

Reconnection of Twisted Flux Tubes

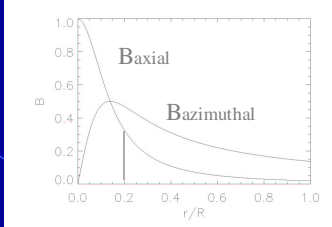
Mark Linton (NRL/ITP)
Russ Dahlburg (NRL)
Spiro Antiochos (NRL)

Motivation

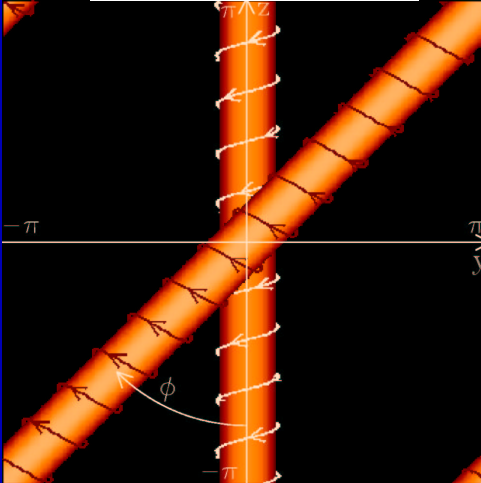
- Sonnerup (1974) 2.5D theory for reconnection of two sheets of equal field strengths:
Field reconnects at any angle,
speed of reconnection slows down as $\sin(\theta/2)$.
Is this true in 3D?
- Flare models for colliding/reconnecting flux tubes: Gold & Hoyle (1960), Heyvaerts et al. (1977), Sturrock et al. (1984), Priest et al. (1994), Farnik et al. (1996).
Are they viable?
- Flux tubes are building blocks of magnetic field:
in convection zone, solar atmosphere, solar wind or Earth's magnetosphere.
Flux tube reconnection should be generally applicable.

Simulation setup

- Two identical Gold-Hoyle magnetic flux tubes
Twist: $qL/2\pi = 10$.
- Back tube R twisted
- Front tube R or L twisted,
- Collision angle $\phi = n\pi/4$:
to give RLn or RRn
- Visco-resistive MHD,
Lundquist numbers ~ 600
- Periodic box, spectral code
- Push tubes together with initial
stagnation point flow at $1/40$ of
Alfven speed on axis



The graph shows the radial profiles of the magnetic field components. The vertical axis is labeled 'B' and ranges from 0.0 to 1.0. The horizontal axis is labeled 'r/R' and ranges from 0.0 to 1.0. Two curves are shown: 'Baxial' and 'Bazimuthal'. The Baxial curve starts at 1.0 at r/R = 0 and decreases to 0 at r/R = 1. The Bazimuthal curve starts at 0 at r/R = 0, peaks at approximately 0.5 around r/R = 0.2, and then decreases to 0 at r/R = 1.



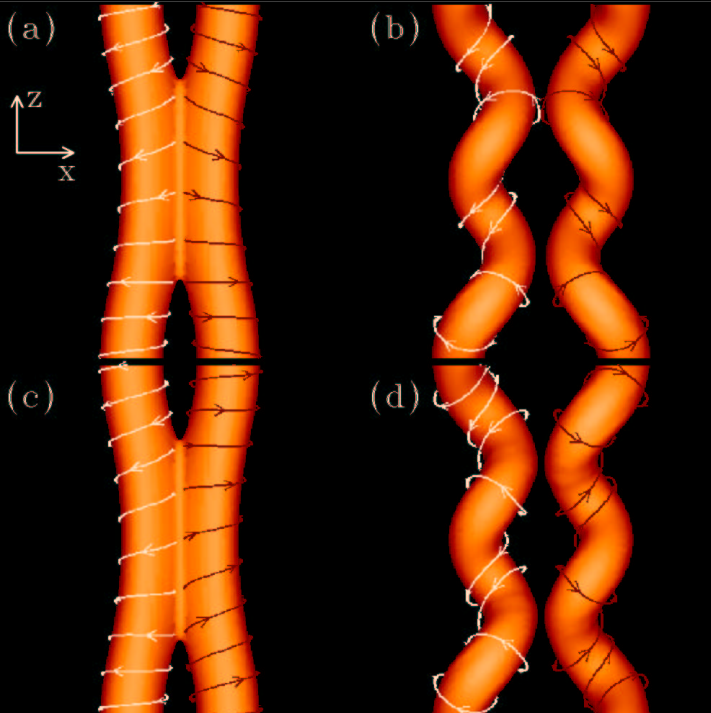
A 3D visualization of two orange magnetic flux tubes. The tubes are twisted, with white lines representing the magnetic field lines. The tubes are oriented along the z-axis. The angle between the tubes is labeled ϕ . The z-axis is labeled with $-\pi$ and π .

