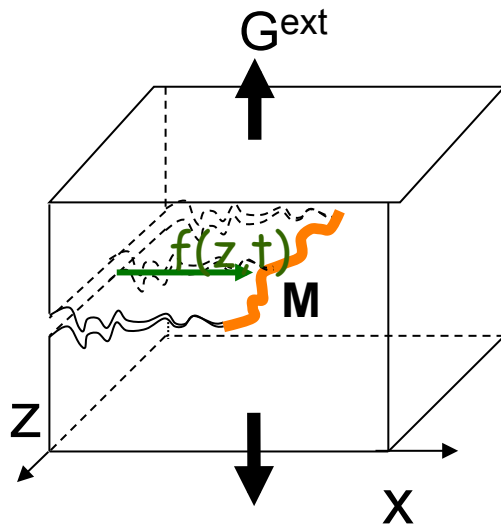
1. Extend the Fracture Mechanics framework to disordered systems

Prediction on the mean crack velocity



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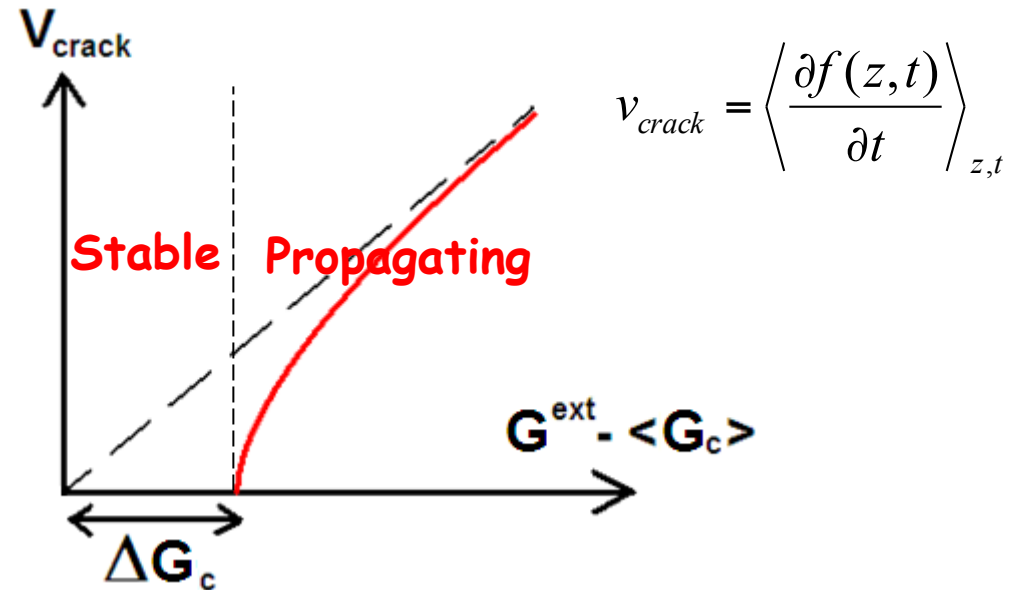
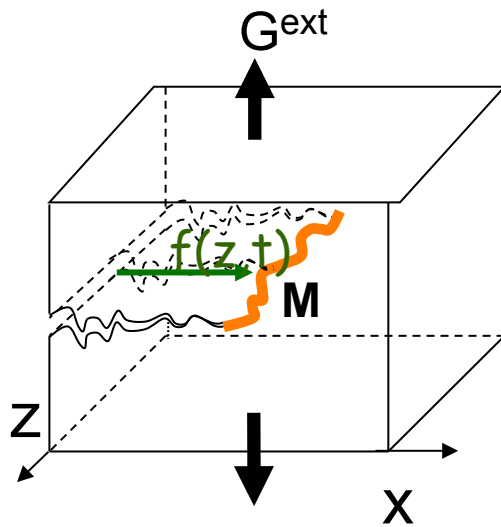


If $G^{ext} < \langle G_c \rangle + \Delta G_c$ \rightarrow The crack front is pinned by the defects

If $G^{ext} > \langle G_c \rangle + \Delta G_c$ \rightarrow The crack propagates

1. Extend the Fracture Mechanics framework to disordered systems

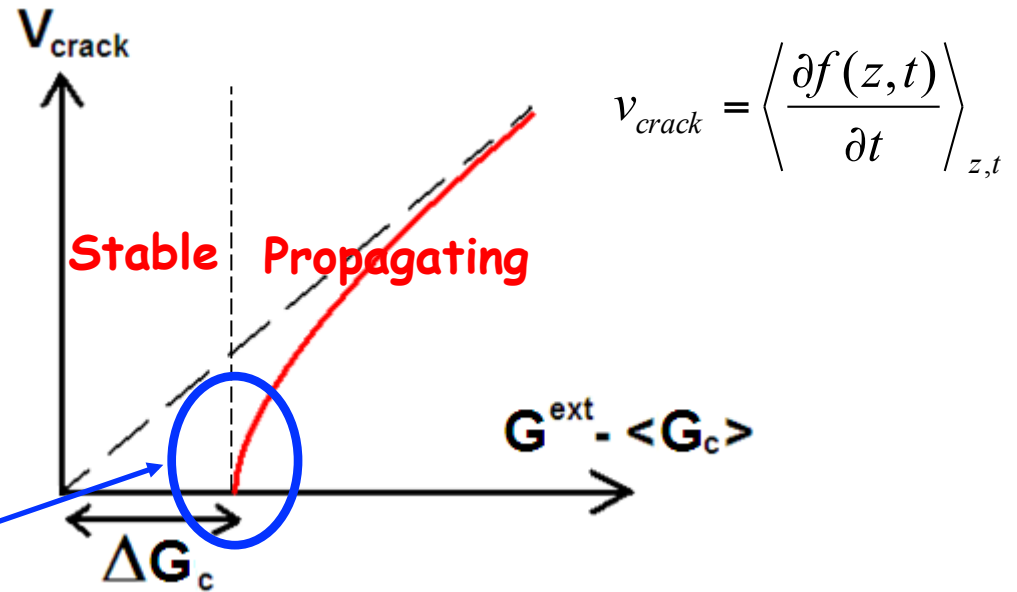
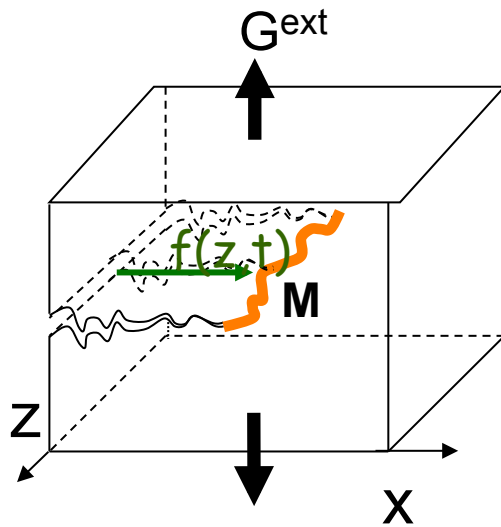
Prediction on the mean crack velocity



