

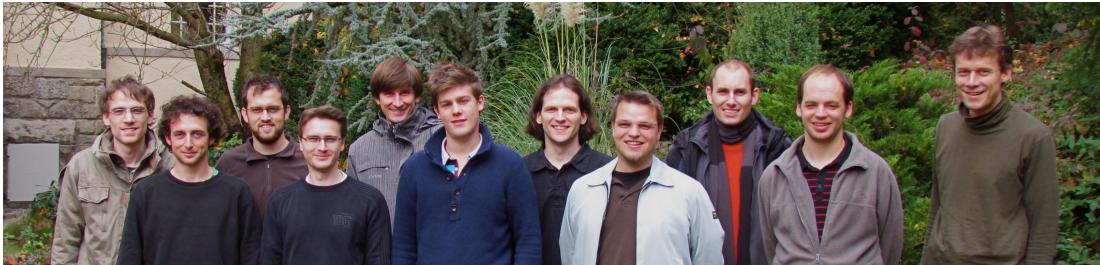
Topological Excitations, Superfluid Turbulence & Non-Thermal Fixed Points in Ultracold Gases

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GSI Helmholtzzentrum für Schwerionenforschung GmbH

Thanks & credits to...



...my work group in Heidelberg:

Sebastian Bock
Sebastian Erne
Martin Gärttner
Roman Hennig
Markus Karl
Steven Mathey
Boris Nowak
Nikolai Philipp
Dénes Sexty
Martin Trappe
Pascal Weckesser

...my former students:

Jan Schole (→ Heidelberg), Jan Zill (→ Queensland), Maximilian Schmidt (→ Jülich), Cédric Bodet (→ NEC), Alexander Branschädel (→ KIT Karlsruhe), Stefan Keßler (→ U Erlangen), Matthias Kronenwett (→ R. Berger), Christian Scheppach (→ Cambridge, UK), Philipp Struck (→ Konstanz), Kristan Temme (→ Vienna)

€€€...

Alexander von Humboldt
Stiftung / Foundation

Heisenberg-
Programm

DFG
Deutsche
Forschungsgemeinschaft



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LGFG BaWue

DAAD
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Dynamics**







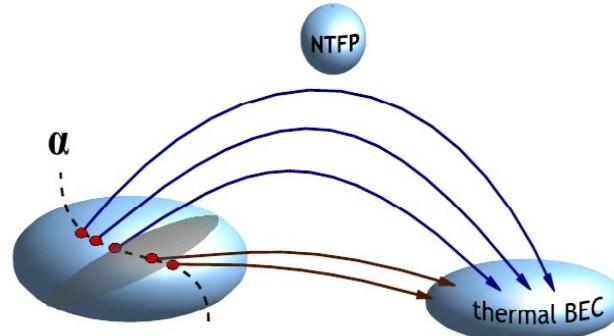
Equilibration



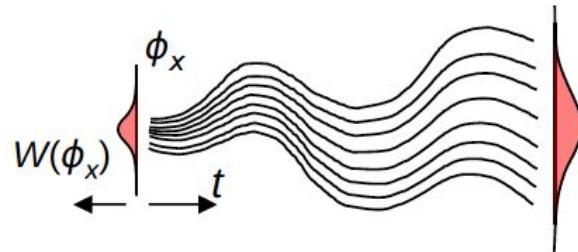
Initial state:
Far from equilibrium

Transient state,
e.g. Turbulence
Non-thermal fixed point

Final state:
Thermal equilibrium



Semi-classical Simulations



Classical field equation for $\phi(x,t)$:

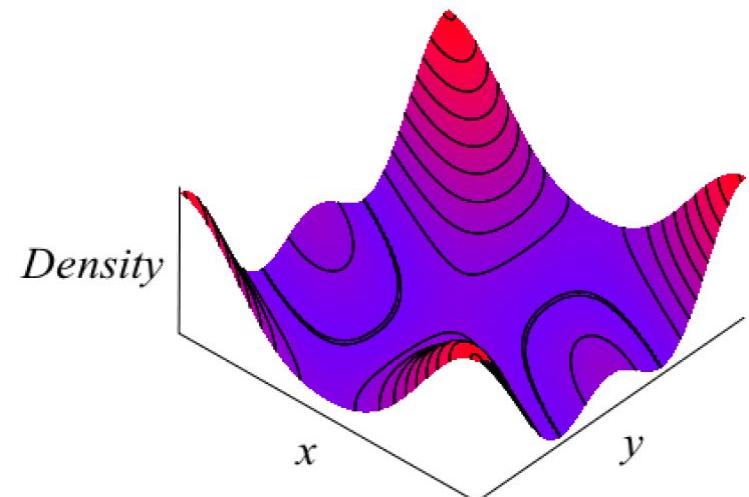
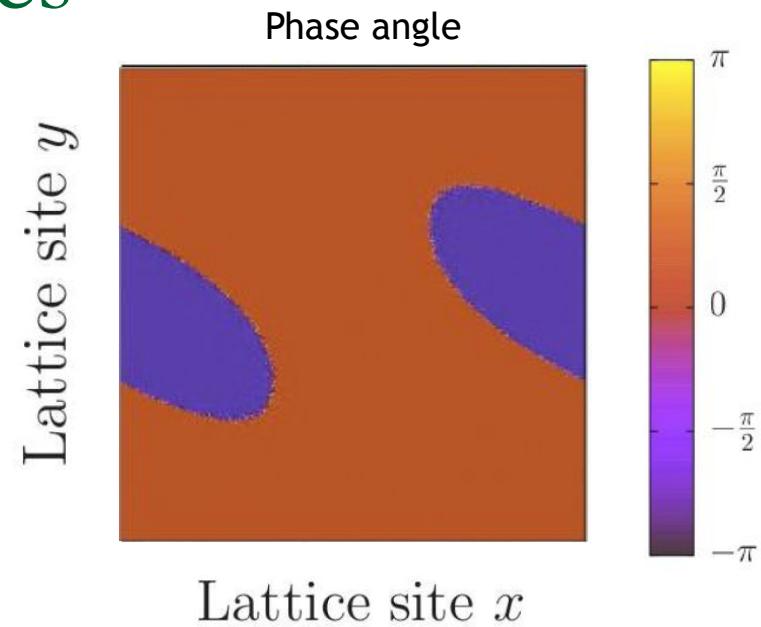
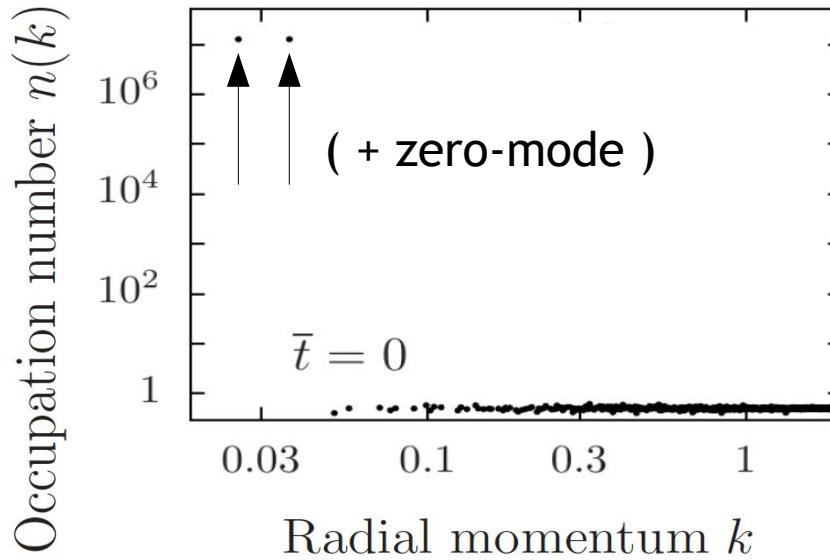
$$i\partial_t \phi(\mathbf{x}, t) = \left[-\frac{\nabla^2}{2m} + g|\phi(\mathbf{x}, t)|^2 \right] \phi(\mathbf{x}, t)$$

Observables: e. g. Momentum distribution

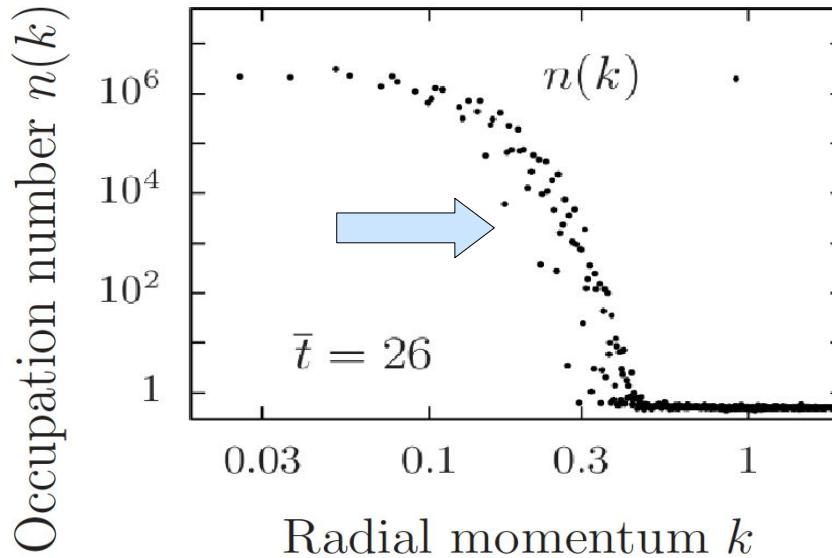
$$n(k) = \int d^{d-1}\Omega_k \langle \phi^*(\mathbf{k})\phi(\mathbf{k}) \rangle_{\text{ensemble}}$$



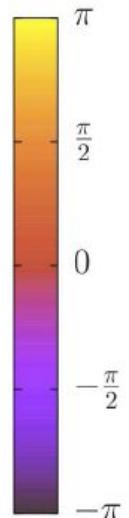
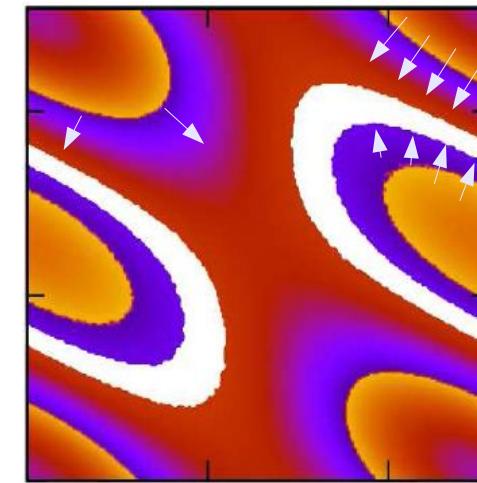
2+1 D: Quench dynamics



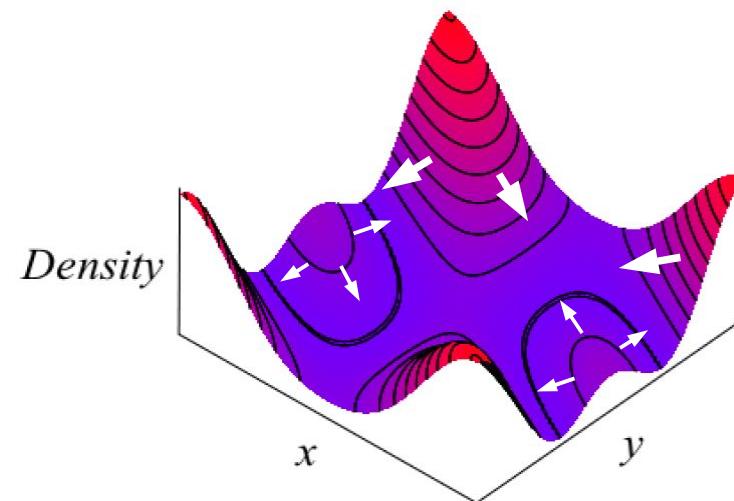
2+1 D: Quench dynamics



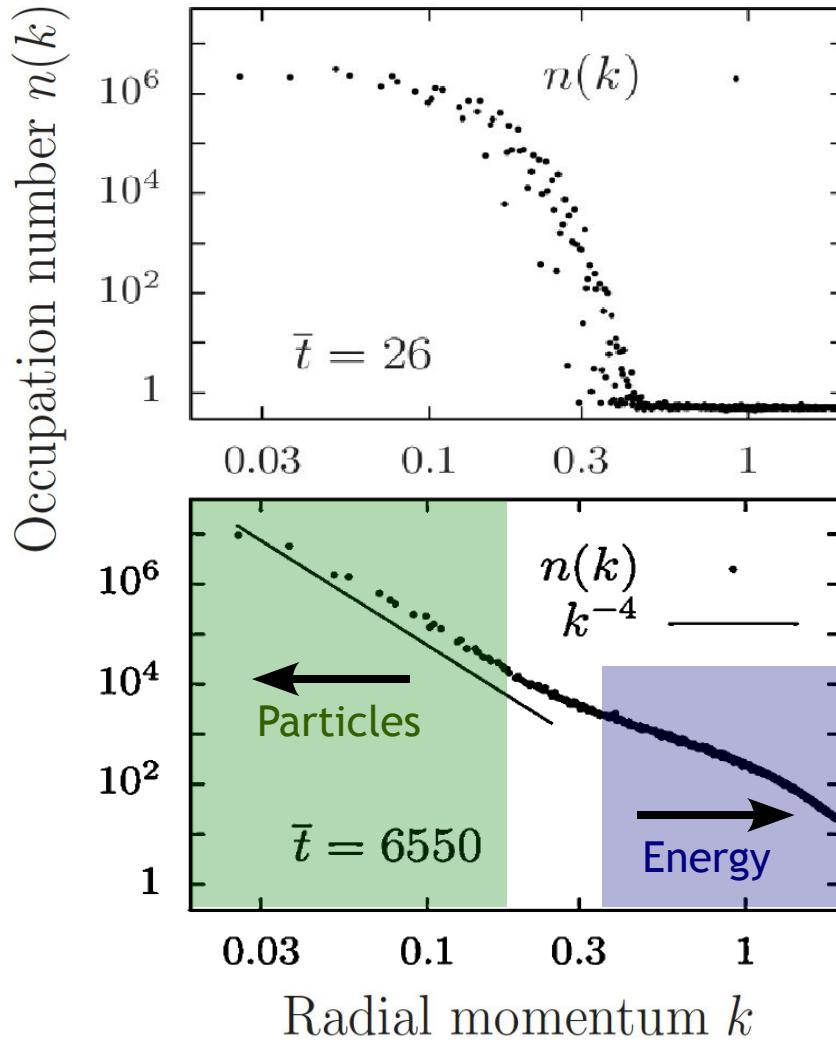
Lattice site y



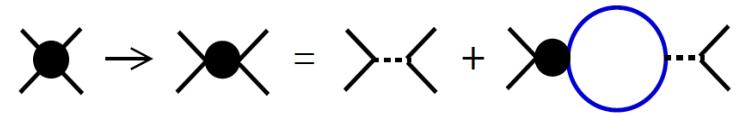
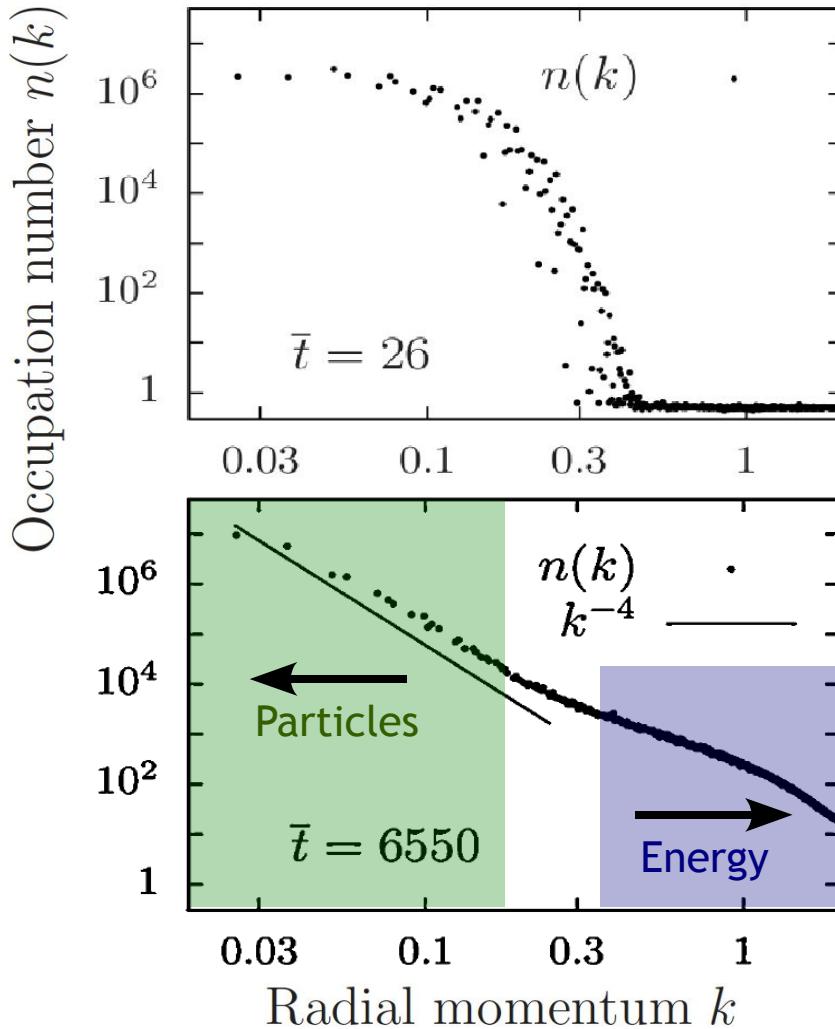
Lattice site x



2+1 D: Quench dynamics



2+1 D: Quench dynamics

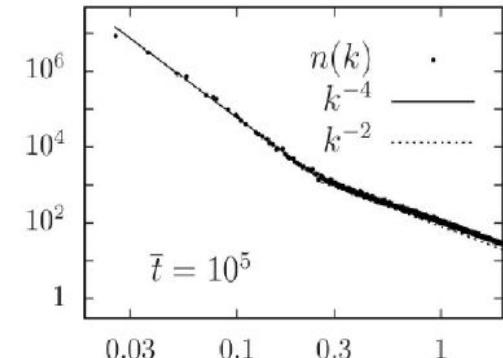
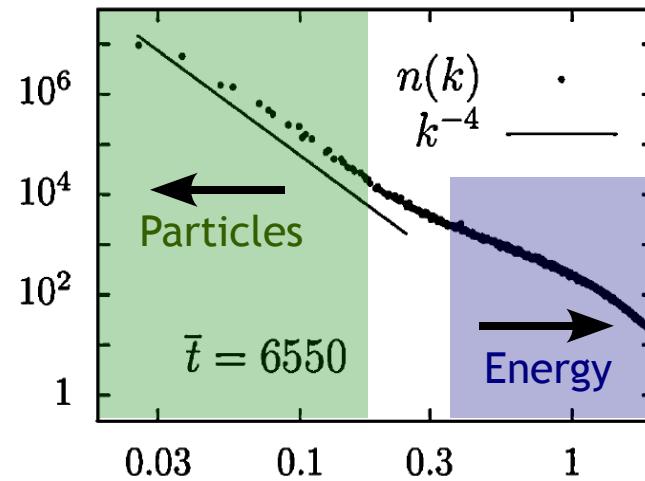
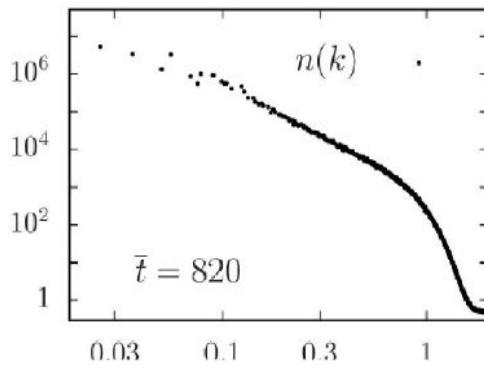


J. Berges, A. Rothkopf, J. Schmidt, PRL 101 (08) 041603,
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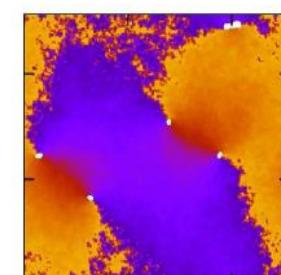
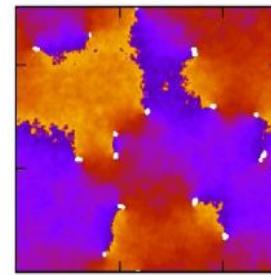
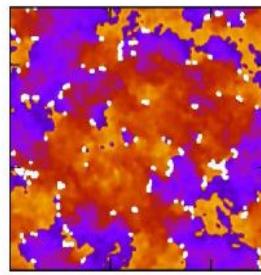
$$n^{\text{IR}} \sim k^{-d-2}$$



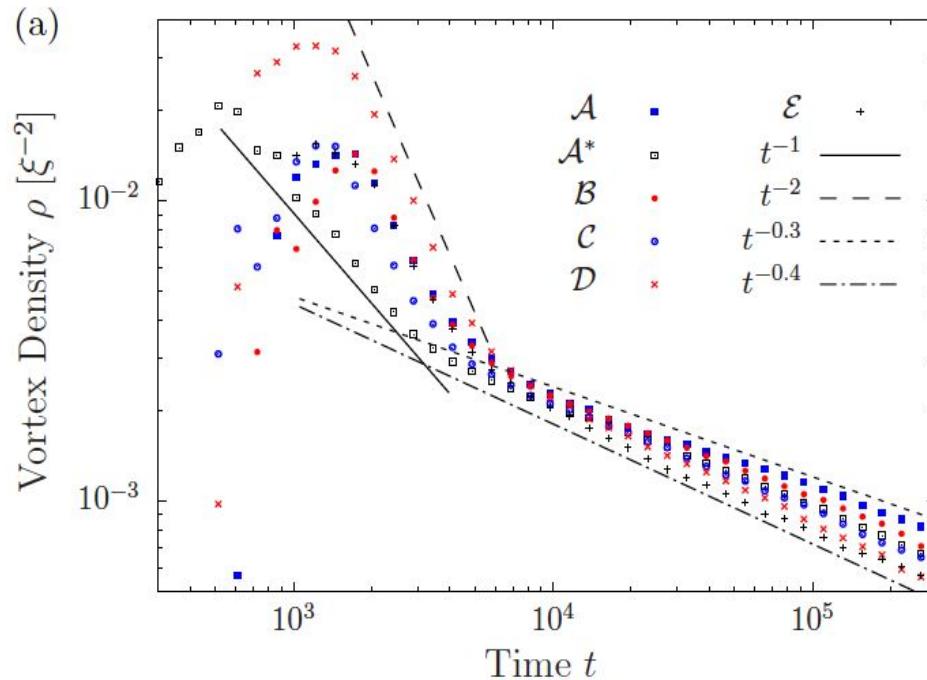
2+1 D: Phase ordering dynamics



Time →



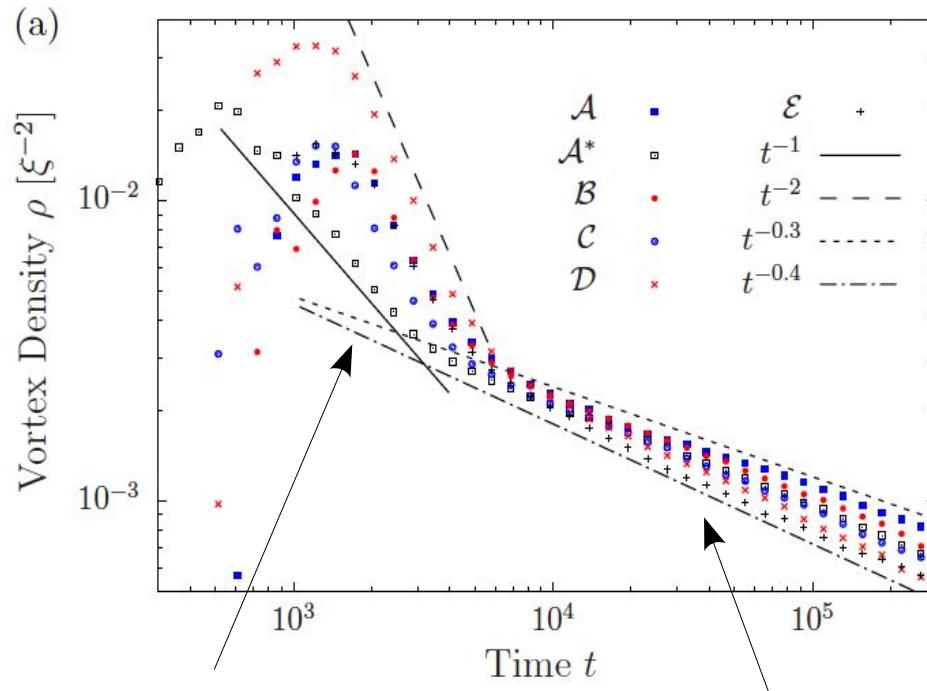
2+1 D: Phase ordering dynamics



J. Schole, B. Nowak, TG, PRA 86 (12) 013624

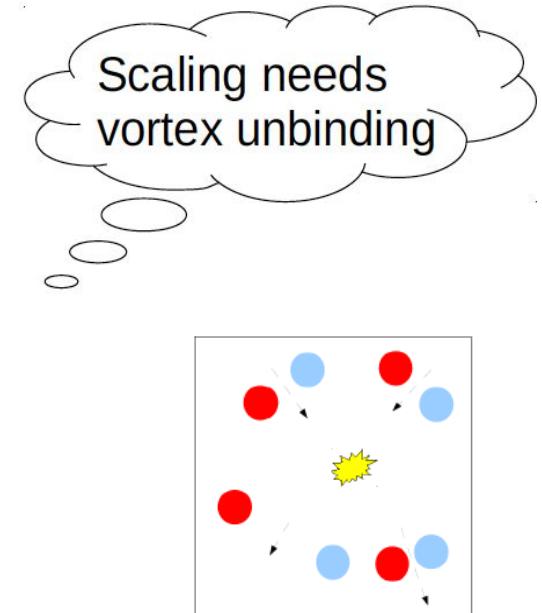


2+1 D: Phase ordering dynamics



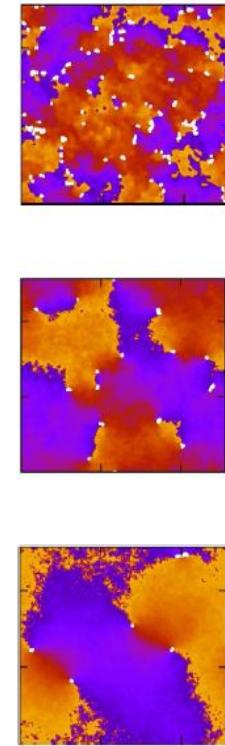
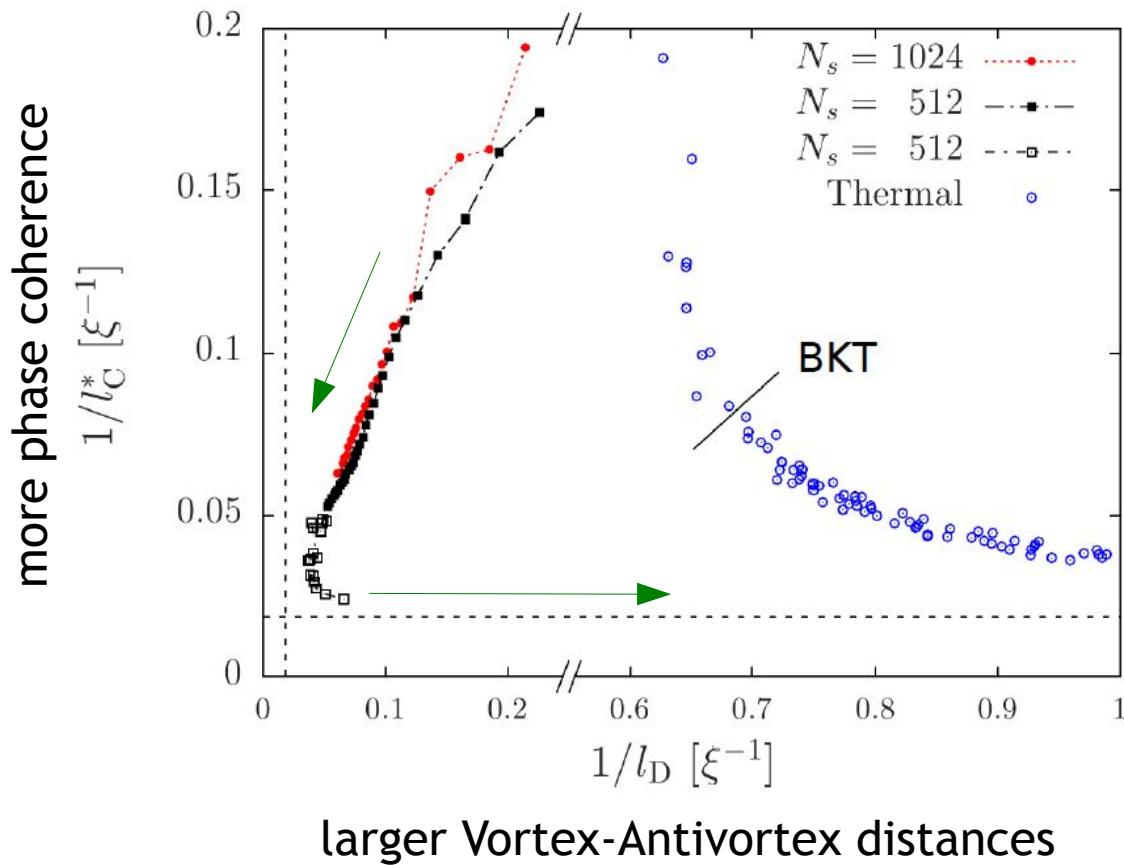
Non-universal decay law
(depends on initial vortex distribution)
Kinetic gas theory for dipoles

Universal decay regime
Strongly correlated dilute vortex gas
Scaling $n(k) \sim k^{-4}$

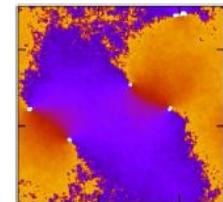
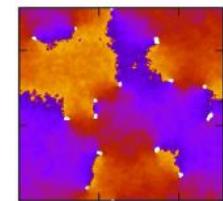
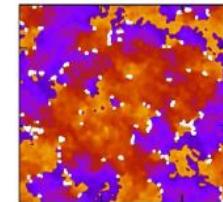
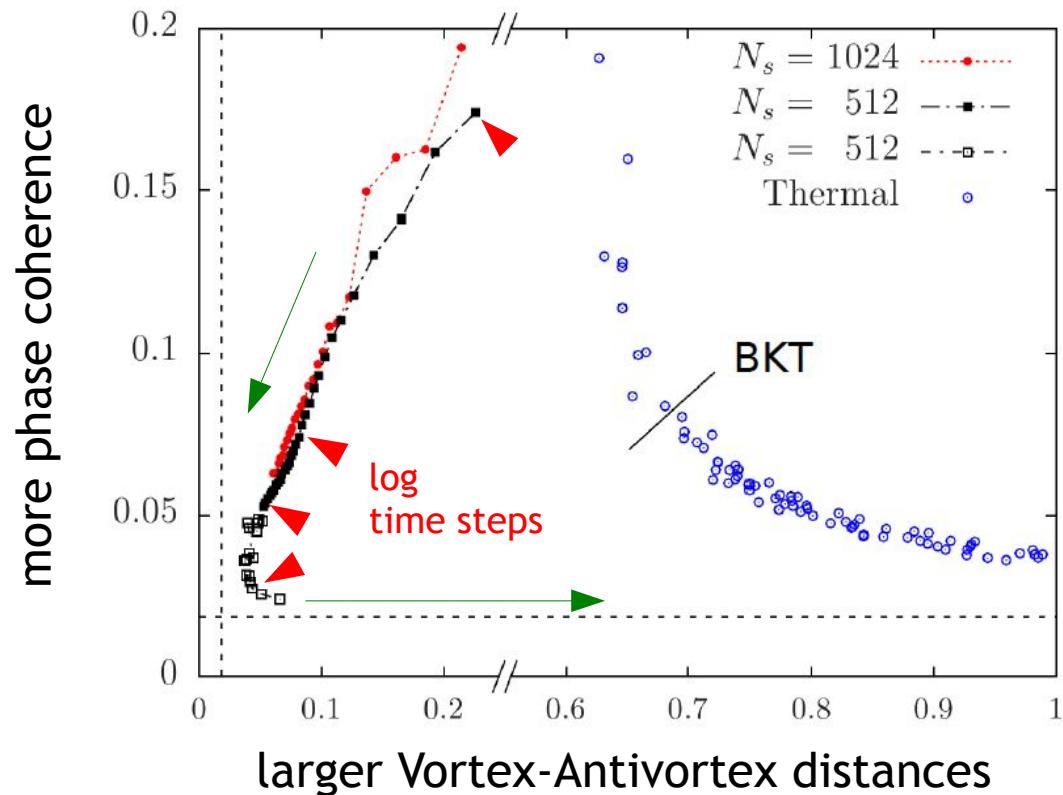


