

# Quantum absorption refrigerator with trapped ions

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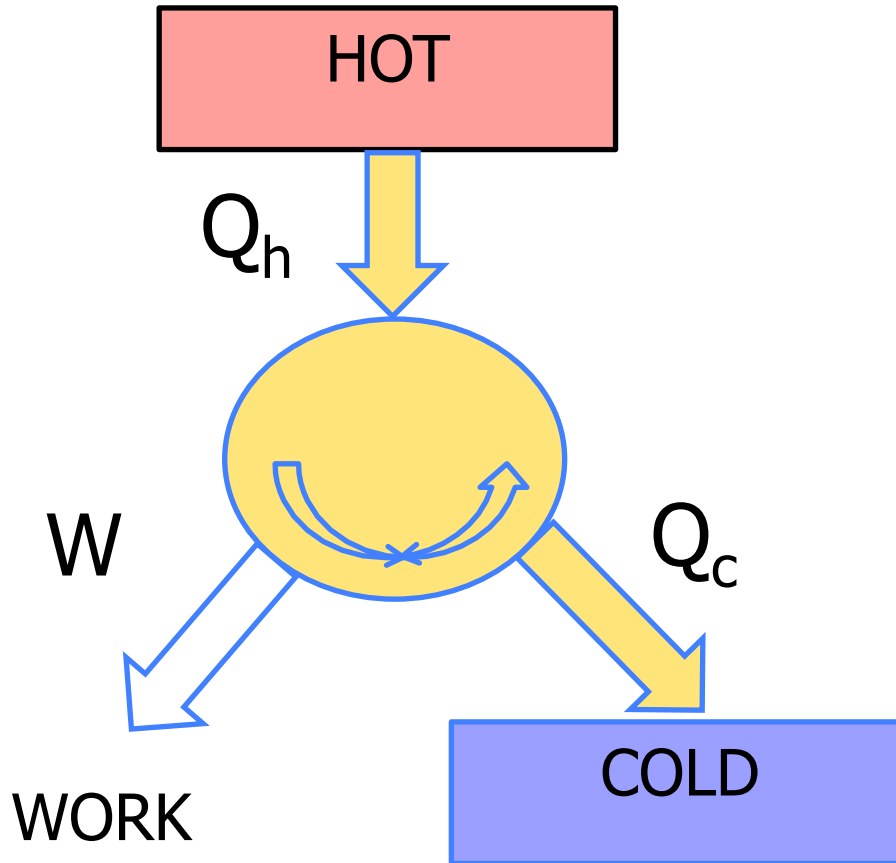
2018, KITP, UC Santa Barbara



# Types of heat machines: Engine

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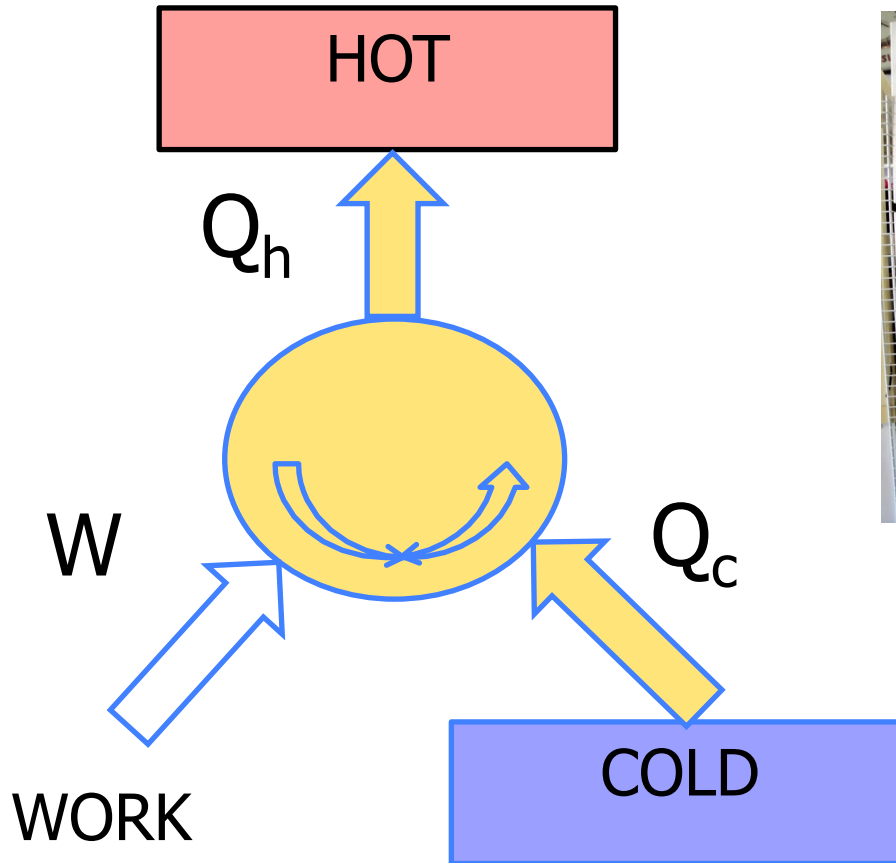
Engine: Uses heat to produce work



# Types of heat machines: Refrigerator

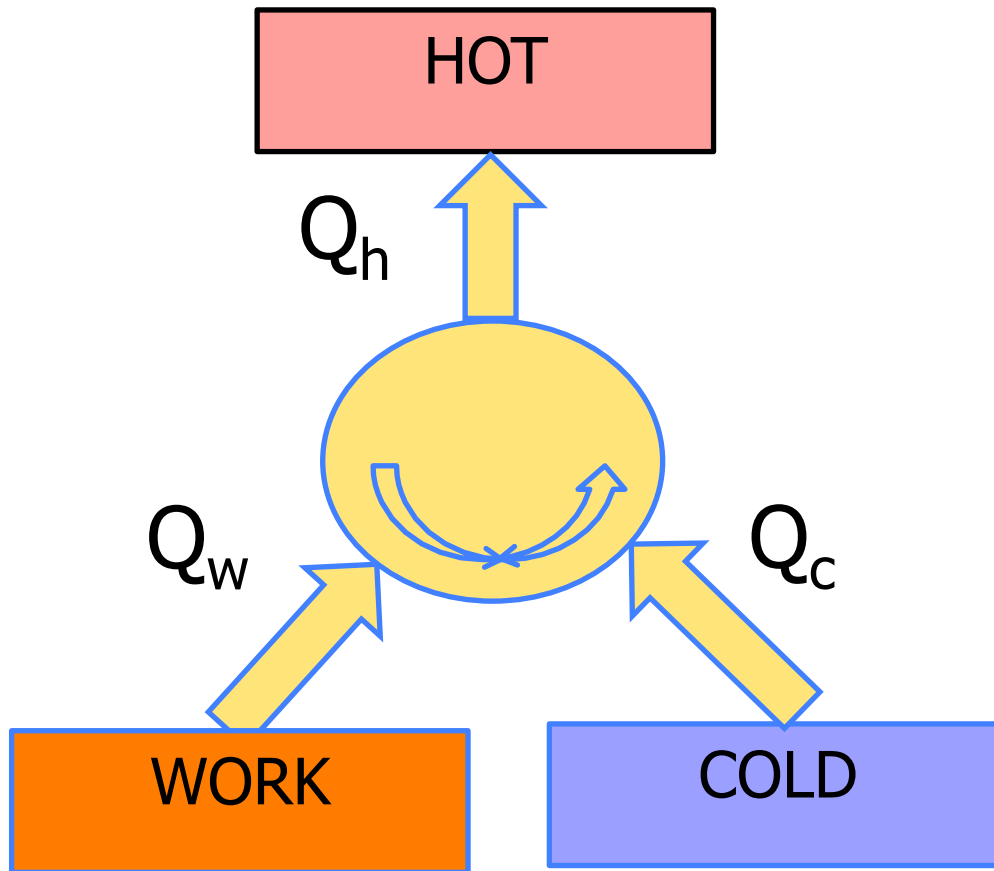
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Refrigerator: Uses work to refrigerate cold body



# Absorption Refrigerator

Absorption Refrigerator: Driven by heat instead of work



$$T_w > T_h > T_c$$



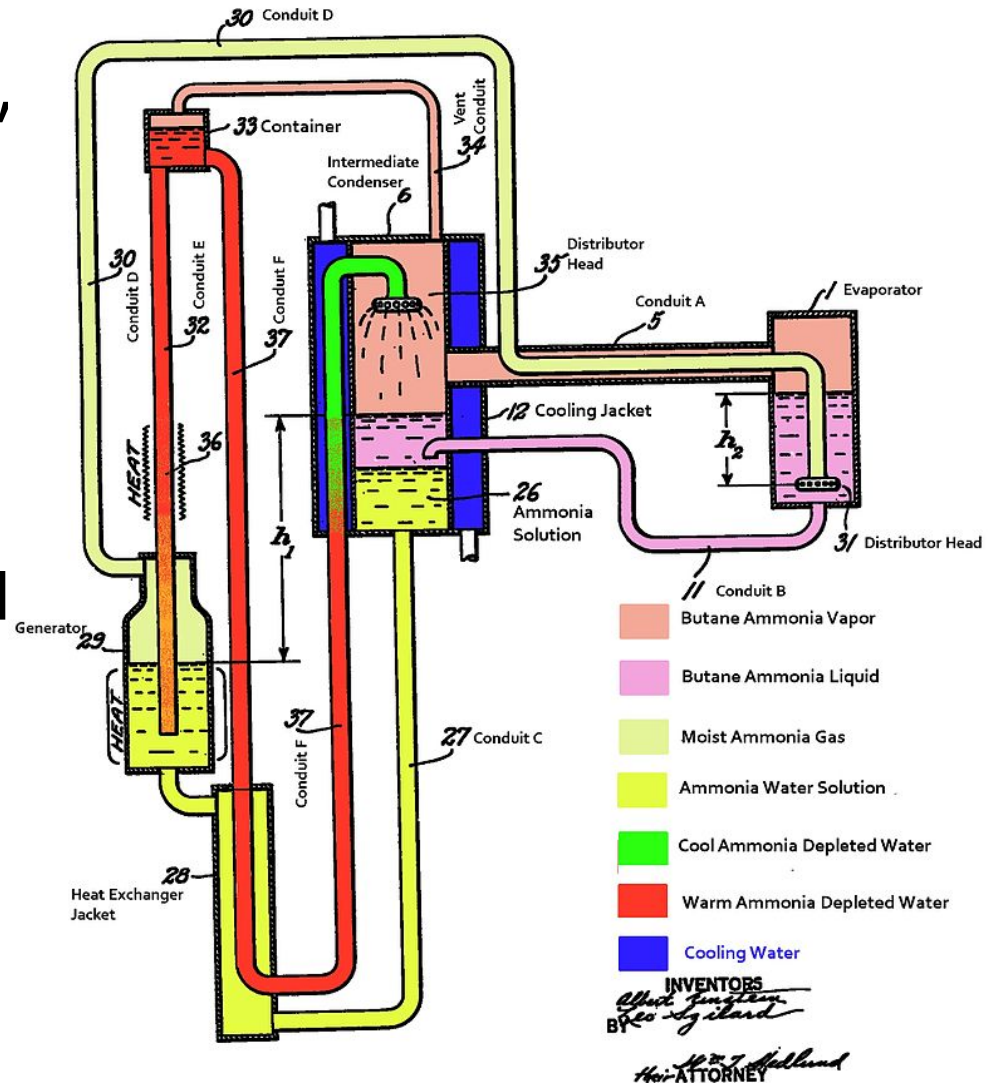
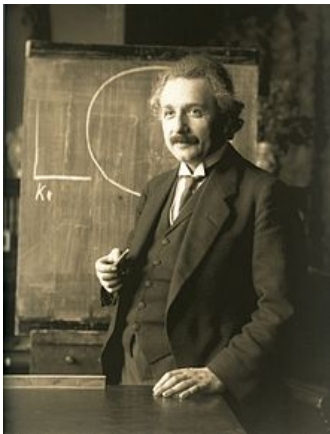


