

# Relaxation, pre-thermalization and diffusion in a noisy Quantum Ising Chain

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# The model and the out of equilibrium protocol

$$H = H_0 + V$$

$$H_0 = -J \sum_i \sigma_i^x \sigma_{i+1}^x + g \sigma_i^z \quad \Delta = |g - 1|$$

Two gapped phases: quantum paramagnetic and ferromagnetic separated by a quantum critical point at  $g=1$

$$V = \sum_i \delta g(t) \sigma_i^z \quad \langle \delta g(t) \delta g(t') \rangle = \frac{\Gamma}{2} \delta(t - t')$$

We perform a quantum quench from the paramagnetic phase at  $g_0$  to the paramagnetic phase at  $g$  and we compute observables of physical interest during the dynamics 



# Dynamics of Thermalization

J. Marino, A. Silva - [arXiv:1203.2108v2](https://arxiv.org/abs/1203.2108v2)

$$\frac{\Gamma}{\Delta} \ll 1$$

- **Prethermalization**

drifts towards GGE (Kollar *et al.* 2010)

- **Dephasing**  $t \approx \frac{1}{\Gamma}$  : coherences killed

- **Thermalization** and population heating towards  $T \rightarrow \infty$



Fig.1



Fig.2



# Order parameter correlations

$$\langle m_x(r, t) m_x(0, t) \rangle = r^{-1/2} e^{-\frac{r}{\xi(t)}}, \quad \frac{1}{\xi(t)} = \frac{1}{\xi} + \frac{\Gamma t}{2(1 + \Delta)^2} \quad \text{No Quench}$$

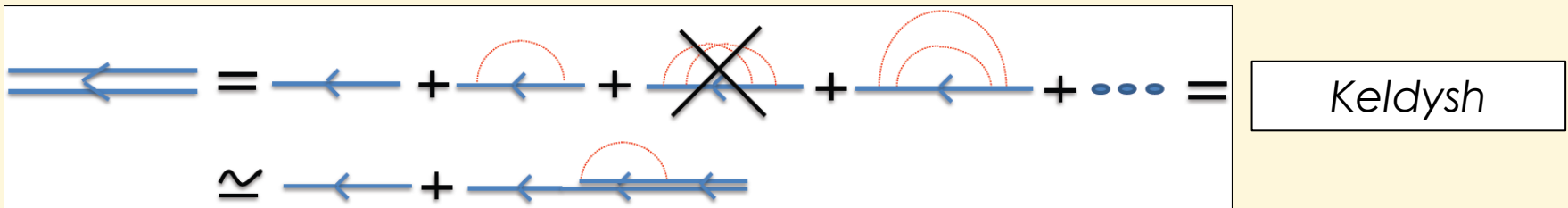
Different signatures in different observables.



# A bit of details ...

$$\rho_k = \begin{pmatrix} \langle \gamma_k^\dagger \gamma_k \rangle & \langle \gamma_k^\dagger \gamma_{-k}^\dagger \rangle \\ \langle \gamma_{-k} \gamma_k \rangle & \langle \gamma_{-k} \gamma_{-k}^\dagger \rangle \end{pmatrix}$$

$$\gamma = \frac{\Gamma}{\Delta} \ll 1$$



$$\partial_t \rho_k = -i[H_k^0, \rho_k] + \frac{\Gamma}{2}(\sigma \rho_k \sigma - \rho_k)$$

