

SMEFT@NLO

Automated one-loop
computations in the SMEFT

[G. Durieux, C. Degrande, F. Maltoni, KM, C. Zhang, E. Vryonidou; arXiv:2008.11743]

[J. Ellis, M. Madigan, KM, V. Sanz & T. You; arXiv:2012.02779]

<http://feynrules.irmp.ucl.ac.be/wiki/SMEFTatNLO>

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New Physics from Precision at High Energies, KITP

8th of April 2021

SMEFT is...

$$\mathcal{L}_{\text{eff}} = \sum_i \frac{c_i \mathcal{O}_i^D}{\Lambda^{D-4}}$$

Model independent

- Underlying assumptions

Heavy new physics: $M > E_{\text{exp}}$

SM field content & gauge symmetries

Linear EWSB: Higgs = doublet

Systematically improvable

- Double expansion

higher dim. $\frac{E^2}{\Lambda^2}$ & $\{g_s, g, g'\}$ more loops

Global

- Model independence: we don't know what operators NP will generate
- Patterns & correlations among observables are key
- Ultimate goal: complete SMEFT likelihood confronted with HEP data

EWPO, Higgs, multiboson, top, DY, flavor, ...

Established part of LHC programme

SMEFT interpretation

Improving sensitivity:

$$\Delta o_n = o_n^{\text{EXP}} - o_n^{\text{SM}} = \sum_i \frac{a_{n,i}^{(6)}(\mu) c_i^{(6)}(\mu)}{\Lambda^2} + \mathcal{O}\left(\frac{1}{\Lambda^3}\right)$$

Global nature
As many observables as possible
Identify patterns & correlations in fits
Exploit energy-growth

Sensitivity
Experiment: Best measurements & understanding of uncertainties and correlations
Theory: Best available predictions for observables (NLO, NNLO, N3LO, ...)

Interpretation
Relies on accurate knowledge of the size & correlation among a_i
Determining $c_i^{(6)}$ requires most precise available SMEFT predictions

SMEFT@NLO

NLO computations for SMEFT: very active field

- Non-universal K-factors in EFT space \Leftrightarrow new information at NLO
- Loop-induced sensitivity
- Control theoretical uncertainties
- Experimental interest in higher precision for SMEFT analyses/interpretations

Challenge: many processes \times many operators

- LO \Rightarrow NLO = more cross-talk/operators/complexity
- Automated tools for fixed-order/NLO+PS are essential to the LHC programme

SMEFT@NLO

[Degrade et al.; arXiv:2008.11743]
<http://feynrules.irmp.ucl.ac.be/wiki/SMEFTatNLO>

- UFO model for `MadGraph5_aMC@NLO`
- Process-independent implementation: SMEFT in top-specific flavor limit

Standard Model Effective Theory at One-Loop in QCD

Céline Degrande, Gauthier Durieux, Fabio Maltoni, Ken Mimasu, Eleni Vryonidou & Cen Zhang, [arXiv:2008.11743](#)

The implementation is based on the Warsaw basis of dimension-six SMEFT operators, after canonical normalization. Electroweak input parameters are taken to be G_F , M_Z , M_W . The CKM matrix is approximated as a unit matrix, and a $U(2)_q \times U(2)_u \times U(3)_d \times (U(1)_l \times U(1)_e)^3$ flavor symmetry is enforced. It forbids all fermion masses and Yukawa couplings except that of the top quark. The model therefore implements the five-flavor scheme for PDFs.

A new coupling order, **NP=2**, is assigned to SMEFT interactions. The cutoff scale **Lambda** takes a default value of 1 TeV^{-2} and can be modified along with the Wilson coefficients in the **param_card**. Operators definitions, normalisations and coefficient names in the UFO model are specified in [definitions.pdf](#). The notations and normalizations of top-quark operator coefficients comply with the LHC TOP WG standards of [1802.07237](#). Note however that the flavor symmetry enforced here is slightly more restrictive than the baseline assumption there (see the [dim6top page](#) for more information). This model has been validated at tree level against the **dim6top** implementation (see [1906.12310](#) and the [comparison details](#)).

Current implementation

UFO model: [SMEFTatNLO_v1.0.tar.gz](#)

The current implementation imposes CP conservation. In the quark sector, it focuses primarily on top-quark interactions. The light-quark current operator, $qq\bar{H}D\bar{H}$, $uu\bar{H}D\bar{H}$, $dd\bar{H}D\bar{H}$, with coefficients **cpq3i**, **cpqMi**, **cpu**, **cpd** are however included. The triple-gluon operator, with coefficient **cG**, is currently not available (see the loop-capable **GGG** implementation). Vertices including more than four scalars or four leptons are not included. Scalar and tensor **QQ11** operators, with coefficients **ct1S3**, **ct1T3**, and **cb1S3**, break our flavor symmetry assumption and are not available for one-loop computations. Top-quark flavor-changing interactions, not compatible with the imposed flavor symmetry, are not included (see the loop-capable [TopFCNC](#) implementation).

Unlike prescribed by the LHC TOP WG, the top quark chromomagnetic-dipole operator coefficient **ctG** is normalized with a factor of the strong coupling, g_S . This normalization factor temporarily ensures compatibility with the 2.X.X series of MadGraph5_aMC@NLO but may be dropped in the future. As with every other appearance of this coupling in MadGraph5_aMC@NLO, its value is renormalisation-group evolved to the QCD renormalization scale (set in the `run_card`).

```
MG5_aMC>import model SMEFTatNLO
MG5_aMC>generate p p > t t~ NP=2 [QCD]
MG5_aMC>output
MG5_aMC>launch
```

'QCD' loops
*coloured particles,
strong coupling or
4-fermion couplings*

What's in the box?

'Warsaw' basis

[Grzadkowski et al.; JHEP 1010 (2010) 085]

X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	Q_φ	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
$Q_{\widetilde{G}}$	$f^{ABC} \widetilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$Q_{\varphi\square}$	$(\varphi^\dagger \varphi) \square (\varphi^\dagger \varphi)$	$Q_{u\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p u_r \widetilde{\varphi})$
Q_W	$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$Q_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^\star (\varphi^\dagger D_\mu \varphi)$	$Q_{d\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p d_r \varphi)$
$Q_{\widetilde{W}}$	$\varepsilon^{IJK} \widetilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{l}_p \gamma^\mu l_r)$
$Q_{\varphi \widetilde{G}}$	$\varphi^\dagger \varphi \widetilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \widetilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi \widetilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \widetilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \widetilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi \widetilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}_p \gamma^\mu u_r)$
$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}_p \gamma^\mu d_r)$
$Q_{\varphi \widetilde{WB}}$	$\varphi^\dagger \tau^I \varphi \widetilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\widetilde{\varphi}^\dagger D_\mu \varphi)(\bar{u}_p \gamma^\mu d_r)$

$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
Q_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	Q_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	Q_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	Q_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	Q_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	Q_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$Q_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	Q_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$Q_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$Q_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$Q_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		B-violating			
Q_{ledq}	$(\bar{l}_p^j e_r)(\bar{d}_s q_t^j)$	Q_{duq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^\gamma)^T C l_t^k]$		
$Q_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	Q_{qqu}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$Q_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	$Q_{qqq}^{(1)}$	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} \varepsilon_{mn} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^m)^T C l_t^n]$		
$Q_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	$Q_{qqq}^{(3)}$	$\varepsilon^{\alpha\beta\gamma} (\tau^I \varepsilon)_{jk} (\tau^I \varepsilon)_{mn} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^m)^T C l_t^n]$		
$Q_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$	Q_{duu}	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		

Some symmetries imposed to control parameter space

- CP, B and flavor conservation
- Top-specific flavour structure of 2 & 4 fermion operators

Flavor symmetry

Approximate flavor symmetry in the SM

- SM: broken by Yukawa interactions
- SMEFT: broken by $\psi^2 X \varphi, \psi^2 \varphi^3, (\bar{L}R)(\bar{L}R), (\bar{L}R)(\bar{R}L) \& \mathcal{O}_{\varphi ud}$
- + any off-diagonal or non-universal entries of other 2F operators

SMEFTatNLO: minimal extension to single out top quark

universal	$U(3)_L \times U(3)_e \times U(3)_Q \times U(3)_u \times U(3)_d$	
top	$U(3)_L \times U(3)_e \times U(2)_Q \times U(2)_u \times U(3)_d$	

*cf. Minimal
flavor violation*

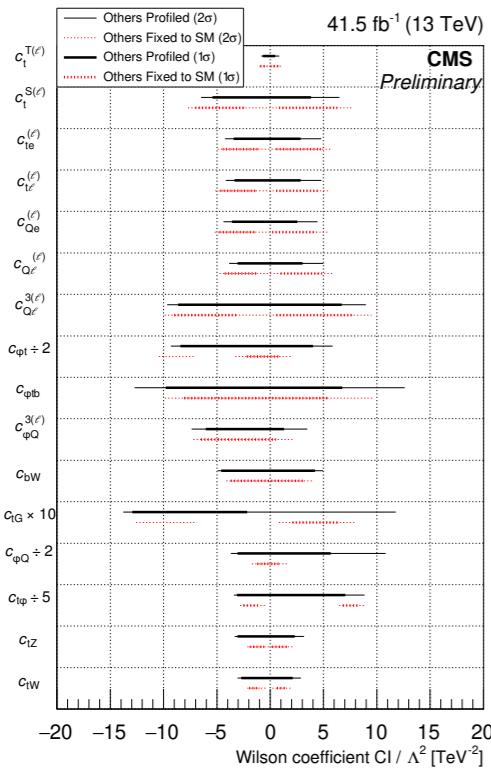
[Buras et al.; PLB 500
(2001) 161]
[D'Ambrosio et al.; NPB
645 (2002) 155]

See **dim6top**

[Aguilar-Saavedra et al.;
arXiv:1802.07237]

Yukawa	$\psi^2 H^3 : (\varphi^\dagger \varphi)^2 (\bar{Q} t \tilde{\varphi})$
Dipoles	$\psi^2 X H : (\bar{Q} \sigma^{\mu\nu} t \tilde{\varphi}) B_{\mu\nu} [W_{\mu\nu}^I, G_{\mu\nu}^a]$
3rd gen. currents	$\psi^2 H^2 D : (\varphi^\dagger \overleftrightarrow{D}_\mu \varphi) (\bar{Q} \gamma^\mu Q) [(\bar{Q} \gamma^\mu \tau^I Q), (\bar{t} \gamma^\mu t), \dots]$
3rd gen. 4F	$\psi^4 : (\bar{Q} \gamma^\mu Q) (\bar{q} \gamma_\mu q), (\bar{Q} \gamma^\mu Q) (\bar{Q} \gamma_\mu Q), \dots$

ttH, ttZ, ttW,
tW, tZ, tH



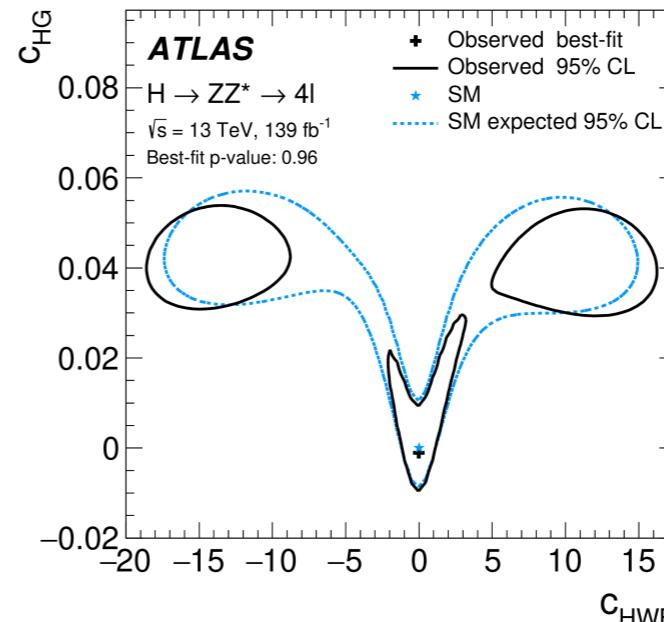
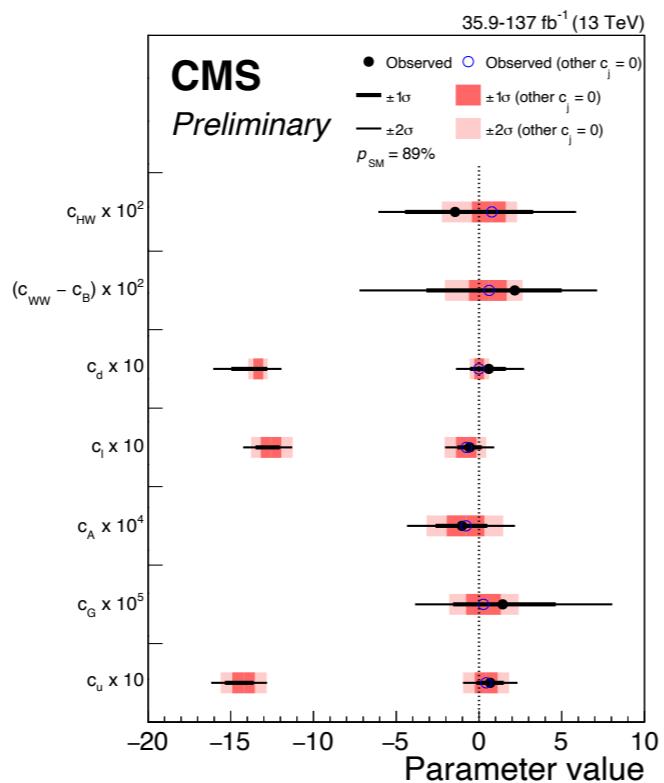
[ATLAS; PRD 99 (2017) 072009]

Coefficients	$\mathcal{C}_{\phi Q}^{(3)}/\Lambda^2$	$\mathcal{C}_{\phi t}/\Lambda^2$	$\mathcal{C}_{tB}/\Lambda^2$	$\mathcal{C}_{tW}/\Lambda^2$
Previous indirect constraints at 68% CL	[-4.7, 0.7]	[-0.1, 3.7]	[-0.5, 10]	[-1.6, 0.8]
Previous direct constraints at 95% CL	[-1.3, 1.3]	[-9.7, 8.3]	[-6.9, 4.6]	[-0.2, 0.7]
Expected limit at 68% CL	[-2.1, 1.9]	[-3.8, 2.7]	[-2.9, 3.0]	[-1.8, 1.9]
Expected limit at 95% CL	[-4.5, 3.6]	[-23, 4.9]	[-4.2, 4.3]	[-2.6, 2.6]
Observed limit at 68% CL	[-1.0, 2.7]	[-2.0, 3.5]	[-3.7, 3.5]	[-2.2, 2.1]
Observed limit at 95% CL	[-3.3, 4.2]	[-25, 5.5]	[-5.0, 5.0]	[-2.9, 2.9]
Expected limit at 68% CL (linear)	[-1.9, 2.0]	[-3.0, 3.2]	—	—
Expected limit at 95% CL (linear)	[-3.7, 4.0]	[-5.8, 6.3]	—	—
Observed limit at 68% CL (linear)	[-1.0, 2.9]	[-1.8, 4.4]	—	—
Observed limit at 95% CL (linear)	[-2.9, 4.9]	[-4.8, 7.5]	—	—

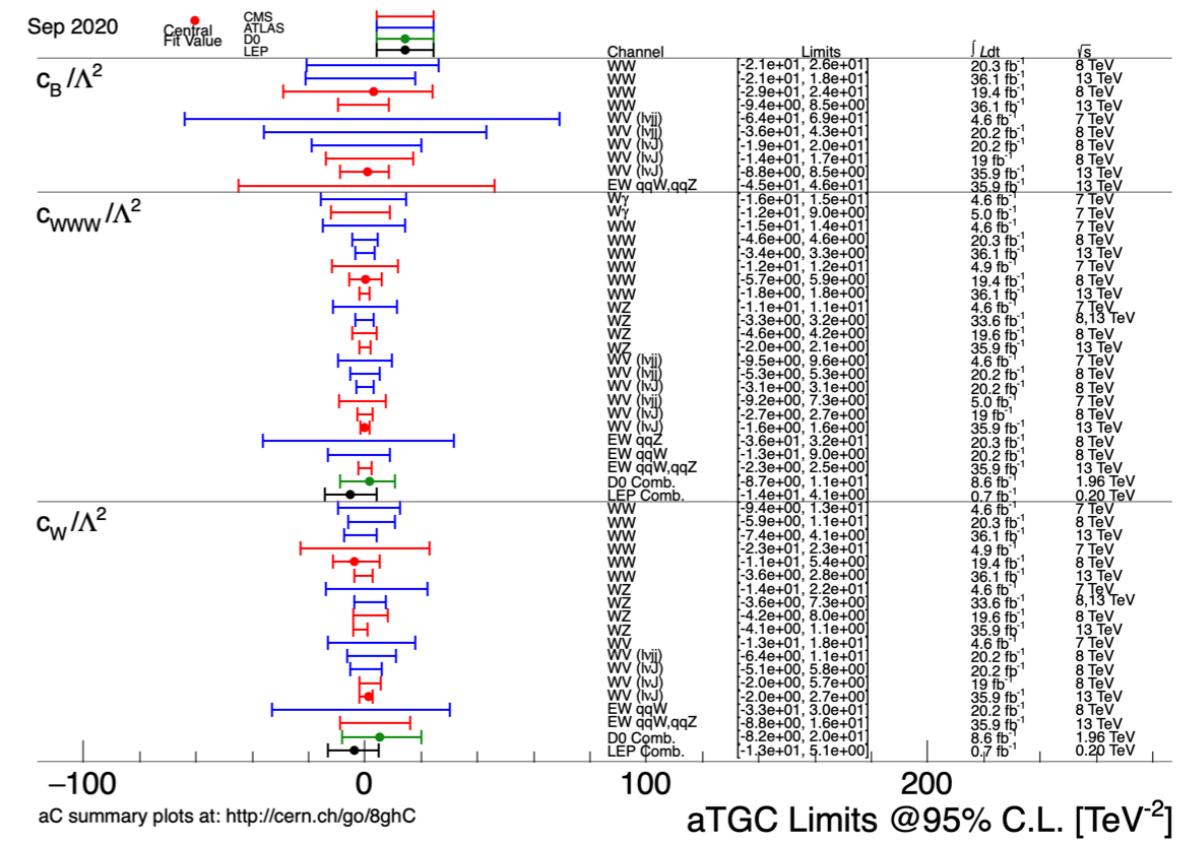
Higgs combination

 $H \rightarrow 4l$

NLO QCD
OK ✓



WW & WZ



hh (LO),
VVV,
VBS,
tt,
4top,
...

Selected results

Some from previous works, superseded by SMEFT@NLO
A few, simple new results presented in 2008.11743

Predictions

Dim-6 SMEFT: $\mathcal{A} = \mathcal{A}_{SM} + \sum_i \mathcal{A}_i \frac{C_i}{\Lambda^2} + \mathcal{O}(\Lambda^{-3})$

$\{\mathcal{A}_i\} \Rightarrow \{\sigma_i, \sigma_{ij}\}$ $\sigma = \sigma_{SM} + \sum_i \sigma_i \frac{C_i}{\Lambda^2} + \sum_{j \geq i} \sigma_{ij} \frac{C_i C_j}{\Lambda^4} + \mathcal{O}(\Lambda^{-4})$

Higher orders (dim > 6) unspecified

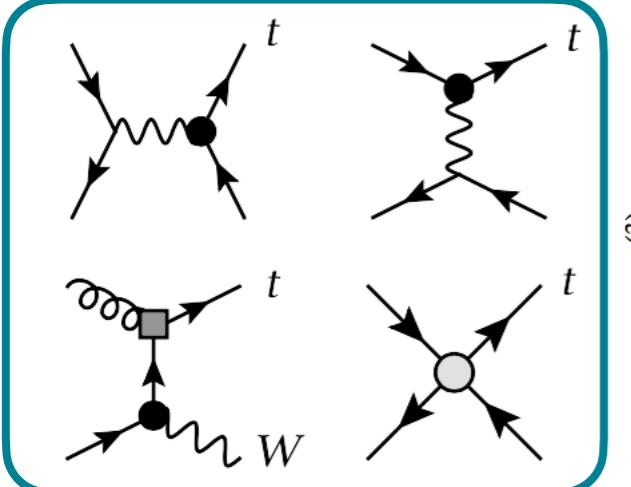
- σ_{ij} formally same order as $\sigma_i^{(8)}$
- Relative importance is model/power-counting dependent
- EFT validity assessment depends further on data sensitivity

We always report both: $\{\sigma_i, \sigma_{ij}\}$

- σ_{ij} contain valuable information
- Can be used in a variety of ways (included in prediction, error estimates,...)

