

Spin Control Based on the Rashba Spin-orbit Interaction

- Rashba spin-orbit interaction
- Competition between Zeeman and Rashba
- Spin interference experiments
- Stern-Gerlach Spin Filter



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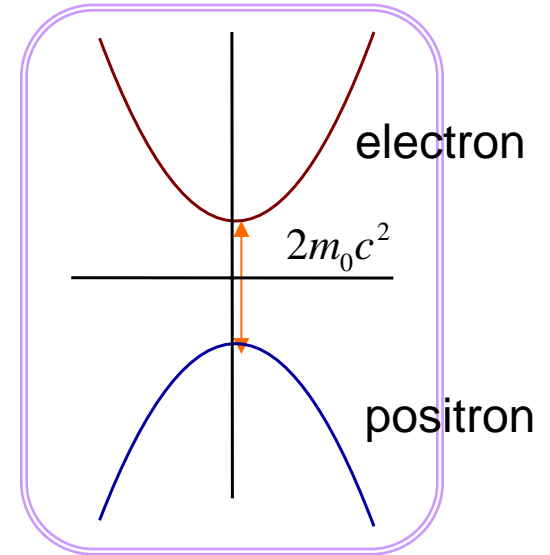
Enhancement of spin-orbit interaction

SOI in vacuum

$$H_{so} = -\mu_B \sigma \cdot \left(\frac{p \times E}{2m_0 c^2} \right) \rightarrow \mathbf{B}_{eff}$$

Dirac gap

$$2m_0 c^2 \approx 1 \text{ MeV}$$

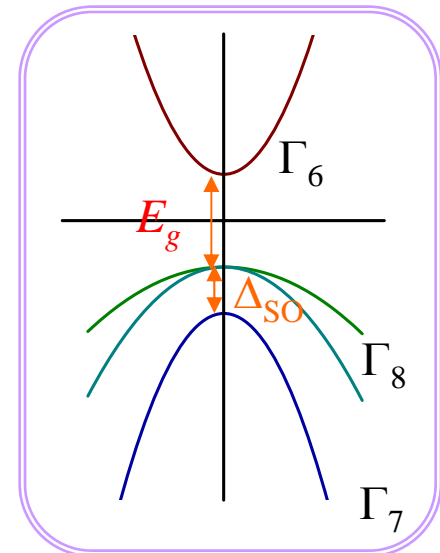


Rashba SOI in semiconductors

$$H_R = \frac{eP^2}{3} \left[\frac{1}{E_g^2} - \frac{1}{(E_g + \Delta_{SO})^2} \right] \sigma \cdot k \times \langle E \rangle_v$$

Energy gap

$$E_g, \Delta_{SO} \approx 1 \text{ eV}$$

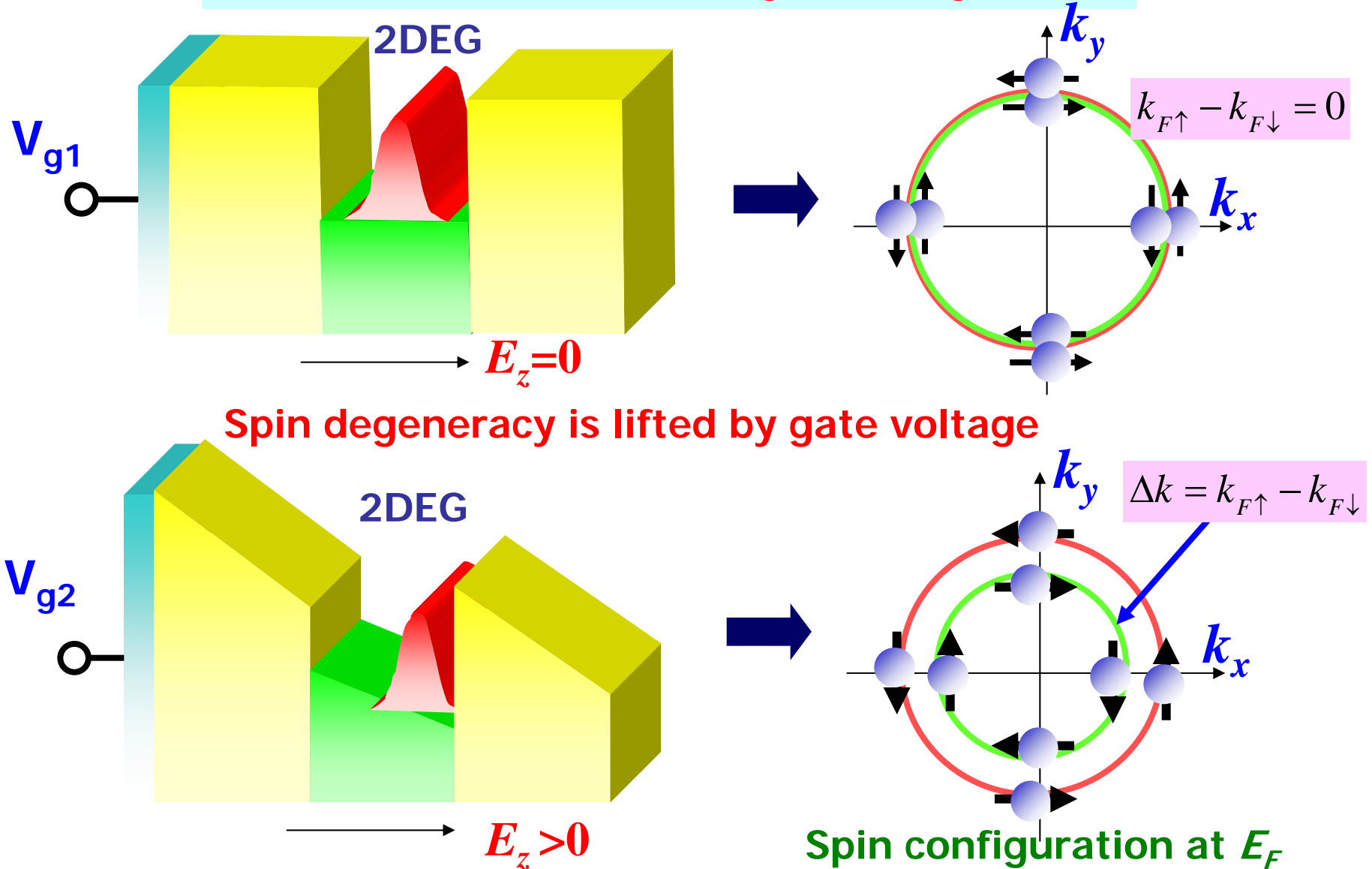


SOI

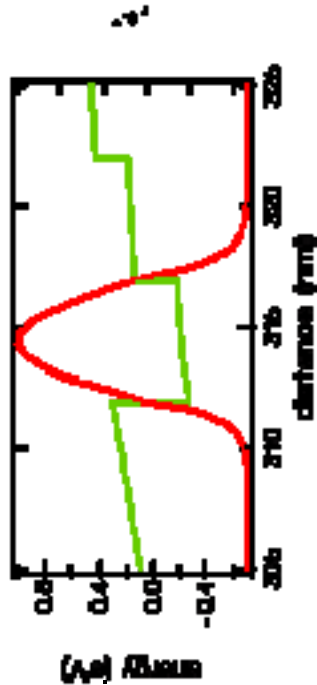
1. Narrow Gap semiconductor
2. Electric Field in QW

Rashba spin-orbit interaction in 2DEG

Structural inversion asymmetry (SIA)



