

Universal “rephasing” dynamics

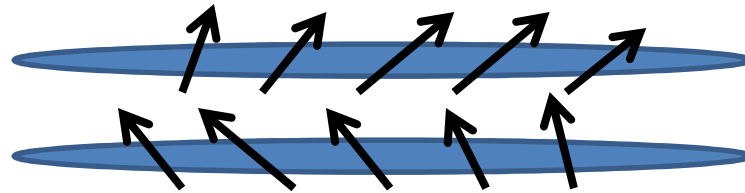
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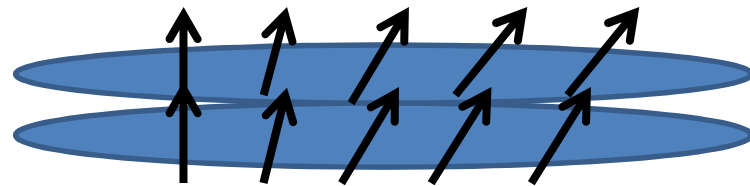
Anatoli Polkovnikov

KITP, October 25, 2012

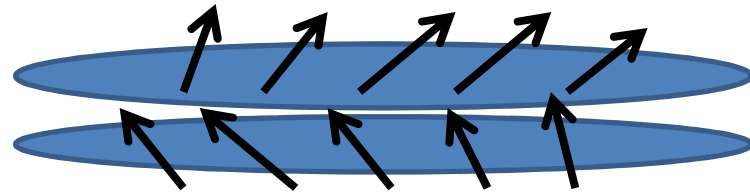
What happens when we couple two independent (quasi)condensates?



Tunneling \rightarrow rephasing



Heating \rightarrow dephasing



Who wins? Is the answer universal?

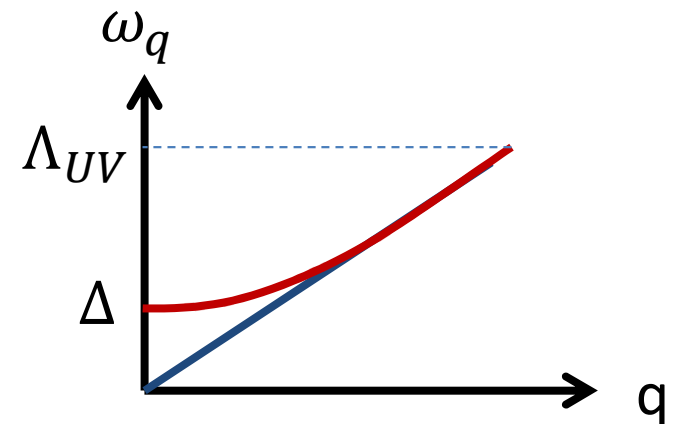
Universality – handwaving argument

$t < 0$

No tunneling: gapless spectrum
(Goldstone mode of the relative phase)

$t > 0$

Finite tunneling : excitation gap Δ



How much heating?

$$\frac{\delta E}{N} = \frac{\text{Total Energy}}{\text{Number of modes}} \approx \Delta \frac{\Delta}{\Lambda_{UV}} \quad \Rightarrow \quad \begin{array}{l} \text{Small coupling:} \\ \Delta \ll \Lambda_{UV} \end{array} \quad \frac{\delta E}{N} \ll \Delta$$

→ universal quantum dynamics protected by the gap

De Grandi, Gritsev, Polkovnikov (PRB, 2010)

Outlook



(a) Mean field



(b) One dimension

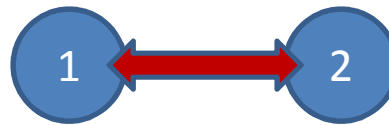


(c) Q&A, Cookies

Mean-field model (Josephson junction)

$$H(t) = \frac{\mu}{N} [(\delta n_1)^2 + (\delta n_2)^2] + j_{\perp}(t)(\psi_1^{\dagger} \psi_2 + H.c.)$$