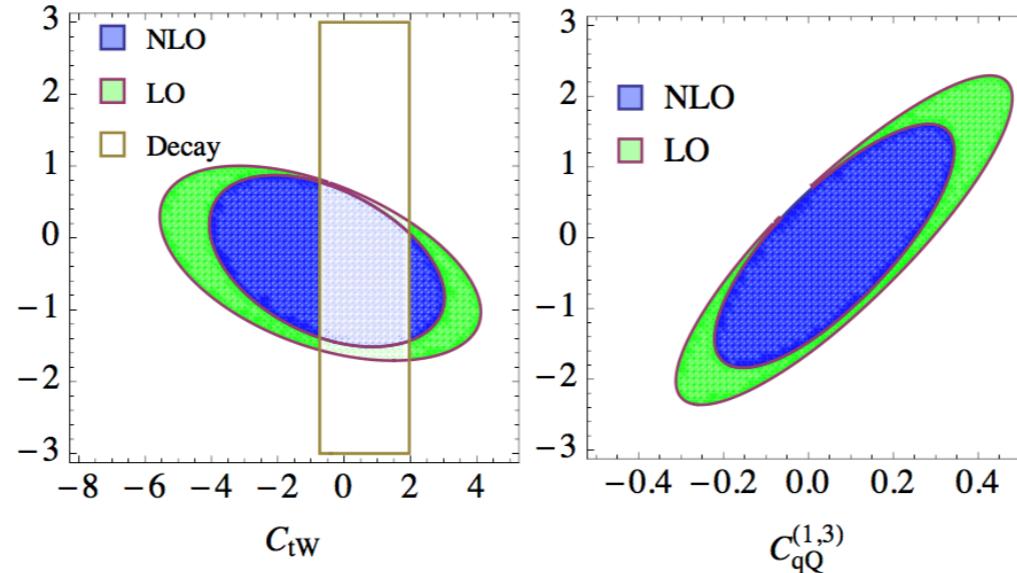
Single top

[Zhang; PRL 116 (2016) 162002]

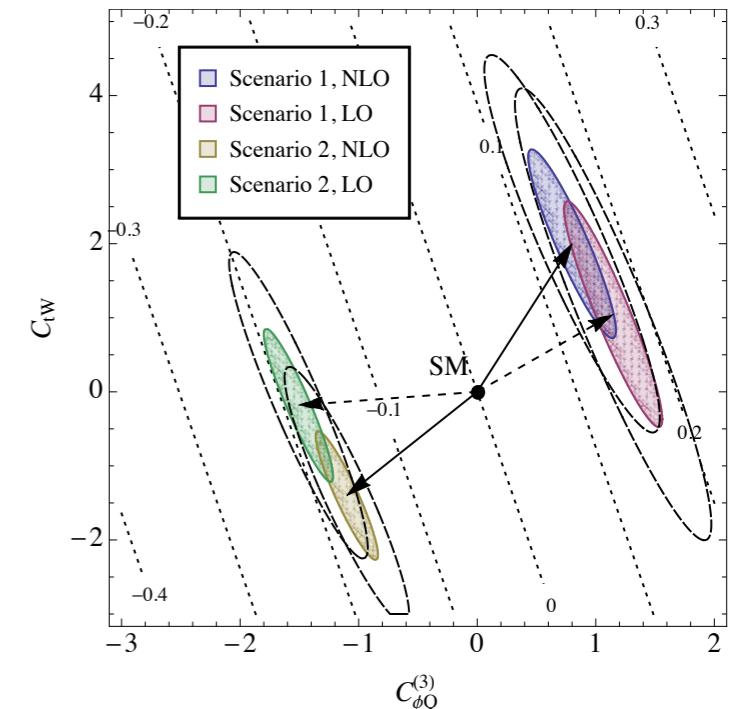


tj, tb, tW

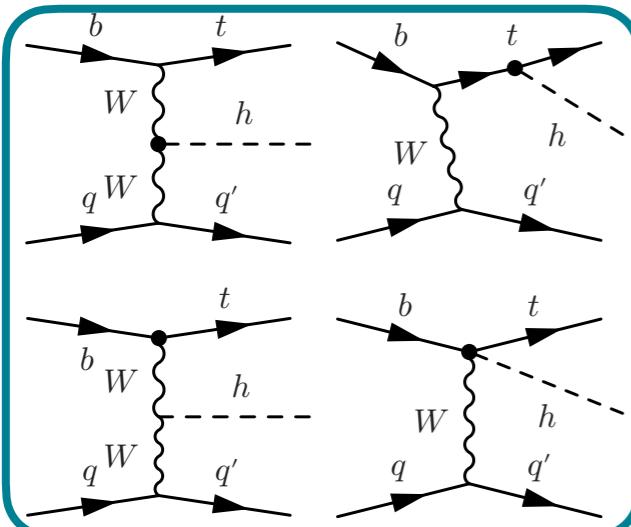
Fit without deviation



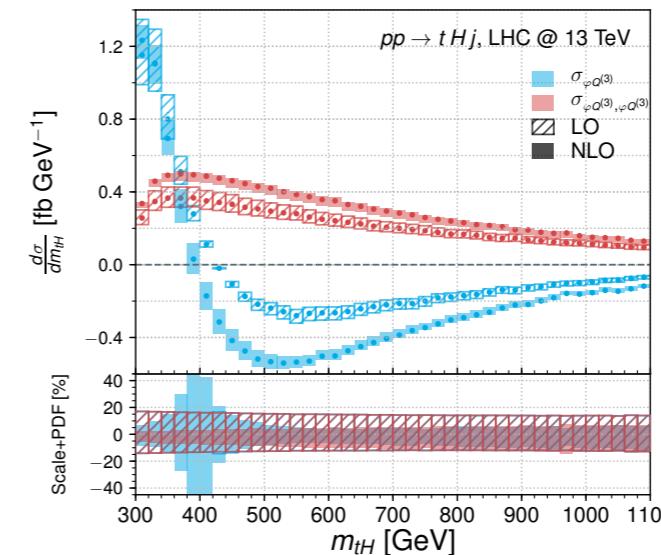
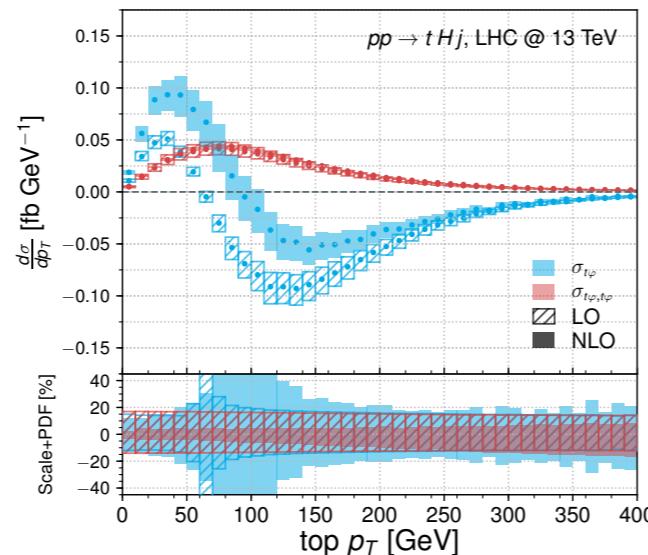
Fit with a
(hypothetical) deviation



[Degrande et al.; JHEP 10 (2018) 005]



$t Z j \& t H j$

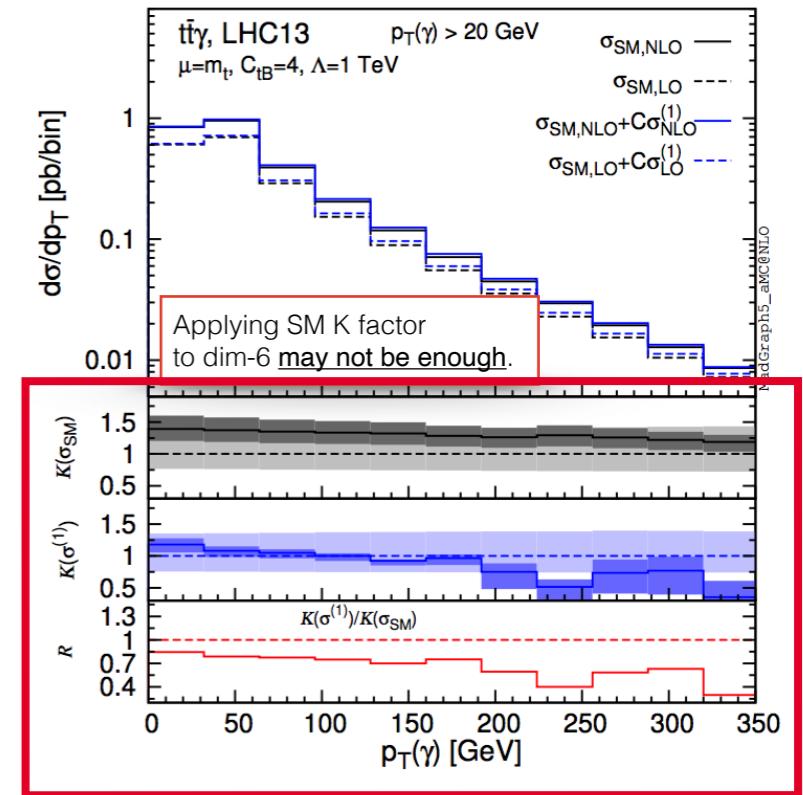
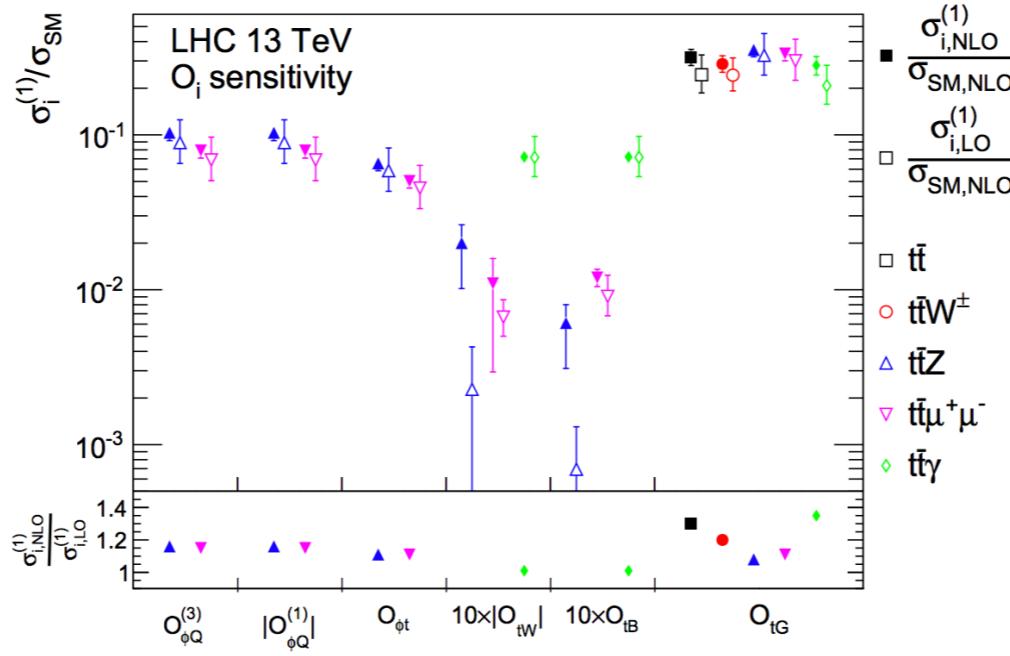
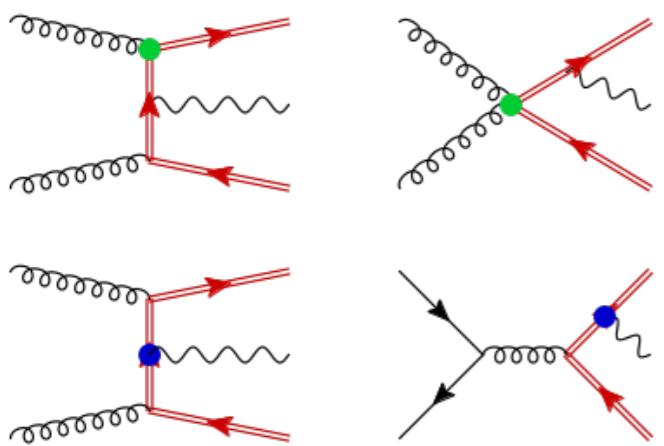


	σ [fb]	K-factor
σ_{SM}		1.32
$\sigma_{\varphi W}$		0.96
$\sigma_{\varphi W, \varphi W}$		1.20
$\sigma_{t\varphi}$	0.20	0.20
$\sigma_{t\varphi, t\varphi}$		1.09
σ_{tW}		1.14
$\sigma_{tW, tW}$		1.54
$\sigma_{\varphi Q^{(3)}}^{(3)}$		3.31
$\sigma_{\varphi Q^{(3)}, \varphi Q^{(3)}}$		1.36

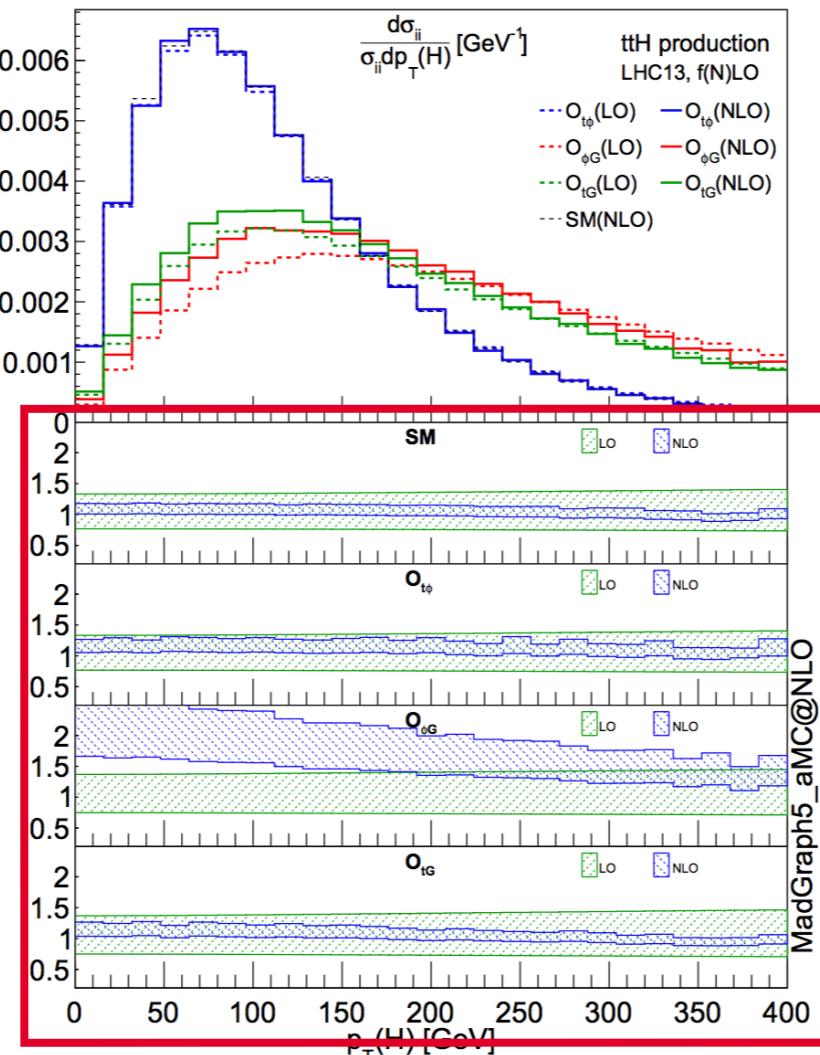
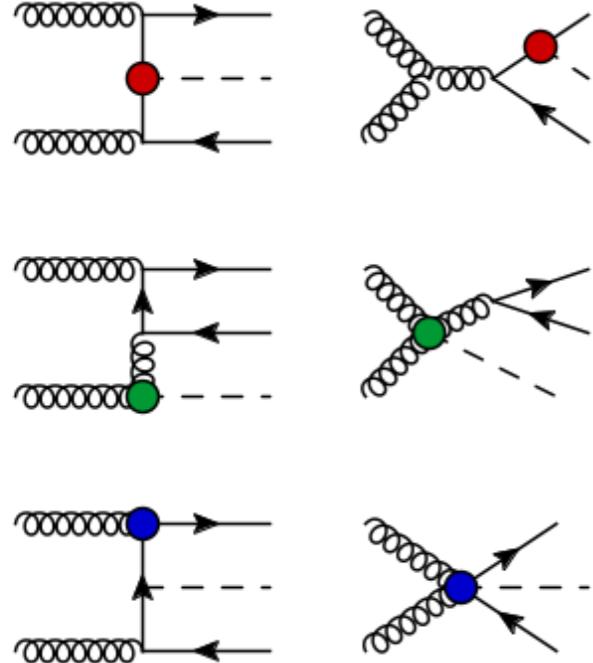
Different patterns of phase-space cancellations at LO/NLO lead to non-trivial & “strange” K factors

NLO > LO

$t\bar{t}\gamma/t\bar{t}Z$



$t\bar{t}H$



13 TeV	$\sigma \text{ LO}$	$\sigma \text{ NLO}$	K
σ_{SM}	$0.464^{+0.16}_{-0.11}$	$0.507^{+0.0}_{-0.0}$	1.09
$\sigma_{t\phi}$	$-0.055^{+0.0}_{-0.0}$	-0.062^{+0}_{-0}	1.13
$\sigma_{\phi G}$	$0.627^{+0.22}_{-0.15}$	$0.872^{+0.1}_{-0.1}$	1.39
σ_{tG}	$0.470^{+0.16}_{-0.11}$	$0.503^{+0.0}_{-0.0}$	1.07
$\sigma_{t\phi, t\phi}$	$0.0016^{+0.0}_{-0.0}$	$0.0019^{+0.0}_{-0.0}$	1.17
$\sigma_{\phi G, \phi G}$	$0.646^{+0.27}_{-0.17}$	$1.021^{+0.2}_{-0.1}$	1.58
$\sigma_{tG, tG}$	$0.645^{+0.27}_{-0.17}$	$0.674^{+0.0}_{-0.0}$	1.04
$\sigma_{t\phi, \phi G}$	$-0.037^{+0.0}_{-0.0}$	-0.053^{+0}_{-0}	1.42
$\sigma_{t\phi, tG}$	$-0.028^{+0.0}_{-0.0}$	-0.031^{+0}_{-0}	1.10
$\sigma_{\phi G, tG}$	$0.627^{+0.25}_{-0.16}$	$0.859^{+0.1}_{-0.1}$	1.37

