Time-Dependent Calculations of Strong Field Ionization and High Harmonics Generation

Anatoli Kheifets and Igor Ivanov



THE AUSTRALIAN NATIONAL UNIVERSITY

Igor Bray



Centre for Antimatter-Matter Studies

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> Interplay between strong field effects and many-electron correlations

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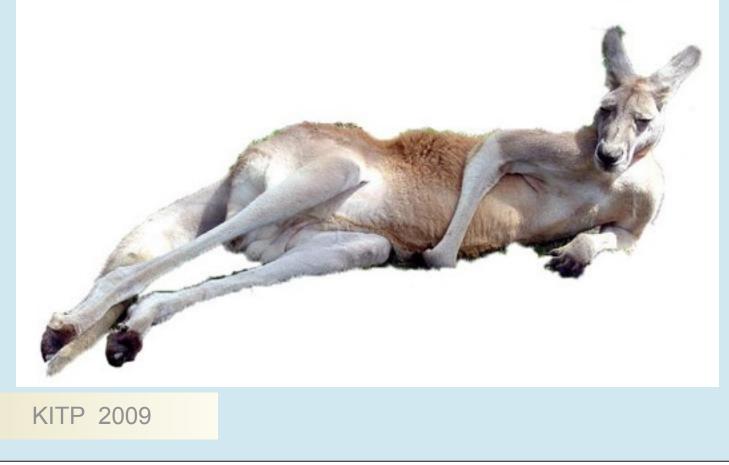


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Interplay between strong field effects and many-electron correlations



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Strong Field Ionization



Strong Field Ionization

- New XUV sources
 - FEL's (FLASH, SCSS, LCLS etc)

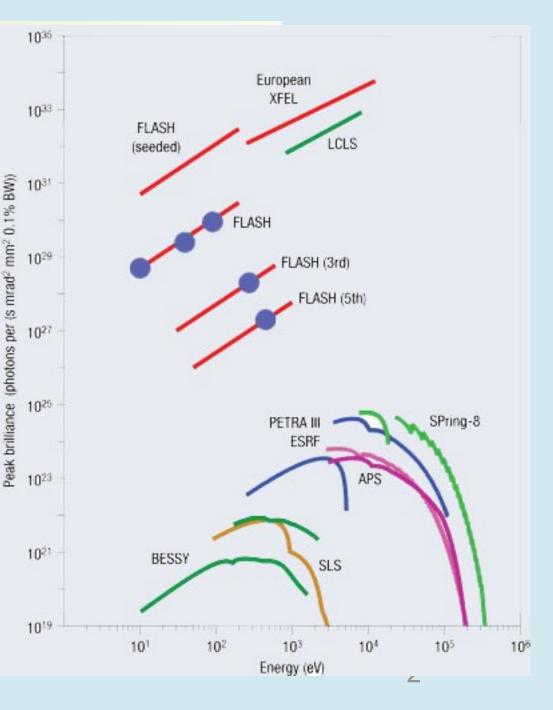


Strong Field Ioniz

New XUV sources

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New XUV sources
 FEL's (FLASH, SCSS, FEL's (FLASH, SCSS)



Strong Field Ionization

- New XUV sources
 - FEL's (FLASH, SCSS, LCLS etc)
 - High harmonic generation



Strong Field Ionization

- New XUV sources
 - FEL's (FLASH, SCSS, LCLS etc)
 - High harmonic generation
- Fundamental importance
 - Interplay of correlations and strong field



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High harmonic generation

- Race for the "table-top synchrotron"
 - Coherent, compact, intense, tunable, source of XUV radiation
 - Neutral gases or ionized plasmas of noble gas atoms



Strong Field Ionization

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 - High harmonic generation
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High harmonic generation

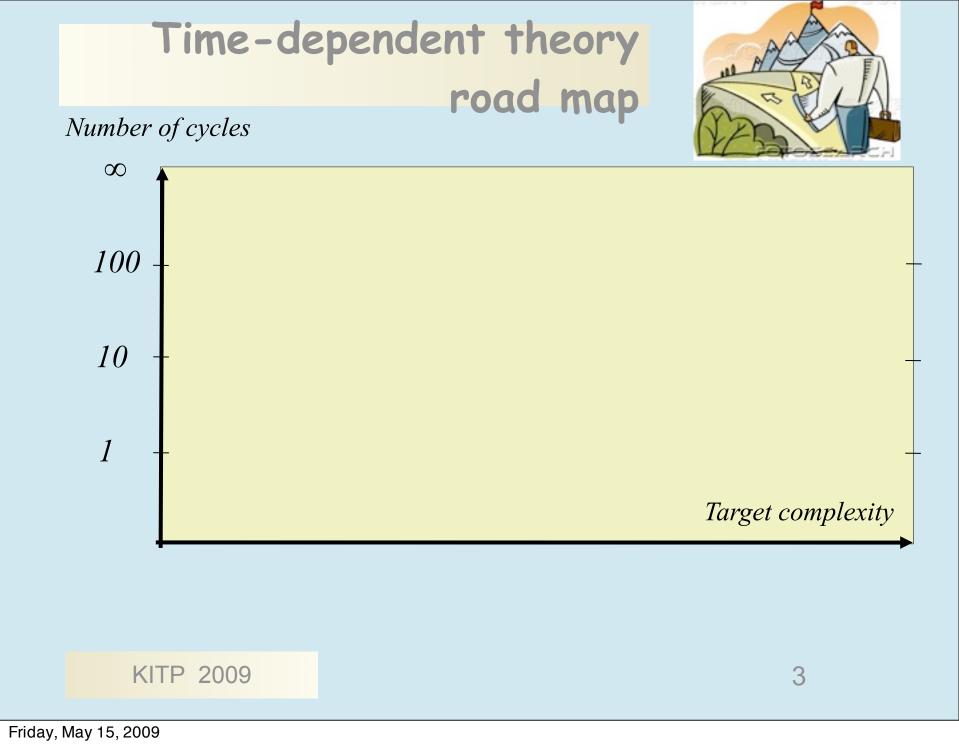
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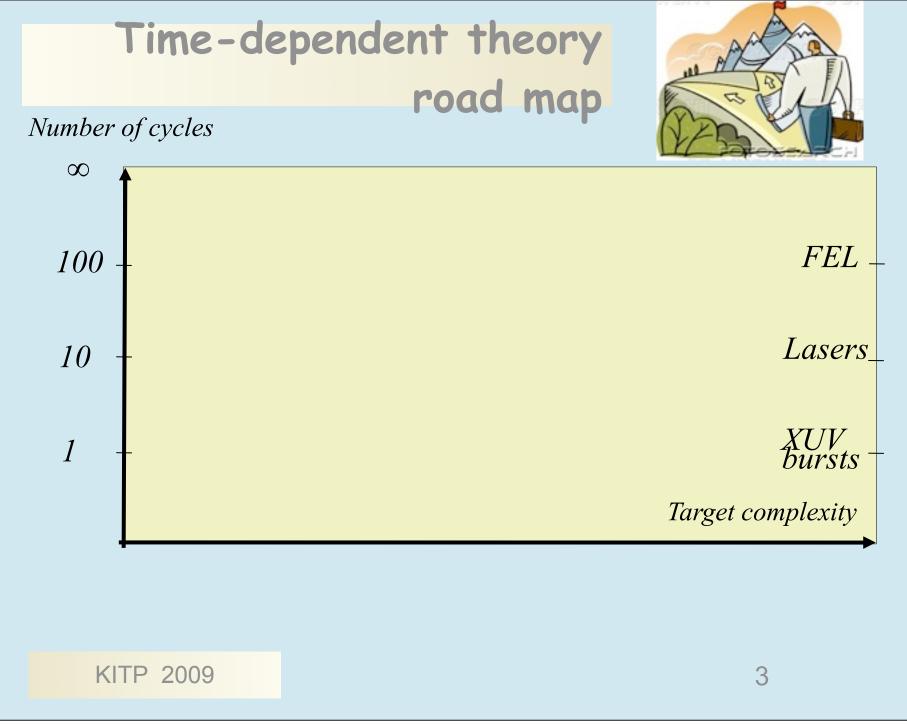
Attosecond control

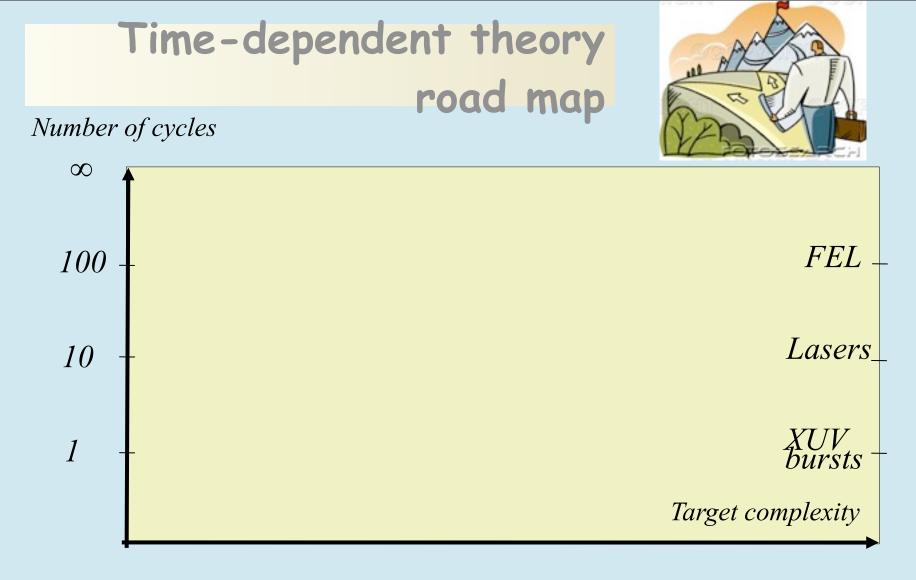
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- Steering of coherent wave packets.
- Time-resolved atomic physics

