

# **Integer quantum Hall edge states far from equilibrium**

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**Work with Dmitry Kovrizhin**

**Phys. Rev. B 84, 085105 (2011) and Phys. Rev. Lett. 109, 106403 (2012)**

**Related papers: D. Kovrizhin + JTC: PRB 81 (2010), PRB 80 (2009)**

**JTC + Y. Gefen and M. Veillette, PRB 76 (2007)**

# Outline

## Motivation

Experiments on evolution of non-equilibrium  
electron distribution in QHE edge states

## Theoretical Approaches

Physical picture: relaxation from mode dispersion

Quantum quench as approximate description

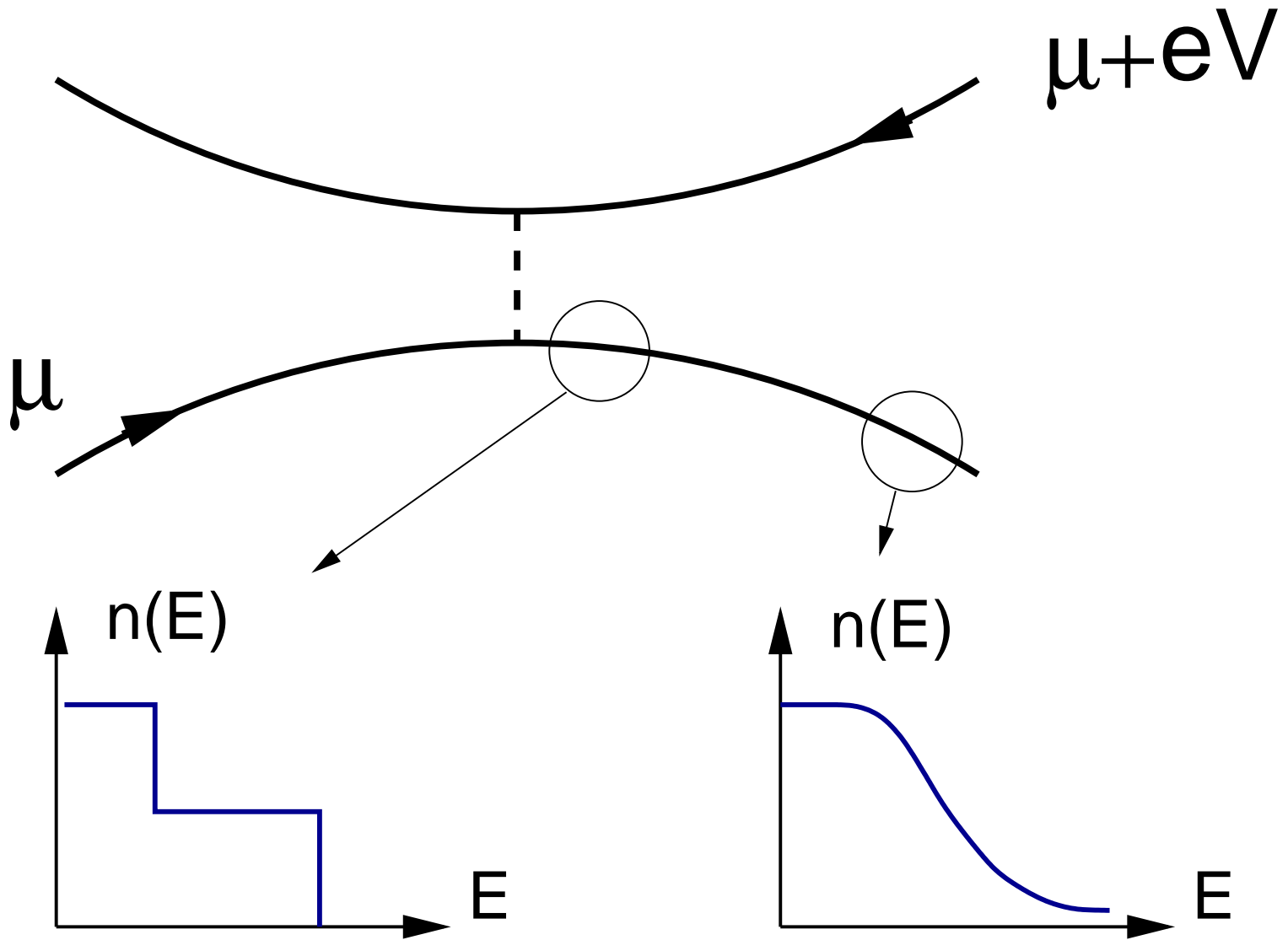
Exact treatment via bosonization + refermionisation

