

Heavy-ion phenomenology vs causality, attractors, non-hydro modes et al

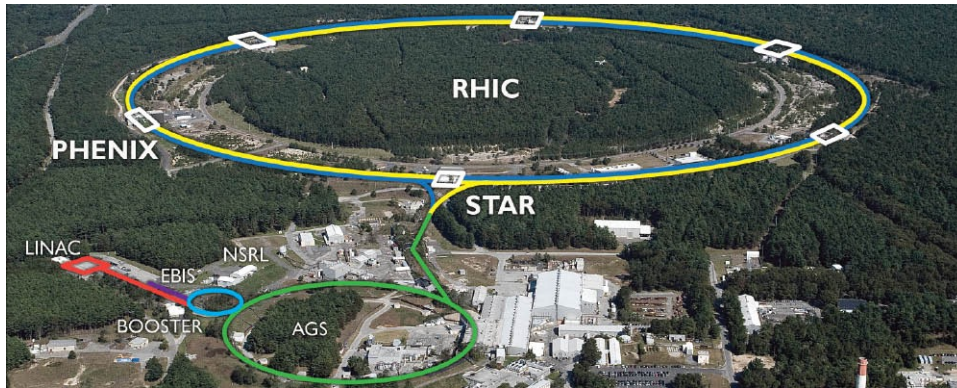
Jean-François Paquet

June 21, 2023



The Many Faces of Relativistic Fluid
Dynamics
Kavli Institute for Theoretical Physics
UC Santa Barbara

Relativistic Heavy Ion Collider (RHIC) [Brookhaven National Lab, Long Island, NY]



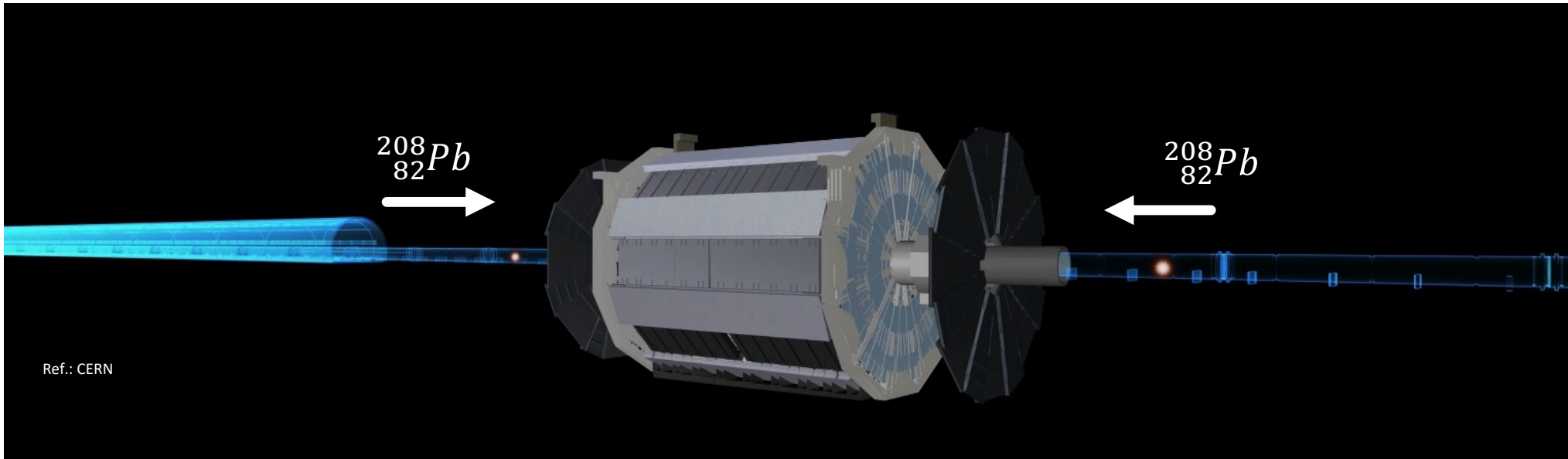
Large Hadron Collider (LHC) [CERN, Geneva, Switzerland/France]



WHAT IS HEAVY-ION PHENOMENOLOGY

Relativistic Heavy Ion Collider (RHIC)
[Brookhaven National Lab, Long Island, NY]

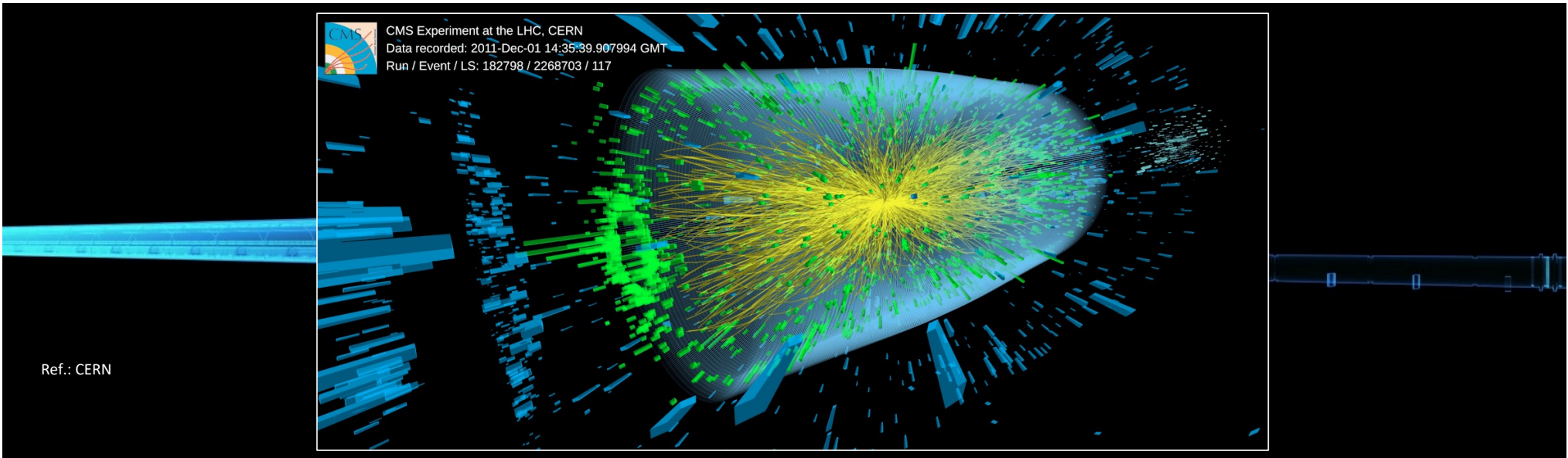
Large Hadron Collider (LHC)
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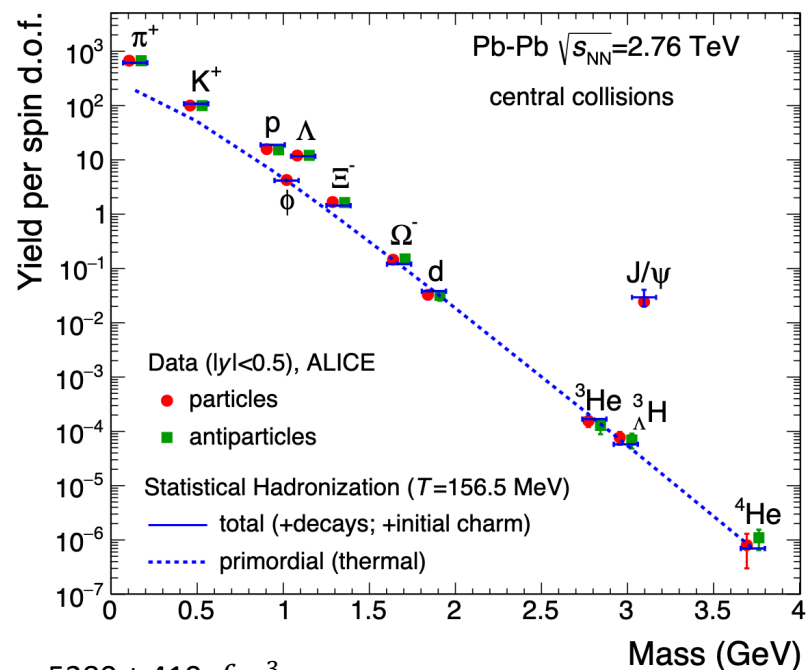
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WHAT IS HEAVY-ION PHENOMENOLOGY

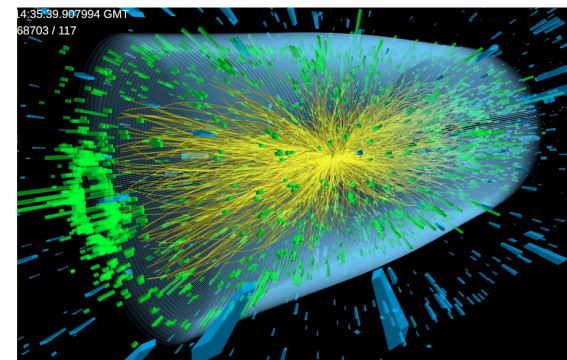
Highlights from measurements

Multiplicity of hadrons



Volume = 5280 ± 410 fm³

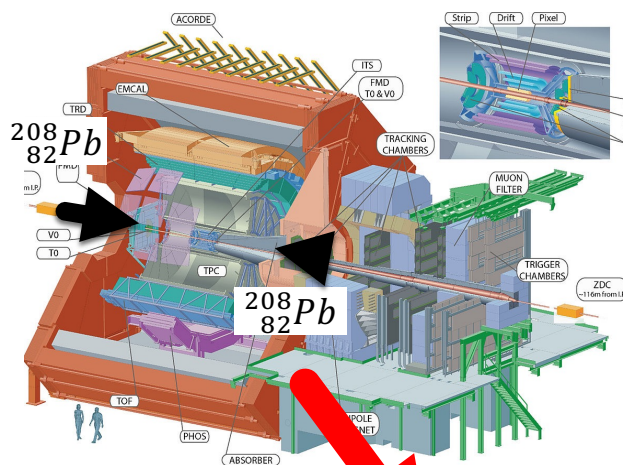
Ref.: Andronic, Braun-Munzinger, Redlich, Stachel (2017) Nature