Approach of the NTFP

l_c^* = Phase coherence length
 l_d = Vortex-Antivortex pair distance



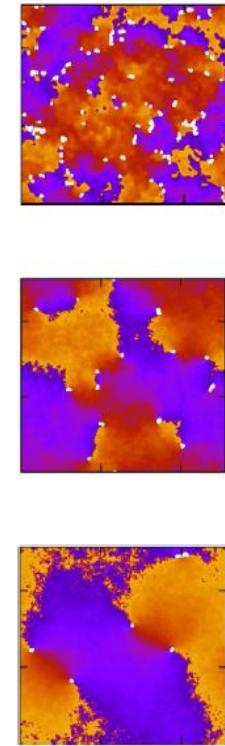
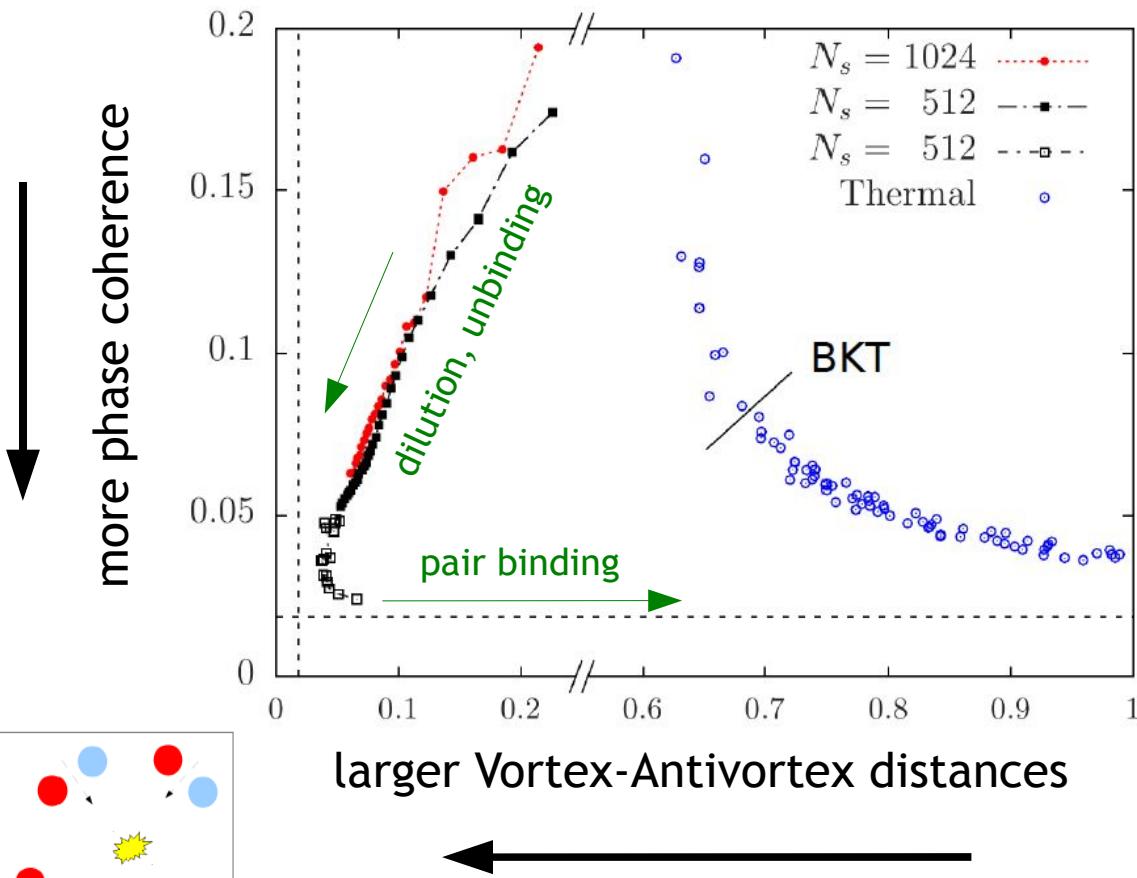
Approach of the NTFP



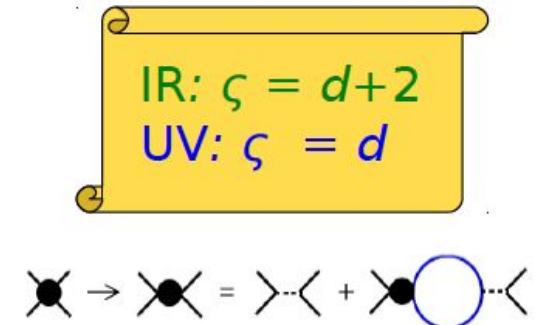
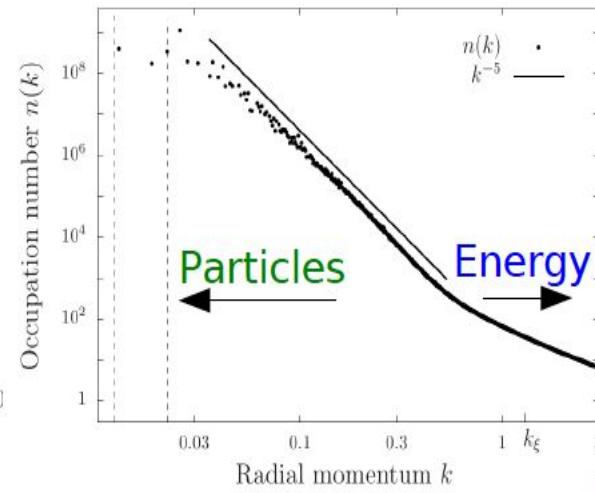
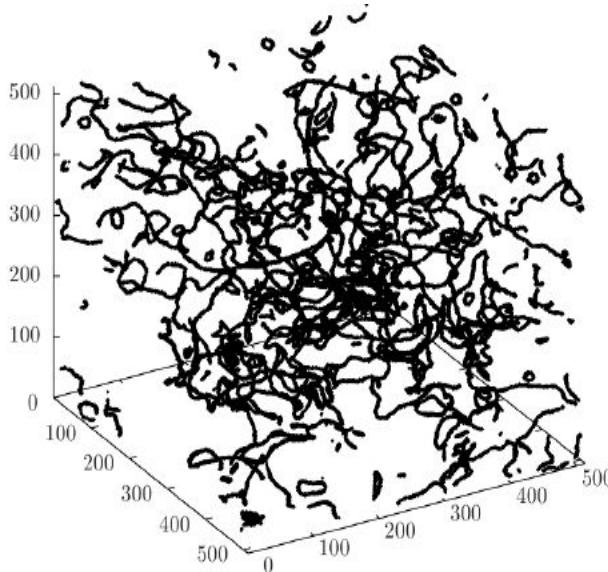
J. Scholz, B. Nowak, TG, PRA 86 (12) 013624



Approach of the NTFP



3D Nonthermal Fixed Point



Vortices



Spectrum



QFT

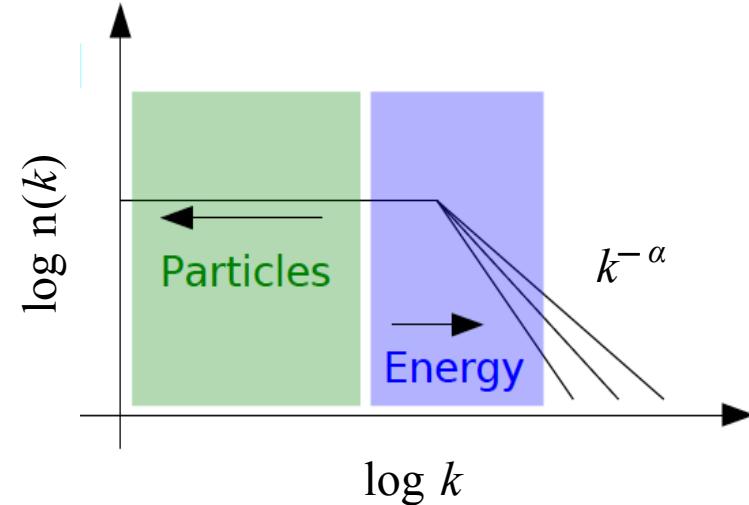
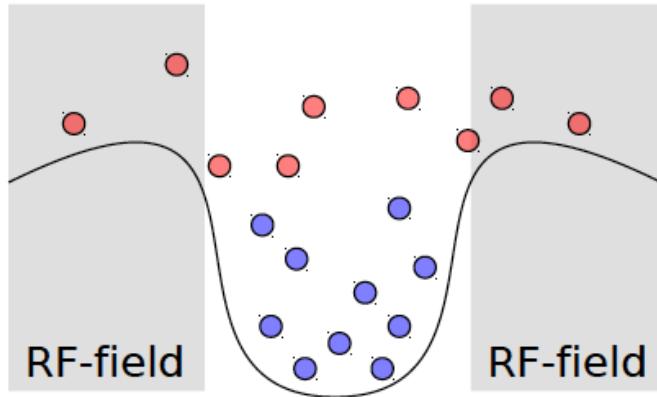
Quantum Turbulence

Strong & Weak Wave Turbulence

B. Nowak, D. Sexty, TG, PRB 84(R) (11); B. Nowak, J. Schole, D. Sexty, TG, PRA 85 (12)



3D: Bose Condensation



Experiments

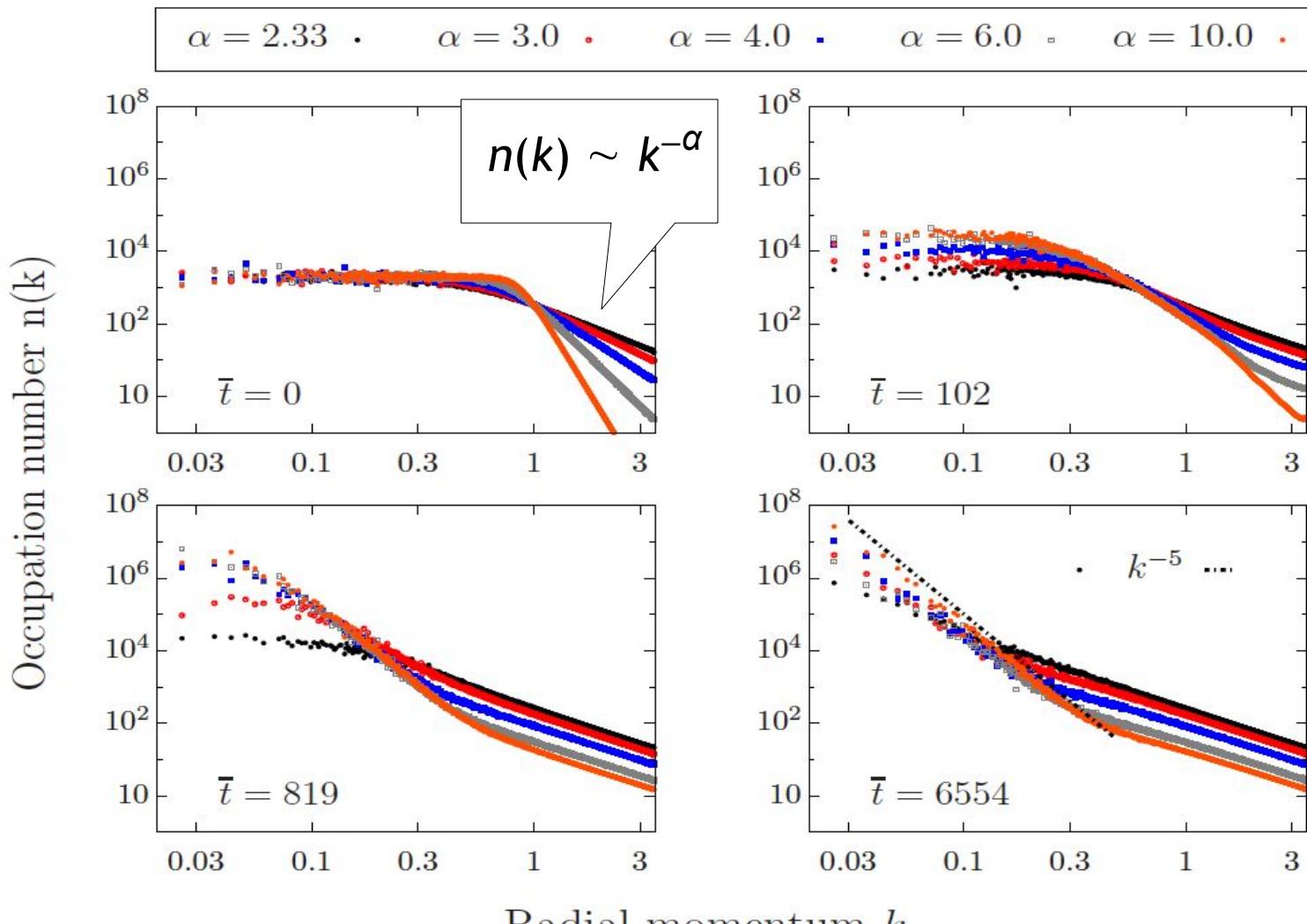
Hänsch, Esslinger (02)
Esslinger (07)
Hadzibabic (12)

Condensation Dynamics

Levich, Yakhot (70s);
Snoke, Wolfe (89);
Kagan, Svistunov, Shlyapnikov (91-94);
Damle, Sachdev (96)
Semikoz, Tkachev (95)
Berloff, Svistunov (02)
Anderson, Davis (08)
Blaizot, McLerran (12)
Berges, Sexty (12)



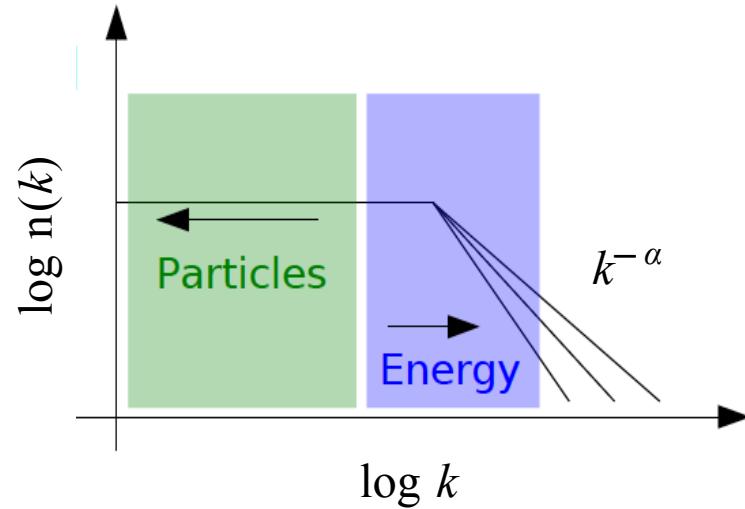
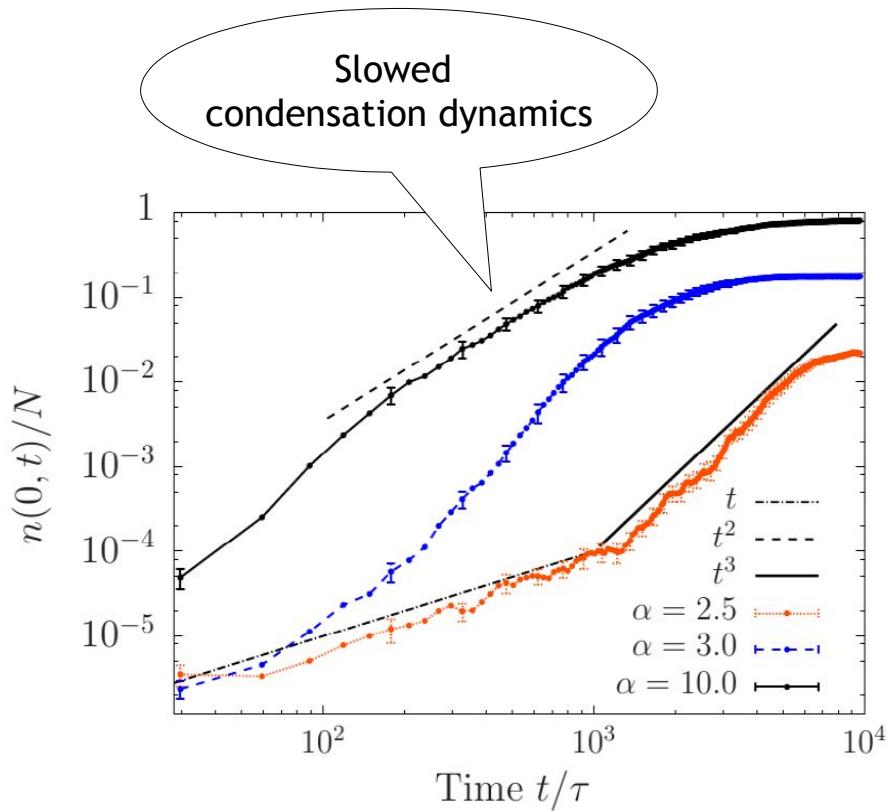
Hydrodynamic vs. kinetic Condensation



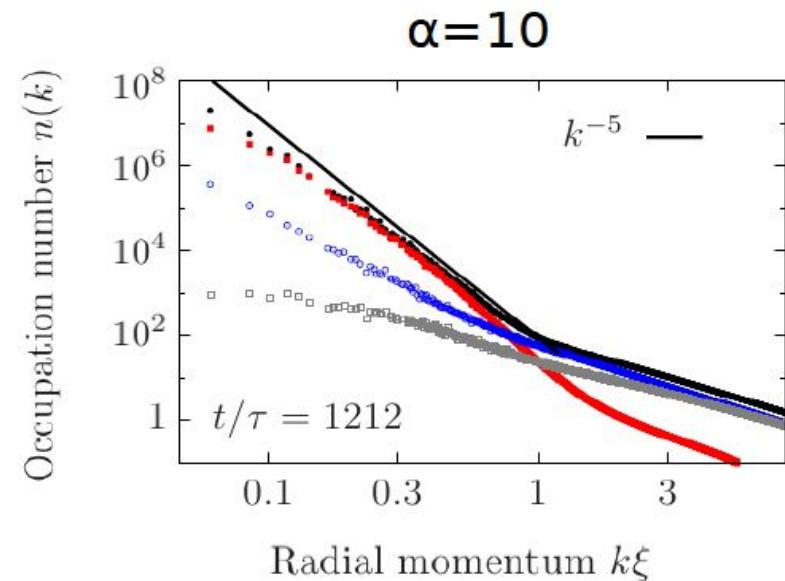
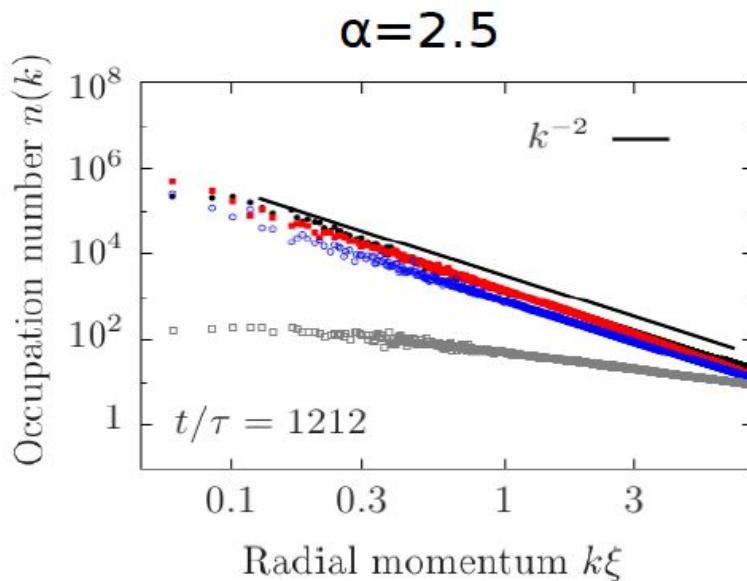
[B. Nowak and TG, unpublished]



3D: Bose Condensation



3D: Bose Condensation



$n^i(k)$

solenoidal
flow

$n^c(k)$

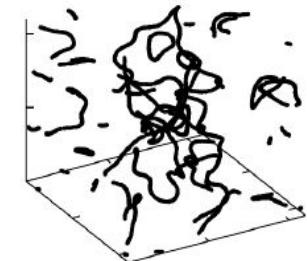
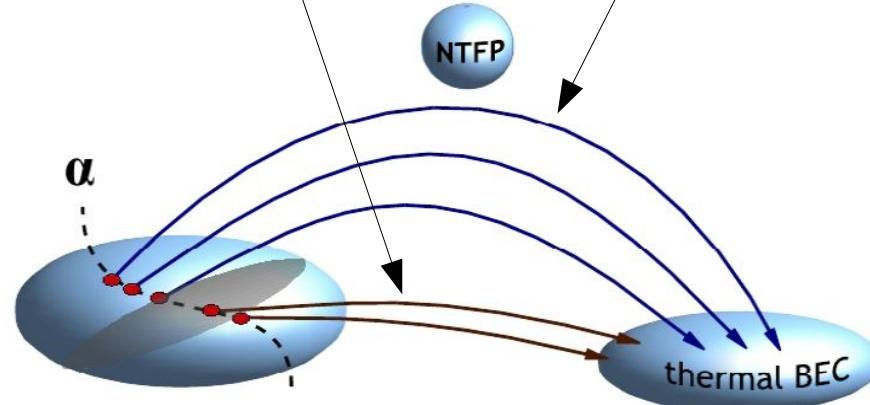
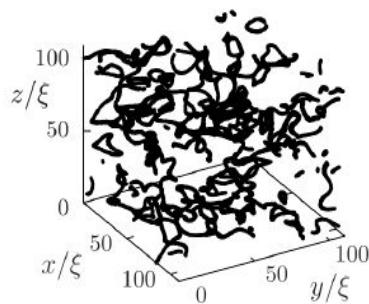
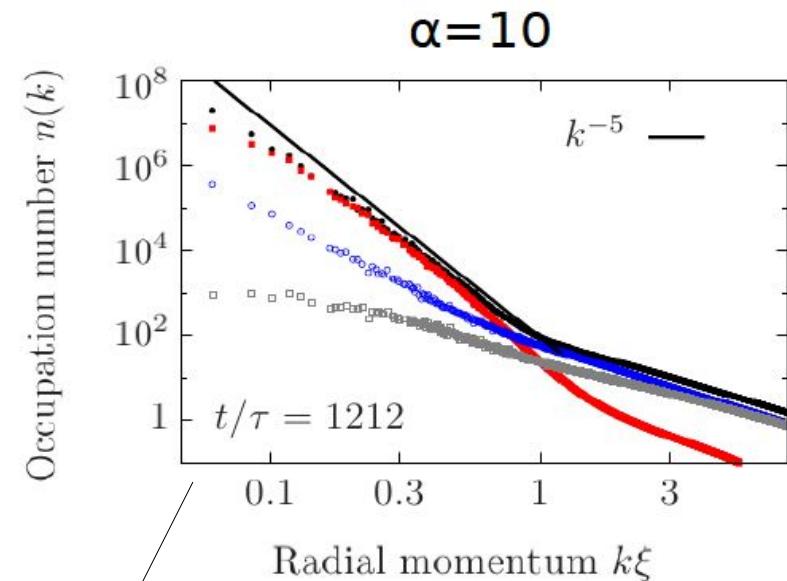
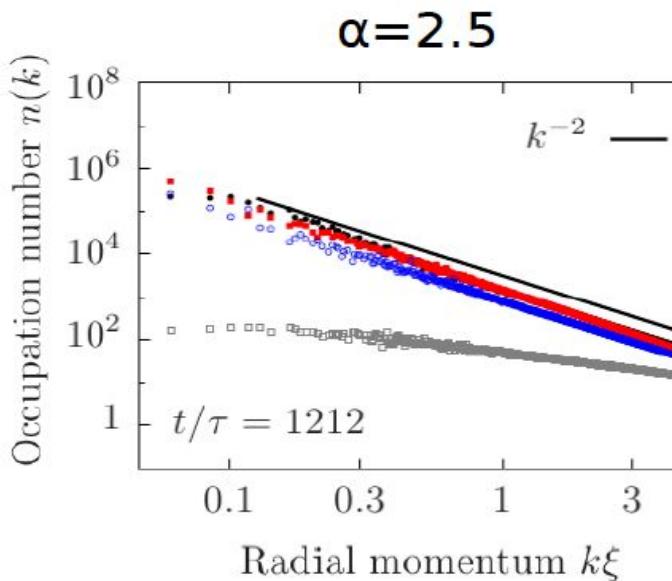
compressible
component

$n^q(k)$

q pressure



3D: Bose Condensation



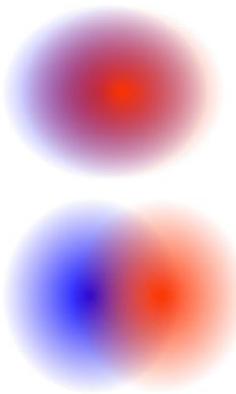
B. Nowak, TG, arXiv: 1206.3181 [cond-mat.quant-gas]



2-component BEC

Bose gas with internal two-level structure

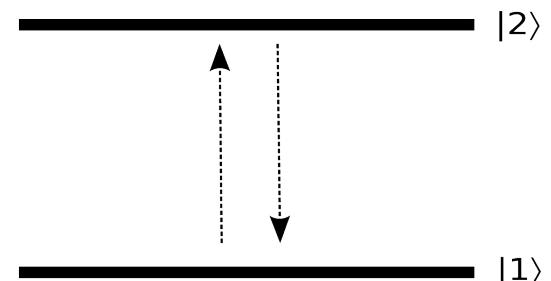
miscible
 $g_{12} < g$



non-linear interactions:

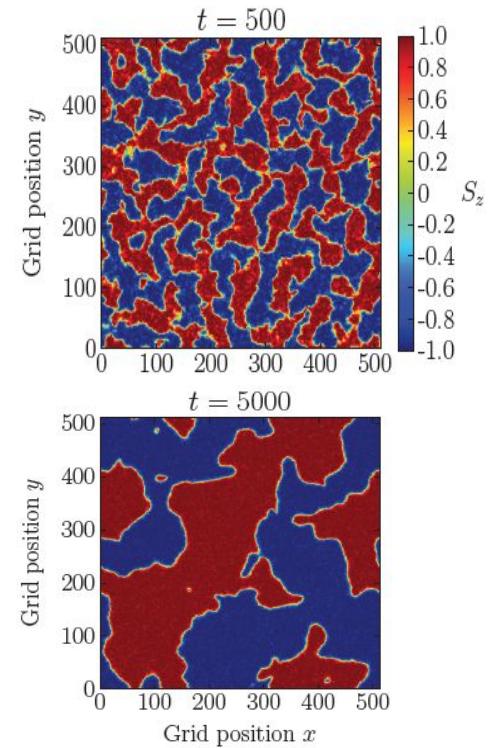
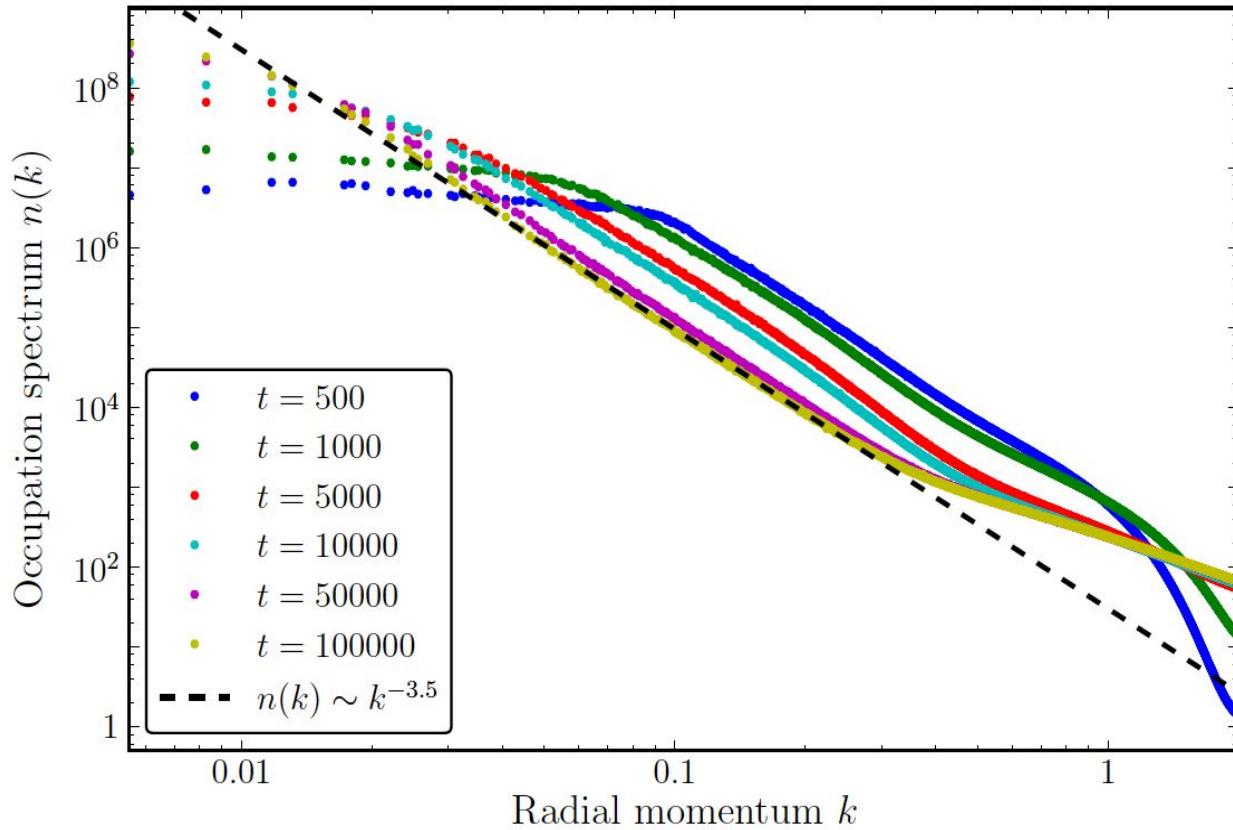
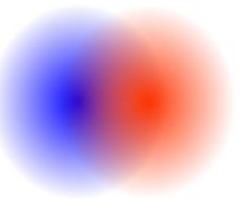
$$\hat{H}_{int} \sim \frac{g}{2}(\hat{n}_1^2 + \hat{n}_2^2) \begin{array}{c} |1\rangle \longleftrightarrow |1\rangle \\ |2\rangle \longleftrightarrow |2\rangle \end{array}$$

$$+ g_{12}\hat{n}_1\hat{n}_2 \quad |1\rangle \longleftrightarrow |2\rangle$$



2-component BEC

immiscible
 $g_{12} > g$



Decomposition of Energy for Spin system

$$\begin{aligned} E_{\text{kin}} &= \frac{1}{2} \int d^2x \nabla \phi_j \nabla \phi_j \\ &= \frac{1}{2} \int d^2x \left[|\nabla \sqrt{\rho_T}|^2 + \frac{\rho_T}{4} \nabla S^a \nabla S^a + |\mathbf{w}|^2 \right] \end{aligned}$$