RL0, RR4 Bounce

Top: RL0

Bottom: RR4

Magnetic
isosurfaces:
 $1/3$ of max
at 85 and 280
Alfven
crossing times

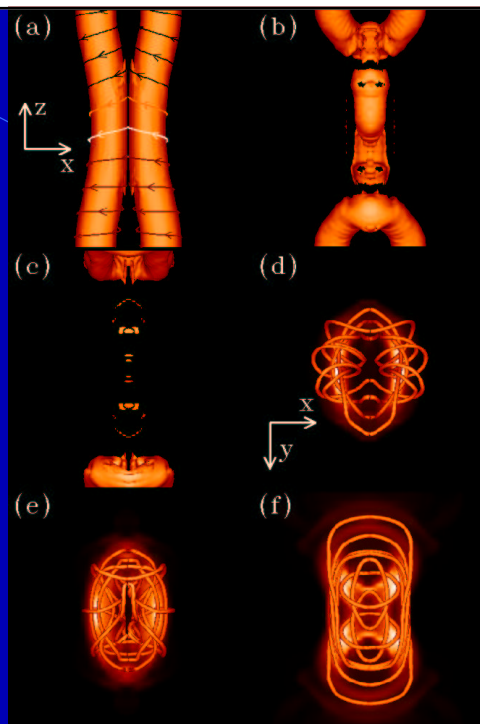


Four panels (a, b, c, d) showing the evolution of magnetic flux tubes. Panel (a) shows two tubes approaching each other. Panel (b) shows the tubes in contact and beginning to reconnect. Panel (c) shows the tubes after reconnection, with a central region where the tubes have merged. Panel (d) shows the tubes after a bounce, with the tubes moving away from each other. A coordinate system with x and z axes is shown in panel (a).

RL4 Slingshot

Similar to that seen by
Yamada et al 1990
Ozaki & Sato 1990
Lau & Finn 1996
Kondrashov et al 1999

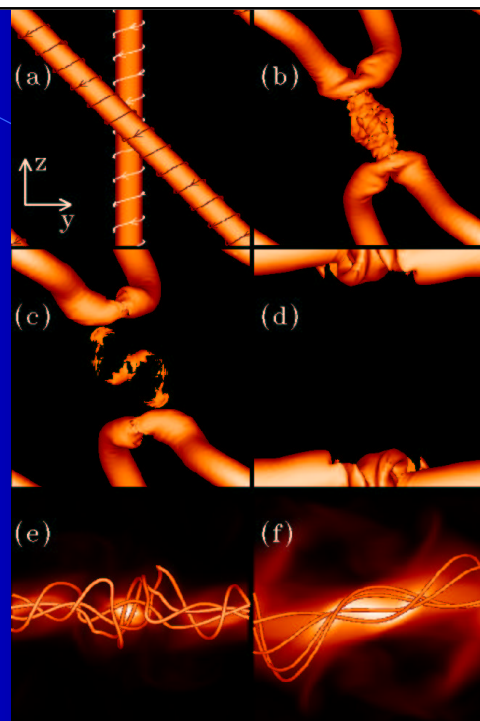
Isosurfaces at times: (85,200,250)
Fieldlines at times:(250,340,830)

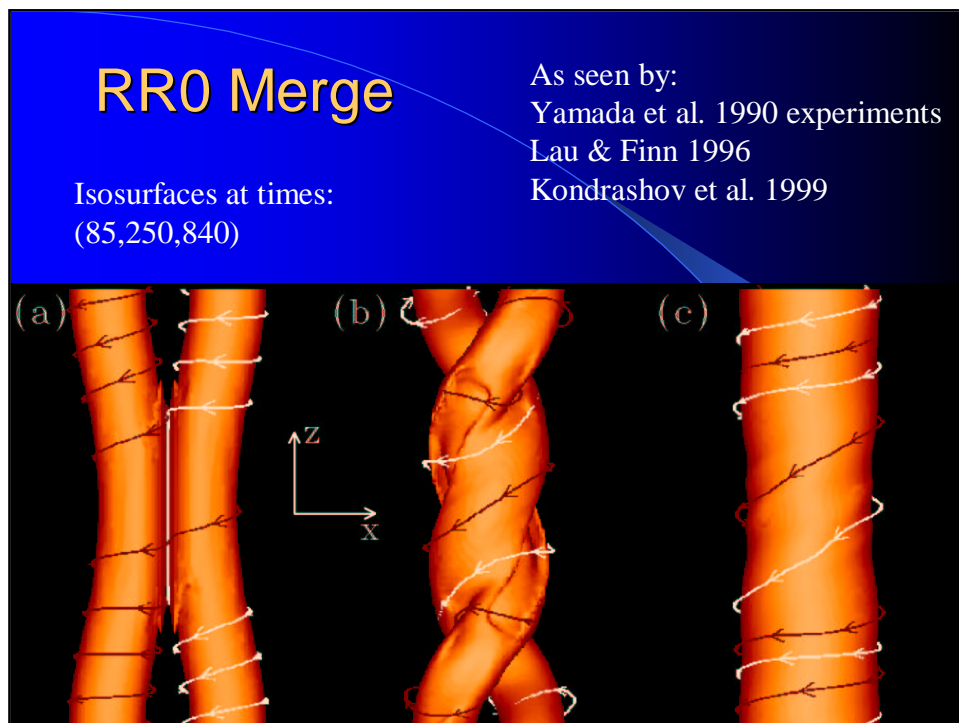
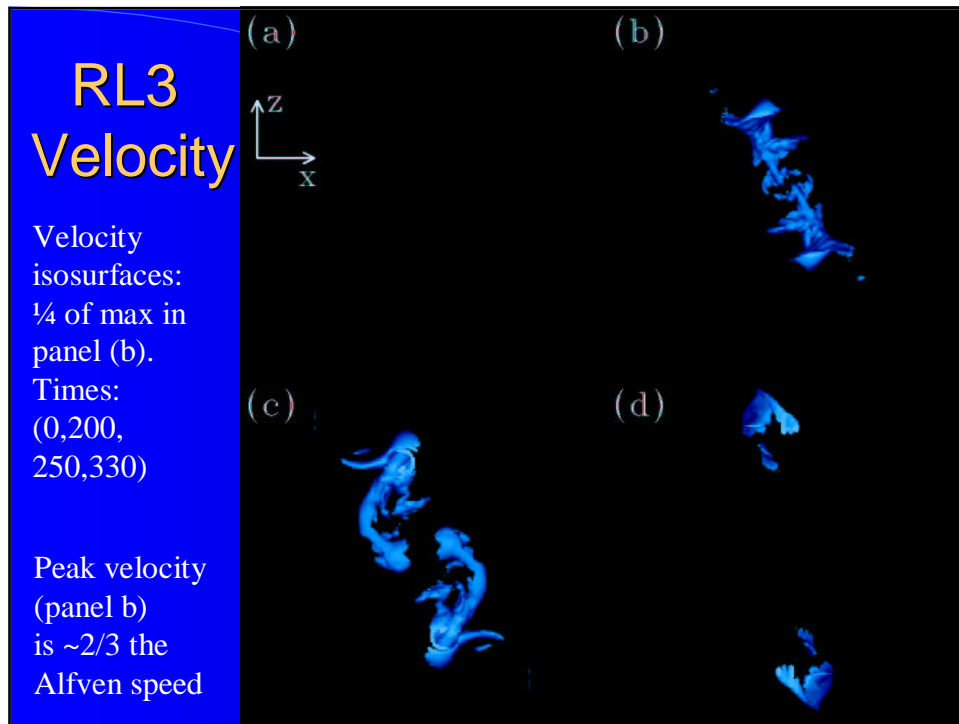


