

# Black Holes and Quantum Field Theory

Warning! Slides have omitted citations to background papers. Apologies to all concerned. See arXiv:1905.00539 for background refs.

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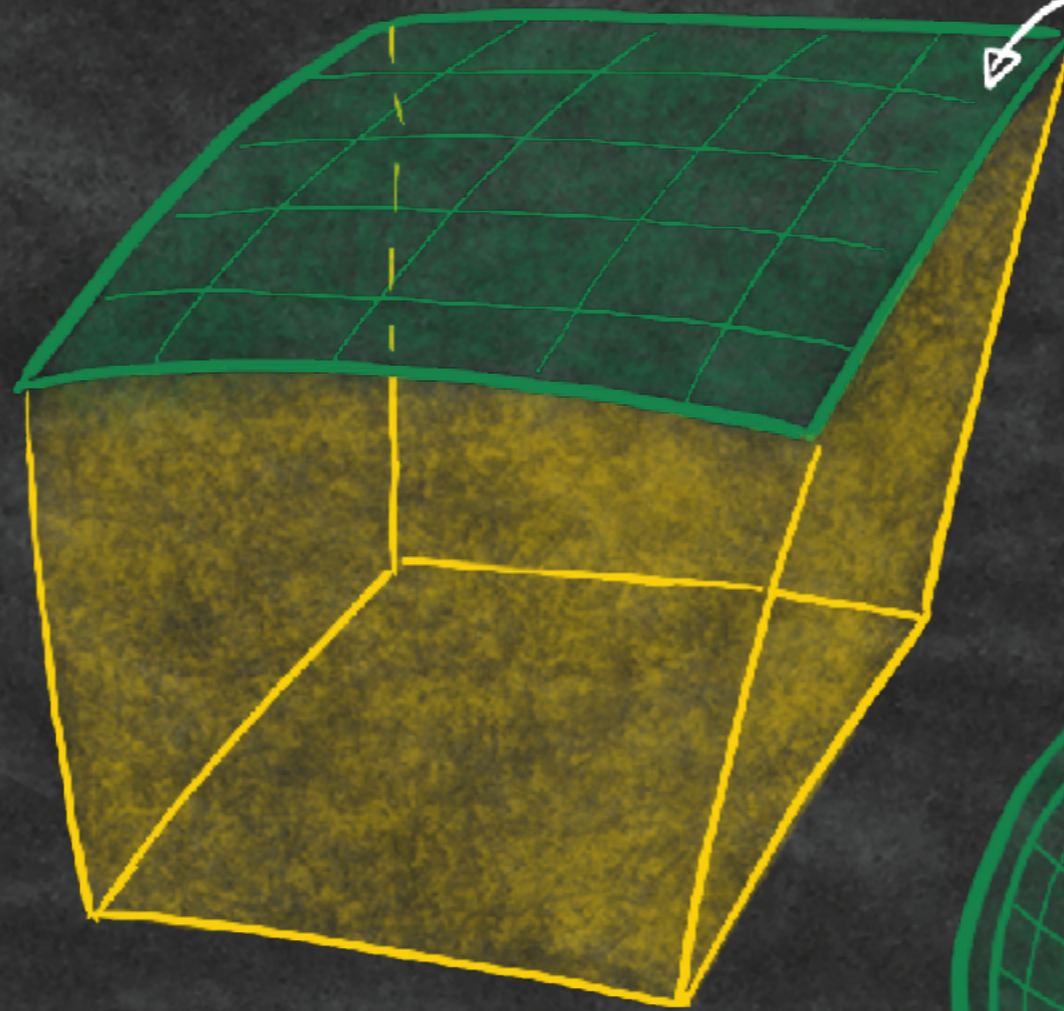
KITP  
2nd May 2019



# Geometrizing Quantum Problems

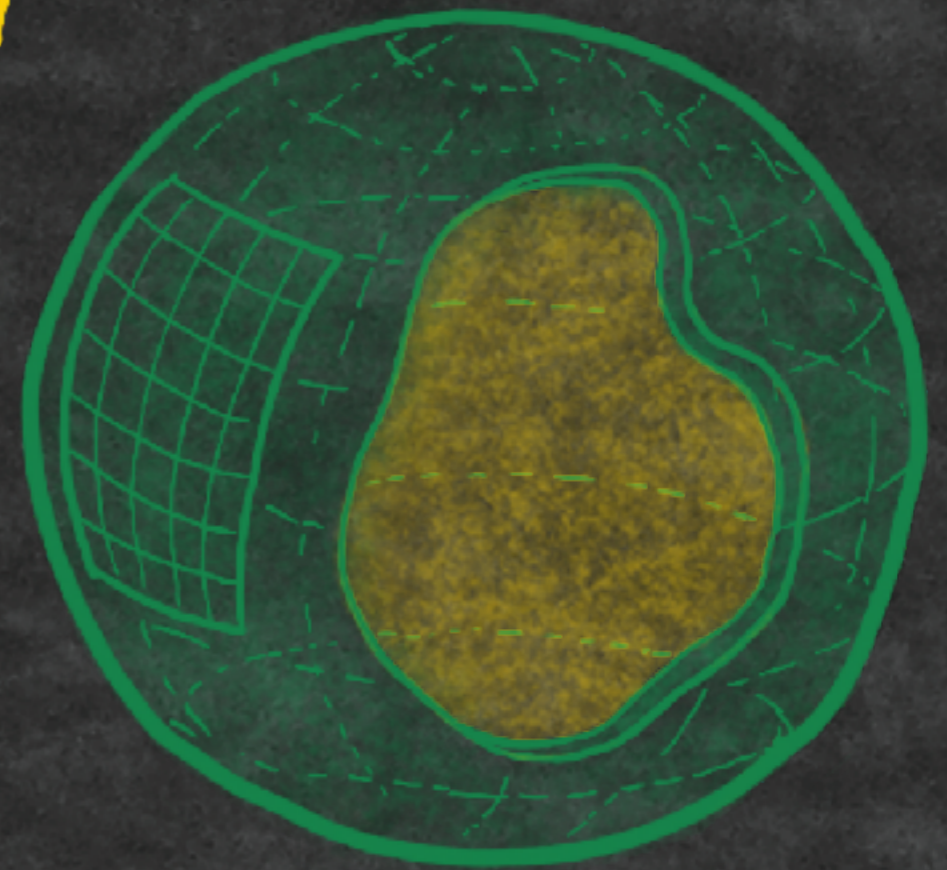
AdS/CFT  
 $\epsilon$   
generalizations

