Large-D Gravity Progress & Prospects

Roberto Emparan ICREA & UBarcelona Geometry from the Quantum - KITP UCSB 16 Jan 2020



Work in 2013-2019 with:

T Andrade, D Grumiller, K Izumi, D Licht, R Luna, M Martínez, T Shiromizu, **R Suzuki, K Tanabe**, T Tanaka Parallel program by *S Bhattacharyya, S Minwalla et al* Effective membrane theory of black holes Same concepts but different implementation



A technique to simplify calculations

A way to gain new insights

D as a parameter

D < 4

UV regulator – good for quantum

But :(No local dof's No GWs No BHs in $R_{\mu\nu} = 0$

D as a parameter

D > 4

IR regulator – bad for quantum?

But :) Has local dof's GWs BHs in $R_{\mu\nu} = 0$

Large N does simplify gauge theories

Reorganize meson dynamics into string worldsheets

Large D: does it simplify gravity?

YES

Reorganizes black hole dynamics into theory of membranes

Large D: does it simplify other aspects of gravity?

Probably yes

Large-D in General Relativity

D² ~ # local degrees of freedom at a point akin to Large-N SU(N) gauge theory also large-c CFT, vector models, Potts models, SYK...

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D ~ # connections between nearby points = directions out of a point akin to Mean Field Theory limit in Stat Mech

average out/dilute long-distance effects

What we've found useful is

D² ~ # local degrees of freedom at a point akin to Large-N SU(N) gauge theory

D ~ # connections between nearby points
= directions out of a point

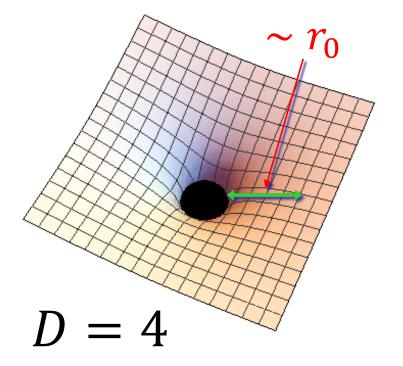
Exploit large gradients of gravitational potential $\frac{1}{r^{D-3}}$

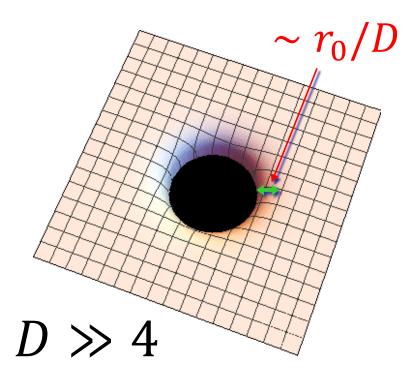
Main findings about BHs @ large D

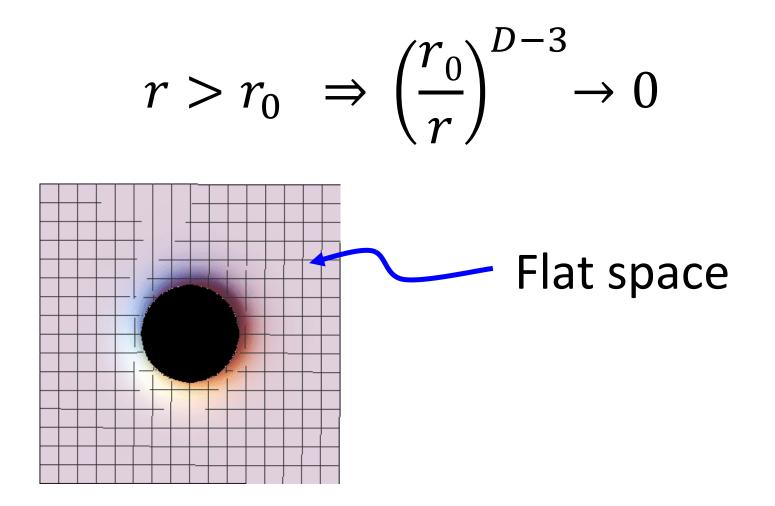
∃ well-defined, universal near-horizon geometry with dynamics that decouples from the far-region

