

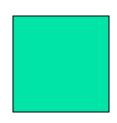
### String-net condensation and topological phases in quantum spin systems

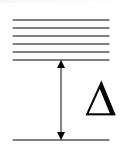
Michael Levin, Xiao-Gang Wen *MIT* 





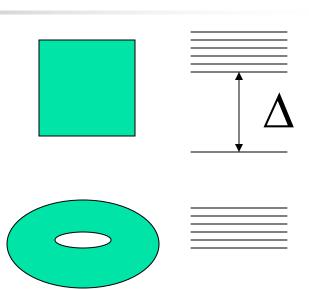
Gapped





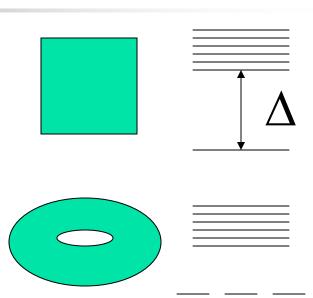


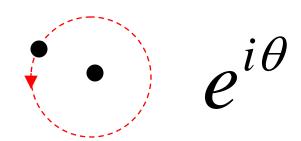
- Gapped
- Degenerate ground state on torus





- Gapped
- Degenerate ground state on torus
- Fractional statistics







#### Real life examples

• FQH liquids.

## 4

#### Real life examples

- FQH liquids.
- Hope: Frustrated magnets
  - Many theoretical models
  - A few candidate materials
    - Cs<sub>2</sub>CuCl<sub>4</sub>
    - $\kappa$ -(BEDT-TTF)<sub>2</sub>Cu<sub>2</sub>(CN)<sub>3</sub>



#### Theory of topological phases



#### Theory of topological phases

- We understand:
  - Low energy/Long distance physics

- We're missing:
  - Connection with microscopics!



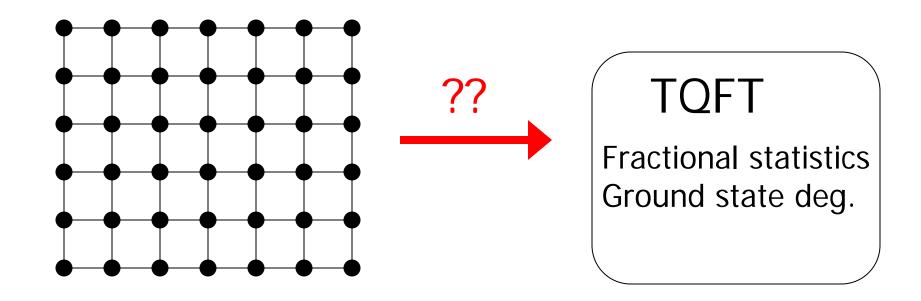
# How do topological phases emerge from microscopic spins?

How can we realize them? What interactions favor them?



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### Outline

I. Physical picture

II. Quantitative results

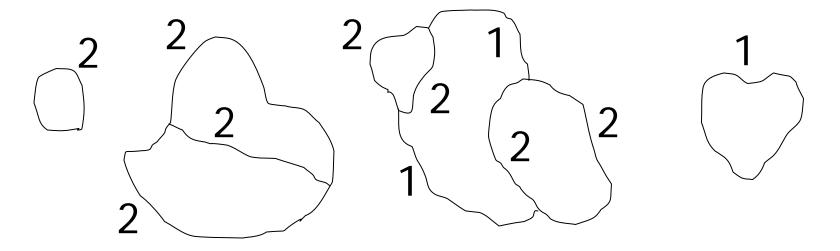
A. Explicit ground state wave functions

B. Exactly soluble Hamiltonians

III. Examples



#### String-net models

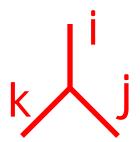


### Data

1. String types: Number of string types N.

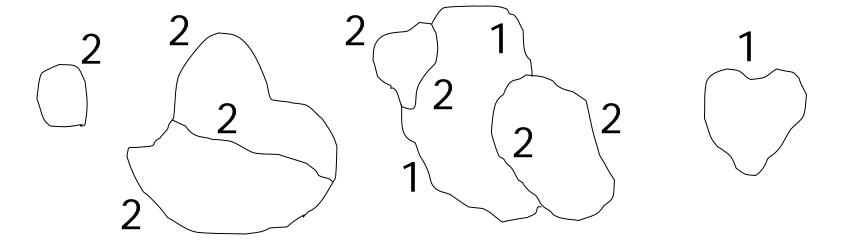
$$_{\underline{\hspace{1cm}}}^{i}$$
 (i = 1,...,N)