extra  
dimension

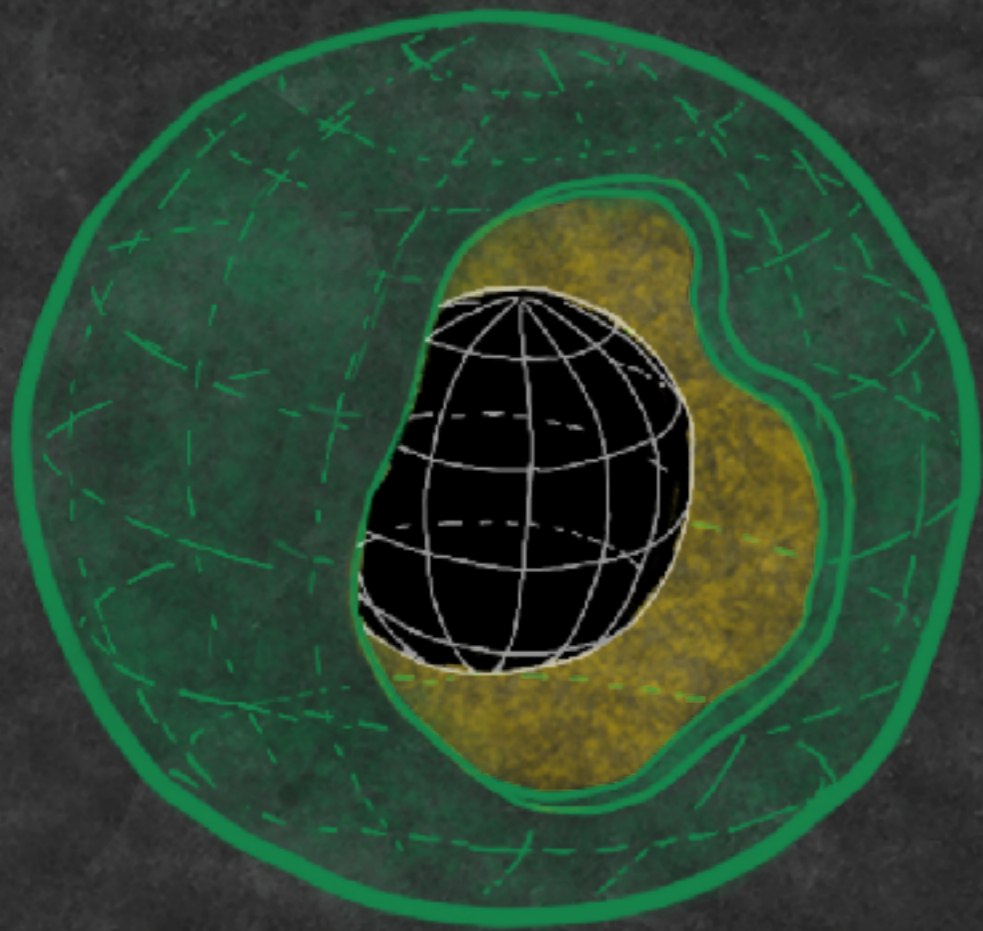


d-dimensional  
QFT  
eg  $SU(N)$  gauge  
theory  
 $g^2, N$

$d+1$  dimensional gravity  
with c.c.  $\Lambda \sim -1/\ell^2$



# Thermal Physics



black hole, radius  $r_h$

$$M = \frac{r_h^3}{2l^2} \rightarrow U$$

$$K = \frac{3r_h}{4\pi l^2} \rightarrow T$$

