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The two-dimensional Ising model and SLE

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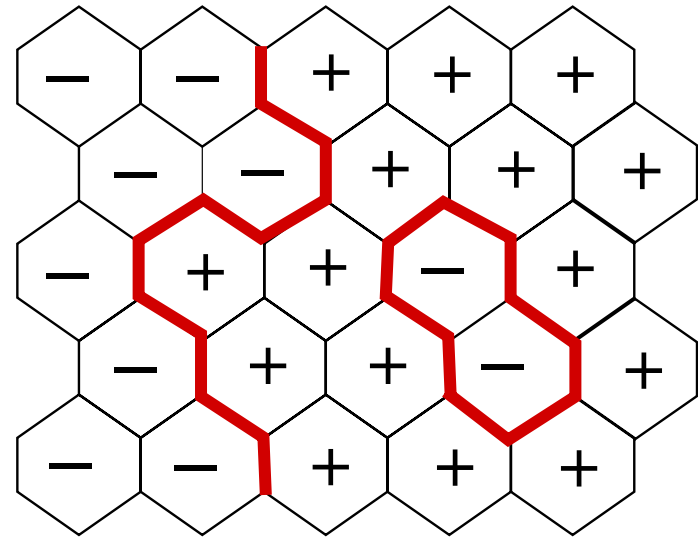
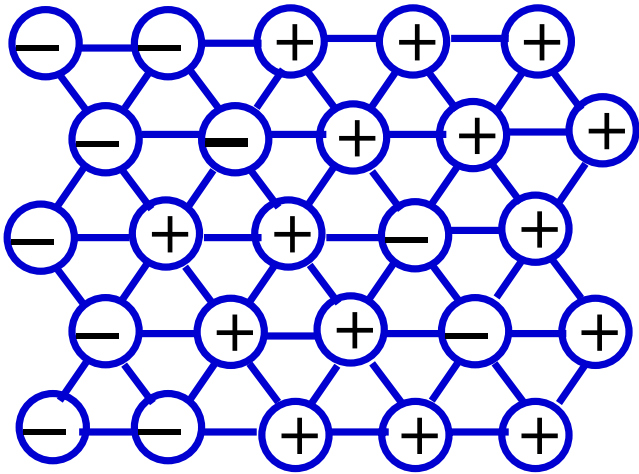


Plan

- 2D Ising model and Conformal Field Theory
- Schramm Loewner Evolution (SLE)
- FK representation and holomorphic fermion
work in collaboration with **John Cardy**

Lattice model

(Ising 1925)



$$H = -J \sum_{\langle ij \rangle} \sigma_i \sigma_j - h \sum_i \sigma_i$$

$$\sigma_i \in \{-1, 1\}$$

$$H = -J \sum_{\langle ij \rangle} \sigma_i \sigma_j - h \sum_i \sigma_i$$

$$\left\{ \begin{array}{l} \langle \mathcal{O} \rangle = \frac{1}{Z} \sum_{\{\sigma\}} \mathcal{O} e^{-\beta H[\sigma]} \\ Z = \sum_{\{\sigma\}} e^{-\beta H[\sigma]} \end{array} \right. \quad \left\{ \begin{array}{l} M = \frac{1}{N} \langle \sum_i \sigma_i \rangle \\ g_{ij} = \langle \sigma_i \sigma_j \rangle \end{array} \right.$$

Exact solution on square lattice at $h=0$ (Onsager 1944)

Critical point: $K \equiv \beta J = K_c, \quad h = 0$

$$\langle \sigma(r) \sigma(0) \rangle \propto \begin{cases} \frac{e^{-r/\xi}}{r^{(d-1)/2}} & K \ll K_c \\ \frac{1}{r^{d-2+\eta}} & K = K_c \end{cases}$$

$$\xi \rightarrow \infty$$

Conformal Field Theory (CFT)

$\xi \rightarrow \infty$
universality



Field Theory description of the scaling limit
(Wilson, Fisher 1960's : RG)

critical exponents:

Specific heat:

$$C \sim |T - T_c|^{-\alpha}$$

Order parameter:

$$M \sim (T_c - T)^\beta$$

Susceptibility:

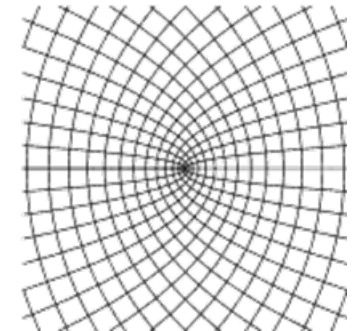
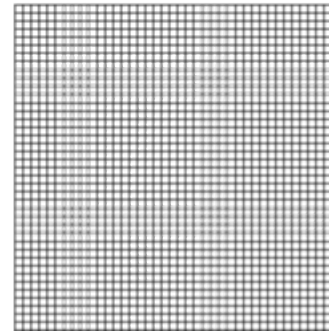
$$\chi \sim |T - T_c|^{-\gamma}$$

...

conformal invariance \longrightarrow Conformal Field Theory

(Belavin-Polyakov-Zamolodchikov 1984)

Conformal transformations:
local translations, rotations and dilatations
preserve angles



2D: analytic functions $z \longrightarrow f(z)$

\longrightarrow equilibrium problems exactly solvable in terms of $\left\{ \begin{array}{l} \text{central charge } c \\ \text{conformal dimensions } \Delta_\phi \end{array} \right.$

\longrightarrow $\left\{ \begin{array}{l} \text{critical exponents: } \Delta_\phi \\ \text{finite size effects: } E_0(L) = -\frac{\pi}{6L} c \end{array} \right.$

Transformation properties in CFT

$$\begin{cases} w = f(z) \\ \bar{w} = \bar{f}(\bar{z}) \end{cases} \longrightarrow \phi(w, \bar{w}) = \left(\frac{df}{dz}\right)^h \left(\frac{d\bar{f}}{d\bar{z}}\right)^{\bar{h}} \phi(z, \bar{z})$$

$$\text{Dilatation: } \begin{cases} w = \lambda z \\ \bar{w} = \lambda \bar{z} \end{cases} \quad \phi(w, \bar{w}) = \lambda^{(h+\bar{h})} \phi(z, \bar{z})$$

$$\text{Rotation: } \begin{cases} w = e^{i\theta} z \\ \bar{w} = e^{-i\theta} \bar{z} \end{cases} \quad \phi(w, \bar{w}) = e^{i(h-\bar{h})\theta} \phi(z, \bar{z})$$

$$\Delta = h + \bar{h} \quad s = h - \bar{h}$$

Holomorphic operators: $\bar{h} = 0$

Ising model

\mathbb{Z}_2 symmetry $\sigma_i \rightarrow -\sigma_i \Rightarrow$ CFT with $c = \frac{1}{2}$

$$\{\phi\} = \{\mathbb{I}, \sigma, \epsilon\}$$

$$\langle \sigma(x)\sigma(0) \rangle = r^{-2\Delta_\sigma} \quad \Delta_\sigma = \frac{1}{8}$$

$$\langle \epsilon(x)\epsilon(0) \rangle = r^{-2\Delta_\epsilon} \quad \Delta_\epsilon = 1$$

$$\langle \sigma_i \sigma_{i+n} \rangle = \frac{1}{|n|^{d-2+\eta}}$$

$$\langle \epsilon_i \epsilon_{i+n} \rangle = \frac{1}{|n|^{2d-2/\nu}}$$

$$\Rightarrow \begin{cases} \eta = 1/4 \\ \nu = 1 \end{cases}$$

scaling laws

$$\Rightarrow \begin{cases} \alpha = 0 \\ \beta = 1/8 \\ \gamma = 7/4 \\ \delta = 15 \end{cases}$$

