



# Wilson Lines in Higher Spin Gravity

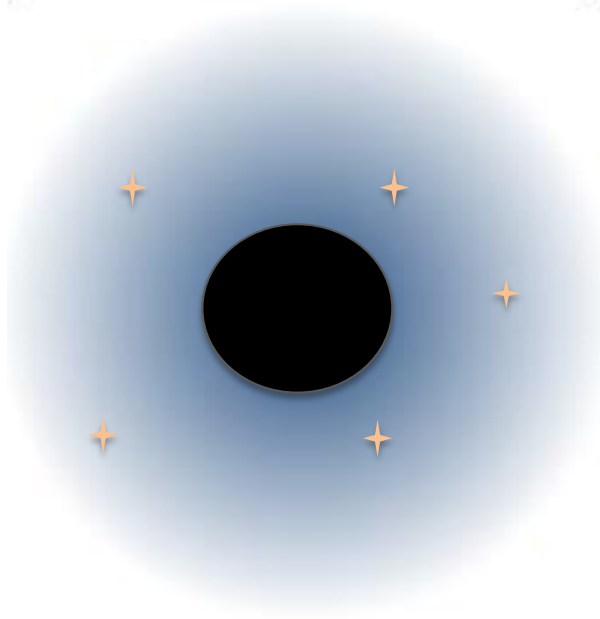
Alejandra Castro  
University of Amsterdam

In collaboration with M. Ammon and N. Iqbal (1306.4338)

“Quantum Fields beyond Perturbation Theory”  
KITP, Santa Barbara  
January 2014

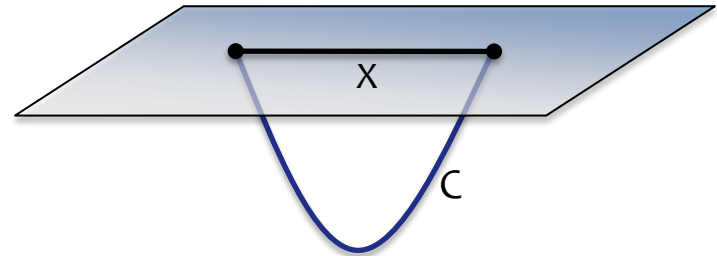
# Universal features of gravity

If Gravity equals Geometry ...



Laws of thermodynamics

[Bekenstein, Hawking]



Quantum information  
(Entanglement Entropy)

[Ryu & Takayanagi]

... Not every gravitational theory admits a geometrical description (or at least not an obvious one)...

# Vasiliev's Higher Spin Gravity

- Infinite number of coupled massless spin fields.
- Gauge group acts non trivially on all fields.

Tractable example of AdS/CFT:  
analytic tractability and intrinsic complexity

[Sezgin & Sundell;  
Klebanov & Polyakov;  
Gaberdiel & Gopakumar]

**Challenge:** Does higher spin gravity know about entanglement?

**Goal:** design a massive probe which will generalize the notion of geodesic for non-local theories.

Our proposal: Wilson line captures the correct dynamics

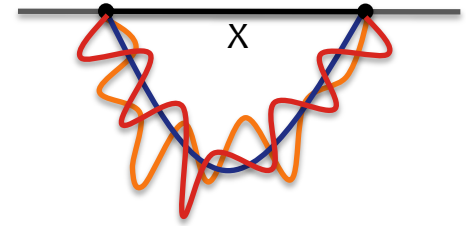
$$W_{\mathcal{R}}(C) = \text{tr}_{\mathcal{R}} \left( \mathcal{P} \exp \int_C \mathcal{A} \right) = \int \mathcal{D}U \exp[-S(U; \mathcal{A})_C]$$

Encodes quantum numbers of the probe. It is an infinite dimensional representation.

Auxiliary field.  
Captures dynamics of the probe.

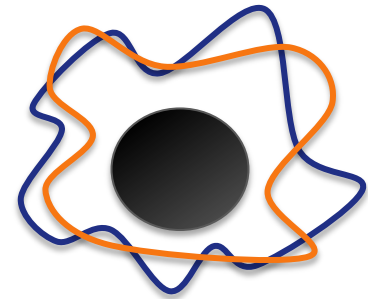
Open Paths

$$S_{\text{EE}} = -\text{Tr}(\rho_X \log \rho_X) = -\log(W_{\mathcal{R}}(C))$$



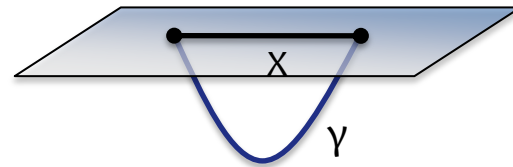
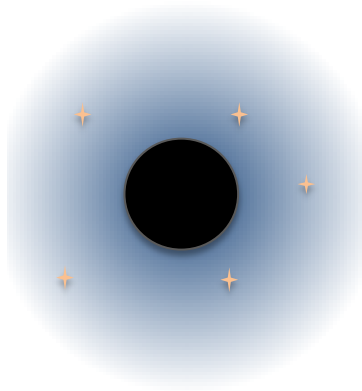
Closed Paths

