

Operator Product Expansion

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based on joint work with J. Holland and Ch. Kopper

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Introduction

Operator Product Expansion [Wilson '69]

Products of composite fields can be expanded as

$$\langle \mathcal{O}_{A_1}(x_1) \cdots \mathcal{O}_{A_N}(x_N) \underbrace{\cdots}_{\text{Spectators}} \rangle \sim \sum_B \underbrace{\mathcal{C}_{A_1 \dots A_N}^B(x_1, \dots, x_N)}_{\text{OPE coefficients}} \langle \mathcal{O}_B(x_N) \cdots \rangle$$

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- ▶ **Asymptotic short distance expansion:**

Difference vanishes in the limit $x_i \rightarrow x_N$ for all $i \leq N$

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- ▶ Asymptotic short distance expansion:
Difference vanishes in the limit $x_i \rightarrow x_N$ for all $i \leq N$
- ▶ Practical application e.g. in deep-inelastic scattering
- ▶ Plays fundamental role in conformal field theory
(Conformal bootstrap, "Vertex operator algebras", ...)
- ▶ Plays fundamental role in QFTCS
(State-independent definition of QFT!)

Topics of today's talk:

1. In what sense does the OPE converge? N -point functions \leftrightarrow 1-point functions & OPE coefficients
2. What are algebraic relations between OPE coefficients?
Vertex algebras in d -dims.
3. A novel recursion scheme for OPE coefficients
New self-consistent construction method

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New self-consistent construction method

Model: Perturbative, Euclidean φ_4^4 -theory

The model: Euclidean φ^4 -theory

- Correlation functions are defined via the path integral

$$\langle \mathcal{O}_{A_1}(x_1) \cdots \mathcal{O}_{A_N}(x_N) \rangle := \mathcal{N} \int \mathcal{D}\varphi \exp[-S] \mathcal{O}_{A_1}(x_1) \cdots \mathcal{O}_{A_N}(x_N),$$

where the action is given by

$$S(\varphi) := \int d^4x \left(\frac{1}{2} (\partial_\mu \varphi)^2(x) + \frac{m^2}{2} \varphi^2(x) + g\varphi(x)^4 - \text{counterterms} \right)$$

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- ▶ Composite operator insertions $\mathcal{O}_A(x) = (\partial^{v_1} \varphi \dots \partial^{v_n} \varphi)(x)$ (renormalized)
- ▶ OPE coefficients can be defined a la Zimmermann or a la Keller-Kopper
- ▶ We use a “renormalization group flow equation” approach [Wilson, Polchinski, Kopper-Keller-Salmhofer]

Outline

- 1 OPE factorisation
- 2 OPE convergence
- 3 Recursive construction of OPE

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2 OPE convergence

3 Recursive construction of OPE

The OPE factorises

Theorem (Holland-SH)

At any arbitrary but fixed loop order:

$$C_{A_1 \dots A_N}^B(x_1, \dots, x_N) = \sum_C C_{A_1 \dots A_M}^C(x_1, \dots, x_M) C_{C A_{M+1} \dots A_N}^B(x_M, \dots, x_N)$$

holds on the domain $\frac{\max_{1 \leq i \leq M} |x_i - x_M|}{\min_{M < j \leq N} |x_j - x_M|} < 1$. (Sum over C absolutely convergent !)

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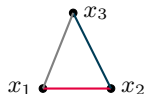
$$\mathcal{C}_{A_1 \dots A_N}^B(x_1, \dots, x_N) = \sum_C \mathcal{C}_{A_1 \dots A_M}^C(x_1, \dots, x_M) \mathcal{C}_{C A_{M+1} \dots A_N}^B(x_M, \dots, x_N)$$

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For $N = 3$: $\varepsilon = \frac{|x_1 - x_2|}{|x_2 - x_3|} < 1$



for $\varepsilon \ll 1$



for $\varepsilon \approx 1$

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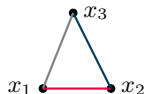
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This shows associativity really holds!

