

ACME

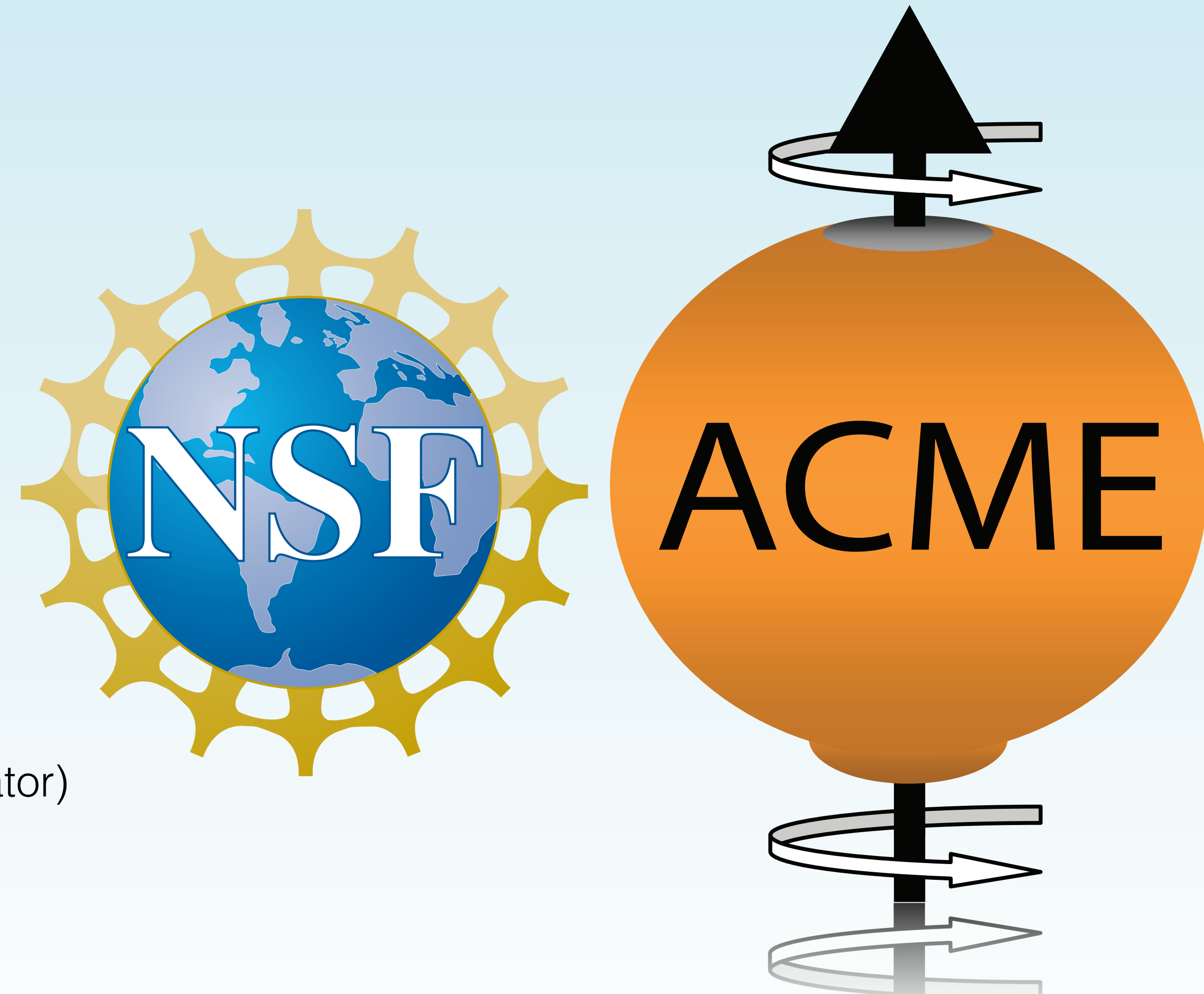
Advanced Cold Molecule Electron EDM Search

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D. DeMille^Y, G. Gabrielse^H, JMD^H

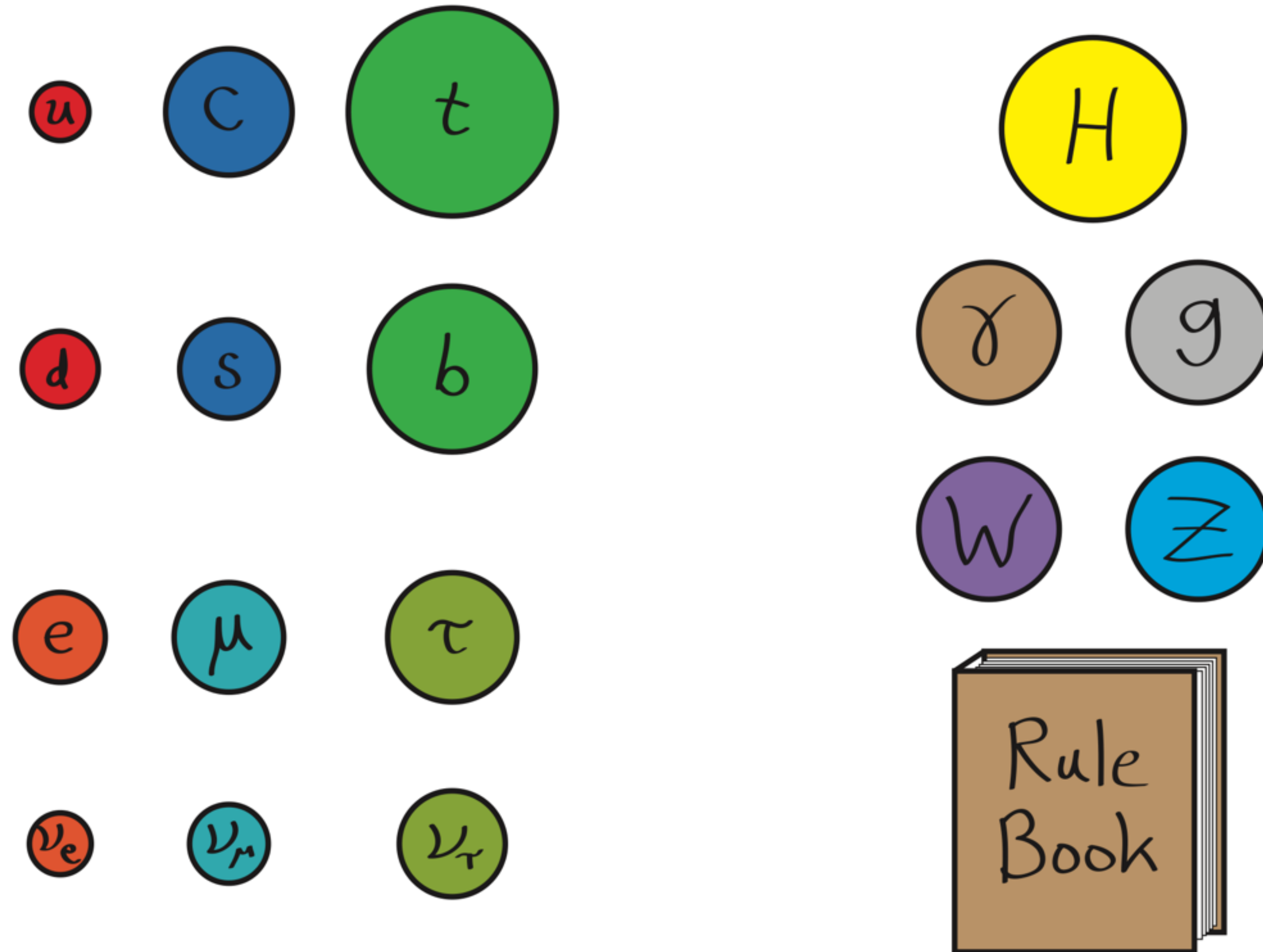
^HHarvard ^YYale

ACME "Graduates":

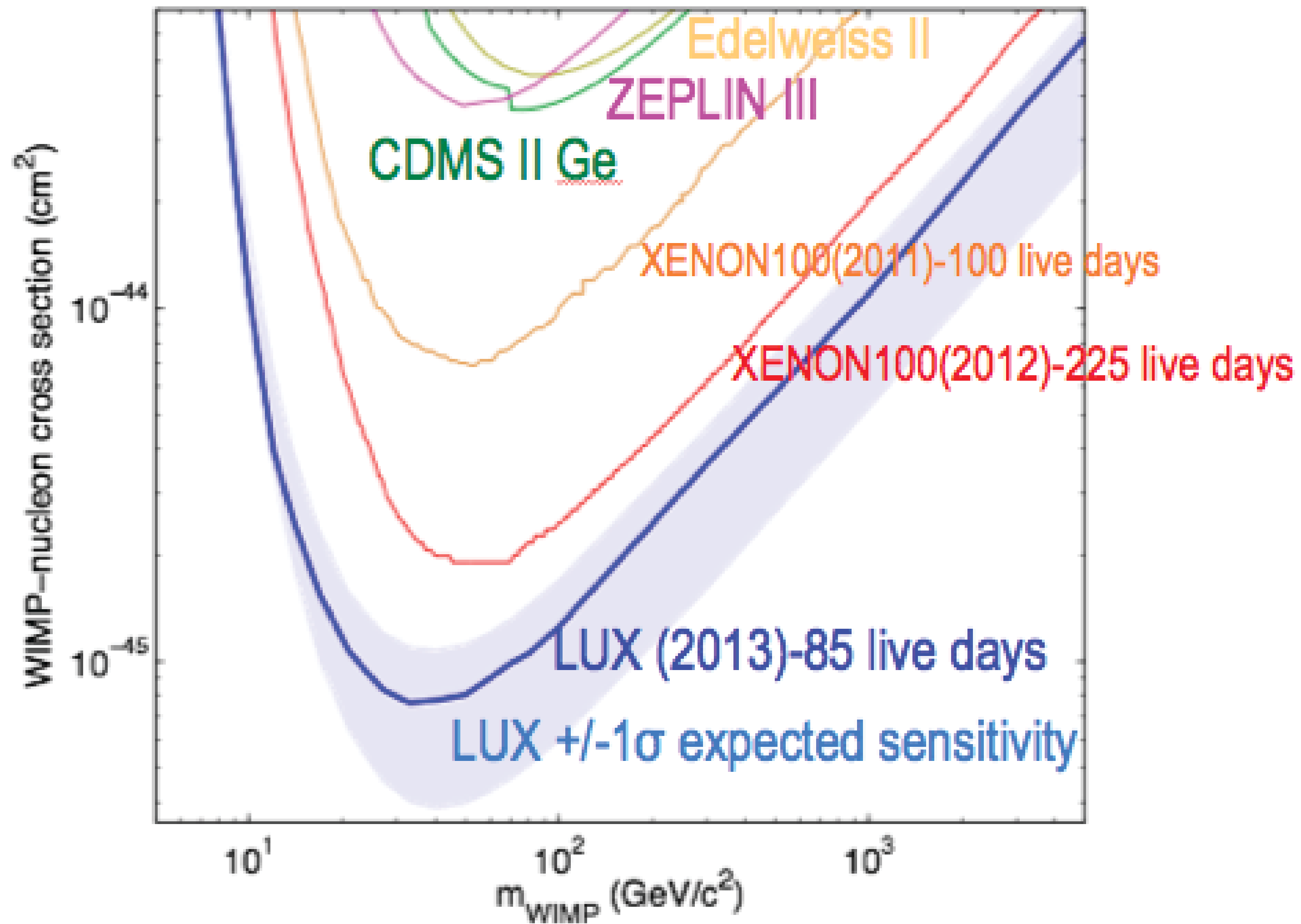
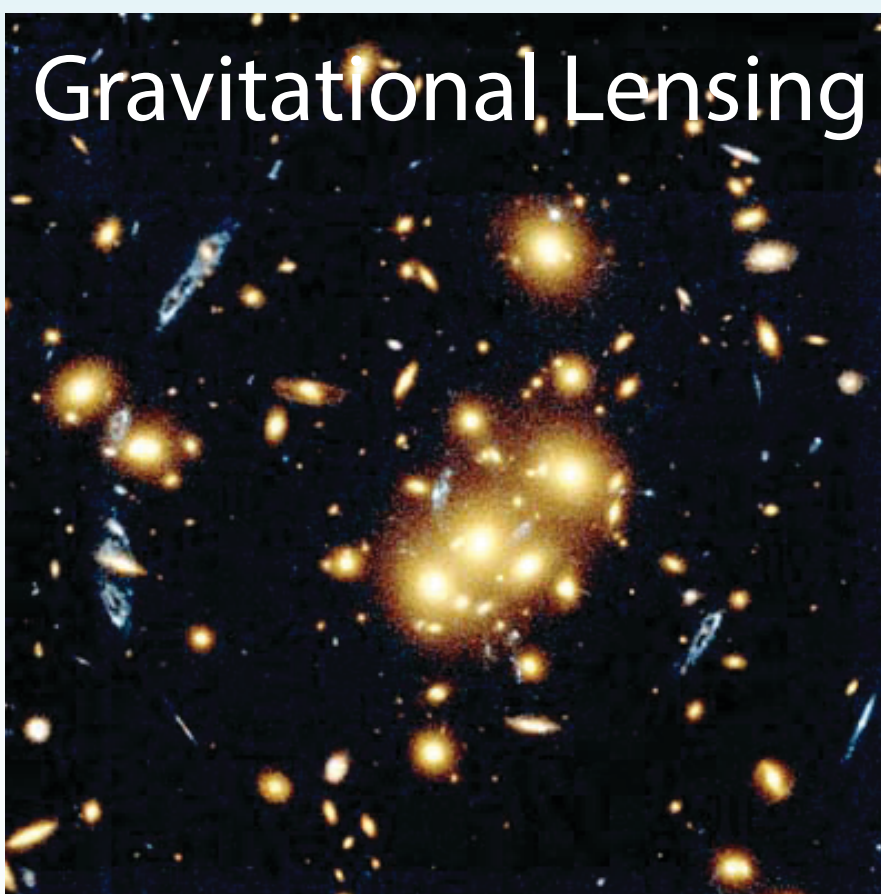
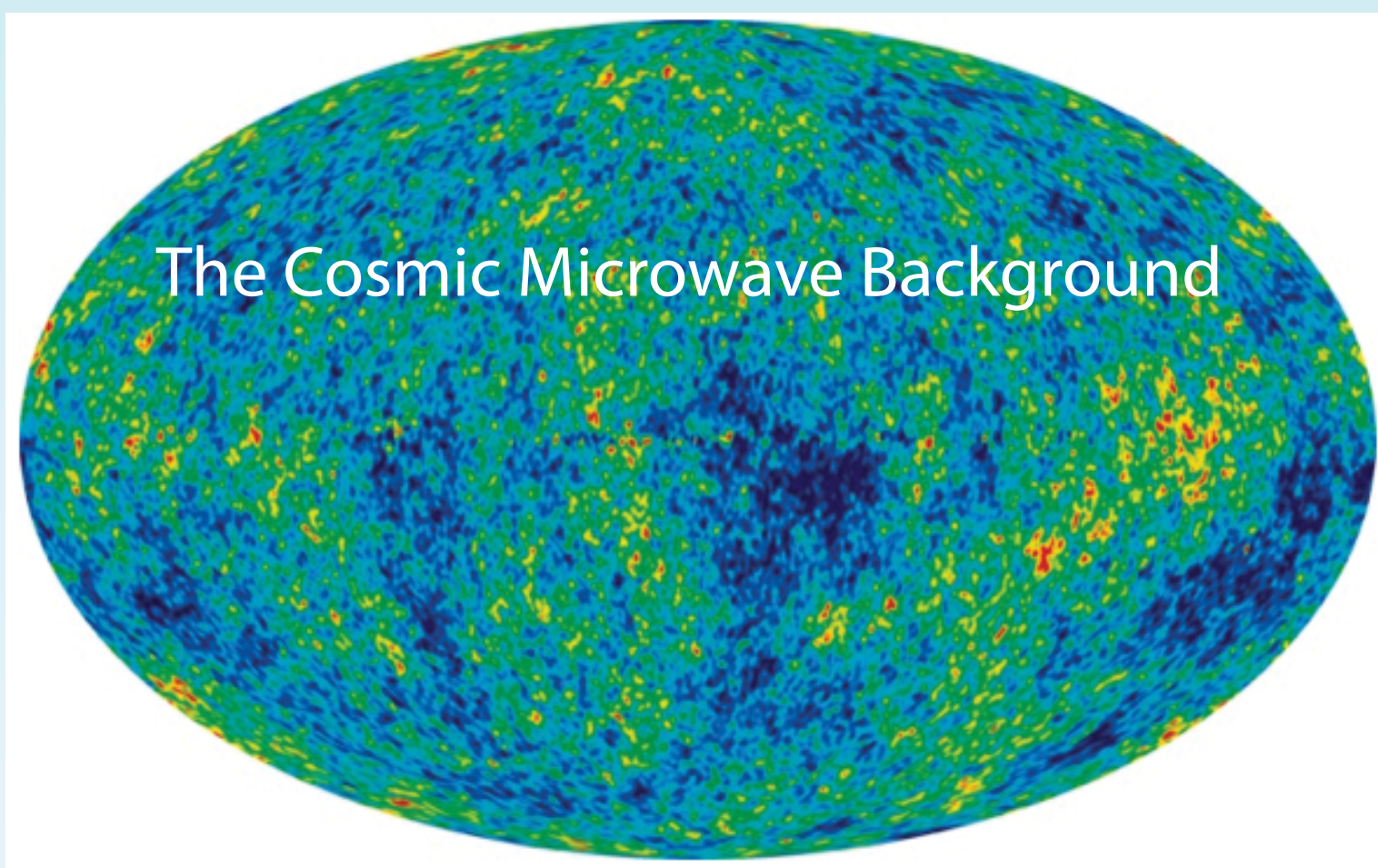
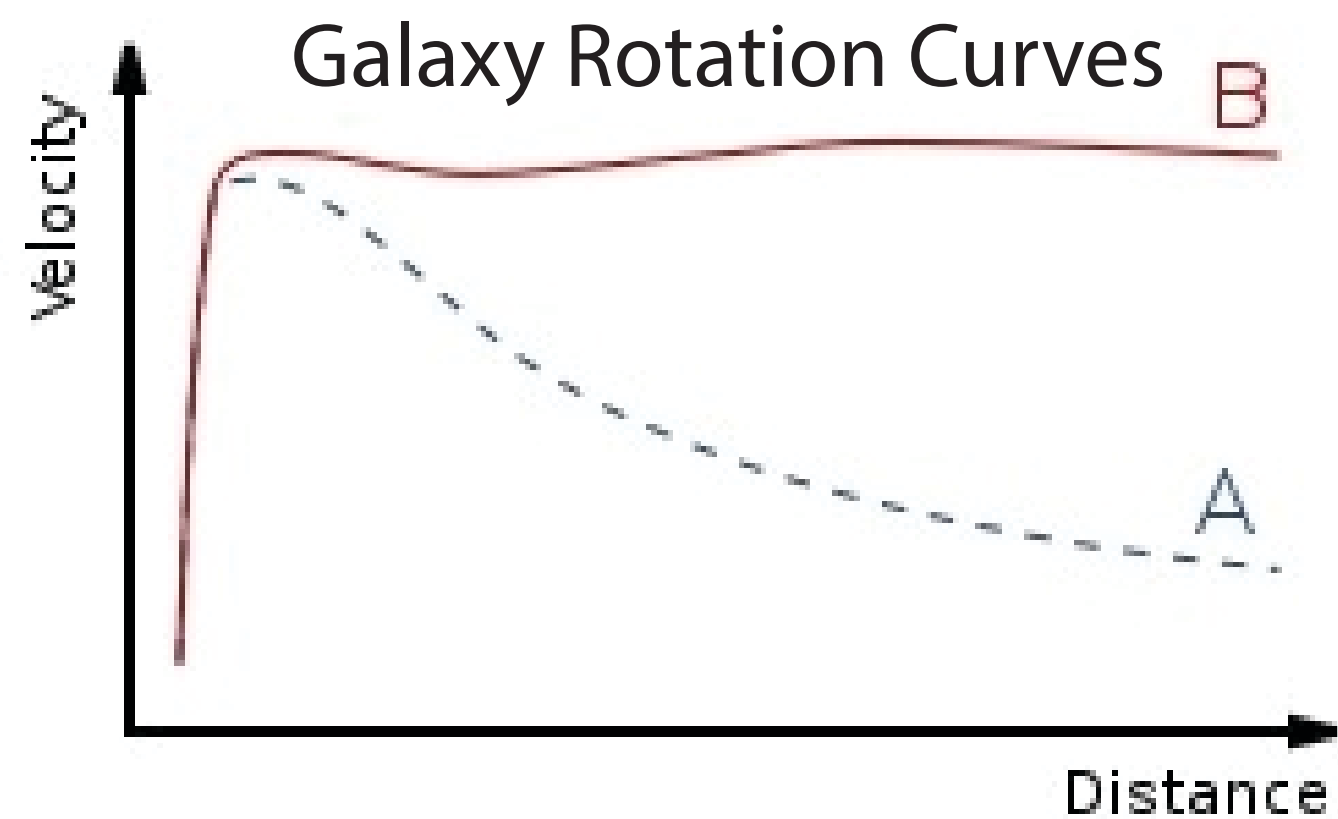
Emil Kirilov (PD) - Innsbruck Scientist
Ivan Kozyryev (U) - Harvard Graduate Student, Doyle Group
Max Parsons (U) - Harvard Graduate Student, Greiner Group
Yulia Gurevich (G) - Yale PD, Lamoreaux Group
Remi Louf (VS) - IPhT Scientist
Amar Vutha (G) - Professor, University of Toronto
Wes Campbell (G) - Professor, UCLA
Chris Overstreet (U) - Stanford Graduate Student, Kasevich Group
Nick Hutzler (G) - Harvard Postdoc, Ni Group (and ACME collaborator)
Ben Spaun (G) - JILA Postdoc, Ye Group
Paul Hess (G) - JQI Postdoc, Monroe Group



Standard Model



Observed Dark Matter — UNEXPLAINED



Direct Detection — so far, no dice

Matter/AntiMatter Asymmetry — UNEXPLAINED

- One key requirement: CP Violation
- CP Violation exists in SM
- Not enough SM ~~CP~~ to explain our Universe

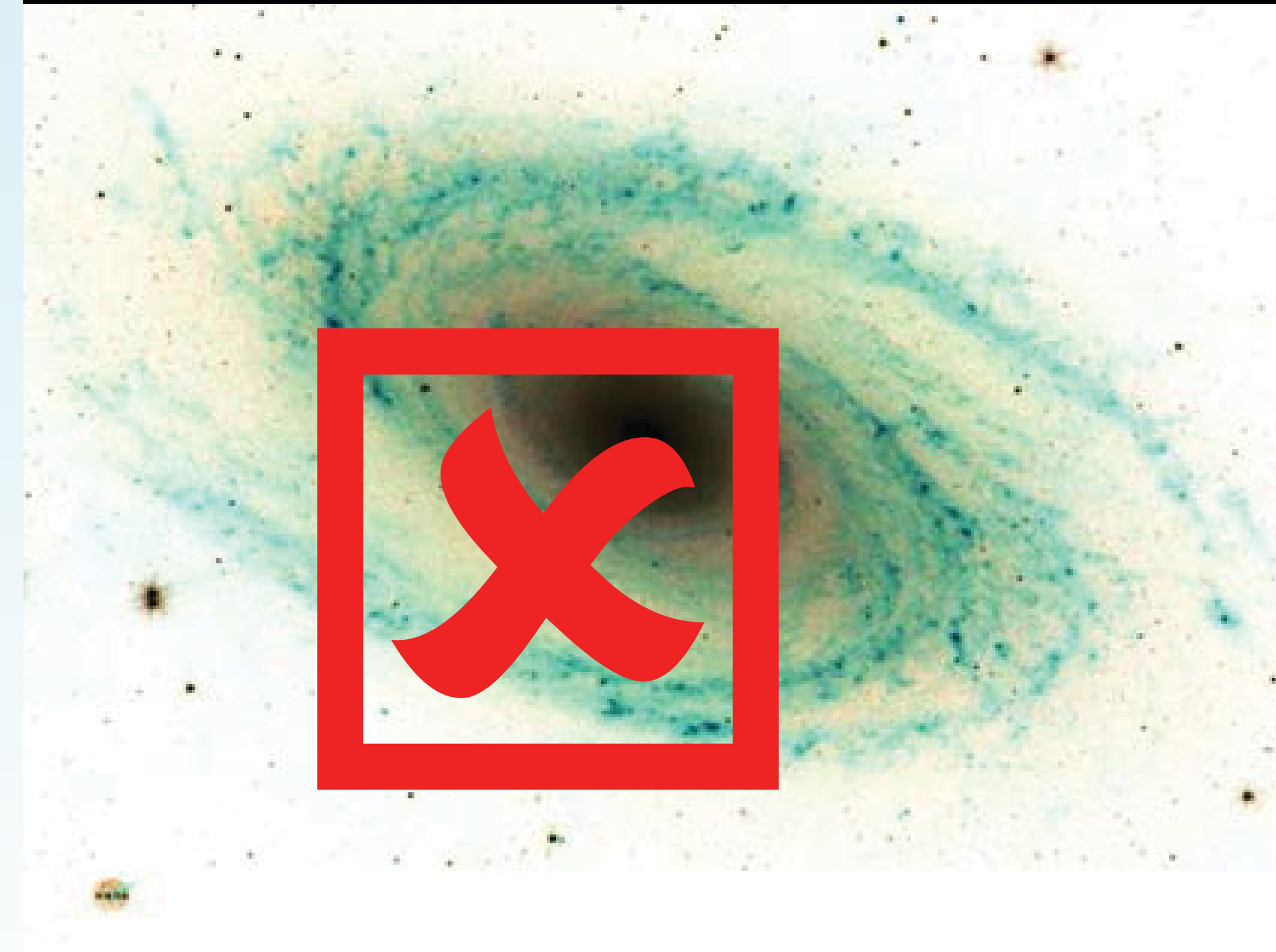
Observed M/AM asymmetry

REQUIRES T-violation

and

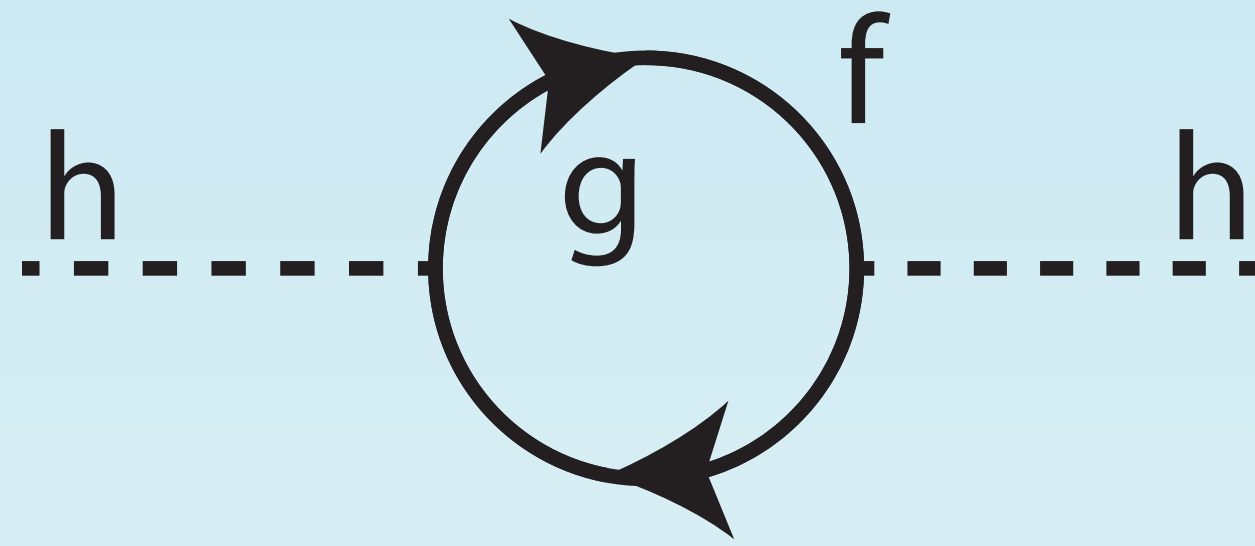
REQUIRES new sources beyond the SM

A. D. Sakharov, JETP Lett. 5, 27 (1967)



Naturalness of the Higgs Mass — *Uncomfortable*

Problem



$$m_h \approx \frac{g}{4\pi} \Lambda$$

$$m_h = 125 \text{ GeV}$$

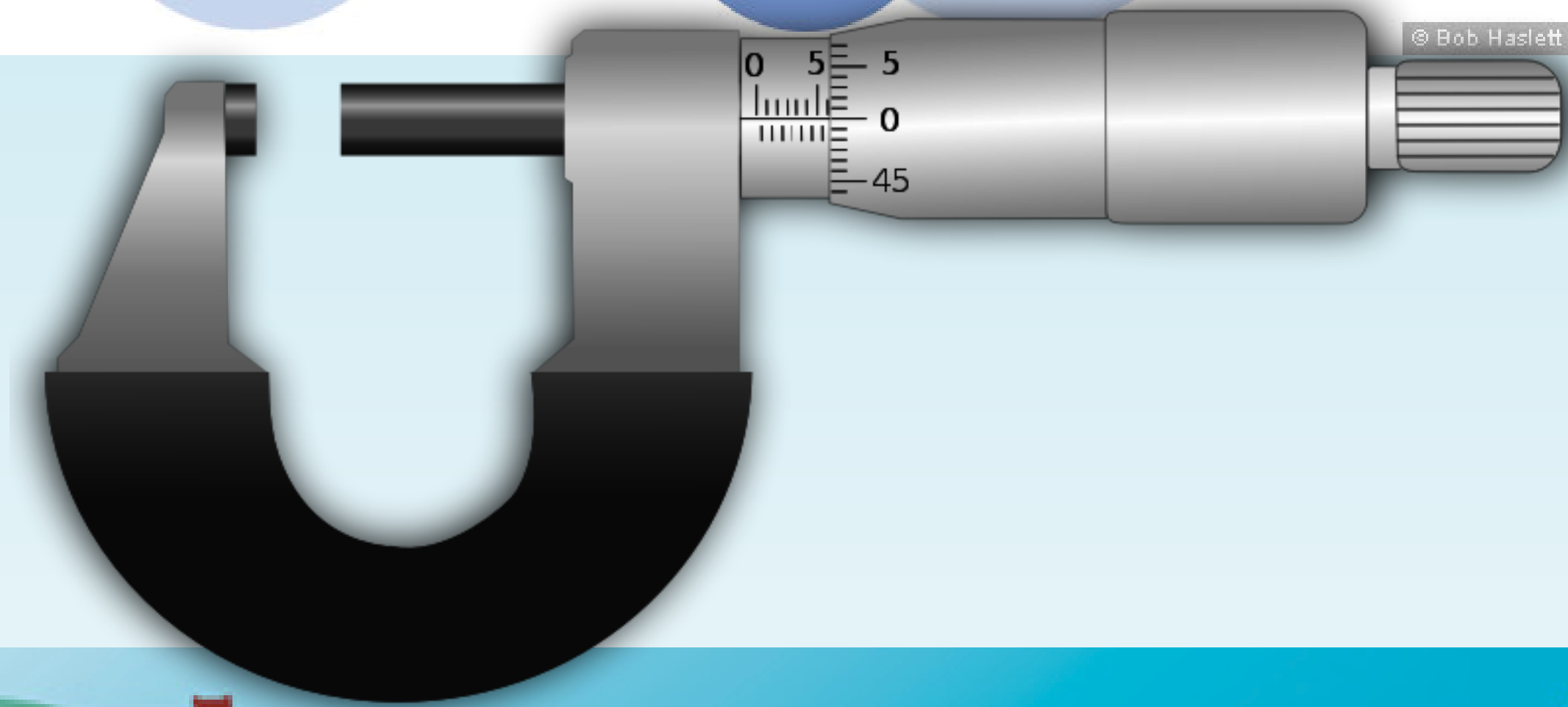
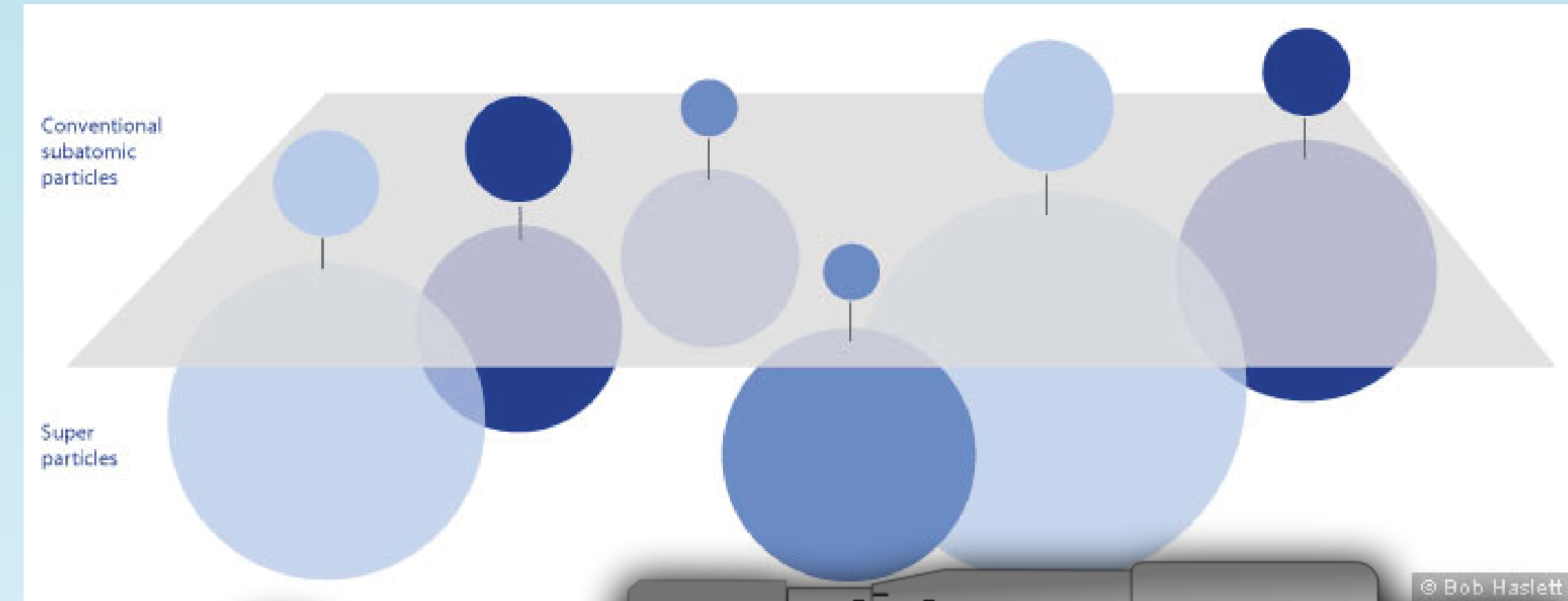
The Higgs Mass² is Quadratically Divergent to Radiative Corrections
So “Naturally” lies near **Lambda**,
UV (New Physics) cutoff

Super
Particles

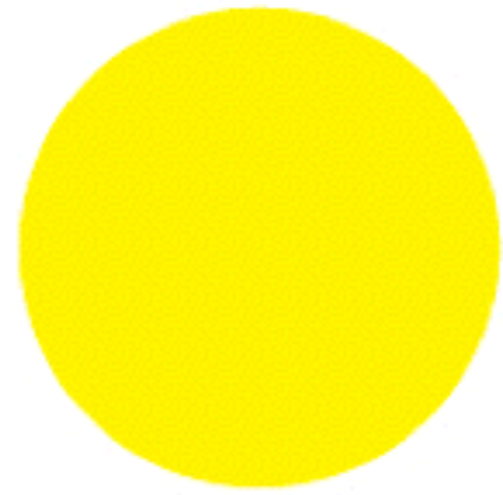
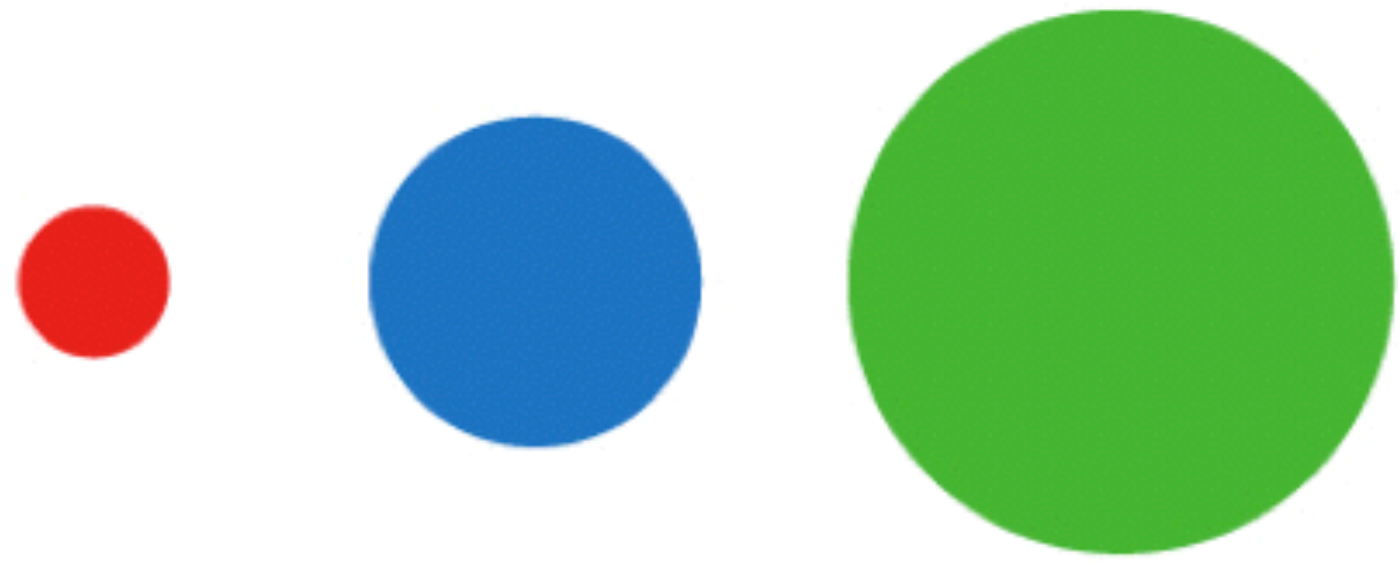
Fine
Tuning

Extra
Dimensions

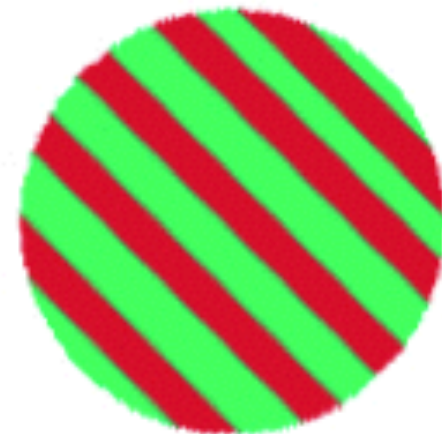
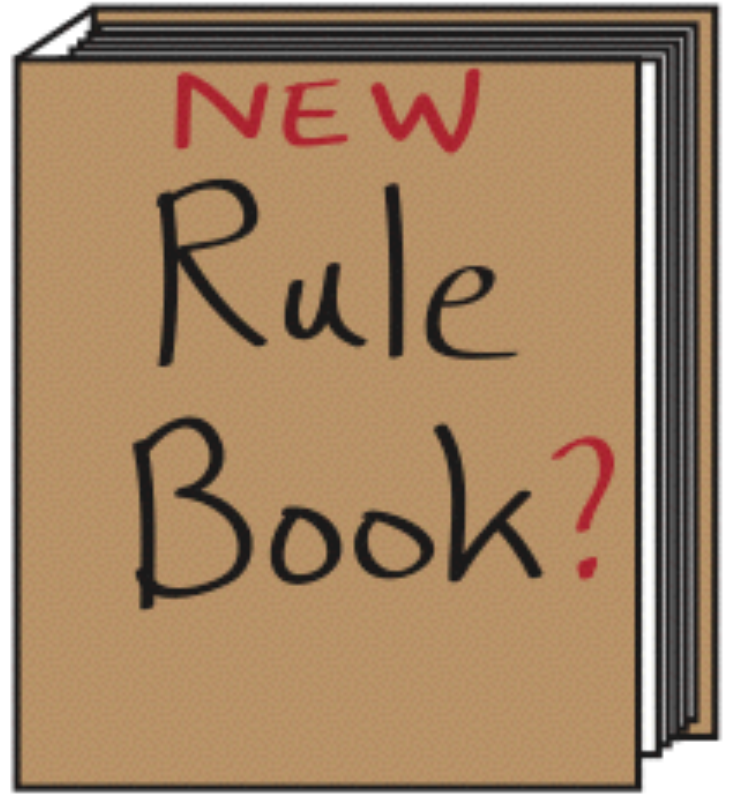
Possible Solutions



Possible Solution

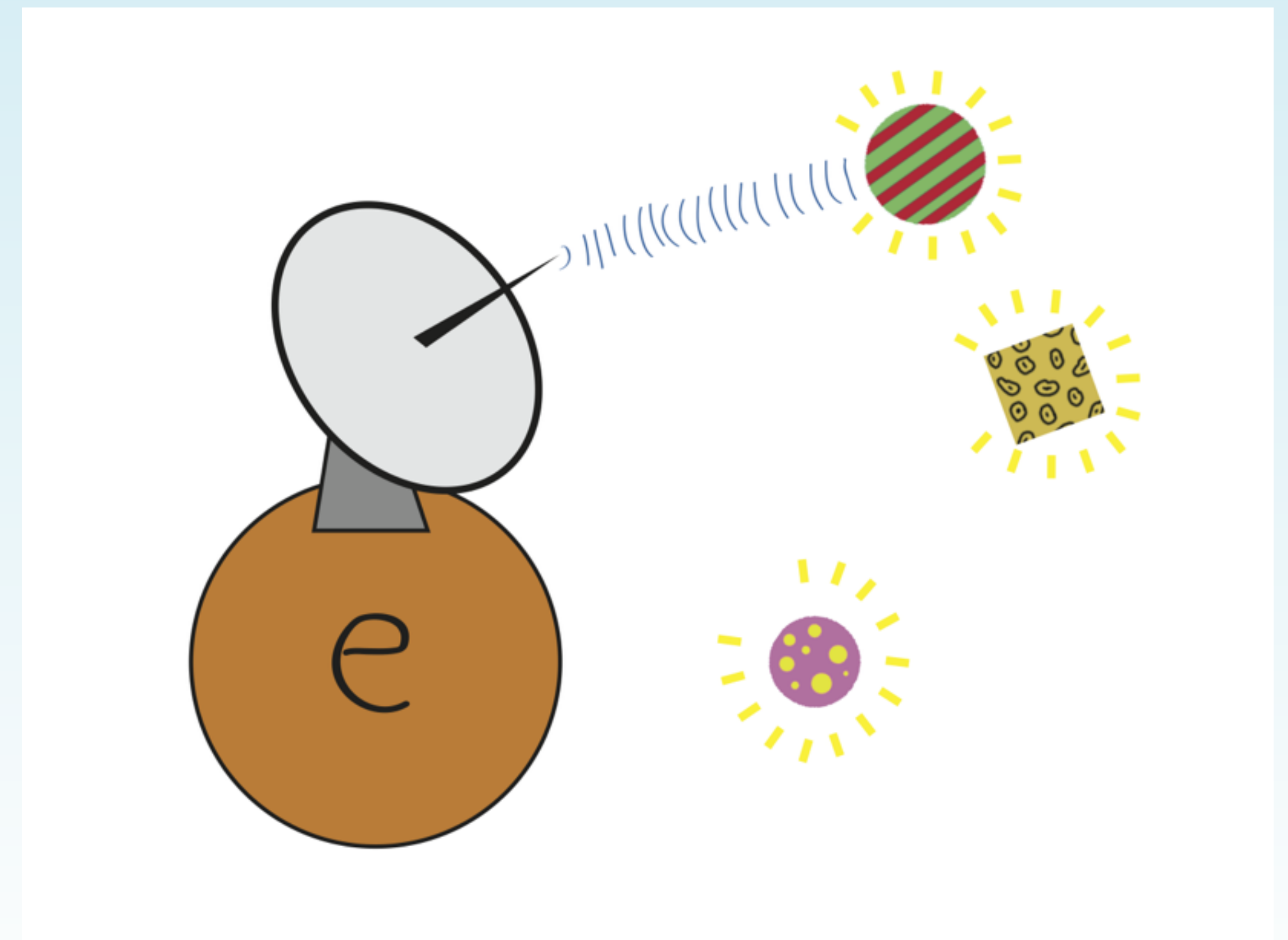
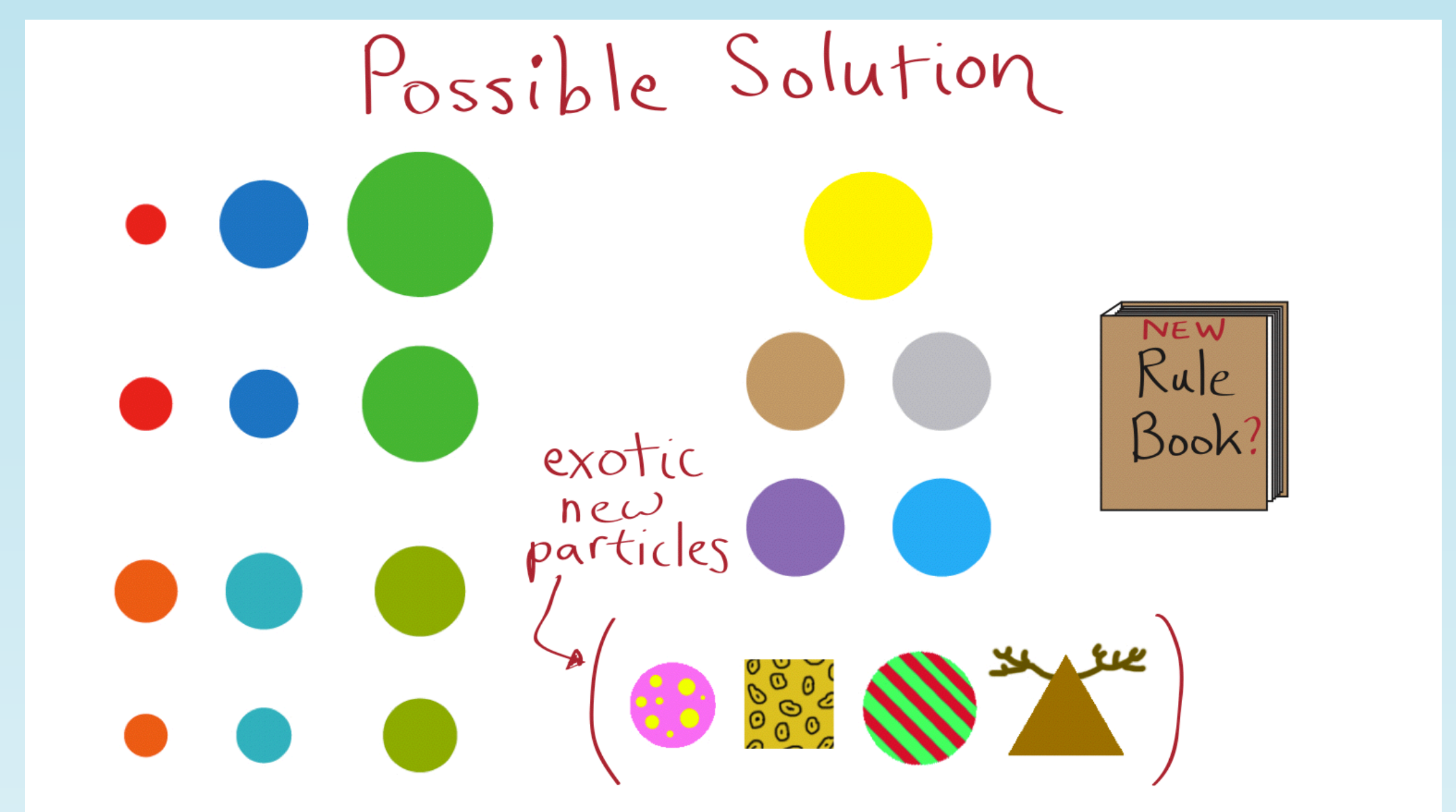


exotic
new
particles

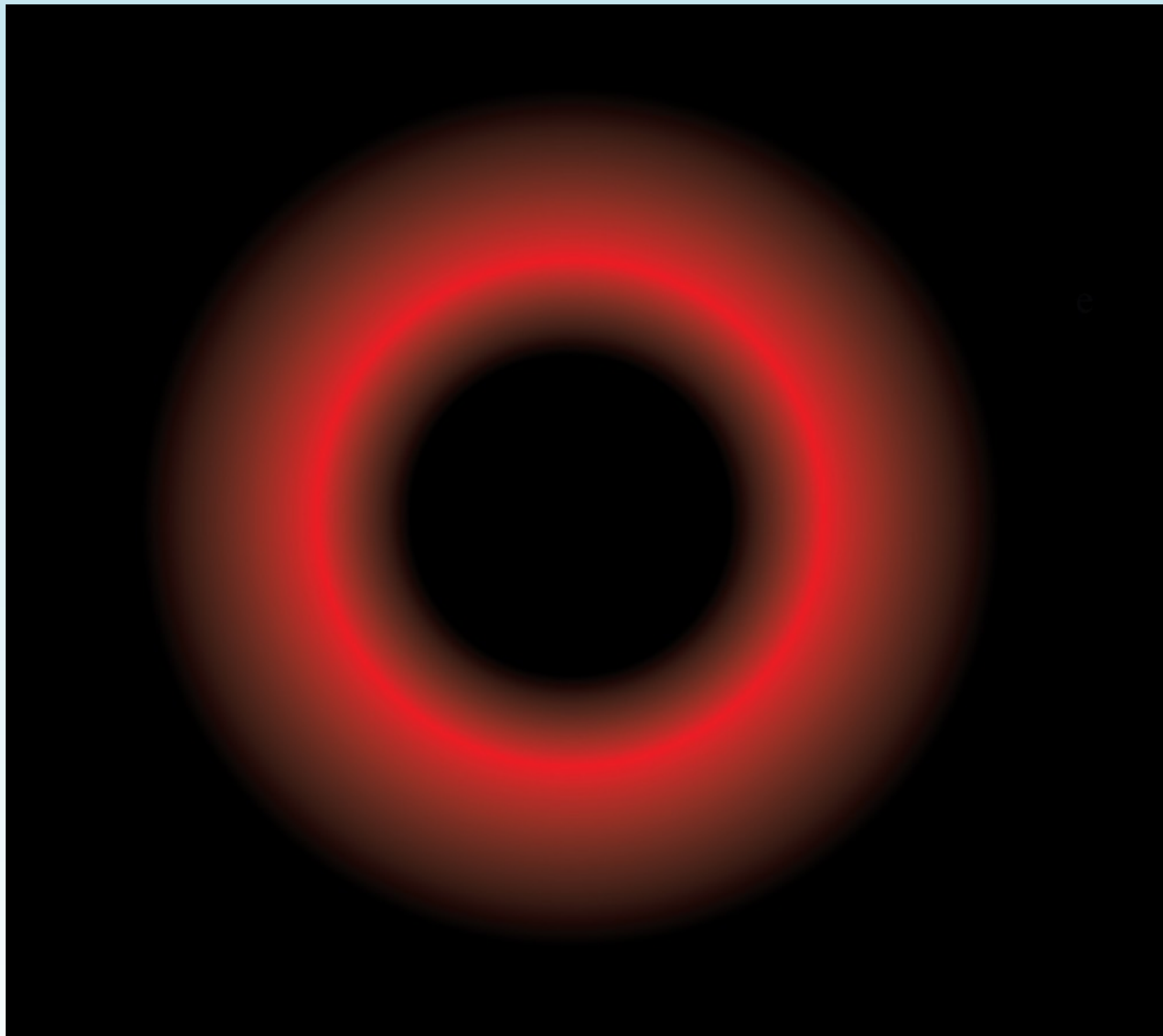


Several SUSY Theories Solve all Problems

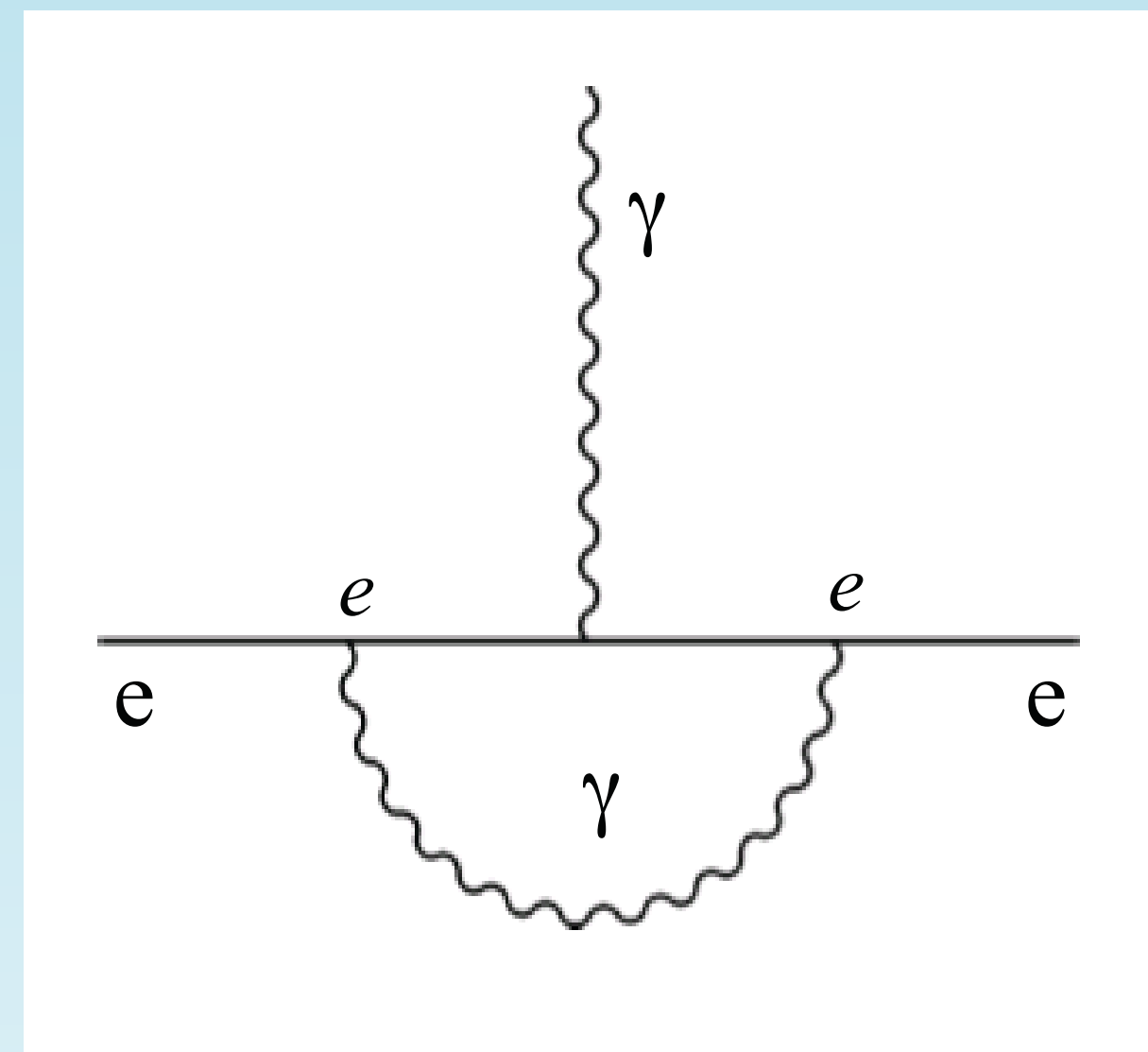
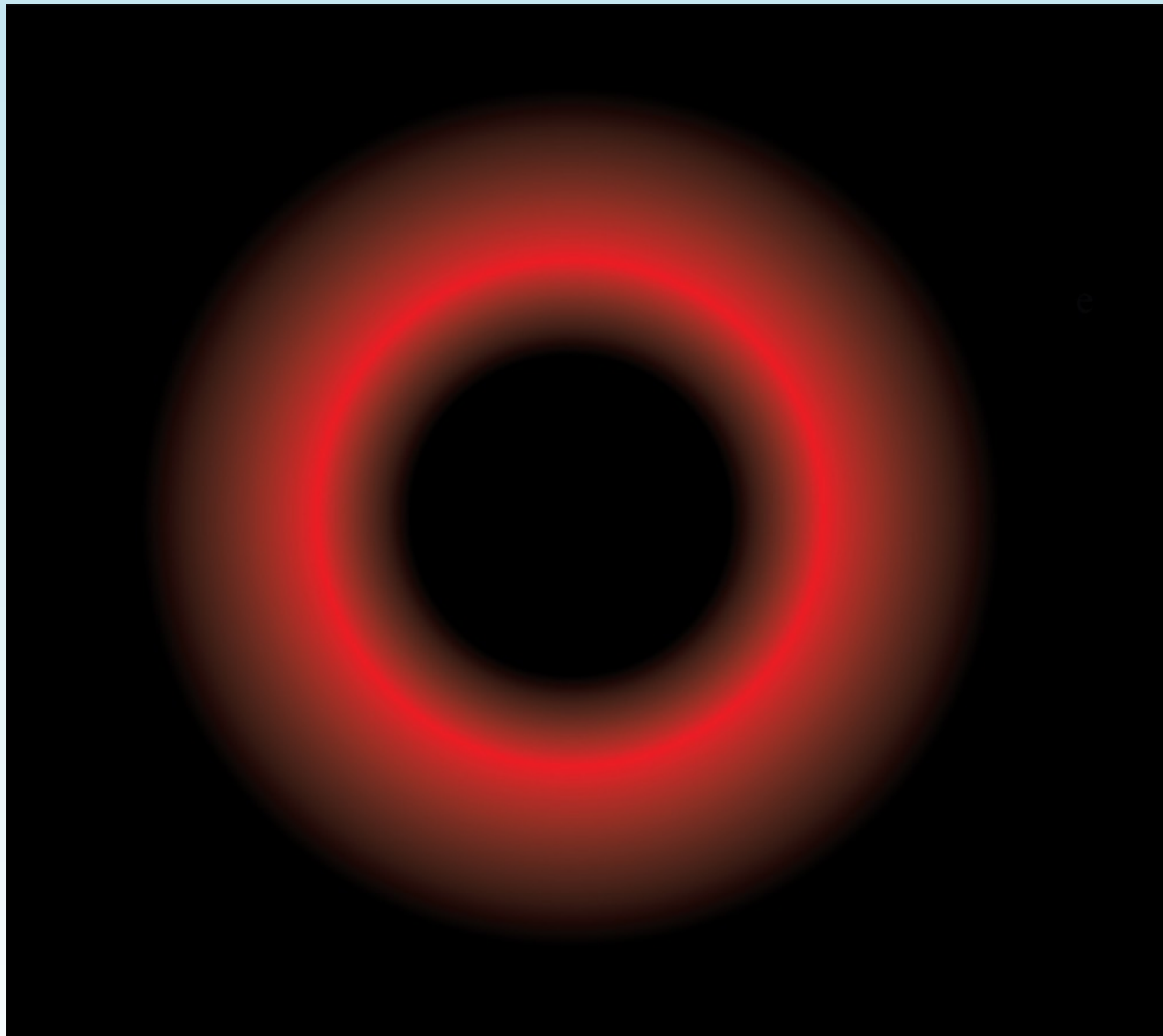
- ★ Matter/Antimatter
 - naturally provides needed T-violation
- ★ Unification/Hierarchy
 - provides needed particles
- ★ Dark Matter
 - provides candidate particle