# Absorption Refrigerator

Invented by:  
Ferdinand Carre,  
US Patent  
30,201 (1860)



Improved design:  
Albert Einstein and Leó Szilárd  
US Patent 1,781,541 (1930)



# Quantum absorption refrigerator

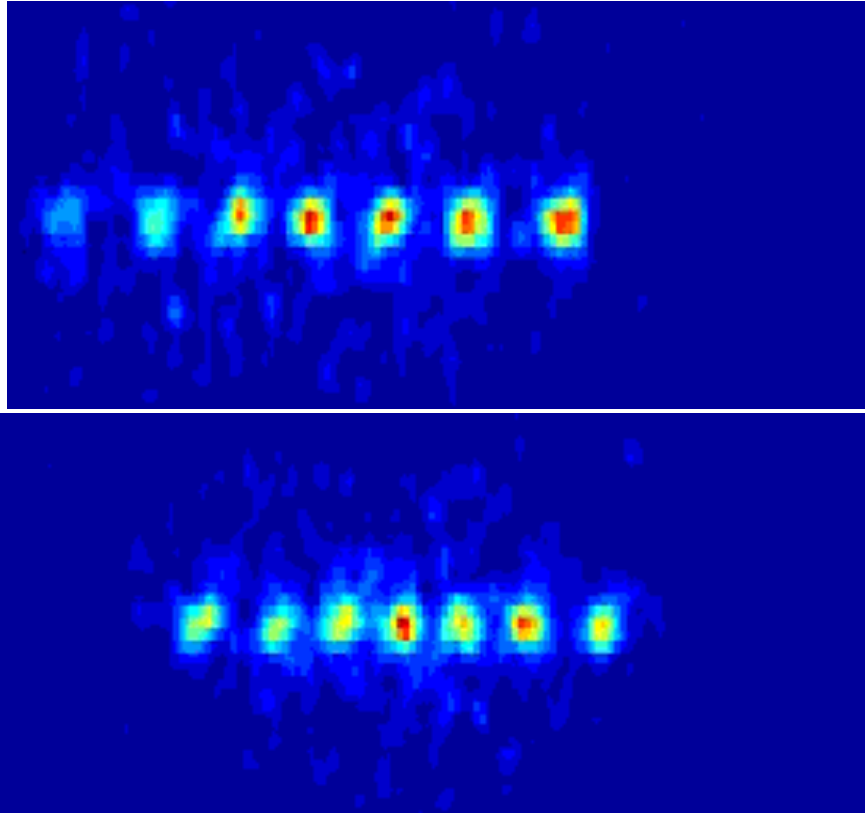
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What is the smallest refrigerator one can build?

How does the quantum effects influence refrigerator performance?

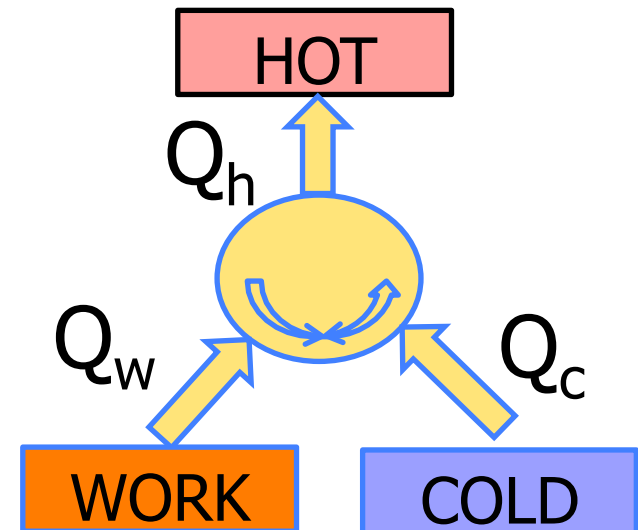


# Trapped Ions

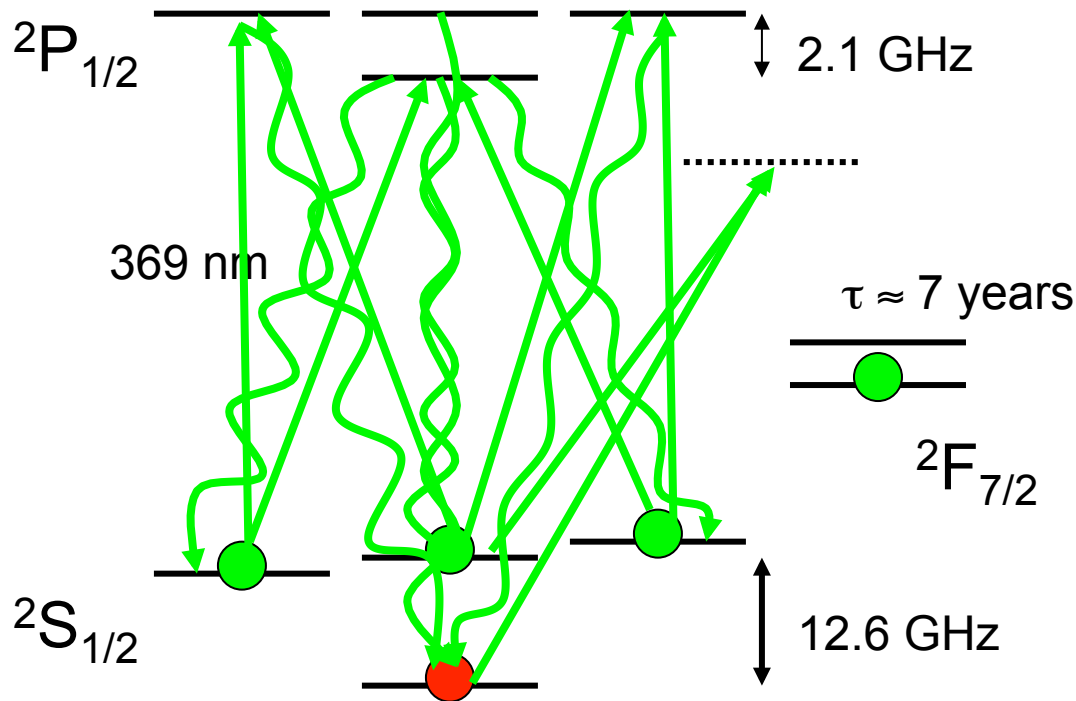


Thanks: R. Blatt, Univ. Innsbruck

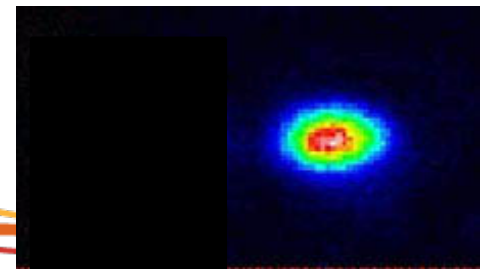
$N$  ions,  $3N$  modes of  
motion,  
(Heat bodies)



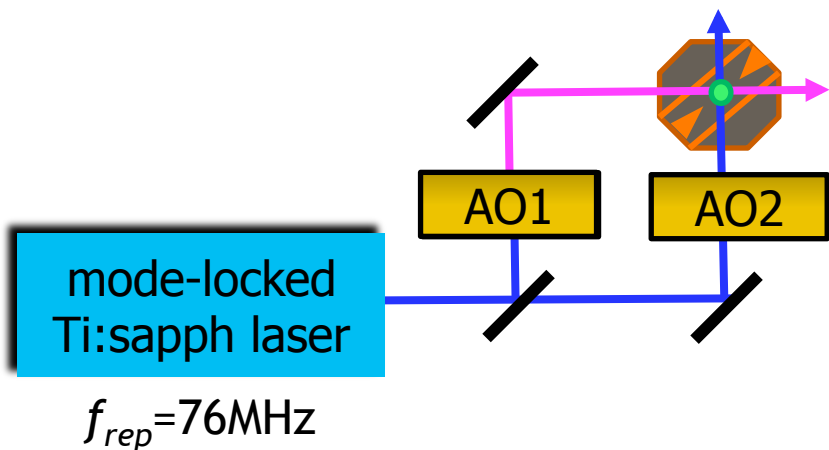
# Experiment: Ytterbium ( $^{171}\text{Yb}^+$ ) Ions



- Optical pumping for state initialization
- Resonance fluorescence for state detection
- Metastable  $2F_{7/2}$  state
- Raman transitions between hyperfine states