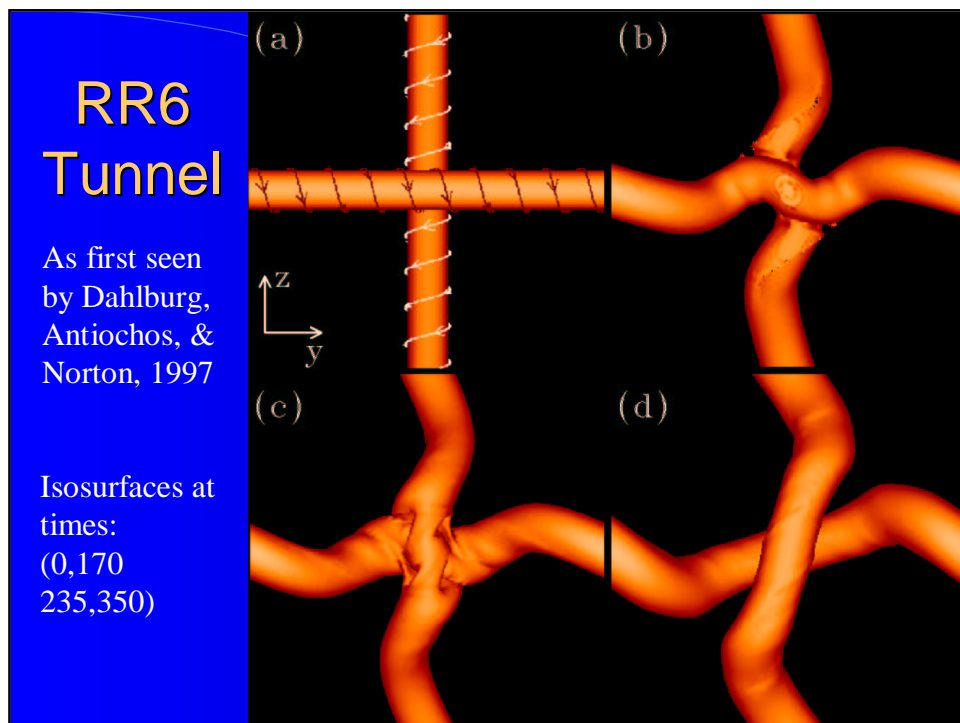
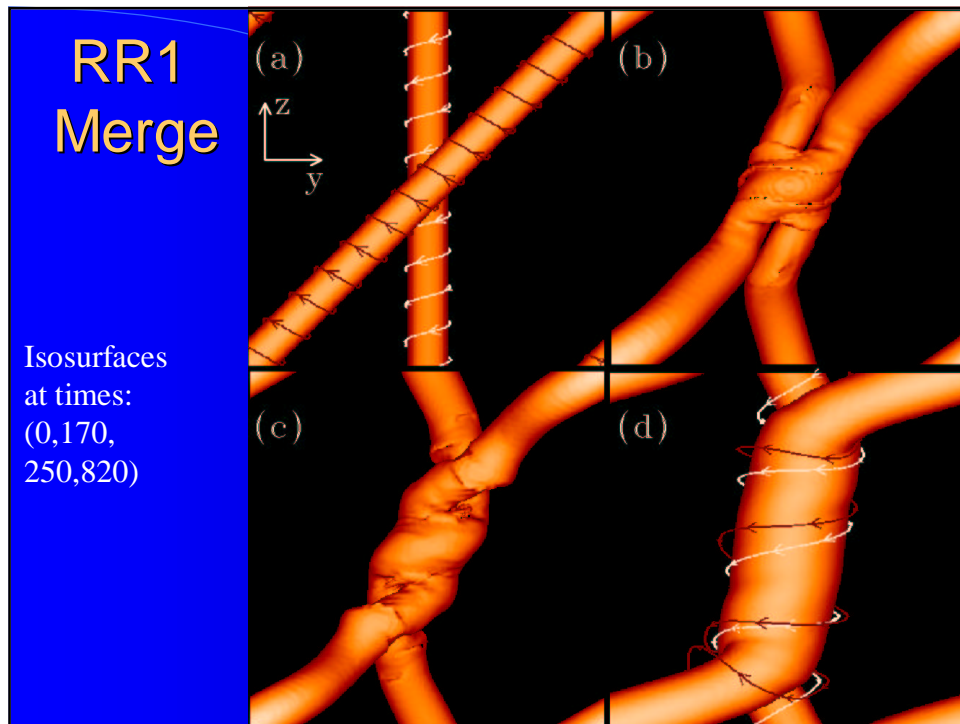
RL3 Slingshot