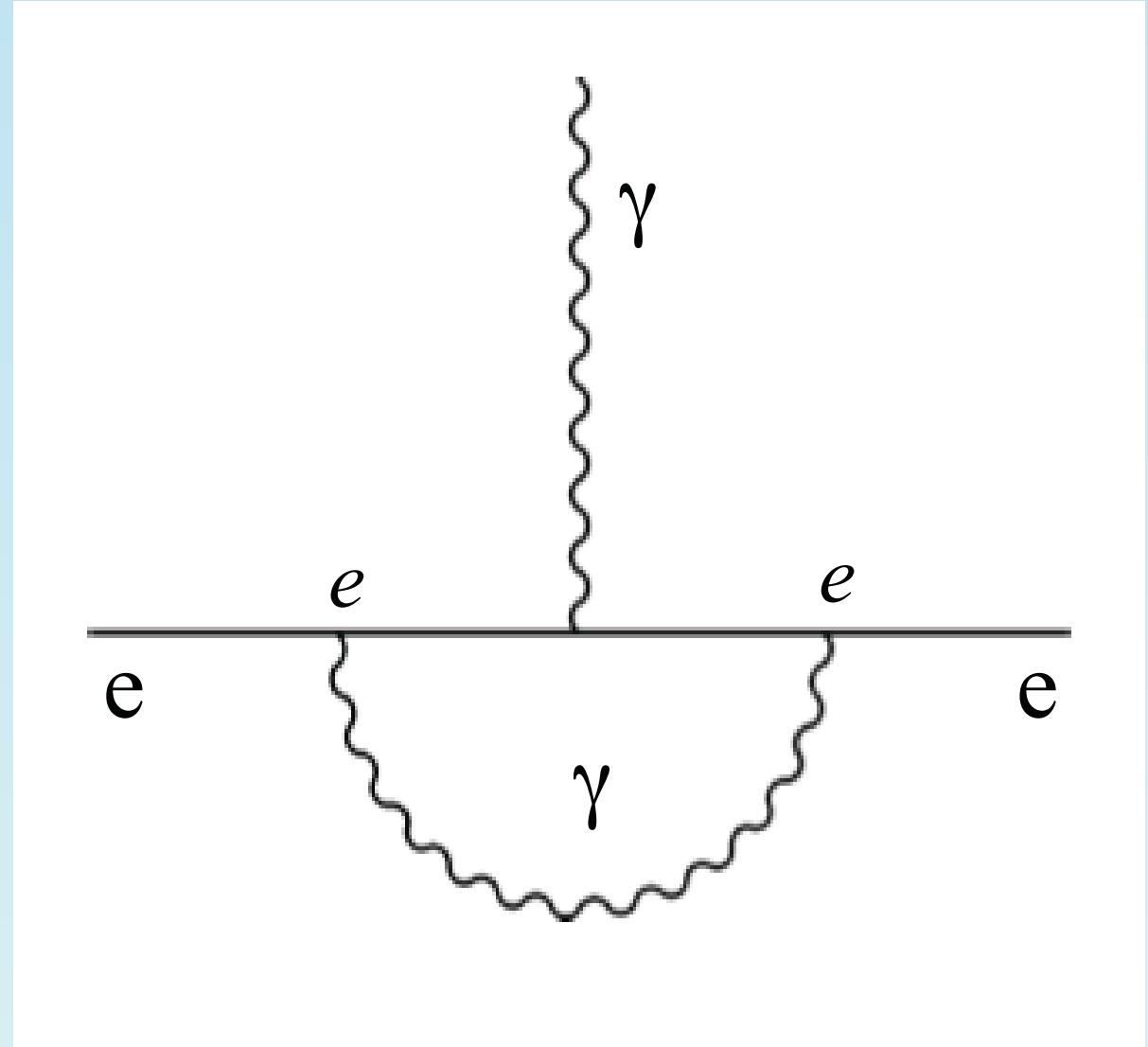
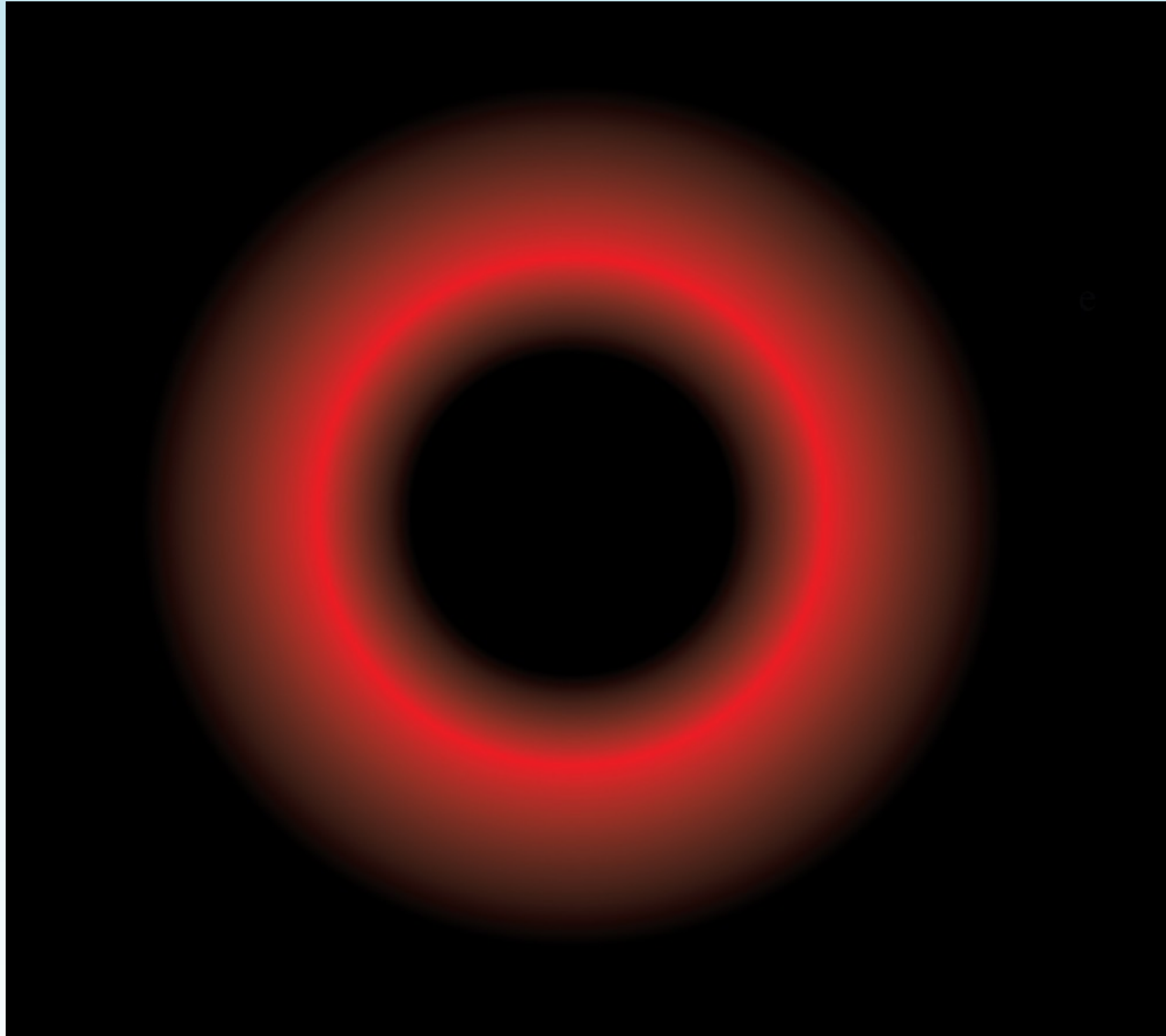
Electron is dressed by Virtual Particles — g-2



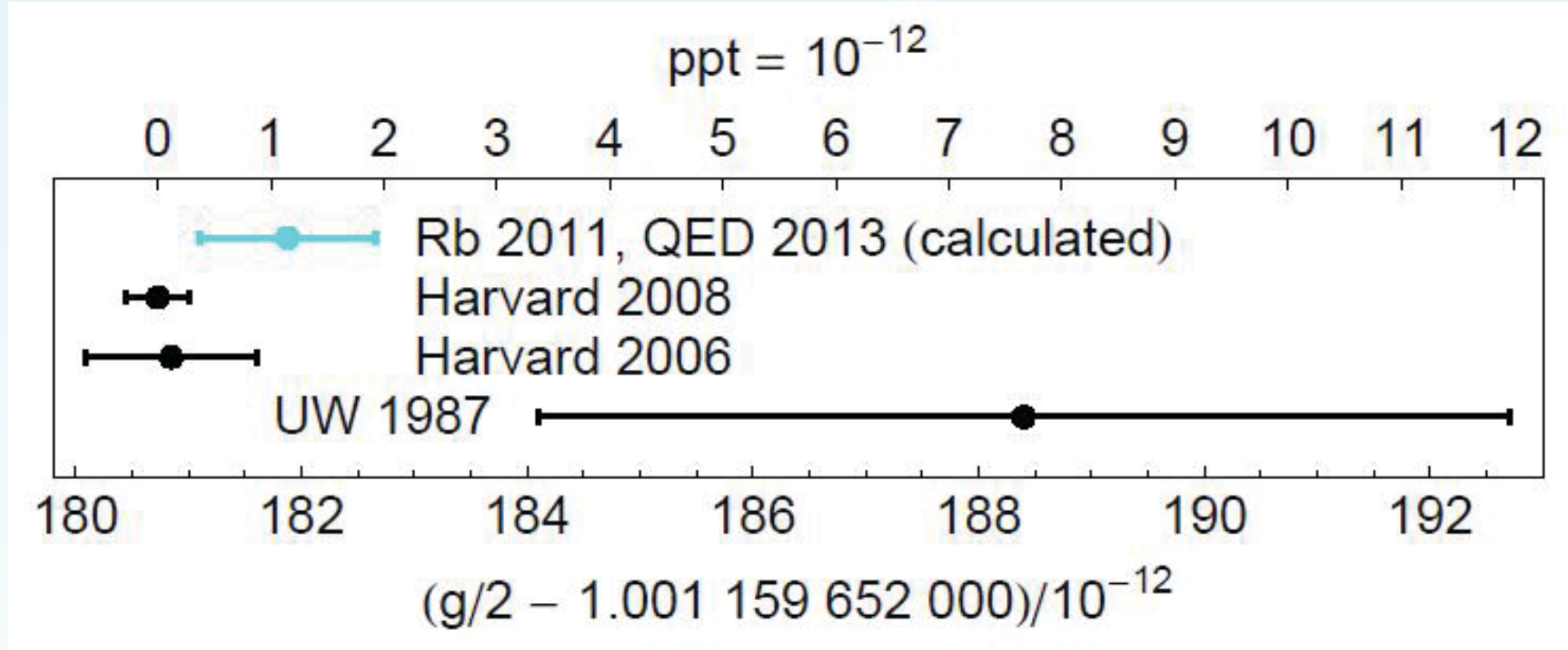
Electron is dressed by Virtual Particles — g-2



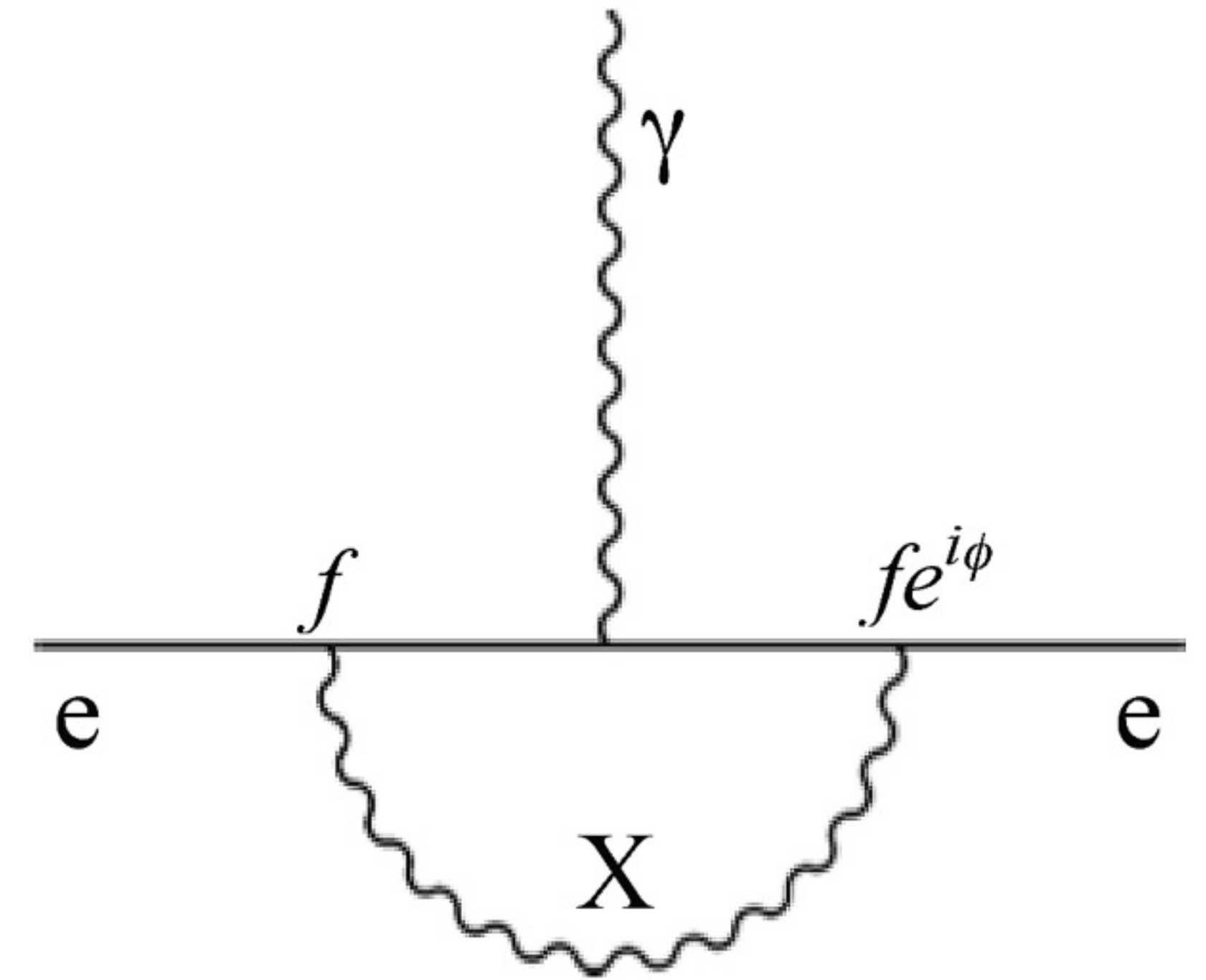
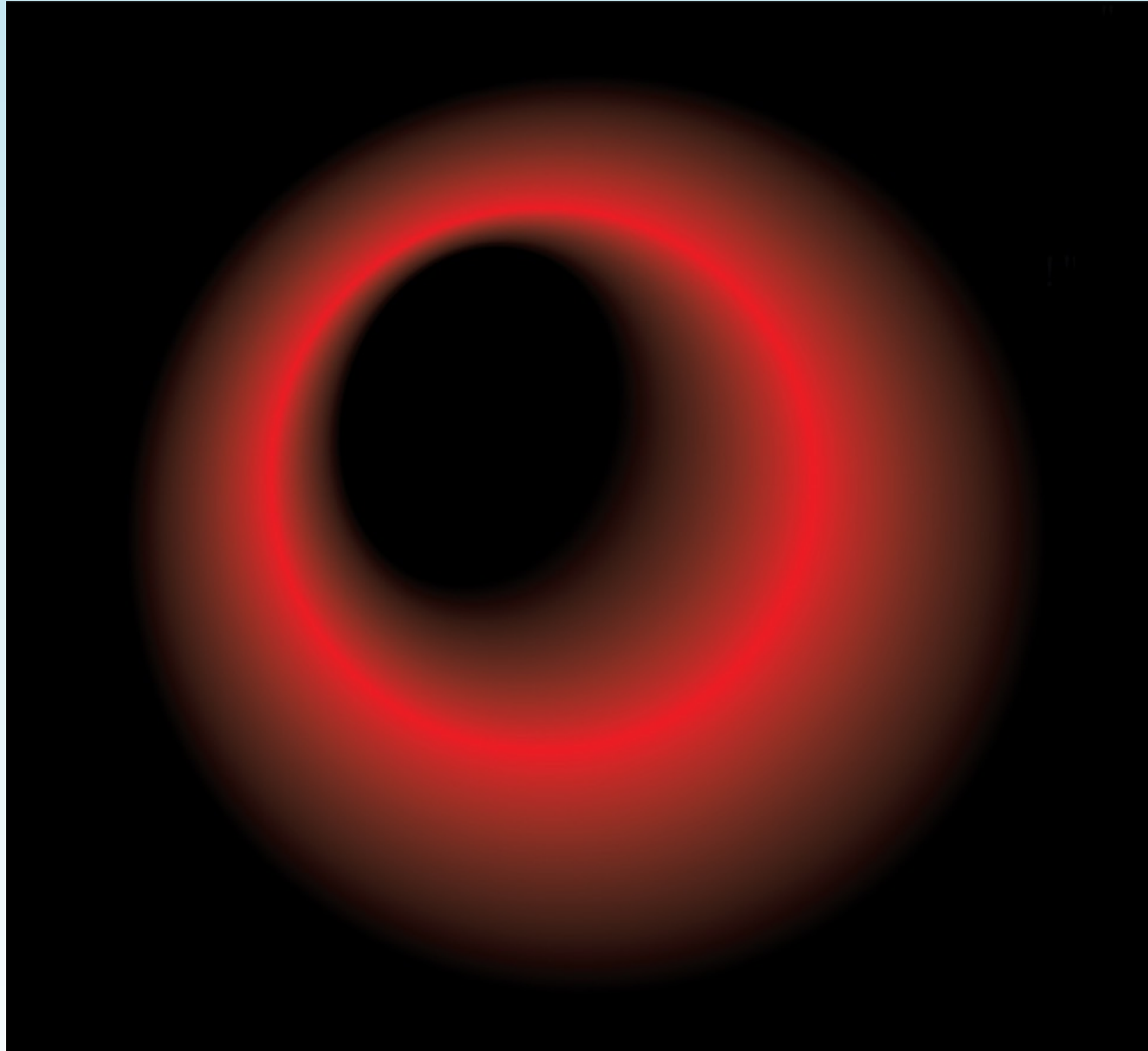
Electron is dressed by Virtual Particles — g-2



Gabriesle group
confirms SM QED
ppt



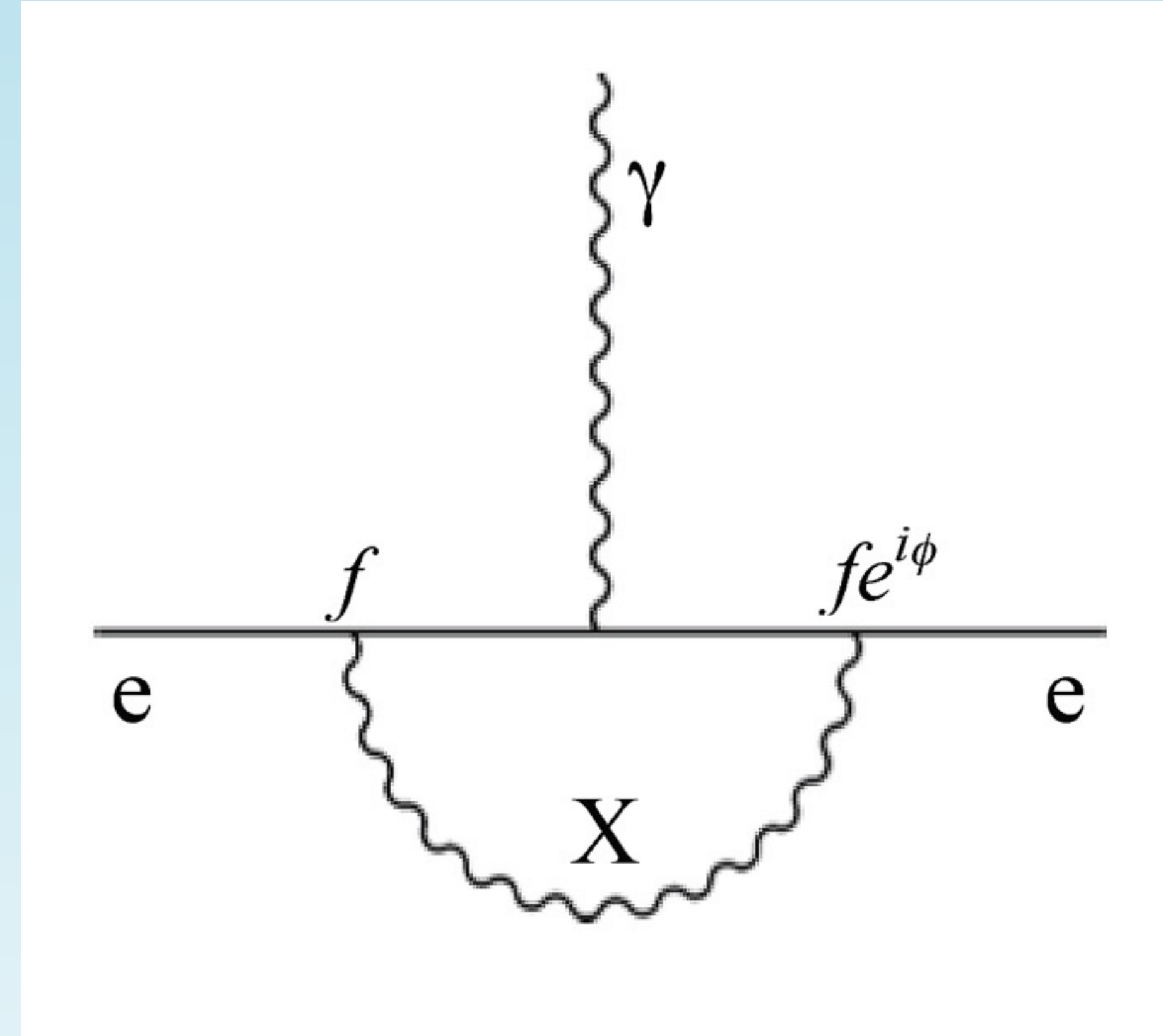
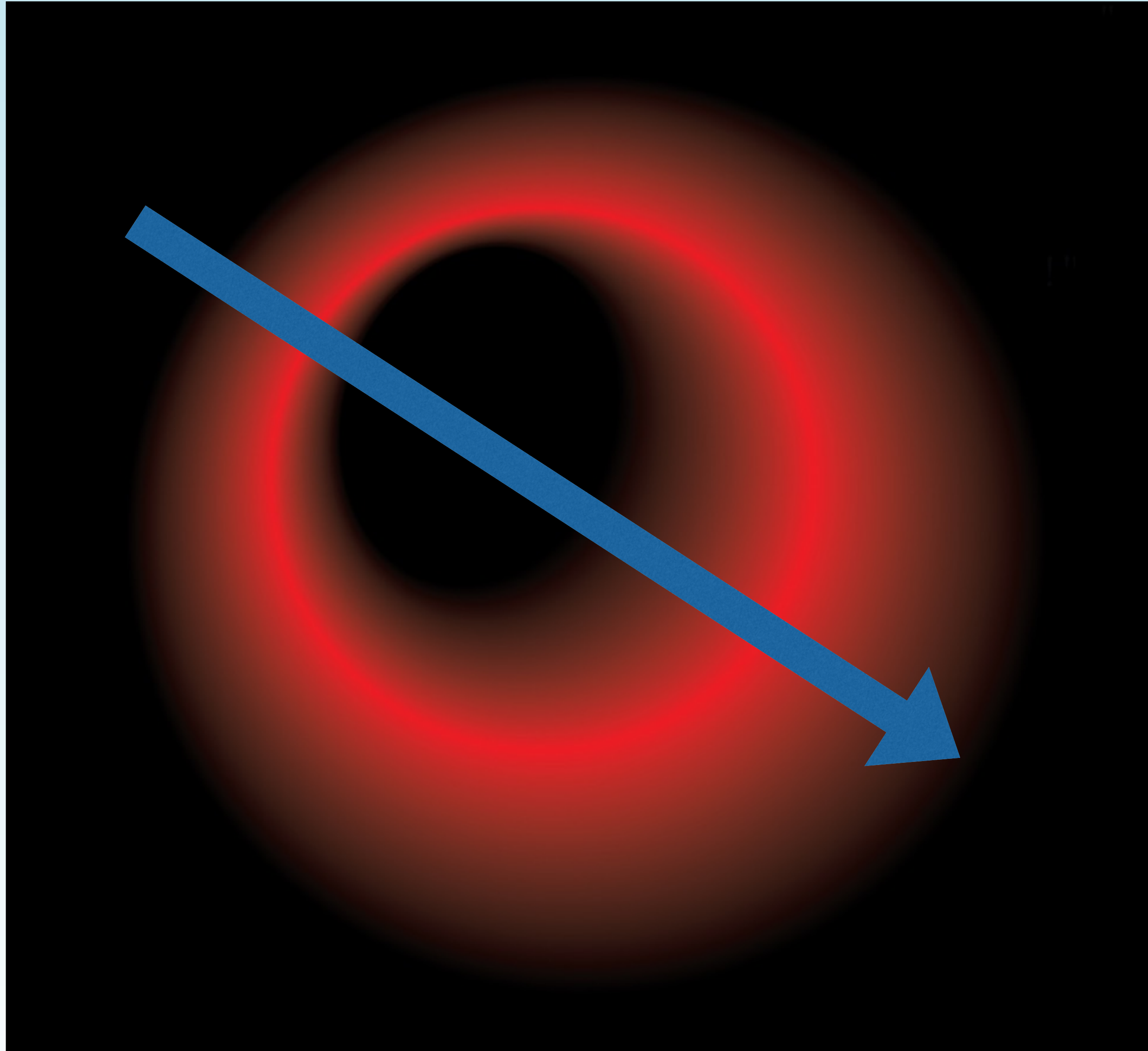
Electron is dressed by Virtual Particles — New Physics



SUSY

1st order perturbation
cancellations not inherent
T-violating phase natural

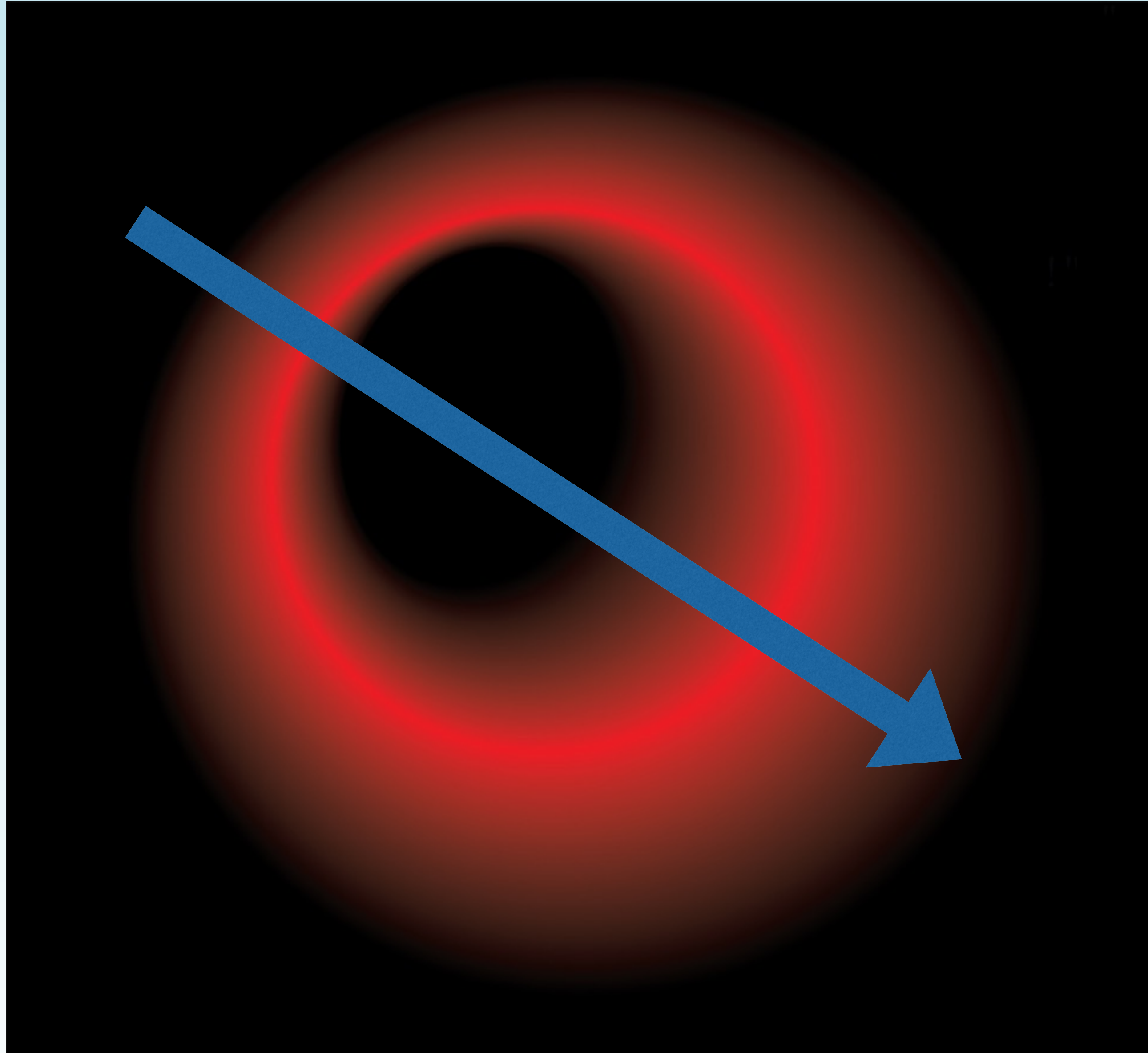
Electron is dressed by Virtual Particles — New Physics



SUSY

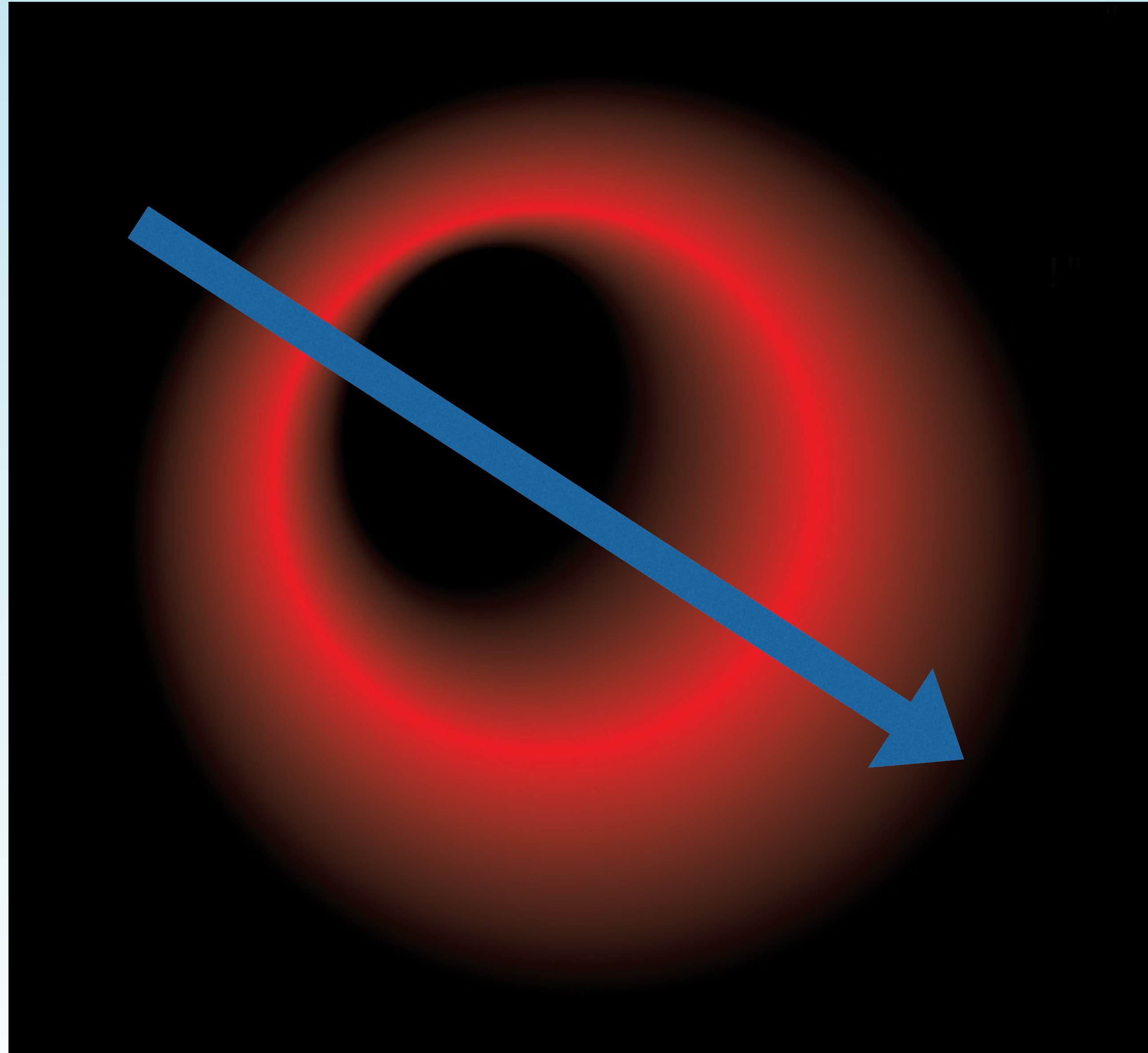
1st order perturbation
cancellations not inherent
T-violating phase natural

Electron is dressed by Virtual Particles — New Physics



EDM is inherently T-violating

Electron is dressed by Virtual Particles — New Physics



EDM is inherently T-violating

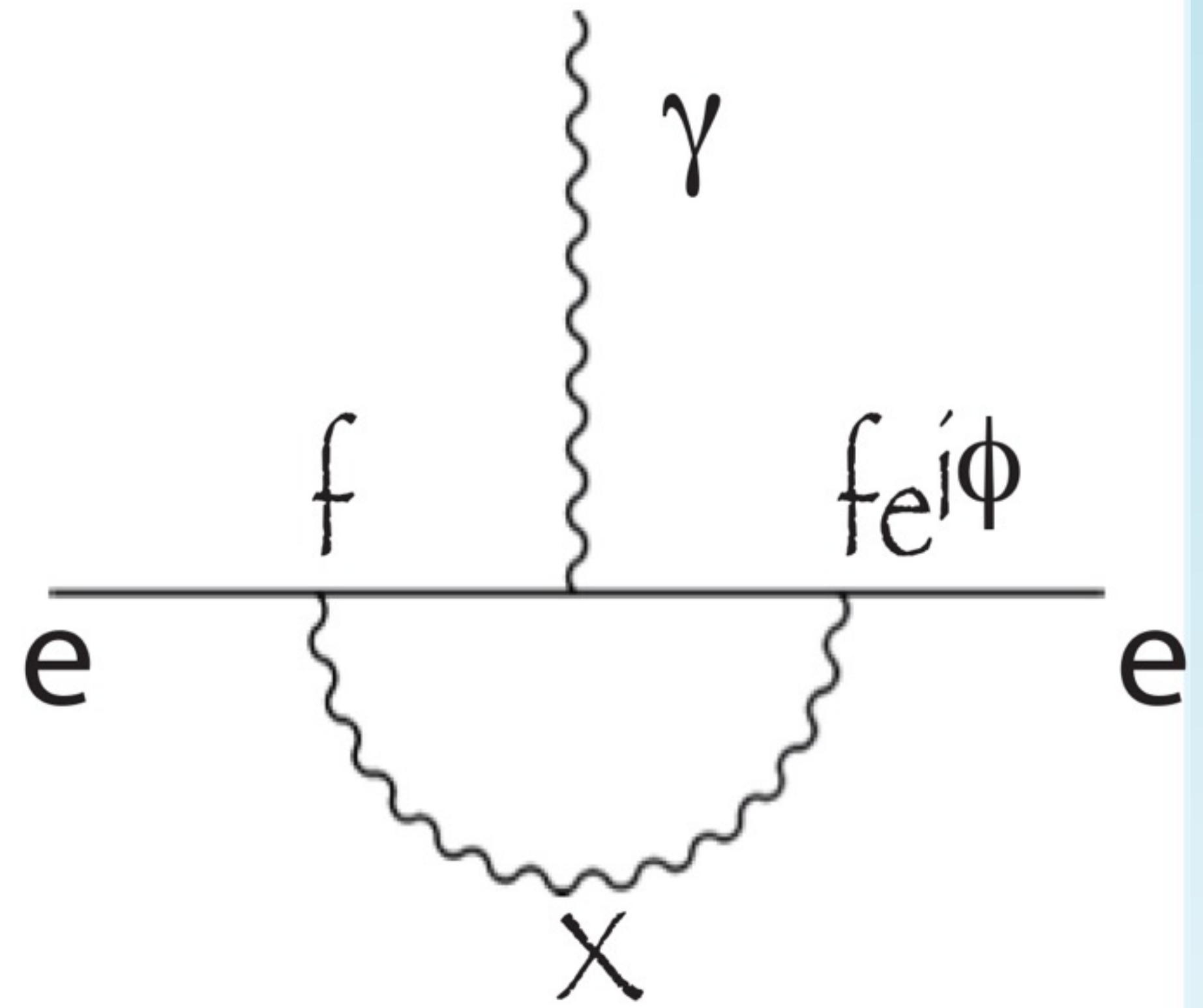
EDM in our experimental range would have enough T-violation to explain Matter/Antimatter asymmetry, indicate new 1-1000 TeV particle

EDM Too Small??

Assume $f^2/hc \approx \alpha$

$\sin(\Phi) \approx 1$

$m_\chi \approx 100 \text{ GeV}$



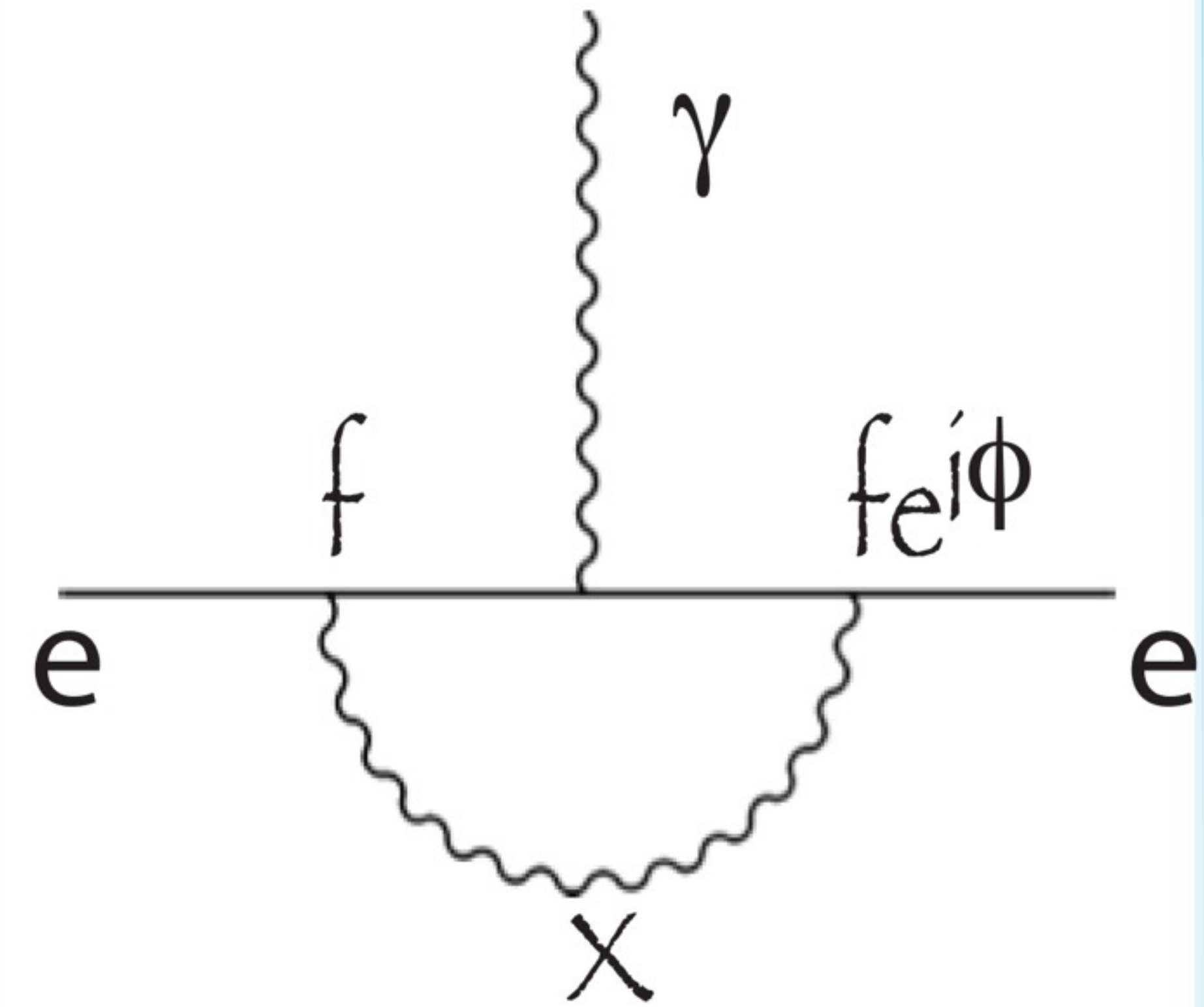
EDM Too Small??

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$$\text{EDM} \approx \mu_B (\alpha/\pi)^N (m_e/m_\chi)^2 \sin(\Phi)$$



EDM Too Small??

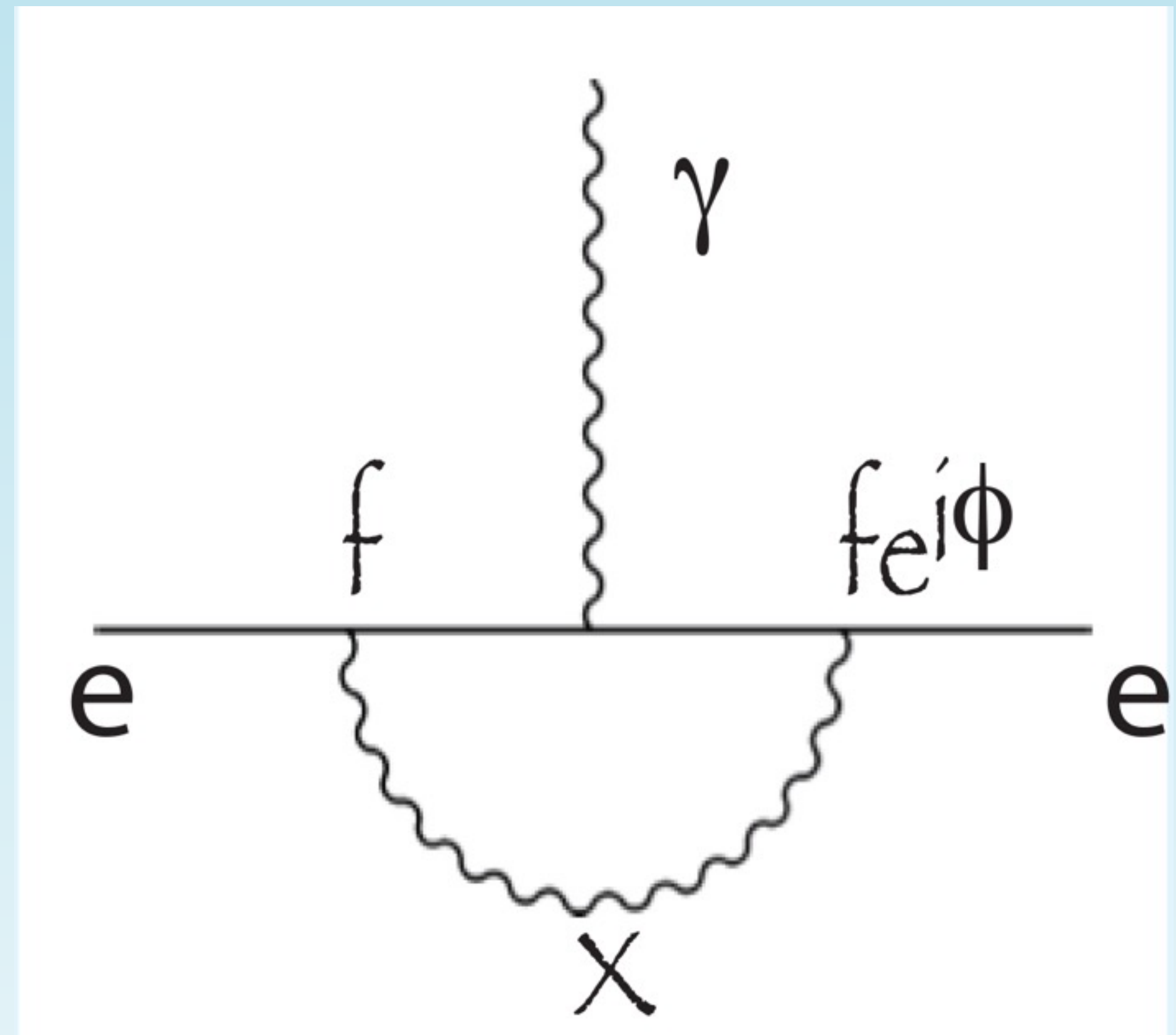
Assume $f^2/hc \approx \alpha$

$$\sin(\Phi) \approx 1$$

$$m_\chi \approx 100 \text{ GeV}$$

$$\text{EDM} \approx \mu_B (\alpha/\pi)^{\overset{\text{number of loops}}{\uparrow} N} (m_e/m_\chi)^2 \sin(\Phi)$$

$$\text{EDM} \approx 10^{-25} \text{ e cm}$$



EDM Too Small??

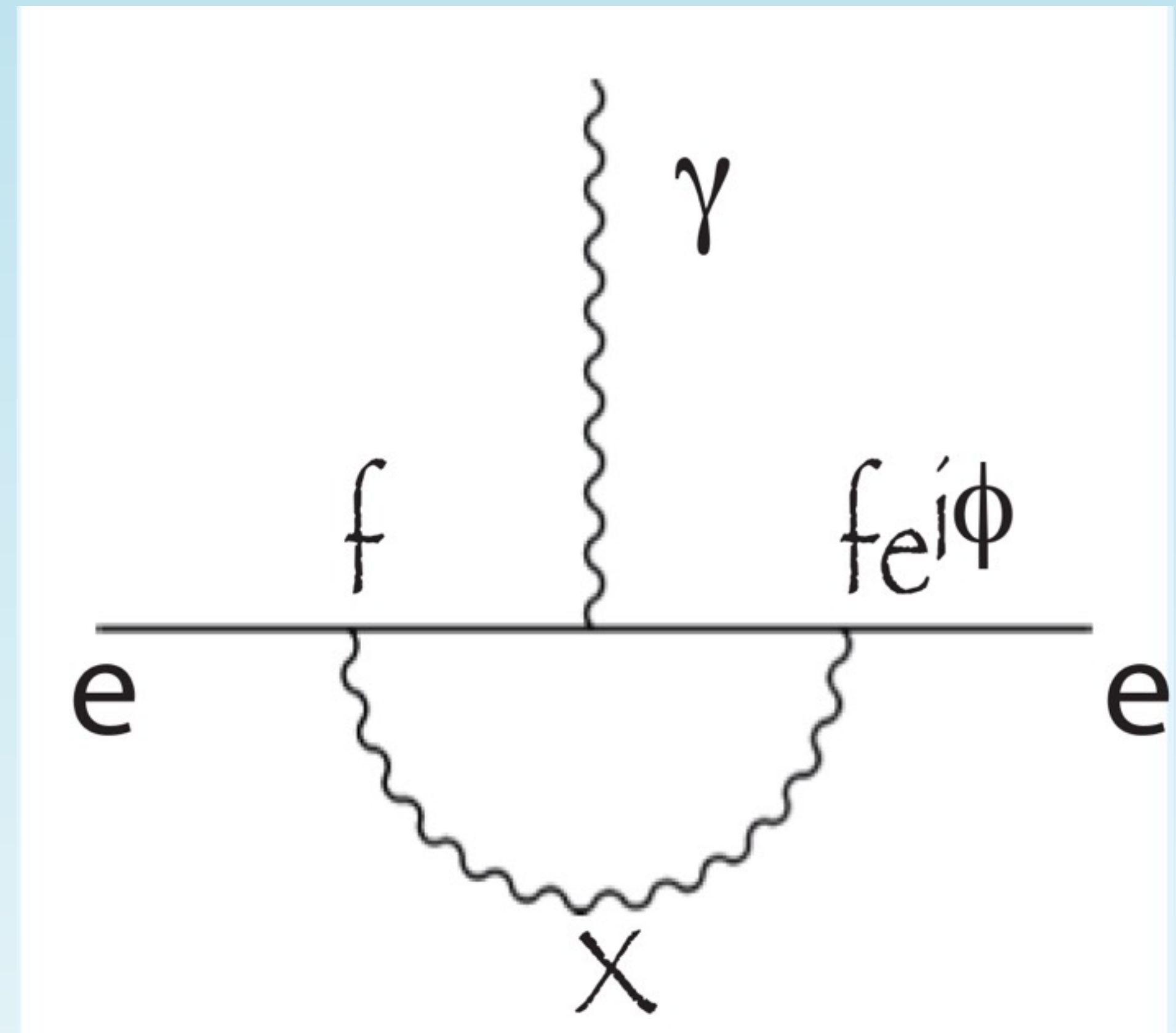
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EDM $\approx 100x$ previous limit



EDM Too Small??

Assume $f^2/hc \approx \alpha$

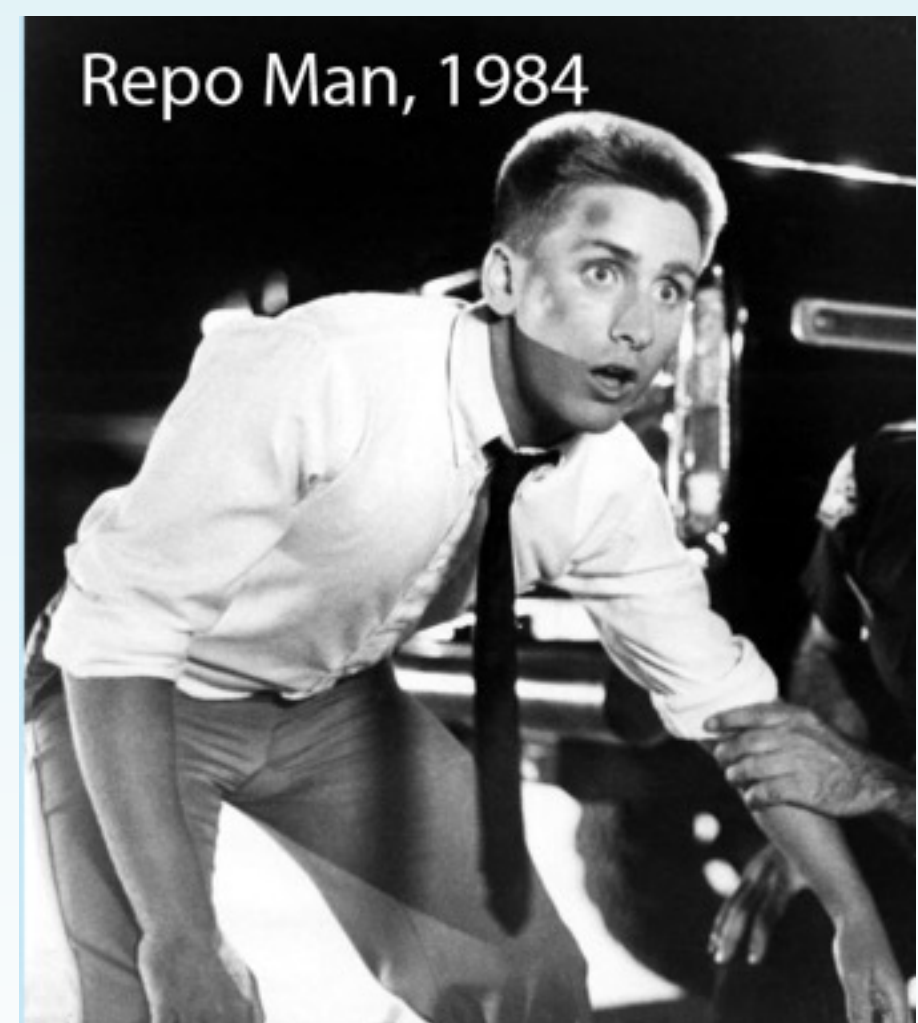
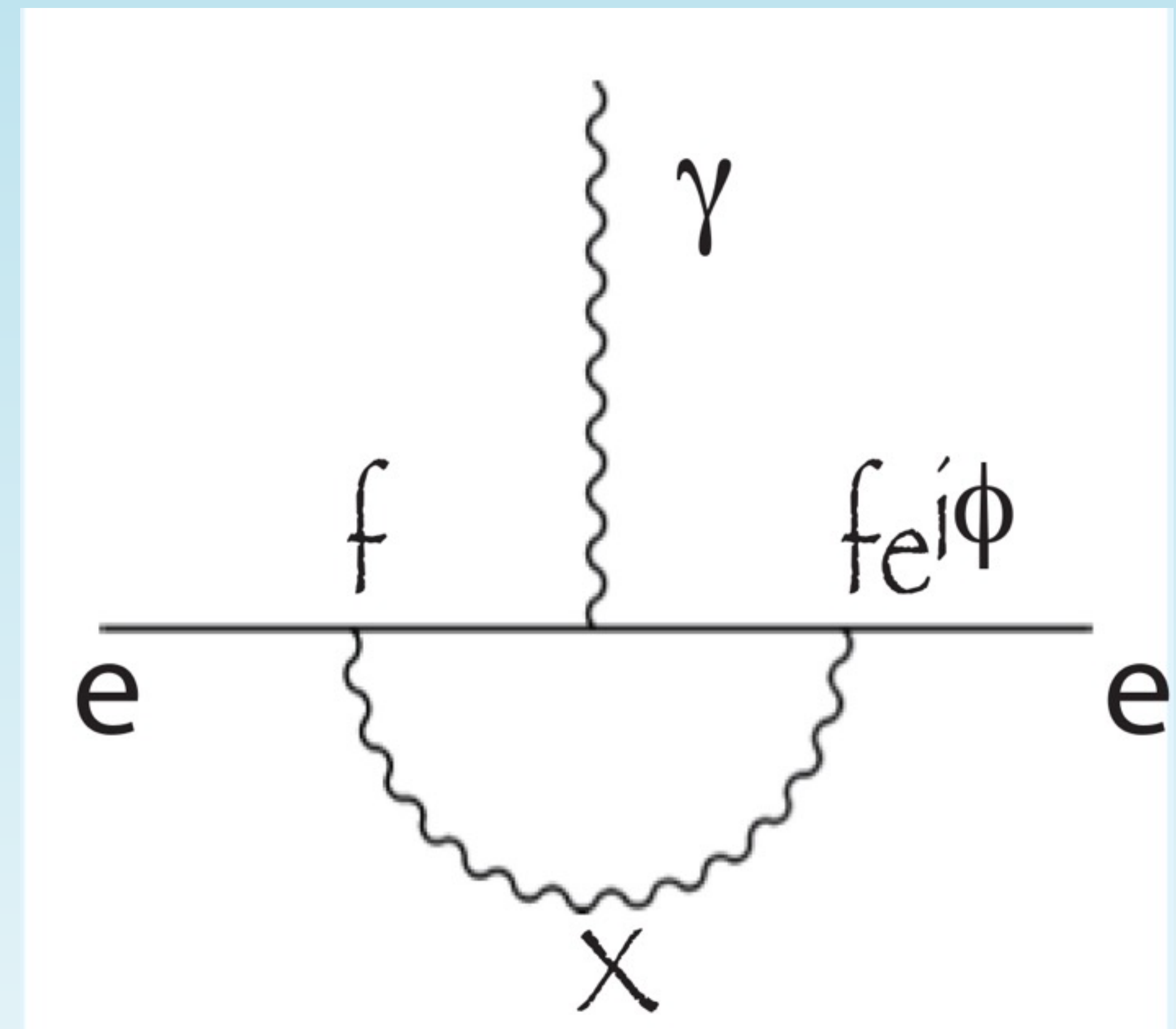
$\sin(\Phi) \approx 1$

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$$\text{EDM} \approx \mu_B (\alpha/\pi)^N (m_e/m_\chi)^2 \sin(\Phi)$$

EDM $\approx 100x$ previous limit

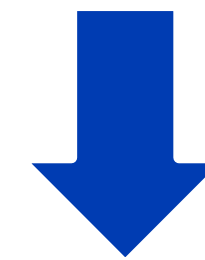
What the heck is going on?



ADMA RITZ TALK 2014 Summary

EDMs are an important class of flavour-diagonal CP-odd observables, testing/limiting new physics (motivated by the need for baryogenesis)

- Disentangling multiple CP-odd operators at 1 GeV requires multiple observables



recent applications

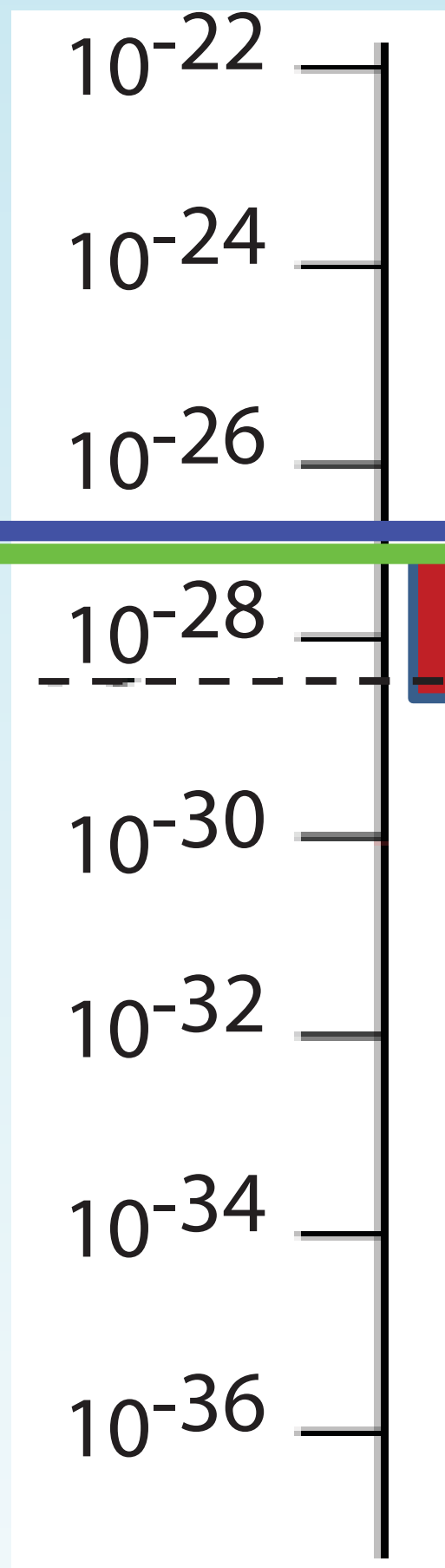
- Useful interplay between EDM constraints and precision

- The SUSY CP problem, hinted at by (1-loop) EDMs for more than 20 years, has been “confirmed” by the LHC, with no squarks seen near the weak scale (thus far). EDMs probe the very high (PeV) sfermion scales characteristic of the “large” observed Higgs mass

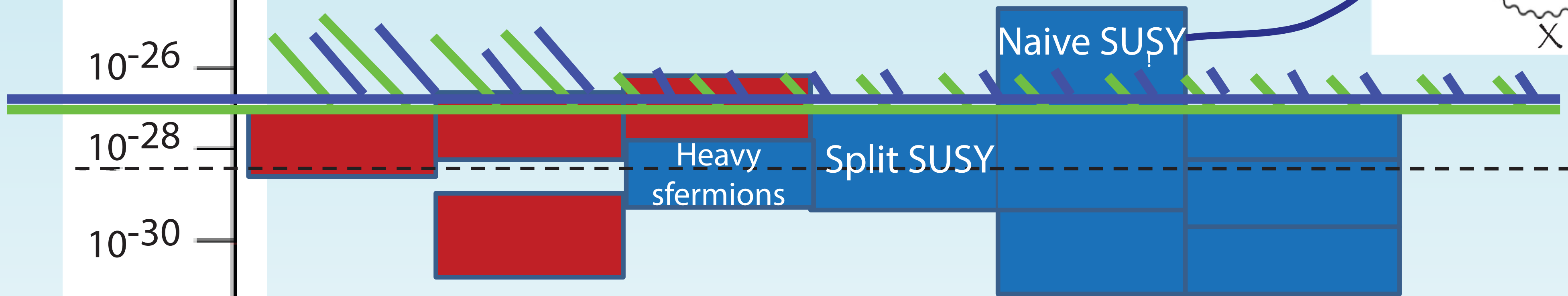
but is it at much higher scales? Or possibly at low scales ($< \text{GeV}$)?