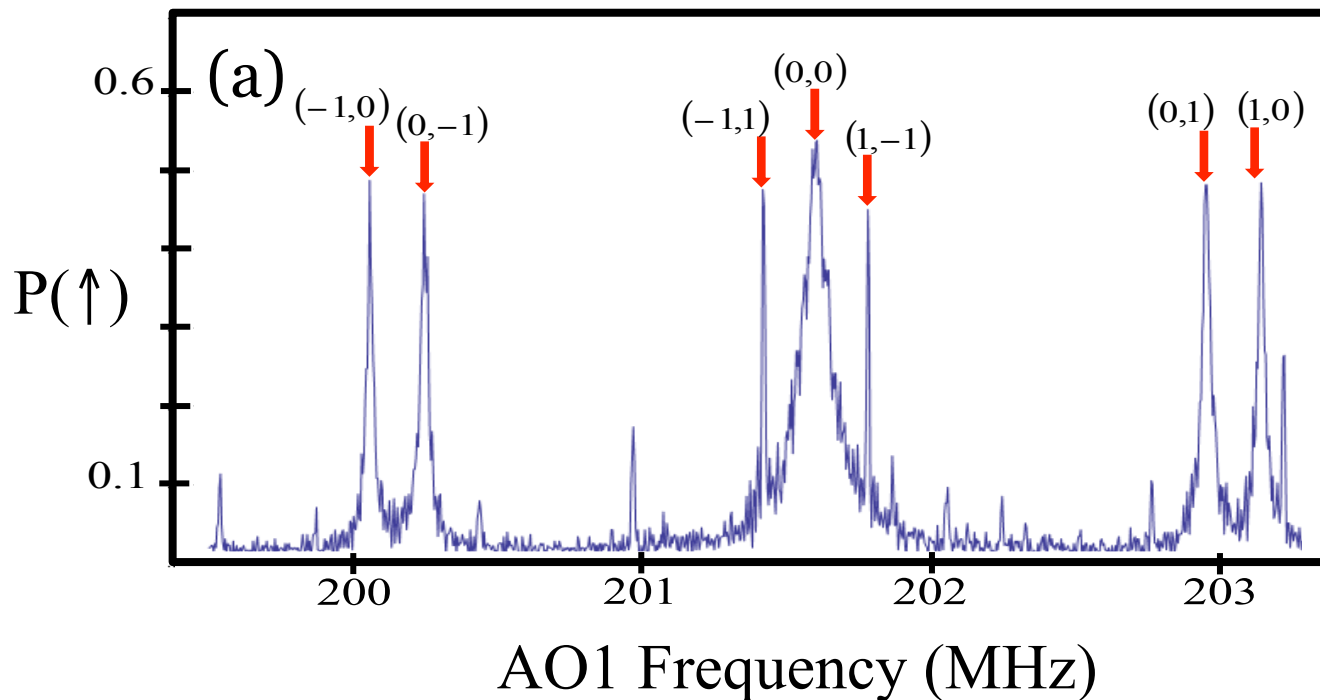
# State preparation




carrier  $H \downarrow c = \hbar\Omega/2 \sigma \downarrow + + h.c.$

blue sideband  $H \downarrow bsb = \hbar\eta\Omega/2 \sigma \downarrow + a \uparrow + +$

red sideband  $H \downarrow rsb = \hbar\eta\Omega/2 \sigma \downarrow + a + +$

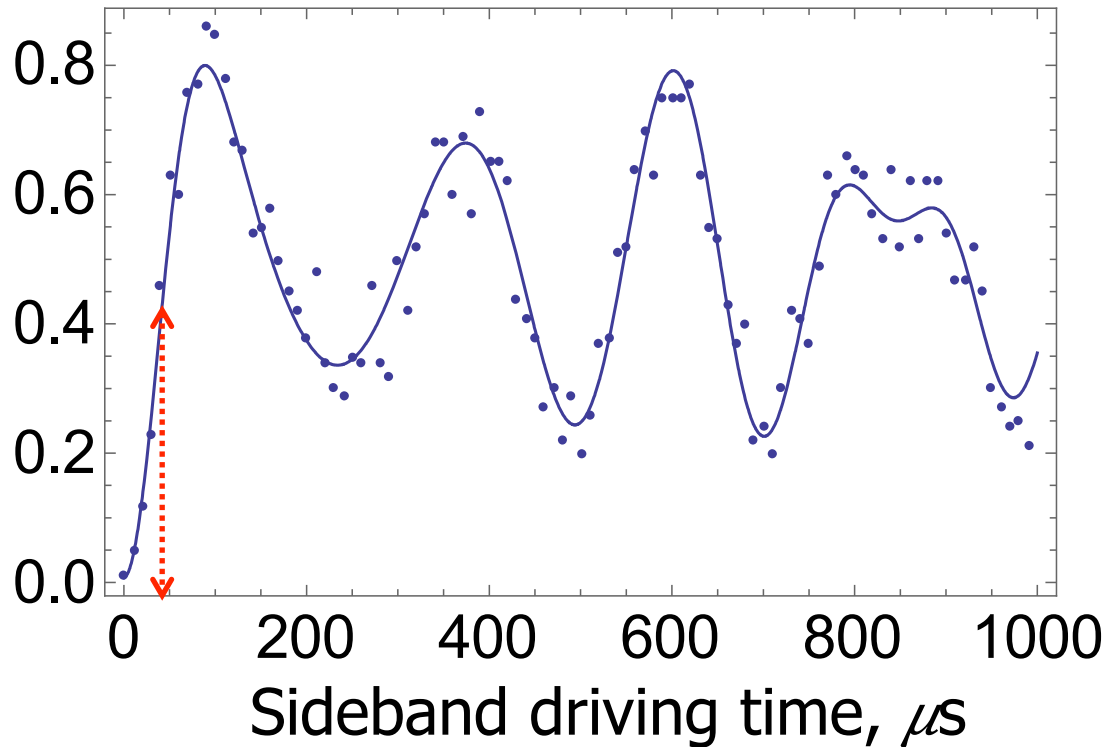


Sisyphus,  
 sideband cooling  
  
 ground state

# How to probe ion motion

$$H_{\text{drive}} = \hbar \eta \Omega / 2 \sigma_{\downarrow} + a \dagger + h.c.$$

$$\Omega_{\downarrow n, n+1} = \sqrt{n+1} \Omega_{\downarrow 0, 1} \quad p_{\downarrow} \uparrow$$



Measure  $p_{\downarrow} \uparrow$  as a function of time

Fourier transform gives  $p_{\downarrow} n$  and  $n$

Requires a lot of data

Measure  $p_{\downarrow} \uparrow$  for fixed  $\tau$

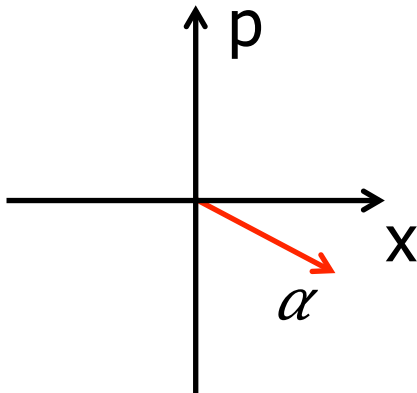
Sensitive to small changes of  $n$

More assumptions

## Coherent state

Force modulated at the mode frequency  $\omega$

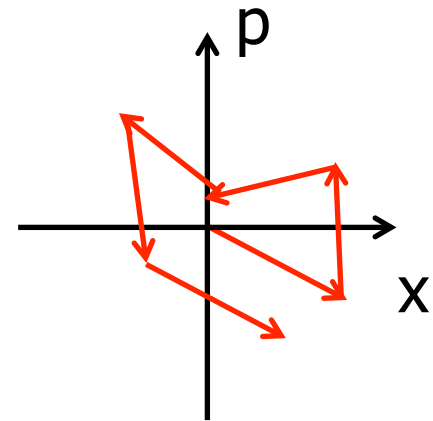
$$|\alpha\rangle = D(\alpha)|0\rangle$$



## Thermal state

Random walk in phase space:  
Random  $\phi$  for each step

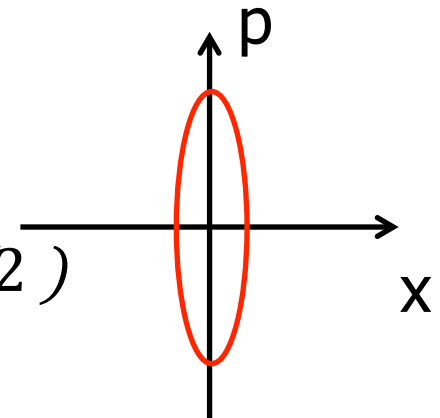
$$D(\alpha e^{i\phi_1}) D(\alpha e^{i\phi_2}) \dots D(\alpha e^{i\phi_n})$$



## Squeezing operator

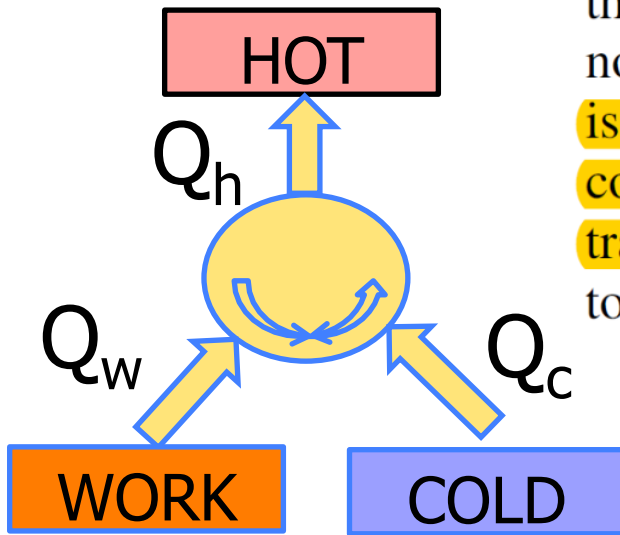
Force modulated at  $2\omega$

$$S(r) = \exp(r a^{\dagger 2} - r^* a^2 / 2)$$



# Can we build refrigerator with ions?

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no movable parts and induce the energy to flow away from a cold reservoir by coupling it with a system that is itself coupled with other (hot) reservoirs. In the quantum regime they were proposed by Kosloff and others [2] using a nonlinear model. Our work implies that, **as nonlinearity is an essential resource, a quantum absorption refrigerator could never be built using harmonic systems such as trapped ions [11] or nanomechanical systems that are ideal to implement other quantum machines [14].** Finally, we

PRL **110**, 130406 (2013)

How can we make motional modes interact with each other ?