2. Branching rules: Triplets {i, j, k} allowed to meet at a point.

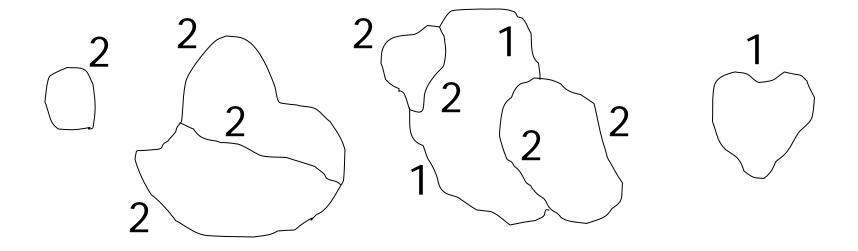




#### Data



### Data



- 1. Number of string types: N = 2.
- 2. Branching rules: {2, 2, 2}, {1, 2, 2}.



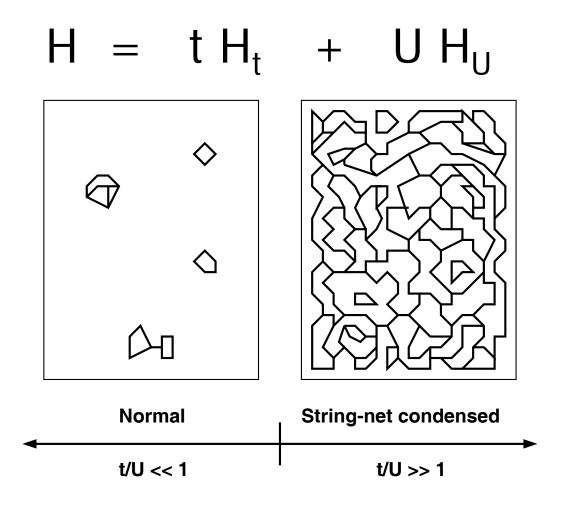
#### String-net Hamiltonian

$$H = t H_t + U H_U$$

$$String String kinetic tension energy$$

### 4

#### String-net Hamiltonian



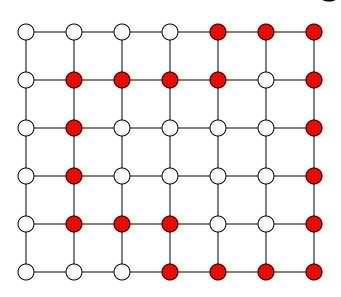


- String-net condensed phases ARE topological phases!
- Mechanism for topological phases
- Very general: all non-chiral topological phases can be realized





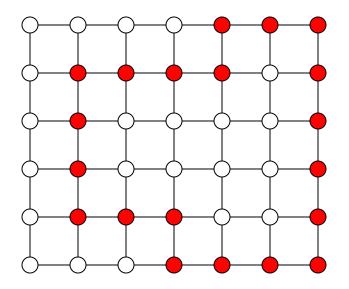
Low energy degrees of freedom can be string-like:

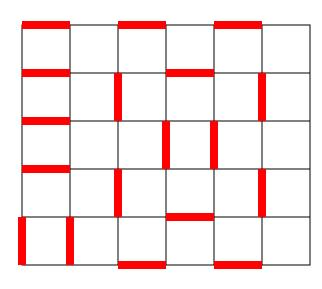




# What does this have to do with spin systems?

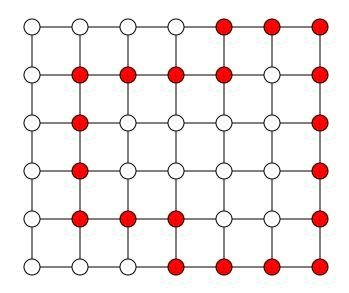
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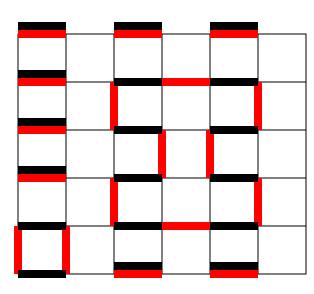