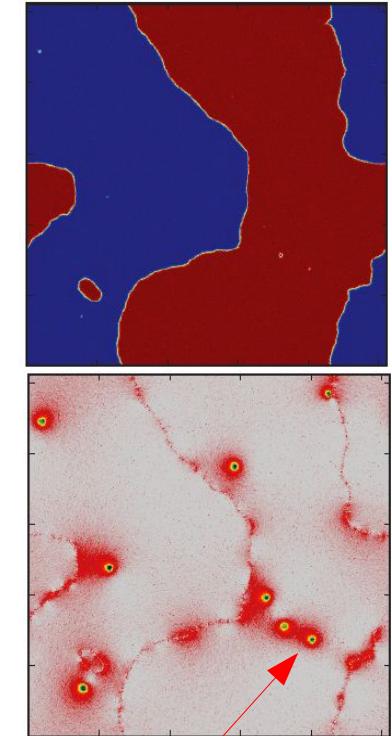
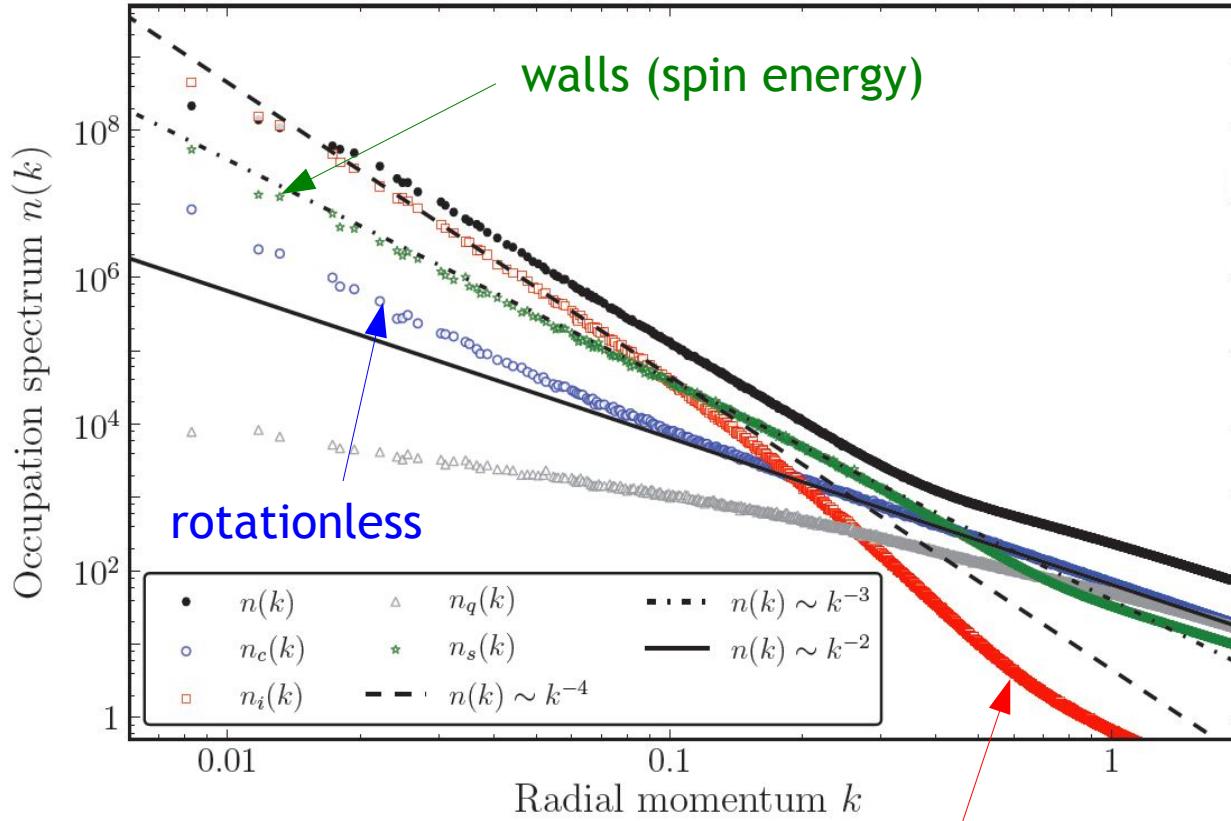
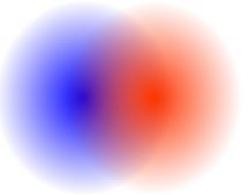
$$S^a = \phi_j \sigma_{ij}^a \phi_i \quad \begin{matrix} \rho_1 + \rho_2 \equiv \rho_T \\ \Theta_T = \varphi_1 + \varphi_2 \end{matrix} \quad S^a \rightarrow \rho_T S^a$$

$$\mathbf{j}_T = \frac{1}{2} \rho_T \nabla \Theta_T + \frac{S^z \rho_T}{2 [(S^x)^2 + (S^y)^2]} (S^y \nabla S^x - S^x \nabla S^y) = \rho_1 \nabla \varphi_1 + \rho_2 \nabla \varphi_2$$



2-component BEC

immiscible
 $g_{12} > g$



solenoidal \leftrightarrow Skyrmions

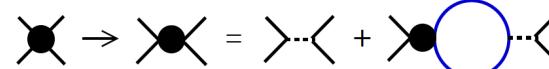
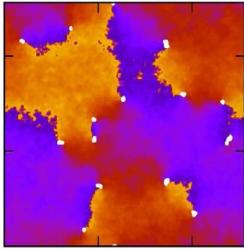
M. Karl, B. Nowak, TG, unpublished (12)

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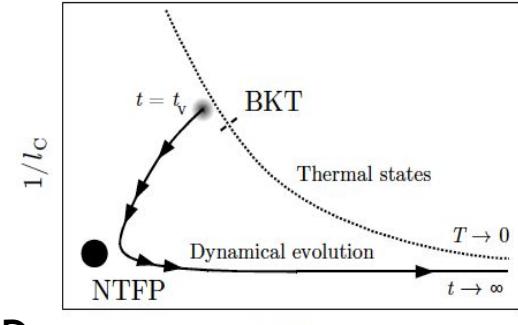
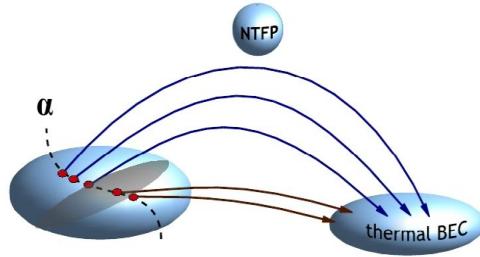
Summary



Non-Thermal Fixed Points

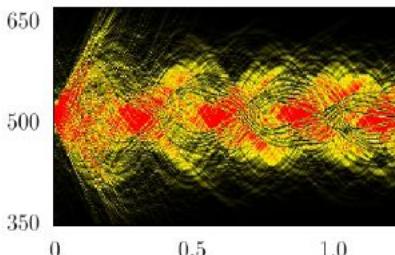


Superfluid Turbulence in 2D

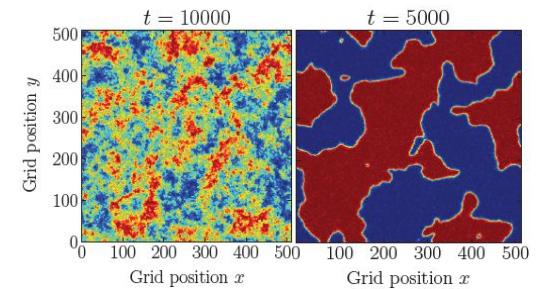


Approach of the NTFP in 2D

Dynamics of BE condensation



Charge separation/Pattern formation



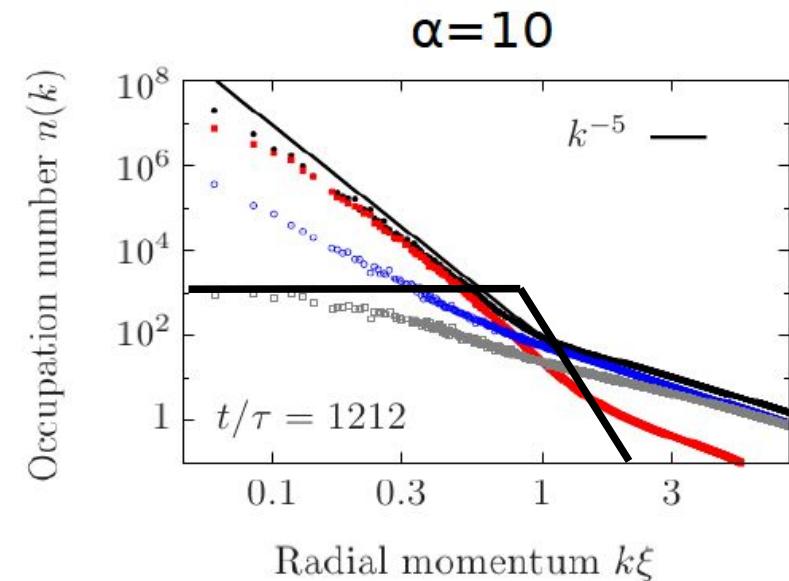
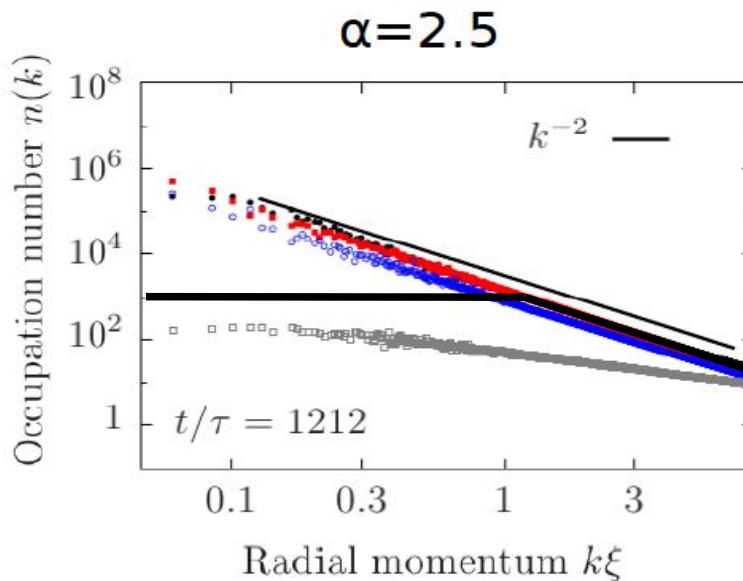
1D Soliton Gas as an NTFP



Supplementary slides

Bose-Einstein Condensation

3D: Bose Condensation



$n^i(k)$

solenoidal
flow

$n^c(k)$

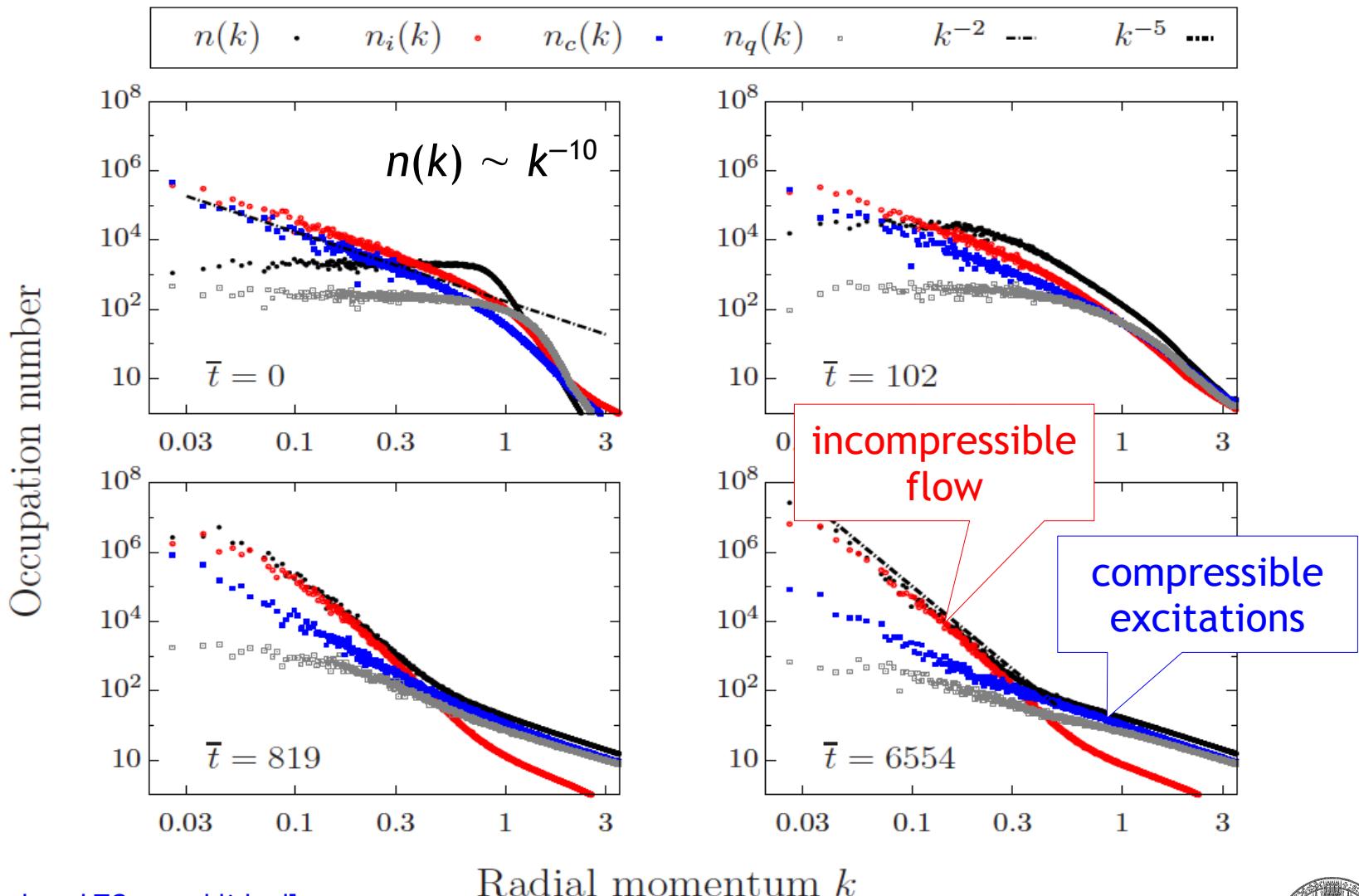
compressible
component

$n^q(k)$

q pressure



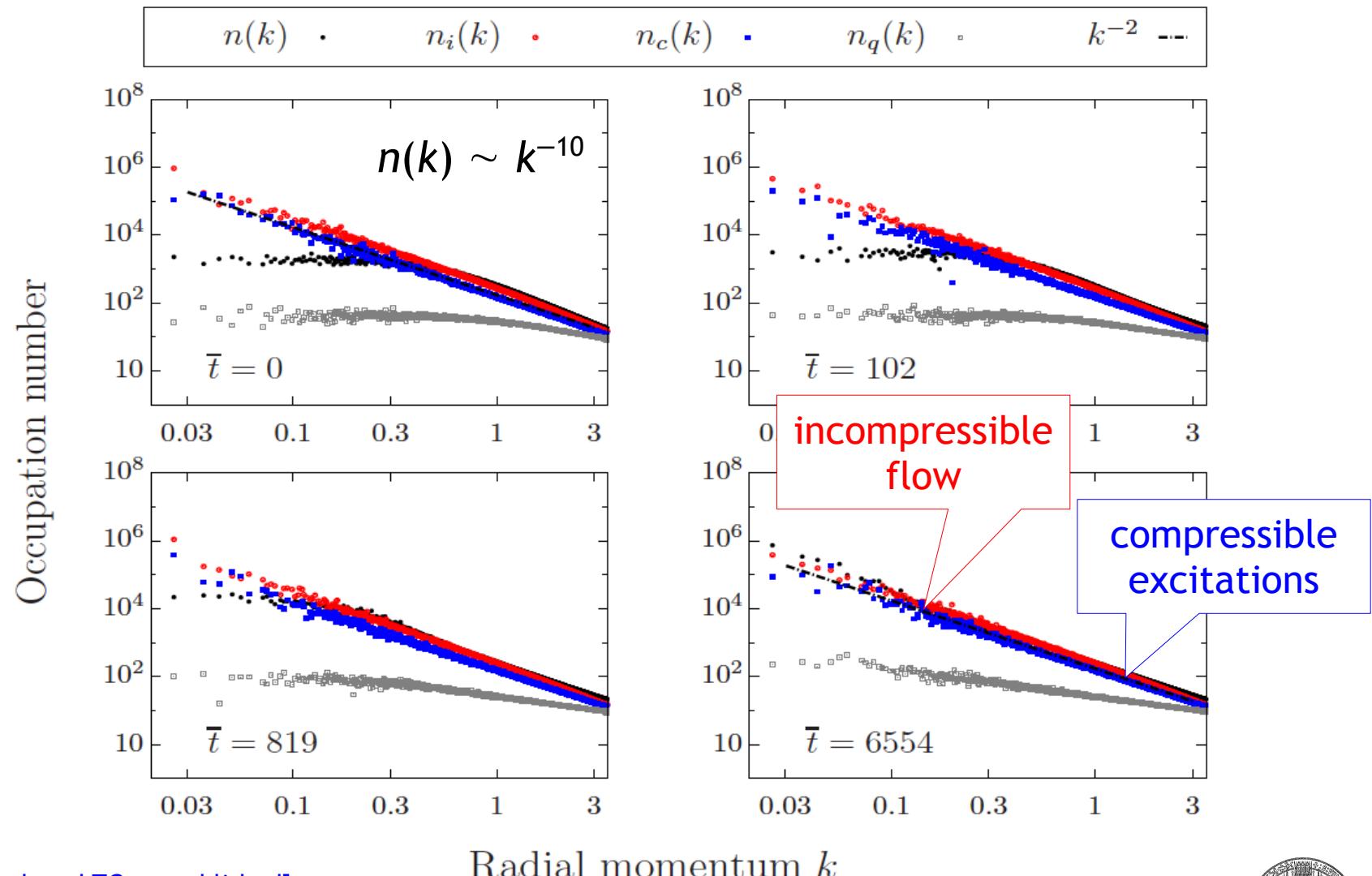
Hydrodynamic vs. kinetic Condensation



[B. Nowak and TG, unpublished]



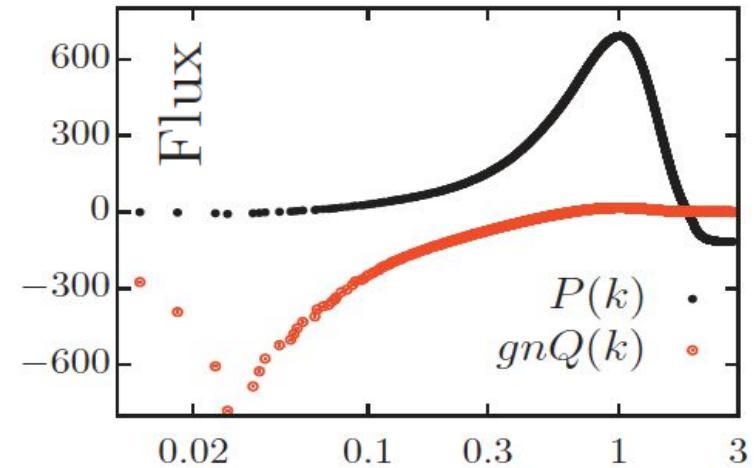
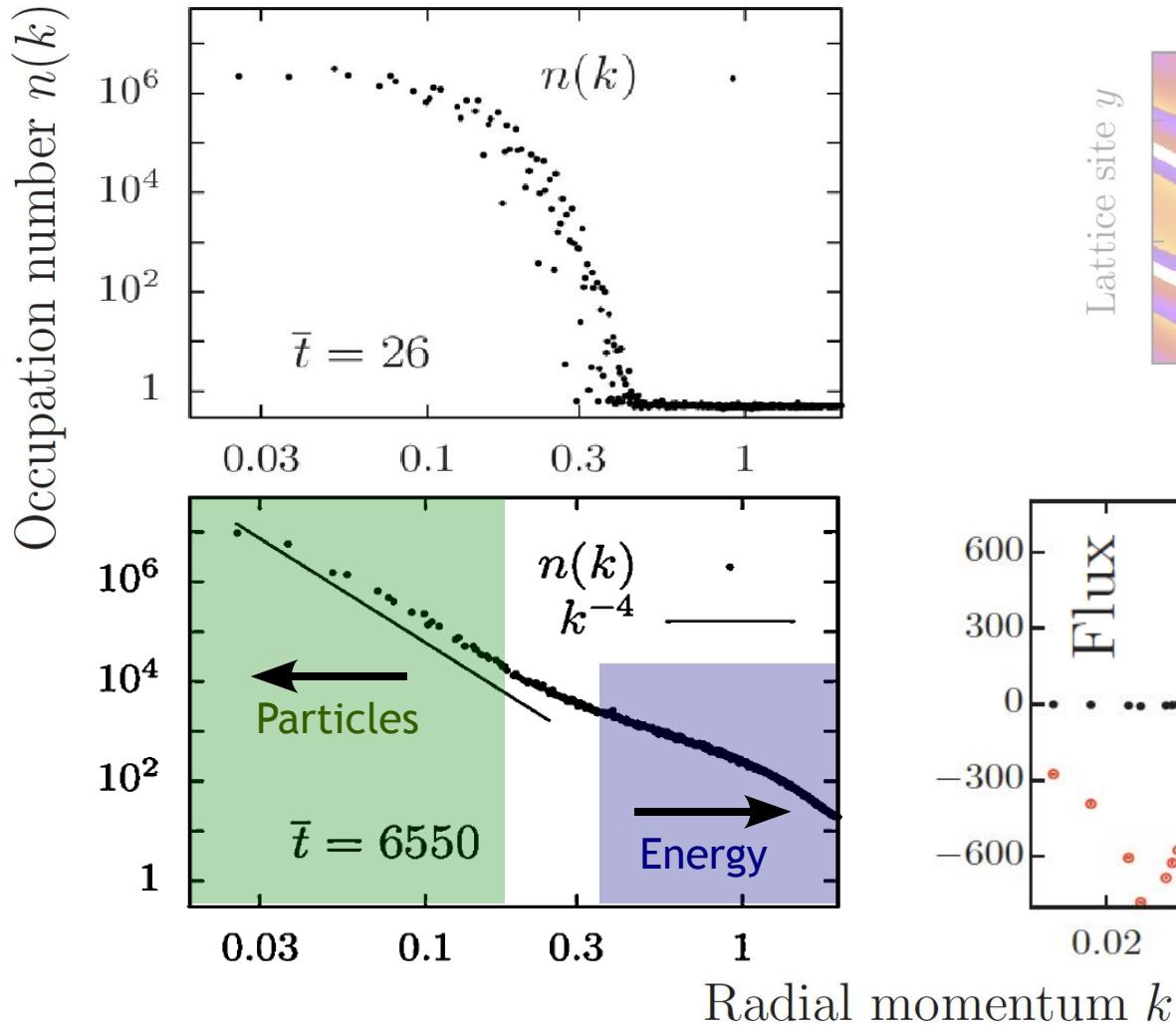
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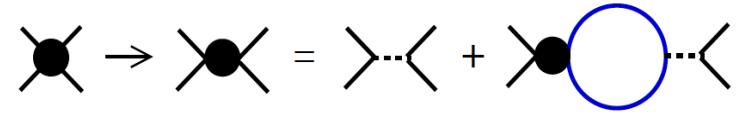
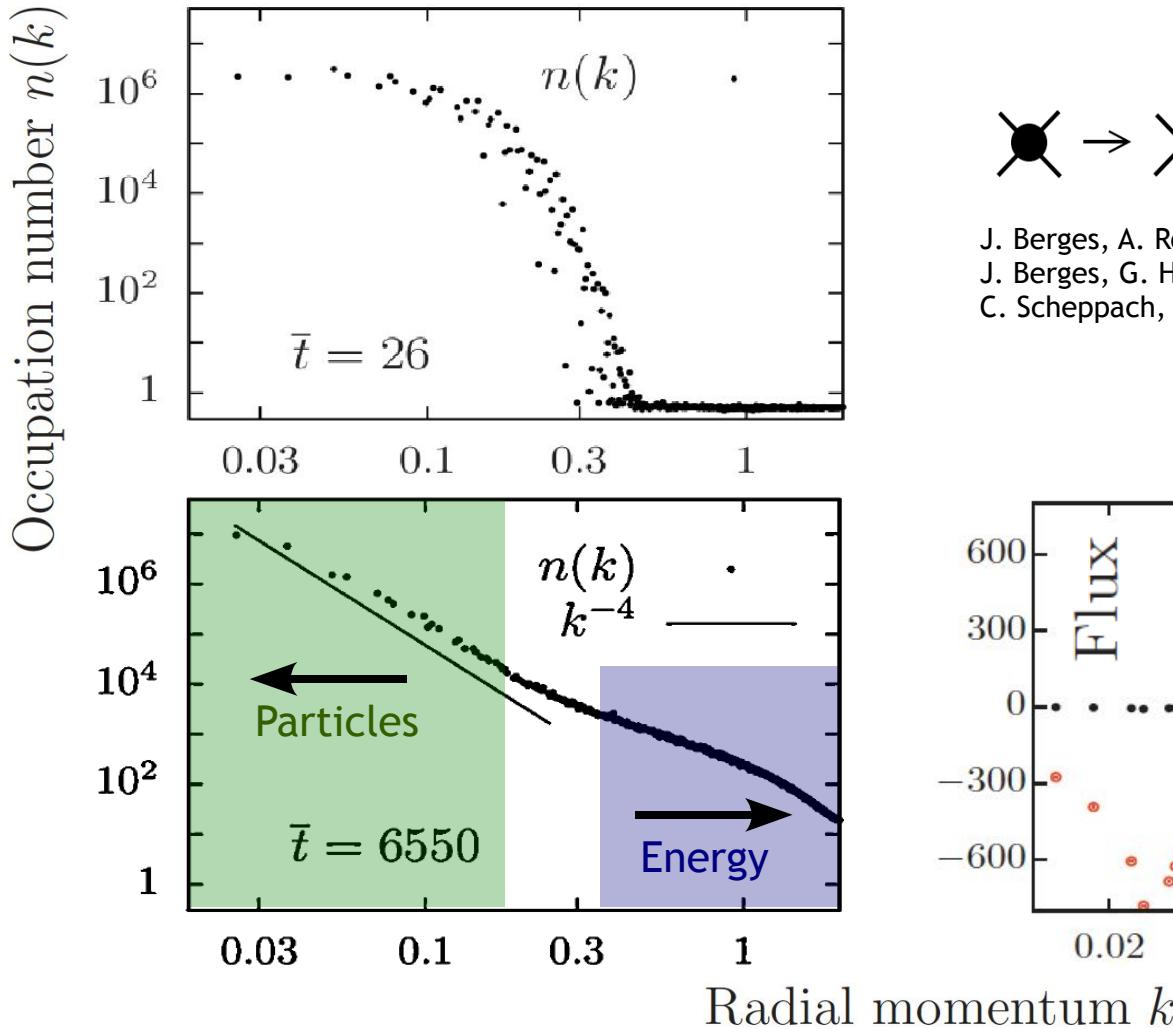
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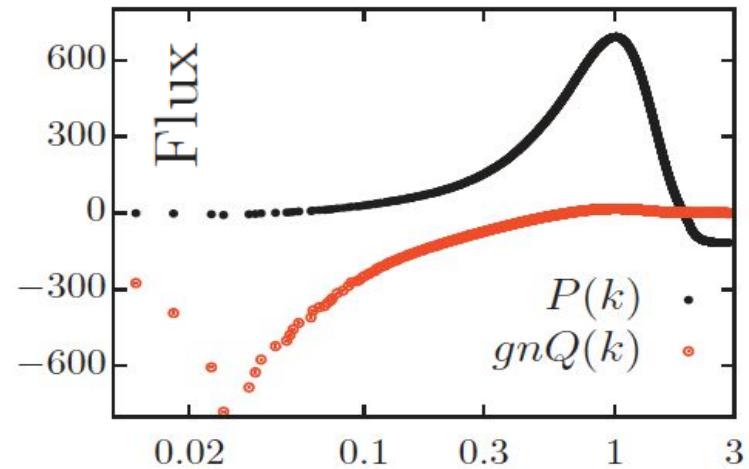
2+1 D: Quench dynamics



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 J. Berges, G. Hoffmeister, NPB 813 (09) 383,
 C. Scheppach, J. Berges, TG PRA 81 (10) 033611,

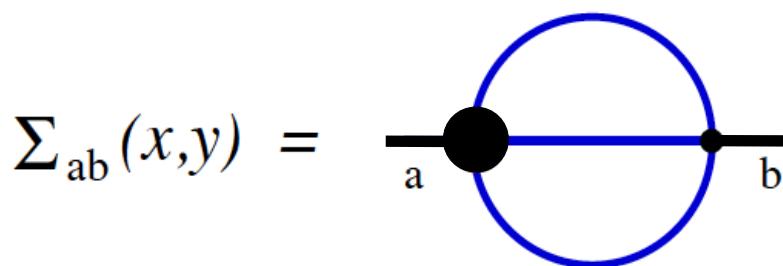
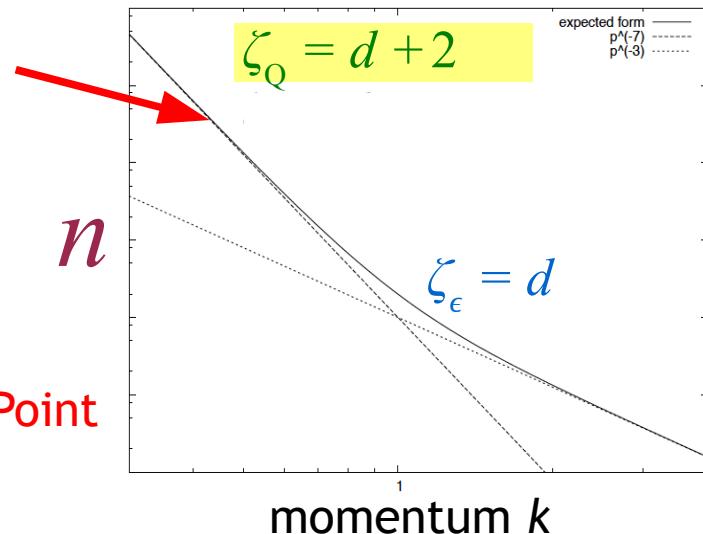


NTFP scaling in d dim^s

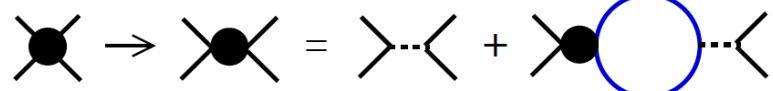
$$n \sim k^{-\zeta}$$

New exponent
beyond
Quantum Boltzmann!