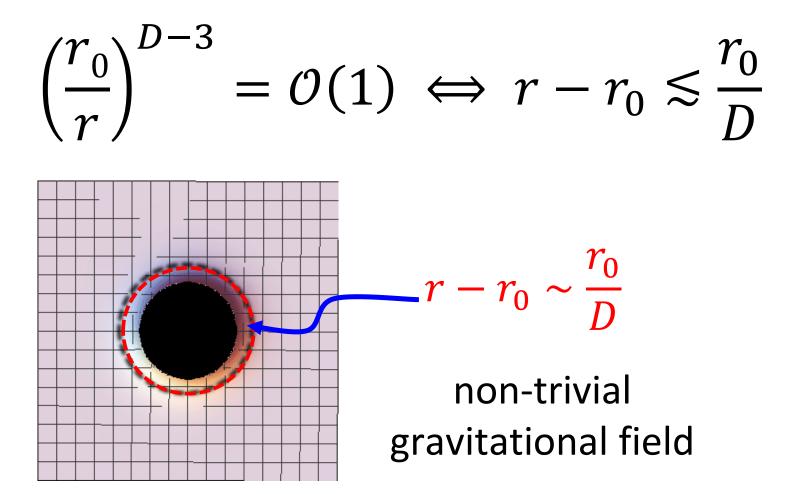
This allows to formulate an effective theory of dynamical black holes for slow fluctuations

Similar to but not the same as membrane paradigm, fluid/gravity







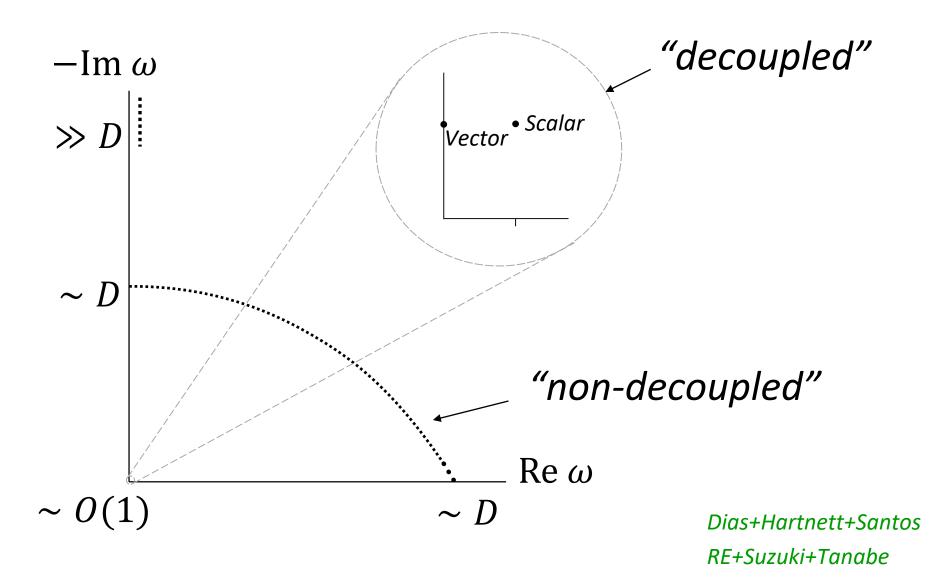


Two kinds of BH fluctuations

• Slow, localized near horizon, decoupled from far zone

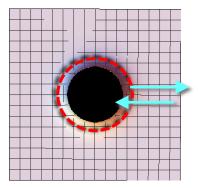
• Fast, everywhere, non-decoupled

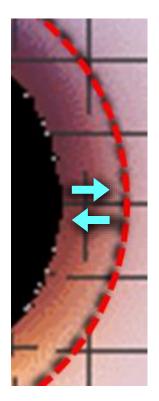
Quasinormal frequencies of Schwarzschild_{D $\gg1$}