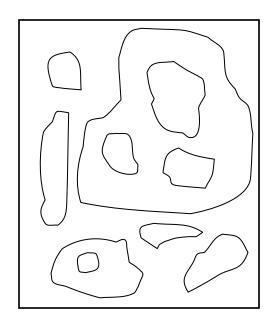
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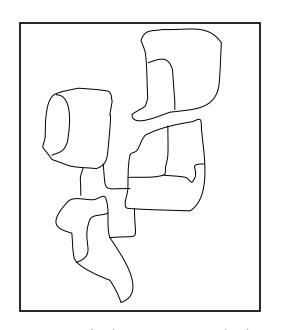




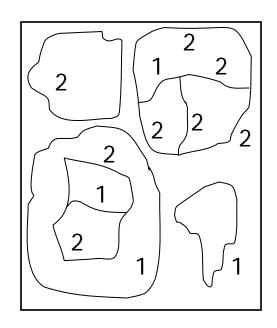
#### Examples



Z<sub>2</sub> gauge theory



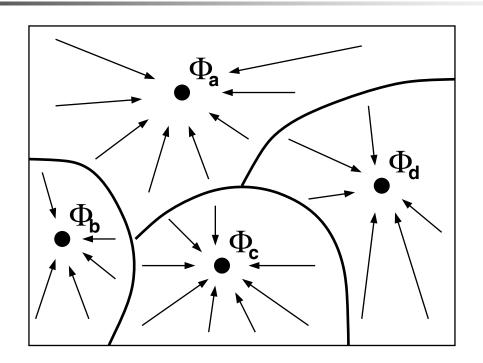
 $SO_3(3) \times SO_3(3)$ Chern-Simons



S<sub>3</sub> gauge theory

### 4

#### Representative wave functions



Want "fixed-point" wave functions:

$$\Phi(\bigcirc\bigcirc\bigcirc\bigcirc) = \dots$$

### Ans

#### Ansatz

- 1. Amplitude of  $\Phi$  only depends on topology of string-net: e.g.,  $\Phi(\ \bigcirc\ ) = \Phi(\ \bigcirc\ )$
- 2. Φ satisfies local constraint equations:

$$\Phi(\bigcirc^{i}) = d_{i} \Phi(\bigcirc)$$

$$\Phi(\frac{i}{i}) = 0 \quad \text{if } i \neq j$$

$$\Phi() = \sum_{n} \operatorname{Fijm}_{k \mid n} \Phi() = \sum_{n \mid k} \Phi()$$

# Local constraints specify $\Phi$ completely

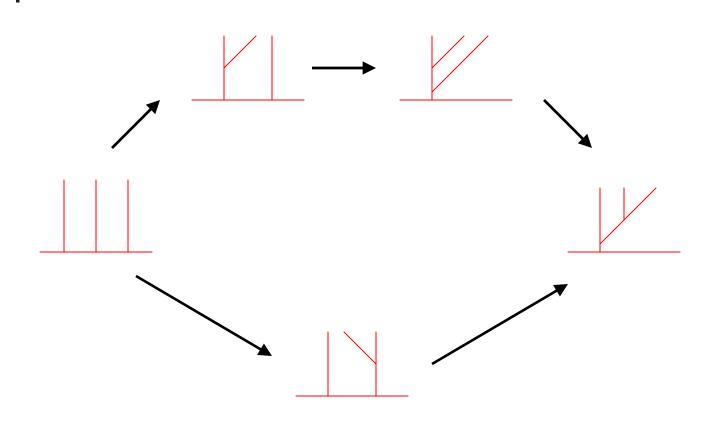
$$\Phi(\underbrace{\downarrow_{k}}) = \sum_{l} F^{ikj}_{kil} \Phi(\underbrace{\downarrow_{k}})$$