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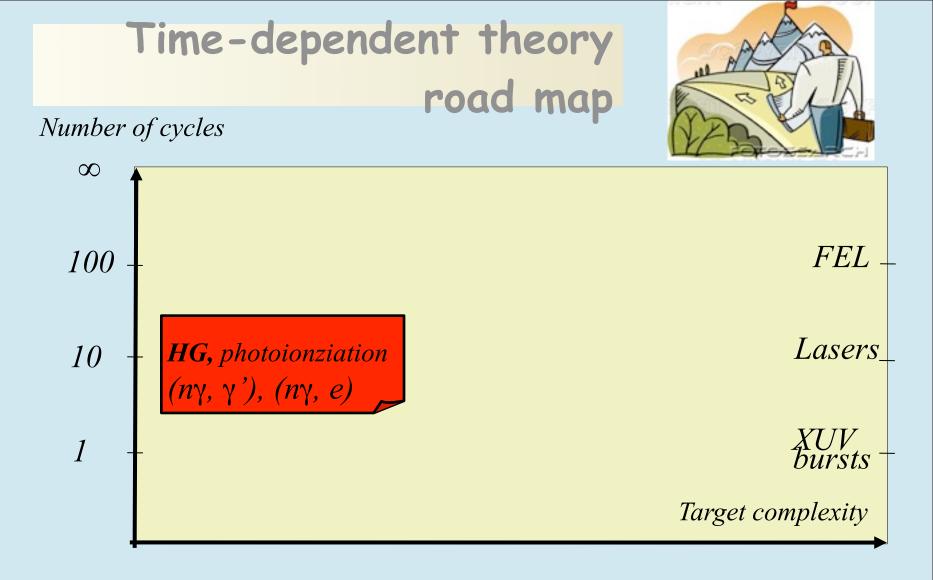






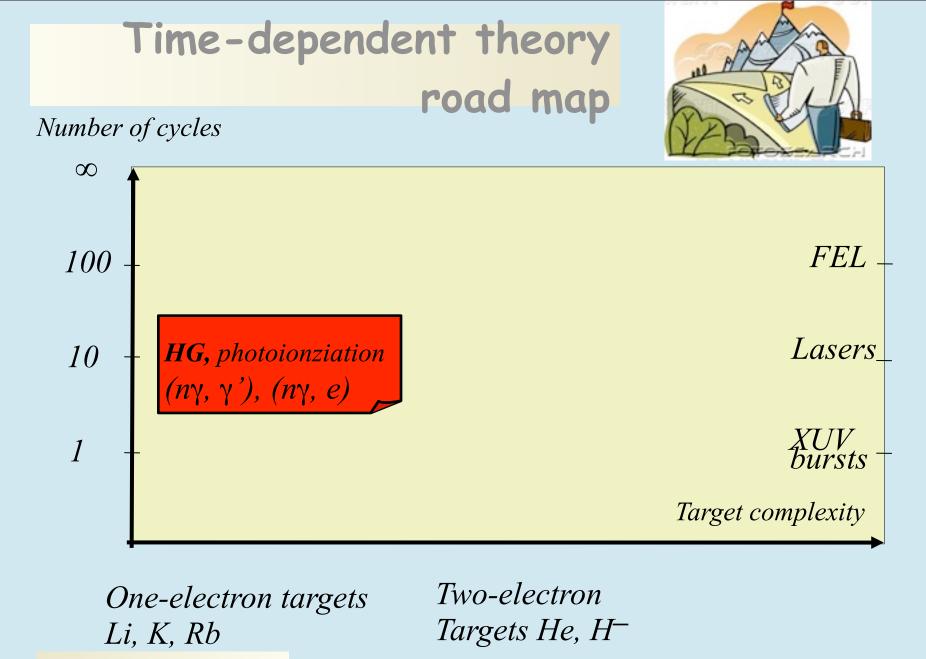
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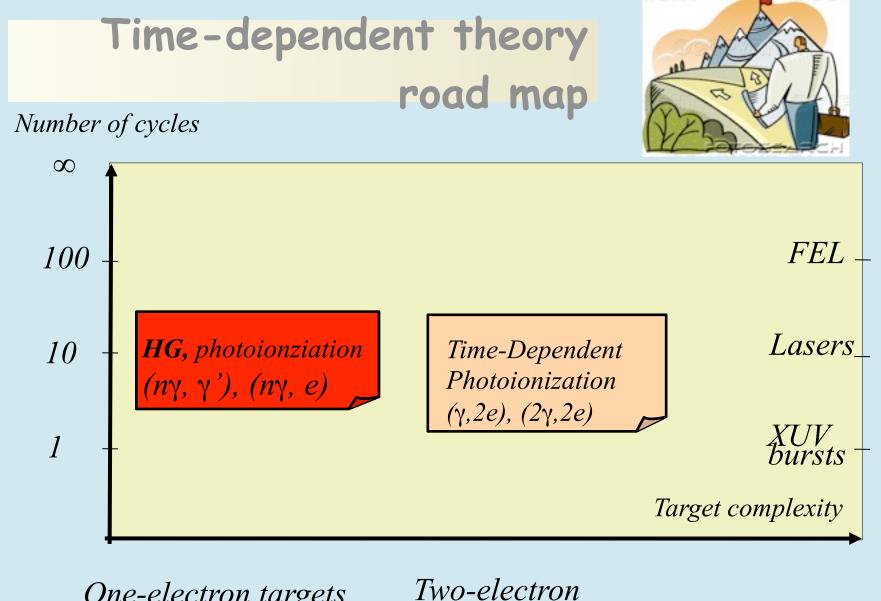
One-electron targets Li, K, Rb KITP 2009



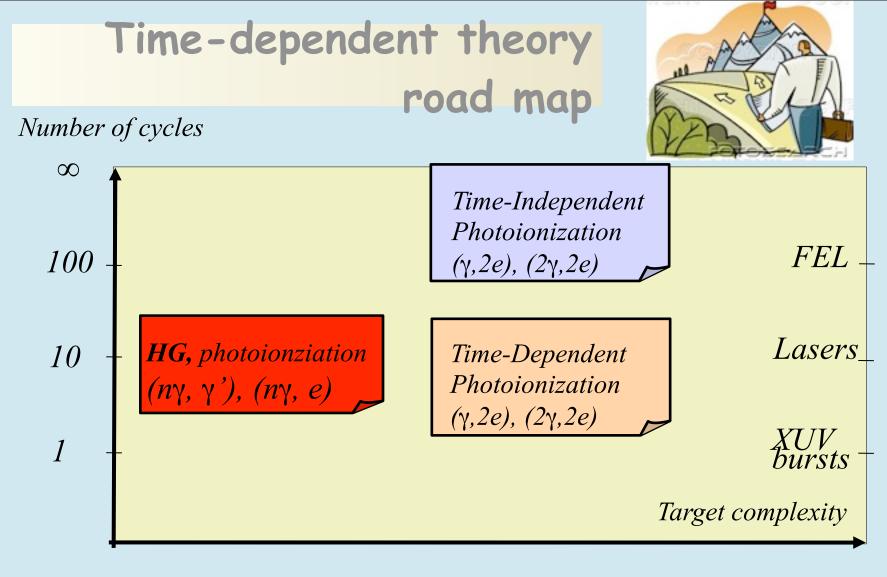
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One-electron targets Li, K, Rb KITP 2009



