

Critical Mechanical Structures: Topology and Entropy

Xiaoming Mao

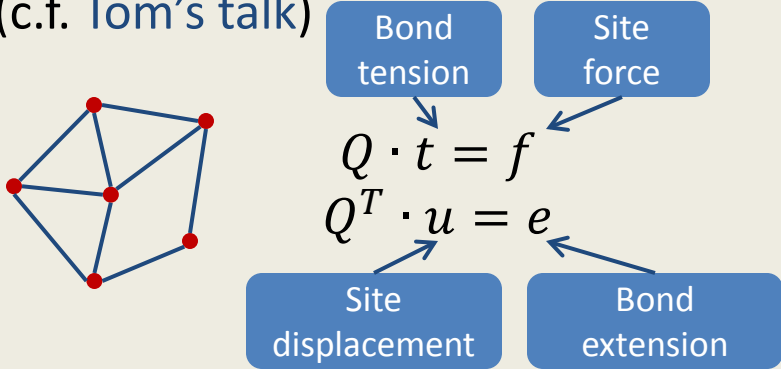
Department of Physics, University of Michigan



KITP workshop: “Geometry, elasticity, fluctuations, and order in 2D soft matter”
Feb 2, 2016

What Are “Critical Mechanical Structures”?

Full machinery of floppy modes counting
(c.f. Tom’s talk)



$$N_B - N_S = N_{d.o.f.} - N_0$$

Rigid-Body

Mechanically Stable: $N_0 = d(d + 1)/2$

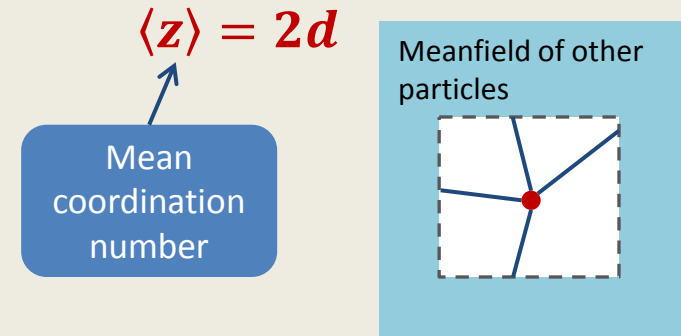
Stress-Free: $N_S = 0$

“Isostaticity = Stable + Stress-free”



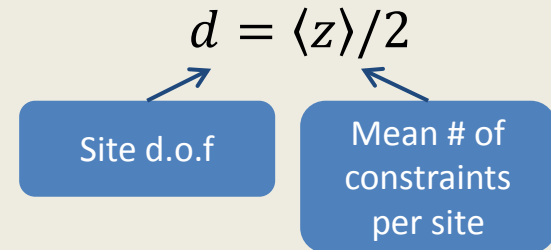
Kindergarten version:

Onset of mechanical stability is at



Look at site averaged quantities

Mechanical stability
(ignoring self-stress)



- J. C. Maxwell, Phil. Mag. **27**, 598 (1864).
- C. R. Calladine, Int. J. Solids Struct. **14**, 161 (1978).
- Thorpe, J. Non-Cryst. Solids, **57**, 355 (1983).
- K. Sun, A. Souslov, X. Mao, and T. C. Lubensky, PNAS **109**, 12369 (2012).
- C. L. Kane and T. C. Lubensky, Nat. Phys. **10**, 39 (2014).
- Lubensky et al, Rep. Prog. Phys., **78**, 073901 (2015)

Why Critical Mechanical Structures Are Interesting?

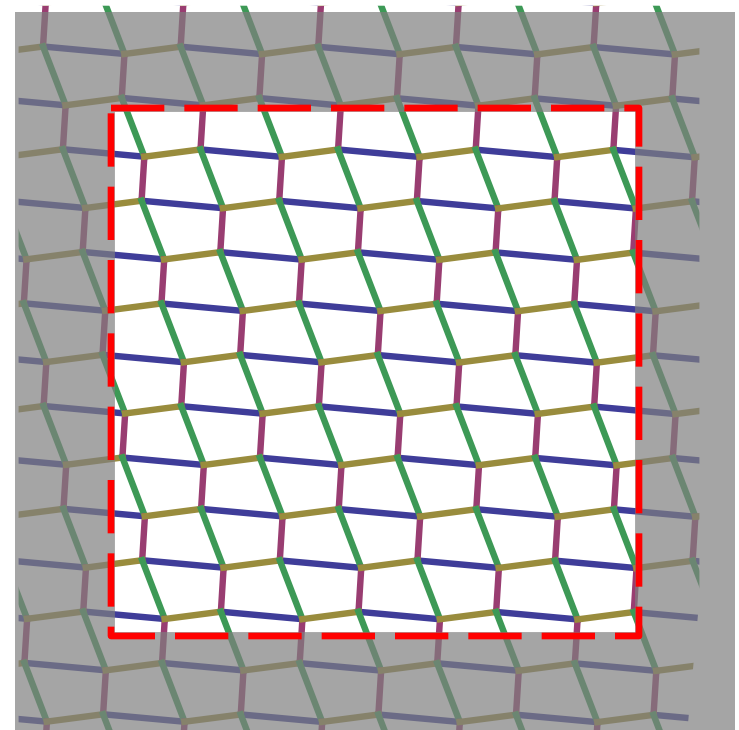
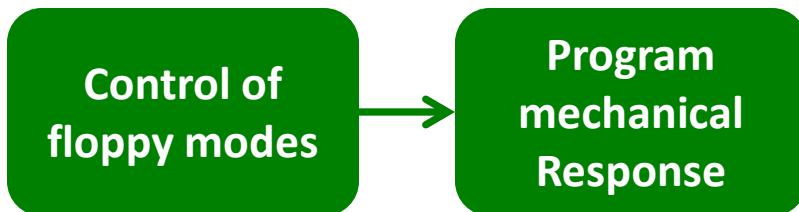
- Central to understanding a wide range of real systems (will discuss later)

Critical mechanical structures



- Interesting fundamental questions:
Start from one simple example:
Cut a finite piece of Maxwell lattice ($z = 2d$)
Deficit of constraints on the boundary
→
of floppy modes \propto **Size of boundary**

But where are these modes located?



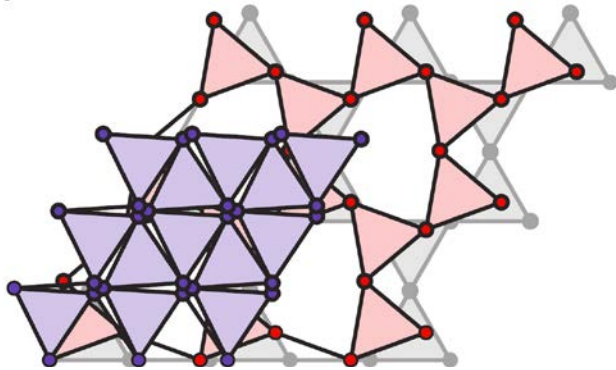
Where Are The Floppy Modes?

Case I. Floppy modes **localized at where you cut**

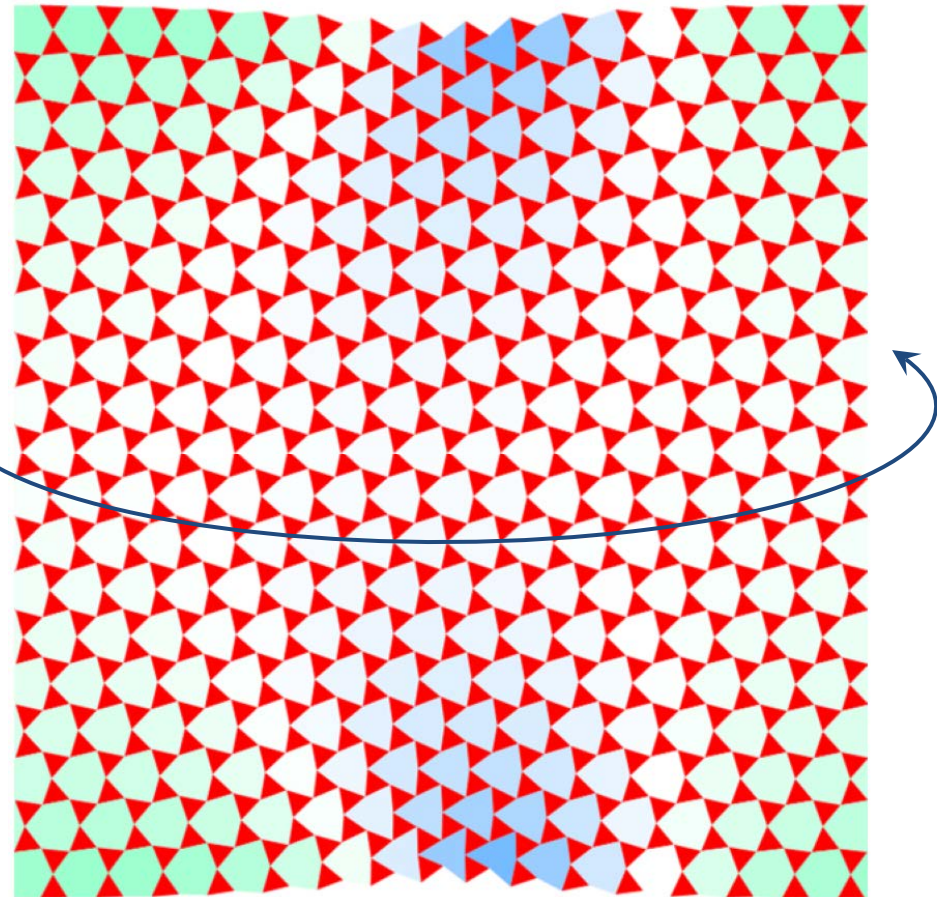
Example: **twisted kagome**

Floppy modes on all edges

Described by conformal transformations



Periodic
BC in x

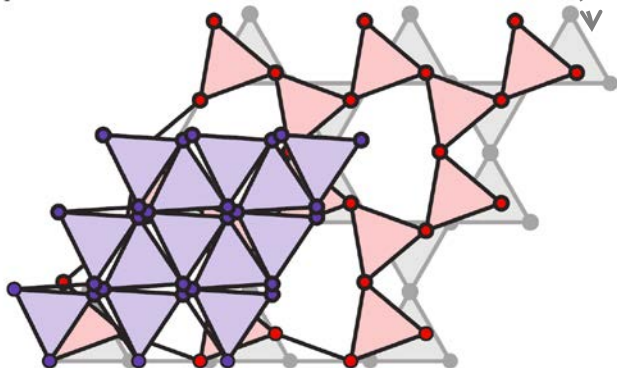


Where Are The Floppy Modes?