$$\frac{A}{4} = \pi r_h^2 \rightarrow S$$

$$u \sim N^{3/2} T^3$$

density  $\nearrow$

$$du = TdS$$

first law

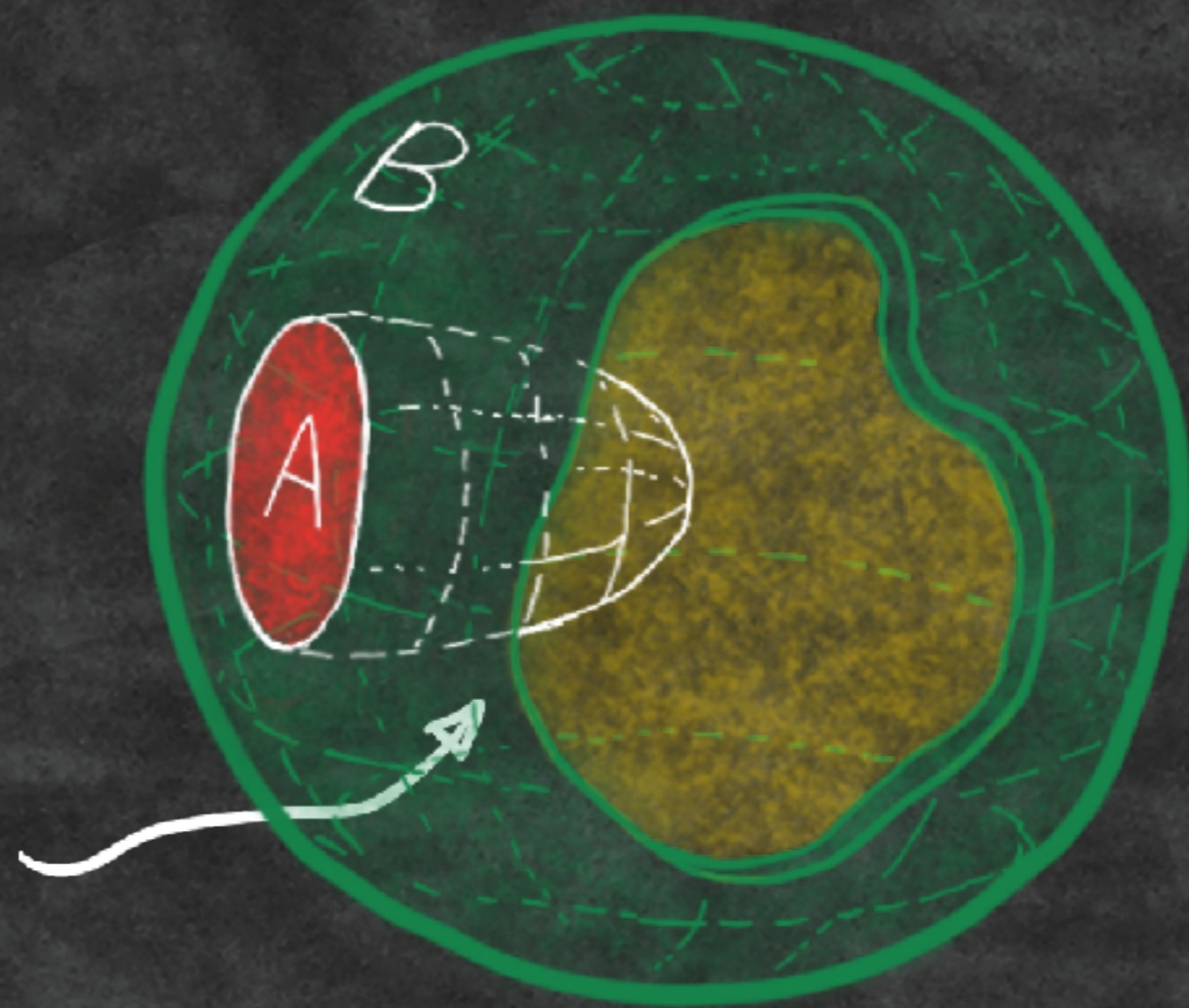
# Quantum Entanglement

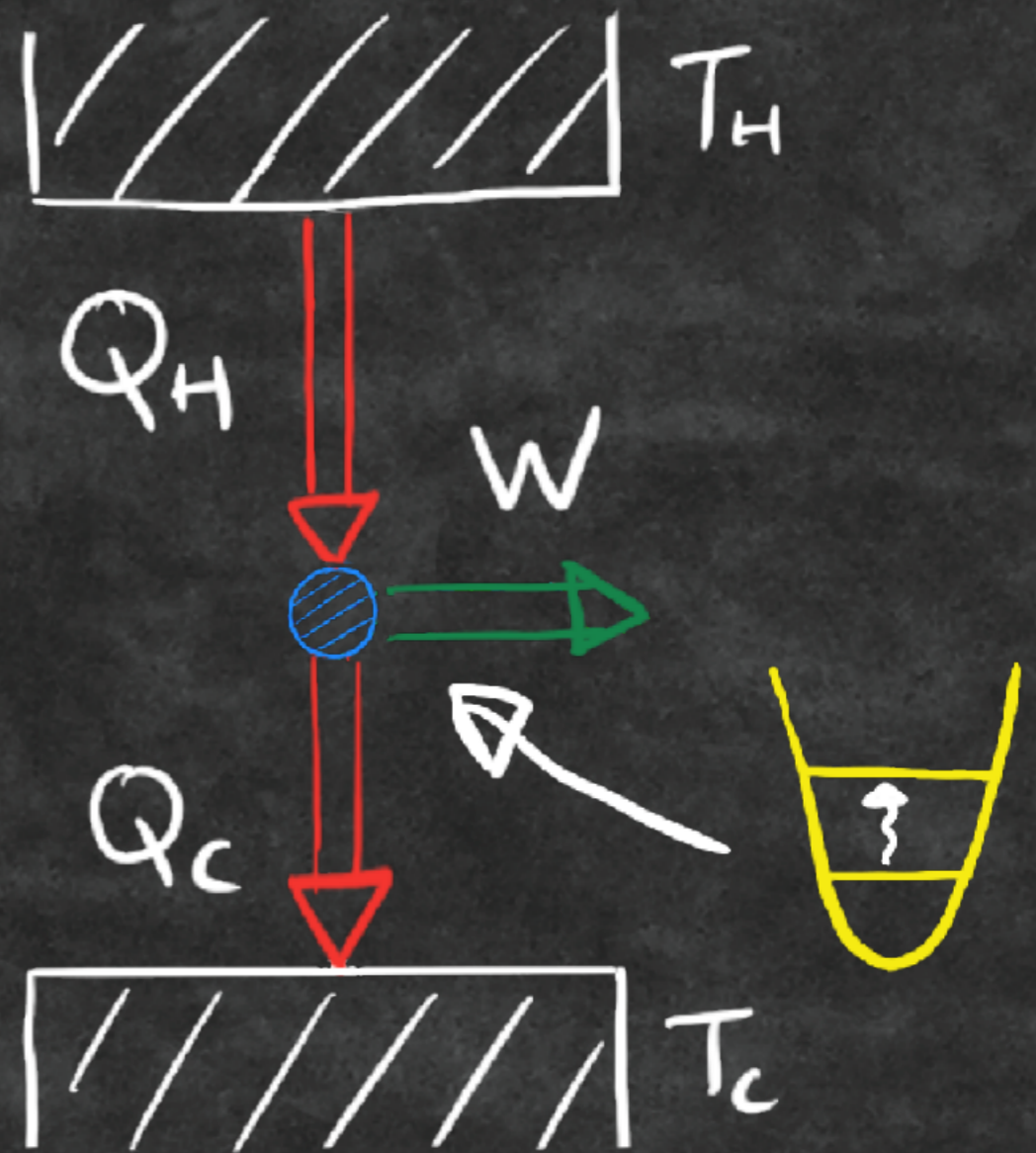
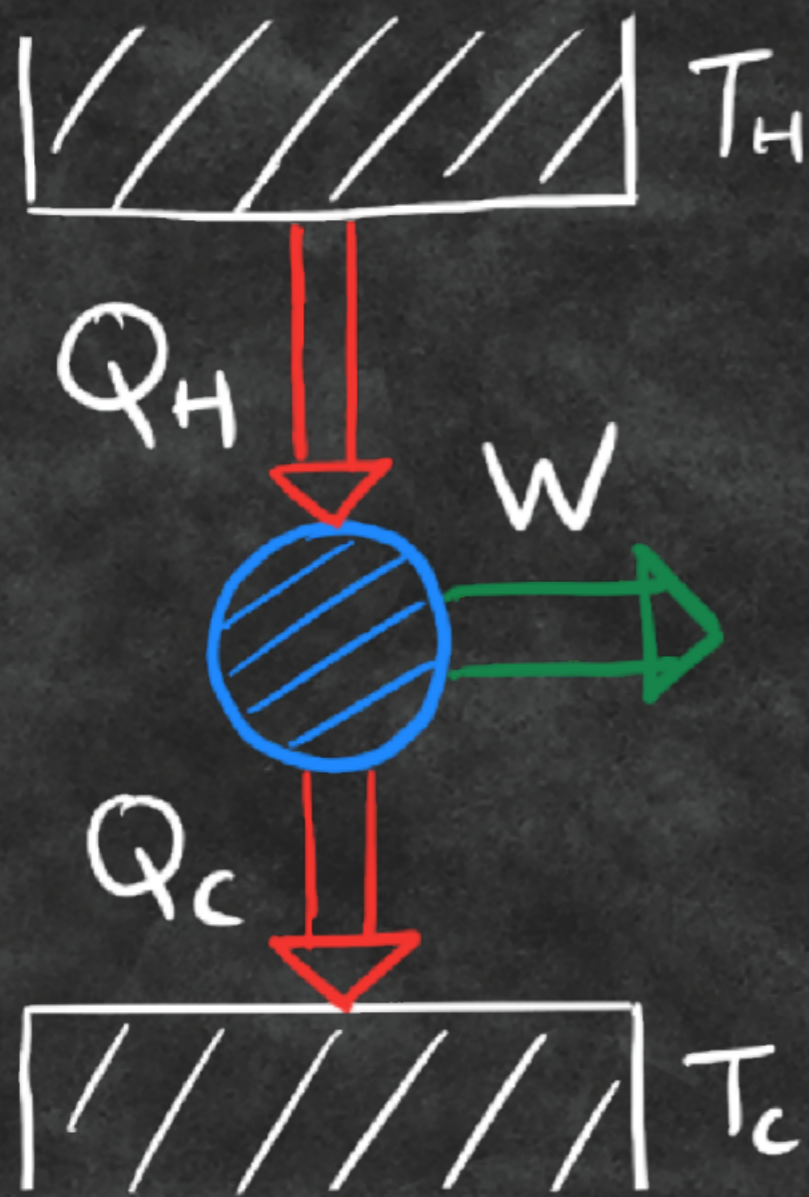
$$\rho_A = \text{Tr}_B \rho$$