$$= F^{ikj}_{ki0} \Phi(\underbrace{\downarrow_{k}})$$

$$= F^{ikj}_{ki0} d_{i} d_{k} \Phi(\text{vacuum})$$

$$= F^{ikj}_{ki0} d_{i} d_{k}$$

# But rules are not usually self-consistent!



### 4

#### Self-consistency conditions

$$\sum_{n} F^{mlq}_{kpn} F^{jip}_{mns} F^{jsn}_{lkr} = F^{jip}_{qkr} F^{riq}_{mls}$$
 (a)

$$F^{ijm}_{kln} = F^{lkm}_{jin} = F^{jim}_{lkn} = F^{imj}_{knl} (d_m d_n / d_j d_l)^{1/2}$$
 (b)

$$F^{ijk}_{ji0} = (d_k/d_id_j)^{1/2} \delta_{ijk}$$
 (c)

(where  $\delta_{ijk} = 1$  if  $\{i,j,k\}$  allowed, 0 otherwise).



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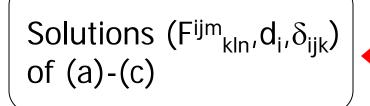
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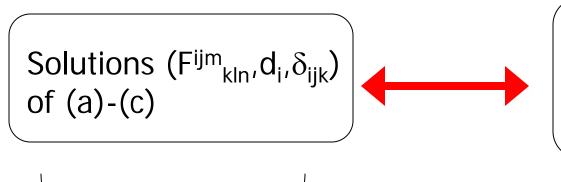
Solutions  $\Leftrightarrow$  fixed point wave functions  $\Phi$ 

# Classification of non-chiral topological phases



String-net condensates/
non-chiral topological
phases

# Classification of non-chiral topological phases

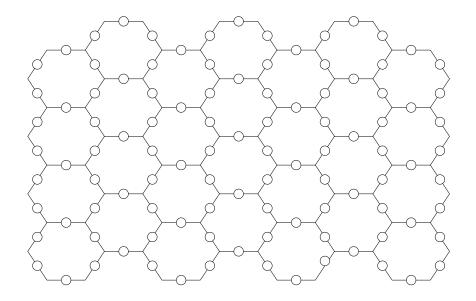


String-net condensates/ non-chiral topological phases

"Tensor categories"



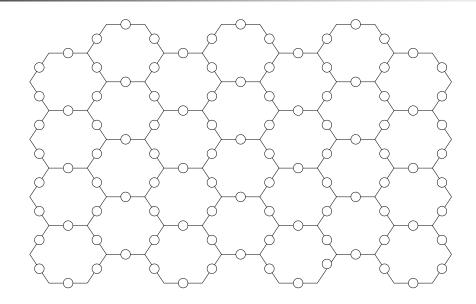
#### Exactly soluble lattice models



Each "spin" can be in N+1 states:  $|0\rangle, |1\rangle, ..., |N\rangle$ 

### 4

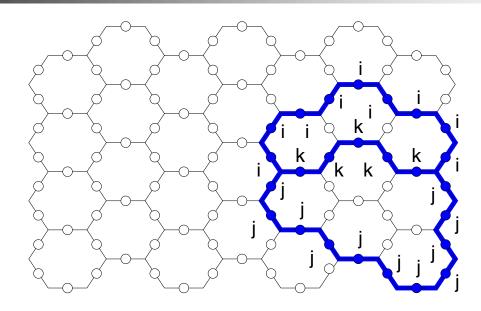
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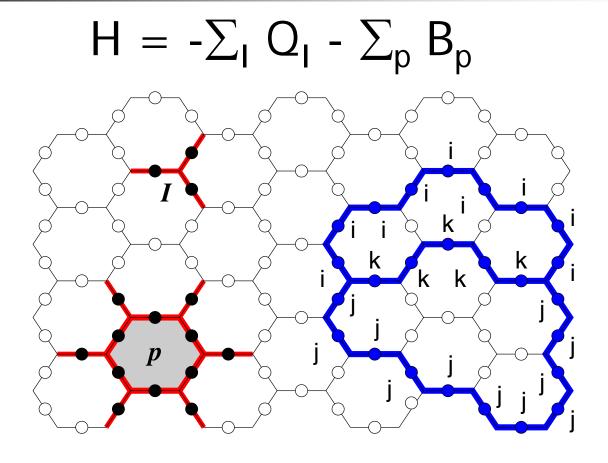
#### Exactly soluble lattice models



