

## Emergent Supersymmetry

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susy  $\rightarrow$  Constraints  $\rightarrow$  Dynamics

## Critical Theories

 $d = 2, 3, 4, 5, 6$ 

Local order Parameter

Exotic

No local order parameter

No Gaussian fixed point

Confinement / Fractionalization

Dualities

Particle - Monopole - Dyon

Particle - Vortex

Emergent

Global Sym, Gauge Sym

Gravity

Non-Renormalization Thms

Exact Relations

state counting, ...

(critical exponents)

Why Study SUSY in CM

Model Systems

Physical Systems

Susy CM - Emergent

If susy interactions irrelevant

↳ Flows to susy in IR

Critical point - RG flows Large

Focus on Emergent Superconformal Sym

Ingredients

Local order parameter  $\phi$

Super-partner  $\psi$

2nd order Phase Transition  $\omega_\phi = \sqrt{K_\phi}$

Space-Time Susy  $\omega_\psi = \sqrt{K_\psi}$   
 $\{Q^\mu, Q^\nu\} = 2p^{\mu\nu}$

X Thermal Fluct.  $n_\phi \neq n_\psi$  Thermal S' b.c. Susy

Quantum Fluct.  $T=0$  - QCP

Most General Set of Interactions  
 (Possibly Restricted by Global Sym)

$d=2$  susy critical points: Tricritical Ising, ...

$d=3$  ?

N=1 SUSY Wilson-Fisher d=3

$\phi$   $\mathbb{R}$   $Z_2: \phi \rightarrow -\phi$   
 $\psi$   $\mathbb{R}$  2 Comp Weyl Spinor

$$Z \partial\phi\partial\phi + \frac{1}{2}m_\phi^2\phi^2 + \frac{\lambda}{3}\phi^3 + \dots$$

$$Z i\psi\cancel{\partial}\psi + \frac{1}{2}m_\psi\psi\psi + h\phi\psi\psi + \dots$$

Two Relevant op.  $m_\phi, m_\psi$   $\mu$ -critical  
 But  $m_\psi=0$  might be obtained everywhere in a phase

Mean Field - 2nd order Phase Transition

•  $\lambda=h$   $m_\phi=m_\psi$  SUSY

$$\delta\phi = \epsilon\psi$$

$$\delta\psi = i\sigma^\mu\epsilon\partial_\mu\phi + h\phi^2$$

Superspace  $(x^\mu, \theta)$   $\Phi = \phi + \theta\psi + \theta^2 F$

$$\int d^3\theta \frac{h}{3} \Phi^3$$

Regulate  $d=4-\epsilon$

Keep explicit spinors after all contractions (external)  
 in  $d=3$ ,  $p^4$  in  $d=4-\epsilon$  -preserves b=f DOF  
 (Motivated by Supergraphs - Bosonic)

One-Loop

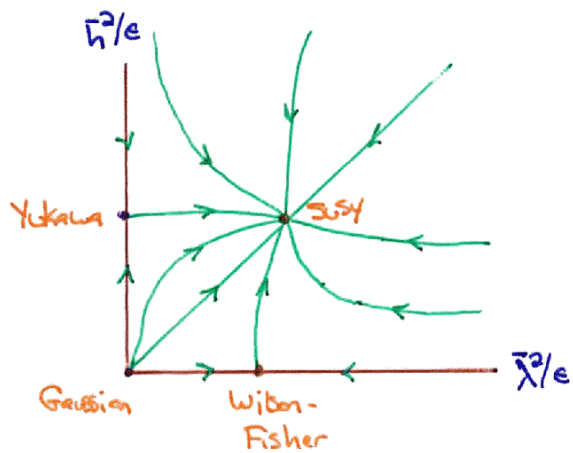
$$\beta_{\bar{h}^2} = -\epsilon\bar{h}^2 + 28\bar{h}^4$$

$$\bar{h}^2 = h^2/16\pi^2$$

$$\beta_{\bar{\lambda}^2} = -\epsilon\bar{\lambda}^2 + 36\bar{\lambda}^2 + 8\bar{h}^2\bar{\lambda}^2 - 16h^4$$

$$\bar{\lambda}^2 = \lambda^2/16\pi^2$$

On  $m_\phi=m_\psi=0$  Emergent Superconformal Sym  
 Conformal Sym  $\Rightarrow$  Accidental Superconformal Sym



