# Quasi-periodic Lattices: memory effects and work statistics



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### THE ENVIRONMENT

Fermi-Hubbard model with a quasi periodic on-site potential

$$\hat{H} = -J\sum_{i} (\hat{a}_{i+1}^{\dagger} \hat{a}_{i} + \hat{a}_{i}^{\dagger} \hat{a}_{i+1}) + \sum_{i} V_{i} \hat{n}_{i}$$

Interplay **quasi-periodic** potential and **kinetic** term

 $V_i = \Delta \cos(2\pi\beta i + \phi)$   $\beta$  Irrational

#### N-M AND LOCALISATION

#### Bose lattice as controllable environment



#### N-M AND LOCALISATION

## Bose lattice as controllable environment superfluid - Markovian





Quantum non-Markovianity induced by Anderson localization

S. Lorenzo, F. Lombardo, F. Ciccarello, and G. M. Palma, Scientific Reports 7, 42729 (2017)

#### N-M AND LOCALISATION

# Bose lattice as controllable environment superfluid - Markovian

1.0

0.8

0.6

0.4

0.2

0.0

0.0

0.5

0.00 0.02 0.04 0.06 0.08 0.10

1.5

2.0

1.0

 $\sigma$ 

 $\overline{\mathcal{N}}$ 



Quantum non-Markovianity induced by Anderson localization

S. Lorenzo, F. Lombardo, F. Ciccarello, and G. M. Palma, Scientific Reports 7, 42729 (2017)

### THE ENVIRONMENT



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#### Aubry-André model $V_i = \Delta \cos(2\pi\beta i + \phi)$ often compared $\beta$ Irrational **Anderson insulator** Eigenstates **localisation** $V_i \in [-\Delta, +\Delta]$ $\Delta/J \ge 2$ Eigenstates $\Delta > 0$ **localisation**

### THE SYSTEM: IMPURITY



#### Impurity atom

- Motional ground state  $\psi_0(x)$
- two lowest internal states  $\left| e 
  ight
  angle \left| g 
  ight
  angle$

$$\hat{H}_S = \frac{\omega_0}{2} \hat{\sigma}_z$$

### THE S-E INTERACTION



#### density-density interaction

#### LOSCHMIDT ECHO OF THE FERMI LATTICE

$$|\rho_{eg}(t)/\rho_{eg}(0)| = |\langle \Psi_0|e^{i\hat{H}_g t}e^{-i\hat{H}_e t}|\Psi_0\rangle| = \sqrt{L(t)}$$

$$\hat{H}_g = \langle g|\hat{H}|g\rangle \quad \hat{H}_e = \langle e|\hat{H}|e\rangle$$

$$probe \qquad \qquad \rho(0) = |\varphi\rangle\langle\varphi| \otimes |\Psi_0\rangle\langle\Psi_0|$$

$$|\varphi\rangle = \frac{1}{2}(|g\rangle + |e\rangle)$$
environment

### IMPURITY DYNAMICS

No initial correlations

$$\hat{\rho}_S(t) = \Lambda_t[\hat{\rho}_S(0)] = \operatorname{Tr}[\hat{U}(t)\rho_S(0) \otimes \rho_E(0)\hat{U}^{\dagger}(t)]$$
$$= \begin{pmatrix} \rho_{gg}(0) & \chi^*(t)\rho_{ge}(0) \\ \chi(t)\rho_{eg}(0) & \rho_{ee}(0) \end{pmatrix}$$

Decoherence function

$$\chi(t) = \langle \Psi_0 | e^{i\hat{H}_g t} e^{-i\hat{H}_e t} | \Psi_0 \rangle$$

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$$\chi(t) = \det(1 - \hat{r} + \hat{r}e^{-i\hat{h}_e t}e^{i\hat{h}_g t})$$

Klich I 2003 An Elementary Derivation of Levitov's Formula BT - Quantum Noise in Mesoscopic Physics ed Nazarov Y V

### IMPURITY DYNAMICS

Uncorrelated initial state

$$\hat{\rho}_S(t) = \Lambda_t[\hat{\rho}_S(0)] = \operatorname{Tr}[\hat{U}(t)\rho_S(0) \otimes \rho_E(0)\hat{U}^{\dagger}(t)]$$
$$= \begin{pmatrix} \rho_{gg}(0) & \chi^*(t)\rho_{ge}(0) \\ \chi(t)\rho_{eg}(0) & \rho_{ee}(0) \end{pmatrix}$$

Decoherence function

$$\chi(t) = \det(1 - \hat{r} + \hat{r}e^{-i\hat{h}_{e}t}e^{i\hat{h}_{g}t})$$
  
Environment  
Single particle counterpart of  
$$\hat{H}_{e/g} = \langle e/g | \hat{H}_{AA} + \hat{H}_{int} | e/g \rangle$$

Klich I 2003 An Elementary Derivation of Levitov's Formula BT - Quantum Noise in Mesoscopic Physics ed Nazarov Y V

### ENVIRONMENT

Initial state of the environment

$$|\Psi_{0}\rangle = \prod_{i \text{ odd}} \hat{a}_{i}^{\dagger}|0\rangle \implies \chi(t) = \det(1 - \hat{n}_{cdw} + \hat{n}_{cdw}e^{-ih_{e}t}e^{ih_{g}t})$$
$$\hat{n}_{cdw} = \sum_{i \text{ odd}} |i\rangle\langle i|$$