Non-universal K-factors
in rates & distributions

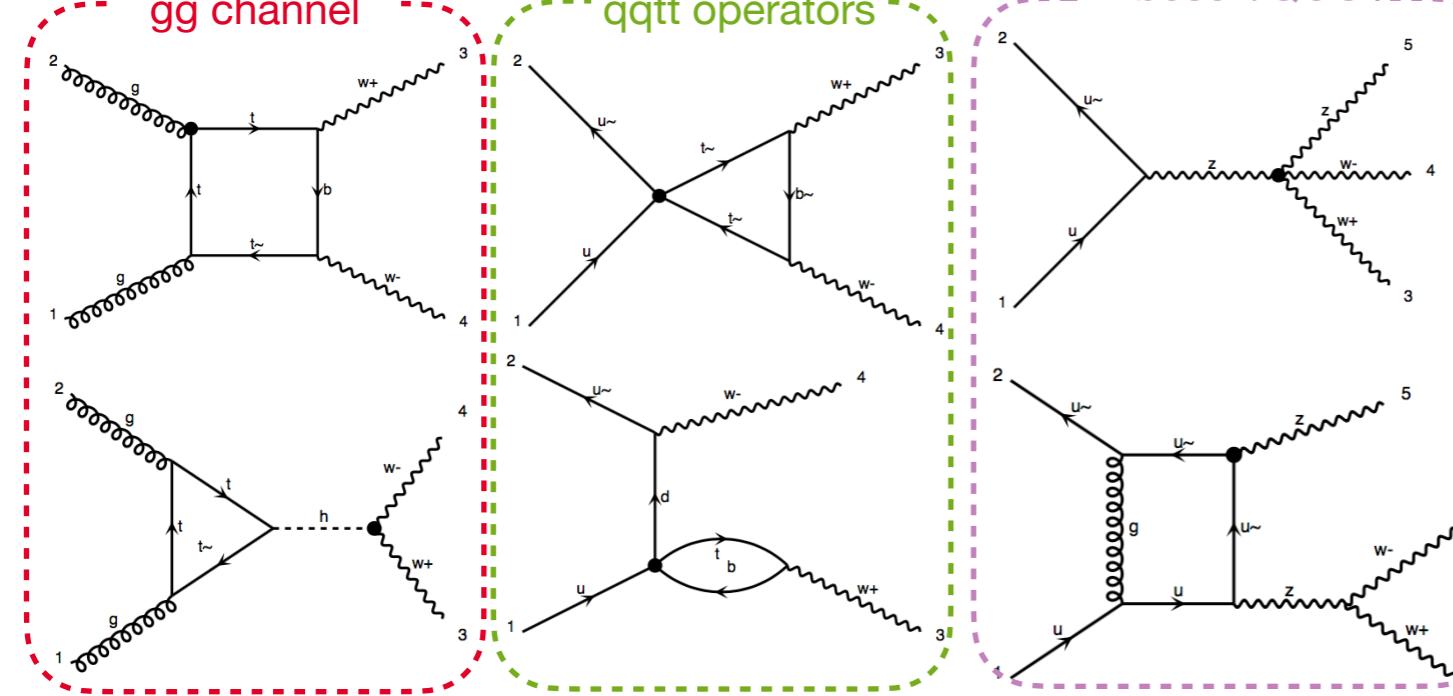
Multiboson

WW/WZ/ZZ/WWW/WWZ/ZZW/ZZZ

Loop-induced
gg channel

Loop-induced
qqt operators

Triboson/QGC



Non-universal NLO corrections, different from SM

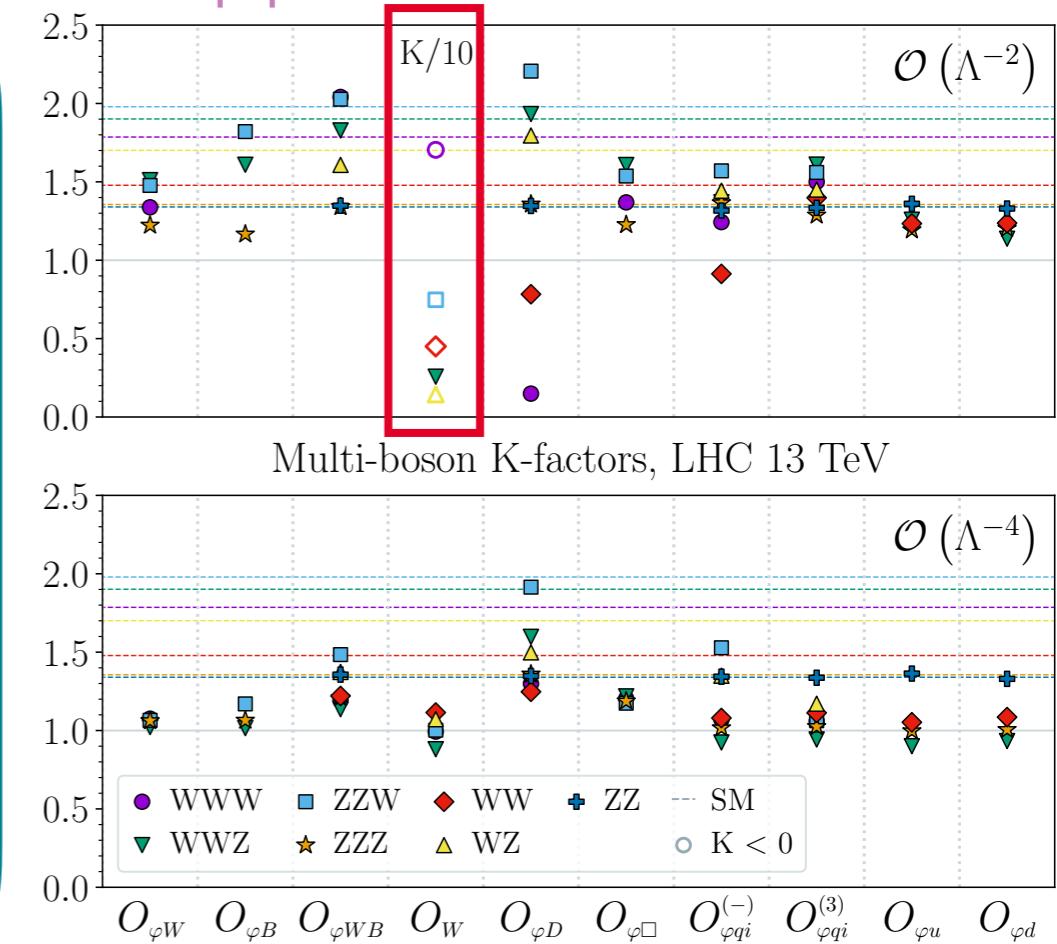
Large, negative K-factors for triple gauge operator, c_w

Non-interference/cancellation at LO broken at NLO

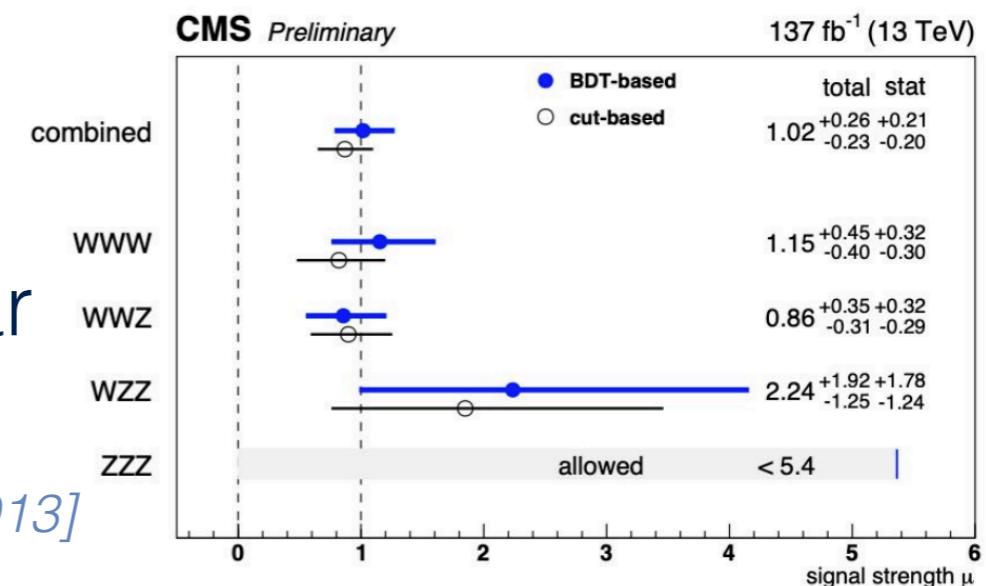
First triboson observation by CMS last year

- Also strong evidence from ATLAS
- New window into SMEFT? [ATLAS; PLB 798 (2019) 134913]

qq-initiated K-factors

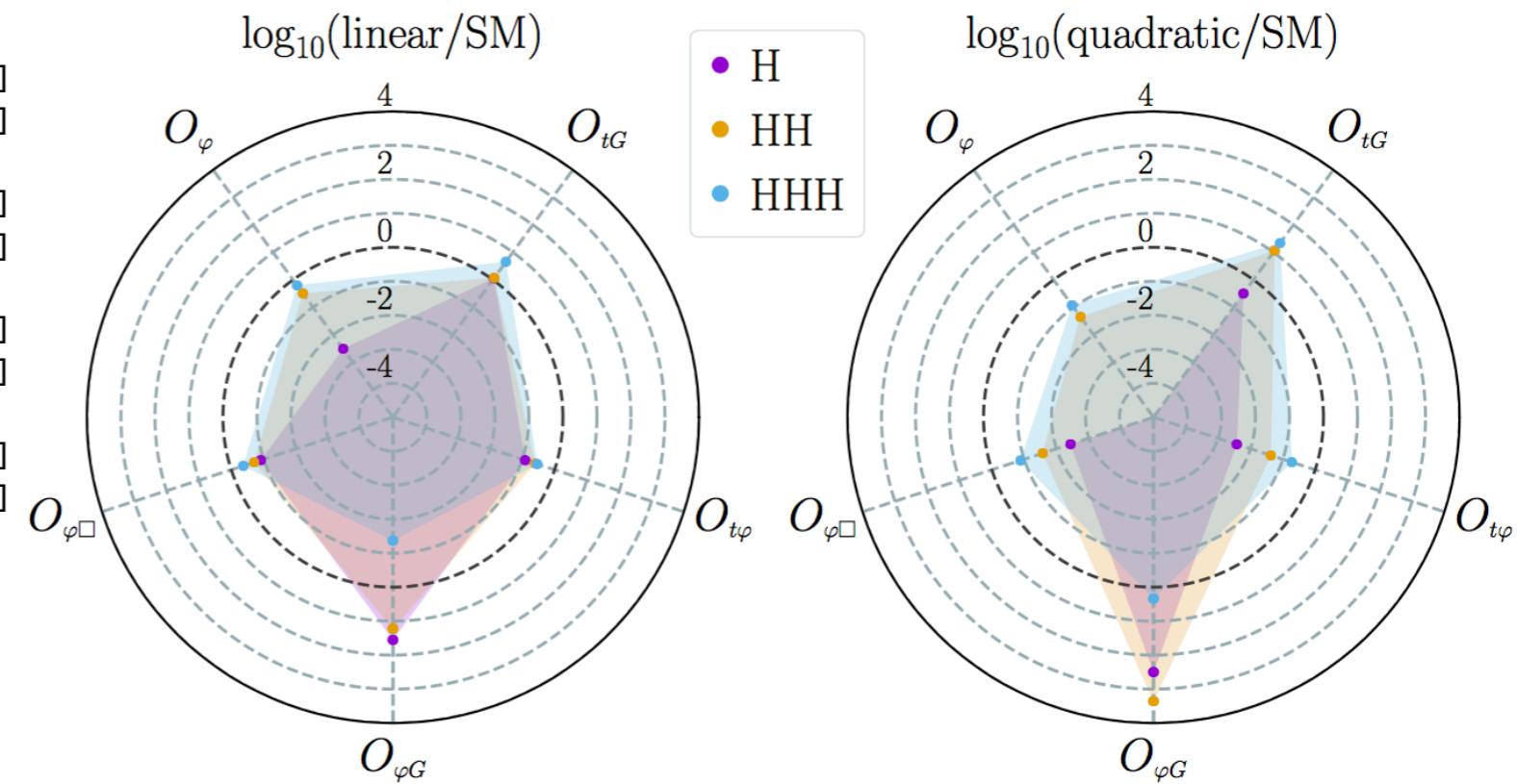
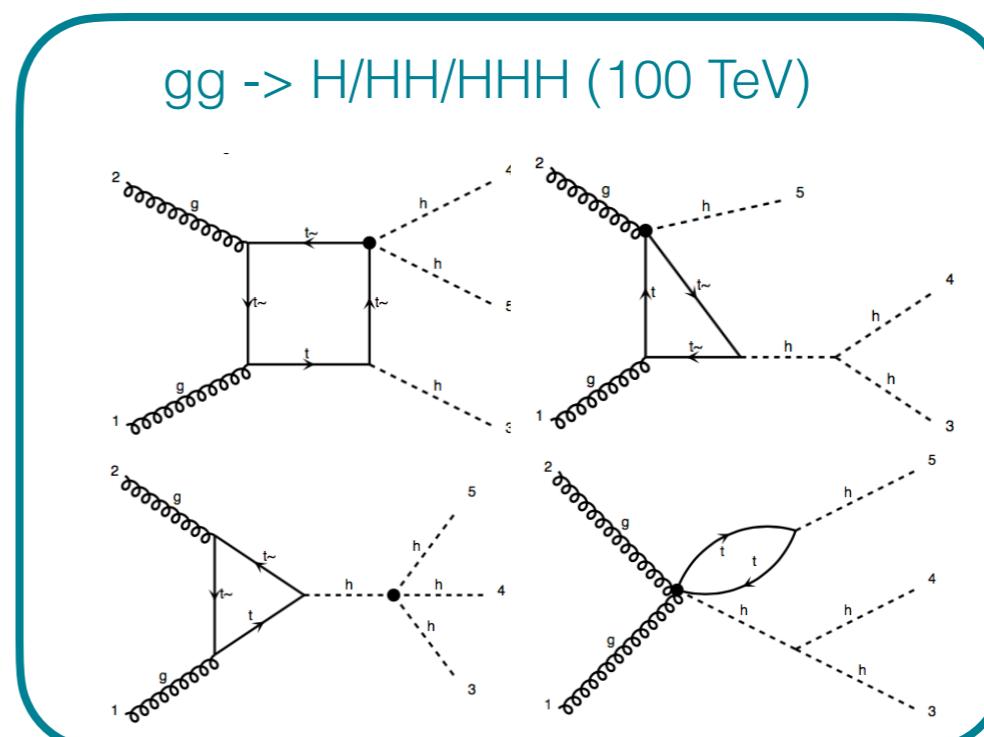
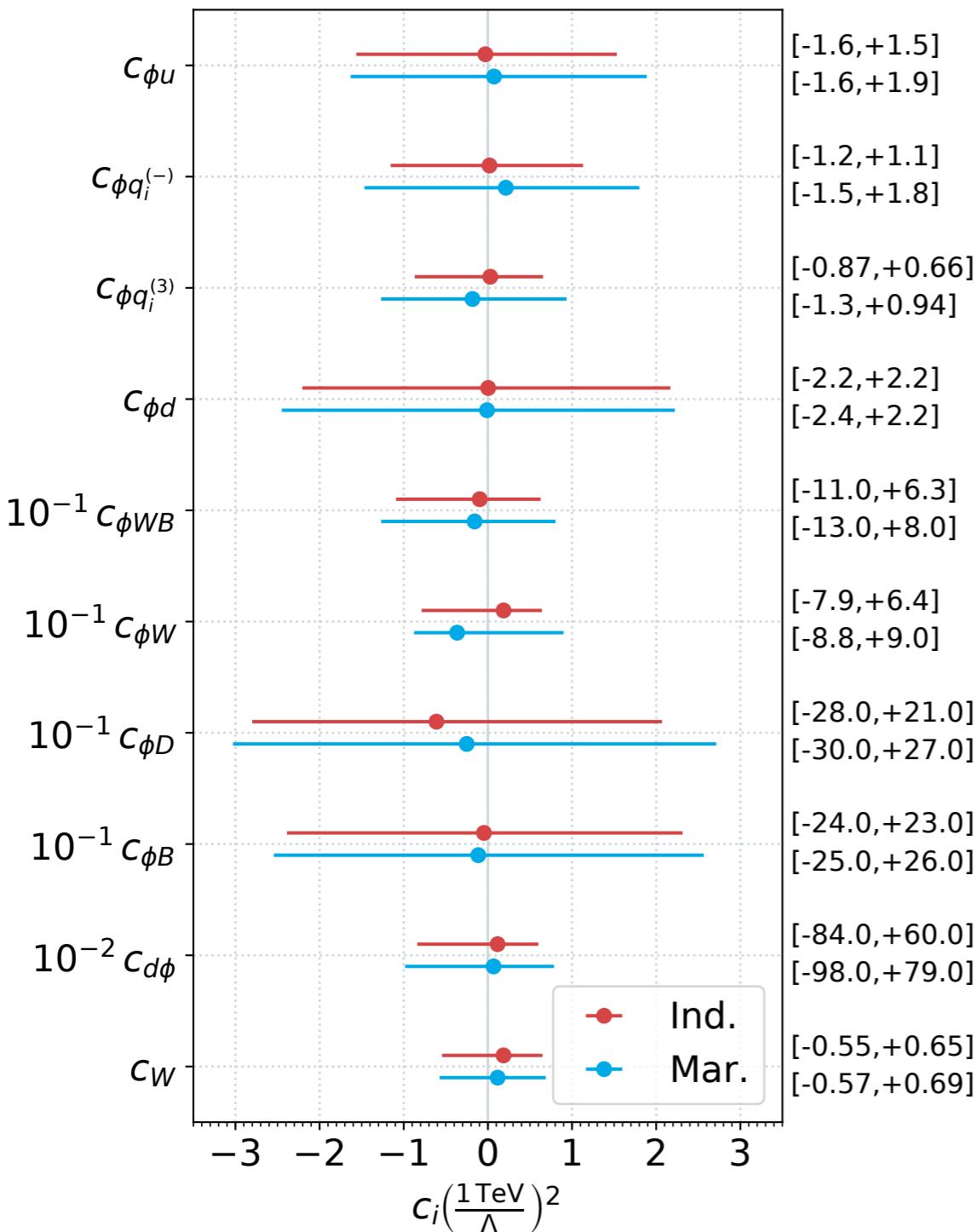


[CMS; PRL 125 (2020) 15, 151802]



Triboson sensitivity

CMS-SMP-19-014 combined sig. str.



- Next: combine with diboson/EWPO

Projected FCC-hh reach: 1%, 5% and 50% on H, HH and HHH

4F in top pair

LHC 13 TeV, SM = 744 pb, K-factor = 1.46, central scale choice = m_t

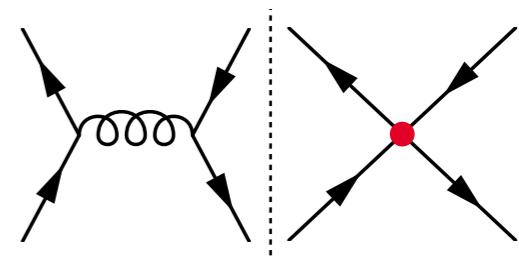
c_i	Interference		K	Square		K
	LO	$\mathcal{O}(\Lambda^{-2})$		LO	$\mathcal{O}(\Lambda^{-4})$	
c_{tu}^8	$4.27^{+11\%}_{-9\%}$	$4.06^{+1\%}_{-3\%}$	0.95	$1.04^{+6\%}_{-5\%}$	$1.03^{+2\%}_{-2\%}$	0.99
c_{td}^8	$2.79^{+11\%}_{-9\%}$	$2.77^{+1\%}_{-3\%}$	0.99	$0.577^{+6\%}_{-5\%}$	$0.611^{+3\%}_{-2\%}$	1.06
c_{tq}^8	$6.99^{+11\%}_{-9\%}$	$6.67^{+1\%}_{-3\%}$	0.95	$1.61^{+6\%}_{-5\%}$	$1.29^{+3\%}_{-2\%}$	0.80
c_{Qu}^8	$4.26^{+11\%}_{-9\%}$	$3.93^{+1\%}_{-4\%}$	0.92	$1.04^{+6\%}_{-5\%}$	$0.798^{+3\%}_{-3\%}$	0.77
c_{Qd}^8	$2.79^{+11\%}_{-9\%}$	$2.93^{+0\%}_{-1\%}$	1.05	$0.58^{+6\%}_{-5\%}$	$0.485^{+2\%}_{-2\%}$	0.84
$c_{Qq}^{8,1}$	$6.99^{+11\%}_{-9\%}$	$6.82^{+1\%}_{-3\%}$	0.98	$1.61^{+6\%}_{-5\%}$	$1.69^{+3\%}_{-3\%}$	1.05
$c_{Qq}^{8,3}$	$1.50^{+10\%}_{-9\%}$	$1.32^{+1\%}_{-3\%}$	0.88	$1.61^{+6\%}_{-5\%}$	$1.57^{+2\%}_{-2\%}$	0.98
c_{tu}^1	$[0.67^{+1\%}_{-1\%}]$	$-0.078(7)^{+31\%}_{-23\%}$	0.61	$4.66^{+6\%}_{-5\%}$	$5.92^{+6\%}_{-5\%}$	1.27
c_{td}^1	$[-0.21^{+1\%}_{-2\%}]$	$-0.306^{+30\%}_{-22\%}$	0.71	$2.62^{+6\%}_{-5\%}$	$3.46^{+5\%}_{-5\%}$	1.32
c_{tq}^1	$[0.39^{+0\%}_{-1\%}]$	$-0.47^{+24\%}_{-18\%}$	1.28	$7.25^{+6\%}_{-5\%}$	$9.36^{+6\%}_{-5\%}$	1.29
c_{Qu}^1	$[0.33^{+0\%}_{-0\%}]$	$-0.359^{+23\%}_{-17\%}$	1.72	$4.68^{+6\%}_{-5\%}$	$5.96^{+6\%}_{-5\%}$	1.27
c_{Qd}^1	$[-0.11^{+0\%}_{-1\%}]$	$0.023(6)^{+114\%}_{-75\%}$	1.72	$2.61^{+6\%}_{-5\%}$	$3.46^{+5\%}_{-5\%}$	1.31
$c_{Qq}^{1,1}$	$[0.57^{+0\%}_{-1\%}]$	$-0.24^{+30\%}_{-22\%}$	0.68	$7.25^{+6\%}_{-5\%}$	$9.34^{+5\%}_{-5\%}$	1.29
$c_{Qq}^{1,3}$	$[1.92^{+1\%}_{-1\%}]$	$0.088(7)^{+28\%}_{-20\%}$	0.55	$7.25^{+6\%}_{-5\%}$	$9.32^{+5\%}_{-5\%}$	1.29

NLO can break degeneracies in fits

- C's enter e.g., $m_{t\bar{t}}$, in fixed combinations at LO

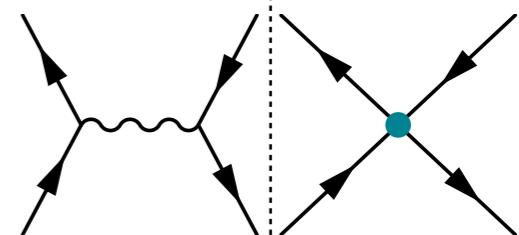
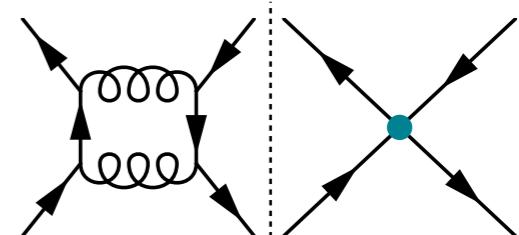
color-octet qqt:

- dominant operators in $t\bar{t}\text{bar}$
- Non SM-like corrections



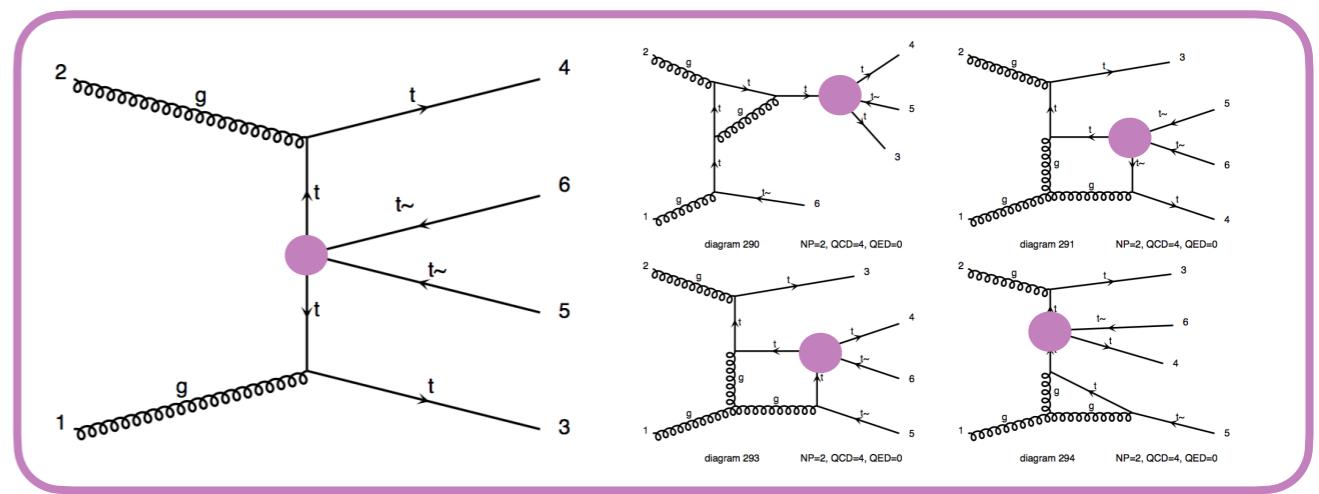
color-singlet qqt:

- int. with QCD $t\bar{t}\text{bar}$ at NLO
- [x] int. with EW $t\bar{t}\text{bar}$
- No error control at LO



4F in 4 top

$$\sigma(pp \rightarrow t\bar{t}t\bar{t}) [\text{fb}], c_i/\Lambda^2 = 1 \text{ TeV}^{-2}$$



! different from arXiv version !

c_i	<i>Interference</i> LO	$\mathcal{O}(\Lambda^{-2})$			<i>Square</i> LO	$\mathcal{O}(\Lambda^{-4})$		
		NLO	K	NLO		K		
c_{QQ}^8	$0.081^{+55\%}_{-33\%}$	$[-0.277]$	$0.090^{+4\%}_{-11\%}$	1.1	$0.115^{+46\%}_{-29\%}$	$0.158^{+4\%}_{-11\%}$	1.37	
c_{Qt}^8	$0.274^{+54\%}_{-33\%}$	$[-0.365]$	$0.311^{+5\%}_{-10\%}$	1.14	$0.342^{+46\%}_{-29\%}$	$0.378^{+4\%}_{-13\%}$	1.10	
c_{QQ}^1	$0.242^{+55\%}_{-33\%}$	$[-0.826]$	$0.24(3)^{+3\%}_{-18\%}$	0.99	$1.039^{+47\%}_{-29\%}$	$1.41^{+4\%}_{-11\%}$	1.36	
c_{Qt}^1	$-0.0098(10)^{+38\%}_{-33\%}$	$[0.852]$	$-0.019(9)^{+63\%}_{-27\%}$	1.9	$1.406^{+46\%}_{-30\%}$	$1.86^{+4\%}_{-10\%}$	1.32	
c_{tt}^1	$0.483^{+55\%}_{-33\%}$	$[-1.38]$	$0.53(8)^{+3\%}_{-10\%}$	1.10	$4.154^{+47\%}_{-29\%}$	$5.61^{+4\%}_{-11\%}$	1.35	

QCD corrections to inclusive 4 top production in SMEFT

- Central scale choice: $\mu = 2m_t$

$$\text{SM} = 11.1^{+25\%}_{-25\%} \text{ fb } (K = 1.83)$$

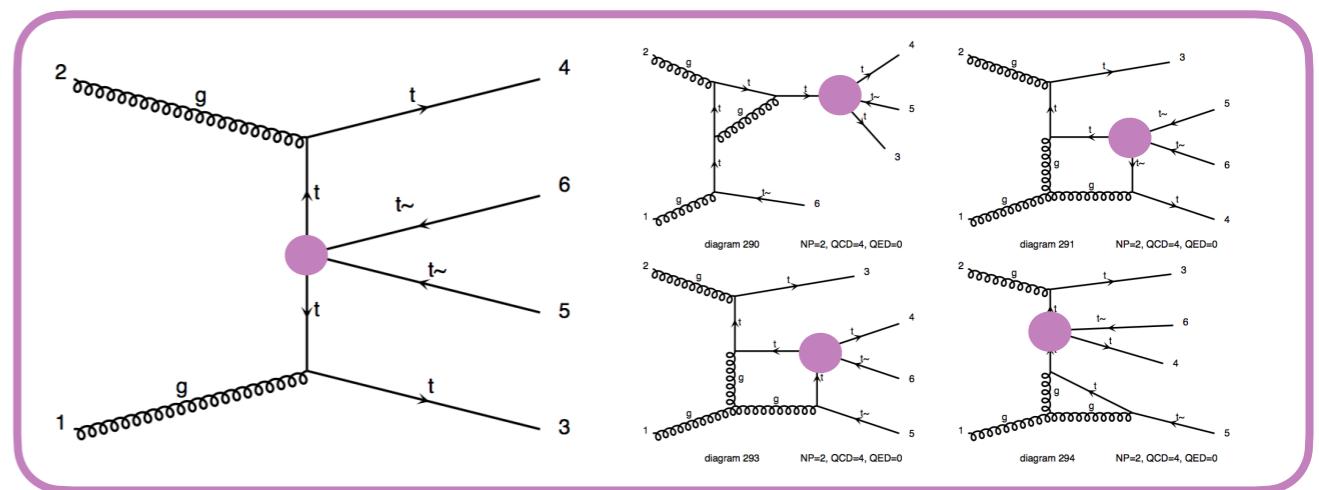
Computationally challenging

- ~1 week per operator run on CP3 computing cluster

4F in 4 top

[Degrande, Durieux, Maltoni, KM,
Vryonidou & Zhang; arXiv:2008.11743]

$$\sigma(pp \rightarrow t\bar{t}t\bar{t}) [\text{fb}], c_i/\Lambda^2 = 1 \text{ TeV}^{-2}$$



! different from arXiv version !

c_i	<i>Interference</i> LO	$\mathcal{O}(\Lambda^{-2})$			<i>Square</i> LO	$\mathcal{O}(\Lambda^{-4})$		
		NLO	K	NLO		K	NLO	K
c_{QQ}^8	$0.081^{+55\%}_{-33\%}$	$[-0.277]$	$0.090^{+4\%}_{-11\%}$	1.1	$0.115^{+46\%}_{-29\%}$	$0.158^{+4\%}_{-11\%}$	1.37	
c_{Qt}^8	$0.274^{+54\%}_{-33\%}$	$[-0.365]$	$0.311^{+5\%}_{-10\%}$	1.14	$0.342^{+46\%}_{-29\%}$	$0.378^{+4\%}_{-13\%}$	1.10	
c_{QQ}^1	$0.242^{+55\%}_{-33\%}$	$[-0.826]$	$0.24(3)^{+3\%}_{-18\%}$	0.99	$1.039^{+47\%}_{-29\%}$	$1.41^{+4\%}_{-11\%}$	1.36	
c_{Qt}^1	$-0.0098(10)^{+38\%}_{-33\%}$	$[0.852]$	$-0.019(9)^{+63\%}_{-27\%}$	1.9	$1.406^{+46\%}_{-30\%}$	$1.86^{+4\%}_{-10\%}$	1.32	
c_{tt}^1	$0.483^{+55\%}_{-33\%}$	$[-1.38]$	$0.53(8)^{+3\%}_{-10\%}$	1.10	$4.154^{+47\%}_{-29\%}$	$5.61^{+4\%}_{-11\%}$	1.35	

Reduction of scale uncertainty, relatively lower than SM

K-factors lower than SM

$$\text{SM} = 11.1^{+25\%}_{-25\%} \text{ fb } (K = 1.83)$$

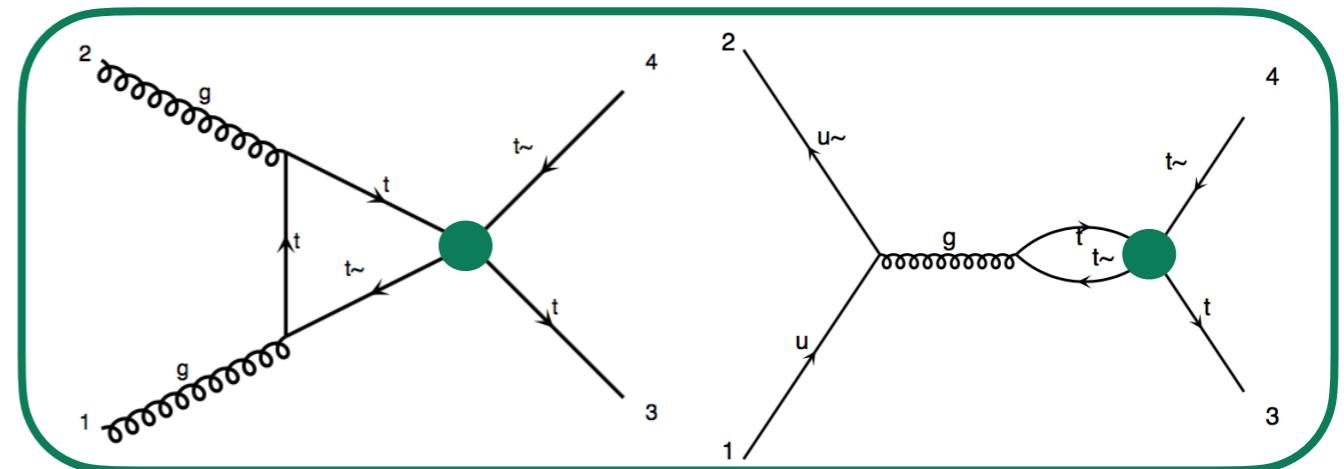
- Relative impact slightly decreases from LO to NLO
- Square typically receives larger corrections

$$K_{QQ}^1 \neq K_{QQ}^8$$

Indirect sensitivity from $t\bar{t}$

Loop-induced effects from 4 top operators in $t\bar{t}$

- $q\bar{q} \rightarrow t\bar{t}$: mixing with $q\bar{q}t\bar{t}$ ops.
 $(\bar{t}\gamma^\mu t)(\bar{t}\gamma_\mu t) \rightarrow (\bar{t}\gamma^\mu T_A D^\nu t) G_{\mu\nu}^A$
- $gg \rightarrow t\bar{t}$: finite contribution
- $b\bar{b} \rightarrow t\bar{t}$: small piece from Q



$gg \rightarrow t\bar{t}$ amplitude: Helicity structure doesn't match SM

- No interference in the massless limit
- Form-factor doesn't grow with energy like $q\bar{q}t\bar{t}$ contact interactions
- Main effects near $t\bar{t}$ threshold

[Craig et al.; JHEP 08 (2020) 086]

Indirect sensitivity from $t\bar{t}$

$\sigma(pp \rightarrow t\bar{t}) [\text{pb}], c_i/\Lambda^2 = 1 \text{ TeV}^{-2}$

Results

- Octet $q\bar{q}t\bar{t}$ for reference
- [EW interference]
- 1-2 orders of magnitude smaller
- Competition/cancellation between gg and qq channels
- Λ^{-4} automatically (loop) suppressed

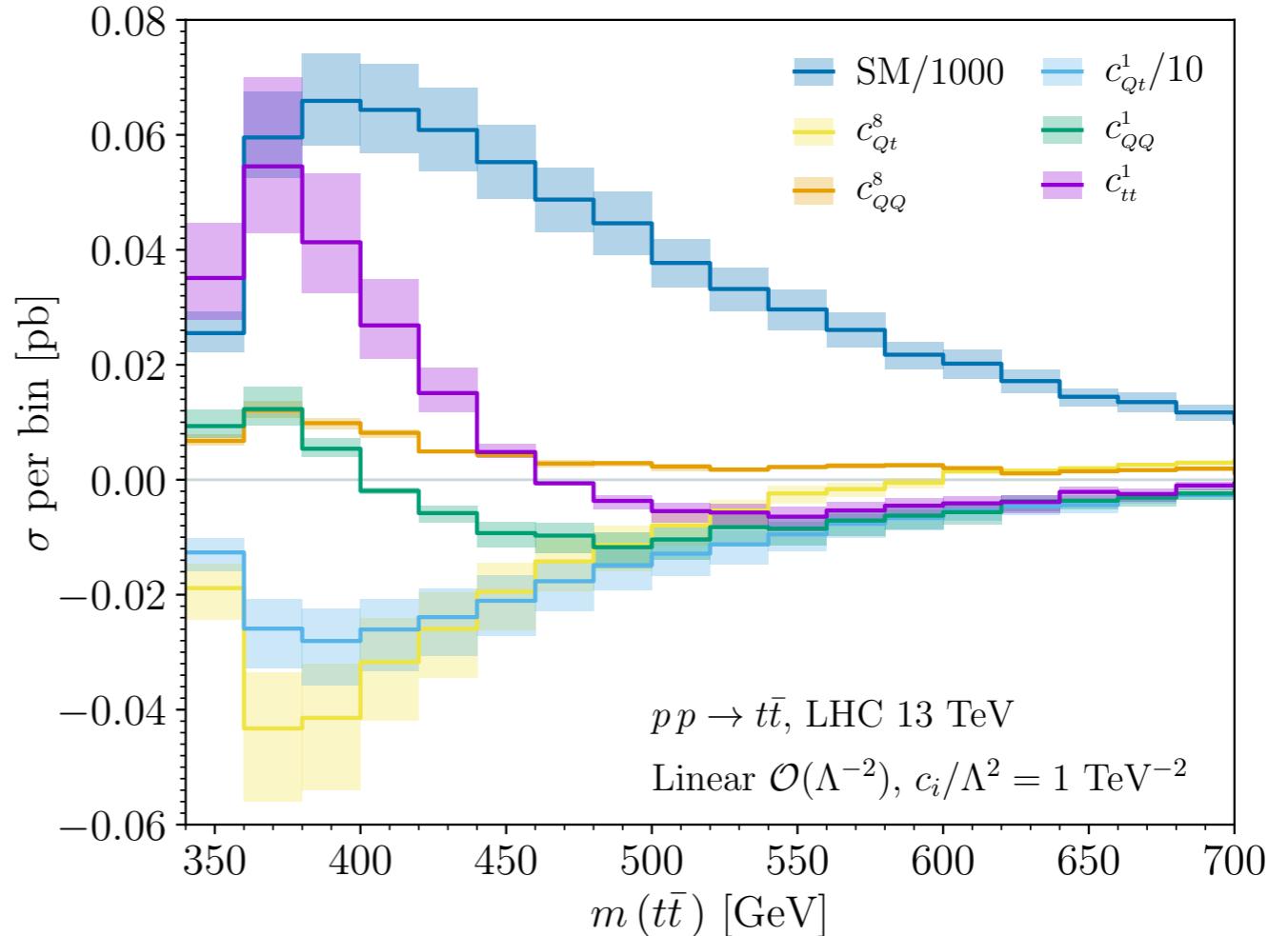
c_i	$\mathcal{O}(\Lambda^{-2})$		$\mathcal{O}(\Lambda^{-4})$	
	LO	NLO	LO	NLO
c_{tu}^8	$4.27^{+11\%}_{-9\%}$	$4.06^{+1\%}_{-3\%}$	$1.04^{+6\%}_{-5\%}$	$1.03^{+2\%}_{-2\%}$
c_{td}^8	$2.79^{+11\%}_{-9\%}$	$2.77^{+1\%}_{-3\%}$	$0.577^{+6\%}_{-5\%}$	$0.611^{+3\%}_{-2\%}$
c_{tq}^8	$6.99^{+11\%}_{-9\%}$	$6.67^{+1\%}_{-3\%}$	$1.61^{+6\%}_{-5\%}$	$1.29^{+3\%}_{-2\%}$
$c_{Q_u}^8$	$4.26^{+11\%}_{-9\%}$	$3.93^{+1\%}_{-4\%}$	$1.04^{+6\%}_{-5\%}$	$0.798^{+3\%}_{-3\%}$
$c_{Q_d}^8$	$2.79^{+11\%}_{-9\%}$	$2.93^{+0\%}_{-1\%}$	$0.58^{+6\%}_{-5\%}$	$0.485^{+2\%}_{-2\%}$
$c_{Qq}^{8,1}$	$6.99^{+11\%}_{-9\%}$	$6.82^{+1\%}_{-3\%}$	$1.61^{+6\%}_{-5\%}$	$1.69^{+3\%}_{-3\%}$
$c_{Qq}^{8,3}$	$1.50^{+10\%}_{-9\%}$	$1.32^{+1\%}_{-3\%}$	$1.61^{+6\%}_{-5\%}$	$1.57^{+2\%}_{-2\%}$
c_{QQ}^8	$0.0586^{+27\%}_{-25\%}$	$0.125^{+10\%}_{-11\%}$	$0.00628^{+13\%}_{-16\%}$	$0.0133^{+7\%}_{-5\%}$
c_{Qt}^8	$0.0583^{+27\%}_{-25\%}$	$-0.107(6)^{+40\%}_{-33\%}$	$0.00619^{+13\%}_{-16\%}$	$0.0118^{+8\%}_{-5\%}$
c_{QQ}^1	$[-0.11^{+15\%}_{-18\%}]$	$-0.039(4)^{+51\%}_{-33\%}$	$[-0.12^{+7\%}_{-5\%}]$	$0.0282^{+13\%}_{-16\%}$
c_{Qt}^1	$[-0.068^{+16\%}_{-18\%}]$	$-2.51^{+29\%}_{-21\%}$	$[-0.12^{+3\%}_{-6\%}]$	$0.0283^{+13\%}_{-16\%}$
c_{tt}^1	x		$0.215^{+23\%}_{-18\%}$	x

- One intriguing number from c_{Qt}^1 , similar in size to $q\bar{q}t\bar{t}$ octets! $\sigma_{\text{int}} \text{ suppressed in } 4t$
- ~ Few percent effect near $t\bar{t}$ threshold assuming current bound ~ 3.5

q^2 dependence

Lack of energy growth

- Sign changes over phase space lead to suppressions
- Quark and gluon channels often have opposite sign
- Optimistic: need few percent precision near threshold



Completely different dependence to $t\bar{t}t\bar{t}$

- Limited prospects but may be at least useful for breaking $t\bar{t}t\bar{t}$ degeneracies
- Further study required