Fast, non-decoupled $\omega \sim D/r_0$

Characteristic crossing time of near-horizon





Slow, decoupled

 $\omega \sim 1/r_0$

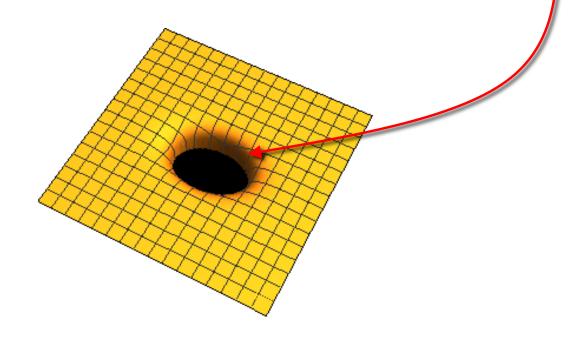
Almost static in near-horizon

Decoupled from far zone

Non-linear dynamics

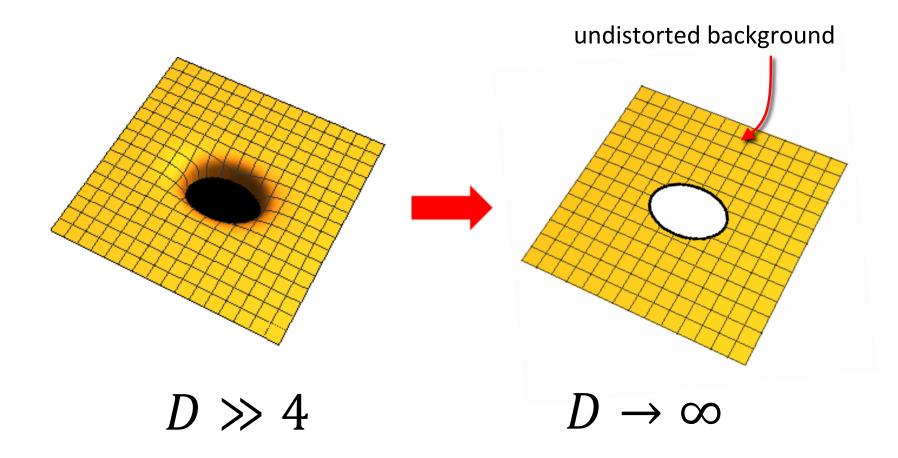
Effective theory of slow decoupled fluctuations

All the black hole dynamics is concentrated here

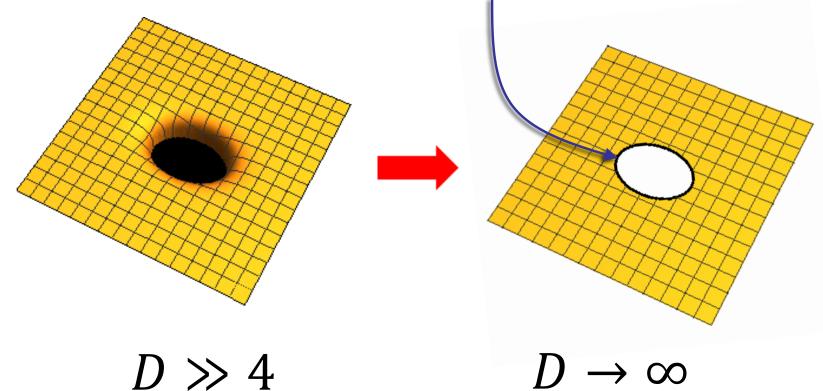


 $D \gg 4$

Replace bh → Surface ('membrane')



What's the dynamics of this membrane?



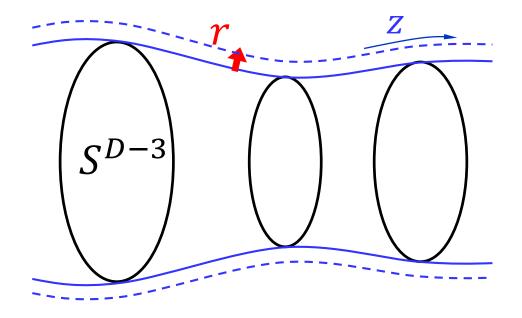
Solve Einstein equations in near-horizon

