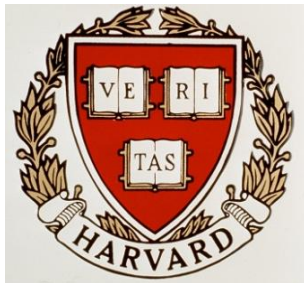


The power of supersymmetry in a lattice model for strongly interacting fermions



Liza Huijse – Harvard University
Feb 2, 2012
KITP

K. Schoutens (UvA),
P. Fendley (UVa), J. Halverson (UVa),
N. Moran (ENS), J. Vala (Maynooth), D. Mehta (Syracuse),
B. Bauer (Station Q), M. Troyer (ETH),
E. Berg (Harvard), B. Swingle (Harvard)

Introduction

Strongly interacting electron systems

Key examples:

- High T_c superconductors
- Heavy fermion compounds

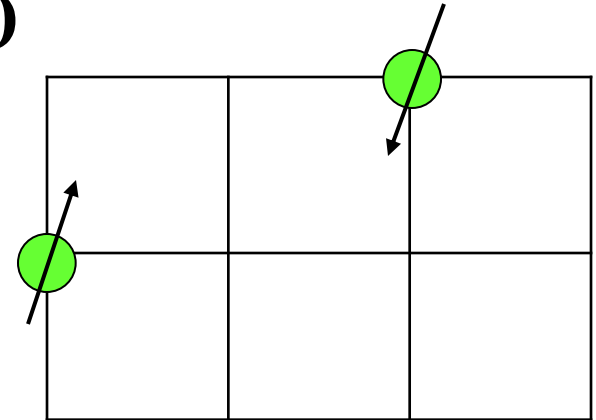
Challenge conventional theoretical techniques

Introduction

Lattice models (time continuous)

configurations:

electrons located on the sites of an ionic lattice in a solid



Hamiltonian:

typically a sum of kinetic (hopping) terms and short range repulsive interactions

Introduction

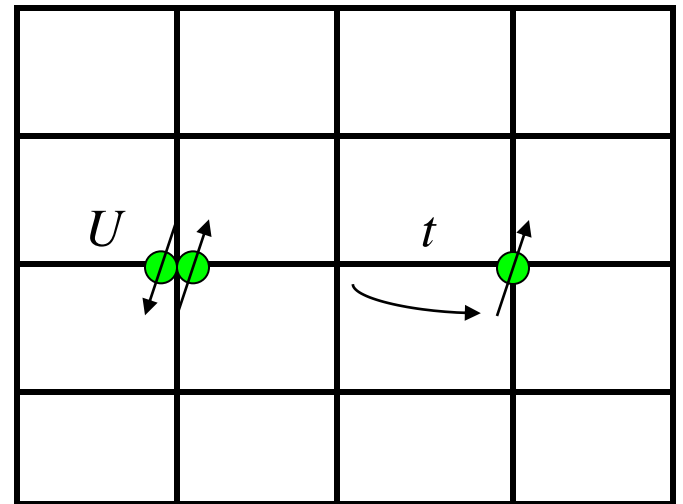
Hubbard model (1963)

Coulomb repulsion \rightarrow onsite repulsion U

$$H = -t \sum_{\langle ij \rangle} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

$$n_{i\sigma} = c_{i\sigma}^\dagger c_{i\sigma} \quad \{c_{i\sigma}^\dagger, c_{j\sigma}\} = \delta_{ij}$$

$$\sigma = \uparrow, \downarrow$$



Introduction

Hubbard model

- Kinetics dominated $n \ll 1$, $U \ll t \rightarrow$ Fermi liquid
- Interaction dominated $U \gg t$
 \rightarrow Mott insulator at half filling



Introduction

Strongly interacting electron systems

Challenge: Intermediate densities

Conventional techniques fail

- Mean field results are unreliable
- Bethe Ansatz does not work in $D > 1$
- Quantum Monte Carlo suffers from sign problem
- ...

→ Too difficult

Our work

A model for strongly interacting fermions

1. Simplifications/adjustments
2. Incorporate supersymmetry

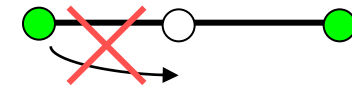
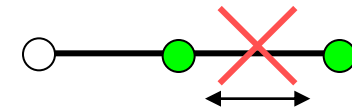
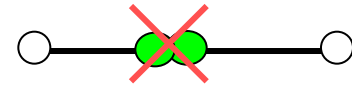
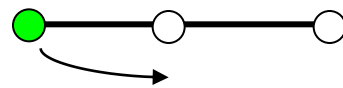
→ Exact result for strongly interacting fermions in $D > 1$
(not accessible via conventional techniques)

The model

Hardcore spinless fermions

- spinless fermions
- hardcore
- hopping t

$$V_1 \rightarrow \infty$$

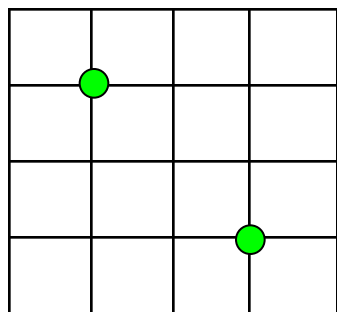
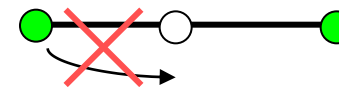
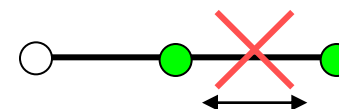
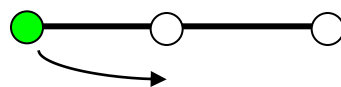


The model

Hardcore spinless fermions

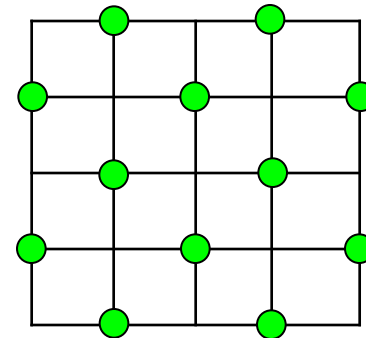
- spinless fermions
- hardcore
- hopping t

$$V_1 \rightarrow \infty$$



Fermi liquid

Stripe phase
[Henley, et. al. '01]



Insulator

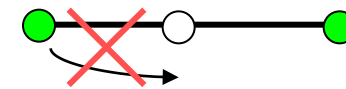
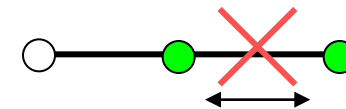
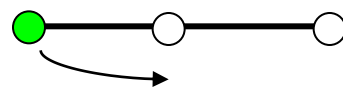
μ

The model

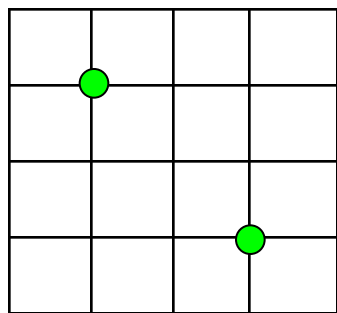
Hardcore spinless fermions

- spinless fermions
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$$V_1 \rightarrow \infty$$



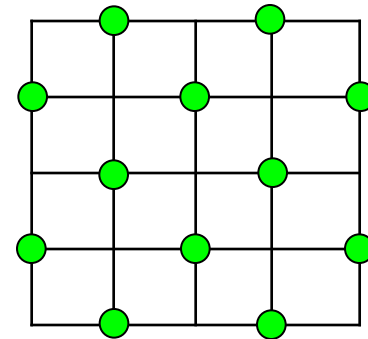
μ



Fermi liquid

Supersymmetry

[Fendley, et. al. '03]



Insulator

Plan of the talk

Benefit of supersymmetry is twofold:

- Powerful tools
- Subtle interplay between kinetic and potential terms leading to quantum criticality and superfrustration
- The model: definition & supersymmetry basics
- Witten index: superfrustration
- Cohomology: quantum ground states as tilings
- Spectral flow: quantum criticality
- Recent developments
- Conclusions

Supersymmetry

Algebraic structure

supercharges Q^+ , $Q^-(Q^+)^\dagger$ and fermion number N_f :

$$(Q^+)^2 = 0, \quad (Q^-)^2 = 0, \quad [N_f, Q^\pm] = \pm Q^\pm$$

Hamiltonian defined as

$$H = \{Q^+, Q^-\}$$

satisfies

$$[H, Q^+] = [H, Q^-] = 0, \quad [H, N_f] = 0$$

Supersymmetry

Spectrum:

- $E \geq 0$ for all states
- $E > 0$ pair into doublets (superpartners)
 $(|\psi\rangle, Q^+ |\psi\rangle), \quad Q^- |\psi\rangle = 0$
- $E = 0$ states are singlets $Q^+ |\psi\rangle = Q^- |\psi\rangle = 0$

High energy physics:

symmetry between bosonic and fermionic particles

Here:

- particles are spinless fermions (f)
- symmetry between “bosonic” (f even) and “fermionic” ($f \pm 1$ odd) states

The model

The supersymmetric lattice model

Supercharges for hardcore spinless fermions:

$$Q^+ = \sum_i c_i^\dagger \prod_{j \text{ next to } i} (1 - n_j), \quad Q^- = (Q^+)^\dagger, \quad n_j = c_j^\dagger c_j$$

Hamiltonian for 1D chain $H = \{Q^+, Q^-\}$

$$H = \sum_i [(1 - n_{i-1})c_i^\dagger c_{i+1}(1 - n_{i+2}) + \text{h.c.}] + \sum_i n_{i-1}n_{i+1} - 2N_f + L$$