@ Nonthermal Fixed Point



Vertex bubble resummation:
(2PI to NLO in $1/N$)



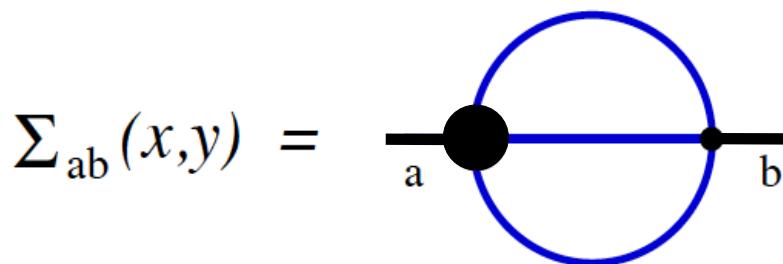
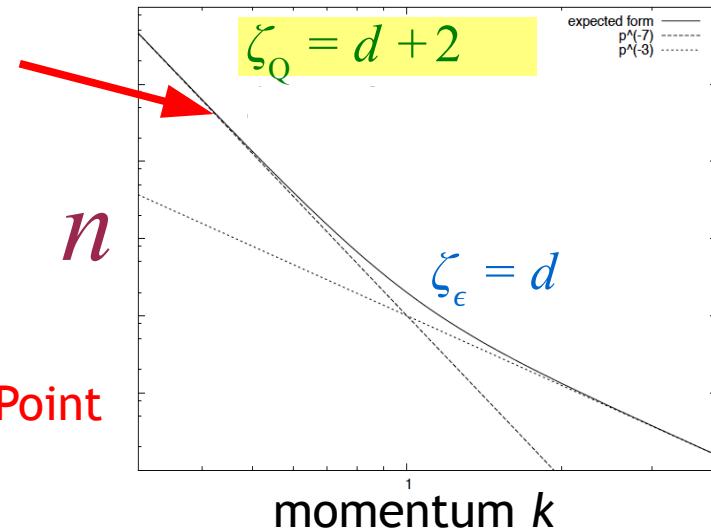
J. Berges, A. Rothkopf, J. Schmidt, PRL 101 (08) 041603, J. Berges, G. Hoffmeister, NPB 813 (09) 383
C. Scheppach, J. Berges, TG PRA 81 (10) 033611



Bose gas in d spatial dimensions $n \sim k^{-\zeta}$

New exponent
beyond
Quantum Boltzmann!

@ Nonthermal Fixed Point



Vertex bubble resummation:
(2PI to NLO in $1/N$)

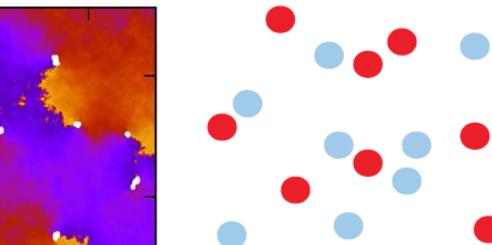
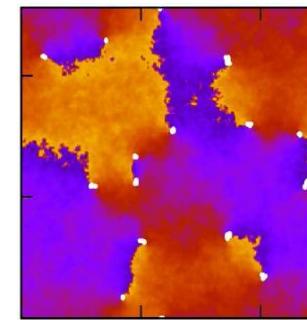
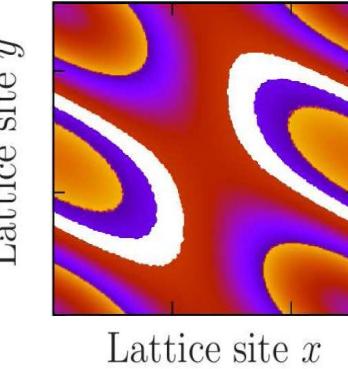
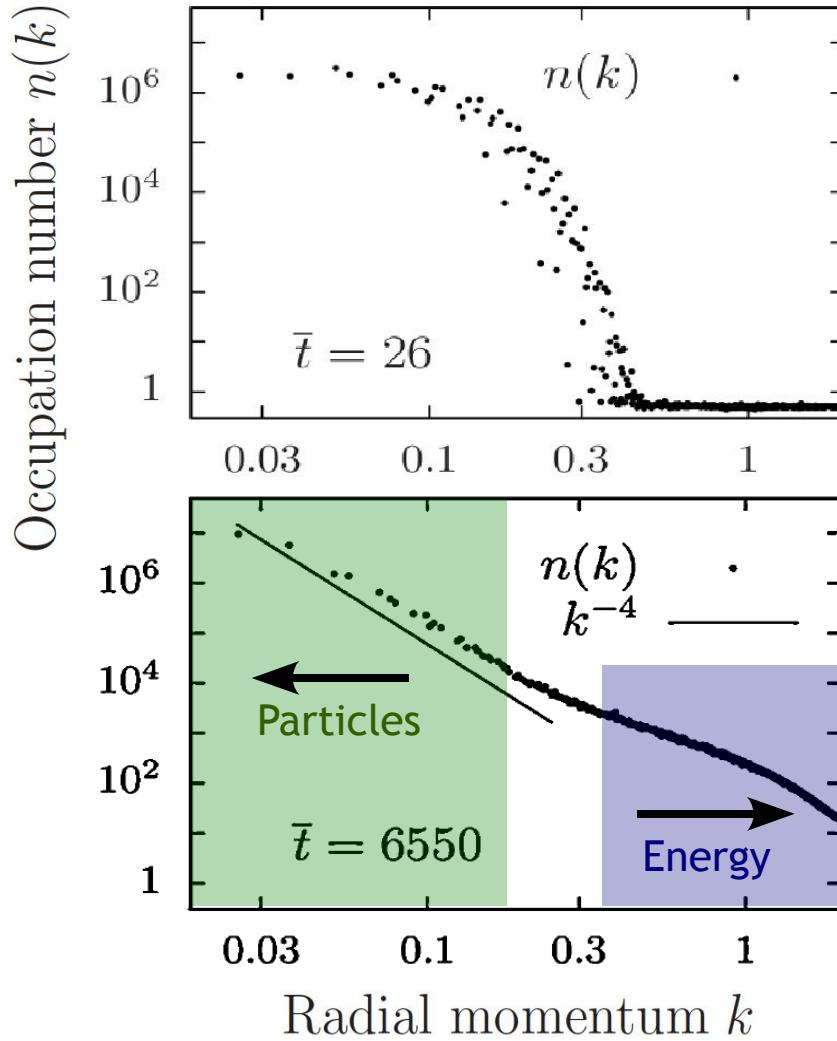
$$\text{---} \rightarrow \text{---} = \text{---} + \text{---}$$

J. Berges, A. Rothkopf, J. Schmidt, PRL 101 (08) 041603, J. Berges, G. Hoffmeister, NPB 813 (09) 383
C. Scheppach, J. Berges, TG PRA 81 (10) 033611



2+1 D: Quench dynamics

B. Nowak, D. Sexty, TG, PRB 84(R) (11);
B. Nowak, J. Scholz, D. Sexty, TG, PRA 85 (12)



$$n(k) \sim k^{-4}$$
$$\Leftrightarrow E(k) \sim k^{-1}$$



Point vortex model

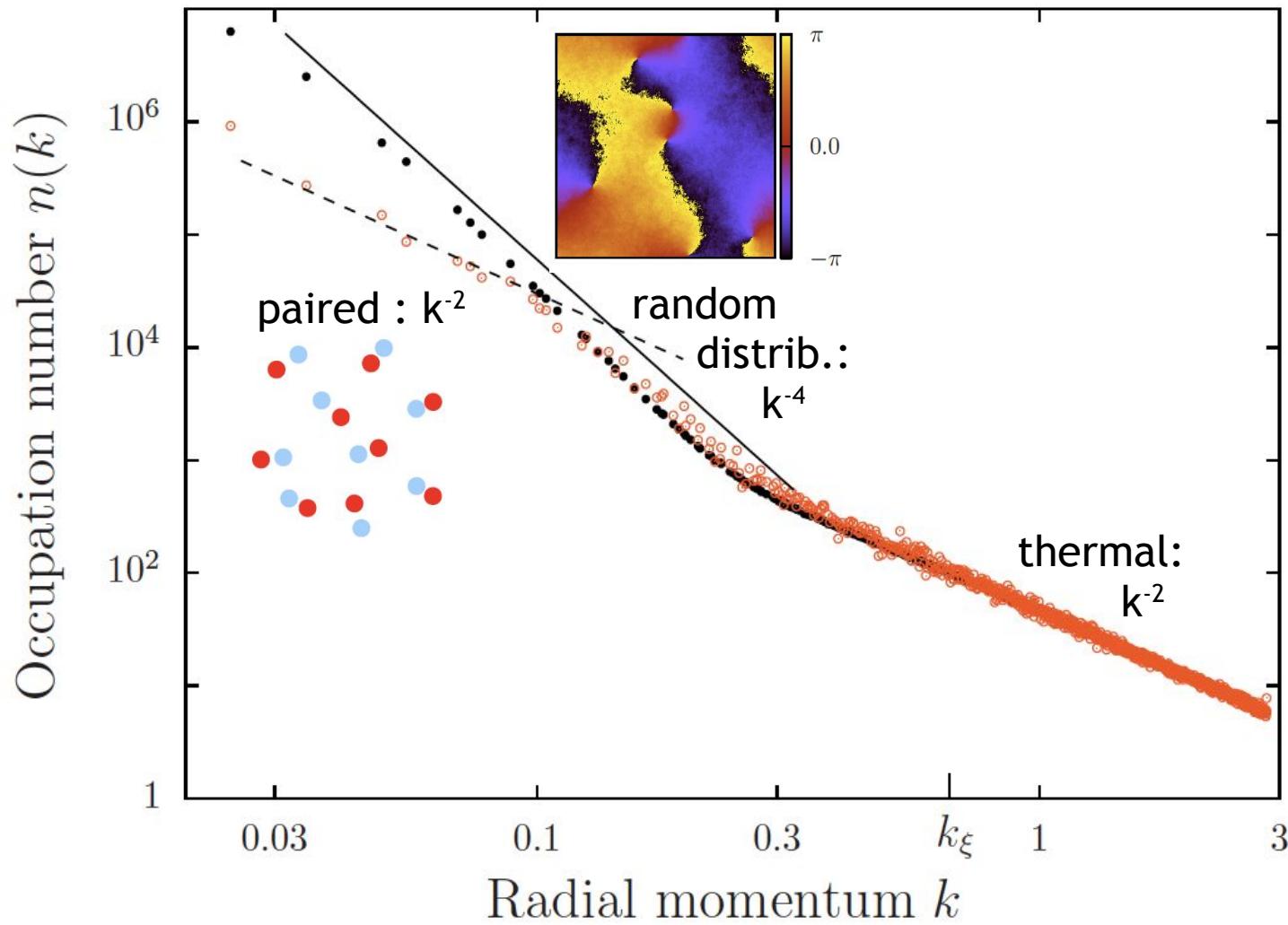
Movie 2: Vortex “gas” & Spectrum

$$n(k) = \langle \Psi^*(\mathbf{k}) \Psi(\mathbf{k}) \rangle \Big|_{\text{angle average}}$$

<http://www.thphys.uni-heidelberg.de/~smp/gasenzer/videos/boseqt.html>



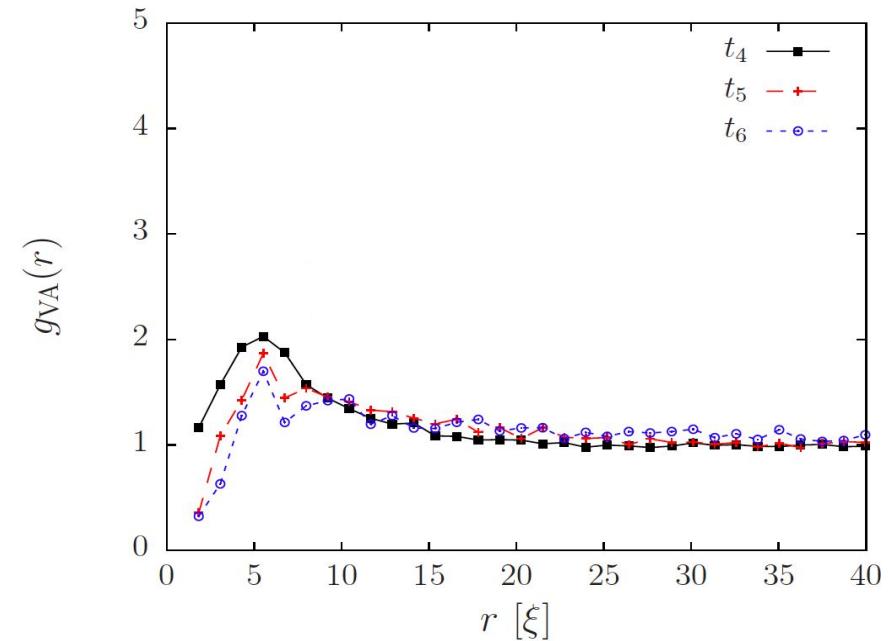
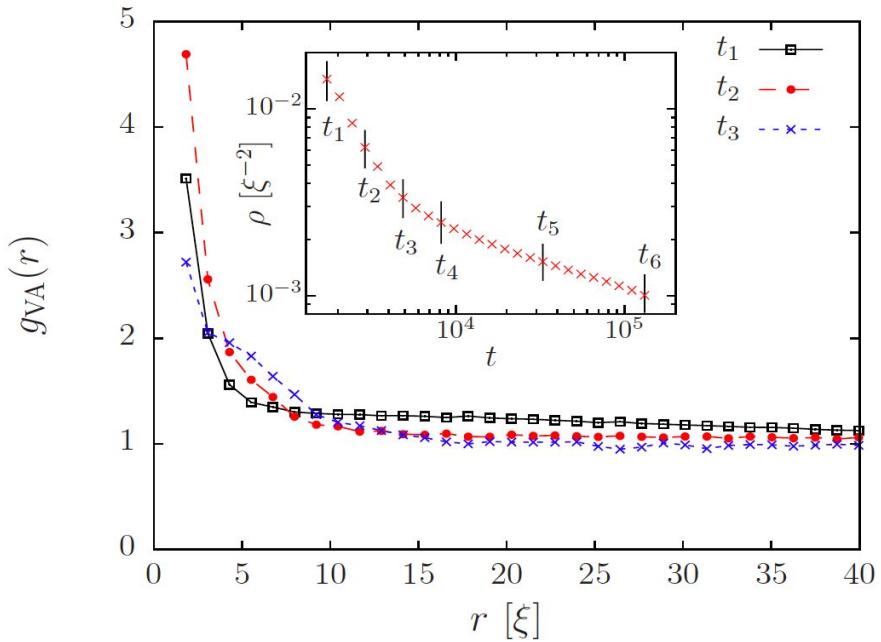
Scaling & Vortex correlations



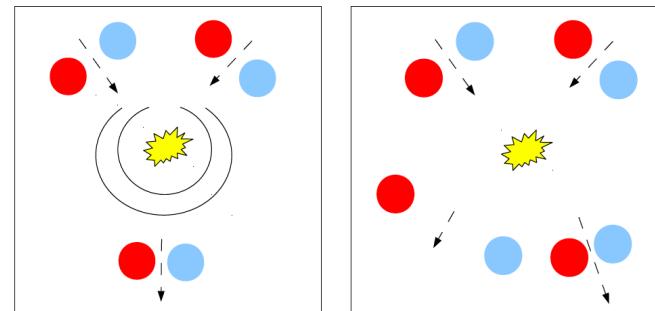
B. Nowak, J. Schole, D. Sexty, TG, arXiv:1111.6127, PRA, to appear (12)



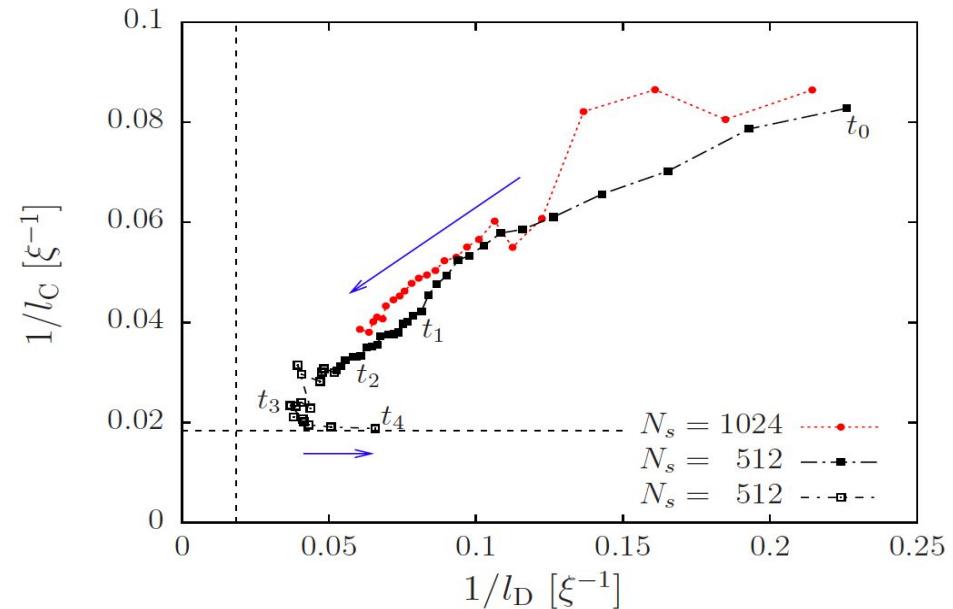
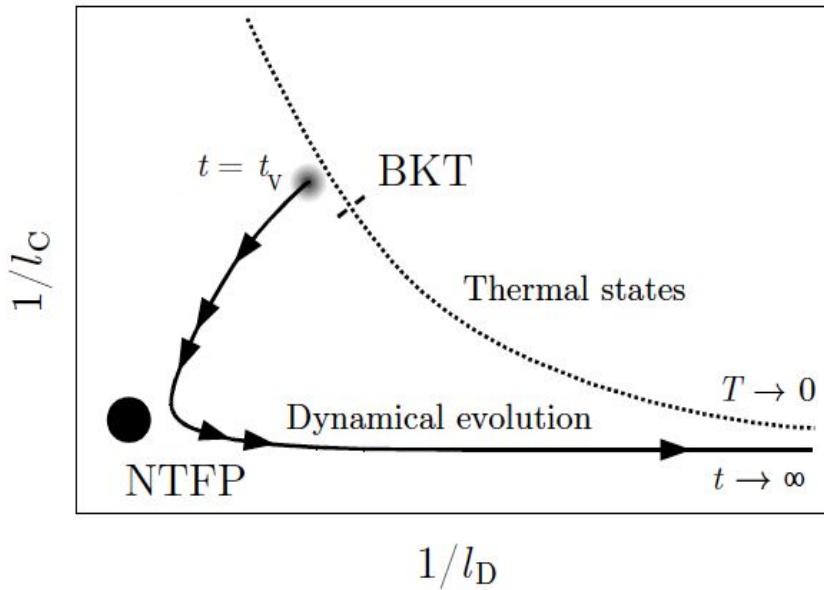
Dynamical vortex unpairing



$$g_{\text{VA}}(\mathbf{x}, \mathbf{x}', t) = \frac{\langle \rho^{\text{V}}(\mathbf{x}, t) \rho^{\text{A}}(\mathbf{x}', t) \rangle}{\langle \rho^{\text{V}}(\mathbf{x}, t) \rangle \langle \rho^{\text{A}}(\mathbf{x}', t) \rangle}$$



Nonthermal fixed point in 2D



J. Schole, B. Nowak, TG, arXiv:1204.2487 [cond-mat.quant-gas]

Perturbative RG for dyn. near BKT: Mathey & Polkovnikov, PRA **80**, 041601R (09), **81**, 033605 (10)
See also: Jelic & Cugliandolo, J. Stat. Mech. P02032 (11)



Spin Systems

Decomposition of Energy

$$\begin{aligned} E_{tot} &= \int \left(\frac{1}{2} |\nabla \sqrt{n} e^{-i\varphi}|^2 + \frac{1}{2} g n^2 \right) d\boldsymbol{\rho} \\ &= E_{kin} + E_q + E_{int} \end{aligned}$$

rotationless

$$E_{kin} = \frac{1}{2} \int |\sqrt{n}\mathbf{u}|^2 d\boldsymbol{\rho} = E_{kin}^i + E_{kin}^c$$

solenoidal

$$E_q = \frac{1}{2} \int (\nabla \sqrt{n})^2 d\boldsymbol{\rho}$$

$\nabla \times (\sqrt{n}\mathbf{u})^c = 0$

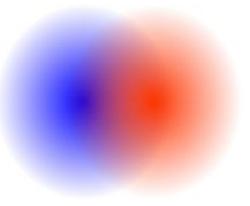
$\nabla \cdot (\sqrt{n}\mathbf{u})^i = 0$

$$E_s(k) = \frac{1}{2} \int d\Omega_k \langle \mathbf{w}_s^a(k) \cdot \mathbf{w}_s^a(k) \rangle \quad \begin{aligned} \mathbf{u}(\boldsymbol{\rho}, t) &= \nabla \varphi(\boldsymbol{\rho}, t) \\ \mathbf{w}_s^a &= \frac{\sqrt{\rho_T}}{2} \nabla S^a \end{aligned}$$

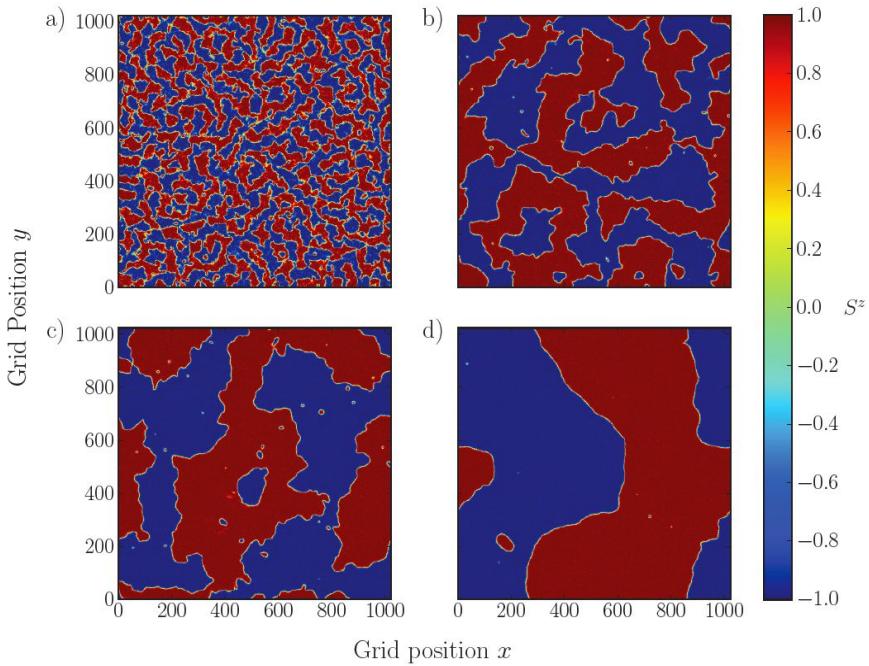


2-component BEC

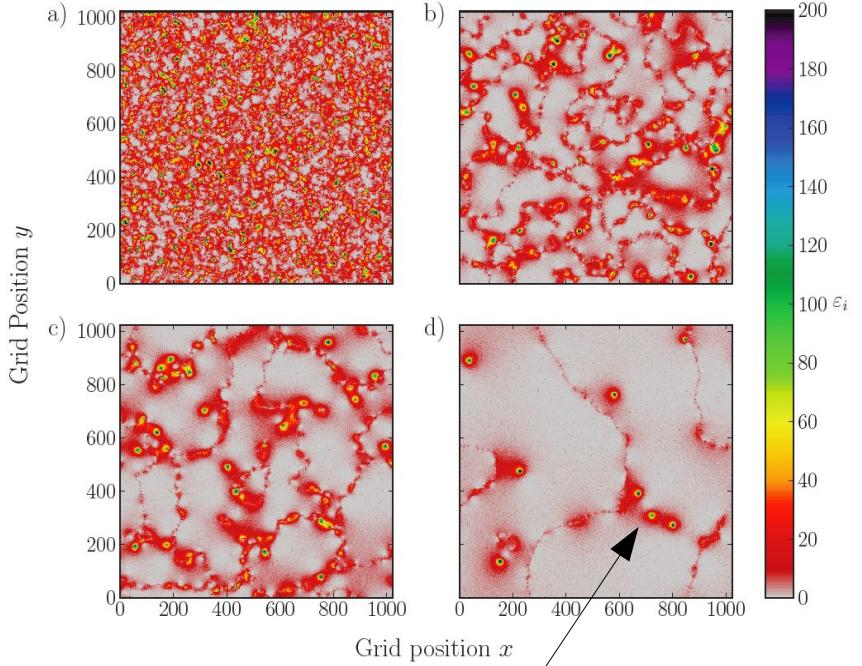
immiscible
 $g_{12} > g$



$S^z(\mathbf{x})$



$\frac{1}{2} |\mathbf{w}_i(\mathbf{x})|^2$

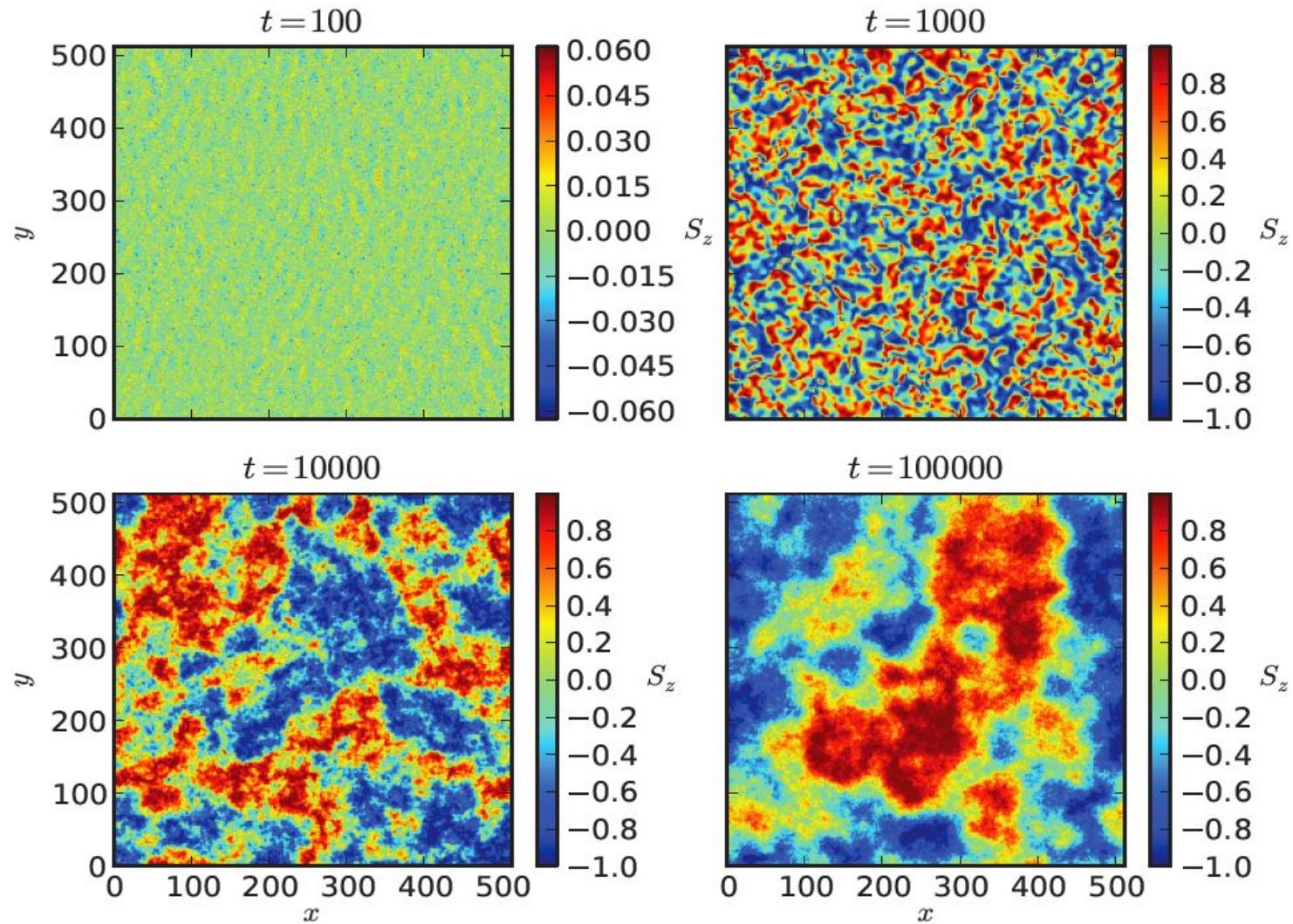


M. Karl, B. Nowak, TG, unpublished (12)

Skyrmions



Domain formation in spin systems

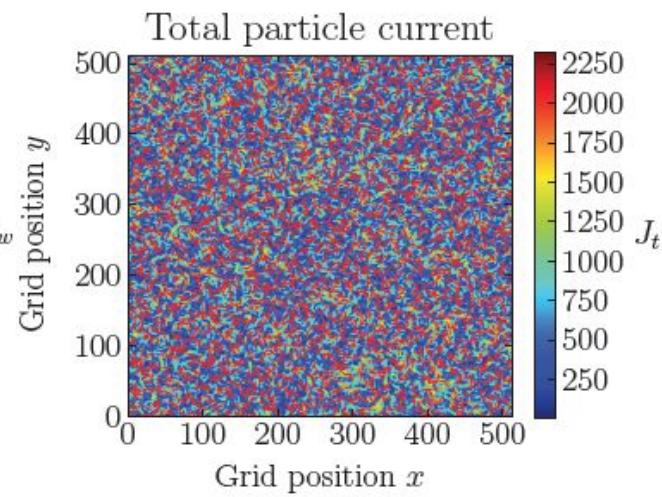
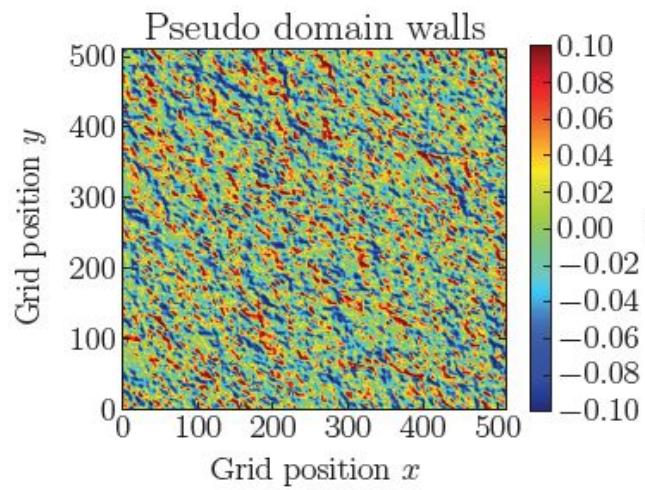
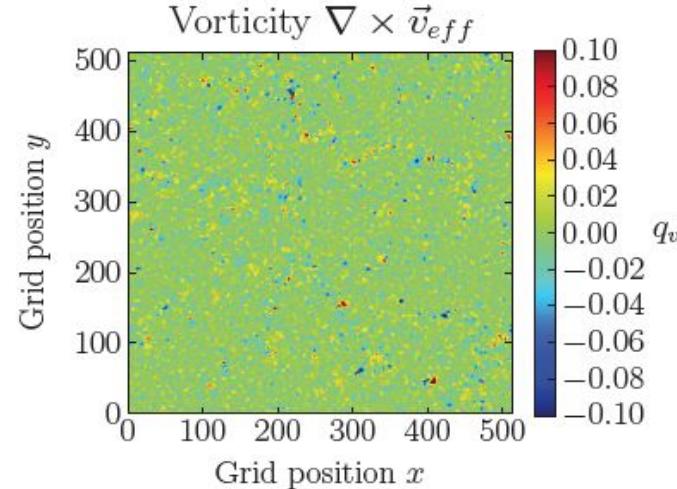
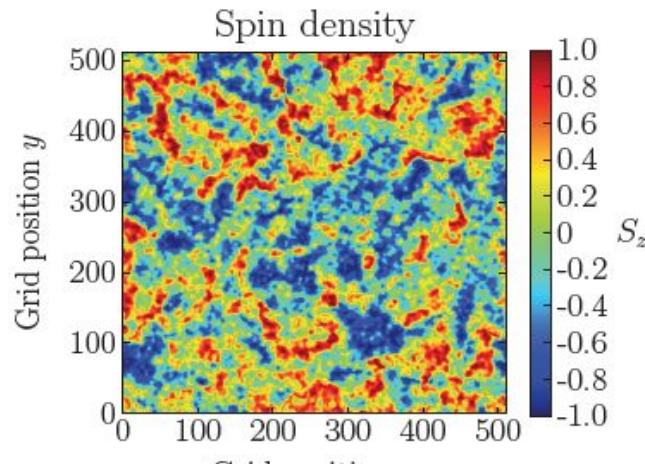


