

Xiuqi  
Ma



Wilbur  
Shirley



Meng  
Cheng



Michael  
Levin



John  
McGreevy



Ho Tat  
Lam

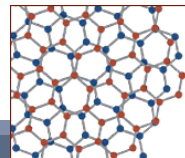
Phys. Rev. B 105, 195124 (2022)  
arXiv:2211.10458

# Fracton and Chern-Simons Theory

XIE CHEN, CALTECH  
KITP, APR. 2023



INSTITUTE FOR QUANTUM INFORMATION AND MATTER

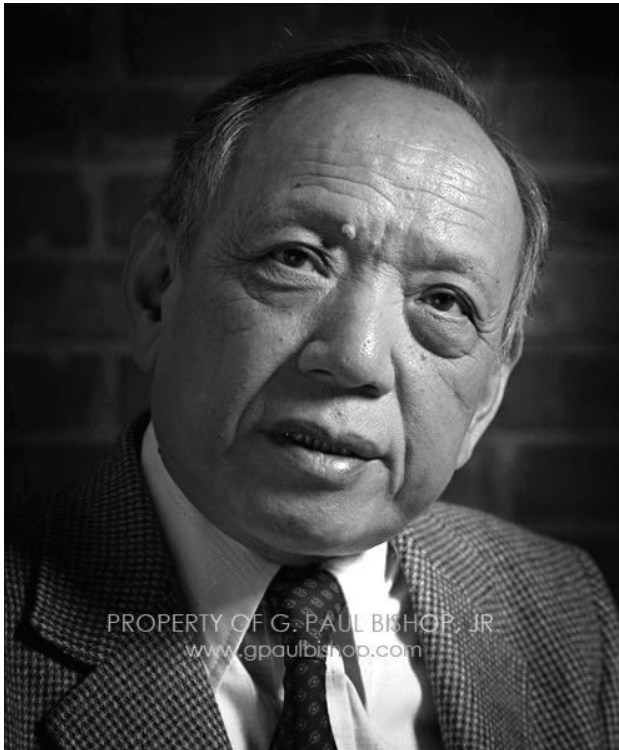


Simons Collaboration on  
Ultra-Quantum Matter

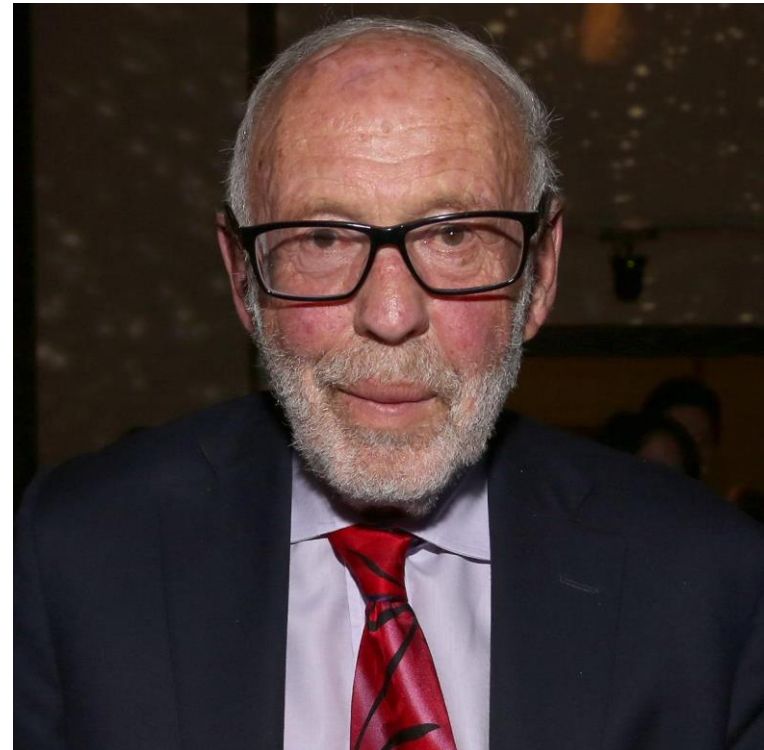


# Chern-Simons Theory

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Shiing-Shen Chern



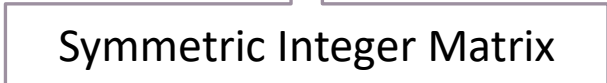
James Simons

# Chern-Simons Theory

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2+1D

$$\mathcal{L} = \frac{1}{4\pi} A \wedge dA \quad \mathcal{L} = \frac{1}{4\pi} \sum_{IJ} K_{IJ} A^I \wedge dA^J$$



$K = 1$      Integer quantum Hall

$K = 3$      Fractional quantum Hall

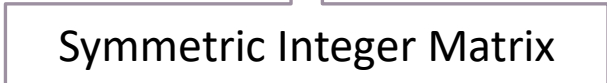
$K = \begin{pmatrix} 0 & 2 \\ 2 & 0 \end{pmatrix}$       $\mathbb{Z}_2$  gauge theory

# Chern-Simons Theory

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2+1D

$$\mathcal{L} = \frac{1}{4\pi} A \wedge dA \quad \mathcal{L} = \frac{1}{4\pi} \sum_{IJ} K_{IJ} A^I \wedge dA^J$$



$\det(K)$       Ground state degeneracy on torus

$K^{-1}$       Fractional statistics, fractional charge

# Chern-Simons Theory

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2+1D

$$\mathcal{L} = \frac{1}{4\pi} A \wedge dA \quad \mathcal{L} = \frac{1}{4\pi} \sum_{IJ} K_{IJ} A^I \wedge dA^J$$

Symmetric Integer Matrix

$$K = 1 \quad \text{Integer quantum Hall} \quad K^{-1} = 1$$

$$K = 3 \quad \text{Fractional quantum Hall} \quad K^{-1} = 1/3$$

$$K = \begin{pmatrix} 0 & 2 \\ 2 & 0 \end{pmatrix} \quad \text{Z}_2 \text{ gauge theory} \quad K^{-1} = \begin{pmatrix} 0 & 1/2 \\ 1/2 & 0 \end{pmatrix}$$