- If $G^{ext} < \langle G_c \rangle + \Delta G_c \rightarrow$ The crack front is pinned by the defects
- If $G^{ext} > \langle G_c \rangle + \Delta G_c \rightarrow$ The crack propagates
- } $G_c^{eff} = \langle G_c \rangle + \Delta G_c$
- Toughening effect**

1. Extend the Fracture Mechanics framework to disordered systems

Prediction on the mean crack velocity



$$v_{crack} \sim (G^{ext} - G_c^{eff})^\theta$$

with $\theta = 0.75$

D. Ertas and M Kardar 1994 (RG, 1st loop)
P. Chauve et al. 2001 (RG, 2nd loop)
O. Duemmer and W. Krauth 2007 (Sim.)

If $G^{ext} < \langle G_c \rangle + \Delta G_c$ → The crack front is pinned by the defects

If $G^{ext} > \langle G_c \rangle + \Delta G_c$ → The crack propagates

} $G_c^{eff} = \langle G_c \rangle + \Delta G_c$

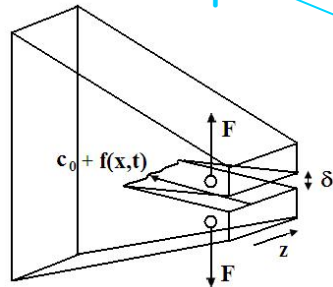
Toughening effect

2. Average crack dynamics: variations of the mean crack velocity with the driving force

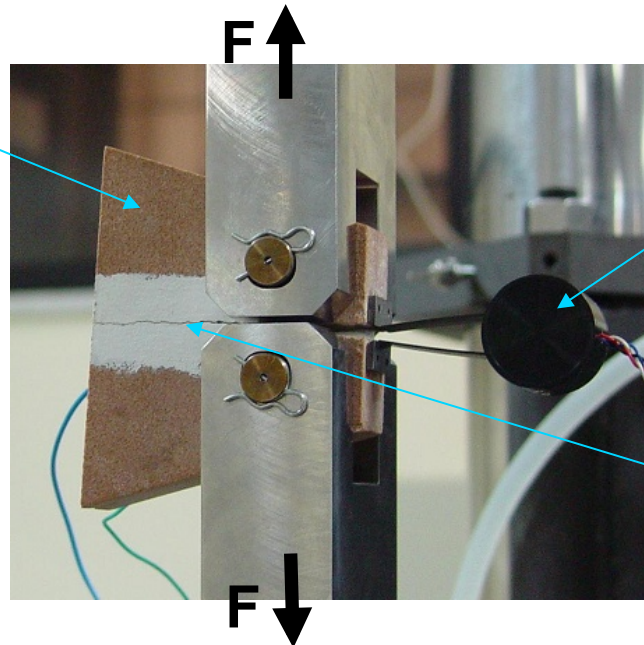
Can we measure the depinning transition?

Experimental setup to measure the curves $v(G^{\text{ext}})$

TDCB Sample



- Slightly unstable geometry that enables to investigate a wide range of velocities
- Sandstone specimen
Brittle + heterogeneous materials



Clip gauge

Crack opening displacement

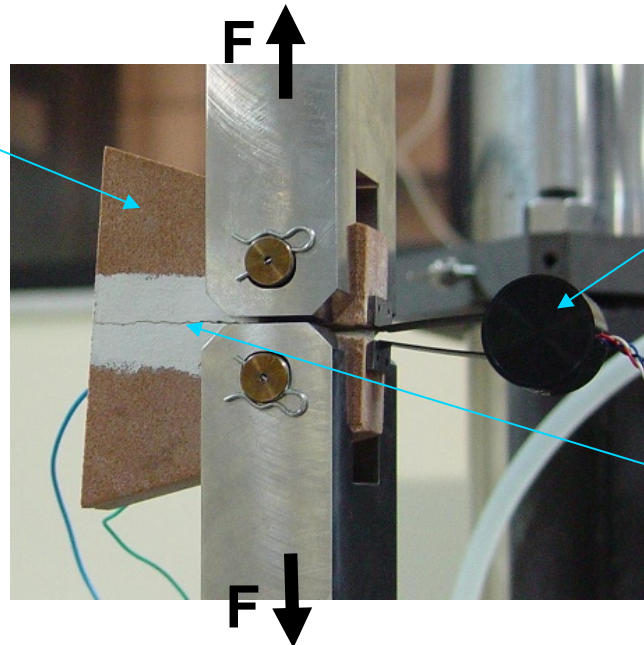
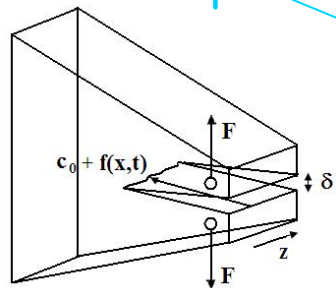
Crack