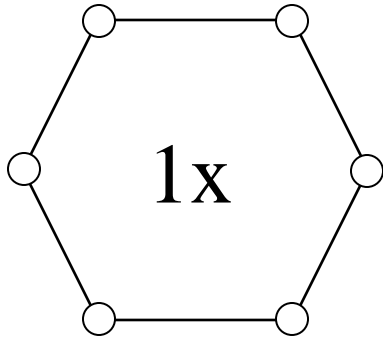
Hamiltonian for general lattice

$$H = \sum_{\langle ij \rangle} P_{\langle i \rangle} c_i^\dagger c_j P_{\langle j \rangle} + \sum_i P_{\langle i \rangle} \quad P_{\langle i \rangle} = \prod_{j \text{ next to } i} (1 - n_j)$$

Supersymmetry: example

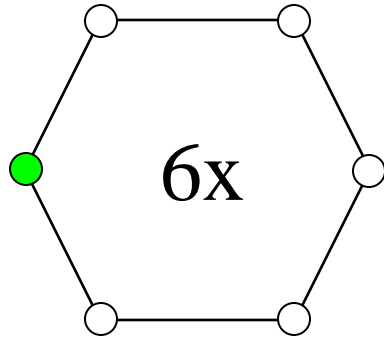
6-site chain

Possible configurations for hardcore spinless fermions



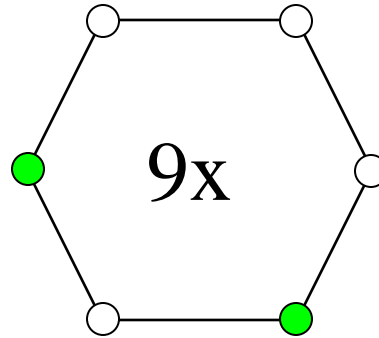
1x

$$|0\rangle$$



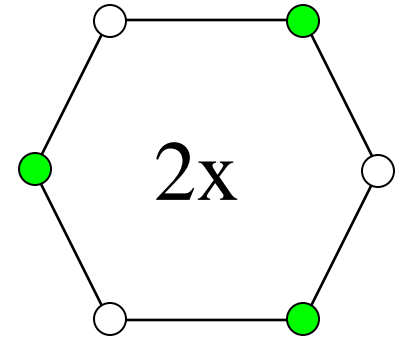
6x

$$c_i^+ |0\rangle$$



9x

$$c_i^+ c_{i+2}^+ |0\rangle$$



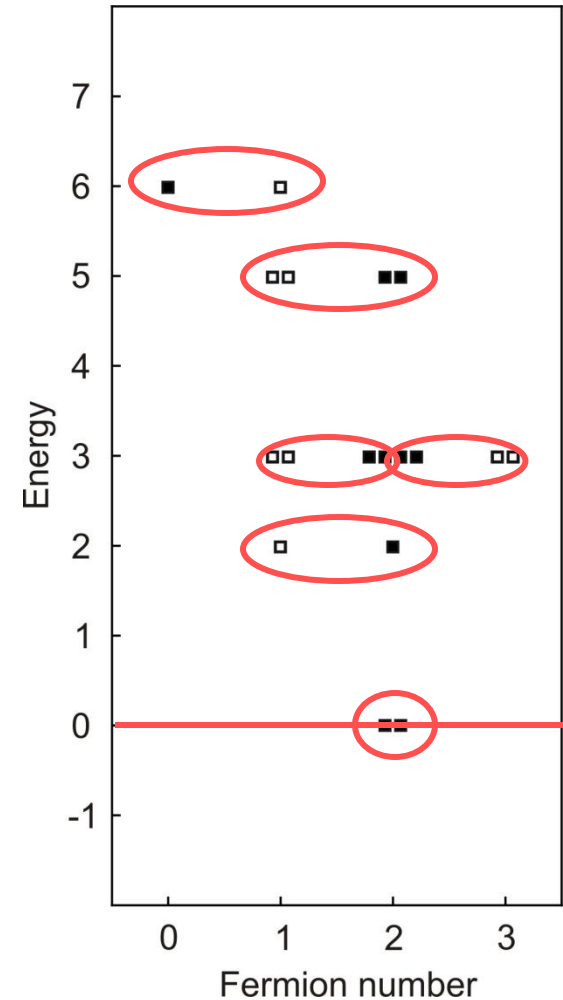
2x

$$c_i^+ c_{i+2}^+ c_{i+4}^+ |0\rangle$$

Supersymmetry: example

Manifestly supersymmetric spectrum

- Energy is positive definite
- $E > 0$ states form pairs between “fermionic” and “bosonic” states
- $E = 0$ states are singlets



Witten index: superfrustration

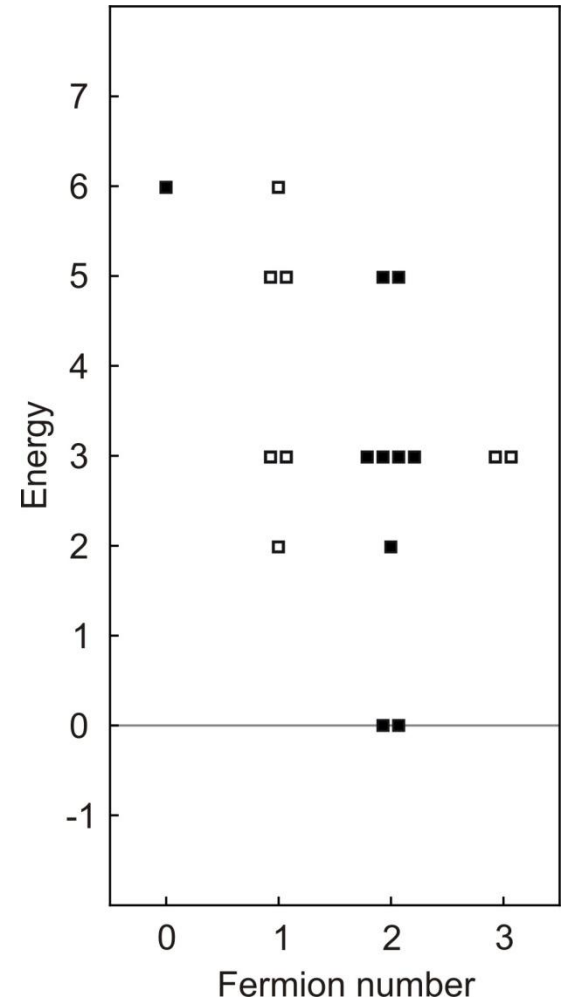
Powerful tool: Witten index

$$W = \text{Tr}(-1)^{N_f}$$

“bosonic” states
contribute +1,
“fermionic” states
contribute -1, so all
superpartners cancel

$$\Rightarrow W = \#GS_B - \#GS_F$$

|W| is lower bound to
number of ground states



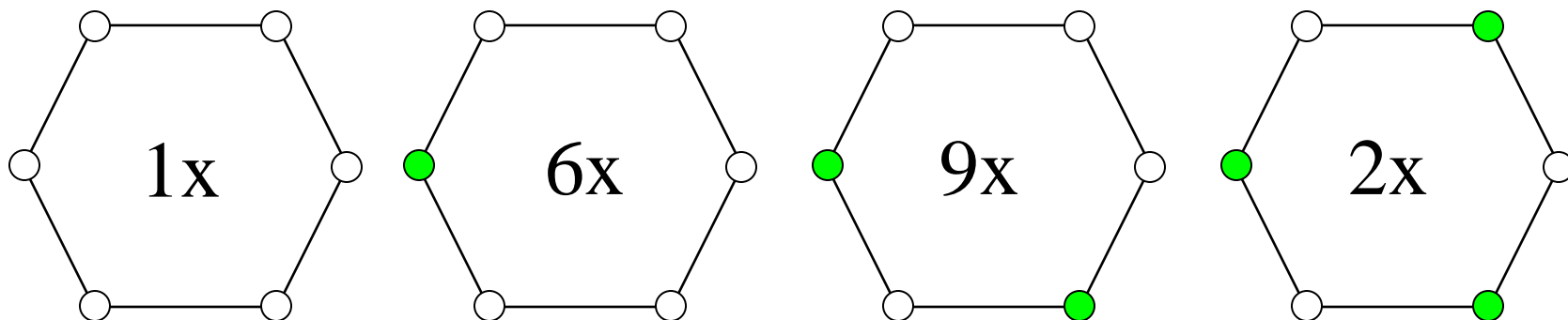
[Witten, '82]