$$\rho = |0\rangle\langle 0|$$

$$S_{EE} = -\text{Tr} \rho_A \ln \rho_A$$
$$= \frac{A_{\text{area}}}{4}$$

minimal  
surface  
area A





Today: Quantum Heat Engines

key: Working substance has small number of d.o.f.

Can gravity + black holes help?



- No  $p$  or  $V$  for mechanical work

- WAY too big a system:

$$\text{d.o.f} \sim N^\#$$

$N$  large for good geometry  
description

Solution: We're in the wrong thermo...

# B.H. Thermodynamics ("extended")

mass  $M$   $\Rightarrow$   $H \equiv U + pV = M$  enthalpy

area  $A$   $\Rightarrow$   $S = \frac{A}{4G}$  entropy

surface gravity  $\kappa$   $\Rightarrow$   $T = \frac{\kappa}{2\pi}$  temp.

Cosmo. Const.  $\Lambda$   $\Rightarrow$   $p = -\frac{\Lambda}{8\pi G}$  pressure

1st Law  $H(p, S)$

$$dH = TdS + Vdp$$

$$T = \left. \frac{\partial H}{\partial S} \right|_p \quad V = \left. \frac{\partial H}{\partial p} \right|_S$$

# Example: Schwarzschild in AdS

$$\Lambda = -\frac{3}{\ell^2} \quad \therefore p = \frac{3}{8\pi G\ell^2}$$

$$ds^2 = -F(r)dt^2 + F(r)^{-1}dr^2 + r^2(d\theta^2 + \sin^2\theta d\phi^2)$$

$$F(r) = \left(1 - \frac{2M}{r} + \frac{r^2}{\ell^2}\right) \quad \begin{array}{l} \text{horizon at largest} \\ \text{root of } F(r) = 0 \end{array}$$

$$\text{i.e. } r_h^3 + r_h\ell^2 - 2M\ell^2 = 0$$

$$S = \frac{A}{4} = \pi r_h^2$$

$$\text{solve: } M = \frac{r_h}{2} \left(1 + \frac{r_h^2}{\ell^2}\right)$$

$$\text{Also, } \beta F' = 4\pi \text{ gives } T = \frac{1}{4\pi} \left( \frac{2M}{r_h^2} + \frac{2r_h}{\ell^2} \right)$$



Now write  $M = H(S, p) = \frac{1}{2} \left( \frac{S}{\pi} \right)^{1/2} \left( 1 + \frac{8Sp}{3} \right)$

check:  $T = \frac{\partial H}{\partial S} \Big|_p$

$$= \frac{1}{2} \cdot \frac{1}{2} \left( \frac{1}{\pi} \right)^{1/2} S^{-1/2} \left( 1 + \frac{8Sp}{3} \right)$$

$$+ \frac{1}{2} \left( \frac{1}{\pi} \right)^{1/2} S^{1/2} \left( 0 + \frac{8p}{3} \right)$$

$$= \frac{1}{4} \frac{1}{\pi^{1/2}} \frac{1}{S^{1/2}} \left( 1 + \frac{8Sp}{3} \right)$$

$$= \frac{1}{4\pi r_h} \left( 1 + 3r_h^2/\ell^2 \right) = \frac{(2M + 2r_h^3/\ell^2)}{4\pi r_h^2}$$