Critical Exponents

$$\eta = 2\gamma_\phi$$

Order Parameter  $G(r)$   
Scaling Regime

$$\nu^{-1} = 2 - \gamma_{m^2}$$

Correlation Length

$$w = \frac{\partial \beta}{\partial t}$$

Approach to Scaling  
Most Relevant irrelevant  $\phi$   
along  $\lambda = h$

$$m_4 = 0$$

N=2 SUSY Wilson Fisher d=3

$\phi \in \mathbb{C}$   
 $\psi^i \in \mathbb{C}$  2 comp Weyl

U(1) - Allow Yukawa

$$Z \partial \phi \partial \phi + m_\phi^2 \phi^* \phi + \lambda^2 (\phi^* \phi)^2 + \dots$$

$$Z i \bar{\psi} \not{\partial} \psi + h \phi \psi \psi + h \phi^* \psi^* \psi^* + \dots$$

Only one Relevant Op.

$Z_2 \times Z_3$  or  $Z_N$  NZS or  $m_\phi = 0$  in phase

$\lambda = h$   $m_\phi = 0$  SUSY -----

Regulate  $d = 4 - \epsilon$

All Spinors in  $d=4$ ,  $p^4 = 4 - \epsilon$  DRED  
 (Supergraphs Bosonic)

One-Loop

$$\beta_{\bar{h}^2} = -\epsilon \bar{h}^2 + 12 \bar{h}^4$$

$$\beta_{\bar{\lambda}^2} = -\epsilon \bar{\lambda}^2 + 20 \bar{\lambda}^4 + 8 \bar{h}^2 \bar{\lambda}^2 - 16 \bar{h}^4$$

On  $m_\phi = 0$  Emergent Superconformal Sym  $U(1) = U(1)_R$

Conformal Sym  $\Rightarrow$  Accidental Superconformal Sym

General  $\int d^d x g_i \sigma_i$        $\sigma_i = \phi\phi \dots \psi\psi \dots$

$$\beta_{g_i} = (d - D_{\sigma_i}^0) g_i + \sum_{\phi} \gamma_{\phi} g_i + \gamma_{\sigma_i}$$

Classical  
Scaling



Wave  
Function



Vertex



Any Dim Reg Scheme  $\gamma_{\phi}, \gamma_{\sigma}$  independent of  $\epsilon$

$\beta=0$        $\epsilon \rightarrow f(g_i^*)$       non-linear; invert

$$g^* = \sum_n g_n^* \epsilon^n$$

Then

$$\gamma(g^*) = \sum_n \gamma_n \epsilon^n$$

Only explicit dep is  $\epsilon!$

N=2 SUSY d=3

Non-Renormalization Thm

Superpotential Receives no Vertex Corrections  
(Exact - All orders and non-perturbatively)

↳ Holomorphy

↳ Hard to see in Components

$$Z F^* F + h F \phi \phi + h F^* \phi^* \phi^* = -\frac{h^2}{Z} (\phi^* \phi)^2$$

Exact →

$$\beta_{\bar{h}^2} = -\epsilon \bar{h}^2 + 6\gamma_{\phi} \bar{h}^2$$

↳ No Vertex Corrections!!