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#### Hamiltonians



Generalization of Kitaev's toric code

# First term: Q

### Defined by:

$$Q_{l} \mid \stackrel{\diamondsuit}{\underset{i \ j}{\Diamond}} \stackrel{k}{\rangle} = \delta_{ijk} \mid \stackrel{\diamondsuit}{\underset{i \ j}{\Diamond}} \stackrel{k}{\rangle}$$

# First term: Q

### Defined by:

$$Q_{l} \mid \stackrel{\diamondsuit k}{\underset{j}{\longleftrightarrow}} \rangle = \delta_{ijk} \mid \stackrel{\diamondsuit k}{\underset{j}{\longleftrightarrow}} \rangle$$

"Electric charge"

# Second term: B<sub>n</sub>

Defined by:  $B_p = \sum_s d_s B_p^s$  where

$$B_{p} \left| \begin{array}{c} b > h < c \\ g & i \\ a < j > d \end{array} \right| =$$

$$\sum_{g'h'\cdots l'} \mathsf{Falg}_{\mathsf{s}g'l'} \mathsf{Fbgh}_{\mathsf{s}h'g'} \cdots \mathsf{Ffkl}_{\mathsf{s}l'k'} \begin{vmatrix} \mathsf{b} \\ \mathsf{g'} \\ \mathsf{a} \\ \mathsf{j'} \\ \mathsf{b} \end{vmatrix}$$

$$\begin{vmatrix}
b & h' < c \\
g' & i' \\
a < & b < d \\
i' & j' \\
f > k' < e
\end{vmatrix}$$



# Second term: B<sub>n</sub>

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$$\begin{vmatrix}
b & h' < c \\
g' & i' \\
a < & b < d \\
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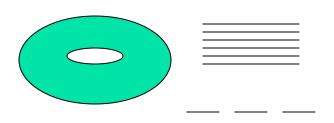
"Magnetic flux"

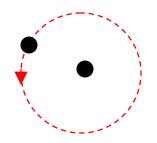


## Properties of Hamiltonian

- 1.  $\{B_p\}$ ,  $\{Q_l\}$  commuting projectors  $\Rightarrow$  H is exactly soluble.
- 2. Ground state wave function is  $\Phi$ .

3. Model describes a topological phase.





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ho}^{i heta}$ 



### Properties of Hamiltonian

- 4. Fixed points: Correlation length  $\xi = 0$ 
  - zero coupling gauge theory

"Right way" to put topological theories on lattice.



## Properties of Hamiltonian

- 4. Fixed points: Correlation length  $\xi$ = 0
  - zero coupling gauge theory

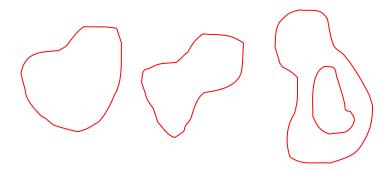
"Right way" to put topological theories on lattice.

Turaev/Viro (1992) Ooguri (1992) Loop quantum gravity: "spin networks"



## Example #1

- 1. String types: N = 1
- 2. Branching rules: No branching



What phase occurs when strings condense?