Fermionic Theory with $c=1/2$

$$S = \frac{1}{2\pi} \int d^2z \left(\psi \partial_{\bar{z}} \psi + \bar{\psi} \partial_z \bar{\psi} + im \bar{\psi} \psi \right) \quad m \propto K - K_c$$

$$\partial_{\bar{z}} \psi = i \frac{m}{2} \bar{\psi}$$

CFT: $m = 0$



$$\psi(z, \bar{z}) = \psi(z)$$

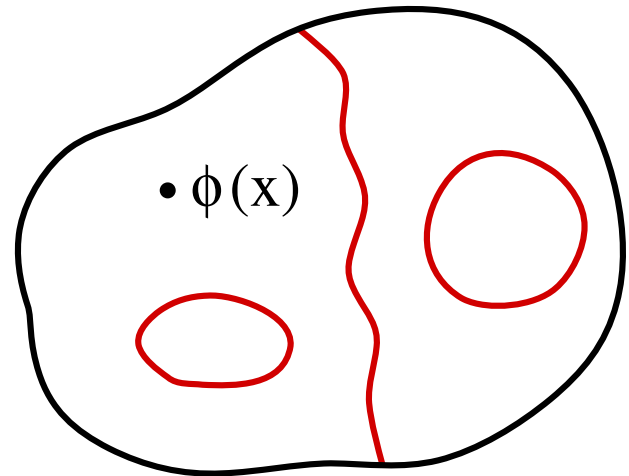
$$\Delta_\psi = s_\psi = \frac{1}{2}$$

What is missing?

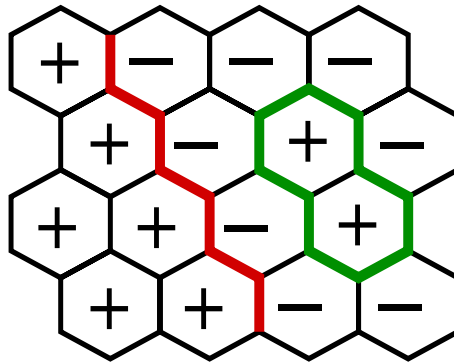
- rigor in relating statistical system to a given model of CFT
- full understanding of non-minimal / non-unitary theories
- satisfactory understanding of **geometrical aspects**

Field Theory is built on local operators
algebraic language

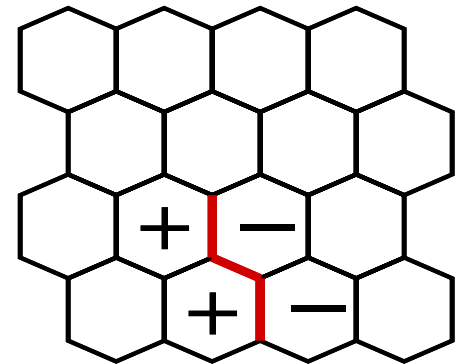
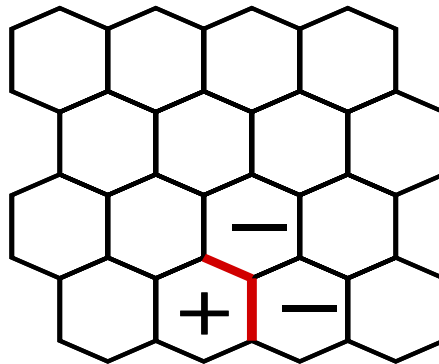
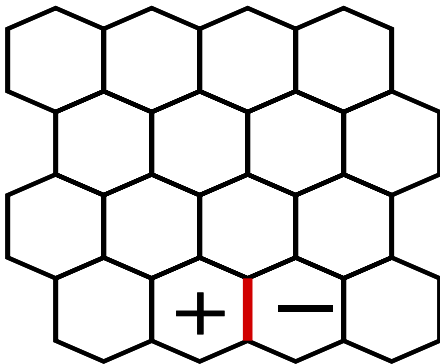
No natural language
for geometric objects like domain walls



Schramm Loewner Evolution (SLE)

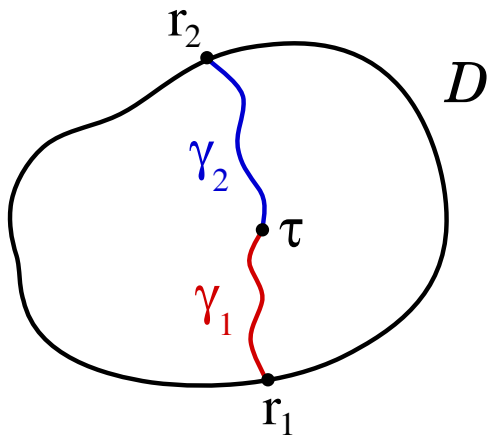


Idea: grow domain wall step by step

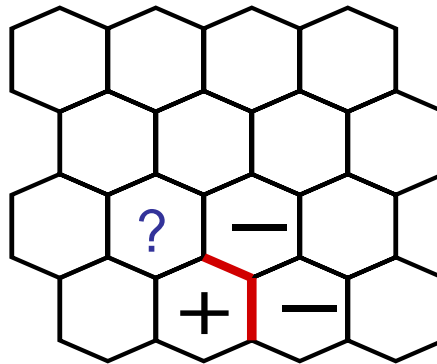


Class of curves in the continuum limit:

Markov property:

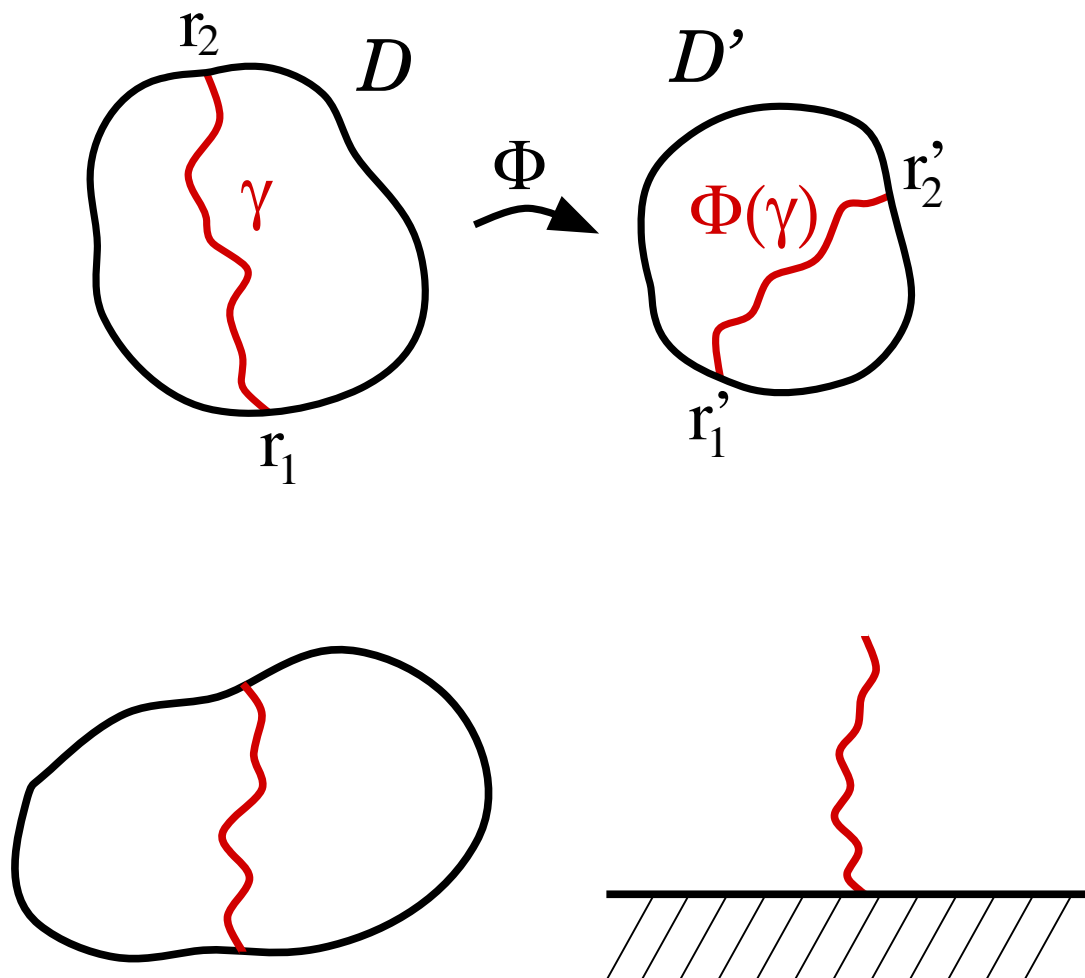


$$\mu(\gamma_2 | \gamma_1; D, r_1, r_2) = \mu(\gamma_2; D \setminus \gamma_1, \tau, r_2)$$

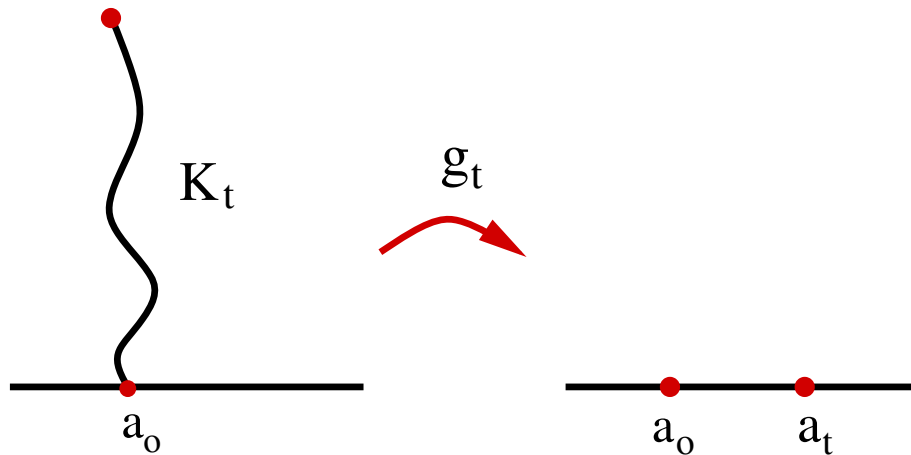


Conformal invariance:

$$(\Phi * \mu)(\gamma; D, r_1, r_2) = \mu(\Phi(\gamma); D', r'_1, r'_2)$$



Non self-crossing curves on the half plane from a_0 to ∞



$K_t =$ curve grown up to time t

$g_t : \mathbb{H} \setminus K_t \rightarrow \mathbb{H}$
2D conformal map

Loewner equation:
 (1923)

$$\frac{dg_t(z)}{dt} = \frac{2}{g_t(z) - a_t}$$

$a_t \in \mathbb{R}$

Schramm:
 (1999)

Markov property
 conformal invariance } \Rightarrow

$$a_t = \sqrt{\kappa} B_t$$

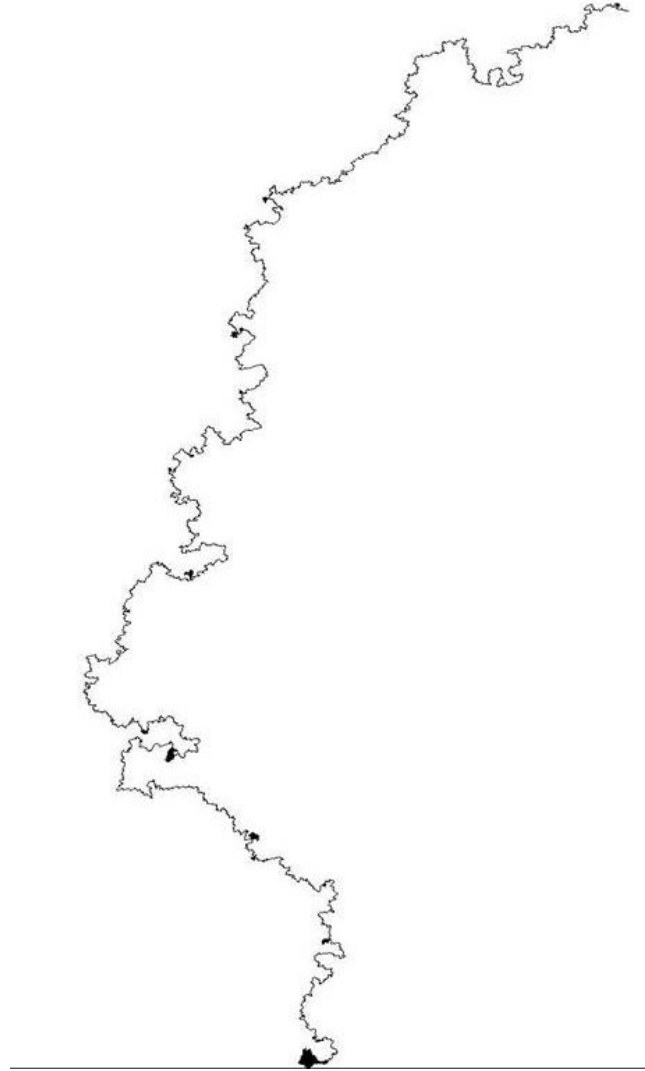
(pictures by Vincent Beffara)