- ▶ Vertex Algebras (Borchers property) also in $4d$.
- ▶ $\mathcal{C}_{A_1 \dots A_N}^B$ uniquely determined in terms of $\mathcal{C}_{A_1 A_2}^B$
- ▶ "Bootstrap construction" of OPE coefficients possible

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1 OPE factorisation

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Bound on OPE remainder I

Theorem (Holland-Kopper-SH)

At any perturbation order r and for any $D \in \mathbb{N}$,

$$\left| \left\langle \left(\mathcal{O}_{A_1}(x_1) \cdots \mathcal{O}_{A_N}(x_N) - \underbrace{\sum_{\dim[B] \leq D} \mathcal{C}_{A_1 \dots A_N}^B(x_1, \dots, x_N) \mathcal{O}_B(x_N)}_{\text{OPE-Remainder}} \right) \underbrace{\hat{\varphi}(p_1) \cdots \hat{\varphi}(p_n)}_{\text{Spectator fields}} \right\rangle \right|$$

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► $M = \begin{cases} m & \text{for } m > 0 \\ \mu & \text{for } m = 0 \end{cases}$ mass or renormalization scale

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- ▶ $|P| = \sup_i |p_i|$: maximal momentum of spectators
- ▶ $\kappa := \inf(\mu, \varepsilon)$, where $\varepsilon = \min_{I \subset \{1, \dots, n\}} |\sum_I p_i|$
 ε : distance of (p_1, \dots, p_n) to “exceptional” configurations

Conclusions from bound on OPE remainder

$$\text{"OPE remainder"} \leq \frac{M^{n-1}}{\sqrt{D!}} \frac{\left(KM \max_{1 \leq i \leq N} |x_i - x_N|\right)^{D+1}}{\min_{1 \leq i < j \leq N} |x_i - x_j|^{\sum_i \dim[A_i] + 1}} \cdot \sup \left(1, \frac{|P|}{\sup(m, \kappa)}\right)^{(D+2)(r+5)}$$

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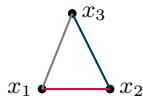
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 - ▶ ratio of max. and min. distances is large, e.g. for $N = 3$



Slow convergence



Fast convergence

Bound on OPE remainder II

Consider now smeared spectator fields $\varphi(f_i) = \int f_i(x)\varphi(x) d^4x$.

Theorem (Holland-Kopper-SH)

At any perturbation order r and for any $D \in \mathbb{N}$,

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M : mass for $m > 0$ or renormalization scale μ for massless fields

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1. Bound is finite for any $f_i \in \mathcal{S}(\mathbb{R}^4)$ (Schwartz space)
OPE remainder is a tempered distribution
2. Let $\hat{f}_i(p) = 0$ for $|p| > |P|$: Bound vanishes as $D \rightarrow \infty$
 \Rightarrow OPE converges at any finite distances!

Outline

1 OPE factorisation

2 OPE convergence

3 Recursive construction of OPE

Motivation for a new construction method

Textbook method (roughly):

- ▶ Write down correlation function with operator insertions
- ▶ Perform short distance/large momentum expansion (in some clever way)
- ▶ Argue that the coefficients obtained this way are state independent

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Not entirely satisfying:

- ▶ Relies on correlation functions \Rightarrow OPE not 'fundamental'
- ▶ State independence not obvious
- ▶ Hard to study general properties of OPE

Recursion formula (for mass $m > 0$)

Theorem (Hollands-JH)

Coupling constant derivatives of OPE coefficients in $g\varphi^4$ -theory can be expressed as

$$\begin{aligned} \partial_g \mathcal{C}_{A_1 \dots A_N}^B(x_1, \dots, x_N) = & - \int d^4 y \left[\mathcal{C}_{\varphi^4 A_1 \dots A_N}^B(y, x_1, \dots, x_N) \right. \\ & - \sum_{i=1}^N \sum_{[C] \leq [A_i]} \mathcal{C}_{\varphi^4 A_i}^C(y, x_i) \mathcal{C}_{A_1 \dots \widehat{A_i} C \dots A_N}^B(x_1, \dots, x_N) \\ & \left. - \sum_{[C] < [B]} \mathcal{C}_{A_1 \dots A_N}^C(x_1, \dots, x_N) \mathcal{C}_{\varphi^4 C}^B(y, x_N) \right]. \end{aligned}$$

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- ▶ Compute OPE coefficients to any perturbation order by iteration.
Initial data: Coefficients of free theory.