# Chern-Simons and Fracton

---

$$\mathcal{L} = \frac{1}{4\pi} \sum_{IJ} K_{IJ} A^I \wedge dA^J$$

**3+1D**

$$I, J = \dots - 1, 0, 1, 2, \dots$$

# Chern-Simons and Fracton

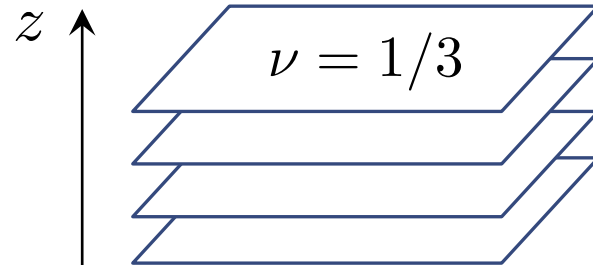
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$$\mathcal{L} = \frac{1}{4\pi} \sum_{IJ} K_{IJ} A^I \wedge dA^J$$

**3+1D**

$$I, J = \dots - 1, 0, 1, 2, \dots$$

$$K = \begin{pmatrix} \ddots & & & & & \\ & 3 & & & & \\ & & 3 & & & \\ & & & 3 & & \\ & & & & \ddots & \end{pmatrix}$$



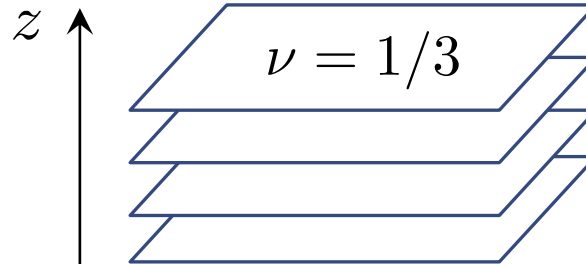
# Chern-Simons and Fracton

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$$\mathcal{L} = \frac{1}{4\pi} \sum_{IJ} K_{IJ} A^I \wedge dA^J$$

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$$K = \begin{pmatrix} \ddots & & & & & \\ & 3 & & & & \\ & & 3 & & & \\ & & & 3 & & \\ & & & & \ddots & \end{pmatrix}$$



- Ground state degeneracy  $3^N$  exponential in height in  $z$
- Fractional point excitations move in  $xy$  planes only – planons

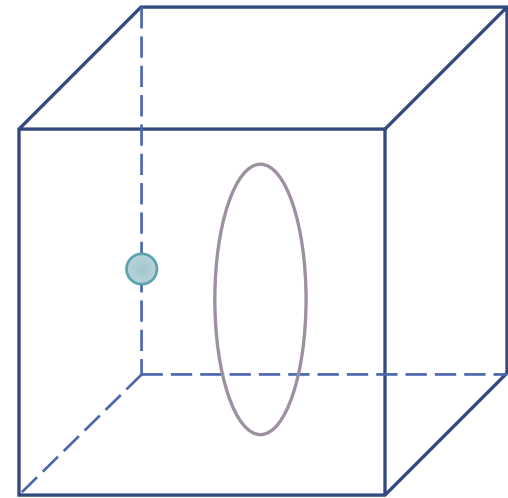


# Compared to 3+1D topo Order

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$Z_2$  gauge theory  $\mathcal{L} = \frac{2}{4\pi} \epsilon^{\mu\nu\rho\sigma} B_{\mu\nu} \partial_\rho A_\sigma$

- Gapped
- Fractional point excitation
- Full motion in 3D space
- Fractional loop excitation
- Ground state degeneracy = 8



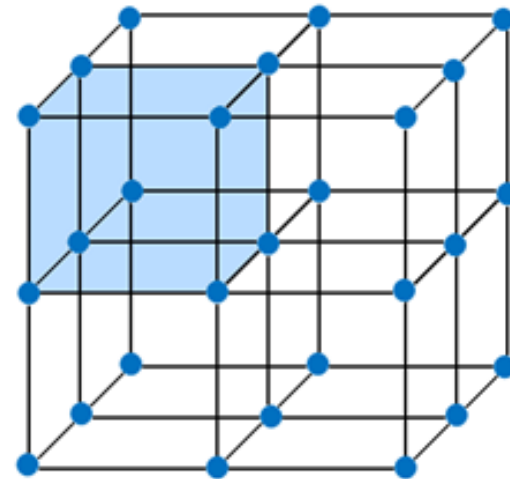
# Compared to X-cube

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$$H = - \sum_c \text{X-cube} - \sum_v \text{Z-cross} - \sum_v \text{Z-cross} - \sum_v \text{Z-cross}$$

$$X = |0\rangle\langle 1| + |1\rangle\langle 0|$$

$$Z = |0\rangle\langle 0| - |1\rangle\langle 1|$$

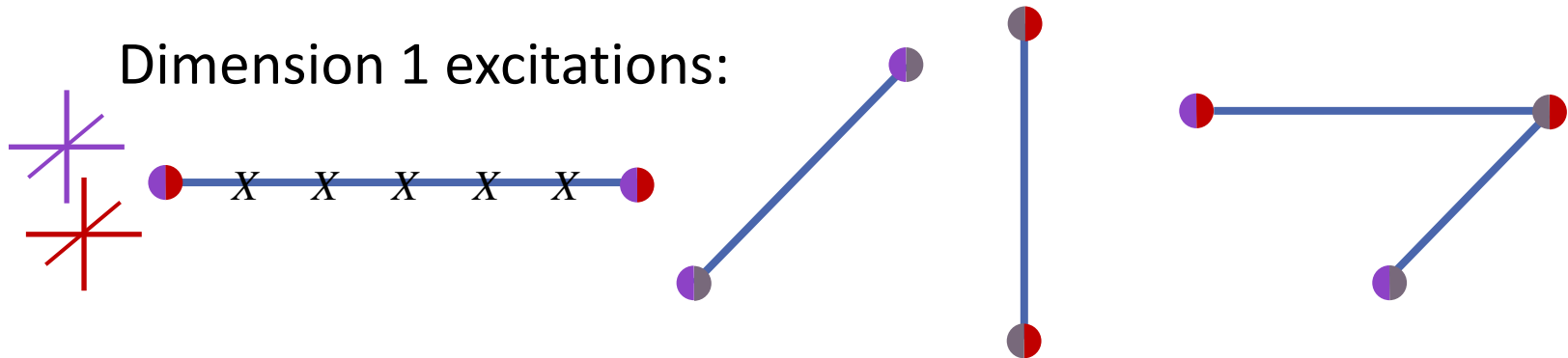


# X-cube

$$H = - \sum_c \sum_{\text{cube}} X - \sum_v \sum_{\text{red}} Z - \sum_v \sum_{\text{purple}} Z - \sum_v \sum_{\text{green}} Z$$