Potential Profile in InGaAs Quantum Well

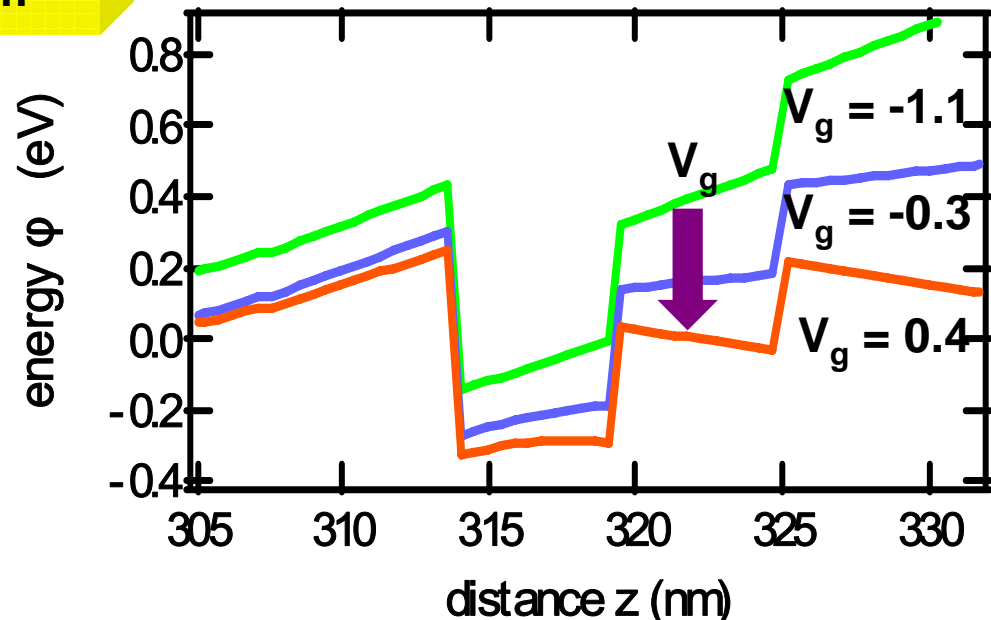


SOI depends on V_g

V_g	$\langle E \rangle$	α
-1.1	large	large
-0.3	medium	medium
0.4	small	small

SOI is proportional to E

$$\alpha \propto \langle E \rangle = -\frac{\partial \phi(z)}{\partial z}$$

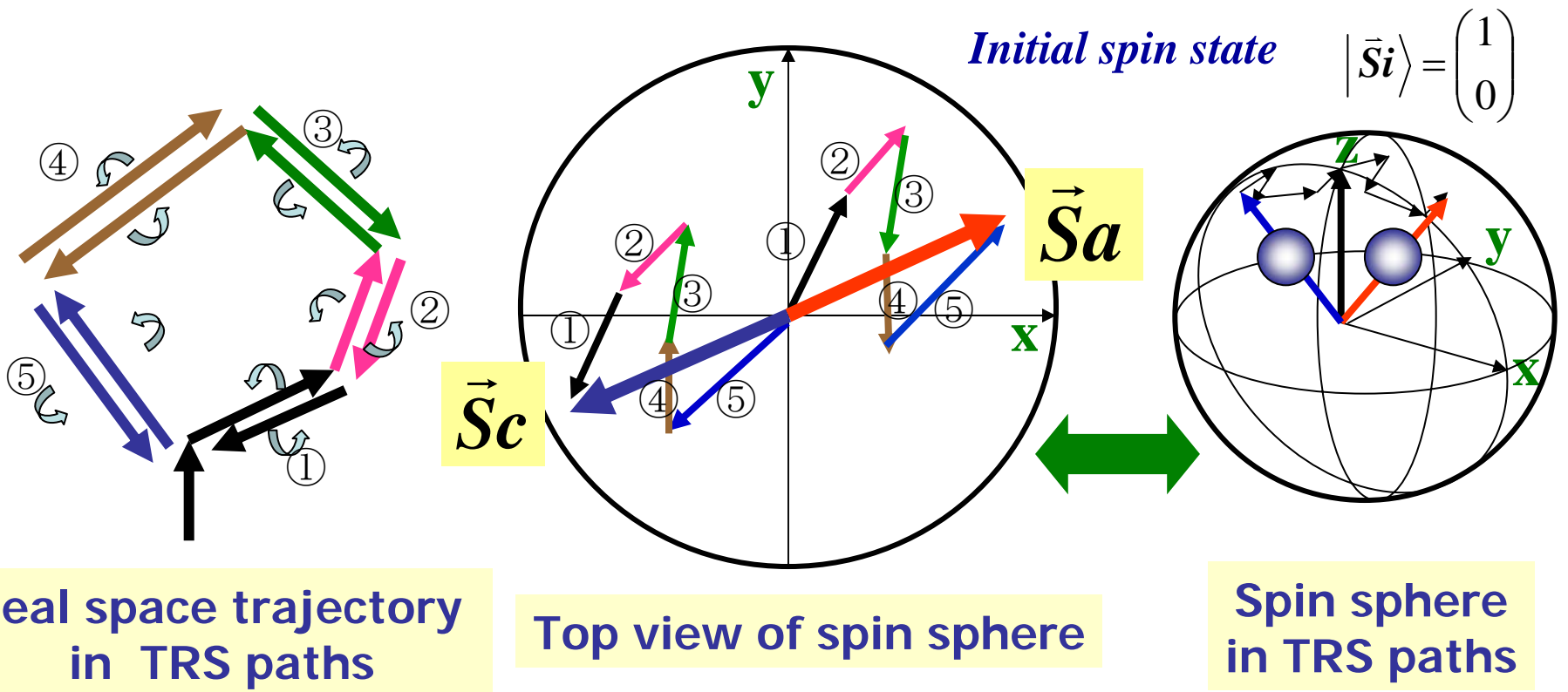


Spin states in Time Reversal Symmetry Paths

SOI does not break TRS

$$TRS \Rightarrow p \rightarrow -p, \sigma \rightarrow -\sigma, E \rightarrow E$$

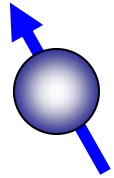
$$H_{SO} = -\mu_B \sigma \cdot (p \times E)$$



Final spin directions are exactly opposite

Destructive interference due to SOI

Final spin states → Exactly opposite



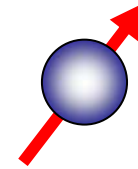
$$|\vec{S}_c\rangle = R|\vec{S}_i\rangle$$

Clockwise



$$|\vec{S}_a\rangle = R^{-1}|\vec{S}_i\rangle$$

Anti-clockwise



$$R(\alpha, \beta, \gamma) = \begin{bmatrix} \cos \frac{\alpha}{2} e^{i(\beta+\gamma)/2} & i \sin \frac{\alpha}{2} e^{-i(\beta-\gamma)/2} \\ i \sin \frac{\alpha}{2} e^{i(\beta-\gamma)/2} & \cos \frac{\alpha}{2} e^{-i(\beta+\gamma)/2} \end{bmatrix}$$

Euler angle (α, β, γ)

Rotational operator

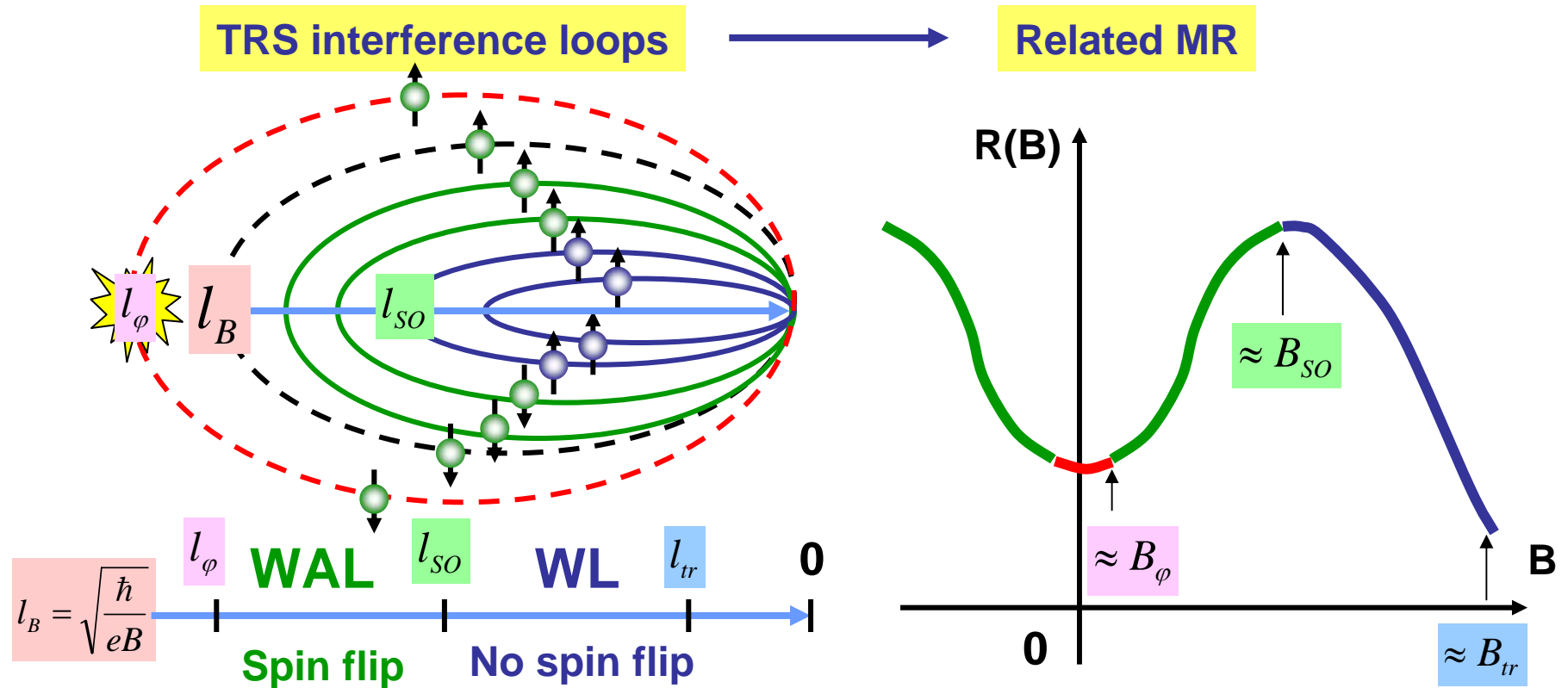
Many different paths with time reversal trajectories
(Final Spin states are statistical)

$$\langle \vec{S}_a | \vec{S}_c \rangle = \langle \vec{S}_i | R^2 | \vec{S}_i \rangle = \cos^2 \frac{\alpha}{2} \cdot e^{i(\beta+\gamma)} - \sin^2 \frac{\alpha}{2} \approx 0 + \frac{\cos \alpha - 1}{2} \approx -\frac{1}{2}$$

Initial spin state $|\vec{S}_i\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ for simplicity

WAL effect

Length scales in quantum interference



l_ϕ ; Maximum interference length $\rightarrow B_\phi = \hbar / 4e \cdot l_\phi^2$

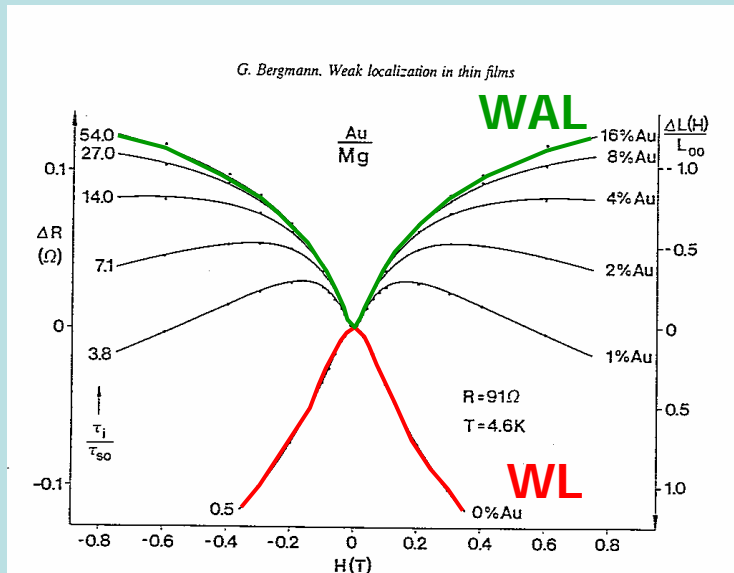
l_{SO} ; Spin flip length $\rightarrow B_{SO} = \hbar / 4e \cdot l_{SO}^2$

l_{tr} ; Minimum interference length $\rightarrow B_{tr} = \hbar / 4e \cdot l_{tr}^2$

l_B replaces l_ϕ as max. interference length when B is applied.