[M. Karl, B. Nowak, and TG, unpublished]



2-component BEC

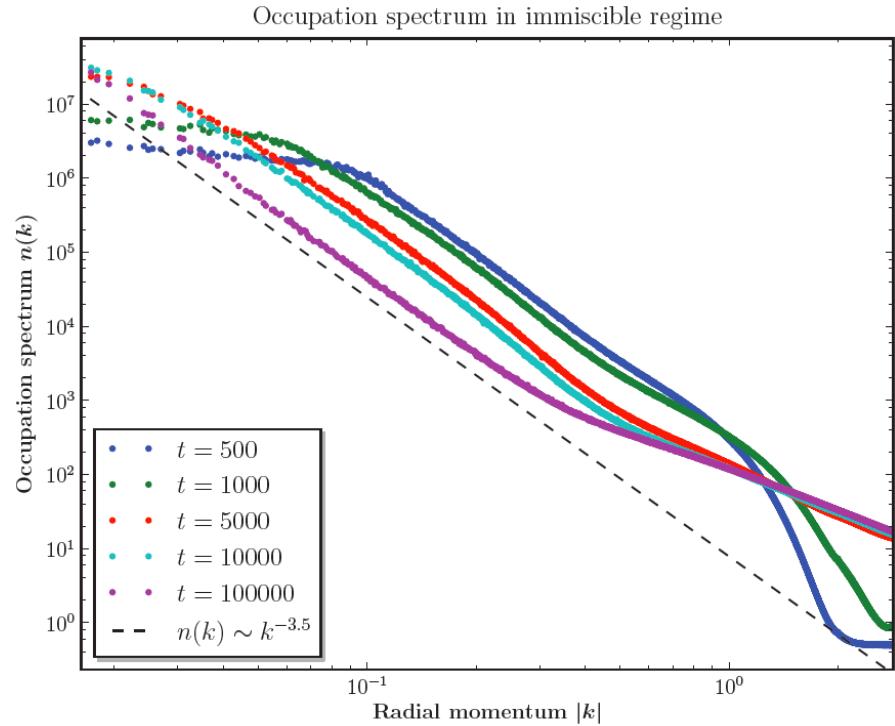
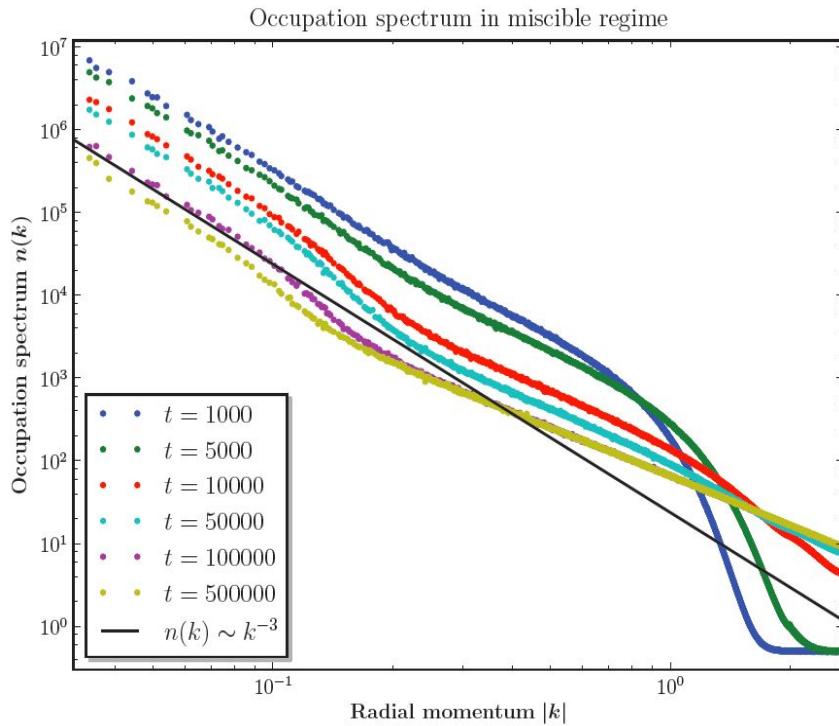
miscible
 $g_{12} < g$



M. Karl, B. Nowak, TG, unpublished (12)



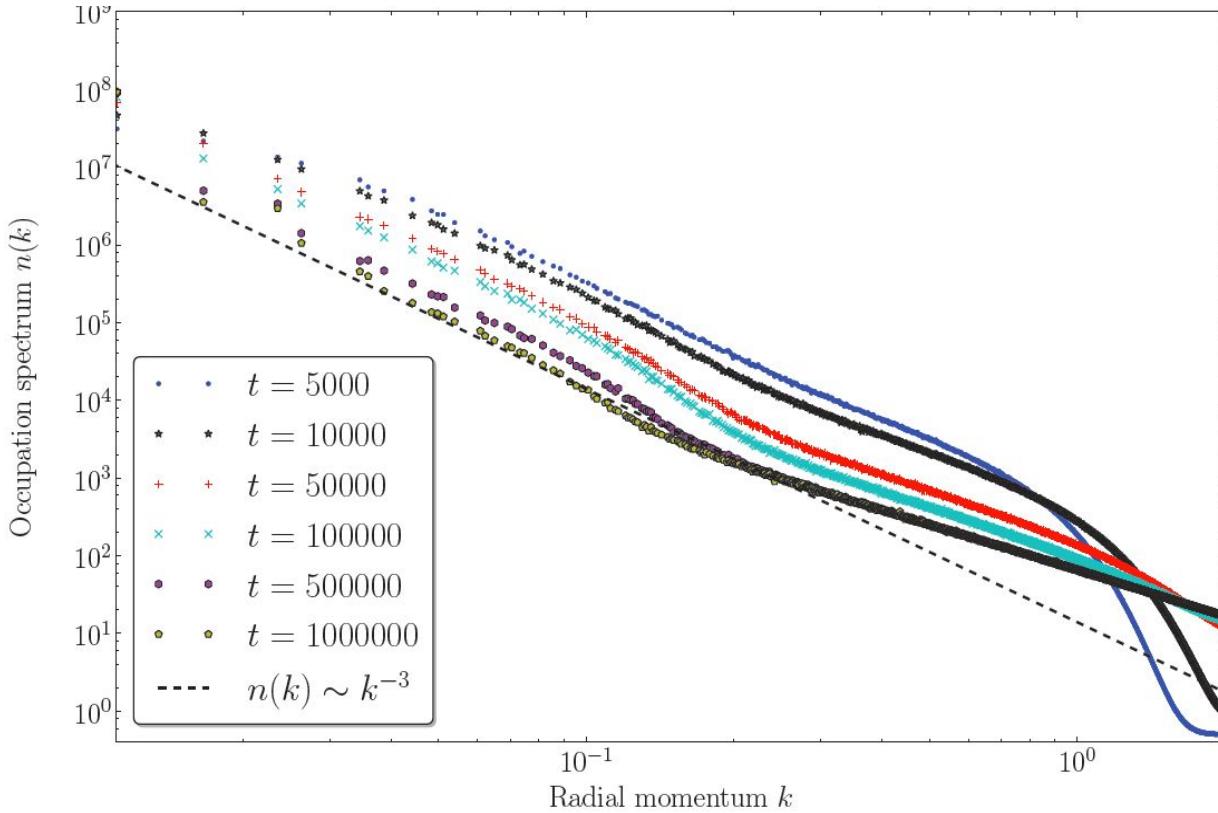
Domain formation in spin systems



[M. Karl, B. Nowak, and T. Gasenzer, unpublished]



2-component BEC



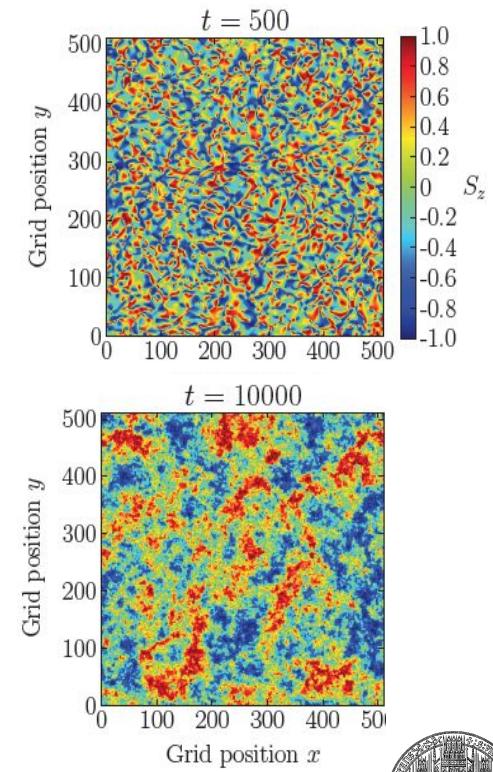
miscible

$$g_{12} = g$$



immiscible

$$g_{12} > g$$

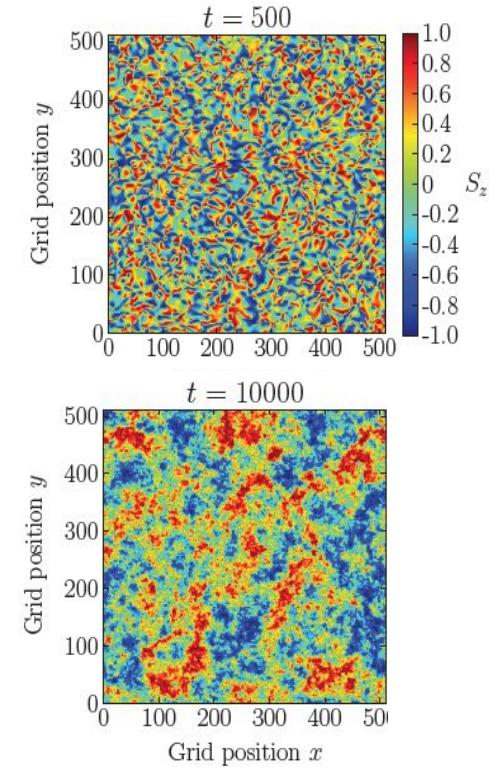
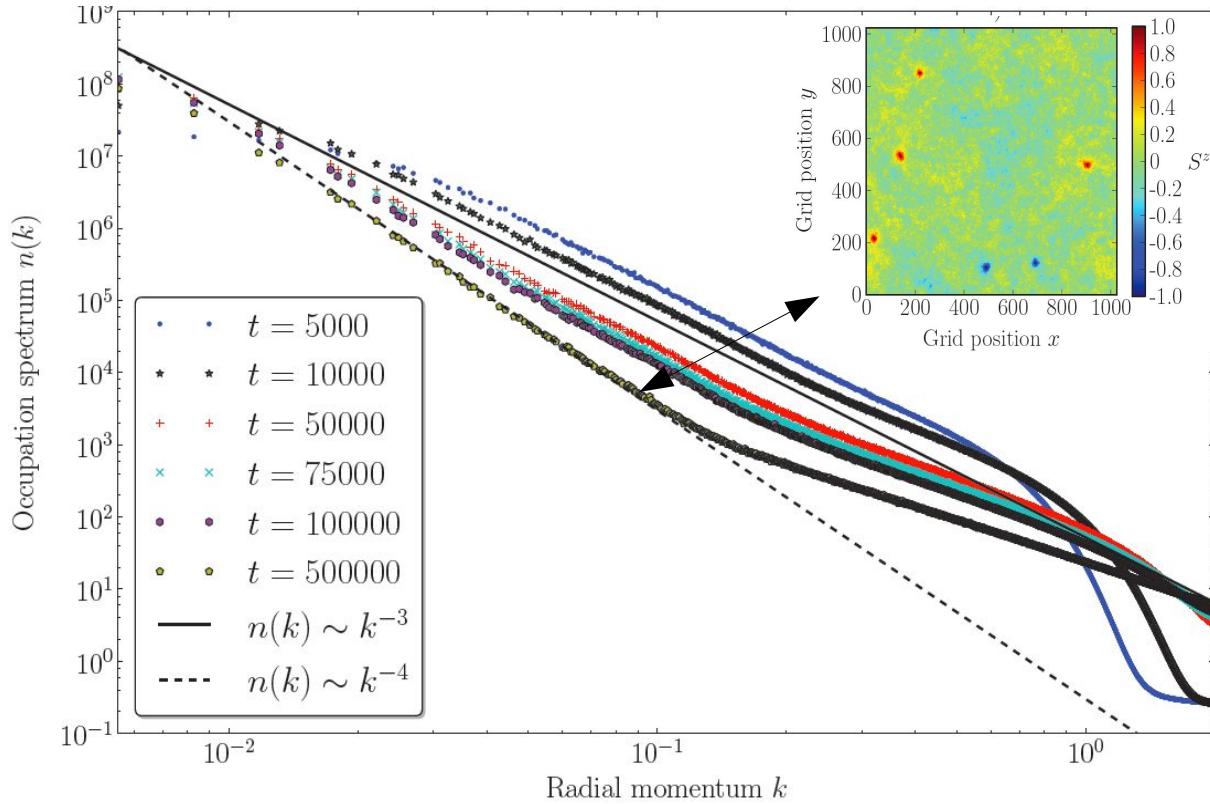


M. Karl, B. Nowak, TG, unpublished (12)



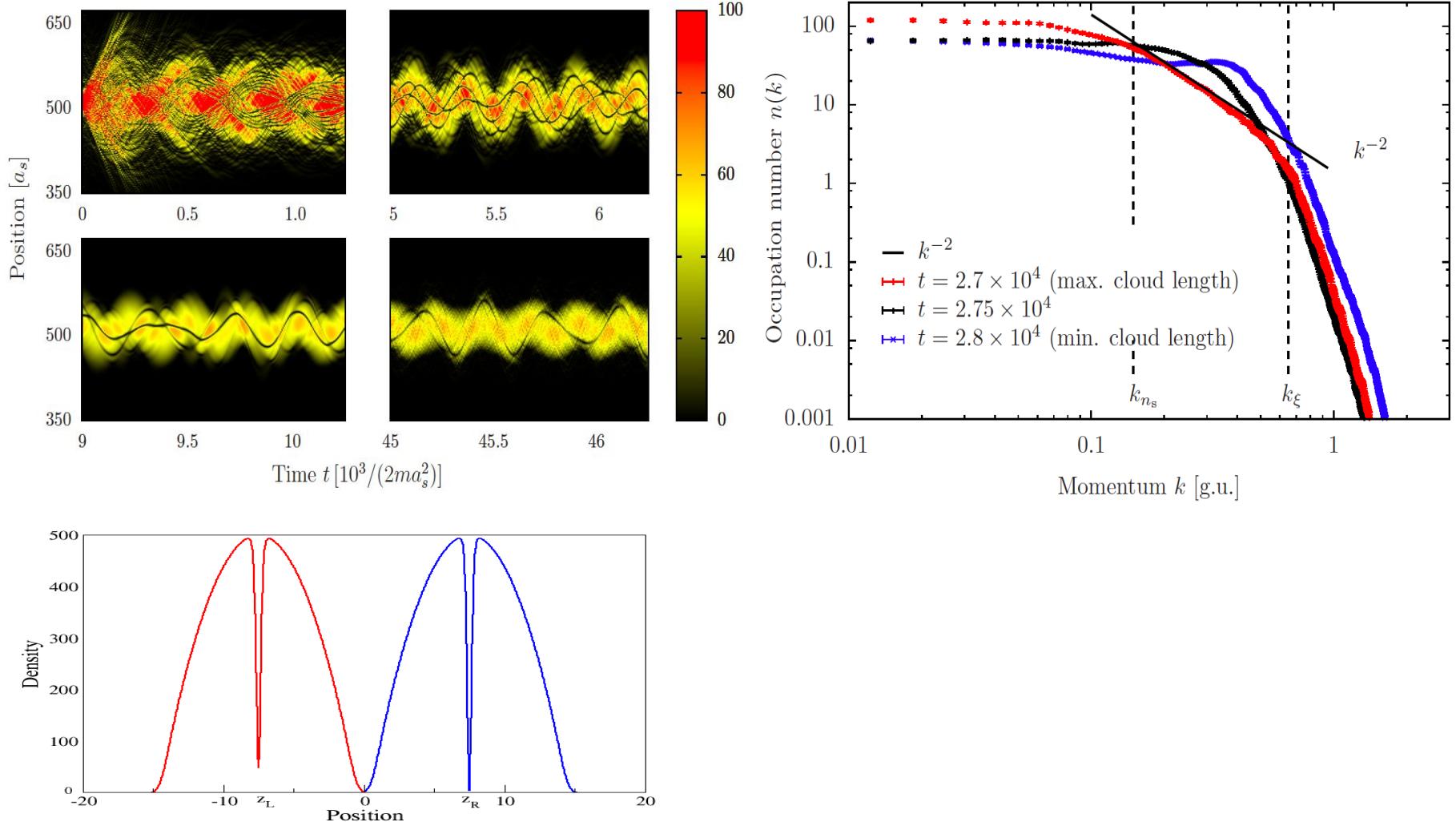
2-component BEC

miscible
 $g_{12} < g$



Solitons in 1D

Solitons in 1 spatial dimension



M. Schmidt, S. Erne, B. Nowak, D. Sexty, and TG, arXiv:1203.3651 [cond-mat.quant-gas]

QDYNAMICS12 · KITP Santa Barbara · 21 Aug 2012

Thomas Gasenzer



Relativistic scalar field

Non-linear Klein-Gordon equation

O(2) symmetry

$$(\partial_t^2 - \partial_x^2)\varphi(x, t) + \lambda\varphi^3(x, t) = 0$$

Initial condition: Highly occupied zero mode, Unoccupied modes with $k>0$

(video)

See also: <http://www.thphys.uni-heidelberg.de/~sextv/videos>

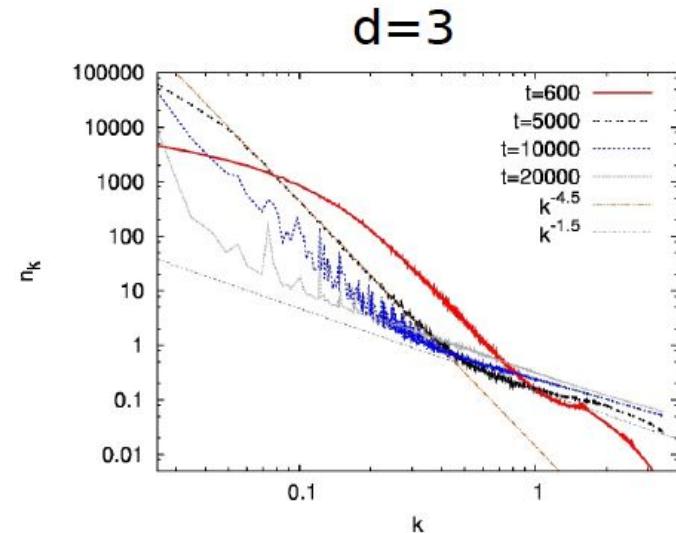
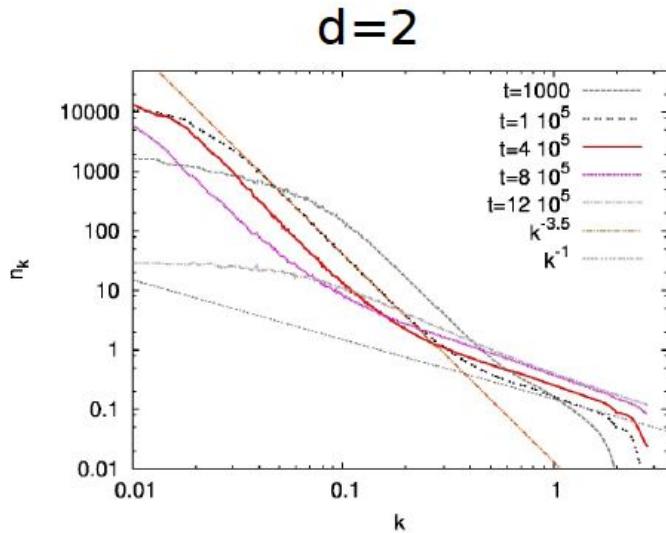
TG, B. Nowak, D. Sexty, arXiv:1108.0541 [hep-ph]



Relativistic simulations

Classical field equation:

$$[\partial_t^2 - \Delta + \Phi^2] \Phi_a = 0$$



reheating after inflation

S. Khlebnikov, I. Tkachev, PRL (96)

R. Micha, I. Tkachev, PRD (04)

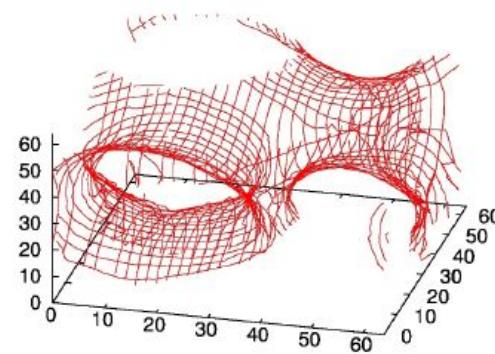
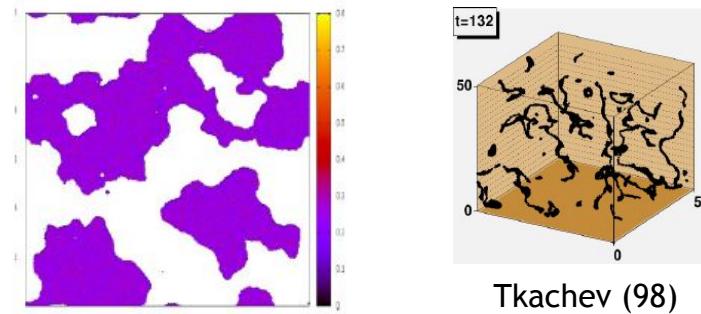
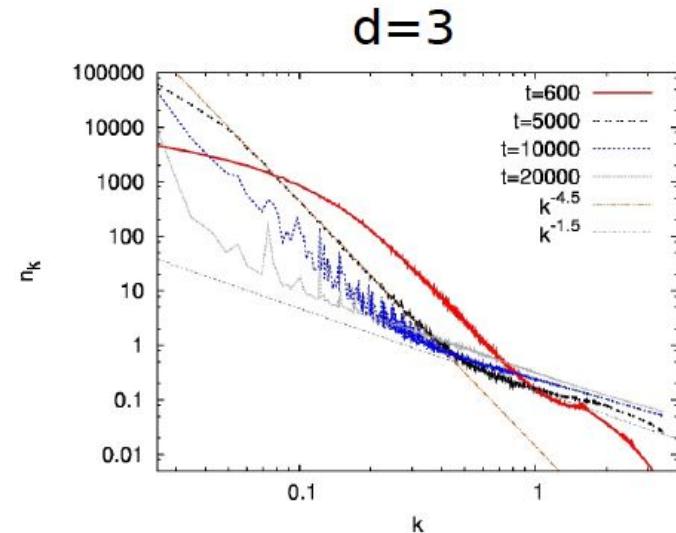
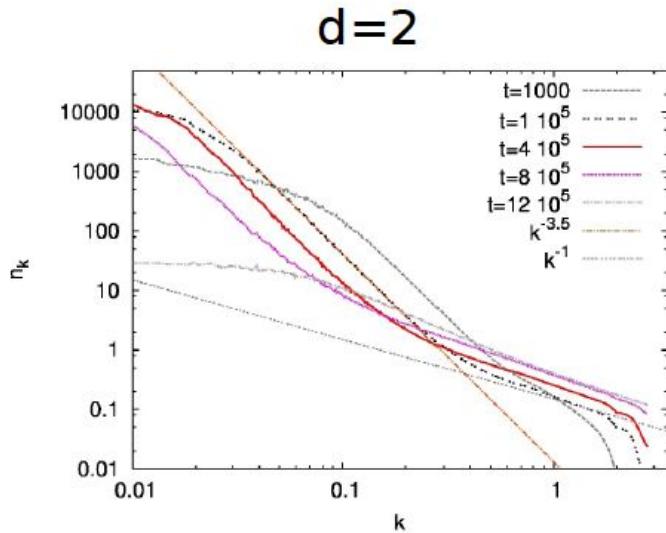
J. Berges, A. Rothkopf, J. Schmidt, PRL (08)

J. Berges, D. Sexty, PRD (11)



Relativistic simulations

Classical field equation: $\left[\partial_t^2 - \Delta + \Phi^2 \right] \Phi_a = 0$



D. Sexty, B. Nowak, TG, PLB 710 (12)



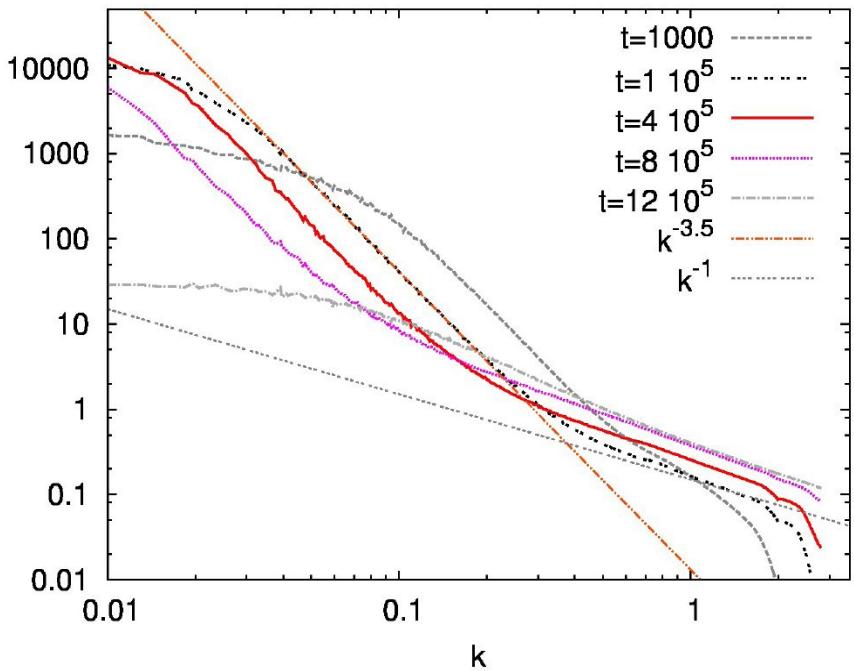
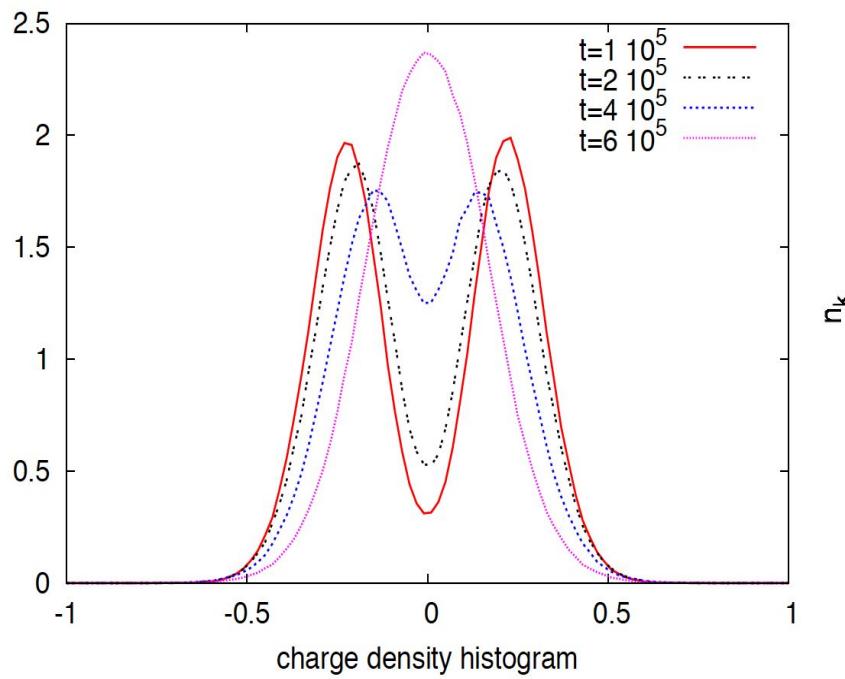
Strong Turbulence = Charge Separation

Charge density distribution

vs.

power spectrum

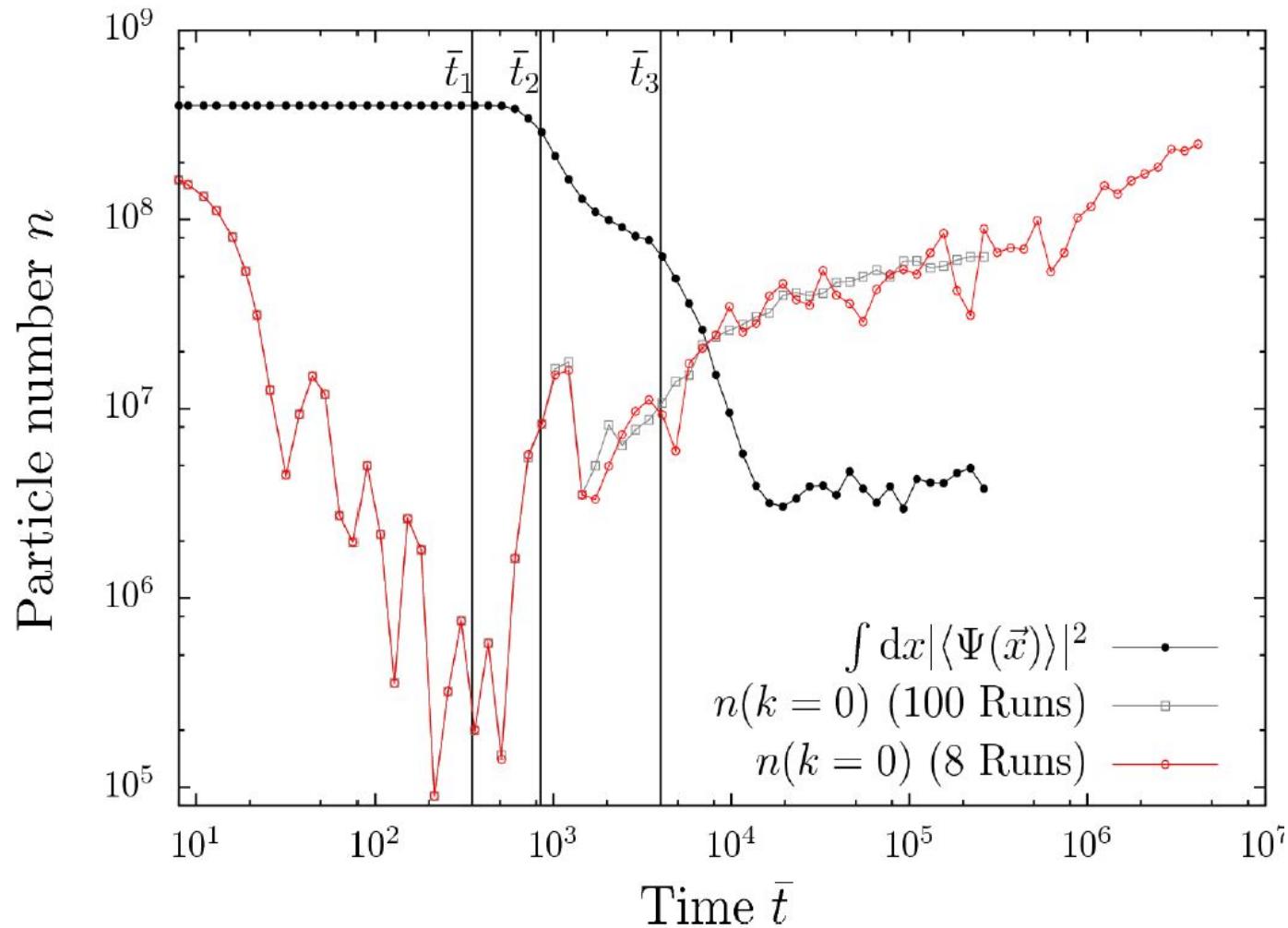
($d = 2, N = 2$)