Charge density wave state (CDW)



H. P. Lüschen, P. Bordia, S. Scherg, F. Alet, E. Altman, U. Schneider, and I Bloch, Phys. Rev. Lett. **119**, 260401 (2017) S. Vardhan, G. De Tomasi, M. Heyl, E. J. Heller, and F. Pollmann, Phys. Rev. Lett. **119**, 016802 (2017)





- Information Flow  $\sigma(t) \equiv \frac{d}{dt} D(\Phi_t \rho_S^1, \Phi_t \rho_S^2)$
- BLP-measure  $\mathcal{N}(\Phi) = \max_{\substack{\rho_{S}^{1,2}}} \int_{\sigma>0} dt \sigma(t)$

Heinz-Peter Breuer, Elsi-Mari Laine, and Jyrki Piilo Phys. Rev. Lett. **103**, 210401 (2009)









Lattice size

 $N_s = 233, 377, 987$ 



Lattice size

 $N_s = 233, 377, 987$ 















### ADIABATIC COUPLING

Asymptotic coherences of the probe

 $\frac{\rho_{eg}(+\infty)}{\rho_{eg}(0)} = |\langle \Psi_0(\epsilon=0)|U(+\infty)|\Psi_0(\epsilon=0)\rangle| = |\langle \Psi_0(\epsilon=0)|\Psi_0(x,\epsilon)\rangle|$ 

Ground states fidelity

$$F(x,\epsilon) = |\langle \Psi_0(\epsilon=0) | \Psi_0(x,\epsilon) \rangle|$$

Statistical Orthogonality catastrophe

$$F_{\text{typ}} \equiv \exp(\overline{\log F}) \sim \exp(-\beta L)$$

V. Khemani, R. Nandkishore, and S. Sondhi, Nature Physics **11** 560–565 (2015) D. L. Deng, J. Pixley, X. Li, and S. D. Sarma, Physical Review B **92** 220201 (2015)

### ADIABATIC COUPLING

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Ground states fidelity

$$F(x,\epsilon) = |\langle \Psi_0(\epsilon=0) | \Psi_0(x,\epsilon) \rangle|$$

Orthogonality catastrophe "probability"

$$\sigma(\epsilon) = \left(\frac{1}{N_s} \sum_{x=1}^{N_s} \theta(\delta - F(x, \epsilon))\right)$$

### GS FIDELITY



#### Half filling

### GS FIDELITY



Half filling

#### PLATEAU-GAP

#### Occupied spectrum



### GAP OPENING



### GAP OPENING



Gaps opening at  $E_{g_2} \simeq \pm \Delta \cos(\pi \beta)$ 

With states on opposite sides of the gap close in space

D. J. Thouless and Q. Niu 1983 J. Phys. A: Math. Gen. 16 1911

#### TWO-LEVEL MODEL

Mechanism **weak** impurity

$$H_2 = \begin{bmatrix} \Delta_2 + \epsilon & -J \\ -J & 0 \end{bmatrix}$$







#### TWO-LEVEL MODEL

Impurity affects neighbouring site  $H_2 = \begin{bmatrix} \Delta_2 & -J \\ -J & \epsilon \end{bmatrix}$ 



![](_page_36_Figure_3.jpeg)

### TWO-LEVEL MODEL

![](_page_37_Figure_1.jpeg)

### ADIABATIC TRANSPORT

![](_page_38_Figure_1.jpeg)

D. L. Deng, J. Pixley, X. Li, and S. D. Sarma, Physical Review B 92 220201 (2015)

#### RADIUS OF DISTURBANCE

$$\mathbf{R}_{x}(\epsilon) = \left\langle \sum_{j=1}^{N_{s}} \left| (j-x) \left[ n(j) - \tilde{n}(j,\epsilon,x) \right] \right| \right\rangle$$

![](_page_39_Figure_2.jpeg)

#### IMBALANCE

![](_page_40_Figure_1.jpeg)

H. P. Lüschen, P. Bordia, S. Scherg, F. Alet, E. Altman, U. Schneider, and I Bloch, Phys. Rev. Lett. 119, 260401 (2017)

### WS AND STOC

![](_page_41_Figure_1.jpeg)

#### CONCLUSIONS

- Oifferent coupling can be used to highlight different properties of a quasi periodic lattice
- Non-trivial connection between energy-space correlation, orthogonality catastrophe and density fluctuations induced by localised impurities
- Connection between non-Markovianity and localisation in quasi-periodic systems (out of equilibrium environment)

#### PEOPLE

Sabrina Maniscalco Elsi-Mari Laine Massimo Borrelli

![](_page_43_Picture_2.jpeg)

#### Antonello Scardicchio

![](_page_43_Picture_4.jpeg)

#### Saverio Pascazio

![](_page_43_Picture_6.jpeg)

![](_page_43_Picture_7.jpeg)

![](_page_43_Picture_8.jpeg)

![](_page_43_Picture_9.jpeg)

arXiv:1803.04382

#### THANKS FOR THE ATTENTION