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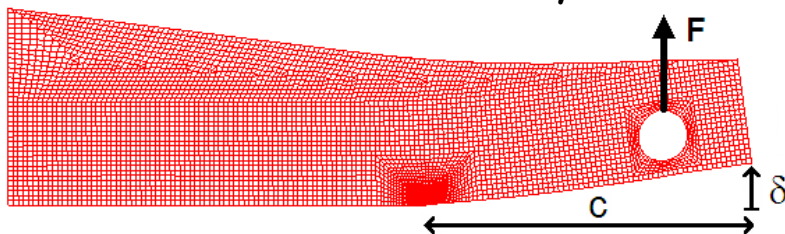
Clip gauge

Crack opening displacement

Crack

- Slightly unstable geometry that enables to investigate a wide range of velocities
- Sandstone specimen
Brittle + heterogeneous materials

Finite element analysis

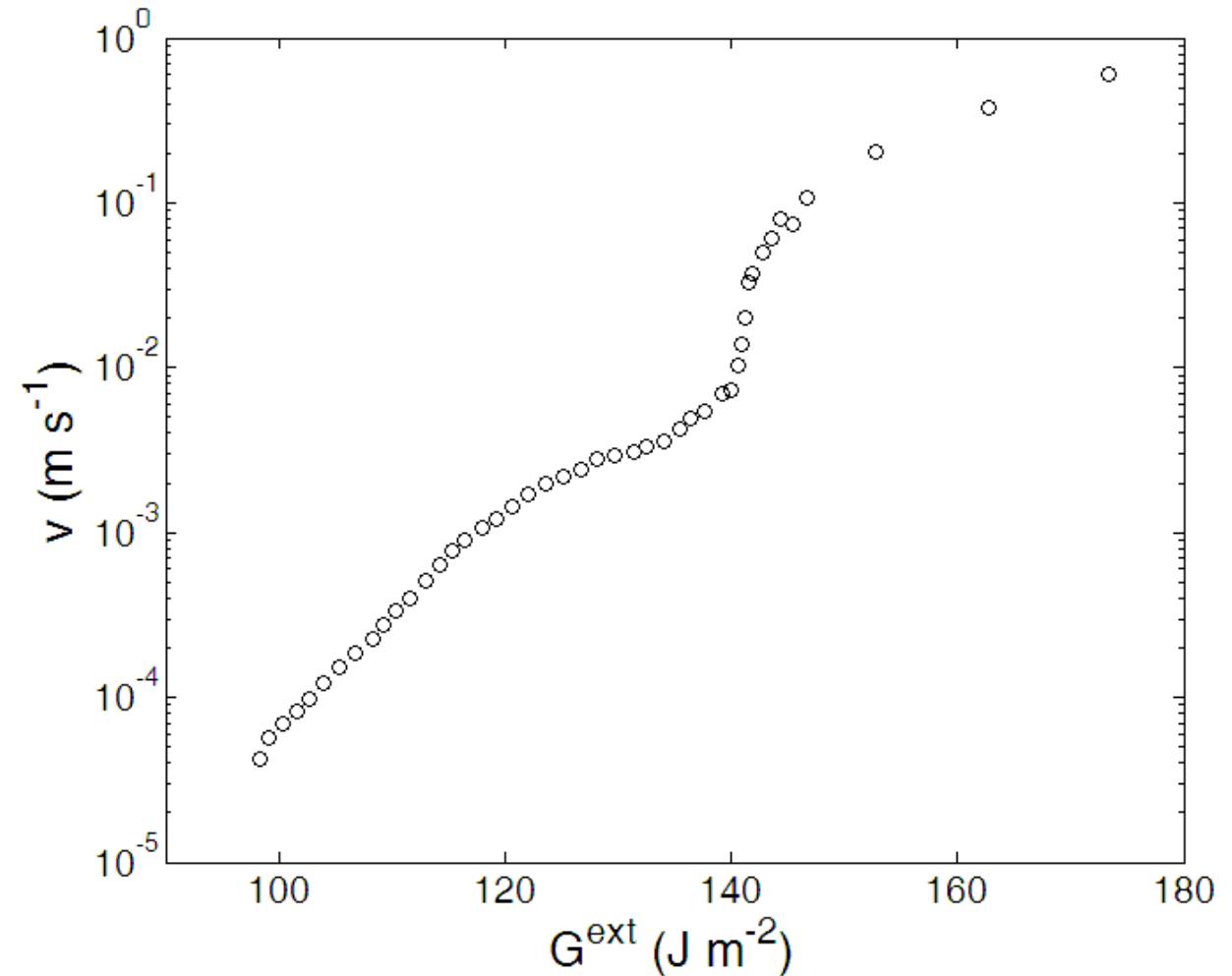
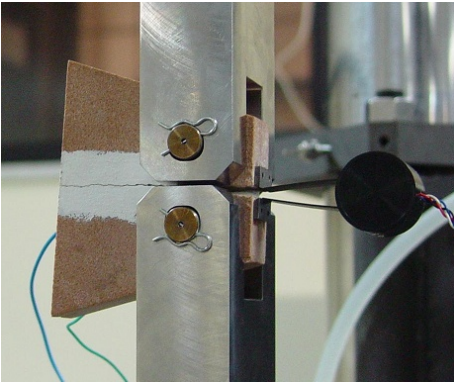


+ Experimental values of
-force
-crack opening displacement
→ $v(G^{\text{ext}})$

2. Average crack dynamics: variations of the mean crack velocity with the driving force

Can we measure the depinning transition?