Experimental Search for the eEDM — *pre-ACME*

Electron EDM
 d_e [e-cm]



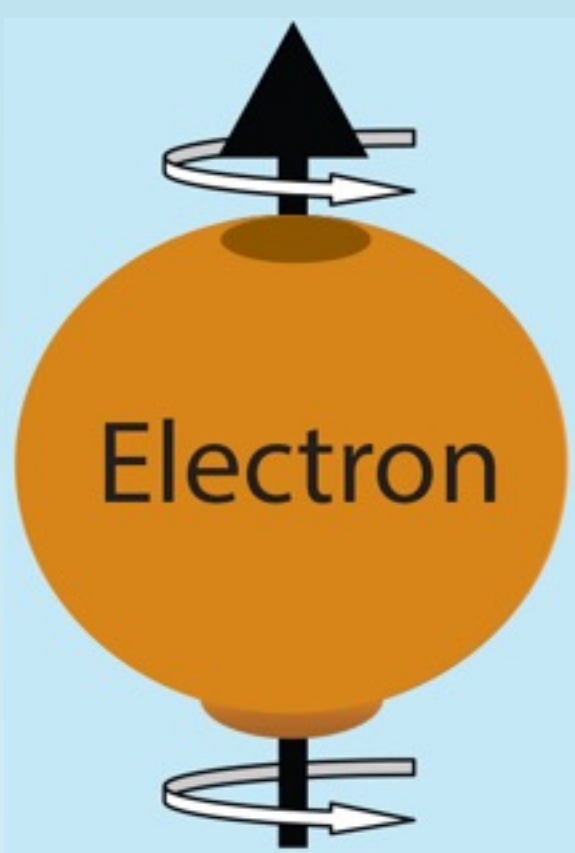
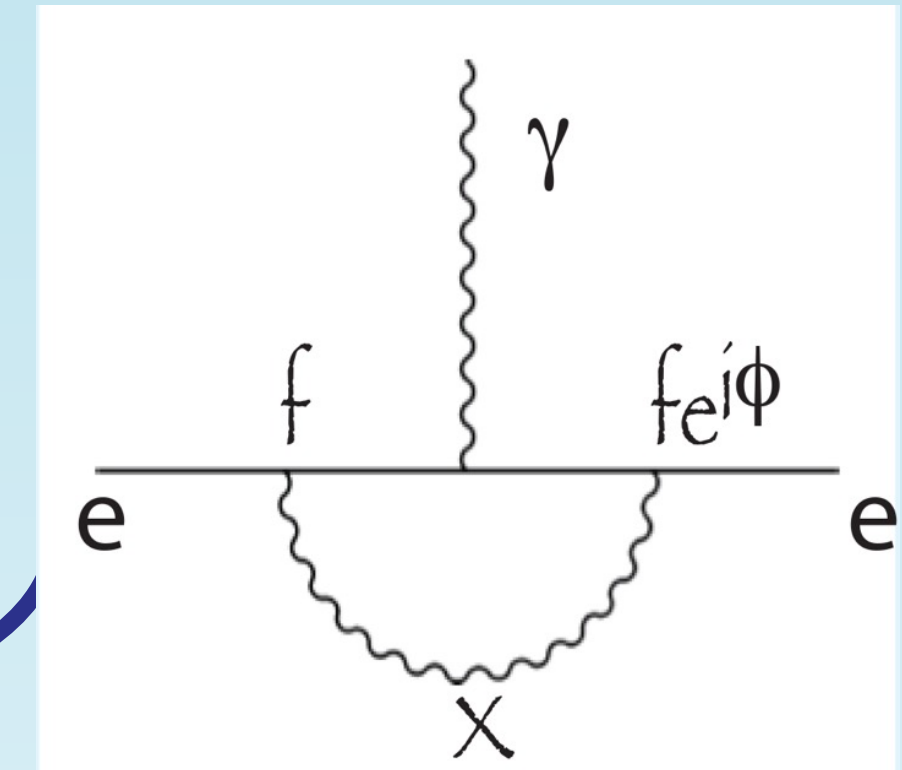
Excluded Region



New Limit on the Electron Dipole Moment, Regan, Commins, Schmidt, DeMille, PRL 88, 2002



Improved measurement of the shape of the electron, Hudson, Kara, Smallman, Sauer, Tarbutt, Hinds, Nature 473, 2011



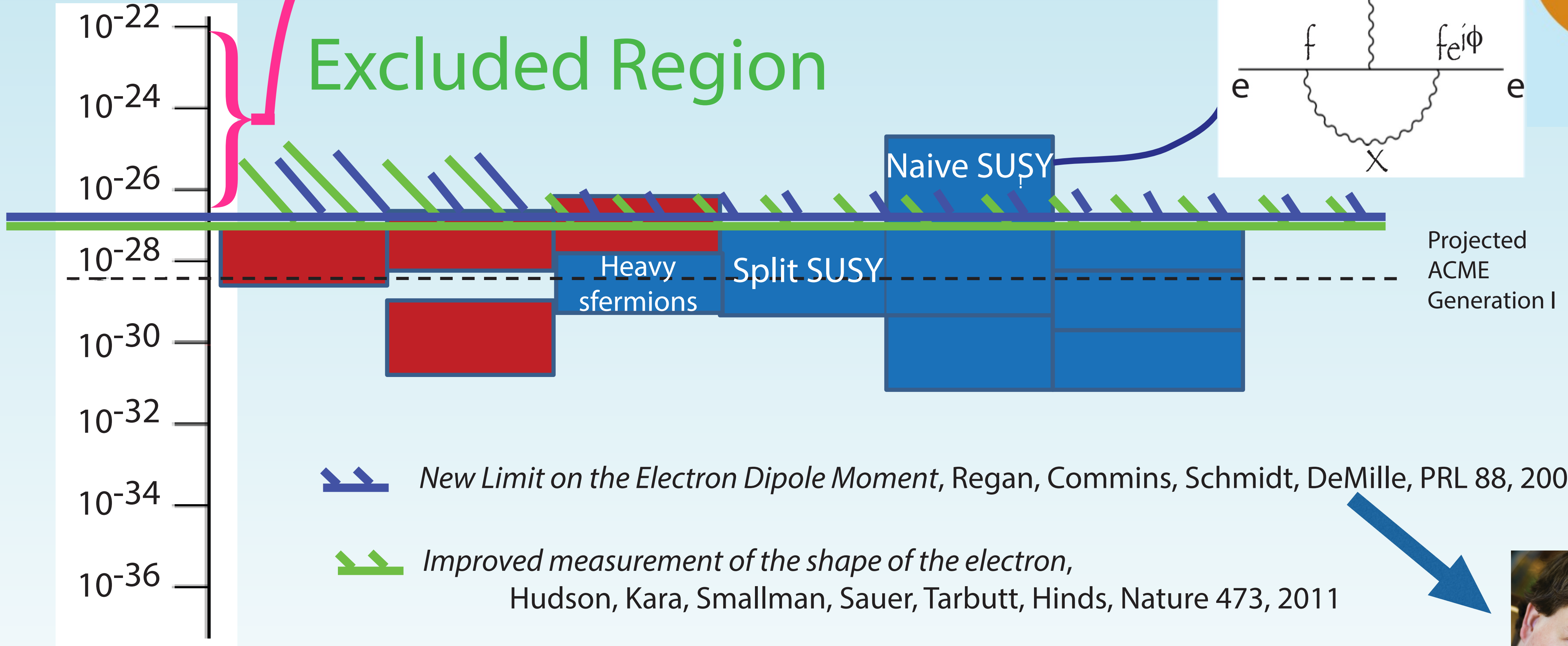
SM

Experimental Search for the eEDM — *pre-ACME*

Electron EDM
 d_e [e-cm]

50 years of experimental progress

Excluded Region

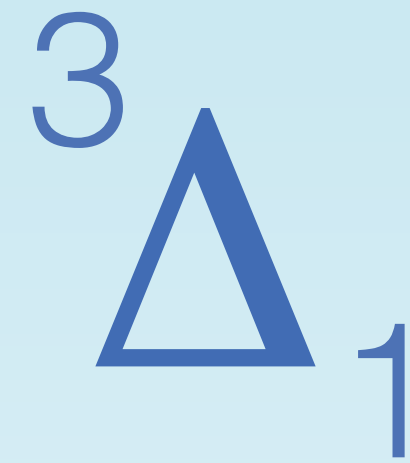


SM

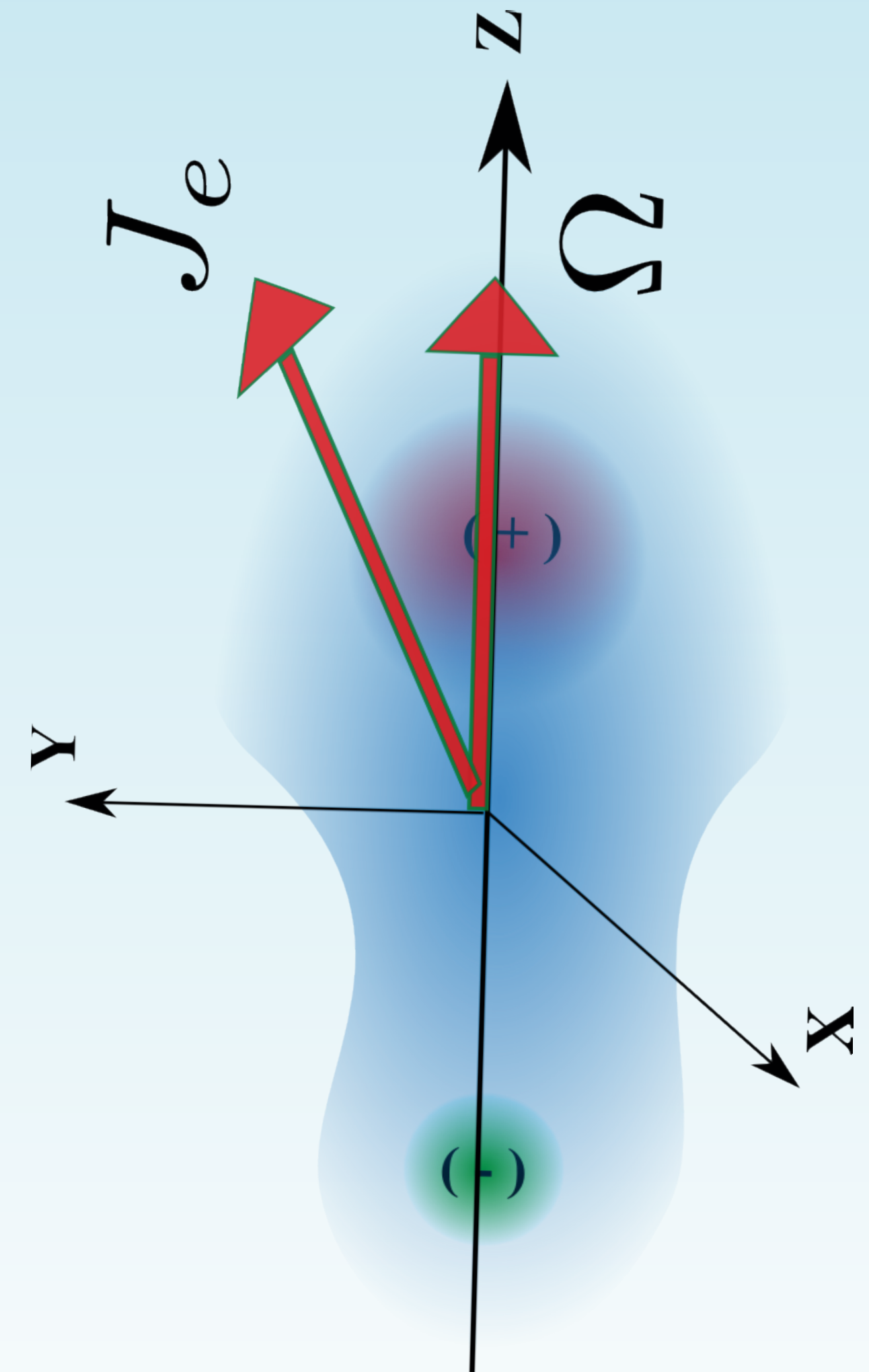


Our “Lab” — the ThO Molecule

ThO “H state”

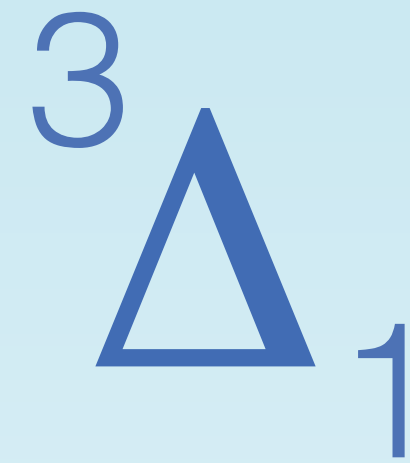


90 Th Thorium 232.0381 [Rn]6d ² 7s ²	8 O Oxygen 15.9994 1s ² 2s ² 2p ⁴
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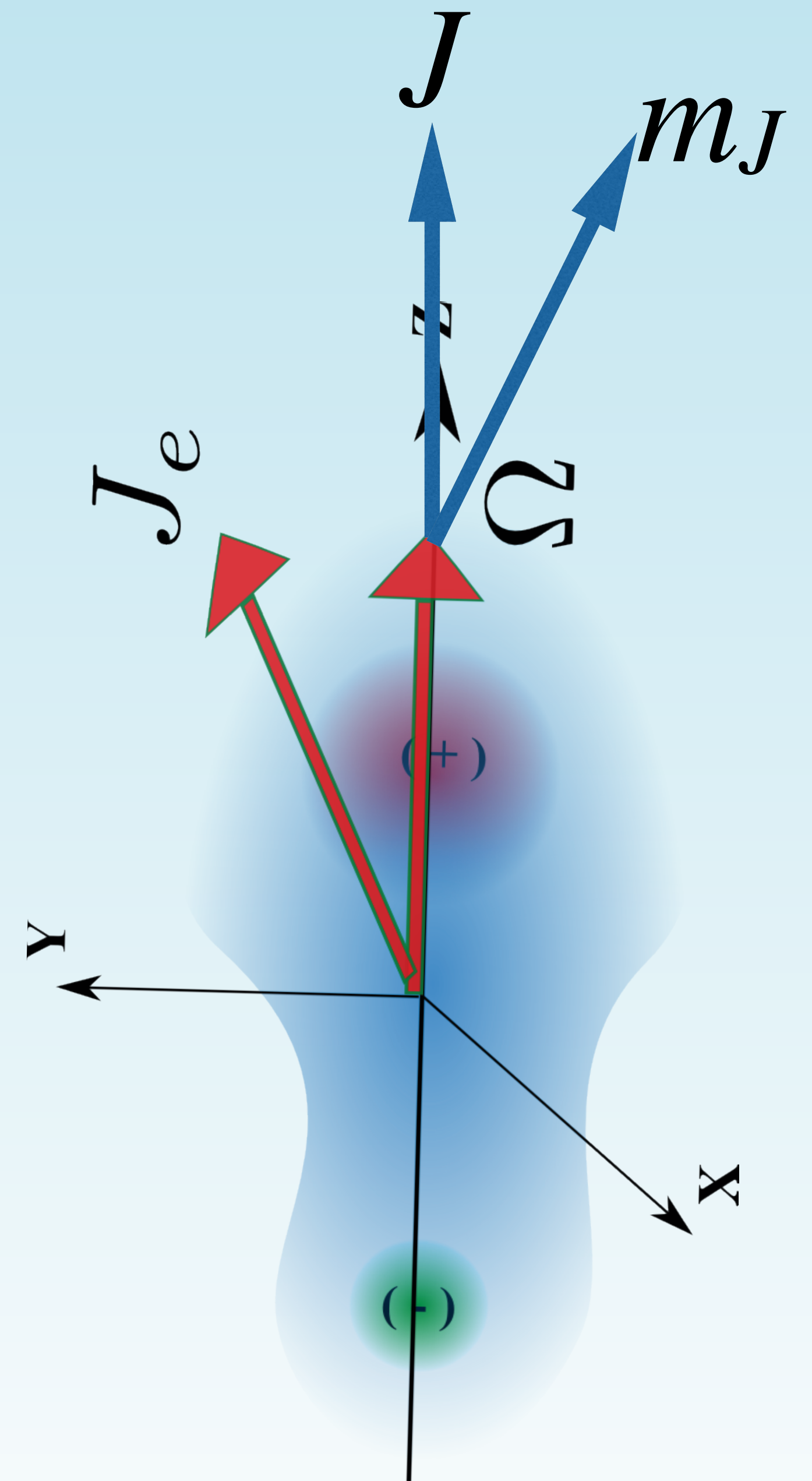


Our “Lab” — the ThO Molecule

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90 Th Thorium 232.0381 [Rn]6d ² 7s ²	8 O Oxygen 15.9994 1s ² 2s ² 2p ⁴
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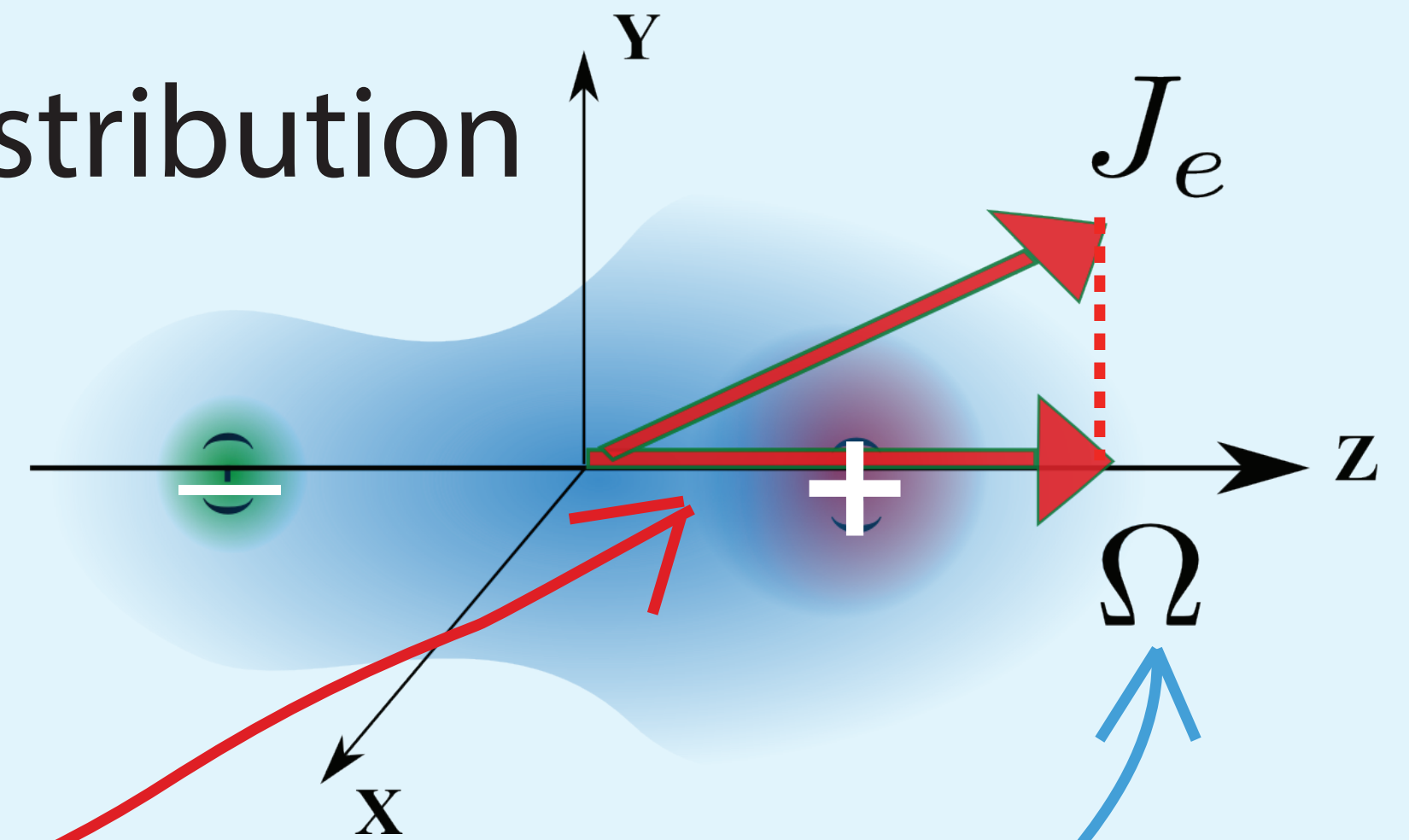
Using Molecules to Search for the EDM

Molecules -- the Good



- "Natural" Asymmetric Electric Field Distribution
(due to chemical bonding)

Strong Electric Field in There

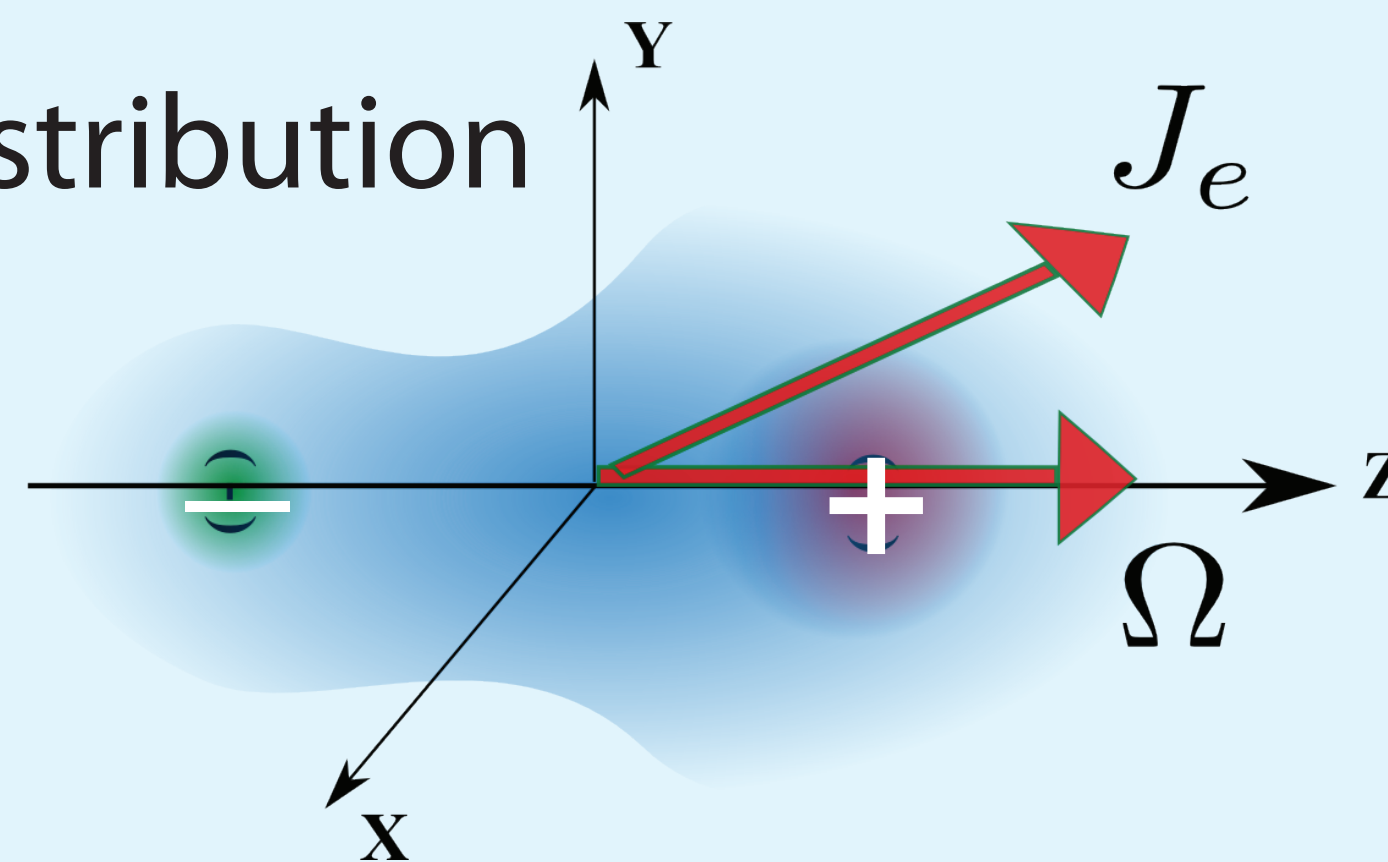


New Quantum Number

Using Molecules to Search for the EDM

Molecules -- the Good

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Molecules -- the Bad



- Go ahead, make my Hamiltonian*...

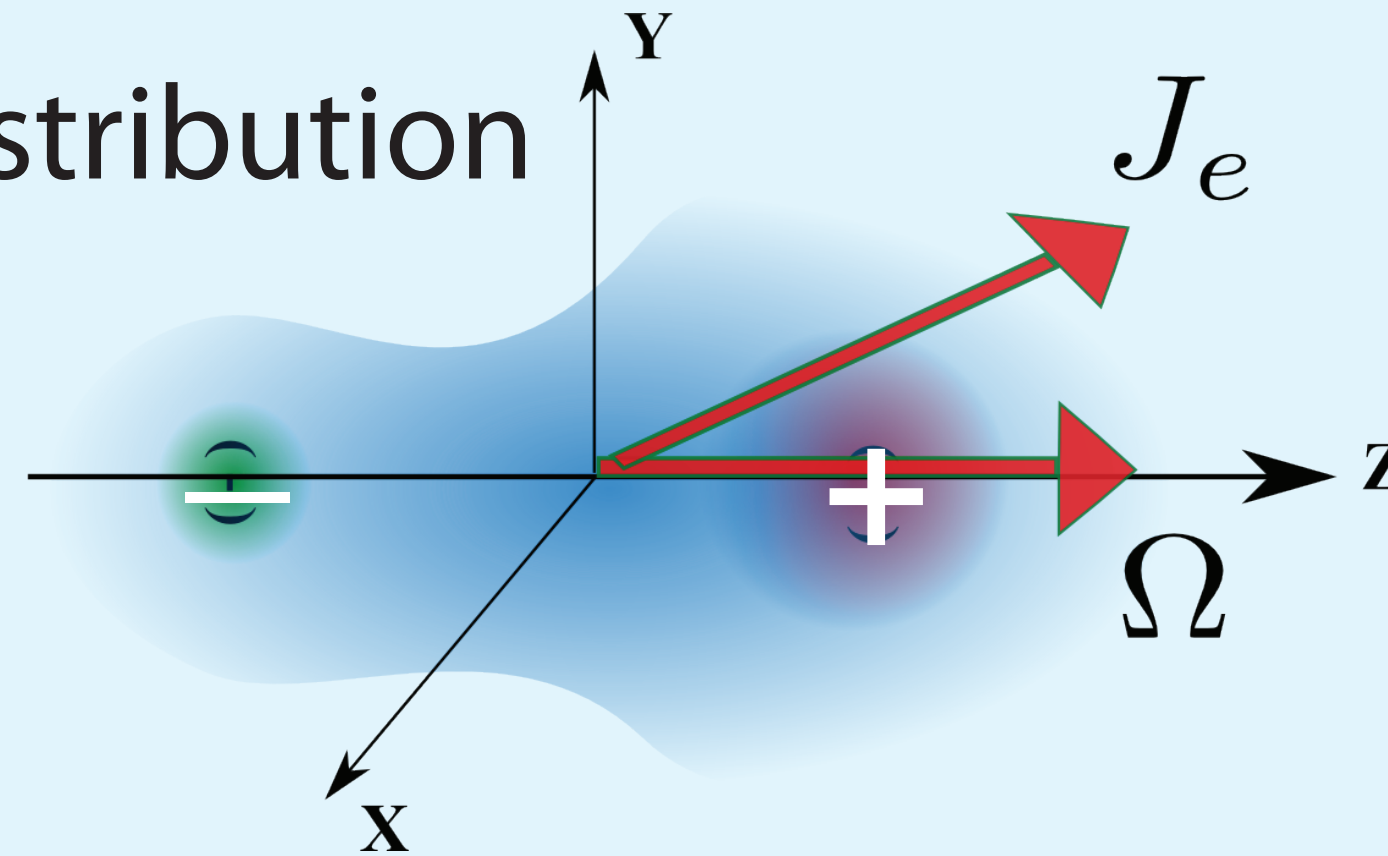
$$\begin{aligned}
 \hat{H} = & B(J^2 - J_z^2) - B(J_-L_+ + J_+L_-) + A_zL_zS_z + \frac{1}{2}(A_\perp + 2B - \gamma)(L_+S_- + L_-S_+) + \left[\frac{2}{3}(B + \lambda_z - 2\gamma) + \lambda_0 \right] S^2 - \frac{1}{3}(B - 2\lambda - 2\gamma)(3S_z^2 - S^2) \\
 & + \left(\lambda_\perp - \frac{1}{4}\gamma_\perp \right) (T_+S_zS_- + T_+S_-S_z + T_-S_+S_z + T_-S_zS_+) + \lambda_{\perp\perp} (T_+^2S_-^2 + T_-^2S_+^2) - \frac{1}{4}\gamma_\perp (L_+T_- + L_-T_+ + T_+L_- + T_-L_+) - (2B - \gamma)(J \cdot S - J_zS_z) \\
 & + \frac{1}{2}\gamma_\perp (J_+T_- + J_-T_+)S_z + \sum_{i=\alpha,\beta} \left[\frac{eQ_i}{2I_i(2I_i - 1)} q_{zi}^e I^2 f_{Lzi} L_z I_{zi} + \frac{1}{2} f_{Lz\perp i} (I_{+i}I_{-i} + L_{-i}I_{+i}) + f_{Szi} S_z I_z + \frac{1}{2} (f_{S\perp i} - c_i) (S_+I_i + S_-I_{+i}) + \frac{1}{2} f_{S\perp zi} (S_+T_- + S_-T_+) I_{zi} \right. \\
 & \left. + \frac{1}{2} f_{S_z\perp i} (I_{+i}T_- + I_iT_+) S_z + \frac{1}{4} f_{S\perp\perp i} (S_+T_-^2 I_{+i} + S_-T_+^2 I_{-i}) + c_i (J \cdot I - J_z I_{zi}) \right] + d(I_\alpha \cdot I_\beta - 3I_{z\alpha} I_{z\beta}) + \delta I_\alpha \cdot I_\beta + \mu_B (g_{lz} - \sigma_{lz}) L_z \mathcal{B}_z + \frac{1}{2} \mu_B (g_{l\perp} + g_z - \sigma_{l\perp}) (L_+ \mathcal{B}_- + L_- \mathcal{B}_+) \\
 & + \frac{1}{4} \mu_B \sigma_{sz\perp} [\mathcal{B}_z (S_+T_- + S_-T_+) + (T_-S_+ + T_+S_-) \mathcal{B}_z + (\mathcal{B}_+T_- + \mathcal{B}_-T_+) S_z + S_z (T_- \mathcal{B}_+ + T_+ \mathcal{B}_-)] + \frac{1}{4} \mu_B \sigma_{s\perp\perp} (\mathcal{B}_+ S_+ T_-^2 + \mathcal{B}_- S_- T_+^2 + T_-^2 S_+ \mathcal{B}_+ T_+^2 S_- \mathcal{B}_-) \\
 & - g_z (J \cdot \mathcal{B} - J_z \mathcal{B}_z) - \sum_{i=\alpha,\beta} \left\{ g_{li} \left[(1 - \sigma_{iz}) I_{zi} \mathcal{B}_z + \frac{1}{2} (1 - \sigma_{i\perp}) (I_{+i} \mathcal{B}_- + I_{-i} \mathcal{B}_+) \right] + \frac{1}{4} \sigma_{i\perp\perp} (\mathcal{B}_+ I_{+i} T_-^2 + \mathcal{B}_- I_{-i} T_+^2 + T_-^2 I_{+i} \mathcal{B}_+ + T_-^2 I_{+i} \mathcal{B}_+ + T_+^2 I_{-i} \mathcal{B}_-) \right\} \\
 & + (\text{terms nondiagonal in } S)
 \end{aligned}$$

*thanks to Roman Krems

Using Molecules to Search for the EDM

Molecules -- the Good

- "Natural" Asymmetric Electric Field Distribution
(due to chemical bonding)



Molecules -- the Bad



- Go ahead, make my Hamiltonian*...

What would possess a person
to
make a precision measurement in such a complicated
"laboratory"
?
?

$$\begin{aligned}
 & + \left(\lambda \right) \\
 & + \frac{1}{2} \gamma_{\perp} (J_{+} T_{-} + J_{-} T_{+}) \\
 & + \frac{1}{2} f_{S_z \perp i} (I_{+i} T_{-} + I_{-i} T_{+}) S_z \\
 & + \frac{1}{4} \mu_B
 \end{aligned}$$

$$\begin{aligned}
 & - 2\gamma) (3S_z^2 - S^2) \\
 & \gamma) (J \cdot S - J_z S_z) \\
 & + T_{-} + S_{-} T_{+}) I_{z i} \\
 & L_{+} B_{-} + L_{-} B_{+}) \\
 & S_{+} B_{+} T_{+}^2 S_{-} B_{-}) \\
 & B_{+} + T_{+}^2 I_{-i} B_{-}) \}
 \end{aligned}$$

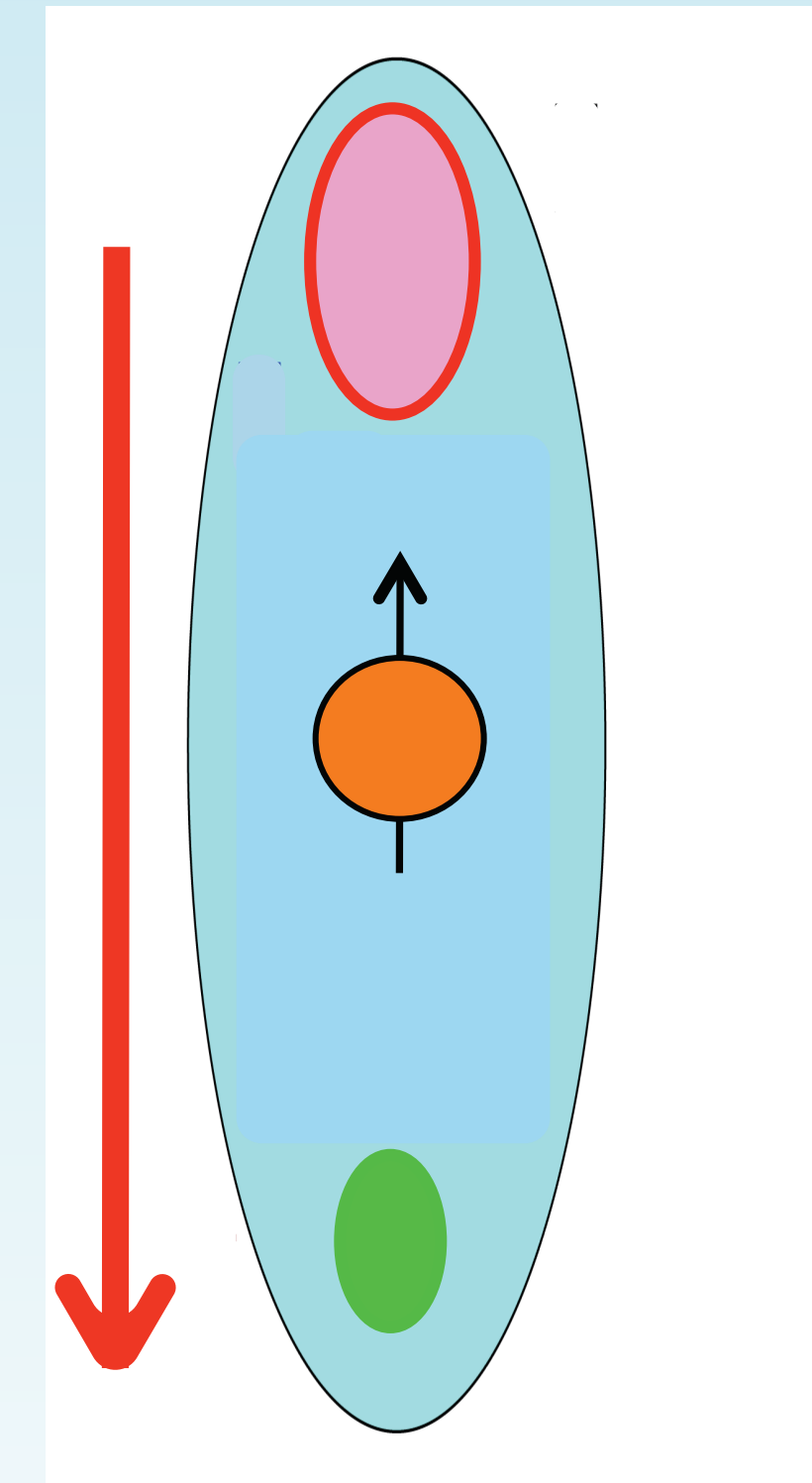
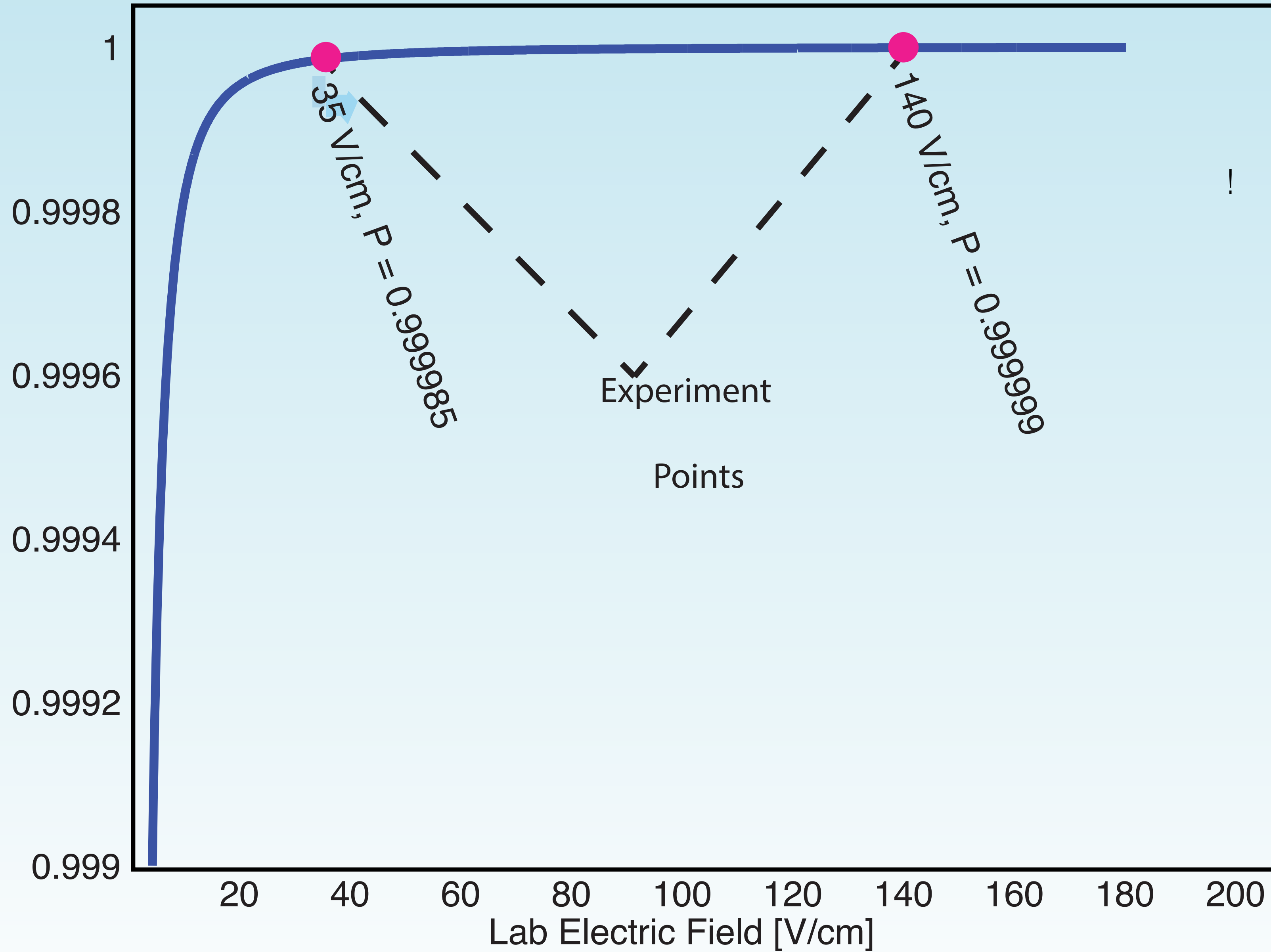
*thanks to Roman Krems

Lets start with the BASICS

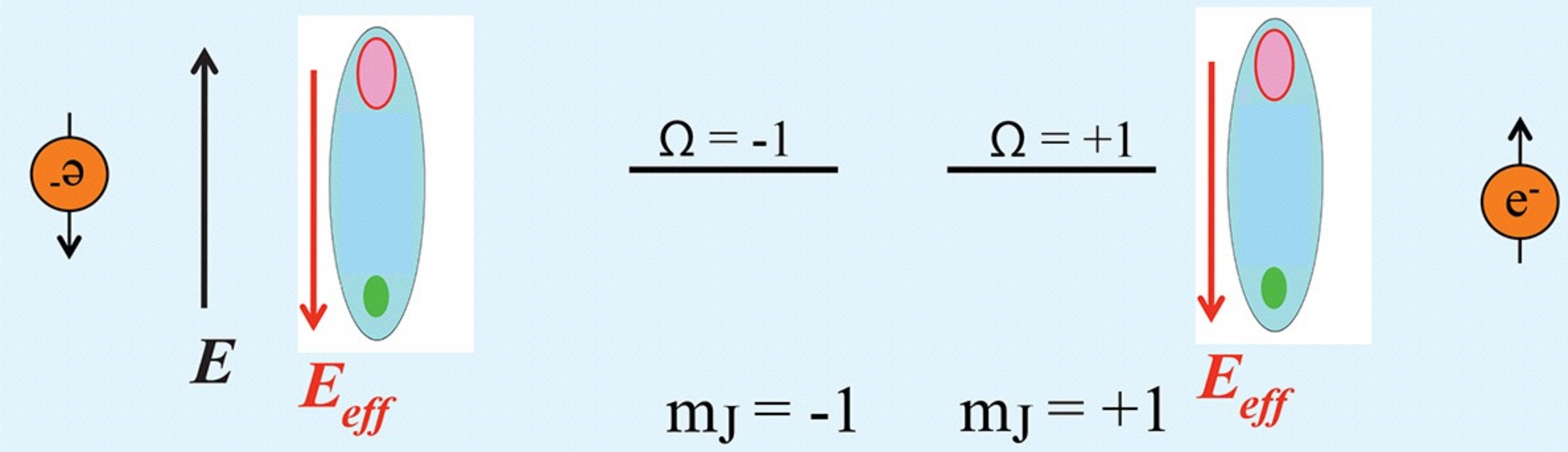
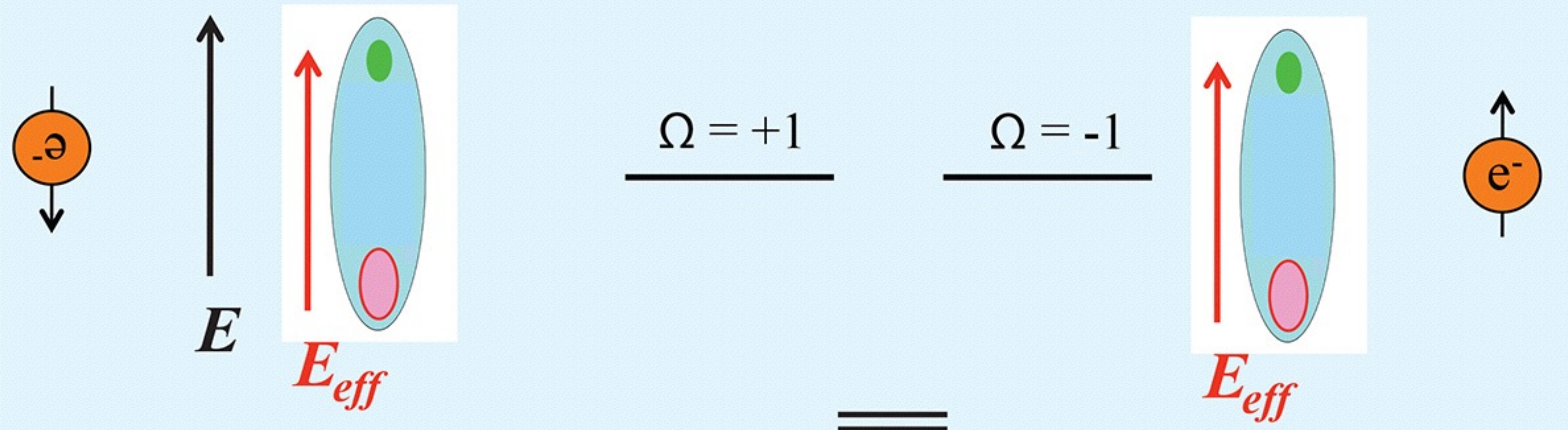
OK, Lets put the molecules in an electric field

Polarization of ThO H,J=1 in a Laboratory Electric Field

Fully Polarized ThO - "Easy"

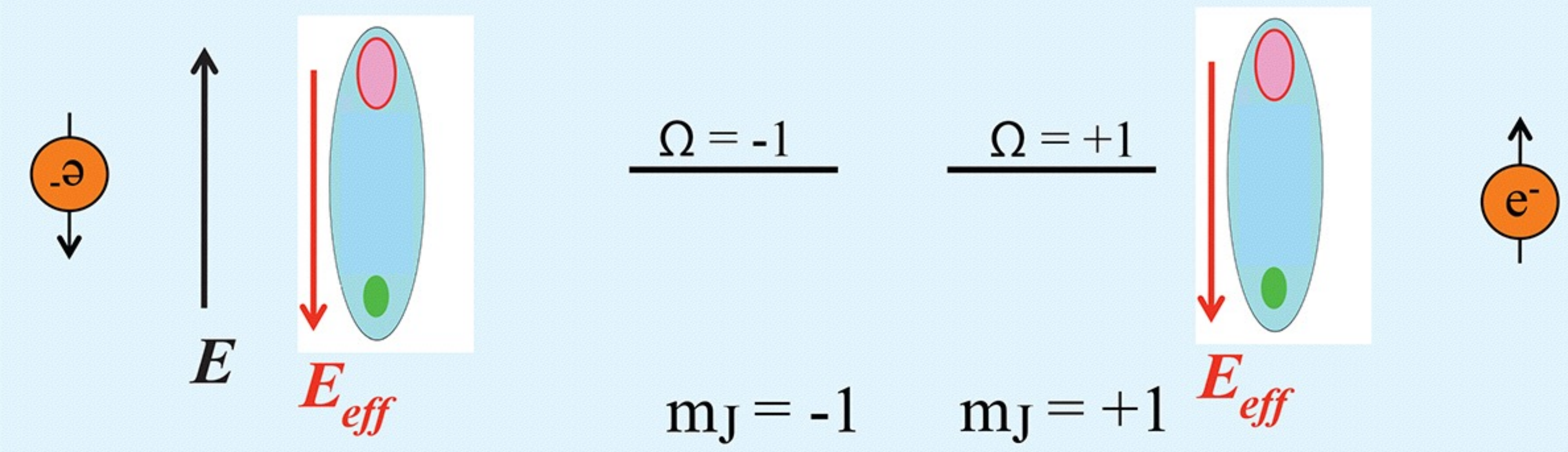
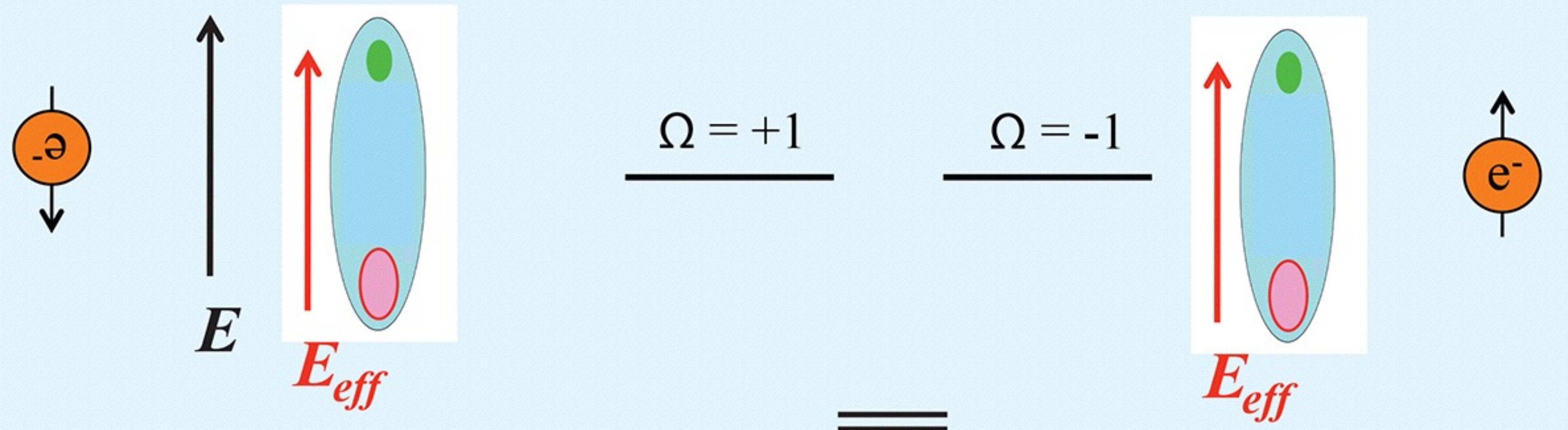


ThO H state in an Electric Field



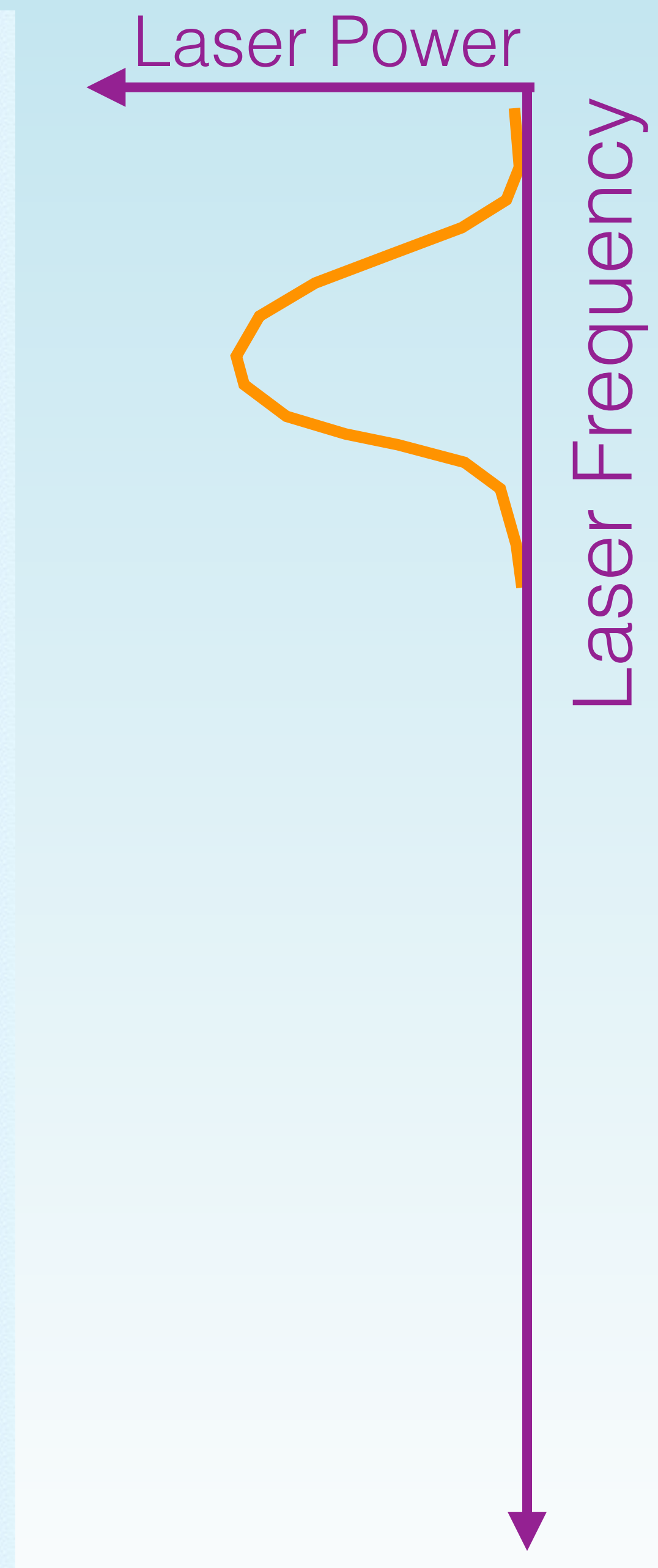
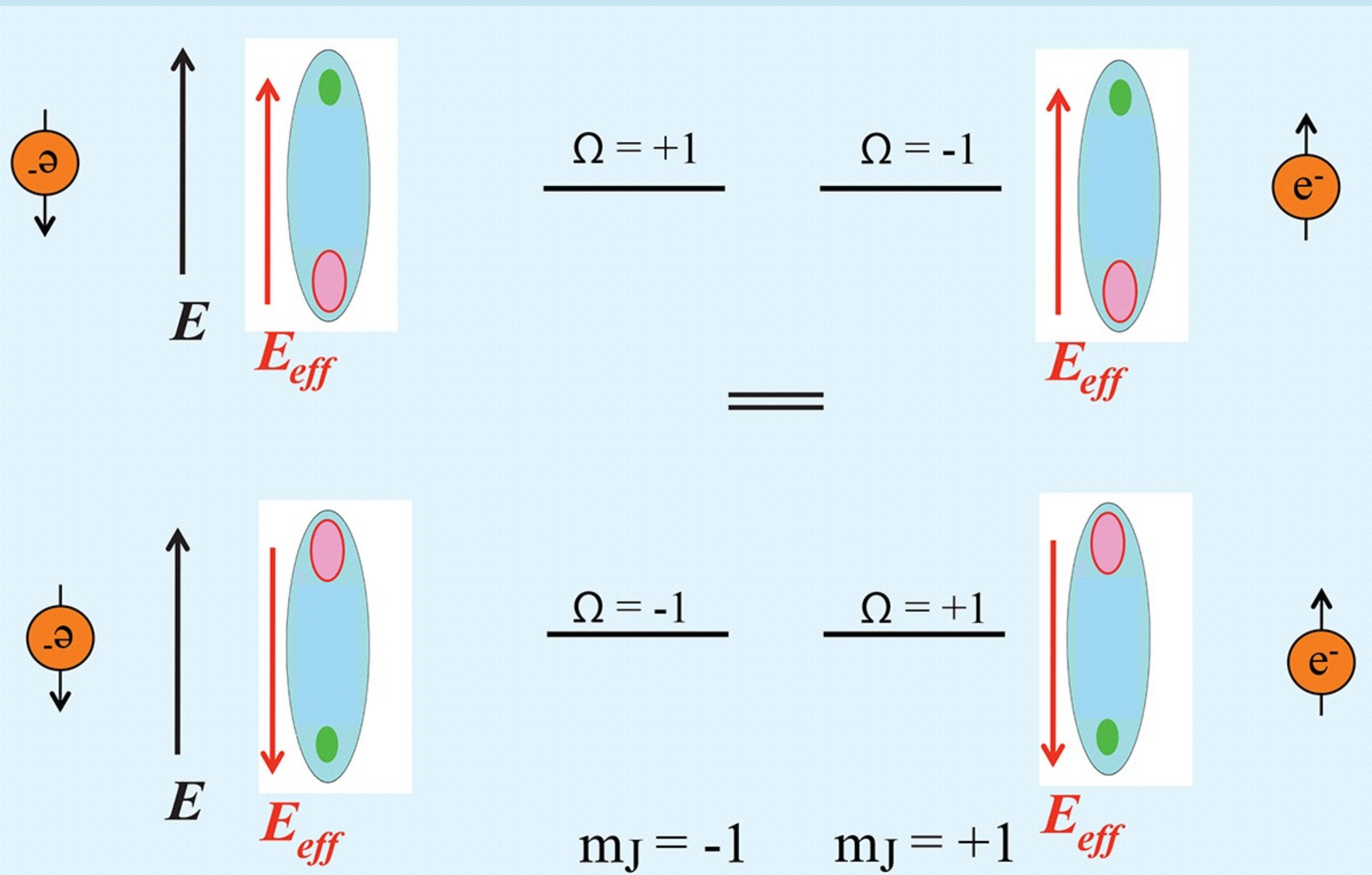
ThO H state in an Electric Field

Frequency Selection



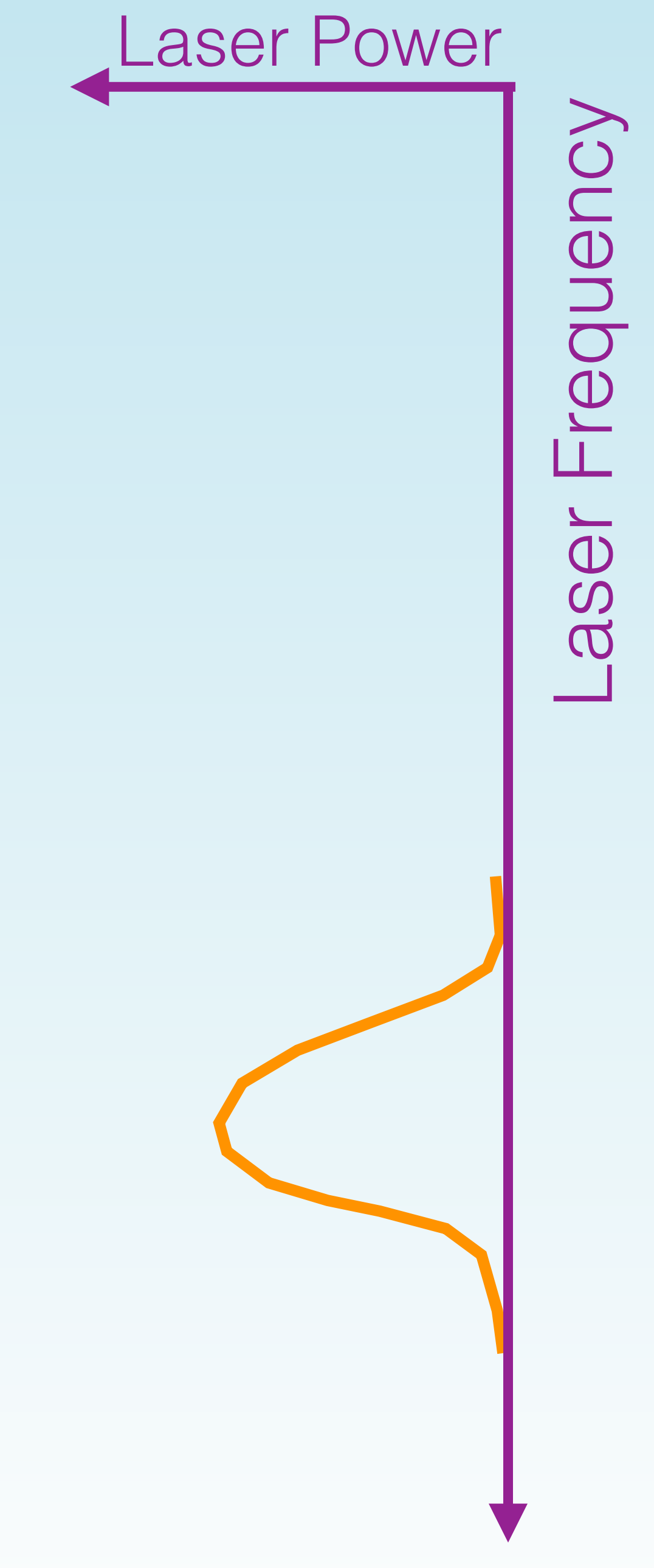
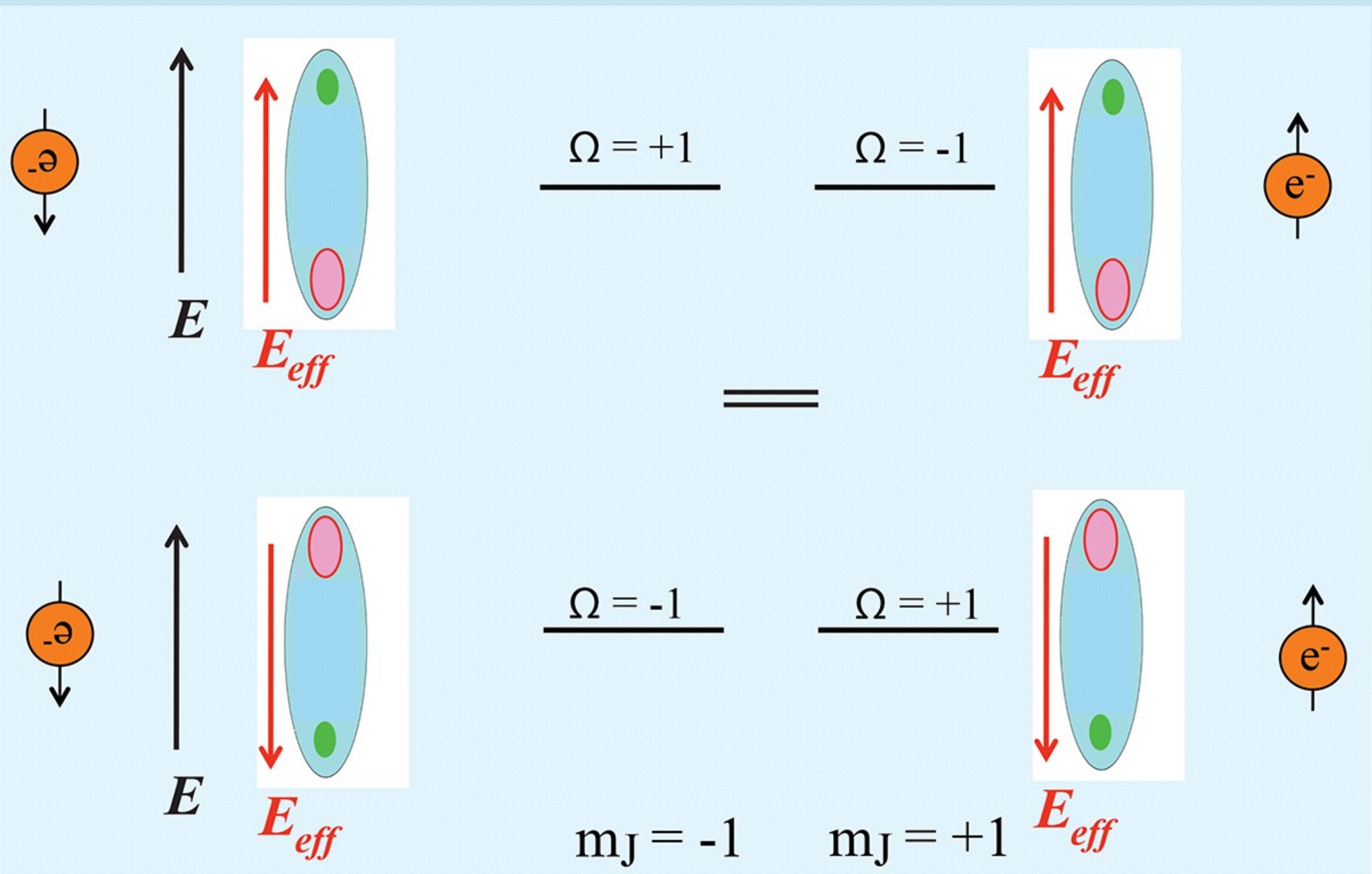
ThO H state in an Electric Field