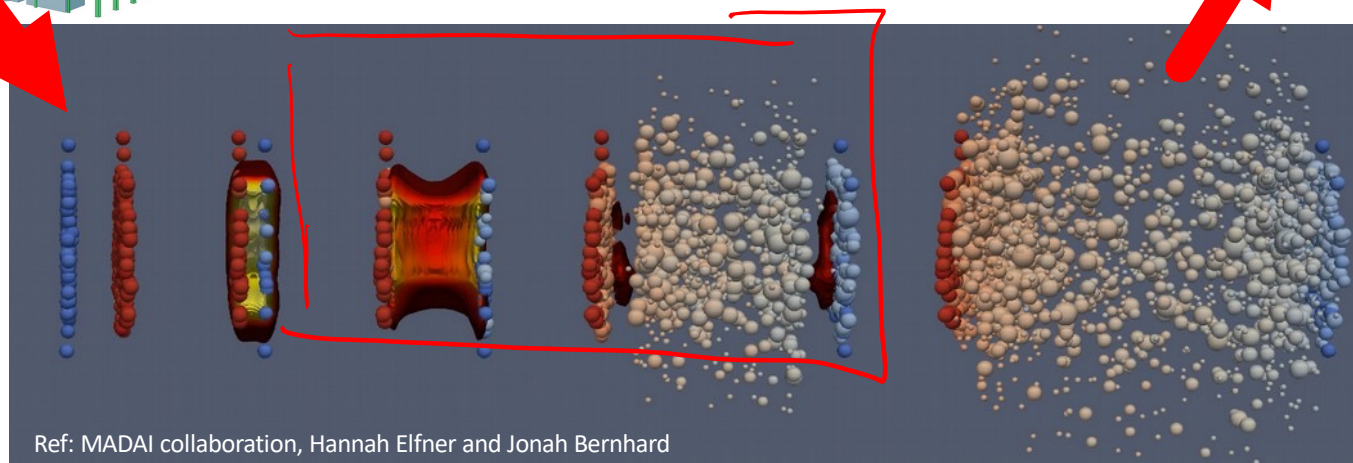
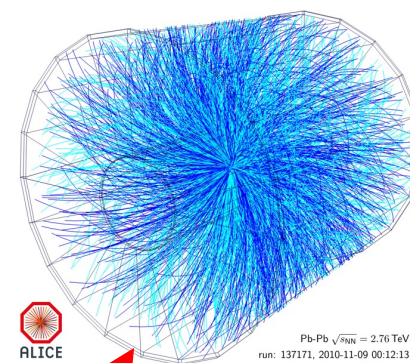
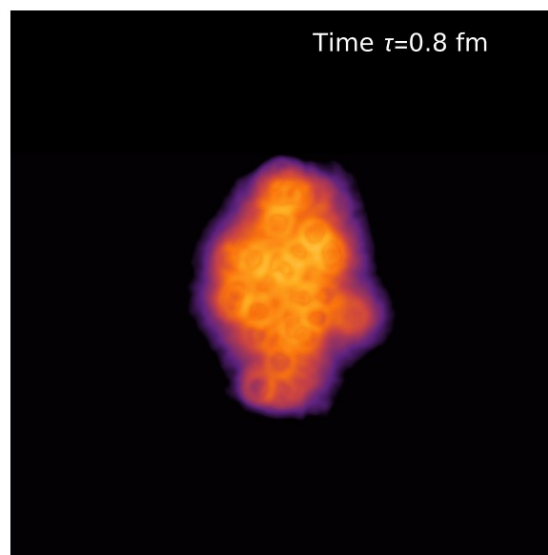
High transverse momentum photon/Z/W[±]/jets

$$d\sigma_{AA \rightarrow \gamma/Z/W} = N_{\text{bin}} d\sigma_{pp \rightarrow \gamma/Z/W}$$

Heavy-ion collisions

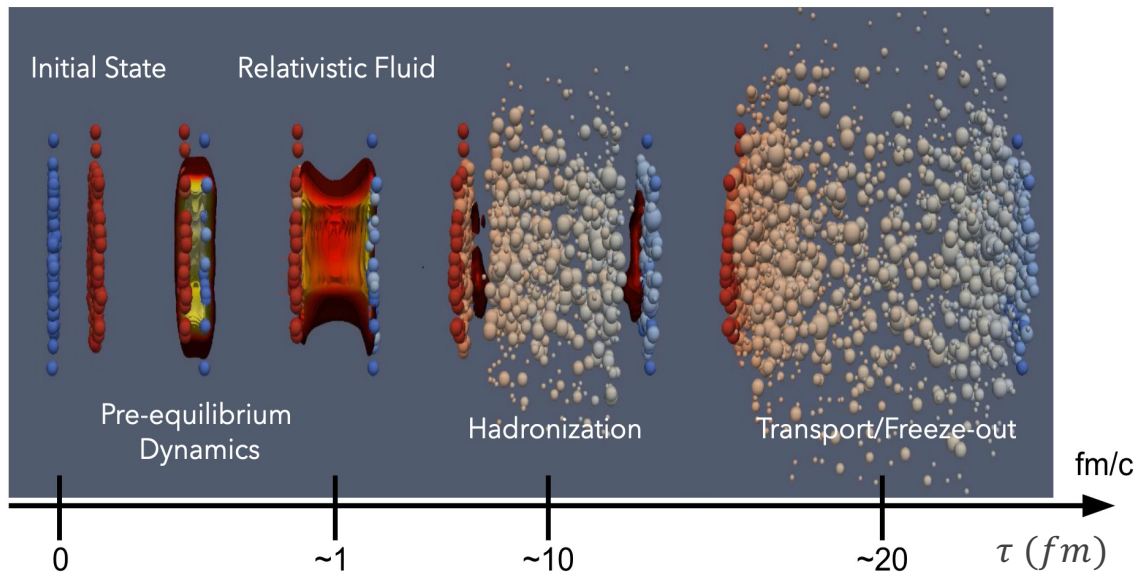


Ref.: ALICE, CERN



JEAN-FRANÇOIS PAQUET (VANDERBILT)