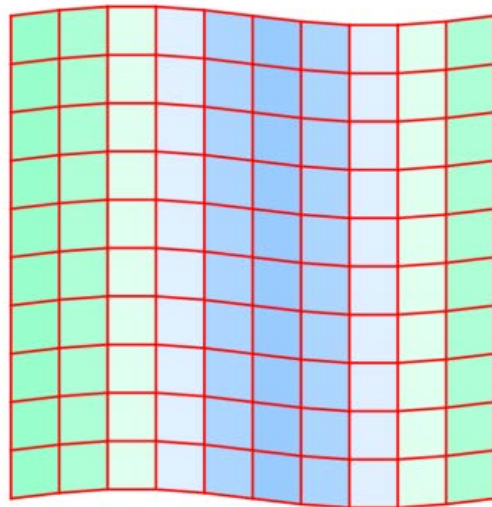
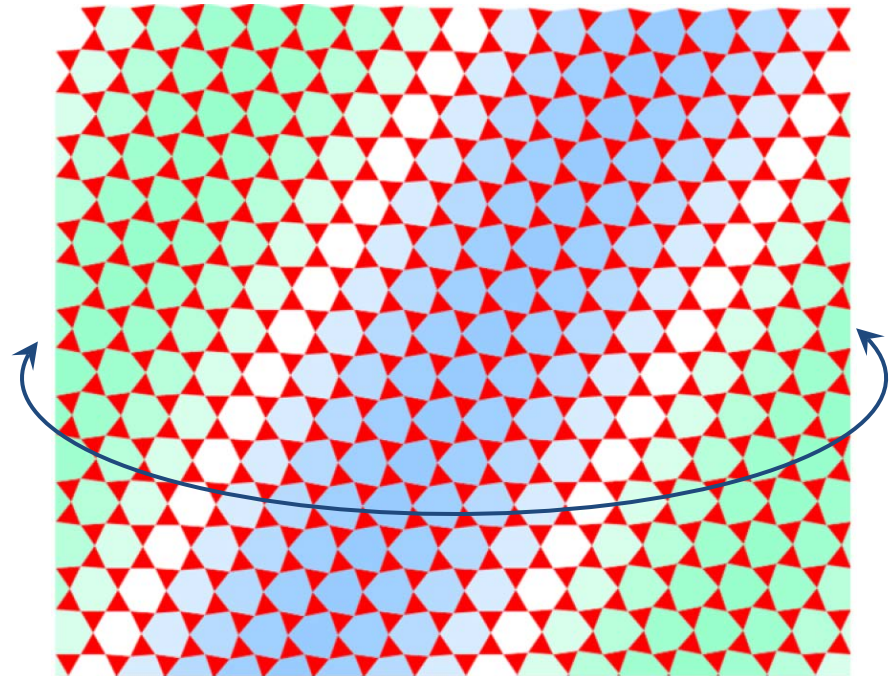
Case II. Floppy modes are **plane waves**

Example: **regular kagome**

regular square



Periodic
BC in x

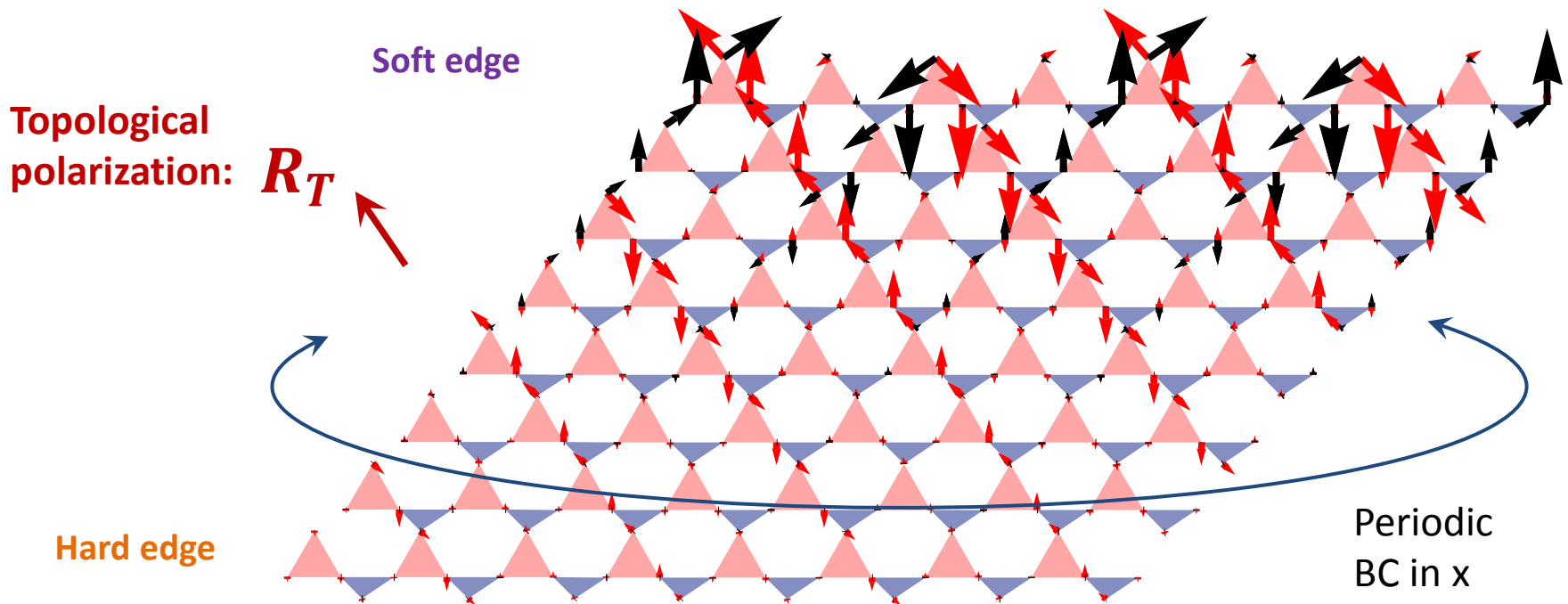


Where Are The Floppy Modes?

Case III. Floppy modes are **topologically protected edge modes on one side**

Example: **deformed kagome**

(arrows show a pair of modes)



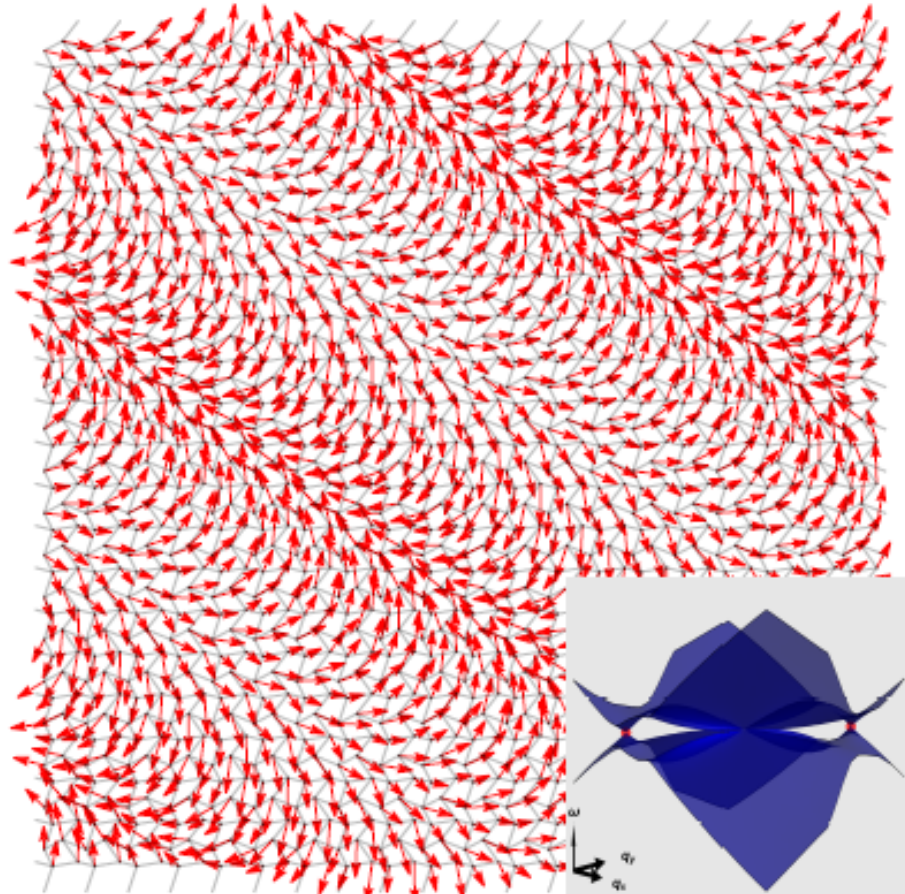
- C. L. Kane and T. C. Lubensky, Nat. Phys. 10, 39 (2014).
- Lubensky, Kane, Mao, Souslov, Sun, Rep. Prog. Phys., 78, 073901 (2015)

Where Are The Floppy Modes?

Case IV. Floppy modes are **topologically protected Weyl modes in the bulk**

Example: **deformed square** (4 site per cell)

(arrows show a Weyl mode)



Talk Outline

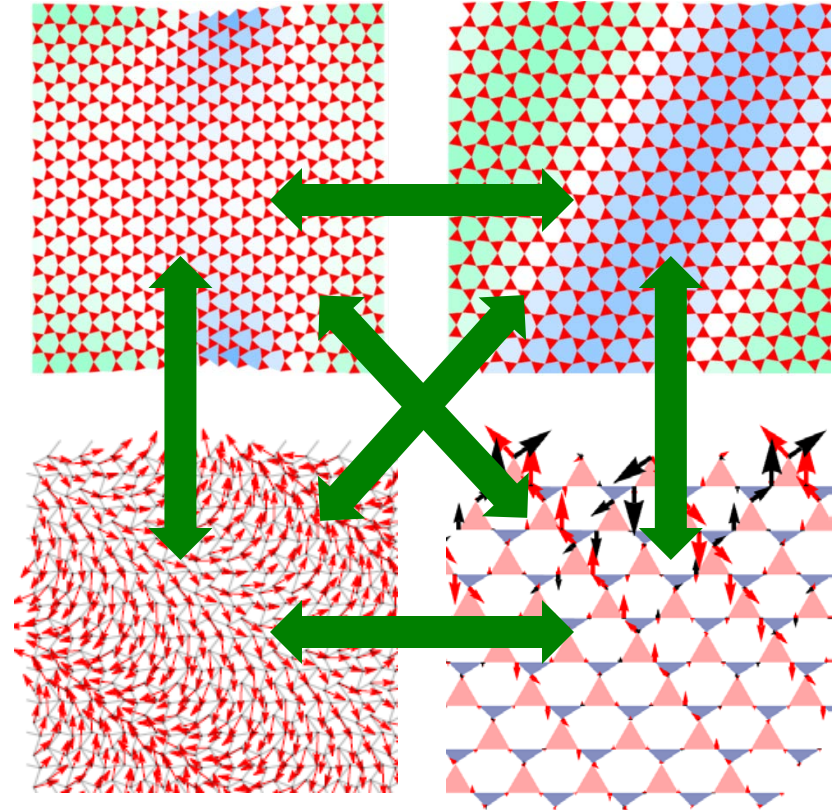
In this talk we will try to answer 2 questions:

1. In a real system can you switch between different classes of floppy mode behaviors?

- Analogous to metal-insulator transitions
- Transformable mechanical metamaterials
- General classification of critical mechanical structures

2. What are the thermodynamic signatures when these transitions occur?

- Large number of competing instabilities
- “Is a mechanism (floppy mode) still a mechanism at finite T ?”