Recursion formula (for mass $m > 0$)

Theorem (Hollands-JH)

OPE coefficients at perturbation order $(r + 1)$ can be expressed as

$$\begin{aligned} (\mathcal{C}_{r+1})_{A_1 \dots A_N}^B(x_1, \dots, x_N) = & - \int d^4 y \left[(\mathcal{C}_r)_{\varphi^4 A_1 \dots A_N}^B(y, x_1, \dots, x_N) \right. \\ & - \sum_{s=0}^r \sum_{i=1}^N \sum_{[C] \leq [A_i]} (\mathcal{C}_s)_{\varphi^4 A_i}^C(y, x_i) (\mathcal{C}_{r-s})_{A_1 \dots \widehat{A}_i C \dots A_N}^B(x_1, \dots, x_N) \\ & \left. - \sum_{s=0}^r \sum_{[C] < [B]} (\mathcal{C}_s)_{A_1 \dots A_N}^C(x_1, \dots, x_N) (\mathcal{C}_{r-s})_{\varphi^4 C}^B(y, x_N) \right]. \end{aligned}$$

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Recursion formula (for mass $m > 0$)

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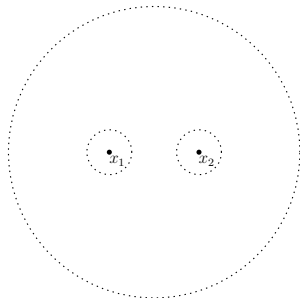
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- ▶ Compute OPE coefficients to any perturbation order by iteration.
Initial data: Coefficients of free theory.
- ▶ State independence obvious.
No other objects enter the construction.
- ▶ The formula depends on the renormalisation conditions.
(Here BPHZ)

Built-in renormalisation (Example: $N = 2$)

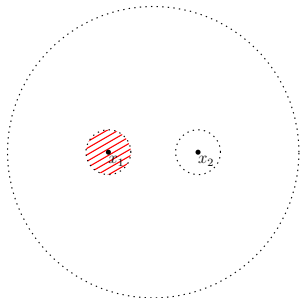
$$\int d^4y \left[\mathcal{C}_{\varphi^4 A_1 A_2}^B(y, x_1, x_2) - \sum_{[C] \leq [A_1]} \mathcal{C}_{\varphi^4 A_1}^C(y, x_1) \mathcal{C}_{C A_2}^B(x_1, x_2) \right. \\ \left. - \sum_{[C] \leq [A_2]} \mathcal{C}_{\varphi^4 A_2}^C(y, x_2) \mathcal{C}_{A_1 C}^B(x_1, x_2) - \sum_{[C] < [B]} \mathcal{C}_{A_1 A_2}^C(x_1, x_2) \mathcal{C}_{\varphi^4 C}^B(y, x_2) \right]$$



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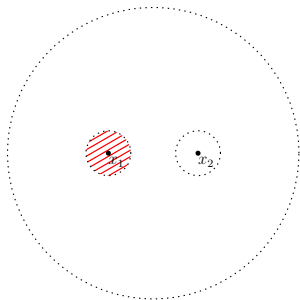
UV-region I ($y \approx x_1$): $\mathcal{C}_{\varphi^4 A_1 A_2}^B$ factorises



Built-in renormalisation (Example: $N = 2$)

$$\int d^4y \left[\sum_{[C]=0}^{\infty} \mathcal{C}_{\varphi^4 A_1}^C(y, x_1) \mathcal{C}_{CA_2}^B(x_1, x_2) - \sum_{[C] \leq [A_1]} \mathcal{C}_{\varphi^4 A_1}^C(y, x_1) \mathcal{C}_{CA_2}^B(x_1, x_2) \right. \\ \left. - \sum_{[C] \leq [A_2]} \mathcal{C}_{\varphi^4 A_2}^C(y, x_2) \mathcal{C}_{A_1 C}^B(x_1, x_2) - \sum_{[C] < [B]} \mathcal{C}_{A_1 A_2}^C(x_1, x_2) \mathcal{C}_{\varphi^4 C}^B(y, x_2) \right]$$