Gate Controlled Spin-orbit Interaction(SOI)

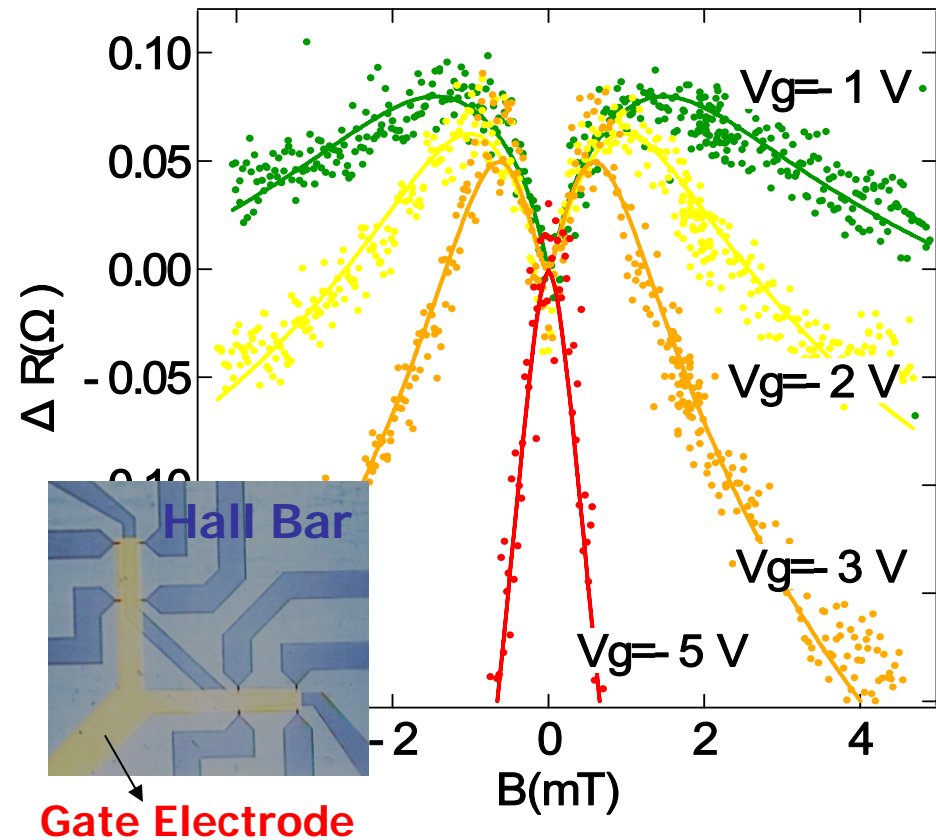
Mg thin film with Au



G. Bergmann,
Phys. Rev. Lett. 48, 1046 (1982)

*Spin-orbit scattering is introduced
by Au (Elliott spin relaxation)*

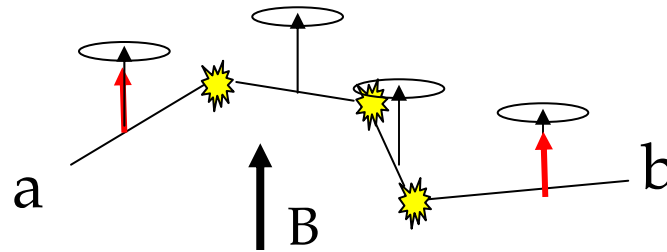
Gate fitted InGaAs 2DEG Hall Bar



Gate voltage can control the Rashba SOI.

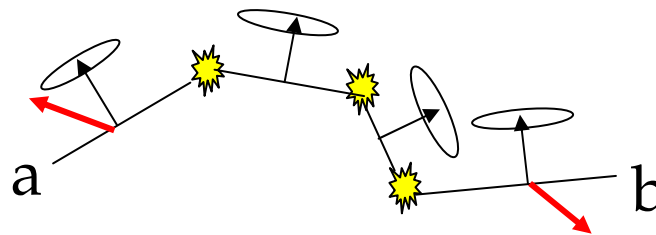
Spin dynamics in diffusive systems

Zeeman



Fixed
precession axis

Rashba SOI



Randomized
precession axis

$$H_R = \frac{\alpha}{\hbar} \cdot (p_x \sigma_y - p_y \sigma_x) \quad \rightarrow \quad \vec{B}_{SOI} \propto \vec{p} \times \hat{z}$$

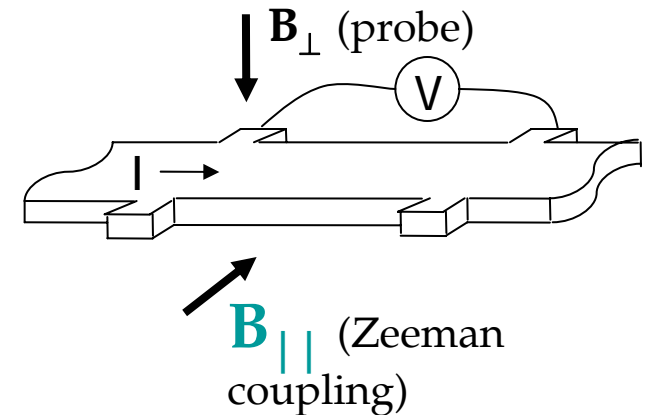
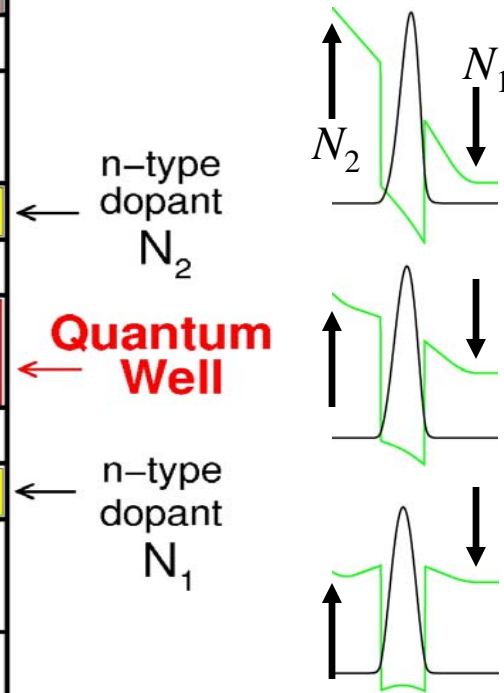
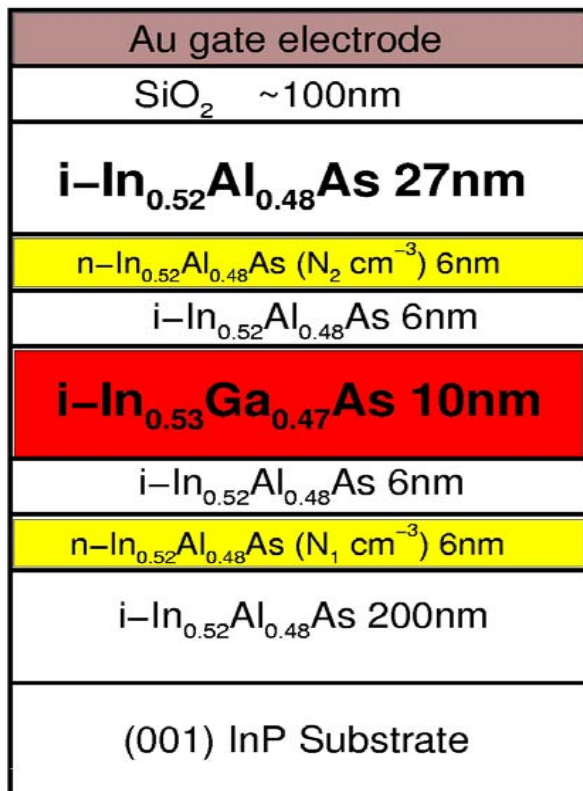
Rashba + scattering:
Spin relaxation (τ_s)

Competition Zeeman and Rashba:

alignment \leftrightarrow randomization

$$E_Z \leftrightarrow E_{SOI} \equiv \hbar/\tau_s$$

Competition between Zeeman and Rashba



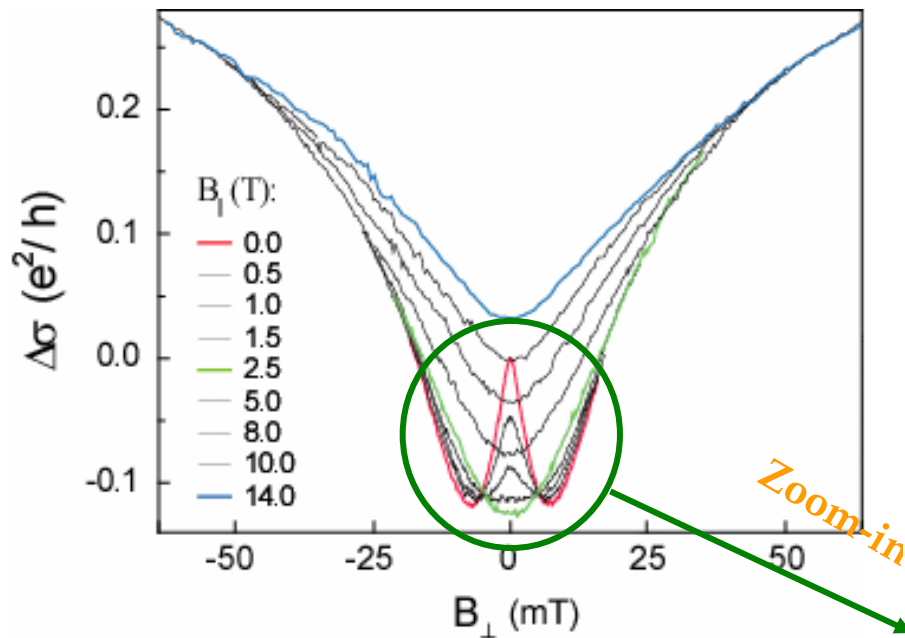
$$N_1 + N_2 = 4 \times 10^{18} \text{ cm}^{-3}$$

N_2/N_1 :	0	1/3	1
Δ (meV):	≈ 2	≈ 1.5	≈ 0.5

Competition Zeeman and Rashba:
alignment \leftrightarrow randomization

$$E_Z \leftrightarrow E_{\text{SOI}} \equiv \hbar/\tau_s$$

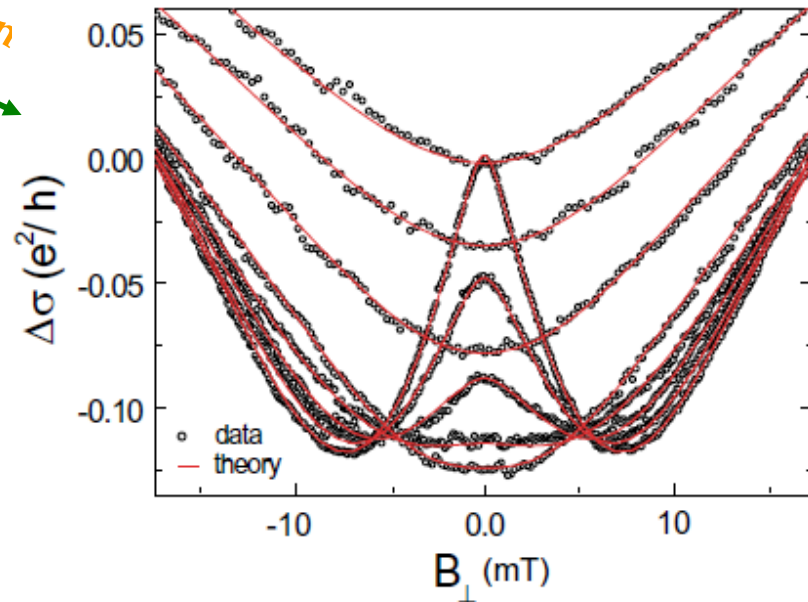
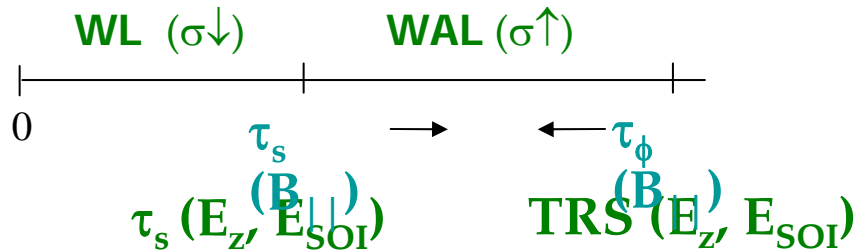
Weak anti-localization and data analysis