Collaborators

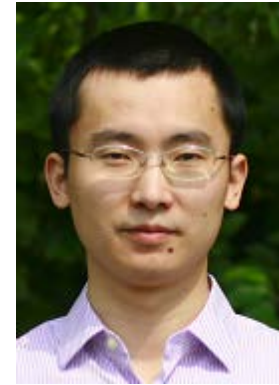
- Topological transitions, mechanical metamaterials, general classifications



Zeb Rocklin
(Michigan)



Shangnan Zhou
(Michigan)



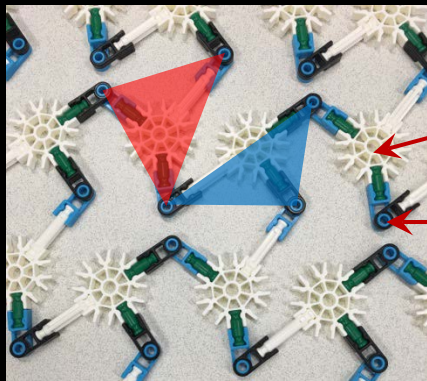
Kai Sun
(Michigan)

Rocklin, Zhou, Sun, and Mao, [arXiv:1510.06389](https://arxiv.org/abs/1510.06389) [cond-mat.soft]

Transformable topological mechanical metamaterials

D. Zeb Rocklin, Shangnan Zhou,
Kai Sun, Xiaoming Mao

University of Michigan



Rigid connector

Hinge

Webpage of video:

<http://www-personal.umich.edu/~maox/research/TTMM/TTMM.html>

What Is The Operation?

On the determinacy of repetitive structures

S.D. Guest^{a,*}, J.W. Hutchinson^b

^a*Department of Engineering, University of Cambridge, Trumpington Street, Cambridge, CB2 1PZ, UK*

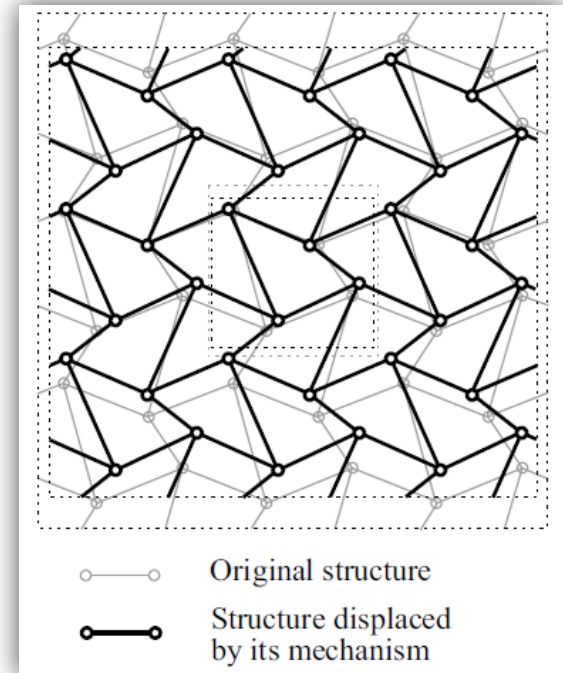
^b*Department of Engineering and Applied Sciences, Harvard University, Cambridge, MA 02138, USA*

Received 5 June 2002; accepted 24 September 2002

“Any $z = 2d$ lattice must have at least $d(d - 1)/2$ homogeneous elastic deformations that are soft (stretch no bonds)”

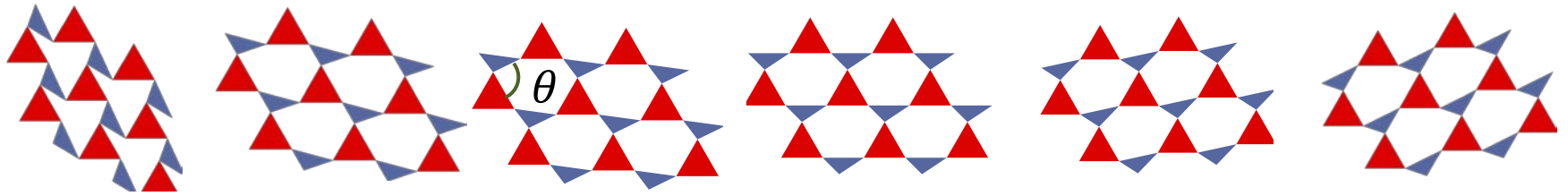
↗
“Guest Modes”

A convenient way to TUNE Maxwell lattices!

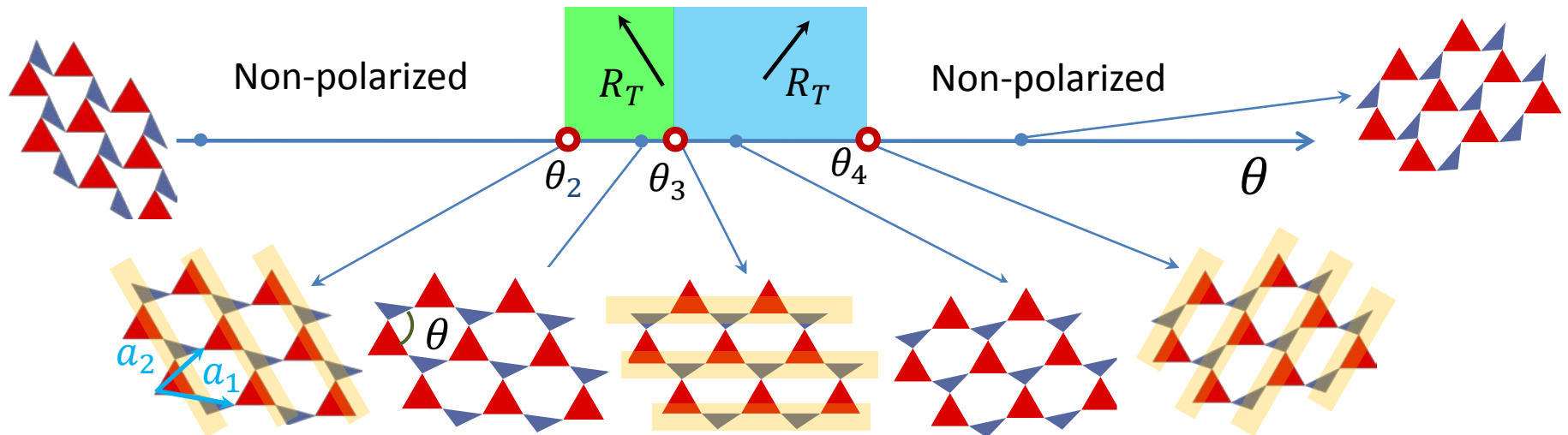


Uniform Soft Twistings of The Deformed Kagome

- Uniform soft twisting is an easy way to transform a lattice



- Does it change the topological structures of the phonons?

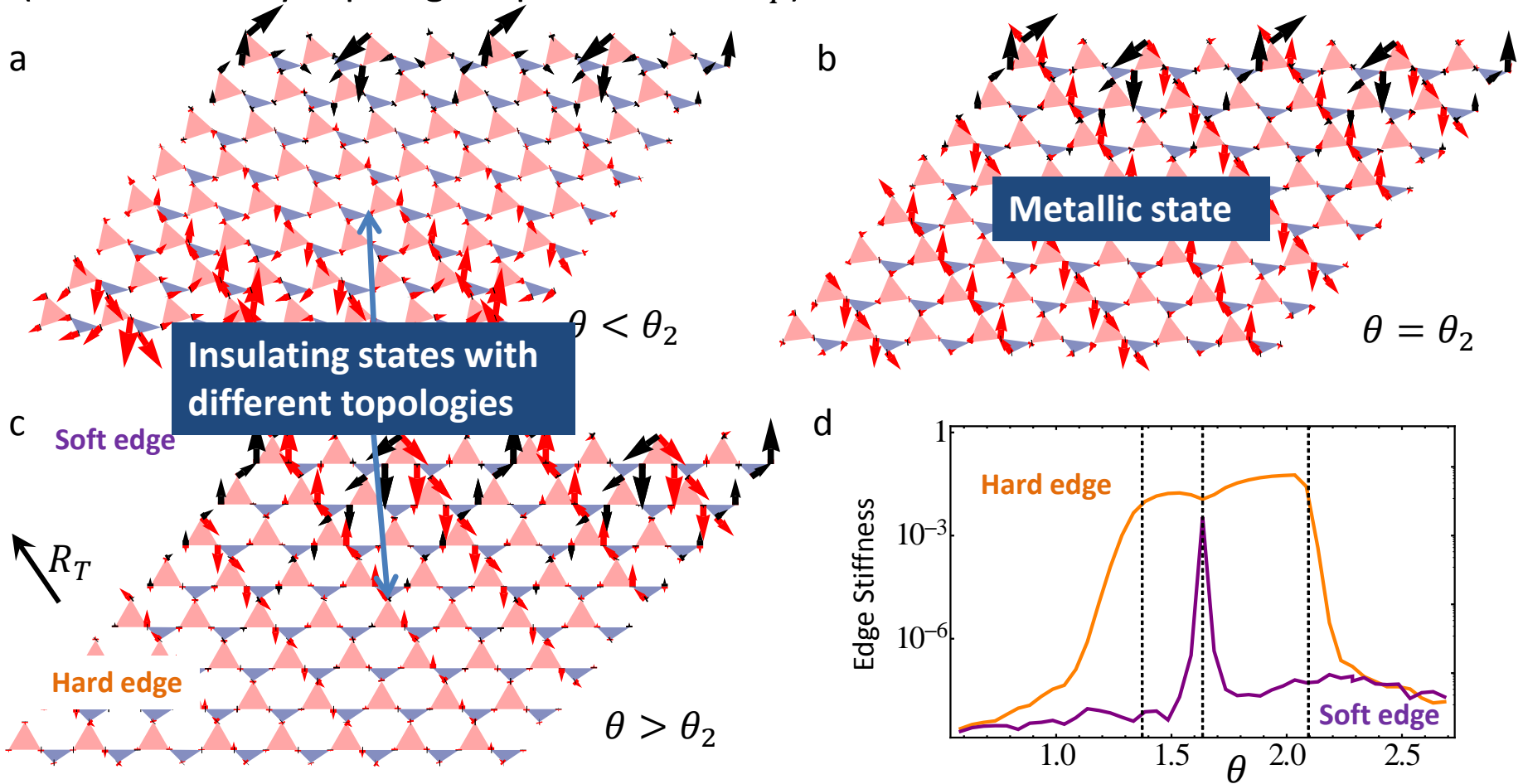


What happens at the topological transitions $\theta_2, \theta_3, \theta_4$?

What happens at the transition?