Witten index: example

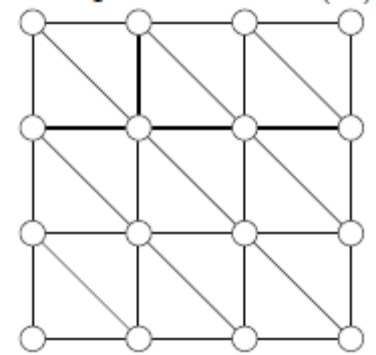
Purely combinatorial problem

$$W = \text{Tr}(-1)^{N_f}$$

$$W = 1 - 6 + 9 - 2 = 2$$



Witten index



Triangular lattice

$N \times M$ sites with periodic BC

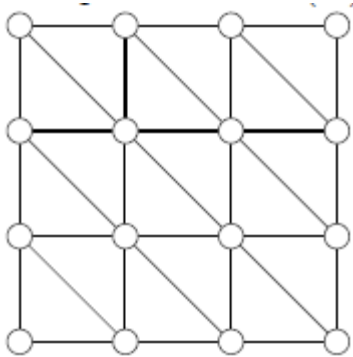
	1	2	3	4	5	6	7	8	9	10
1	1	1	1	1	1	1	1	1	1	1
2	1	-3	-5	1	11	9	-13	-31	-5	57
3	1	-5	-2	7	1	-14	1	31	-2	-65
4	1	1	7	-23	11	25	-69	193	-29	-279
5	1	11	1	11	36	-49	211	-349	811	-1064
6	1	9	-14	25	19	188	-19	-415	1462	-4911
7	1	-13	1	-69				3403	-7055	5237
8	1	-31	31	193				881	-28517	50849
9	1	-5	-2	-29	881	1462	-7055	-28517	31399	313315
10	1	57	-65	-279	-1064	-4911	5237	50849	313315	950592
11	1	67	1	859	1651	12607	32418	159083	499060	2011307
12	1	-47	130	-1295	-589	-26006	-152697	-535895	-2573258	-3973827
13	1	-181	1	-77	-1949	67523	330331	-595373	-10989458	-49705161
14	1	-87	-257	3641	12611	-139935	-235717	5651377	4765189	-232675057
15	1	275	-2	-8053	-32664	272486	-1184714	-1867189	134858383	-702709340

$$|W| \sim 1.14^{NM}$$

[van Eerten, '05]

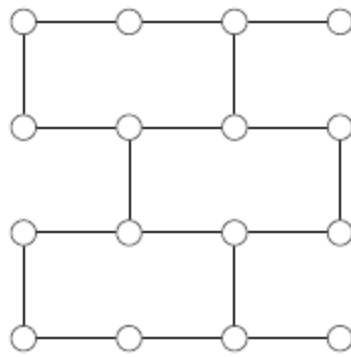
Superfrustration - examples

Tri



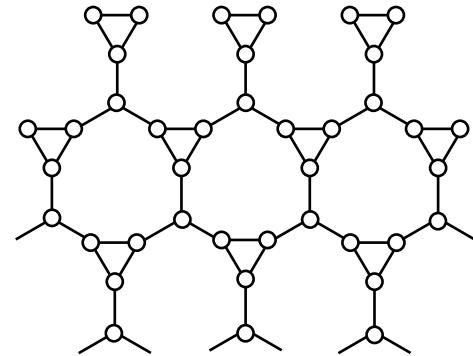
$$|W| \sim 1.14^V$$

Hex



$$|W| \sim 1.2^V$$

Martini



$$|W| \sim 1.17^V$$

V is number of sites (2D volume)

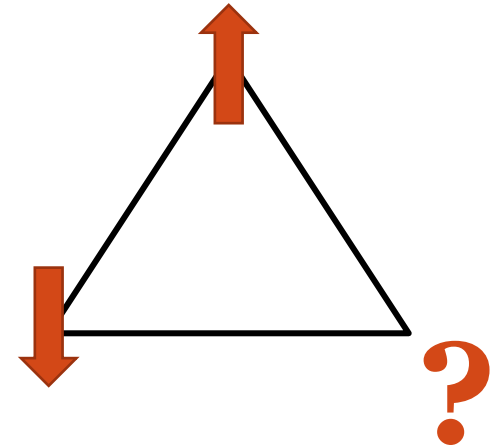
[van Eerten, '05;
Fendley, Schoutens, '05]

Superfrustration

Frustration

Competing terms in hamiltonian

→ multiple ground states



Supersymmetry

Subtle competition between kinetic and potential terms

→ for 2D lattices exponential ground state degeneracy

Violation of 3rd law of thermodynamics

Exponential number of ground states

→ finite zero temperature entropy

Superfrustration

'3-rule'

- repulsive interactions favor 3-site interparticle distance
- chemical potential favors higher densities

Combined with kinetic terms

→ quantum charge frustration at intermediate densities

**Cohomology:
quantum ground states as tilings**

Powerful tool: Cohomology

- Cohomology of Q
 - GS are in 1-1 correspondence with cohomology elements
- More difficult to compute than Witten index
- But gives more information:
 - gives total number of gs
 - gives fermion number of gs
 - often gives relation between gs and geometric object

Cohomology technique

Lemma

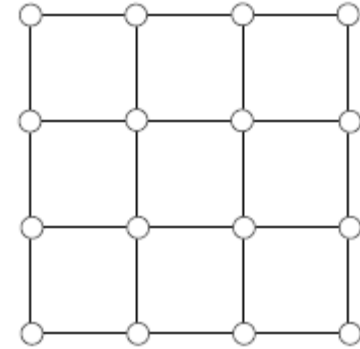
Susy ground states are in 1-1 correspondence with the cohomology

$$H_{Q,N_f} = \text{Ker}[Q^+]_{N_f} / \text{Im}[Q^+]_{N_f-1}$$

of Q^+ in the complex

$$\dots \xrightarrow{Q^+} H_{N_f} \xrightarrow{Q^+} H_{N_f+1} \xrightarrow{Q^+} \dots$$

Square lattice: Witten index



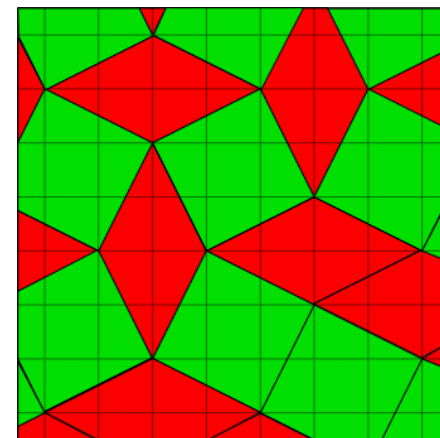
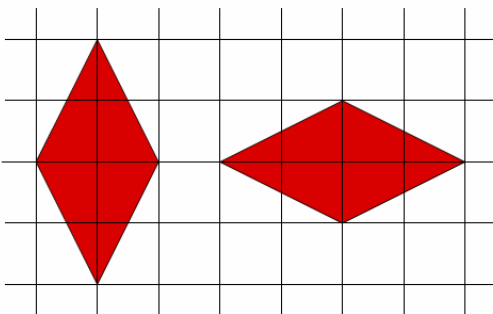
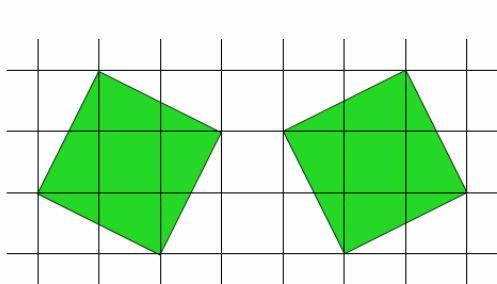
$N \times M$ sites with periodic BC

	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	19	20
1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
2	1	-1	1	3	1	-1	1	3	1	-1	1	3	1	-1	1	3	1	-1	1	3
3	1	1	4	1	1	4	1	1	4	1	1	4	1	1	4	1	1	4	1	1
4	1	3	1	7	1	3	1	7	1	3	1	7	1	3	1	7	1	3	1	7
5	1	1	1	1	-9	1	1	1	1	11	1	1	1	1	-9	1	1	1	1	11
6	1	-1	4	3	1	14	1	3	4	-1	1	18	1	-1	4	3	1	14	1	3
7	1	1	1	1	1	1	1	1	1	1	1	1	1	-27	1	1	1	1	1	1
8	1	3	1	7	1	3	1	7	1	43	1	7	1	3	1	7	1	3	1	47
9	1	1	4	1	1	4	1	1	40	1	1	4	1	1	4	1	1	76	1	1
10	1	-1	1	3	11	-1	1	43	1	9	1	3	1	69	11	43	1	-1	1	13
11	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
12	1	3	4	7	1	18	1	7	4	3	1	166	1	3	4	7	1	126	1	7
13	1	1	1	1	1	1	1	1	1	1	1	1	-51	1	1	1	1	1	1	1
14	1	-1	1	3	1	-1	-27	3	1	69	1	3	1	55	1	451	1	-1	1	73
15	1	1	4	1	-9	4	1	1	4	11	1	4	1	1	174	1	1	4	1	11

[Fendley - Schoutens - van Eerten '05;
Jonsson '06]

Square lattice (periodic BC)

- Cohomology of Q gives direct relation between ground states and rhombus tilings
- # of tiles = # of fermions



[LH - Schoutens '10]

Square lattice (periodic BC)