$$\delta n_1 = \psi_1^{\dagger} \psi_1 - \frac{N}{2}$$

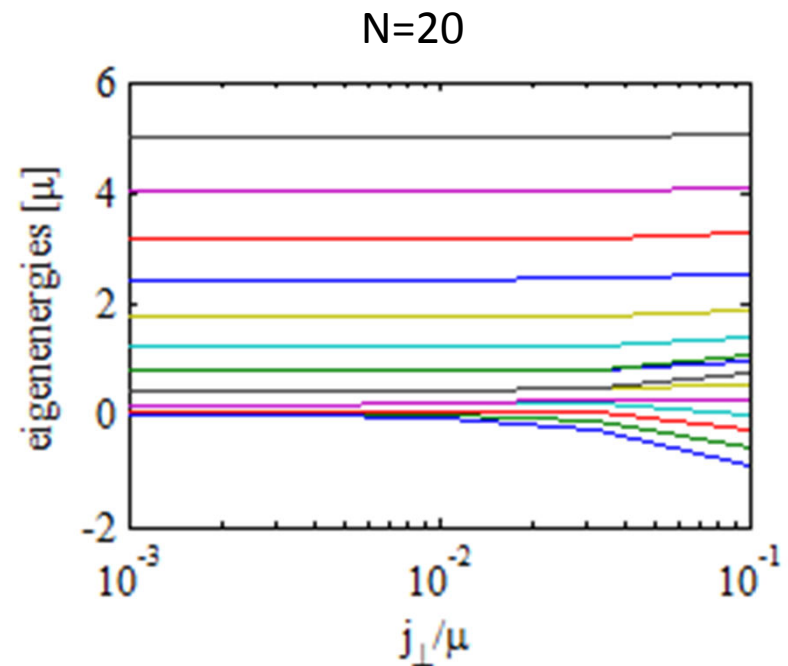


- ✓ Initial Hamiltonian is gapless:

$$\Delta = \frac{2\mu}{N} \rightarrow 0$$

- ✓ Final Hamiltonian is gapped:

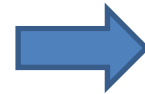
$$\Delta = \omega_p = \sqrt{4j_{\perp}\mu}$$



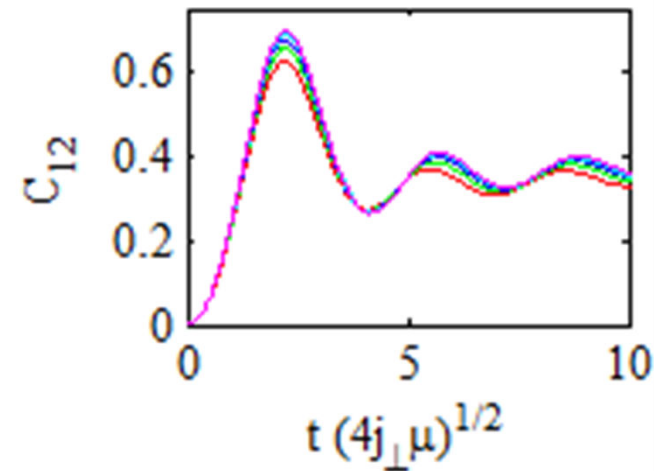
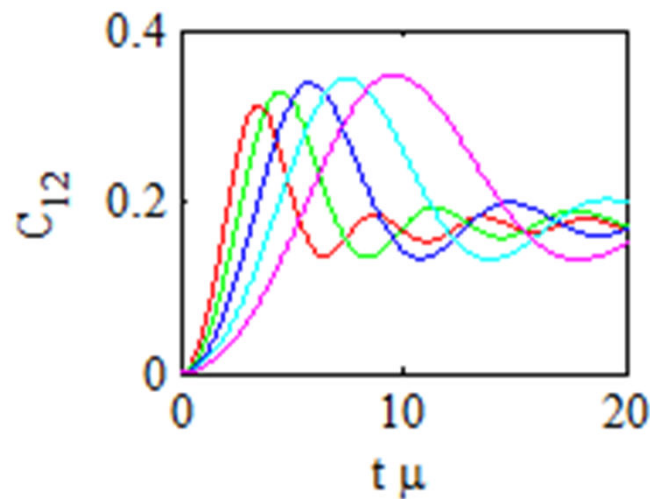
Mean field – observing universality

$$C_{12}(t) = \langle \psi_1^\dagger \psi_2 + H.c. \rangle$$

$$\frac{j_\perp}{\mu} = 0.1, 0.06, 0.037, 0.022, 0.013$$



Rescaling



What is the universal function?