$$0 < \kappa \leq 4$$

simple curve

$$d_f = 1 + \frac{\kappa}{8}$$

$$\kappa = 2$$

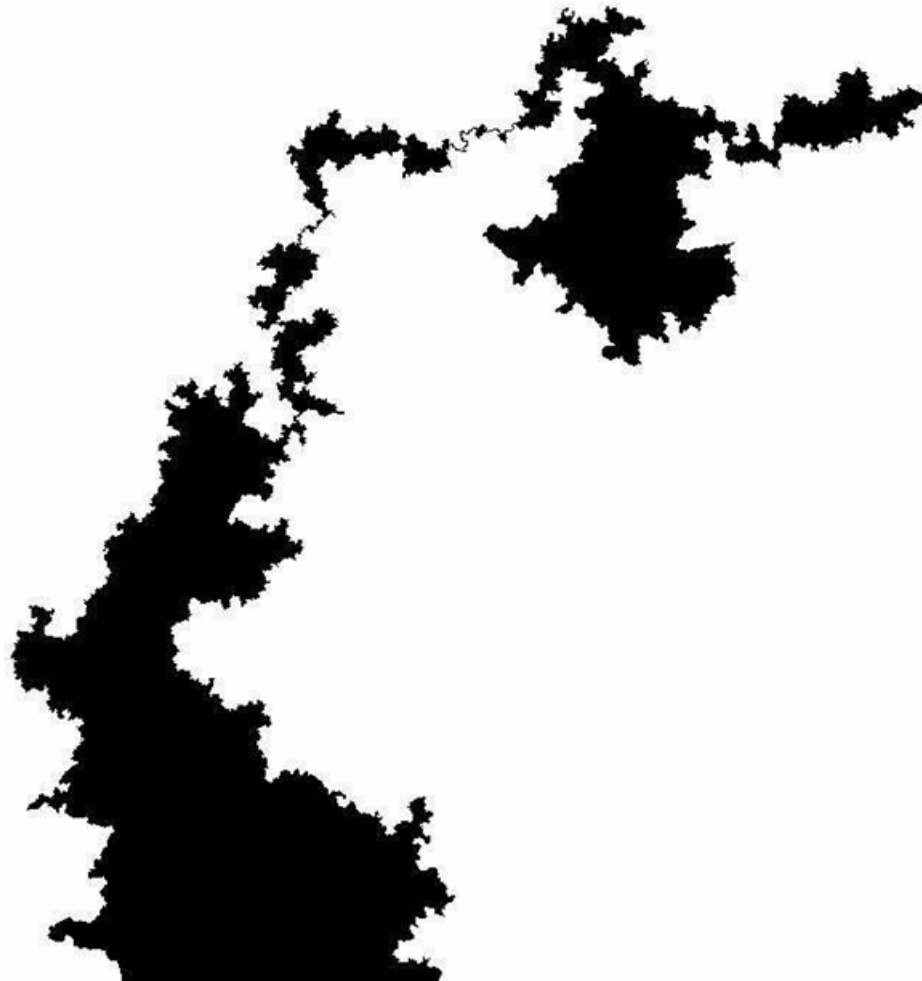


$$4 < \kappa \leq 8$$

self-intersecting curve

$$\kappa = 6$$

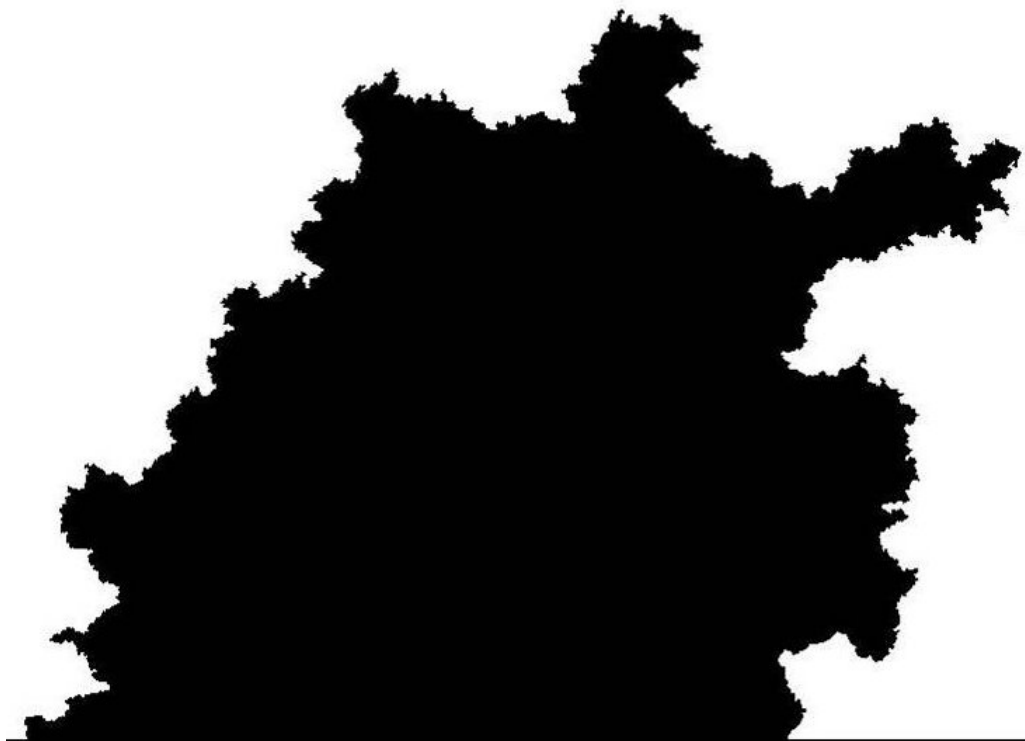
$$d_f = 1 + \frac{\kappa}{8}$$



$$\kappa > 8$$

space filling

$$d_f = 2$$



$$\kappa = 9$$

Proven results

- LERW $\kappa = 2$, spanning trees $\kappa = 8$ (LSW 2000)
- percolation $\kappa = 6$ (Smirnov 2001)
- Ising (FK) $\kappa = \frac{16}{3}$ (Smirnov 2006)

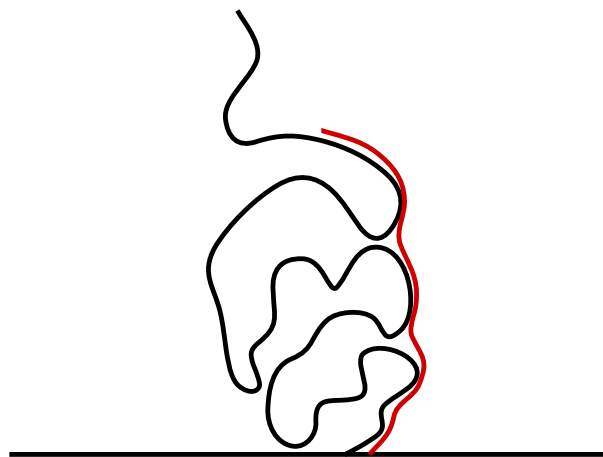
Conjectures

- self-avoiding walks $\kappa = \frac{8}{3}$
- Ising $\kappa = 3$
- Q-state Potts (FK) $\sqrt{Q} = -2 \cos \frac{4\pi}{\kappa}$

SLE duality:

(Duplantier, Belfara)

$$\kappa \longleftrightarrow \frac{16}{\kappa}$$



frontier of $\overline{K_t}$

Ising:

$$\kappa = 3$$

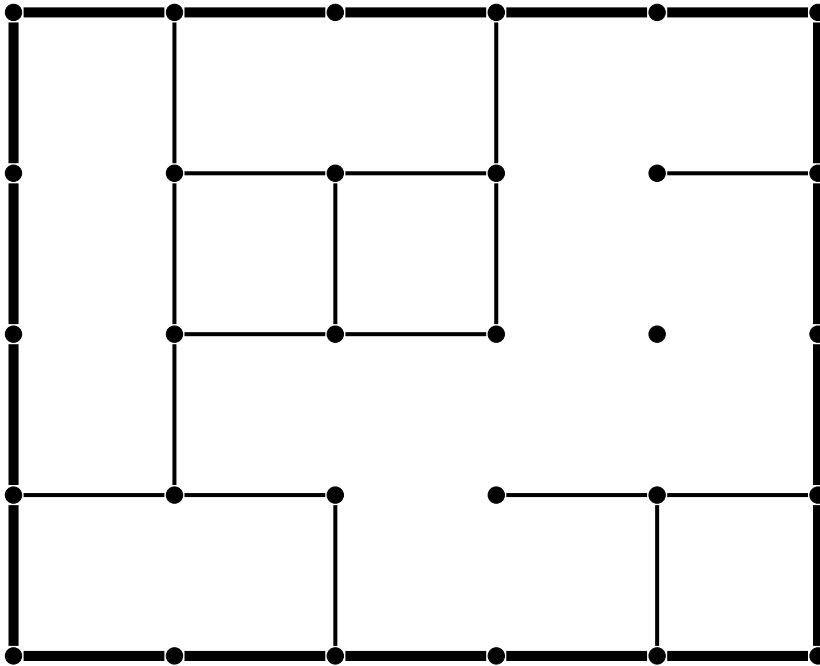


$$\kappa = \frac{16}{3}$$

FK representation

(Fortuin-Kasteleyn 1972)

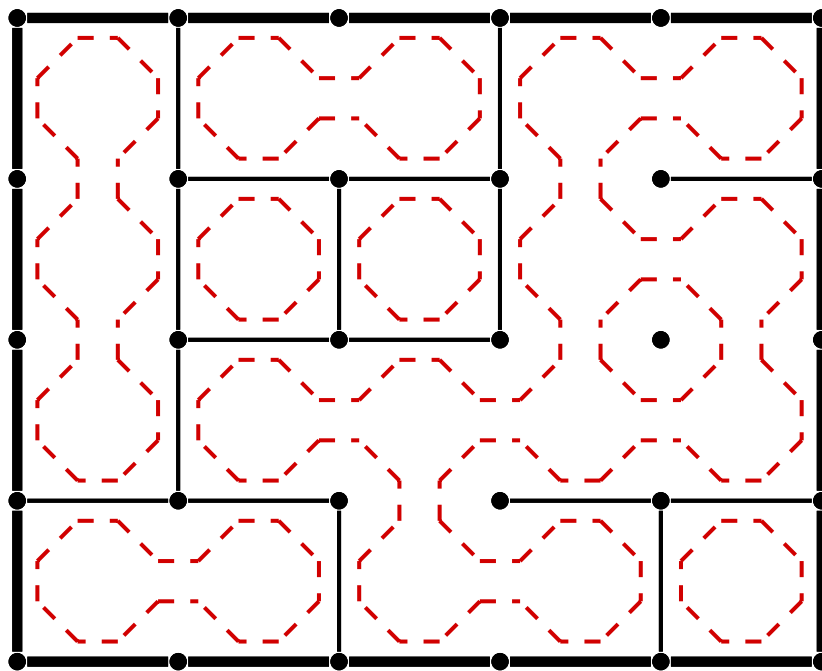
$$Z = \sum_{\{\sigma\}} e^{\sum_{\langle ij \rangle} K \sigma_i \sigma_j} = \sum_{\{\sigma\}} \prod_{\langle ij \rangle} [1 + u \delta(\sigma_i, \sigma_j)] \quad u = e^{2K} - 1$$



$$Z = \sum_G u^b 2^c$$

$$Z = \sum_G u^b 2^c = 2^{\frac{N}{2}} \sum_G \left(\frac{u}{\sqrt{2}} \right)^b \sqrt{2}^d$$

$$c - \ell = N - b$$

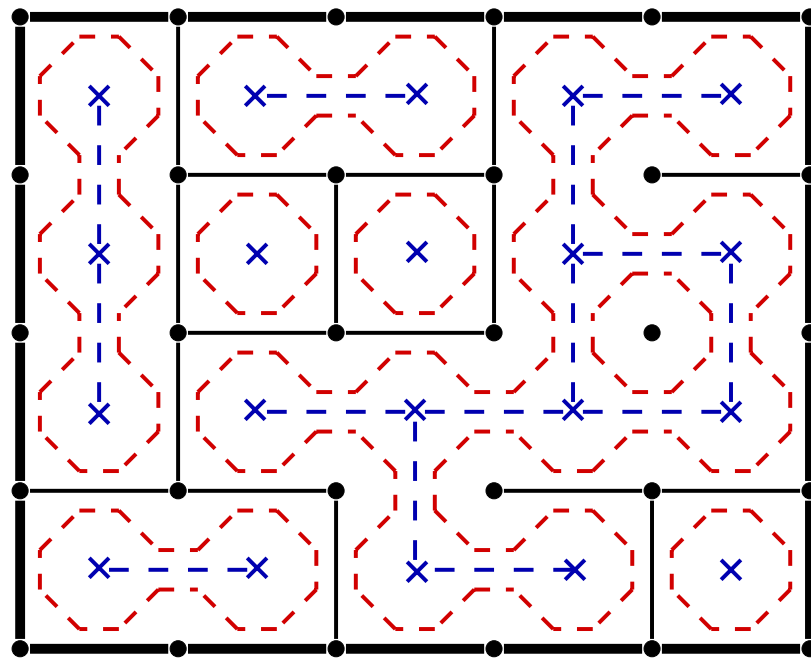


$$d = \ell + c$$

$$Z = \sum_G u^b 2^c = 2^{\frac{N}{2}} \sum_G \left(\frac{u}{\sqrt{2}} \right)^b \sqrt{2}^d$$