TG, B. Nowak, D. Sexty, arXiv:1108.0541 [hep-ph], PLB, to appear



Bose-Einstein condensation

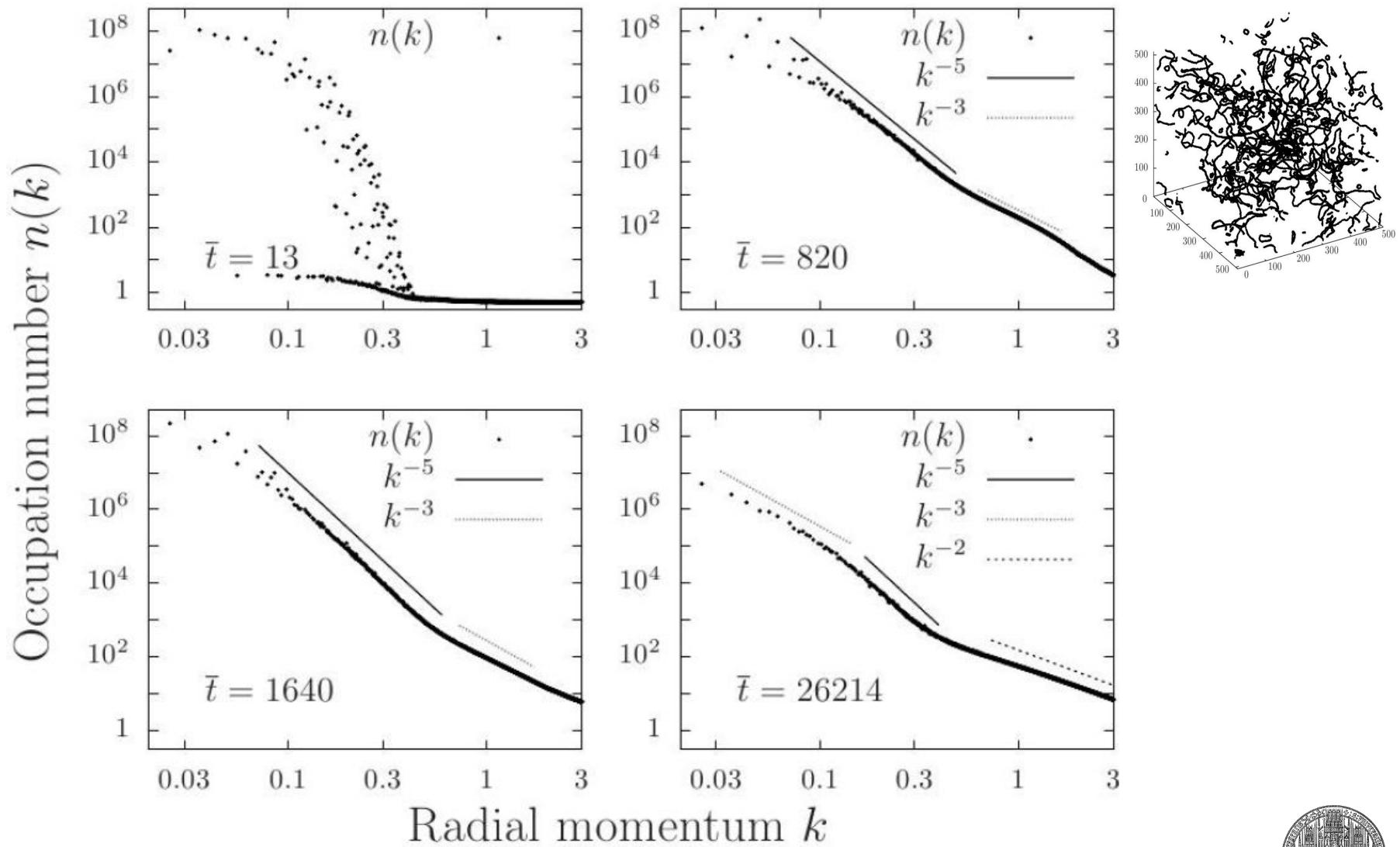


J. Schole, B. Nowak, D. Sexty, TG (unpublished)

For 3D see also N. Berloff & B. Svistunov, PRA 66 (02)



3+1 D simulations



Acoustic Turbulence

Decomposition of Energy

$$E_{tot} = \int \left(\frac{1}{2} |\nabla \sqrt{n} e^{-i\varphi}|^2 + \frac{1}{2} g n^2 \right) d\rho$$

$$= E_{kin} + E_q + E_{int}$$

$$\mathbf{u}(\rho, t) = \nabla \varphi(\rho, t)$$

$$E_{kin} = \frac{1}{2} \int |\sqrt{n} \mathbf{u}|^2 d\rho = E_{kin}^i + E_{kin}^c$$

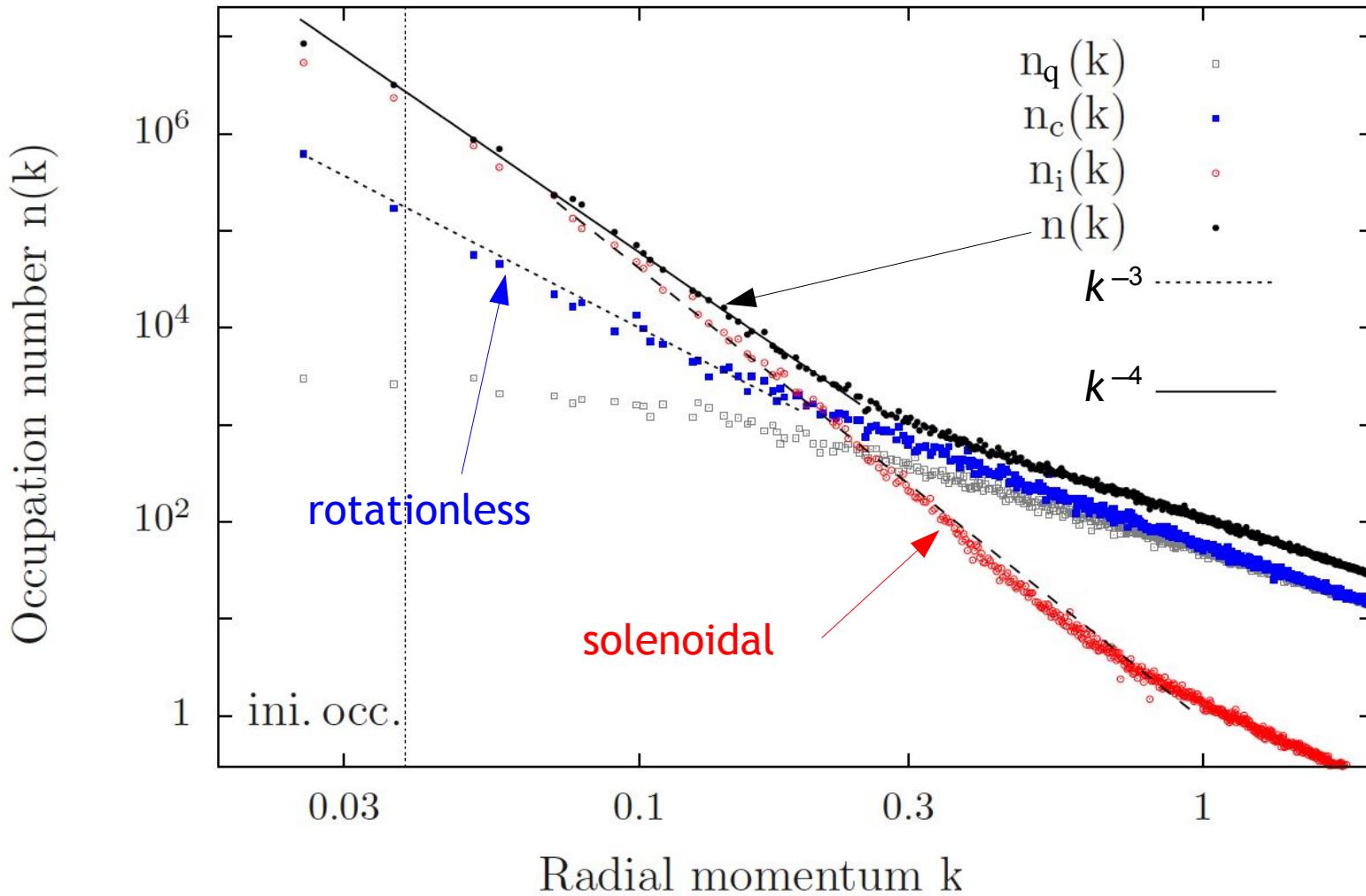
$$\nabla \times (\sqrt{n} \mathbf{u})^c = 0$$

$$\nabla \cdot (\sqrt{n} \mathbf{u})^i = 0$$

$$E_q = \frac{1}{2} \int (\nabla \sqrt{n})^2 d\rho$$

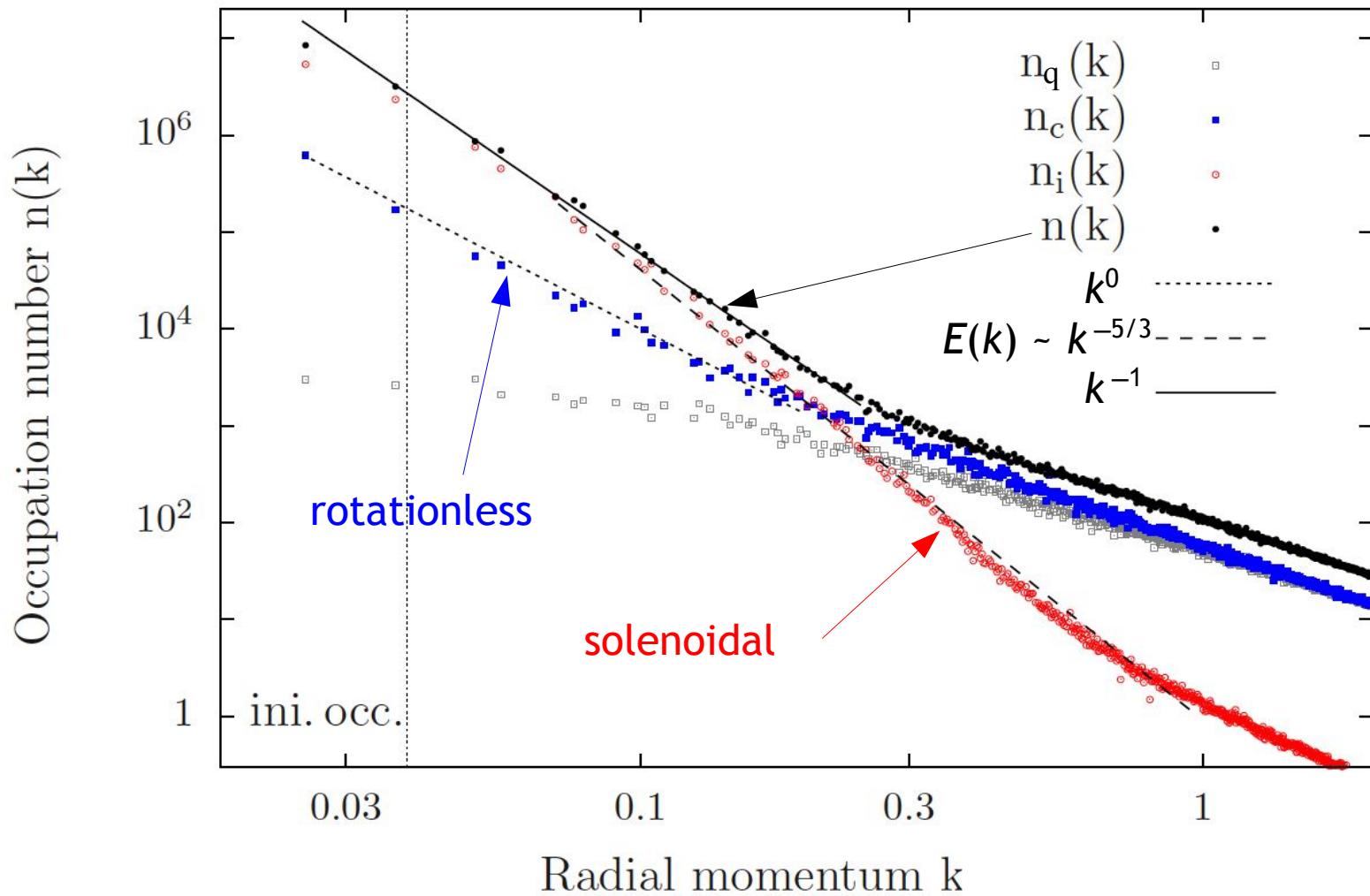


Decomposition of flow



Simulations in 2+1 D

$$E(k) = \omega(k) k^{d-1} n(k)$$

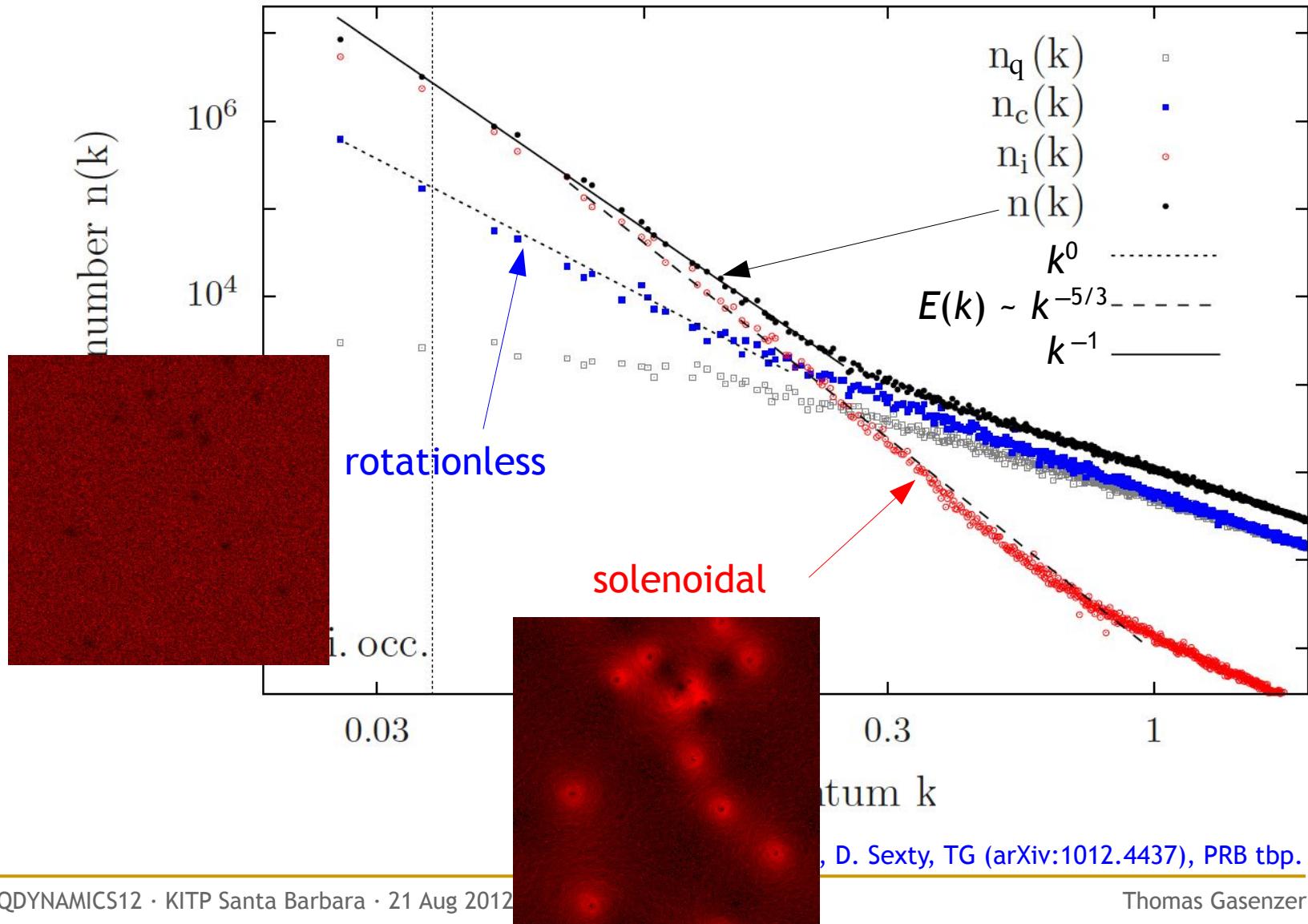


B. Nowak, D. Sexty, TG (arXiv:1012.4437), PRB tbp.



Simulations in 2+1 D

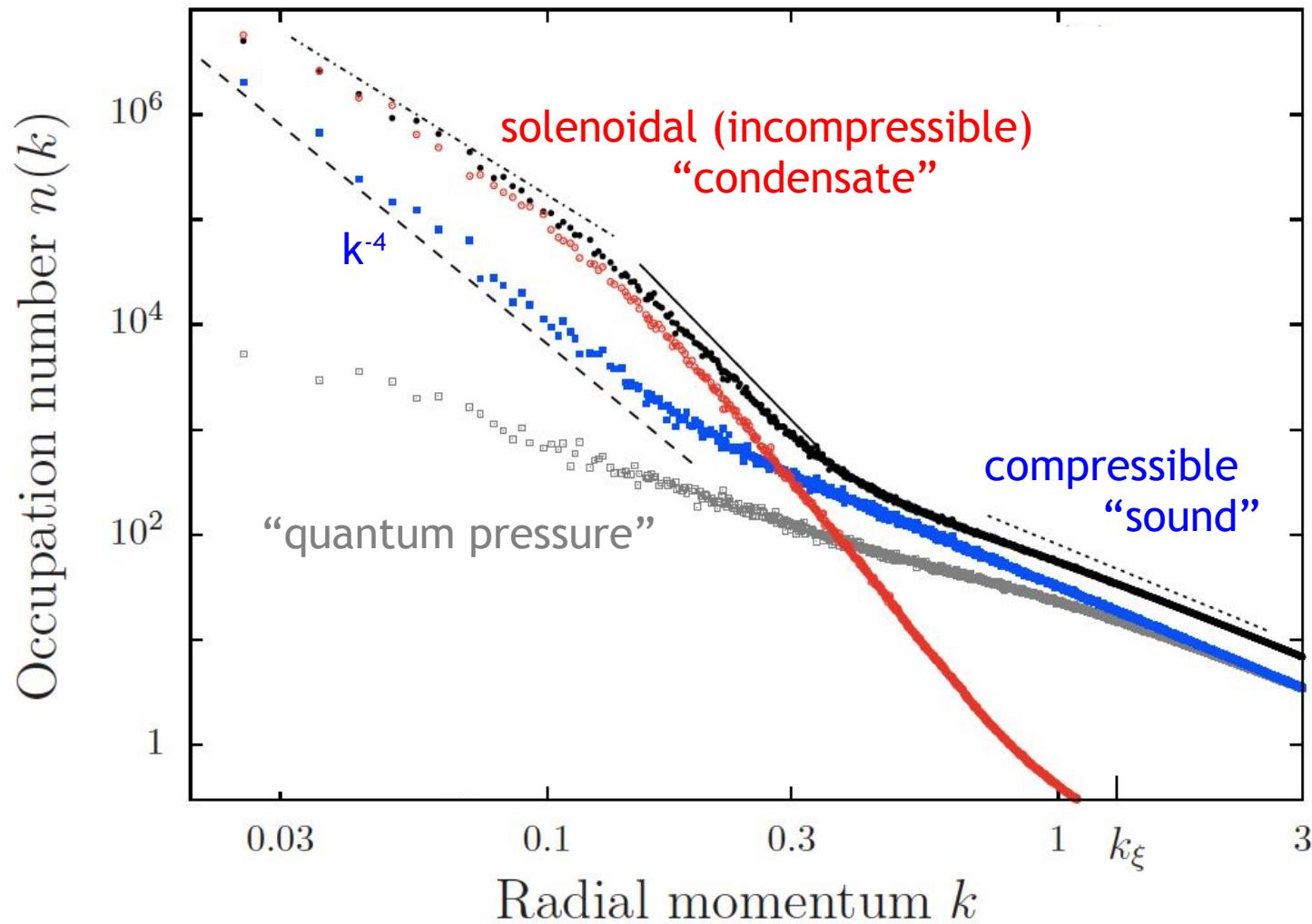
$$E(k) = \omega(k) k^{d-1} n(k)$$



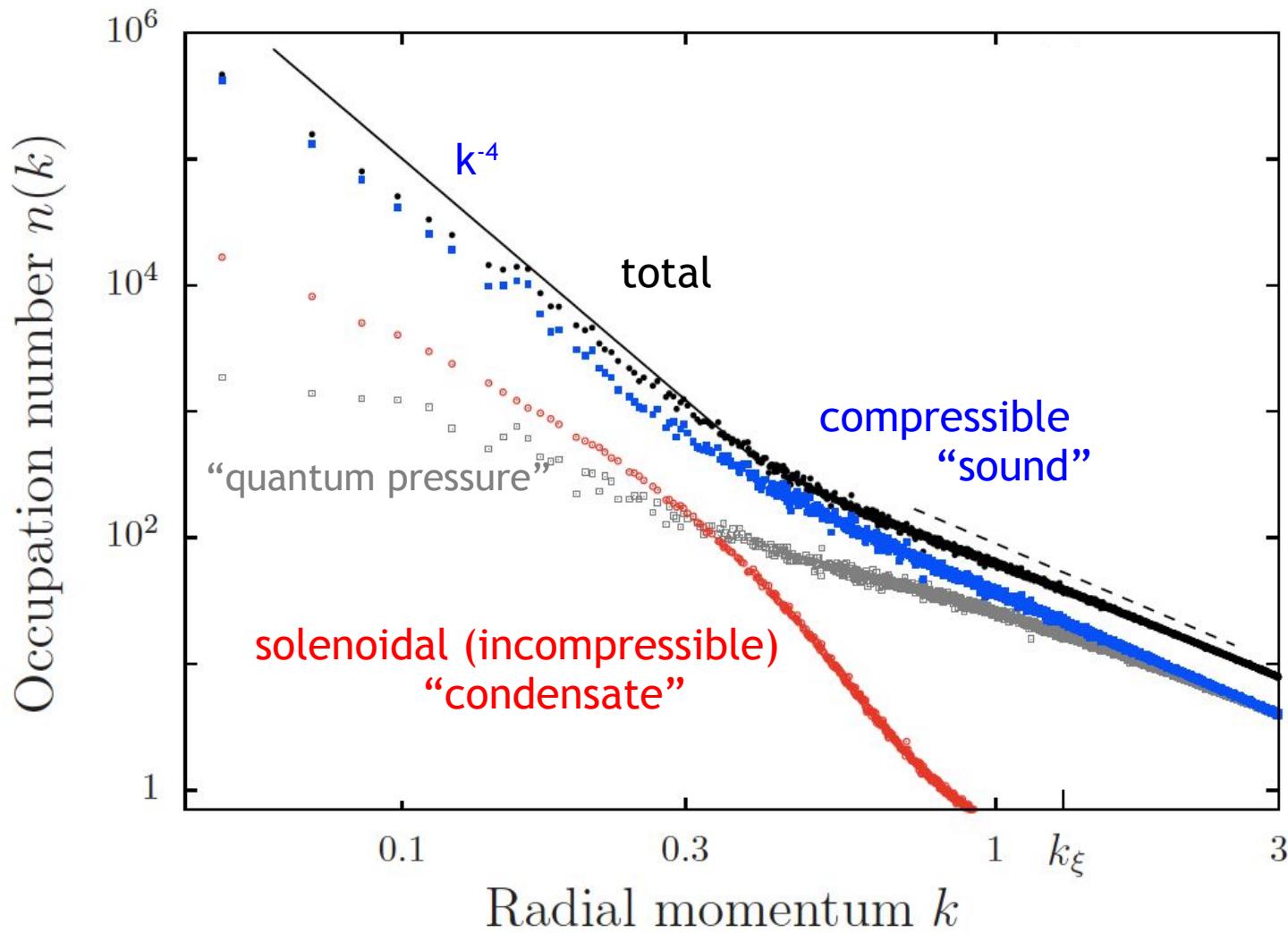
, D. Sexty, TG (arXiv:1012.4437), PRB tbp.



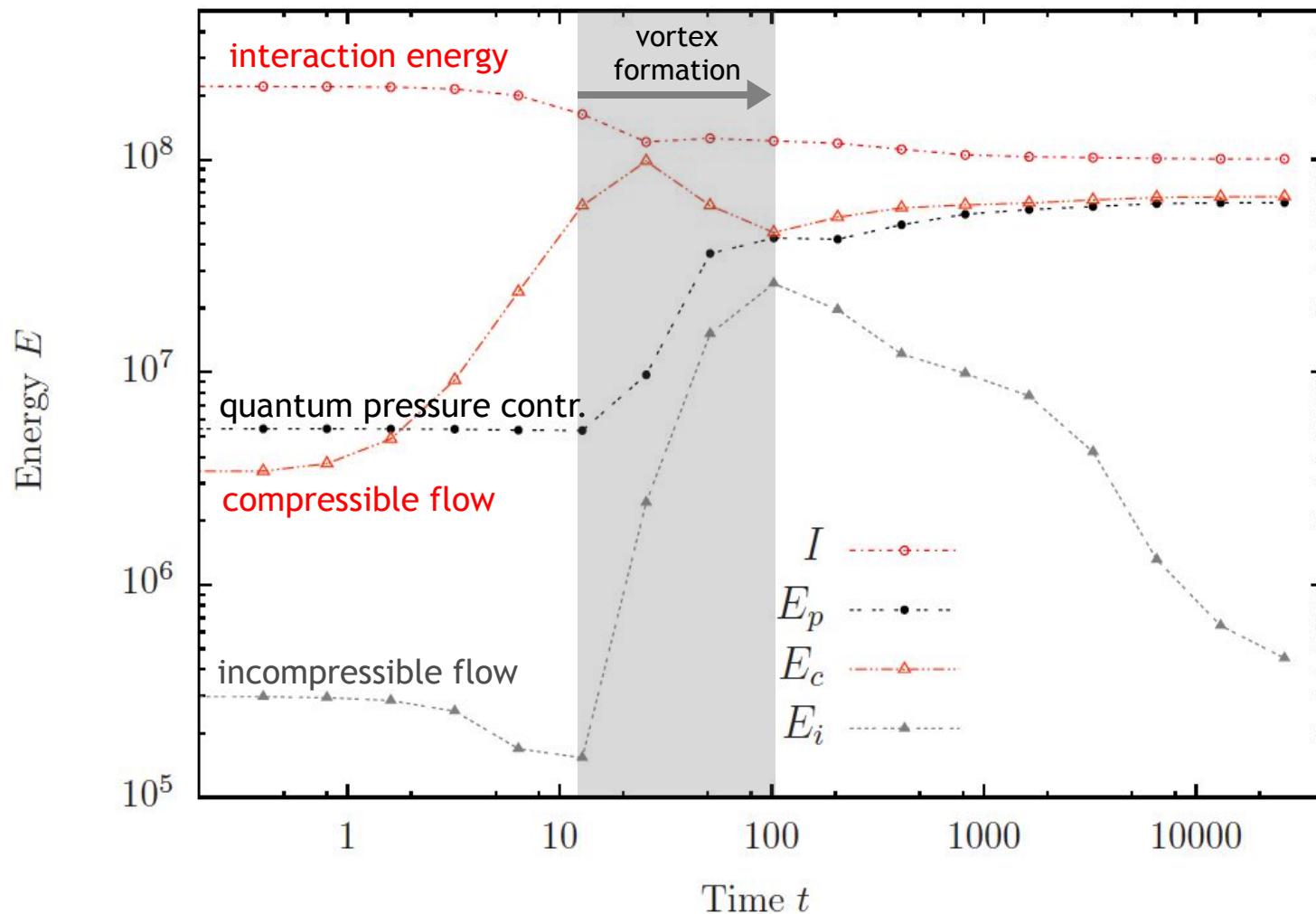
Decomposition of flow



Acoustic turbulence



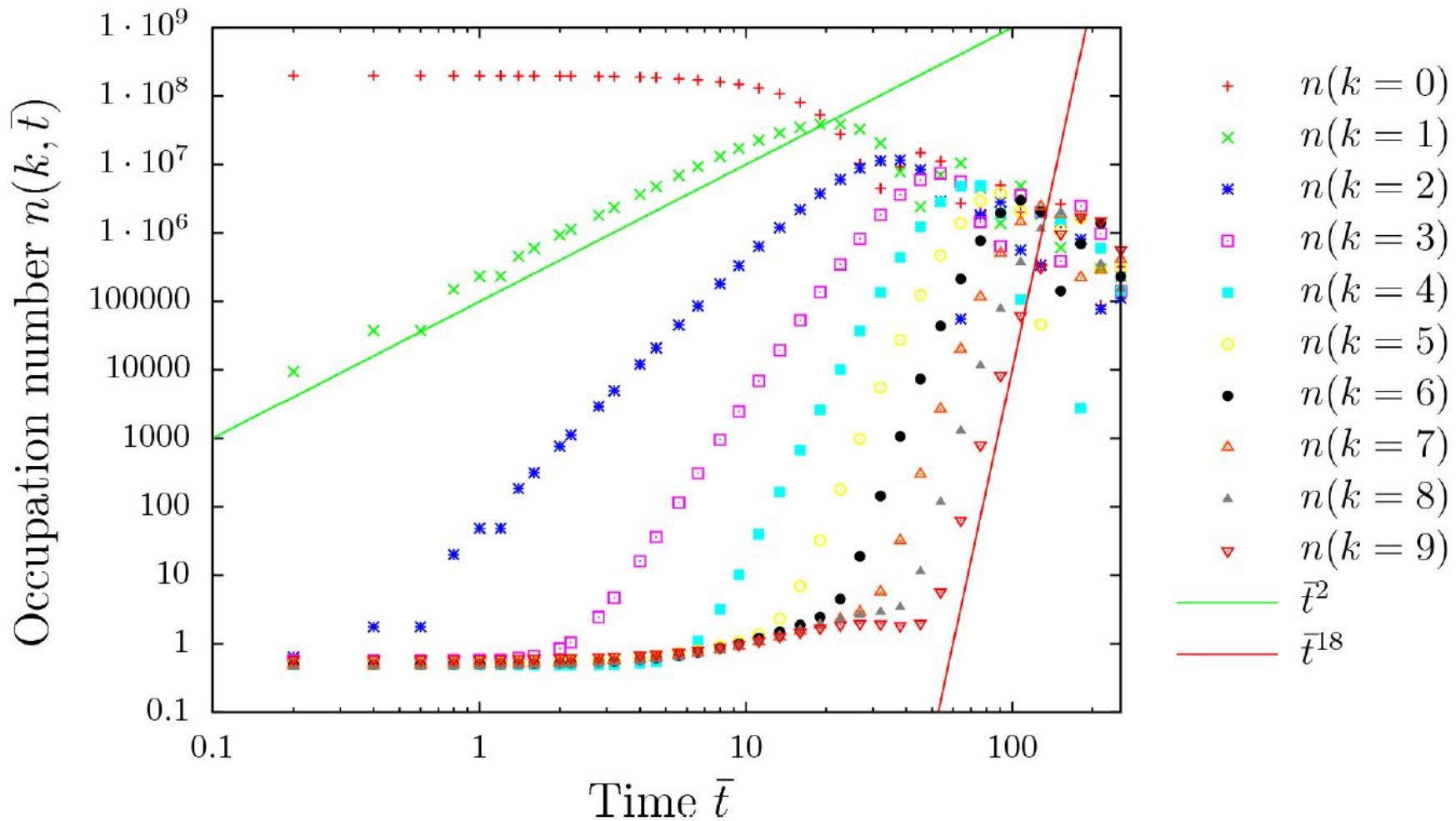
Time evolution of Energy Components (3+1 D)