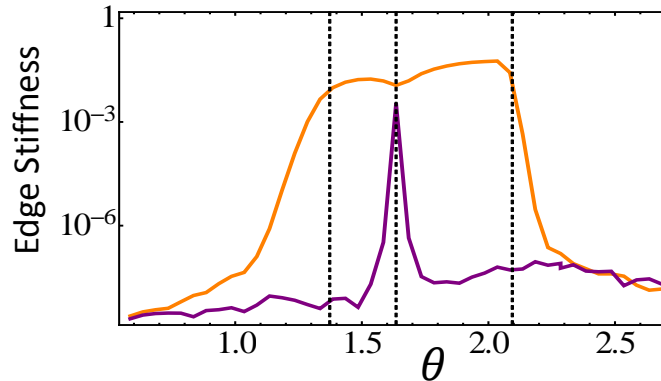
Evolution of a pair of floppy edge modes

(determined by topological polarization R_T)



→ Transformable mechanical metamaterial

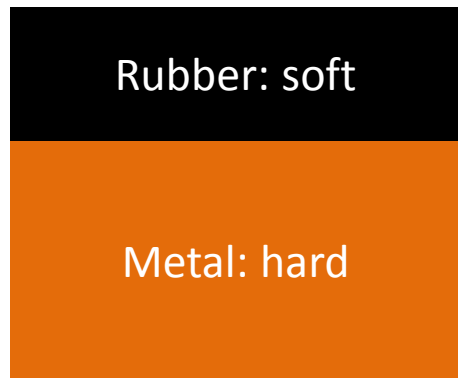
Switch of Stiffness Is Protected by Topology



→ Transformable mechanical metamaterial

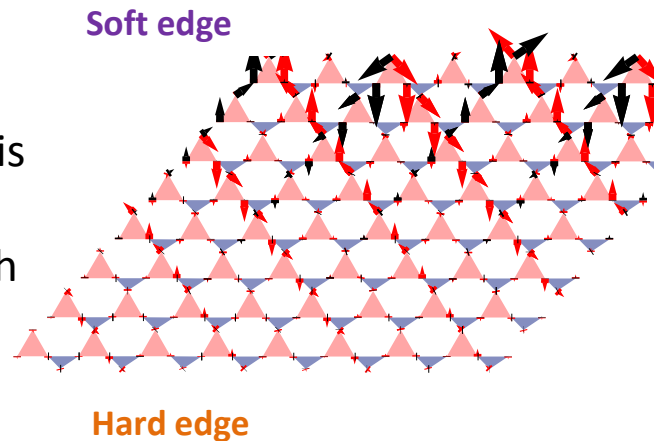
Compare 2 approaches to get contrasting edge stiffness

(1) composite



- If rubber wears out, protection is lost
- Can't switch

(2) topological



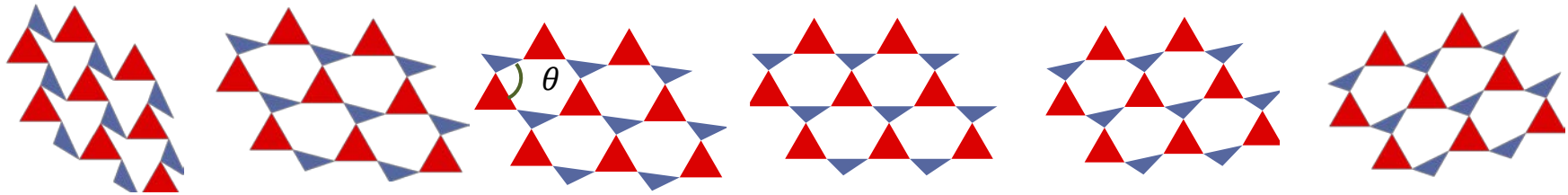
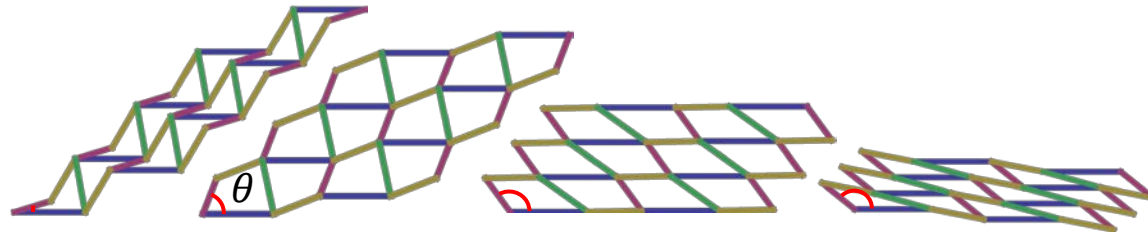
- If top layers wear off, newly exposed surface becomes soft
- Can switch

Which Lattices Exhibit Topological Transitions?

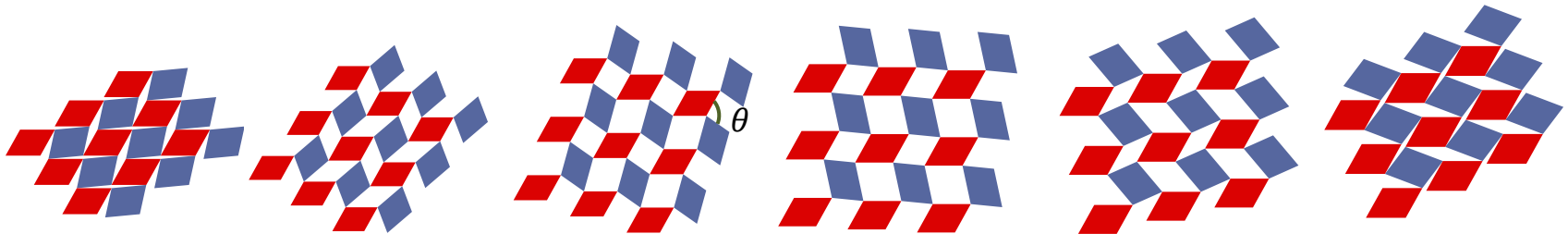
Many structures exhibit soft elastic deformations:

“Uniform soft twistings”

- ALL $\langle z \rangle = 2d$ lattices (guaranteed by the Guest modes)

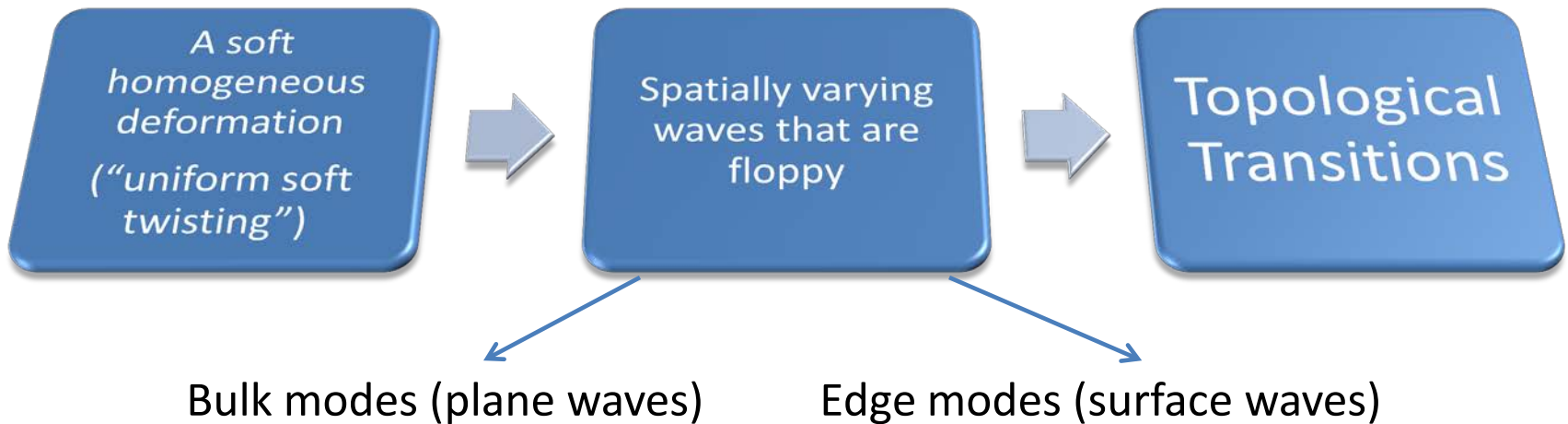


- Some $\langle z \rangle > 2d$ lattices (fine tuning needed)

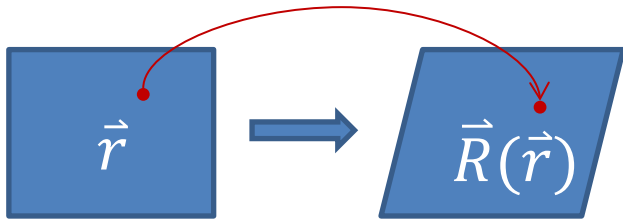


How many “universality classes” are there for mechanical instability?

Our Approach: General Classification



Floppy Modes



Deformation gradient: $\Lambda_{ij} \equiv \frac{\partial R_i}{\partial r_j}$

Metric Tensor: $g = \Lambda^T \Lambda$

(Right Cauchy-Green) strain tensor: $\epsilon \equiv (g - I)/2$

- General form of elastic energy (assume no stress)

Elastic constants

$$F = \int [c_{ijkl} \epsilon_{ij} \epsilon_{kl} + O(\epsilon^3)] d^2 r$$

- The **uniform soft twisting** can be written as

$$\tilde{\epsilon} = \begin{pmatrix} \epsilon_{xx} & \epsilon_{xy} \\ \epsilon_{xy} & \epsilon_{yy} \end{pmatrix} \quad \text{and they satisfy} \quad c_{ijkl} \tilde{\epsilon}_{ij} \tilde{\epsilon}_{kl} = 0$$

- So all strain tensors of the form $\epsilon(\mathbf{r}) = \mathbf{f}(\mathbf{r}) \tilde{\epsilon}$ must also be soft

Arbitrary Scalar

$$F = \int [f(\mathbf{r})^2 c_{ijkl} \tilde{\epsilon}_{ij} \tilde{\epsilon}_{kl} + O(\tilde{\epsilon}^3)] d^2 r = O(\tilde{\epsilon}^3)$$

Floppy Modes

- Not all arbitrary scalar function $f(r)$ correspond to physical deformations
- Extra condition on $g(r)$ for deformations in **flat space**

$$\Gamma_{kij} = \frac{1}{2}(\partial_j g_{ki} + \partial_i g_{kj} - \partial_k g_{ij})$$

No curvature

$$R_{ijkl} = \partial_k \Gamma_{ilj} - \partial_l \Gamma_{ikj} + g^{mn} \Gamma_{ikm} \Gamma_{nlj} - g^{mn} \Gamma_{ilm} \Gamma_{nkj} = 0$$

- Two families of floppy modes must also exist

$$\epsilon_+(\mathbf{r}) = \tilde{\epsilon} f_+(x + \lambda_+ y)$$

$$\lambda_+ = (\tilde{\epsilon}_{xy} + \sqrt{-\det \tilde{\epsilon}}) / \tilde{\epsilon}_{xx}$$

$$\epsilon_-(\mathbf{r}) = \tilde{\epsilon} f_-(x + \lambda_- y)$$

with

$$\lambda_- = (\tilde{\epsilon}_{xy} - \sqrt{-\det \tilde{\epsilon}}) / \tilde{\epsilon}_{xx}$$