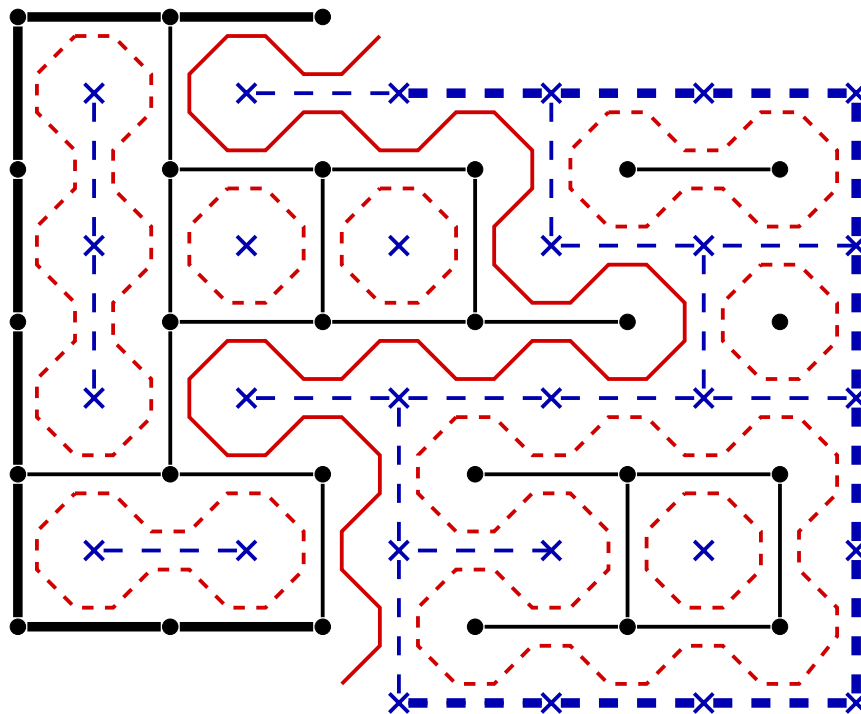
$$\propto \sum_{G^*} \left(\frac{2}{u} \right)^{b^*} 2^{c^*}$$

$u_c = \sqrt{2}$



$$d = l + c$$

Wired / free boundary conditions \longrightarrow SLE with $\kappa = \frac{16}{3}$

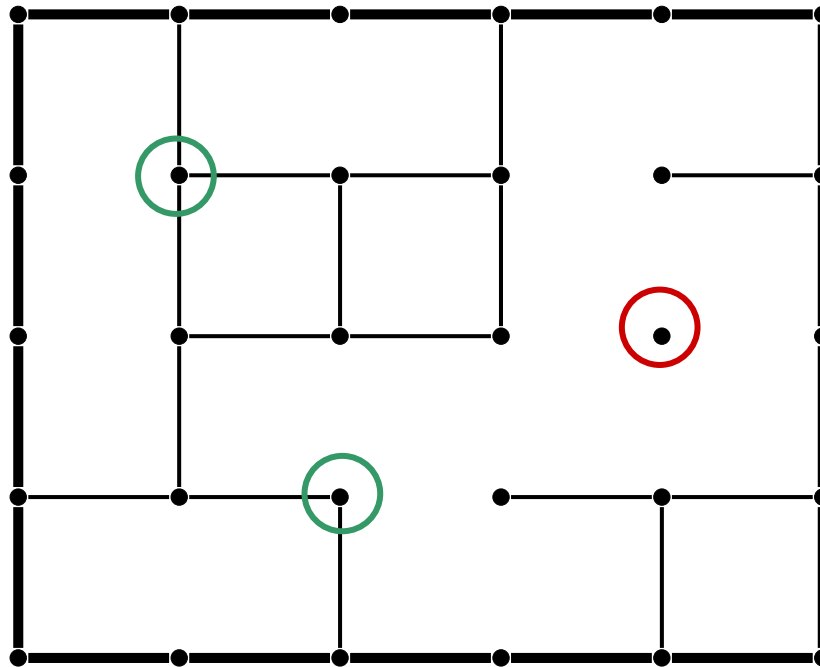


Ising fermion

Order operator

$$Z = \sum_G u^b 2^c$$

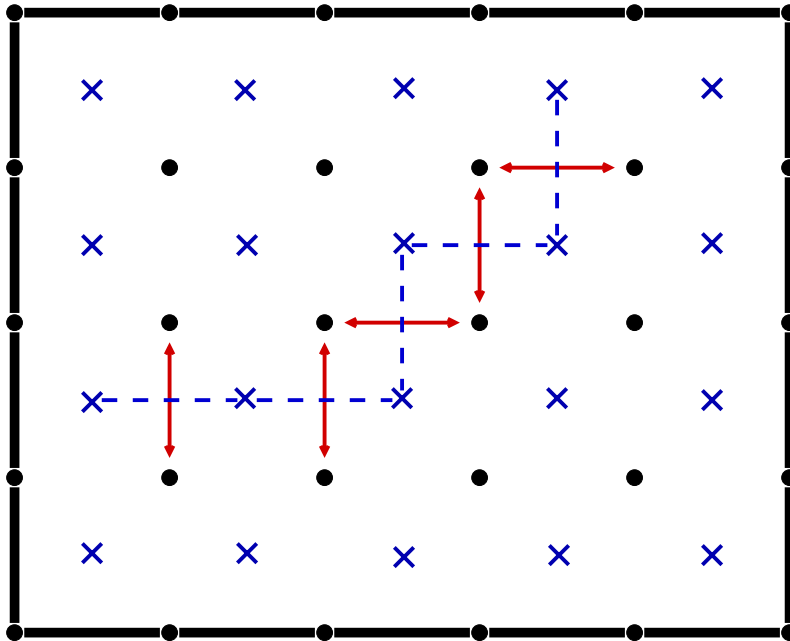
$$\langle \sigma_i \sigma_j \rangle = \frac{1}{Z} \sum_{G(i,j)} u^b 2^c$$



Disorder operator (Kadanoff, Ceva, Fradkin)