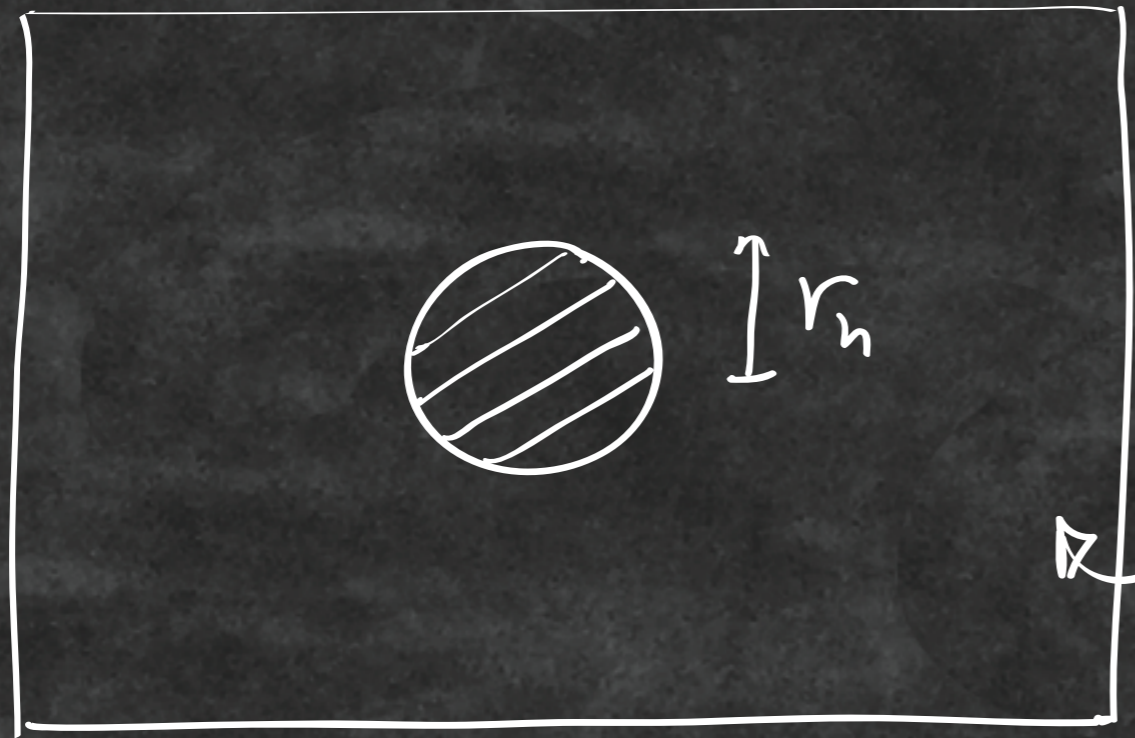
Now for something new...

$$H = \frac{1}{2} \left( \frac{S}{\pi} \right)^{1/2} \left( 1 + \frac{8Sp}{3} \right)$$

$$V = \left. \frac{\partial H}{\partial p} \right|_S = \frac{1}{2} S^{1/2} \frac{1}{\pi^{1/2}} \frac{8S}{3}$$

$$= \frac{4}{3} \frac{S^{3/2}}{\pi^{1/2}} = \frac{4}{3} \pi r_h^3 \quad !!$$

The 'naive' volume!



energy density

$$\rho = \frac{\Lambda}{8\pi G}$$

so cost is  $\rho V$

What about the degrees of freedom?

Let's check: use specific heat

$$C(T) = T \frac{\partial S}{\partial T}$$

$$dU = TdS - pdV$$

- $C_V(T)$  tells us about available d.o.f.
- $C_P(T)$  also includes energy used for work.

For black holes, usual  $C \rightarrow C_P$  in extended.  
So what is  $C_V$ ? (largely overlooked)

In fact, for our favourite black holes:

$$C_V = 0 \quad \checkmark$$

follows from  $V$  and  $S$  both being geometrical

$$V = \frac{4\pi}{3} r_h^3 \quad S = \pi r_h^2$$

So, **No degrees of freedom!**

(Well, no "traditional" ones.)

So, what to do? Use black holes w  $C_V \neq 0$

eg. Kerr



$$V = \frac{2}{3\pi H} \left( S \left( S + 8\rho S^2 \right) + 2\pi^2 J^2 \right) \quad (\text{Dolan})$$

$$H = \frac{1}{2} \sqrt{\frac{\left( S + 8\rho S^2 \right)^2 + 4\pi \left( 1 + \frac{8\rho S}{S} \right) J^2}{\pi S}}$$

So, what to do?