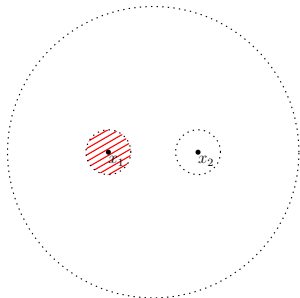
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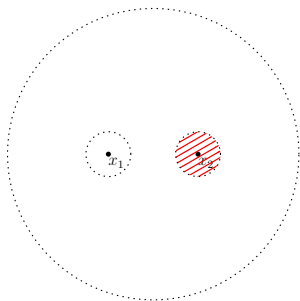
UV-region I ($y \approx x_1$): $\mathcal{C}_{\varphi^4 A_1 A_2}^B$ factorises \Rightarrow divergences cancel



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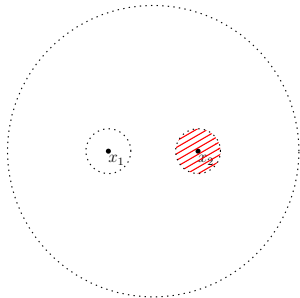
UV-region II ($y \approx x_2$): $\mathcal{C}_{\varphi^4 A_1 A_2}^B$ factorises



Built-in renormalisation (Example: $N = 2$)

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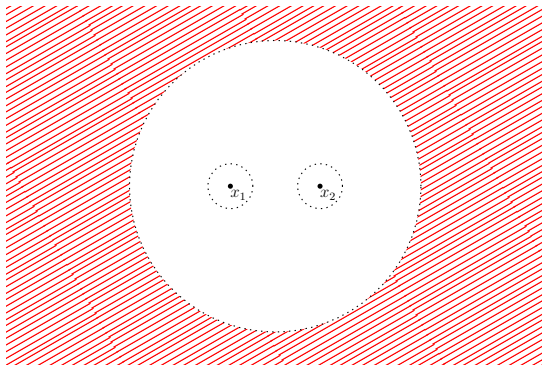
UV-region II ($y \approx x_2$): $\mathcal{C}_{\varphi^4 A_1 A_2}^B$ factorises \Rightarrow divergences cancel



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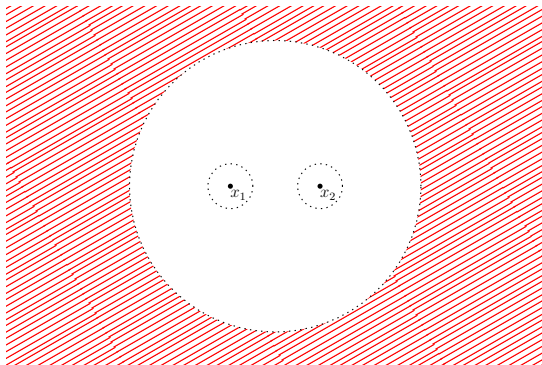
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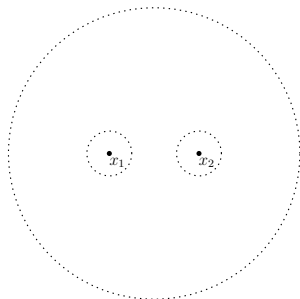
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The integral is absolutely convergent due to the factorisation property.



Conclusions & Outlook

In Euclidean perturbation theory, we found that:

1. The OPE converges at finite distances.
2. The OPE factorises (associativity).
3. The OPE satisfies a recursion formula.

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Possible Generalisations

- ▶ Gauge theories (in progress)
- ▶ Curved manifolds
- ▶ Minkowski space
- ▶ ...

Applications of the Recursion Formula

- ▶ Does the algorithm facilitate computations?
- ▶ Does the perturbation series for OPE coefficients converge?