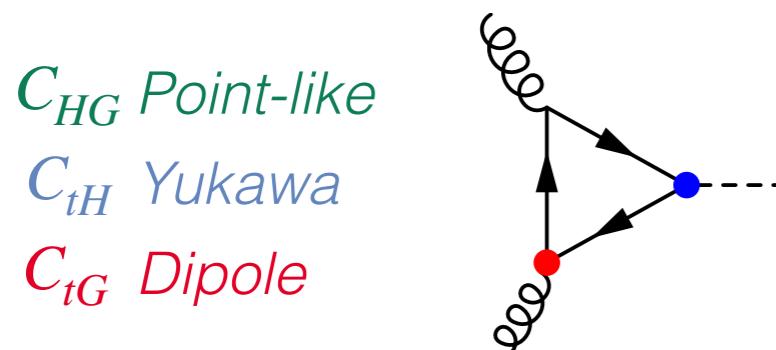
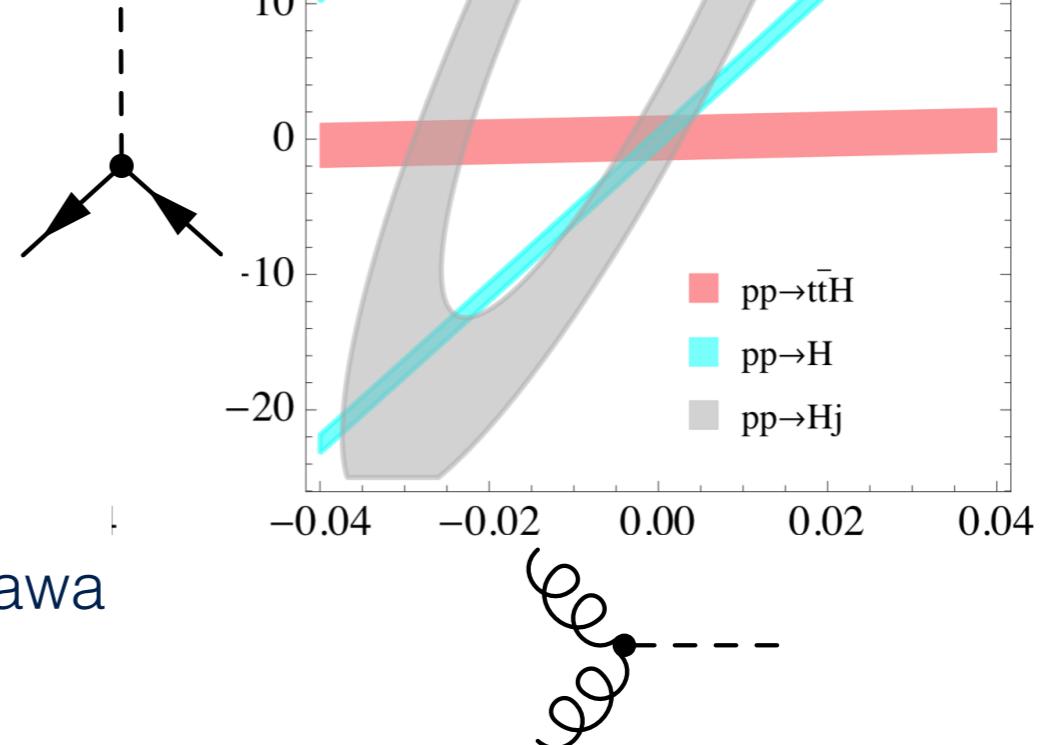
Top/Higgs interplay

Inextricably linked in the SM

- Yukawa interaction controls ggF
- Strong BSM motivation to study tops

ggF is well measured now

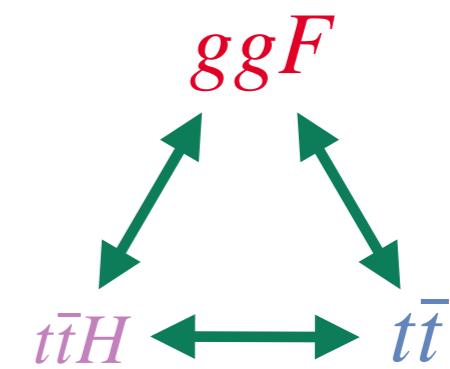
- Cannot exclude top partners/anomalous Yukawa



Need more data to break degeneracy

- $t\bar{t}H$ production for direct Yukawa measurement
- $t\bar{t}$ data to constrain dipole

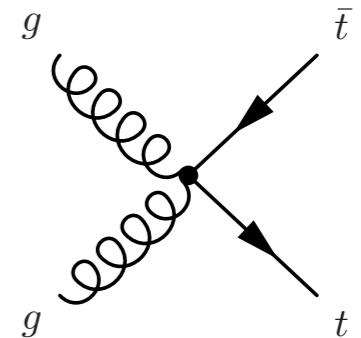
Blind direction in BSM scenarios
Effective coupling degeneracy



The role of top data

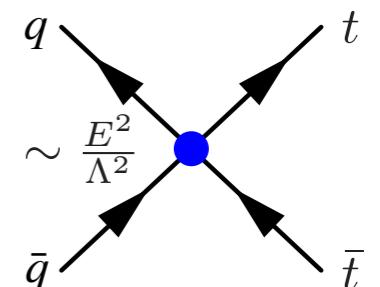
$t\bar{t}$ cross section measurements constrain C_{tG}

- Indirectly improve bounds on C_{HG} and C_{tH}



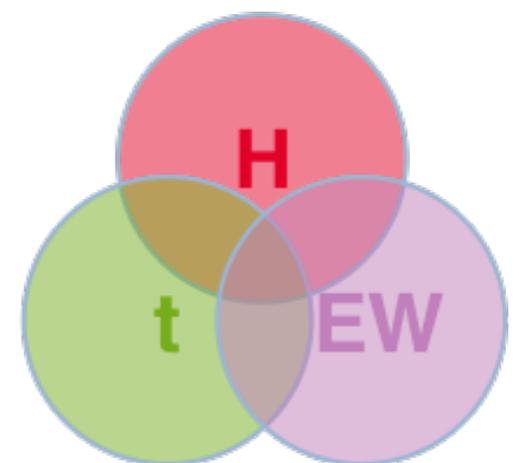
Several other new interactions can affect $t\bar{t}$

- Notably $q\bar{q}t\bar{t}$ operators, of which there are many (14)
- To what extent do these limit ultimate NP sensitivity in top/Higgs sector?



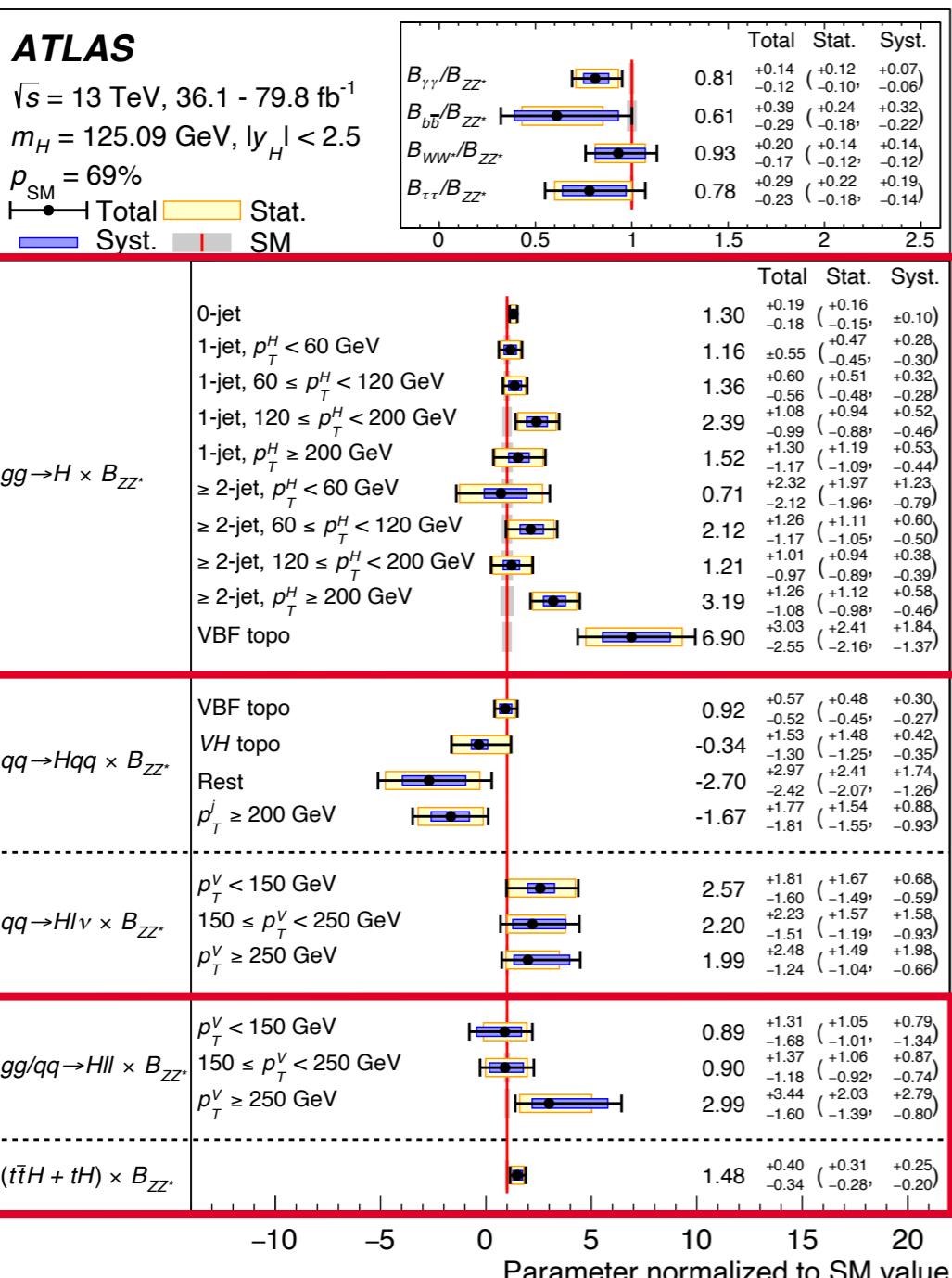
Can only be addressed in combined fit

- Beyond tree-level (at least for ggF)
- Identify other cross-talk (non-trivial correlations)
- Broaden range of applicability to UV models



STXS

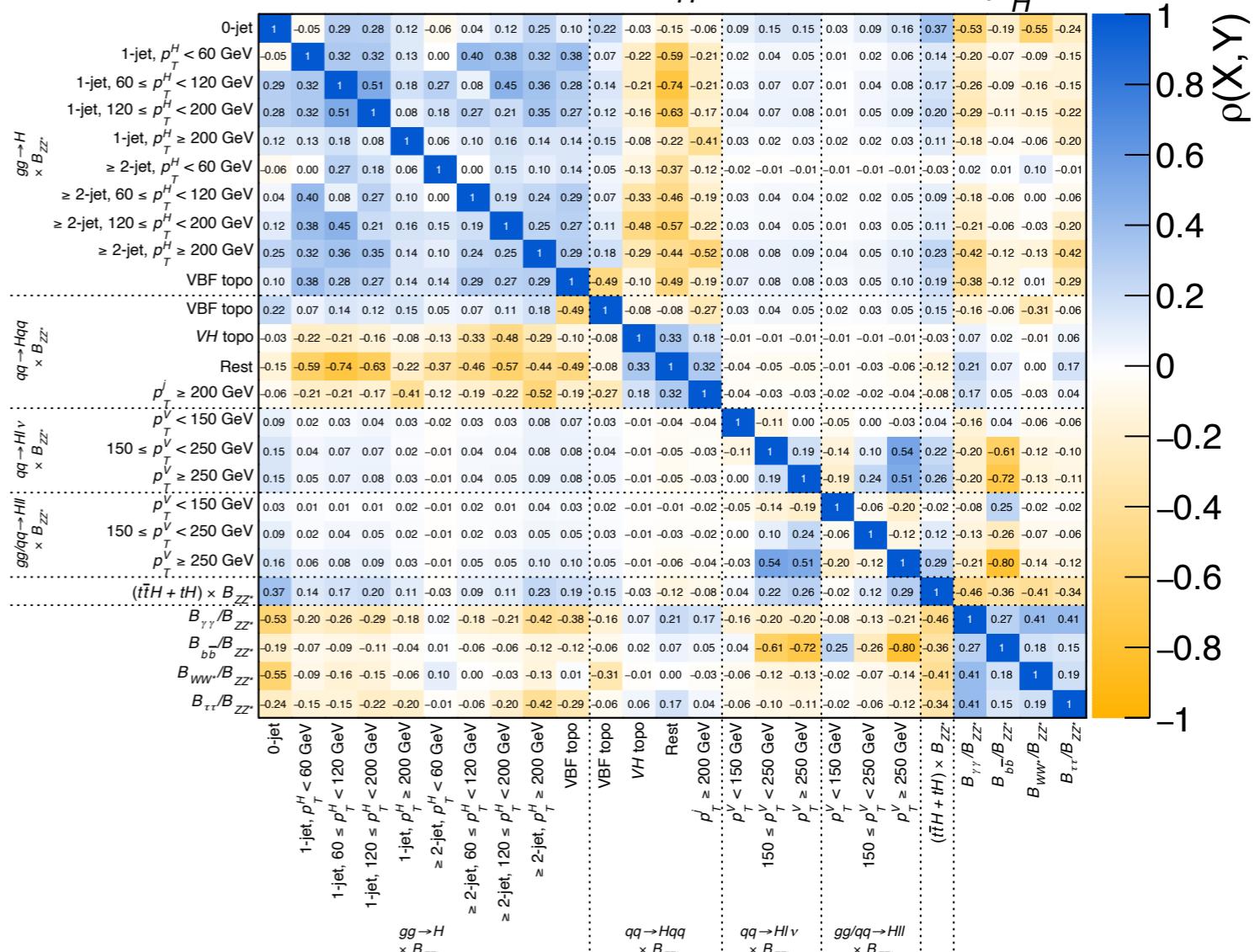
(Simplified Template Cross Sections)



[ATLAS; PRD 101 (2020) 012002]

ATLAS

$\sqrt{s} = 13 \text{ TeV}, 36.1 - 79.8 \text{ fb}^{-1}$
 $m_H = 125.09 \text{ GeV}, |\eta_H| < 2.5$



See also:

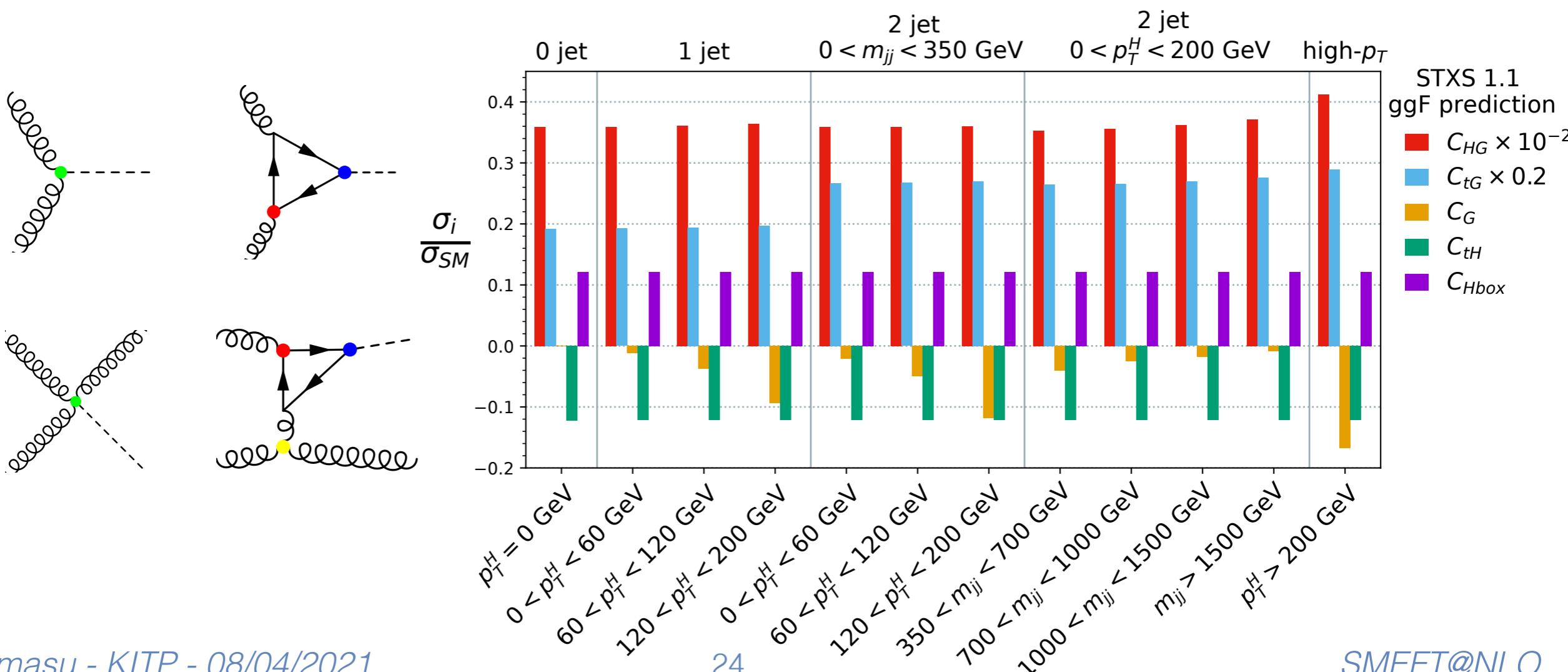
- [CMS-HIG-19-015]
- [CMS-PAS-HIG-19-010]
- [CMS-HIG-19-001]

- [ATLAS-CONF-2020-053]
- [ATLAS; EPJC 80 (2020) 10]
- [ATLAS-CONF-2020-026]
- [ATLAS; EPJC 81 (2021) 178]

Improving fits

STXS \leftrightarrow gluon fusion in the SMEFT

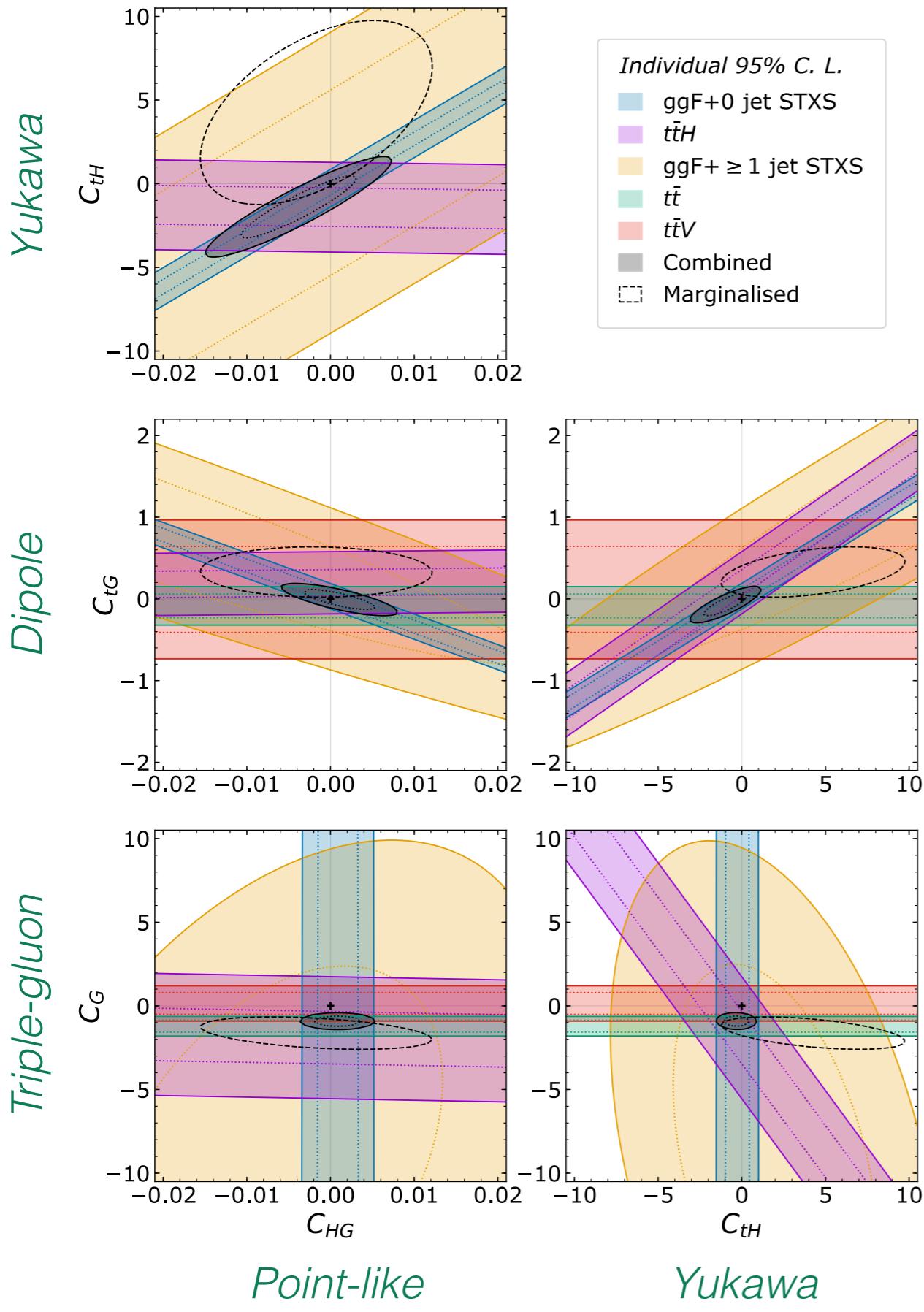
- LO in the SM is one-loop
- Tree-EFT x loop-SM + loop-EFT x loop-SM interference terms
- Heavy top limit is OK for 0-jet, breaks down at high- p_T