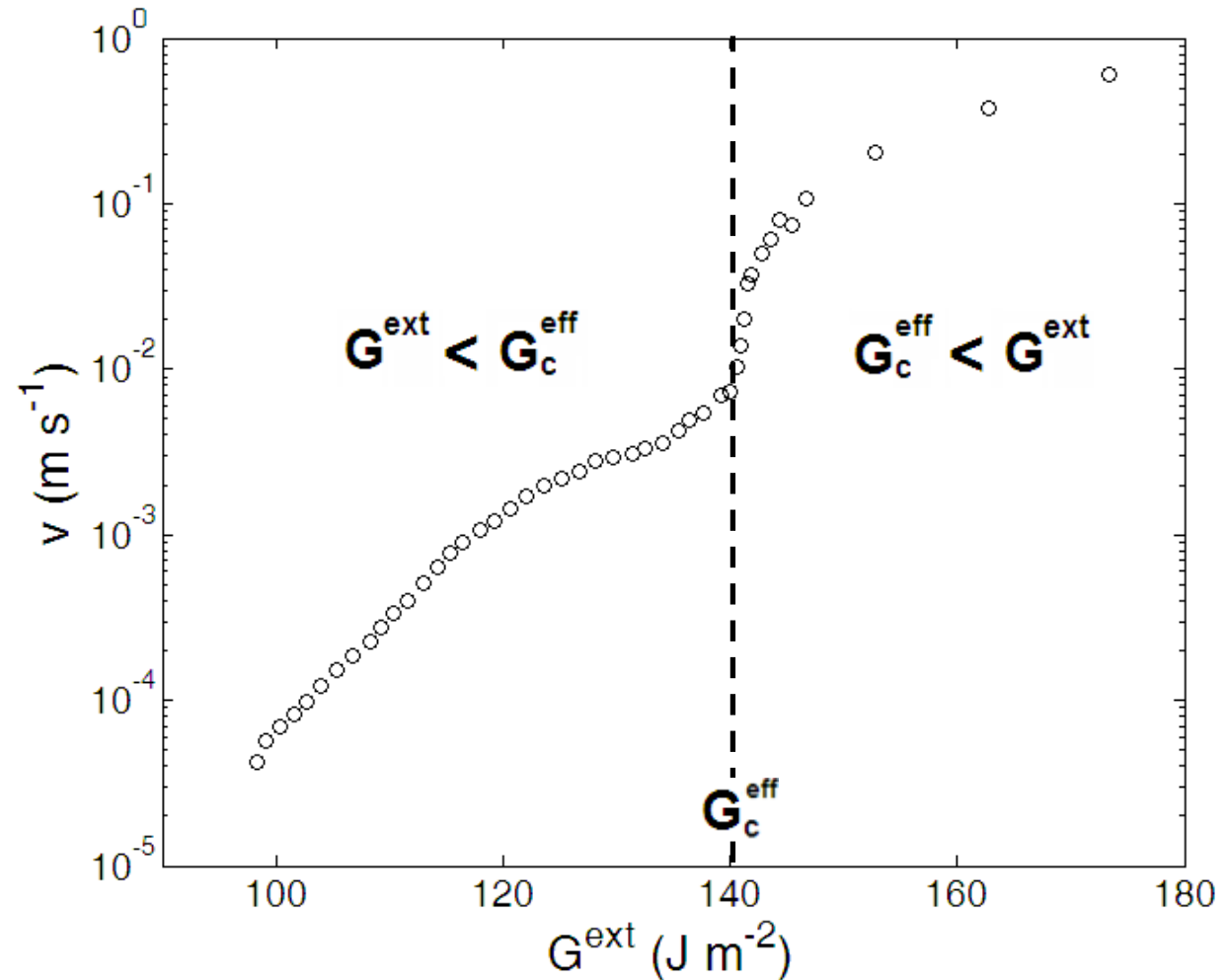
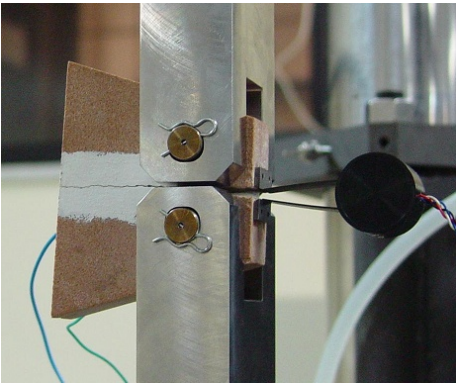
Fracture tests of brittle rocks