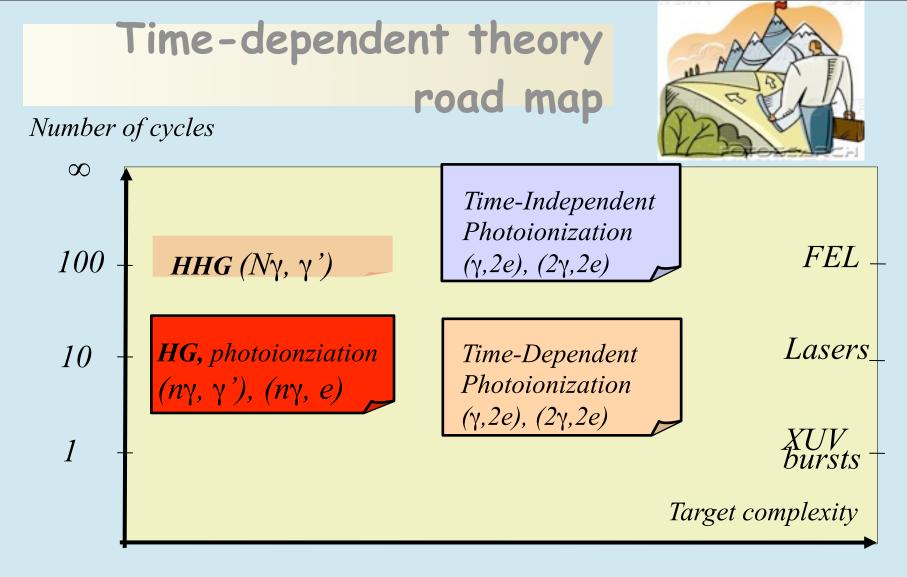


Two-electron Targets He, H⁻

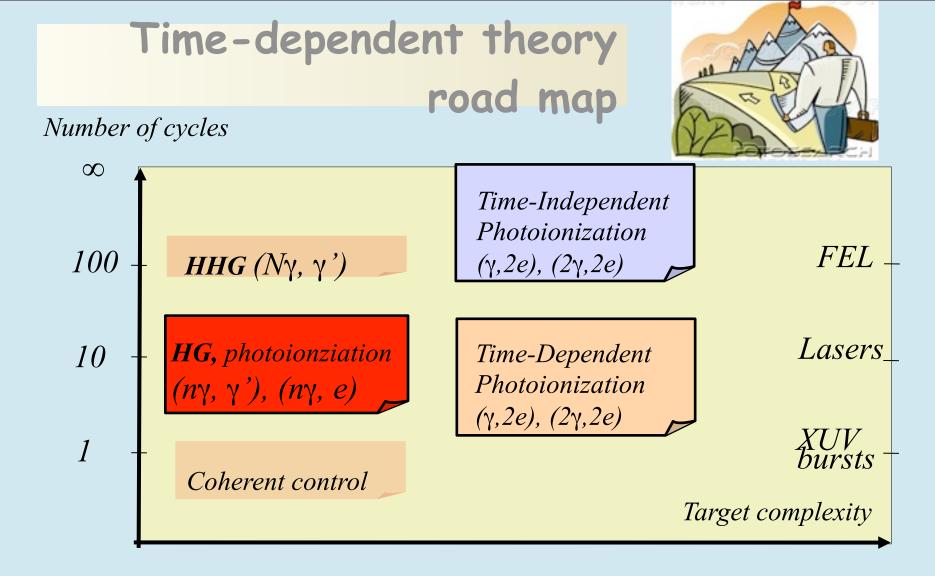


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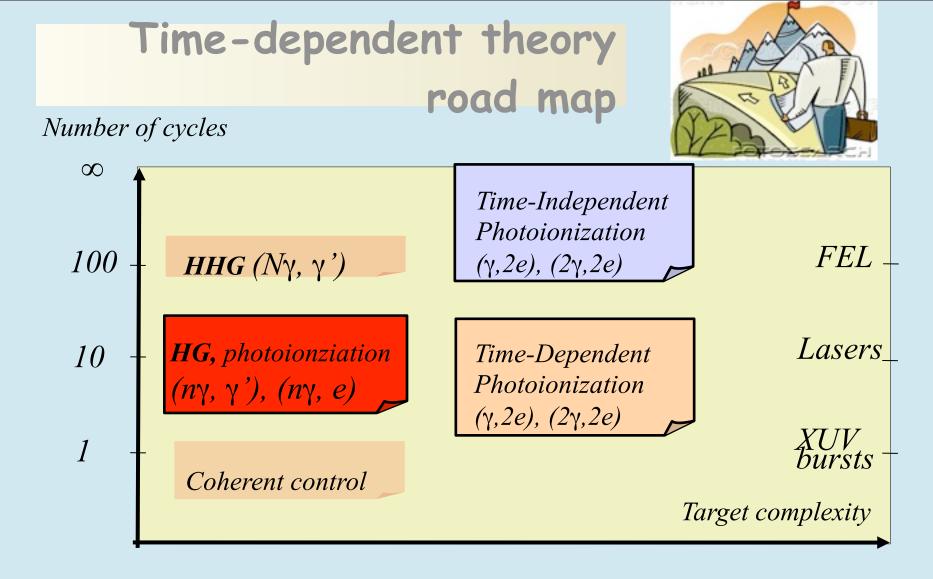
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Two-electron Targets He, H⁻⁻



Two-electron Targets He, H⁻⁻



Two-electron Targets He, H⁻ Many-electron Targets Ne, Ar, Xe

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- Theoretical model
 - Time-dependent Schrödinger equation



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- Theoretical model
 - Time-dependent Schrödinger equation
 - Basis based integration vs coordinate space integration



- Theoretical model
 - Time-dependent Schrödinger equation
 - Basis based integration vs coordinate space integration
 - Field-free atomic states



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 - Discretization of continuum
 - Pseudostates vs continuum states



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Applications



- Applications
 - Multi-photon ionization



- Applications
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 - Two-photon double ionization of He



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 - Four-photon ionization of Li



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Outline

- Applications
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 - Two-photon double ionization of He
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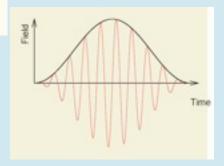
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 - Optimal control

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Field on:





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Field on:

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Time-Dependent Schrödiger equation $i\frac{\partial\Psi}{\partial t}$

$$i\frac{\partial\Psi}{\partial t} = \hat{H}\Psi$$

$$\hat{H} = \frac{p_1^2}{2} + \frac{p_2^2}{2} - \frac{2}{r_1} - \frac{2}{r_2} + \frac{1}{|r_1 - r_2|} + \frac{D \cdot F_{AC} \cos \omega t}{|r_1 - r_2|}$$



Field on:

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$$\hat{H} = \frac{p_1^2}{2} + \frac{p_2^2}{2} - \frac{2}{r_1} - \frac{2}{r_2} + \frac{1}{|r_1 - r_2|} + \frac{D \cdot F_{AC} \cos \omega t}{|r_1 - r_2|}$$

Solution:

$$\Psi(\boldsymbol{r}_{1}, \boldsymbol{r}_{2}, t) = \sum_{j \equiv \left\{ {n_{1}n_{2}N \atop l_{1}l_{2} \ L} \right\}} a_{j}(t) \ \phi_{n_{1}l_{1}}^{N}(r_{1})\phi_{n_{2}l_{2}}^{N}(r_{2}) \ |l_{1}(1)l_{2}(2)L\rangle$$

$$\mathsf{Pseudostate \ basis:} \ \langle \phi_{i}^{N}|\hat{H}|\phi_{j}^{N}\rangle = E_{i}\delta_{ij}$$

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Field on:

Time-Dependent Schrödiger equation $i \frac{\partial \Psi}{\partial t} = \hat{H} \Psi$

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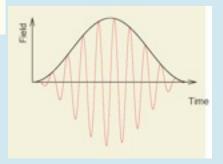
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Pseudostate basis: $\langle \phi_{i}^{N} | \hat{H} | \phi_{j}^{N} \rangle = E_{i}\delta_{ij}$

Matrix notations:

 $i\mathbf{R}\cdot\dot{\mathbf{a}}(t) = \mathbf{H}\cdot\mathbf{a}(t)$, with initial condition $\Psi(\mathbf{r}_1,\mathbf{r}_2,t=0) = \Psi_0(\mathbf{r}_1,\mathbf{r}_2)$

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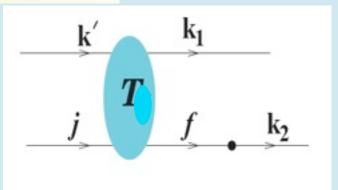
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Field off:



Field off:

CCC expansion for two-electron continuum

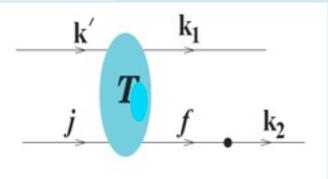




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Field off:

CCC expansion for two-electron continuum



$$\Psi_{k_1k_2}(r_1, r_2) = \Psi_{k_1f}(r_1, r_2) \langle k_2 | f \rangle \quad , \ \epsilon_f = k_2^2/2$$

$$\Psi_{kf}(\boldsymbol{r}_1, \boldsymbol{r}_2) = \chi_{k}(\boldsymbol{r}_1)\phi_f(\boldsymbol{r}_2) + \sum_{\boldsymbol{k}'j} \frac{\langle \boldsymbol{k}f|T_J|\boldsymbol{k}'j\rangle}{E - k'^2/2 - \varepsilon_j + i0} \ \chi_{\boldsymbol{k}'}(\boldsymbol{r}_1)\phi_j(\boldsymbol{r}_2)$$

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Field off:

CCC expansion for two-electron continuum

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$$\Psi_{kf}(\boldsymbol{r}_1, \boldsymbol{r}_2) = \frac{\chi_{k}(\boldsymbol{r}_1)}{\kappa_{k'j}} \phi_f(\boldsymbol{r}_2) + \sum_{\boldsymbol{k'}j} \frac{\langle \boldsymbol{k}f | T_J | \boldsymbol{k'}j \rangle}{E - k'^2/2 - \varepsilon_j + i0} \ \chi_{\boldsymbol{k'}}(\boldsymbol{r}_1) \phi_j(\boldsymbol{r}_2)$$
Coulomb wave

Field off:

CCC expansion for two-electron continuum

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Coulomb wave
Pseudostate

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Coulomb wave
Pseudostate

Fully differential cross-section

$$\frac{d\sigma}{dk_1 dk_2} = \left| \langle \Psi(\boldsymbol{r}_1, \boldsymbol{r}_2, t = T) | \Psi_{\boldsymbol{k}_1 \boldsymbol{k}_2}(\boldsymbol{r}_1, \boldsymbol{r}_2) \rangle \right|^2 \quad T = N \frac{2\pi}{\omega} \ , \ N \gg 1$$

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Harmonics Generation:



Harmonics Generation:

Harmonics spectrum

$$|d(\omega)|^{2} = \left| \frac{1}{t_{2} - t_{1}} \int_{t_{1}}^{t_{2}} e^{-i\omega t} d(t) dt \right|^{2}$$

