

# **Quantum Correction to Conductivity Close to Ferromagnetic Quantum Critical Point in 2d**

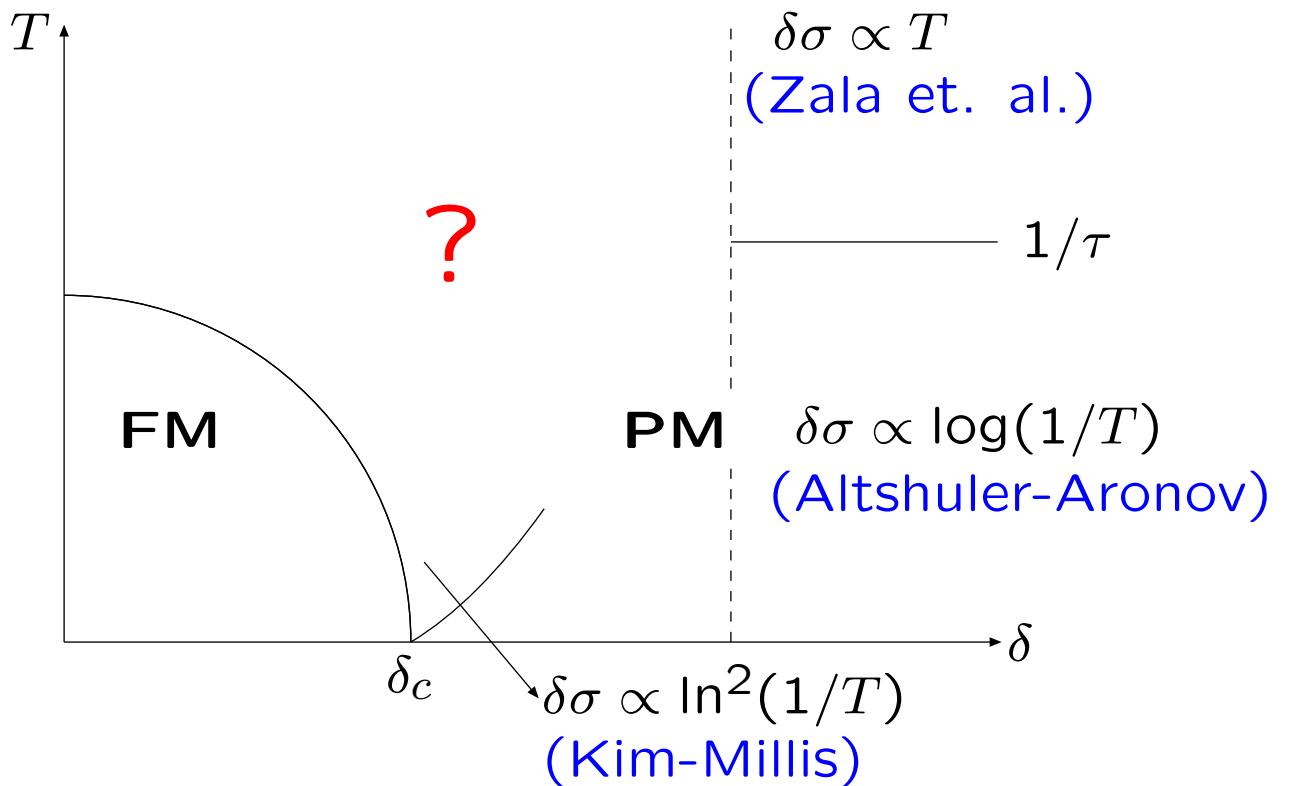
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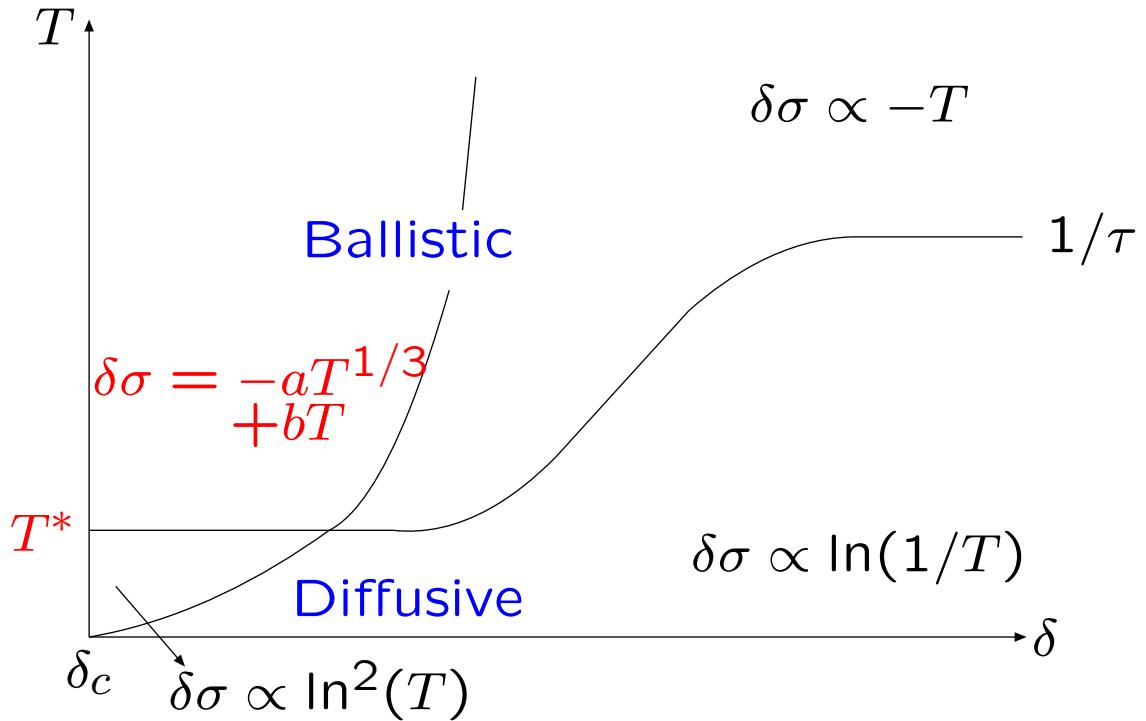
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## Effect of quantum interference on conductivity