2. Average crack dynamics: variations of the mean crack velocity with the driving force

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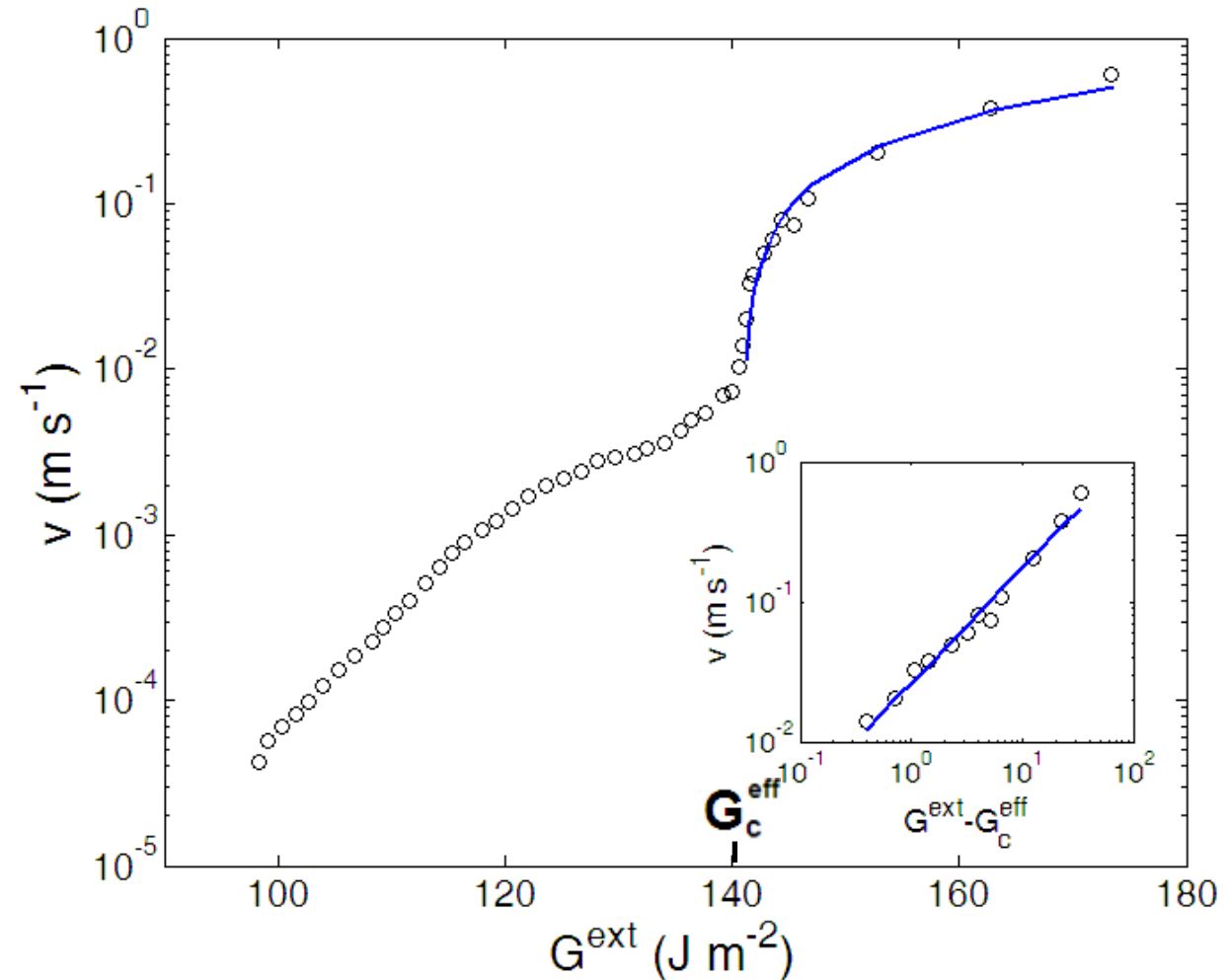
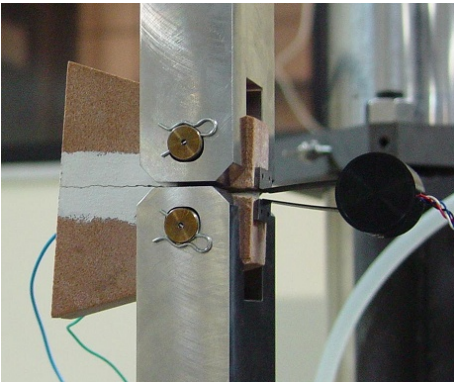
Fracture tests of brittle rocks



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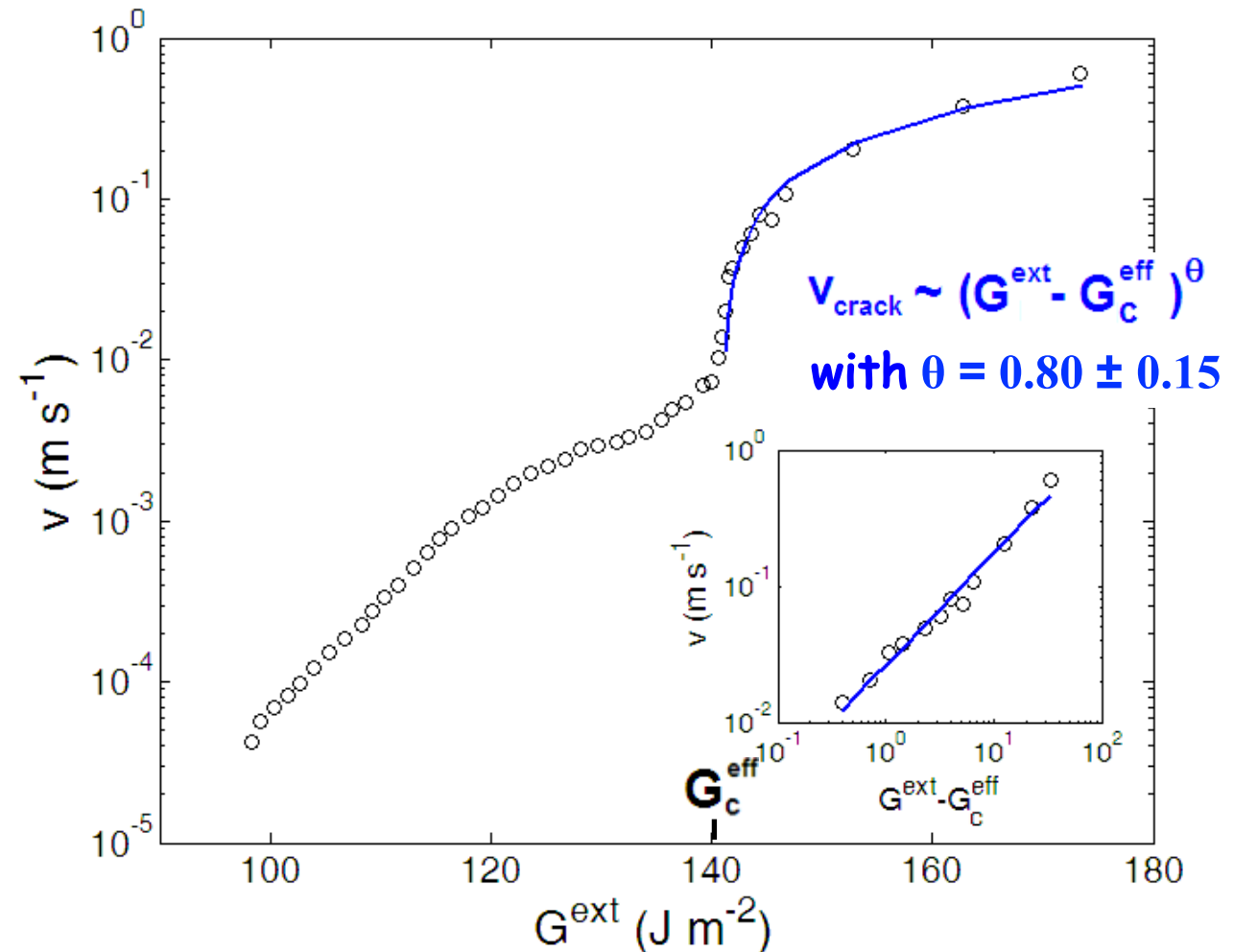
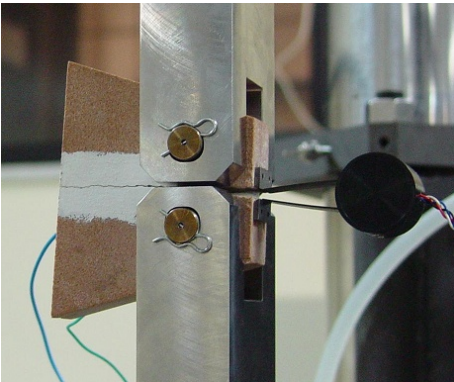
Fracture tests of brittle rocks