$$\langle \mu_k \mu_l \rangle \equiv \frac{Z'}{Z}$$

$$Z = \sum_{\{\sigma\}} e^{\sum_{\langle ij \rangle} K \sigma_i \sigma_j}$$



$$Z \rightarrow Z'$$

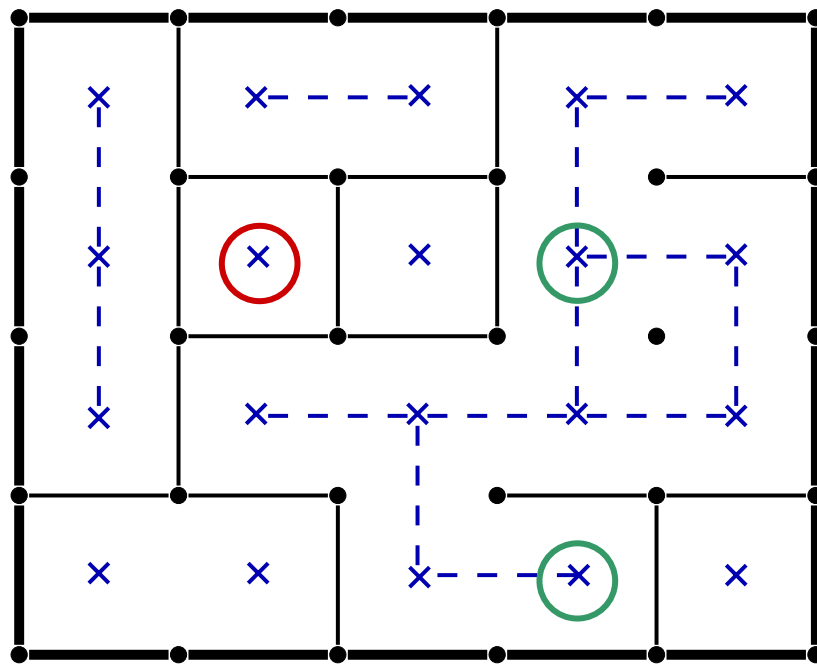
$$\uparrow$$

$$K \rightarrow -K$$

$$Z = \sum_G u^b 2^c$$

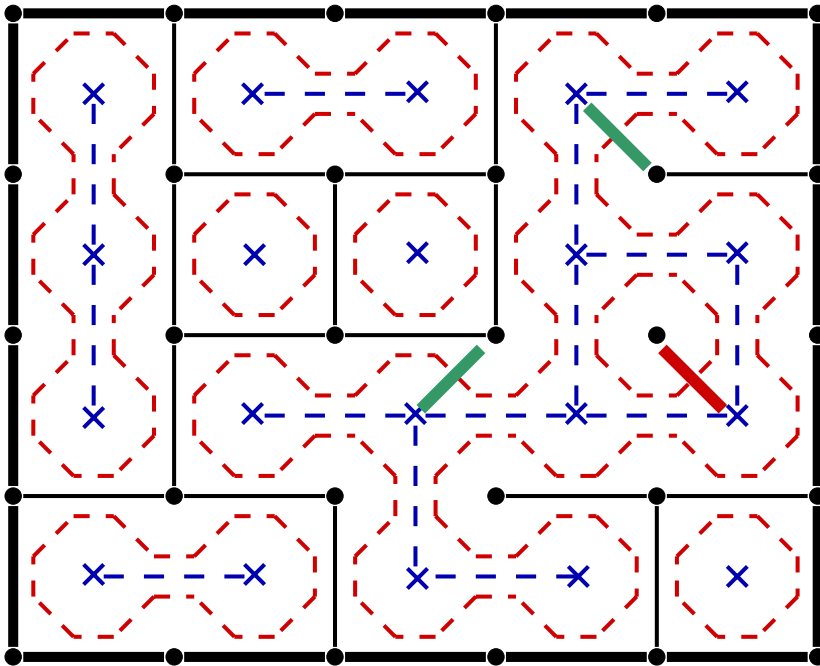
$$\langle \sigma_i \sigma_j \rangle = \frac{1}{Z} \sum_{G(i,j)} u^b 2^c$$

$$\langle \mu_k \mu_l \rangle = \frac{1}{Z} \sum_{G^*(k,l)} u^b 2^c$$



Fermion operator

$$\psi_p(e) = \sigma(e)\mu(e) e^{-ip\theta(\gamma,e)}$$



$\langle \psi_p(e_1) \psi_p(e_2) \rangle$
observable of the loops

$$\arg[\psi_p(e)] = -p\theta(\gamma, e)$$

(VR, J. Cardy)

Lattice holomorphicity

$$p = \frac{1}{2}$$

$$\sum_{e \in C} \langle \psi_p(e_1) \psi_p(e) \rangle \delta z_e = 0 \quad \Leftrightarrow \quad 2 \sin\left(p \frac{\pi}{2}\right) = \sqrt{2}$$

discretized version of $\oint_C dz \langle \psi_p(z_1) \psi_p(z) \rangle = 0$

$$\partial_{\bar{z}} \langle \psi_p(z_1) \psi_p(z) \rangle = 0$$

Fermionic Theory

$$S = \frac{1}{2\pi} \int d^2z \left(\psi \partial_{\bar{z}} \psi + \bar{\psi} \partial_z \bar{\psi} + im \bar{\psi} \psi \right) \quad m \propto K - K_c$$

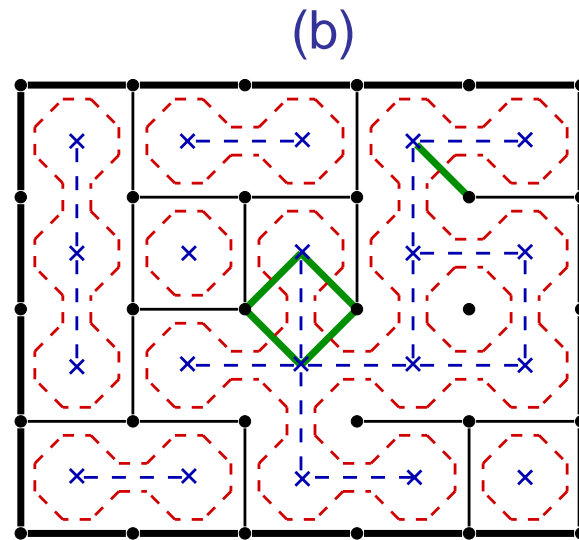
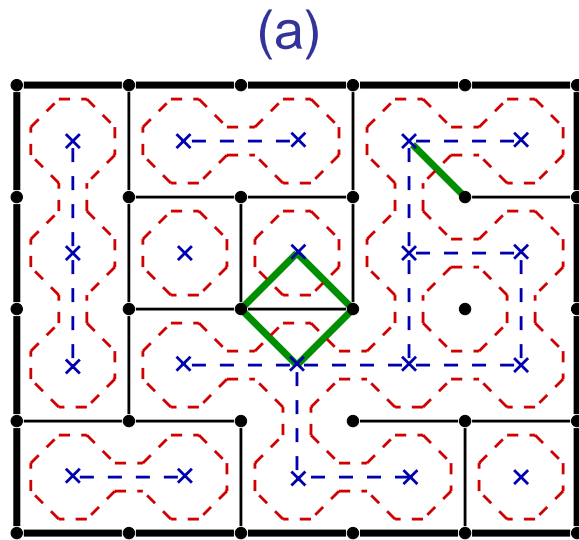
$$\partial_{\bar{z}} \psi = i \frac{m}{2} \bar{\psi}$$

$$\Delta_\psi = s_\psi = \frac{1}{2}$$

Proof:

$$p = \frac{1}{2}$$

$$\sum_{e \in C} \langle \psi_p(e_1) \psi_p(e) \rangle \delta z_e = 0 \quad \Leftrightarrow \quad 2 \sin\left(p \frac{\pi}{2}\right) = \sqrt{2}$$



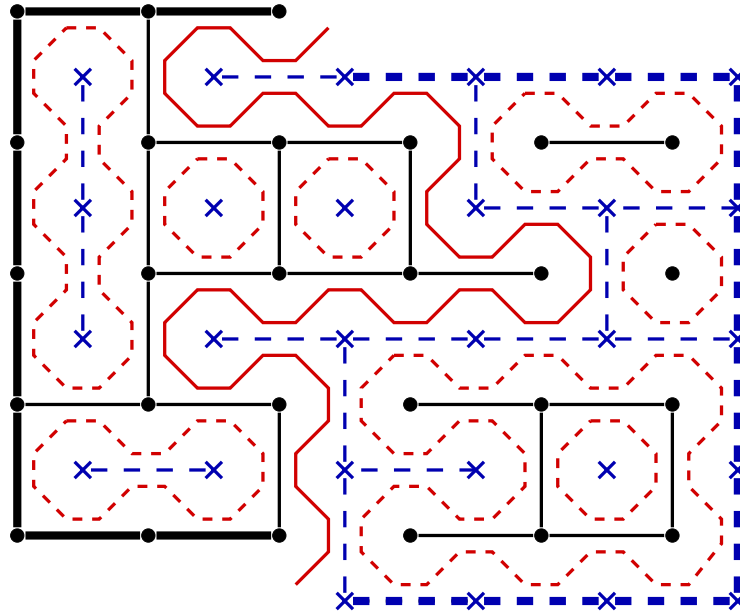
$$\left[1 + (i)e^{ip\frac{\pi}{2}} \right] \underbrace{P(a)}_{\sqrt{2}P(b)} + \left[1 + (i)e^{ip\frac{\pi}{2}} + (-1)e^{ip\pi} + (-i)e^{-ip\frac{\pi}{2}} \right] P(b) = 0$$