also, STU



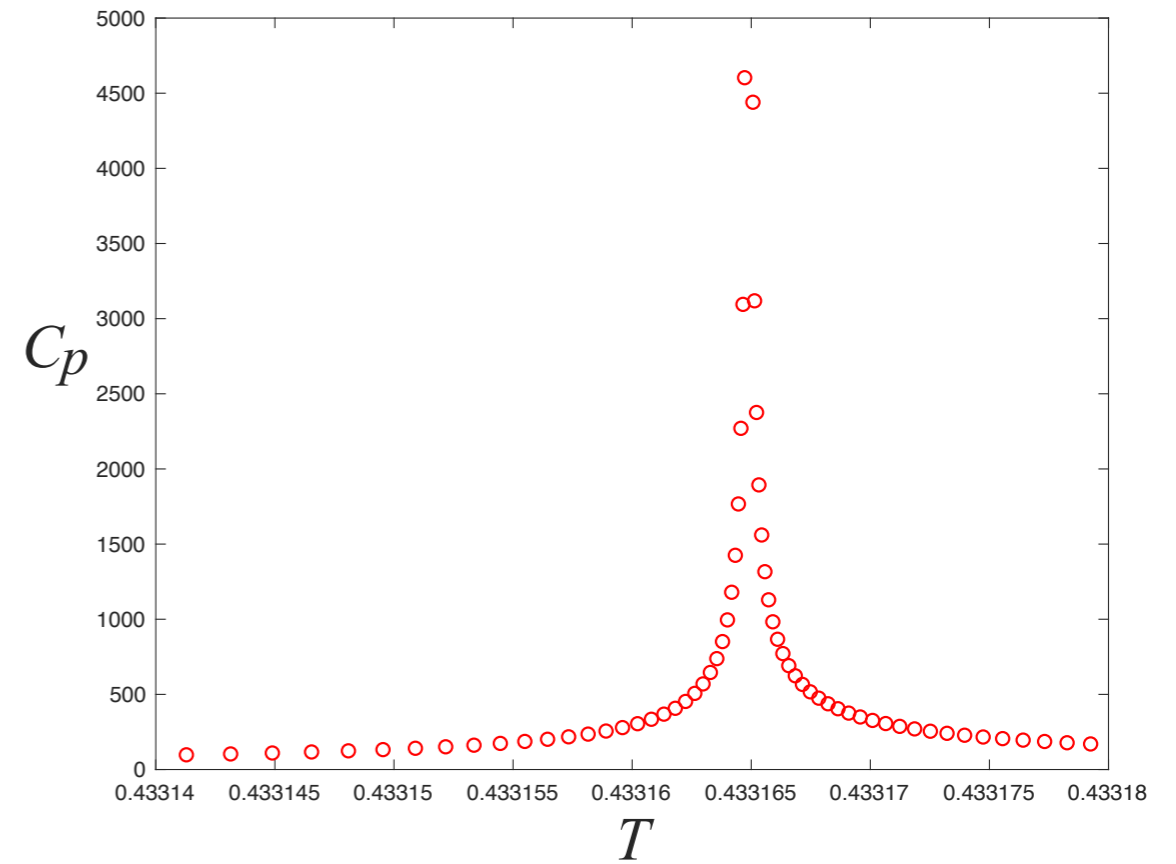
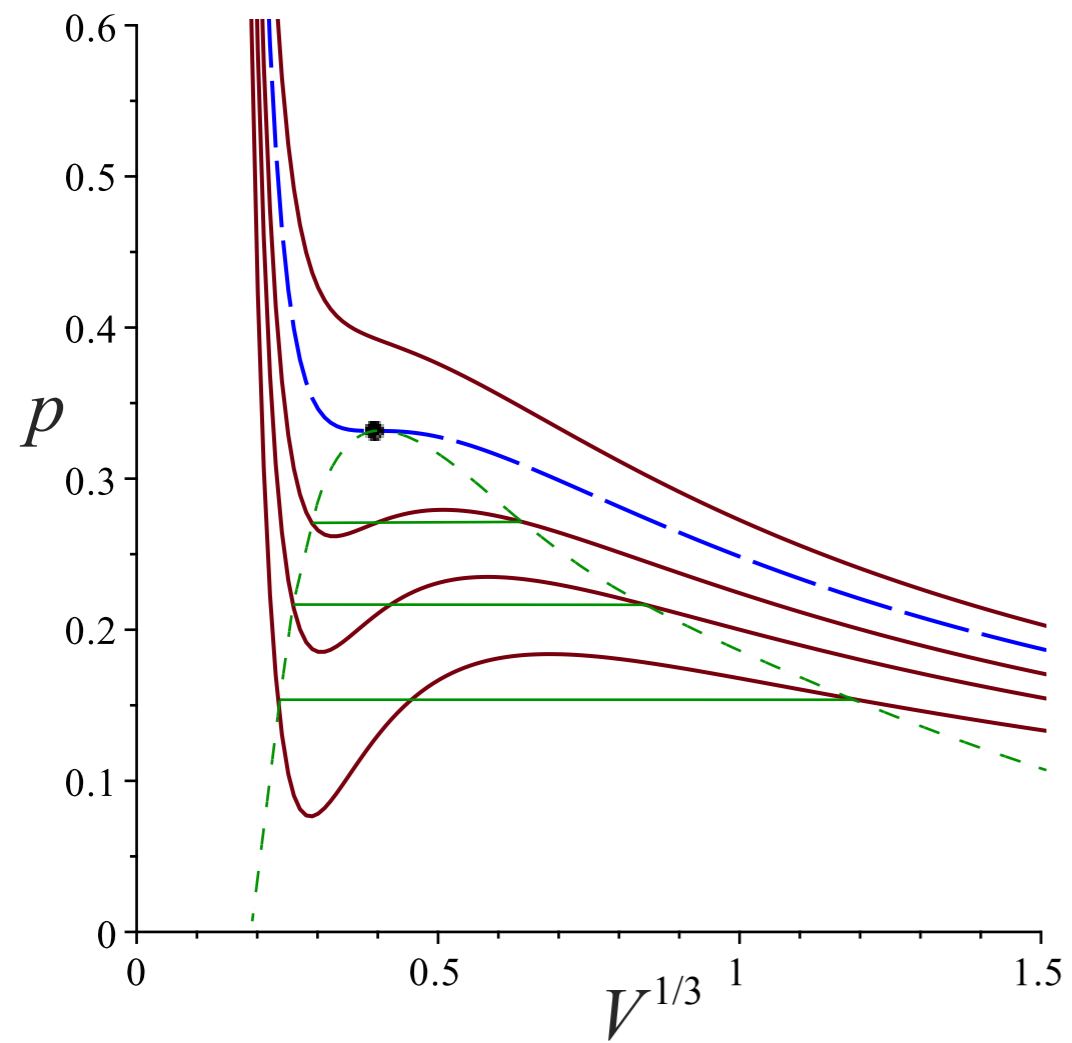
b.h. Charged non-diag.