Ref: MADAI collaboration, Hannah Elfner and Jonah Bernhard

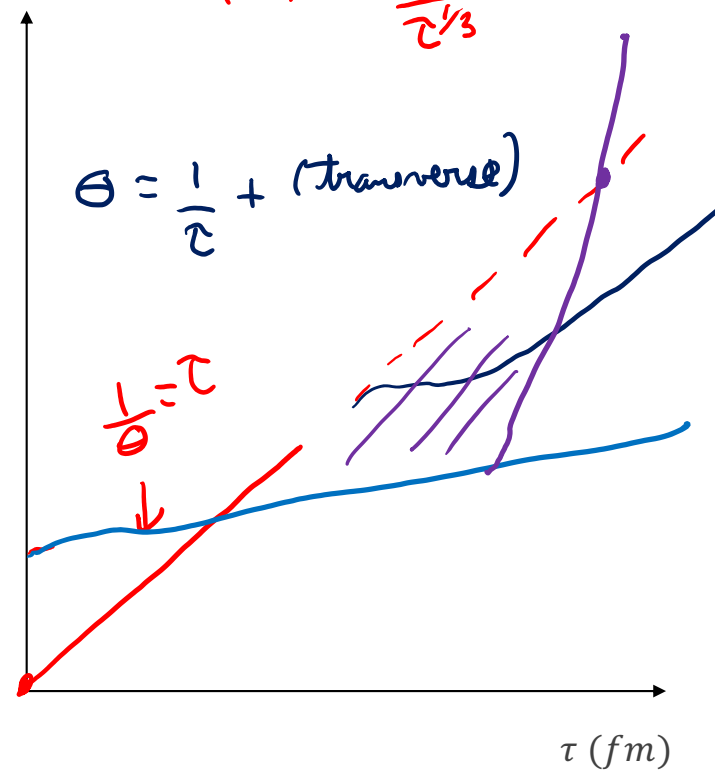


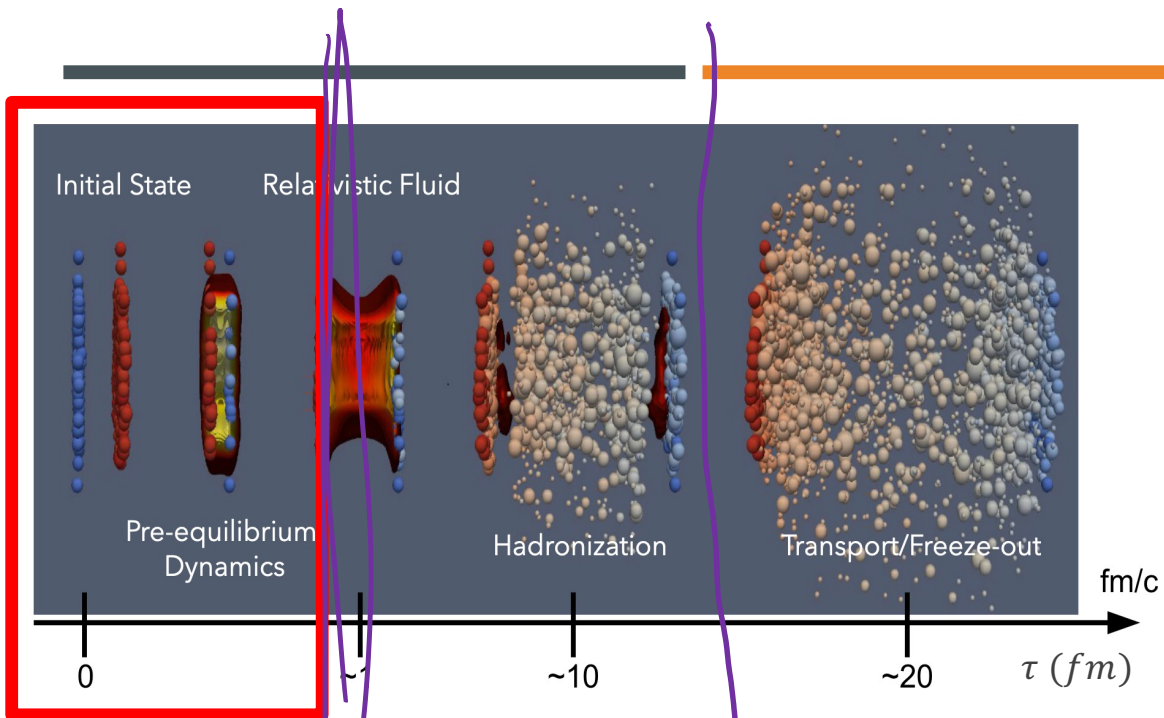
Hydro: $\tau_{\text{relax}} \ll \frac{1}{\text{expansion rate}}$

$\Theta = \partial_{\mu} u^{\mu} = \frac{1}{\tau}$; $\tau = \sqrt{t^2 - z^2}$

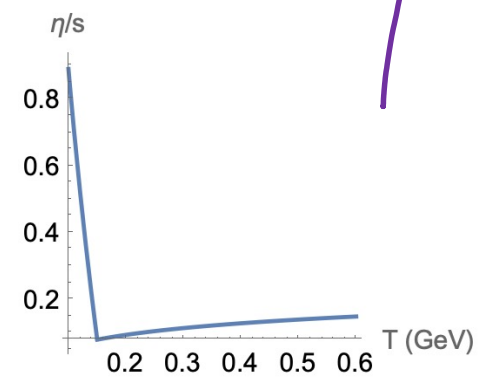
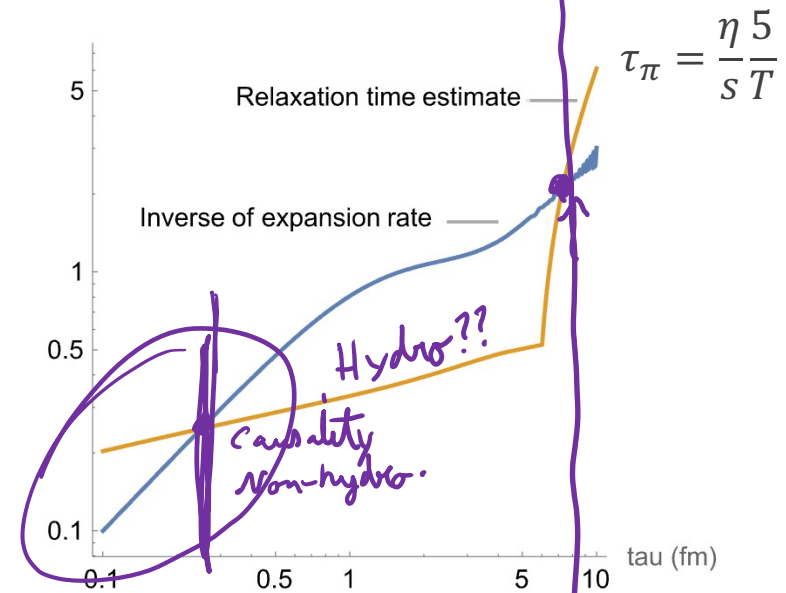
$$\tau_{\pi} = \frac{\eta}{s} \frac{1}{T} \quad (3-5-10?) \approx \frac{\eta}{s} \tau^{1/3} \neq$$

$$\uparrow T \propto \frac{1}{\tau^{1/3}}$$



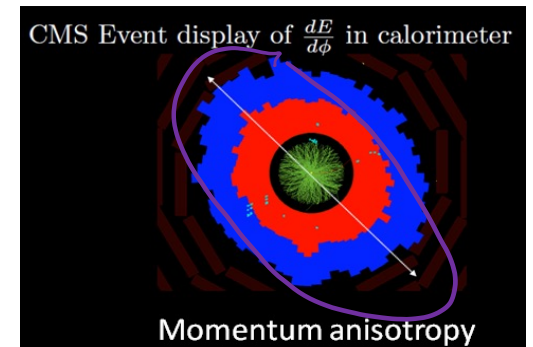
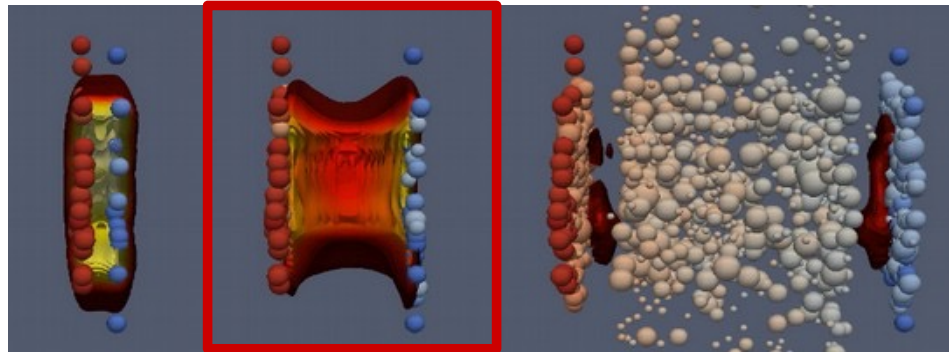
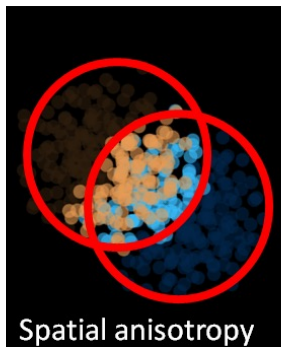


- Causality?
- Attractors?
- Non-hydrodynamic modes & transient hydrodynamics?



Multistage simulations of heavy ion collisions

Based on figures by Derek Teaney, CMS Coll., MADAI, H. Elfner and J. Bernhard



Energy deposition

Early dynamics

Hydrodynamics

Hadronic transport

- Energy-momentum tensor: $T^{\mu\nu} = \epsilon u^\mu u^\nu - (P(\epsilon) + \Pi)(g^{\mu\nu} - u^\mu u^\nu) + \pi^{\mu\nu}$
- Conservation of energy/momentum: $\partial_\nu T^{\mu\nu} = 0$
- DNMR transient relativistic viscous hydrodynamics

$$\tau_\pi \Delta_{\alpha\beta}^{\mu\nu} \dot{\pi}^{\alpha\beta} + \boxed{\pi^{\mu\nu} = 2\eta(T)(\partial^\mu u^\nu + \dots)} + (2^{\text{nd}} \text{ order}); \quad \tau_\Pi \dot{\Pi} + \boxed{\Pi = -\zeta(T) \partial_\mu u^\mu} + (2^{\text{nd}} \text{ order});$$

DNR

- Energy-momentum tensor:

$$T^{\mu\nu} = \underline{\epsilon} u^\mu u^\nu - (\underline{P(\epsilon)} + \underline{\Pi}) \Delta^{\mu\nu} + \underline{\pi^{\mu\nu}}$$

- Conservation of energy and momentum:

$$\partial_\nu T^{\mu\nu} = 0$$

- Bulk pressure Π equation of motion

$$\underline{\tau_\Pi} \dot{\Pi} + \Pi = \underline{-\zeta} \theta - \underline{\delta_{\Pi\Pi}} \Pi \theta + \underline{\lambda_{\Pi\pi}} \pi^{\mu\nu} \sigma_{\mu\nu}$$

- Shear tensor $\pi^{\mu\nu}$ equation of motion:

$$\tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = 2\eta \sigma^{\mu\nu} - \delta_{\pi\pi} \pi^{\mu\nu} \theta + \phi_7 \pi_\alpha^{\langle\mu} \pi^{\nu\rangle\alpha} - \tau_{\pi\pi} \pi_\alpha^{\langle\mu} \sigma^{\nu\rangle\alpha} + \lambda_{\pi\Pi} \Pi \sigma^{\mu\nu}$$