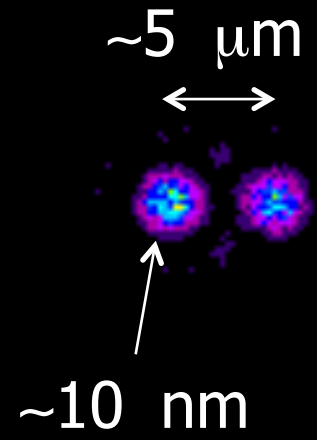
**We need anharmonicity !!!**



# Anharmonicity

Trap size  $\sim$  mm,  
Amplitude of ion motion  $\sim$  10 nm

Harmonic trapping potential  
is a very good approximation



0.5mm

2mm

Coulomb force gives anharmonicity

Push ions closer  $\rightarrow$  weaker radial confinement

Trap frequency  $\sim$  MHz

# Anharmonicity

$$\frac{1}{2} m \omega_x^2 \Delta x^2 + \frac{1}{2} m \omega_y^2 \Delta y^2 + \frac{1}{2} m \omega_z^2 \Delta z^2 + \frac{e^2}{8 \pi \epsilon_0} \sum_{n,m=1,2,3, n \neq m} \dots$$

Harmonic terms

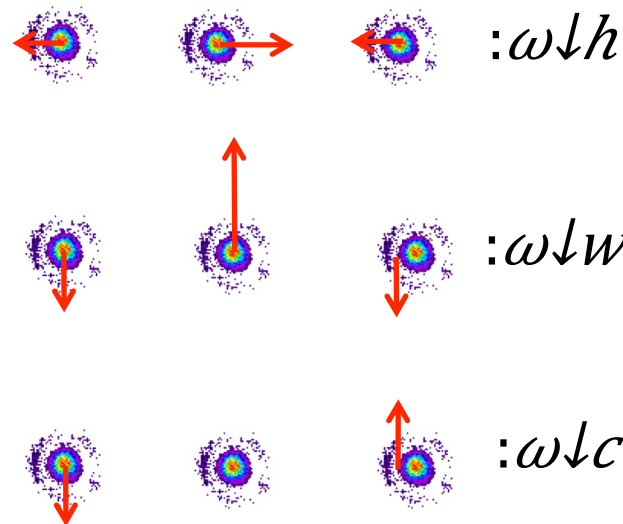
Anharmonicity

Normal mode coordinates -> Taylor expansion -> Quantization of motion

$$H_{int} = \hbar \xi (h^\dagger w c + h w^\dagger c^\dagger)$$

$$\omega \hbar = \omega \hbar_w + \omega \hbar_c$$

$$\frac{1}{2} m \omega^2 \Delta x^2 \approx \sqrt{\hbar / \omega} \omega \hbar_w \quad \text{For } \omega \approx 1 \text{ MHz, } \xi = 2-3 \text{ kHz}$$

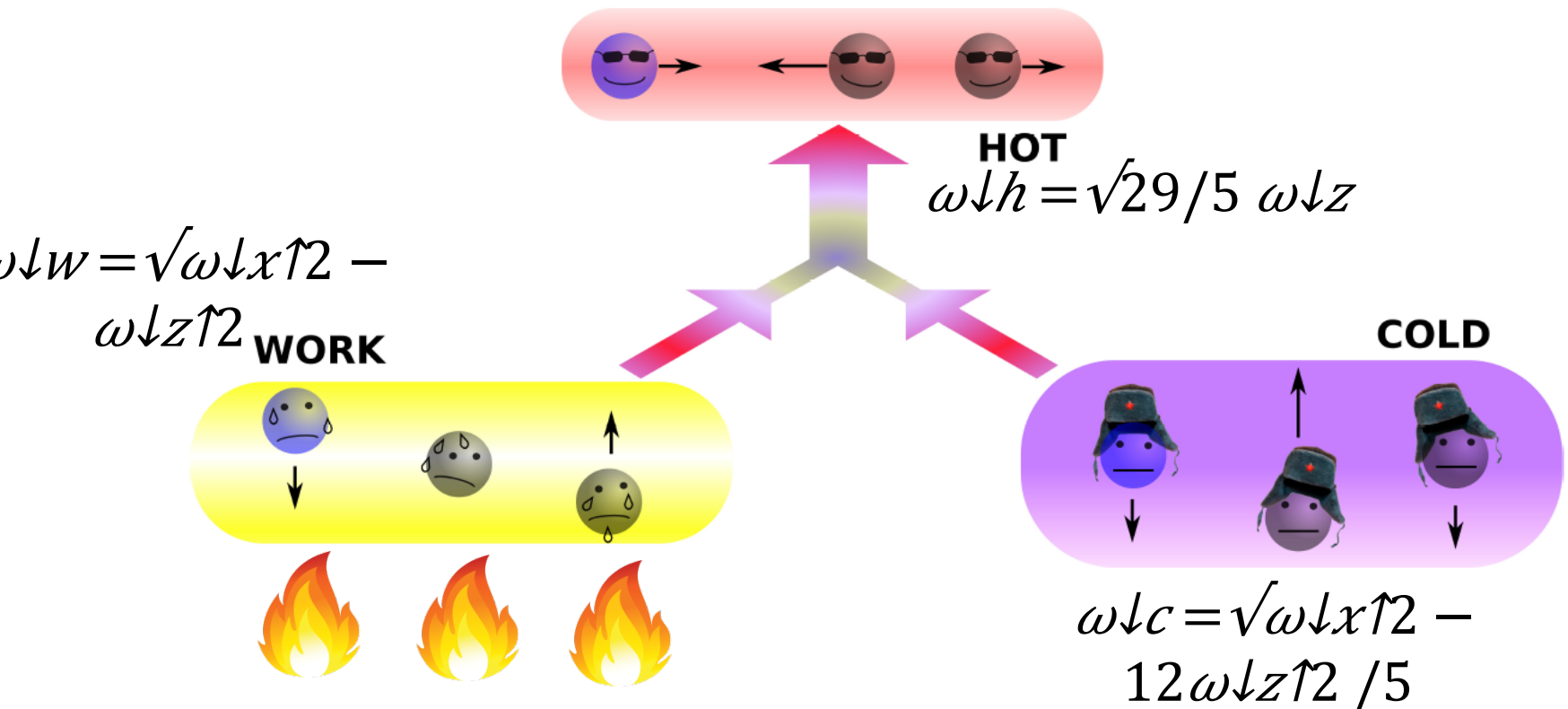


# Refrigerator with trapped ions

Harmonic oscillators interacting via trilinear Hamiltonian

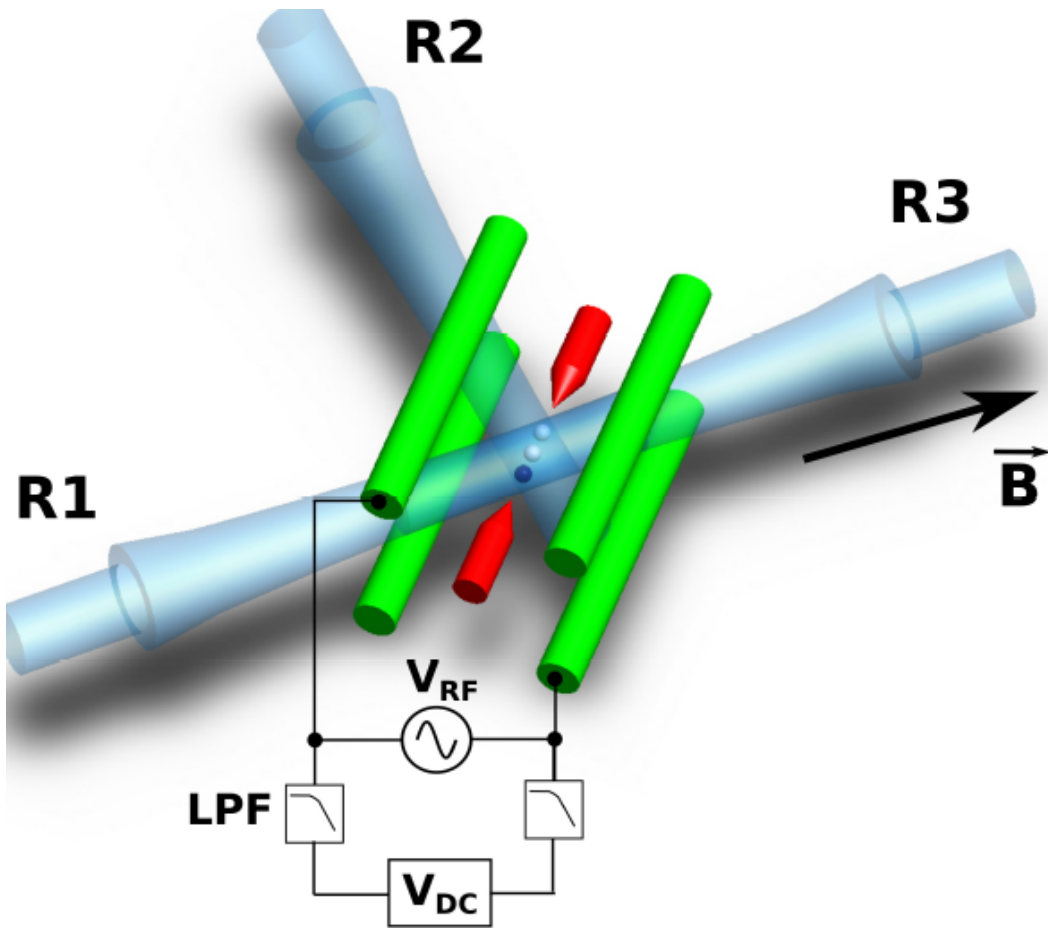
$$H_{\text{int}} = \hbar \xi (h \dagger + w c + h w \dagger + c \dagger)$$

$$\omega_h = \omega_w + \omega_c$$



# Experimental setup

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Three  $^{171}\text{Yb}$  ions in the trap

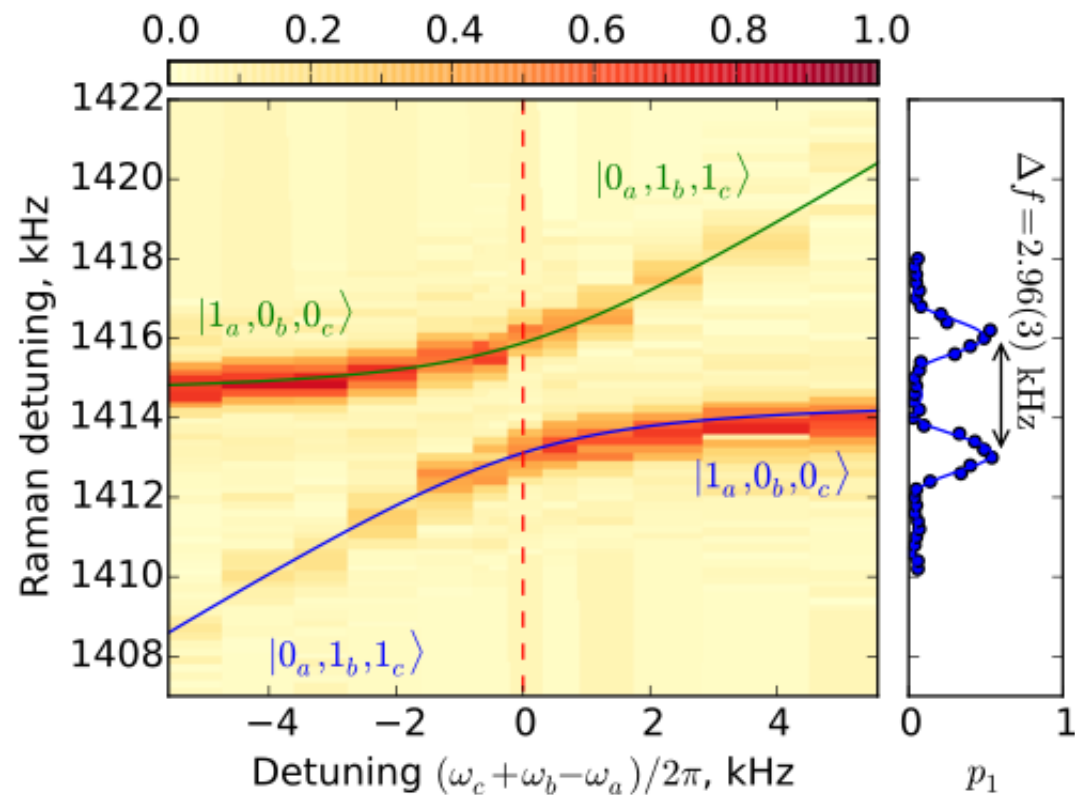
Two pairs of Raman beams to address motion from axial and radial directions