## Results

Relaxation in an integrable system

Non-thermal steady state

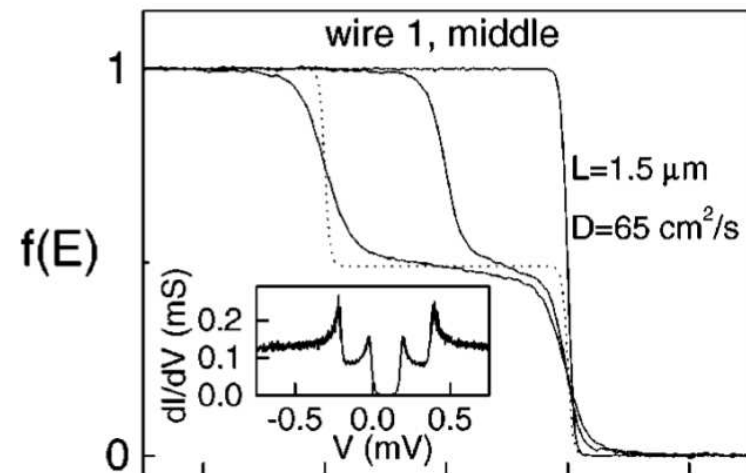
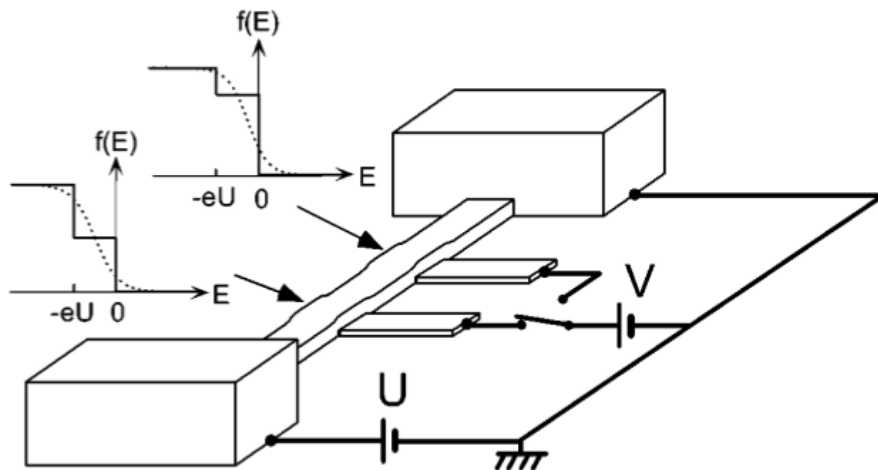
# Schematic view of experiment



le Sueur, Altimiras, Gennser, Cavanna, Mailly & Pierre, PRL (2010)

# Analogue of earlier experiment in wires

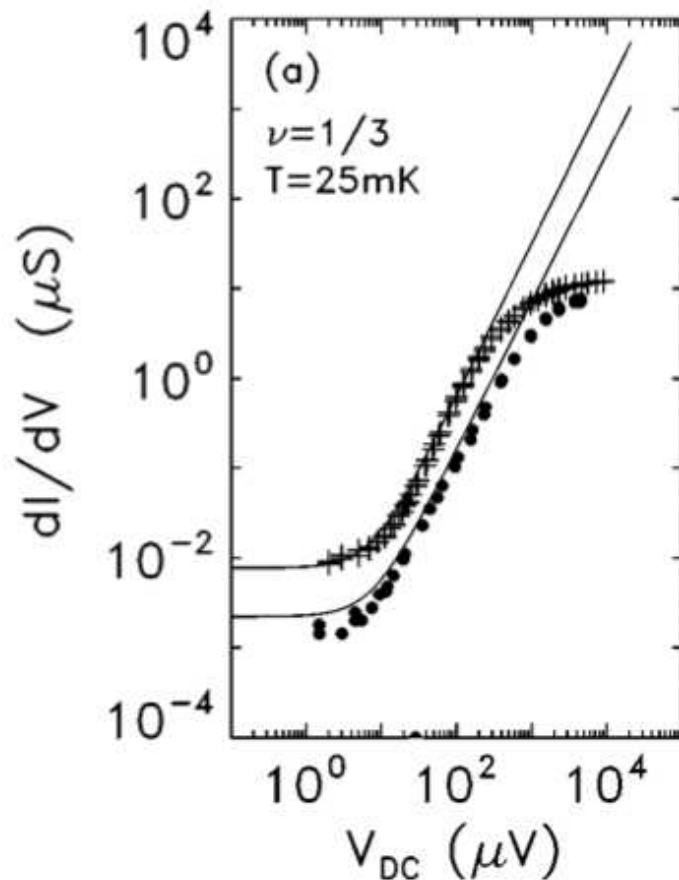
## Electron distribution in biased diffusive wires



Poithier, Guzron, Birge, Esteve, & Devoret, PRL (1997)

# Contrast with tunnelling probe of fractional quantum Hall edge states

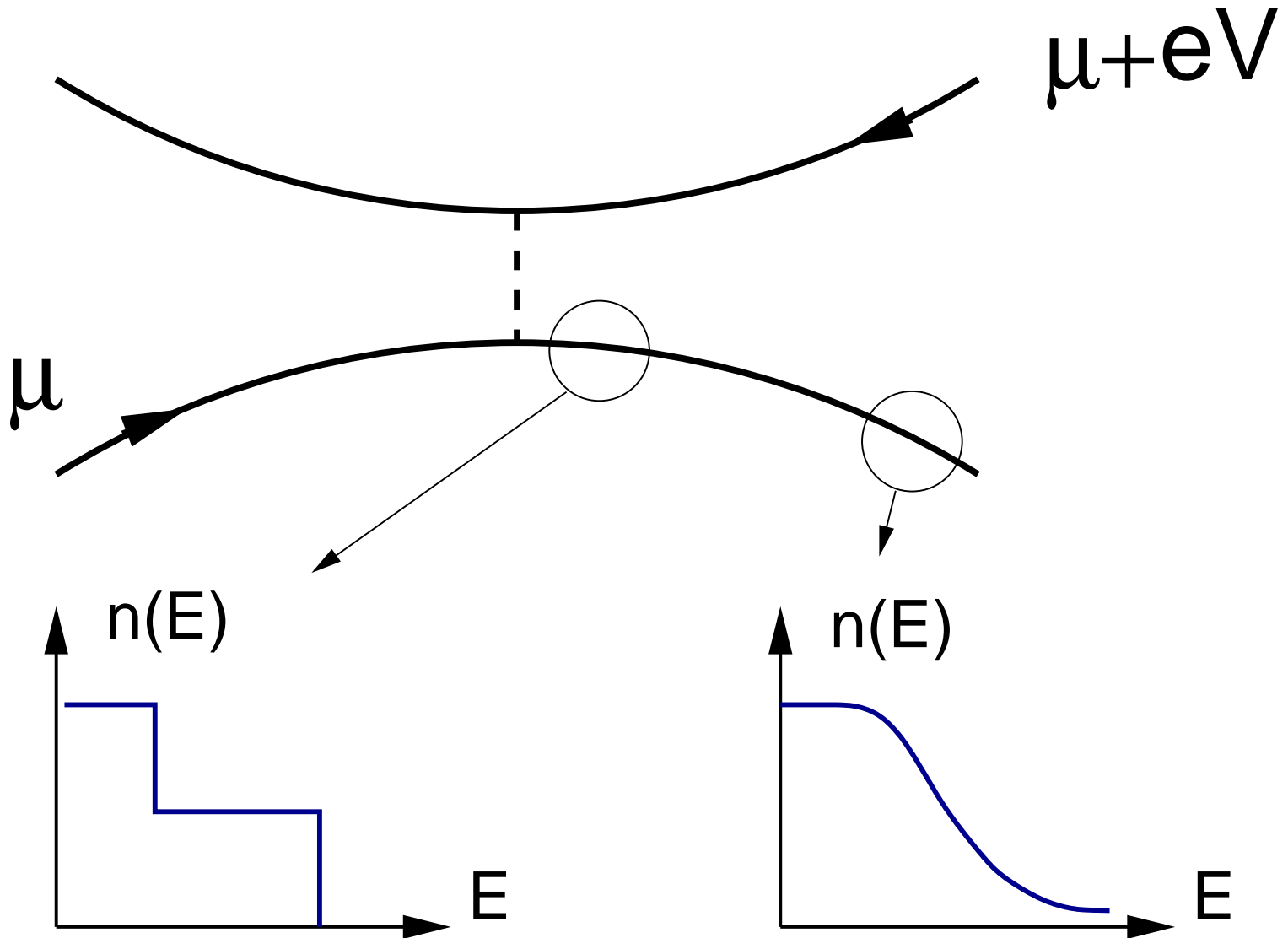
Voltage-dependent differential conductance from electron correlations