- What is the influence of FM QCP on quantum interference processes?
- which is more important:  
elastic or inelastic scattering?

## Main Results



1. near QCP ballistic-diffusive cross-over at  $T^* \ll 1/\tau$ .  
 $T^* \approx 1/((E_F\tau)^2\tau)$
2. near QCP, ballistic regime:  
 $\delta\sigma = \sigma(T) - \sigma(T=0) = -aT^{1/3}$ .

## Spin-Fermion Model

$$\begin{aligned} S = & \frac{1}{\beta} \sum_{\omega_n} \int d^2r \psi_\alpha^\dagger(\mathbf{r}, \omega_n) \left( i\omega_n + \frac{\nabla^2}{2m} + \mu \right) \psi_\alpha(\mathbf{r}, \omega_n) \\ & + \frac{1}{\beta \chi_0} \sum_{\Omega_n, \mathbf{q}} D^{-1}(\mathbf{q}, \Omega_n) \mathbf{S}(\mathbf{q}, \Omega_n) \cdot \mathbf{S}(-\mathbf{q}, -\Omega_n) \\ & + \left( \frac{\alpha}{\chi_0 \nu} \right)^{1/2} \int d^2r \int_0^\beta d\tau \psi_\alpha^\dagger(\mathbf{r}, \tau) \psi_\beta(\mathbf{r}, \tau) [\mathbf{S}(\mathbf{r}, \tau) \cdot \boldsymbol{\sigma}_{\alpha\beta}] \\ & + \int_0^\beta d\tau \int d^2r \psi_\alpha^\dagger(\mathbf{r}, \tau) \mathbf{V}(\mathbf{r}) \psi_\alpha(\mathbf{r}, \tau). \end{aligned}$$