Log(Ground state degeneracy) = 6L-3

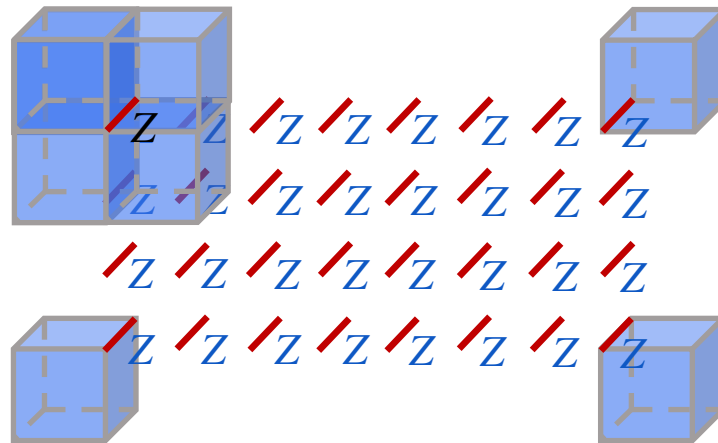
Dimension 1 excitations:



# X-cube

$$H = - \sum_c \sum_{\text{cube}} X - \sum_v \sum_{\text{red}} Z - \sum_v \sum_{\text{purple}} Z - \sum_v \sum_{\text{green}} Z$$

Fracton Excitations

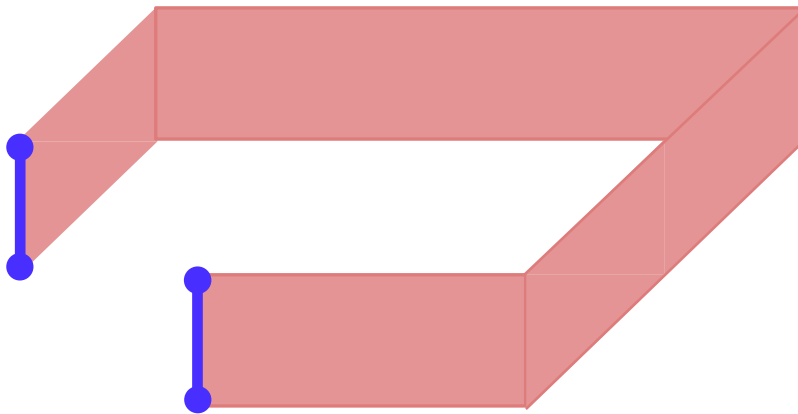


# Compare to X-cube

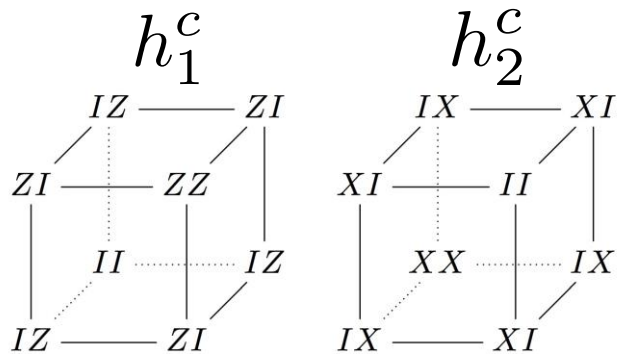
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$$H = - \sum_c \text{[cube with X's]} - \sum_v \text{[red star with Z's]} - \sum_v \text{[purple star with Z's]} - \sum_v \text{[green star with Z's]}$$

Planon excitations:



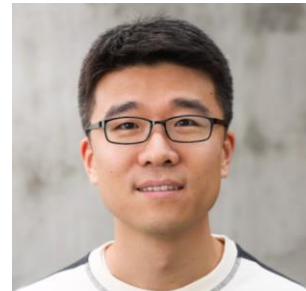
# Compared to Haah's code



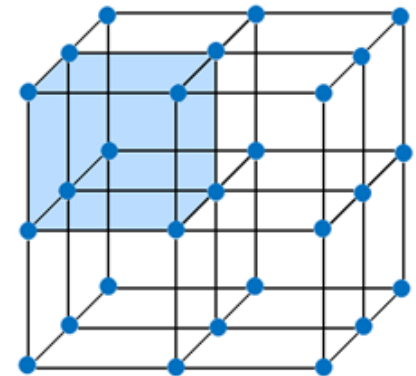
$$X = |0\rangle\langle 1| + |1\rangle\langle 0|$$

$$Z = |0\rangle\langle 0| - |1\rangle\langle 1|$$

$$H = - \sum_c h_1^c - \sum_c h_2^c$$

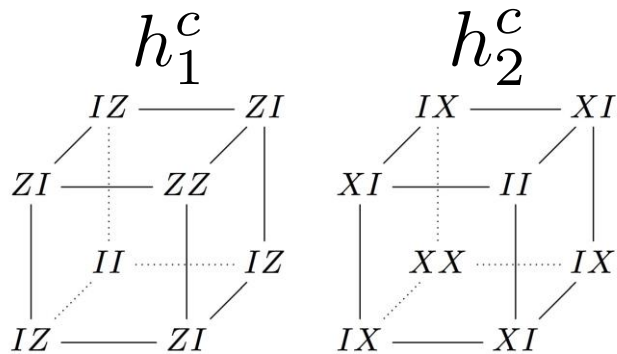


Jeongwan Haah  
2011



- Cubic lattice model
- Two qubits per site
- All Hamiltonian terms commute

# Compared to Haah's code



Stable ground state degeneracy

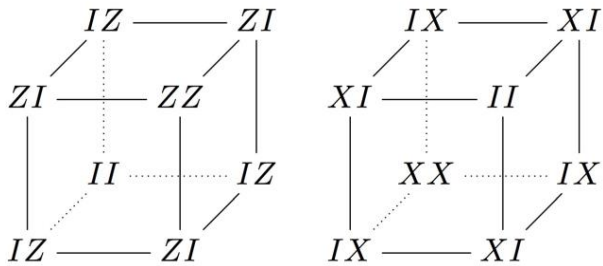
$$\log \text{GSD} =$$

$$\frac{k+2}{4} = \deg_x \text{gcd} \left[ \begin{array}{l} 1 + (1+x)^L, \\ 1 + (1+\omega x)^L, \\ 1 + (1+\omega^2 x)^L \end{array} \right]_{\mathbb{F}_4}$$