Homogenous in  $\bar{h}^2$

$$\gamma_{\phi}^* = \frac{\epsilon}{6} \quad \eta = \frac{\epsilon}{3}$$

$\epsilon$ -expansion for  $\eta$  exact at  $\sigma(\epsilon)$

Superconformal Algebra

$J_0^*, J_R^*$       Same Supermultiplet

$$D_{\phi}^* = \frac{d-1}{2} R_{\phi} \quad R_{\psi\psi\phi} = -2 \quad R_{\phi} = 2/3$$

$$D_{\phi}^* = \frac{1}{3}(d-1) = 1 - \frac{\epsilon}{3} \quad D_{\phi}^c = 1 - \frac{\epsilon}{6}$$

$$\therefore \gamma_{\phi}^* = \frac{\epsilon}{6} \quad \eta = \frac{\epsilon}{3} \quad \text{by superconformal Sym (emergent)}$$

$$v^{-1} = 2 - \gamma_m^2 \quad : \quad m^2 \phi^* \phi \quad \text{susy}$$

- Extend  $\beta$ -fun in superspace  
Exact Relation  $v, w$

$$h^2 = \frac{h_0^+ h_0}{z^2} \rightarrow H^2 \text{ Background superfield}$$

$$\text{with } z = Z(1 - m^2 \theta^2 \bar{\theta}^2), \quad \int d^4 \theta z \Phi^\dagger \Phi \subset m^2 \phi^* \phi$$

$$\beta_{H^2} = -\epsilon H^2 + 6\gamma_\phi H^2 \quad \gamma_\phi = \gamma_\phi(H^2)$$

$$\gamma_{H^2} = 6\gamma_\phi$$

$$\gamma_{H^2} \Big|_{\theta^2 \bar{\theta}^2} = 3m\gamma_m^2$$

$\therefore$

$$\gamma_m^2 = 6\bar{h}^2 \frac{\partial \gamma}{\partial \bar{h}^2}$$

Exact on  $\lambda = h$   
Away from critical pt.

But

$$w = \frac{\partial \beta^*}{\partial \bar{h}^2} = 6\bar{h}^2 \frac{\partial \gamma^*}{\partial \bar{h}^2}$$

$$\therefore w = \gamma_m^2 \quad w = 2 - v^{-1} \quad \text{Exact!}$$

Susy

$N=2$

$\frac{w}{v}$  (Exact)

$\frac{1}{2} + \frac{\epsilon}{4} + O(\epsilon^2)$

$2 - v^{-1}$  (Exact)

Wilson Fisher  
N=0

SWY  
N=1

n  $0 + \mathcal{O}(\epsilon^2)$

$\frac{\epsilon}{7} + \mathcal{O}(\epsilon^2)$

v  $\frac{1}{2} + \frac{\epsilon}{12} + \mathcal{O}(\epsilon^2)$

$\frac{1}{2} + \frac{\epsilon}{7} + \mathcal{O}(\epsilon^2)$

w  $\epsilon + \mathcal{O}(\epsilon^2)$

$\epsilon + \mathcal{O}(\epsilon^2)$

Mean Field: 2nd order Phase Transition

Q-Fluctuations of other Fields  
coupled to order Parameter might ruin  $\rightarrow$  1st order

$\rightarrow$  continuing  $\beta_{\phi}$  into 2 space  
Inherits Homogeneity in  $m^2$   
 $\uparrow$   
Can't develop Tachyon

i)  $\beta_{m^2} = -2m^2 + 2\gamma_{\phi} m^2$

ii)  $\bar{\lambda}^* > 0$

stable

$\therefore$  Phase Transition Remains 2nd order  
Quantum Super Critical Point Stable



• Generalizations

$U(1)_R$  c.f.  $SO(2)$  Sym

Lorentz viol.  $V_\psi \neq V_\phi$

$\mu$ -Component

More SUSY (Requires Gauge Fields)



Entire Region  
Flows to SUSY in IR

Simple Models - Retain 2nd order  
General Model - Fluctuation induced 1st order  
by tachyon -  
QSCP unstable  
↳ Effects  $t = \text{finite}$

Summary

Emergent Superconformal Sym  
Possible in Simple  $d=3$  Systems

Single Relevant Direction Possible

Stable

Provide new Examples Quantum Super Critical Points

With  $N=2$   $d=3$  Exact Relations Among  $\nu, \nu, \omega$

Lessons

Exact Results - SUSY sufficient

Non-SUSY Exact Results