. 0



Harmonics Generation:

Harmonics spectrum

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~

Expectation value of dipole operator

$$d(t) = \langle \Psi(t) | z | \Psi(t) \rangle = \int d\mathbf{r} \, dt \, \Psi(\mathbf{r}, t) \, (\mathbf{r} \cdot \mathbf{e}) \Psi(\mathbf{r}, t)$$
$$= \sum_{n,m} z_{mn} \exp[-i(m-n)\omega t] = \sum_{N} z_{N} \exp[-i(2M+1)\omega t]$$

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Harmonics Generation:

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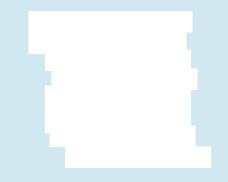
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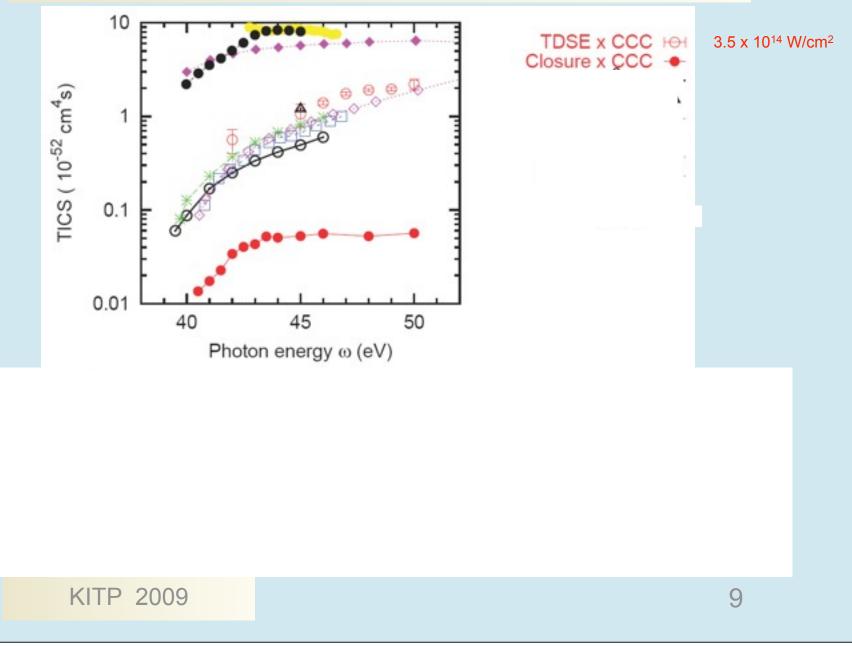
Parity conservation $\Pi = -1 = (-1)^{2M+1}$

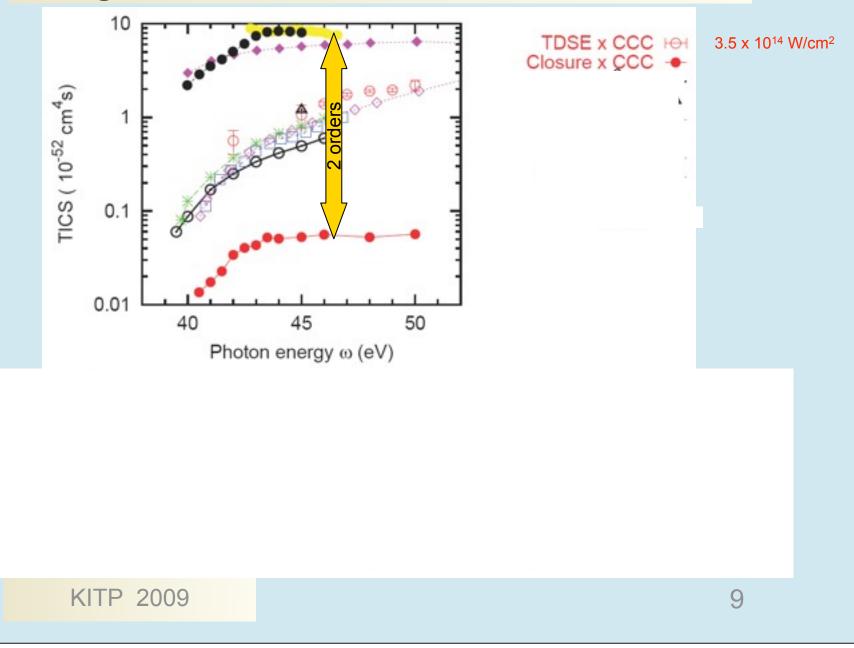
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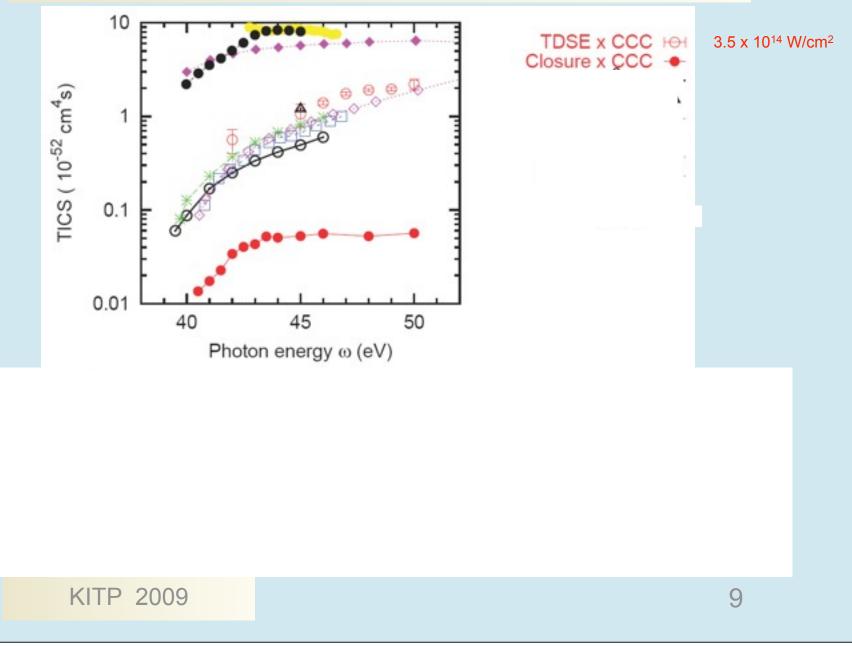


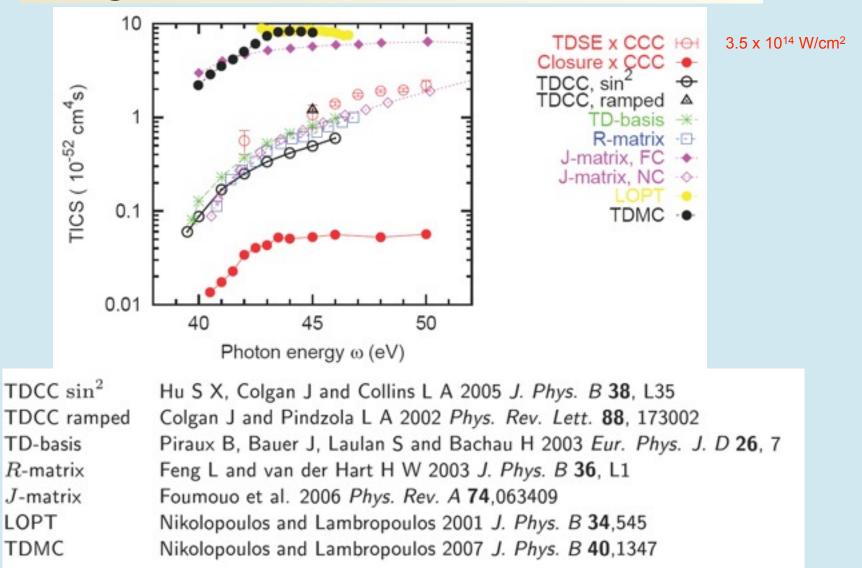


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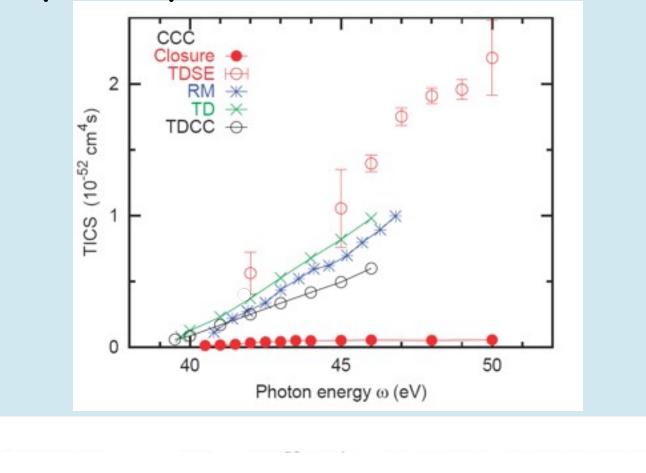
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Theory vs. experiment:

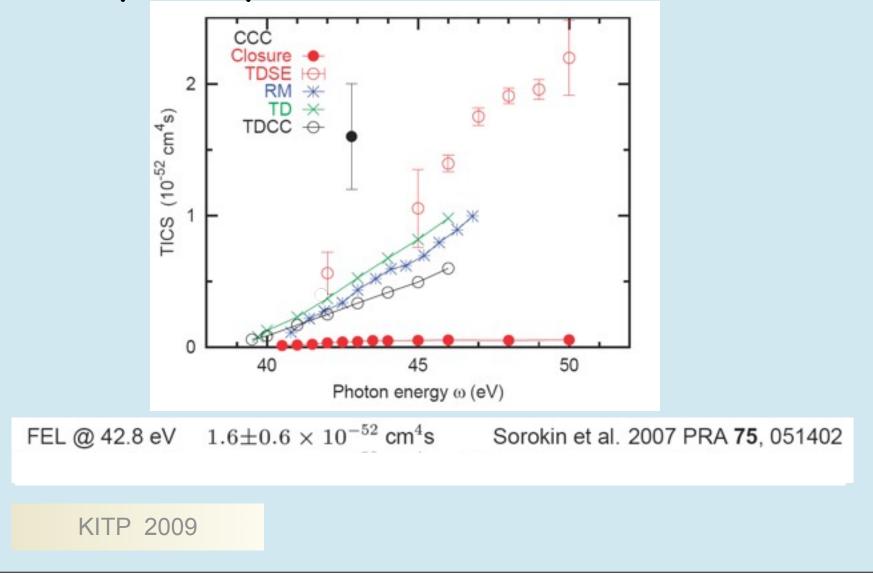


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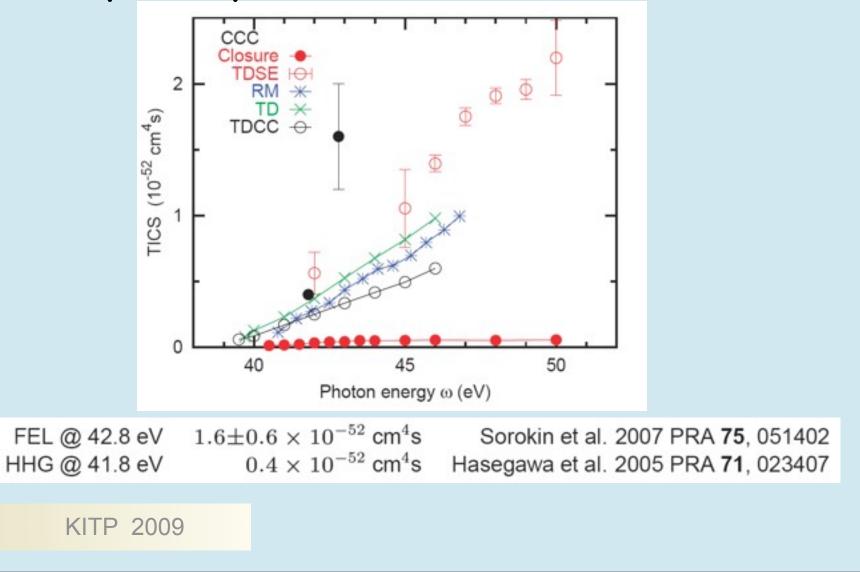


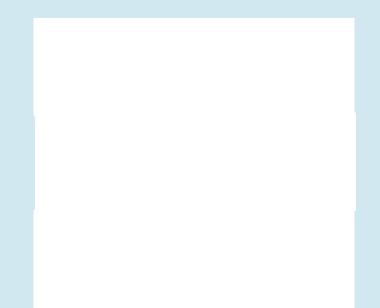
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Theory vs. experiment:



Theory vs. experiment:





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| | | | / | |
|--------|--------------|-------------------------------|------|---------------|
| | 1 st IP [eV] | $I_{\rm OBI} [{\rm W/cm}^2]$ | N | $U_p \; [eV]$ |
| Li(2s) | 5.39 | $3.38 \cdot 10^{12}$ | 5.29 | 0.20 |
| Li(2p) | 3.54 | $6.28 \cdot 10^{11}$ | 9.76 | 0.037 |
| He | 24.59 | $1.46 \cdot 10^{15}$ | 0.53 | 86.1 |
| Ne | 21.57 | $8.65 \cdot 10^{14}$ | 0.65 | 50.5 |
| Ar | 15.76 | $2.47 \cdot 10^{14}$ | 1.0 | 14.6 |
| Xe | 12.13 | $8.66 \cdot 10^{13}$ | 1.5 | 5.11 |

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| | | | / | |
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Low over-the-barrier intensities as compared to noble gases, absence of ponderomotive effects.

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| Li(2s) Li(2p) | 5.39 3.54 | $\begin{array}{c} 3.38 \cdot 10^{12} \\ 6.28 \cdot 10^{11} \end{array}$ | 5.29 9.76 | 0.20 0.037 | High Keldysh-parameter γ, i.e. deep within multiphoton regime, | |
| He | 24.59 | $1.46 \cdot 10^{15}$ | 0.53 | 86.1 | but not perturbative. | |
| Ne | 21.57 | $8.65 \cdot 10^{14}$ | 0.65 | 50.5 | Only 4 (2s) / 3 (2p) Ti:Sa-Photon | |
| Ar | 15.76 | $2.47 \cdot 10^{14}$ | 1.0 | 14.6 | required to ionize. Low non-linearity $(\sigma \propto I^{N})$. | |
| Xe | 12.13 | $8.66 \cdot 10^{13}$ | 1.5 | 5.11 | (0 ~ 1). | |

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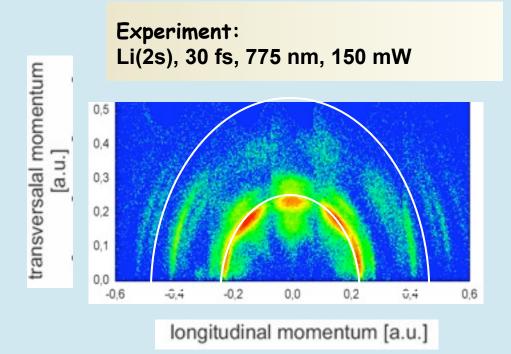
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Photoelectron momentum distribution

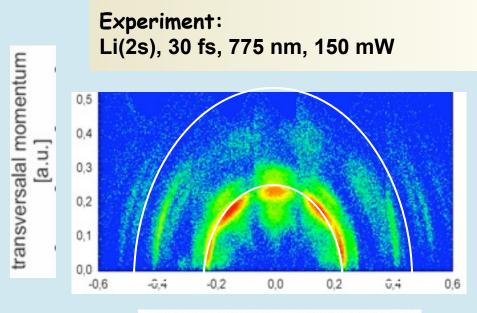


Photoelectron momentum distribution



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Photoelectron momentum distribution



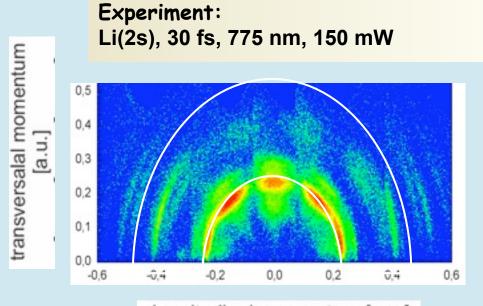
longitudinal momentum [a.u.]

Momentum space multi-photon rings:

 $P_{x}^{2}+P_{z}^{2}=N\omega - IP, N \ge 4$ for Li 2s

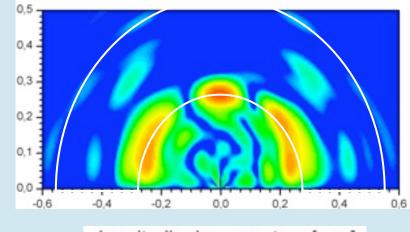
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Photoelectron momentum distribution



longitudinal momentum [a.u.]

Calculation: Li(2s), 25 fs, 800 nm, 3x10¹¹ W/cm²



longitudinal momentum [a.u.]

Momentum space multi-photon rings:

 $P_x^2 + P_z^2 = N \odot - IP$, N ≥ 4 for Li 2s

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Resonant Enhancement of Harmonics Generation

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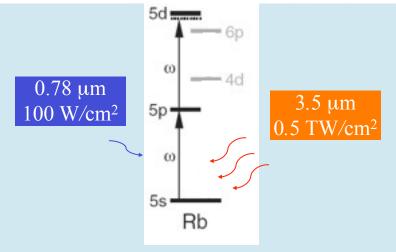
Enhanced High Harmonic Generation from an Optically Prepared Excited Medium



We investigate high harmonics generated from rubidium atoms irradiated simultaneously by an intense 3.5 μ m fundamental field and a weak cw diode laser. When 5p, 5d, and 4d excited states are populated through cascade excitation or deexcitation, orders-of-magnitude increases in harmonic yield as compared with the ground state are observed.