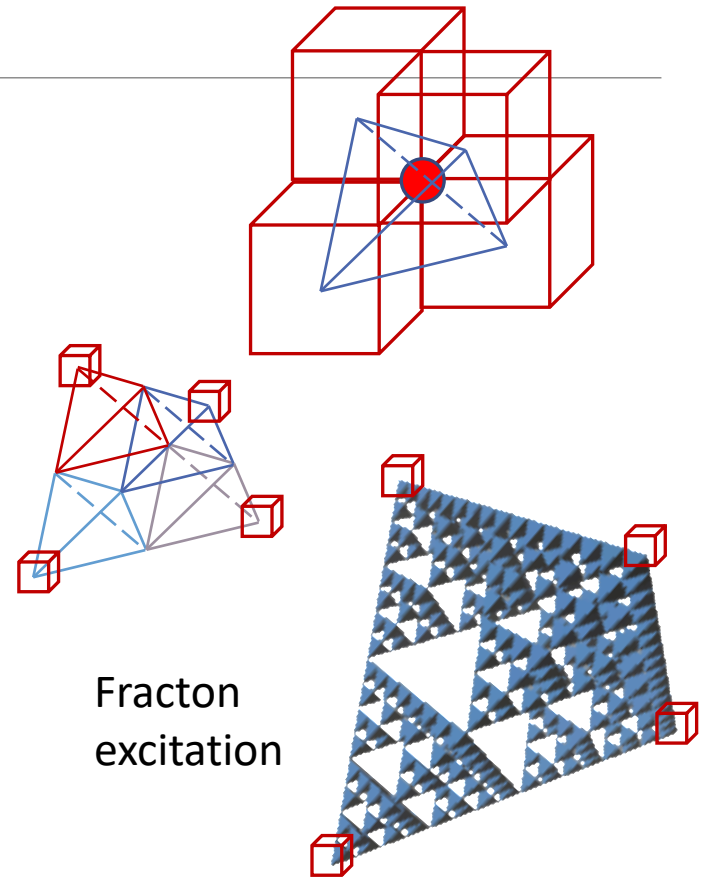
$$= \begin{cases} 1 & \text{if } L = 2^p + 1 \ (p \geq 1), \\ L & \text{if } L = 2^p \ (p \geq 1) \end{cases}$$

$$H = - \sum_c h_1^c - \sum_c h_2^c$$

# Haah's code



- Cube excitations created four at a time
- Individual cube excitation cannot move
- Cube excitations separated through a fractal process





# Fracton

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## Type I

X-cube

- Some fractional point excitations can move
- No fractal structure

- exponential growth of ground state degeneracy
- fractional point excitation with restricted motion

## Type II

Haah's code

- None of the fractional point excitations can move
- Fractal structure

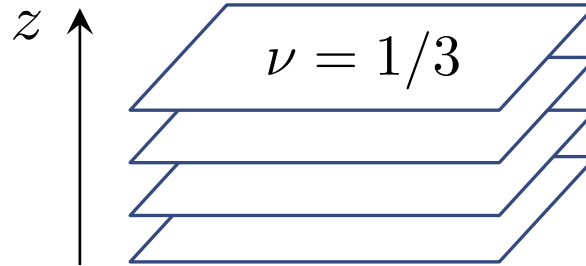
# Chern-Simons and Fracton

---

$$\mathcal{L} = \frac{1}{4\pi} \sum_{IJ} K_{IJ} A^I \wedge dA^J$$

$$I, J = \dots - 1, 0, 1, 2, \dots$$

$$K = \begin{pmatrix} \ddots & & & & & \\ & 3 & & & & \\ & & 3 & & & \\ & & & 3 & & \\ & & & & \ddots & \\ & & & & & \ddots \end{pmatrix}$$



- Ground state degeneracy  $3^N$  exponential in height in  $z$
- Fractional point excitations move in  $xy$  planes only – planons

# Fracton

---

Type I

X-cube

diagonal CS

Type II

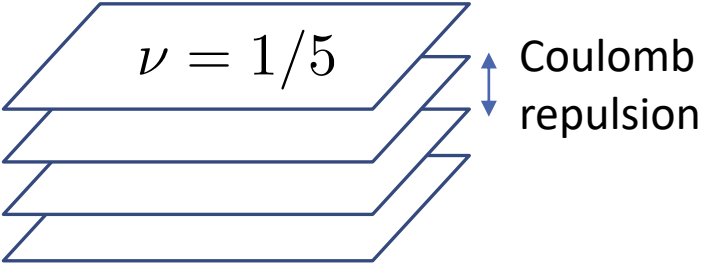
Haah's  
code

# Tri-diagonal K matrix

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$$K(131) = \begin{pmatrix} 3 & 1 & & & 1 \\ 1 & 3 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & 1 & 3 & 1 \\ 1 & & & 1 & 3 \end{pmatrix}_{N \times N}$$

$z$  ↑



$\nu = 1/5$

Coulomb repulsion

The diagram shows a stack of four horizontal layers. The top layer is labeled with  $\nu = 1/5$ . A vertical arrow labeled  $z$  points upwards from the layers. A double-headed vertical arrow to the right of the layers is labeled "Coulomb repulsion".

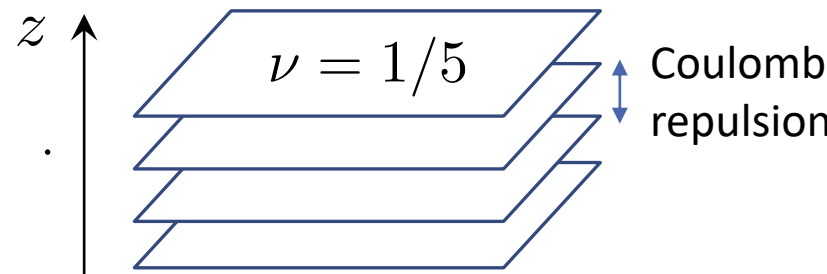
# Tri-diagonal K matrix

---

$$K(131) = \begin{pmatrix} 3 & 1 & & & & & 1 \\ 1 & 3 & 1 & & & & \\ & 1 & 3 & 1 & & & \\ & & 1 & 3 & 1 & & \\ & & & 1 & 3 & 1 & \\ & & & & 1 & 3 & 1 \\ 1 & & & & & 1 & 3 \end{pmatrix}$$

$$K(131)^{-1} = \frac{1}{65} \begin{pmatrix} 29 & -11 & 4 & -1 & -1 & 4 & -11 \\ -11 & 29 & -11 & 4 & -1 & -1 & 4 \\ 4 & -11 & 29 & -11 & 4 & -1 & -1 \\ -1 & 4 & -11 & 29 & -11 & 4 & -1 \\ -1 & -1 & 4 & -11 & 29 & -11 & 4 \\ 4 & -1 & -1 & 4 & -11 & 29 & -11 \\ -11 & 4 & -1 & -1 & 54 & -11 & 29 \end{pmatrix}$$