Frequency Selection



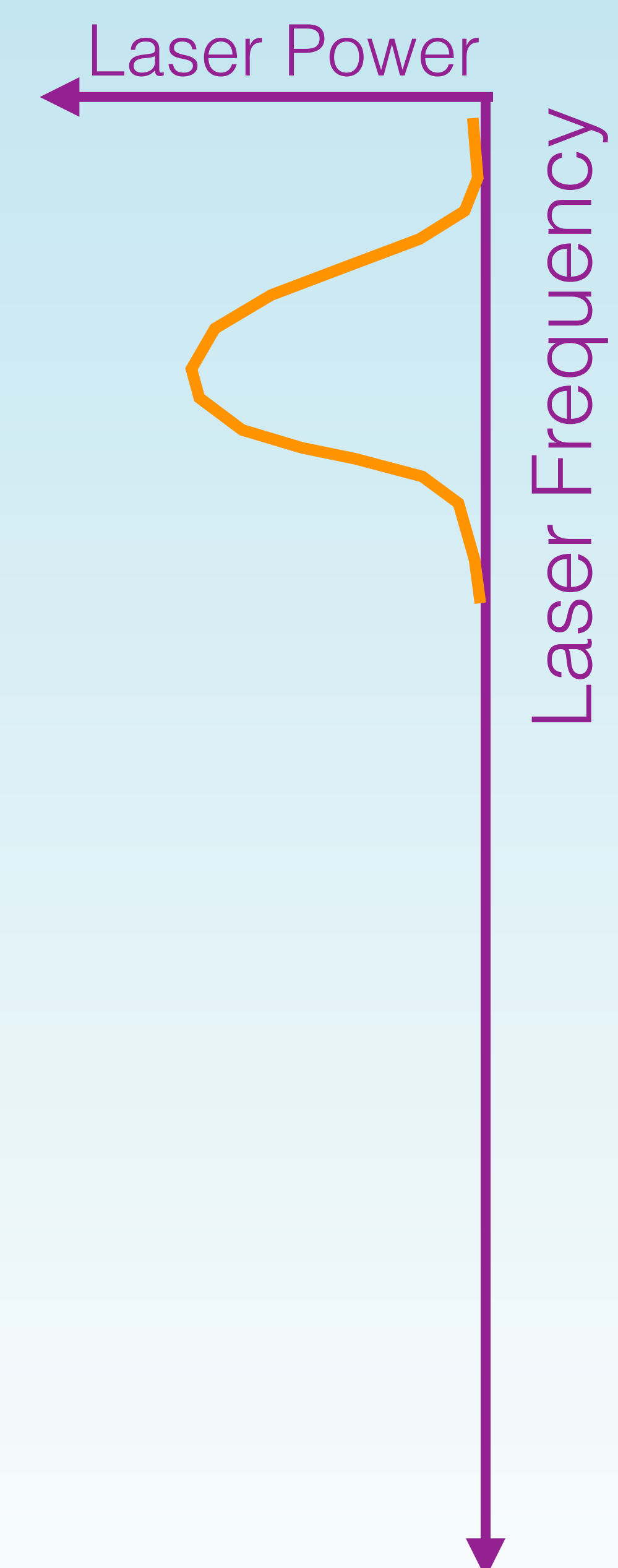
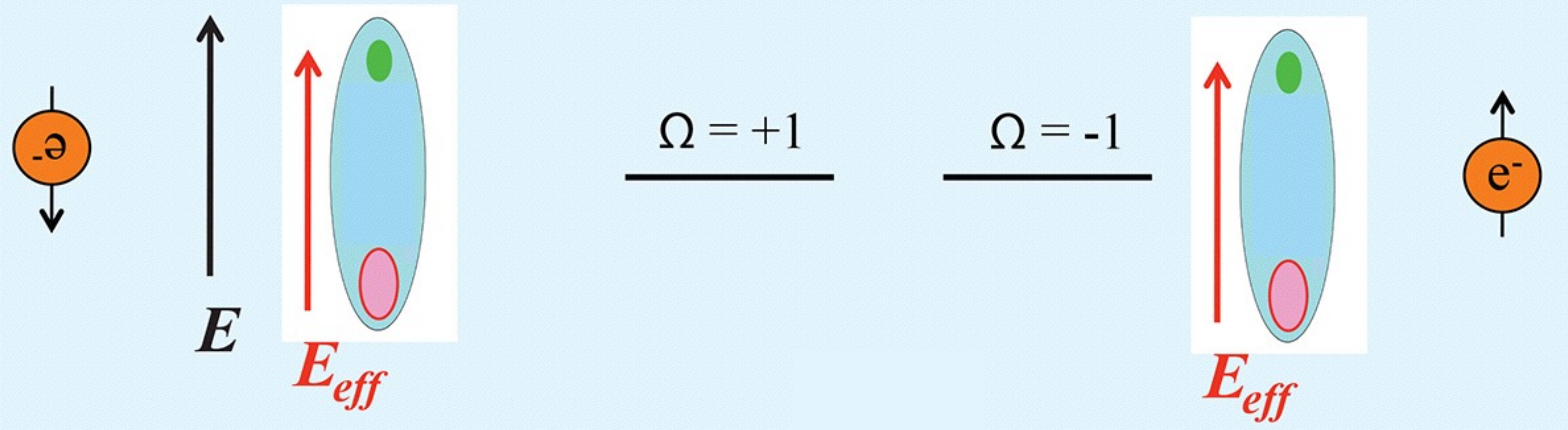
ThO H state in an Electric Field

Frequency Selection

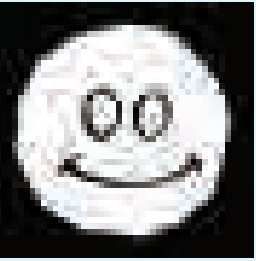


ThO H state in an Electric Field

Frequency Selection

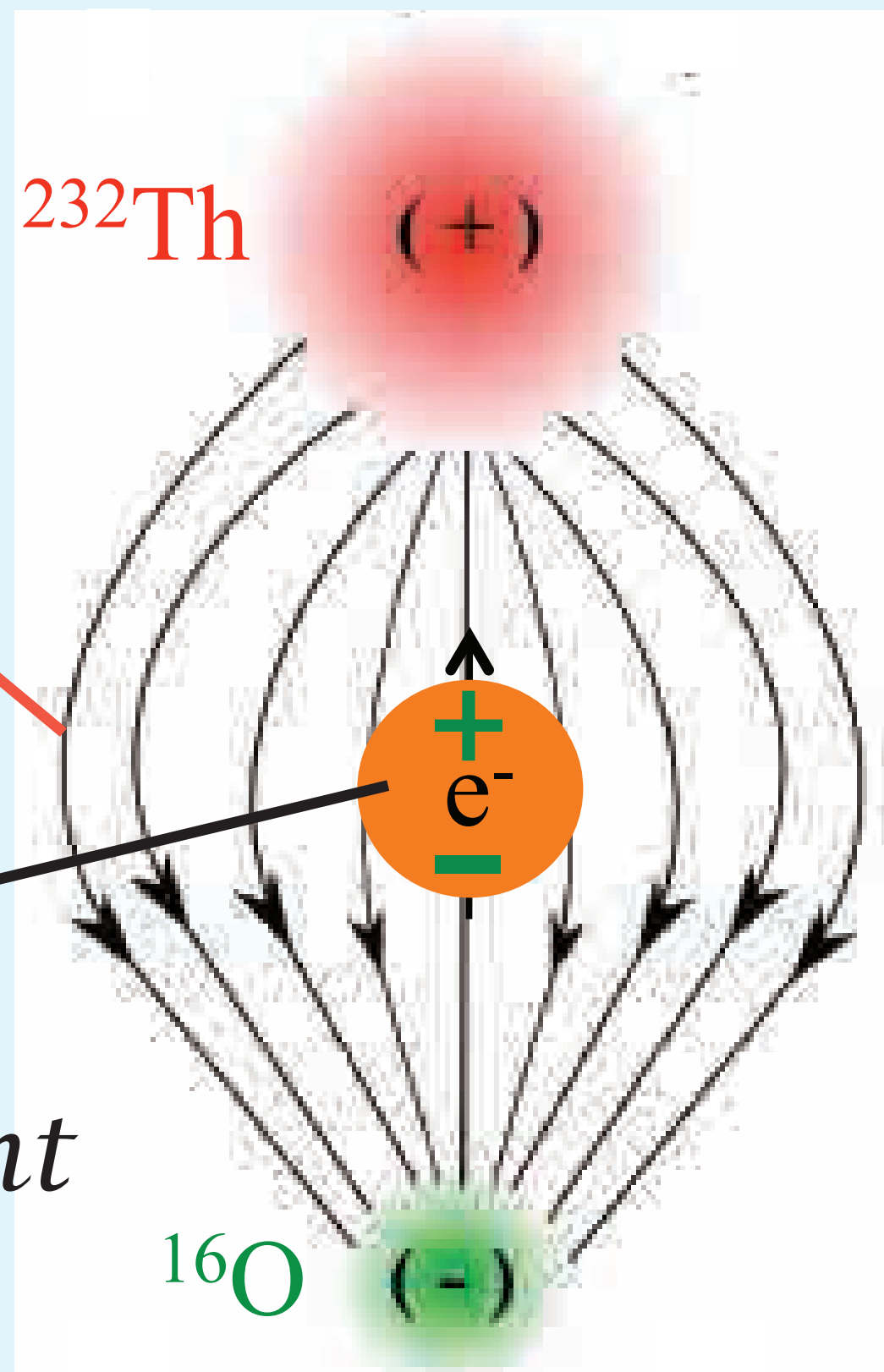


Ignore these levels for the moment...



Add the eEDM -- Effective Field Interaction

\mathcal{E}_{eff}
Effective
Electric Field
Inside the Molecule



d_e

Electric Dipole Moment
of the Electron

$\mathcal{E}_{lab} \sim 100 \text{ V/cm}$

$$H'_{de} = -d_e \cdot \mathcal{E}_{eff}$$

d_e interacts with \mathcal{E}_{eff}

$\mathcal{E}_{eff} \sim 10^{11} \text{ V/cm}$

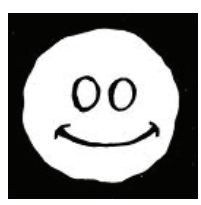
THE JOURNAL OF CHEMICAL PHYSICS 139, 221103 (2013)

Communication: Theoretical study of ThO for the electron electric dipole moment search

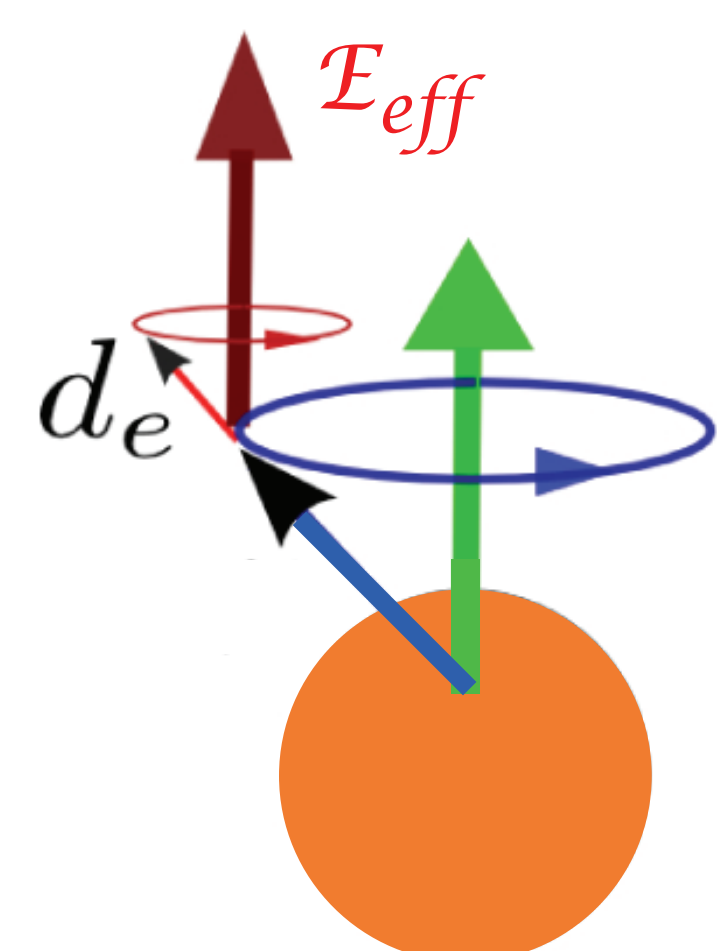
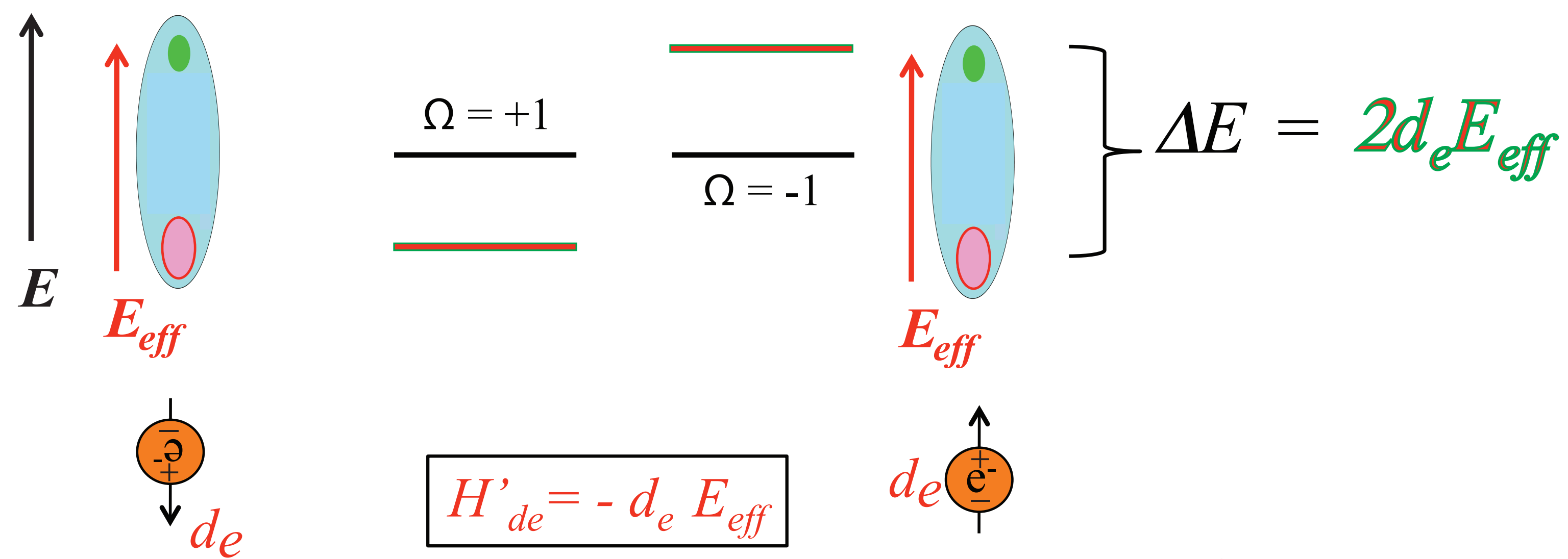
L. V. Skripnikov,^{1,2,a)} A. N. Petrov,^{1,2} and A. V. Titov^{1,2,b)}

¹Federal State Budgetary Institute "Petersburg Nuclear Physics Institute," Gatchina, Leningrad district 188300, Russia

Later, more detailed calculation



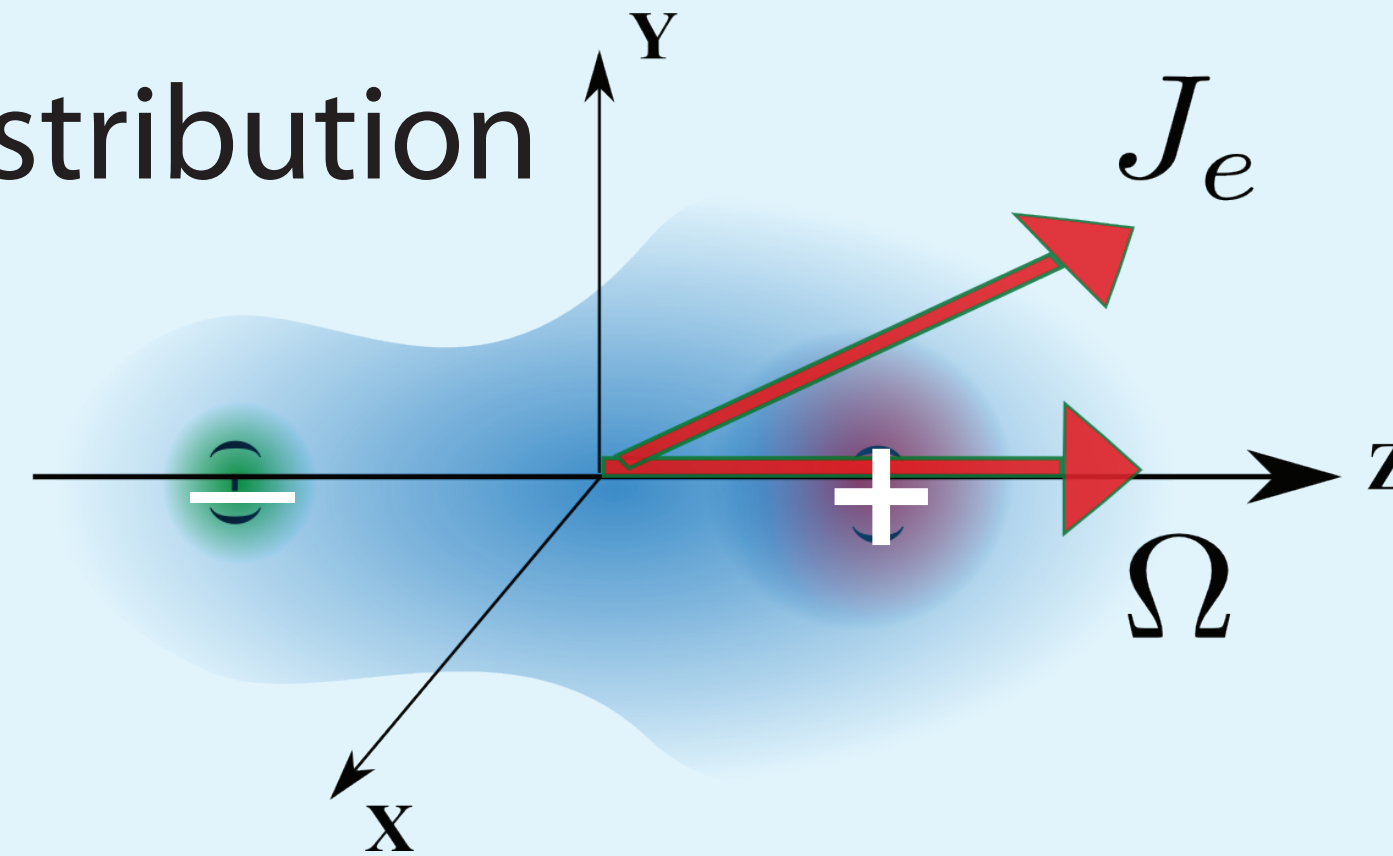
Add the eEDM - Levels Shift, Electron Spin Precesses



Using Molecules to Search for the EDM

Molecules -- the Good

- "Natural" Asymmetric Electric Field Distribution
(due to chemical bonding)



Molecules -- the Bad



- Go ahead, make my Hamiltonian*...

$$\begin{aligned}
 & + \left(\lambda \right) \\
 & + \frac{1}{2} \gamma_{\perp} (J_{+} T_{-} + J_{-} T_{+}) \\
 & + \frac{1}{2} f_{S_z \perp i} (I_{+i} T_{-} + I_{-i} T_{+}) S_z \\
 & + \frac{1}{4} \mu_E
 \end{aligned}$$

$$\begin{aligned}
 & - 2\gamma) (3S_z^2 - S^2) \\
 & \gamma)(J \cdot S - J_z S_z) \\
 & + T_{-} + S_{-} T_{+}) I_{zi} \\
 & L_{+} B_{-} + L_{-} B_{+}) \\
 & S_{+} B_{+} T_{+}^2 S_{-} B_{-}) \\
 & B_{+} + T_{+}^2 I_{-i} B_{-}) \}
 \end{aligned}$$

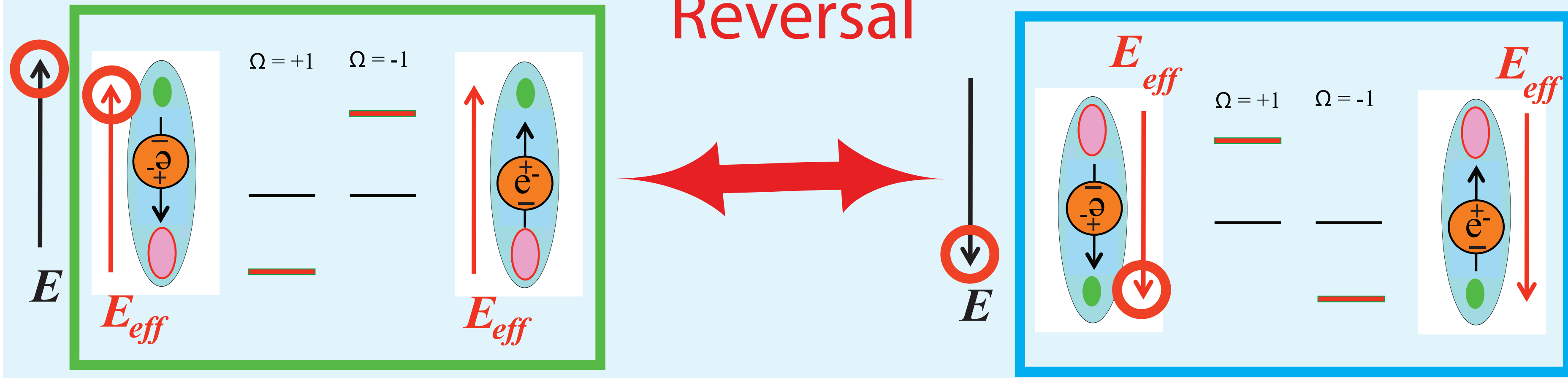
Here is how we handle these complications...

*thanks to Roman Krems



Reverse Direction of E field, Keep Laser Frequency Fixed

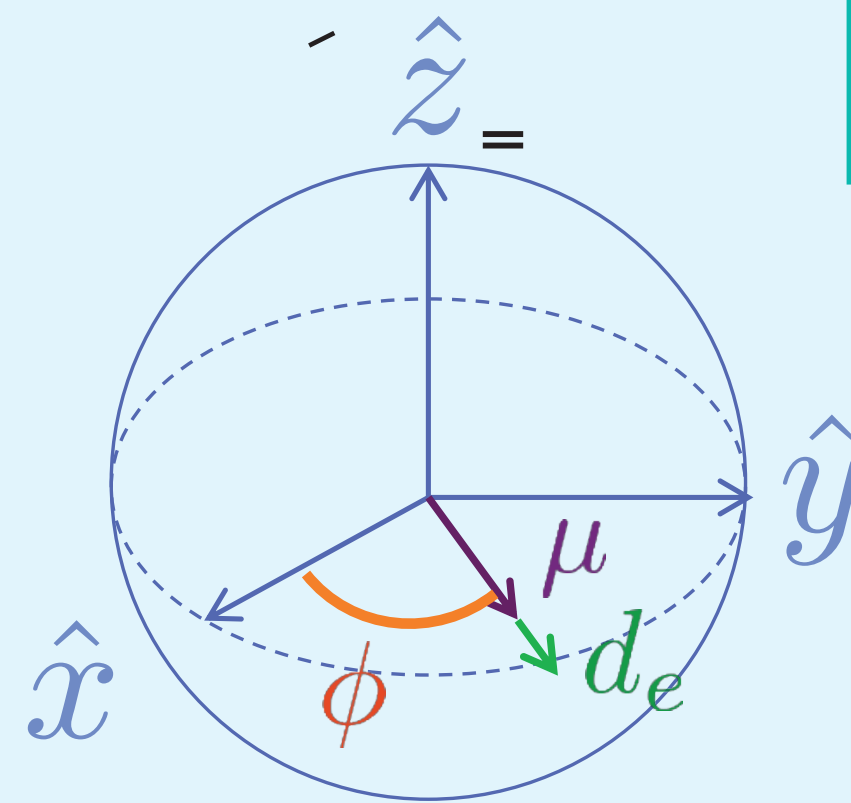
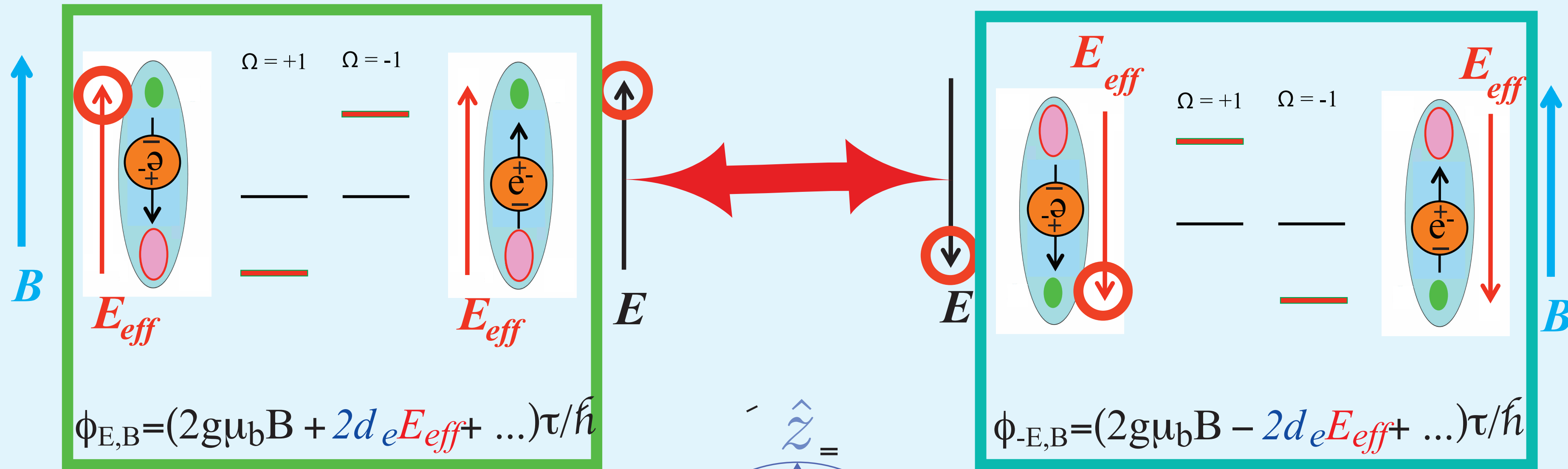
E field Reversal





Reverse Direction of E field, Keep Laser Frequency Fixed

E field Reversal

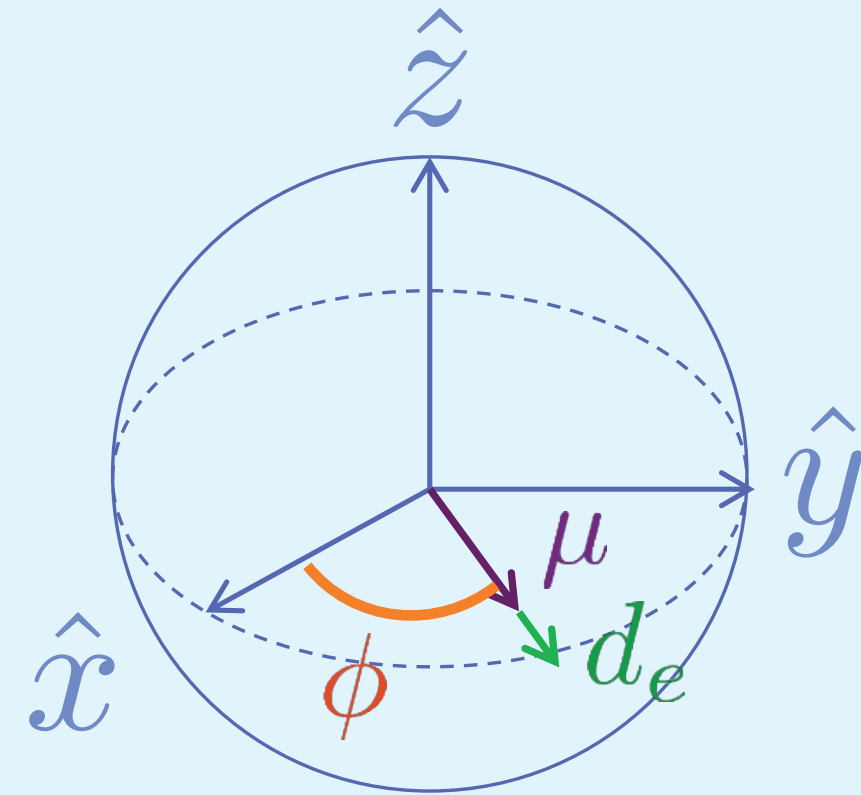


$$P_- \equiv \boxed{\phi_{E,B} = (2g\mu_b B + 2d_e E_{eff} + \dots)\tau/\hbar} - \boxed{\phi_{-E,B} = (2g\mu_b B - 2d_e E_{eff} + \dots)\tau/\hbar} = (4d_e E_{eff} + \dots)\tau/\hbar$$



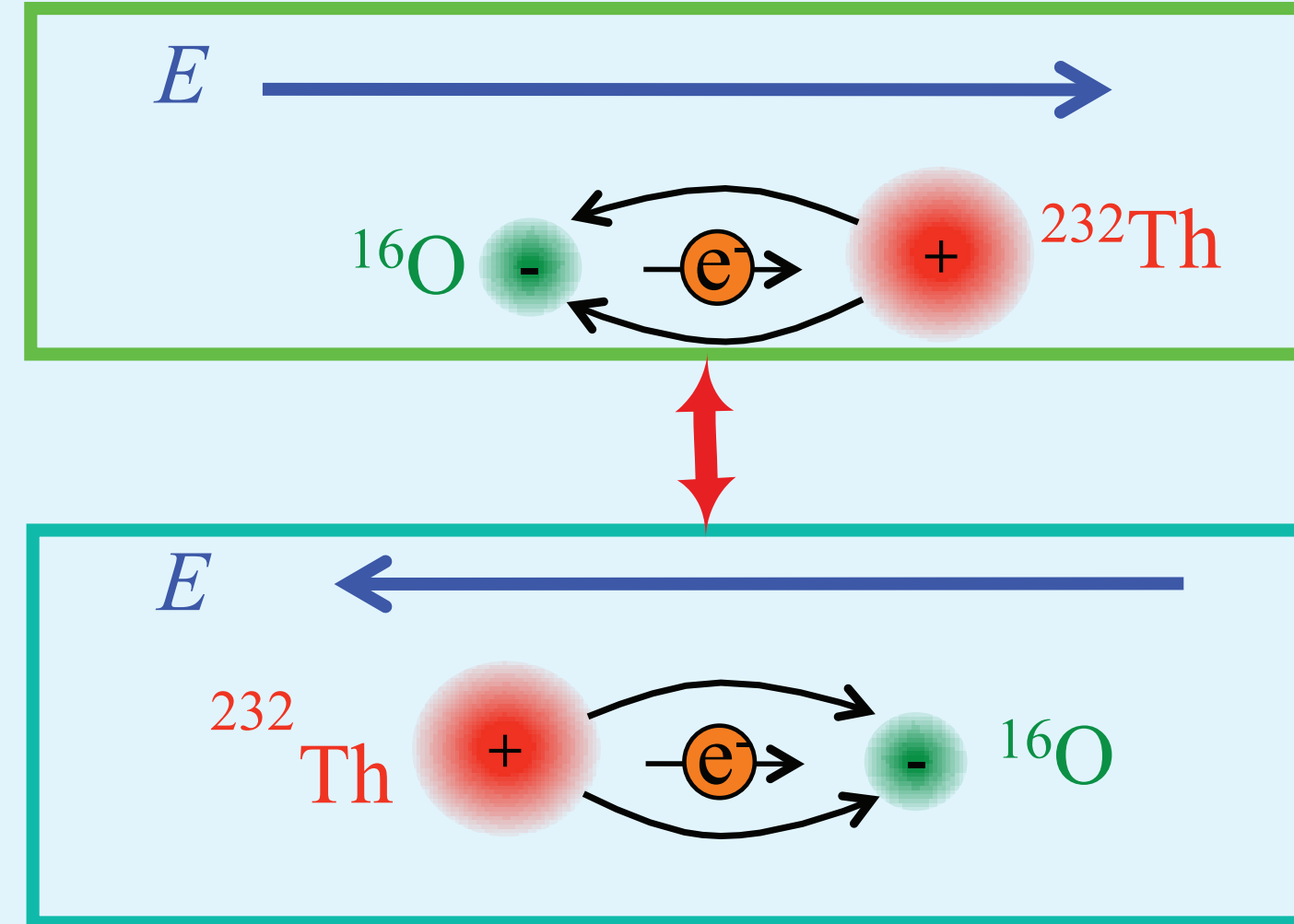
Determining eEDM, now with Efield Reversal

Measure the precession of the spin with E field reversal, subtract the two values



$$H = H_0 + H_1 + H_2 + H_3 + H_4 + H_5 + H_6 + H_7 + H_8 + H_9 + H_{10} + H_{11} + H_{12} + H_{13} + H_{14} + H_{15} + H_{16} + H_{17} + H_{18} + H_{19} + H_{20} + H_{21} + H_{22} + H_{23} + H_{24} + H_{25} + H_{26} + H_{27} + H_{28} + H_{29} + H_{30} + H_{31} + H_{32} + H_{33} + H_{34} + H_{35} + H_{36} + H_{37} + H_{38} + H_{39} + H_{40} + H_{41} + H_{42} + H_{43} + H_{44} + H_{45} + H_{46} + H_{47} + H_{48} + H_{49} + H_{50} + H_{51} + H_{52} + H_{53} + H_{54} + H_{55} + H_{56} + H_{57} + H_{58} + H_{59} + H_{60} + H_{61} + H_{62} + H_{63} + H_{64} + H_{65} + H_{66} + H_{67} + H_{68} + H_{69} + H_{70} + H_{71} + H_{72} + H_{73} + H_{74} + H_{75} + H_{76} + H_{77} + H_{78} + H_{79} + H_{80} + H_{81} + H_{82} + H_{83} + H_{84} + H_{85} + H_{86} + H_{87} + H_{88} + H_{89} + H_{90} + H_{91} + H_{92} + H_{93} + H_{94} + H_{95} + H_{96} + H_{97} + H_{98} + H_{99} + H_{100} + \dots$$

+ terms due to applied fields + EDM terms



$$P_- = (4d_e E_{eff} + g\mu_b \mathbf{B}_{E-corr} + \eta\mu_b E_{nr} |\mathbf{B}| + \dots) \tau / \hbar$$

\mathbf{B} field correlated with \mathcal{E} reversal

Need $B_{E-corr} < 10^{-7}$ gauss, for 10^{-28} e-cm

B_{E-corr} measured with fluxgates, $< 10^{-8}$ gauss sensitivity

“Non-reversing \mathcal{E} field”

-- g factor is linear in E, due to rotational mixing

-- Need $E_{nr} < 1$ V/cm,

-- E_{nr} measured by raman spectroscopy, \sim few mv/cm

-- AND η determined from \mathcal{N} , \mathcal{B} reversals (1 nm/volt)

*\mathcal{E} field reversal alone could give greatly improved measurement of eEDM
But we don't do just \mathcal{E} reversals.....*

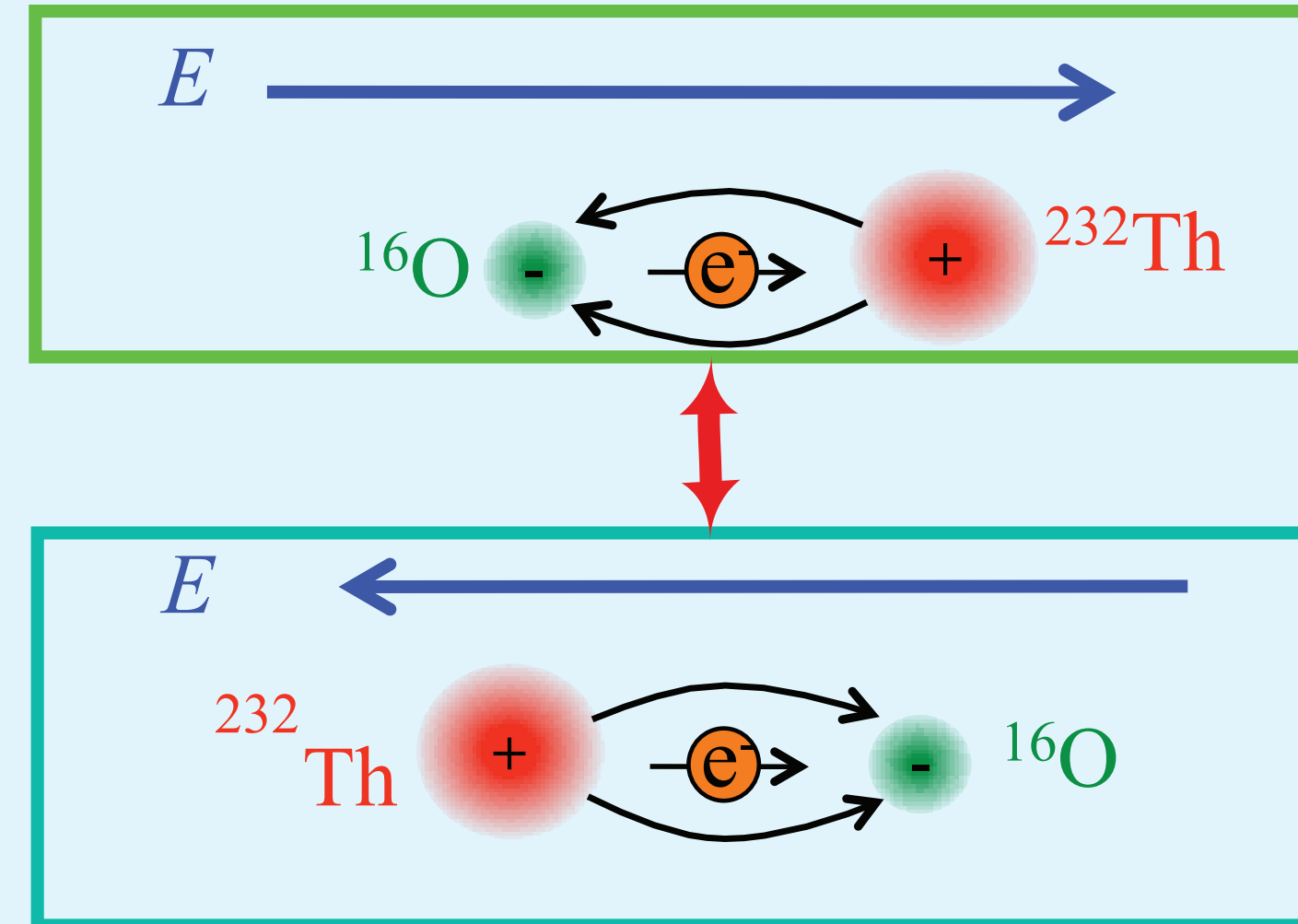


Determining eEDM, now with Efield Reversal

Measure the precession of the spin with E field reversal, subtract the two values



$$\begin{aligned}
 & \left[\frac{1}{2} \mu_B \frac{d}{dt} \langle S_z \rangle + \frac{1}{2} \mu_B \frac{d}{dt} \langle L_z \rangle + \frac{1}{2} \mu_B \frac{d}{dt} \langle S_x \rangle + \frac{1}{2} \mu_B \frac{d}{dt} \langle L_x \rangle + \frac{1}{2} \mu_B \frac{d}{dt} \langle S_y \rangle + \frac{1}{2} \mu_B \frac{d}{dt} \langle L_y \rangle + \frac{1}{2} \mu_B \frac{d}{dt} \langle S_z \rangle + \frac{1}{2} \mu_B \frac{d}{dt} \langle L_z \rangle + \dots \right] \\
 & + \text{terms due to applied fields} + \text{EDM terms}
 \end{aligned}$$



$$P_- = (4d_e E_{eff} + g\mu_b \mathbf{B}_{E\text{-corr}} + \eta\mu_b E_{nr} |\mathbf{B}| + \dots) \tau / \hbar$$

B field correlated with *E* reversal

Need $B_{E\text{-corr}} < 10^{-7}$ gauss, for 10^{-28} e-cm

$B_{E\text{-corr}}$ measured with fluxgates, $< 10^{-8}$ gauss sensitivity

“Non-reversing *E* field”

- g factor is linear in E, due to rotational mixing
- Need $E_{nr} < 1$ V/cm,
- E_{nr} measured by raman spectroscopy, ~ few mv/cm
- *AND* η determined from *N*, *B* reversals (1 nm/volt)

E field reversal alone could give greatly improved measurement of eEDM
But we don't do just *E* reversals.....

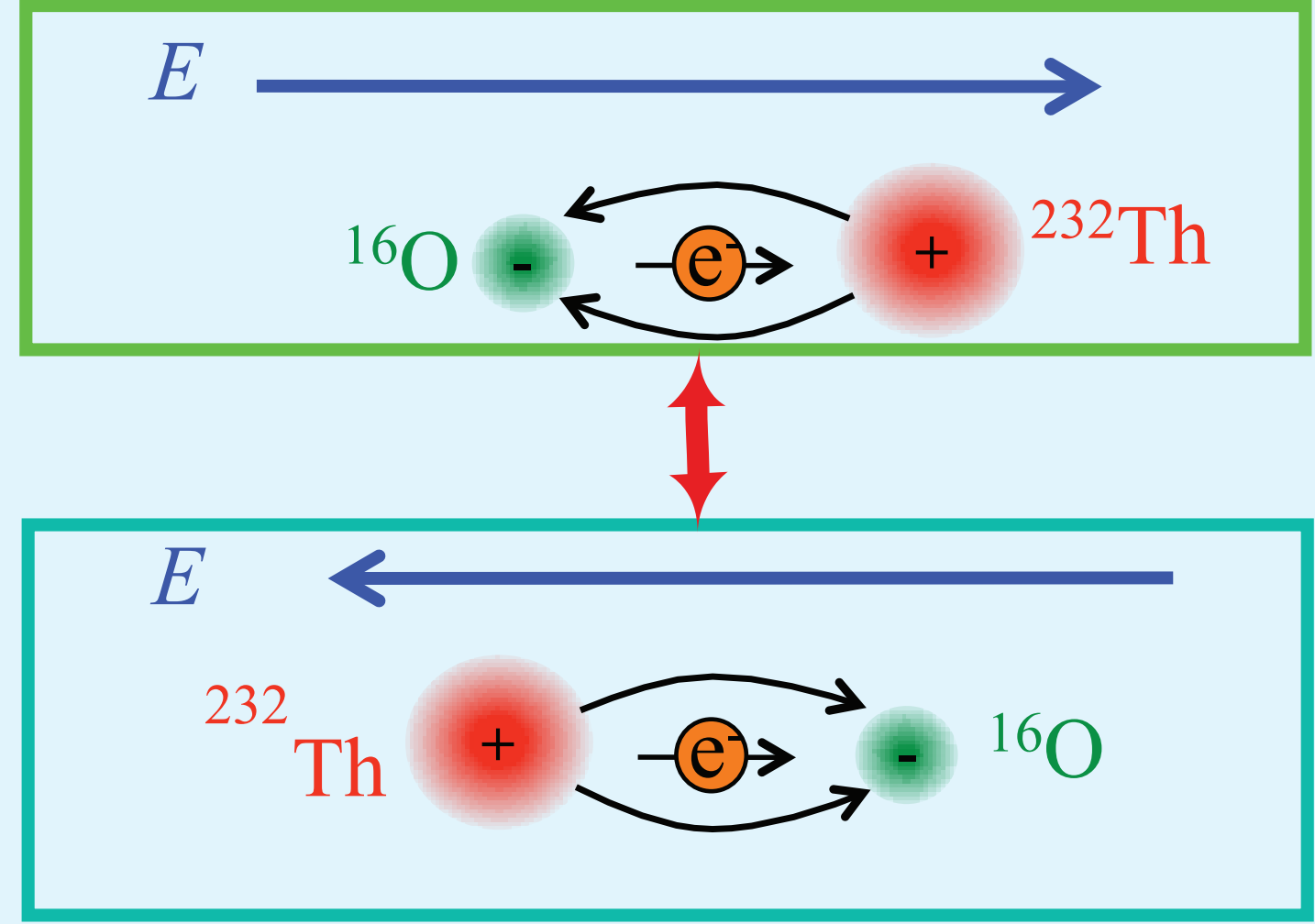
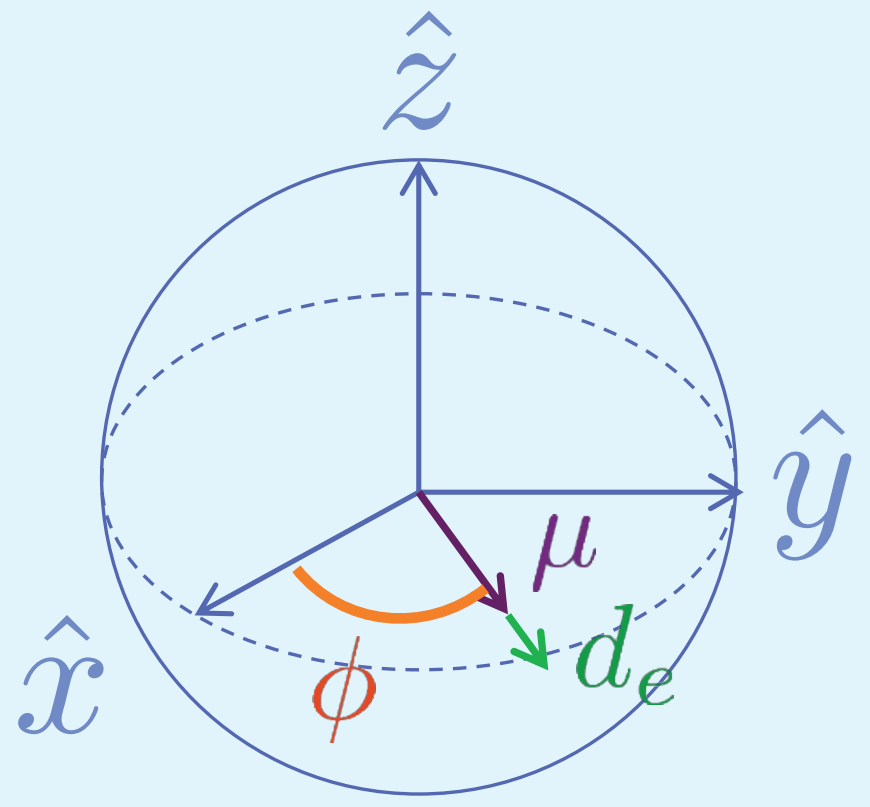


Determining eEDM, now with Efield Reversal

Measure the precession of the spin with E field reversal, subtract the two values

$$H = \mu_B \mathbf{B} \cdot \mathbf{J} + \mu_N \mathbf{B} \cdot \mathbf{I} + \mu_B \mathbf{B} \cdot \mathbf{L} + \mu_N \mathbf{B} \cdot \mathbf{S} + \mu_B \mathbf{B} \cdot \mathbf{J} + \mu_N \mathbf{B} \cdot \mathbf{I} + \mu_B \mathbf{B} \cdot \mathbf{L} + \mu_N \mathbf{B} \cdot \mathbf{S} + \dots + \text{terms due to applied fields} + \text{EDM terms}$$

We call this a "parity sum"



$$P_- = (4d_e E_{eff} + g\mu_b \mathbf{B}_{E\text{-corr}} + \eta\mu_b E_{nr} |\mathbf{B}| + \dots) \tau / \hbar$$

\mathbf{B} field correlated with \mathcal{E} reversal

Need $B_{E\text{-corr}} < 10^{-7}$ gauss, for 10^{-28} e-cm

$B_{E\text{-corr}}$ measured with fluxgates, $< 10^{-8}$ gauss sensitivity

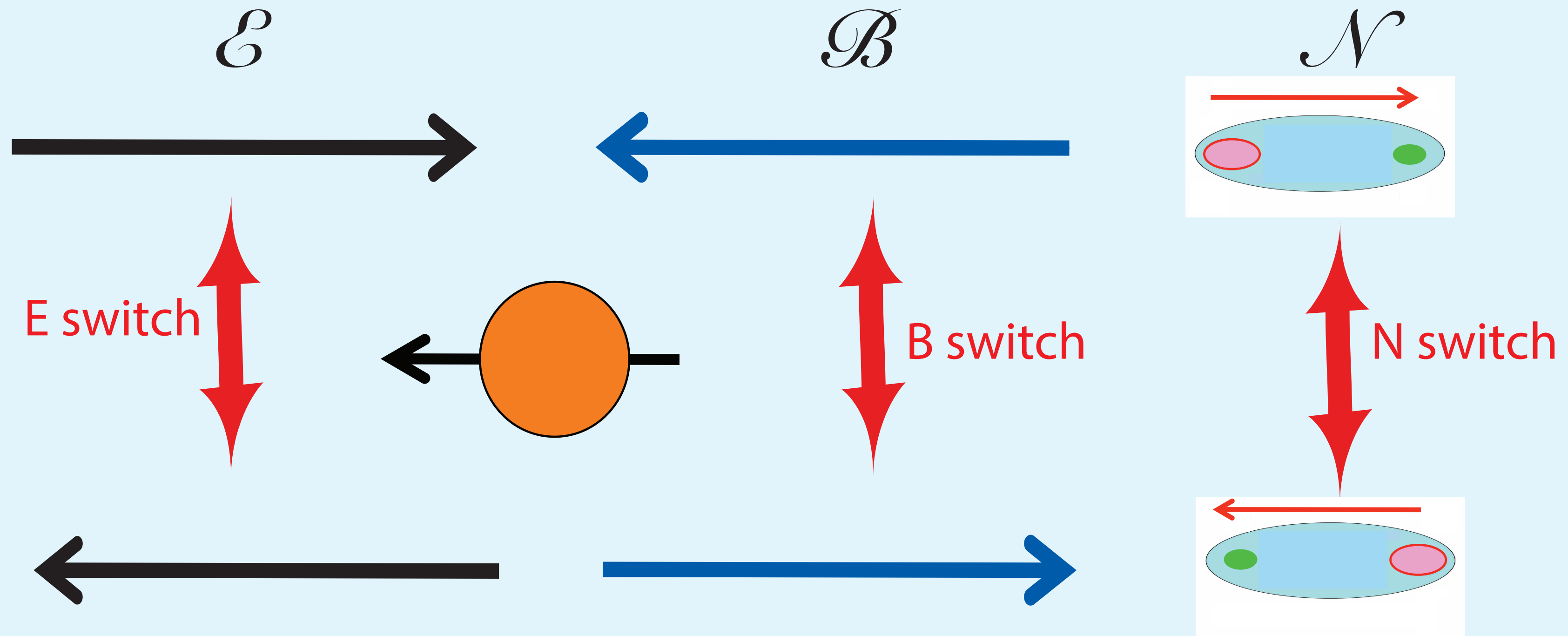
"Non-reversing \mathcal{E} field"

- g factor is linear in E, due to rotational mixing
- Need $E_{nr} < 1$ V/cm,
- E_{nr} measured by raman spectroscopy, ~ few mv/cm
- AND η determined from \mathcal{N} , \mathcal{B} reversals (1 nm/volt)

*\mathcal{E} field reversal alone could give greatly improved measurement of eEDM
But we don't do just \mathcal{E} reversals.....*



THREE reversals -- EIGHT combinations



These reversals together generate
8 parity sums

Each contain information about the system

ACME ThO EDM Experiment Apparatus

ThO Molecular Beam Path

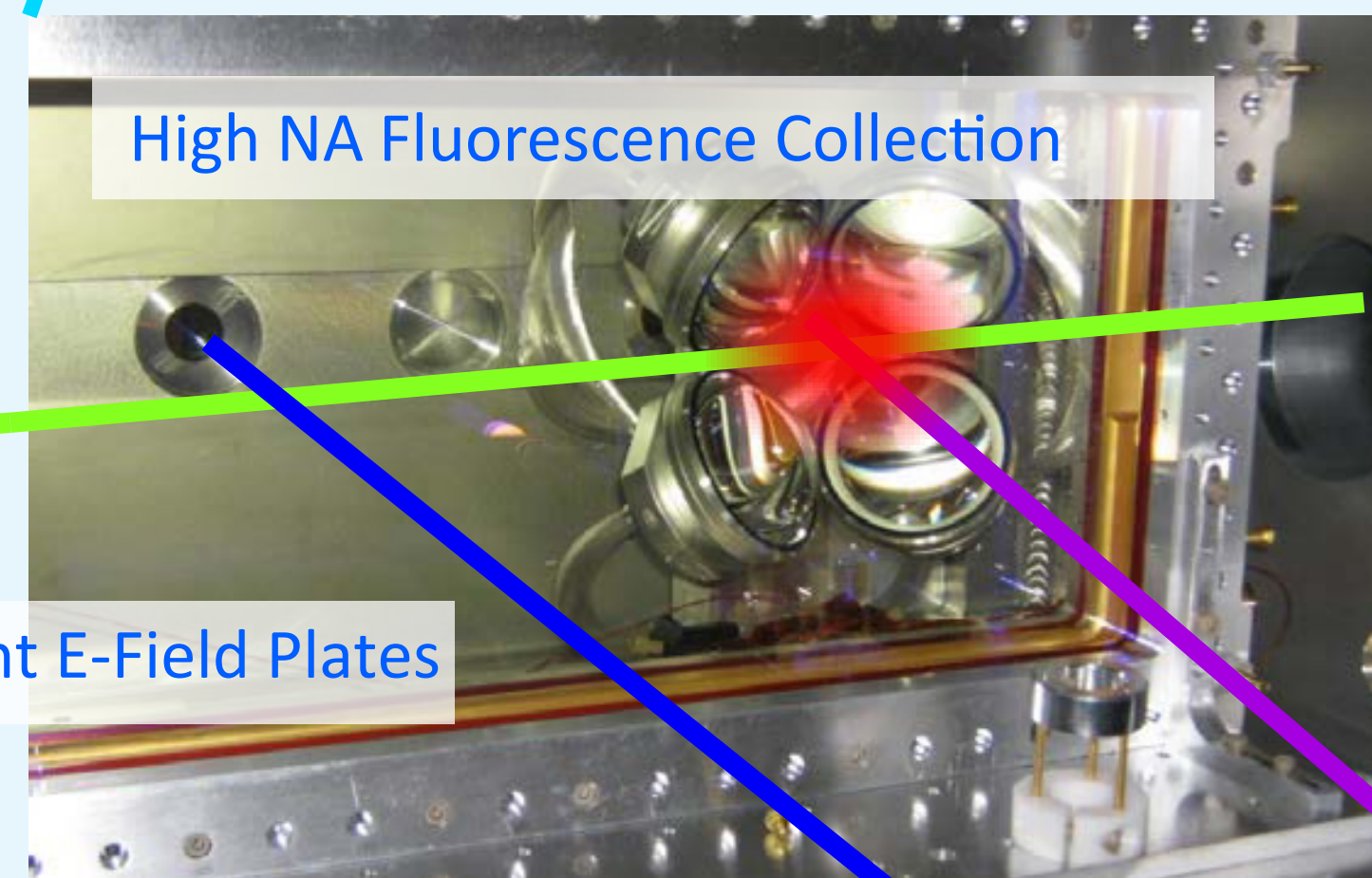
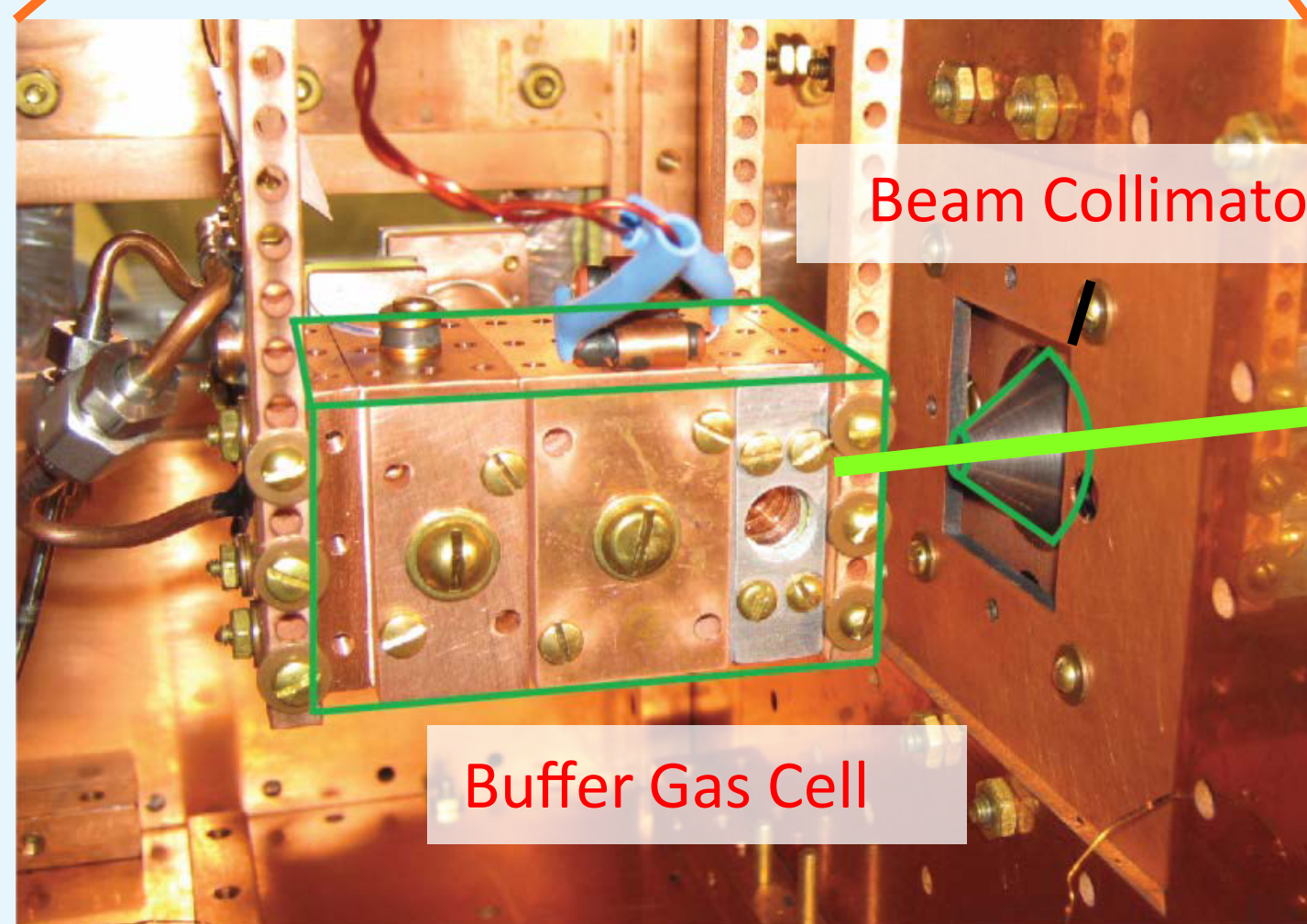
Pulsed Beam Source