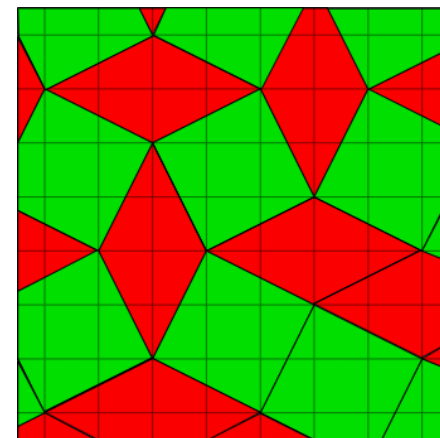
- Cohomology of Q gives direct relation between ground states and rhombus tilings
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- GS at intermediate filling

$$\nu = \frac{N_f}{NM} \in [1/5, 1/4] \cap \mathbb{Q}$$

- Sub-extensive GS entropy

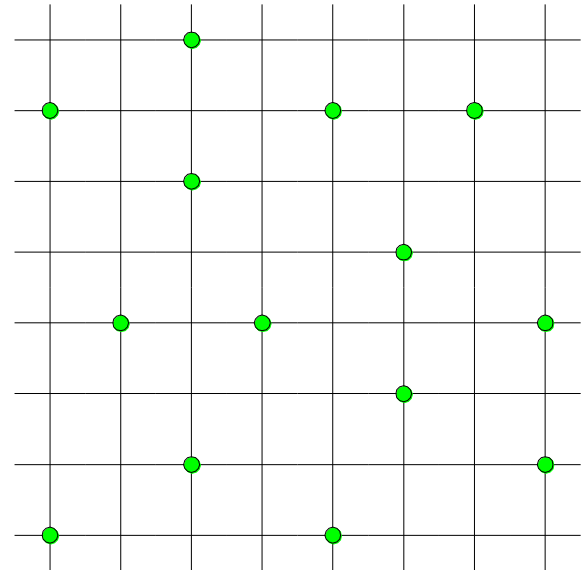
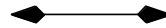
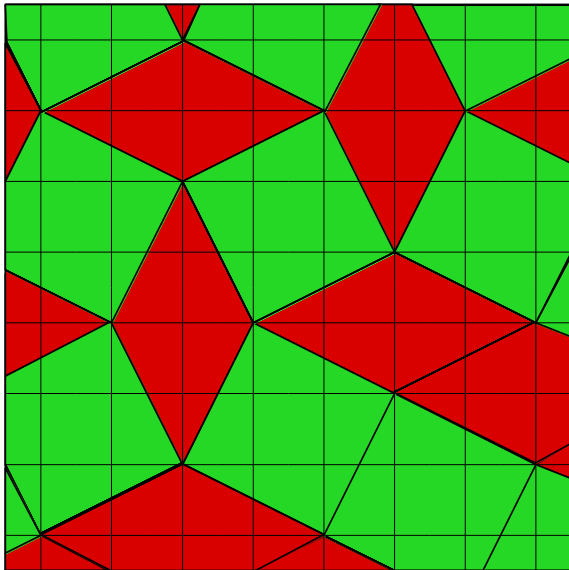
$$S_{gs} \sim 0.46(N+M)$$



[LH - Schoutens '10]

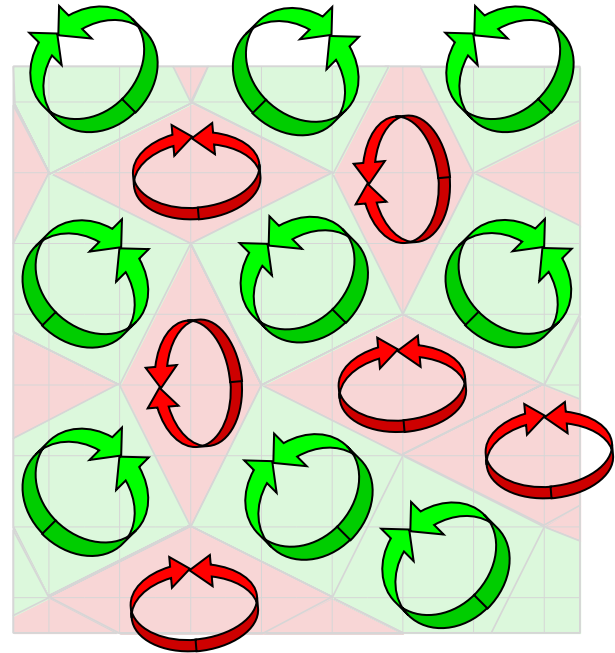
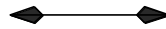
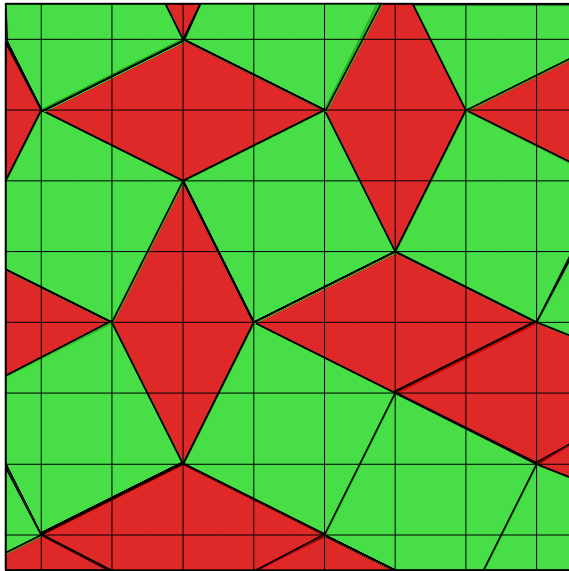
Square lattice: ground states

- # gs grows exponentially with the **linear** size of the system
- zero energy ground states found at **intermediate** filling
- compelling evidence for **critical edge modes**
- what is the nature of these states?

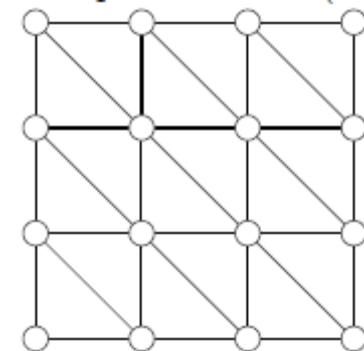


Square lattice: ground states

- # gs grows exponentially with the **linear** size of the system
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Witten index



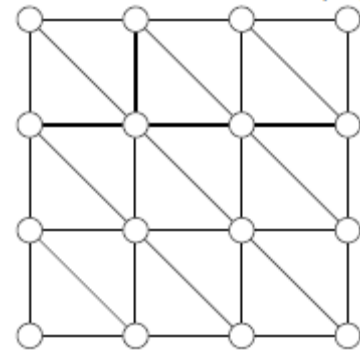
Triangular lattice

$N \times M$ sites with periodic BC

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6	1	9	-14	25	-49	-102	-13	-415	1462	-4911
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[van Eerten, '05]

Triangular lattice: ground states

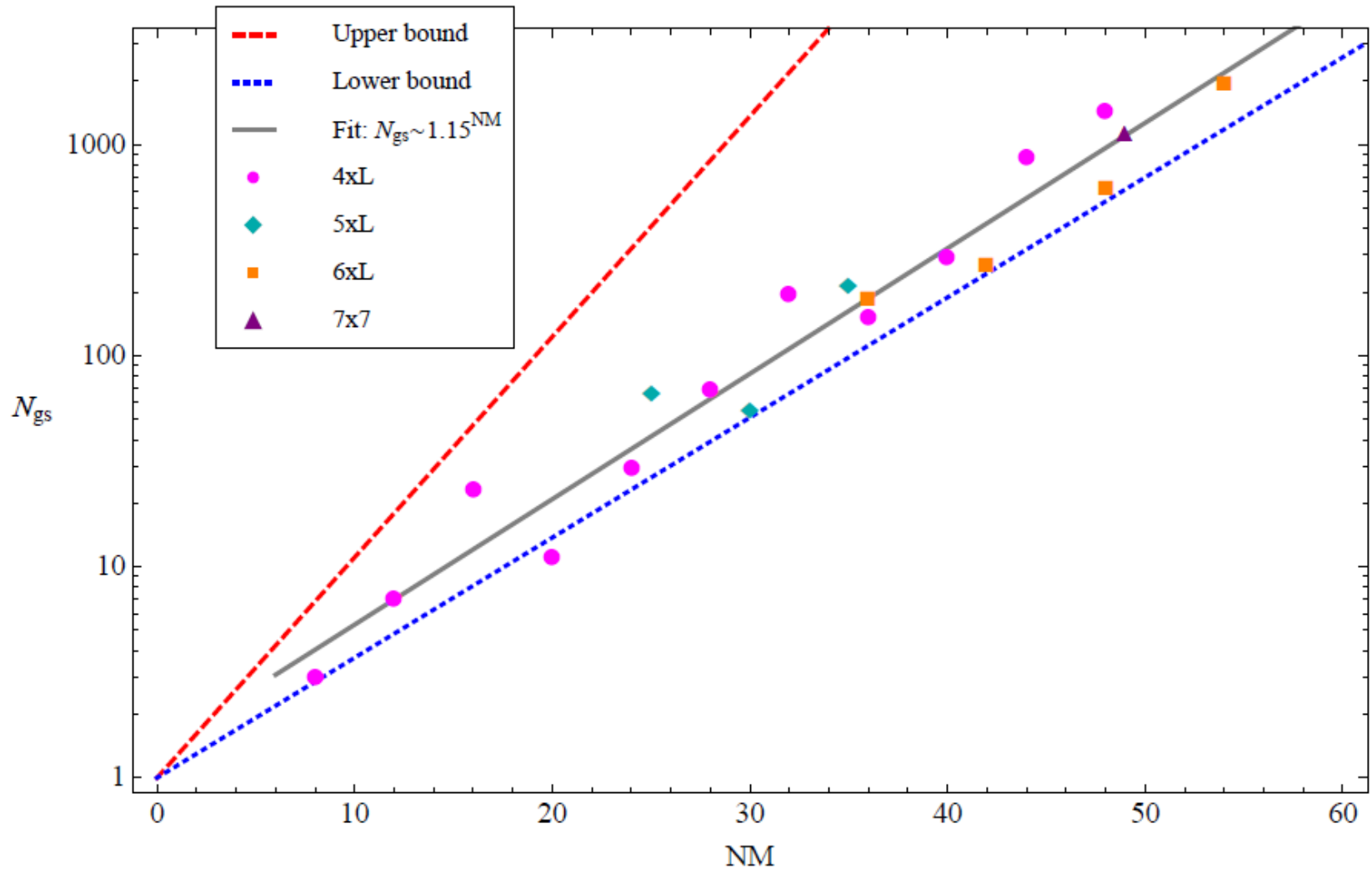


Full cohomology problem is very hard, but intermediate results:

- Upper bound on #gs $\#gs \lesssim (\sqrt{\phi})^{MN} \sim 1.27^{MN}$
- Lower bound on filling $\nu \in [1/7, 1/5] \cap \mathbb{Q}$
- Upper bound on filling $\nu \in [1/8, 1/4] \cap \mathbb{Q}$

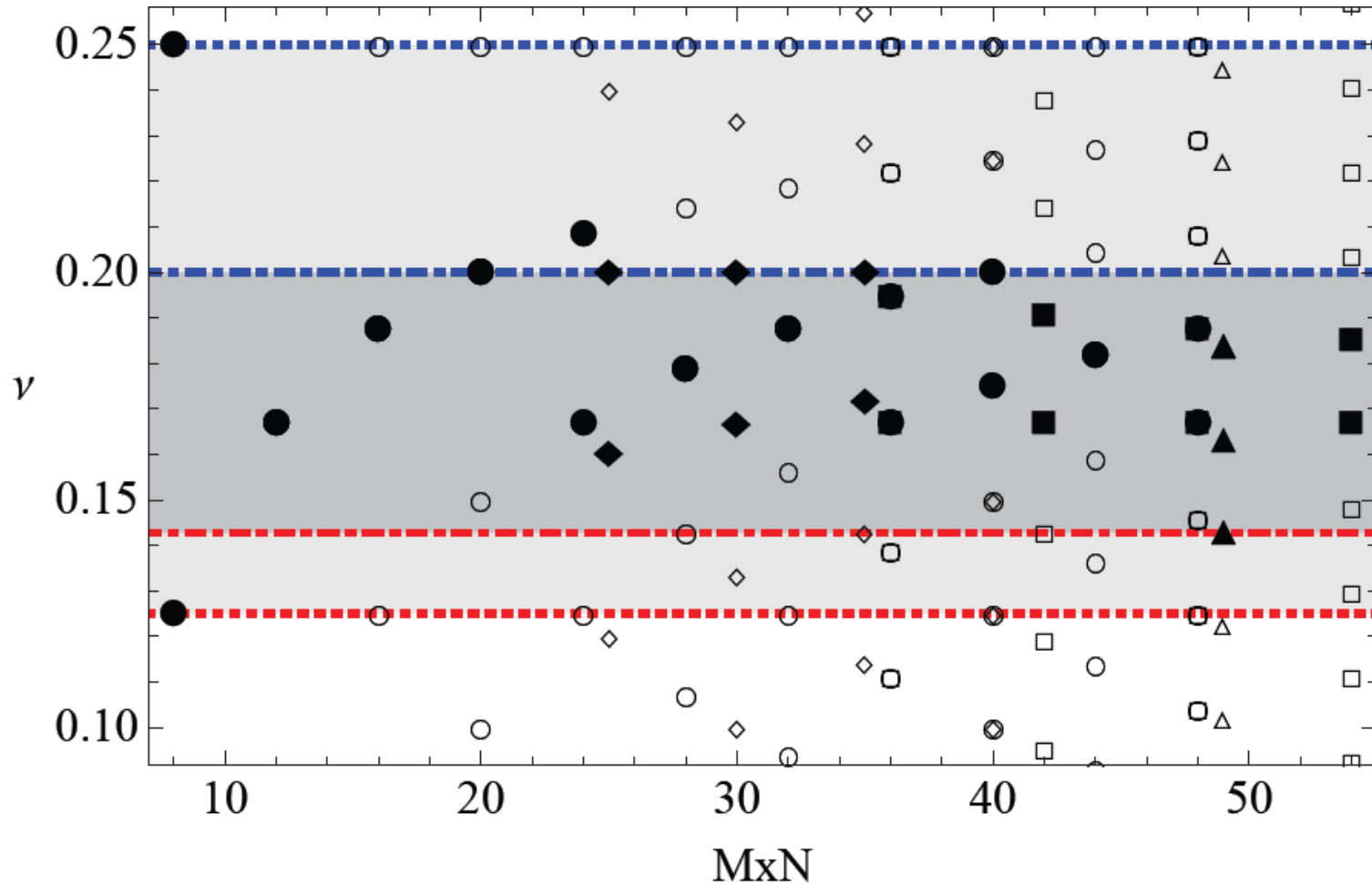
[Jonsson, '10; Engström, '09;
LH, Mehta, Moran, Schoutens, Vala, '11]

Numerical results: Ground state degeneracy



[LH, Mehta, Moran, Schoutens, Vala, '11]

Numerical results: Filling fraction



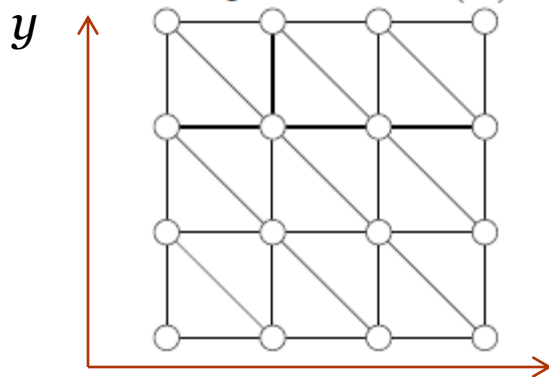
[LH, Mehta, Moran, Schoutens, Vala, '11]

Numerical results: Momentum

**Flatband dispersion:
Zero energy states at
all momenta**

Eigenvalues of translations:

$$t_x = e^{2\pi i k_x / N}, \quad t_y = e^{2\pi i k_y / M}$$



x

[LH, Mehta, Moran, Schoutens, Vala, '11]

6×9

$N_{\text{gs}} = 1926, W = 1462$

$f = 9$						$f = 10$							
(k_x, k_y)	0	1	2	3	4	5	(k_x, k_y)	0	1	2	3	4	5
0	9	3	3	9	3	3	0	33	31	31	33	31	31
1	4	5	3	4	5	3	1	32	31	31	32	31	31
2	4	3	5	4	3	5	2	32	31	31	32	31	31
3	7	4	4	7	3	4	3	32	32	31	32	30	31
4	4	5	3	4	5	3	4	32	31	31	32	31	31
5	4	3	5	4	3	5	5	32	31	31	32	31	31
6	7	4	3	7	4	4	6	32	31	30	32	31	32
7	4	5	3	4	5	3	7	32	31	31	32	31	31
8	4	3	5	4	3	5	8	32	31	31	32	31	31

Spectral flow: quantum criticality

1D chain

$$H = \sum_i [(1 - n_{i-1})c_i^\dagger c_{i+1}(1 - n_{i+2}) + \text{h.c.}] + \sum_i n_{i-1}n_{i+1} - 2N_f + L$$

- Periodic chain of length L :
2 gs for $L \bmod 3 = 0$;
1 gs otherwise
- Fermion number in ground state: $f = [L/3]$
- Bethe Ansatz solution (integrable)
- Continuum limit: $\mathcal{N}=(2,2)$ SCFT with central charge $c=1 \rightarrow$ quantum critical, emergent spacetime supersymmetry

[Fendley-Schoutens-deBoer '03,
Fendley-Nienhuis-Schoutens '03,
Beccaria-DeAngelis '05,
Fendley-Hagendorf '10 & '11, LH '11]

Relations/extensions

- XXZ spin chain – exact mapping
- SUSY Matrix models of Veneziano-Wosiek (via mapping to XXZ spin chain)
- Generalize hard-core constraint
 - Allow k particles to be nearest neighbors, but not $k+1$:
 M_k susy model \leftrightarrow k -th SCFT minimal model

[Fendley , Nienhuis , Schoutens '03
Veneziano, Wosiek, '06]

1D Quantum criticality

Numerical techniques to identify continuum CFT

- Finite size scaling of energy gap: $E \sim 1/L$
- Entanglement entropy: $S \sim c \log L$
- Superconformal field theories: Spectral flow:
Energy depends parabolically on boundary twist
 - + Accurate for small systems
 - + No scaling required

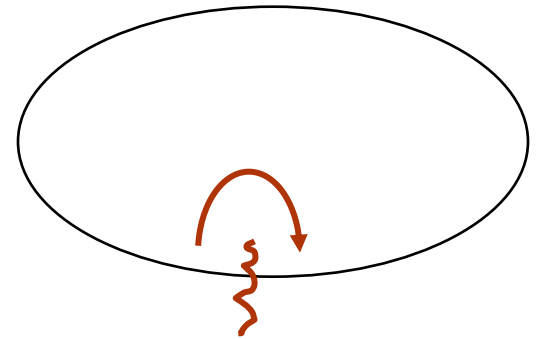
Boundary twist: spectral flow

wave function picks up a phase $\exp(2\pi i\alpha)$
as a particle hops over a “boundary”:

$$e^{2\pi i\alpha} c_L^\dagger c_1 + \text{h.c.}$$

twist: $\alpha: 0 \leftrightarrow 1/2$

“pbc \leftrightarrow apbc” = “R \leftrightarrow NS sector”