# Tri-diagonal K matrix

$$K(131) = \begin{pmatrix} 3 & 1 & & & 1 \\ 1 & 3 & 1 & & \\ & \ddots & \ddots & \ddots & \\ & & 1 & 3 & 1 \\ 1 & & & 1 & 3 \end{pmatrix}_{N \times N}$$


$N \rightarrow \infty$  Irrational statistics

$$\theta_{IJ} = \frac{1}{\sqrt{5}} \left( \frac{\sqrt{5} - 3}{2} \right)^{|I-J|}$$

Ground state degeneracy

$$\left( \frac{3+\sqrt{5}}{2} \right)^N + \left( \frac{3-\sqrt{5}}{2} \right)^N - 2(-1)^N \sim \left( \frac{3+\sqrt{5}}{2} \right)^N$$

Planon Fusion group  $G_N = \mathbb{Z}_{F_N} \times \mathbb{Z}_{5F_N}$   $F_N$  Nth number in the Fibonacci sequence

# Fracton

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Type I

X-cube

Diagonal CS

tri-Diagonal CS

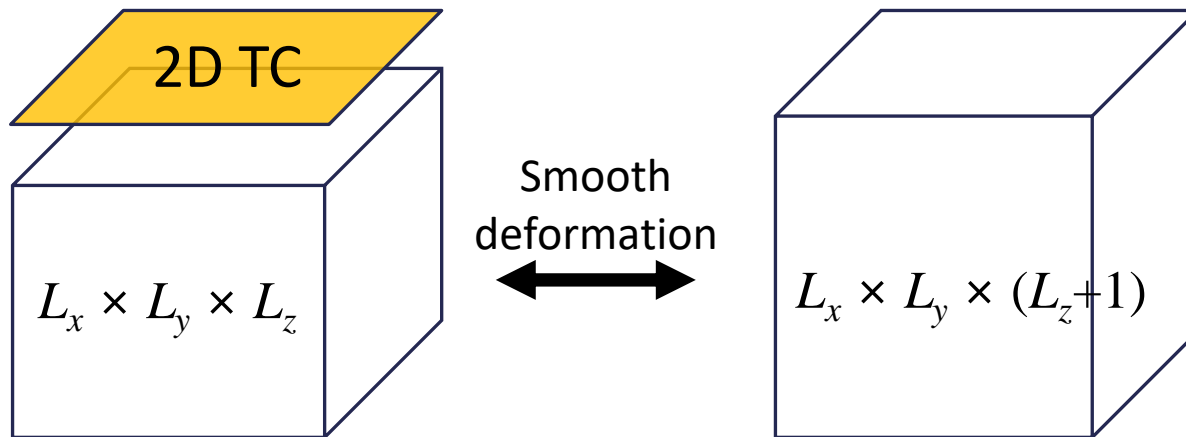
Type II

Haah's  
code

# X-cube and foliation

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## Foliation

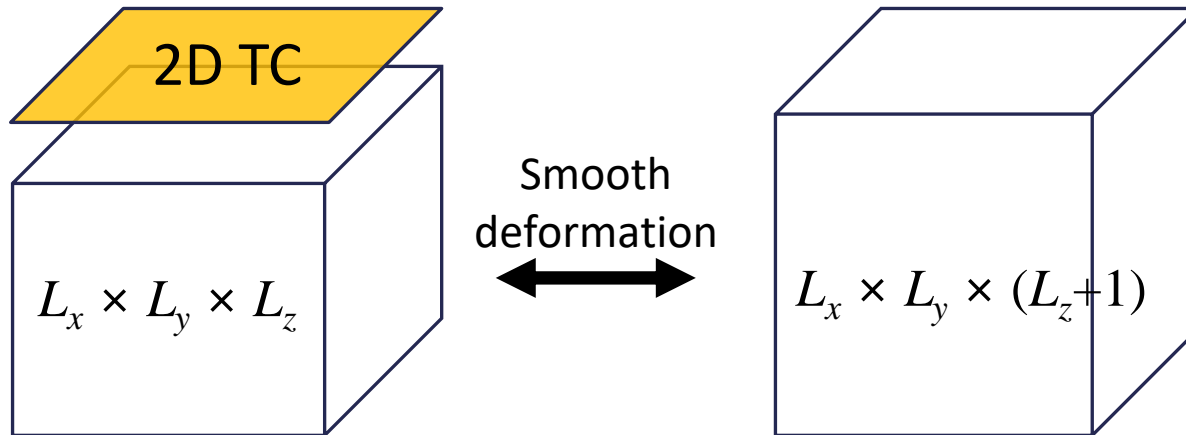




# X-cube and foliation

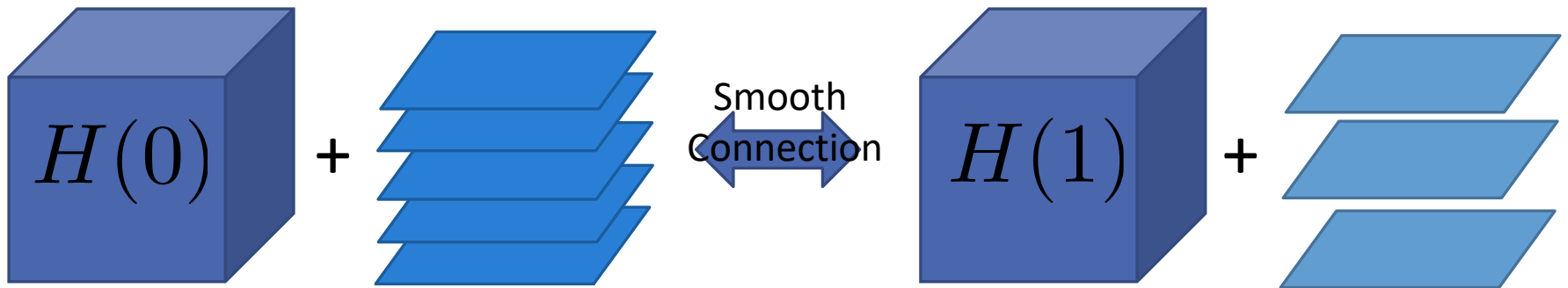
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## Foliation



- Exponential scaling of ground state degeneracy
- Planon excitation

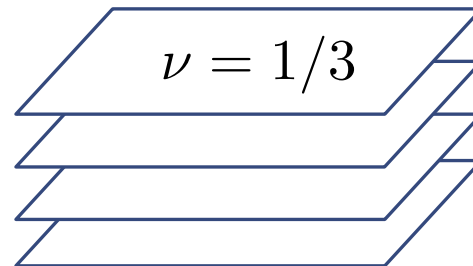
# Foliated Fracton Phase



Comparing different models: mod contribution from 2D layers  
Excitations, entanglement ...