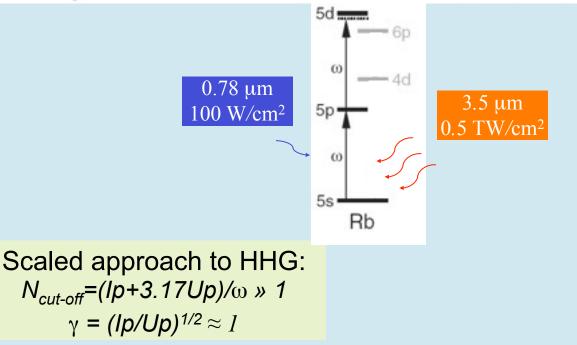


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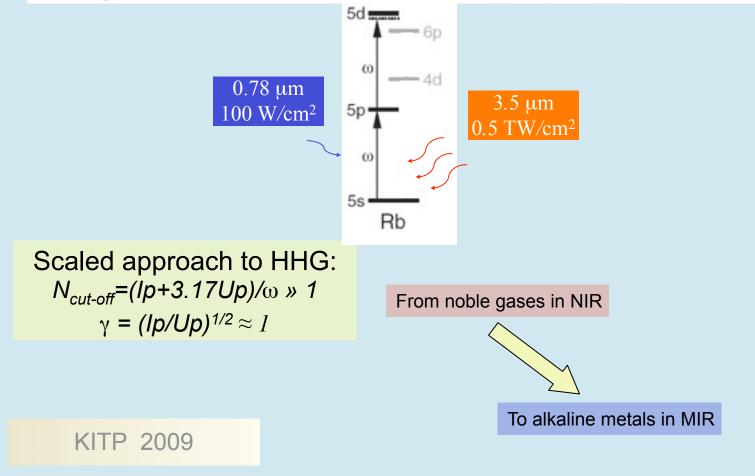


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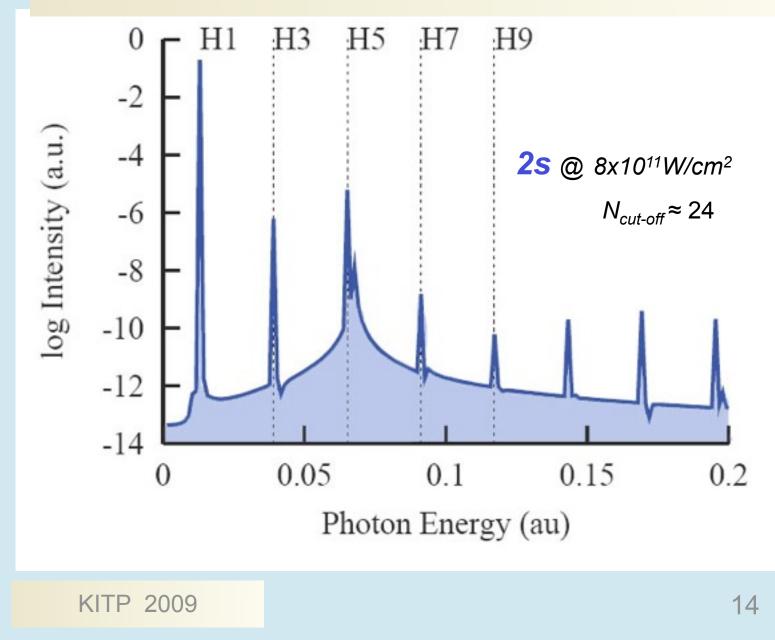
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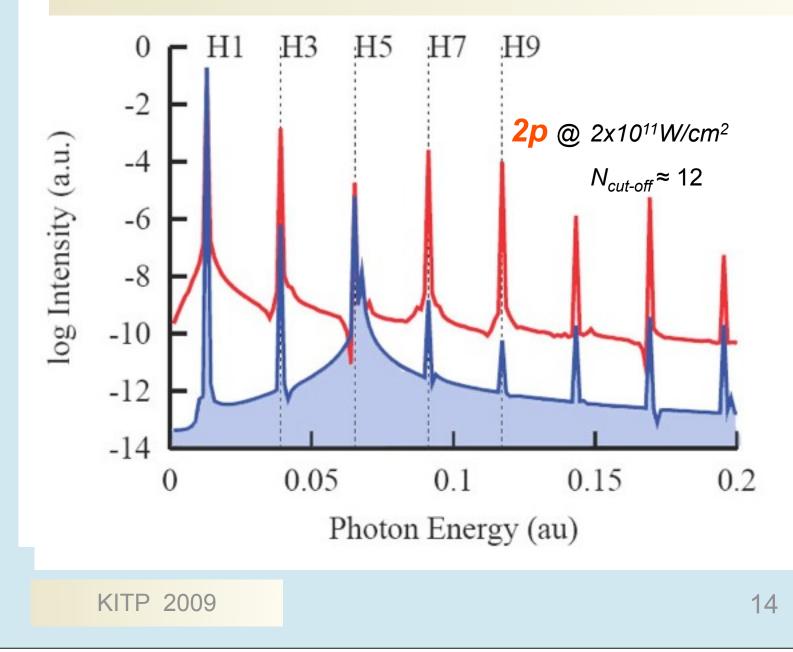
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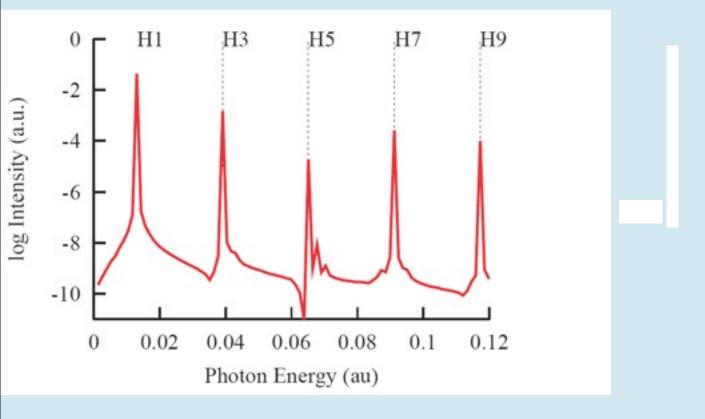


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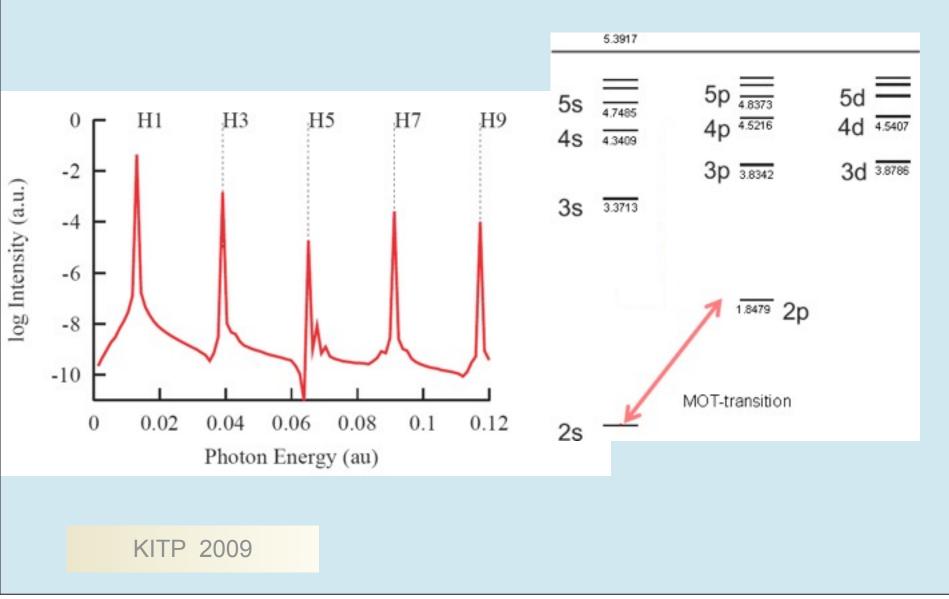
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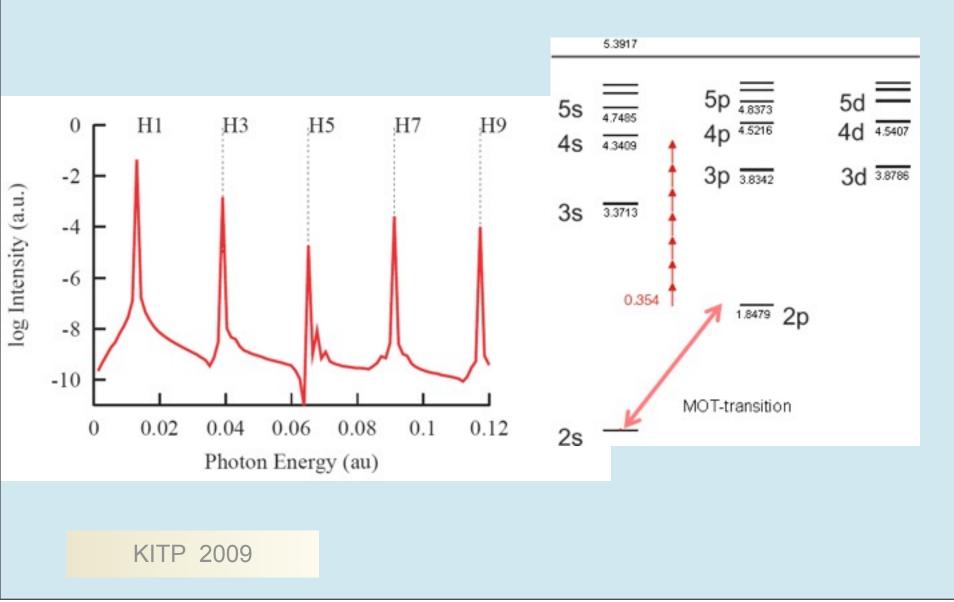