- Two classes of behaviors

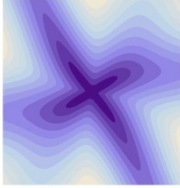
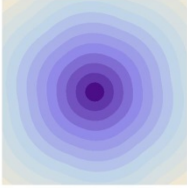
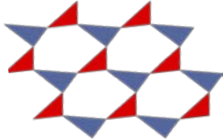



$\det \tilde{\epsilon} > 0$ Dilation dominant, λ_{\pm} imaginary, floppy modes are **edge waves**

$\det \tilde{\epsilon} < 0$ Shear dominant, λ_{\pm} real, floppy modes are **bulk waves**

Auxetic (negative Poisson's ratio)

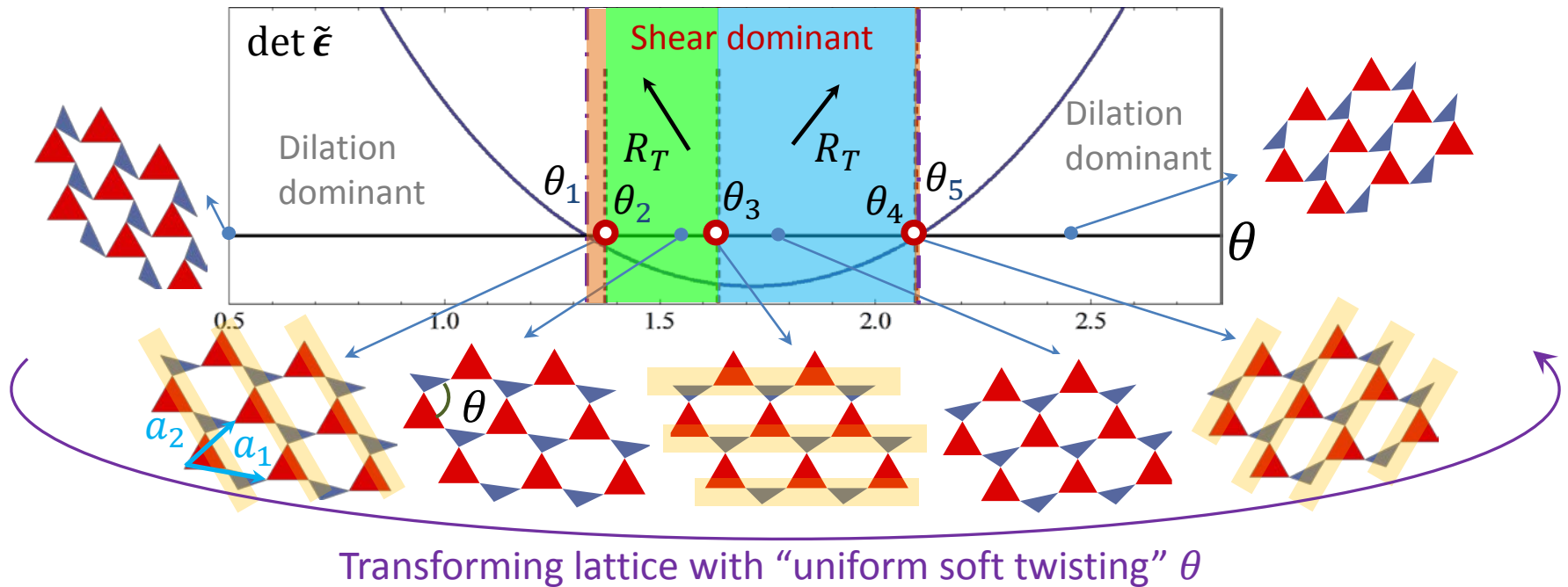
Independent of
microscopic details

Classification of Structures with “Uniform Soft Twisting”

Soft twisting characteristic		Shear dominant $\det \tilde{\epsilon} < 0$		Dilation dominant $\det \tilde{\epsilon} > 0$	
Spatially varying floppy modes		$k_y = \lambda_{\pm} k + O(k^2)$ with $\lambda_{\pm} \in \mathbb{R}$		$k_y = \text{Re}(\lambda_{\pm}) k + i \text{Im}(\lambda_{\pm}) k + O(k^2)$	
Bulk phonon spectra		Vanishing speed of sound in two directions 		Positive speed of sound in all directions 	
Edge modes		Additional floppy edge modes may arise on some edges and can be topological		Floppy edge modes on ALL edges and can be described by conformal transformations	
Floppy edge modes (when present)	Example lattices	$z = 2d$	$z > 2d$	$z = 2d$	$z > 2d$
					
	Frequency	$\omega = 0$	$\omega = O(k^2)$	$\omega = 0$	$\omega = O(k^2)$
	Decay rate	$O(k^2)$	$O(k^2)$	$O(k)$	$O(k)$

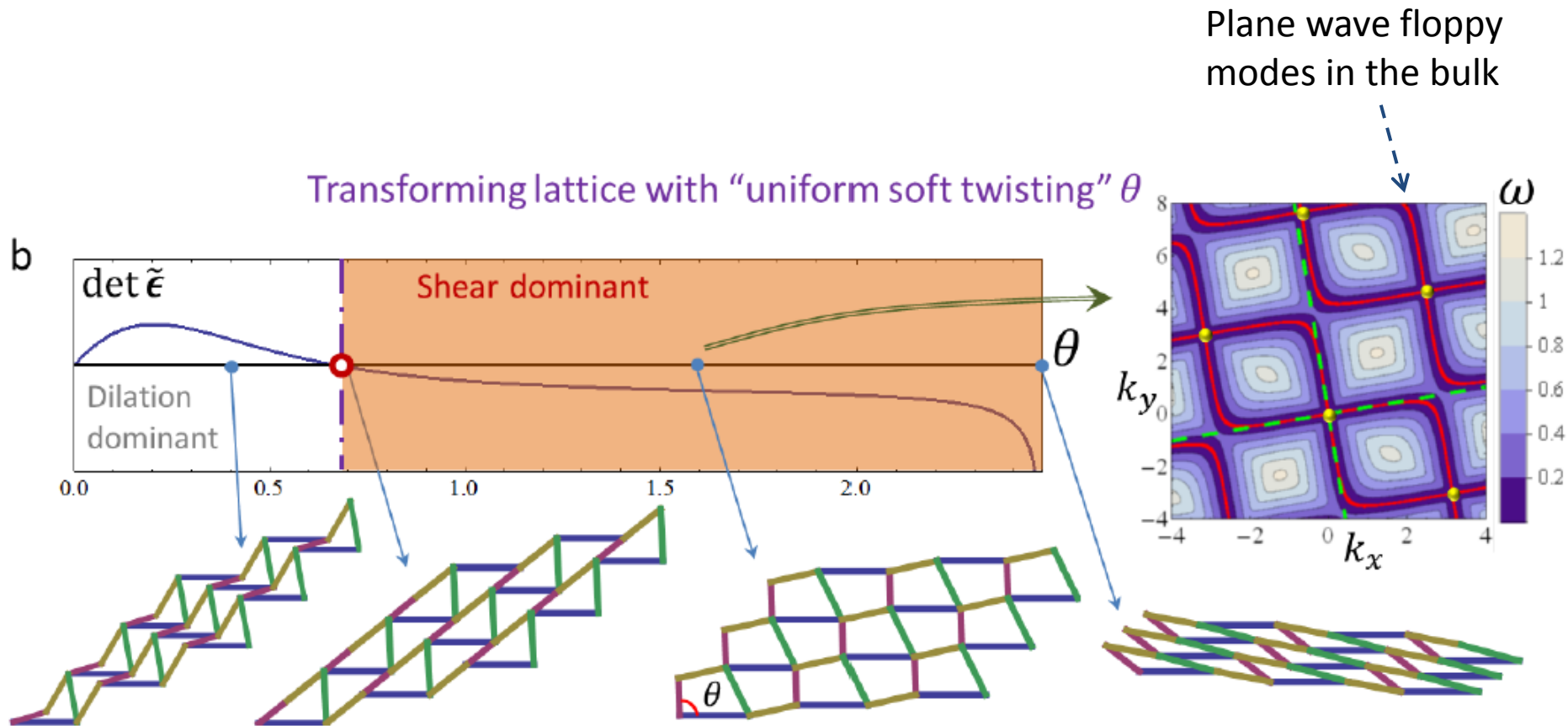
Transitions In The Deformed Kagome

The “uniform soft twisting” can transform a structure between different class



Shear dominant is a necessary but not sufficient condition for topological states

Shear Dominant But Not Topologically Polarized



Talk Outline

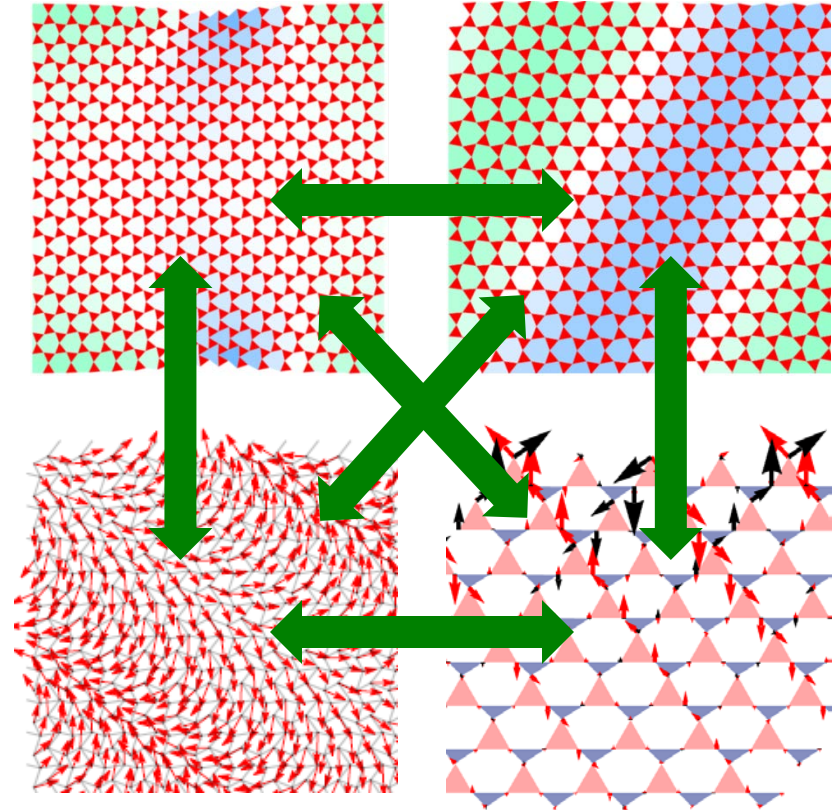
In this talk we will try to answer 2 questions:

1. In a real system can you switch between different classes of floppy mode behaviors?

- Transformable mechanical metamaterials
- New types of phase transitions
- General classification of critical mechanical structures

2. What are the thermodynamic signatures when these transitions occur?

- Large number of competing instabilities
- “Is a mechanism (floppy mode) still a mechanism at finite T ?”