- Buffer gas cooled
- Best source for reactive radicals

Ramsey Region

Probe Laser

Preparation Laser





Harvard University

ACME

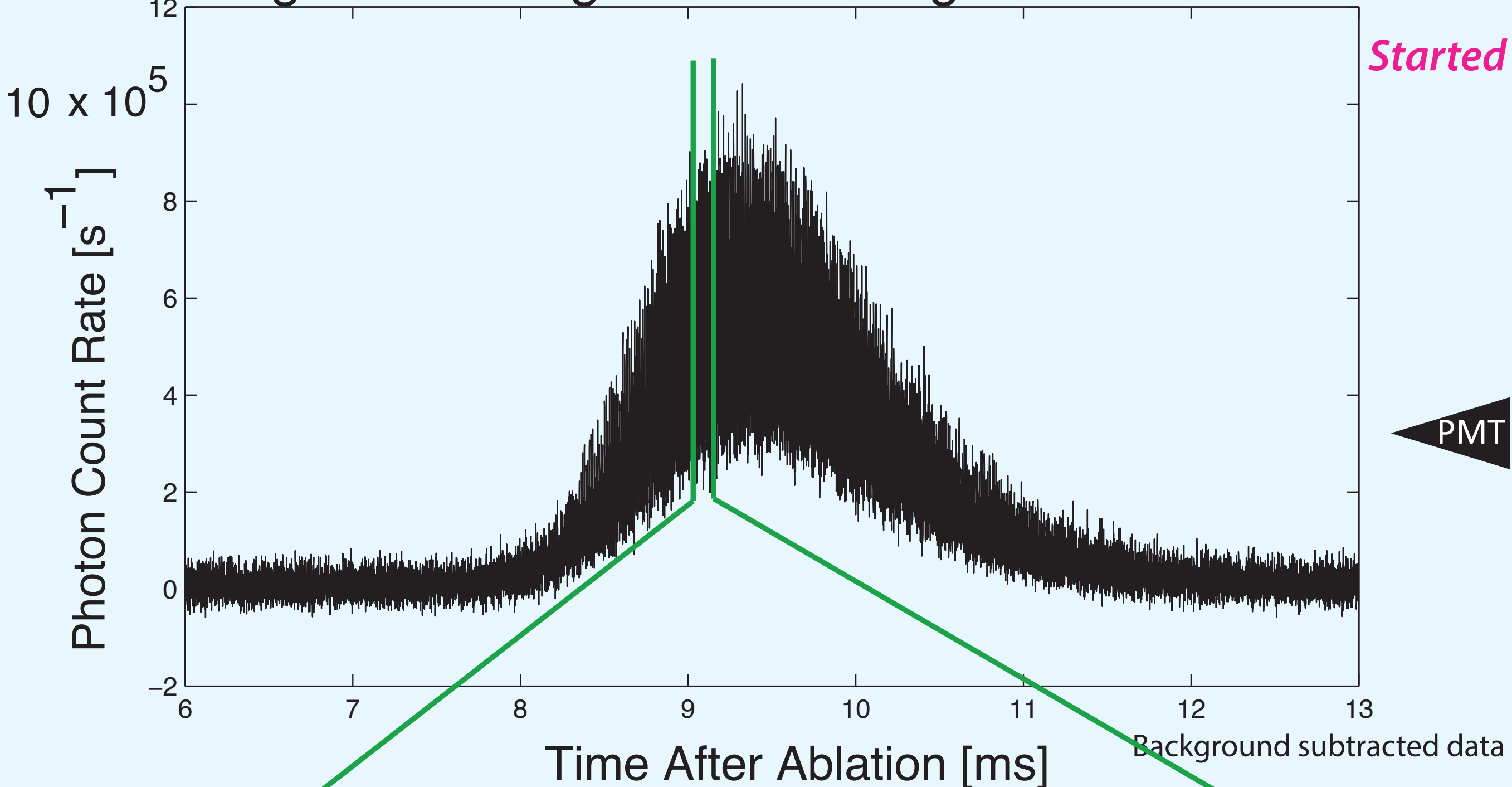
Advanced Cold Molecule Electron EDM



Yale University

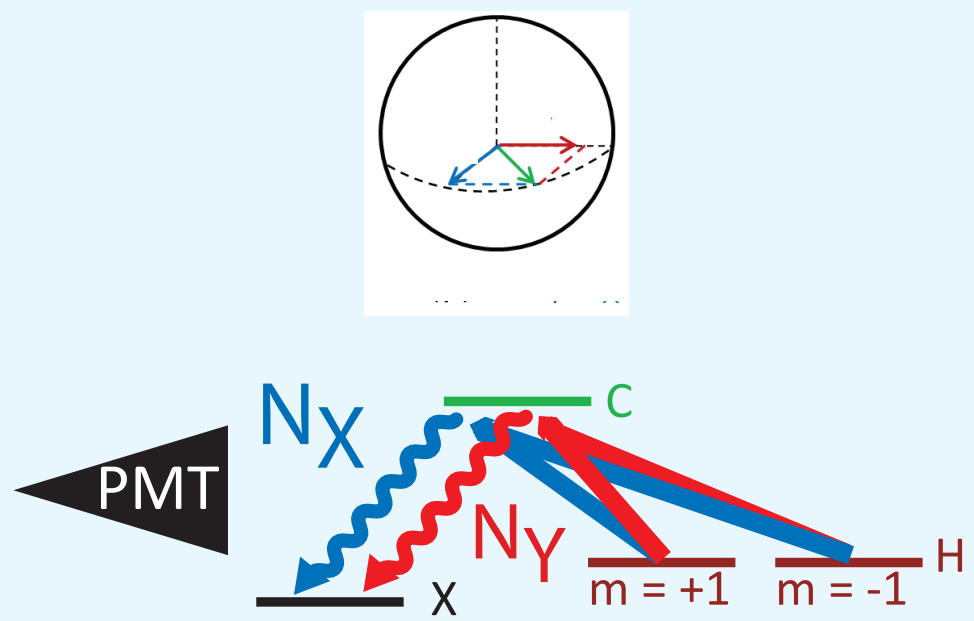
Typical EDM Data -- Pulses of Molecules Through the Ramsey Region

Single-block Signal Photon Signal from EDM Run

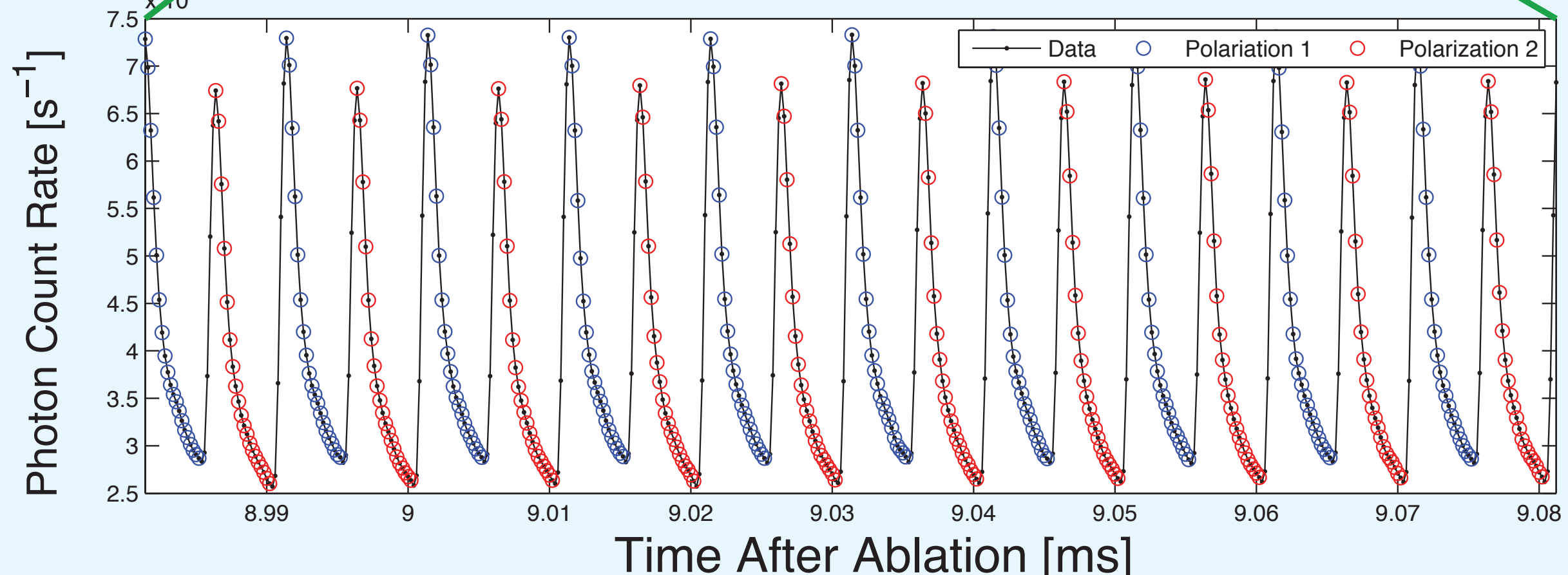


Started Taking EDM Data May 2012

Read out quantum state



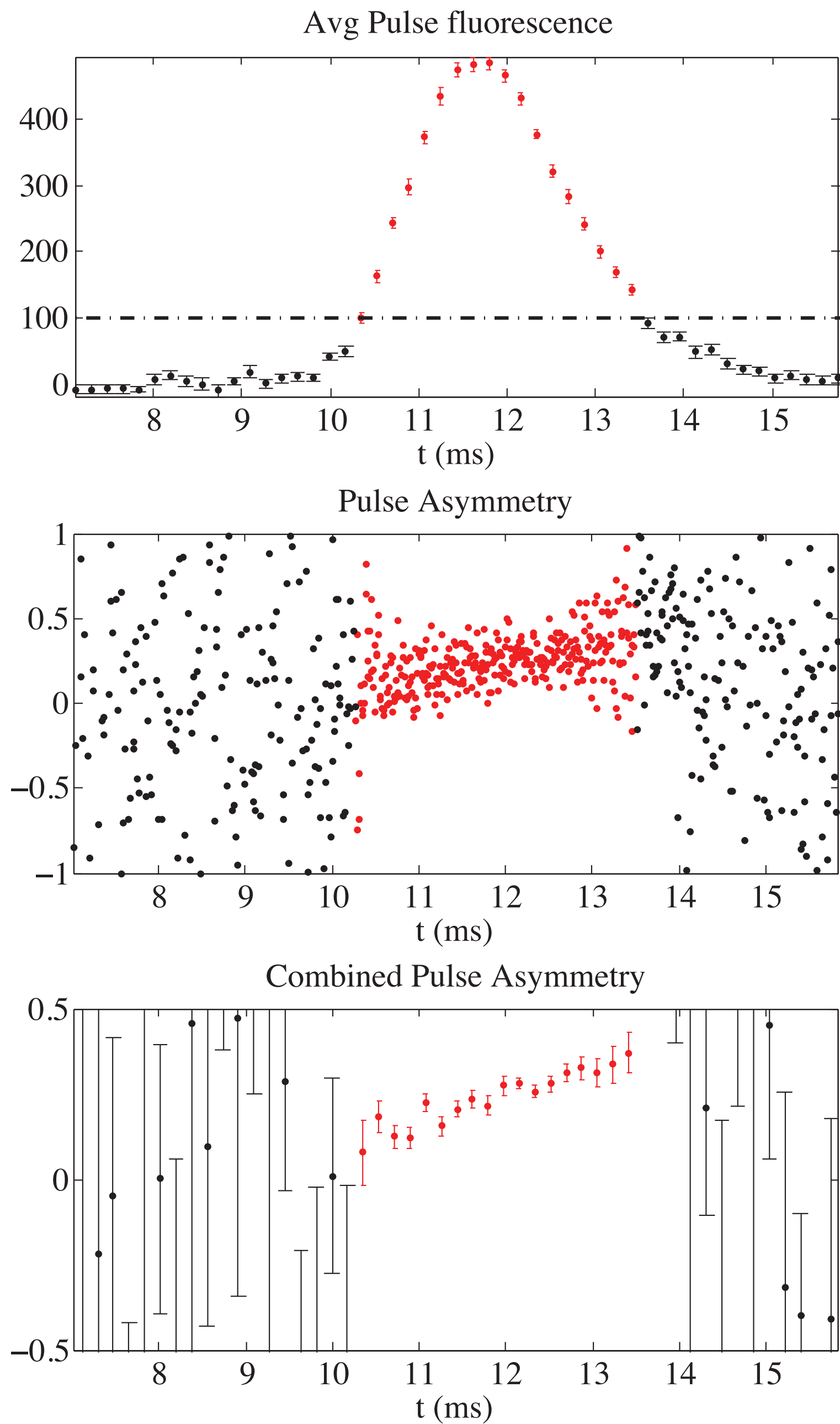
Zoom in Near Maximum Signal Area, Long time Average



Red and Blue indicate fluorescence collected under application of orthogonal polarizations of probe light.

Ratio of Red to Blue areas gives phase angle ϕ

Data Analysis



signal
threshold cut
ensures
gaussian
asymmetry

binning of
asymmetries
(18-30 pts)

compute the standard error in the mean, and
use standard error propagation to obtain the
result.

- Computation of Uncertainty
- Uncertainty propagation

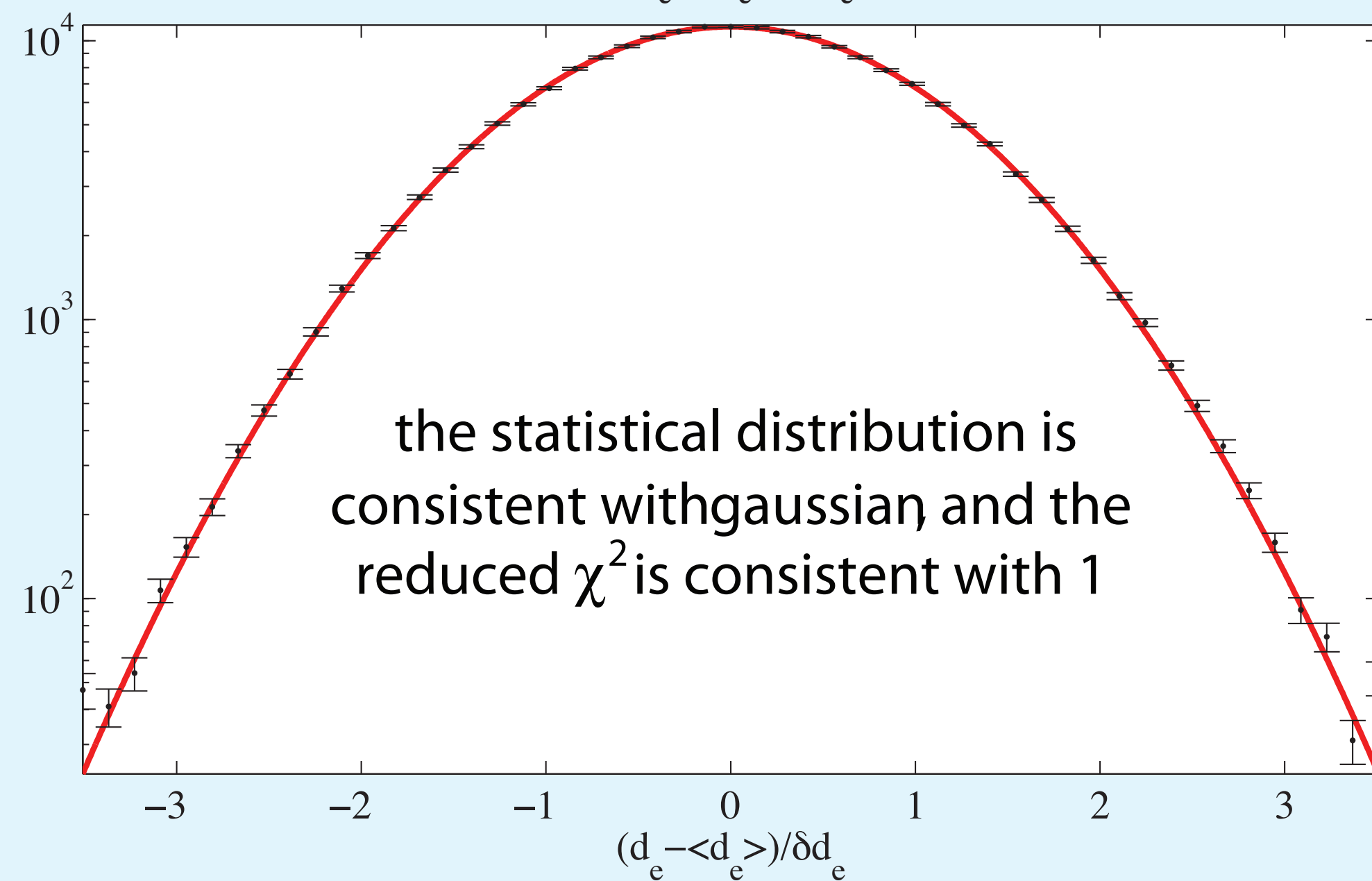
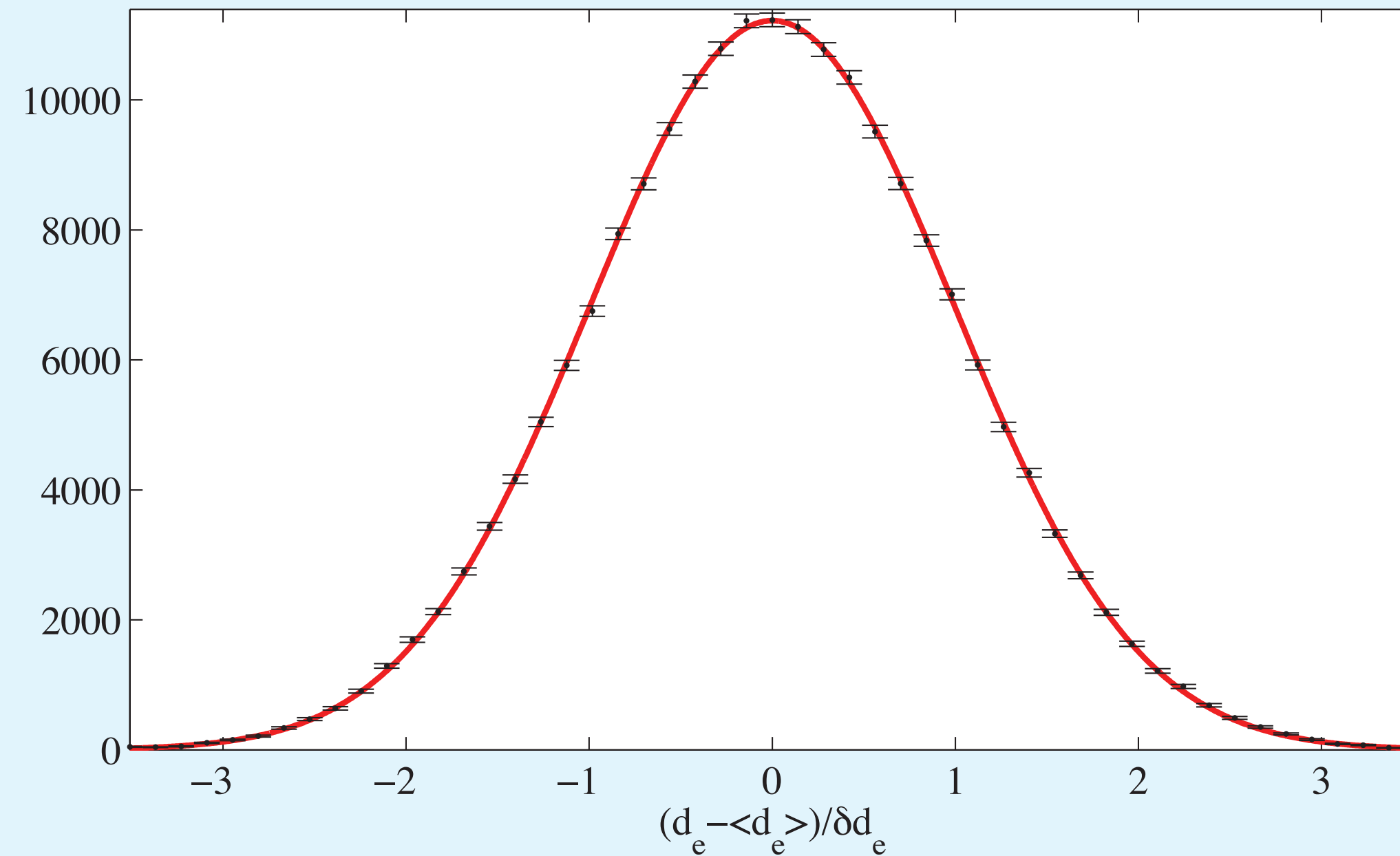
• Statistical Distribution

• Blind Offset

• Feldman-Cousins
Confidence Intervals

Data Analysis

Sigma=0.999±0.002, Center=0.000±0.002



• Computation of Uncertainty

• Uncertainty propagation

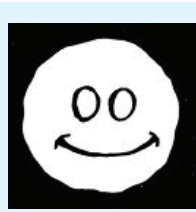
• Statistical Distribution

• Blind Offset

• Feldman-Cousins Confidence Intervals

the uncertainty is limited by photon shot noise in the photodetectors

$$\delta\omega^{NE} \approx 1.15 \times \frac{1}{C\tau \sqrt{N}}$$



Statistical Sensitivity to the eEDM

Shot noise limited error in measurement of the electron EDM using our method

$$\delta d_e = \frac{1}{2 E_{eff}} \frac{\hbar}{\tau \sqrt{\dot{N} T}}$$

Effective
Electric Field

Coherence
Time

(Photon)
Counting Rate

Integration
Time

↪ = (molecular flux) X (detection efficiency)

Statistics OK

$\pm 4 \times 10^{-29}$ e-cm

best previous measurements $d_e < 1 \times 10^{-27}$ e-cm

Statistics OK

$\pm 4 \times 10^{-29}$ e-cm

best previous measurements $d_e < 1 \times 10^{-27}$ e-cm

Got Systematics ?

Systematic Error Search - I

Turn Knobs -- see what happens

These Knobs Should NOT change EDM, but let's check!

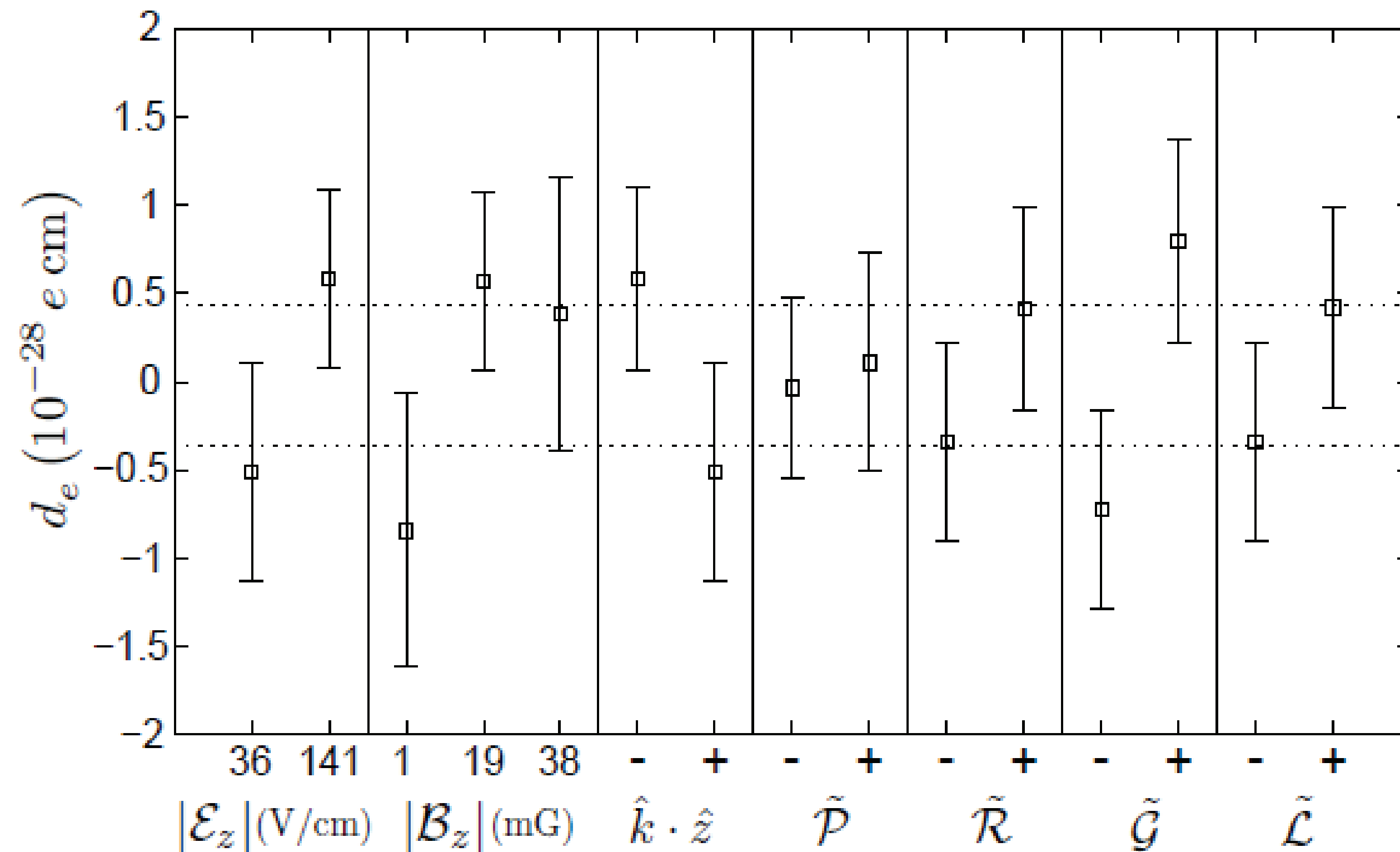
Extra "Switches"

TWO Different
Pump-probe

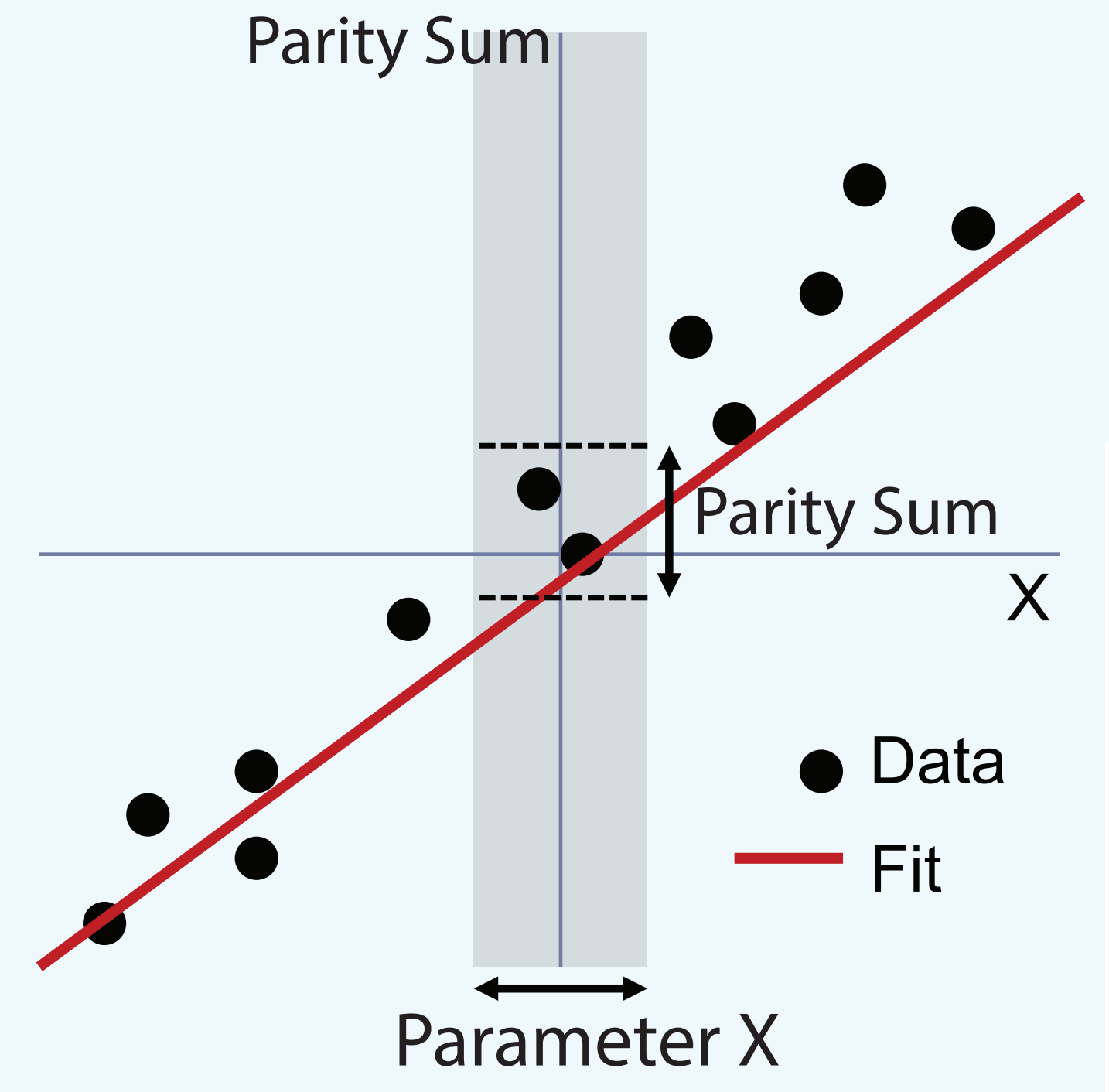
- relative polarizations
- global polarizations

Efield plate lead positions
Laser propagation directions
Probe upper states
Electric field (E) magnitudes

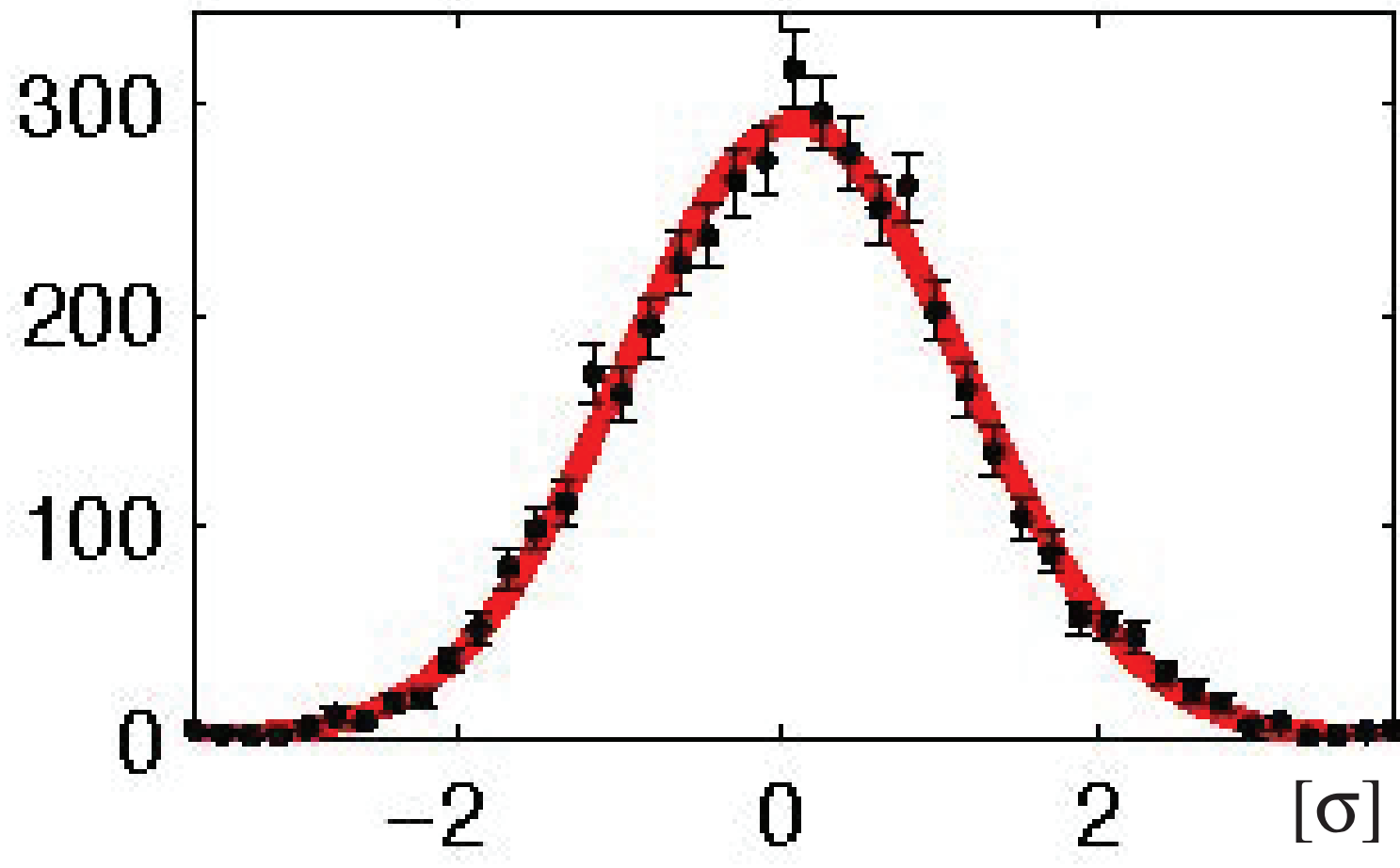
THREE different
Magnetic field (B) values



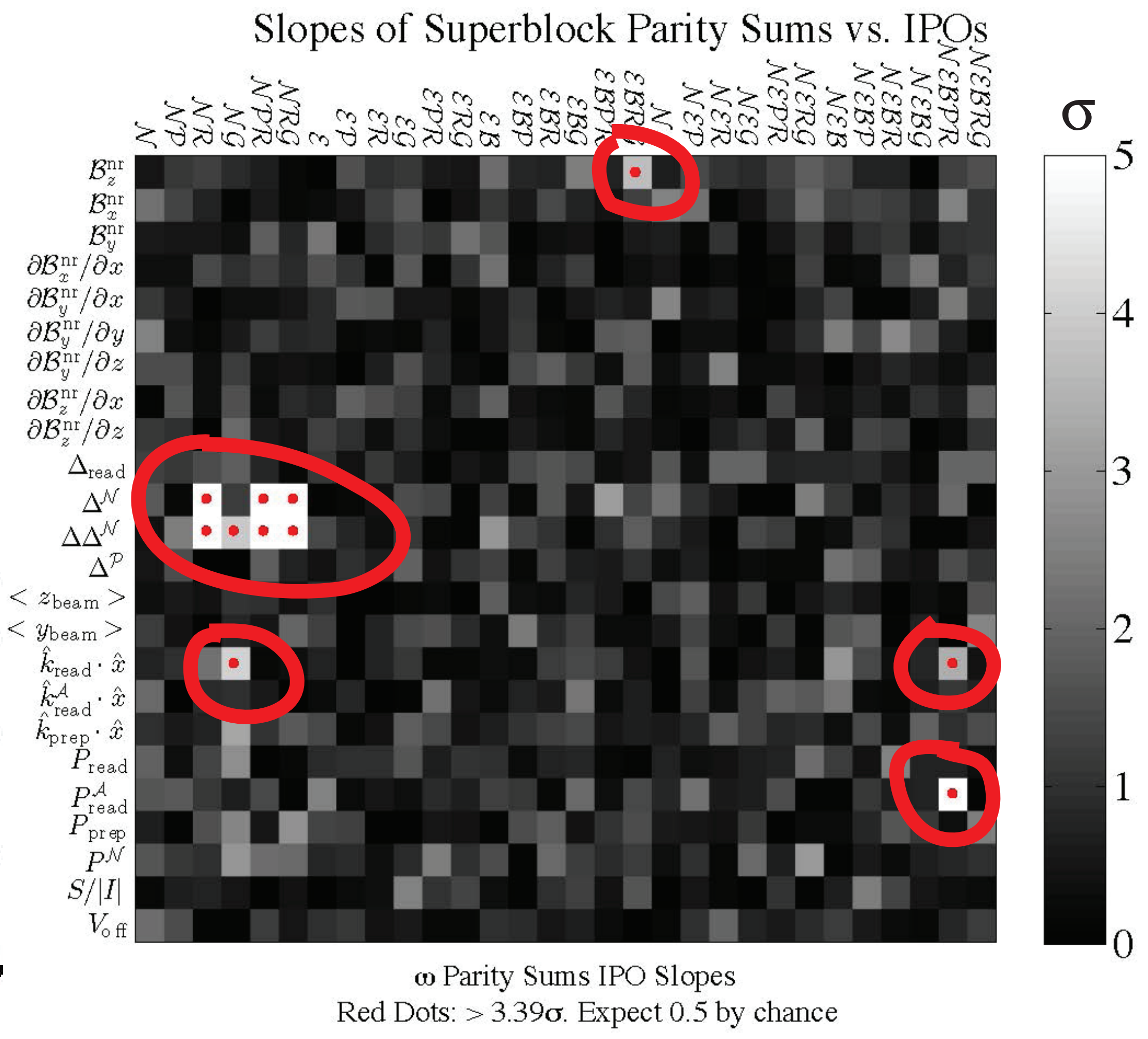
Systematic Checks - III, Correlation Search Method



Statistical Distribution of Diagnostic Signals from Systematic Error Search



Pixel Plot Identifies Correlations

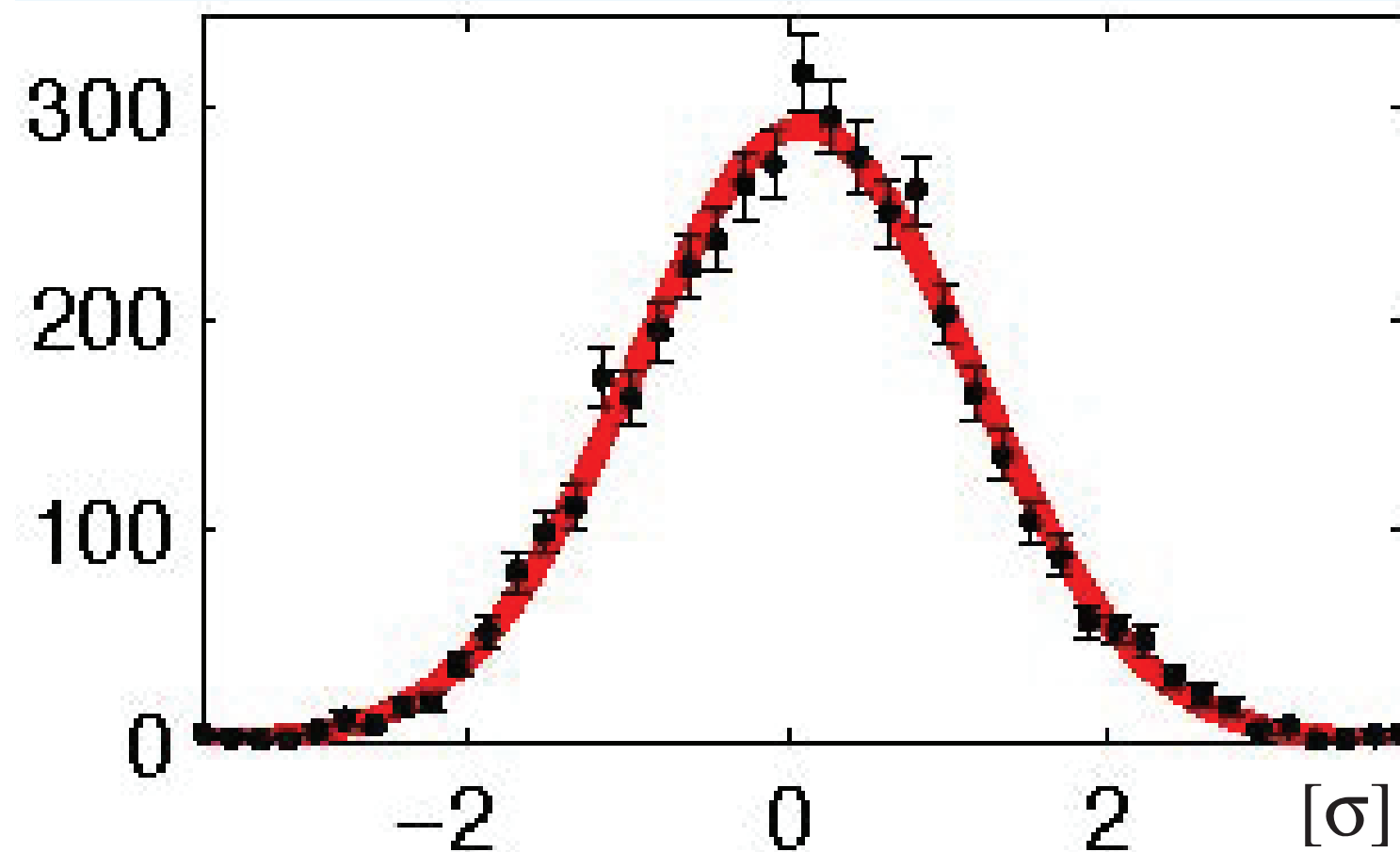


Systematic Checks - III, Correlation Search Method

Few "outliers",
one "switch" (or so)
from the EDM channel

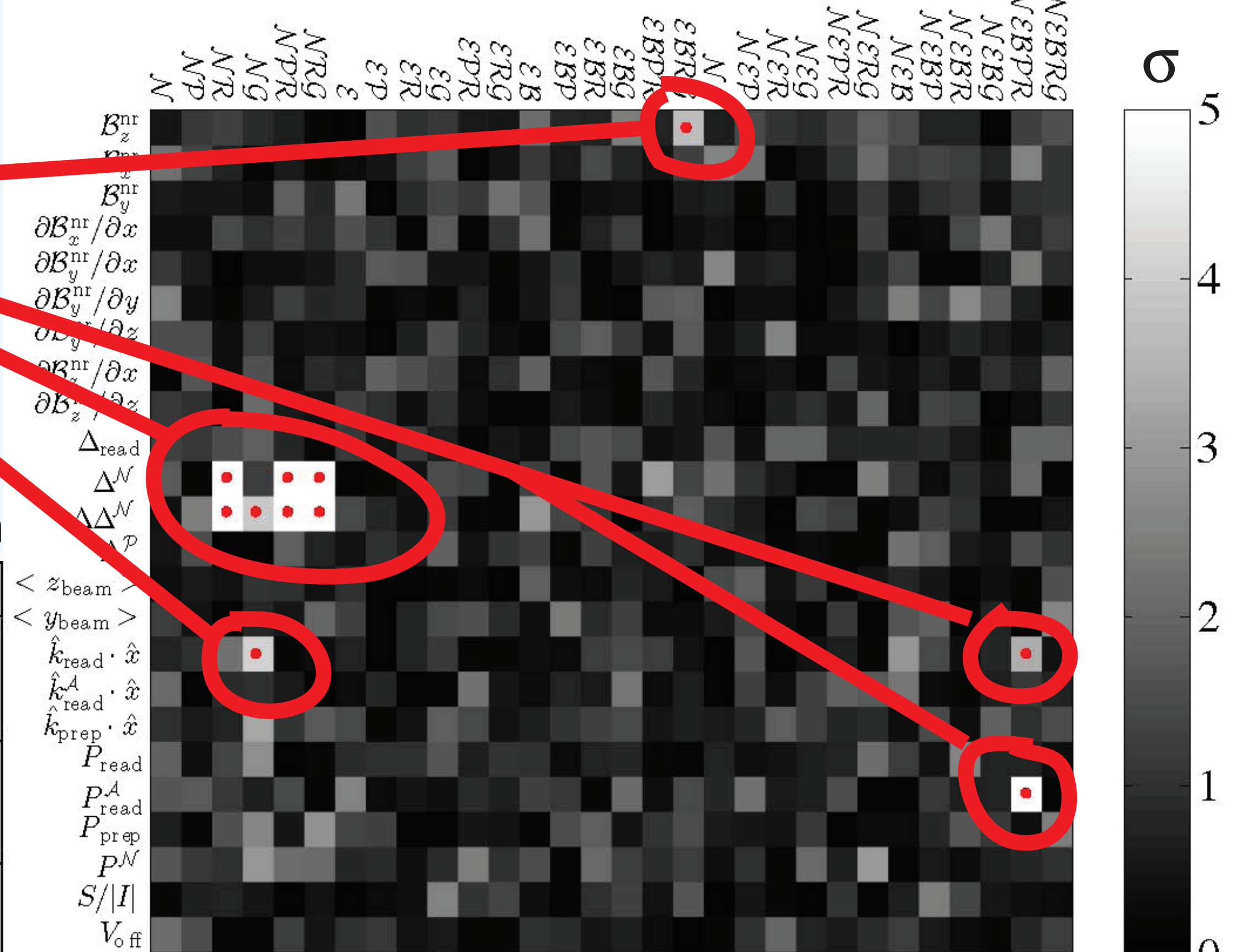
All understood and/or
controlled

Statistical Distribution of Diagnostic
Signals from Systematic Error Search



Pixel Plot Identifies Correlations

Slopes of Superblock Parity Sums vs. IPOs



Parity Sums IPO Slopes
Red Dots: $> 3.39\sigma$. Expect 0.5 by chance

EDM Measurement Error Budget Sheet

Systematic and statistical errors for rate of angular precession of the electron spin in units of mrad/second.

$$1 \text{ mrad/s} \sim 10^{-29} \text{ e cm}$$

Parameter	Shift	Uncertainty
\mathcal{E}^{nr} correction	-0.81	0.66
$\Omega_r^{\mathcal{N}\mathcal{E}}$ correction	-0.03	1.58
$\phi^{\mathcal{E}}$ -correlated effects	-0.01	0.01
$\phi^{\mathcal{N}}$ correlation		1.25
Nonreversing B -field (B_z^{nr})		0.86
Transverse B -fields ($B_x^{\text{nr}}, B_y^{\text{nr}}$)		0.85
B -field gradients		1.24
Prep./read laser detunings		1.31
$\tilde{\mathcal{N}}$ -correlated detuning		0.90
\mathcal{E} -field ground offset		0.16
Total systematic	-0.85	3.24
Statistical		4.80
Total uncertainty		5.79

Systematics

Statistics

$$d_e = \left(\pm 3.7_{\text{stat}} \pm 2.5_{\text{syst}} \right) \times 10^{-29} \text{ e}\cdot\text{cm}$$

using $E_{\text{eff}} = 84 \text{ GV/cm}$, calculated by Skripnikov, Petrov and Titov JCP (2013) and Meyer and Bohn PRA (2008)

EDM Measurement Error Budget Sheet

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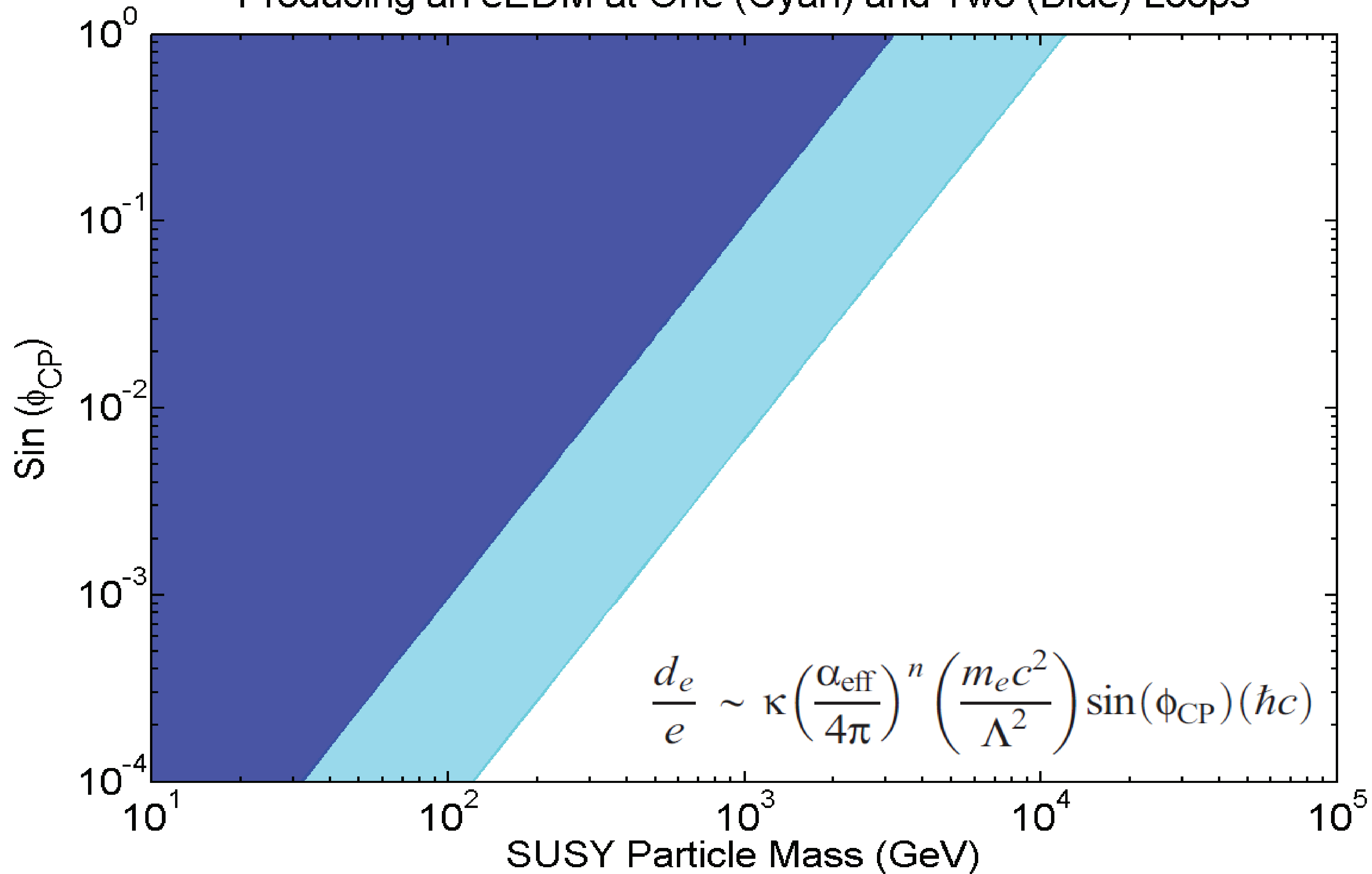
Systematics

Statistics

$$d_e = (-2.1 \pm 3.7_{\text{stat}} \pm 2.5_{\text{syst}}) \times 10^{-29} \text{ e}\cdot\text{cm}$$

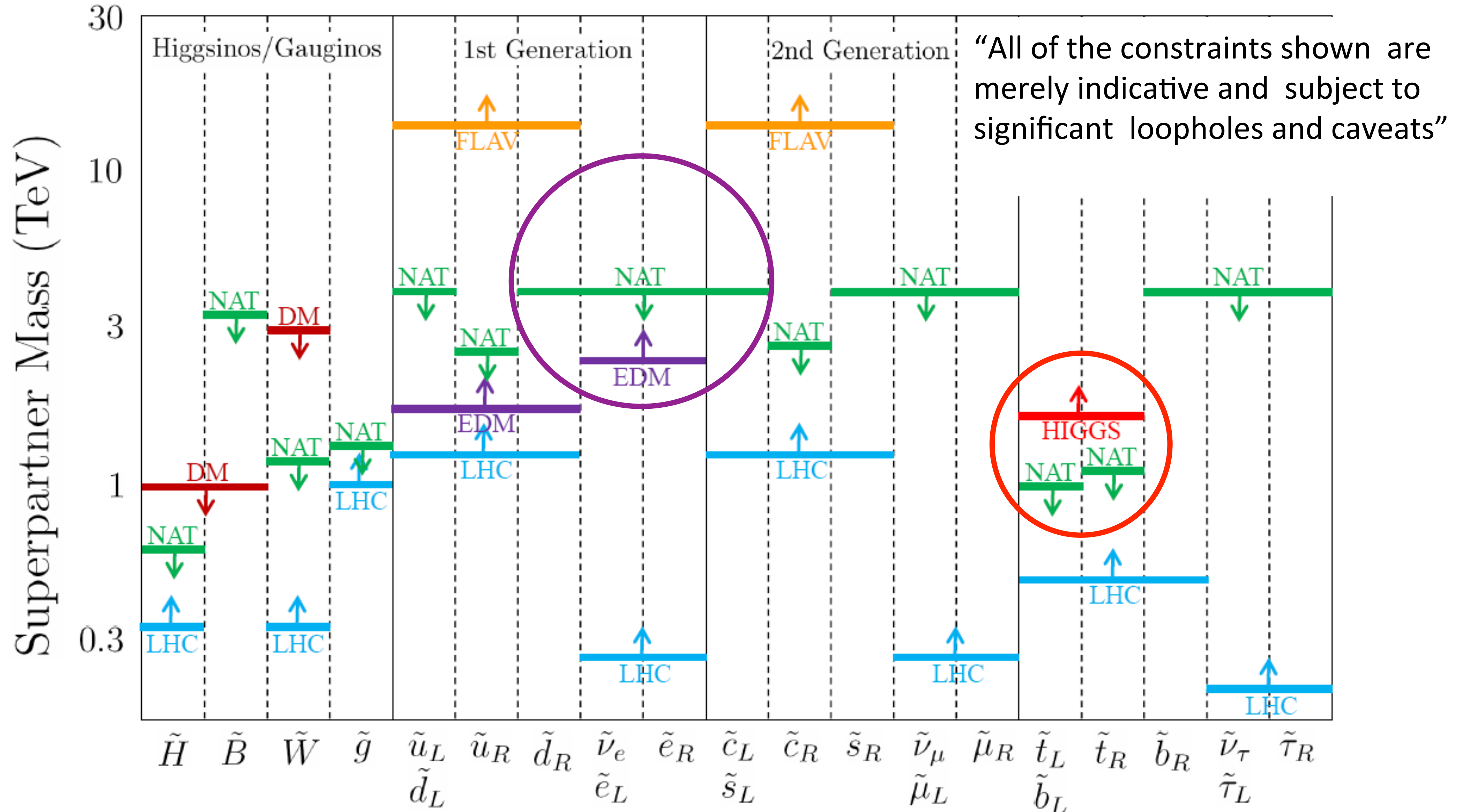
using $E_{\text{eff}} = 84 \text{ GV/cm}$, calculated by Skripnikov, Petrov and Titov JCP (2013) and Meyer and Bohn PRA (2008)

ACME Excluded Regions for Generic SUSY Model
Producing an eEDM at One (Cyan) and Two (Blue) Loops



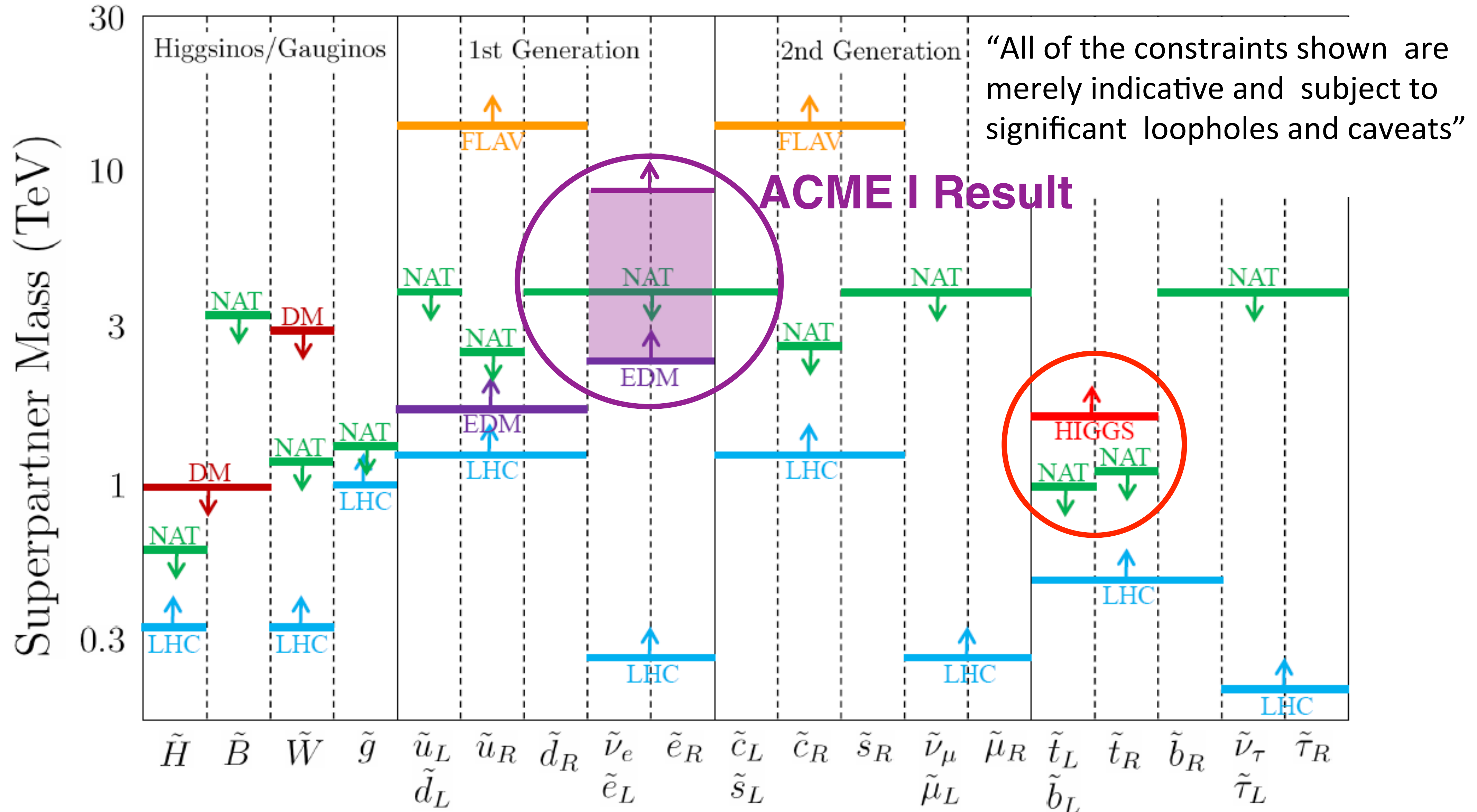
What does our limit mean for particle physics?

J. Feng: "Naturalness and the status of SUSY", Ann. Rev. Nucl. Part. Sci. (2013)

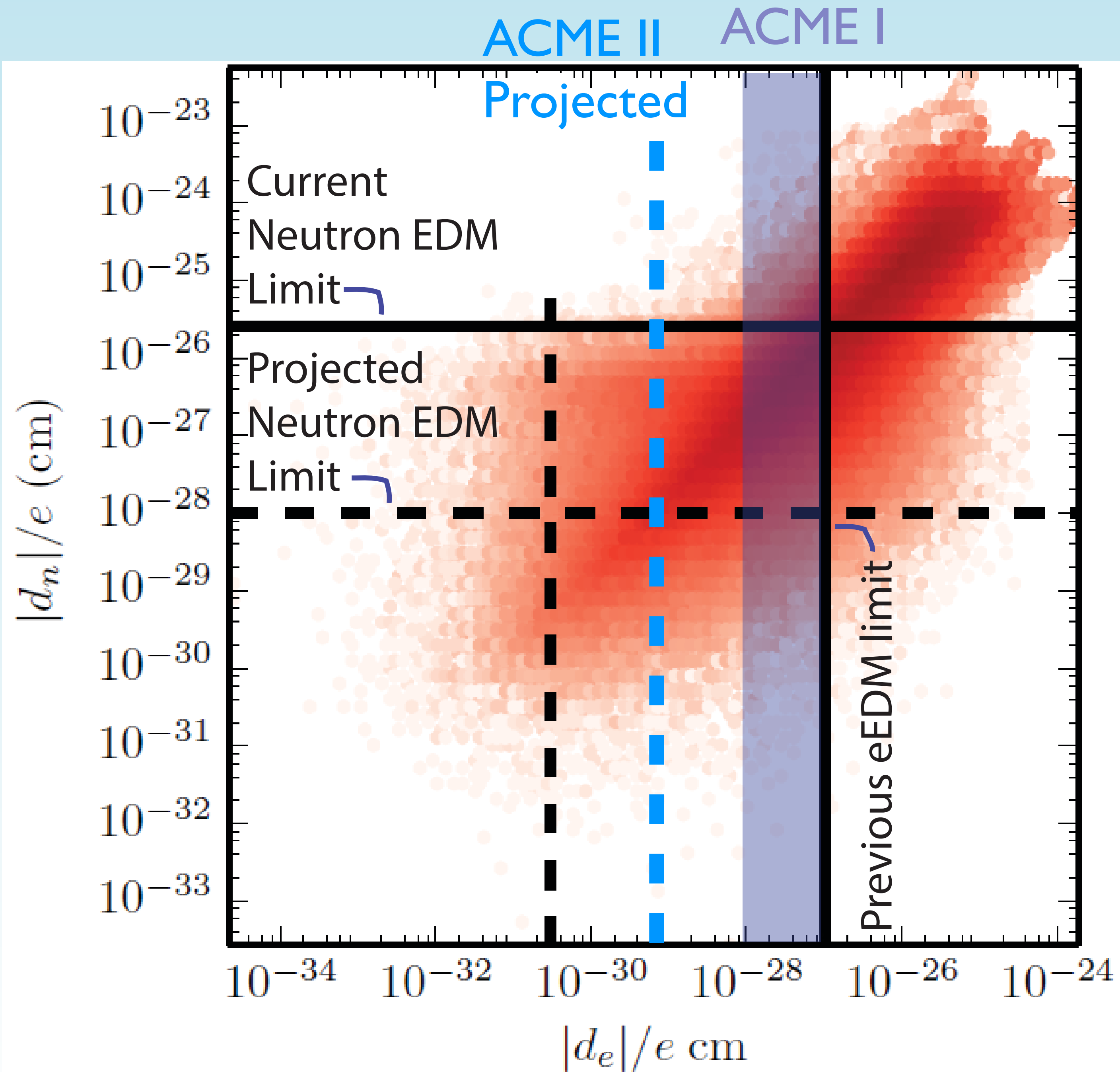


What does our limit mean for particle physics?