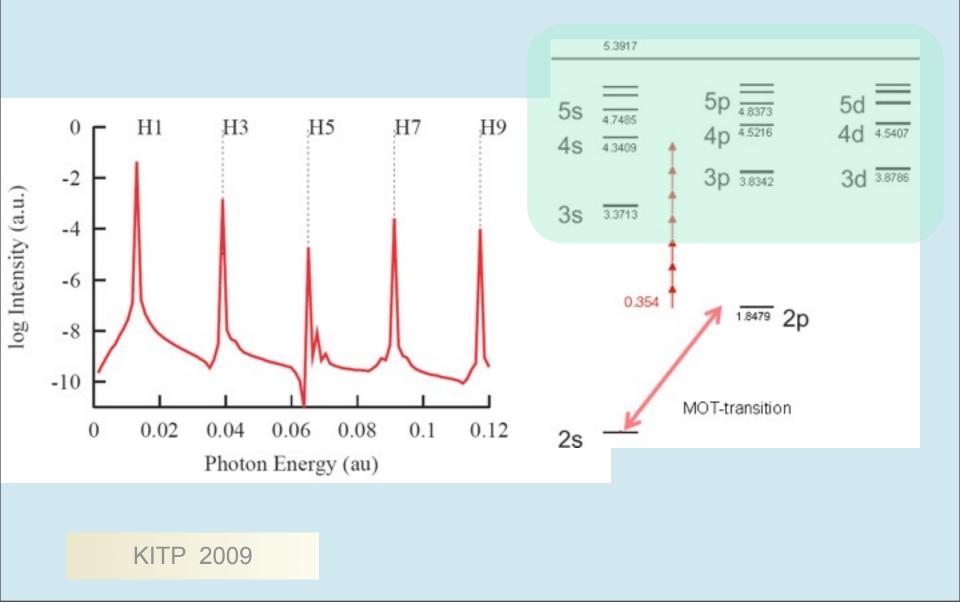


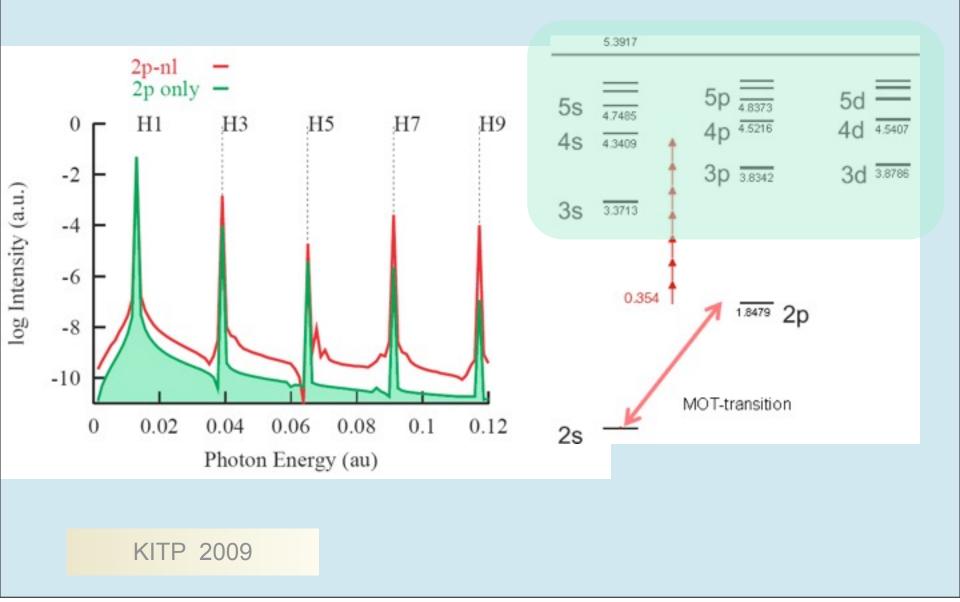


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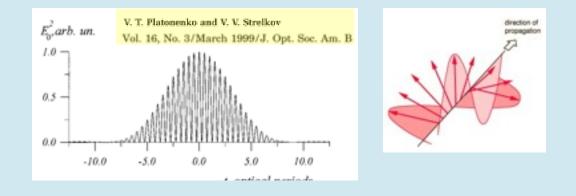






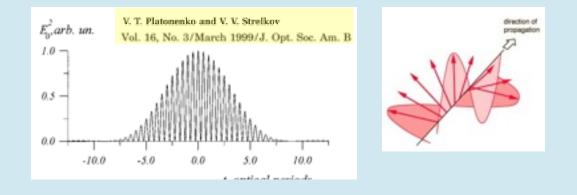
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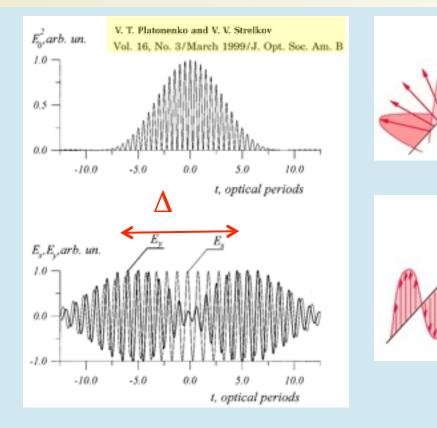


$$\hat{H}_{int}(t) = E_0 \left[(xf_1(t)\cos\omega t + yf_2(t)\sin\omega t) \right]$$

$$f_1(t) = f(t) + f(t - \Delta)), f_2(t) = f(t) - f(t - \Delta))$$

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direction of

propagation

Z

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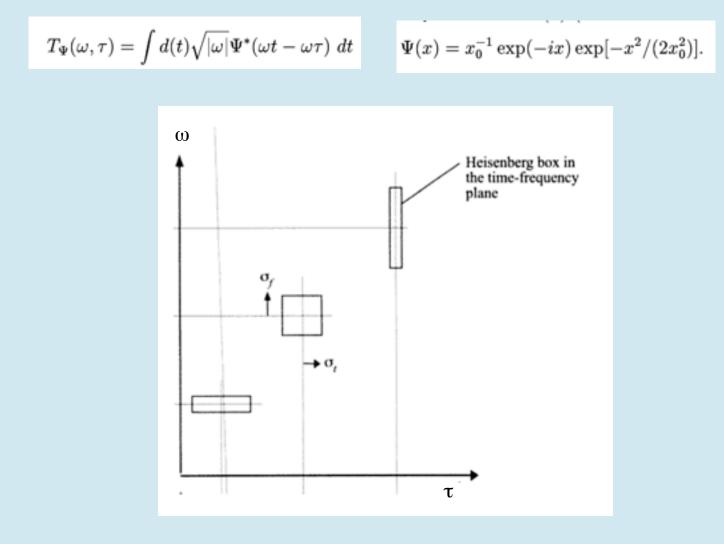
17

$$T_{\Psi}(\omega,\tau) = \int d(t) \sqrt{|\omega|} \Psi^*(\omega t - \omega \tau) \ dt$$

 $\Psi(x)=x_0^{-1}\exp(-ix)\exp[-x^2/(2x_0^2)].$



17

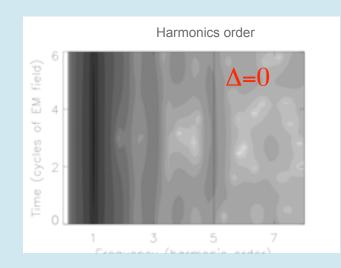


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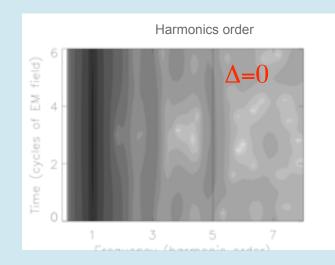


 $\Psi(x) = x_0^{-1} \exp(-ix) \exp[-x^2/(2x_0^2)].$

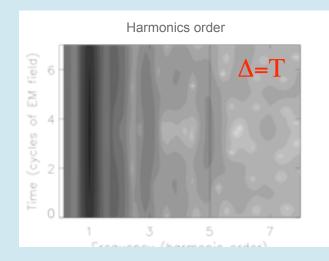
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Rb
$$F = 0.002 \ a.u. \ \omega = 0.2 \ eV$$

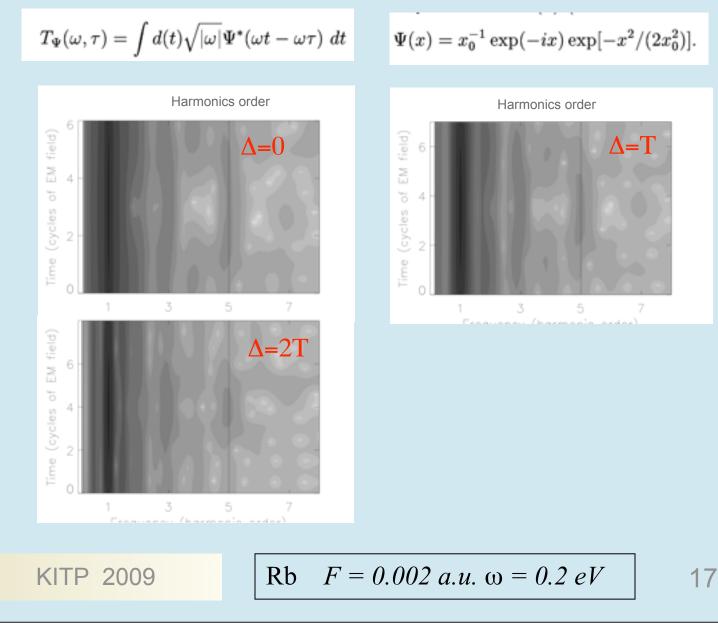
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Harmonics order
$$\int Harmonics order$$