Collaborators

- Finite-T mechanical instability in an ordered lattice



Anton Souslov
(Georgia Tech)



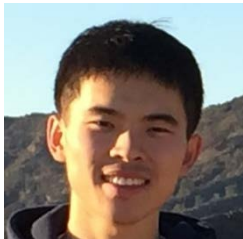
Carlos Mendoza
(Universidad Nacional
Autonoma de Mexico)



Tom Lubensky
(Penn)

- Mao, Souslov, Mendoza, and Lubensky, *Nature Communications* 6, 5968 (2015).

- Finite-T mechanical instability in disordered lattices



Leyou Zhang
(University of Michigan)

- Zhang and Mao, [arXiv:1503.05274](https://arxiv.org/abs/1503.05274) (2015).

- Mechanisms, topological phase transitions and fluctuation effects

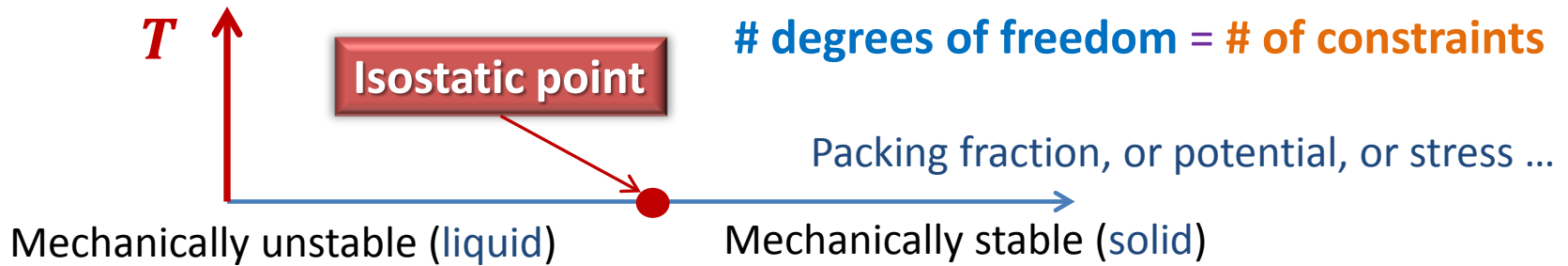


Zeb Rocklin
(Michigan)

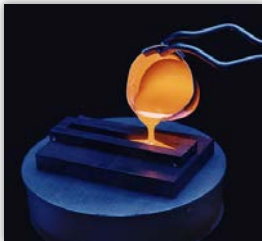


Vincenzo Vitelli
(Leiden)

Motivation of The Original “Finite-T Instability” Project



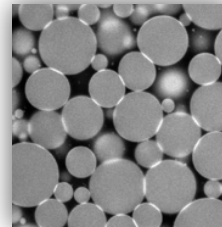
Examples of soft matter systems:



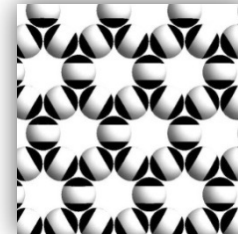
Metallic glasses



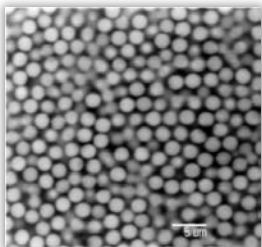
Granular matter



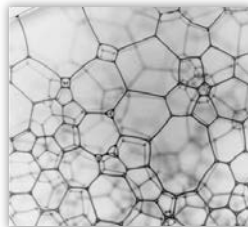
Emulsions



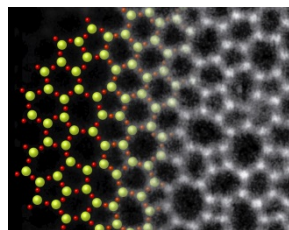
Self-assembled open lattices



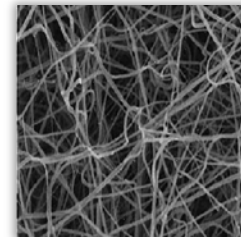
Colloids



Foams



Network glasses



Fiber networks

Finite-T Transitions in Critical Lattices

Ordinary continuous phase transitions:

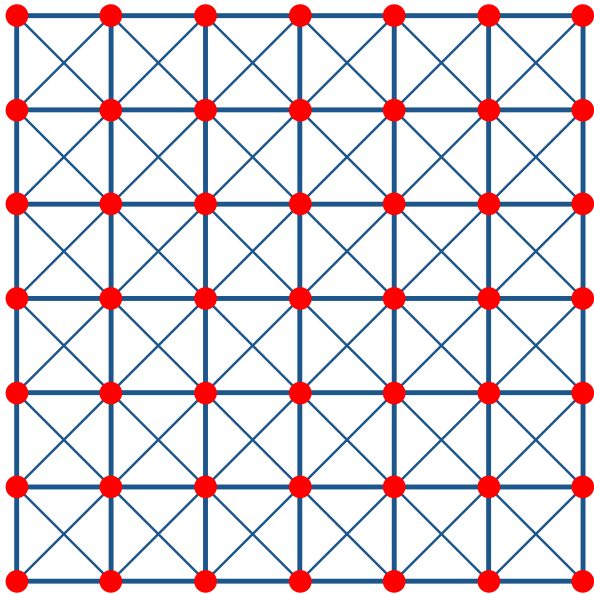


Critical mechanical structures:

Competition between **a large number of instability**

Square Lattice Model

A simple lattice model with a **controlled mechanical instability**



- NN bond: harmonic spring

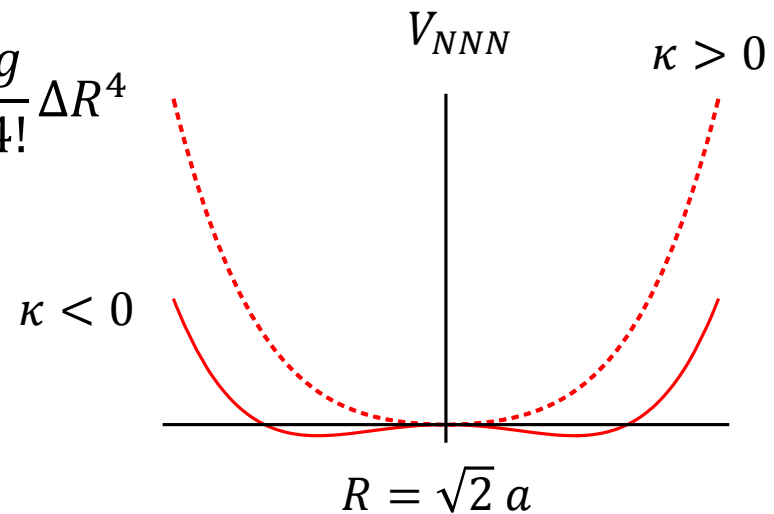
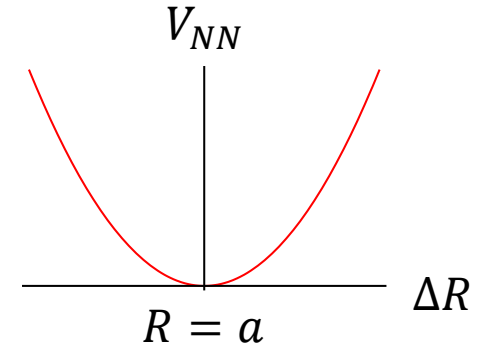
$$V_{NN} = \frac{k}{2} \Delta R^2$$

- NNN bond: anharmonic spring

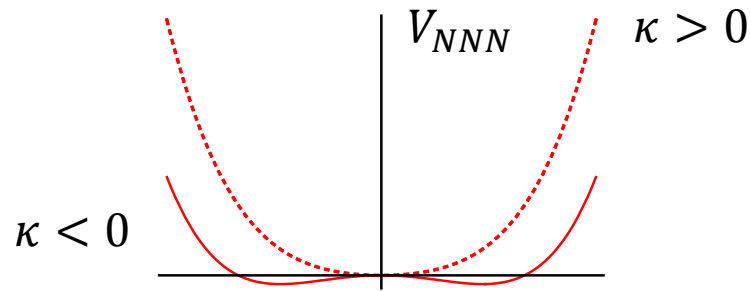
- Harmonic κ

- Quartic g

$$V = \frac{\kappa}{2} \Delta R^2 + \frac{g}{4!} \Delta R^4$$

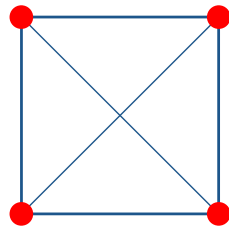


The $T = 0$ Mechanical Instability

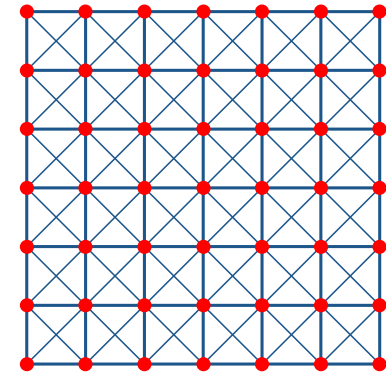


Consider one plaquette:

- When $\kappa > 0$ stable state ($T = 0$): **square**

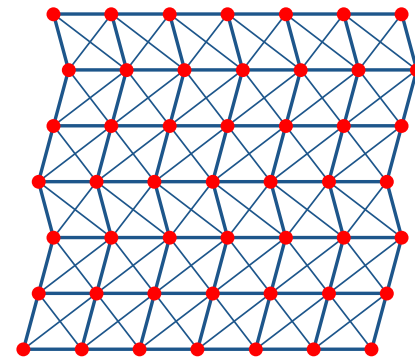
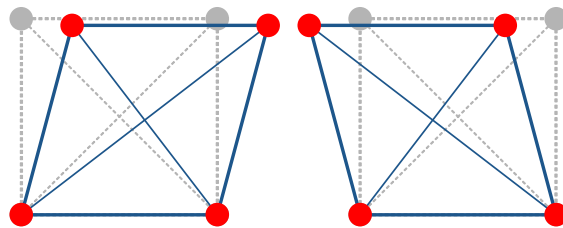