Two ions are pumped in the "dark"  $^2F_{7/2}$  state

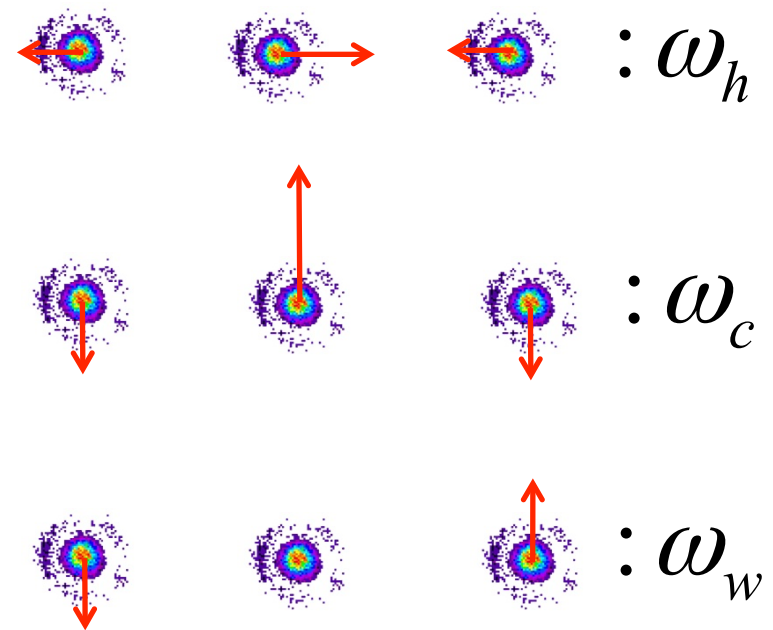
Offset voltages to tune the radial trap frequencies in and out of resonance

# Coupling between three modes

$$H_{\text{int}} = \hbar \xi (h \uparrow \uparrow w c + h w \uparrow \uparrow c \uparrow \uparrow)$$



$$\omega \downarrow h = \omega \downarrow w + \omega \downarrow c$$



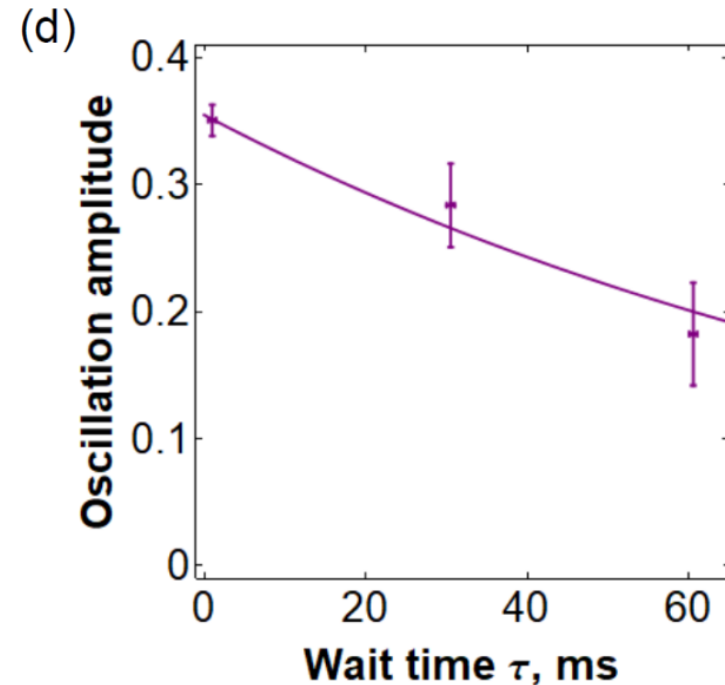
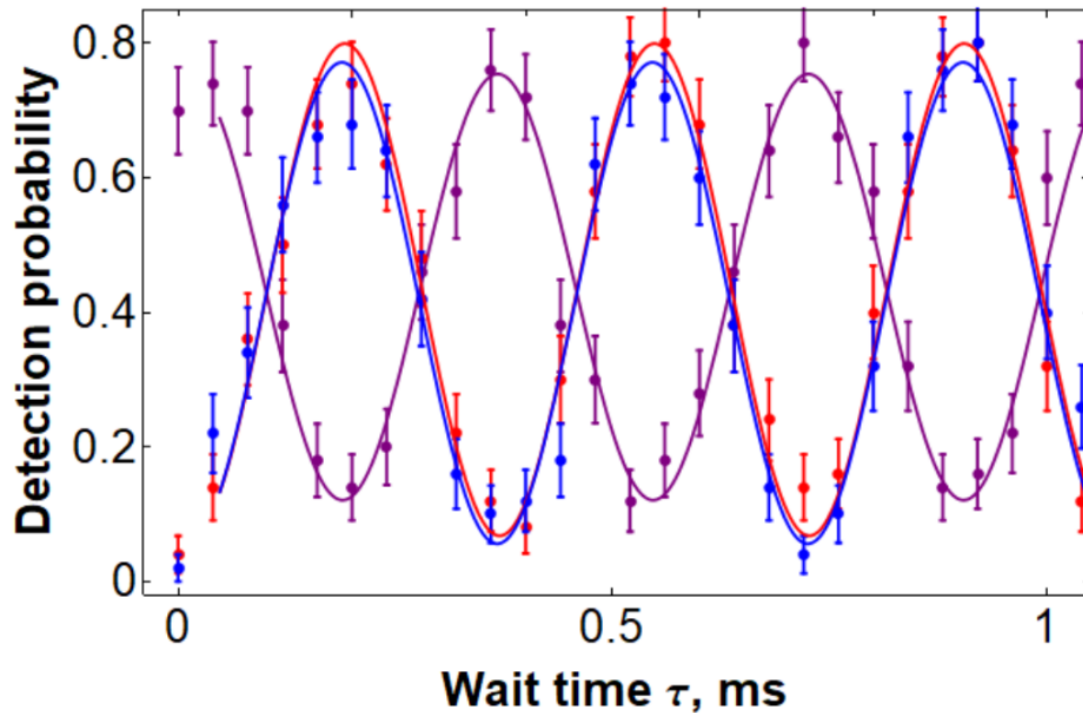


# Coupling of 3 modes, state evolution

$$H_{\text{int}} = \hbar \xi (h \hat{a}^\dagger w c + \hbar w \hat{a}^\dagger c \hat{a}^\dagger)$$

$|h 0 \downarrow w 0 \downarrow c \rangle \rightarrow |0 \downarrow h 1 \downarrow w 1 \downarrow c \rangle \rightarrow |1 \downarrow h 0 \downarrow w 0 \downarrow c \rangle$

→



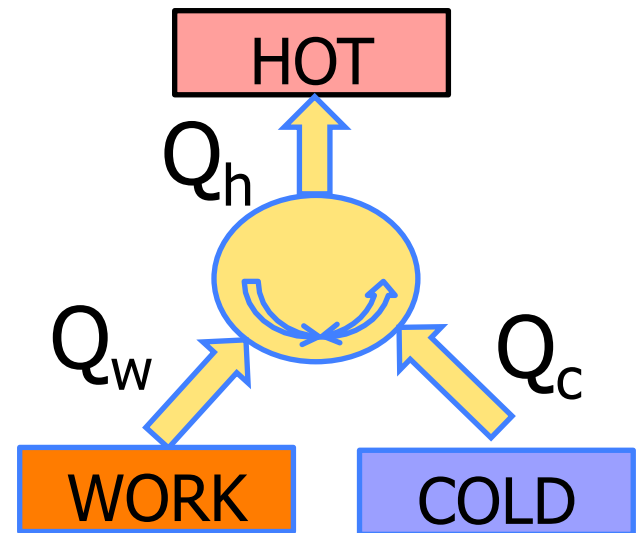
Coherence time > 8ms

S. Ding et. al. arXiv:1805.11193

# Testing refrigerator

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- Prepare initial states
- Tune the trap frequencies to resonance
- Let the modes interact for time  $\tau$
- Turn off interaction and measure mode energies



# Equilibrium

2<sup>nd</sup> law  $\Delta S = Q_h / T_h + Q_w / T_w + Q_c / T_c = 0$

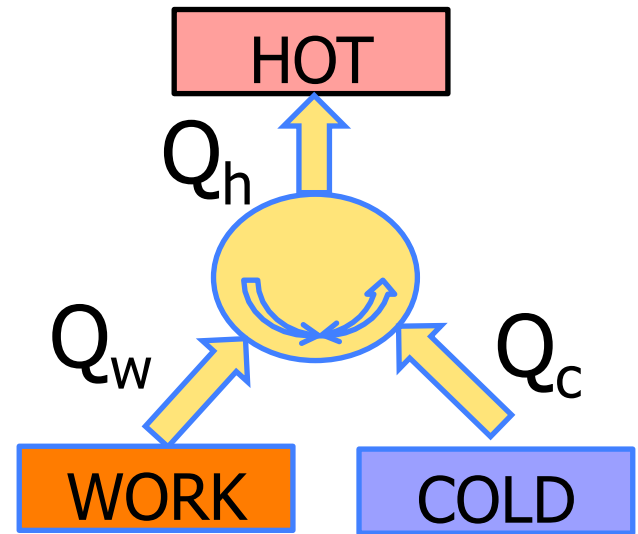
1<sup>st</sup> law  $Q_h = \hbar \omega n$

Coupling Hamiltonian  $n \hbar = -n \omega = -n c$

In Thermal Equilibrium



$n \hbar = (1 + 1/n \omega) (1 + 1/n c)$   
 $n = 1 / (\exp(\hbar \omega / kT) - 1)$