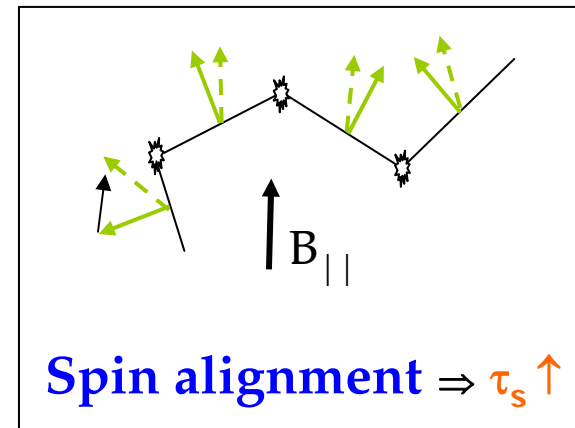
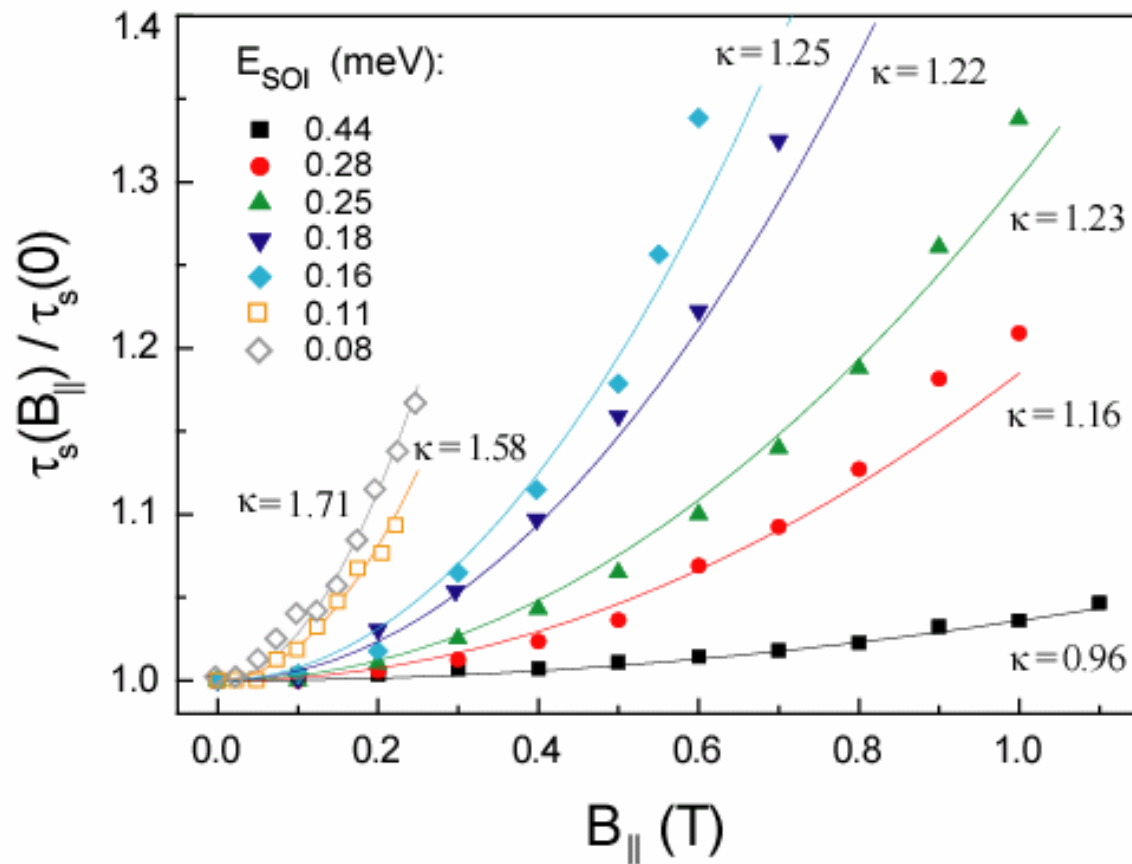
Good agreement with ILP theory

- $B_{||}$ (i.e. E_z/E_{SOI})
- ne (i.e. τ)
- all samples

WAL data Analysis $\rightarrow \tau_s(B_{||}), \tau_\phi(B_{||})$



Increase in spin relaxation time: $\tau_s(B_{||})$



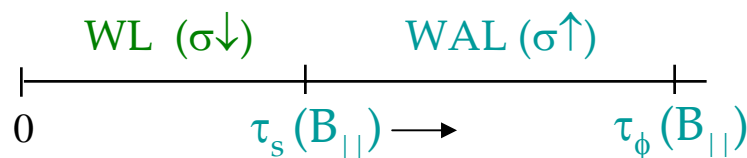
Energy Scale:

- **Zeeman** E_Z
- **Rashba** $E_{SOI} \equiv \hbar/\tau_s(0)$

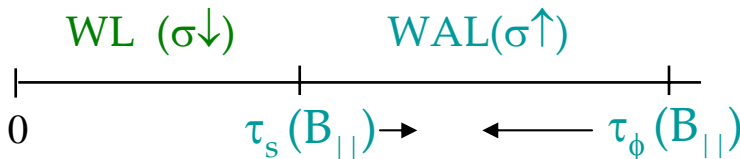
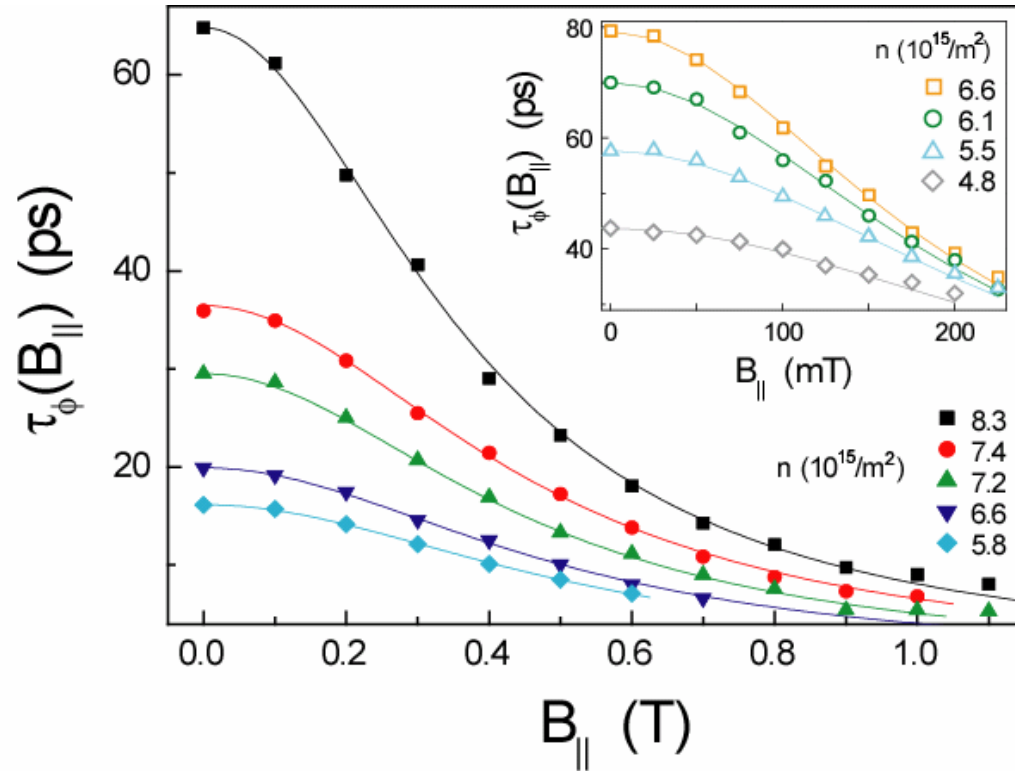
Theory:

$$\frac{\tau_s(B_{||})}{\tau_s(0)} \approx 1 + \frac{1}{2} (E_Z / E_{SOI})^2$$

(V.A. Frolov, PRB '01)



Decrease in dephasing time: $\tau_\phi(B_{||})$



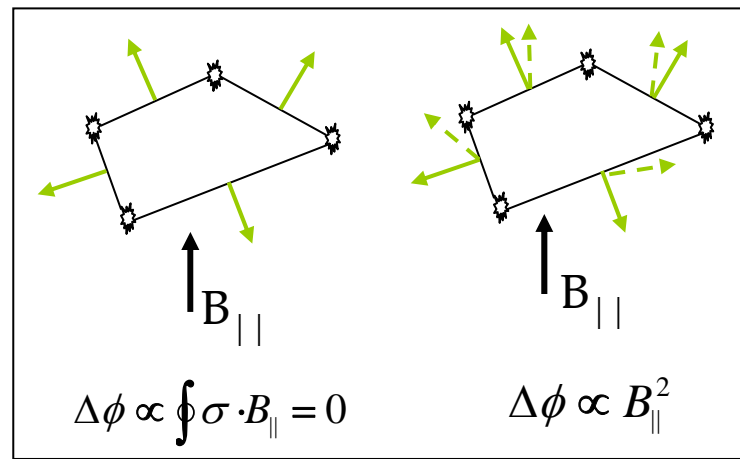
Theory ($E_z/E_{SOI} \ll 1$) :

$$\frac{1}{\tau_\phi(B_{||})} = \frac{1}{\tau_\phi(0)} + \frac{1}{T_\phi(B_{||})}$$

Spin induced dephasing

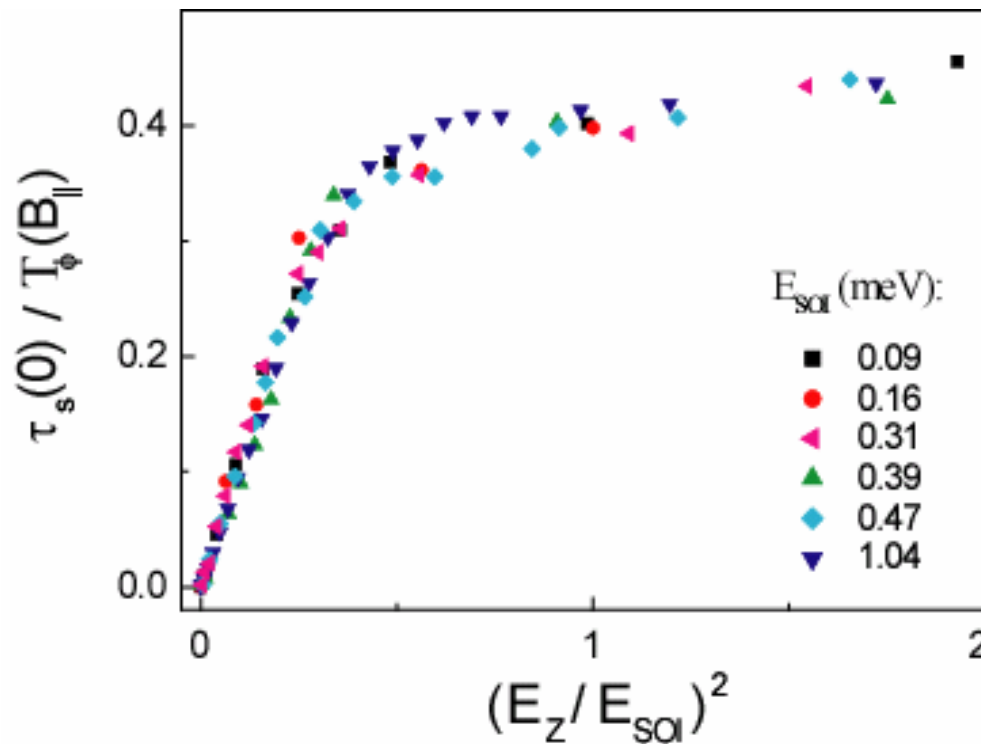
$$\frac{1}{T_\phi(B_{||})} \propto (E_z / E_{SOI})^2$$