### Example #1

Two solutions to self-consistency equations:

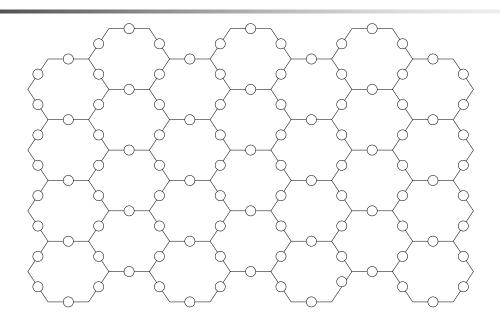
$$\begin{aligned} &d_0 = 1 \\ &d_1 = F^{110}_{110} = \pm 1 \\ &F^{000}_{000} = F^{101}_{101} = F^{011}_{011} = 1 \\ &F^{000}_{111} = F^{110}_{001} = F^{101}_{010} = F^{011}_{100} = 1 \end{aligned}$$

Two sets of local rules:

$$\Phi(\bigcirc) = \pm \Phi(\bigcirc)$$
  
 $\Phi(\bigcirc) = \pm \Phi(\bigcirc)$ 

Two solutions:  $\Phi_{+}(X) = (\pm 1)^{N_{loops}(X)}$ 

### Lattice realization



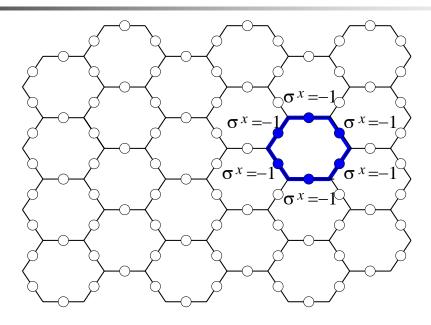
Each "spin" can be in 2 states:  $|0\rangle$ ,  $|1\rangle$ 

Convenient to use spin-1/2 notation:

$$|0\rangle = |\sigma^x = +1\rangle$$

$$|1\rangle = |\sigma^x = -1\rangle$$

### Lattice realization



Each "spin" can be in 2 states:  $|0\rangle$ ,  $|1\rangle$ 

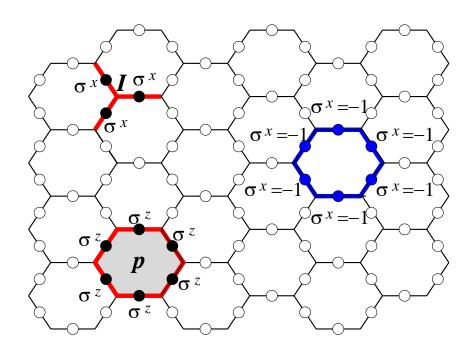
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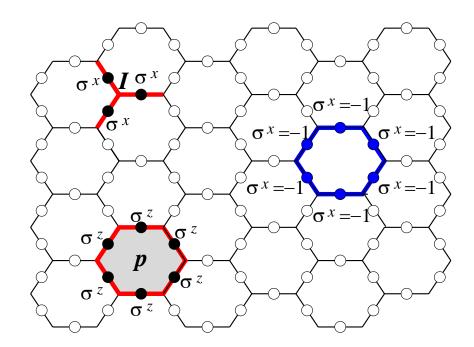
# Hamiltonian: $\Phi_+$

$$H_{+} = -\sum_{I} \prod_{a} \sigma^{x}_{a} - \sum_{p} \prod_{b} \sigma^{z}_{b}$$



## Hamiltonian: $\Phi_{+}$

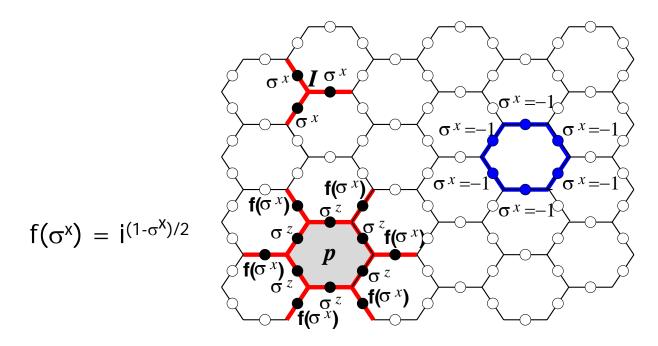
$$H_{+} = -\sum_{I} \prod_{a} \sigma^{x}_{a} - \sum_{p} \prod_{b} \sigma^{z}_{b}$$



Toric code: Lattice model for Z<sub>2</sub> gauge theory!