$$\int A=0$$

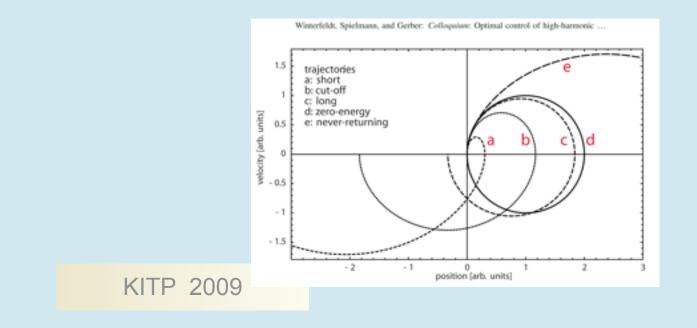
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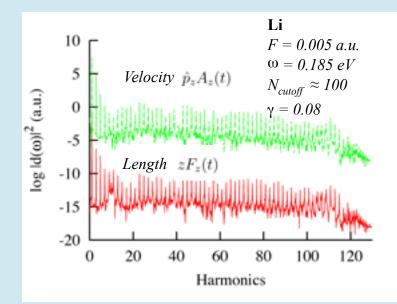
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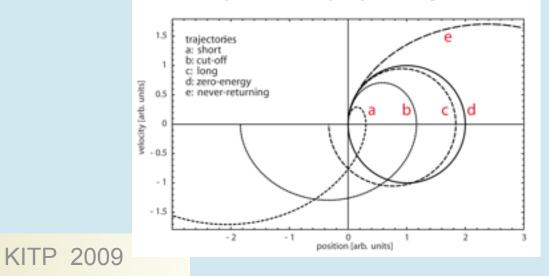
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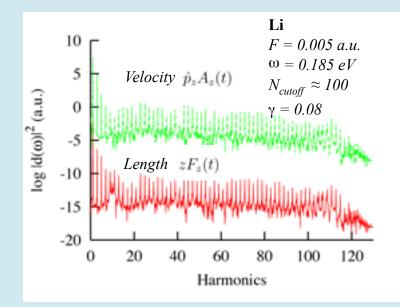
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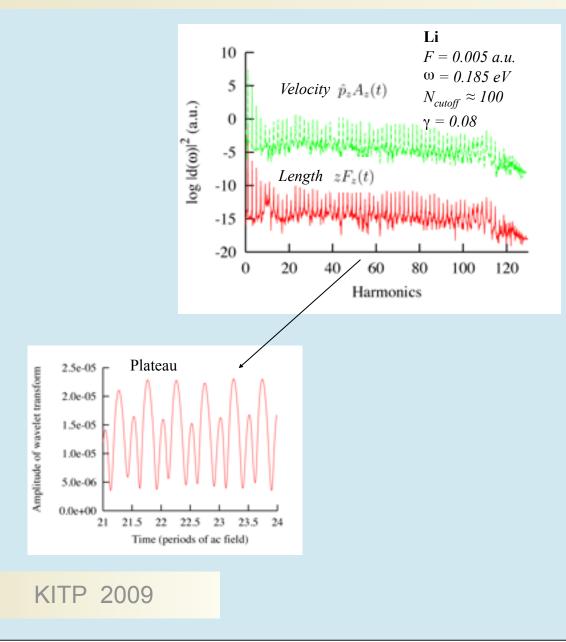


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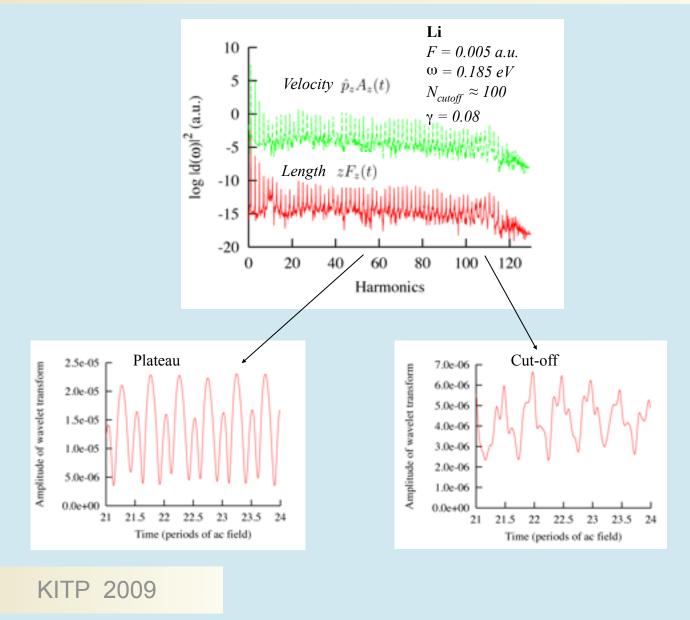




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Juan J. Carrera and Shih-I Chu PHYSICAL REVIEW A 75, 033807 (2007)

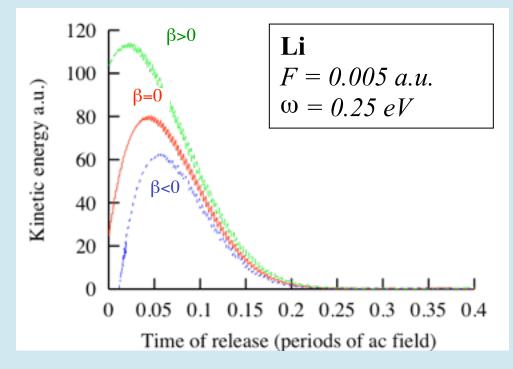
$$\hat{V}(\mathbf{r},t) = -\mathbf{F} \cdot \mathbf{r} E(t) = -Fzf(t)\cos[\omega t + \delta(t)],$$

$$\delta(t) = -\beta \tanh\left(\frac{t-t_0}{\tau}\right).$$

$$\omega(t) = \omega + d\delta(t)/dt = \omega - \frac{\beta}{\tau} \frac{1}{\cosh^2\left(\frac{t-t_0}{\tau}\right)}.$$

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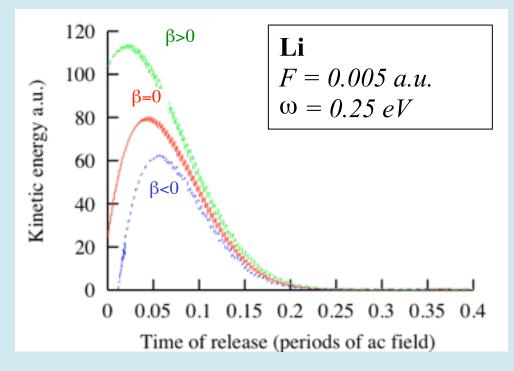
Juan J. Carrera and Shih-I Chu PHYSICAL REVIEW A 75, 033807 (2007)

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Juan J. Carrera and Shih-I Chu PHYSICAL REVIEW A 75, 033807 (2007)

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Kinetic energy at the moment of return as function of the time of release field without chirp (red), and positive and negative chirps

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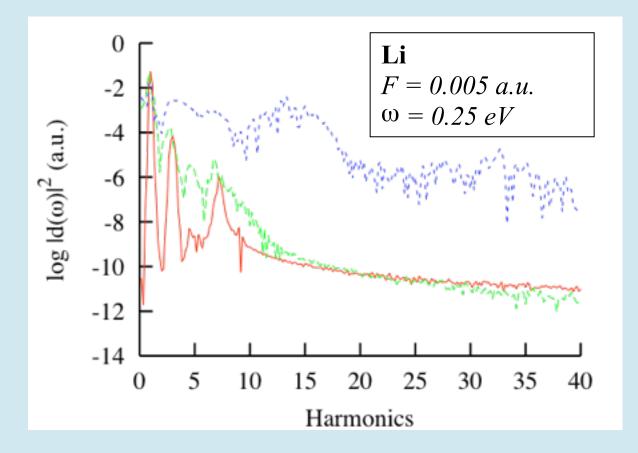
Li

$$F = 0.005 \ a.u.$$

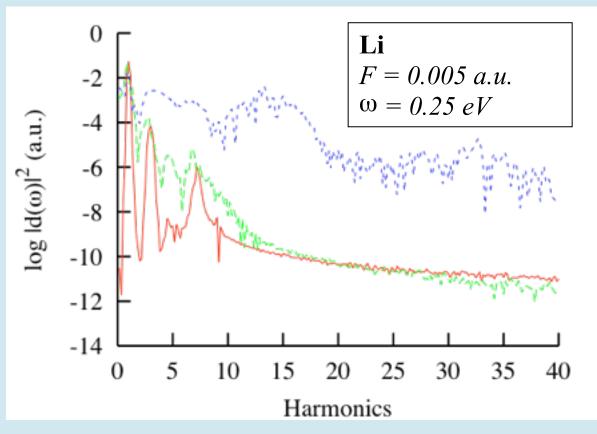
 $\omega = 0.25 \ eV$



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Three calculations for Li for a short pulse with increasing chirp (red, green blue progressively)

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Basis-based solution of TDSE applied to



Basis-based solution of TDSE applied to

- Ionization
 - (2γ,2e) on He
 - (4γ,e) on Li



• Basis-based solution of TDSE applied to

- Ionization
 - (2γ,2e) on He
 - (4γ,e) on Li
- HHG
 - Li $N_{cut-off} \approx 10$ optical regime
 - Li, Rb N_{cut-off} ≈100 XUV regime
 - Resonant enhancement
 - Polarization gating
 - Optimal control
- Further directions
 - Multi-photon many-electron processes on complex targets
 - More HHG control

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