(A.G. Malshukov *et al.*, PRB '97)



Universal Spin-Induced Time Reversal Symmetry Breaking

Spin induced dephasing rate: $1/T_\phi(B)$



$$\frac{1}{T_\phi(B_{||})} = \frac{1}{\tau_\phi(B_{||})} - \frac{1}{\tau_\phi(0)}$$

Total dephasing rate (points to $\tau_\phi(B_{||})$)
Inelastic scattering rate (points to $\tau_\phi(0)$)

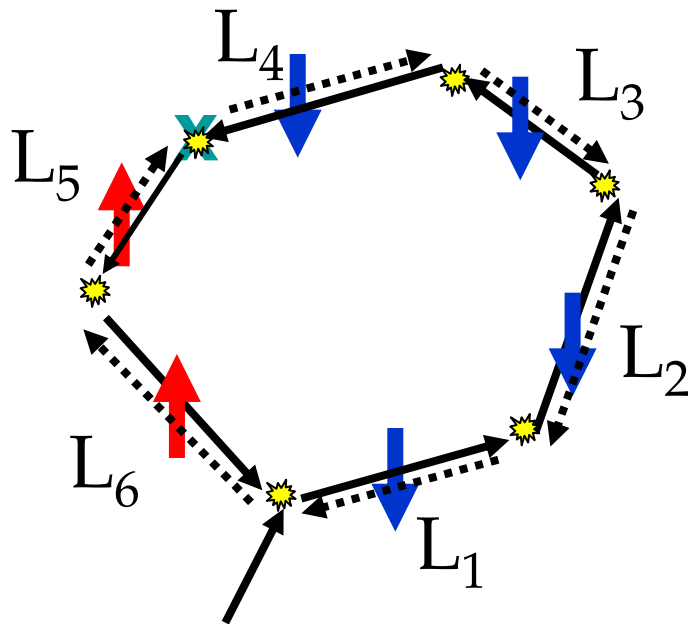
Upper limit $\sim \tau_s$

- Saturation ($E_Z/E_{SOI} \gg 1$)
- Universal behavior

\Rightarrow No available theory

Spin-induced Time Reversal Symmetry Breaking

Time-Reversal Symmetric Interference

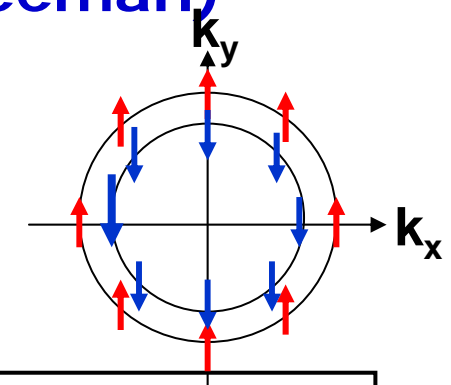
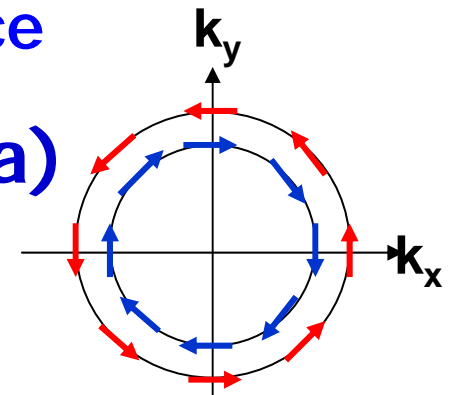


Spin flip (Rashba)

$$E_{SOI} \longrightarrow \tau_S$$

Difference in k_F (Zeeman)

$$Ez \longrightarrow |k_{F\uparrow}| \neq |k_{F\downarrow}|$$



Phase shift

CW $\phi_c = k_{F\uparrow}(L_6 + L_5) + k_{F\downarrow}(L_4 + L_3 + L_2 + L_1)$

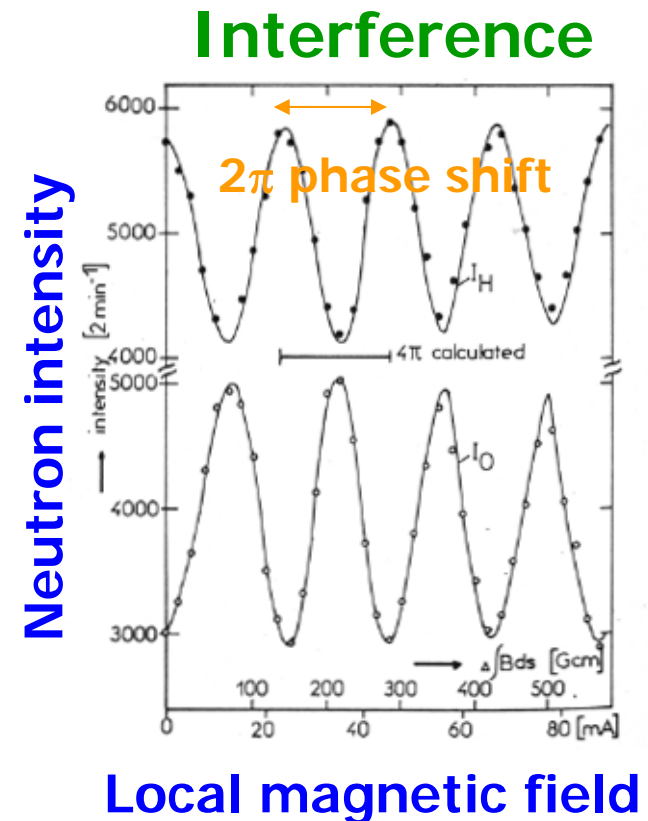
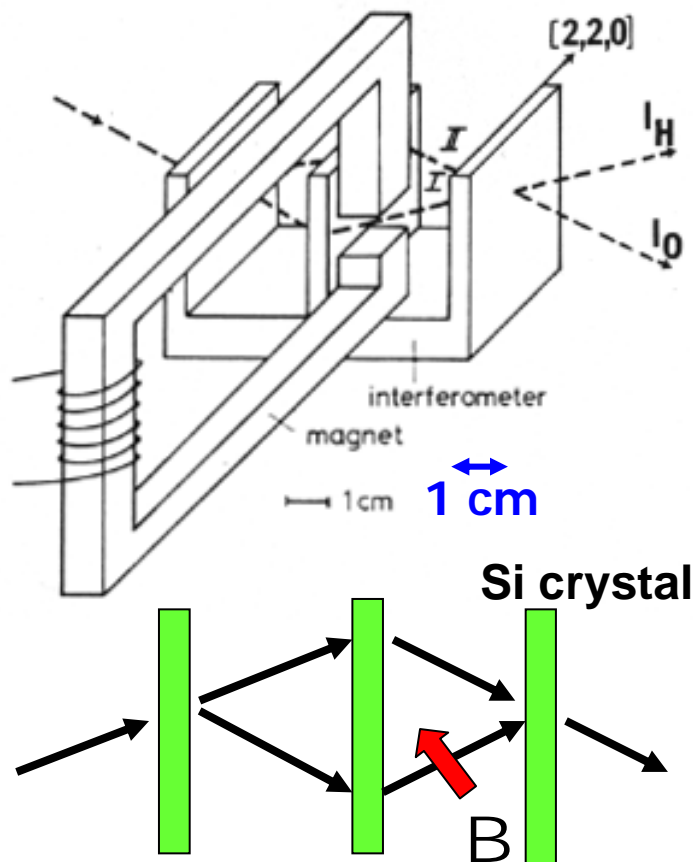
CCW $\phi_{ac} = k_{F\downarrow}(L_1 + L_2 + L_3 + L_4) + k_{F\uparrow}(L_5 + L_6)$

Breaking of interference

$$\Delta\phi = \phi_c - \phi_{ac} \approx 1 \quad \because |k_{F\uparrow}| \neq |k_{F\downarrow}|$$

Neutron Spin-interference Exp.

4π -spin precession = 2π -phase shift



Spin precession: Local magnetic field

H. Rauch *et al* Phys. Lett. 54A (1975)

Operational Principle of Spin Interferometer

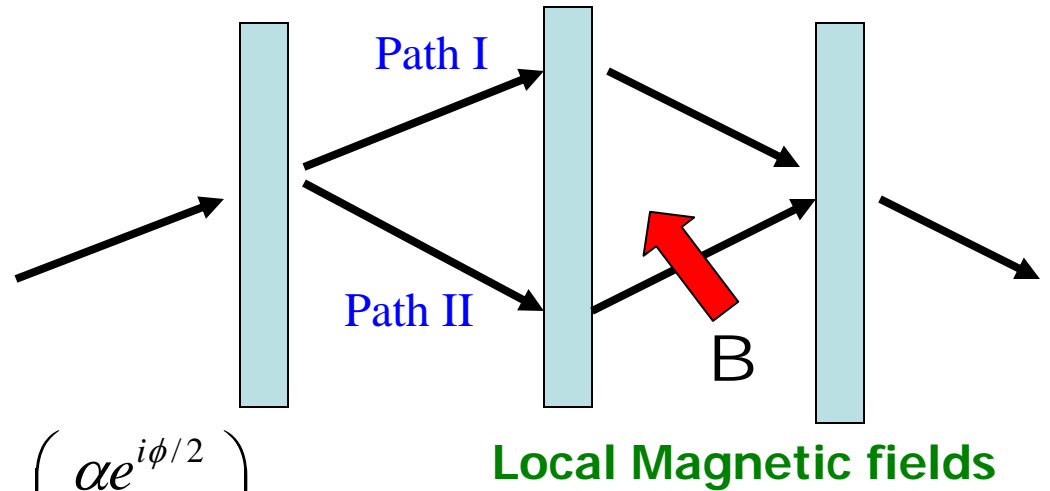
Non-spin polarized beam

Path I

$$\begin{pmatrix} \alpha \\ \beta \end{pmatrix} \rightarrow \begin{pmatrix} \alpha \\ \beta \end{pmatrix}$$