Theory: Kane & Fisher (1992)

Expt: Chang, Pfeiffer & West (1996)

# Analogy with quantum quench

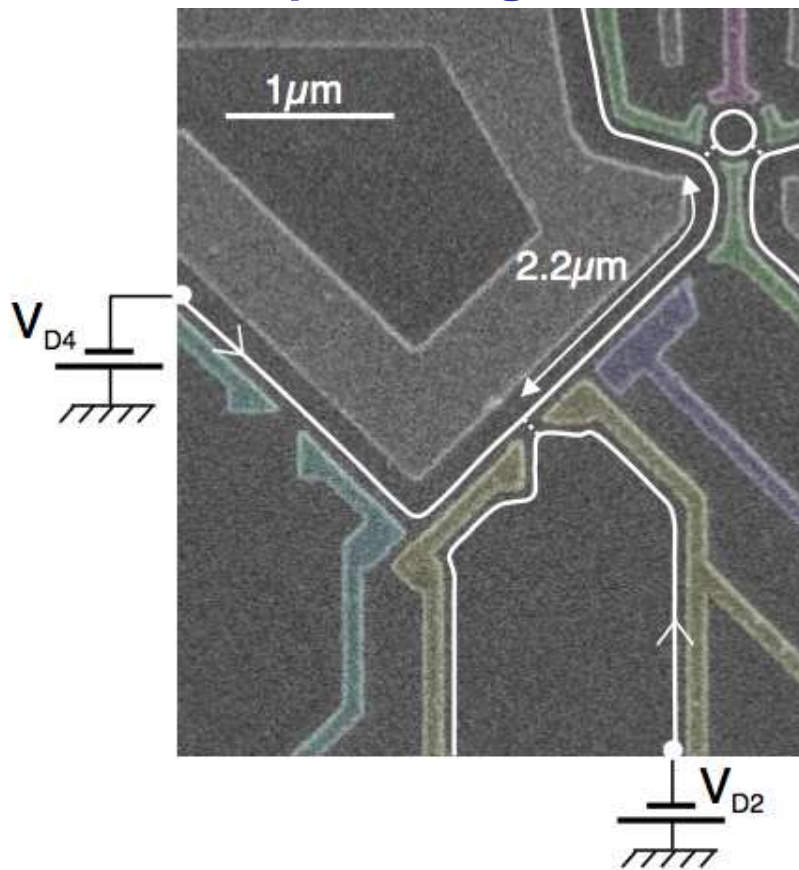


**Distance from contact  $\propto$  time after 'quench'**

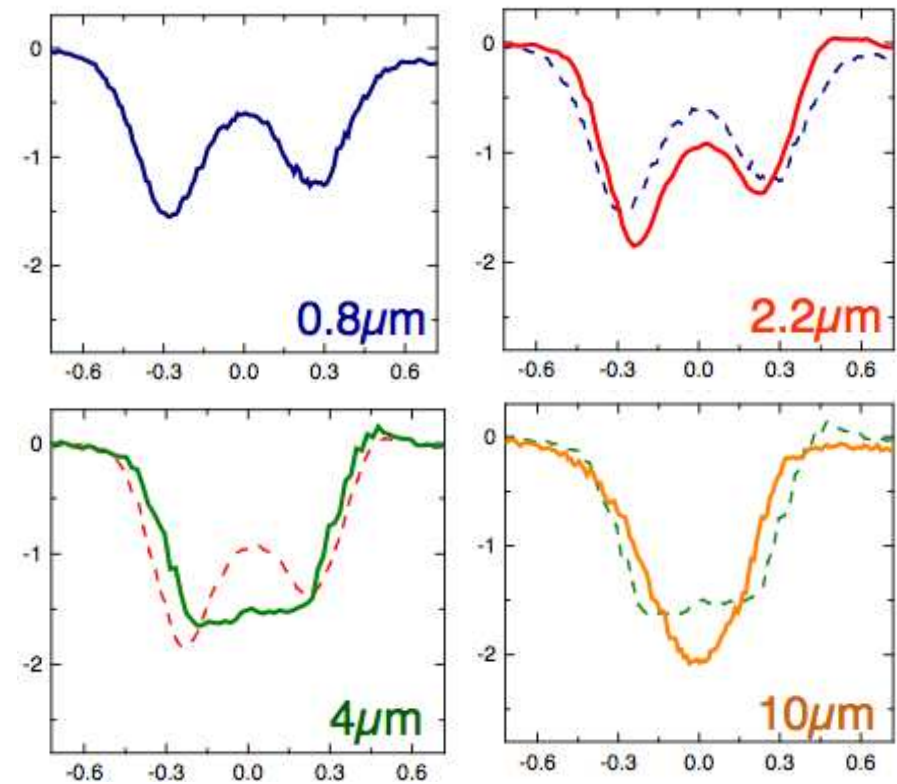
# Experiment – Actual

le Sueur, Altimiras, Gennser, Cavanna, Mailly & Pierre, PRL (2010)