Mean field – universal curves

In the thermodynamic limit $N \rightarrow \infty$ we can map the problem to
classical equation of motion

$$H(t) = \frac{\mu}{N} [\delta n_1^2 + \delta n_2^2] + j_{\perp}(t) (\psi_1^{\dagger} \psi_2 + H.c.)$$

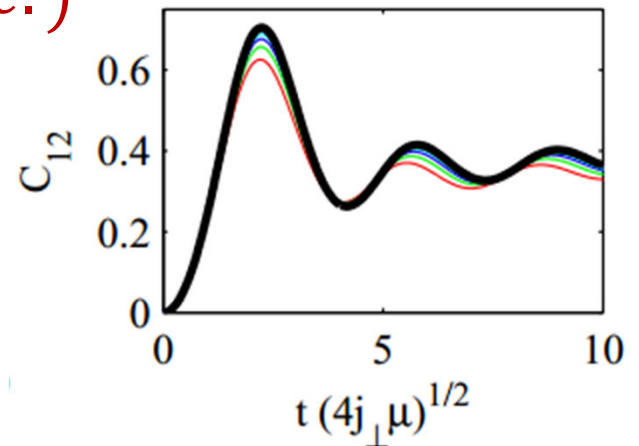
$$\psi_1 \rightarrow \sqrt{N} e^{i\phi_1}$$

$$\delta n_1 \rightarrow \frac{N}{2\mu} \partial_t \phi_1$$

$$\phi = \phi_1 - \phi_2$$

$$\hat{\phi} = \phi_1 + \phi_2$$

$$\frac{1}{2\mu} \partial_t^2 \phi = 2j_{\perp}(t) \cos(\phi)$$



Non-linear simple pendulum with
period $\Delta = \sqrt{4j_{\perp}\mu}$: analytic solution
via elliptic functions (**black curve**)

One dimension

The initial state does not have a well defined phase

$$H(t) = H_{LL}[\phi_1] + H_{LL}[\phi_2] + g \Theta(t) \int dx \cos(\phi_1 - \phi_2)$$

Relative and COM coordinates:

$$\phi = \phi_1 - \phi_2$$

$$\hat{\phi} = \phi_1 + \phi_2$$

$$H(t) = H_{LL}[\hat{\phi}] + H_{LL}[\phi] + g \Theta(t) \int dx \cos(\phi)$$

Quench in a Sine-Gordon – many different cases!

Iucci & Cazalilla (2010)

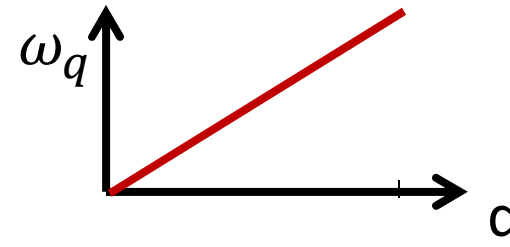
Mitra (2012)

Quench in a Sine-Gordon - review

(1) Turning **off** the coupling:

Long time dynamics is **classical** (“Dephasing”)

Calabrese&Cardy, Bistritzer&Altman (2006)

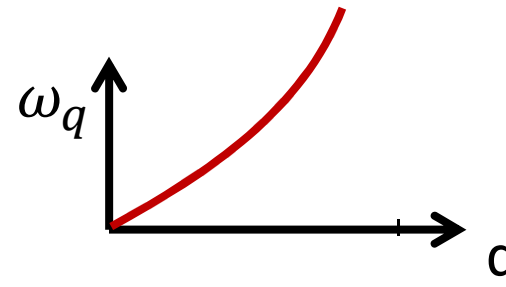


$$\Delta = 0$$

(2) Turning **on** an irrelevant coupling: ($K > 2$)

Long time dynamics is **not-universal**

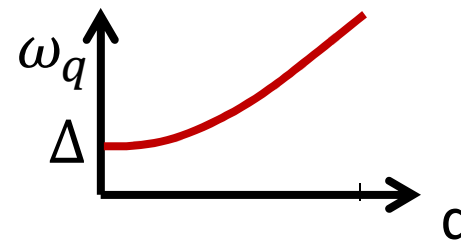
Mitra and Giamarchi (2011)



$$\frac{\delta E}{N} \sim \Lambda_{UV}$$

(3) Turning **on** a relevant coupling: ($K < 2$)