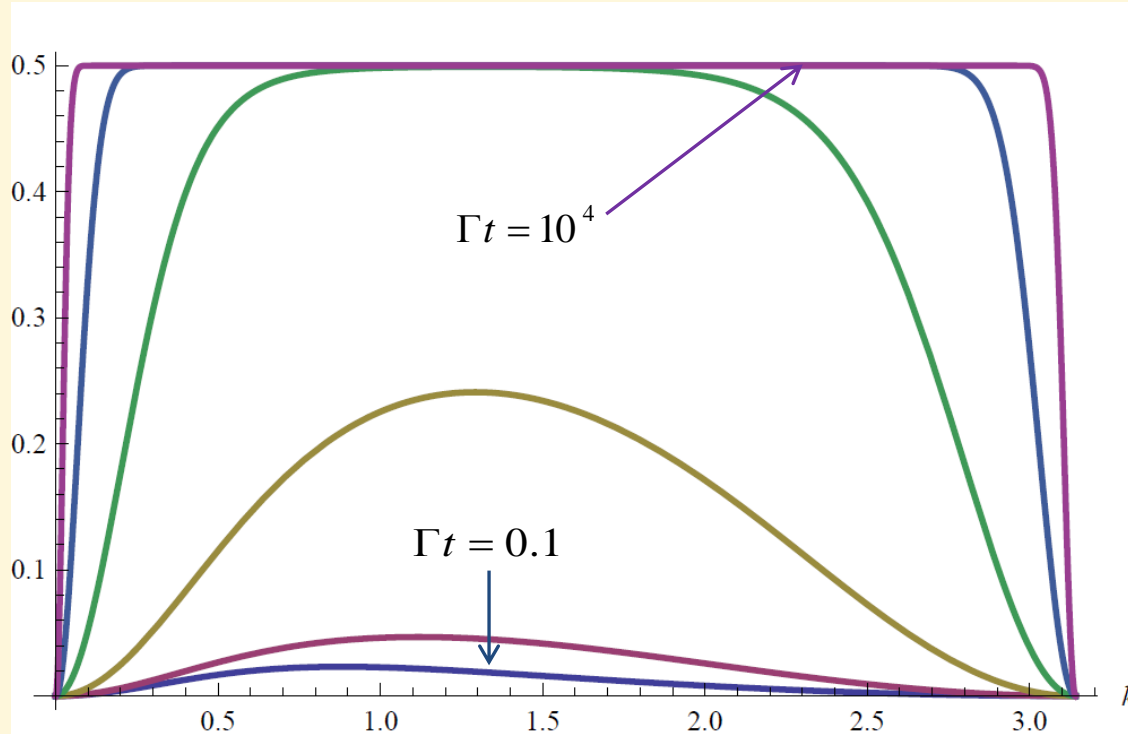
$$\sigma \equiv \cos 2\theta_k \sigma_z + \sin 2\theta_k \sigma_y$$



# Populations

$$\langle \gamma_k^+(t) \gamma_k(t) \rangle = \frac{1}{2} + (\sin^2(\Delta \alpha_k) - \frac{1}{2}) e^{-\Gamma \sin^2 \theta_1 t} \quad \text{Populations}$$

$$\theta(g) = \frac{1}{2} \arctan \left( \frac{\sin k}{g - \cos k} \right) \quad \Delta \alpha_k = \theta(g_1) - \theta(g_0)$$



The relaxation rates tend to vanish close to the band edges ( $k = 0, \pm \pi$ ) for  $\Gamma t \gg 1$

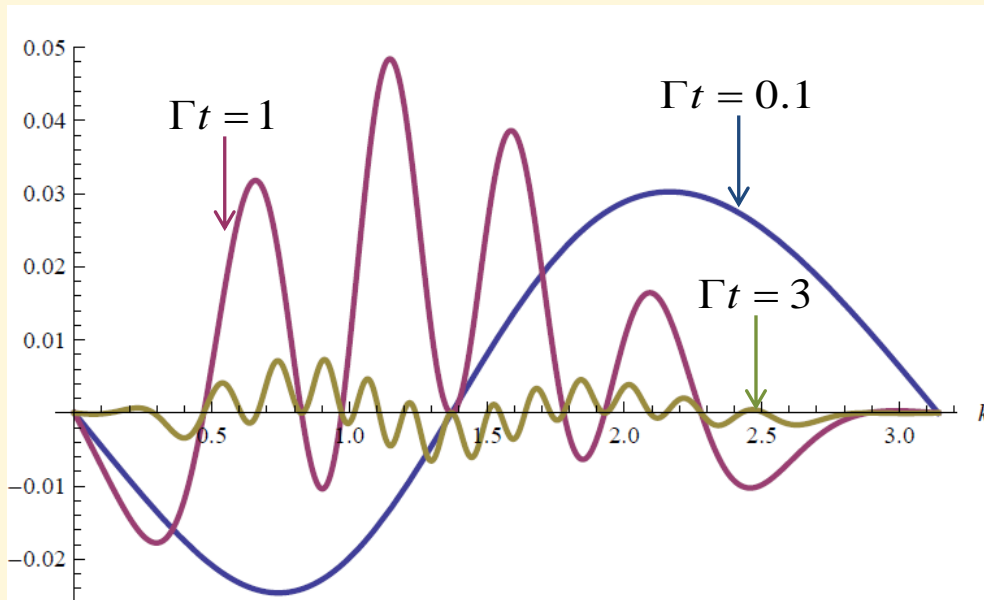
$$\Gamma_k \approx_{(k \approx 0)} \Gamma \frac{k^2}{\Delta^2}$$