## Sample Design



## Evolution of Distribution



$\partial n(E)/\partial E$  vs.  $E$

# Theoretical description of edge states

As electrons:

$$H = -i\hbar v \int dx \psi^\dagger(x) \partial_x \psi(x) + \int dx \int dx' U(x-x') \rho(x) \rho(x')$$

$$\rho(x) = \psi^\dagger(x) \psi(x)$$

As collective modes:

$$H = \sum_q \hbar \omega(q) b_q^\dagger b_q$$

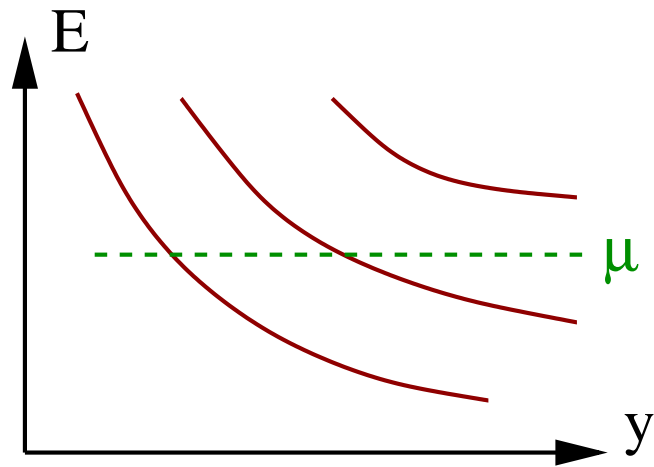
$$\omega(q) = [v + u(q)] q \quad u(q) = (2\pi\hbar)^{-1} \int dx e^{iqx} U(x)$$

— related via bosonization



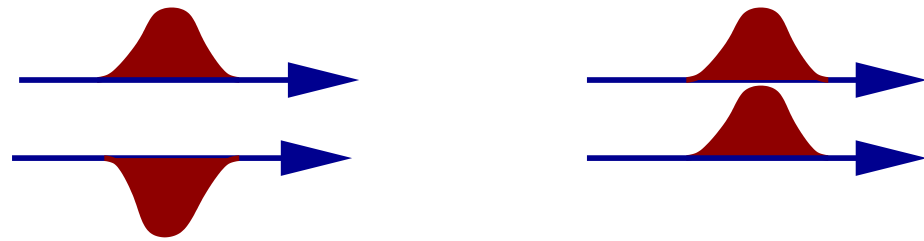
# Fractionalisation at $\nu = 2$

## Two filled Landau levels



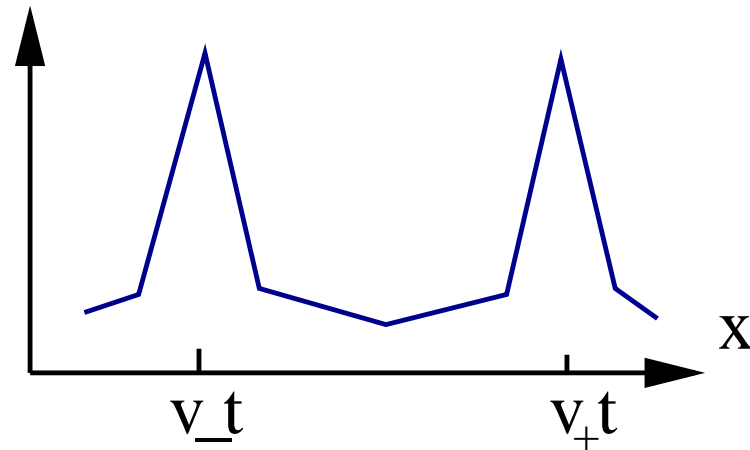
Modes mixed by interaction  $g$

$$v_{\pm} = v \pm g$$



## Injected electron fractionalises

$$|\langle \psi_1^\dagger(x, t) \psi_1(0, 0) \rangle|^2 \sim$$



# Description of edge state tunnelling experiment

## Difficulties

Interactions treated most simply via bosonization

Tunneling at QPC simplest in fermionic language

## Previous work on theory for relaxation expt

Boltzmann Eqn: Lunde *et al* (2010)

QPC as source of plasmon noise: Degiovanni *et al* (2010)

Approx theory of tunnelling: Levkivskyi & Sukhorukov (2012)