→ *Effective membrane theory*

Non-linear effective theory of lightest quasinormal modes

Gradient hierarchy

 $\perp \text{ Horizon: } \partial_{\gamma} \sim D$ $\parallel \text{ Horizon: } \partial_{z} \sim 1 \quad (\text{or} \sim \sqrt{D})$



Effective equations

Most general & elegant formulation by Bhattacharyya+Minwalla et al

$$\left(\frac{\nabla^2 u}{\mathcal{K}} - \frac{\nabla \mathcal{K}}{\mathcal{K}} + u \cdot K - (u \cdot \nabla)u\right) \cdot \mathcal{P} = 0$$

$$\mathcal{K} = \eta^{AB} K_{AB}$$
$$\nabla \cdot u = 0, \qquad n \cdot u = 0$$
$$\mathcal{P}_{AB} = \eta_{AB} - n_A n_B + u_A u_B$$

n, *K*_{AB}: normal & extrinsic curvature of membrane

u: **velocity** field on membrane

Simplifies -conceptually and technicallyin two important cases:

1. Stationary black holes

2. Black branes, AdS or AF

Stationary solution

Soap-bubble equation (redshifted)

$$K = 2\gamma\kappa$$

K = trace **extrinsic curvature** of membrane $\gamma =$ **redshift** factor on membrane

 $\kappa = surface gravity$

Effective equations for fluctuating black brane

effective fields

 $m(t, z^i)$: mass and area density of black brane $p_i(t, z^j)$: pressure/momenta along brane

$$\partial_{t}m - \nabla^{2}m = -\nabla_{i}p^{i}$$
$$\partial_{t}p_{i} - \nabla^{2}p_{i} = \underbrace{\pm}_{f}\nabla_{i}m - \nabla_{j}\left(\frac{p_{i}p^{j}}{m}\right)$$
$$\underset{\text{AF/AdS}}{\overset{f}{\longrightarrow}}$$

$$\begin{split} \partial_t m - \nabla^2 m &= -\nabla_i p^i \\ \partial_t p_i - \nabla^2 p_i &= \pm \nabla_i m - \nabla_j \left(\frac{p_i p^j}{m} \right) \end{split}$$

Black brane:
$$m = 1, p_i = 0$$

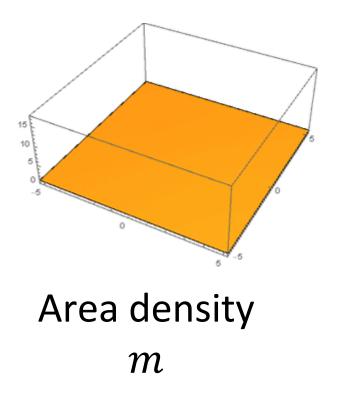
Fluctuate: $m(t, z^i), p_i(t, z^j)$

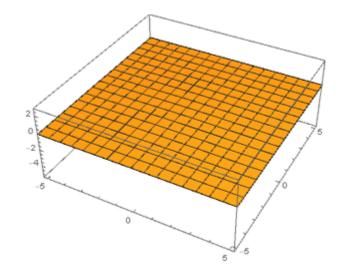
$$\partial_t m - \nabla^2 m = -\nabla_i p^i$$
$$\partial_t p_i - \nabla^2 p_i = \pm \nabla_i m - \nabla_j \left(\frac{p_i p^j}{m}\right)$$
almost linear diffusion equations, except for