J. Feng: "Naturalness and the status of SUSY", Ann. Rev. Nucl. Part. Sci. (2013)



MSSM Parameter Scatter Plot

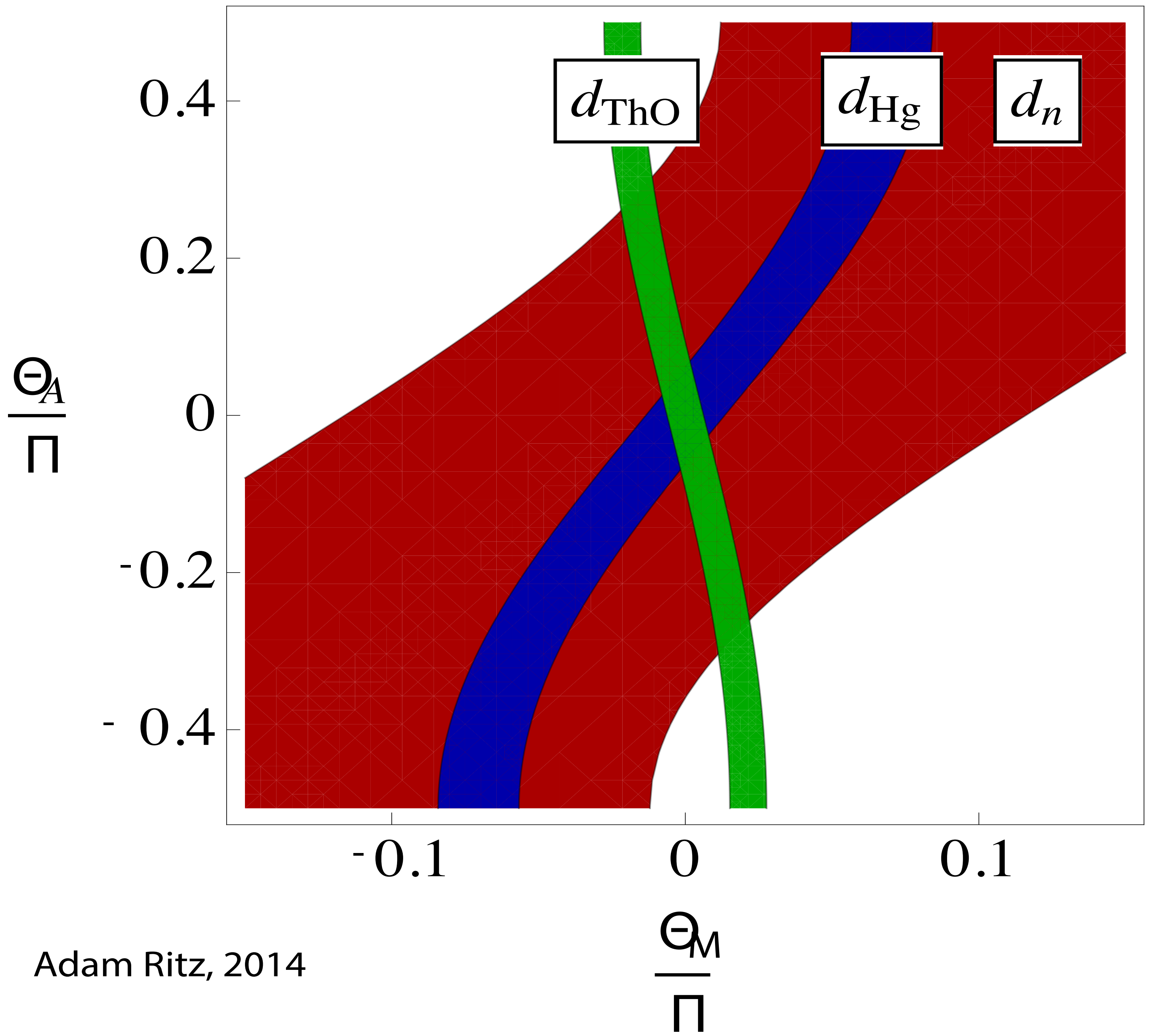


Color Scale indicates 10^4 Range of "Probability Density"

2/3 of parameter space excluded by *previous* EDM limits
ACME I excluded 2/3 of Remaining

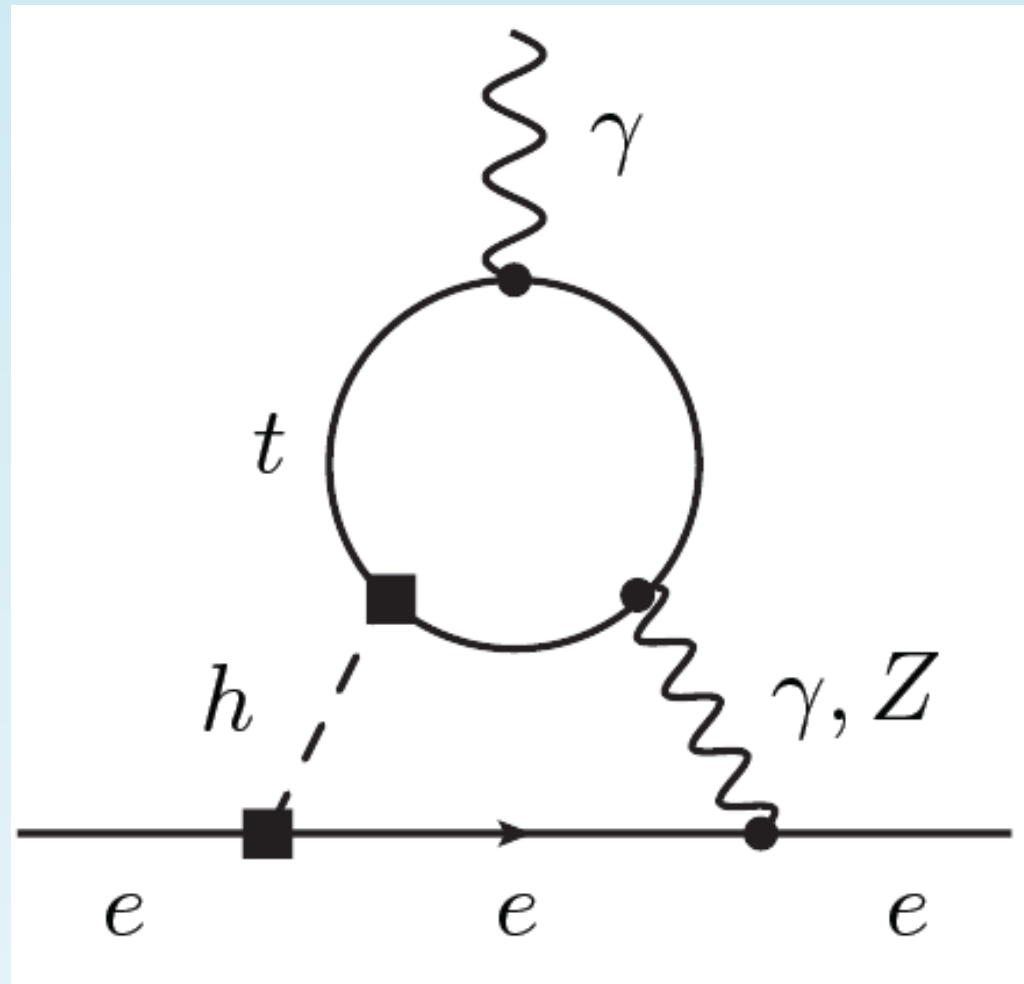
Flat log distribution of SUSY CP Phases

J. Berger et al., 1309.7653 (2013)



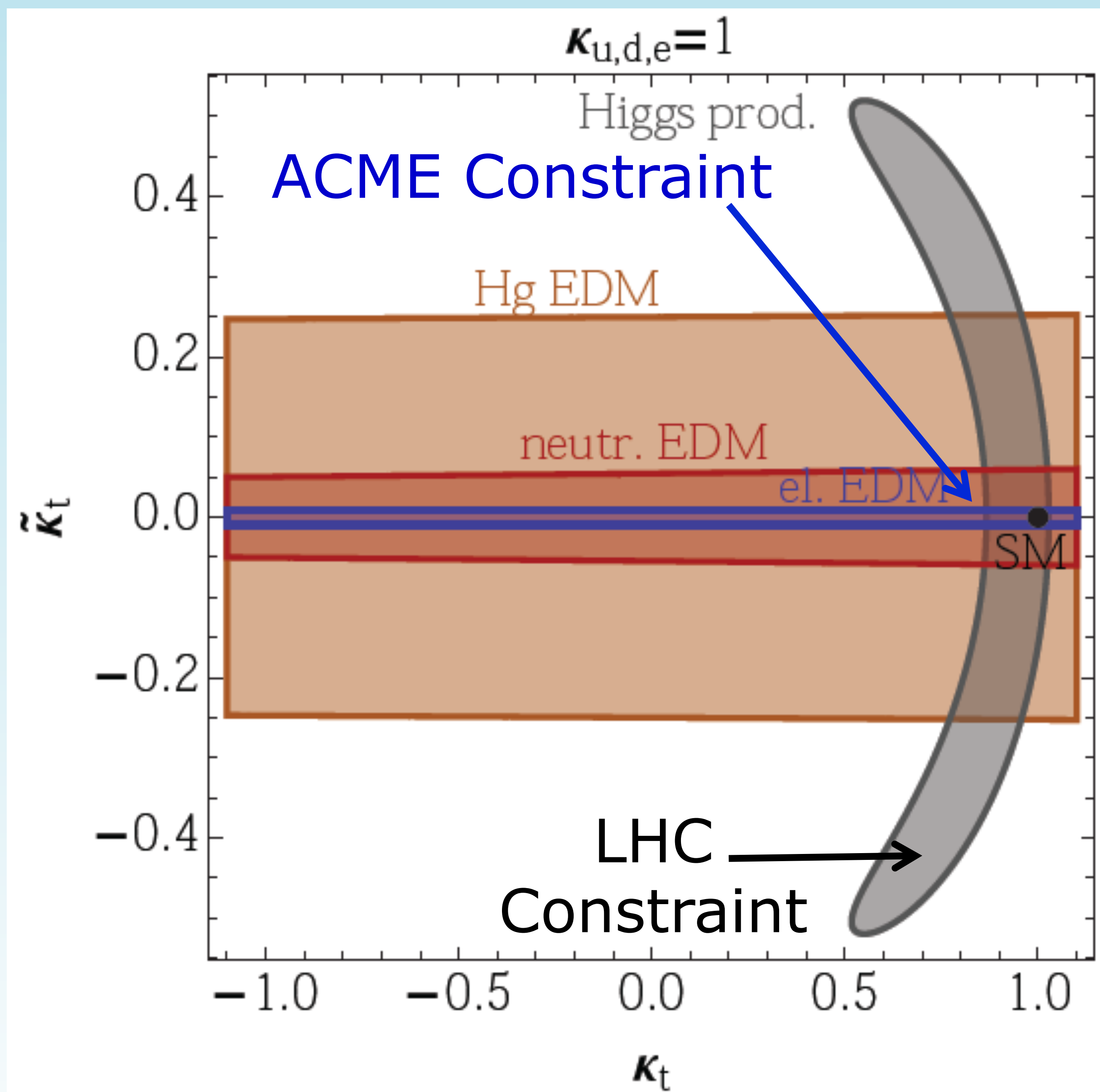
Adam Ritz, 2014

Diagrams with **known** SM particles can rule out non-SM couplings



Example:
CP-violating
Higgs-top coupling

Brod *et al.*,
arXiv: 1310.1385

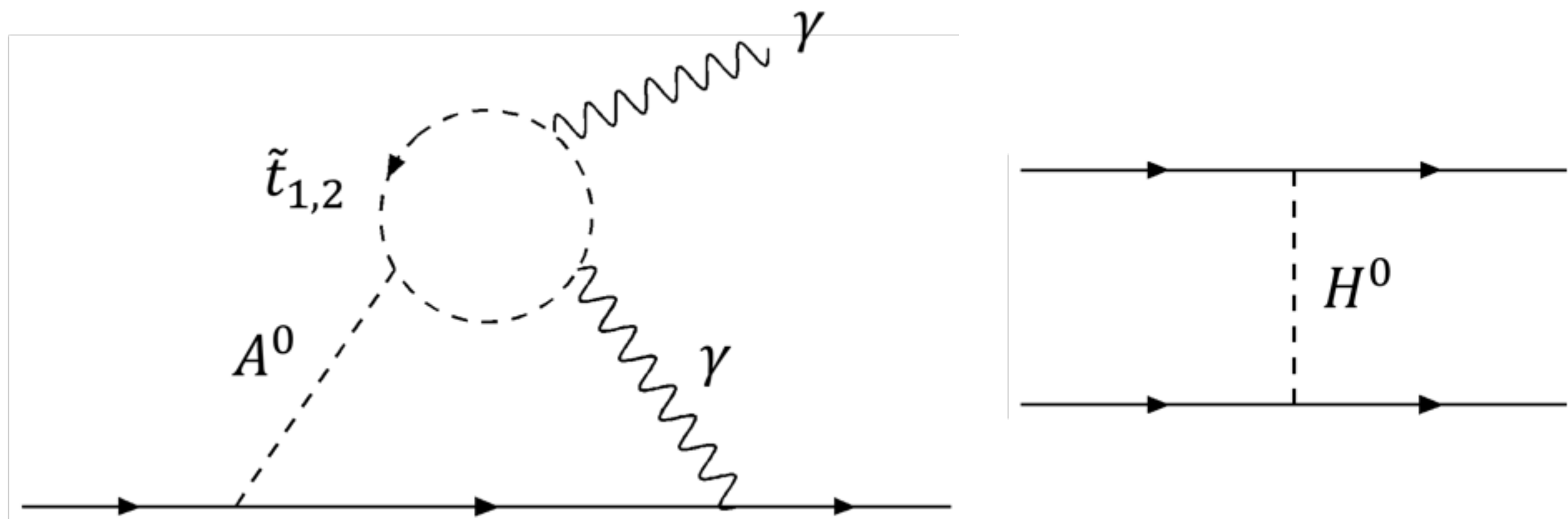


CP-odd/CP-even Higgs-top coupling $< 1\%$ from ACME

Present bounds on stop masses in the MSSM

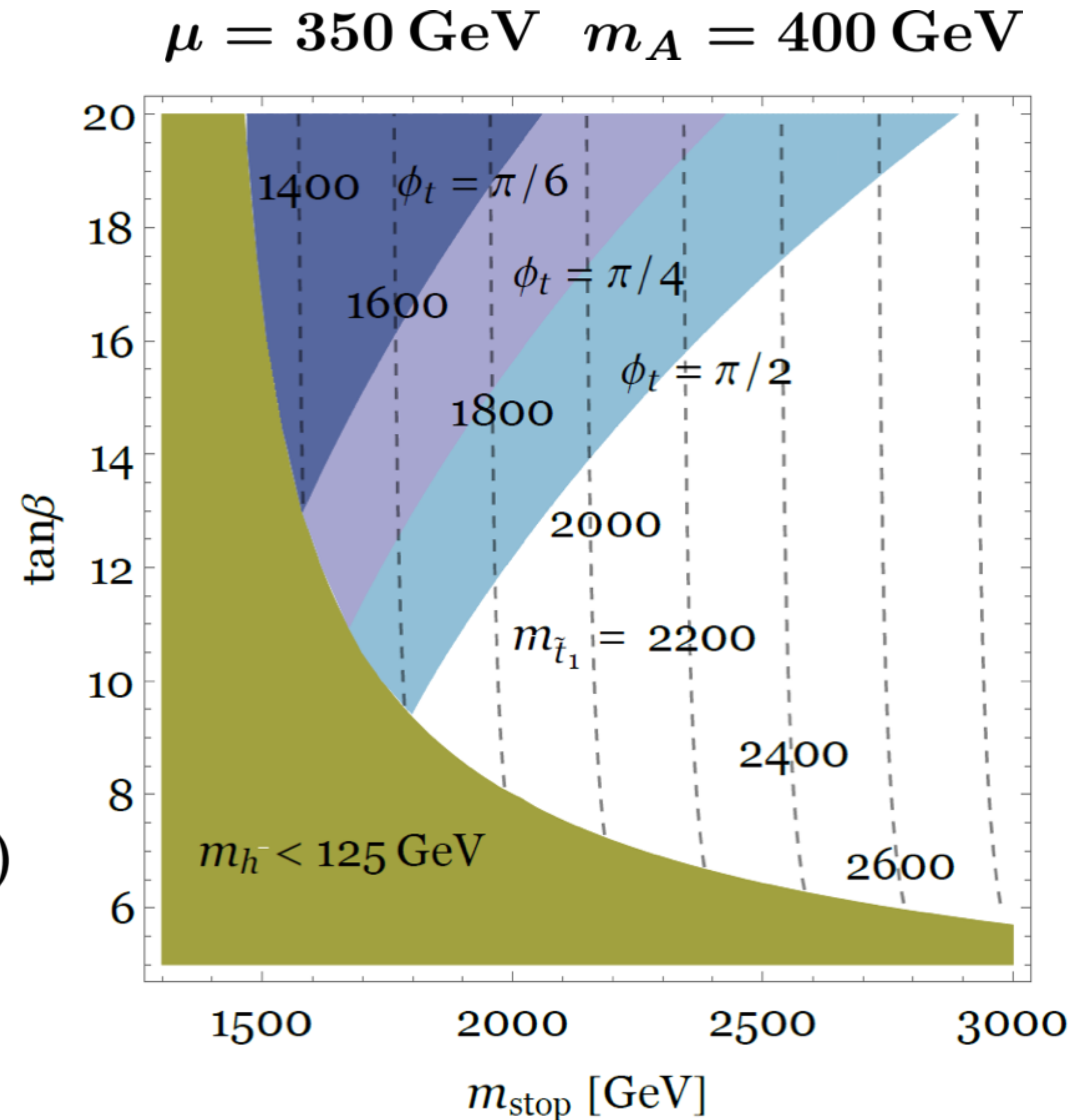
Yuichiro Nakai and Matthew Reece: **preliminary!**

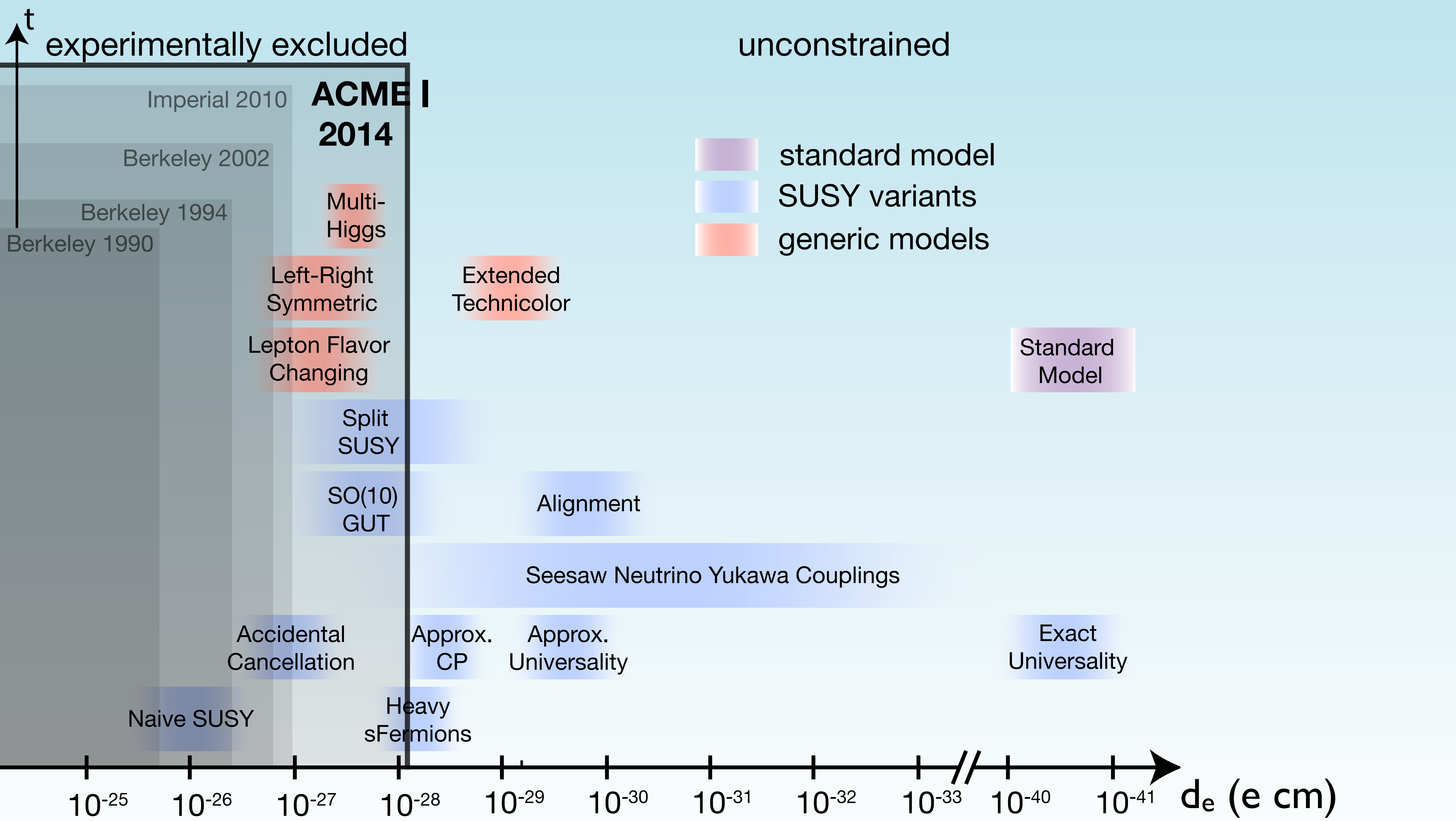
Superpartners of top quarks (stops), essential for the natural EWSB, can generate a sizable EDM.



For the maximal CP phase $\phi_t = \arg(A_t \mu)$

$$m_{\tilde{t}_1} > 1.6 \text{ TeV}$$





What led to statistical sensitivity improvements in ThO over YbF?

1) ThO (change of molecule)

÷1.5 ÷2 radiative decay of metastable EDM state
x6 x6 higher effective field
2.5 6.5 lower state preparation
= x2 better

2) Hydrodynamic Buffer Gas Beam Source (New Source)

÷3.7 x14 useful molecular flux
x1.7 x1.7 coherence time
= x6.5 better

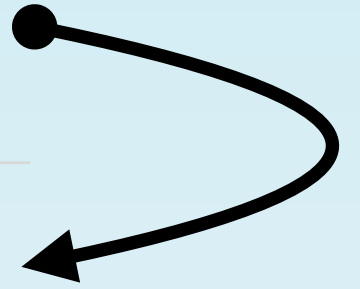
3) Technical Changes

÷1.2 1.5 longer length
x1.7 x3 higher collection efficiency
÷1.4 ÷2 shorter running time
x1.4 x2 optical rotational cooling
=x1.5 better

x15, better overall statistical sensitivity

Experiment	One Day Statistical Sensitivity e-cm day ^{-1/2}	Published Limit $ d_e <$ in e-cm
------------	---	--------------------------------------

Berkeley Tl	0.5×10^{-27}	1.6×10^{-27}
Imperial YbF	2×10^{-27}	1.5×10^{-27}
ACME I ThO	1×10^{-28}	0.9×10^{-28}



Improvement 1	Improvement 2	Improvement 3
---------------	---------------	---------------

EDM Sensitivity Gain over Previous Experiment

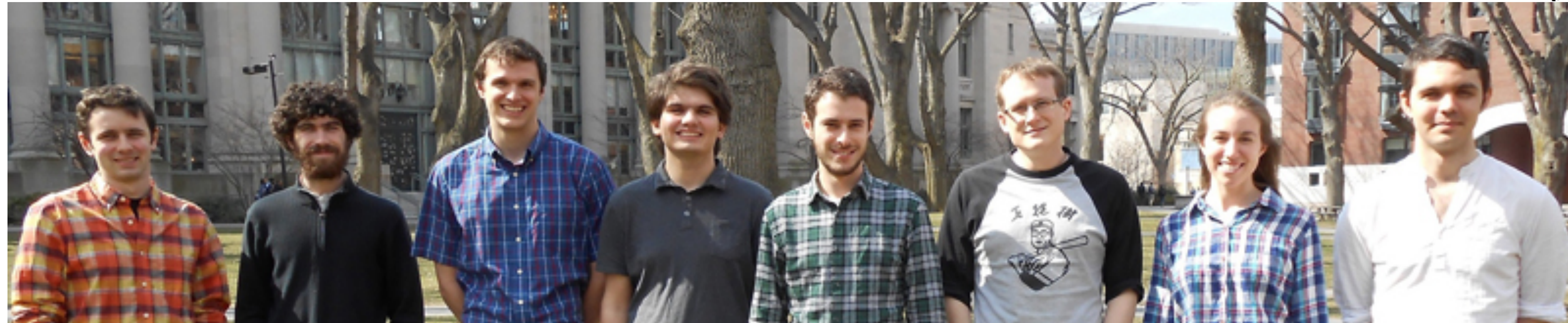
used a	molecule	
--------	----------	--

~x1

Beam Source x6.5	Molecule to ThO x2	Technical x1.5
------------------	--------------------	----------------

x15

The ACME team



Paul
Hess

Brendon
O'Leary

Ben
Spaun

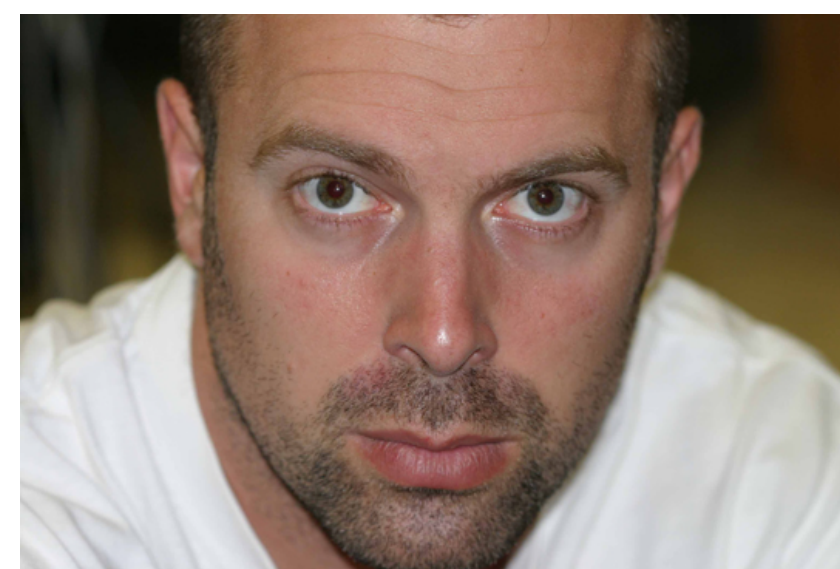
Cris
Panda

Jacob
Baron

Nick
Hutzler

Elizabeth
Petrik

Adam
West



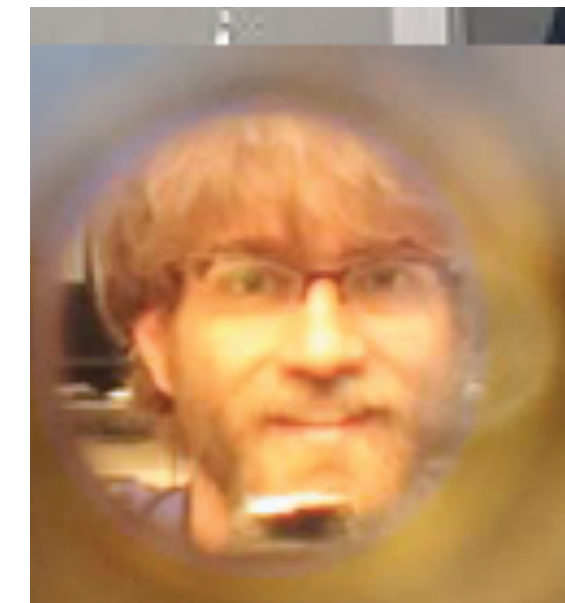
Emil
Kirilov



Amar
Vutha



Yulia
Gurevich



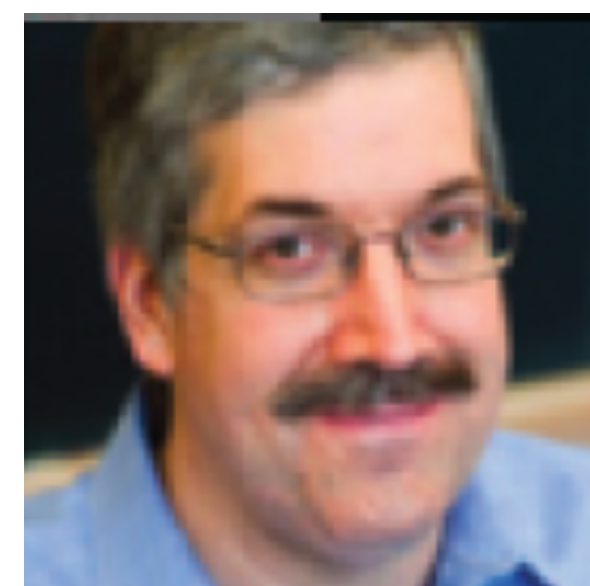
Wes
Campbell



Ivan
Kozyryev



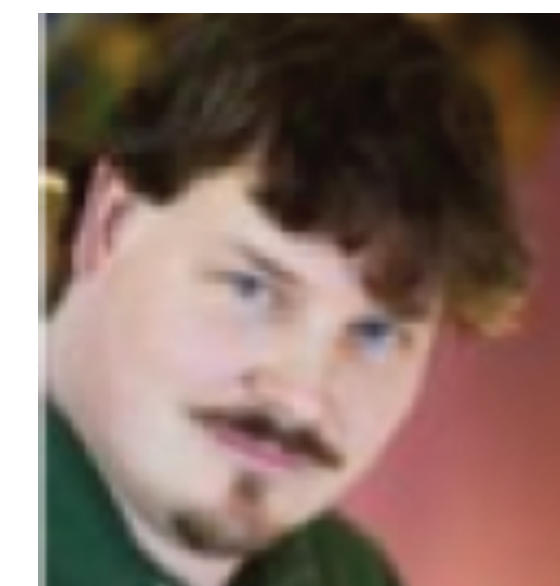
Max
Parsons



JMD



Gabrielse

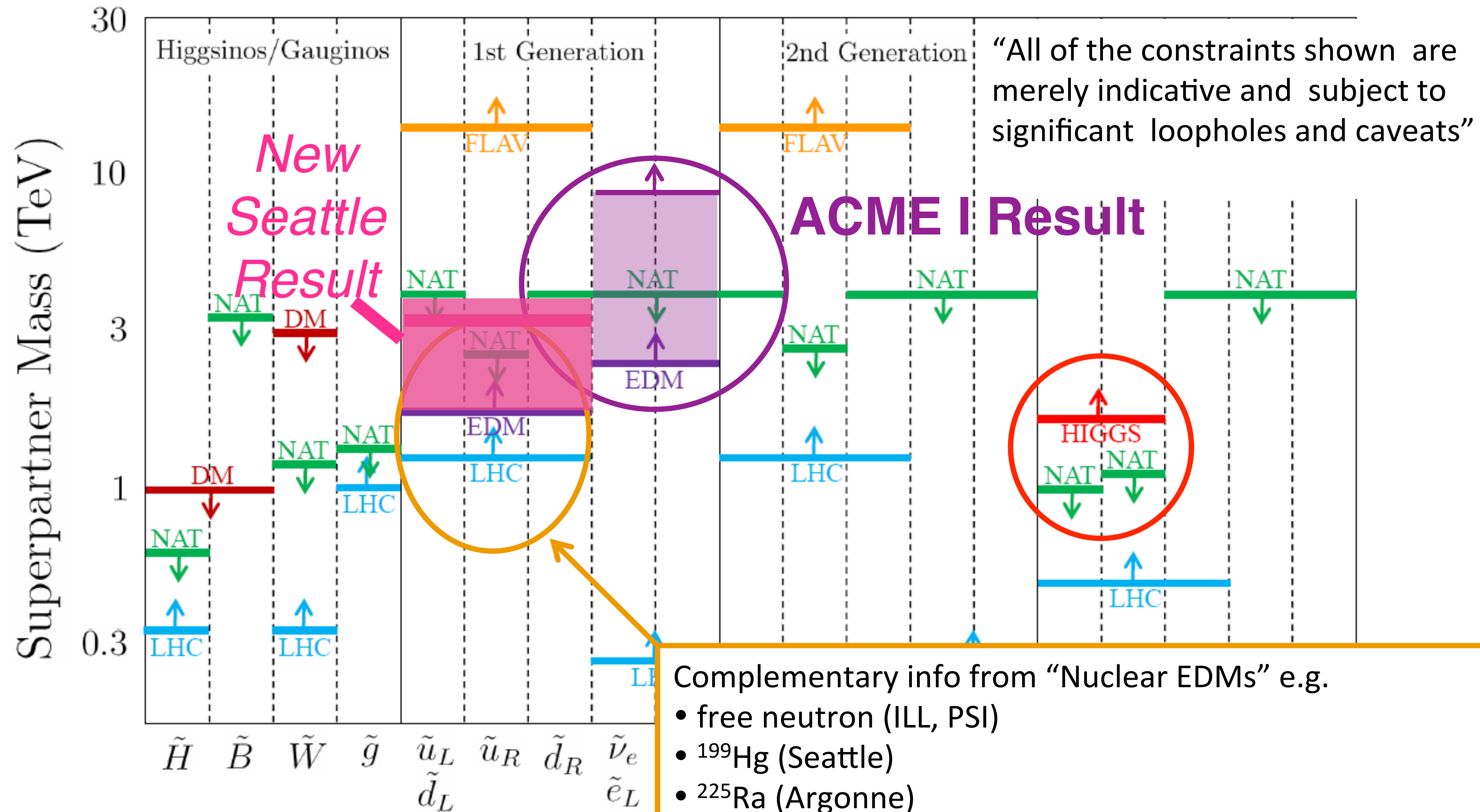


DeMille



Nuclear EDMs and Naturalness

J. Feng: "Naturalness and the status of SUSY", Ann. Rev. Nucl. Part. Sci. (2013)

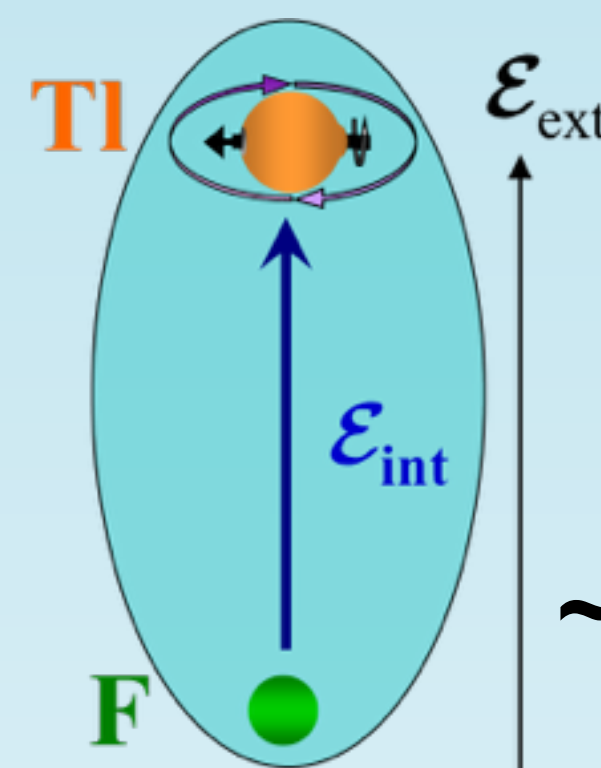


Complementary info from "Nuclear EDMs" e.g.

- free neutron (ILL, PSI)
- ^{199}Hg (Seattle)
- ^{225}Ra (Argonne)
- ^{129}Xe (TUM/Mich/PTB): Mo-127, Tu-126
- ^{129}Xe (Mainz)
- ^{221}Ra (Michigan)

ACME method will be applied to Nuclear EDM search

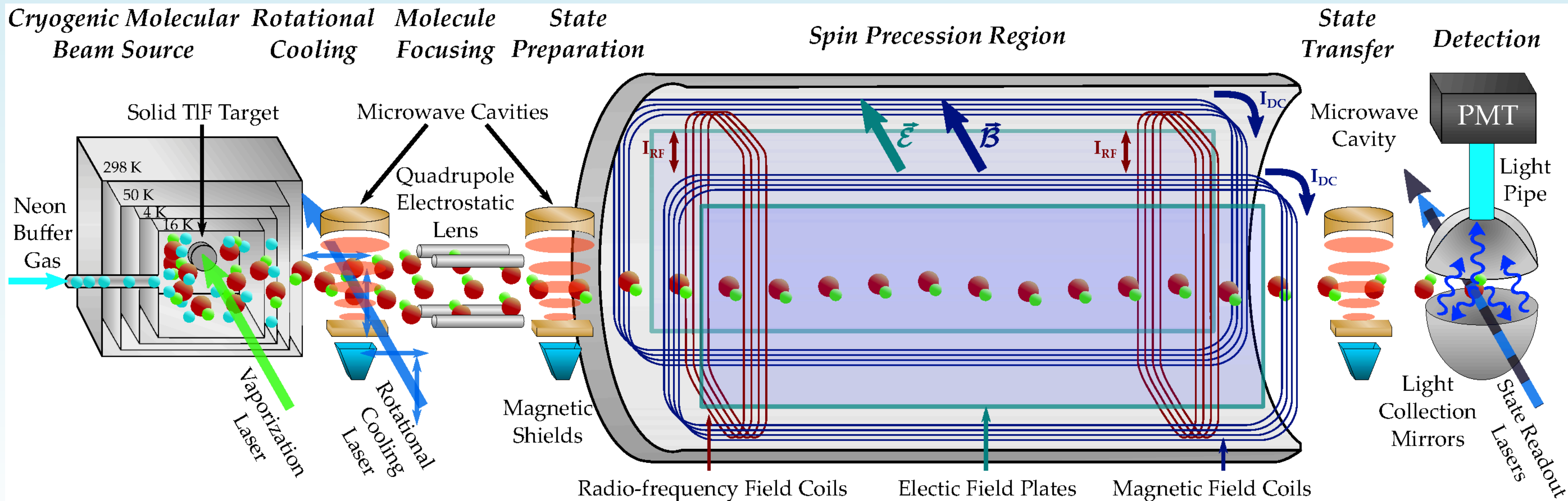
Current state-of-the-art Schiff moment limit from ^{199}Hg atom experiment (Seattle): **Already sensitive to new physics at $> \text{TeV}$ scale**



CeNTREX

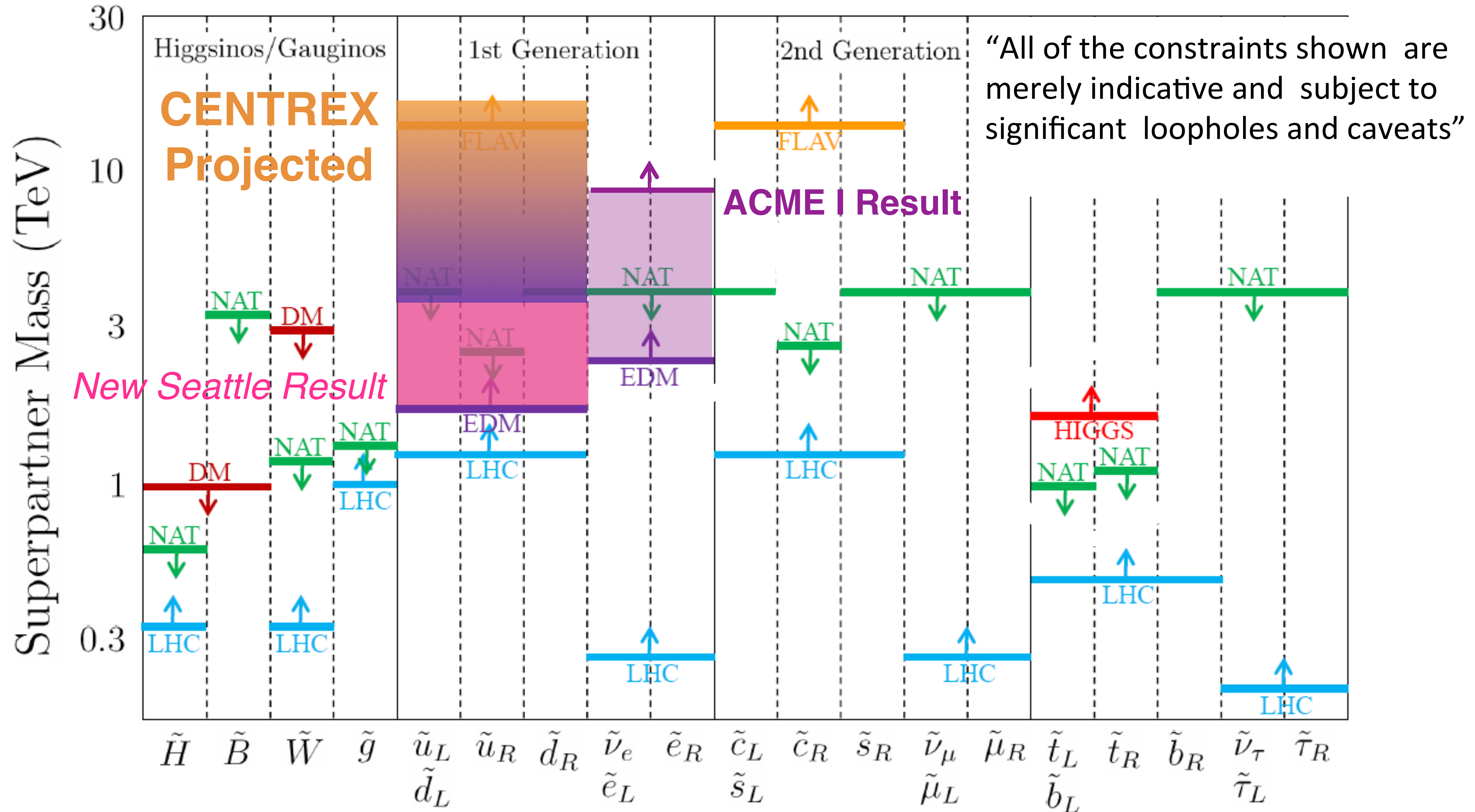
Similar to e-EDM, huge intra-molecular E-field \Rightarrow nuclear spin precession rate $\sim 10^4 \times$ larger than in ^{199}Hg atoms

D. DeMille, D. Kawall, S. Lamoreaux, T. Zelevinsky



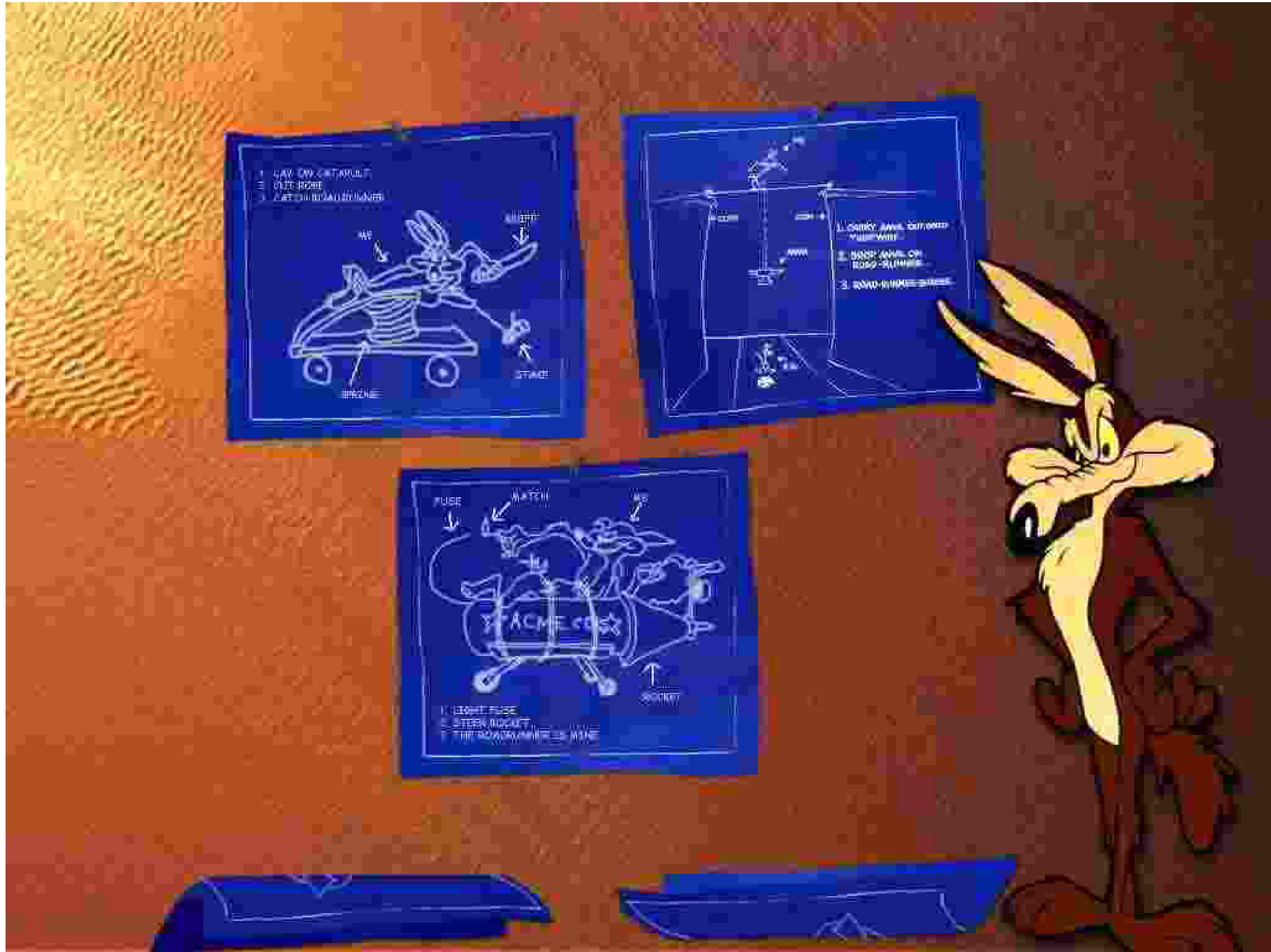
Nuclear EDMs and Naturalness

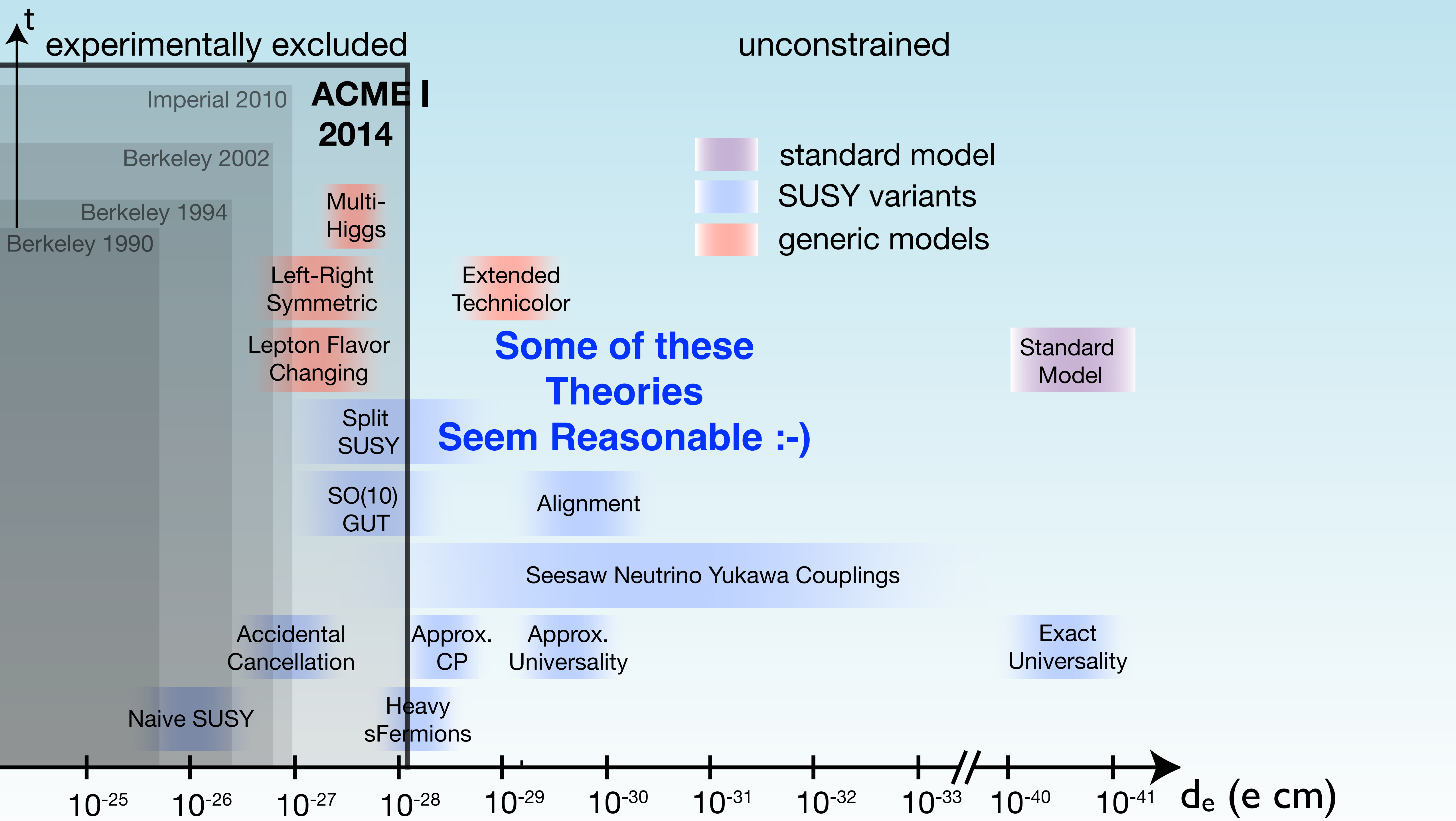
J. Feng: "Naturalness and the status of SUSY", Ann. Rev. Nucl. Part. Sci. (2013)



We are still hard at work...

....with ACME II !

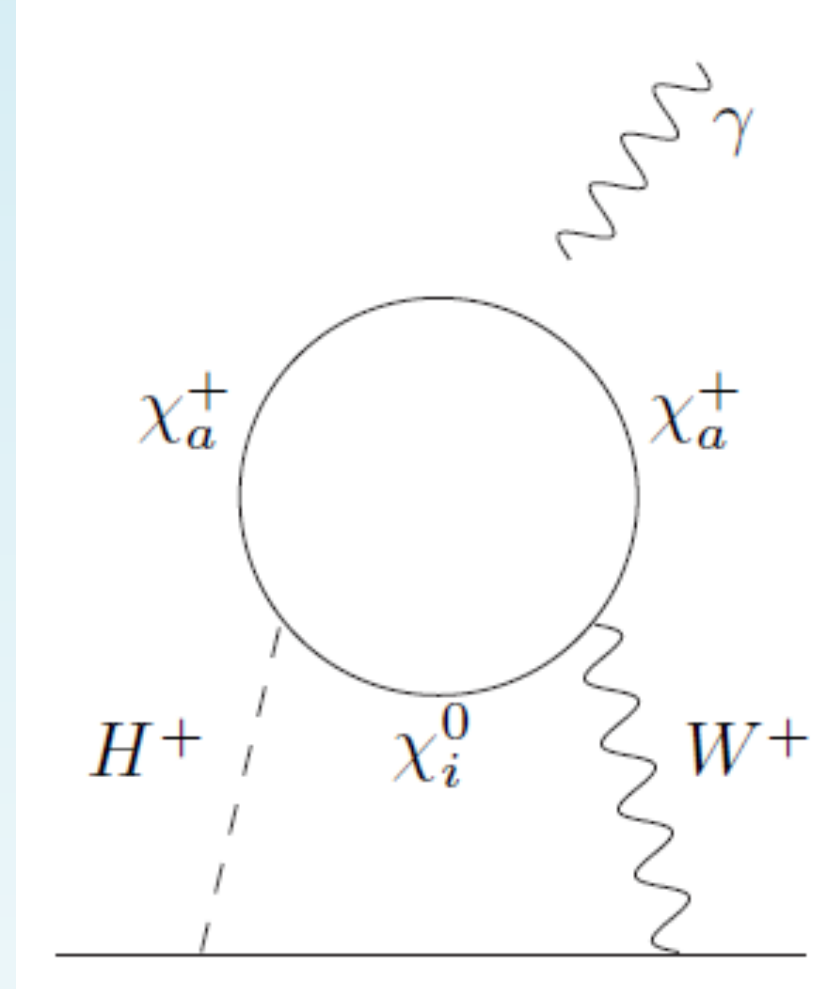




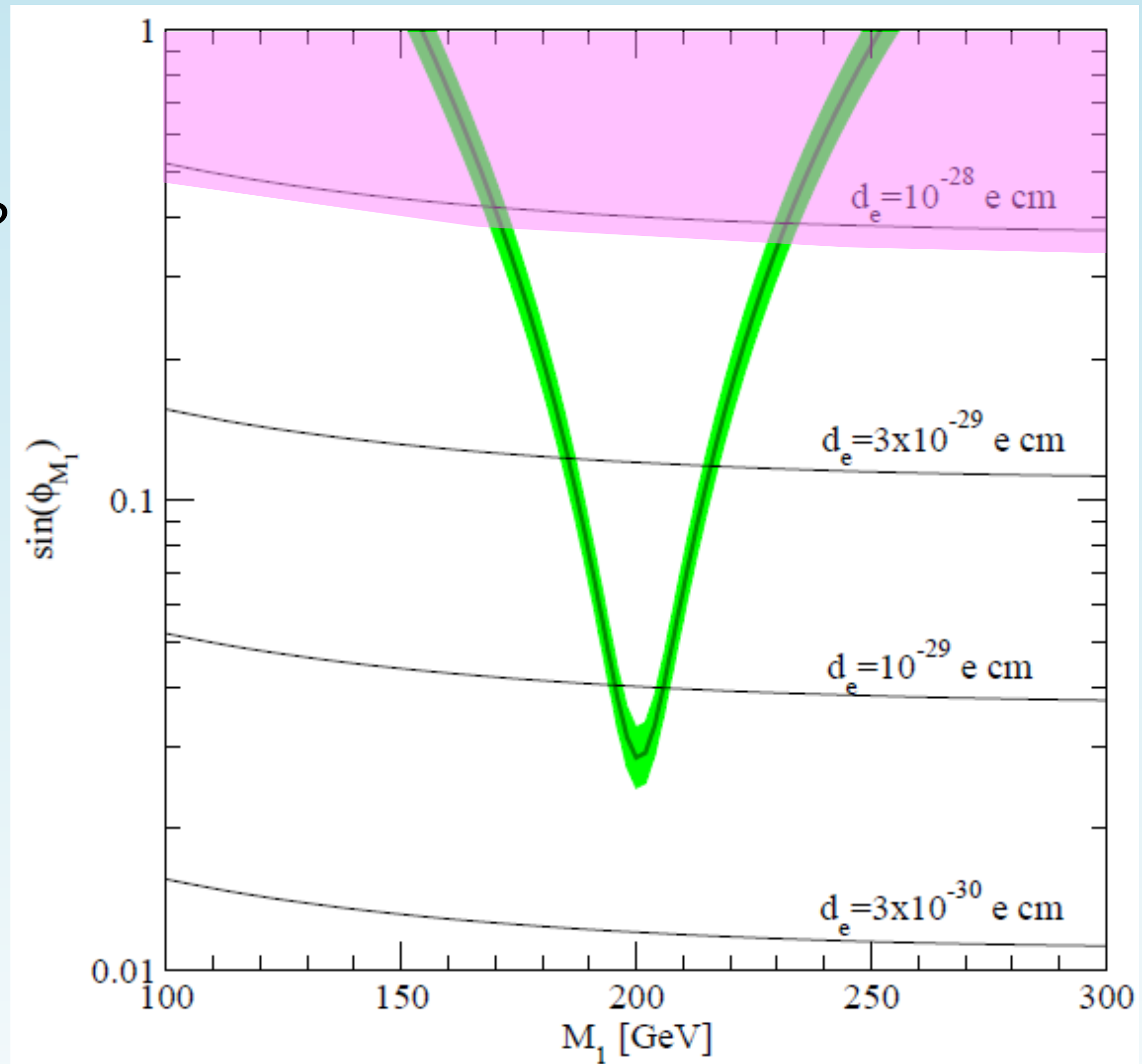
Implications for baryon asymmetry...?

Last viable corner for
Electroweak Baryogenesis
(a testable model for
matter/antimatter asym)...?

“Bino-driven EWBG”
can elude ACME limit, but...



requires non-universal
SUSY CP phases ($\varphi_2=0$)



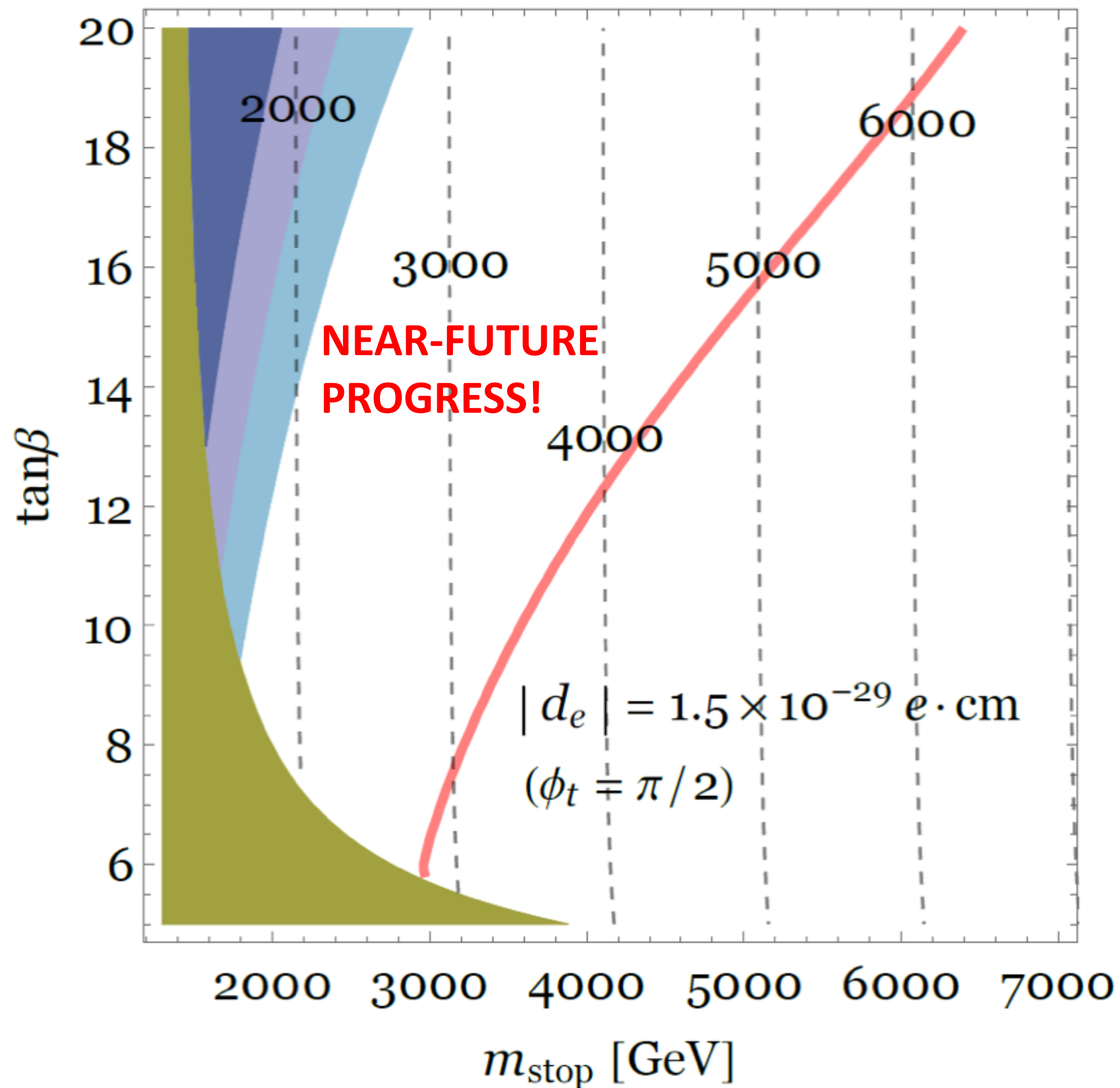
Li *et al.*, Phys. Lett. (2009)

$\sim 10x$ improvement may rule out Electroweak Baryogenesis...?