Top-Higgs interplay

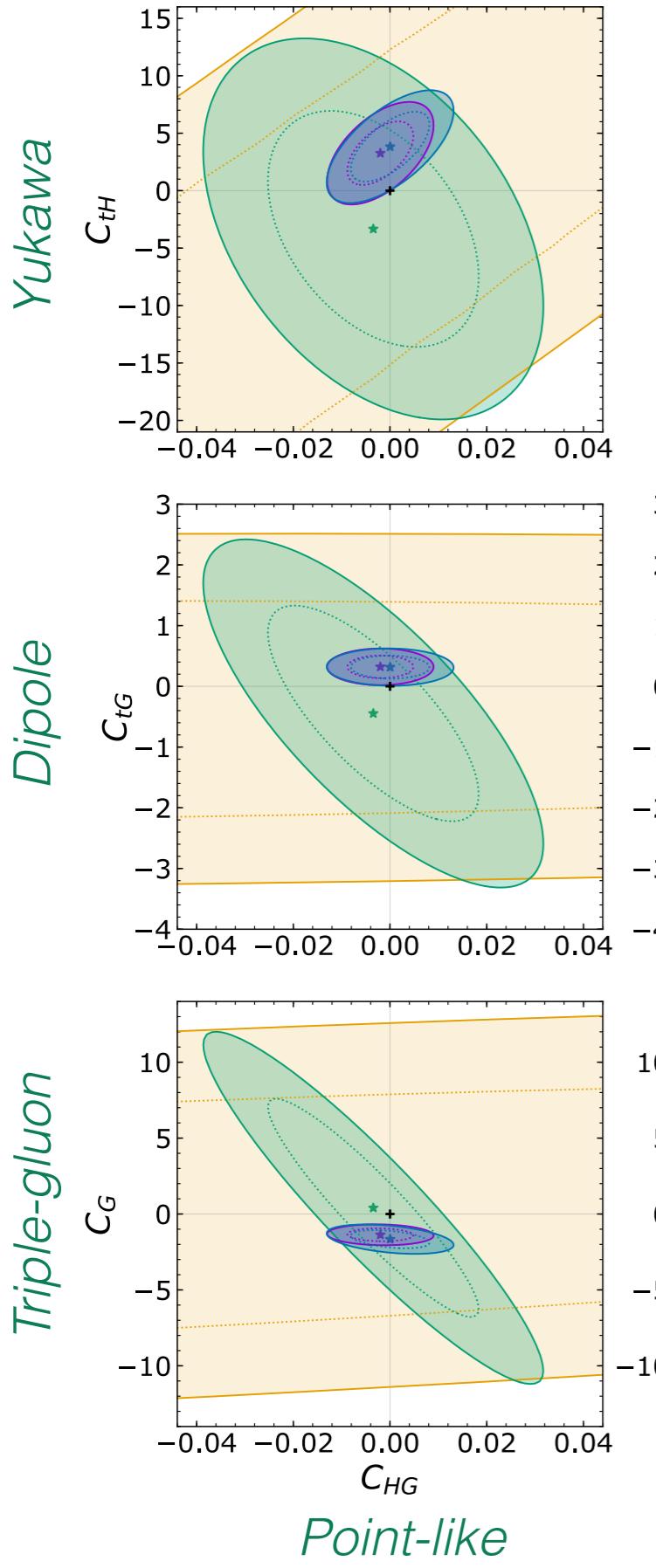
2D individual constraints

- All others set to 0
- $ggF/t\bar{t}H$ complementarity for (C_{HG}, C_{tH})
- H+jets STXS & $t\bar{t}V$ not yet competitive
- Strong impact of $t\bar{t}$ evident for (C_{tG}, C_G)
- Tension with SM $> 2\sigma$
- Significant correlations remain
- Large marginalisation effects (including 4F)



What is the concrete impact of 4F?

4F impact



Marginalised

- Marginalised 95% C. L.
- Higgs data (no $t\bar{t}H$)
- Higgs data
- Higgs & Top data
- Higgs & Top data (+4F)
- SM

Fit to ‘Higgs-only’ subspace

$$C_{H\square}, C_{HG}, C_{HW}, C_{HB}, C_{tH}, C_{bH}, C_{\tau H}, C_{\mu H} \\ + C_{tG} \& C_G$$

- Allow a closed fit to Higgs data only
- Emphasises impact of $t\bar{t}H$ & $t\bar{t}$

Now add in $t\bar{t}$ 4F operators

$$+ C_{Qq}^{3,8}, C_{Qq}^{1,8}, C_{Qu}^8, C_{Qd}^8, C_{tq}^8, C_{tu}^8, C_{td}^8$$

- Relatively mild impact
- Preferred $t\bar{t}$ phase space is different

C_{tG} : low $m_{t\bar{t}}$

4F : high $m_{t\bar{t}}$

- Able to constrain them independently

Conclusions & future plans

SMEFT@NLO is a milestone in tools for SMEFT predictions

- Automated, fully-differential computations up to one-loop
- NLO+PS, loop-induced, tree-loop interference
- Crucial for inputs to global SMEFT likelihood for LHC & beyond

Planned extensions

- Generalise flavor structure: $U(2)^5$ (b chirality flipping operators)
- 4 light fermion operators (qqqq & qql \bar{l})
- CP violation
- Open to suggestions/requests!

Work in progress for running of Wilson coefficients in MG5

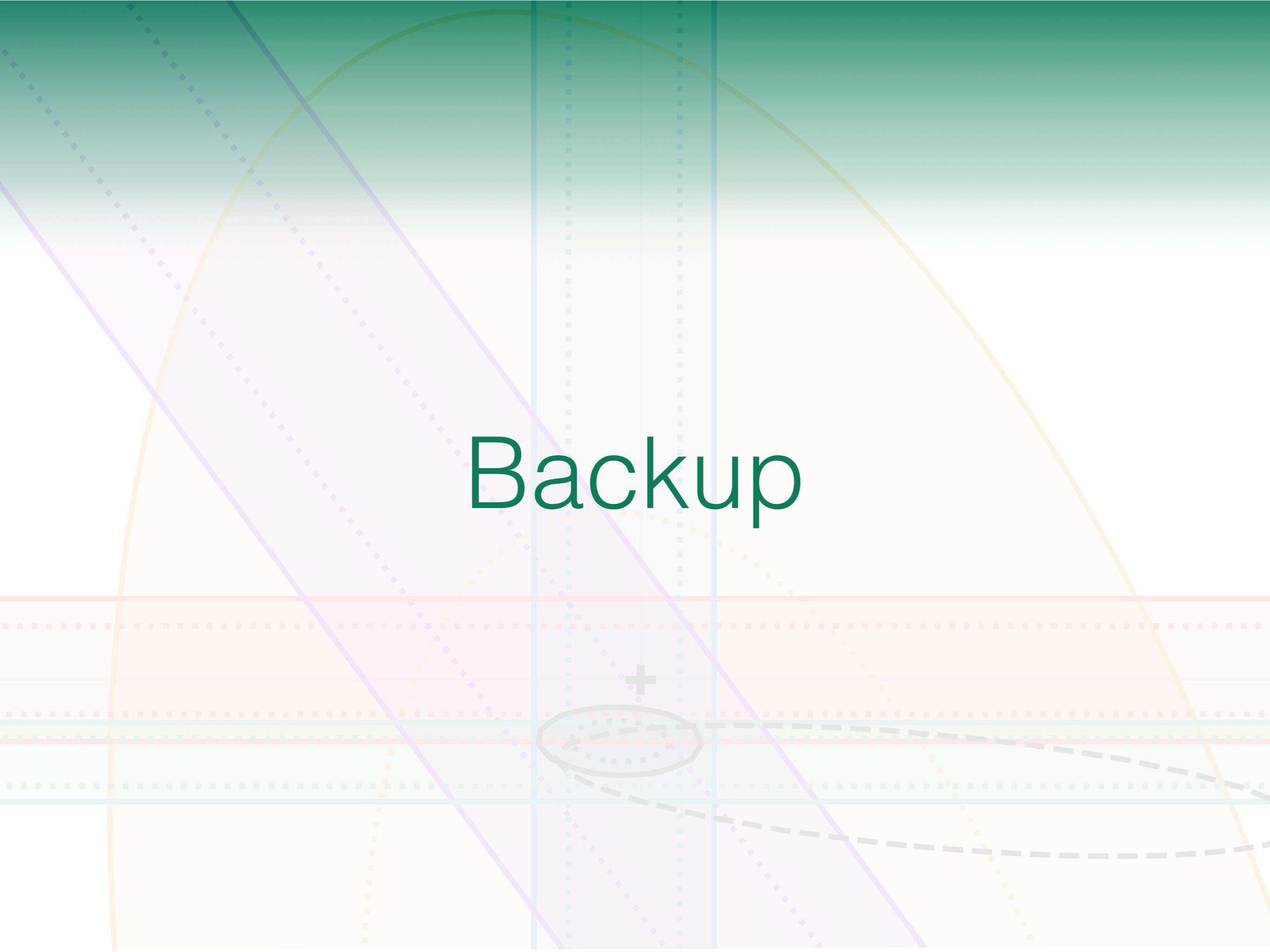
Long term: EW loops, already possible for the SM in MG5

Next week: HEFT 2021

USTC Hefei, China

<https://indico.ihep.ac.cn/event/13632/>

Backup



Technical details

Lepton sector: $[U(1)_L \times U(1)_e]^3$, flavor diagonal (e, μ, τ)

5-flavor scheme (massless b) & CKM=1

EW input scheme: $\{G_F, m_W, m_Z\}$

- Relevant field redefinitions & EW parameter shifts performed

EFT (\overline{MS}) renormalisation scale: `mueft`

- Separate, fixed renormalisation scale for Wilson coefficients
- MG5 does not run the Wilson coefficients (yet)
- Usual `muR` & `muF` are kept for α_S & PDFs

Validated at LO against existing implementations

- `dim6top` & `SMEFTsim`

[Aguilar-Saavedra et al.; arXiv:1802.07237]

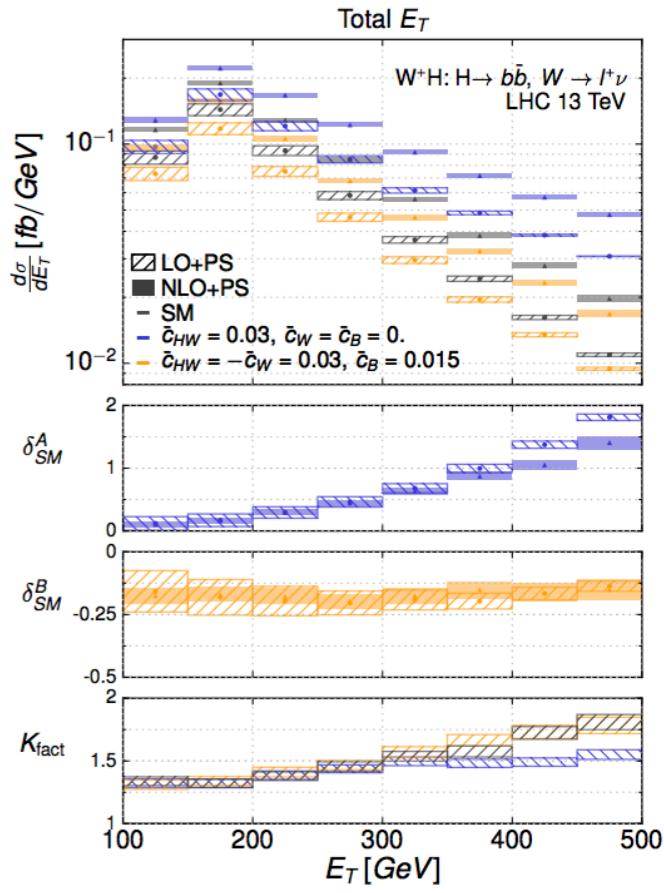
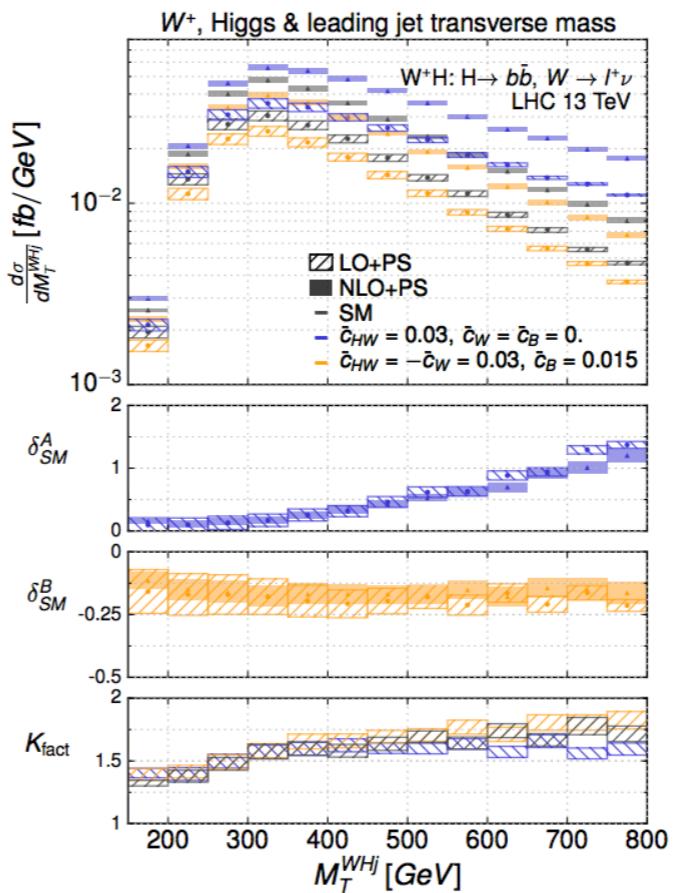
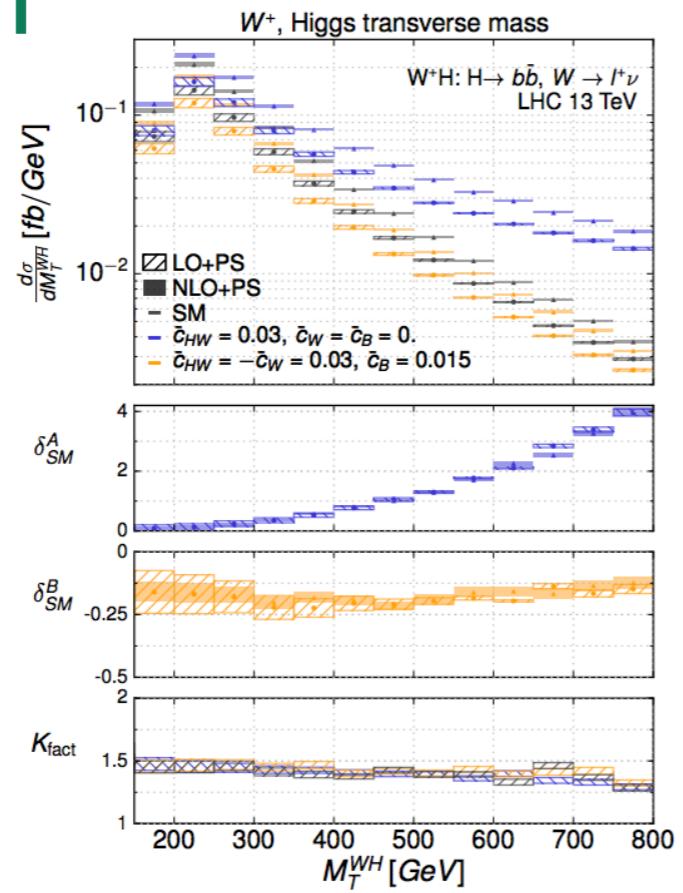
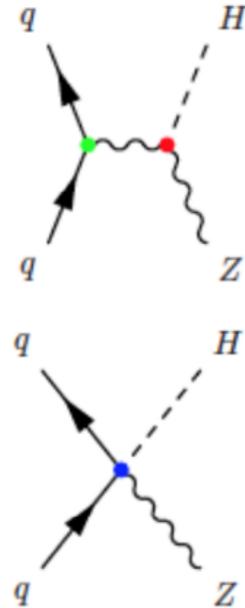
[Brivio Jiang & Trott; JHEP 12 (2017) 070]

[Brivio; arXiv:2012.11343]

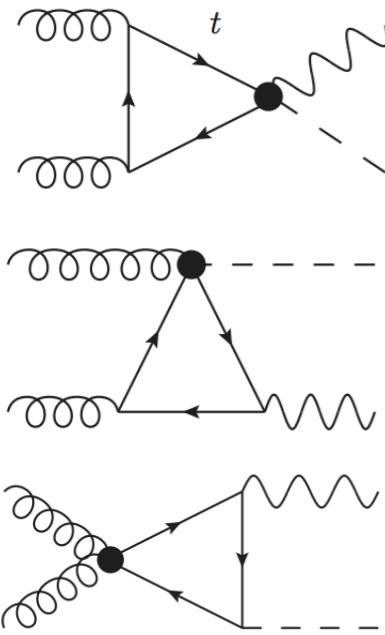
$pp \rightarrow ZH$

[Degrande, et al.; EPJC 77 (2017) 4, 262]

Quark-initiated

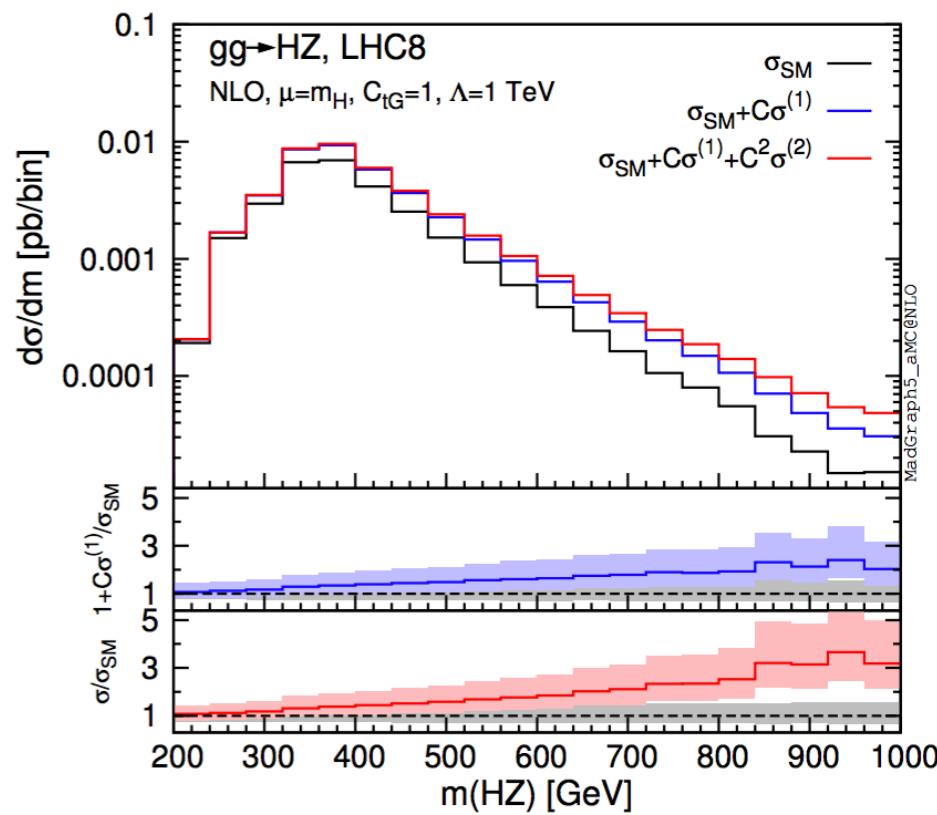


gg, loop-induced



[Bylund et al.; JHEP 1605 (2016) 052]

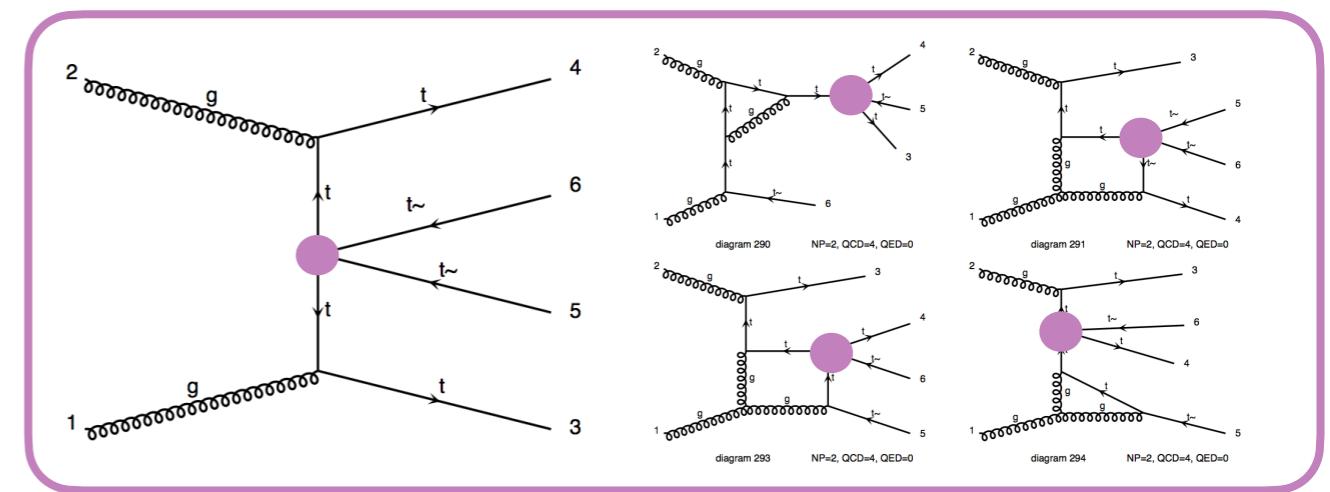
$[\text{fb}]$	SM	\mathcal{O}_{tG}	$\mathcal{O}_{\phi Q}^{(1)}$
8TeV	$29.15^{+40.0\%}_{-26.6\%}$	$\sigma_i^{(1)}$	$10.37^{+41.3\%}_{-27.2\%}$
		$\sigma_i^{(2)}$	$1.621^{+45.1\%}_{-28.7\%}$
		$\sigma_i^{(1)}/\sigma_{SM}$	$0.356^{+0.9\%}_{-0.8\%}$
		$\sigma_i^{(2)}/\sigma_i^{(1)}$	$0.156^{+2.6\%}_{-2.0\%}$
13TeV	$93.6^{+34.3\%}_{-23.8\%}$	$\sigma_i^{(1)}$	$34.6^{+35.2\%}_{-24.5\%}$
		$\sigma_i^{(2)}$	$6.09^{+39.2\%}_{-26.1\%}$
		$\sigma_i^{(1)}/\sigma_{SM}$	$0.370^{+0.7\%}_{-0.9\%}$
		$\sigma_i^{(2)}/\sigma_i^{(1)}$	$0.176^{+2.9\%}_{-2.1\%}$