$$S_{\text{thermal}} = -\log(W_{\mathcal{R}}(C))$$



# AdS<sub>3</sub> Gravity à la Chern-Simons

Reminder: 3D gravity has no local degrees of freedom...



As a topological theory,  
how does 3d gravity know about thermodynamics and entanglement?

## AdS<sub>3</sub> Gravity à la Chern-Simons

$$\mathcal{A} \in so(2, 2) = SL(2, \mathbb{R}) \times SL(2, \mathbb{R})$$

$$S_{\text{EH}} = S_{\text{CS}}[A] - S_{\text{CS}}[\bar{A}]$$

Dictionary to gravity

$$A = \omega + \frac{e}{\ell}$$

$$\bar{A} = \omega - \frac{e}{\ell}$$

Gauge transformations

$$A_\mu \rightarrow L(x) (A_\mu + \partial_\mu) L^{-1}(x)$$

$$\bar{A}_\mu \rightarrow R^{-1}(x) (\bar{A}_\mu + \partial_\mu) R(x)$$

$$L, R \in SL(2, \mathbb{R})$$

## Explicit construction of the Wilson line

$$W_{\mathcal{R}}(C) = \text{tr}_{\mathcal{R}} \left( \mathcal{P} \exp \int_C \mathcal{A} \right) = \int \mathcal{D}U \exp[-S(U; \mathcal{A})_C]$$

[Witten; Carlip]

Effective acting along the curve

$$\begin{aligned} S(U, P; \mathcal{A})_C &= S(U)_{C, \text{free}} + S(U; \mathcal{A})_{C, \text{int}} \\ &= \int_C ds \left( \text{Tr} (PU^{-1} D_s U) + \lambda(s) (\text{Tr}(P^2) - c_2) \right) \end{aligned}$$

**Lagrange multiplier**      **Casimir representation**

Auxiliary field and its conjugate momenta

$$U(s) \rightarrow LU(s)R \quad P(s) \rightarrow R^{-1}P(s)R, \quad L, R \in SL(2, \mathbb{R})$$



## Explicit construction of the Wilson line

$$W_{\mathcal{R}}(C) = \text{tr}_{\mathcal{R}} \left( \mathcal{P} \exp \int_C \mathcal{A} \right) = \int \mathcal{D}U \exp[-S(U; \mathcal{A})_C]$$

Effective acting along the curve

$$\begin{aligned} S(U, P; \mathcal{A})_C &= S(U)_{C,\text{free}} + S(U; \mathcal{A})_{C,\text{int}} \\ &= \int_C ds \left( \text{Tr} (PU^{-1} D_s U) + \lambda(s) (\text{Tr}(P^2) - c_2) \right) \end{aligned}$$

Covariant derivative couples probe to background

$$D_s U = \frac{d}{ds} U + A_s U - U \bar{A}_s, \quad A_s \equiv A_\mu \frac{dx^\mu}{ds}$$

Recovering the geodesic equation: Set  $U(s) = 1$

$$\frac{d}{ds} \left( (A - \bar{A})_{\mu} \frac{dx^{\mu}}{ds} \right) + [\bar{A}_{\mu}, A_{\nu}] \frac{dx^{\mu}}{ds} \frac{dx^{\nu}}{ds} = 0 \quad \text{eom reduces to geodesic eqn!}$$

$$\begin{aligned} S_{C,\text{on-shell}} &= \sqrt{c_2} \int_C ds \sqrt{\text{Tr} \left( (A - \bar{A})_{\mu} (A - \bar{A})_{\nu} \frac{dx^{\mu}}{ds} \frac{dx^{\nu}}{ds} \right)} \\ &= \sqrt{2c_2} \int_C ds \sqrt{g_{\mu\nu}(x) \frac{dx^{\mu}}{ds} \frac{dx^{\nu}}{ds}} \quad \text{Action is the proper distance!} \end{aligned}$$

$$W_{\mathcal{R}}(x_i, x_f) \sim \exp \left( -\sqrt{2c_2} L(x_i, x_f) \right)$$

But please remember: the path is not important! This is a little miracle of  $SL(2, \mathbb{R})$