- Definition:

$$\theta = \partial_\mu u^\mu; \Delta^{\mu\nu} = (\text{tensor perpendicular to } u^\mu);$$

$$\Delta_{\alpha\beta}^{\mu\nu} = (\text{symmetric traceless projection operator}); \sigma^{\mu\nu} = \Delta_{\alpha\beta}^{\mu\nu} \partial^\alpha u^\beta$$

$$T^{\mu\nu} = \epsilon u^\mu u^\nu - (P(\epsilon) + \Pi)\Delta^{\mu\nu} + \pi^{\mu\nu}; \partial_\nu T^{\mu\nu} = 0$$

- Bulk pressure Π equation of motion

$$\tau_\Pi \dot{\Pi} + \Pi = -\zeta\theta - \delta_{\Pi\Pi}\Pi\theta + \lambda_{\Pi\pi}\pi^{\mu\nu}\sigma_{\mu\nu}$$

- Shear tensor $\sigma^{\mu\nu}$ equation of motion:

$$\tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi = 2\eta\sigma^{\mu\nu} - \delta_{\pi\pi}\pi^{\mu\nu}\theta + \phi_7\pi_\alpha^{\langle\mu}\pi^{\nu\rangle\alpha} - \tau_{\pi\pi}\pi_\alpha^{\langle\mu}\sigma^{\nu\rangle\alpha} + \lambda_{\pi\Pi}\Pi\sigma^{\mu\nu}$$

$$\phi_7 = \frac{9}{70\mathcal{P}}$$

$$\tau_\Pi = \frac{\zeta}{15\left(\frac{1}{3} - c_s^2\right)^2(\epsilon + \mathcal{P})}$$

$$\delta_{\Pi\Pi} = \frac{2}{3}\tau_\Pi$$

$$\lambda_{\Pi\pi} = \frac{8}{5}\left(\frac{1}{3} - c_s^2\right)\tau_\Pi$$

$$\tau_\pi = \frac{5\eta}{(\epsilon + \mathcal{P})}$$

$$\delta_{\pi\pi} = \frac{4}{3}\tau_\pi$$

$$\tau_{\pi\pi} = \frac{10}{7}\tau_\pi$$

$$\lambda_{\pi\Pi} = \frac{6}{5}$$

$$T^{\mu\nu} = \epsilon u^\mu u^\nu - (P(\epsilon) + \Pi)\Delta^{\mu\nu} + \pi^{\mu\nu}; \partial_\nu T^{\mu\nu} = 0$$

- Bulk pressure Π equation of motion

$$\tau_\Pi \dot{\Pi} + \Pi = -\zeta\theta - \delta_{\Pi\Pi}\Pi\theta + \lambda_{\Pi\pi}\pi^{\mu\nu}\sigma_{\mu\nu}$$

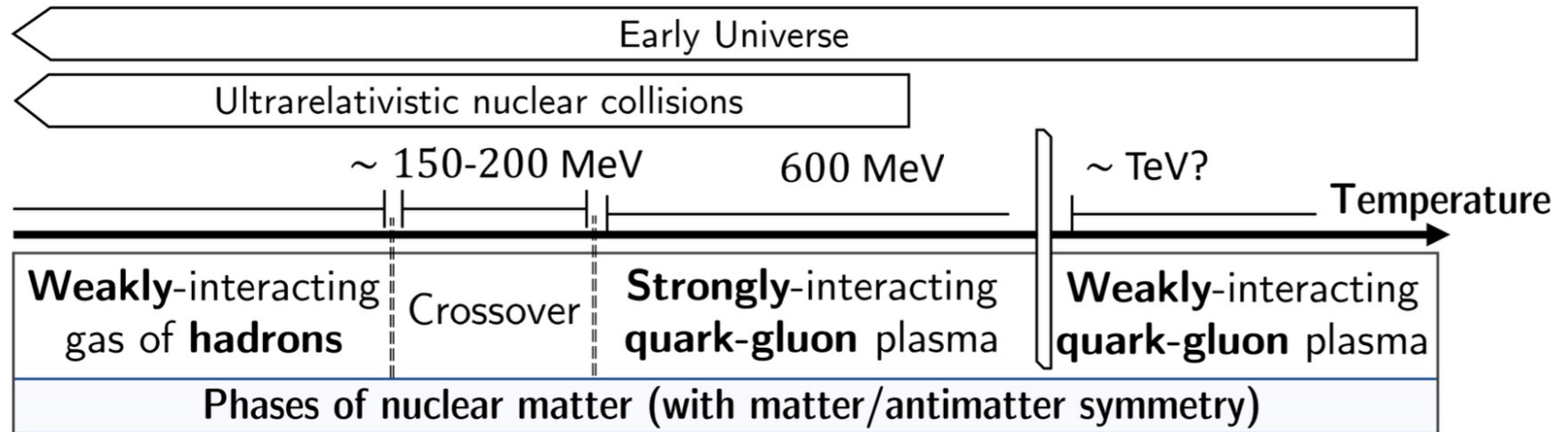
- Shear tensor $\sigma^{\mu\nu}$ equation of motion:

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$$\tau_\pi = \frac{\eta}{5} \frac{1}{T} \frac{1}{2(2-\ln 2)}$$

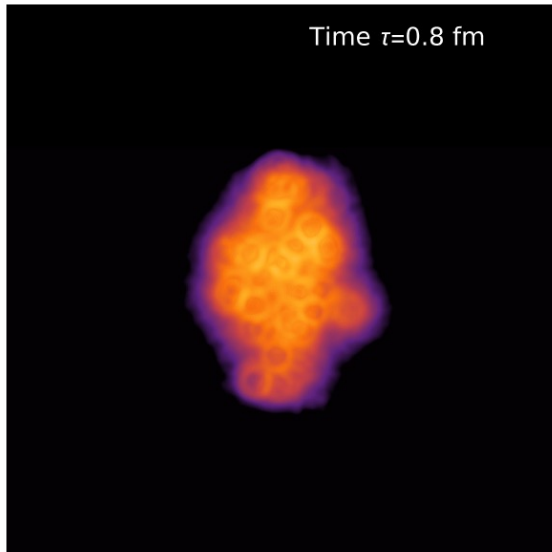
	Gauge/Gravity	Kinetic (BGK)	pQCD
$\epsilon(P)$	3 P	$\frac{\Delta T}{(\epsilon+P)\tau_R}$	3 P
η	$\frac{\epsilon+P}{4\pi T}$ $\frac{\eta}{5} = \frac{1}{4\pi}$	$\frac{5}{\tau_R}$	$\frac{3.85(\epsilon+P)}{g^4 \ln(2.765g^{-1})T}$
τ_π	$\frac{1}{4\pi} \frac{2-\ln 2}{2T}$	τ_R	$\frac{5.9\eta}{\epsilon+P}$