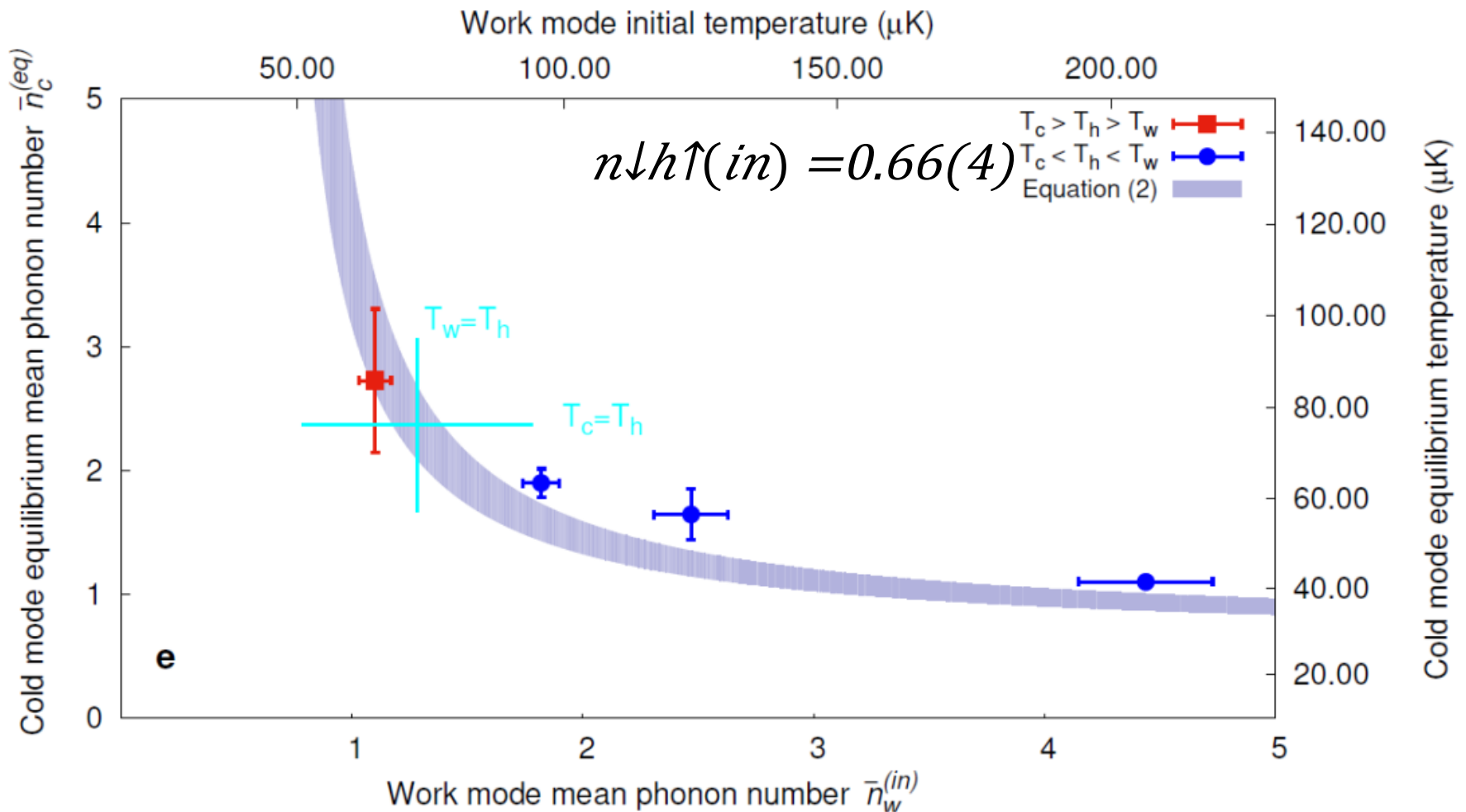
For Harmonic oscillator

Quantum mechanics:  $[H_{int}, \rho_h \otimes \rho_w \otimes \rho_c] = 0$



# Equilibrium temperatures

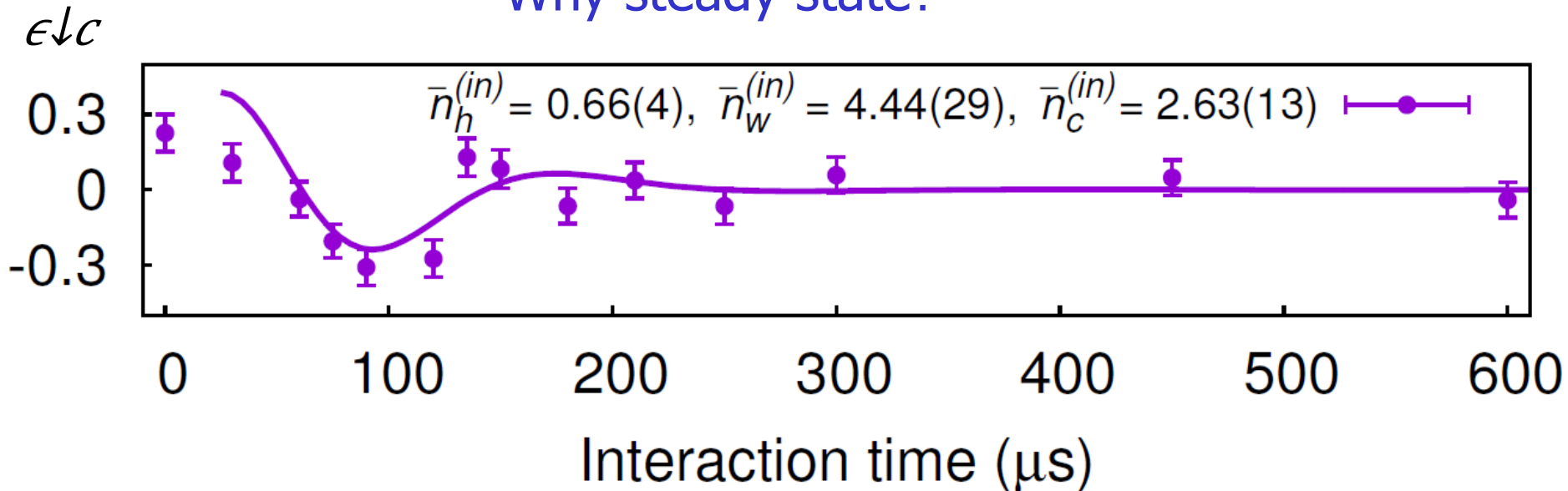
Equilibrium conditions  $(1 + 1/n \downarrow h \uparrow (eq)) = (1 + 1/n \downarrow w \uparrow (eq)) (1 + 1/n \downarrow c \uparrow (eq))$



# Away from equilibrium

Start away from equilibrium. .... Evolution is unitary.

Why steady state?

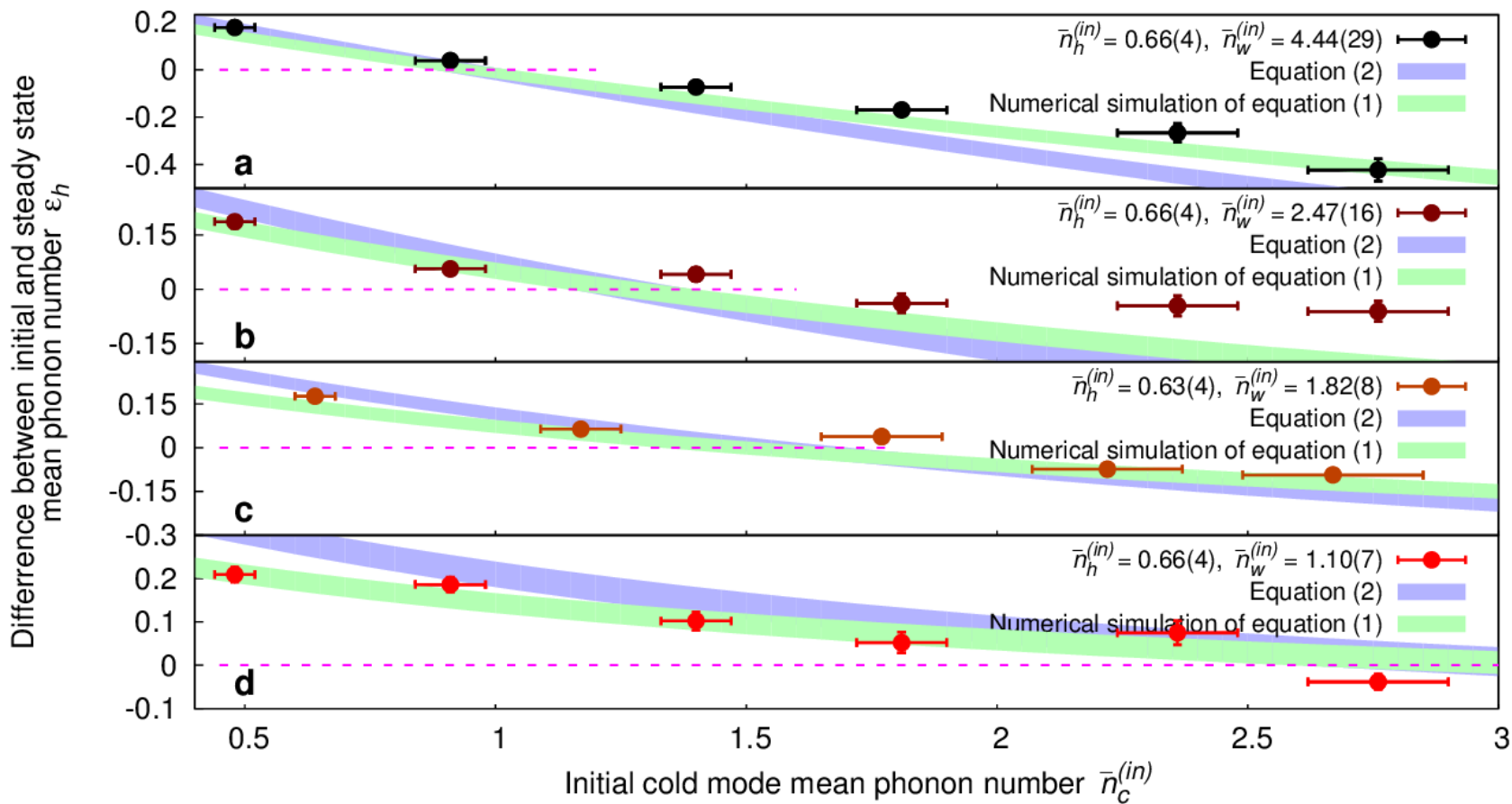


Effective equilibration:

- Large "effective" Hilbert space
- Many Rabi frequencies which are incommensurable