## Wilson line and Entanglement Entropy

$$W_{\mathcal{R}}(x_i, x_f) \sim \exp\left(-\sqrt{2c_2}L(x_i, x_f)\right)$$

Casimir = mass<sup>2</sup> probe = (dimension twist field in CFT)<sup>2</sup>


$$m = h + \bar{h}$$

$$h = \bar{h} = \frac{c}{24} \left( n - \frac{1}{n} \right)$$

[Cardy & Calabrese]

$$c = \frac{3\ell}{2G_3}$$

[Brown & Henneaux]


$$2c_2 = 4h(h - 1) \xrightarrow{n \rightarrow 1} \left(\frac{c}{6}\right)^2$$

$$S_{\text{EE}} = -\log(W_{\mathcal{R}}(C))$$

Follows from  
Lewkowycz & Maldacena

## Entanglement Entropy

(open path; end points at boundary)

$$S_{\text{EE}} = -\log(W_{\mathcal{R}}(C)) = \frac{c}{3} \log\left(\frac{\Delta X}{\epsilon}\right)$$

Infinite system, zero temperature

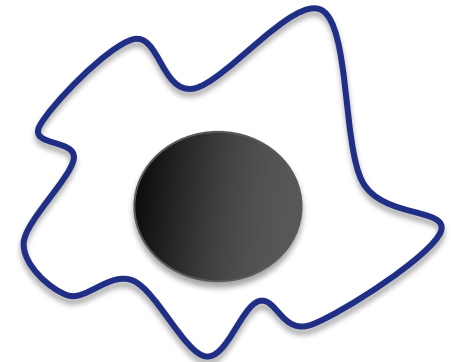
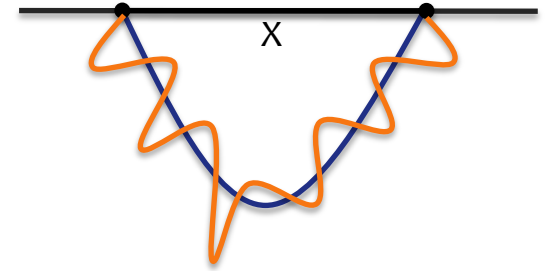
$$S_{\text{EE}} = -\log(W_{\mathcal{R}}(C)) = \frac{c}{3} \log\left(\frac{\beta}{\pi\epsilon} \sinh\left(\frac{\pi\Delta X}{\beta}\right)\right)$$

Infinite system, finite temperature

## Thermal entropy

(closed spatial path)

$$S_{\text{thermal}} = -\log(W_{\mathcal{R}}(C)) = \frac{A_H}{4G_3}$$



# SL(3) Higher Spin Gravity

$$S_{\text{HS}} = S_{\text{CS}}[A] - S_{\text{CS}}[\bar{A}] \quad A, \bar{A} \in SL(3, \mathbb{R})$$

[Bergshoeff, Blencowe & Stelle; Campoleoni, Fredenhagen, Pfenninger & S. Theisen]

Gravitational interpretation: Interacting theory  
of a graviton with a massless spin-3 field

Interesting solutions:

- AdS<sub>3</sub> and BTZ black hole  
[Bañados, Teitelboim & Zanelli]
- Higher Spin Black holes: non-zero spin-3 potential  $(\beta, \mu)$   
[Gutperle & Kraus]
- Domain walls, Janus solutions, conical defects, ...  
[Ammon, Gutperle, Kraus & Perlmutter; Gutperle; AC, Gopakumar, Gutperle & Raeymaekers]

## Constructing a massive probe

$$W_{\mathcal{R}}(C) = \text{tr}_{\mathcal{R}} \left( \mathcal{P} \exp \int_C \mathcal{A} \right) = \int \mathcal{D}U \exp[-S(U; \mathcal{A})_C]$$

Effective acting along the curve

$$S(U, P; A, \bar{A})_C = \int_C ds \left( \text{Tr}(PU^{-1}D_s U) + \lambda_2(\text{Tr}(P^2) - c_2) + \lambda_3(\text{Tr}(P^3) - c_3) \right)$$

Cubic and quadratic casimir

Covariant derivative couples probe to background

$$D_s U = \frac{d}{ds} U + A_s U - U \bar{A}_s, \quad A_s \equiv A_\mu \frac{dx^\mu}{ds}$$

Auxiliary field and its conjugate momenta

$$U(s) \rightarrow LU(s)R, \quad P \rightarrow R^{-1}P(s)R, \quad L, R \in SL(3, \mathbb{R})$$

## Wilson line and Entanglement Entropy

Follow the same logic... Wilson line is a massive probe that induces a conical deficit on the background.