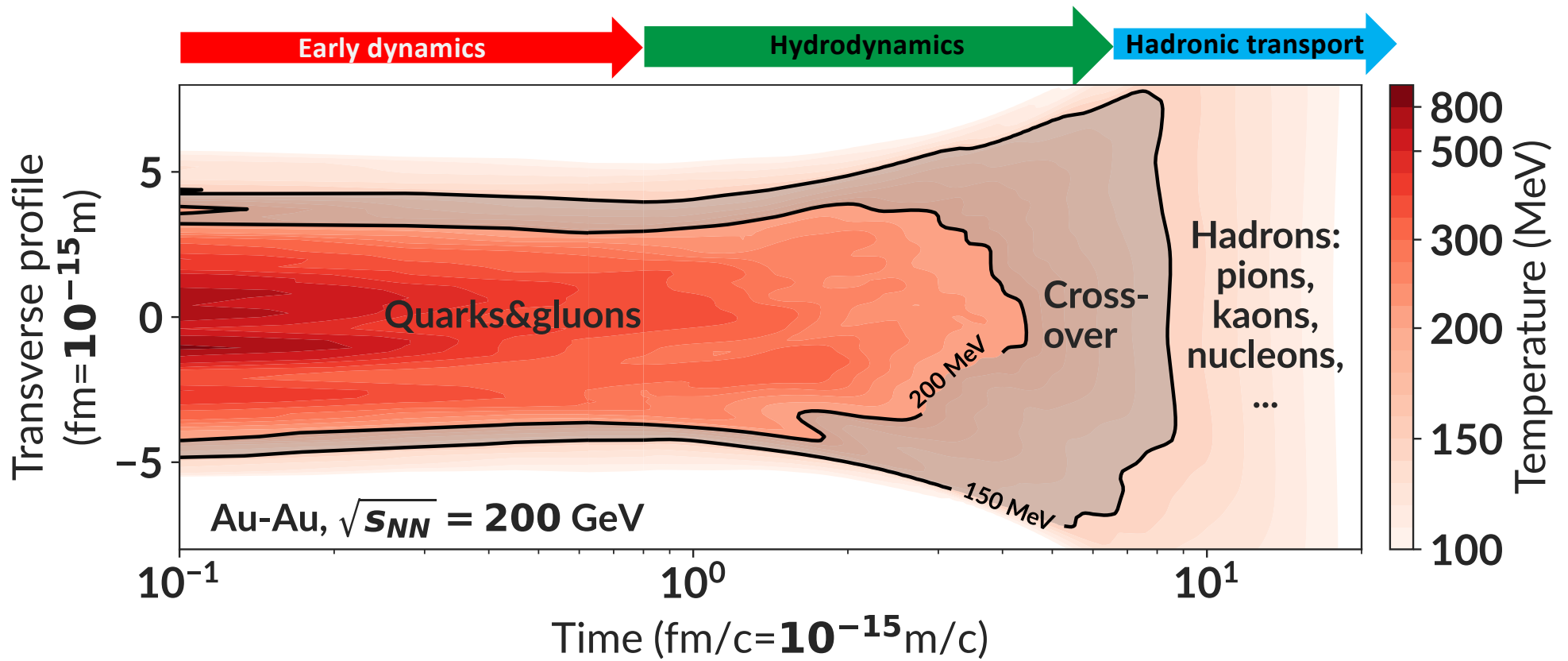
Ref.: Romatschke&Romatschke



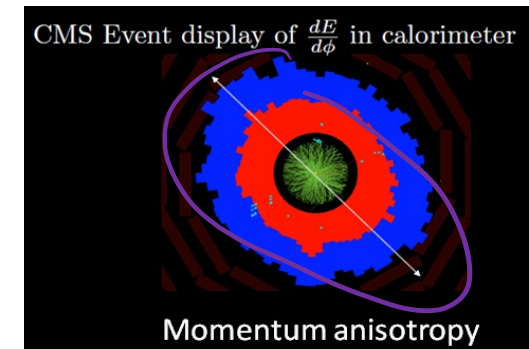
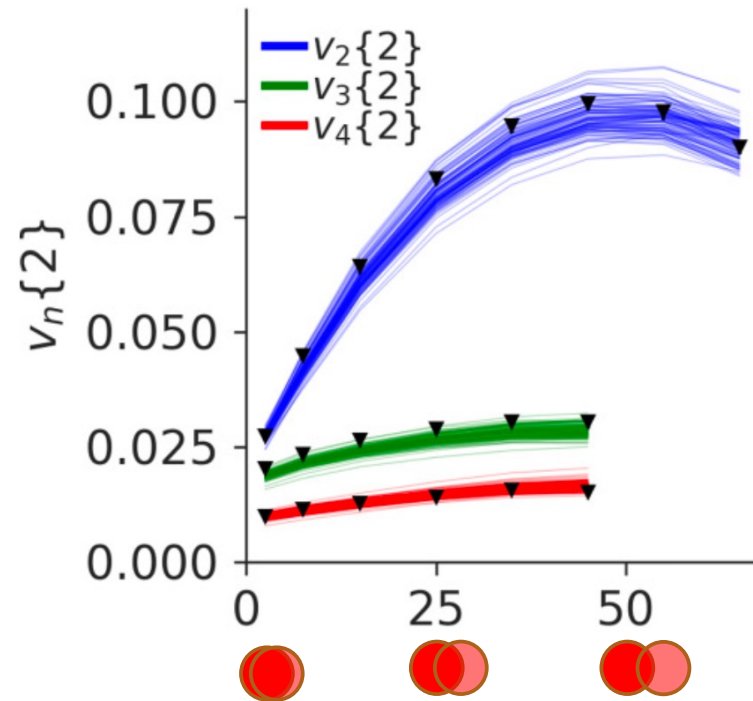
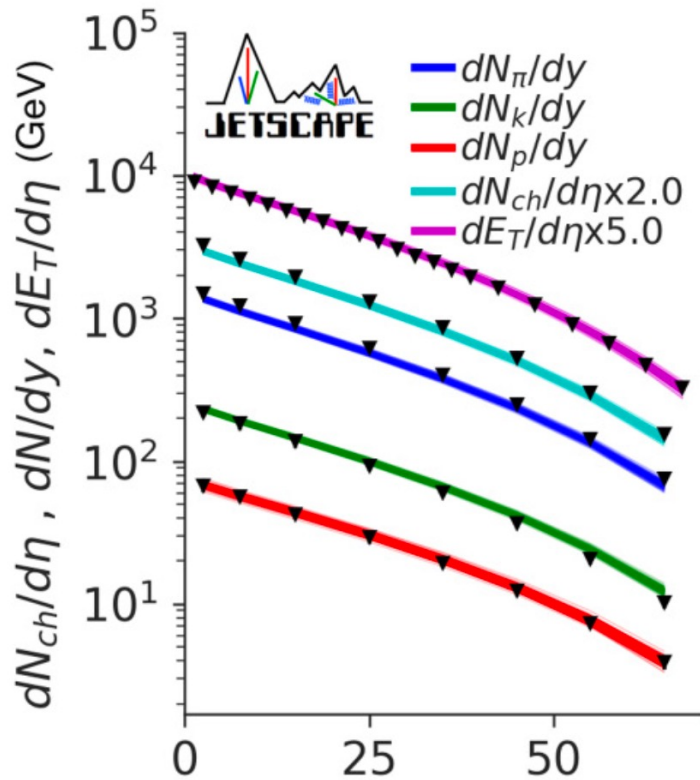
	Gauge/Gravity	Kinetic (BGK)	pQCD
$\epsilon(P)$	$3 P$		$3 P$
η	$\frac{\epsilon+P}{4\pi T}$	$\frac{(\epsilon+P)\tau_R}{5}$	$\frac{3.85(\epsilon+P)}{g^4 \ln(2.765g^{-1})T}$
τ_π	$\frac{2-\ln 2}{2\pi T}$	τ_R	$\frac{5.9\eta}{\epsilon+P}$

Ref.: Romatschke&Romatschke

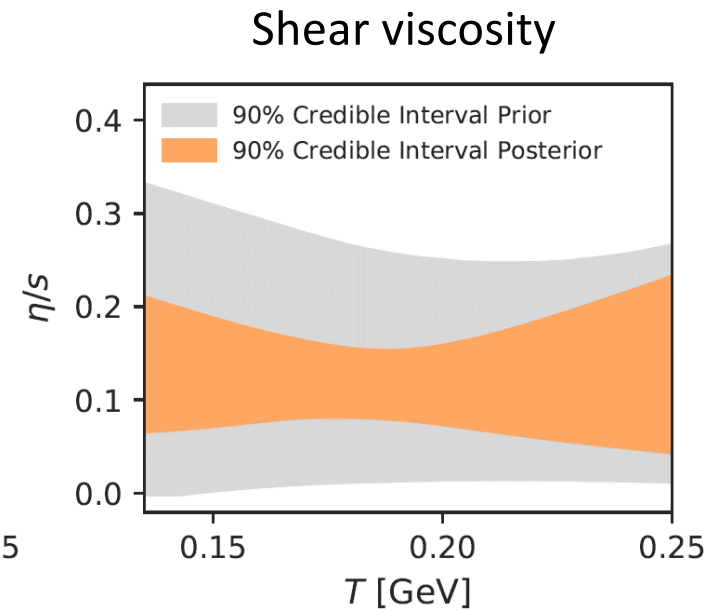
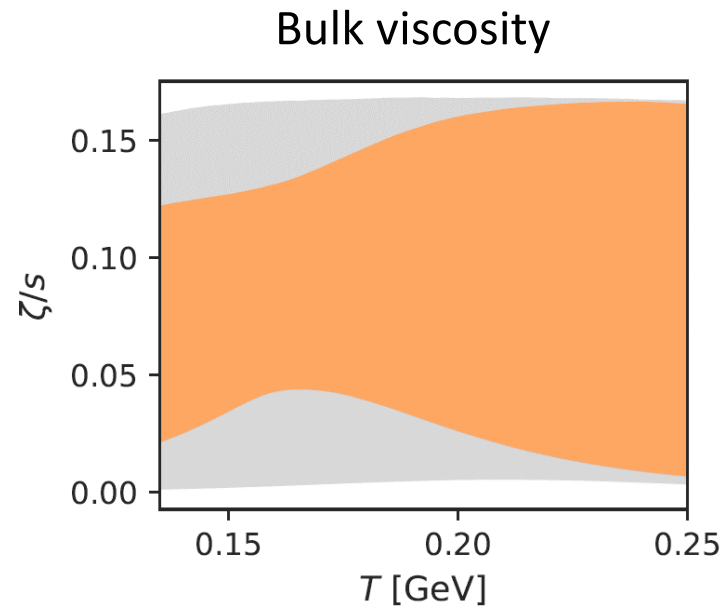
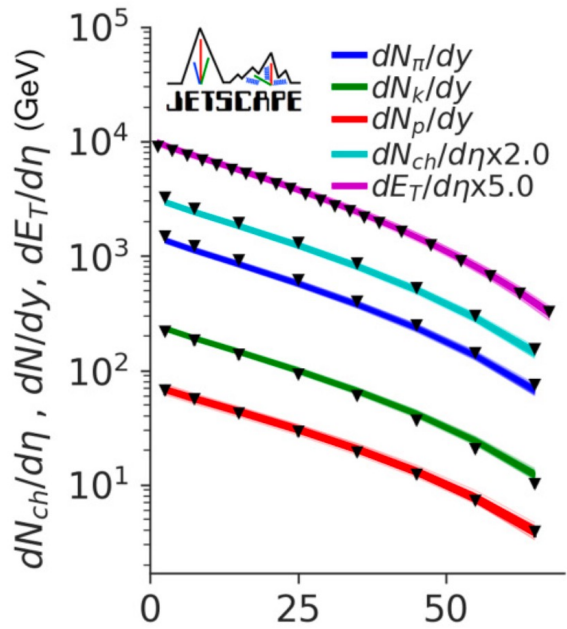