BH collisions and Cosmic Censorship violation in Hi-D



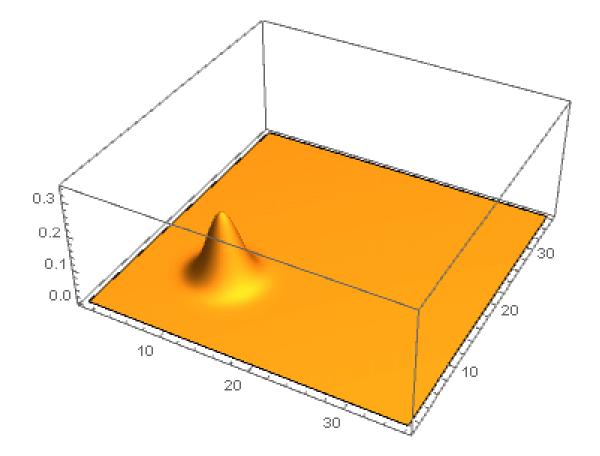
"Black hole blob" in a black membrane





Area radius $r_H = \ln m$

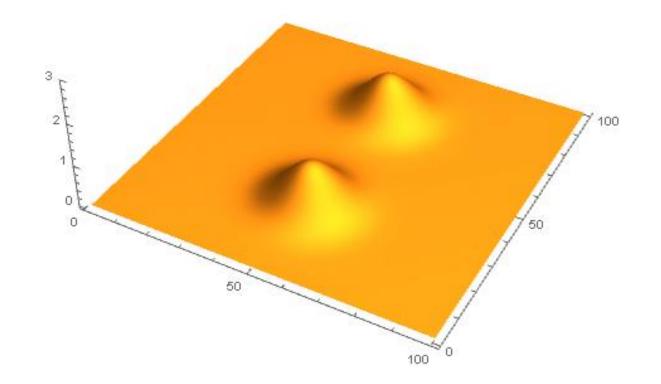
Moving black hole



Collisions of black hole blobs

Brane acts as a "regulator": continuous horizon

BHs never really merge nor split: smooth evolution



Large D and the Quantum

Large D gravity and the Quantum

Still at preliminary stage

A few studies, some general observations

Not yet a systematic understanding nor "effective theory" Short-distance quantum fluctuations strongly enhanced

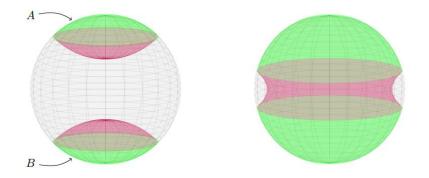
 Long-distance quantum fluctuations average out

Holo-dual to behavior of gravitational field at large D

Entanglement structure

Colin-Ellerin, Hubeny, Niehoff, Sorce 2019

Phase transition in holographic mutual information



Analytical solution @ large D

Entanglement structure

At $D \to \infty$

Entanglement between neighbors vanishes too many neighbors

Spatial decoupling of holographic correlations

Entanglement between neighbors vanishes at $D \rightarrow \infty$

In the bulk this does not involve actual large field gradients

RT/HRT surfaces explore full dimensionality

Can this be exploited/encoded in some effective theory?

Holographic renormalization RE+Suzuki

UV structure of holographic CFTs (planar) simplifies at $D \rightarrow \infty$

Fefferman-Graham expansion

$$ds^{2} = \frac{1}{z^{2}} (dz^{2} + g_{ij}(z, x) dx^{i} dx^{j})$$
$$g(z, x) = g^{0}(x) + z^{2}g^{1}(x) + \dots + z^{D-1} \langle T(x) \rangle + \dots$$

UV, state-independent Counterterms Quantum effective action IR, state-dependent

vevs

Fefferman-Graham expansion

$$ds^{2} = \frac{1}{z^{2}} (dz^{2} + g_{ij}(z, x) dx^{i} dx^{j})$$
$$g(z, x) = g^{0}(x) + z^{2}g^{1}(x) + \dots + z^{D-1}(T(x)) + \dots$$

UV and IR decoupled at all 1/D orders

IR: effective theory of black branes UV: boundary expansion

Bulk reconstruction at $D \rightarrow \infty$

$$ds^{2} = \frac{1}{z^{2}}(dz^{2} + g_{ij}(z, x)dx^{i}dx^{j} + d\Sigma_{D-p-1})$$

Bulk is reconstructed as Ricci flow from boundary geometry

$$\frac{\partial g_{ij}}{\partial \log z} = 2R_{ij}$$

Counterterm action truncates finitely at each order in 1/D

Large D gravity and the Quantum

Scales in Hawking radiation emission time, scrambling time, evaporation time Hod

Holdt-Sørensen, McGady, Wintergerst

String theory @ large D: an old story

Alvarez, Ambjorn, Gubser...

Ferrari

Large-N + large-D tensor model X_{μ}^{ij}

Large D gravity and the Quantum

still lacking basic concepts what problems are good? much to explore... Thank you