Quadratic Casimir  $\sim \text{mass}^2 \text{ probe} = (\text{dimension twist field in CFT})^2$   
Cubic Casimir  $\sim (\text{spin-3 charge})^3 = (\text{spin-3 of twist field})^3$

$$m = h + \bar{h}$$

$$h = \bar{h} = \frac{c}{24} \left( n - \frac{1}{n} \right)$$

$$c = \frac{3\ell}{2G_3}$$

[Campoleoni, Fredenhagen, Pfenninger & Theisen; Henneaux & Rey]

$$2c_2 = 4h(h-1) \xrightarrow{n \rightarrow 1} \left( \frac{c}{6} \right)^2$$

$$c_3 = \frac{3}{8}w \left( h^2 - \frac{1}{4}w^2 \right) = 0$$



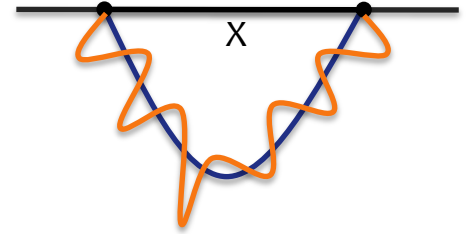
$$S_{\text{EE}} = -\log(W_{\mathcal{R}}(C))$$

## Entanglement Entropy

(open path; end points at boundary)

$$S_{\text{EE}} = \Delta X \left( 2\sqrt{2\pi\mathcal{L}k} \frac{\sqrt{1 - \frac{3}{4C}}}{1 - \frac{3}{2C}} \right)$$

( $\Delta X \gg \beta$ ;  $\mu$  : fixed;  $\beta$  : fixed)



$$S_{\text{EE}} = \frac{c}{3} \log \left( \frac{\beta}{\pi\epsilon} \sinh \left( \frac{\pi\Delta X}{\beta} \right) \right)$$

( $\mu \rightarrow 0$  ;  $\beta$  : fixed)

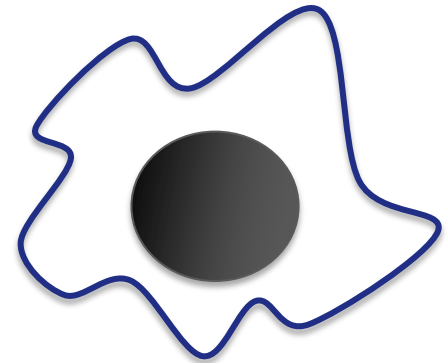
$$S_{\text{EE}} = \frac{c}{3} \log \left( \frac{\Delta X}{\epsilon} \left| 1 - \frac{16\mu^2}{\Delta X^2} \right|^{1/4} \right)$$

( $\beta \rightarrow \infty$  ;  $\mu$  : fixed)

## Thermal entropy

(closed spatial path)

$$S_{\text{thermal}} = 2\pi \sqrt{\frac{c_2}{2}} \text{tr}_f((\lambda_\phi - \bar{\lambda}_\phi)L_0)$$





## Comments

- If the background is AdS (or BTZ) we recover all previous formulas
- Setting  $U=1$  is not allowed. Geodesic equation is not recovered (which is good!).
- Good agreement with thermodynamics of higher spin black holes.
- Agreement with J. de Boer and J.I. Jottar [[1306.4347\[hep-th\]](#)]
- Puzzles for domain wall backgrounds. Not every background satisfy monotonicity of entanglement entropy.
- Preliminary results for conical defects: non-trivial agreements with CFT!
- Extending results of Cardy-Calabrese for non-vanishing spin-3 potential is needed in the CFT.
- Still more to explore!

## Summary

$$W_{\mathcal{R}}(C) = \text{tr}_{\mathcal{R}} \left( \mathcal{P} \exp \int_C \mathcal{A} \right) = \int \mathcal{D}U \exp[-S(U; \mathcal{A})_C]$$

- A new way to cast RT formula for AdS<sub>3</sub> gravity.
- For closed paths, our proposal computes thermal entropy.
- Our formula reproduces non-trivial results of entanglement entropy.

# Outlook

- Generalization to Vasiliev's theory with infinite number of higher spin fields.
- Interpretation of quantum corrections to the Wilson line.
- Renyi entropies in Chern-Simons formulation. [w J. de Boer and J. I. Jottar]
- Entanglement entropy in the presence of diff anomalies. [w M. Ammon, S. Detournay, N. Iqbal and E. Perlmutter]
- Puzzles regarding sub-additivity. Better understanding in the CFT is needed.