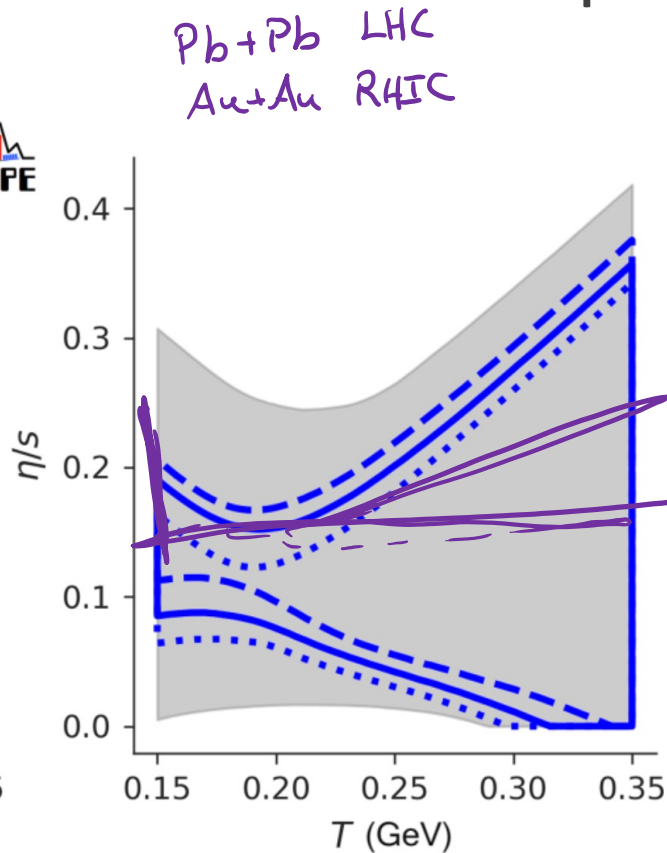
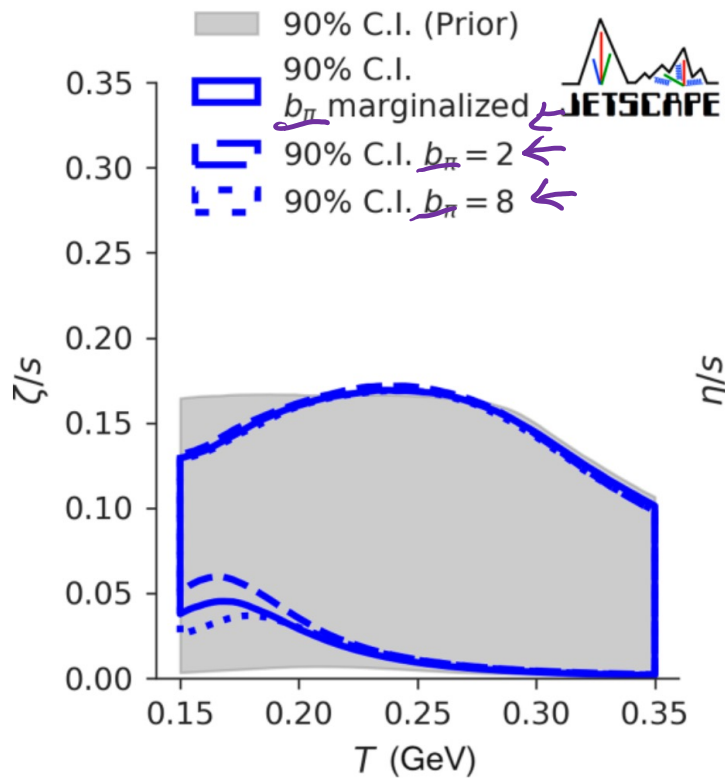
Comparison with data and constraints on viscosity



Comparison with data and constraints on viscosity

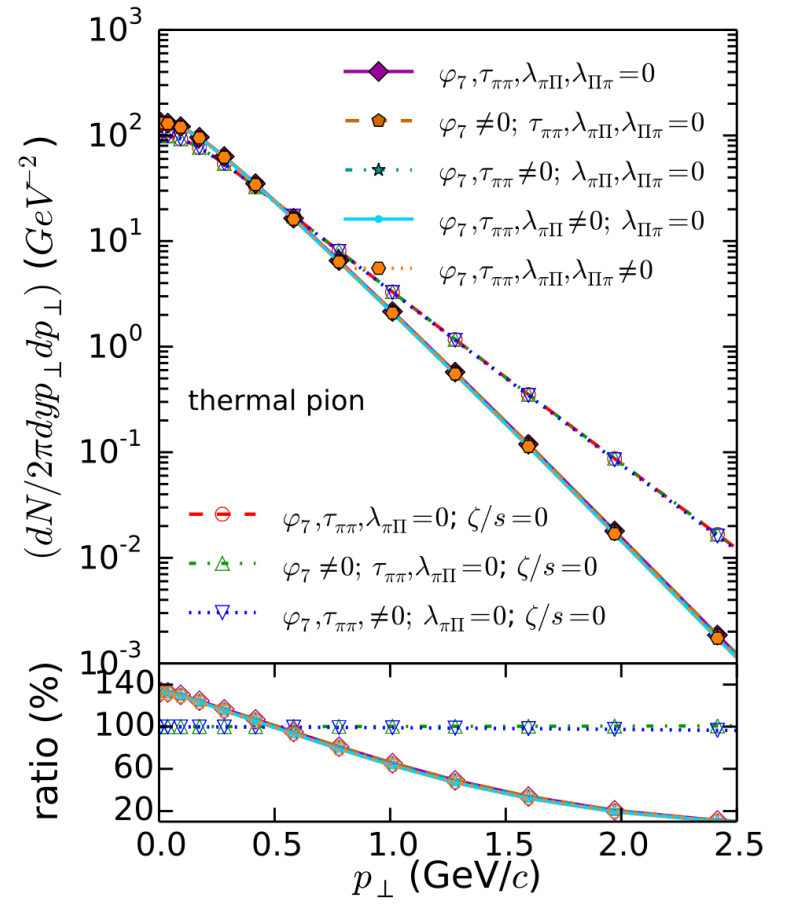


Dependence on second-order transport coefficients

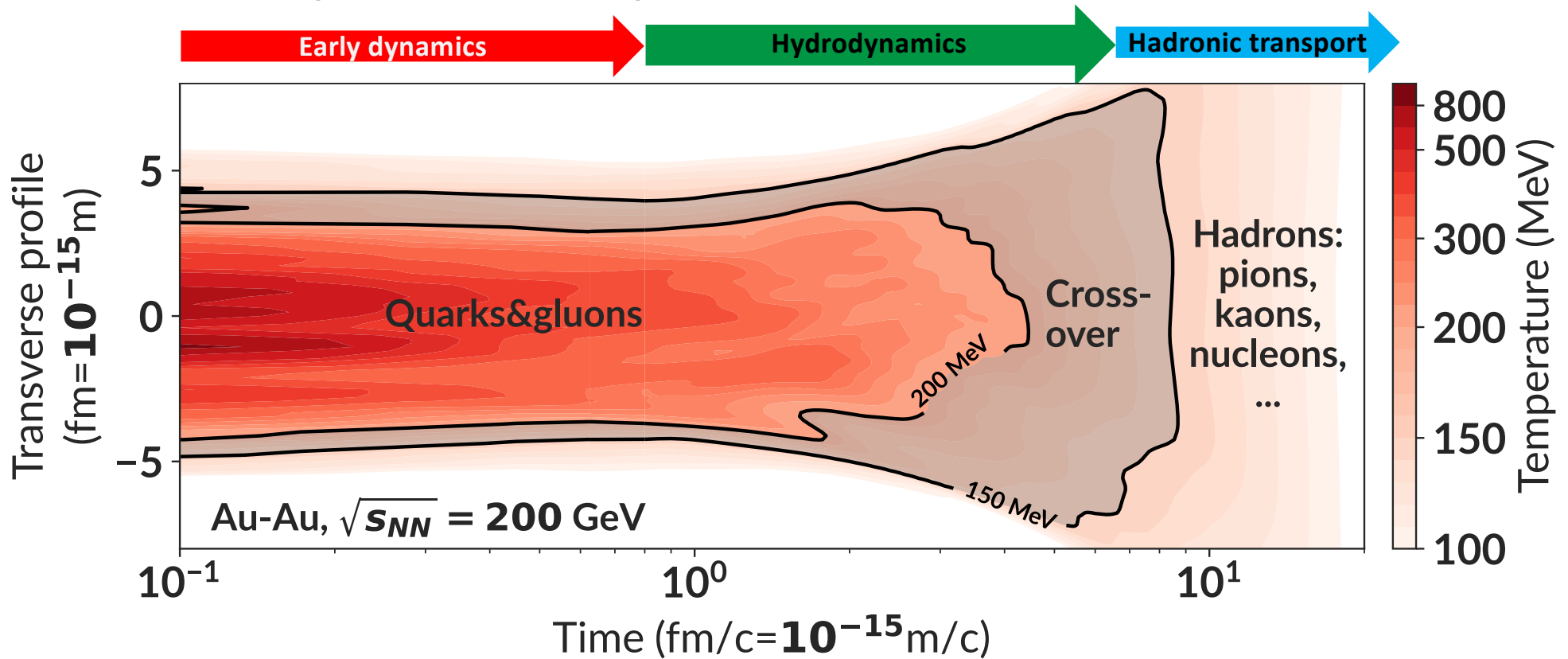


$$\begin{aligned}
 & \tau_\pi = \frac{\eta}{s} \frac{1}{T} b_\pi \quad \Leftarrow \\
 & \downarrow \\
 & \tau_\pi \dot{\pi}^{\langle \mu \nu \rangle} + \pi \\
 & = 2\eta \sigma^{\mu \nu} \\
 & - \delta_{\pi\pi} \pi^{\mu \nu} \theta \\
 & + \phi_7 \pi_\alpha^{\langle \mu} \pi^{\nu \rangle \alpha} \\
 & - \tau_{\pi\pi} \pi_\alpha^{\langle \mu} \sigma^{\nu \rangle \alpha} \\
 & + \lambda_{\pi\Pi} \Pi \sigma^{\mu \nu}
 \end{aligned}$$