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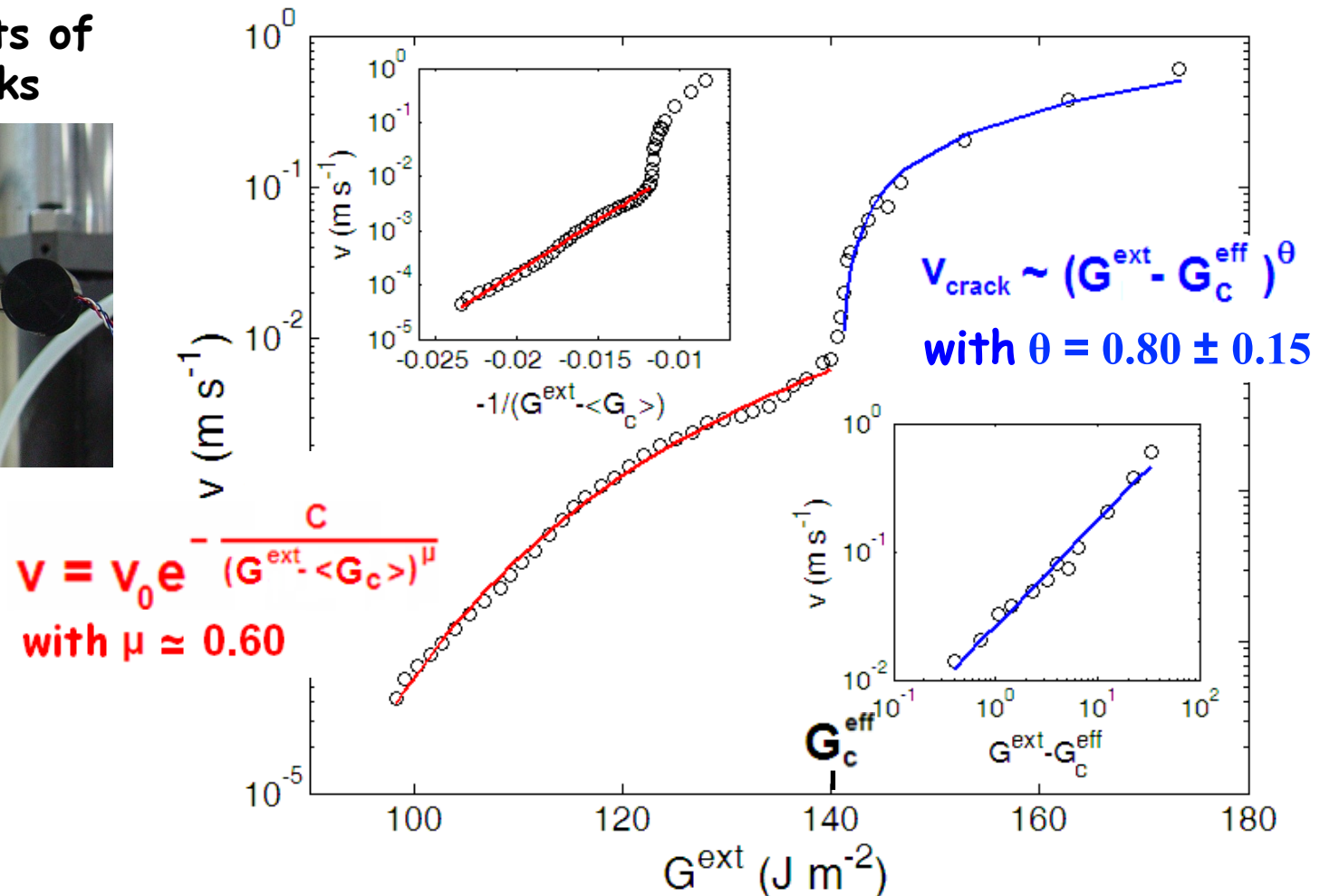
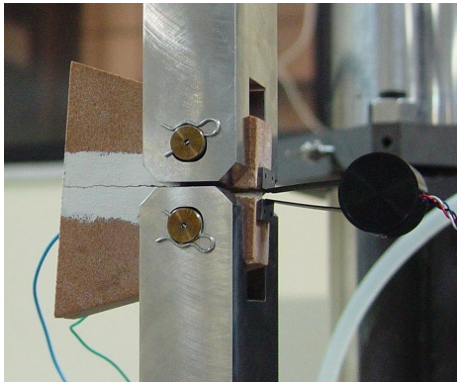
Fracture tests of brittle rocks