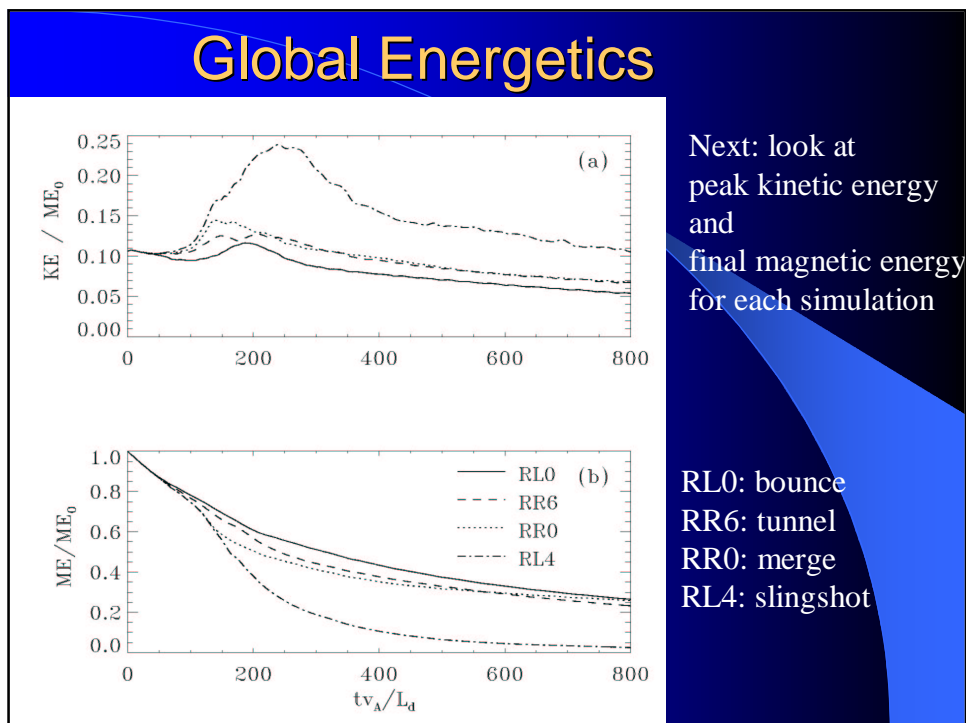
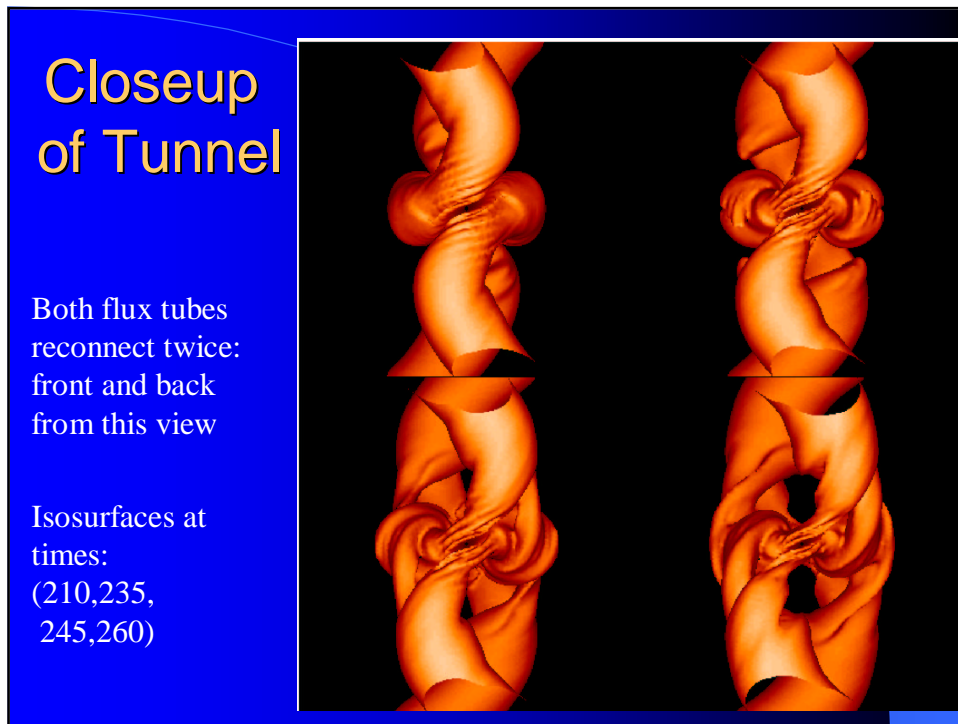
Similar to the flare
model proposed by Gold
& Hoyle (1960)