Path II

$$\begin{pmatrix} \alpha \\ \beta \end{pmatrix} \rightarrow \begin{pmatrix} e^{i\phi/2} & 0 \\ 0 & e^{-i\phi/2} \end{pmatrix} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} = \begin{pmatrix} \alpha e^{i\phi/2} \\ \beta e^{-i\phi/2} \end{pmatrix}$$



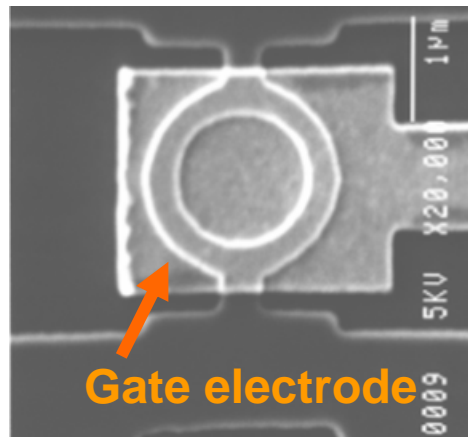
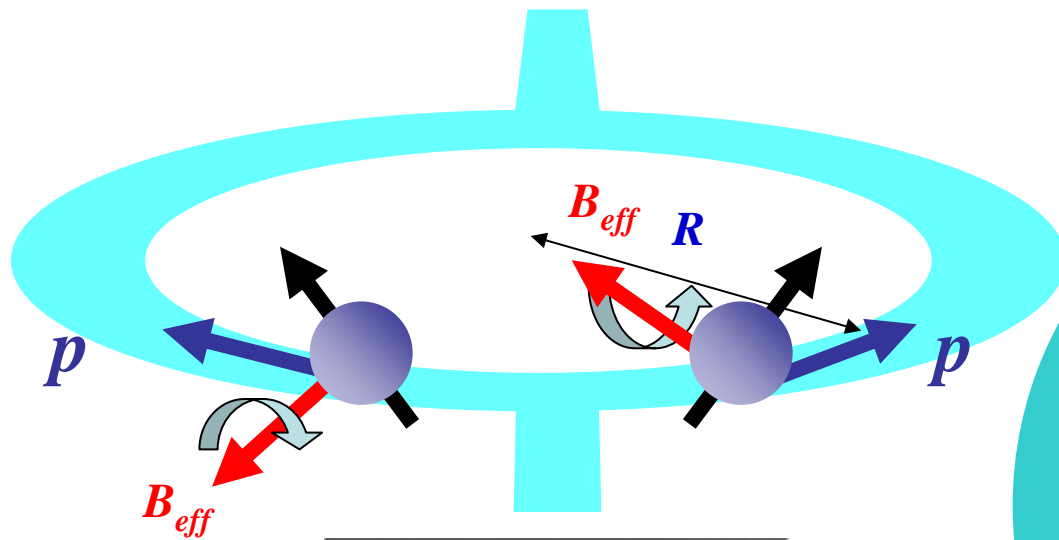
Local Magnetic fields

Interference

$$\frac{\left| \begin{pmatrix} \alpha e^{i\phi/2} \\ \beta e^{-i\phi/2} \end{pmatrix} + \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \right|^2}{\left| \begin{pmatrix} \alpha \\ \beta \end{pmatrix} + \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \right|^2} = \frac{(1 + \cos \phi / 2)}{2}$$

$$\phi = \gamma B \tau$$

Spin interferometer by the Rashba SOI

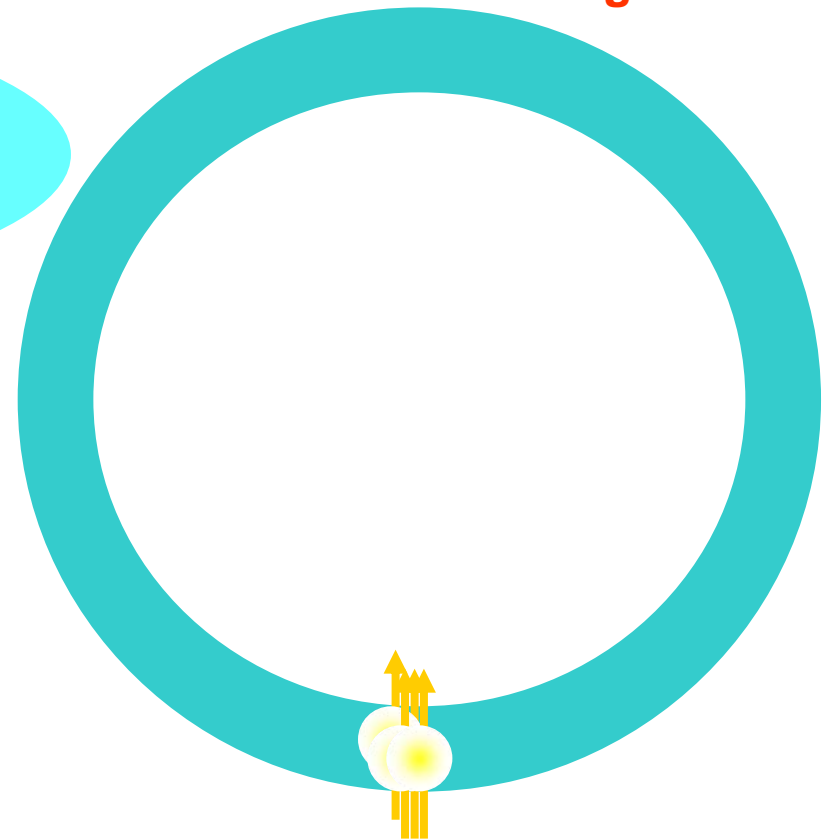


Spin interference device

Appl. Phys. Lett. 75, 695 (1999)

$V_g = 0$ V

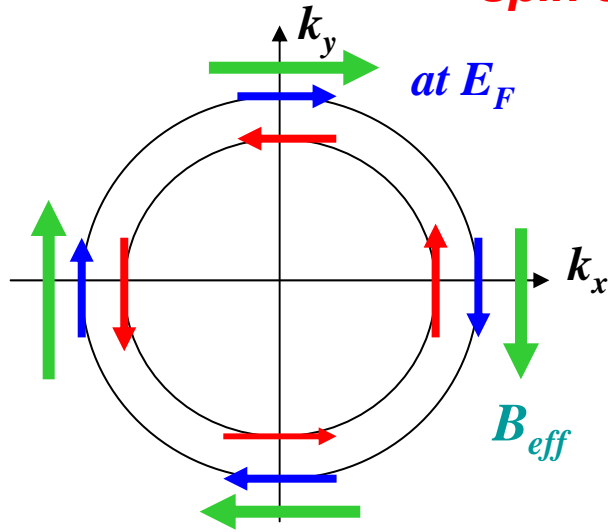
$V_g = 1$ V



Ring conductance depends on the precession angle

Spin precession by the Rashba SOI

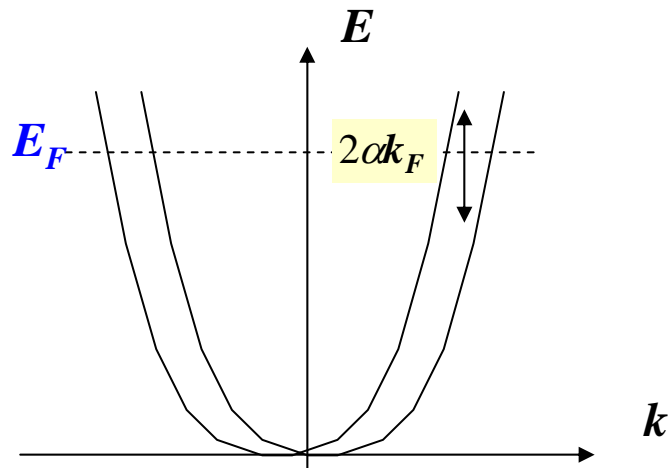
Spin splitting as if a spin feels an effective magnetic field perpendicular to its momentum direction



$$E(k) = \frac{\hbar^2 k^2}{2m^*} \pm \alpha k$$

$$\frac{\hbar^2 k_{\uparrow}^2}{2m^*} + \alpha k_{\uparrow} = \frac{\hbar^2 k_{\downarrow}^2}{2m^*} - \alpha k_{\downarrow} \quad \therefore k_{\downarrow} - k_{\uparrow} = \frac{2\alpha m^*}{\hbar^2}$$

$$k_{\downarrow} \approx k_F + \frac{\alpha m^*}{\hbar^2}, \quad k_{\uparrow} \approx k_F - \frac{\alpha m^*}{\hbar^2}$$