J. Schole, B. Nowak, D. Sexty, TG (unpublished)



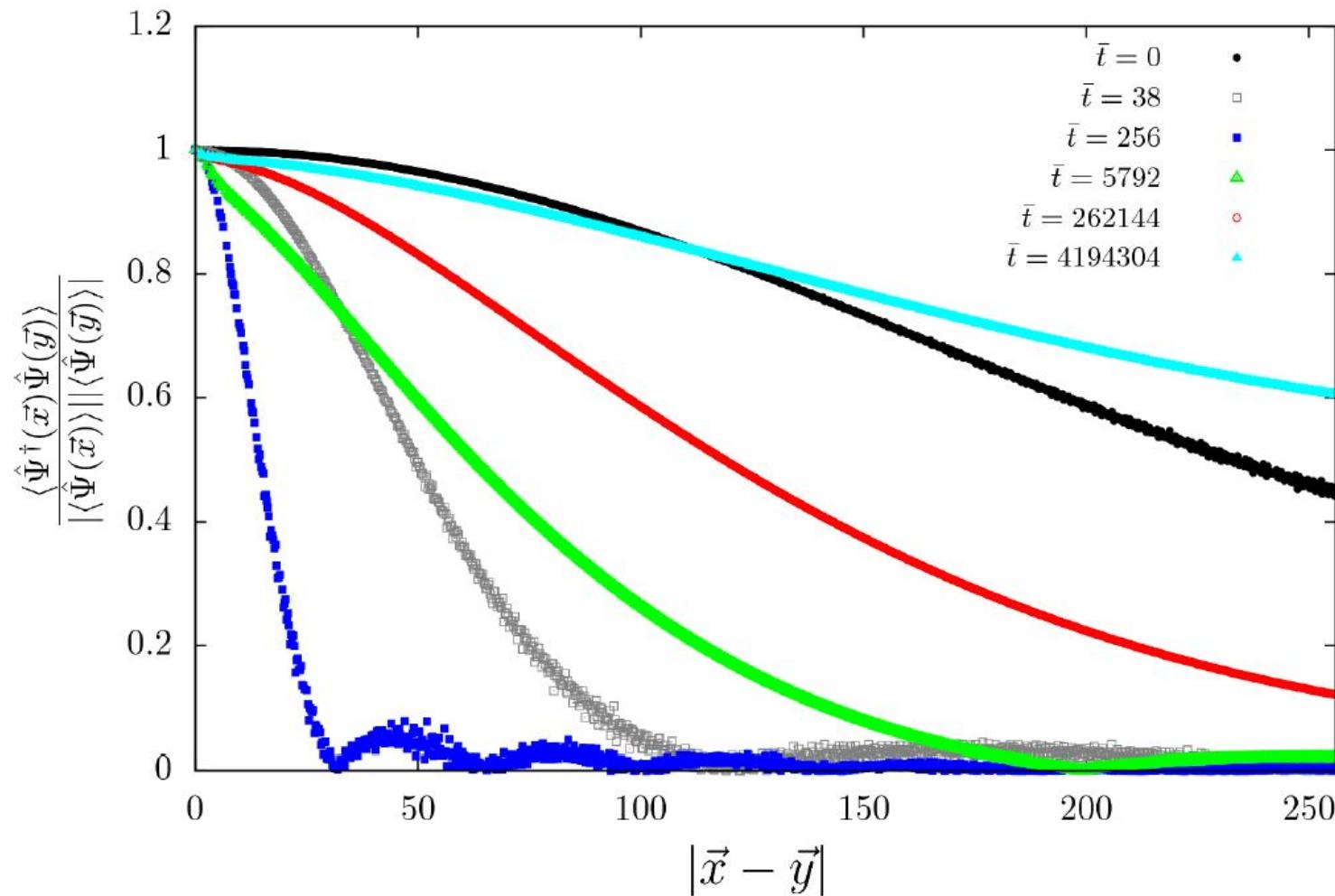
Mode Occupations



J. Schole, B. Nowak, D. Sexty, TG (unpublished)



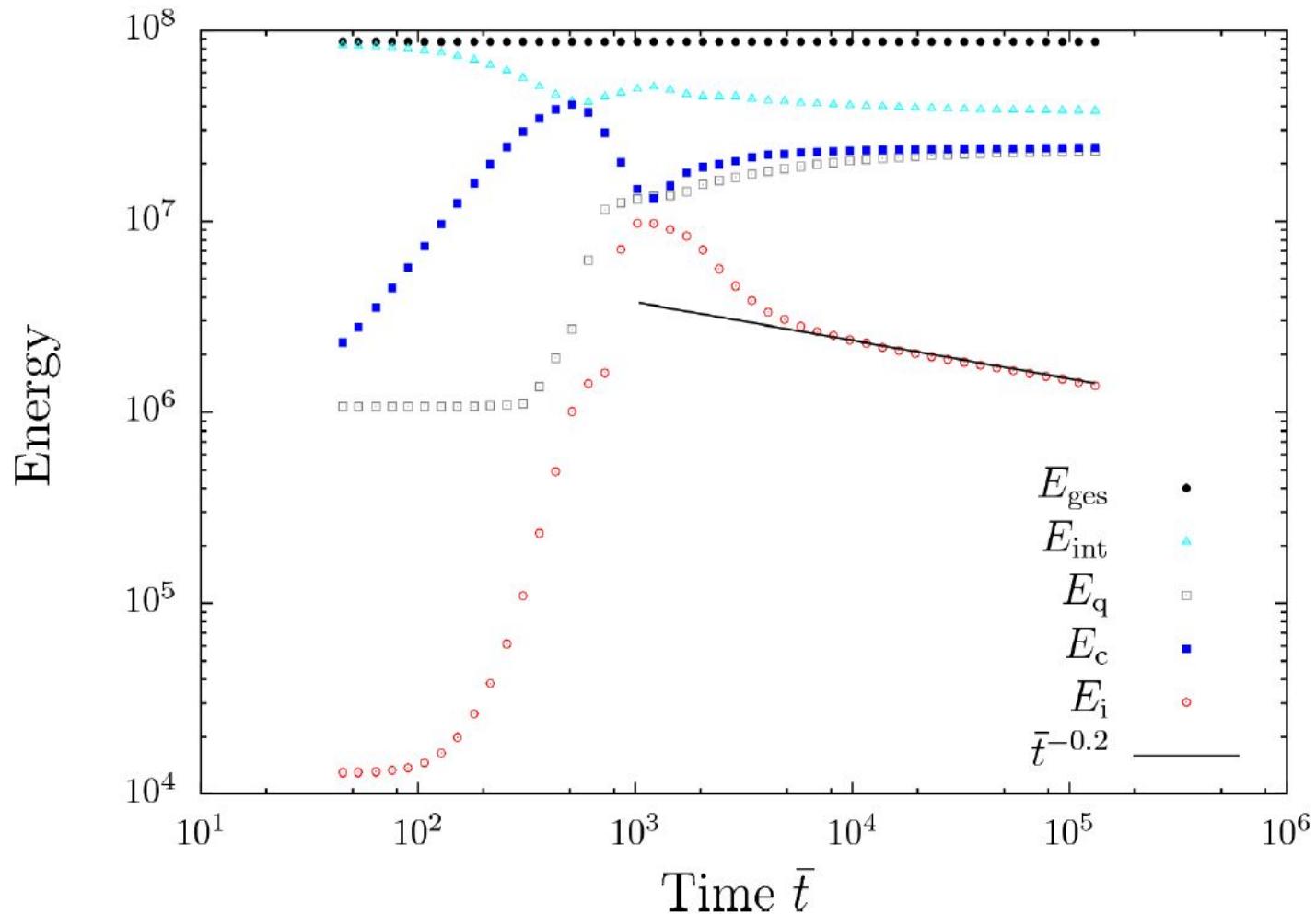
1st-order Coherence



J. Schole, B. Nowak, D. Sexty, TG (unpublished)



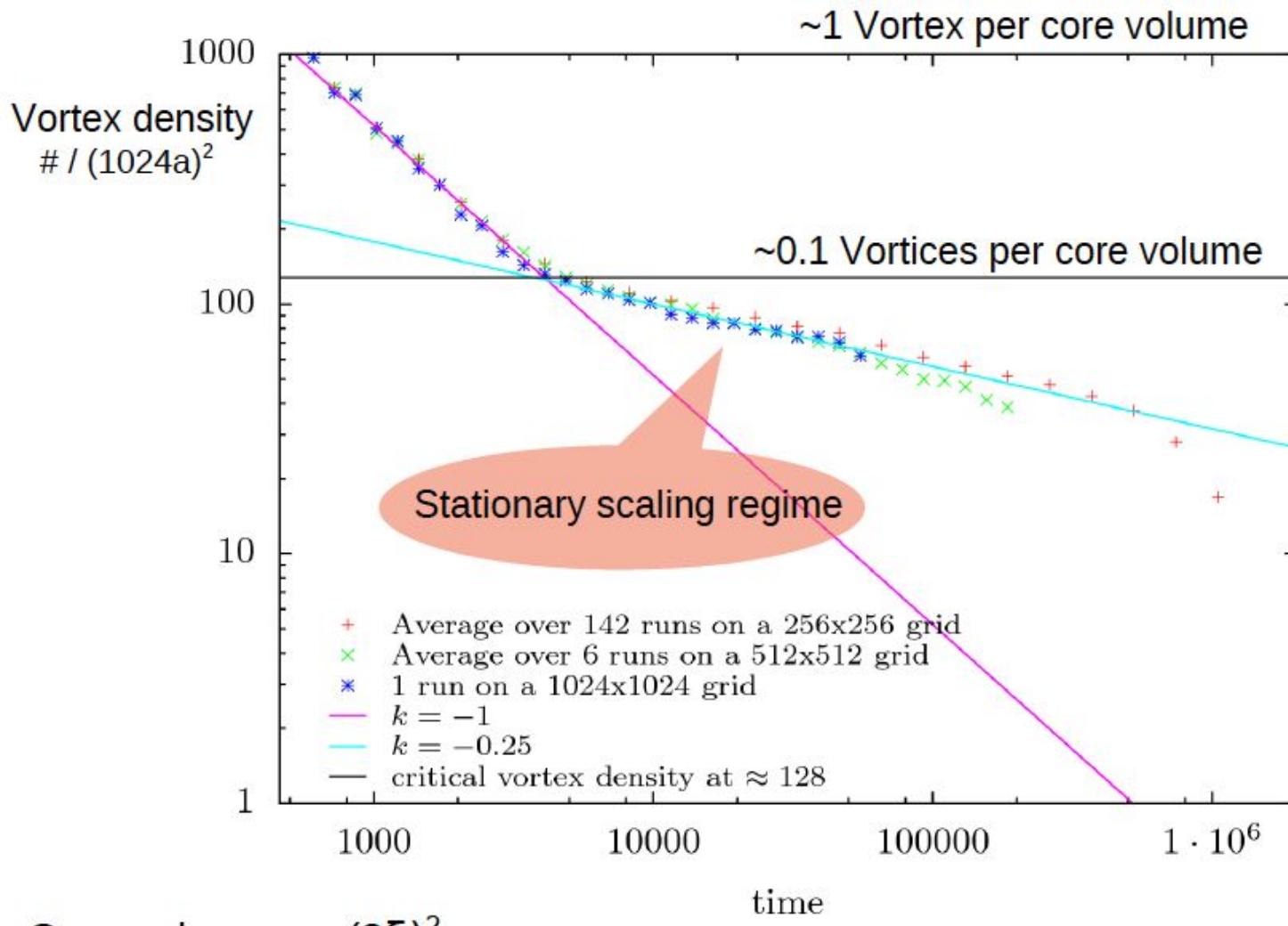
Time-Evolution of Energy-components (2+1 D)



J. S., B. Nowak, D. Sexty, T. Gasenzer (unpublished)



Time evolution of vortex density



Core volume $\sim \pi(3\xi)^2$

J. Schole, B. Nowak, D. Sexty, TG (unpublished)



Enstrophy in classical turbulence

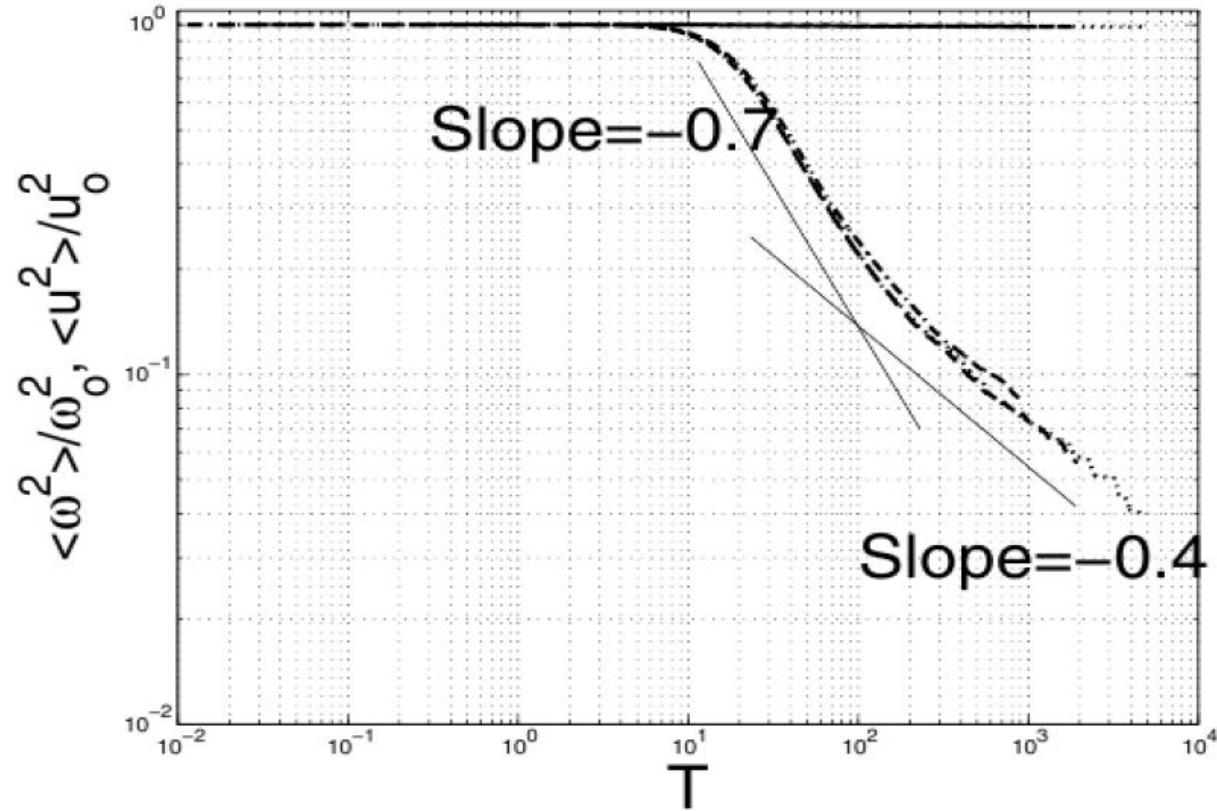
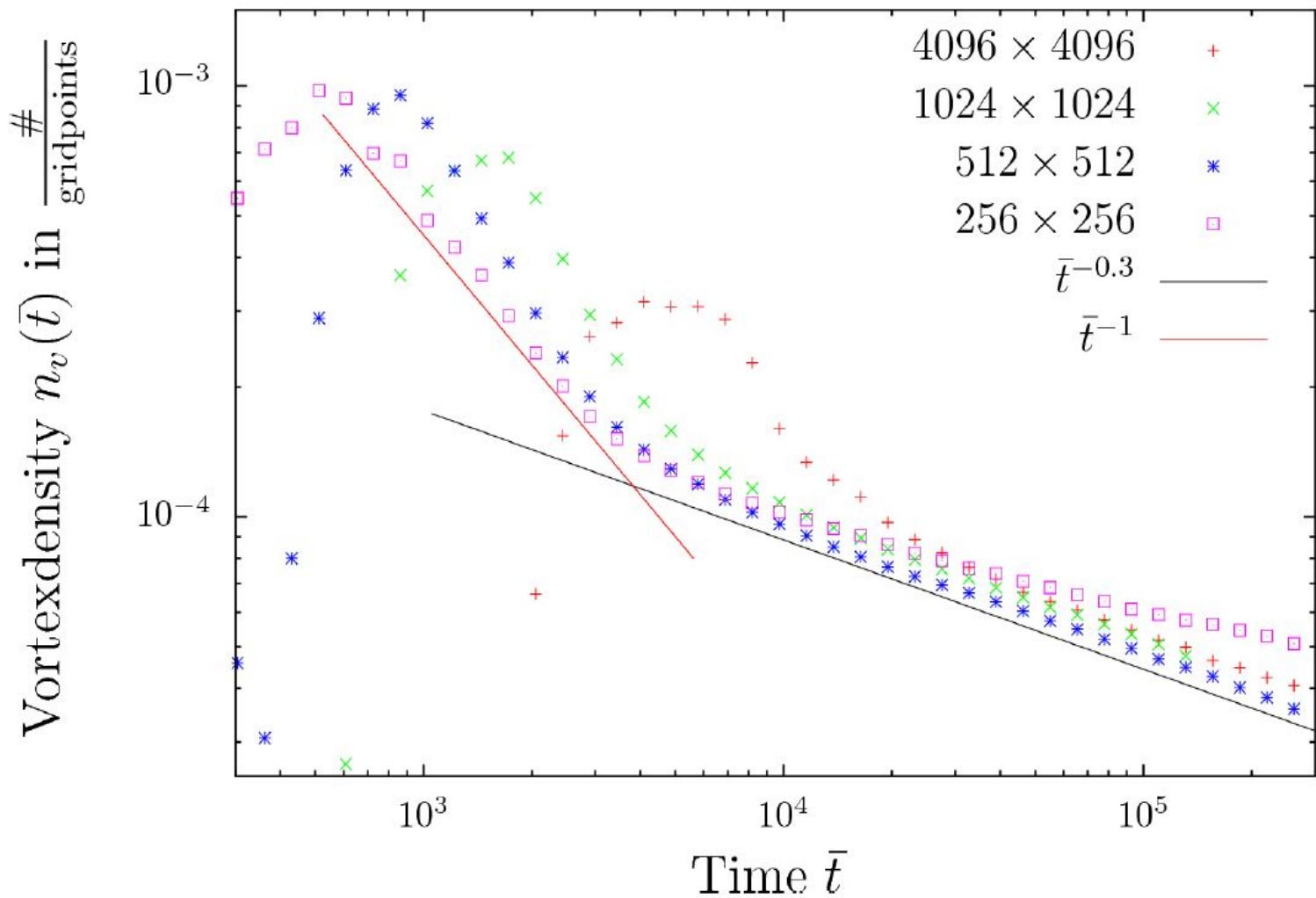


FIG. 1. Time evolution of energy (upper horizontal curve) and enstrophy. Resolution: dotted line (512^3); dashed line (1024^3); dash-dotted line (2048^3).

V. Yakhot, J. Wanderer, PRL 93:154502



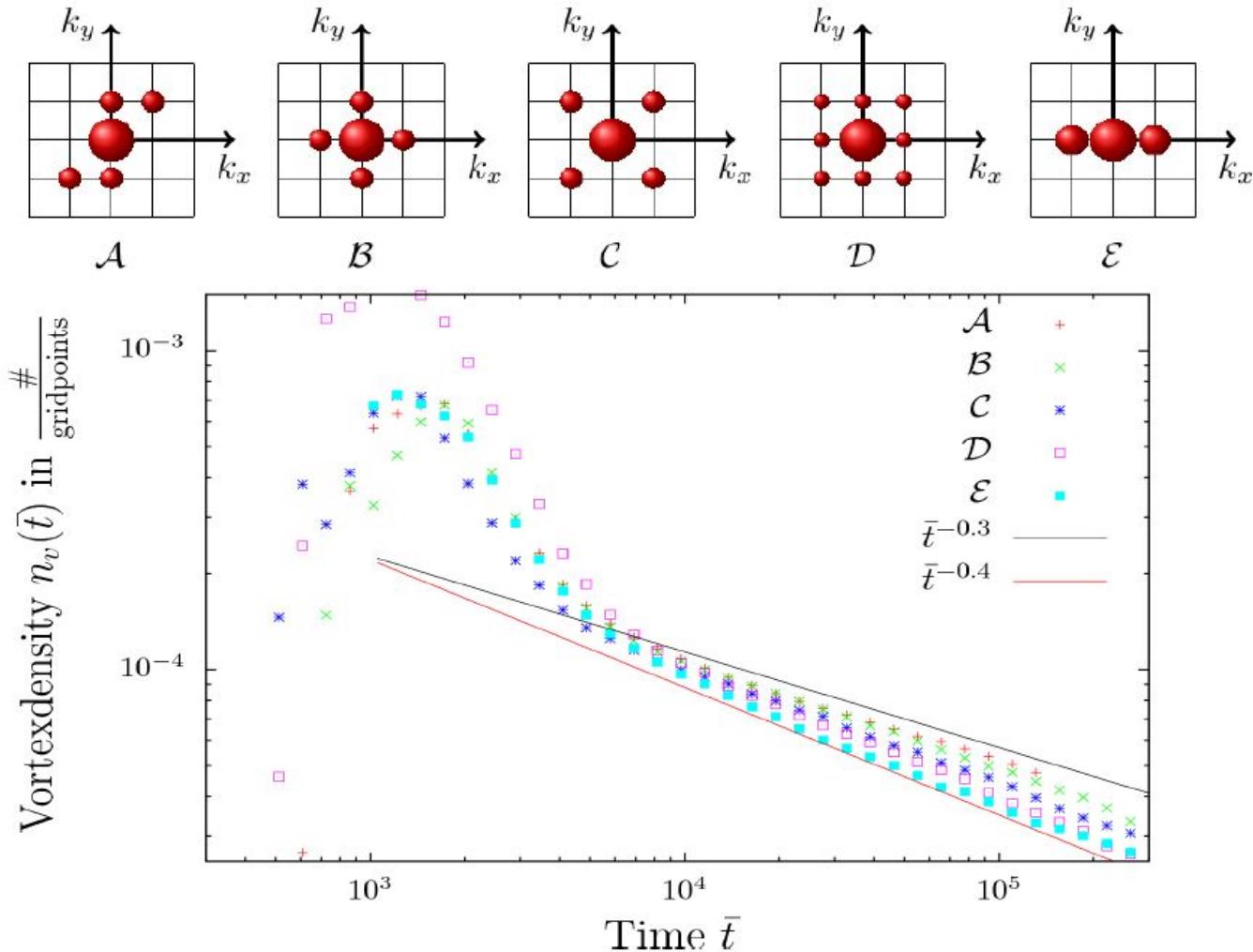
Vortex-Density Decay in 2d



J. S., B. Nowak, D. Sexty, T. Gasenzer (unpublished)



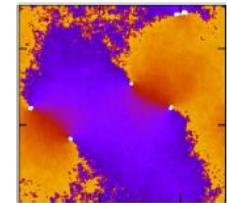
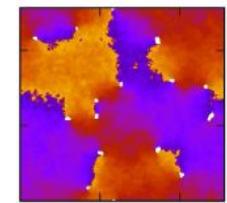
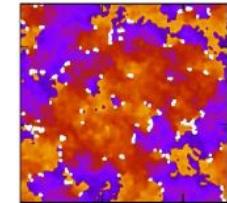
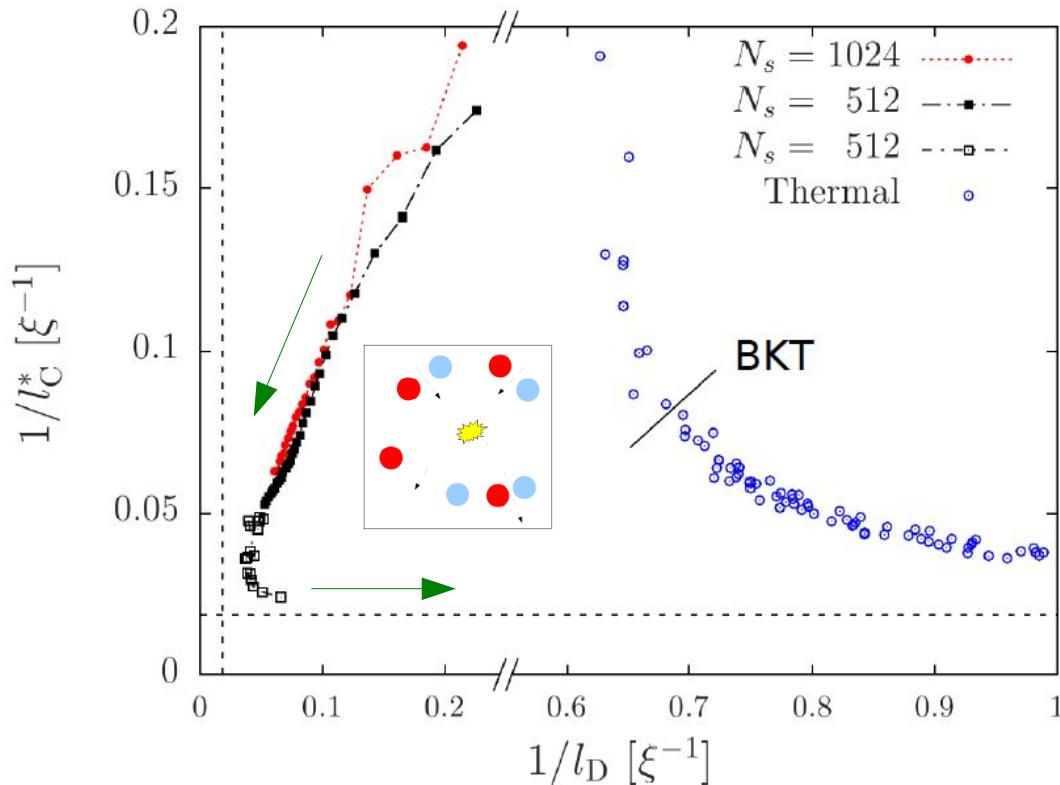
Vortex-Density Decay in 2d



J. S., B. Nowak, D. Sexty, T. Gasenzer (unpublished)



Approach of the NTFP



l_c^* = Phase coherence length

l_d = Vortex-Antivortex pair distance

J. Schole, B. Nowak, TG, PRA 86 (12) 013624



Kinetic Theory

One of the power laws can be explained by a kinetic theory:

$$\partial_t n_v(t) = -\frac{n_{\text{dip}}}{\tau_{\text{ann}}}$$

$$n_{\text{dip}} \sim n_v$$

$$\sigma \sim d$$

d : average pair distance

$$\tau_{\text{ann}} = \tau_{\text{coll}} \alpha$$

$$\bar{v} = \frac{1}{d}$$

\bar{v} : average pair velocity

$$\tau_{\text{coll}} = \frac{l}{\bar{v}}$$

$$d = \frac{1}{\sqrt{n_v}}$$

l : mean free path

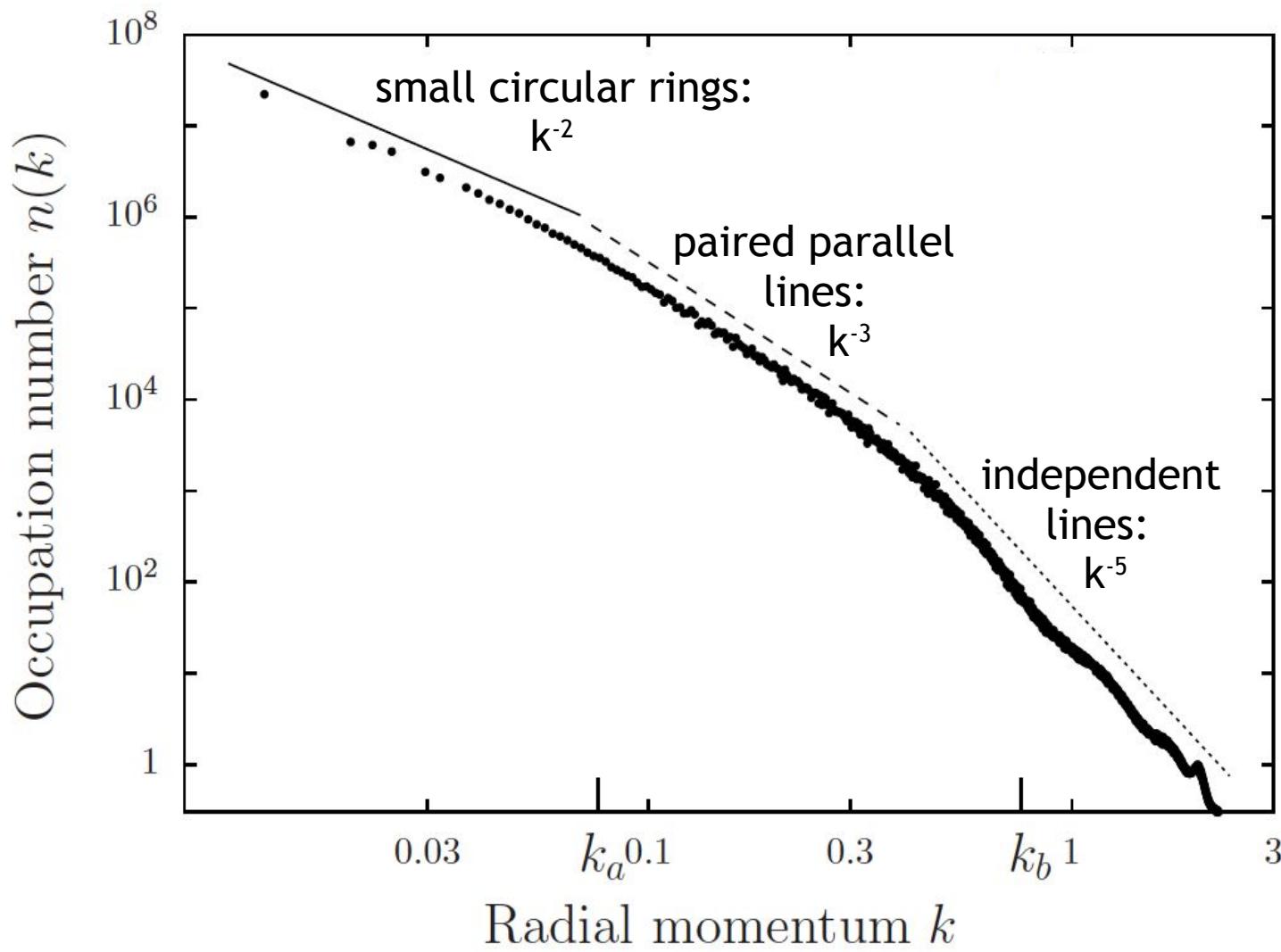
$$l \sim \frac{1}{n_v \sigma}$$

$$\Rightarrow \partial_t n_v(t) \sim -n_v^2 \Rightarrow n_v(t) \sim t^{-1}$$

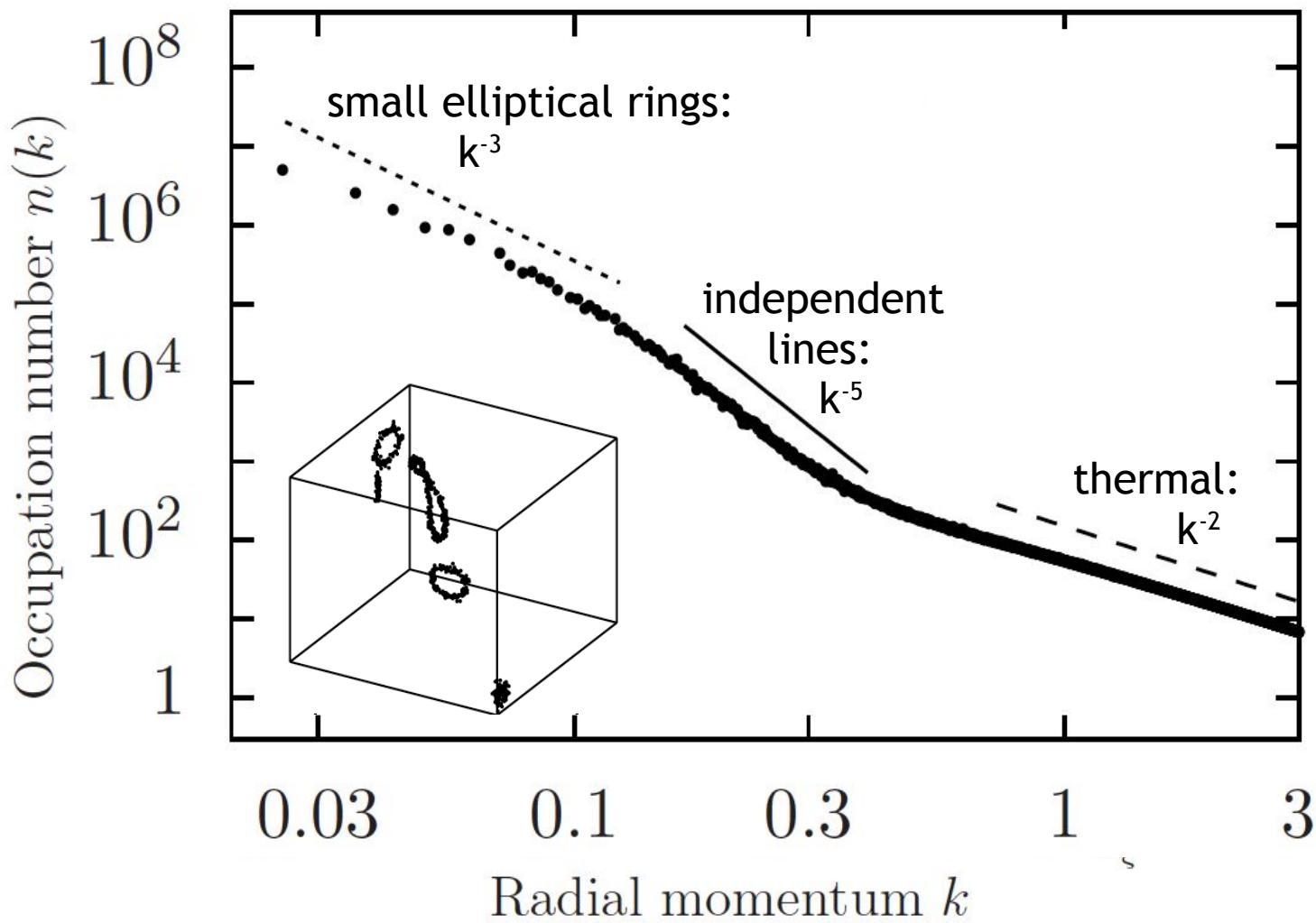
This result is valid under the assumption that the vortices are moving in pairs and that the pairs are homogeneously distributed.