Isosurfaces at times:
(0,200,250,790)
Fieldlines at times:
(790,1950)





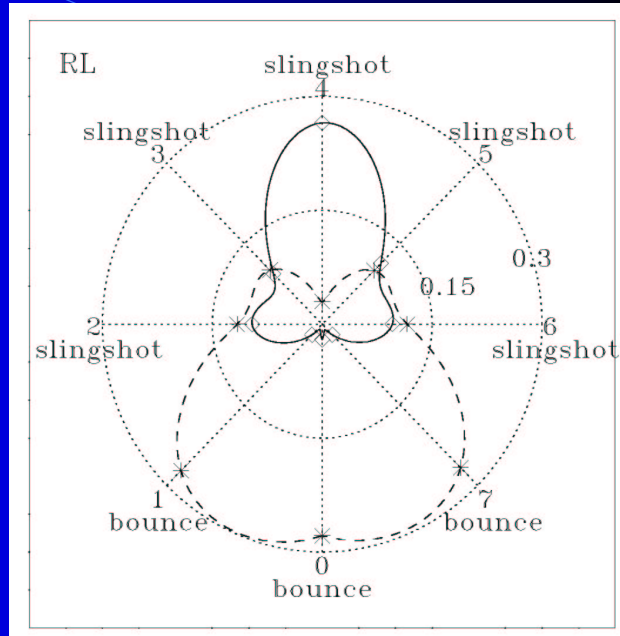




RL Reconnection Energy

Solid curve:
twice the peak
kinetic energy

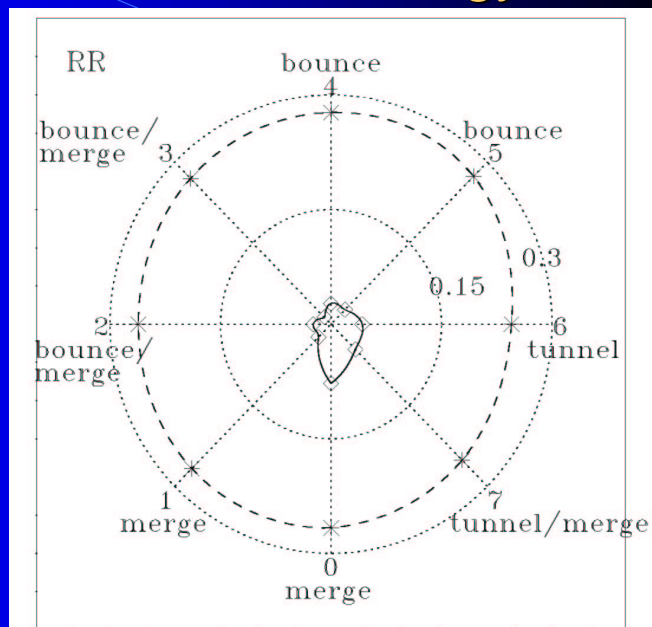
Dashed curve:
magnetic energy
remaining at end of
simulation
(750 Alfvén times)



RR Reconnection Energy

Solid: kinetic

Dashed: magnetic



Relative Helicity for a pair of flux tubes

Following Berger 1984, Wright & Berger 1989

$$H_r = \int_V (\mathbf{A} + \mathbf{A}') \cdot (\mathbf{B} - \mathbf{B}') d^3x = 2H_m + 2H_t = \text{const.}$$

$$H_m = \frac{\delta}{2\pi} \Phi^2$$

Mutual helicity depends on the angle by which the tubes rotate about each other.

$$H_t = \frac{qL}{2\pi} \Phi^2 = N\Phi^2$$

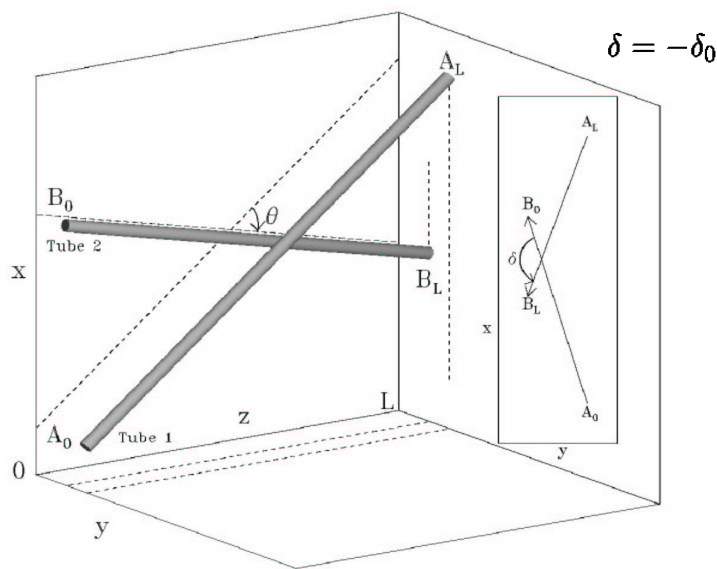
Twist helicity depends on the angle by which the fieldlines rotate about the tube axis.