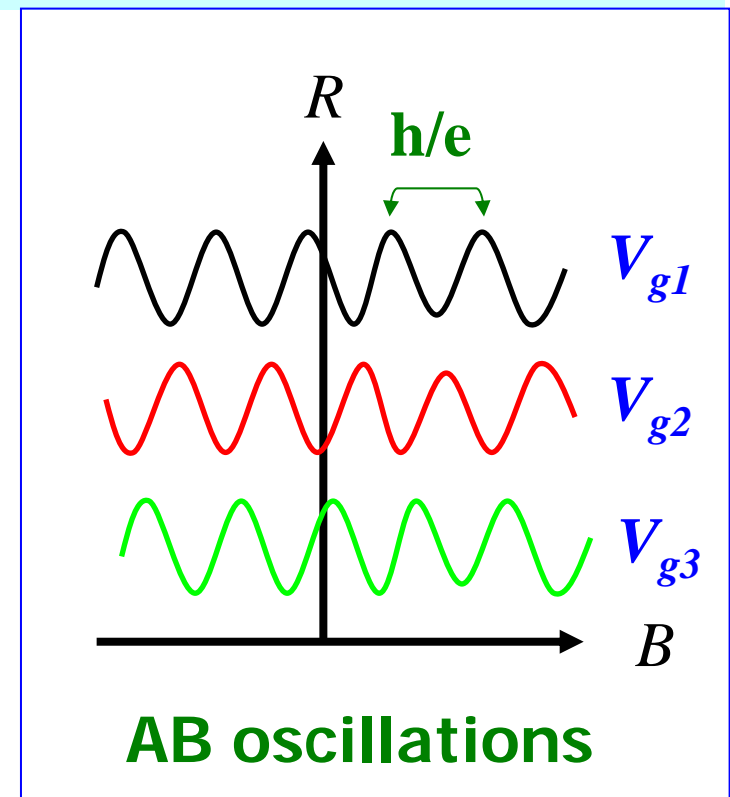
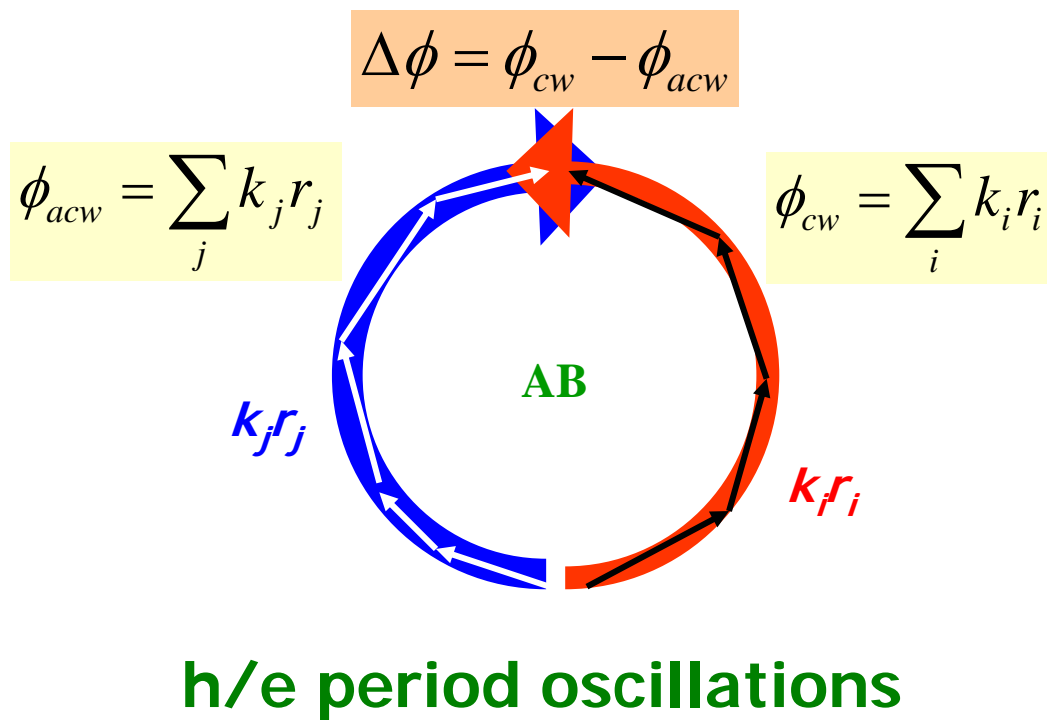
Spin precession and precession angle

$$\psi(r) = \frac{1}{\sqrt{2}} e^{ik_F r} \begin{pmatrix} e^{-\frac{\alpha m^*}{\hbar^2} r} \\ e^{\frac{\alpha m^*}{\hbar^2} r} \end{pmatrix} \quad \theta = \frac{2\alpha m^* L}{\hbar^2}$$

Aharonov-Bohm Oscillations

Sample specific feature

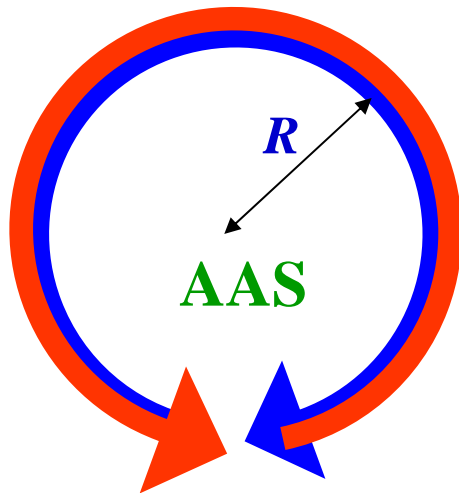
Detail of the trajectory and k_F affect the interference



Gate voltage changes k_F , and its interference pattern !

AAS effect; Time-reversal symmetry interference

AAS oscillation does not depend on **wave-vector k_F**
 But it depends on **spin precession angle**



$h/2e$ period oscillations

Trajectory: Same length
between cw and anti-cw

Wave-function

$$\Phi_{k\sigma} = \frac{1}{\sqrt{2}} e^{ikx} \begin{pmatrix} e^{\frac{i\alpha m^* x}{\hbar^2}} \\ e^{\frac{i\alpha m^* x}{\hbar^2}} \end{pmatrix}$$

Orbital phase	Spin phase
$\phi_{cw} = \sum_{i=1}^n k_i r_i = \phi_{acw}$	$\varphi_{cw} = \frac{2\pi\alpha m^* R}{\hbar^2} = -\varphi_{acw}$
<p>Phase difference= 0</p>	<p>Phase difference \neq 0</p>
$\Delta\phi = \phi_{cw} - \phi_{acw} = 0$	$\Delta\varphi = \varphi_{cw} - \varphi_{acw} = \frac{4\pi\alpha m^* R}{\hbar^2}$
<p>Always constructive</p>	<p>Constructive/Destructive</p>

Ensemble averaging of AB oscillations

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PHYSICAL REVIEW LETTERS

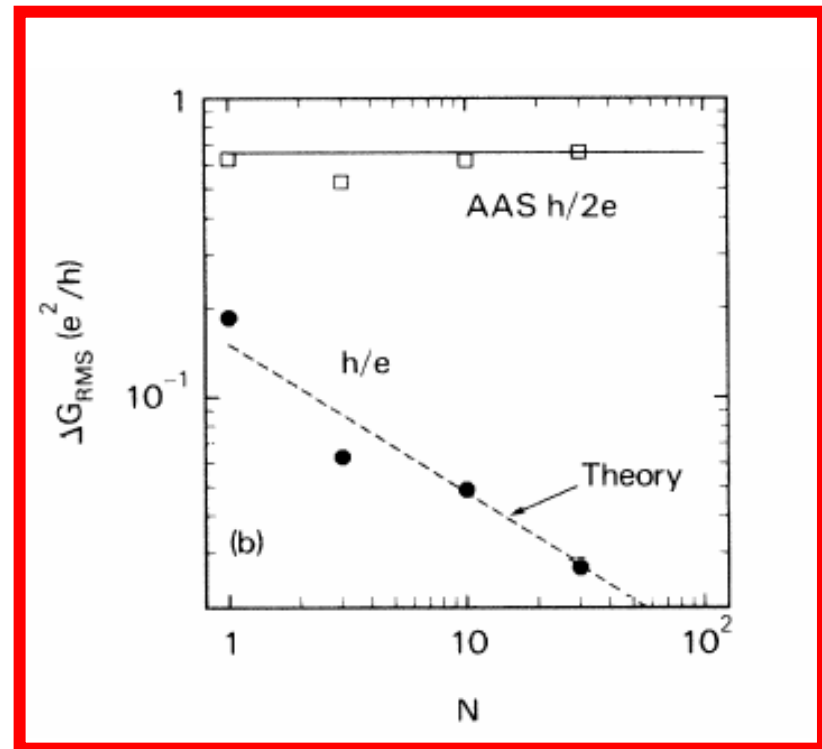
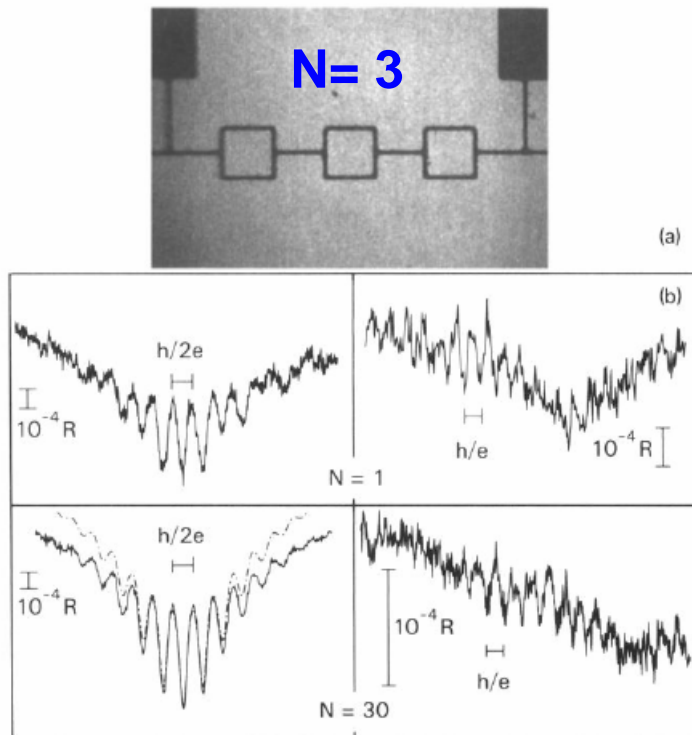
27 JANUARY 1986

Direct Observation of Ensemble Averaging of the Aharonov-Bohm Effect in Normal-Metal Loops