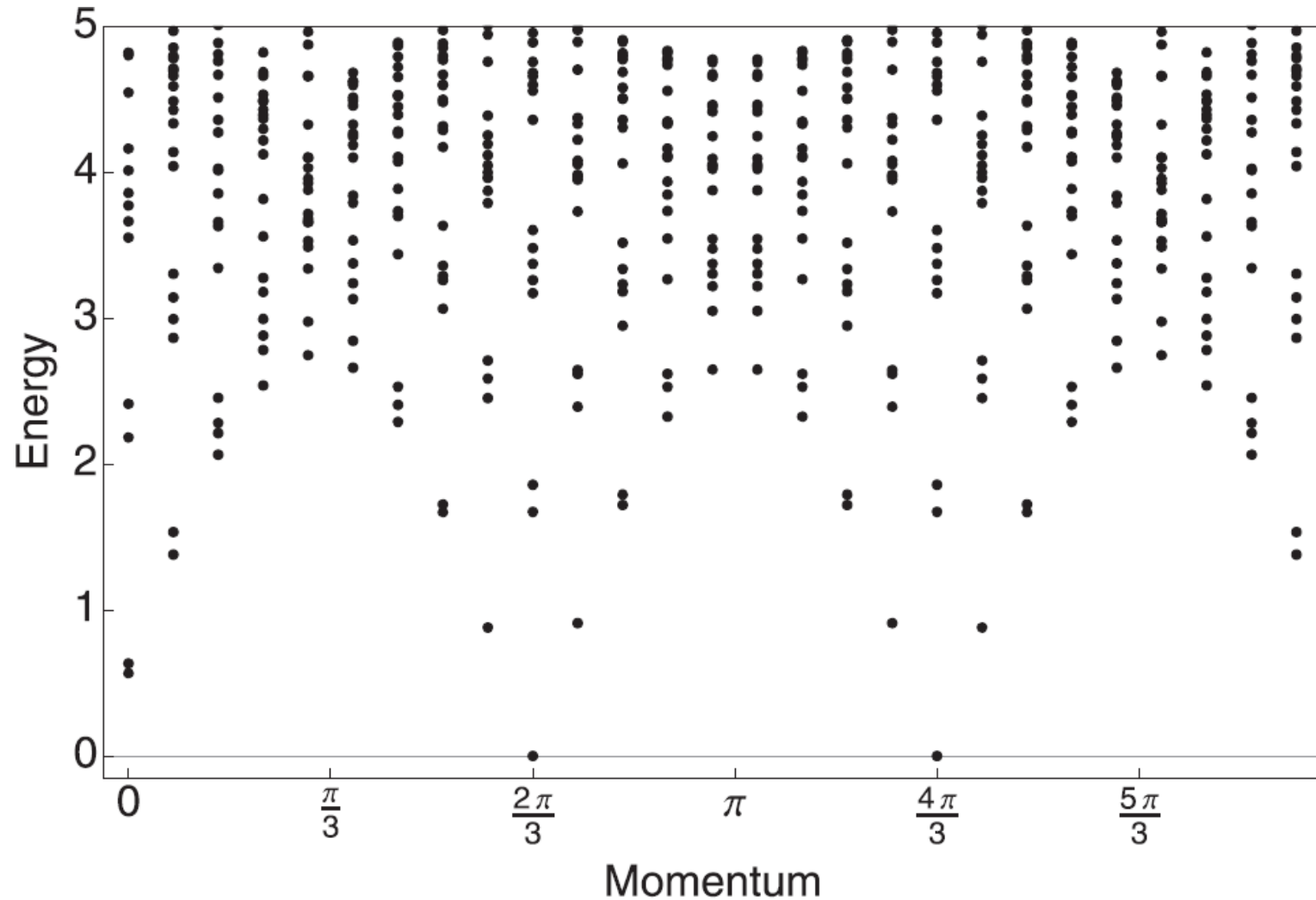
SCFT: energy is parabolic function of twist parameter