## Approaches here

Physical picture

Quantum quench as toy problem

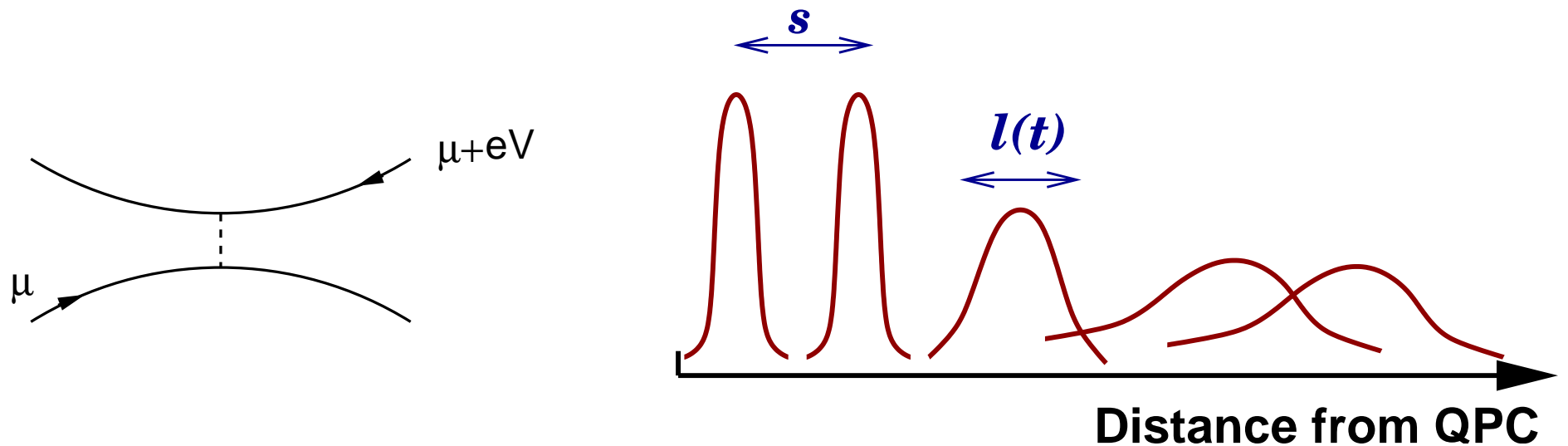
Exact treatment via bosonization + refermionisation

# Physical picture of relaxation

Collective mode Hamiltonian  $\mathcal{H} = \sum_{nq} \hbar\omega_n(q) b_{nq}^\dagger b_{nq}$

Edge magnetoplasmon dispersion  $\rightarrow$  electron equilibration?

Initial quasi-particle separation  $s \propto \hbar v / eV$



Relaxation when wavepacket spread  $l(t) \gtrsim s$

# Relaxation from two mode velocities

## Two edge modes with short-range interactions

**Two linearly dispersing modes**  $\omega_1(q) = v_+q$  &  $\omega_2(q) = v_-q$

**Initial quasi-particle separation**  $s = \hbar v / eV$

**Relaxation when wavepacket spread**  $l(t) \gtrsim s$

**Spread**  $l(t) = [v_+ - v_-]t$

**Relaxation time:**  $t_{\text{eq}} \sim \frac{\hbar}{eV} \cdot \frac{v}{v_+ - v_-}$

**Relaxation distance:**  $v t_{\text{eq}}$

# Theoretical Idealisation: Quantum quench

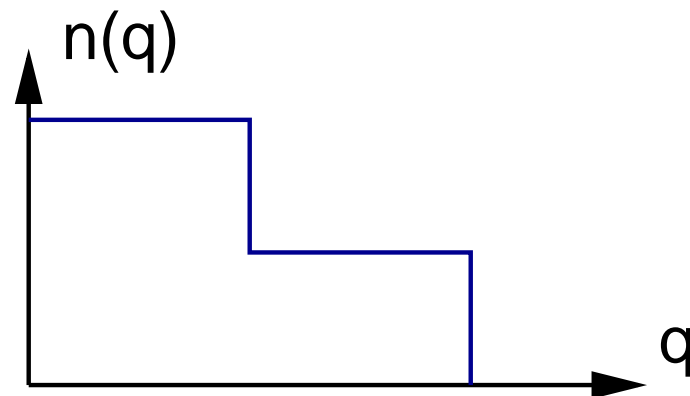
Evade treatment of point contact

– study time evolution in translationally-invariant edge

Initial state

$$|\Psi_0\rangle$$

with



Time evolution

$$|\Psi(t)\rangle = e^{i\mathcal{H}t} |\Psi_0\rangle$$

Properties of  $|\Psi(t)\rangle$  ?

Energies of collective modes conserved  
— consequences for equilibration?

# What is the equilibrium state?

## Characterise via one-electron correlations

**Calculate**  $G(x, t) = \langle \psi^\dagger(x, t) \psi(0, t) \rangle$

**in thermal state**  $G(x, t) = [-2i\beta\hbar v \sinh(\pi[x + i0]/\beta\hbar v)]^{-1}$

## Approach to calculations:

### Alternate between fermionic and bosonic descriptions

$$\Psi(x) \sim e^{-i\varphi(x)}$$

$$\varphi(x) \propto \sum q^{-1/2} [b_q^\dagger e^{-iqx} + \text{h.c.}]$$

$$b_q^\dagger \propto \sum c_{k+q}^\dagger c_k$$

**Initial state**

**simple for fermions**

**Time evolution**

**simple for bosons**

**Observables**

**fermion fields**

# Comparison with thermal state

## Short-distance correlations

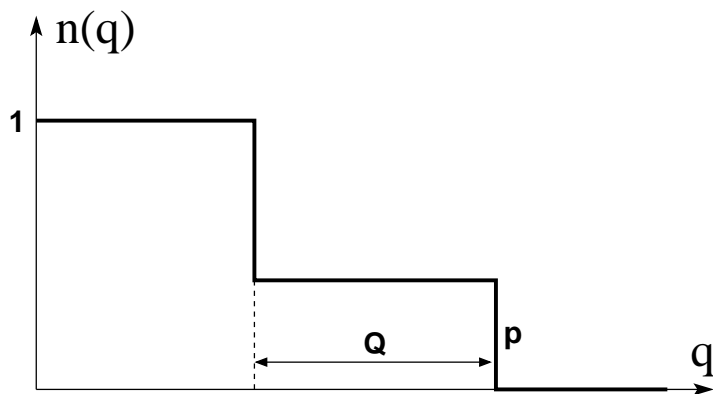
As in thermal state at same energy density