Future prospects

Yuichiro Nakai and Matthew Reece: **preliminary!**

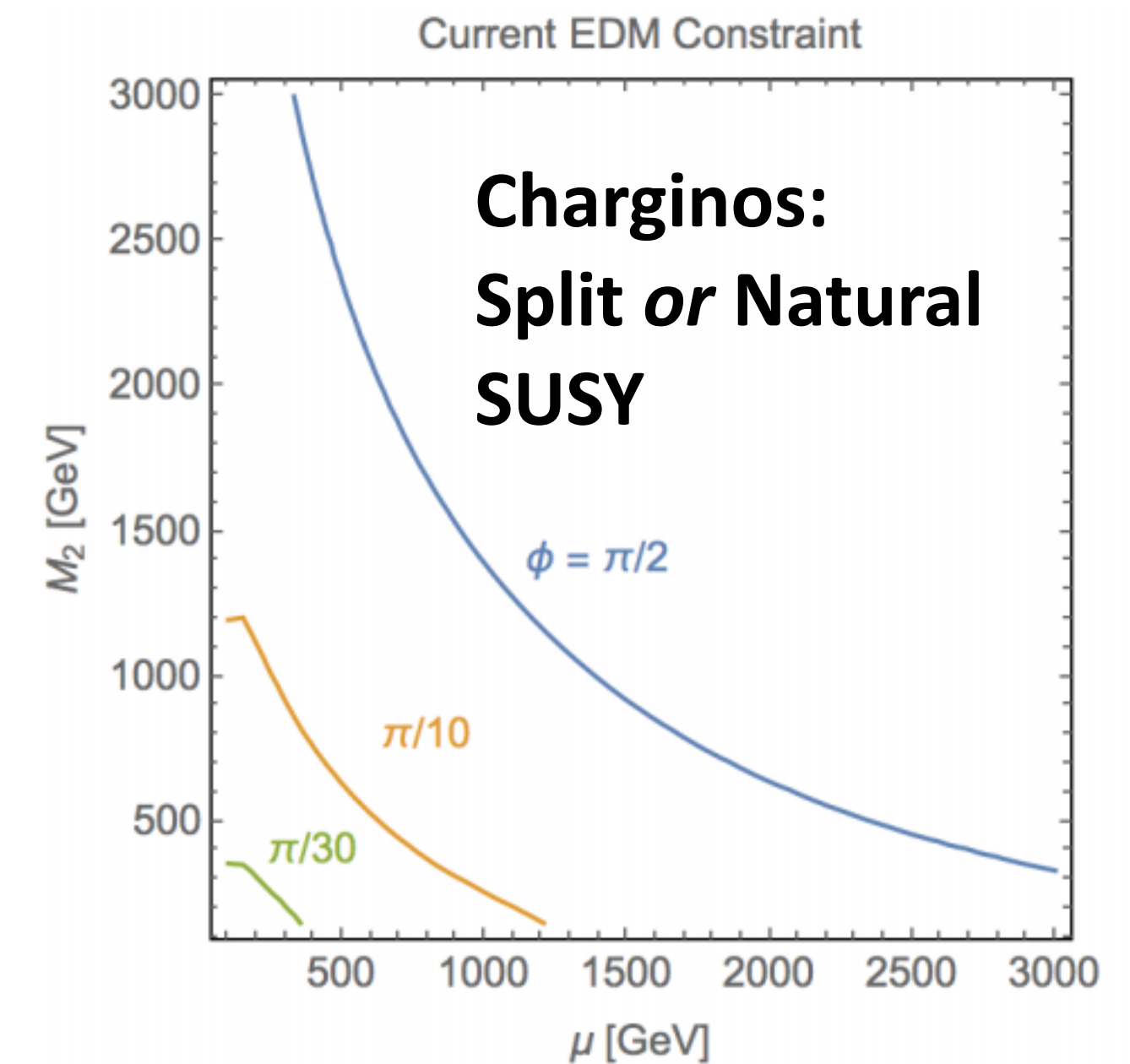
$$\mu = 350 \text{ GeV} \quad m_A = 400 \text{ GeV}$$



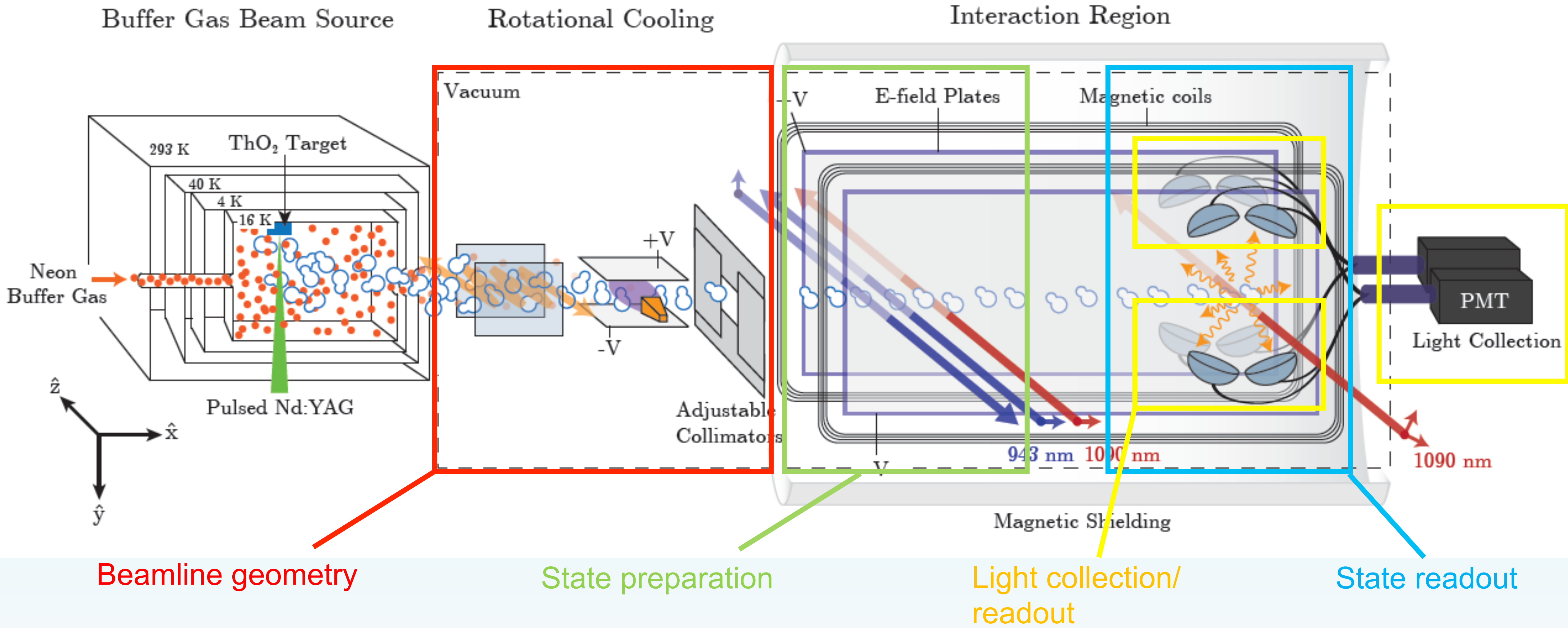
Large parameter space of stop masses unconstrained by the Higgs mass can be explored. Factor of 10 improvement will cover large part of interesting parameter space w/ CP phase.

More to the story: charginos, Higgs physics beyond MSSM, ...

Under study



ACME I → ACME II Upgrades



The Path to Improvements

This slide is about five years old....

BEFORE

the ACME I result

Increases in statistical sensitivity already demonstrated

Fully Demonstrated:

- ✓ Electrostatic Focussing } 10x
- ✓ Rotational Cooling } 10x
- ✓ STIRAP State Preparation 5x



Emil Kirilov

Mostly Demonstrated:

- ✓ Photon Cycling/collection geometry
(photon shot noise ---> molecule shot noise) 10x
- ✓ Thermochemical Production 10x

Not Demonstrated:

- More cryogenic cooling 4x
- PMT-->Cooled Photodiode or PMT 2x

Possible combined increase in signal: 100 -- 10,000

Possible combined increase in statistical sensitivity: 10 -- 100

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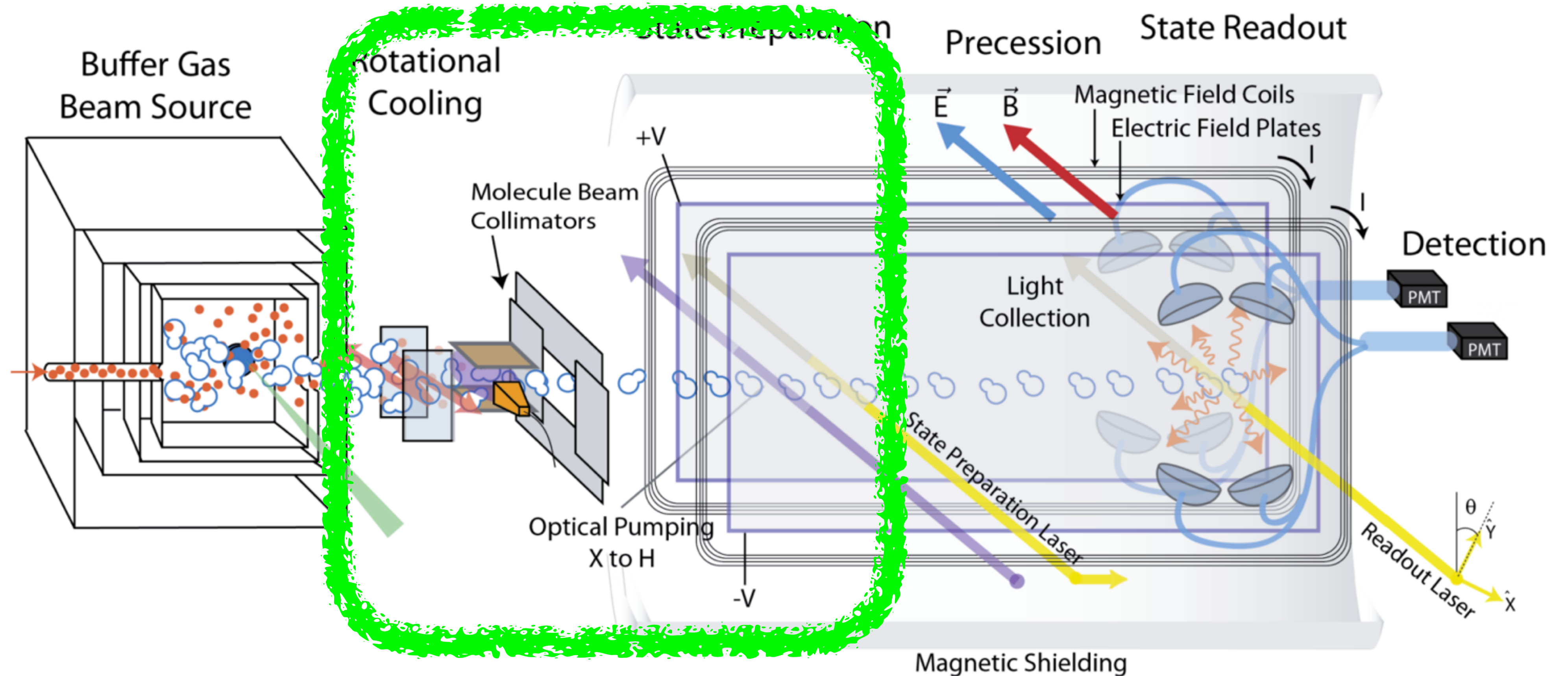
Not Demonstrated:

- More cryogenic cooling 4x
- PMT-->Cooled Photodiode or PMT 2x

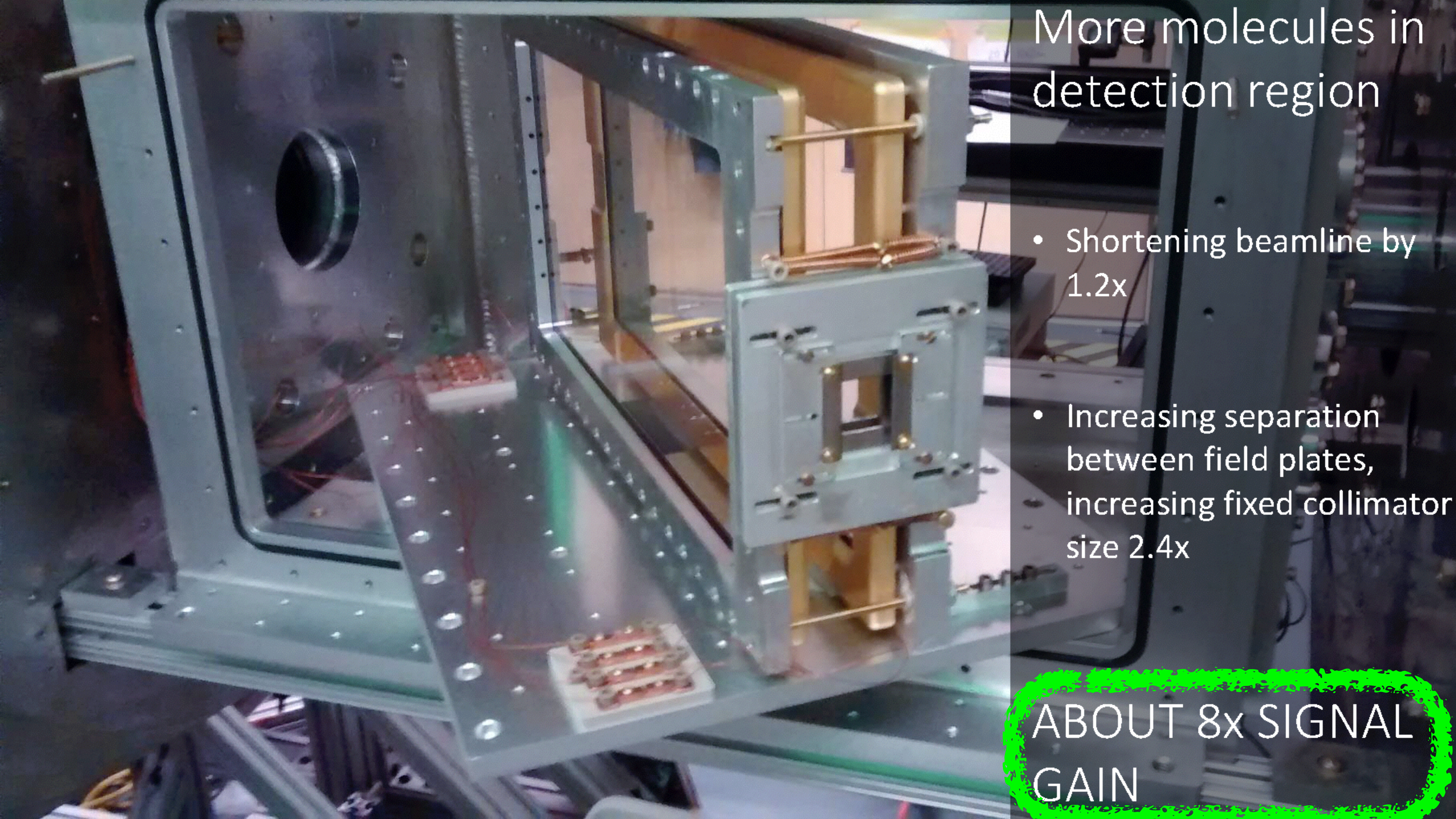
Possible combined increase in signal: 100 -- 10,000

Possible combined increase in statistical sensitivity: 10 -- 100

ACME Generation II Apparatus Improvements



Increasing Detected Solid Angle of Molecule Beam

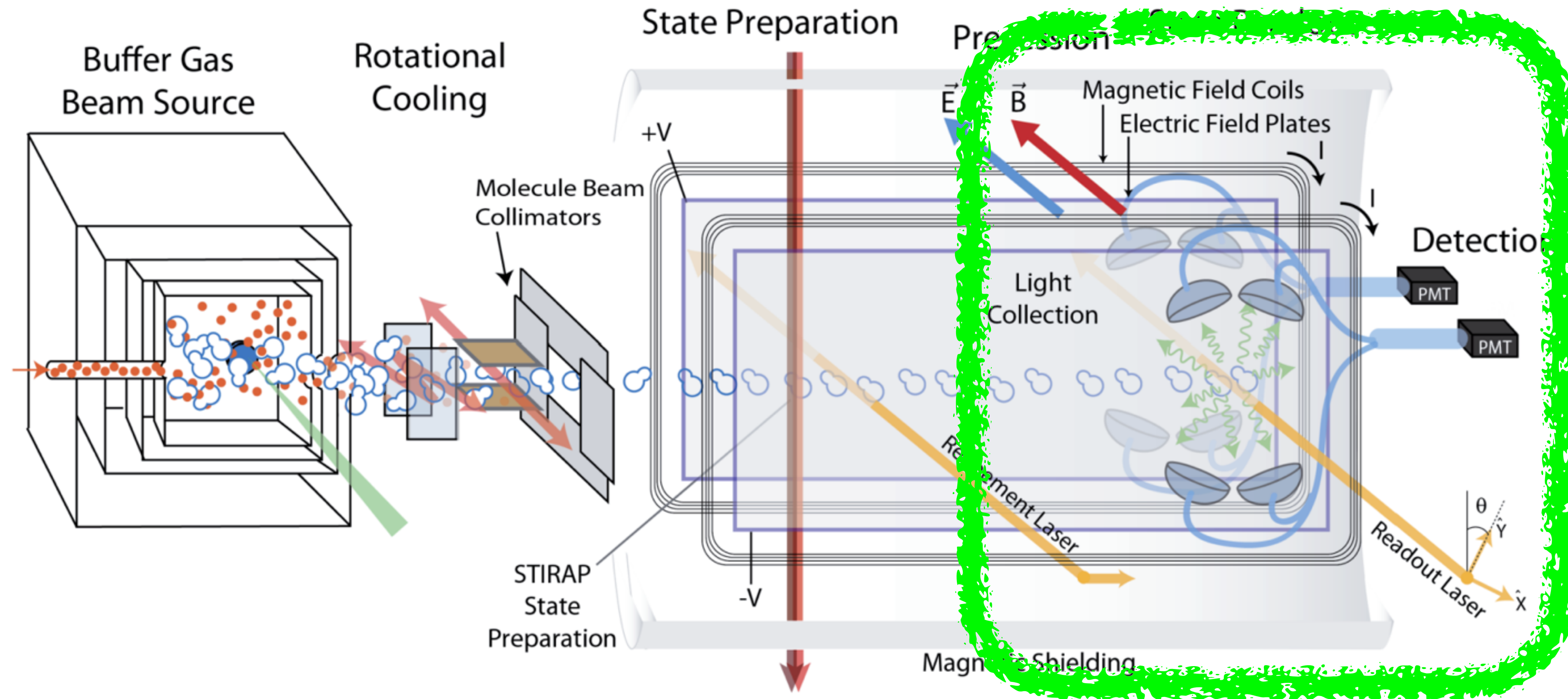


More molecules in
detection region

- Shortening beamline by 1.2x
- Increasing separation between field plates, increasing fixed collimator size 2.4x

ABOUT 8x SIGNAL
GAIN

ACME Generation II Apparatus Improvements



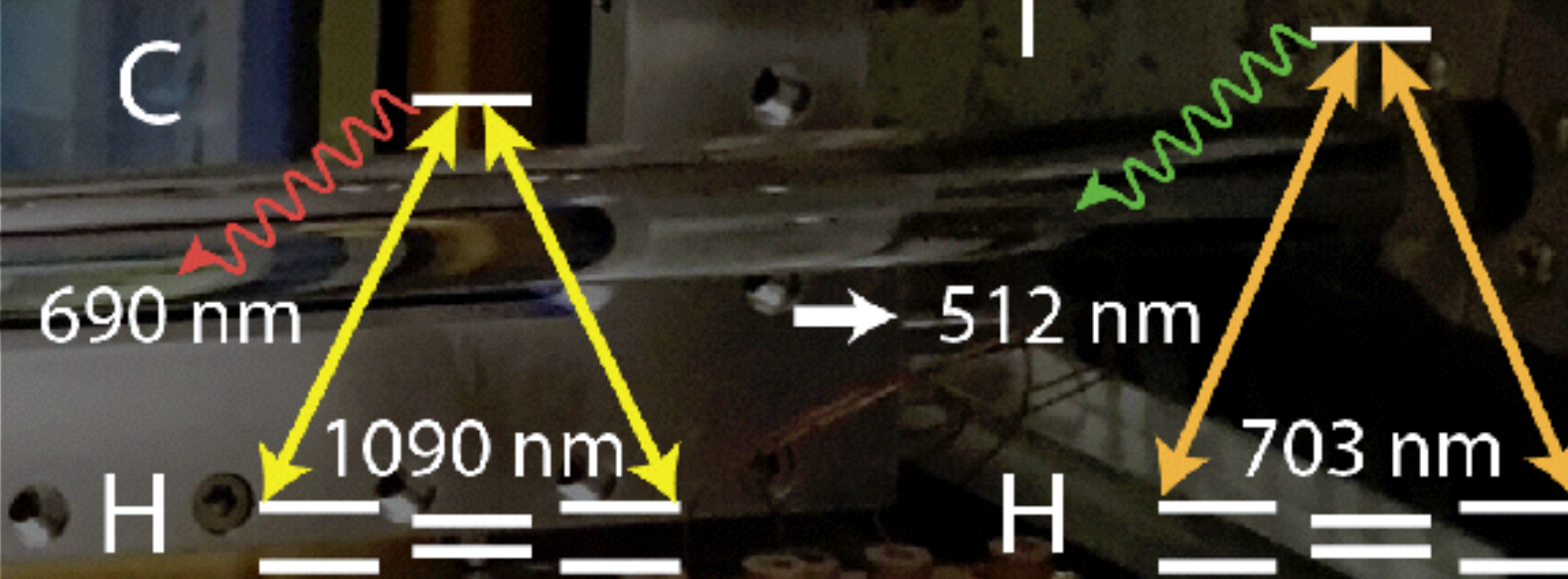
Increasing Molecule Detection Efficiency with New Transition

Improving Detection Efficiency

Curved light pipes and other light collection optimizations
About 1.8x Gain

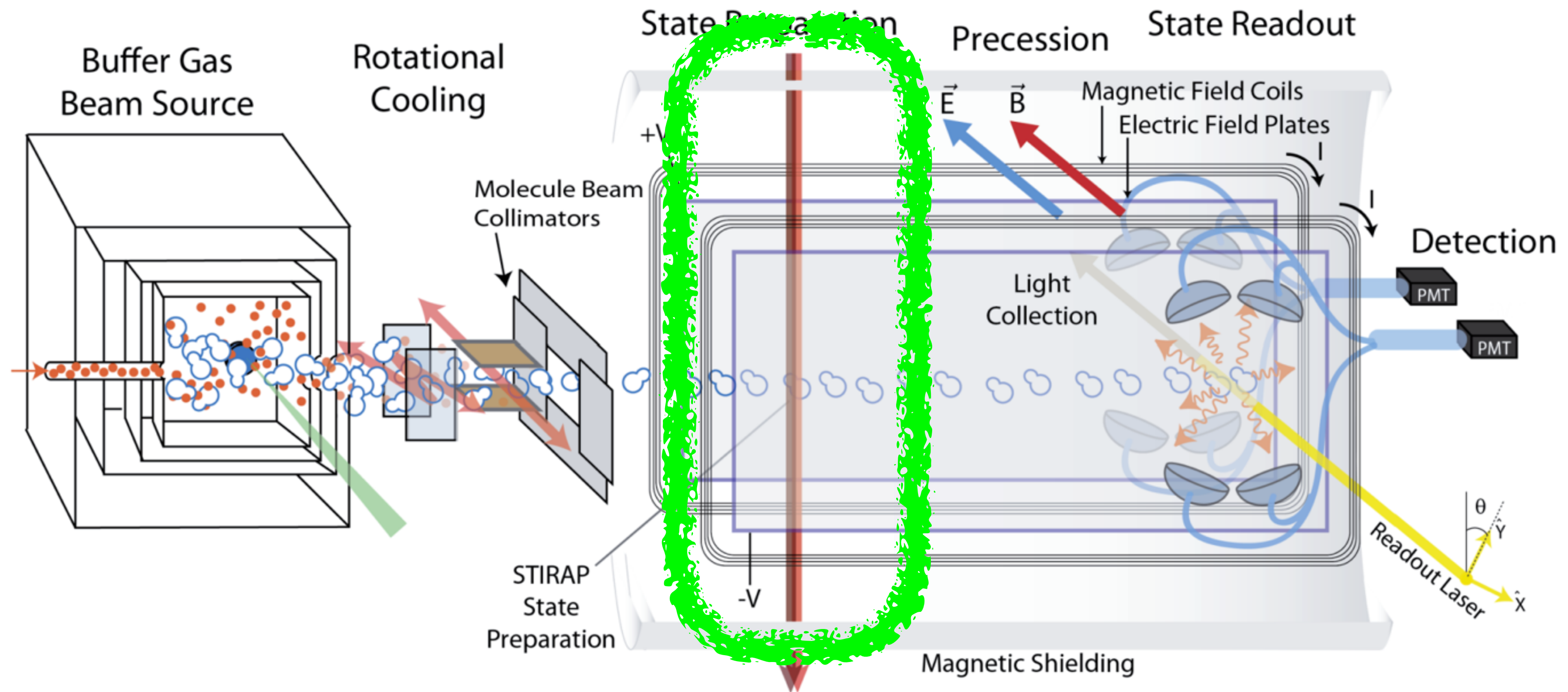
Detecting on new transition provides more light/molecule due to larger branching ratio, About 1.2x Gain

Higher quantum efficiency at new detection wavelength 512 nm, About 2.5x Gain



ABOUT 5x SIGNAL GAIN

ACME Generation II Apparatus Improvements

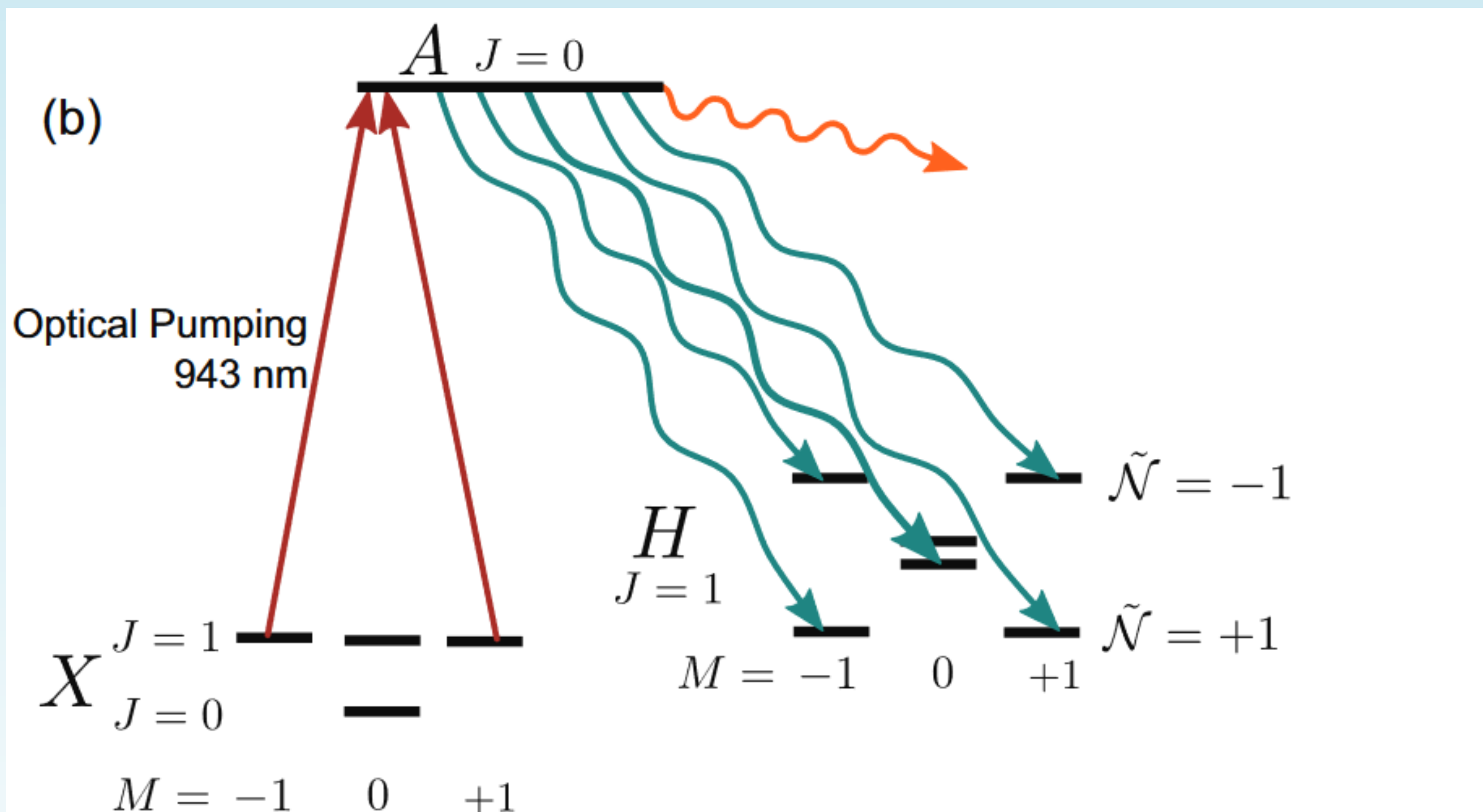


Increasing the Efficiency of State Preparation

ACME I \longrightarrow ACME II

STIRAP State Preparation

ACME I

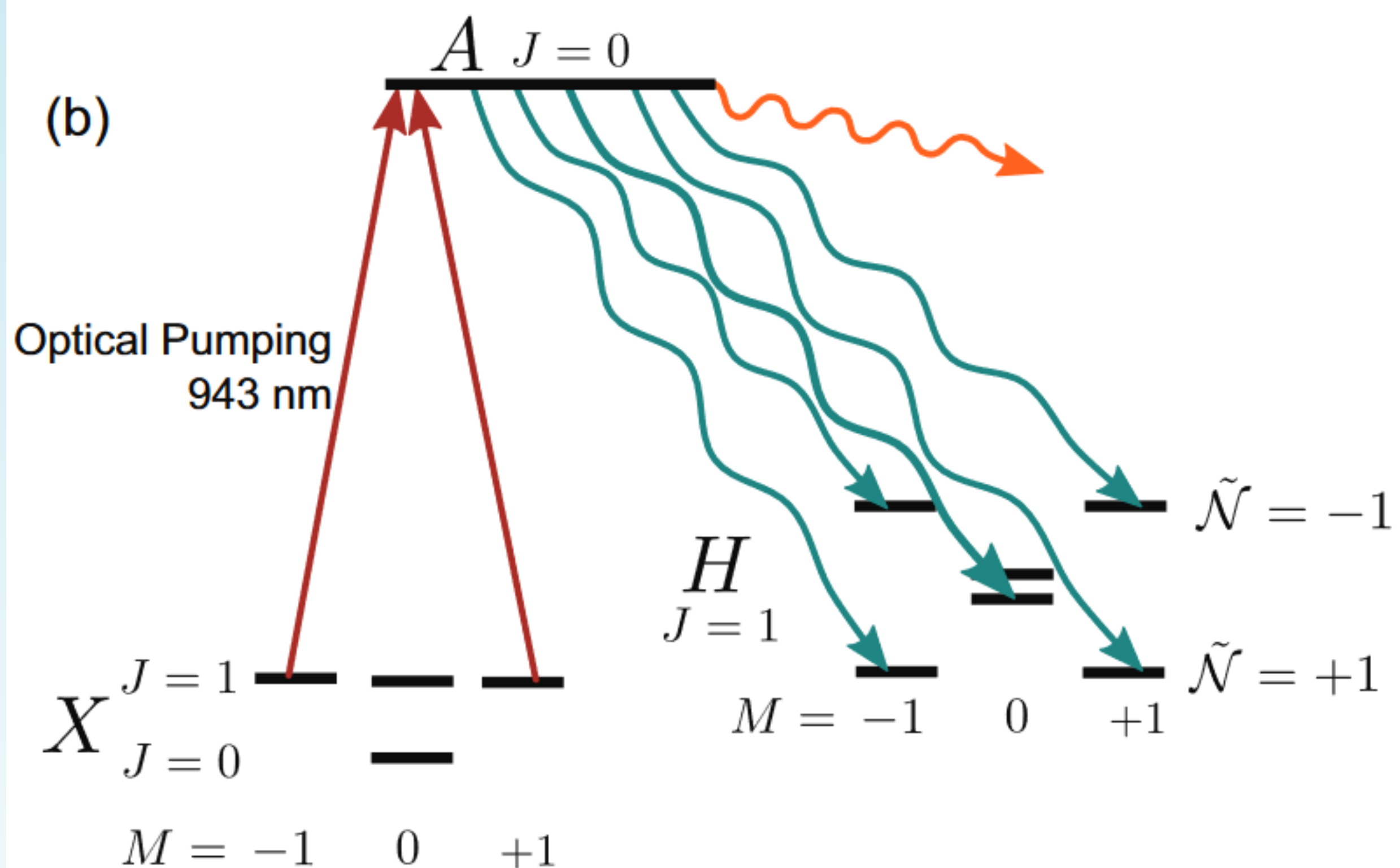


Optical Pumping (incoherent)

ACME I \longrightarrow ACME II

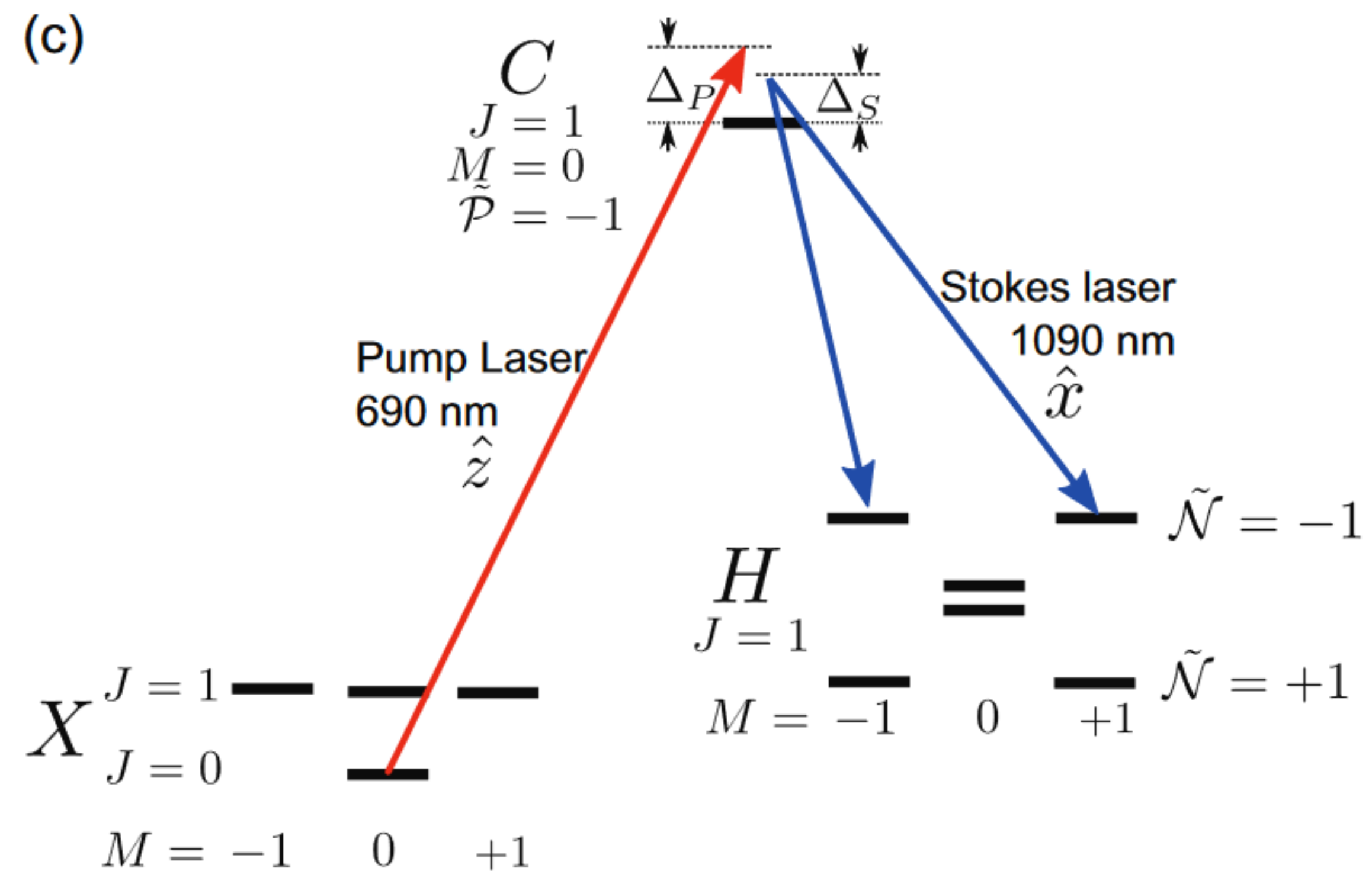
STIRAP State Preparation

ACME I



Optical Pumping (incoherent)

ACME II



STIRAP (coherent)

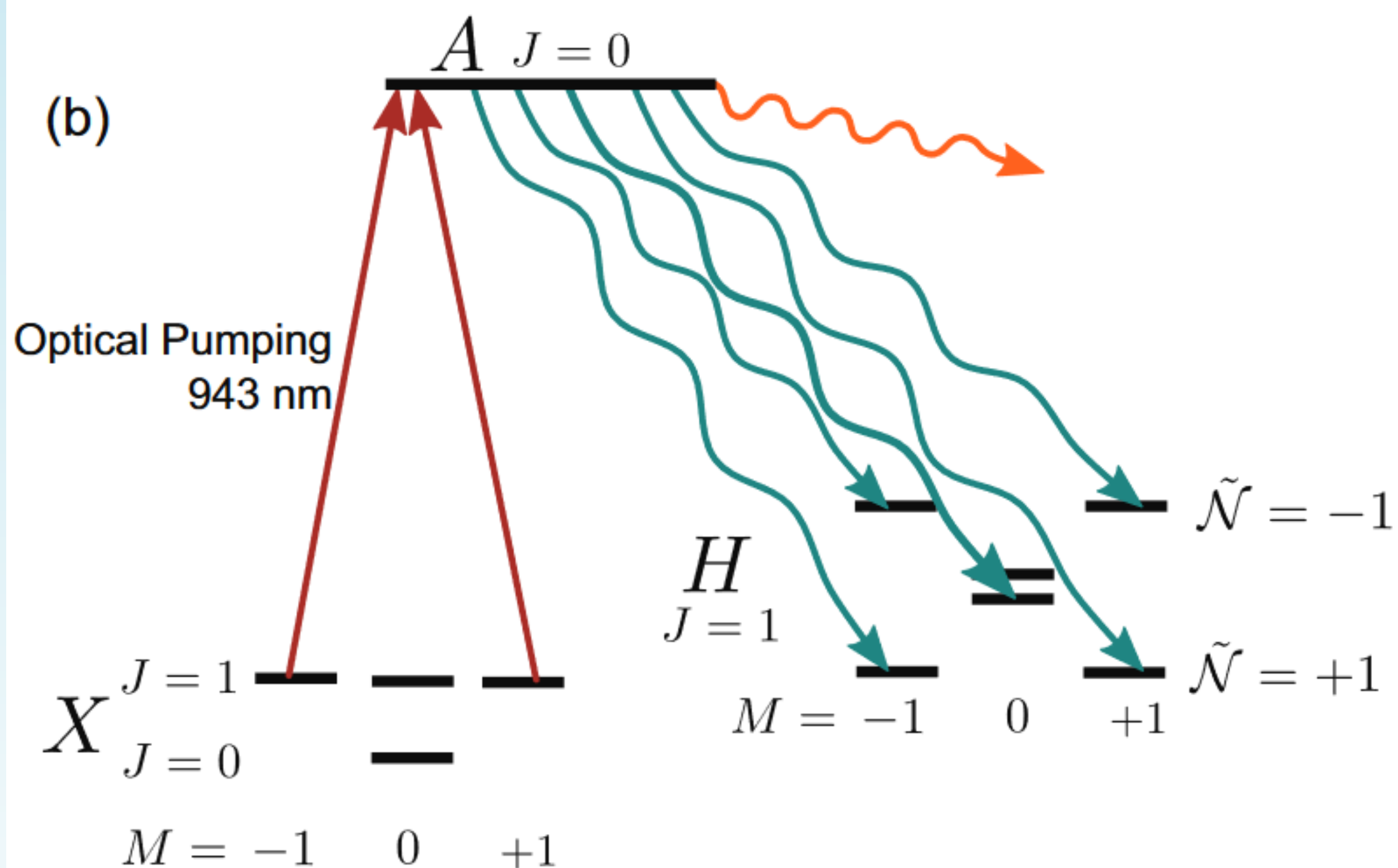
ACME I \longrightarrow ACME II

STIRAP State Preparation

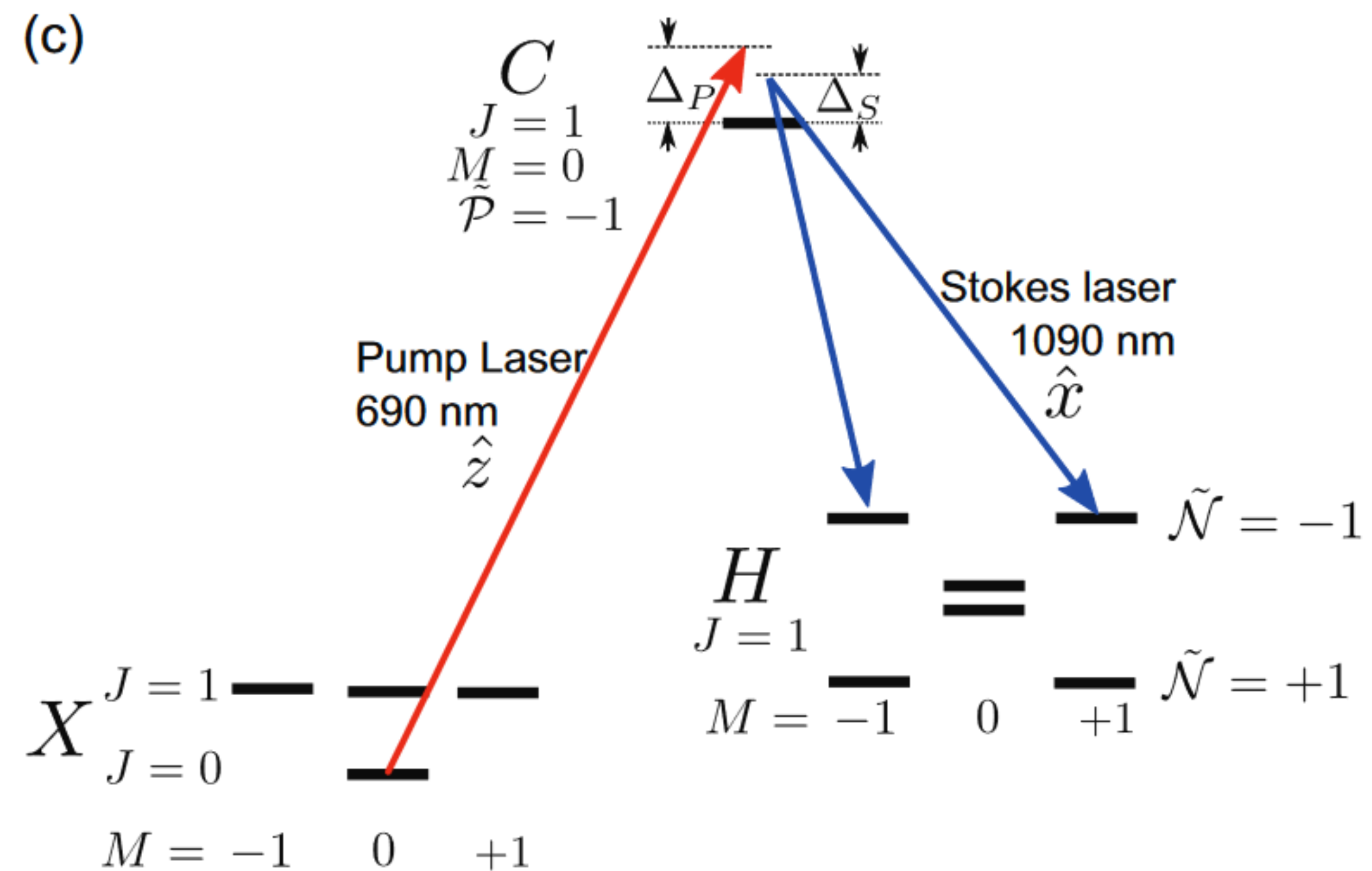
Known Technique

ACME I

ACME II



Optical Pumping (incoherent)



STIRAP (coherent)

ACME I \rightarrow ACME II

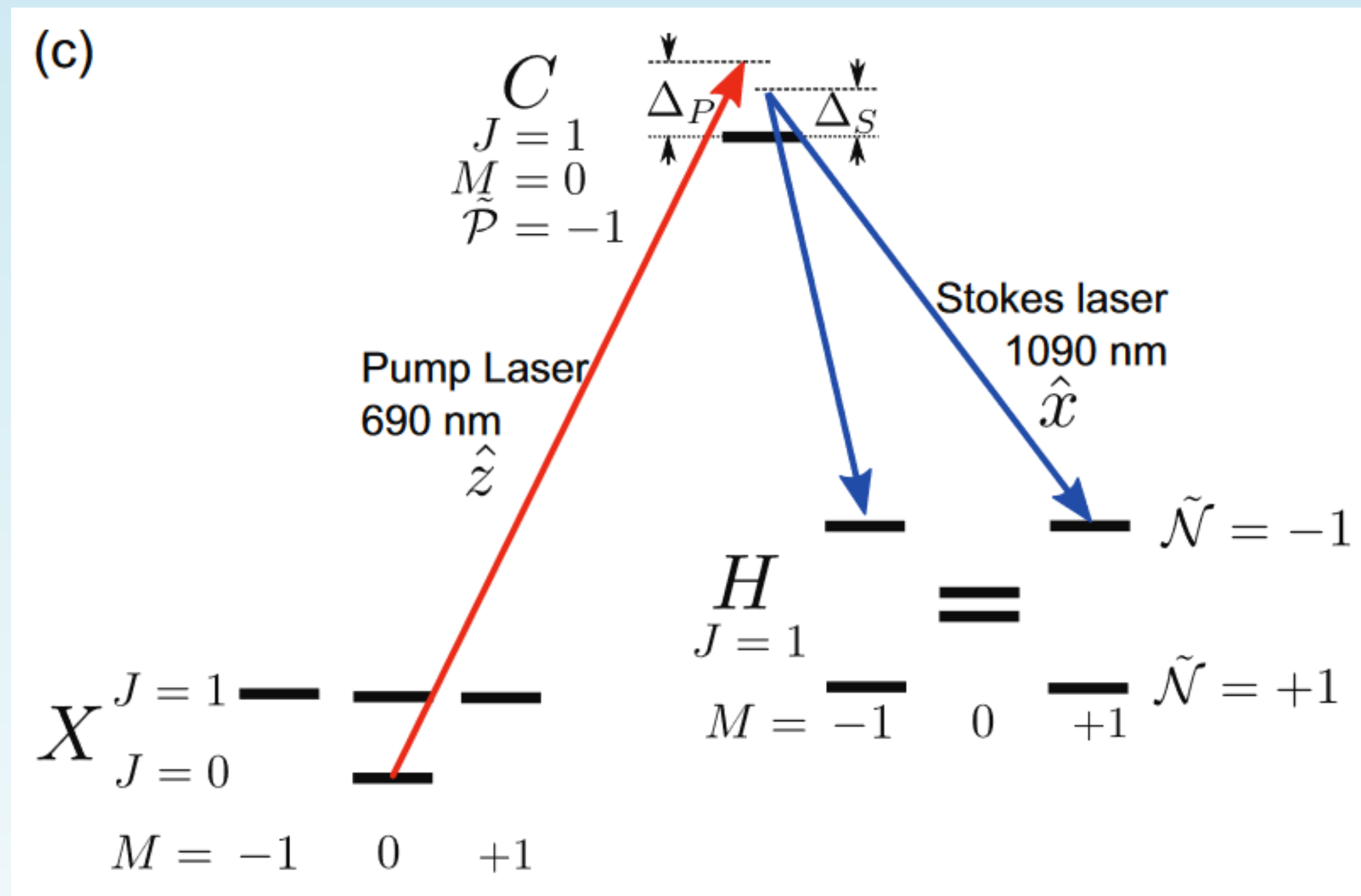
STIRAP State Preparation

BUT Challenges

- Weak transition
- Wavelengths differ by x2
- Large doppler width
- Large diameter molecular beam
- Interaction region about 1 m away

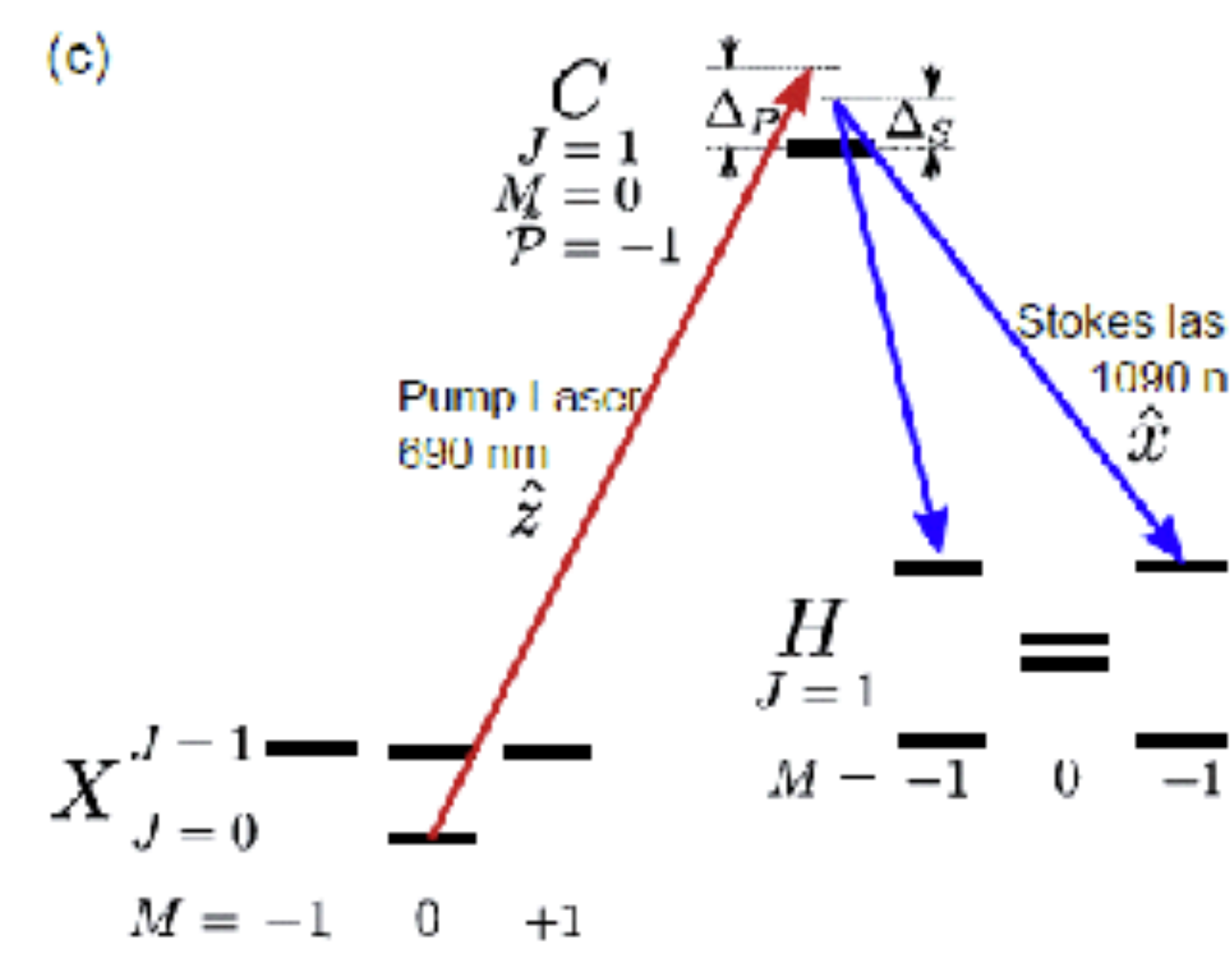
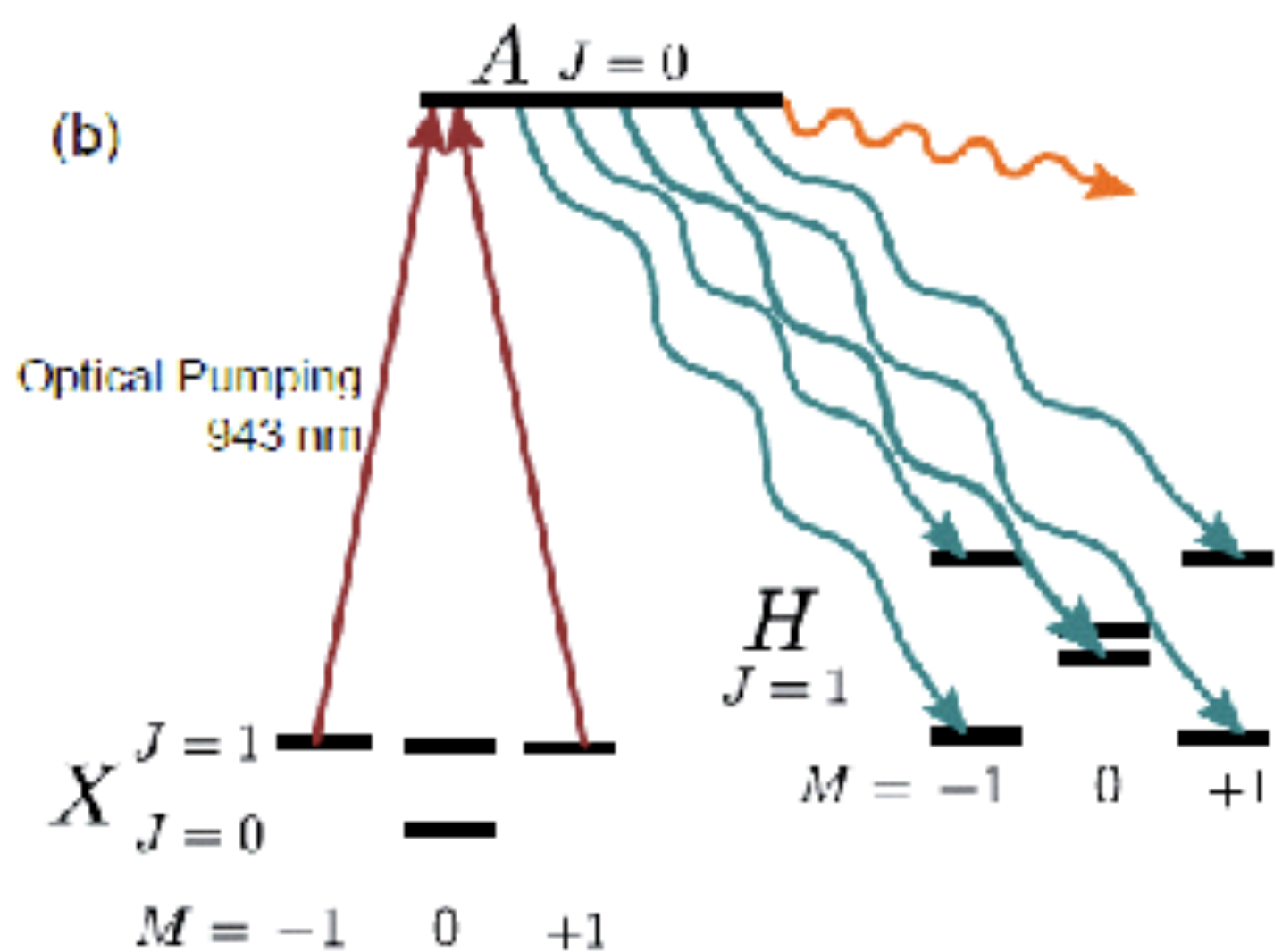
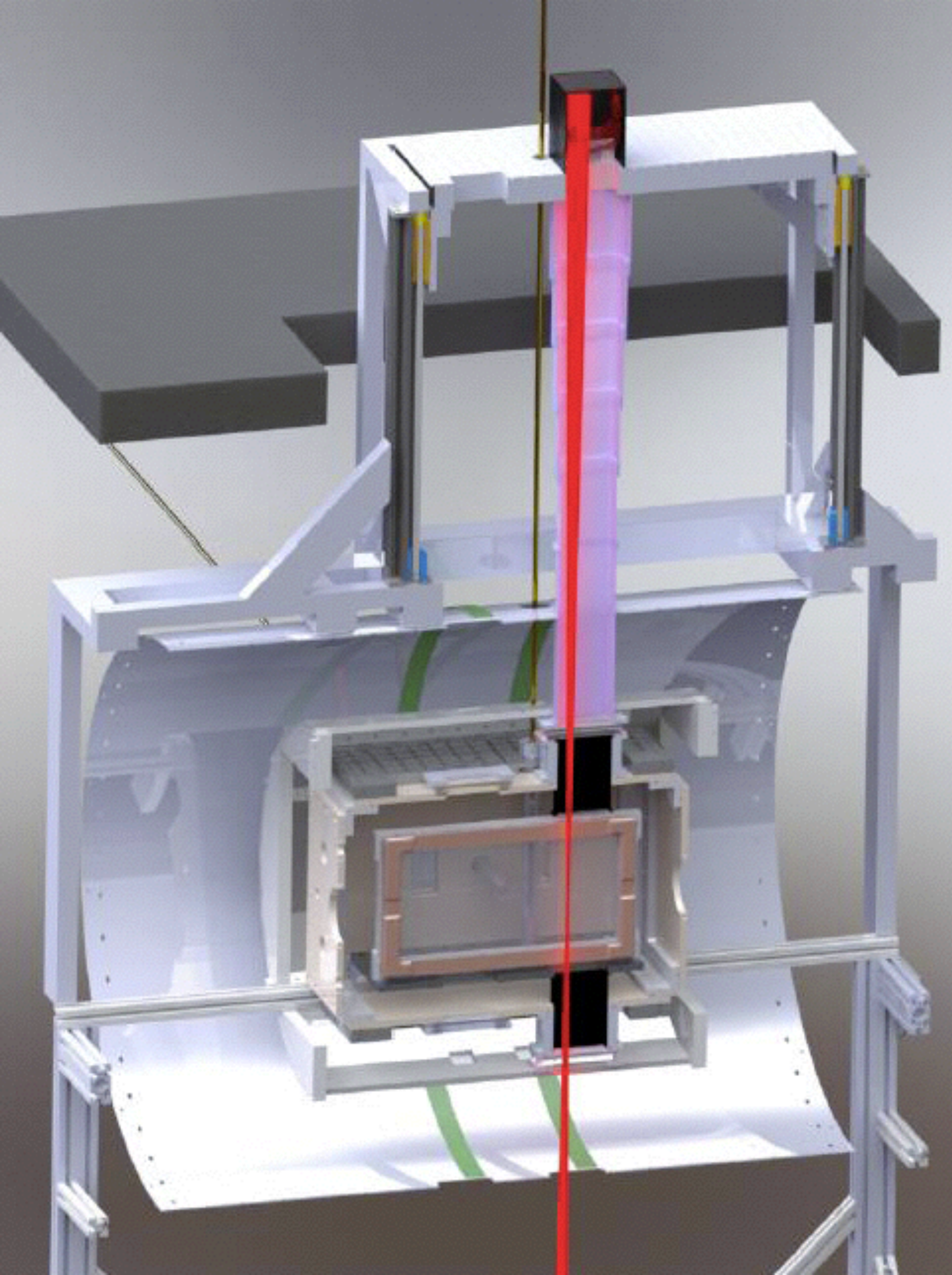
Known Technique

ACME II



STIRAP (coherent)

More Efficient Molecule Manipulation



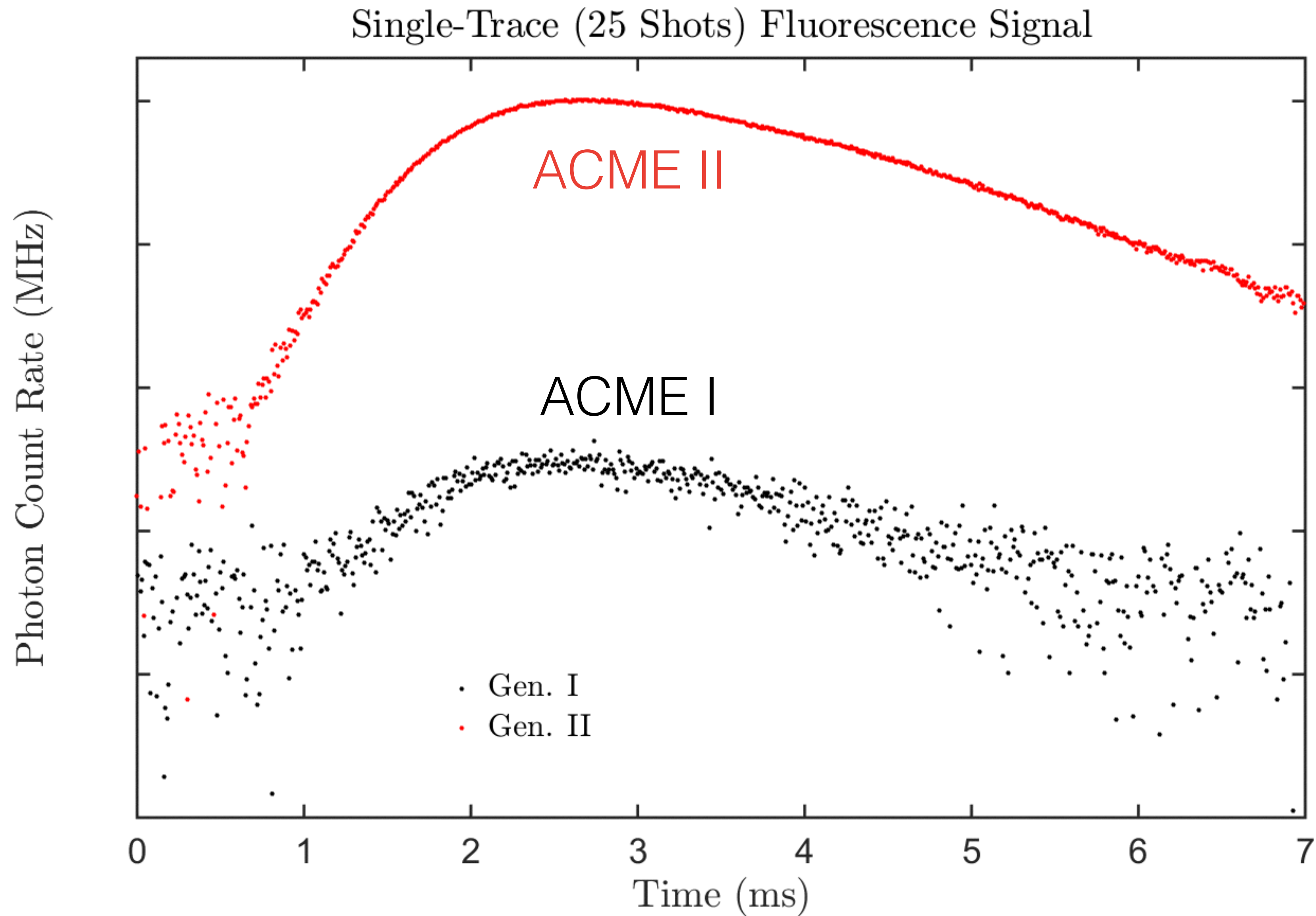
- Demonstrated STIRAP with 75% efficiency
- Compare to Optical Pumping with 6% efficiency
- New systematic errors require refinement beam
- Requires vertical launch of lasers