$$K = \begin{pmatrix} \ddots & & & & \\ & 3 & & & \\ & & 3 & & \\ & & & 3 & \\ & & & & \ddots \end{pmatrix}$$

$z$  ↑



A trivial foliated  
fracton model

# Foliated K matrix

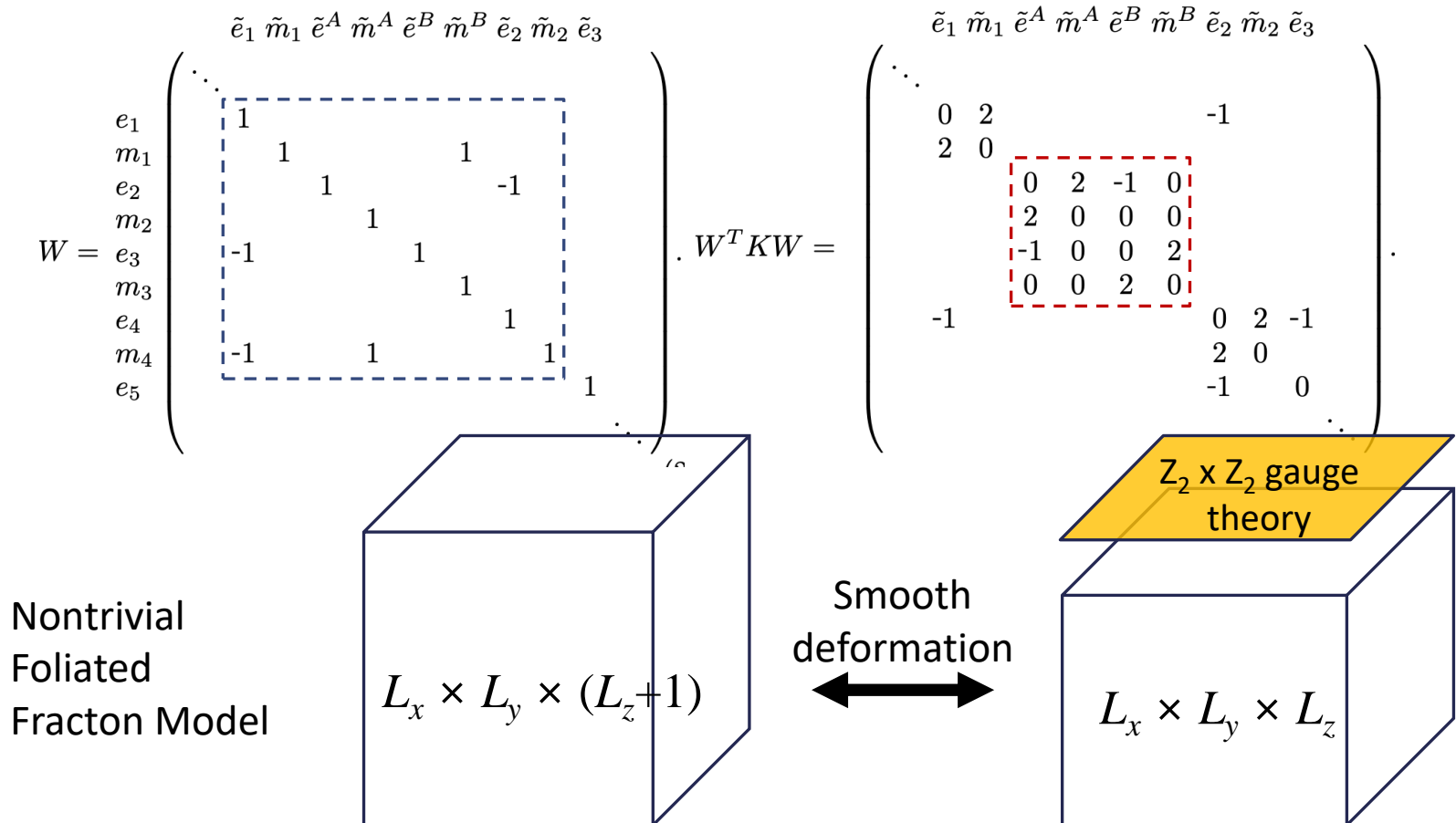
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$$K = \begin{pmatrix} \ddots & & & & & & & & \\ & 0 & 2 & -1 & & & & & \\ & 2 & 0 & & & & & & \\ & -1 & & 0 & 2 & -1 & & & \\ & & & 2 & 0 & & & & \\ & & & -1 & & 0 & 2 & -1 & \\ & & & & & 2 & 0 & & \\ & & & & & -1 & & 0 & \\ & & & & & & & \ddots & \end{pmatrix}.$$

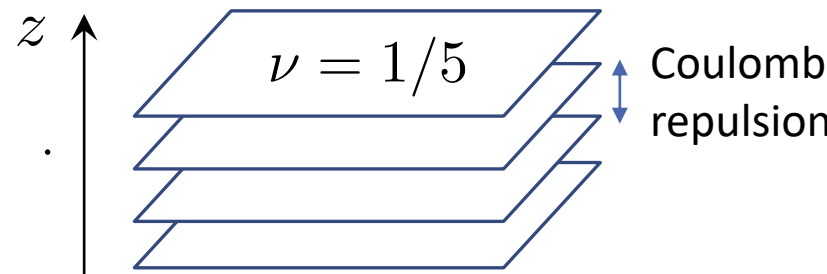
$$K^{-1} = \frac{1}{4} \begin{pmatrix} \ddots & & & & & & & & \\ & 0 & & 1 & & & & & \\ & & 0 & 2 & & & & & \\ & 1 & 2 & 0 & & 1 & & & \\ & & & & 0 & 2 & & & \\ & & & 1 & 2 & 0 & & 1 & \\ & & & & & & 0 & 2 & \\ & & & & & & 1 & 2 & 0 & \\ & & & & & & & \ddots & \end{pmatrix}.$$