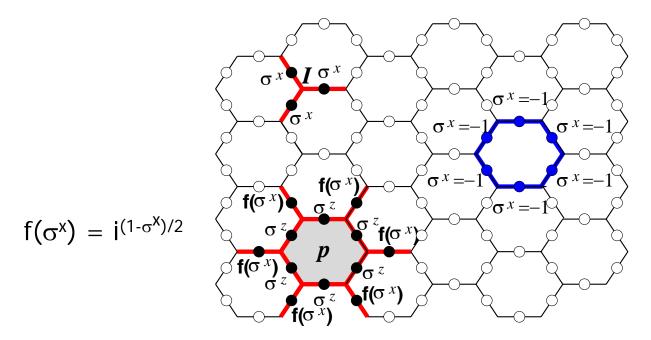
### Hamiltonian: Φ<sub>-</sub>

$$H_{-} = -\sum_{i} \prod_{a} \sigma_{a}^{x} - \sum_{p} \prod_{b} \sigma_{b}^{z} \cdot \prod_{c} i^{(1-\sigma_{c}^{x})/2}$$



### Hamiltonian: Φ<sub>-</sub>

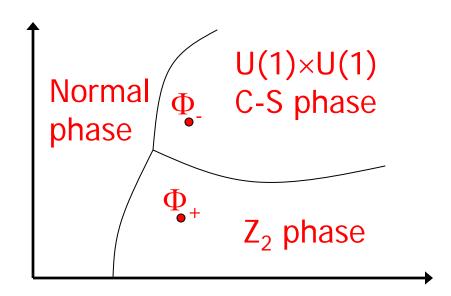
$$H_{-} = -\sum_{i} \prod_{a} \sigma_{a}^{x} - \sum_{p} \prod_{b} \sigma_{b}^{z} \cdot \prod_{c} i^{(1-\sigma_{c}^{x})/2}$$



U(1)×U(1) Chern-Simons theory with semions!



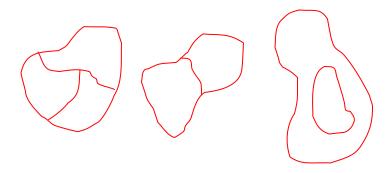
### Two string condensed phases





## Example #2

- 1. String types: N = 1
- 2. Branching rules: {1,1,1}



What phase occurs when string-nets condense?

### Example #2

Only one set of self-consistent local rules:

$$\Phi(\bigcirc) = \tau \Phi(\bigcirc)$$

$$\Phi(\bigcirc) = 0$$

$$\Phi(\bigcirc) = \tau^{-1} \Phi(\bigcirc) + \tau^{-1/2} \Phi(\bigcirc)$$

$$\Phi(\bigcirc) = \tau^{-1/2} \Phi(\bigcirc) - \tau^{-1} \Phi(\bigcirc)$$

$$\tau = (1+5^{1/2})/2$$

# Example #2

Wave function: No closed form!

Hamiltonian: Spin-1/2 model (complicated)

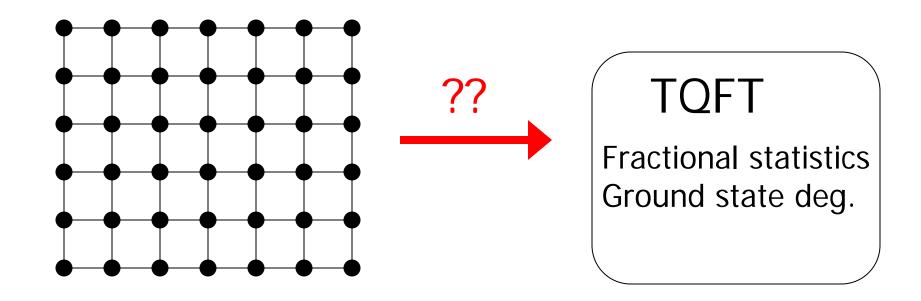
Topological phase:  $SO_3(3) \times SO_3(3)$  Chern-Simons theory

- "Fibonacci theory"
- Non-abelian anyons



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How can we realize them? What interactions favor them?