System $T = 0$
stable state



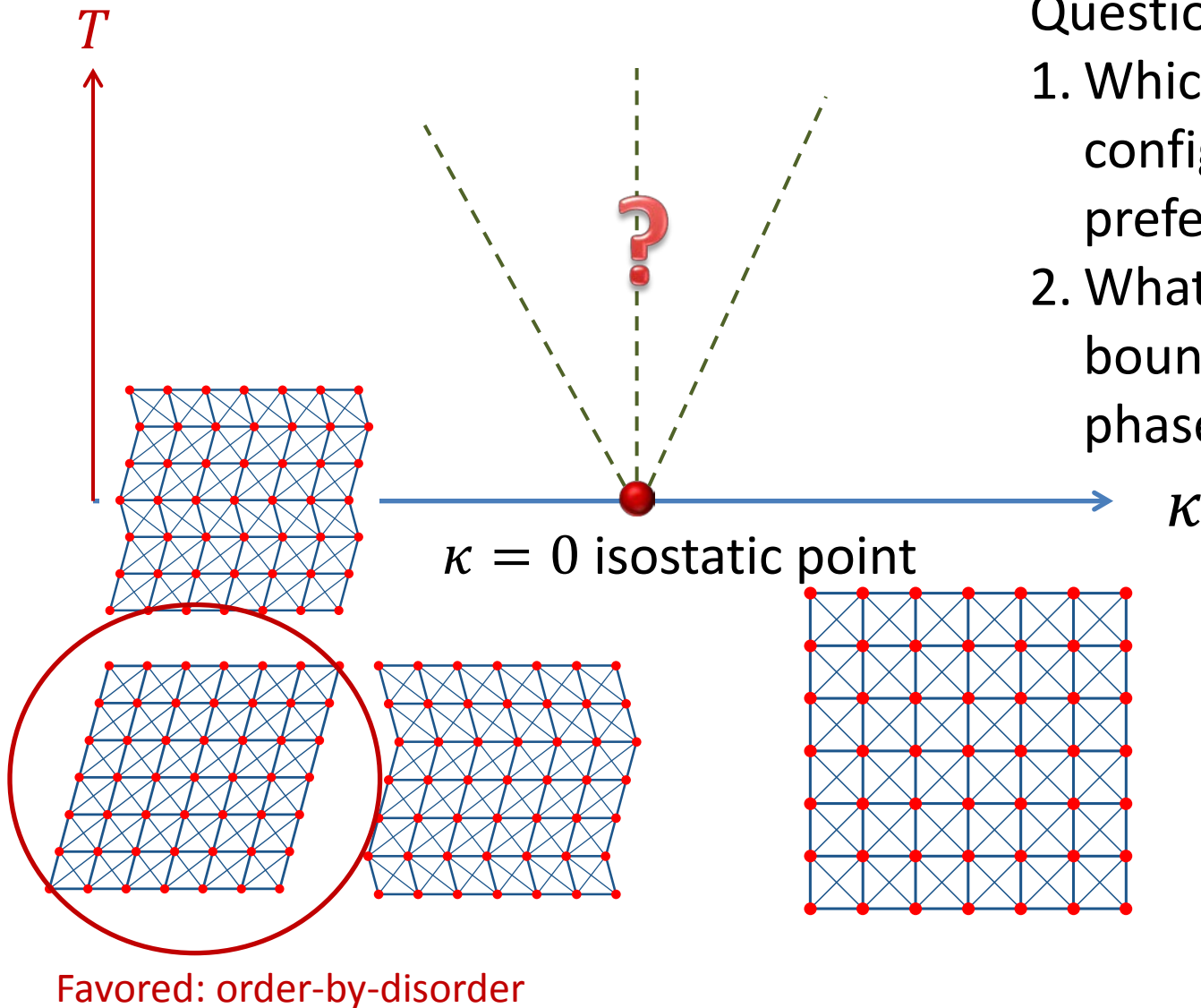
Unique ground state

- When $\kappa < 0$ stable state ($T = 0$): **rhombus**



Multiple degenerate ground state: $e^{\sqrt{N}}$

Effect of Temperature?



Questions:

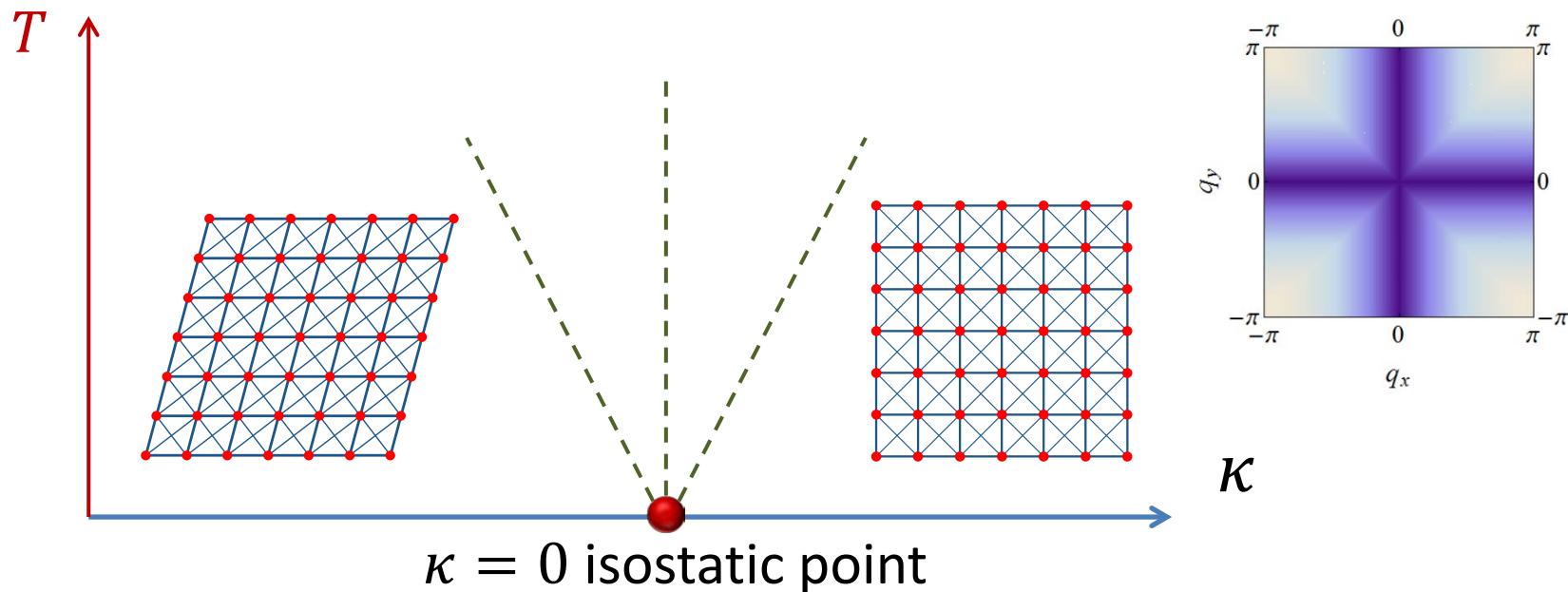
1. Which “zig-zagging configuration” is more preferred? ✓ ($F = U - TS$)
2. What is the stability boundary of the square phase?

Entropic Stabilization of the Square Lattice

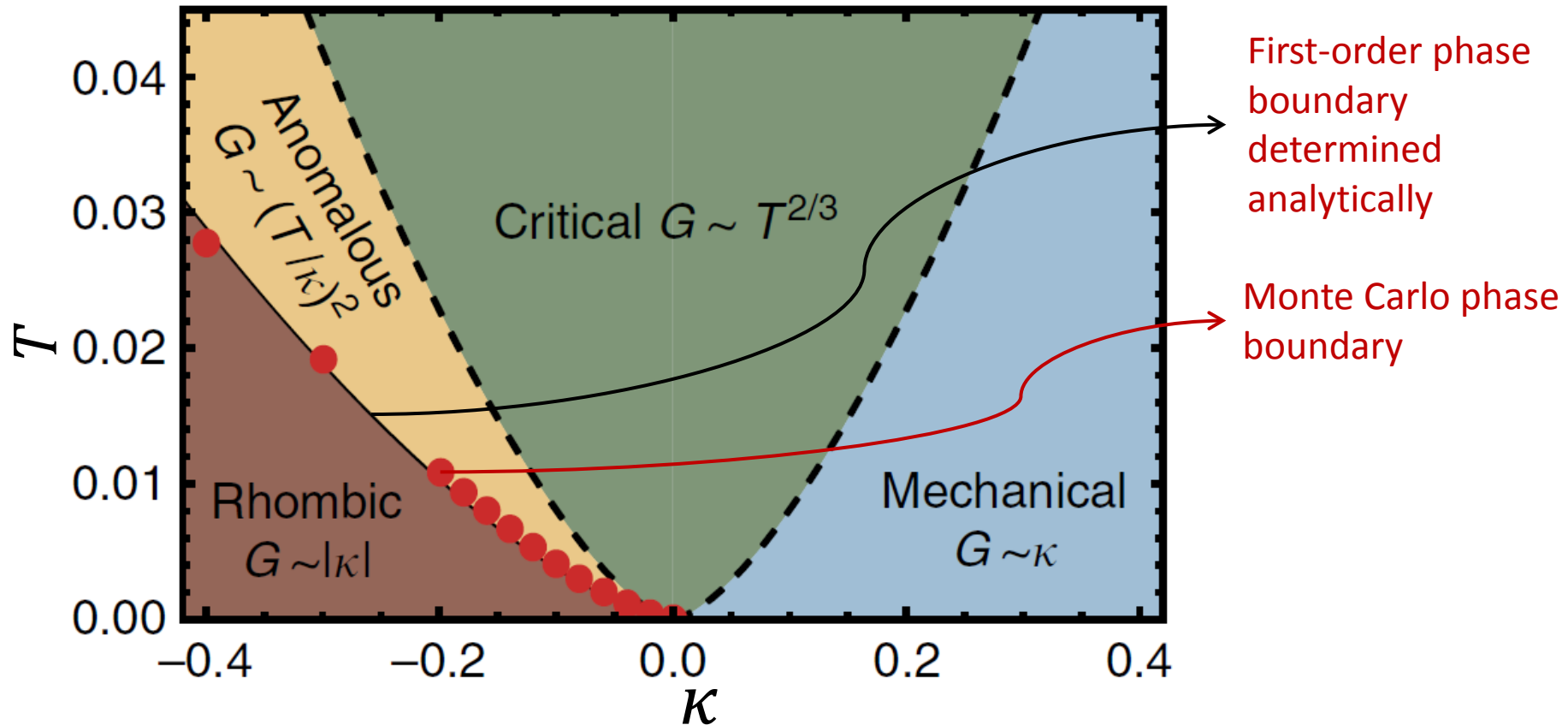
2. What is the stability boundary of the square phase?

Method: Calculate fluctuation corrected shear rigidity (integrate out finite wavelength fluctuations)

What's special: floppy modes living on a 1D manifold in p -space and give a divergent contribution to shear modulus \rightarrow fluctuation driven first order transition (similar to Brazovskii '75 theory for finite wavelength instability in pattern formation)



Phase Diagram

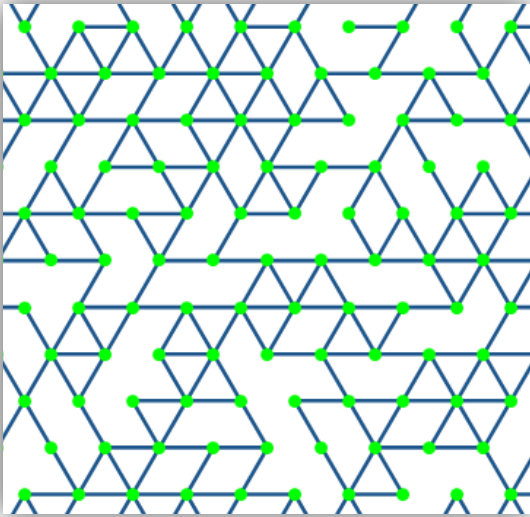


- Transitions near mechanical instability can become first-order if a large number of modes **simultaneously** become critical at mechanical instability (connection to glass transitions)
- This is a mechanical transition in an ordered lattice, which can be seen as a special type of structural phase transition (crystal - crystal), what about liquid – (disordered) solid?