# Foliated K matrix



# Tri-diagonal K matrix

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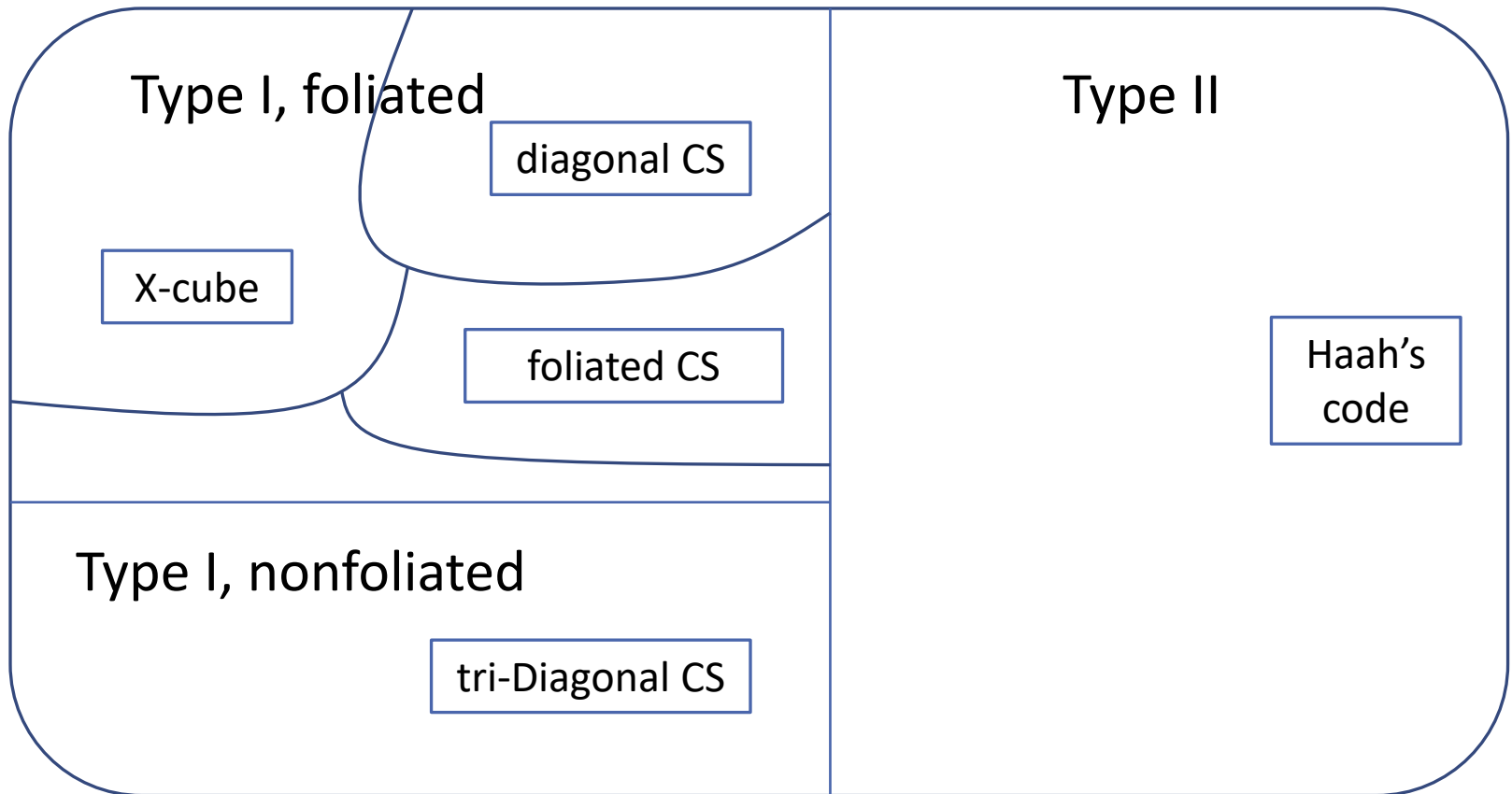
Ground state degeneracy

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Planon Fusion group  $G_N = \mathbb{Z}_{F_N} \times \mathbb{Z}_{5F_N}$   $F_N$  Nth number in the Fibonacci sequence

# Fracton

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# Spectrum

---

$$\mathcal{L} = -\frac{1}{4e^2} \sum_I F_{\mu\nu}^I F^{I,\mu\nu} + \frac{1}{4\pi} \sum_{IJ} K_{IJ} \epsilon^{\mu\nu\lambda} A_\mu^I \partial_\nu A_\lambda^J$$

Gap of each mode is proportional to |eigenvalues of K|^2



# Spectrum

---

$$\mathcal{L} = -\frac{1}{4e^2} \sum_I F_{\mu\nu}^I F^{I,\mu\nu} + \frac{1}{4\pi} \sum_{IJ} K_{IJ} \epsilon^{\mu\nu\lambda} A_\mu^I \partial_\nu A_\lambda^J$$

Gap of each mode is proportional to |eigenvalues of K|^2

Gapped:  $K(131), K(141), K(151)$

Gapless:  $K(101), K(111), K(121)$

# Spectrum

---

$$\mathcal{L} = -\frac{1}{4e^2} \sum_I F_{\mu\nu}^I F^{I,\mu\nu} + \frac{1}{4\pi} \sum_{IJ} K_{IJ} \epsilon^{\mu\nu\lambda} A_\mu^I \partial_\nu A_\lambda^J$$