4F in 4 top

[Degrande, Durieux, Maltoni, KM,
Vryonidou & Zhang; arXiv:2008.11743]

$$\sigma(pp \rightarrow t\bar{t}t\bar{t}) [\text{fb}], c_i/\Lambda^2 = 1 \text{ TeV}^{-2}$$



! different from arXiv version !

c_i	Interference $\mathcal{O}(\Lambda^{-2})$			K	Square $\mathcal{O}(\Lambda^{-4})$		
	LO	NLO	K		LO	NLO	K
c_{QQ}^8	$0.081^{+55\%}_{-33\%}$	[-0.277]	$0.090^{+4\%}_{-11\%}$	1.1	$0.115^{+46\%}_{-29\%}$	$0.158^{+4\%}_{-11\%}$	1.37
c_{Qt}^8	$0.274^{+54\%}_{-33\%}$	[-0.365]	$0.311^{+5\%}_{-10\%}$	1.14	$0.342^{+46\%}_{-29\%}$	$0.378^{+4\%}_{-13\%}$	1.10
c_{QQ}^1	$0.242^{+55\%}_{-33\%}$	[-0.826]	$0.24(3)^{+3\%}_{-18\%}$	0.99	$1.039^{+47\%}_{-29\%}$	$1.41^{+4\%}_{-11\%}$	1.36
c_{Qt}^1	$-0.0098(10)^{+38\%}_{-33\%}$	[0.852]	$-0.019(9)^{+63\%}_{-27\%}$	1.9	$1.406^{+46\%}_{-30\%}$	$1.86^{+4\%}_{-10\%}$	1.32
c_{tt}^1	$0.483^{+55\%}_{-33\%}$	[-1.38]	$0.53(8)^{+3\%}_{-10\%}$	1.10	$4.154^{+47\%}_{-29\%}$	$5.61^{+4\%}_{-11\%}$	1.35

Current limits $\sim \mathcal{O}(\text{few}) \text{ TeV}^{-2}$: square > interference

c_i & normalisation independent measure: $\xi_i \equiv \frac{1}{2} |\sigma_{\text{int.}}^i| / (\sigma_{\text{SM}} \sigma_{\text{sq.}}^i)^{\frac{1}{2}}$

	C_{4t}	QQ^8	Qt^8	QQ^1	Qt^1	tt
ξ_{4t}	LO	0.048	0.095	0.048	0.002	0.048
	NLO	0.034	0.075	0.03	-	0.034

	C_{2t}	$Qq^{8,1}$	tq^8	tu^8
$\xi_{t\bar{t}}$	LO	0.12	0.12	0.092
	NLO	0.096	0.11	0.073

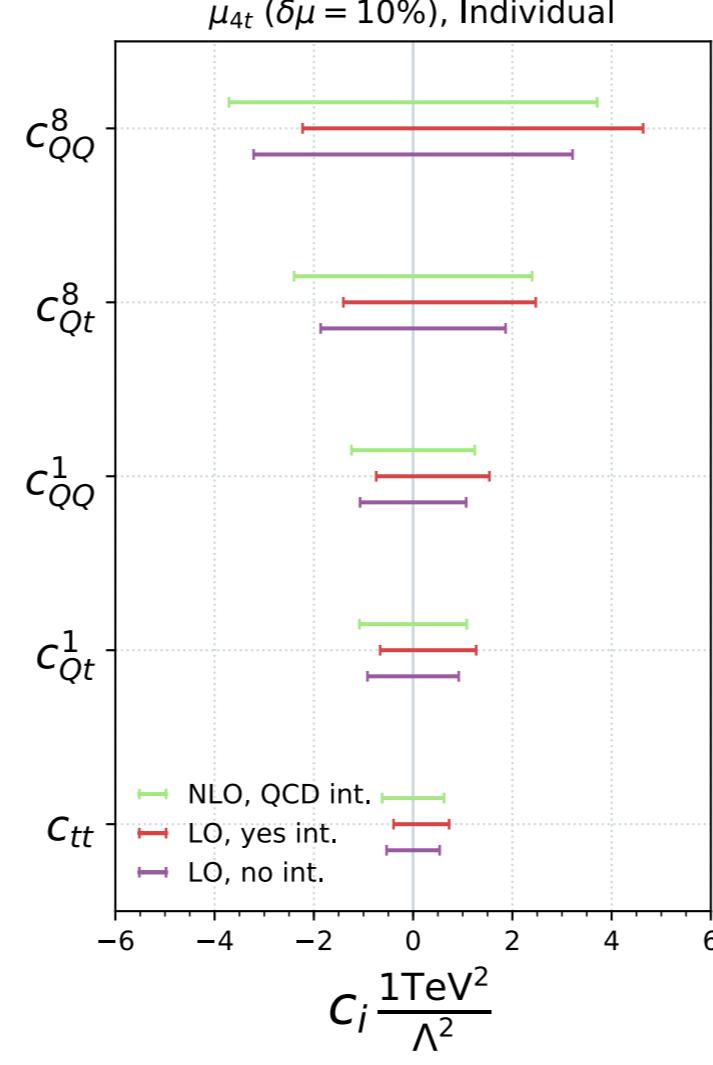
4 top

[Degrande, Durieux, Maltoni, KM,
Vryonidou & Zhang; arXiv:20

$$\sigma(pp \rightarrow t\bar{t}t\bar{t}) [\text{fb}], c_i/\Lambda^2 = 1$$

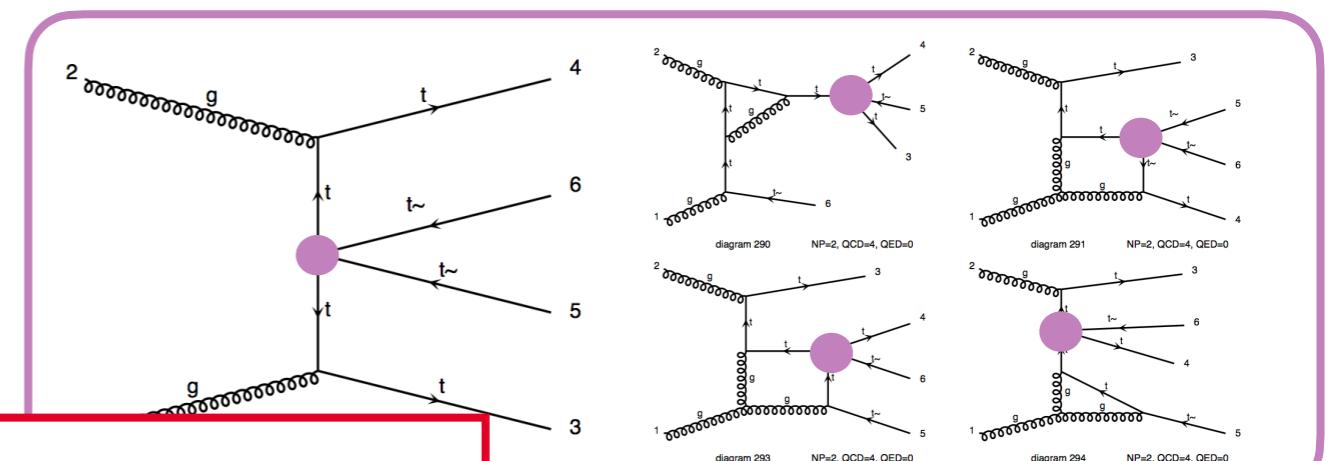
c_i	Interference LO	...
c_{QQ}^8	$0.081^{+55\%}_{-33\%}$	[...]
c_{Qt}^8	$0.274^{+54\%}_{-33\%}$	[...]
c_{QQ}^1	$0.242^{+55\%}_{-33\%}$	[...]
c_{Qt}^1	$-0.0098(10)^{+38\%}_{-33\%}$	[...]
c_{tt}^1	$0.483^{+55\%}_{-33\%}$	[...]

Interference boosted



	C_{4t}	QQ^8	Qt^8	QQ^1	Qt^1	tt
ξ_{4t}	LO	0.048	0.095	0.048	0.002	0.048
	NLO	0.034	0.075	0.03	-	0.034
	+EW	0.33	0.25	0.33	0.29	0.27

	$\sigma_{\text{sq.}}$	σ_{NLO}	σ_{Limit}
QQ^8	2.4	1.1	0.8
Qt^8	1.8	1.0	0.6
QQ^1	0.8	0.6	0.5
Qt^1	0.6	0.5	0.25
tt	0.33	0.25	0.19



Different from arXiv version !

square	$\mathcal{O}(\Lambda^{-4})$	K
LO	NLO	
$0.5^{+46\%}_{-29\%}$	$0.158^{+4\%}_{-11\%}$	1.37
$1.2^{+46\%}_{-29\%}$	$0.378^{+4\%}_{-13\%}$	1.10
$3.9^{+47\%}_{-29\%}$	$1.41^{+4\%}_{-11\%}$	0.93
$9.6^{+46\%}_{-30\%}$	$1.86^{+4\%}_{-10\%}$	1.32
$54^{+47\%}_{-29\%}$	$5.61^{+4\%}_{-11\%}$	1.35

$t\bar{t}t\bar{t}$ amplitude

Status in a nutshell

Global new physics searches via high precision/energy

- Z & W-pole data: handle on the EW gauge sector [Han & Skiba; PRD 71 (2005) 075009]
[Falkowski & Riva; JHEP 02 (2015) 039]
- LHC: thriving Higgs & top programmes
- Probing gauge interactions at high energy (VV , VBS , VVV , ...)

How much cross-talk? Where does being global matter?

We know that Higgs physics greatly complements LEP data

- Access to parameter directions not probed at LEP
- Allows for a closed fit to flavor-universal SMEFT
- Crucial to combine EWPO, Diboson & Higgs data

[Corbett et al.; PRD 87 (2013) 015022]

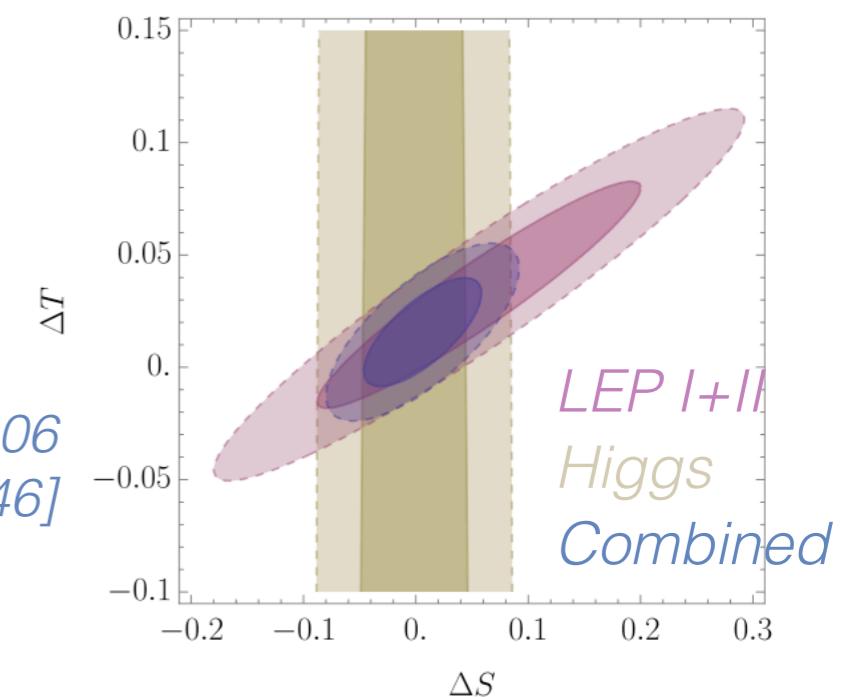
[Pomarol & Riva; JHEP 01 (2014) 151]

[Ellis, Sanz & You; JHEP 03 (2015) 157]

[Biekötter Corbett & Plehn; SciPost Phys 6 (2019) 6, 064]...

[Ellis et al.; JHEP 06

(2018) 146]



The fit

arXiv:2012.02779 **Top, Higgs, Diboson and Electroweak Fit to the Standard Model Effective Field Theory**

John Ellis,^{a,b,c} Maeve Madigan,^d Ken Mimasu,^a Veronica Sanz^{e,f} and Tevong You^{b,d,g}

Global SMEFT interpretation of 4 categories of data

- Based on
- 14 • Electroweak Precision Observables (EWPO): Z-pole & W-mass [Ellis et al.; JHEP 06 (2018) 146]
 - 118 • LEP2 & LHC diboson production: differential WW, WZ, Zjj
 - 72 • Higgs measurements: signal strengths & STXS
 - 124 • Top data: single-top, ttbar & asymmetries, ttV, tZ, tW

Big thanks to authors of
SMEFiT analysis
[JHEP 04 (2019) 100]
for sharing some of their
top predictions

328 measurements across categories

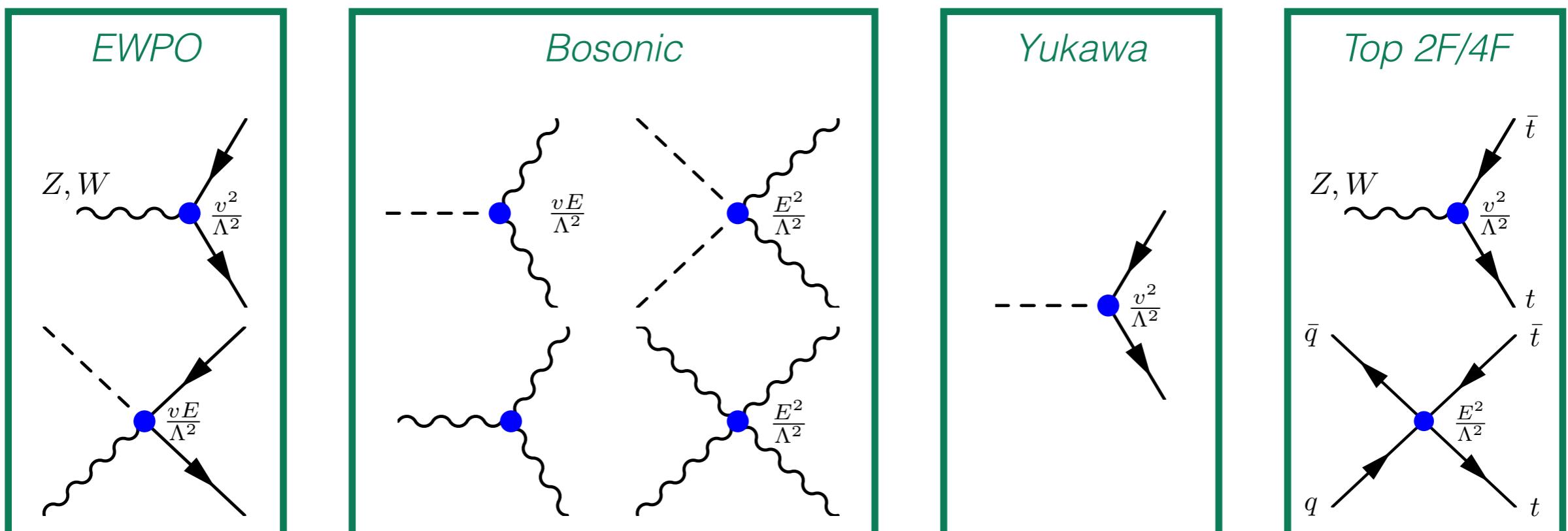
- Chosen to be statistically independent & maximise reach
- Correlations included when publicly available (mostly are)

Linear EFT approximation: $\mu_X \equiv \frac{X}{X_{SM}} = 1 + \sum_i a_i^X \frac{C_i}{\Lambda^2} + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$

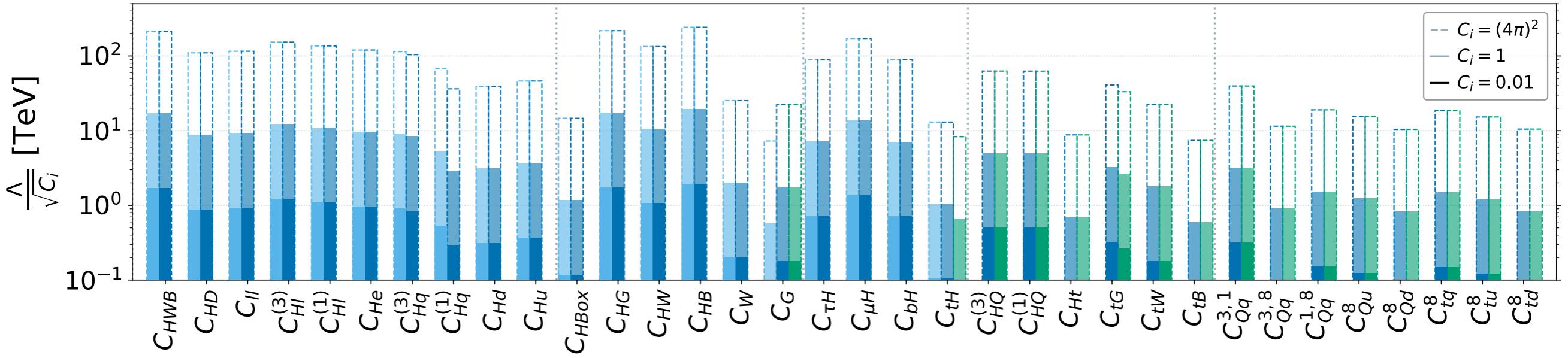
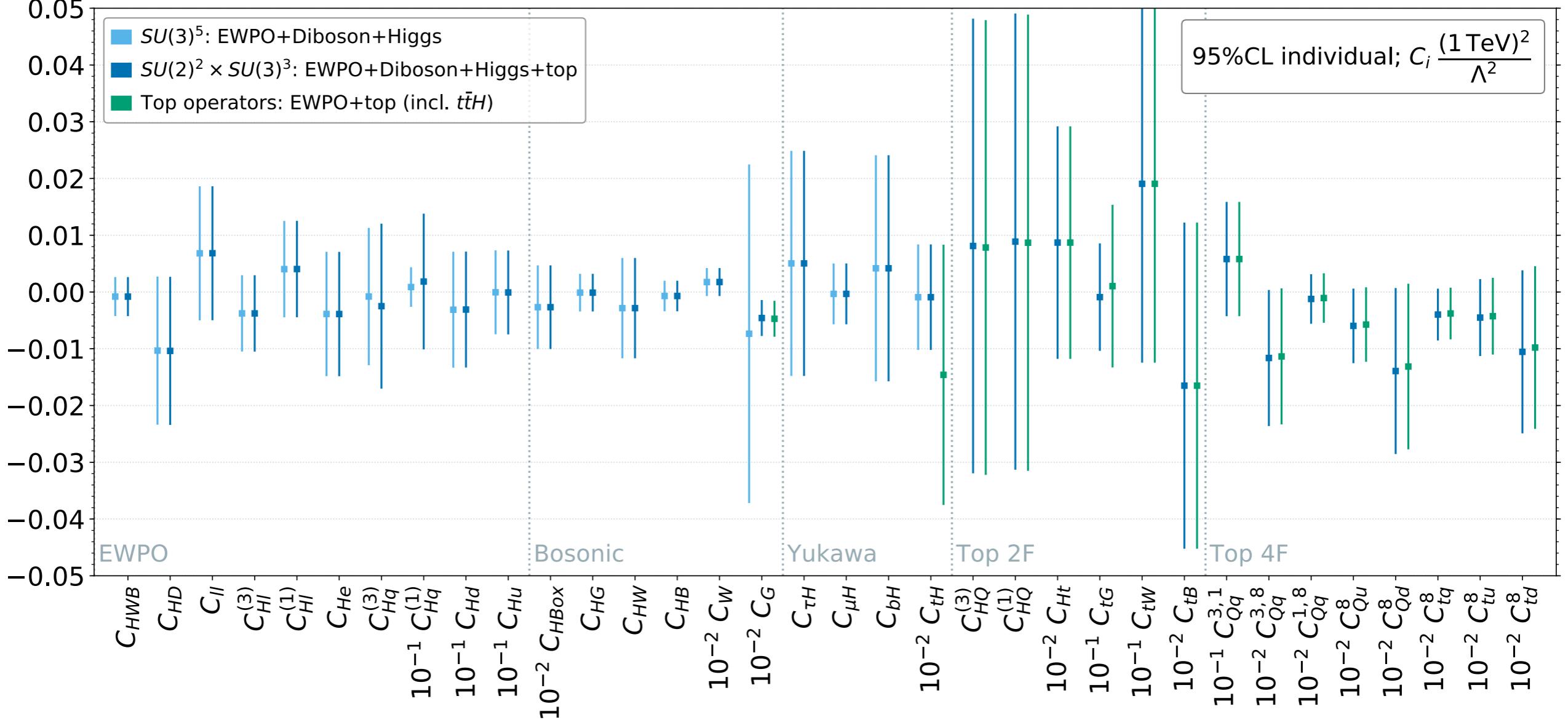
Degrees of freedom