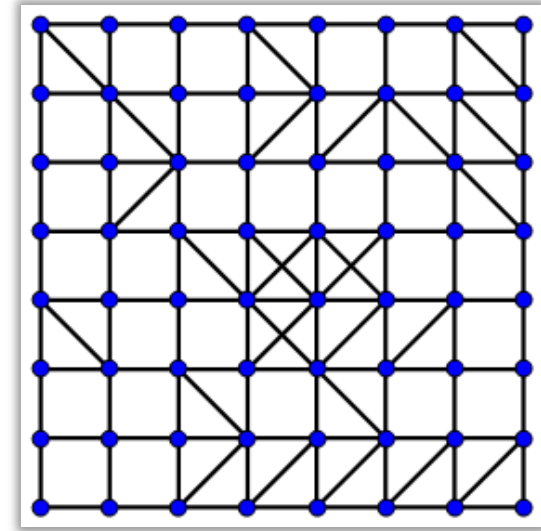
Two CF Model Lattices

Finite T?

Rigidity percolation in a triangular lattice



Square lattice with random NNN bonds



$T = 0$ transition

$$p_c \approx 2/3$$
$$\xi \sim (p - p_c)^{-1.21}$$

Non-meanfield

Isostatic point

$$p_c = 0$$
$$\xi \sim (p - p_c)^{-1}$$

Meanfield,
Agrees with **Jamming**

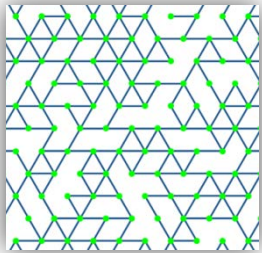


Leyou Zhang

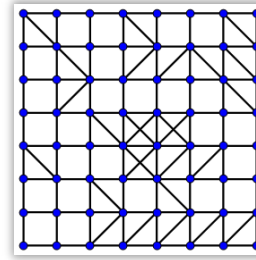
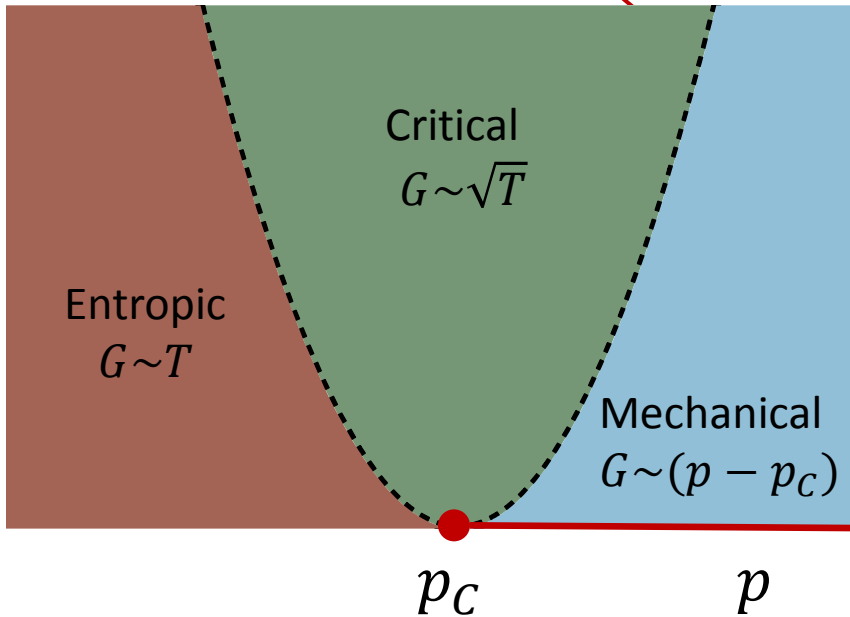
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- Zhang and Mao, [arXiv:1503.05274](https://arxiv.org/abs/1503.05274) (2015).
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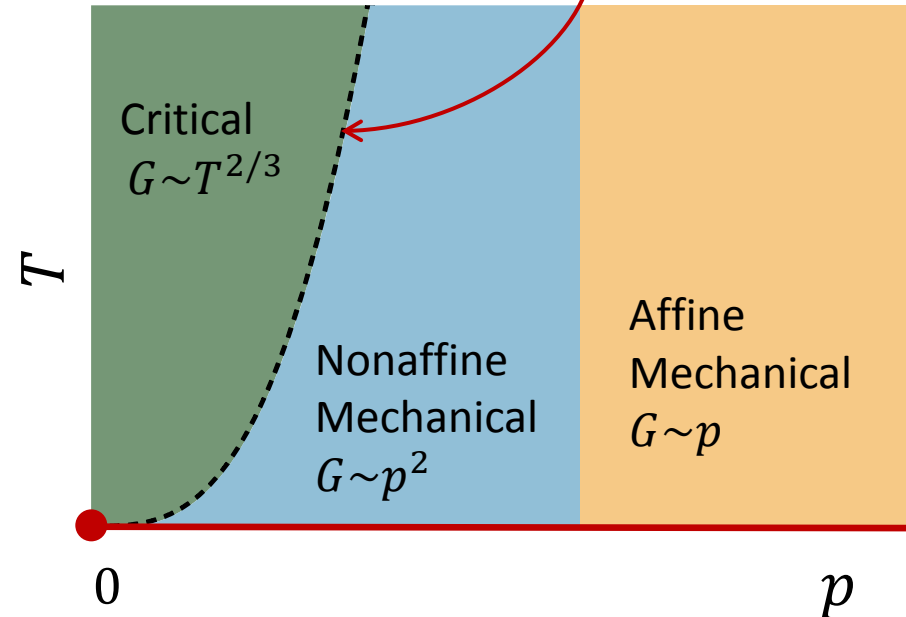
Finite-T EMT Phase Diagrams



Boundary: $T \sim (p - p_c)^2$



Boundary: $T \sim p^3$

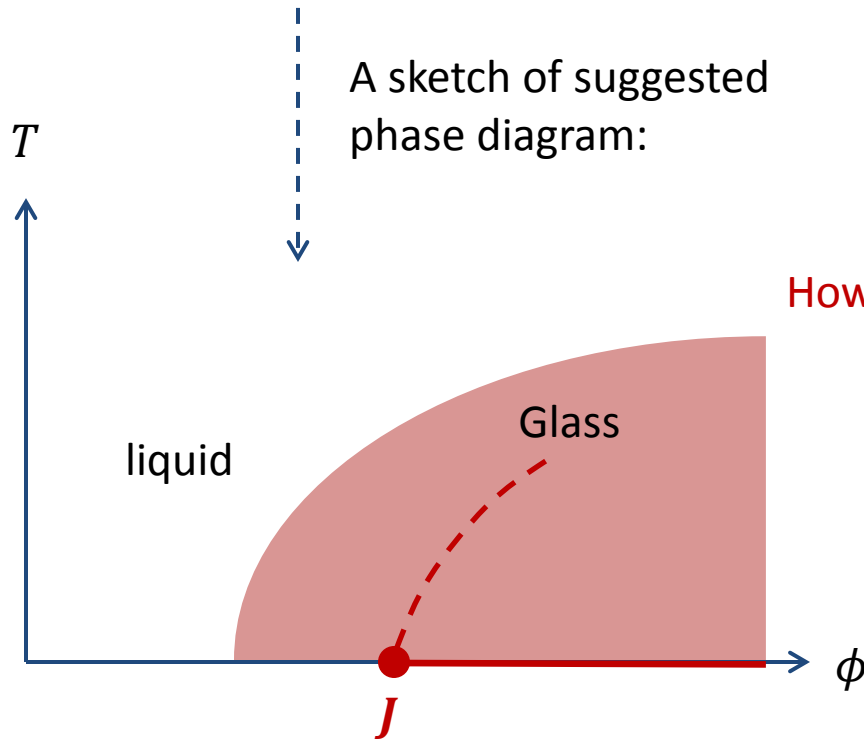


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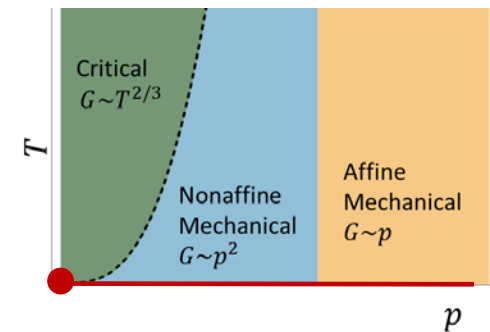
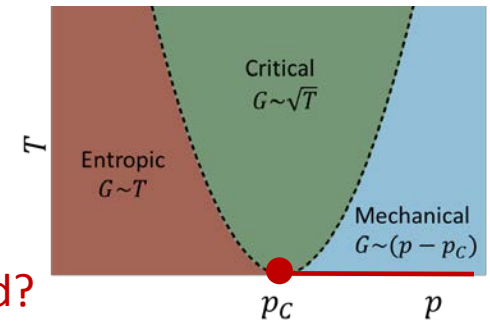
Relation to Glasses?

What happens when thermal fluctuations are introduced near **point J**?

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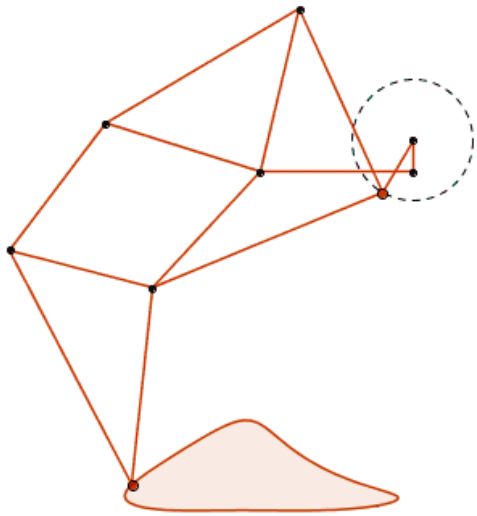


How are they related?

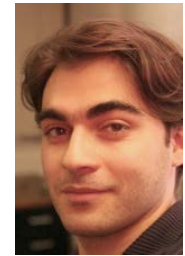


Nonlinear Mechanisms At Finite T?

Macroscopic machines and robots



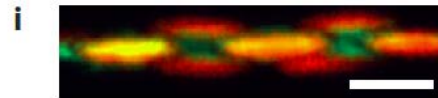
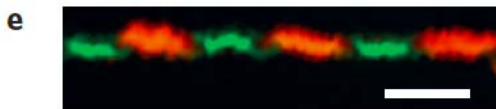
Zeb Rocklin
(Michigan)



Vincenzo Vitelli
(Leiden)

Theo Jansen's walking mechanism and "Strandbeest" (wind powered robot using this mechanism)

How small can a mechanism go?



Summary

- Topological transitions in critical mechanical structures → dramatic change of mechanical properties → “transformable topological mechanical metamaterials” (TTMM)
- General classification of critical mechanical structures
Floppy edge modes vs. Floppy bulk modes
- Interplay between entropy – floppy modes -- topology

THANK YOU!