## Spin-Fermion Model

$$\overrightarrow{\text{---}} = (\epsilon - \epsilon_{\mathbf{k}} \pm i/(2\tau))^{-1}$$

$$\overbrace{\text{---}} = D(\mathbf{q}, \Omega)$$

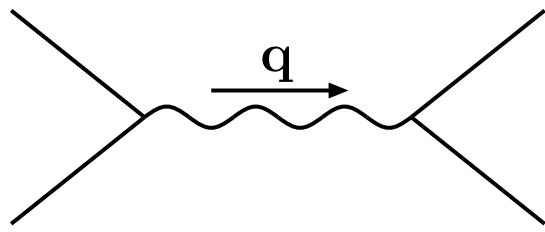
$$\overline{\text{-----}} = 1/(2\pi\nu\tau)$$

$$\begin{array}{c} \diagup \\ \diagdown \end{array} \text{---} = \sqrt{\alpha}$$

$$D(\mathbf{q}, \Omega_n) = \left[ \delta + \left( \frac{q}{p_F} \right)^2 + \frac{\gamma |\Omega_n|}{v_F q} \right]^{-1} \quad (\text{ballistic})$$

- $\alpha, \gamma, 1/(E_F\tau)$  are phenomenological parameters.
- theory is controlled for  $\gamma \gg \alpha$ .

## Ballistic-Diffusive Cross-over



- distance traveled by electrons during interaction  $\Delta r \approx 1/q$ .
- distance traveled between two impurity scattering  $v_F\tau$ .
- ballistic:  $\Delta r \ll v_F\tau$   
single impurity scattering is important.
- in FL,  $v_F q \sim T \Rightarrow$   
cross-over temperature is  $1/\tau$ .
- near QCP,  $v_F q \sim (\gamma T)^{1/3} E_F^{2/3} \Rightarrow$   
cross-over is at  $T^* = 1/(\gamma \tau^3 E_F^2)$ .

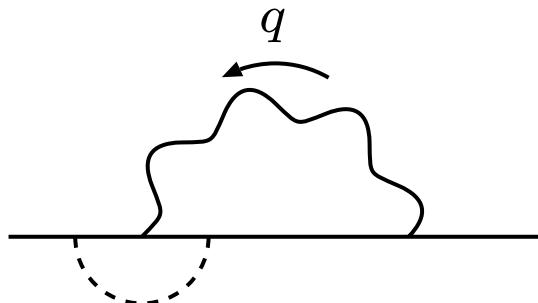
# T-dependent Lifetime for Elastic Scattering

FL

QCP



$$\frac{1}{\tau} \propto T^2$$



$$\frac{1}{\tau} \propto T^{2/3}$$

$$\frac{1}{\tau_{\text{el}}} \propto \left( \frac{1}{v_F q \tau} \right) T^2$$

$$\frac{1}{\tau_{\text{el}}} \propto \left( \frac{1}{v_F q \tau} \right) T^{2/3}$$

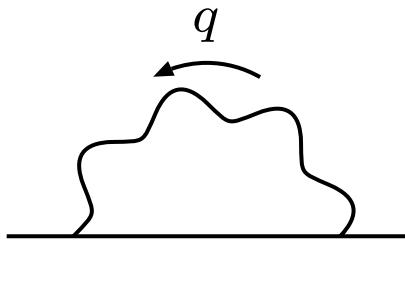
$$q \propto T$$

$$q \propto T^{1/3}$$

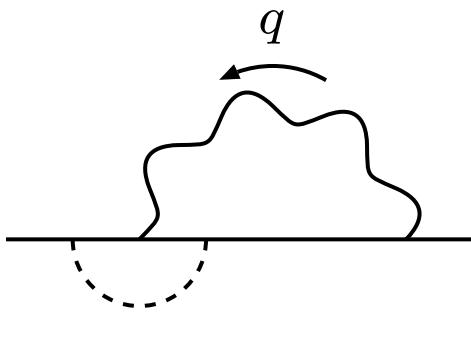
$$\frac{1}{\tau_{\text{el}}} \propto T$$