## Long-distance correlations

$G(x, t) \sim \exp(-\alpha|x|)$  with  $\alpha$  not fixed by energy density

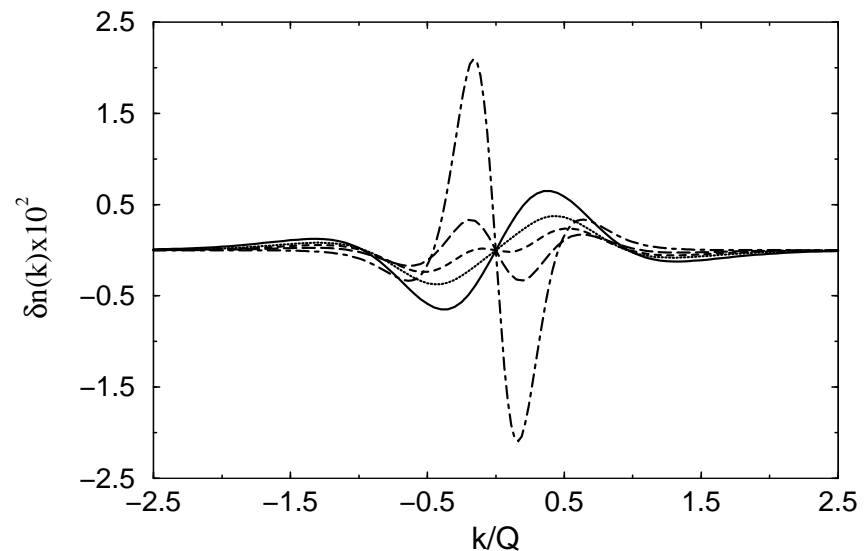
## Example

### Initial momentum distribution

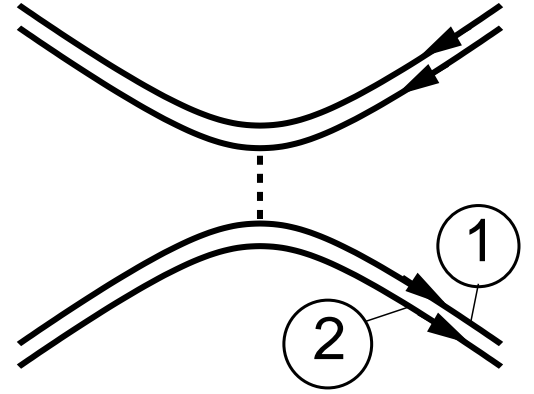


### Difference from thermal in steady state

$p = 0.1, 0.2, 0.25, 0.3, 0.5$



# Solvable model



Each edge

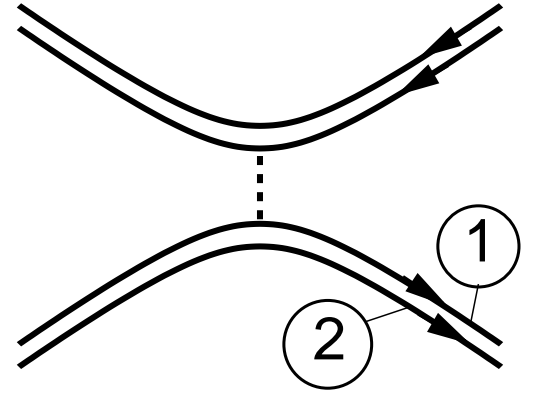
$$\mathcal{H} = \mathcal{H}_{\text{kin}} + \mathcal{H}_{\text{int}}$$

$$\mathcal{H}_{\text{kin}} = -i\hbar v \int \left( \Psi_1^\dagger(x) \partial_x \Psi_1(x) + \Psi_2^\dagger(x) \partial_x \Psi_2(x) \right) dx$$

$$\mathcal{H}_{\text{int}} = 2\pi\hbar g \int \rho_1(x) \rho_2(x) dx, \quad \rho_n(x) = \Psi_n^\dagger(x) \Psi_n(x)$$



# Solvable model



## Each edge

$$\mathcal{H} = \mathcal{H}_{\text{kin}} + \mathcal{H}_{\text{int}}$$

$$\mathcal{H}_{\text{kin}} = -i\hbar v \int \left( \Psi_1^\dagger(x) \partial_x \Psi_1(x) + \Psi_2^\dagger(x) \partial_x \Psi_2(x) \right) dx$$

$$\mathcal{H}_{\text{int}} = 2\pi\hbar g \int \rho_1(x) \rho_2(x) dx, \quad \rho_n(x) = \Psi_n^\dagger(x) \Psi_n(x)$$

## Bosonize and diagonalise

$$\Psi_{1,2}(x) \sim e^{-i\varphi_{1,2}(x)}$$

$$\mathcal{H} = \frac{\hbar v_+}{2} \int [\partial_x \varphi_+(x)]^2 \frac{dx}{2\pi} + \frac{\hbar v_-}{2} \int [\partial_x \varphi_-(x)]^2 \frac{dx}{2\pi}$$

## Modes

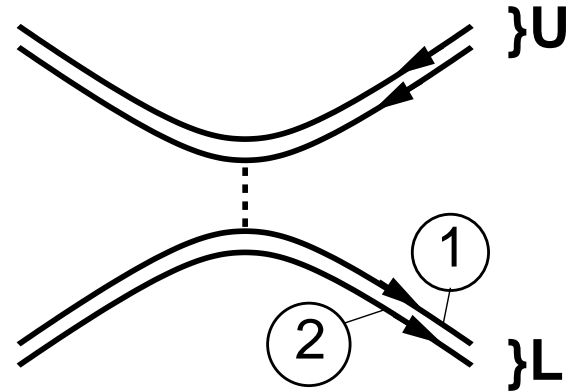
$$v_\pm = v \pm g \quad \varphi_\pm(x) = \frac{1}{\sqrt{2}} [\varphi_1(x) \pm \varphi_2(x)]$$