# Steady state

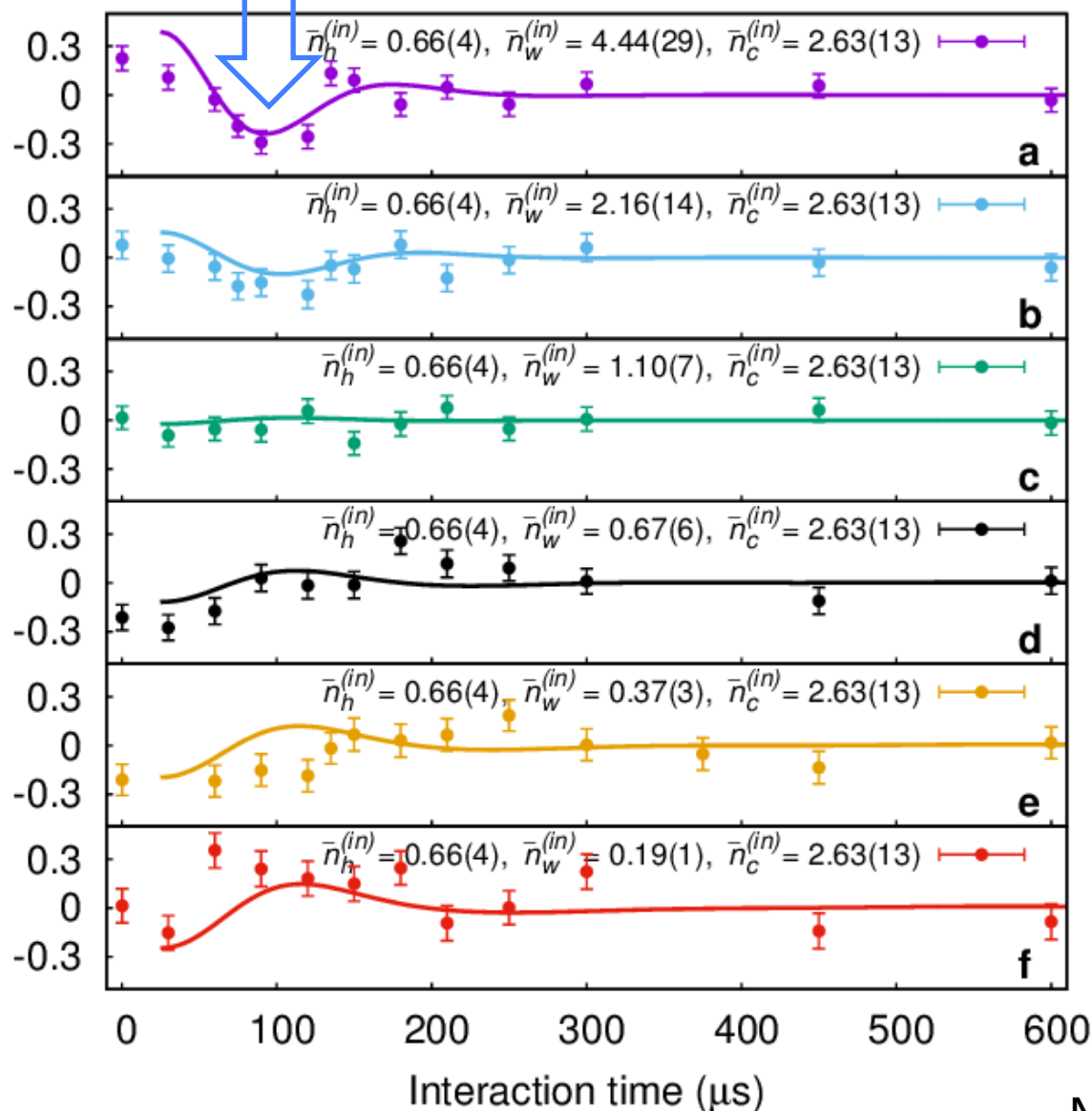
Blue: thermal equilibrium  $\langle n \rangle = (1 + 1/n \downarrow w \uparrow (eq) - \epsilon) (1 + 1/n \downarrow w \uparrow (eq) + \epsilon)$   
 Green: "quantum" steady state



Final states are not thermal

It is cold here !!!

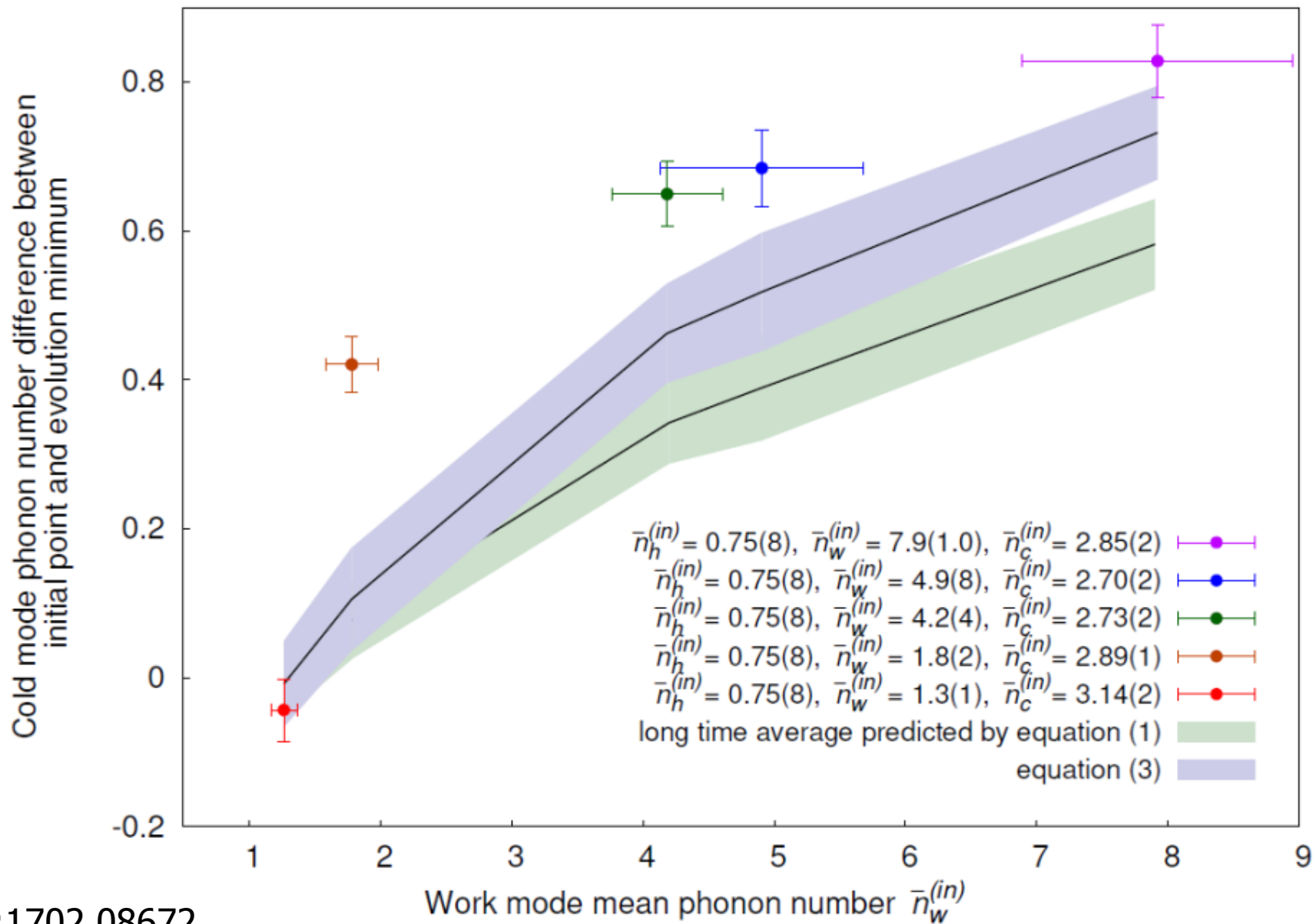
# Single shot cooling



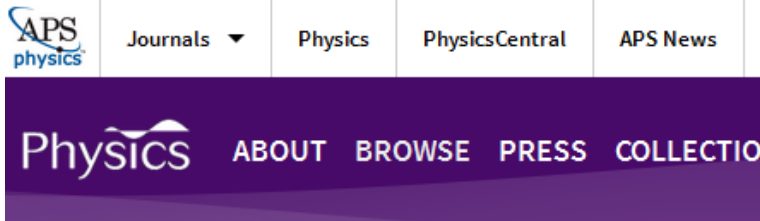
The cold mode  
"overshoots"  
before the steady  
state is reached.

Stop evolution  
when the mode is  
coldest !!!

# Single shot cooling



# Squeezing as a thermal resource



## Synopsis: “Squeezed” Engine Could Break Thermodynamic Limits

January 22, 2014

A theoretical analysis shows that nanoengines based on quantum squeezed states can perform several times more efficiently than classical engines.

Squeezed refrigerator also  
can ???

Prepare work mode in  
thermal state

Apply squeezing operation

$$S(r) = \exp\left(\frac{1}{2} r a \hat{\dagger} \hat{\dagger} - \frac{1}{2} r \hat{*} a \hat{\dagger}^2\right)$$

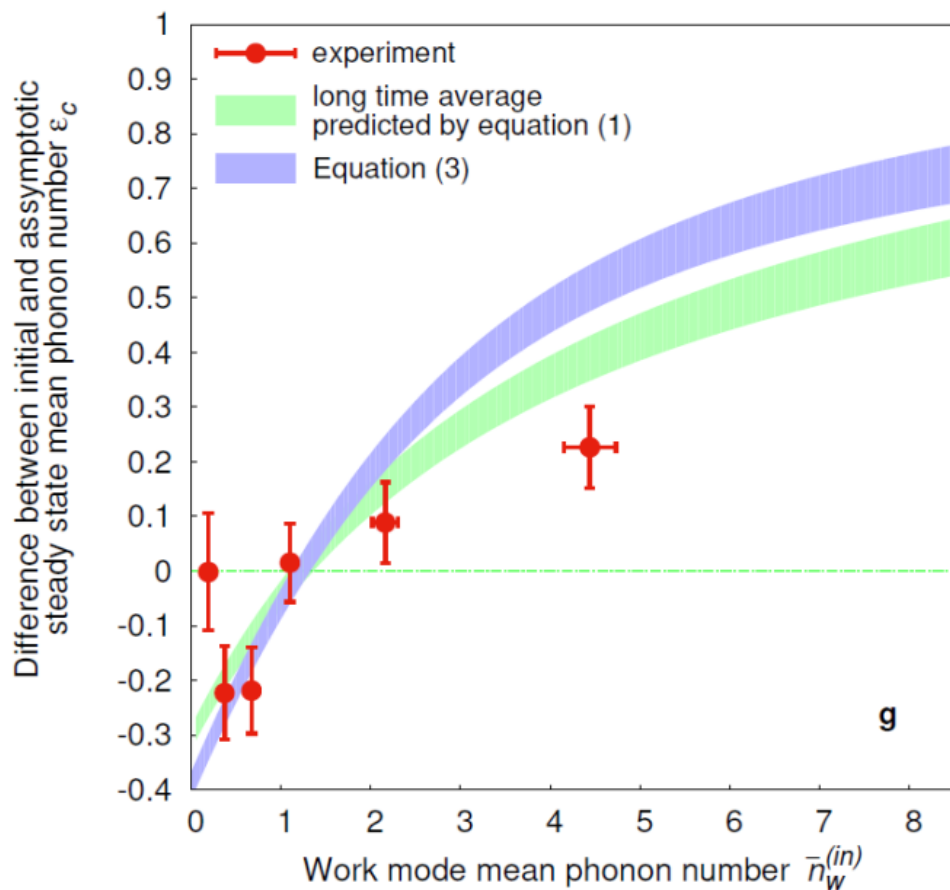
$n$  increases to:

$$n \hat{\dagger}(\text{in,eq}) = n \hat{\dagger}(\text{in}) \cosh(2r) - \sinh \hat{\dagger}^2$$