Gap of each mode is proportional to |eigenvalues of K|^2

Gapped:  $K(131), K(141), K(151)$

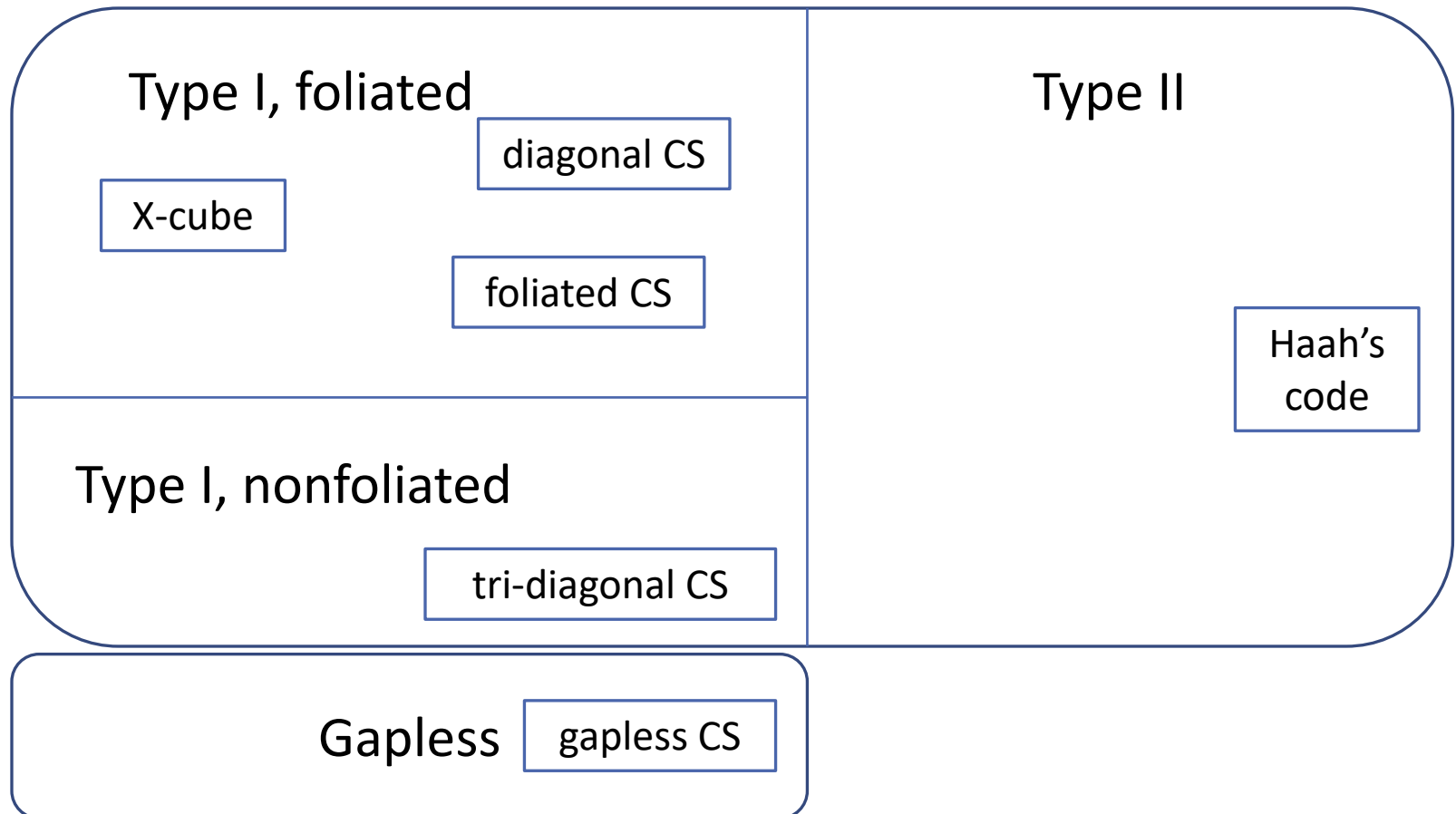
Gapless:  $K(101), K(111), K(121)$

Sullivan, Dua, Cheng, arXiv: 2109.13267

Lam, Ma, Chen, arXiv:2211.10458

# Fracton

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# Gapless 2+1D CS

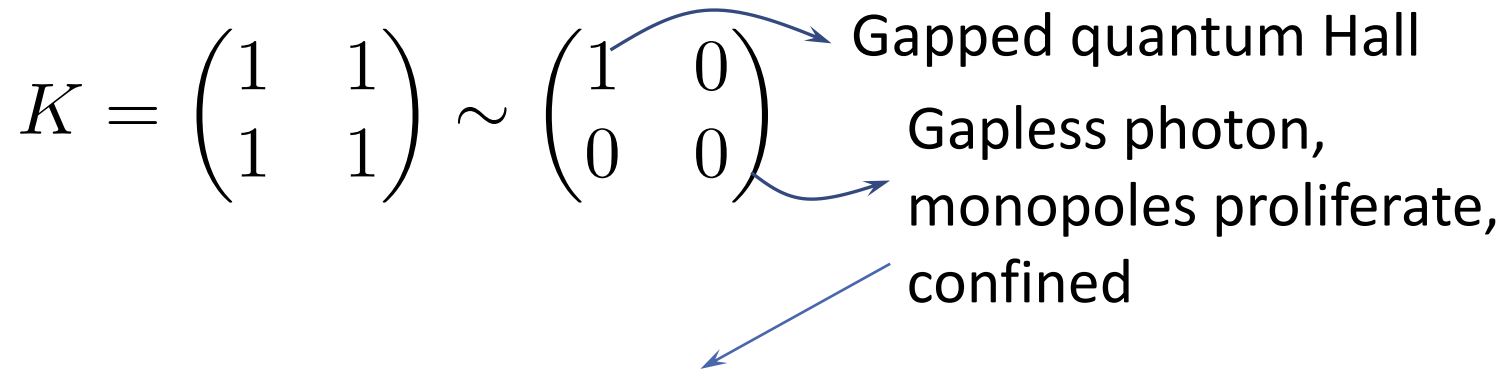
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$$\mathcal{L} = -\frac{1}{4e^2} \sum_I F_{\mu\nu}^I F^{I,\mu\nu} + \frac{1}{4\pi} \sum_{IJ} K_{IJ} \epsilon^{\mu\nu\lambda} A_\mu^I \partial_\nu A_\lambda^J$$

For example

$$K = \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \sim \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Gapped quantum Hall  
Gapless photon,  
monopoles proliferate,  
confined

The diagram shows the equivalence of two K matrices. The first matrix is a 2x2 matrix with all elements equal to 1. The second matrix is a 2x2 matrix with the top-left element equal to 1 and all other elements equal to 0. A blue arrow points from the top-right element of the second matrix to the text "Gapped quantum Hall". Another blue arrow points from the bottom-right element of the second matrix to the text "Gapless photon, monopoles proliferate, confined". A third blue arrow points from the bottom-left element of the second matrix to the text "Explicit U(1) symmetry breaking, Goldstone mode gapped".

Explicit U(1) symmetry breaking,  
Goldstone mode gapped

# Gapless infinite Chern Simons

---

K(121)

Contains mode with  
exactly zero energy  
 $k_x = k_y = 0, k_z = \pi$

Quadratic  
dispersion

K(232)

no mode with exactly  
zero energy

Linear  
dispersion

- Stability?
- Topological feature?

# Gapless infinite Chern Simons

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## □ Stability

no gauge invariant local monopole operators  
no relevant term to gap out the theory  
robust gapless-ness

## □ Deconfined point excitation

polynomially decaying potential

## □ Topological feature (linear dispersion)

finite braiding phase between point excitations

# Summary

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- Generalizing 2+1D Chern Simons theory to 3+1D by including an infinite number of 2+1D gauge fields
- Type I fracton
- Foliated
- Non-foliated
- Stable gapless phases
- Continuous field theory description (linear dispersion)

# Open questions

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- Conditions for gapped foliated / gapped non-foliated
- Equivalence classes of foliated models (as equivalence class of  $K$  matrices)
- General properties of gapped non-foliated models (degeneracy, fusion group, statistics, RG, etc.)
- Building block for 3-foliated models (like x-cube)?
- Experiment?