C. P. Umbach, C. Van Haesendonck,^(a) R. B. Laibowitz, S. Washburn, and R. A. Webb

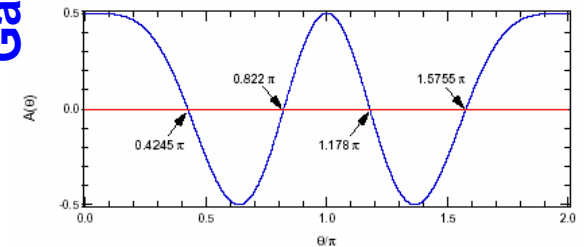
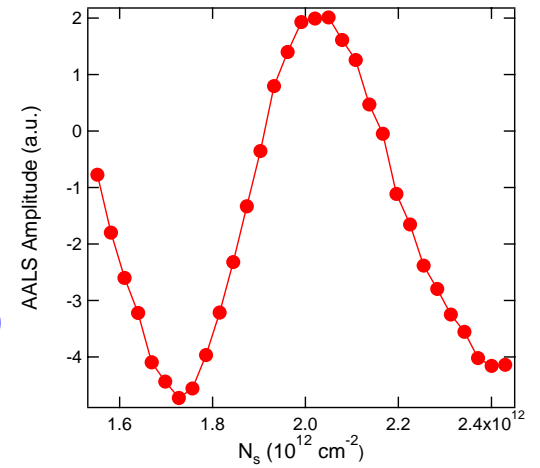
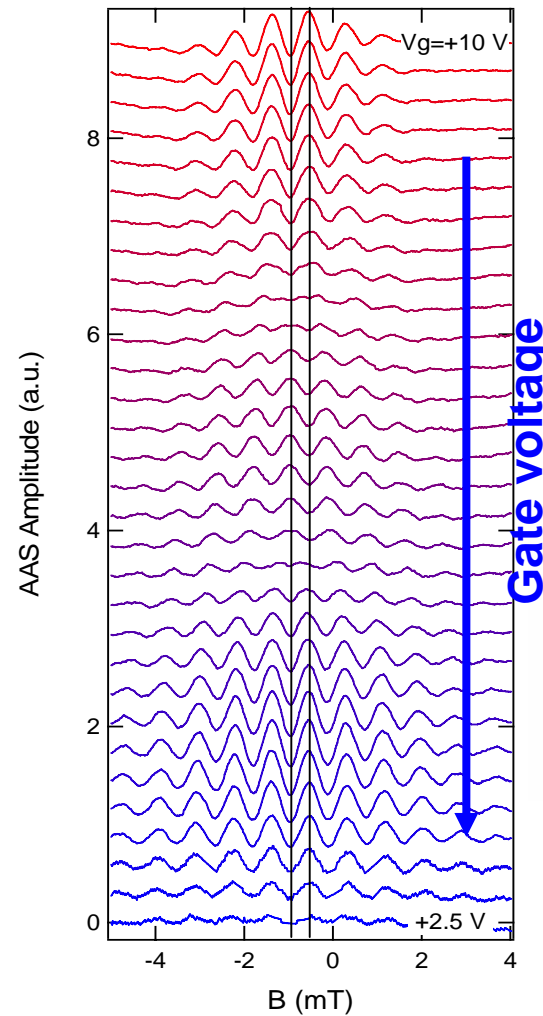
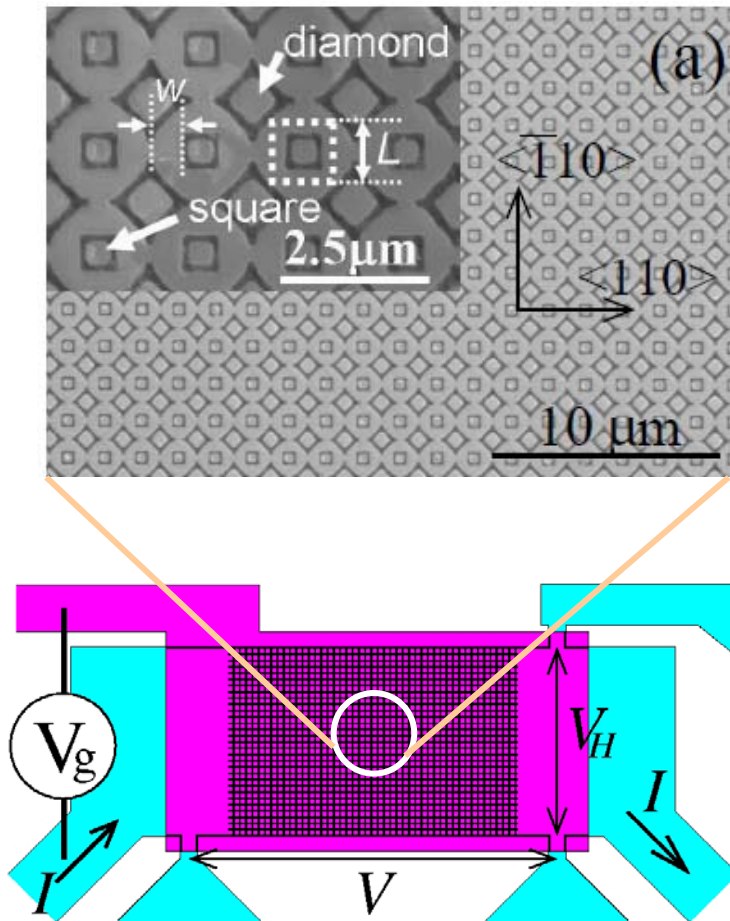
IBM T.J. Watson Research Center, Yorktown Heights, New York 10598

(Received 6 November 1985)



V_g dependence of AAS oscillations

Array of 7700 loops

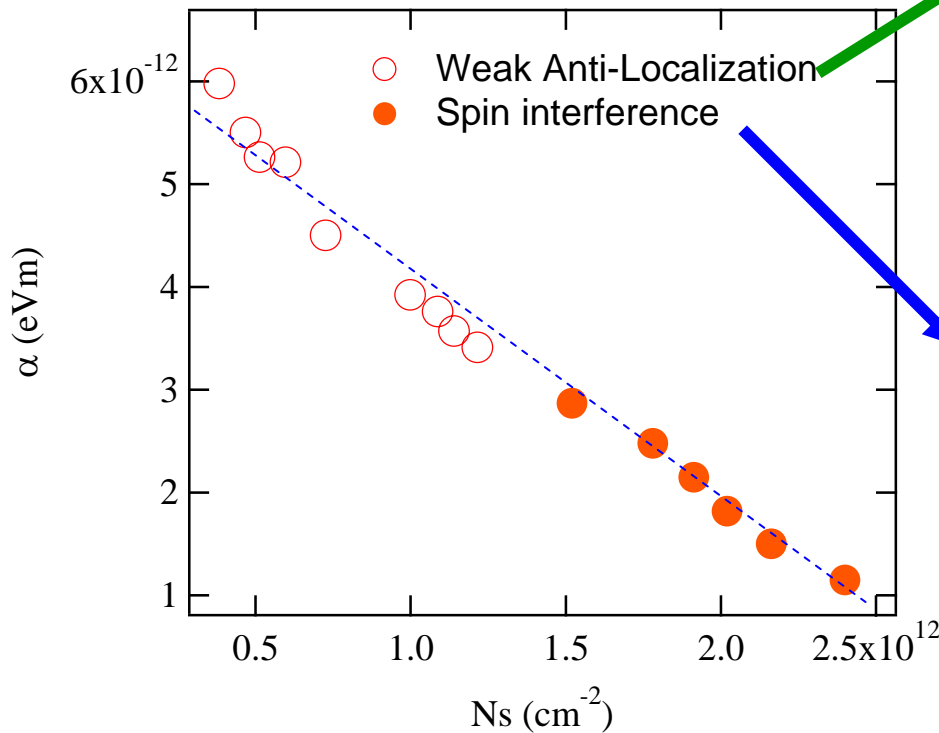


$$\alpha = \frac{\hbar^2 \theta}{2m^* L}$$

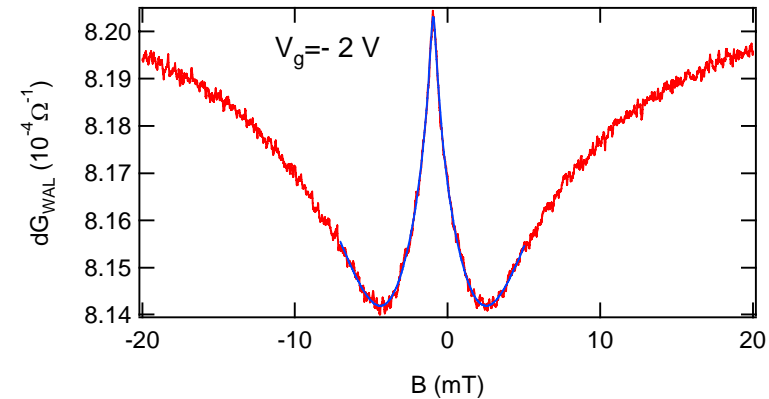
Array of loops is covered with gate to control SOI

Gate controlled SOI and spin precession

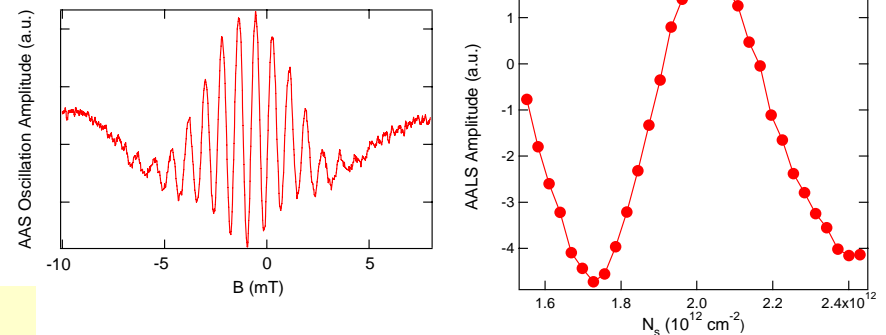
Gate voltage dependence of SOI



Weak anti-localization



Spin interference

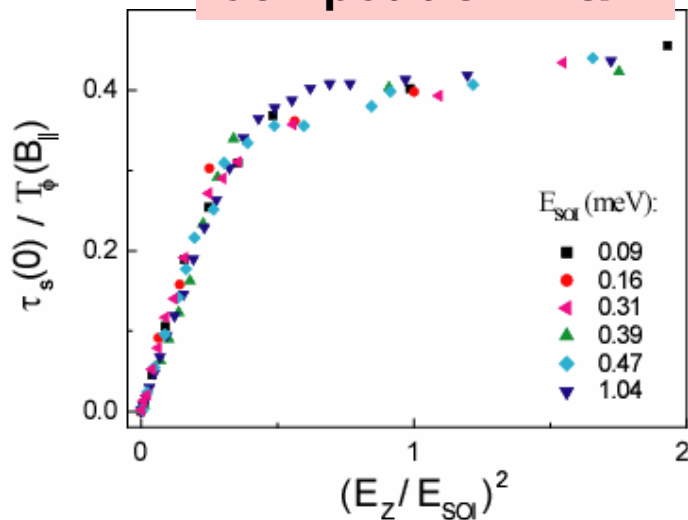


Experimental demonstration of Aharonov-Casher Effect

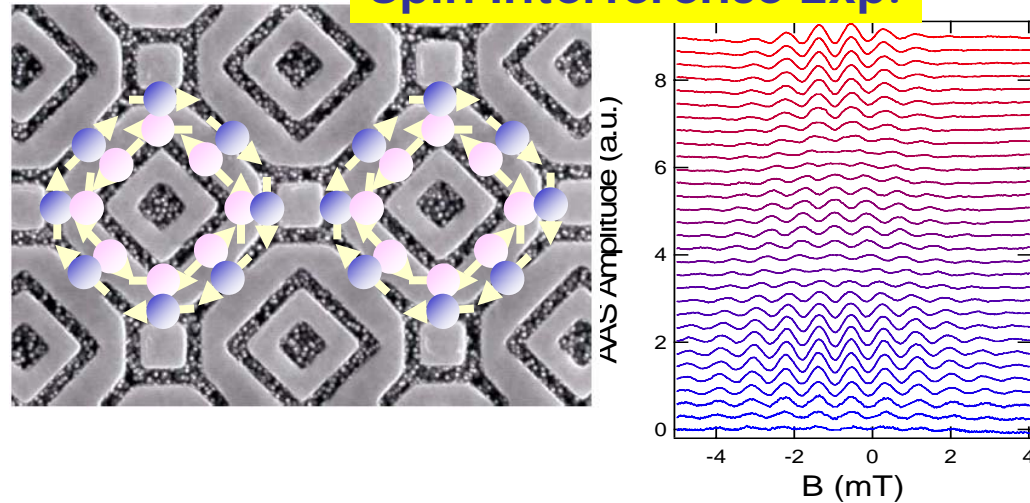
Summary

1. The origin of Rashba SOI and gate control
2. Competition between Zeeman and Rashba
3. Spin interference device

Competition R & Z



Spin interference Exp.



F. Meijer, A. Morpurgo, T. Klapwijk (TUD)

T. Koga, Y. Sekine, T. Bergsten (NTT)

J. Ohe, T. Ohtsuki (Sophia Univ.)