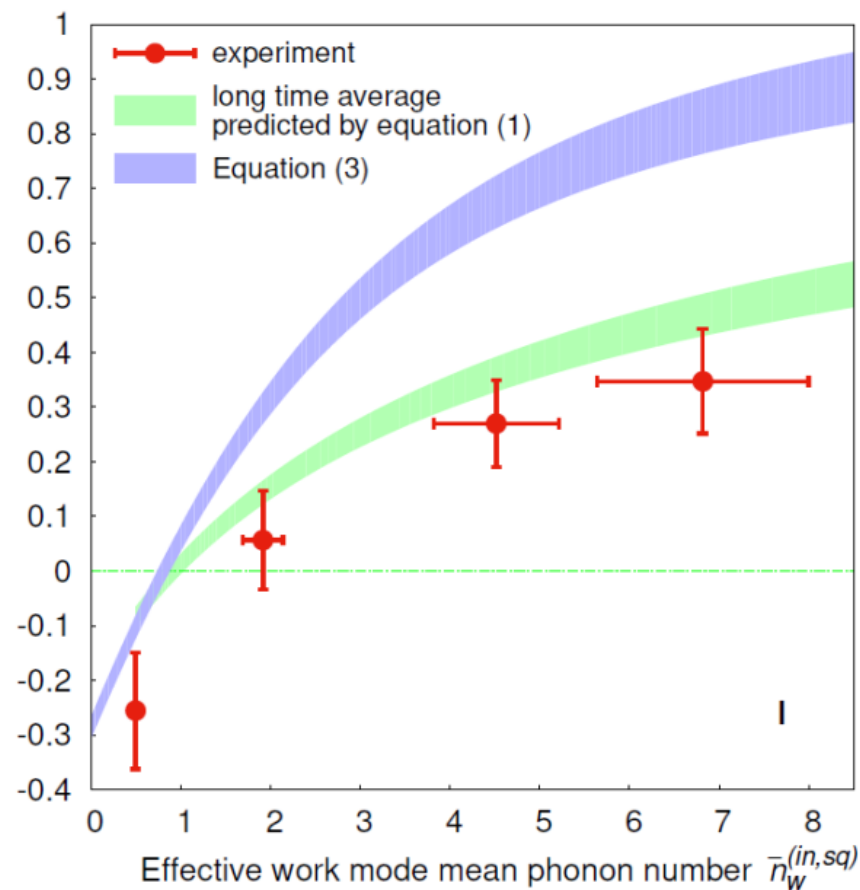
Compare to the thermal state  
with the same  $n$

# Which state of work mode is more efficient

## Thermal



## Squeezed thermal

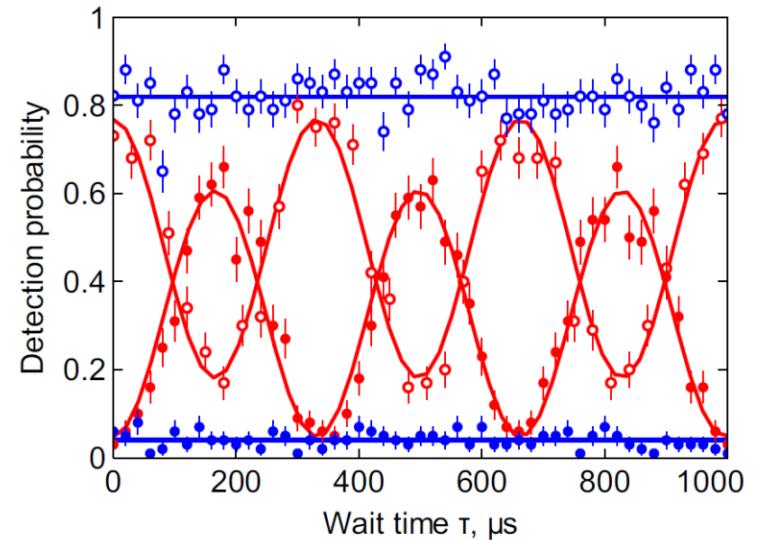


# Other types of nonlinear coupling

Parametric oscillator:

$$H_{\text{int}} = \hbar \xi (a^\dagger \dagger \dagger b + a \dagger \dagger \dagger b^\dagger)$$

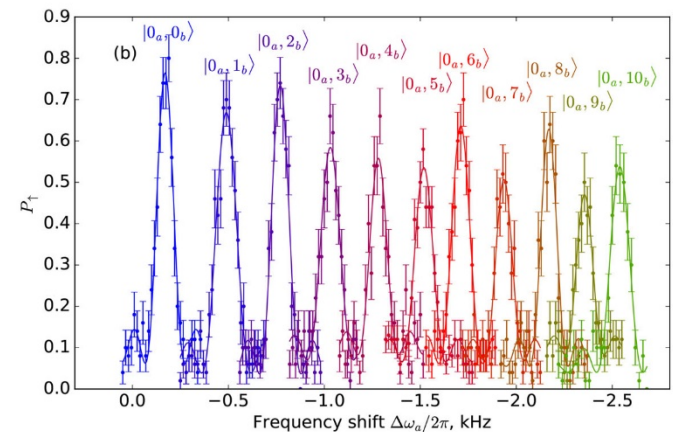
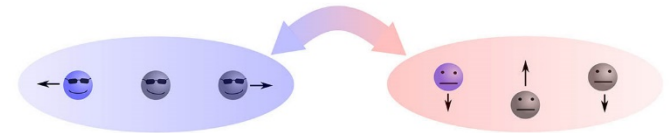
Phys. Rev. Lett. 119, 150404 (2017)



Cross Kerr interaction:

$$H_{\text{int}} = \hbar \xi n \dagger a n \dagger b$$

Phys. Rev. Lett. 119, 193602 (2017)



Any thermodynamic applications ??



- Implemented absorption refrigerator with trapped ions
  - Equilibrium properties of the refrigerator
  - Single shot cooling
  - Effects of squeezing on the fridge performance

Nonlinear coupling enables simulation of other systems

- Jaynes-Cummings model
- Parametric down conversion
- Hawking radiation

$$H_{int} = \hbar \xi (h \hat{a}^\dagger + h \omega c + h \omega \hat{a}^\dagger c^\dagger)$$



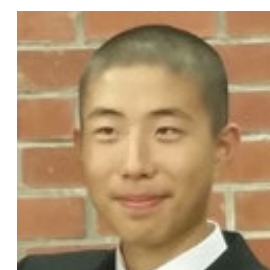
Jaren  
Gan



Gleb  
Maslennikov



Ko Wei  
Tseng



Beicheng  
Lou

## Former members



Shiqian  
Ding  
(JILA)



Huanqian  
Loh  
(NUS)



Roland  
Hablutzel  
(Waterloo)



**NATIONAL  
RESEARCH  
FOUNDATION**

PRIME MINISTER'S OFFICE  
SINGAPORE

Theory : Alexandre Roulet, Stefan Nimmrichter,  
Jibo Dai, Valerio Scarani

# Singapore

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We have atoms ...

ions ...

molecules ...

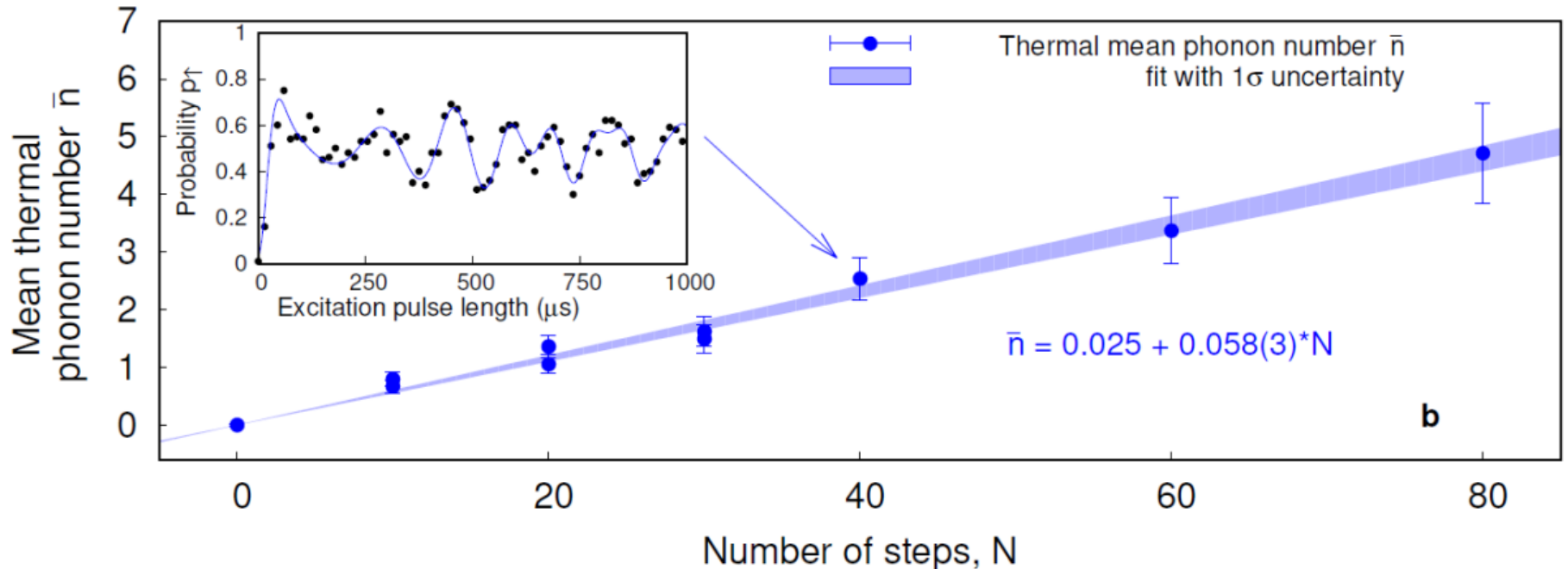
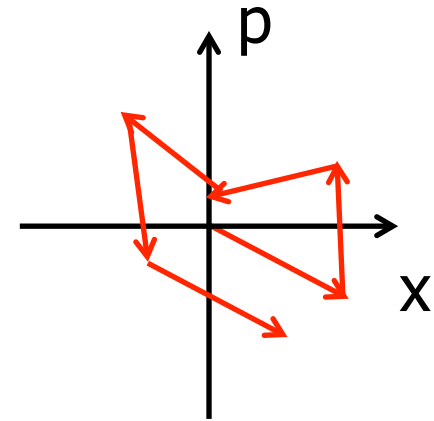


# Thermal state

Random walk in phase space:

$$(\alpha e^{i\phi} \downarrow 1) D (\alpha e^{i\phi} \downarrow 2) \dots D (\alpha e^{i\phi} \downarrow n) |0\rangle$$

Phase chosen randomly for each step:



# Thermal state preparation