Most of the modes relax on time scales of the order of  $1/\Gamma$



# Coherences

$$\langle \gamma_k^+(t) \gamma_{-k}^+(t) \rangle$$

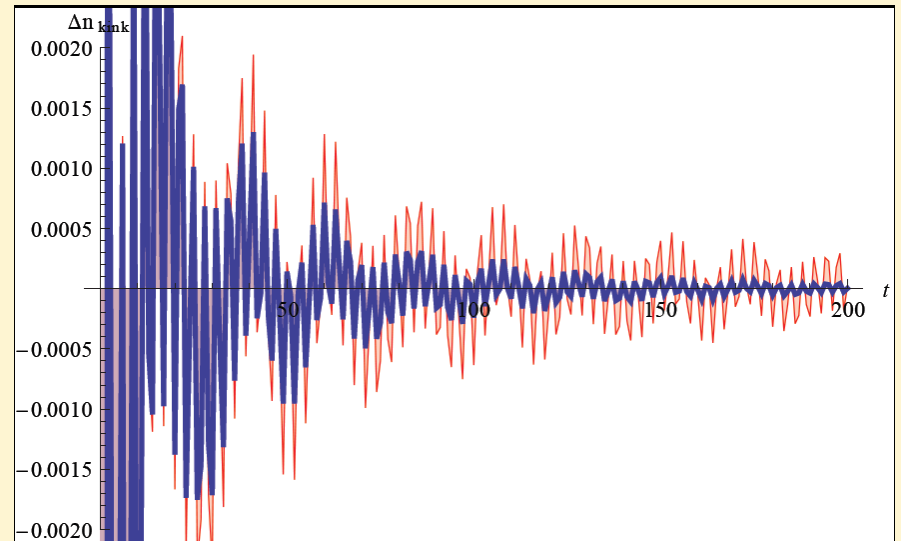
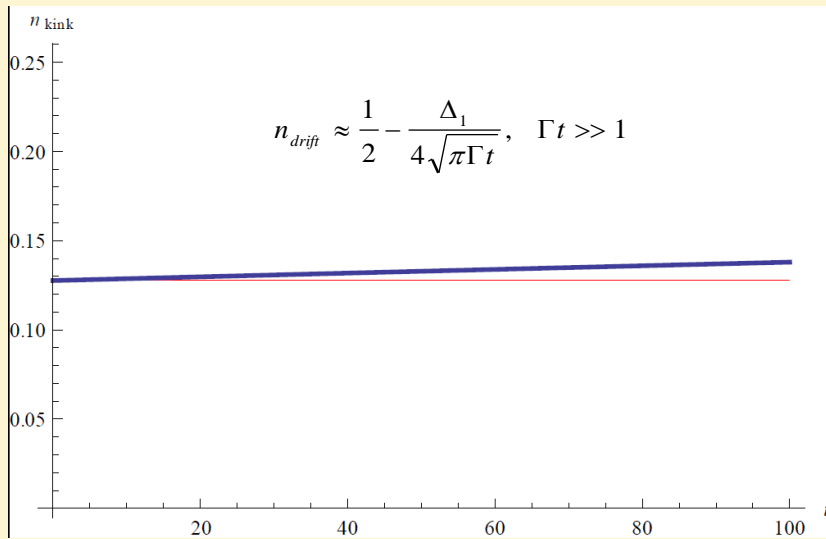


Coherences decay exponentially fast as  $e^{-\Gamma t}$ , and they are subleading compared to populations for  $\Gamma t \gg 1$ .



# Number of kinks

$$n_{kink}(t) = n_{drift}(t) + \Delta n(t)$$



Number of excitations

$$N = \int_0^\pi dk \langle \gamma_k^+ \gamma_k \rangle \approx \frac{\pi}{2} - \frac{\alpha}{\sqrt{\Gamma t}}$$





# Conclusions

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- We study the dynamics of thermalization resulting from a time-dependent noise in a Quantum Ising Chain subject to a sudden quench of  $g$ .
  - For weak noises the dynamics shows a pre-thermalized state at intermediate time scales, eventually drifting towards an asymptotic infinite temperature state, characterized by a diffusive behaviour in the transverse magnetization correlations.
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# Future issues

- Different times correlation functions?
- Fluctuation-dissipation theorem?
- Effect of spatial inhomogeneous noise?
- Integrability breaking in a closed system?

$$V = \sum_i \delta g_i(t) \sigma_i^z$$