Long time dynamics is **quantum** (“Rephasing”)



$$\frac{\delta E}{N} \ll \Delta$$

Independent modes approximation

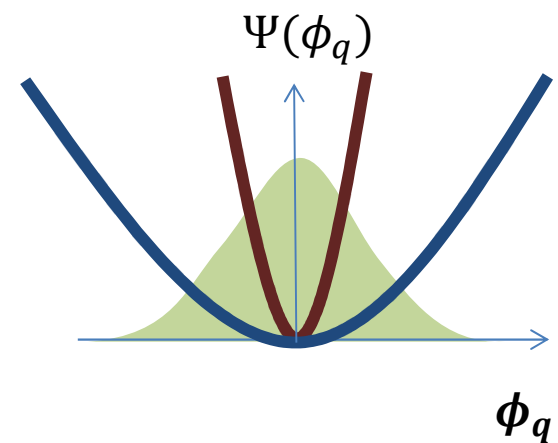
Iucci & Cazalilla (2010)

Fluctuations of the q mode oscillate

$$\langle |\phi_q(t)|^2 \rangle \approx \frac{K}{|q|} \cos^2(\Delta t) \rightarrow \frac{K}{2|q|}$$

Total fluctuations $\langle \phi^2(x, t) \rangle$ diverge

$$\langle \cos(\phi(t)) \rangle = e^{-\frac{1}{2}\langle \phi^2 \rangle} \rightarrow 0$$



→ The system does not rephase!

Why is this wrong?

Non-perturbative series expansion

$$\begin{aligned}
 \langle \cos \phi(x,t) \rangle &= \int \mathcal{D}\phi \cos \phi e^{iS_0 + g \int_0^t dt' dx' \cos \phi'} \\
 &= \sum_{N=0}^{\infty} \frac{1}{N!} (ig)^N \langle \cos \phi \left(\int_0^t dt' dx' \cos \phi' \right)^N \rangle_0 \\
 &\sim \sum_{N=0}^{\infty} C_N g^N (\mu t)^{2N - (N+1)k/2} \\
 &\sim \left(\frac{\Delta}{\mu} \right)^k \underbrace{\sum_{N=0}^{\infty} C_N (\Delta t)^{2N - (N+1)k/2}}_{R(\Delta t)} \quad \Delta = \mu \left(\frac{g}{\mu} \right)^{\frac{1}{2-k}}
 \end{aligned}$$

Expectation value - scaling result

$$\langle \cos(\phi(t)) \rangle = \left(\frac{\Delta}{\Lambda_{UV}} \right)^K R(\Delta t)$$

Same scaling exponent
as in the ground state

Lukyanov & Zamolodchikov (1997)

Universal dynamics –
connecting short times
to the steady state

Theory: Janssen et al (1989)

Numerical: Zheng (1998)

Review: Albano et al (RPP 2011)

Non universal contributions

From the UV cutoff \leftrightarrow integral over short times

$$\langle \cos(\phi(t)) \rangle = \left(\frac{g}{\Lambda_{UV}} \right)^{\frac{K}{2-K}} R(\Delta t) + A \frac{g}{\Lambda_{UV}} + o(g^2)$$

Must be an analytic function of “g”

 The universal part is dominant for $\frac{K}{2-K} < 1 \Leftrightarrow K < 1$

At K=1 Exact solution (re-fermionization)

$$\langle \cos(\phi(t)) \rangle = \frac{g}{\Lambda_{UV}} [\text{CosIntegral}(2gt) + A]$$

both terms scale like g

Iucci & Cazalilla (2010)

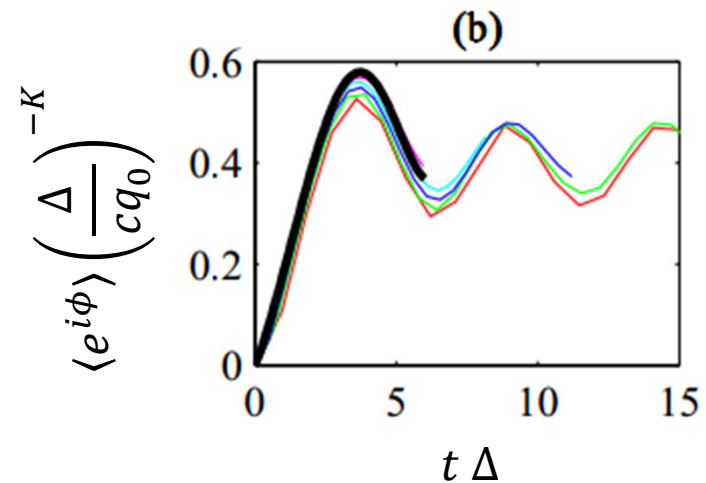
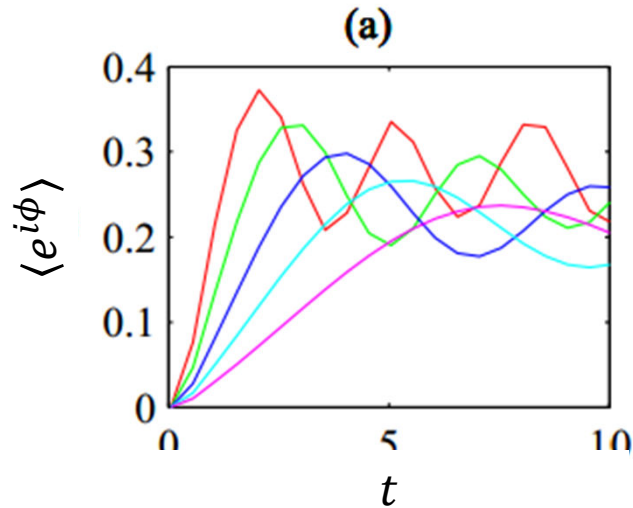
Lattice model (TEBD)