under

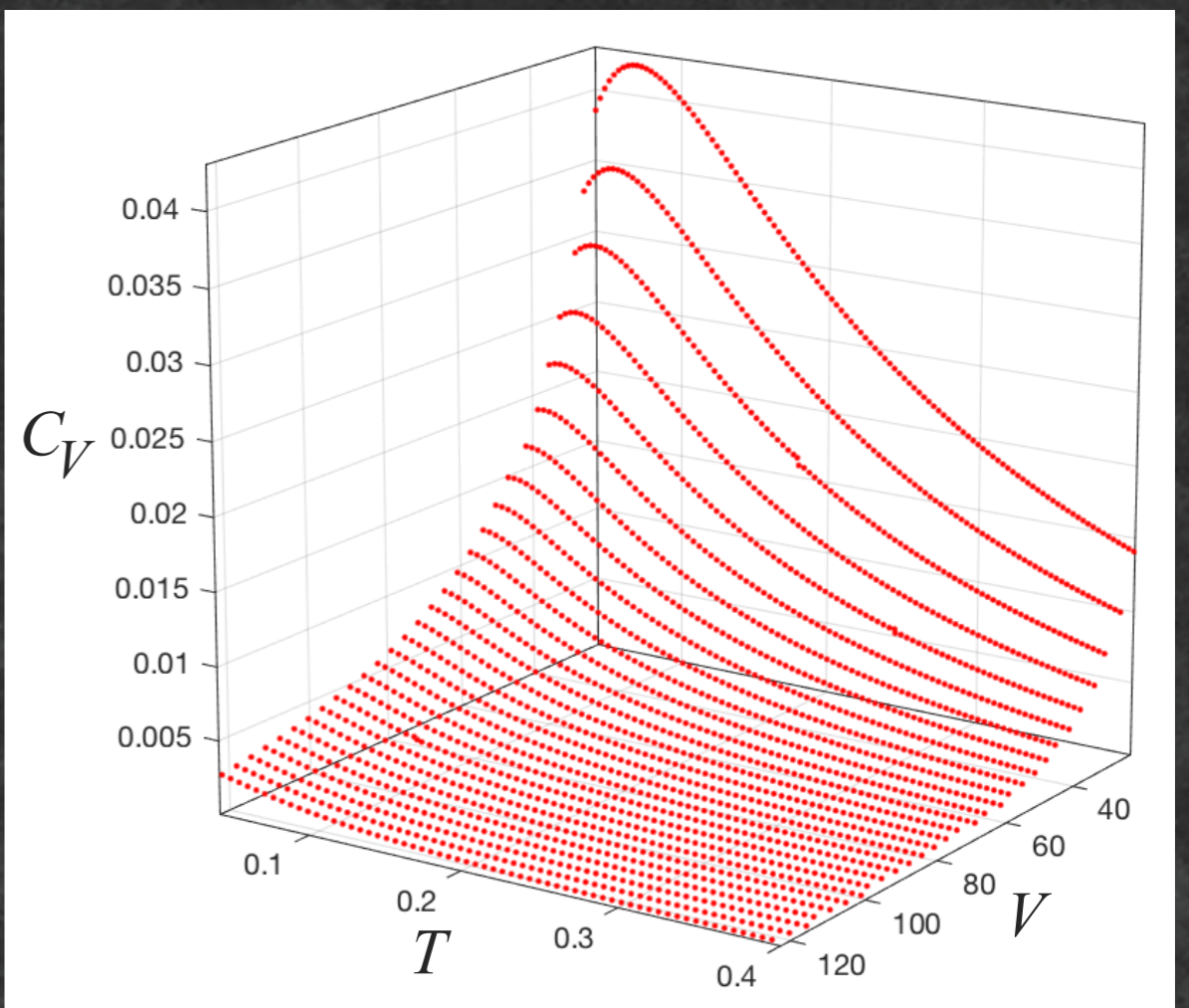
$$U(1) \times U(1) \times U(1) \times U(1)$$