$$E_\alpha = E_0 - \alpha Q_0 + \alpha^2 c/3$$

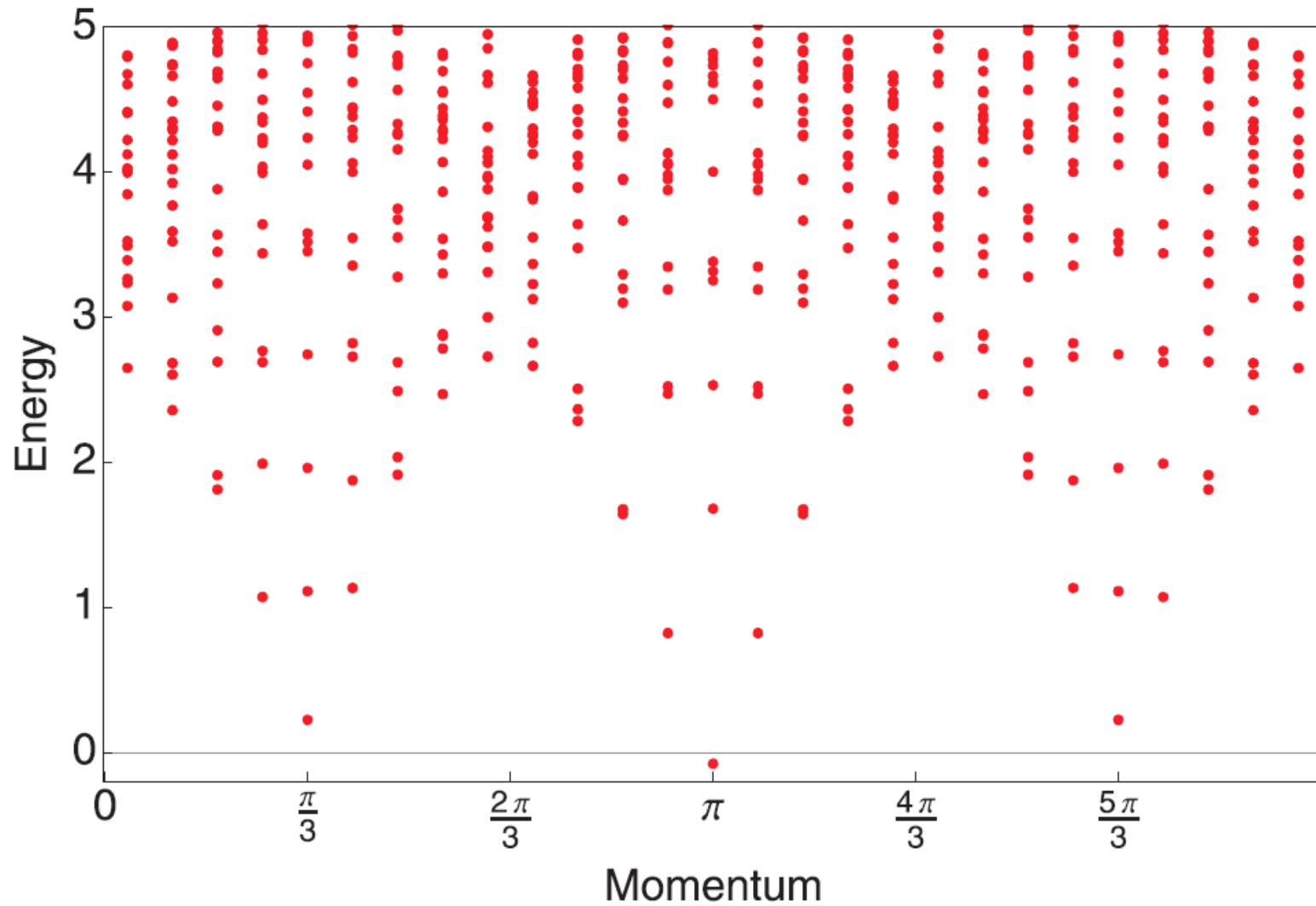
$$\text{R} : \alpha = 0$$

$$\text{NS} : \alpha = 1/2$$

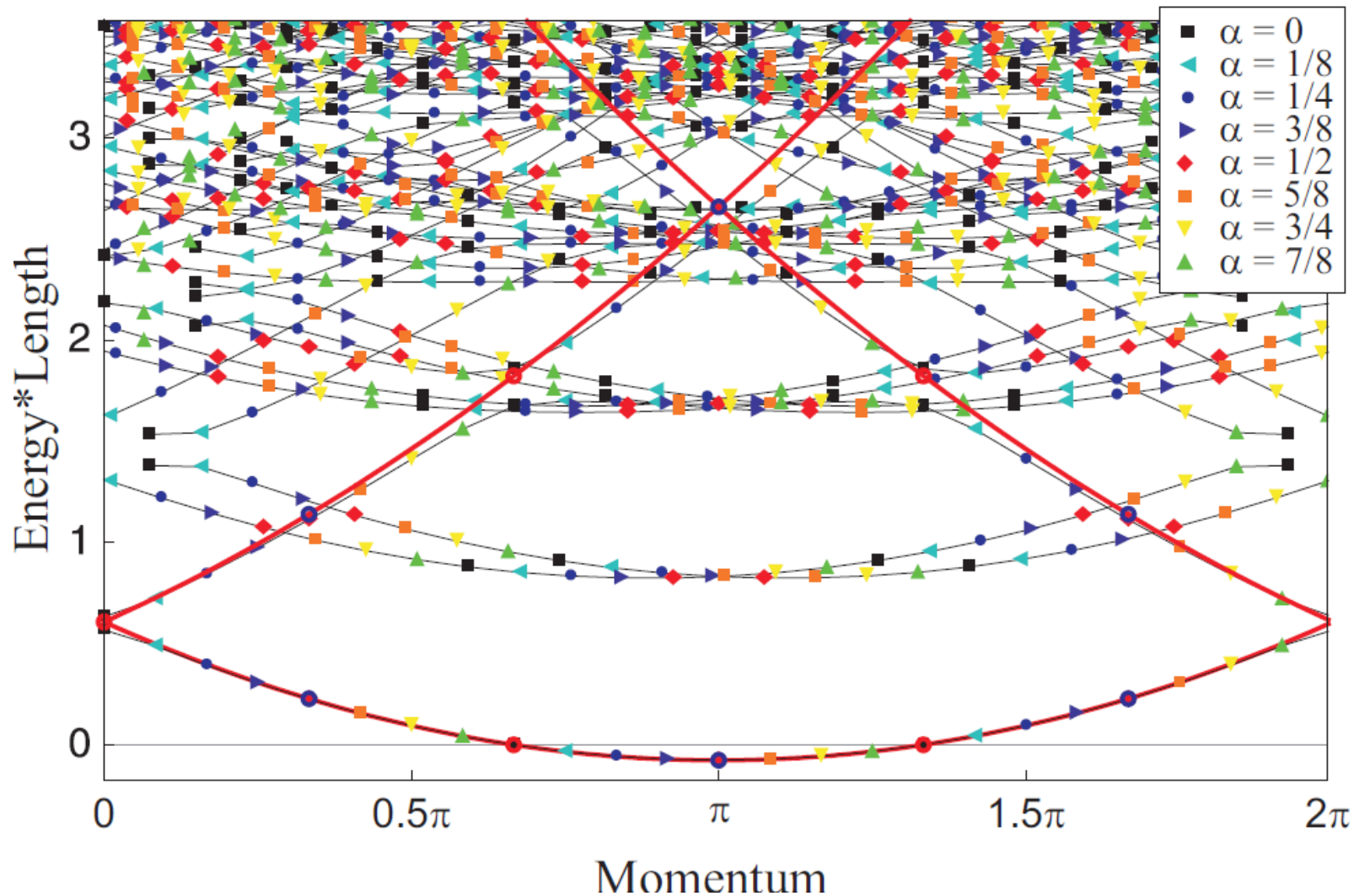
Chain Spectrum, $L=27$, $N_f=9$, PBC ($\alpha=0$)



Chain Spectrum, $L=27$, $N_f=9$, APBC

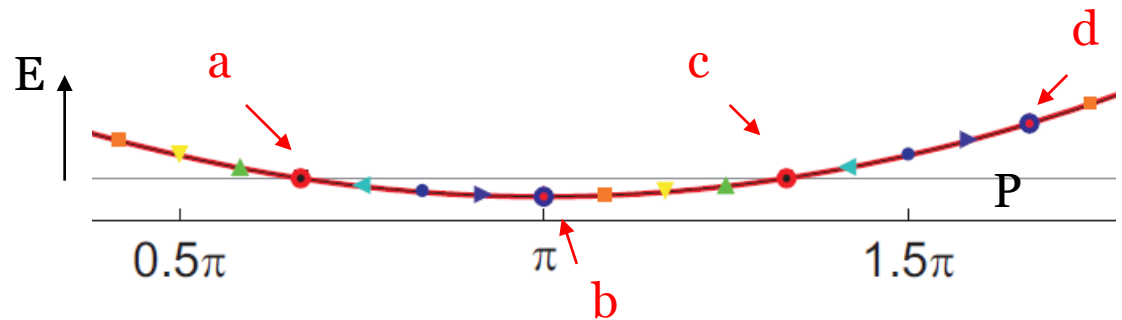


Spectral flow chain, $L=27$, $N_f=9$



What can we learn from spectral flow?

- 3 fit parameters
- 4 unknowns:
 E , Q_0 , c and v_F
- \rightarrow ratios
- for 1D chain we extract:



numerics

SCFT

state	E/c	Q_0/c	E/c	Q_0/c
a	0	-0.334	0	-1/3
b	-0.083	0	-1/12	0
c	0	0.342	0	1/3
d	0.254	0.675	1/4	2/3

Spectral flow for the square lattice

- square ladder

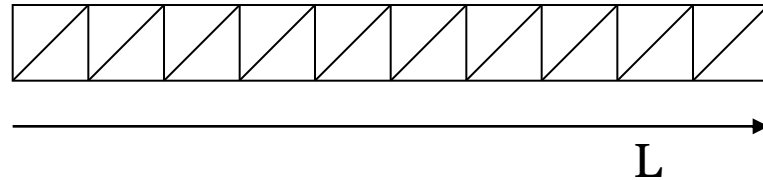
$(2,0) \times (0,L)$



- zigzag ladder

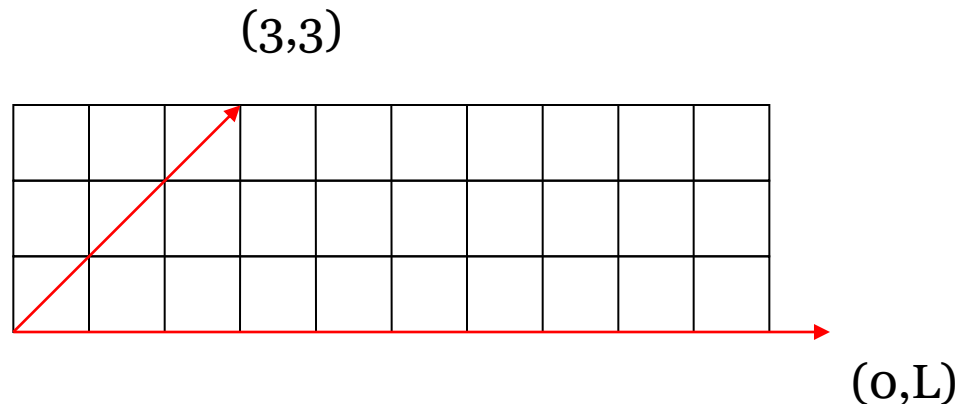
$(2,1) \times (0,L)$

GS for $\nu \in [1/5, 1/4]$

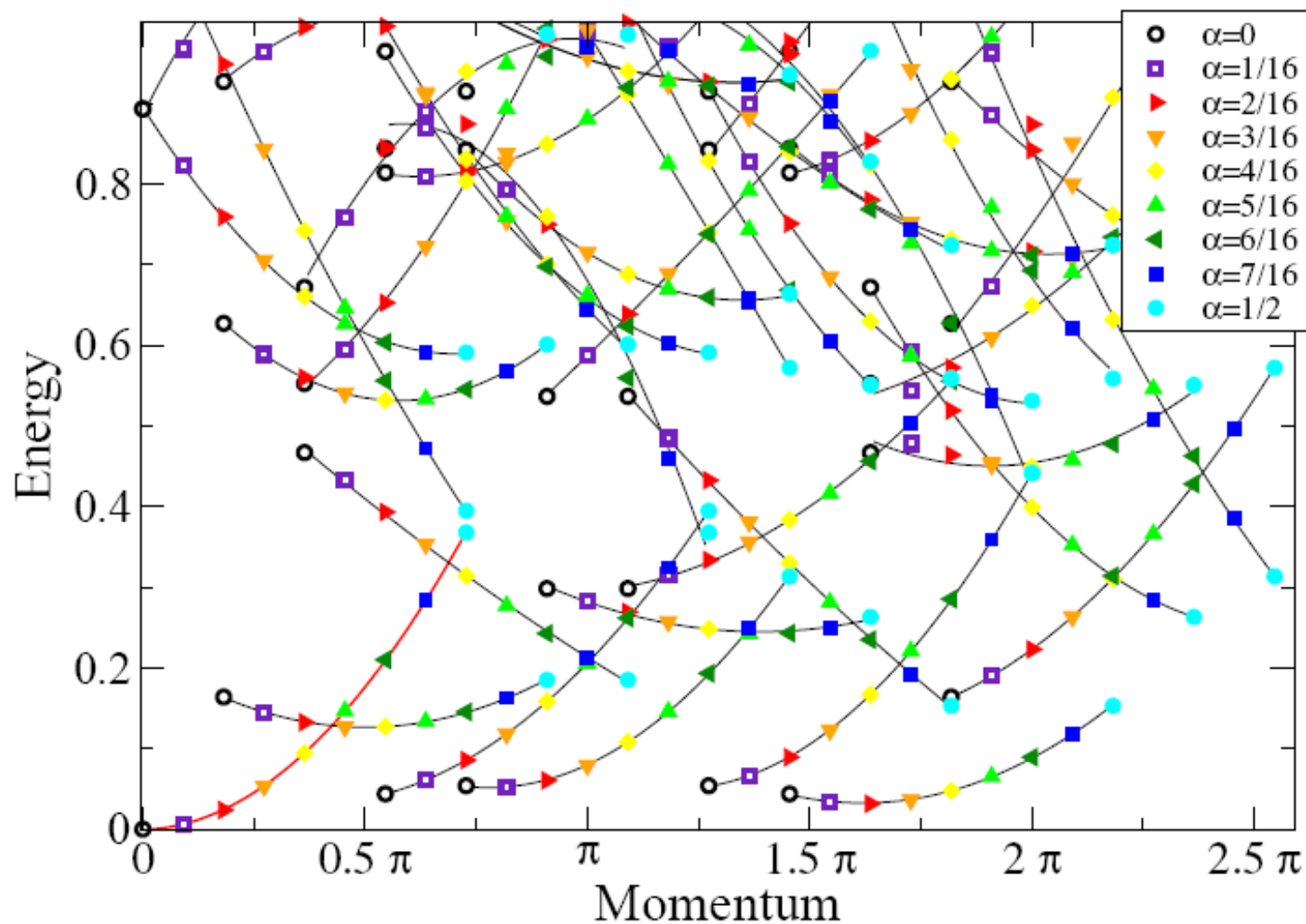


- $(3,3) \times (0,L)$

fermions can hop
past each other



Spectral flow results $(3,3) \times (0,11)$, $N_f=8$



Spectral flow results

$(L, 0) \times (3, 3)$

N	f	E/c	Q/c
18	4	-0.0851	0.004
36	8	-0.0841	-0.002
15	4	0.0898	0.349
21	4	0.0850	0.337
24	5	0.0850	0.337
30	7	0.0853	0.338
33	8	0.0855	0.338