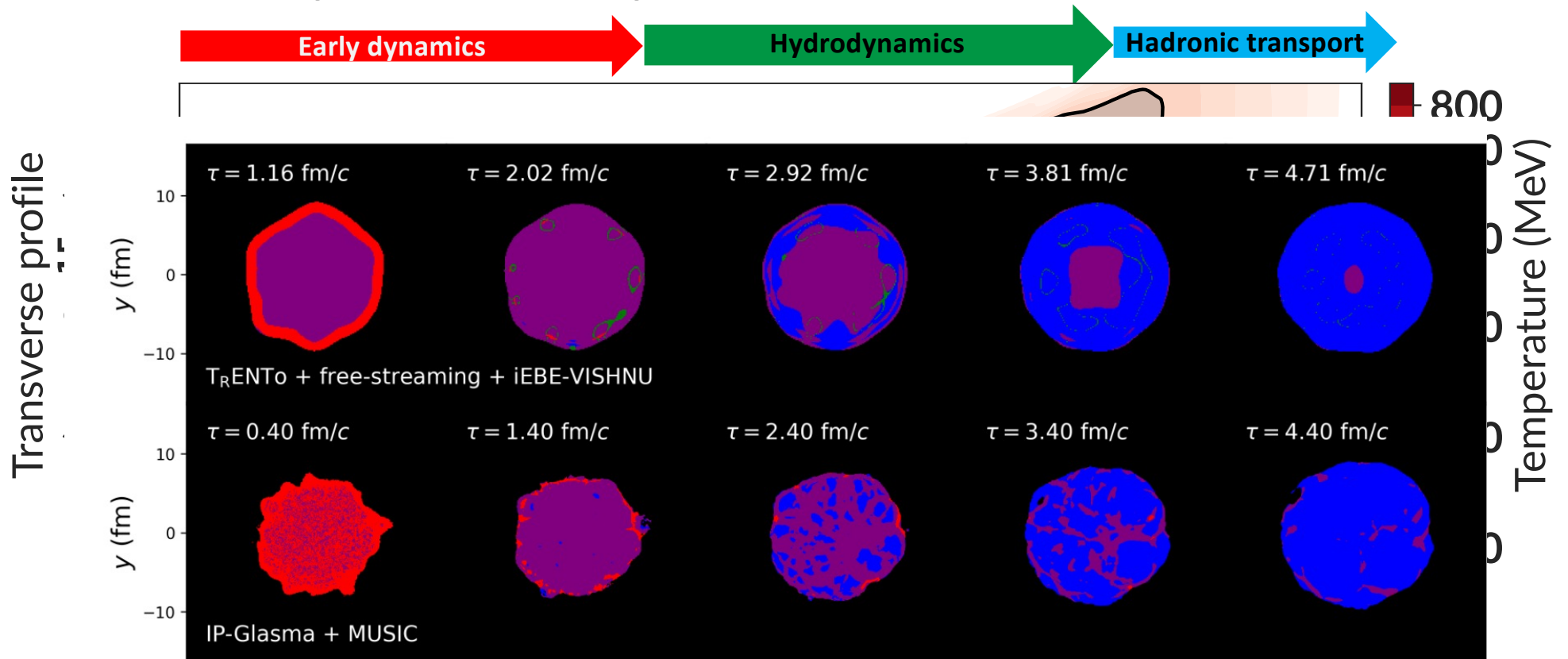
■ Jia Liu's PhD thesis



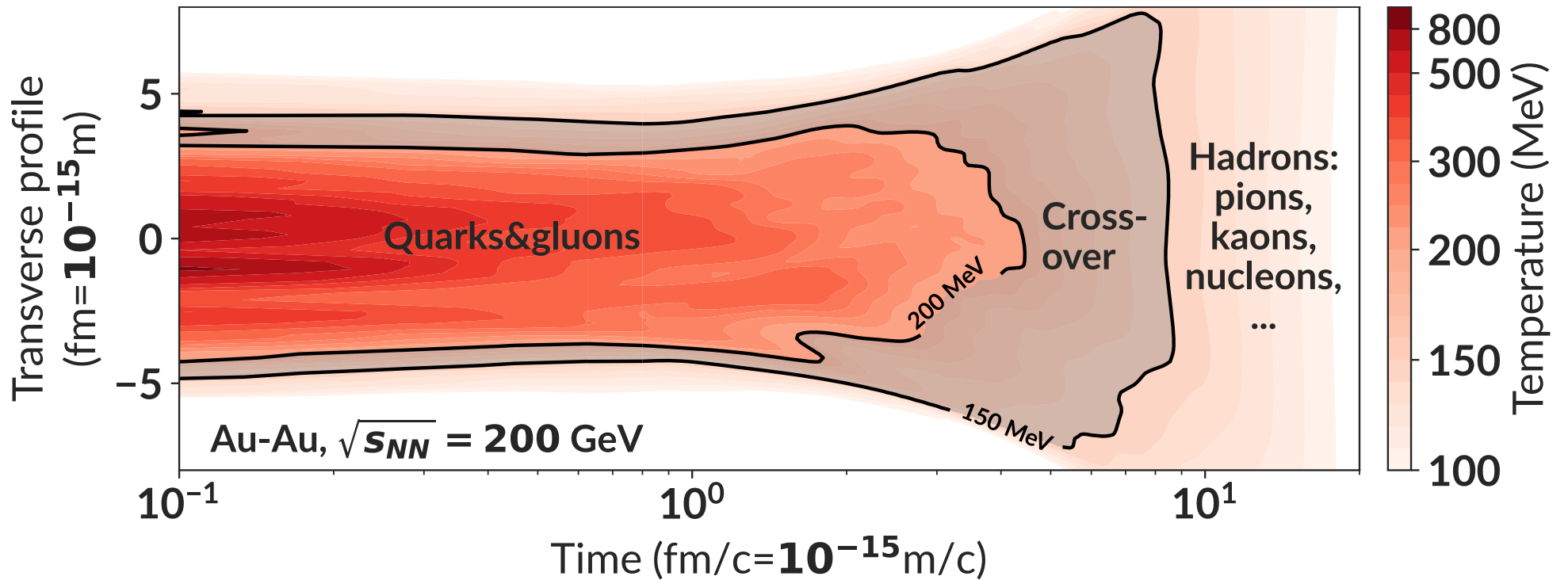
Acausality (c.f. Dekrayat's talk)



Acausality (c.f. Dekrayat's talk)

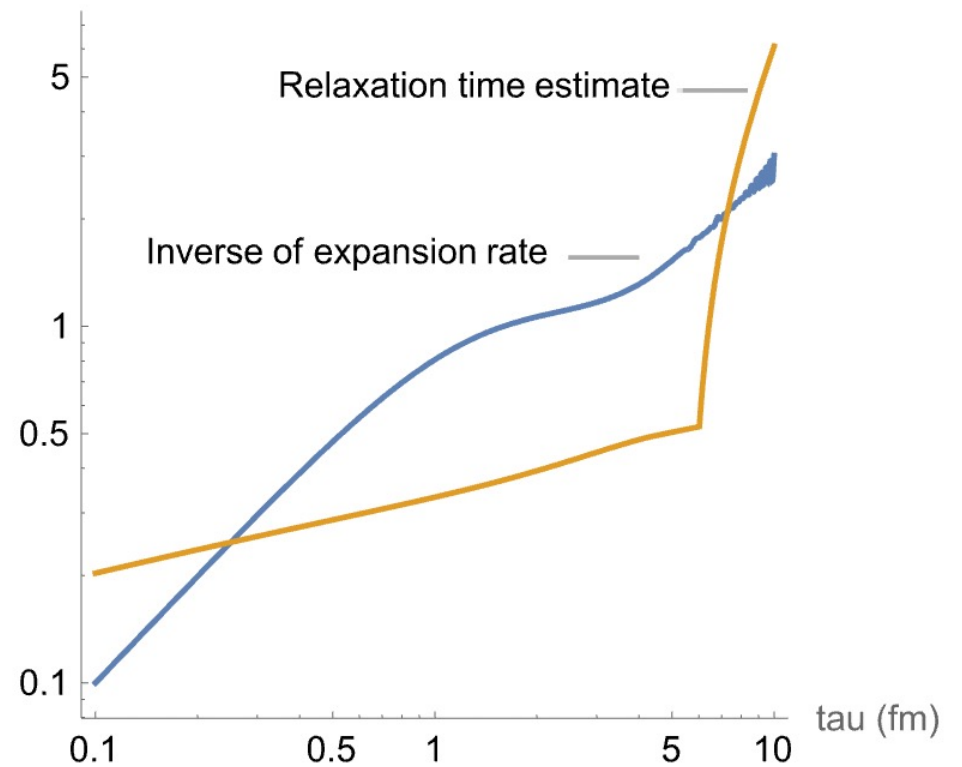


Non-hydro modes + acausality = problem?