Hard core bosons ladder:



J, V — blue line
 $J_{\perp}(t)$ — red line

$K = 0.4$



Black curve: short time perturbative result

More about universality

How sudden is “sudden”?

$$\tau_{\text{ramp}} \ll \frac{1}{\Delta}$$

How cold is “cold”?

$$T_0 \ll \Delta$$

Spatially integrated coherence?

$$Z = \int_{-L/2}^{L/2} dx e^{i\phi(x)}$$

$$P(|Z|) = f(K, \Delta L, \Delta t)$$

Equilibrium case:

V. Gritsev, E. Altman, E. Demler, A. Polkovnikov, Nat.Phys. (2006)

A. Imambekov, V. Gritsev, E. Demler, PRA (2008)

Summary

1. The “rephasing” process is dominated by the low-frequency modes and is thus **universal**.

$$\langle \cos(\phi(t)) \rangle = \left(\frac{\Delta}{c q_0} \right)^K R(\Delta t)$$

2. The **scaling** ansatz relates the perturbative short-time dynamics to the non-perturbative steady states
3. The full distribution of the **interference** phase and contrast show respectively mean-field and beyond mean-field dynamics

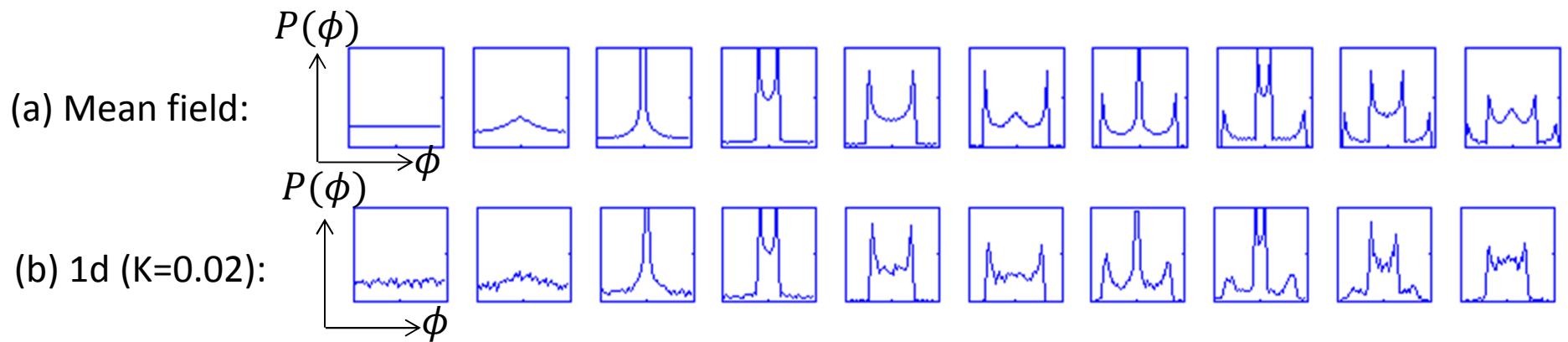
Non-equal time correlations?

Analytic expression for the one dimensional case?

Extension to other systems: Kondo physics from short-time QMC?

Q.: Why integrated coherence?

Problem: In atom-on-a-chip $K \ll 1 \rightarrow$ the dynamics is **very** mean field



Solution: use $|Z|$ to observe 1d physics

$$Z = \int_{-L/2}^{L/2} dx e^{i\phi(x)}$$