$(L, 0) \times (1, 2)$

N	f	E/c	Q/c
9	2	-0.0858	-0.005
18	4	-0.0842	-0.002
27	6	-0.0839	-0.001
17	4	0.0844	0.336
26	6	0.0840	0.335
35	8	0.0839	0.335
14	3	0.2666	0.701
23	5	0.2458	0.657
32	7	0.2432	0.652

$(L, 0) \times (0, 2)$

N	f	E/c	Q/c
16	4	-0.0897	-0.014
24	6	-0.0889	-0.012
32	8	-0.0885	-0.011
12	3	0.0911	0.350
20	5	0.0900	0.348
28	7	0.0894	0.347
14	4	0.0855	0.338
22	6	0.0849	0.337
30	8	0.0847	0.336

Spectral flow results

$(L, 0) \times (3, 3)$

N	f	E/c	Q/c
18	4	-0.0851	0.004
36	8	-0.0841	-0.002
15	4	0.0898	0.349
21	4	0.0850	0.337
24	5	0.0850	0.337
30	7	0.0853	0.338
33	8	0.0855	0.338

$(L, 0) \times (1, 2)$

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minimal models in SCFT: $c = \frac{3k}{k+2}$

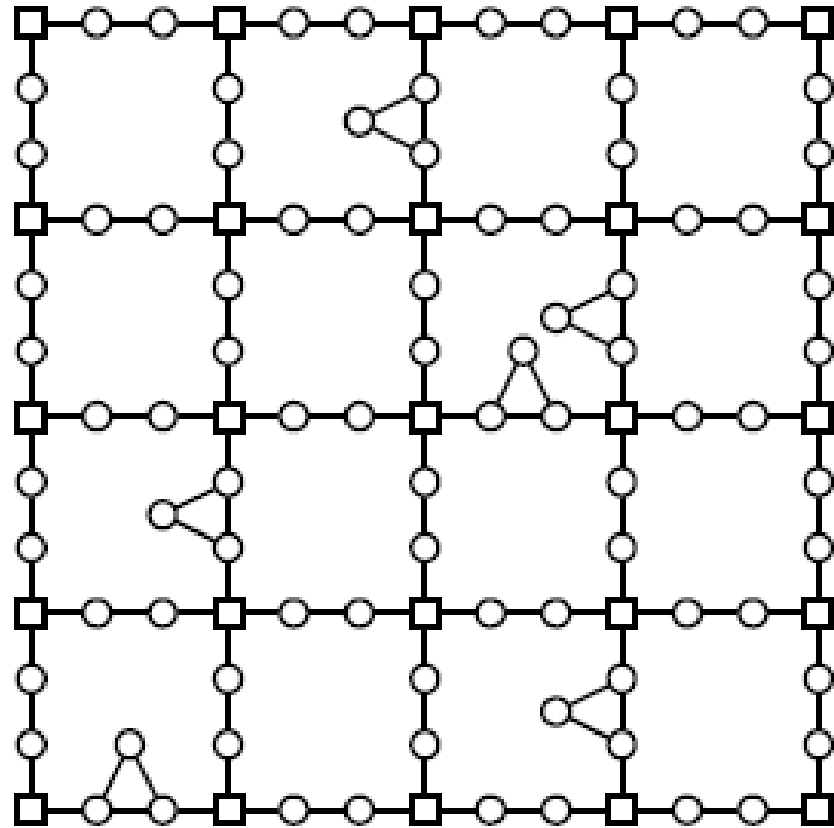
$$E/c = \frac{4l-k}{12k} \text{ and } Q_0/c = \frac{2l}{3k}$$

$l = 0 : (-1/12, 0), l = k/2 : (1/12, 1/3), l = k : (1/4, 2/3)$

Open problems/recent developments

Controllable gs degeneracy

- Triangles (t) as impurities
- $\#gs = 2^t$
- Anomalous scaling of entanglement entropy



[LH-Swingle (to appear)]

Staggering: beyond gs counting in 2D

$$Q^+ = \sum \lambda_i^* c_i^\dagger P_{\langle i \rangle}$$
$$Q^- = \sum \lambda_i c_i P_{\langle i \rangle}$$
$$\lambda_i = \begin{cases} 1 & \text{for } i \in S_1 \\ y & \text{for } i \in S_2 \end{cases}$$

- Exact ground states from cohomology for lattices with reduced frustration
- Towards understanding the phases in the square lattice: system of decoupled chains in infinite staggering limit

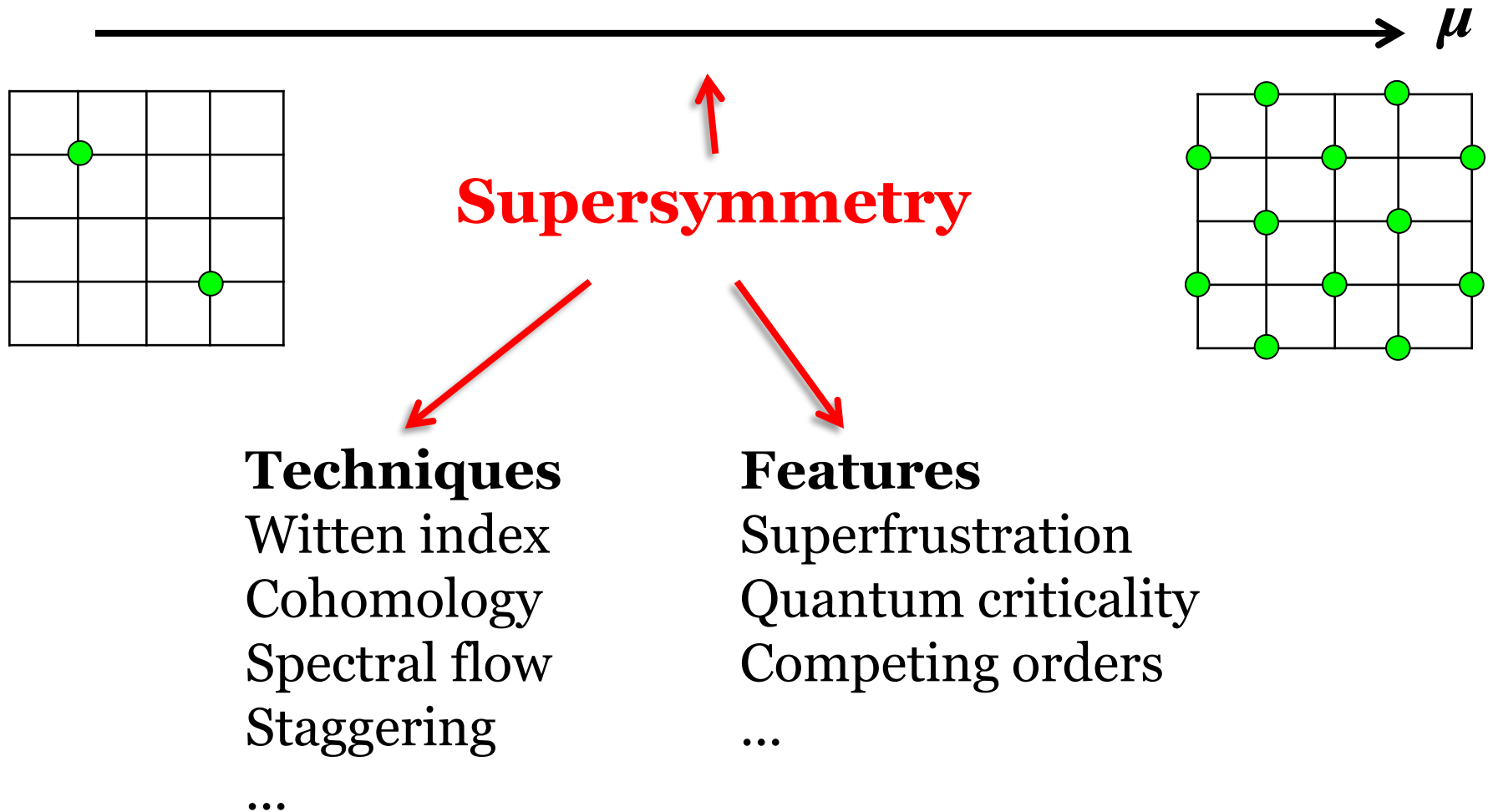
[LH - Moran - Schoutens - Vala '11 ;
Fendley – Hagendorf '10; '11;
LH-Berg (work in progress)]

Open questions/future directions

- AdS/CMT: apply gauge/gravity duality to condensed matter systems
- Toy model for AdS/CMT??
- Typical features: extensive ground state entropy, supersymmetry and large N
- What is continuum theory in 2D?
- Is there an emergent gauge symmetry?
- Can we include gauge symmetry explicitly on the lattice?
- ...

Conclusions

Exact results for strongly interacting fermions



Thank you