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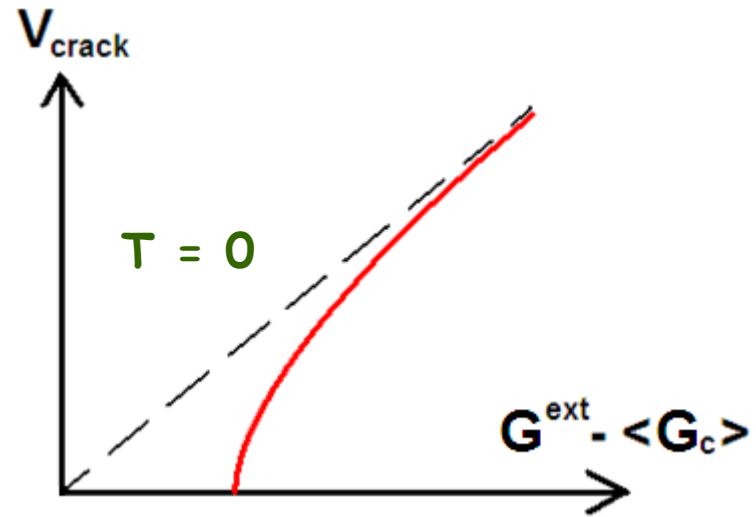
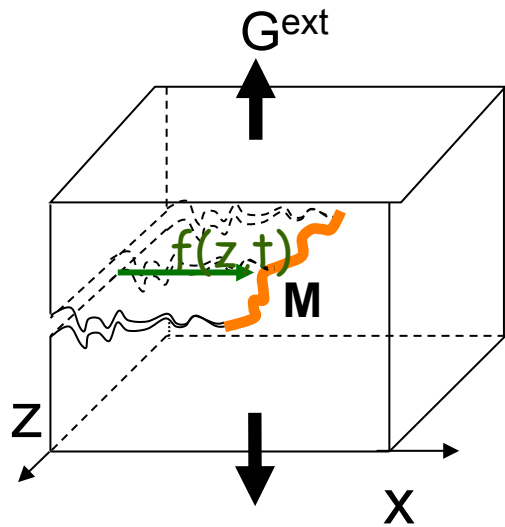
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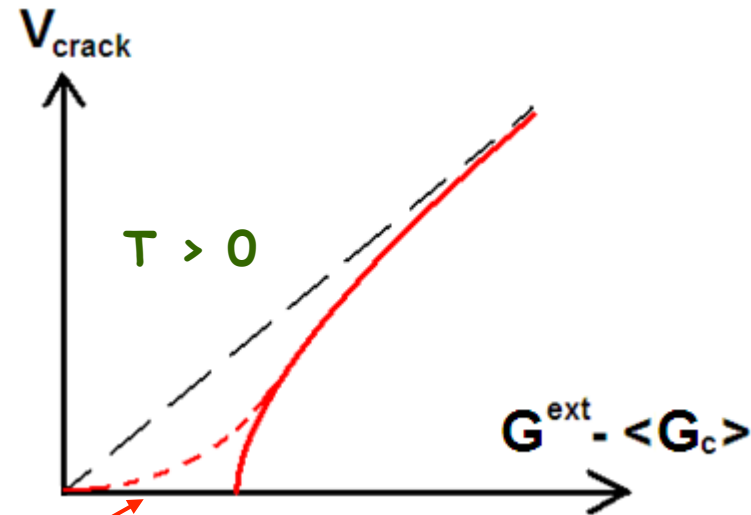
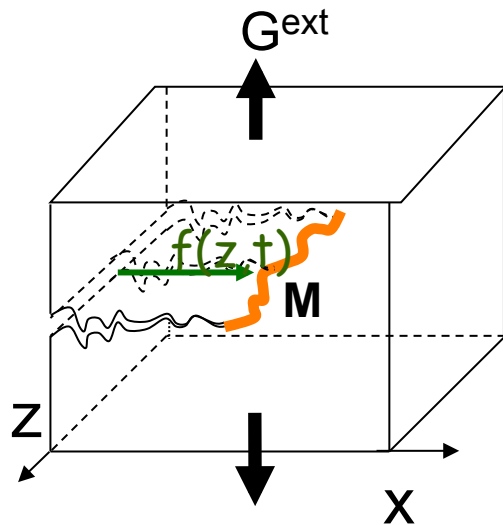
2. Average crack dynamics: variations of the mean crack velocity with the driving force

Coming back to the model..



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Coming back to the model..