ABOUT 12x
SIGNAL GAIN

STIRAP State Preparation

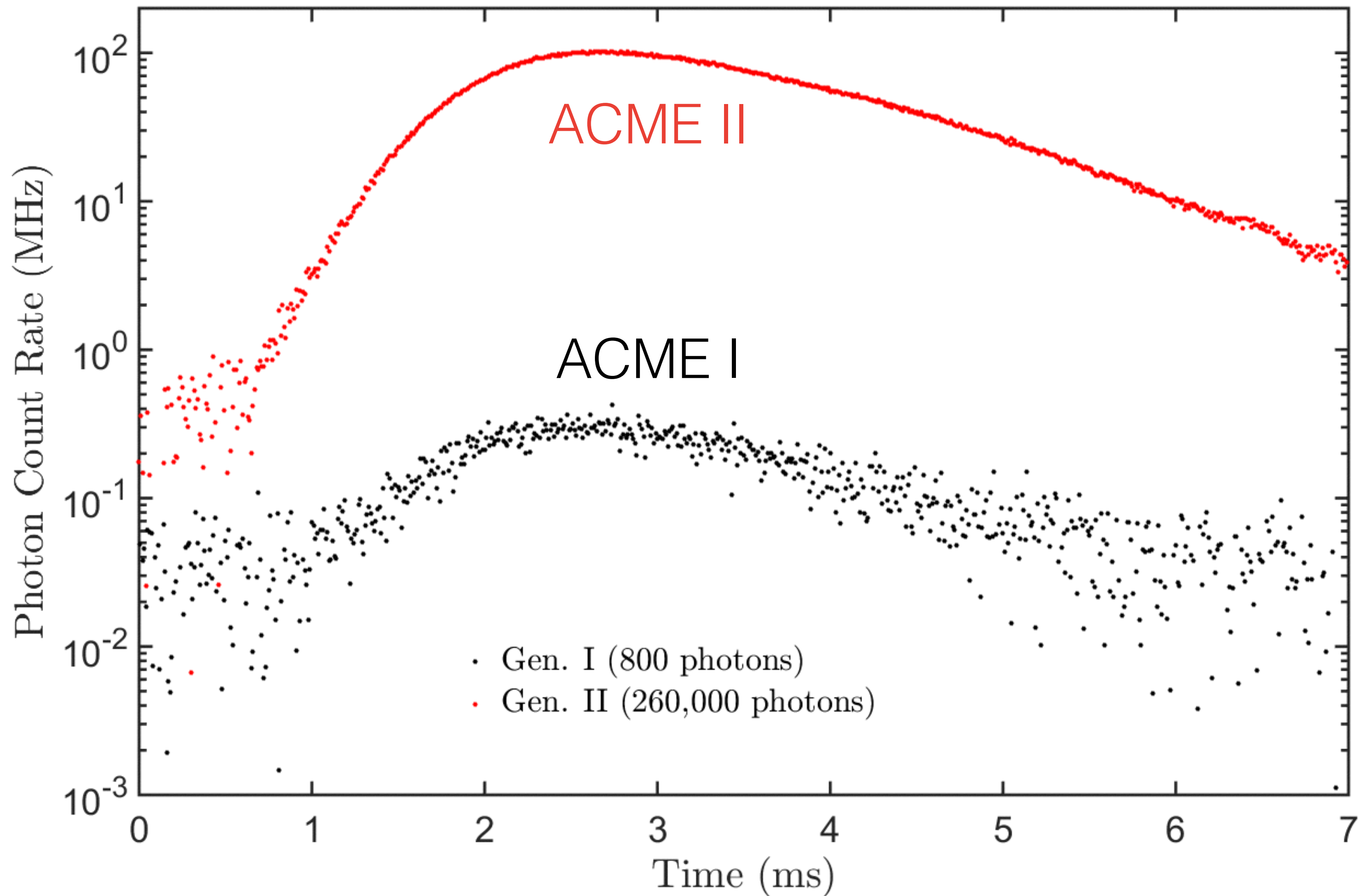
C. Panda et al. Phys. Rev. A. 93 052110 (2016)

ACME II Preliminary Signals



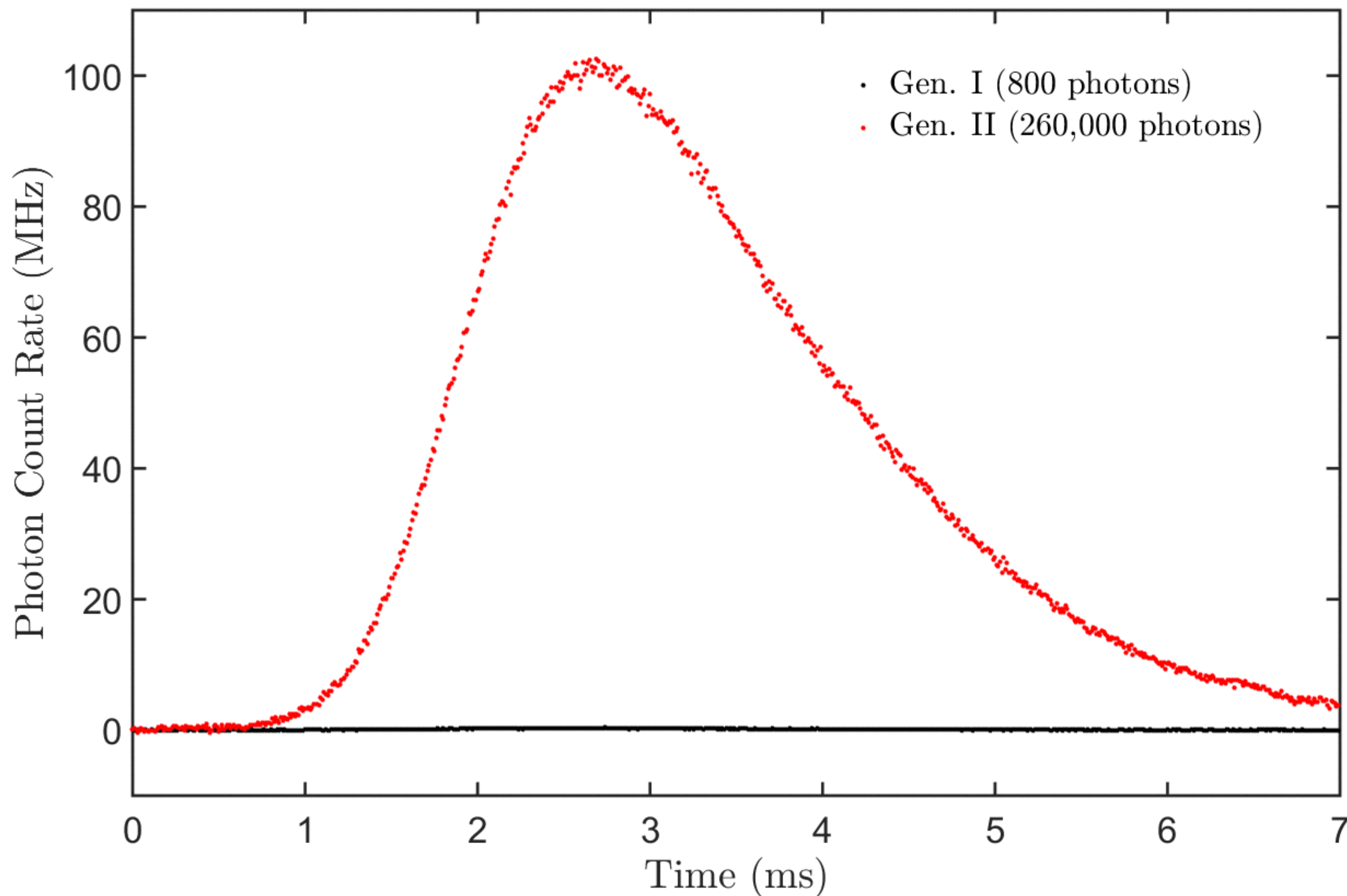
ACME II Preliminary Signals

Single-Trace (25 Shots) Fluorescence Signal

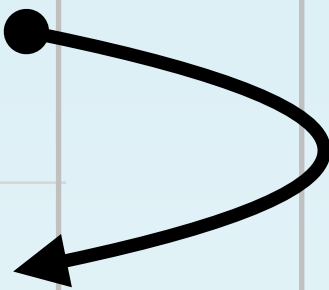


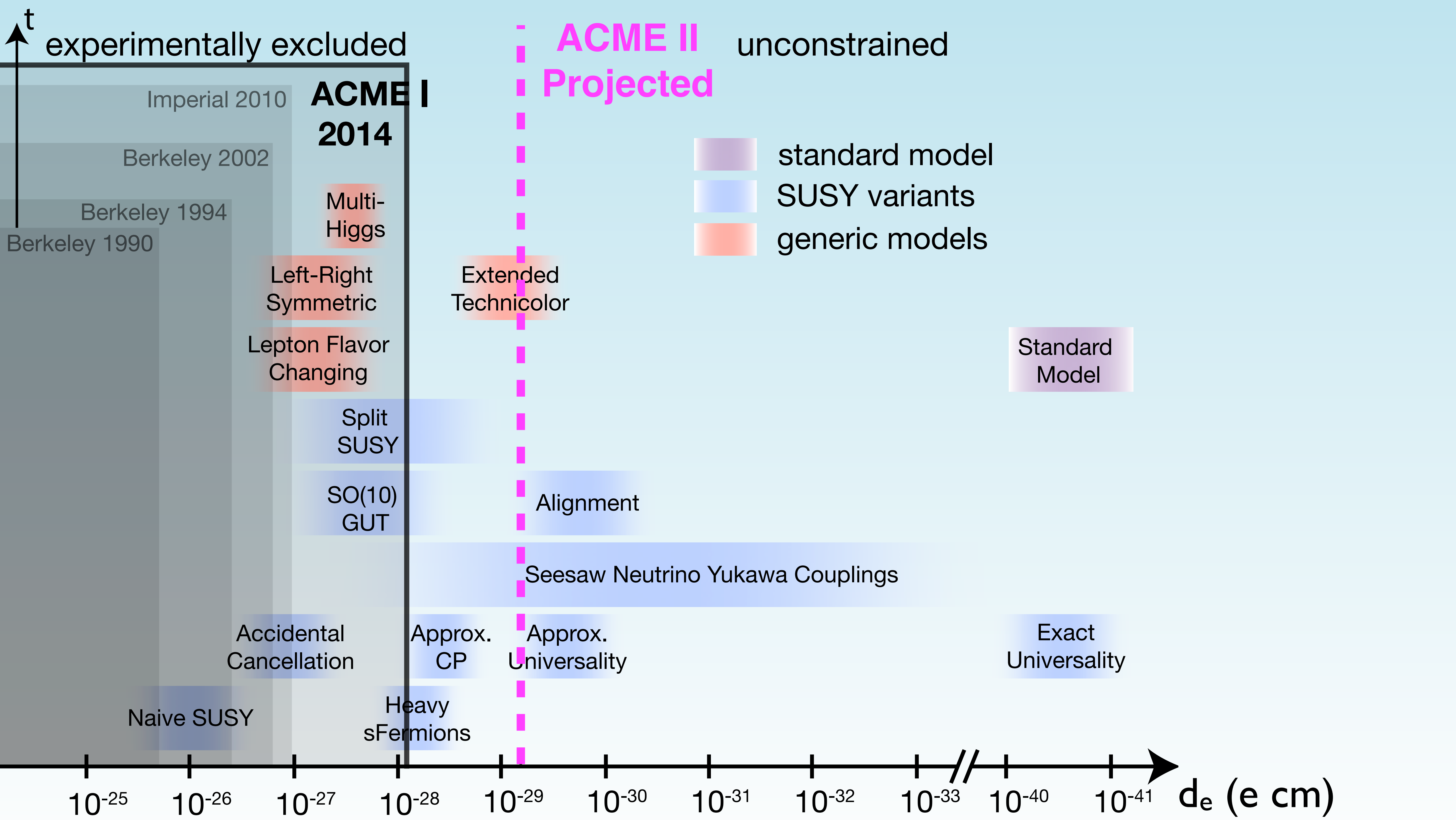
ACME II Preliminary Signals

Single-Trace (25 Shots) Fluorescence Signal



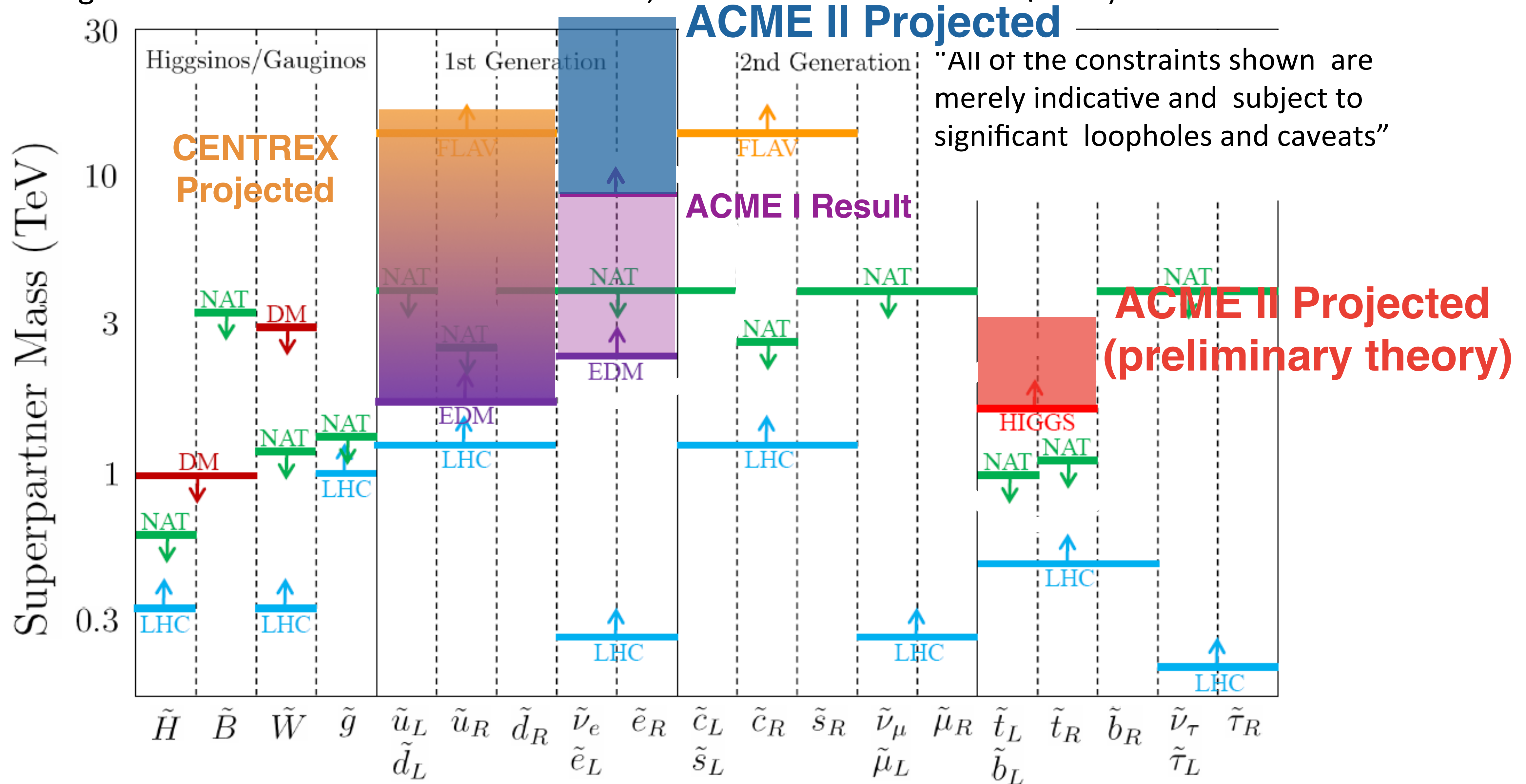
Experiment	One Day Statistical Sensitivity e-cm day ^{-1/2}	Published Limit $ d_e <$ in e-cm	Improvement			EDM Sensitivity Gain over Previous Experiment
			1	2	3	
Berkeley TI	0.5×10^{-27}	1.6×10^{-27}	used a	molecule		~x1
Imperial YbF	2×10^{-27}	1.5×10^{-27}	Beam Source x6.5	Molecule to ThO x2	Technical x1.5	x15
ACME I	1×10^{-28}	0.9×10^{-28}	Geometry x3	STIRAP x3.5	Detection x2.2	x20 <i>Projected</i>
<i>ACME II</i>	0.5×10^{-29}	<i>Projected</i>				





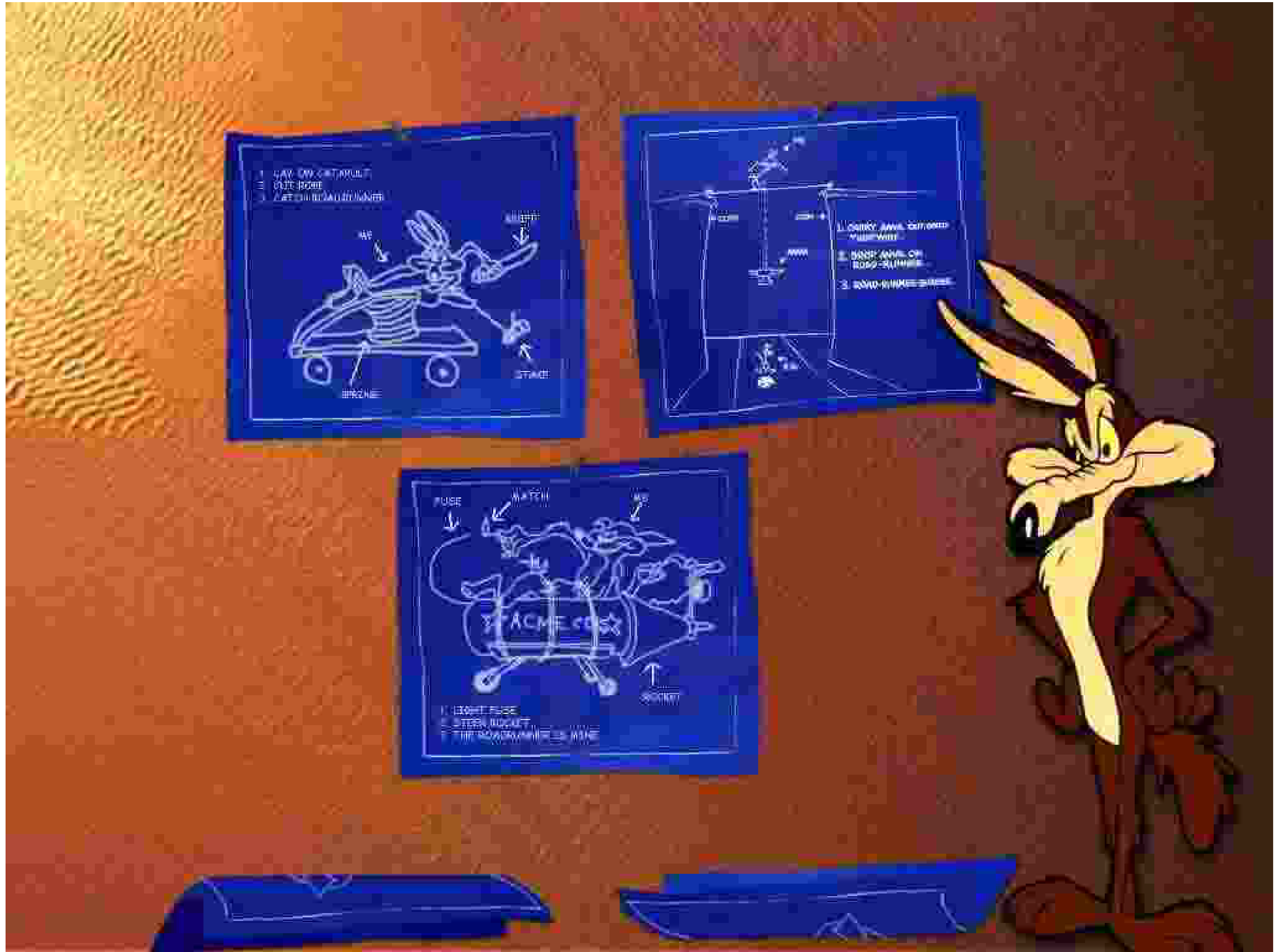
ACME II Projected Impact - 2017

J. Feng: "Naturalness and the status of SUSY", Ann. Rev. Nucl. Part. Sci. (2013)



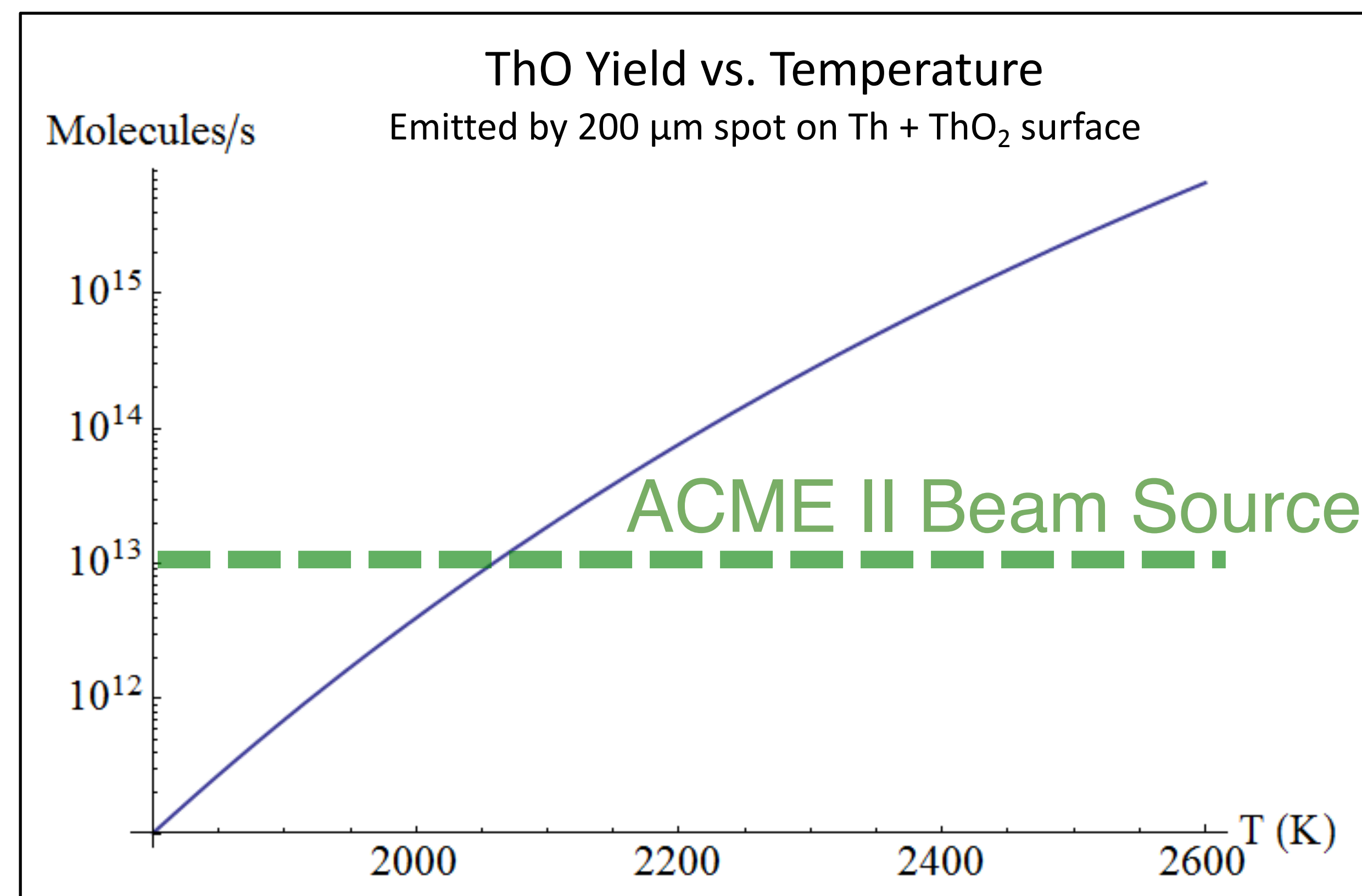
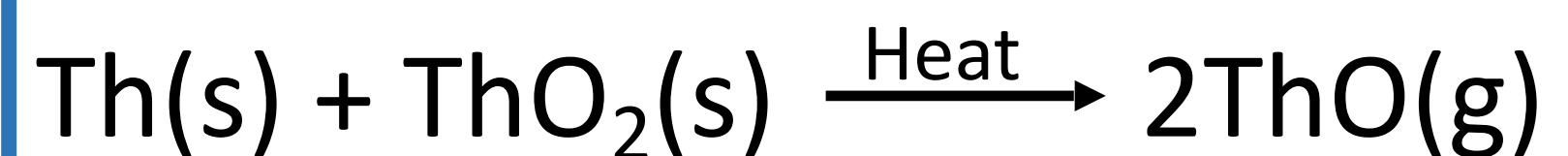
We are still hard at work...

....with ACME II ! **And ACME III !**



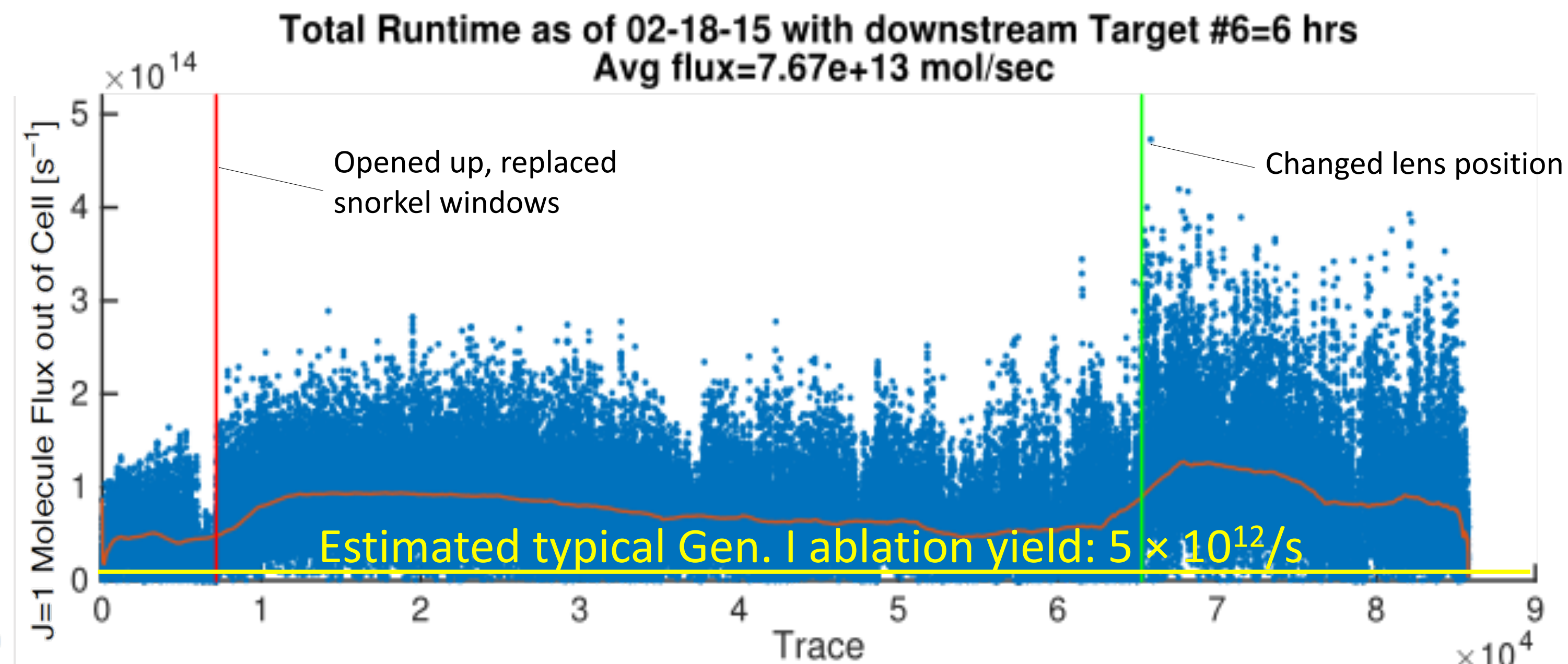
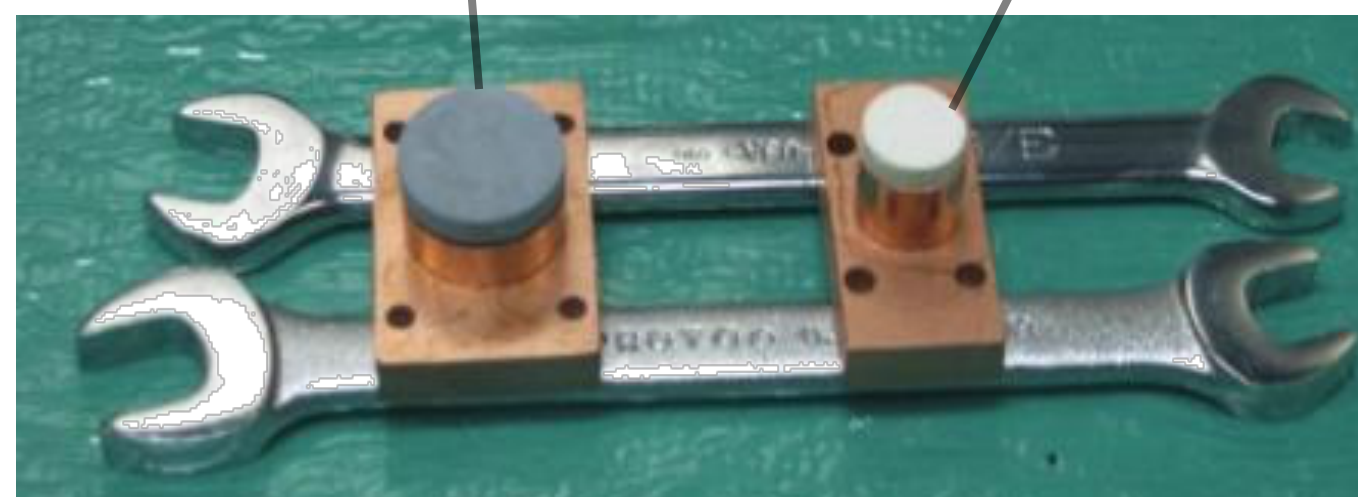
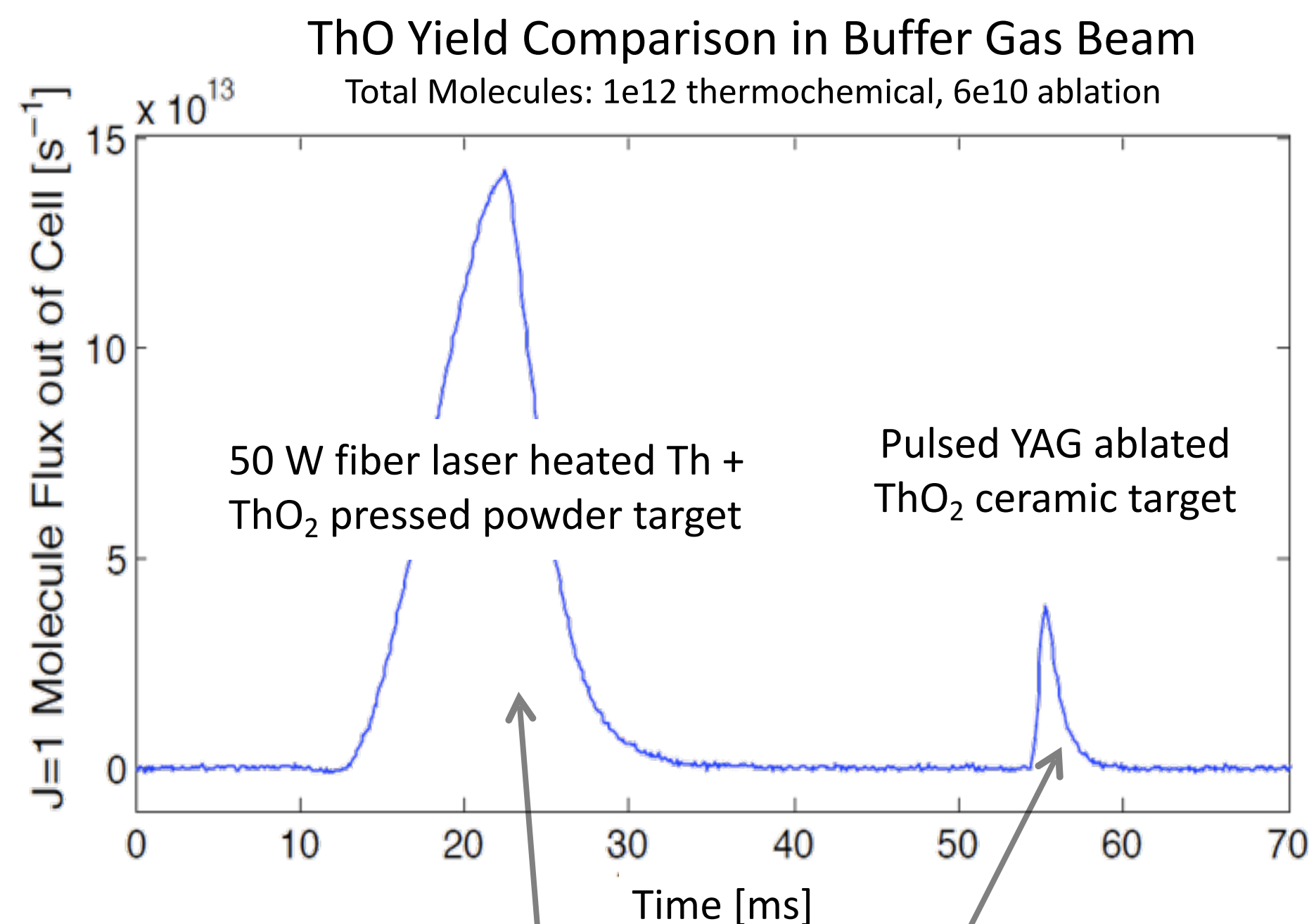
High-Temperature Thorium-Oxygen Chemistry

Chemical reaction with favorable yields >2000 K



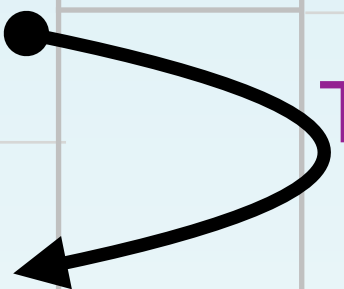
1. Darnell and McCollum, "High Temperature Reactions of Thorium and Thoria and the Vapor Pressure of Thoria," *Atomics International*, September 1961.
2. Hildenbrand and Murad, "Mass Spectrometric Studies of Gaseous ThO and ThO₂," *J. Chem. Phys.*, August 1974.
3. Rand et al., *Chemical Thermodynamics of Thorium*, OECD 2007.

Thermochemical Source Yields



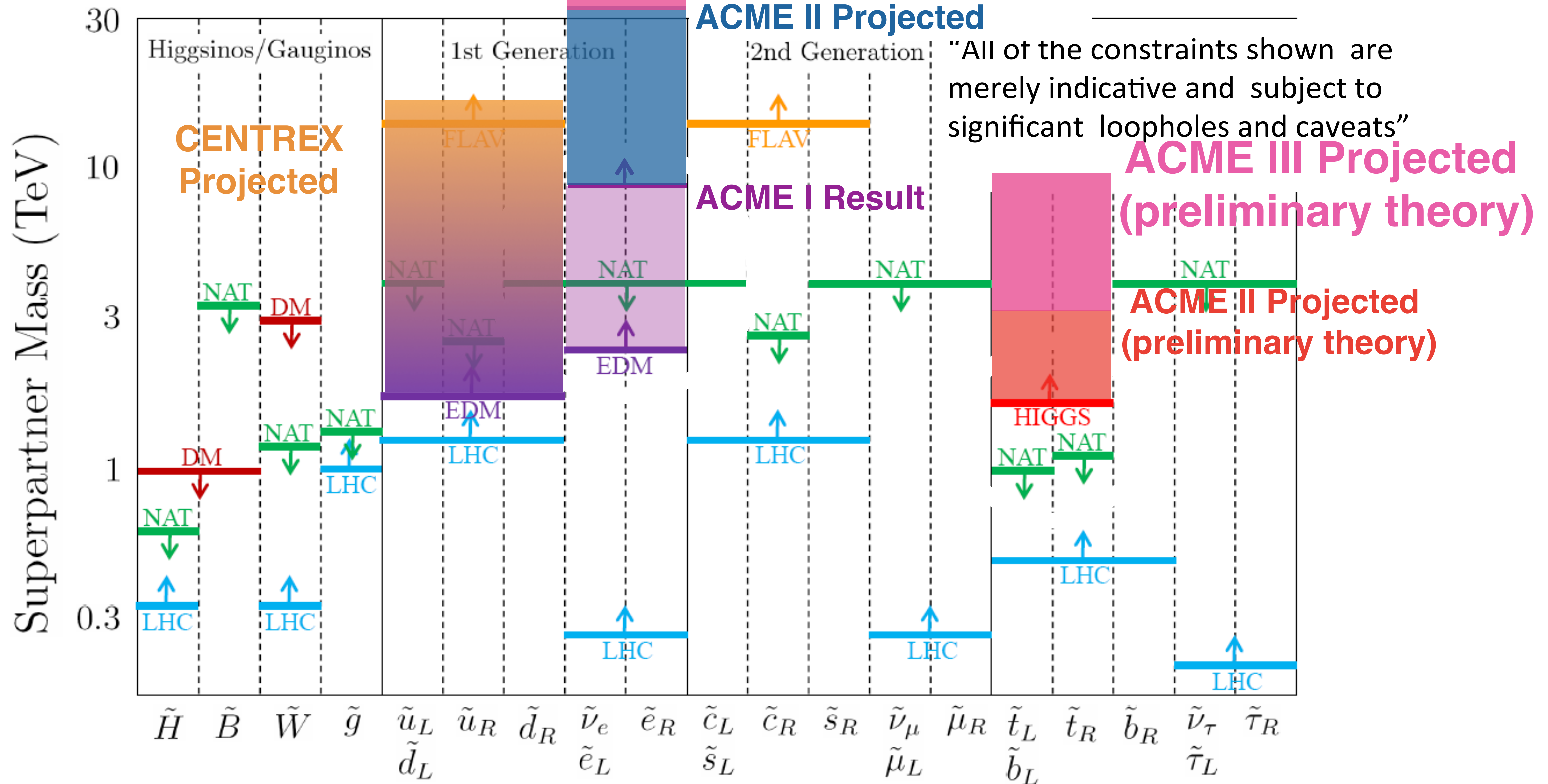
$\sim 10\times$ larger flux than ablation source

Experiment	One Day Statistical Sensitivity e-cm day ^{-1/2}	Published Limit $ d_e <$ in e-cm	Improvement			EDM Sensitivity Gain over Previous Experiment
			1	2	3	
Berkeley Tl	0.5×10^{-27}	1.6×10^{-27}	used a	molecule		~x1
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ACME I	1×10^{-28}	0.9×10^{-28}	Geometry x3	STIRAP x3.5	Detection x2.2	x20 <i>Projected</i>
ACME II	0.5×10^{-29} <i>Projected</i>		Thermochemical Source x3	<i>Electrostatic Lens x1.5</i>	<i>Detection x5</i>	x20 <i>Projected</i>
ACME III	0.3×10^{-30} <i>Projected</i>					



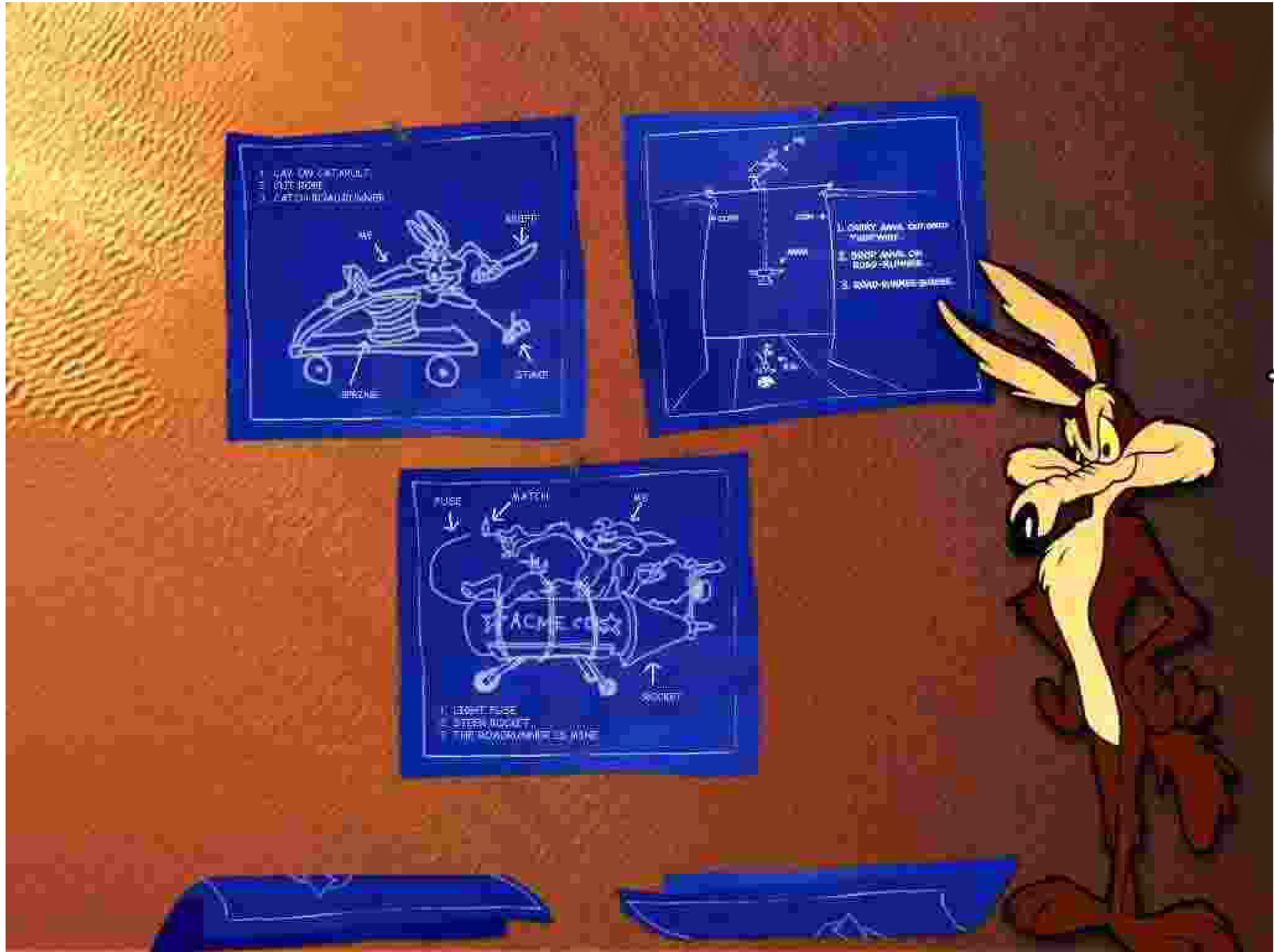
ACME III Projected

J. Feng: "Naturalness and the status of SUSY", Ann. Rev. Nucl. Part. Sci. (2013)



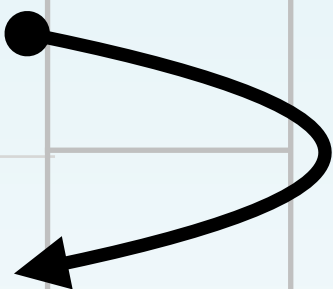
We are still hard at work....

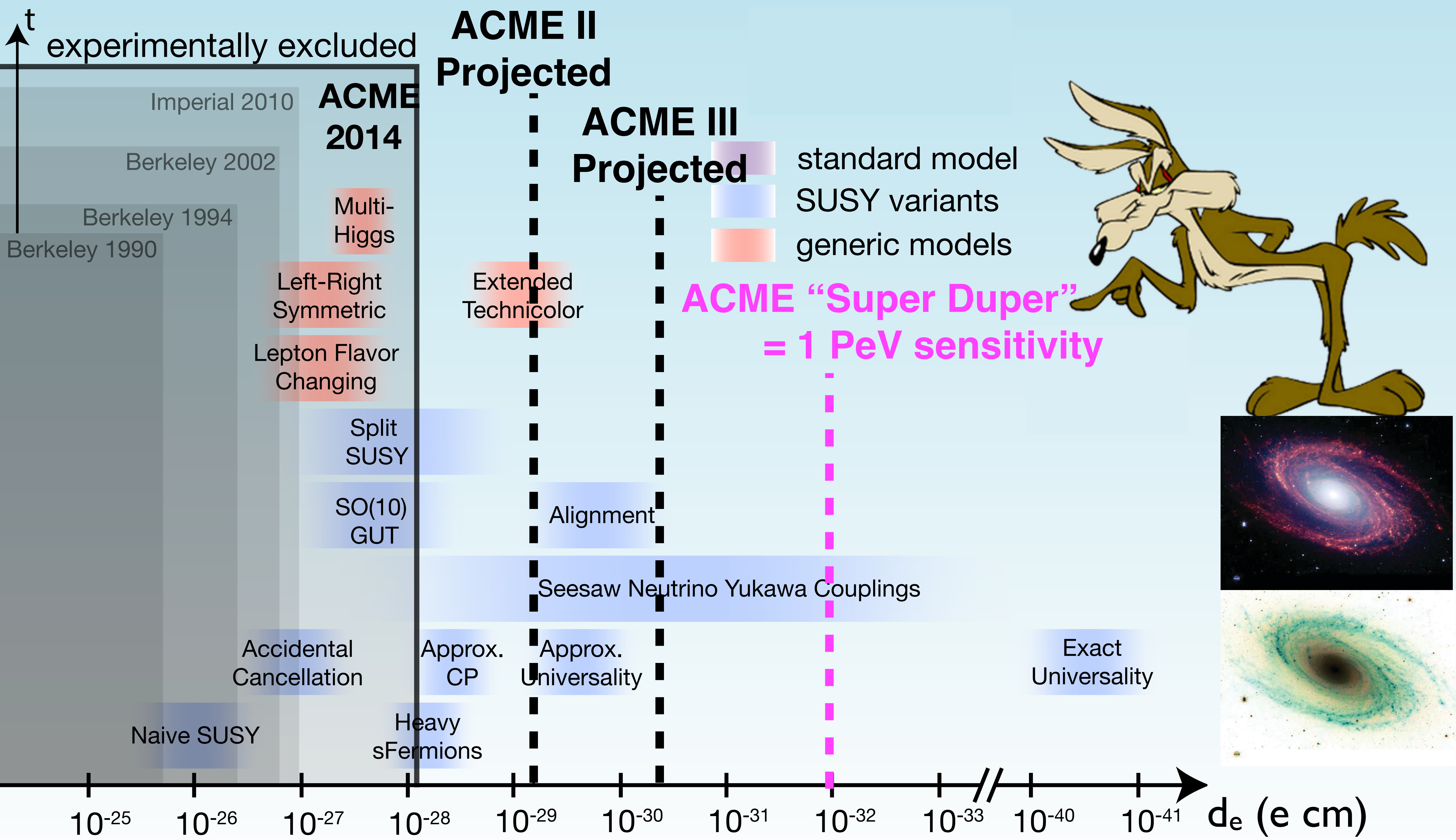
....with ACME II ! And ACME III !

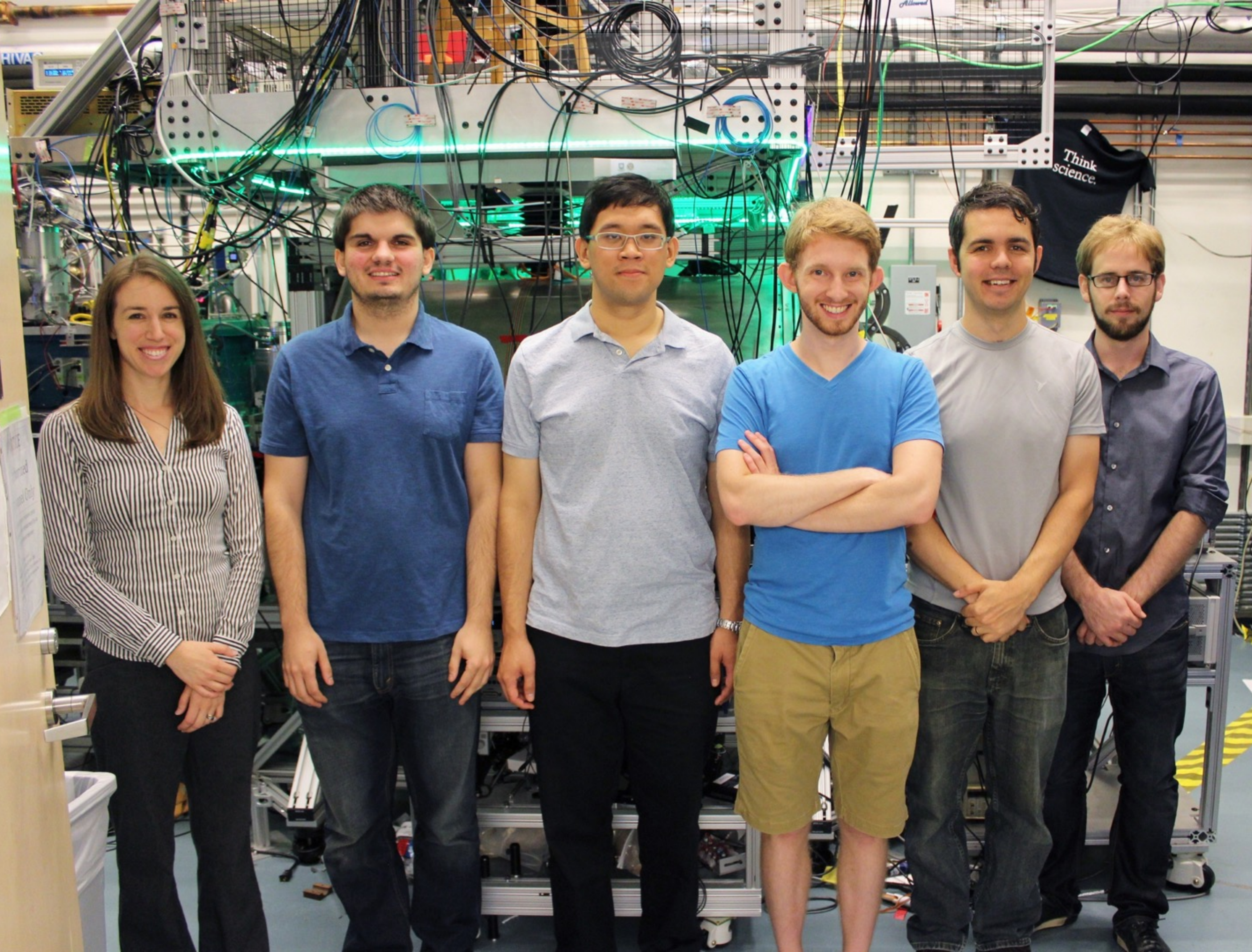


and
ACME
Super-duper!

Experiment	One Day Statistical Sensitivity e-cm day ^{-1/2}	Published Limit $ d_e <$ in e-cm	Improvement			EDM Sensitivity Gain over Previous Experiment
			1	2	3	
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ACME I	1×10^{-28}	0.9×10^{-28}	Geometry x3	STIRAP x3.5	Detection x2.2	x20 <i>Projected</i>
<i>ACME II</i>	0.5×10^{-29} <i>Projected</i>		Thermochemical Source x3	<i>Electrostatic Lens x1.5</i>	<i>Detection x5</i>	x20 <i>Projected</i>
<i>ACME III</i>	0.3×10^{-30} <i>Projected</i>		<i>Advanced Beam Phase Space</i>	<i>Optimized Detection x1.5</i>	?	?
<i>ACME SD</i>	1×10^{-32} (?) <i>Projected</i>					



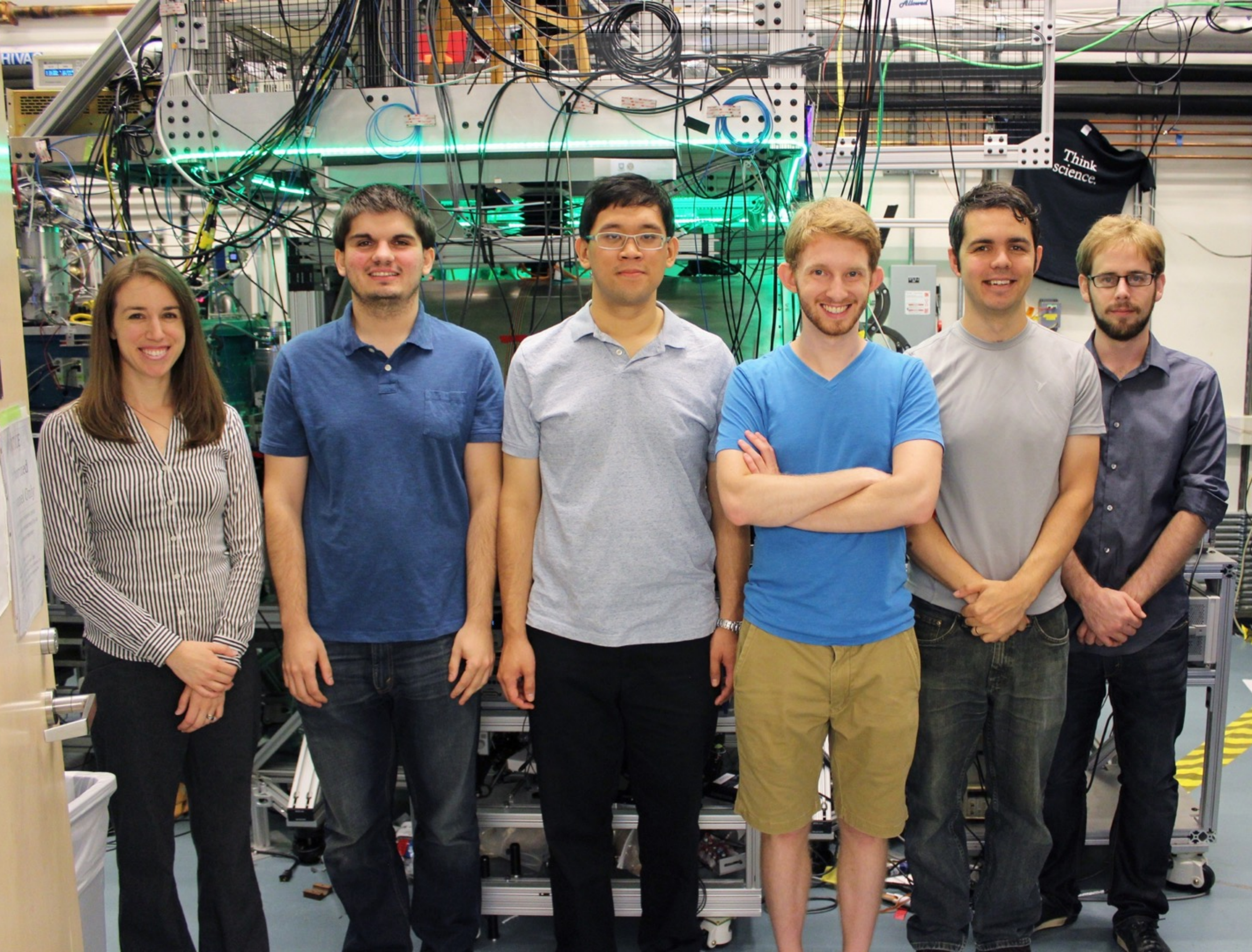




Brendon
O'Leary

Nick
Hutzler

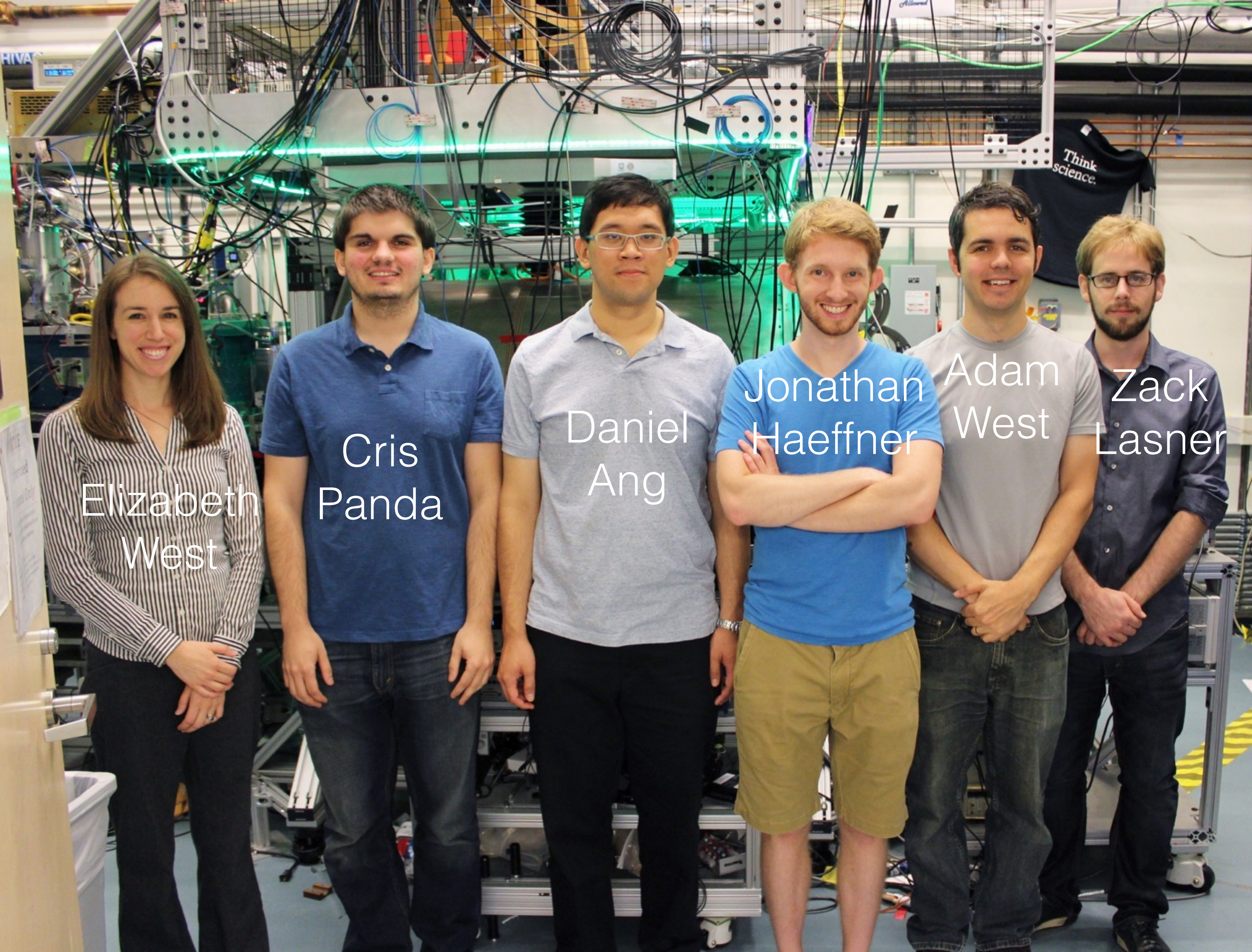




Dave DeMille

Gerald Gabrielse





Elizabeth West

Cris Panda

Daniel Ang

Jonathan Haeffner

Adam West

Zack Lasner