Line vortex model in 3+1 D

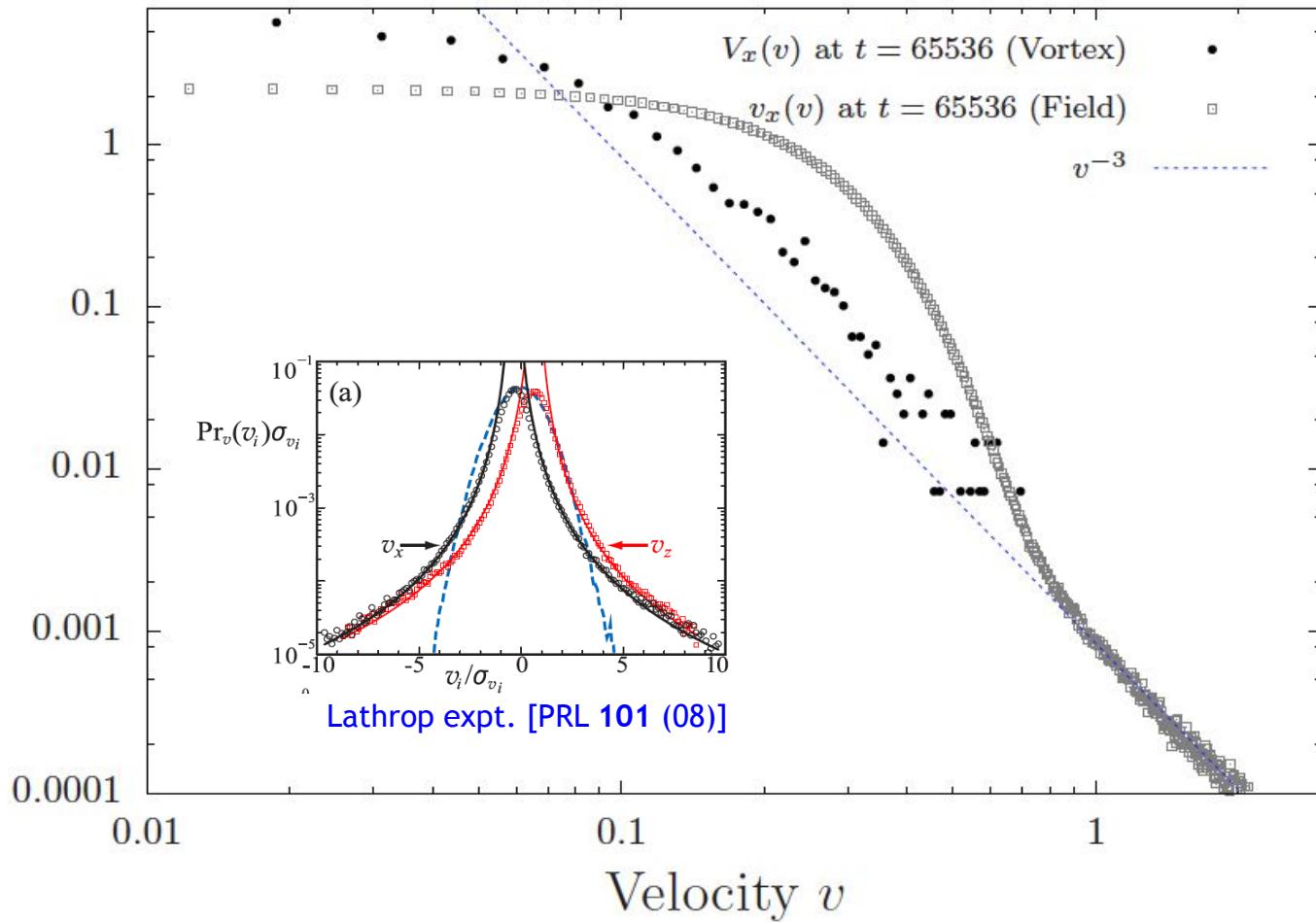


Simulations in 3+1 D



Vortex velocity distribution

Probability distribution $P(v)$



J. Schole, B. Nowak, D. Sexty, TG (unpublished)
s. also C.F. White et al., PRL 104 (10); I.A. Min, Phys. Fluids 8 (96)



Velocity distributions

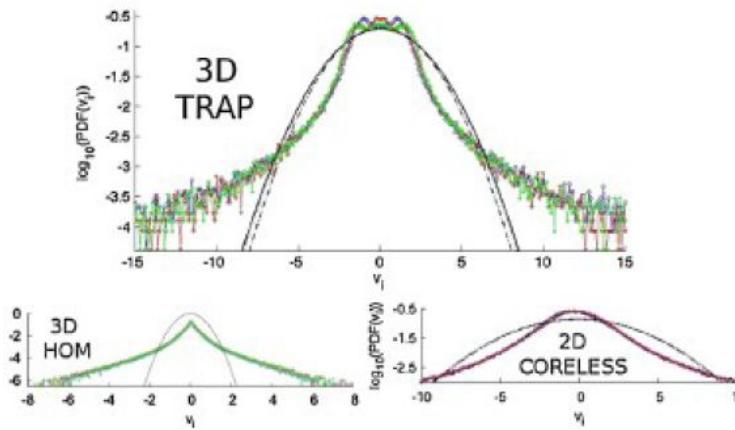
Paoletti et al. PRL 101, 154501 (2008):

Power law tails distinguish classical turbulence from classical turbulence.

Min et al. Phys. Fluids 8, 1169 (1996), White et al. PRL 104, 075301 (2010):

Point vortices: Power law tails

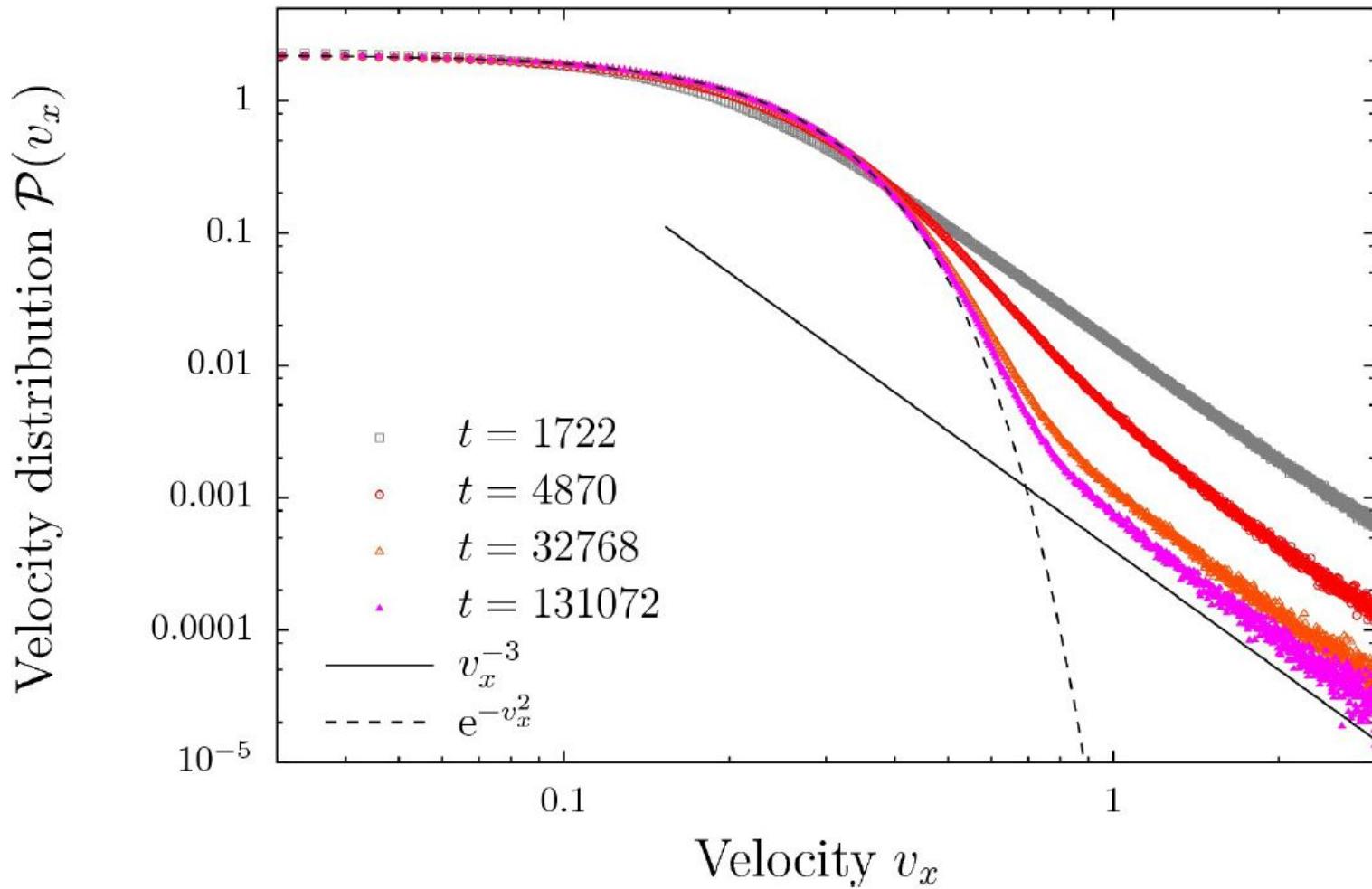
Vorticity patches: Gaussian distributions



White et al. PRL 104, 075301 (2010)



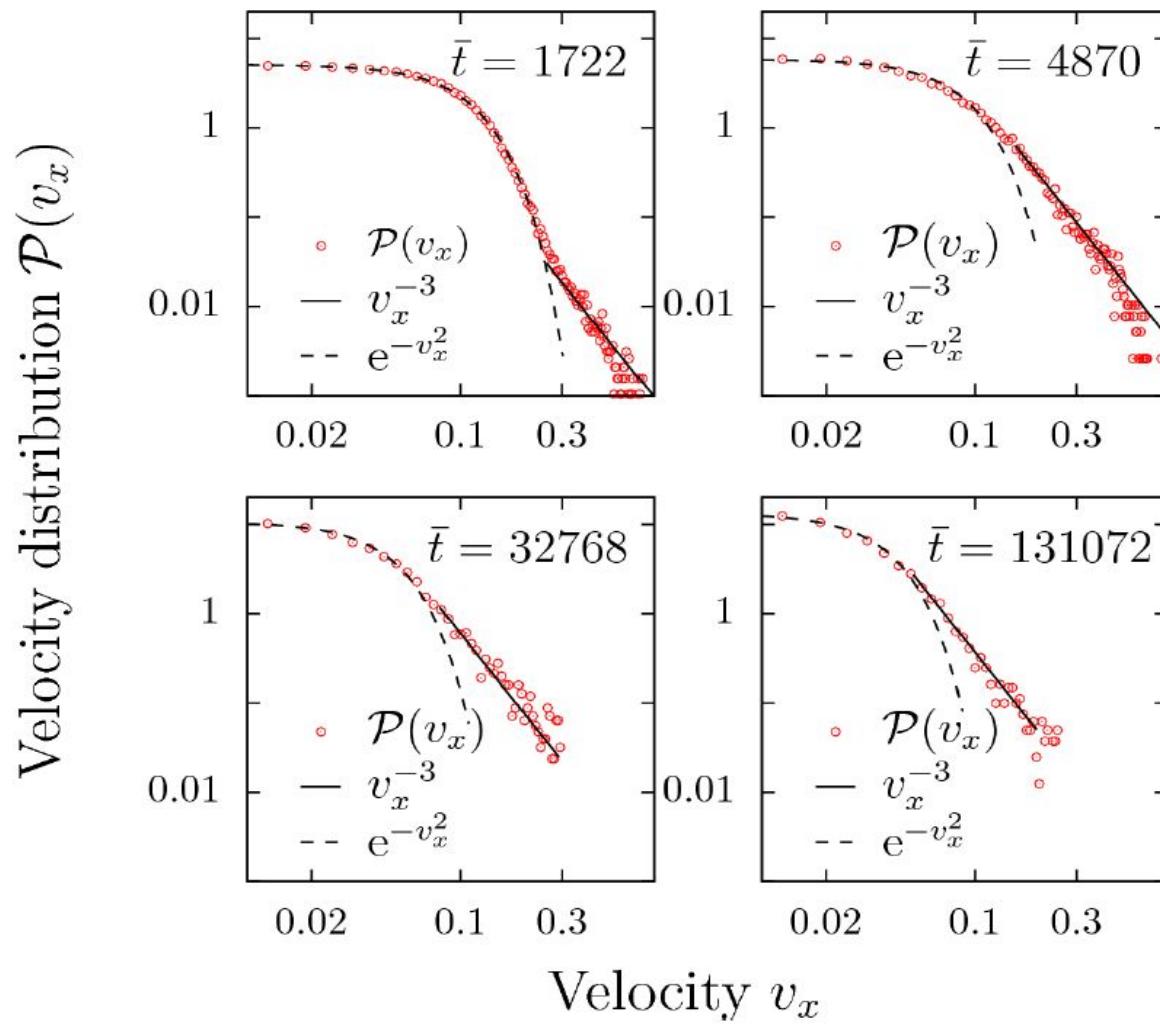
Velocity distributions (Field)



J. S., B. Nowak, D. Sexty, T. Gasenzer (unpublished)



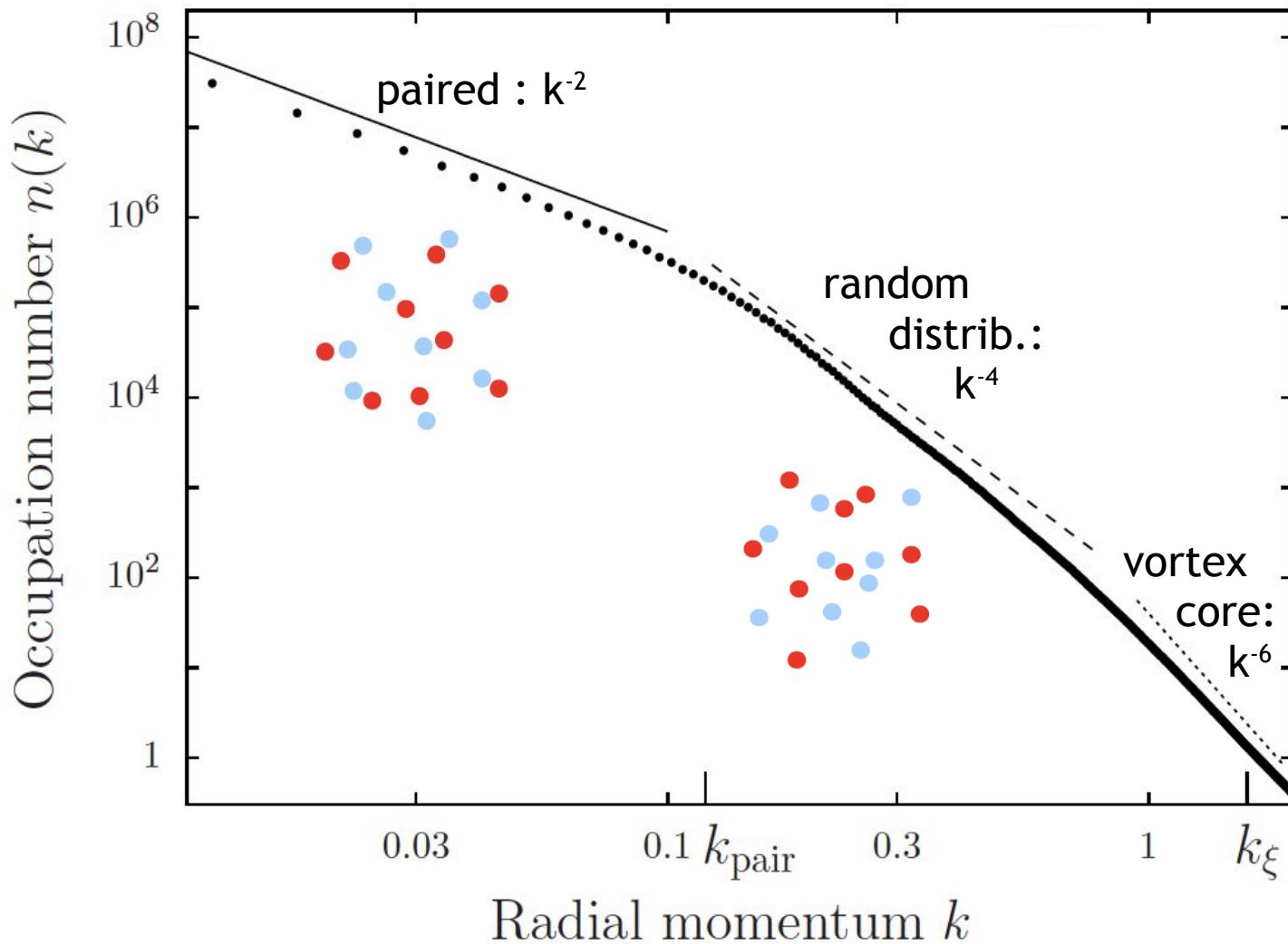
Velocity distributions (Vortices)



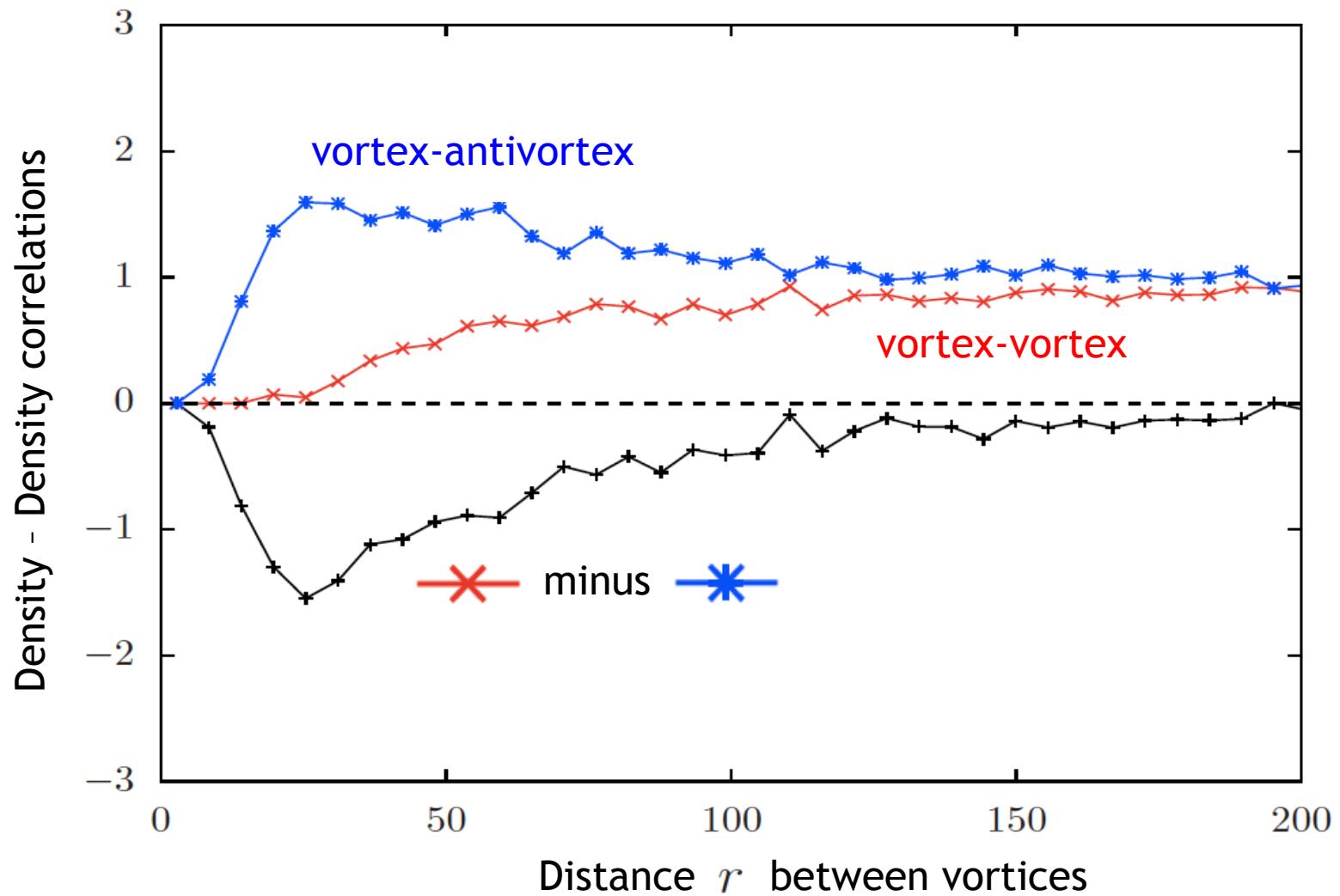
J. S., B. Nowak, D. Sexty, T. Gasenzer (unpublished)



Point vortex model in 2+1 D



Vortex position correlations



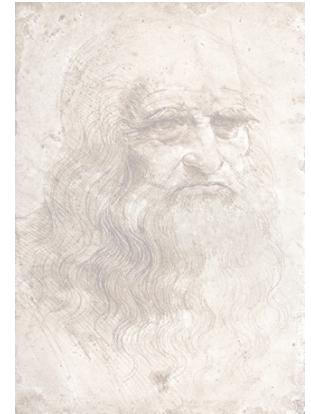
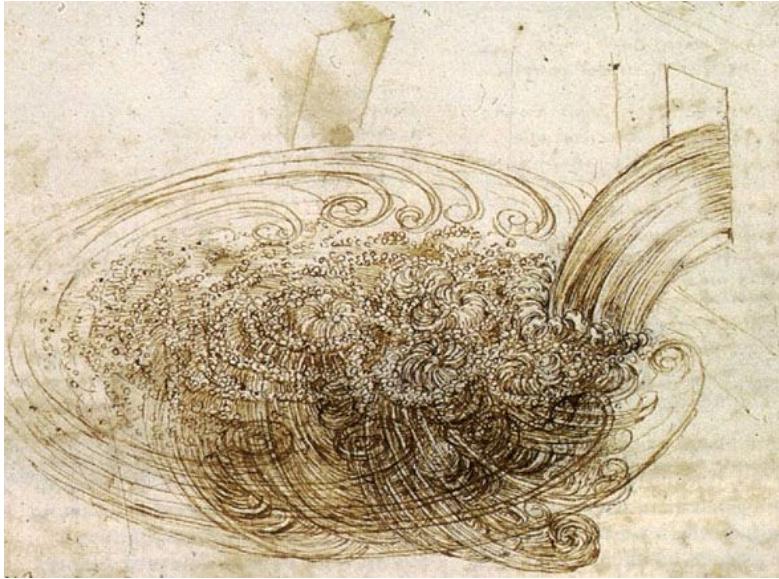
B. Nowak, J. Schole, D. Sexty, TG, arXiv:1111.6127

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Thomas Gasenzer



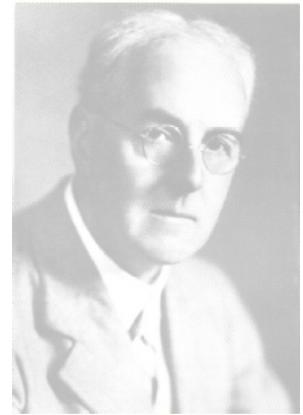
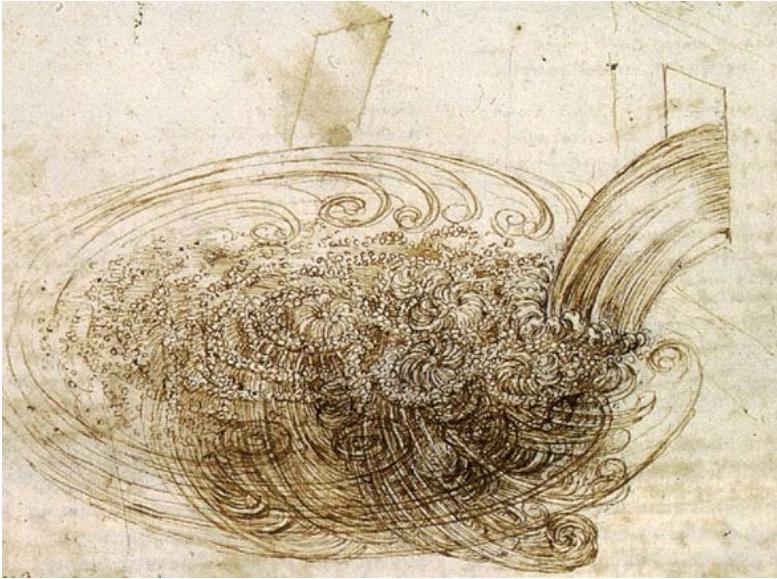
Classical Turbulence



Leonardo da Vinci
(1452-1519)



Classical Turbulence



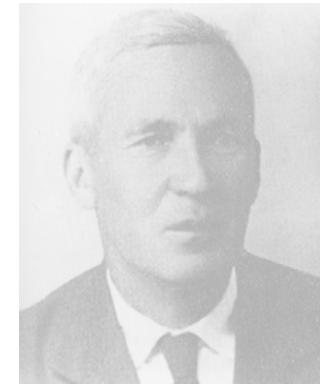
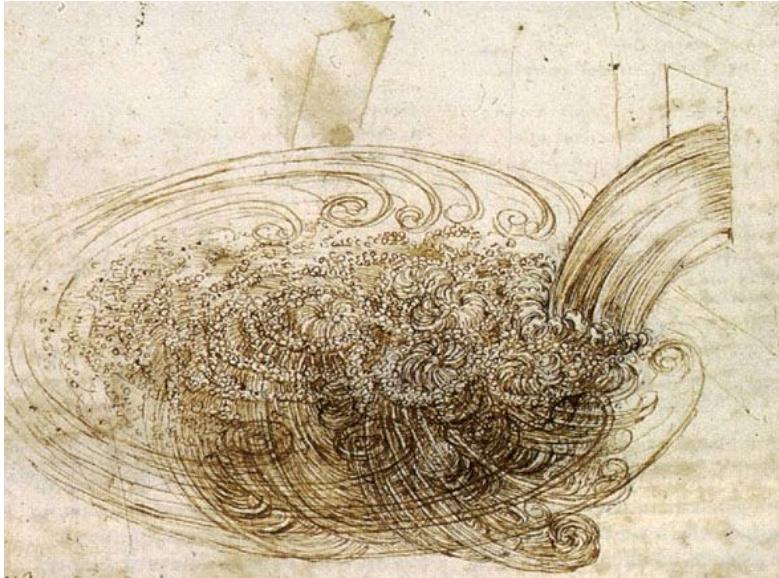
Richardson cascade
large scales (source)
→ small scales (sink)

*“Big whirls have little whirls that feed on their velocity,
and little whirls have lesser whirls and so on to viscosity.”*

(Richardson, 1920)



Classical Turbulence



Andrey N. Kolmogorov
(1903-1987)

Richardson cascade
large scales (source)
→ small scales (sink)

*“Big whirls have little whirls that feed on their velocity,
and little whirls have lesser whirls and so on to viscosity.”*

(Richardson, 1920)

Kolmogorov (1941)

$$E(k) \sim k^{-5/3} \quad (\text{for incompressible fluids})$$