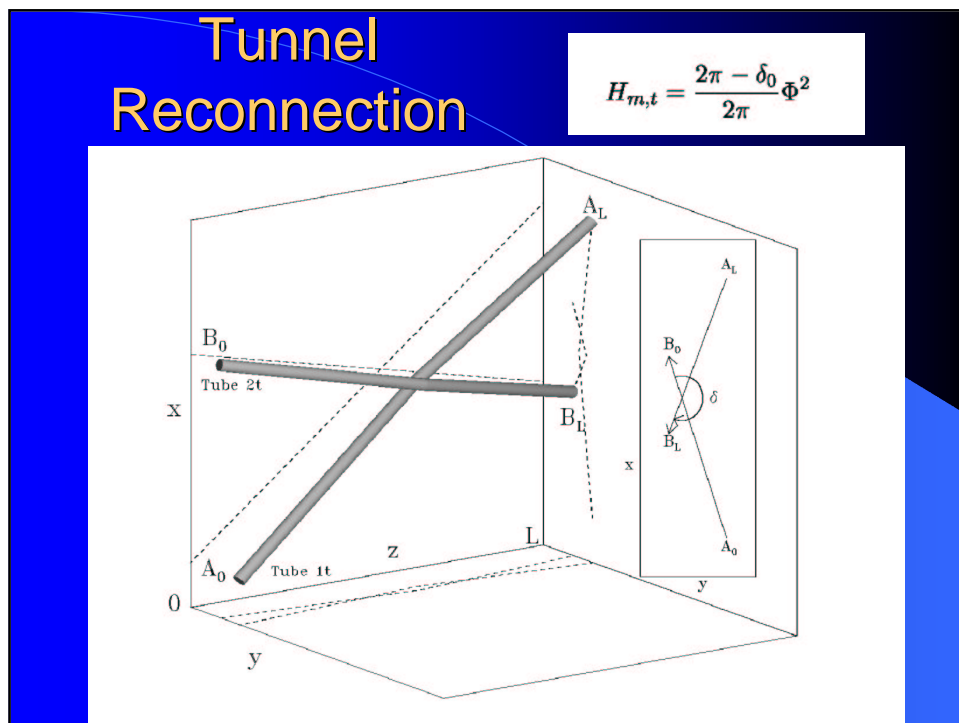
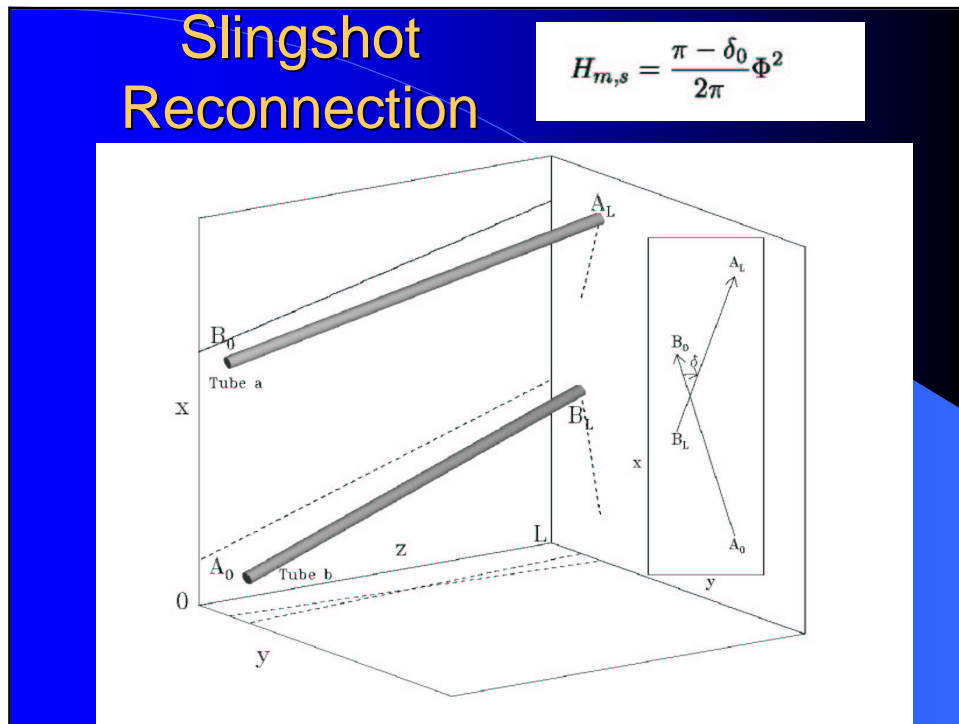
So that:
$$H_r = \Phi^2 \left(2N + \frac{\delta}{\pi} \right)$$

Assume helicity is conserved in reconnection:
sum of mutual and twist helicity is constant.