Connection with SLE

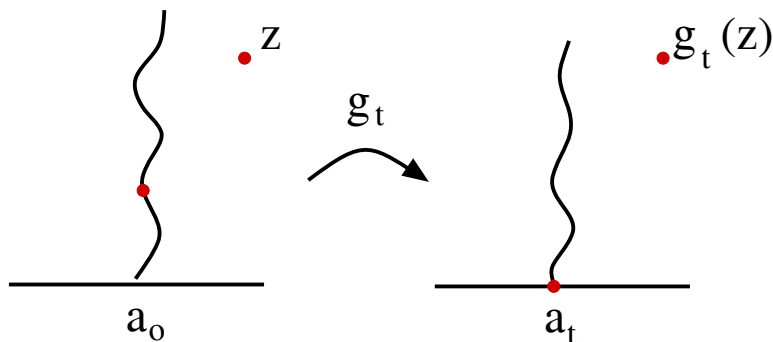
$\langle \psi_p(e) \rangle$ is an observable of the domain wall

SLE has holomorphic observable of spin 1/2 at $\kappa = \frac{8}{p+1} = \frac{16}{3}$

(Smirnov)



$P(a_0, z, \bar{z})$ = probability that the curve passes to the left / right of $z \in \mathbb{H}$

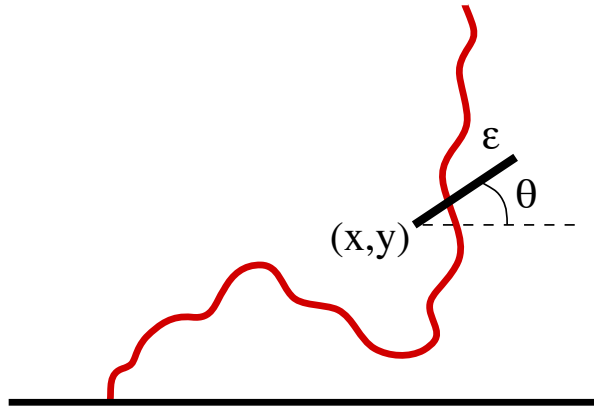


$$\frac{dg_t(z)}{dt} = \frac{2}{g_t(z) - a_t}$$

$$P(a_0, z, \bar{z}) = \left\langle P \left(a_0 + \sqrt{\kappa} dB_t, z + \frac{2dt}{z - a_0}, \bar{z} + \frac{2dt}{\bar{z} - a_0} \right) \right\rangle$$

$$\Rightarrow \left(\frac{\kappa}{2} \frac{\partial^2}{\partial a_0^2} + \frac{2}{z - a_0} \frac{\partial}{\partial z} + \frac{2}{\bar{z} - a_0} \frac{\partial}{\partial \bar{z}} \right) P(a_0, z, \bar{z}) = 0$$

(B. Doyon, V.R., J. Cardy)

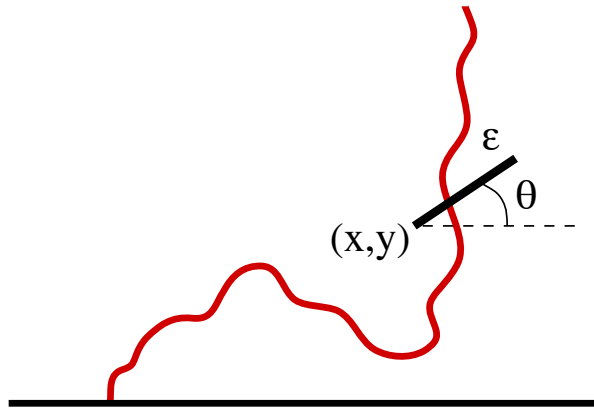


$$P(x, y, \epsilon, \theta)$$

probability that the SLE
curve passes between 2 points

$$\left\{ \frac{\kappa}{2} (\partial_w + \partial_{\bar{w}})^2 + \frac{2}{w} \partial_w + \frac{2}{\bar{w}} \partial_{\bar{w}} - \left(\frac{1}{w^2} + \frac{1}{\bar{w}^2} \right) \epsilon \partial_\epsilon + \left(\frac{1}{w^2} - \frac{1}{\bar{w}^2} \right) i \partial_\theta \right\} P(w, \bar{w}, \epsilon, \theta) = 0$$

(B. Doyon, V.R., J. Cardy)

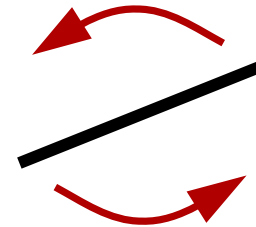


$$P(x, y, \epsilon, \theta)$$

probability that the SLE curve passes between 2 points

$$Q_p(x, y, \epsilon) = \int d\theta e^{-ip\theta} P(x, y, \epsilon, \theta)$$

p has physical meaning of spin



expansion in small ϵ : $Q_p \sim \epsilon^{\Delta_p}$

$$\left\{ \frac{\kappa}{2} (\partial_w + \partial_{\bar{w}})^2 + \frac{2}{w} \partial_w + \frac{2}{\bar{w}} \partial_{\bar{w}} - \left(\frac{1}{w^2} + \frac{1}{\bar{w}^2} \right) \epsilon \partial_\epsilon + \left(\frac{1}{w^2} - \frac{1}{\bar{w}^2} \right) i \partial_\theta \right\} P(w, \bar{w}, \epsilon, \theta) = 0$$

$$\left\{ \frac{\kappa}{2} (\partial_w + \partial_{\bar{w}})^2 + \frac{2}{w} \partial_w + \frac{2}{\bar{w}} \partial_{\bar{w}} - \left(\frac{1}{w^2} + \frac{1}{\bar{w}^2} \right) \epsilon \partial_\epsilon - p \left(\frac{1}{w^2} - \frac{1}{\bar{w}^2} \right) \right\} Q_p(w, \bar{w}, \epsilon) = 0$$

suppose $Q_p \sim \epsilon^p$

$Q_p(w, \bar{w}, \epsilon) = \text{const} \times \left(\frac{\epsilon}{w} \right)^p$ is solution if

$$\kappa = \frac{8}{p+1}$$

Conclusions

- Other models (Q-state Potts,...): parafermions
- Off-critical $\partial_{\bar{z}} \psi = i \frac{m}{2} \bar{\psi}$
- Loops \longrightarrow CLE (Sheffield, Werner)