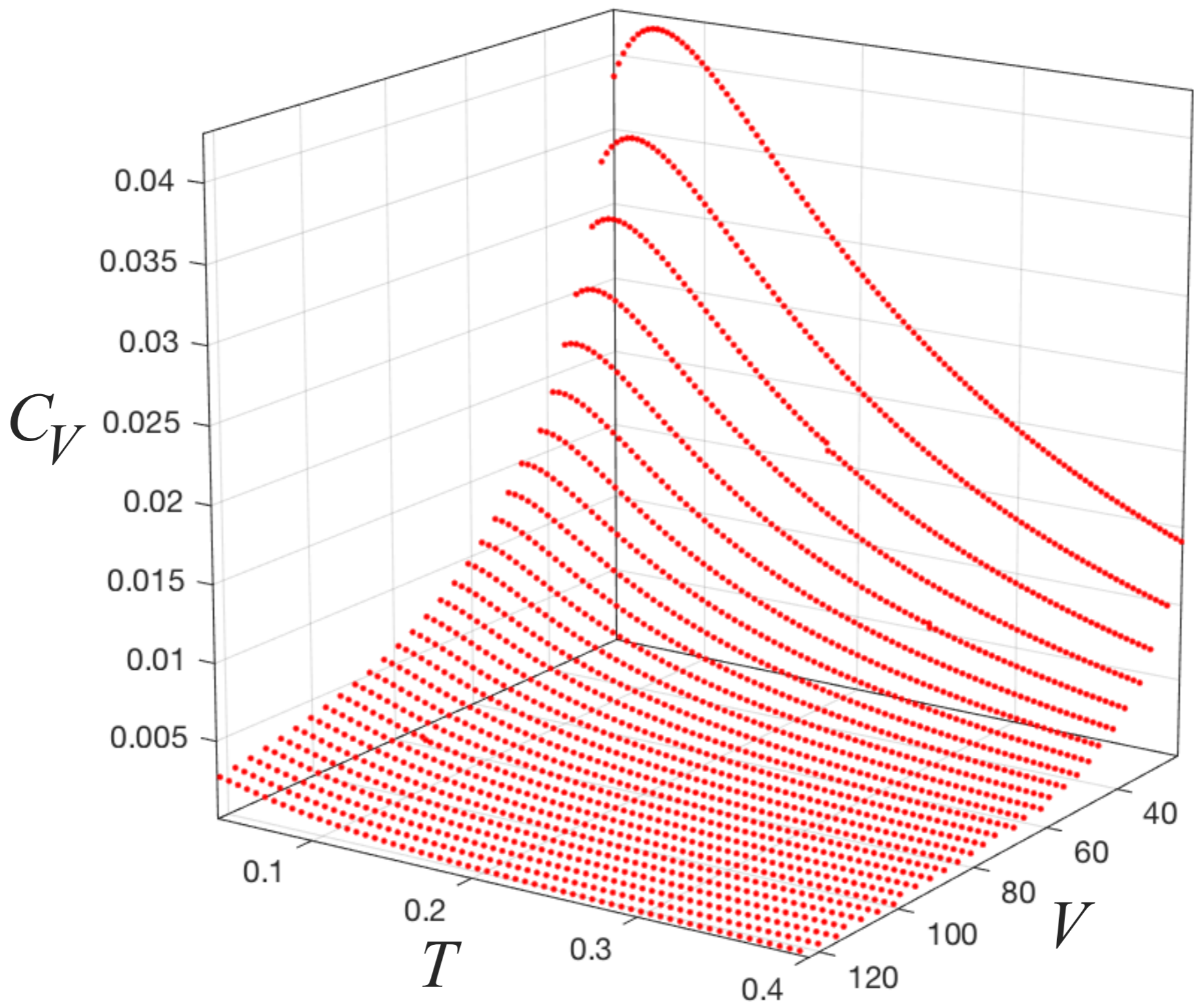
$Q_1 \quad Q_2 \quad Q_2 \quad Q_3$

(RN is  
all  $Q$ 's equal)



New exciting results!  
See: [arXiv:1905.00539](https://arxiv.org/abs/1905.00539)







## Schottky-like Peaks!

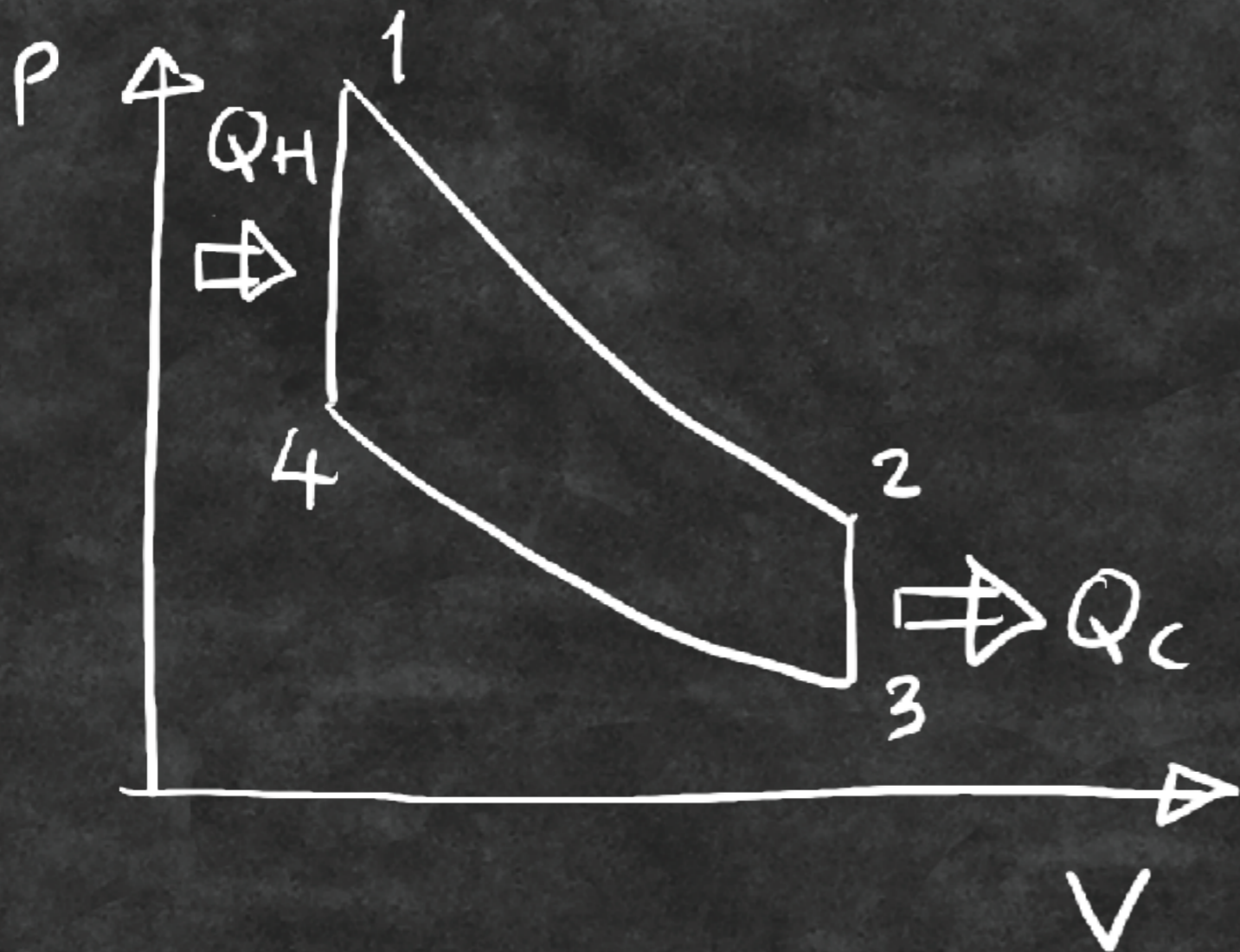
- A finite window of energies.
- Height, position set by  $Q \approx J$ .



This gives us the "small"  
system we need!

Holographic heat engines  
as

Quantum heat engines



Otto cycle

In each case  
a cycle is  
a tour through  
a family  
of models

(SHO,  $\approx$  QFT)

## Summary + Discussion

- Are these solvable models useful?
- What issues in experiment + theory can they help with?  
(Power vs. efficiency?  
Running near criticality?)
- Can we model friction/noise?
- Models of heat baths?
- Do QHE do other cycles?

Thank you!