Initial State:

$$H_{m,0} = \frac{-\delta_0}{2\pi} \Phi^2$$





Helicity conservation determines the post-reconnection twist:

Initial helicity

$$H_0 = \left(2N_0 - \frac{\delta_0}{\pi}\right) \Phi^2$$

Helicity conservation:

$$H_0 = H_s = H_t$$

Slingshot helicity

$$H_s = \left(2N_s - \frac{\delta_0}{\pi} + 1\right) \Phi^2$$

Tubes each lose half a turn of twist

$$N_s = N_0 - \frac{1}{2}$$

Tunnel helicity

$$H_t = \left(2N_t - \frac{\delta_0}{\pi} + 2\right) \Phi^2$$

Tubes each lose a full turn of twist

$$N_t = N_0 - 1$$

Assumptions for equilibrium states

Flux tube field profile:
Same form before and after reconnection, but constants can change.

$$B_{axial} = bf(\lambda r)$$

$$B_{azimuthal} = qr B_{axial}$$

Helicity conservation

$$N_0 = N_s - \frac{1}{2} = N_t - 1$$

Mass conservation

$$\rho L \pi R^2 = \rho L_0 \pi R_0^2$$

Flux conservation

$$\begin{aligned} \Phi = \Phi_0 &= b_0 \int_0^{R_0} 2\pi f(\lambda_0 r) r dr = \frac{2\pi b_0}{\lambda_0^2} \Gamma_{11}(\lambda_0 R_0) \\ &= \frac{2\pi b}{\lambda^2} \Gamma_{11}(\lambda R) \end{aligned}$$

For identical initial tubes can show:

$$\lambda R = \lambda_0 R_0,$$

Predicted Energies for RR6

Initial state:

$$E'_0 = \Gamma_{12} + \Gamma_{32}$$

Slingshot:

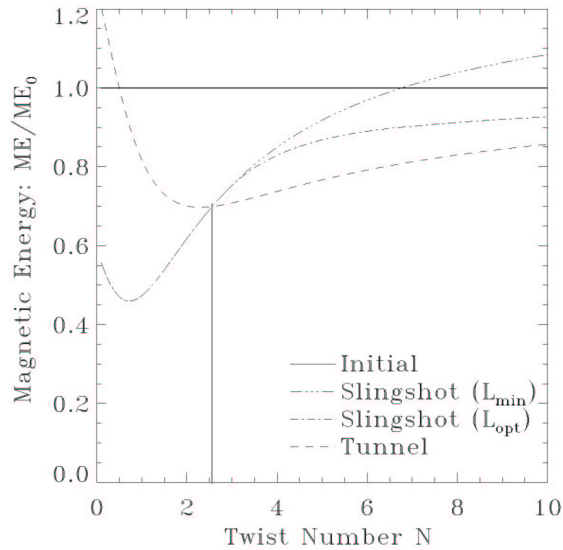
$$E'_s = \frac{L_s^2}{L_0^2} \Gamma_{12} + \frac{L_0 (N_0 - 1/2)^2}{L_s N_0^2} \Gamma_{32}$$

Tunnel:

$$E'_t = \Gamma_{12} + \frac{(N_0 - 1)^2}{N_0^2} \Gamma_{32}$$

Normalization:

$$E' = 2q_0^2 E / (L_0 b_0^2)$$



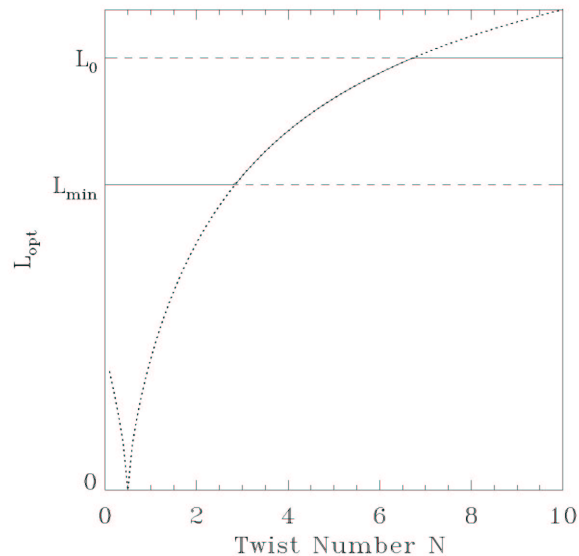
Optimal Length for Slingshot

1) Shortest length allowed by geometry:

$$L_{min} = L_0 / \sqrt{2}$$

2) Optimal length, minimizing energy

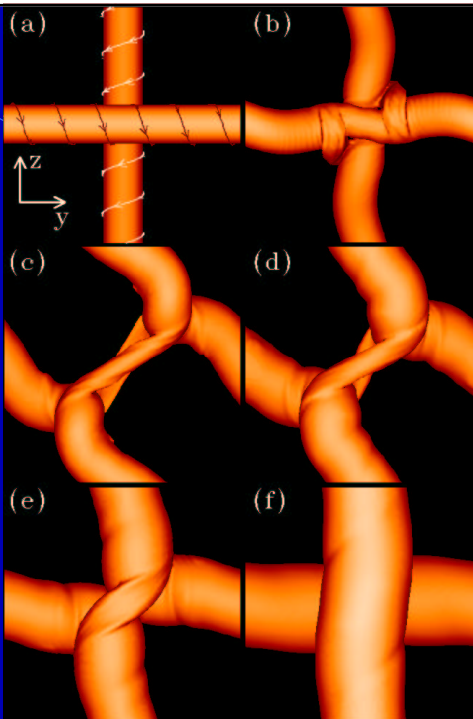
$$L_{opt} = L_0 \left(\frac{N_s^2 \Gamma_{32}}{2N_0^2 \Gamma_{12}} \right)^{1/3}$$