# Refermionize

Combine bosons from opposite edges

$$\varphi_{S\pm} = \frac{1}{\sqrt{2}}[\varphi_{U\pm}(x) + \varphi_{L\pm}(x)]$$

$$\varphi_{A\pm} = \pm \frac{1}{\sqrt{2}}[\varphi_{U\pm}(x) - \varphi_{L\pm}(x)]$$



**Tunneling**  $\mathcal{H}_{\text{tun}} = t_{\text{QPC}}[\Psi_{U1}^\dagger(0)\Psi_{L1}(0) + \Psi_{L1}^\dagger(0)\Psi_{U1}(0)]$

$$\Psi_1^\dagger(0)\Psi_2(0) \sim e^{i[\varphi_{U1}(0)-\varphi_{L1}(0)]} \sim e^{i[\varphi_{A+}(0)-\varphi_{A-}(0)]} \sim \Psi_{A+}^\dagger(0)\Psi_{A-}(0)$$

**Edges**

$$\mathcal{H}_{\text{kin}} + \mathcal{H}_{\text{int}} = -i\hbar \int [v_+ \Psi_{A+}^\dagger(x) \partial_x \Psi_{A+}(x) + v_- \Psi_{A-}^\dagger(x) \partial_x \Psi_{A-}(x)] dx + [A \leftrightarrow S]$$

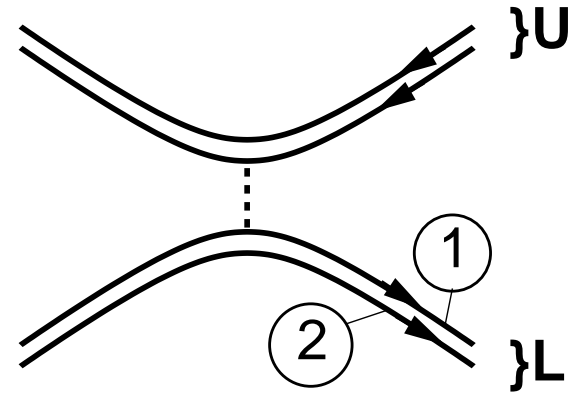
# Observables

Electron energy distribution at  $x$  from

$$\langle \Psi_{L1}^\dagger(x, t) \Psi_{L1}(x, 0) \rangle$$

or

$$\langle \Psi_{L2}^\dagger(x, t) \Psi_{L2}(x, 0) \rangle$$



Transforms to  $\langle e^{i\pi(n_- - n_+)} \rangle$  or  $\langle e^{i\pi(n_- + n_+)} \rangle$

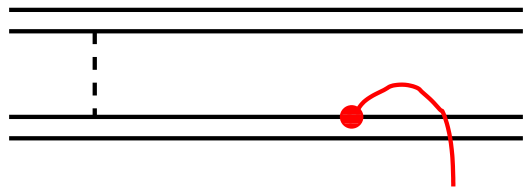
with  $n_{\pm} = \int_x^{x+v_{\pm}t} \Psi_{A_{\pm}}^\dagger(y) \Psi_{A_{\pm}}(y) : dy$

Energy exchange between channels — analytic result

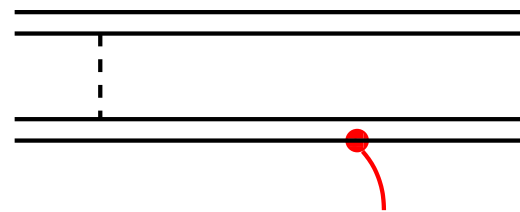
Electron distribution — evaluate free fermion averages numerically

# Results: energy exchange between channels

Measurement in channel  
coupled at QPC

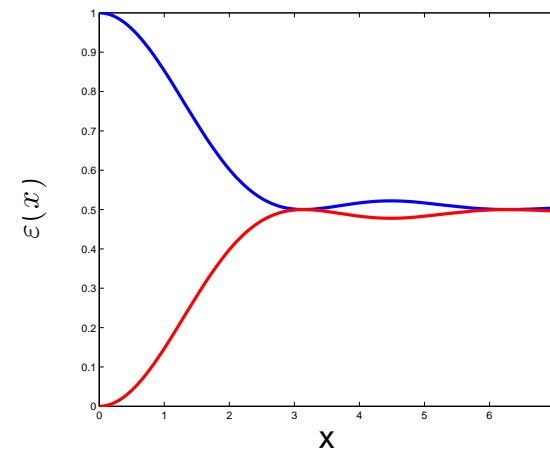


Measurement in channel  
coupled by interactions



Energy density

$$\varepsilon(x) = \varepsilon_T + \frac{\varepsilon_V}{2} \left[ 1 \pm \frac{\ell_V^2 \sin^2(x/2\ell_V)}{\ell_T^2 \sinh^2(x/2\ell_T)} \right]$$



Thermal contribution  $\varepsilon_T$

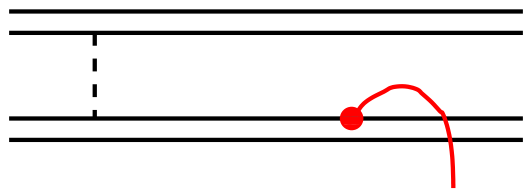
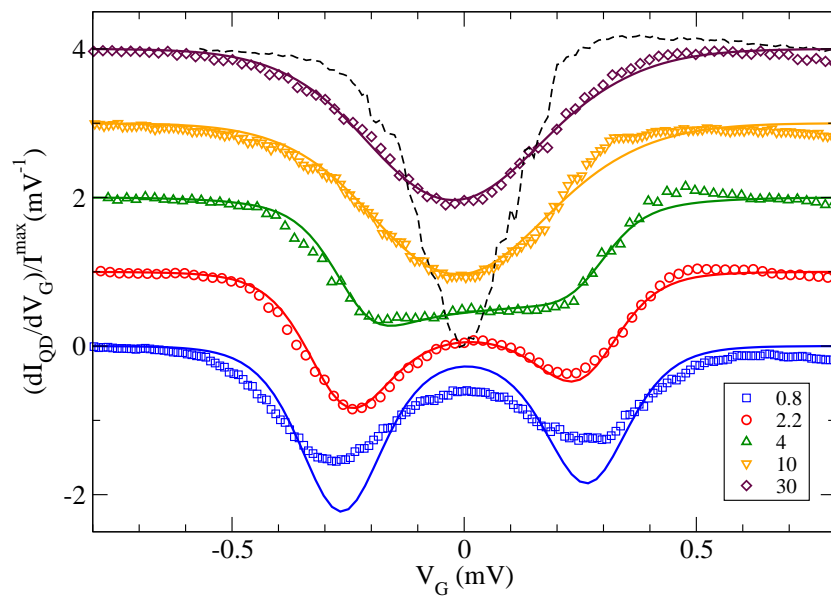
Contribution from bias voltage  $\varepsilon_V$