Attractors and phenomenology

- Expansion-dominated phase ("free-streaming")
- Approach to hydrodynamics



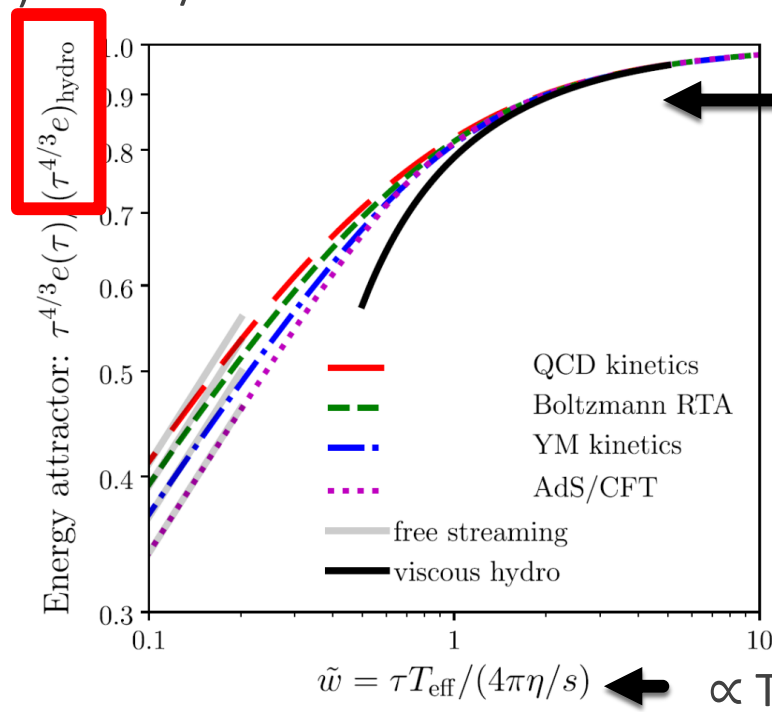
Matching to 0+1D (boost-invariant) hydrodynamics

- In 0+1D hydro, $T^{\mu\nu}$ characterized by single component: energy density
- 0+1D dynamical models with smooth transition to hydrodynamics:
 - Kinetic theory (gluons, QCD, RTA) or AdS/CFT

(note: not all systems have the same final energy density, but this can be rescaled)

$$\text{Free-streaming: } \propto \tilde{w}^{4/9}$$

$$\text{Navier-Stokes: } 1 - \frac{2}{3\pi} \frac{1}{\tilde{w}}$$



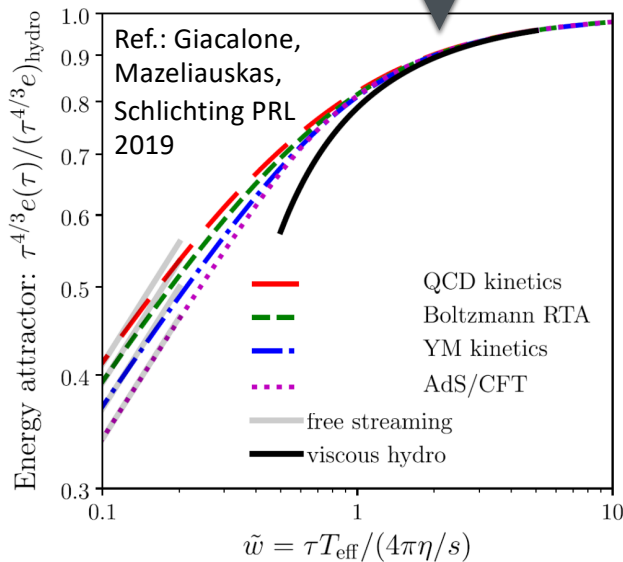
All systems converge to 0+1D Navier-Stokes

Ref.: Giacalone, Mazeliauskas, Schlichting PRL 2019

Matching to 2+1D hydrodynamics: the “KøMPøST” approach

- Better approximation [KøMPøST]: decompose $T^{\mu\nu}$ in 0+1D background + linear perturbation

$$T^{\mu\nu}(\tau_{\text{hydro}}, \mathbf{x}) = \overline{T}_{\mathbf{x}}^{\mu\nu}(\tau_{\text{hydro}}) + \frac{\overline{T}_{\mathbf{x}}^{\tau\tau}(\tau_{\text{hydro}})}{\overline{T}_{\mathbf{x}}^{\tau\tau}(\tau_{\text{EKT}})} \int d^2\mathbf{x}' G_{\alpha\beta}^{\mu\nu}(\mathbf{x}, \mathbf{x}', \tau_{\text{hydro}}, \tau_{\text{EKT}}) \delta T_{\mathbf{x}}^{\alpha\beta}(\tau_{\text{EKT}}, \mathbf{x}')$$

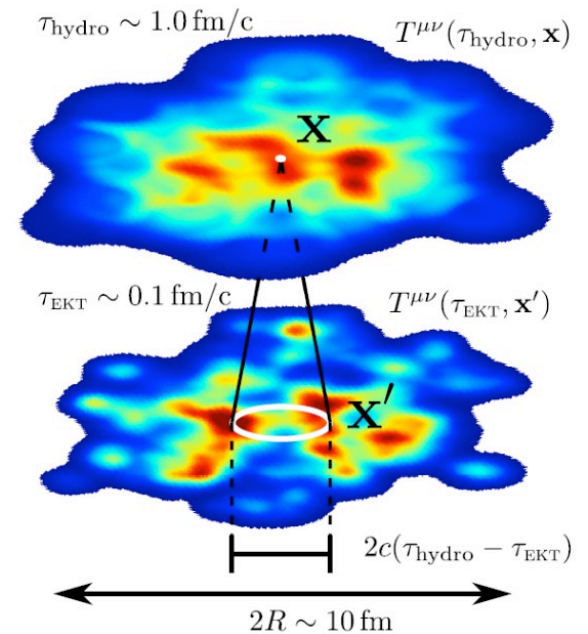


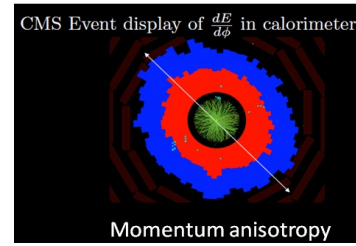
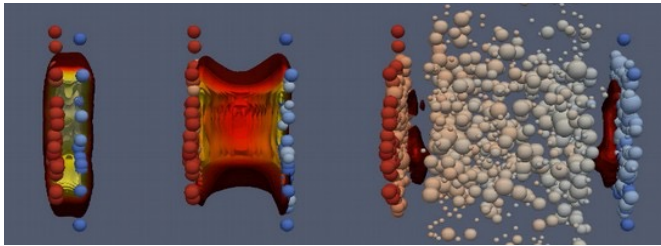
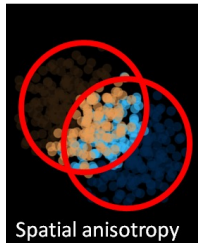
Response functions describing evolution of perturbations

- Also scale with relaxation time

[See also Kamata, Martinez, Plaschke, Ochsenfeld, Schlichting, PRD 2020]

Ref.: Kurkela, Mazeliauskas, Paquet, Schlichting, Teaney PRL2019, PRC2019

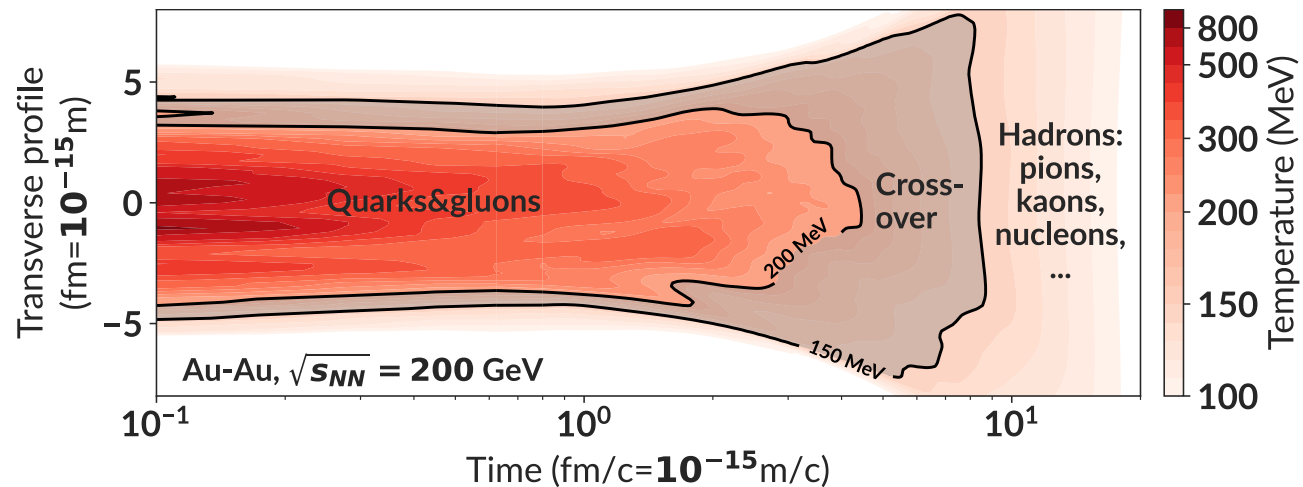




SUMMARY

Summary

- Need transient hydrodynamics or need early-time model



- Early-time challenges are subtle because we “solve hydrodynamics ~backward”