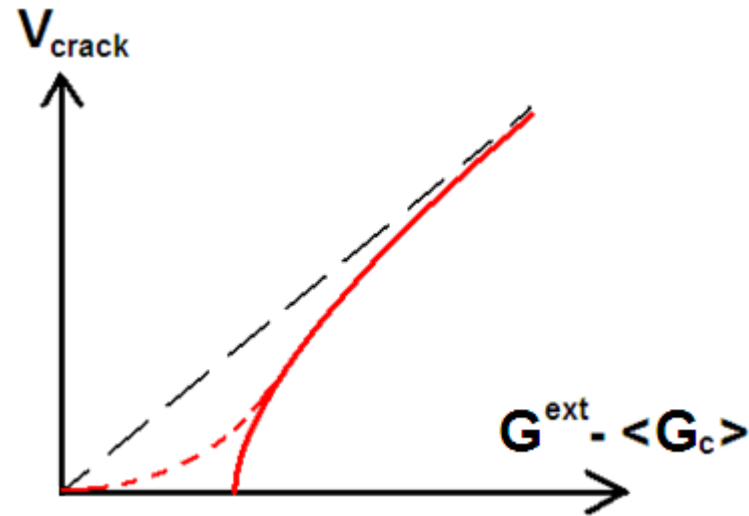
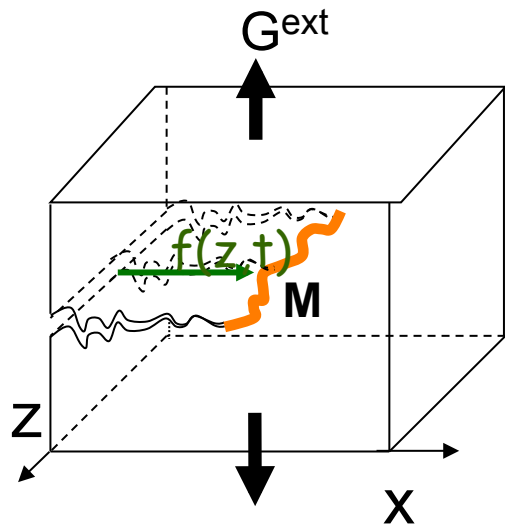
$$v = v_0 e^{-\frac{c}{(G^{ext} - \langle G_c \rangle)^\mu}}$$

with $\mu \approx 0.60$ (long range elasticity)

➔ Crack propagation is possible below the critical threshold through thermal activation processes

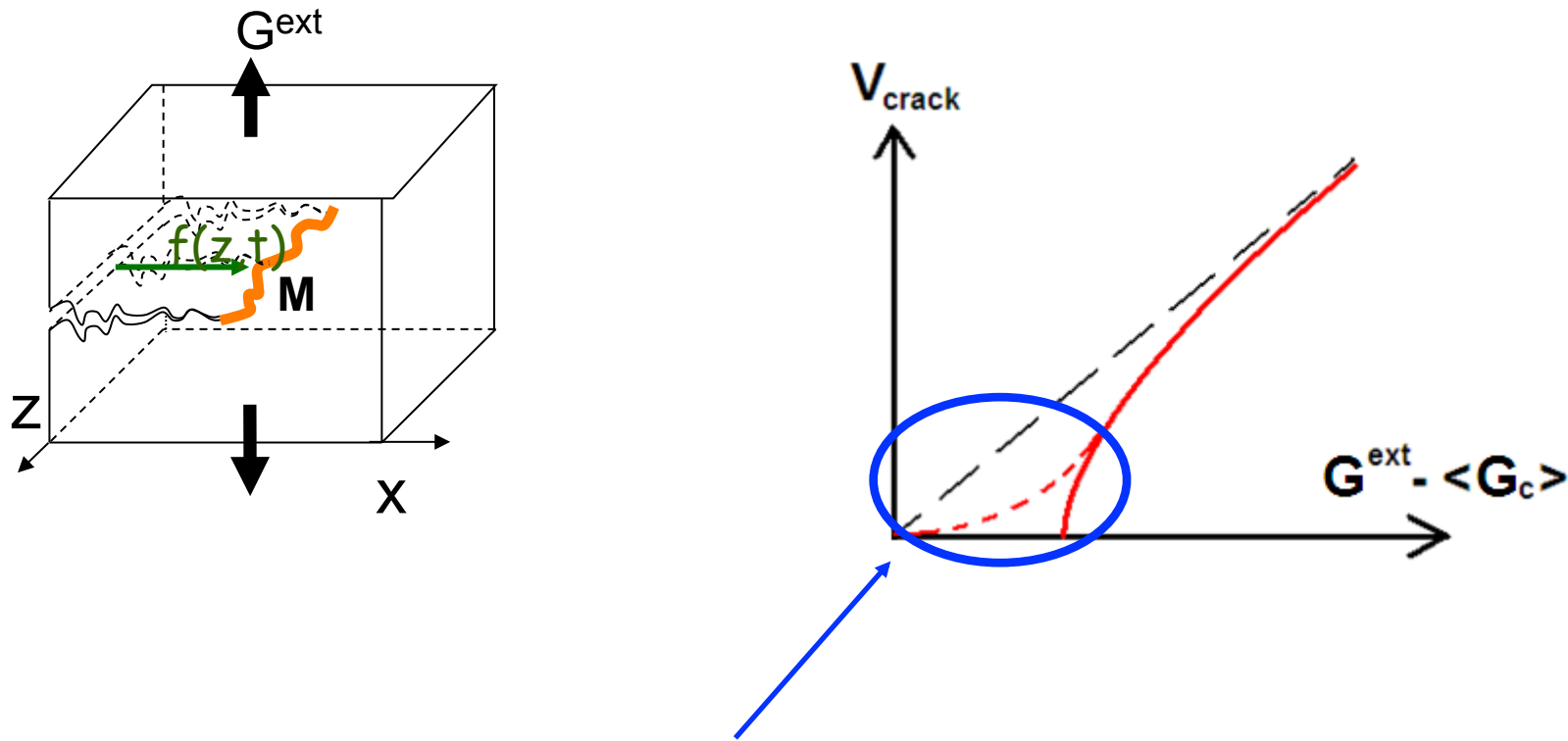
3. Intermittency and avalanches during crack propagation

Intermittent crack propagation at the transition and below



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Intermittent crack propagation at the transition and below



At the transition and below, the crack propagates through sudden jumps (avalanches) with no characteristic size (power law distributed)