Lengthscales  $\ell_V = \frac{\hbar}{eV} \frac{v_+ v_-}{v_+ - v_-}$  and  $\ell_T = \frac{\hbar}{2\pi k_B T} \frac{v_+ v_-}{v_+ - v_-}$

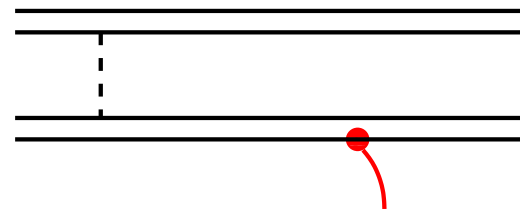
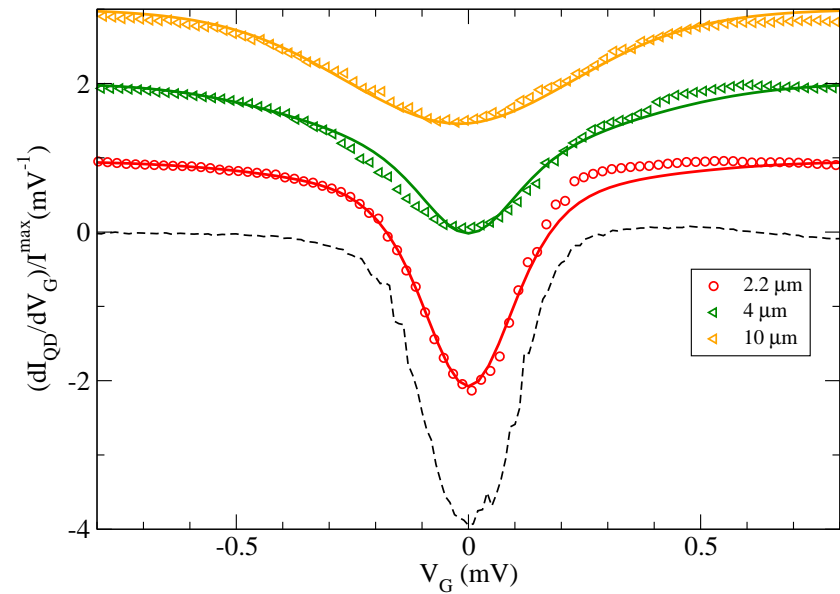
see also Degiovanni *et al.* 2010

# Comparison with experiment

Measurement in channel  
coupled at QPC



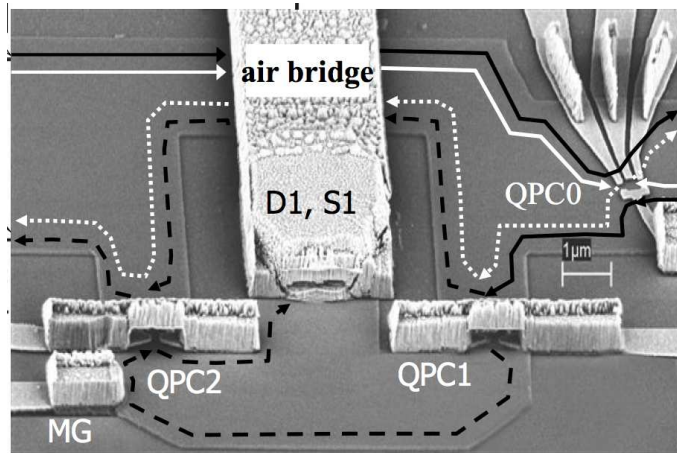
Measurement in channel  
coupled by interactions



Fitting parameter: interaction strength  $(v^2 + g^2)/2g = 6.5 \times 10^4 \text{ms}^{-1}$

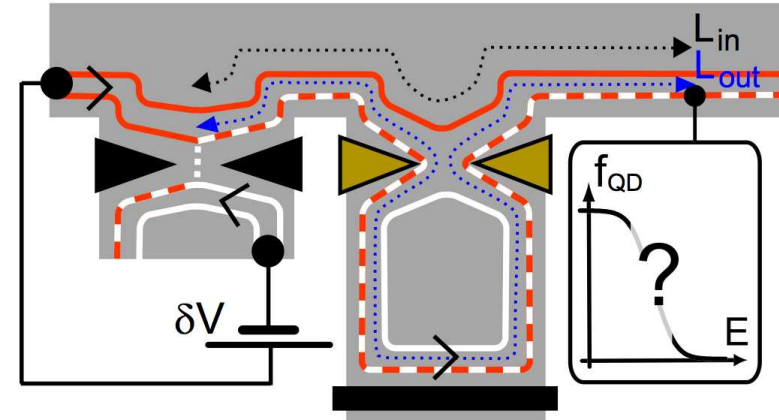
# Related experiments

## Interferometers out of equilibrium



Heiblum group (2005)

## Controlling relaxation



Altimiras et al (2011)

# Summary

## Experiment probes relaxation in an integrable system

- Interactions bring system into steady state

## Full problem with QPC solvable at $\nu = 2$

- Asymptotic state is non-thermal
- Calculated evolution of tunneling density of state with distance matches experiment well