**RR6 N=6:
Tunnel**

- Still tunnels at N=6
- Comes close to slingshot, however

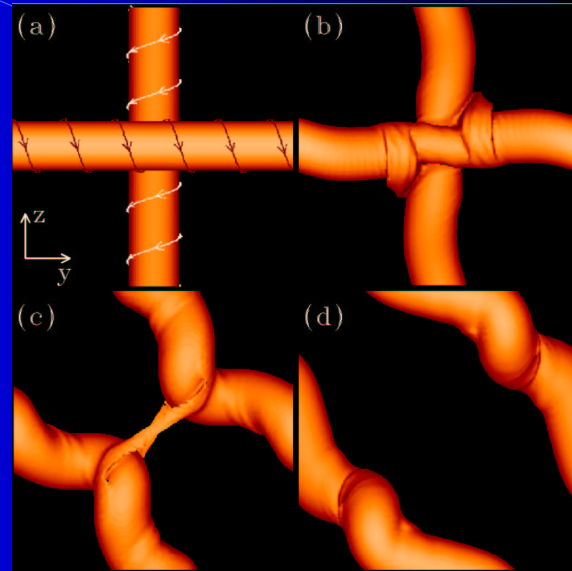
Isosurfaces at times:
(0,150,
480,789,
1115,2700)



**RR6 N=5.5:
Slingshot**

- Transition at N=5.5 higher than prediction of N=2.5
- Final state kinked: minimum energy length is longer than the shortest possible length

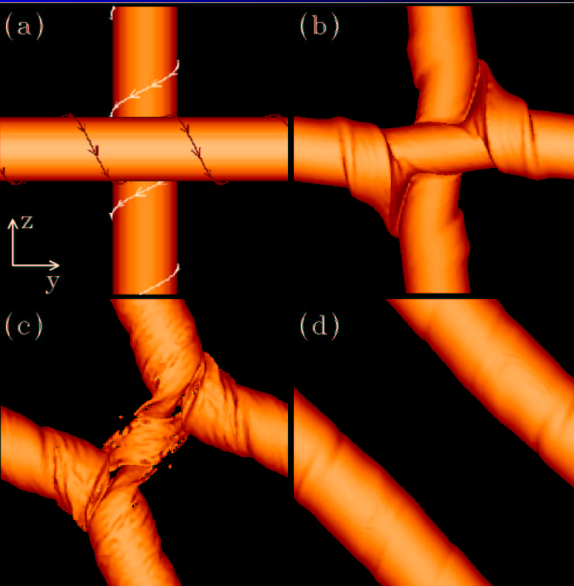
Isosurfaces at times:
(0,150,290,450)



RR6 N=3: Slingshot

- Tubes no longer kink at this twist: optimal length is shortest length

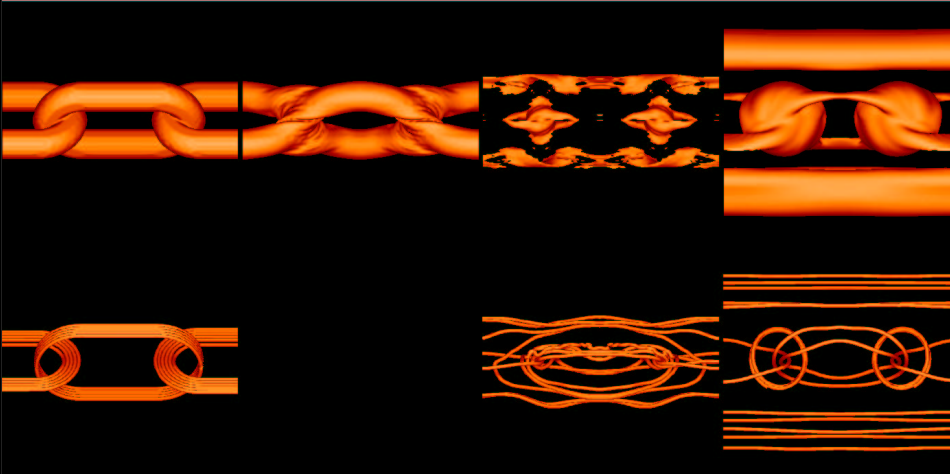
Isosurfaces at times:
(0,140,220,500)



Zero twist reconnection

- Energetically, tubes should tunnel
- In fact $\sim 2/3$ of flux slingshots
- Tunnel is inaccessible to fieldlines?

Isosurfaces at times
 $\sim(0,2,12,60)$



Conclusions

- Four types of twisted flux tube reconnection: bounce, slingshot, merge, and tunnel.
- Slingshot is most energetic.
- Helicity conservation allows one to predict the twist, and therefore the energy of reconnected flux tubes.
- Energy calculation predicts the tunnel will happen at large twist for RR6, while the slingshot will occur at low twist.
- RR6 transition occurs, but at higher twist than predicted: due to helicity loss?
- Tunnel does not occur for zero twist flux tubes, even where is energetically advantageous: fieldlines are unable to reconnect twice in untwisted tubes?