EWPO:	$\mathcal{O}_{HWB}, \mathcal{O}_{HD}, \mathcal{O}_{ll}, \mathcal{O}_{Hl}^{(3)}, \mathcal{O}_{Hl}^{(1)}, \mathcal{O}_{He}, \mathcal{O}_{Hq}^{(3)}, \mathcal{O}_{Hq}^{(1)}, \mathcal{O}_{Hd}, \mathcal{O}_{Hu},$
Bosonic:	$\mathcal{O}_{H\square}, \mathcal{O}_{HG}, \mathcal{O}_{HW}, \mathcal{O}_{HB}, \mathcal{O}_W, \mathcal{O}_G,$
Yukawa:	$\mathcal{O}_{\tau H}, \mathcal{O}_{\mu H}, \mathcal{O}_{b H}, \mathcal{O}_{t H},$
Top 2F:	$\mathcal{O}_{HQ}^{(3)}, \mathcal{O}_{HQ}^{(1)}, \mathcal{O}_{Ht}, \mathcal{O}_{tG}, \mathcal{O}_{tW}, \mathcal{O}_{tB},$
Top 4F:	$\mathcal{O}_{Qq}^{3,1}, \mathcal{O}_{Qq}^{3,8}, \mathcal{O}_{Qq}^{1,8}, \mathcal{O}_{Qu}^8, \mathcal{O}_{Qd}^8, \mathcal{O}_{tQ}^8, \mathcal{O}_{tu}^8, \mathcal{O}_{td}^8.$
	20
	+ 14

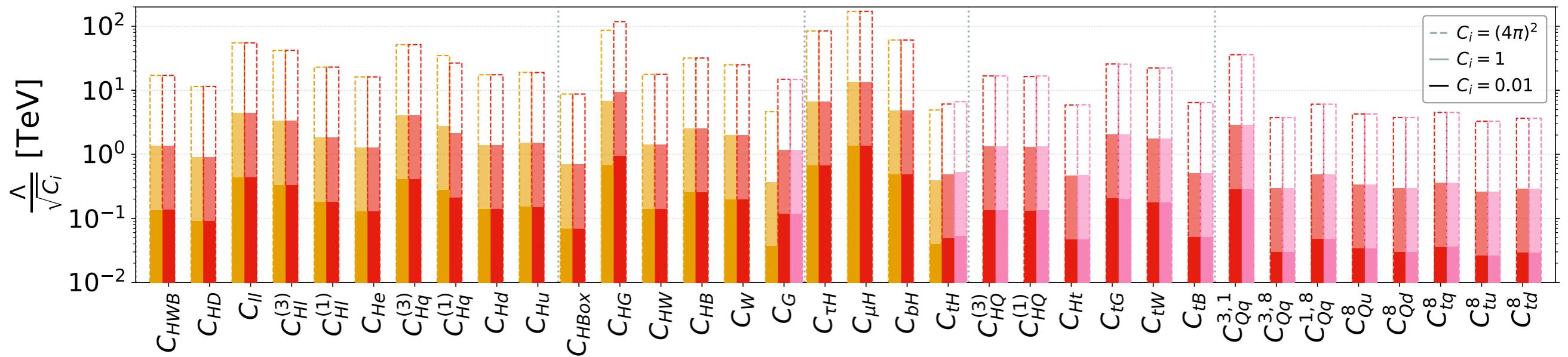
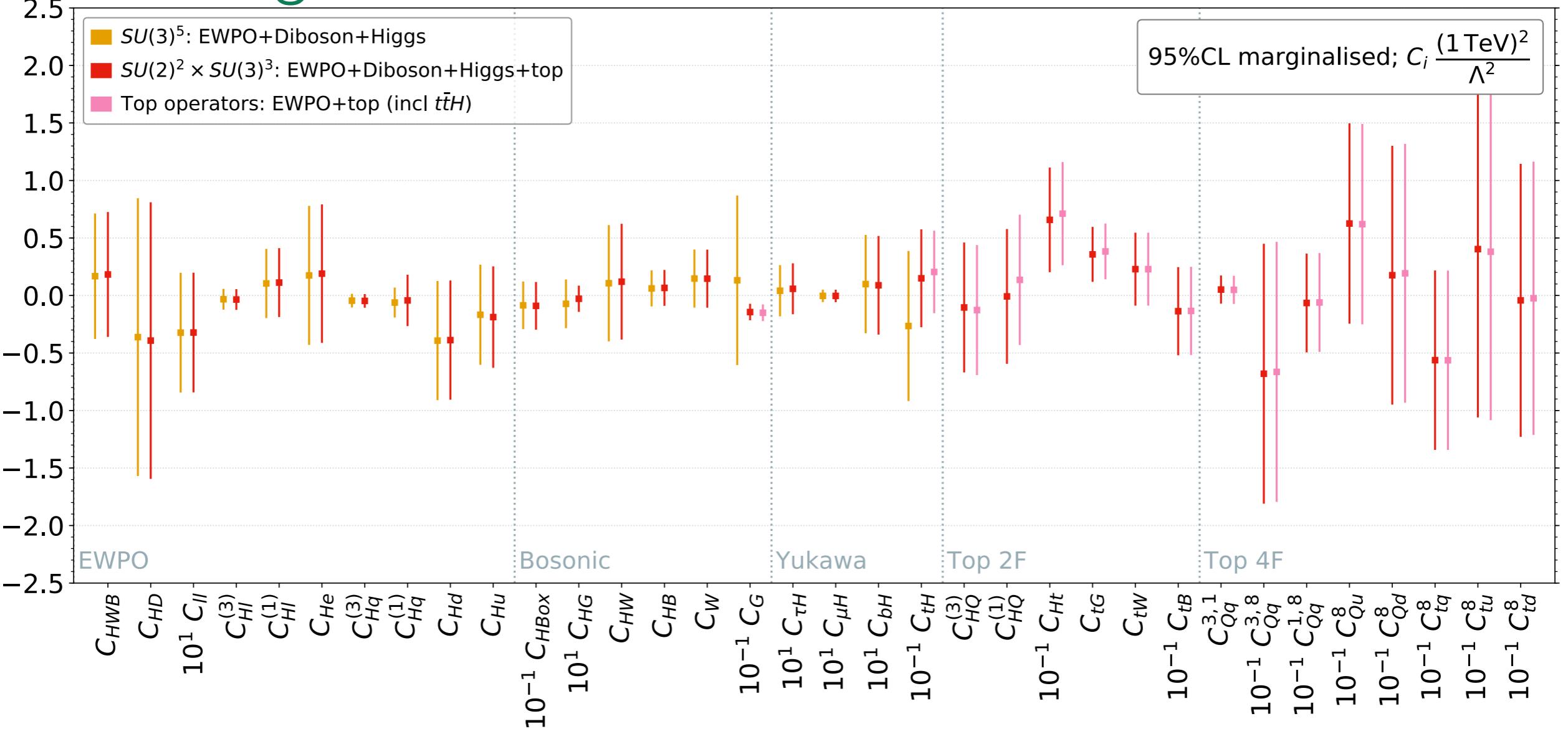
In total: 20(34) d.o.f. for the two flavor scenarios



Full fit: individual



Full fit: marginalised



Correlations

Block diagonal:
correlations *within*
'sector'

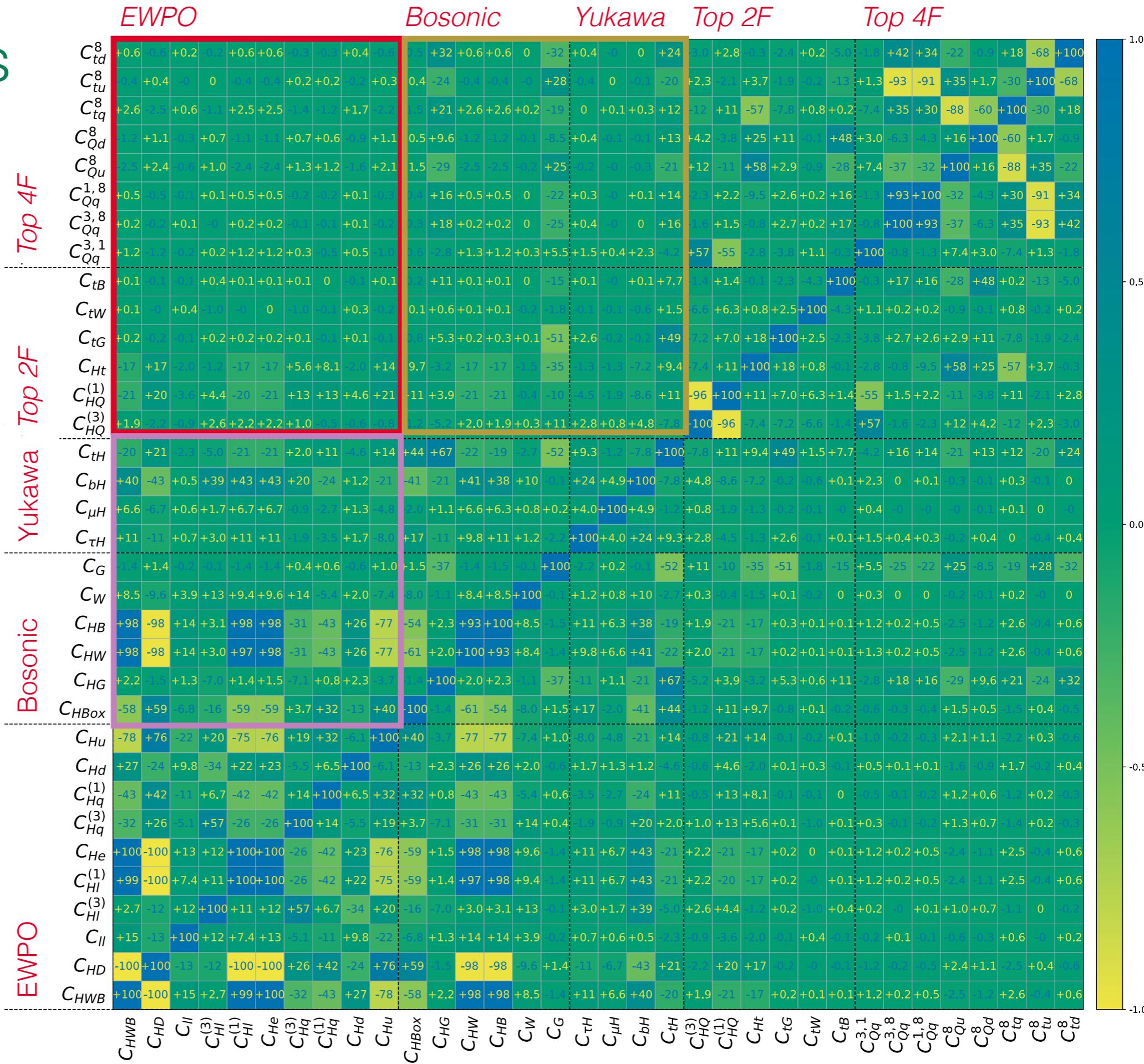
Block off-diagonal:
correlations *among*
'sectors'

EWPO & top ~uncorrelated

EWPO-Higgs $C_{HB}, C_{HW}, C_{H\square}$ & Yukawa with EWPO

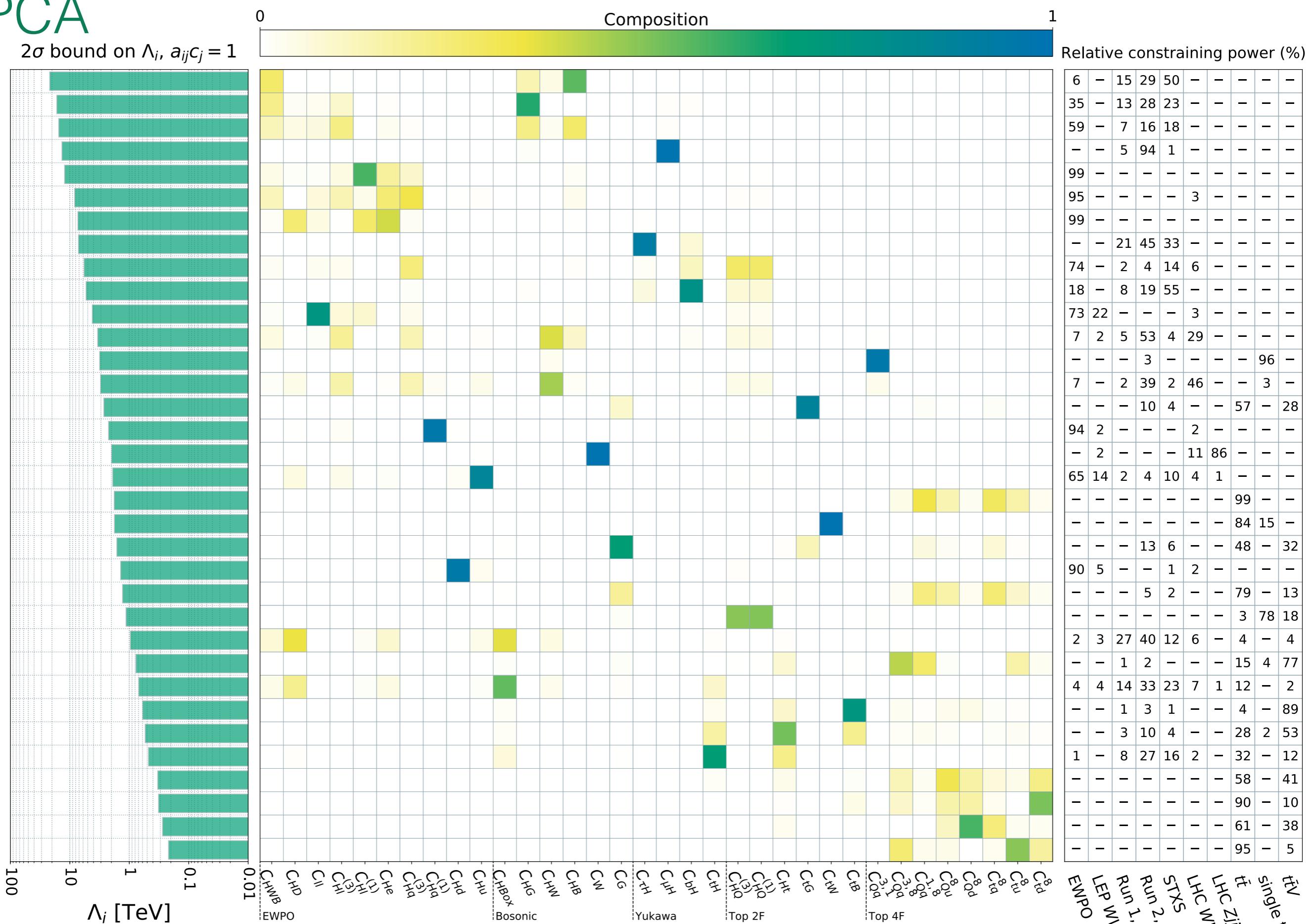
Higgs precision rivalling LEP

Top-Higgs C_{HG}, C_G, C_{tH} with 4F



PCA

2σ bound on Λ_i , $a_{ij}c_j = 1$



Single field extensions

Name	Spin	SU(3)	SU(2)	U(1)	Name	Spin	SU(3)	SU(2)	U(1)
S	0	1	1	0	Δ_1	$\frac{1}{2}$	1	2	$-\frac{1}{2}$
S_1	0	1	1	1	Δ_3	$\frac{1}{2}$	1	2	$-\frac{1}{2}$
φ	0	1	2	$\frac{1}{2}$	Σ	$\frac{1}{2}$	1	3	0
Ξ	0	1	3	0	Σ_1	$\frac{1}{2}$	1	3	-1
Ξ_1	0	1	3	1	U	$\frac{1}{2}$	3	1	$\frac{2}{3}$
B	1	1	1	0	D	$\frac{1}{2}$	3	1	$-\frac{1}{3}$
B_1	1	1	1	1	Q_1	$\frac{1}{2}$	3	2	$\frac{1}{6}$
W	1	1	3	0	Q_5	$\frac{1}{2}$	3	2	$-\frac{5}{6}$
W_1	1	1	3	1	Q_7	$\frac{1}{2}$	3	2	$\frac{7}{6}$
N	$\frac{1}{2}$	1	1	0	T_1	$\frac{1}{2}$	3	3	$-\frac{1}{3}$
E	$\frac{1}{2}$	1	1	-1	T_2	$\frac{1}{2}$	3	3	$\frac{2}{3}$
T	$\frac{1}{2}$	3	1	$\frac{2}{3}$	TB	$\frac{1}{2}$	3	2	$\frac{1}{6}$

Considered single field extensions of the SM

- Complete tree-level matching dictionary is known [de Blas et al.; JHEP 03 (2018) 109]
- Interpret in terms of simplified 1 & 2 parameter versions of the models

One parameter models

Model	C_{HD}	C_{ll}	C_{Hl}^3	C_{Hl}^1	C_{He}	$C_{H\square}$	$C_{\tau H}$	C_{tH}	C_{bH}
S						-1			
S_1		1							
Σ			$\frac{5}{8}$	$\frac{3}{16}$			$\frac{y_\tau}{4}$		
Σ_1			$-\frac{5}{8}$	$-\frac{3}{16}$			$\frac{y_\tau}{8}$		
N			$-\frac{1}{4}$	$\frac{1}{4}$					
E			$-\frac{1}{4}$	$-\frac{1}{4}$			$\frac{y_\tau}{2}$		
Δ_1					$\frac{1}{2}$		$\frac{y_\tau}{2}$		
Δ_3					$-\frac{1}{2}$		$\frac{y_\tau}{2}$		
B_1	1					$-\frac{1}{2}$	$-\frac{y_\tau}{2}$	$-\frac{y_t}{2}$	$-\frac{y_b}{2}$
Ξ	-2					$\frac{1}{2}$	y_τ	y_t	y_b
W_1	$-\frac{1}{4}$					$-\frac{1}{8}$	$-\frac{y_\tau}{8}$	$-\frac{y_t}{8}$	$-\frac{y_b}{8}$
φ							$-y_\tau$	$-y_t$	$-y_b$
$\{B, B_1\}$						1	y_τ	y_t	y_b
$\{Q_1, Q_7\}$								y_t	

Model	C_{HG}	C_{Hq}^3	C_{Hq}^1	$(C_{Hq}^3)_{33}$	$(C_{Hq}^1)_{33}$	C_{Hu}	C_{Hd}	C_{tH}	C_{bH}
U		$-\frac{1}{4}$	$\frac{1}{4}$	$-\frac{1}{4}$	$\frac{1}{4}$			$\frac{y_t}{2}$	
D		$-\frac{1}{4}$	$-\frac{1}{4}$	$-\frac{1}{4}$	$-\frac{1}{4}$				$\frac{y_b}{2}$
Q_5							$-\frac{1}{2}$		$\frac{y_b}{2}$
Q_7						$\frac{1}{2}$		$\frac{y_t}{2}$	
T_1		$-\frac{5}{8}$	$-\frac{3}{16}$	$-\frac{5}{8}$	$-\frac{3}{16}$			$\frac{y_t}{4}$	$\frac{y_b}{8}$
T_2		$-\frac{5}{8}$	$\frac{3}{16}$	$-\frac{5}{8}$	$\frac{3}{16}$			$\frac{y_t}{8}$	$\frac{y_b}{4}$
T	$-\frac{M_T^2}{v^2} \frac{\alpha_s(0.02)}{8\pi}$			$-\frac{1}{2} \frac{M_T^2}{v^2}$	$\frac{1}{2} \frac{M_T^2}{v^2}$			$y_t \frac{M_T^2}{v^2}$	

$$\times \frac{\lambda^2}{M^2}$$

One parameter models