$$\frac{1}{\tau_{\text{el}}} \propto T^{1/3}$$

## Comparing Elastic and Inelastic Lifetimes



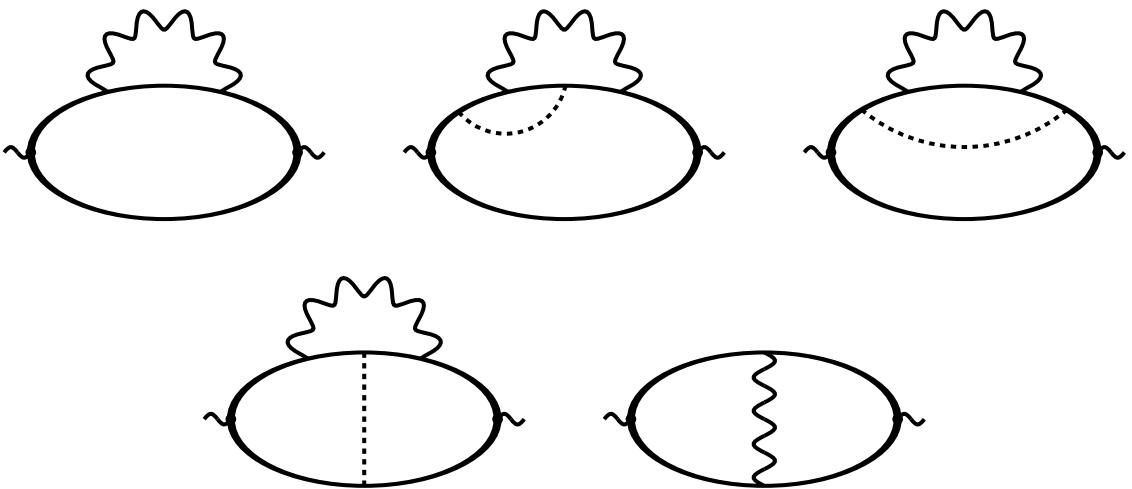
$$\frac{1}{\tau_{\text{in}}} = \underbrace{(1 - \cos \theta)}_{(q/p_F)^2} \underbrace{\text{Im} \Sigma_0}_{T^{2/3}} \underbrace{}_{T^{4/3}}$$



$$\frac{1}{\tau_{\text{el}}} = \underbrace{\left( \frac{1}{v_F q \tau} \right)}_{T^{1/3}} \underbrace{\text{Im} \Sigma_0}_{T^{2/3}}$$

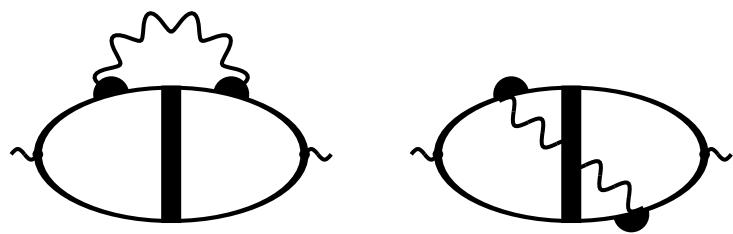
- Elastic scattering is important below the temperature  $1/(\gamma\tau)$ .
- In  $\text{Sr}_3\text{Ru}_2\text{O}_7$  this temperature is very low.
- Increasing disorder will raise this temperature.

## Ballistic Regime



$$\begin{aligned}\delta\sigma &= -(3\pi e^2 v_F^2 \tau \alpha) \int_{-\infty}^{\infty} \frac{d\Omega}{4\pi^2} \left[ \frac{\partial}{\partial \Omega} \left( \Omega \coth \frac{\Omega}{2T} \right) \right] \\ &\times \text{Im} \int \frac{d^2 q}{(2\pi)^2} D^A(\mathbf{q}, \Omega) B(\mathbf{q}, \Omega).\end{aligned}$$

## Diffusive Regime



## Conclusion

