



## Collaborators:

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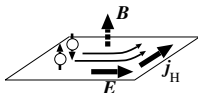
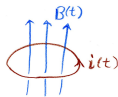
# Electromagnetism for charge and spin

- Electromagnetism

$$\nabla \times \mathbf{E} + \dot{\mathbf{B}} = 0$$

$$\nabla \times \mathbf{B} = \mu \mathbf{j} + \epsilon \mu \dot{\mathbf{E}}$$

$$\mathbf{j}_{\text{Hall}} = \sigma_{\text{H}}(\mathbf{E} \times \mathbf{B})$$



- Spin electromagnetism    Ferromagnetic metals

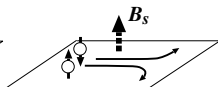
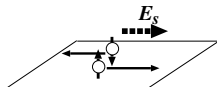
*Volovik'87*

- Drives conduction electron spin
- *sd* exchange interaction

$$\nabla \times \mathbf{E}_{\text{S}} + \dot{\mathbf{B}}_{\text{S}} = 0$$

$$\nabla \times \mathbf{B}_{\text{S}} = \mu_{\text{S}} \mathbf{j} + \epsilon_{\text{S}} \mu_{\text{S}} \dot{\mathbf{E}}_{\text{S}}$$

$$\mathbf{j}_{\text{Hall}} = \sigma_{\text{SH}}(\mathbf{E} \times \mathbf{B}_{\text{S}})$$



- $E_{\text{S}}$  : Spin motive force,  $B_{\text{S}}$  : Spin Berry's phase

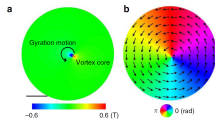
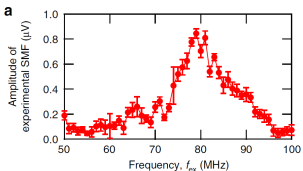
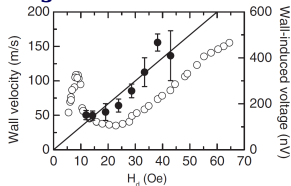
*Berger'86, Stern'92, Barnes&Maekawa'07*

- Real field    Detectable by electric measurements

$$\mathbf{j}_{\text{S}} = P \mathbf{j} \text{ in ferromagnetic metals} \quad P \sim O(1)$$

# Spin electromagnetic field

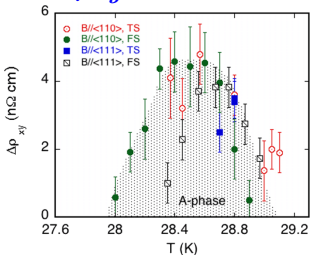
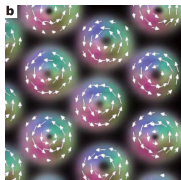
- $E_s$  from motion of domain wall, vortex  $V \sim \mu V \propto E_s \propto v_{dw}$



Domain wall *Yang'09*

Vortex *Tanabe'12*

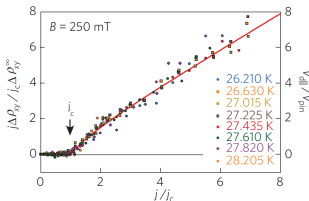
- Skyrmion lattice  $\rho_{xy} \sim 4n\Omega cm \propto B_s$



Topological Hall effect  $B_s$

*Yu'10*

*Neubauer'09*



Voltage,  $E_s \propto v_{skym}$

*Schulz'12*

# Spin electromagnetic field

*Volovik'87, Stern'92, Barnes&Maekawa'07*

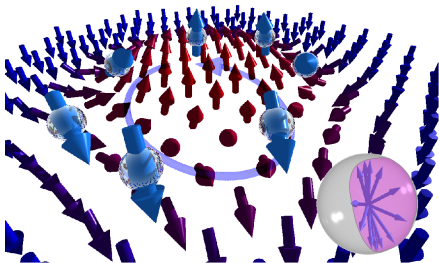
Adiabatic limit

- Electron spin rotation

⇒ Phase  $e^{i\varphi}$

$$\varphi = \int_C dr \cdot A_s$$

$$A_s = \frac{1}{2}(1 - \cos \theta) \partial \phi$$



- Spin magnetic field

$$\varphi = \int_S dS \cdot B_s$$

- Faraday's law is satisfied

- Spin electric field (dynamics)

$$\dot{\varphi} = - \int_C dr \cdot E_s$$

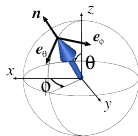
$$\nabla \times E_s = - \frac{\partial B_s}{\partial t}$$

Electromagnetic field coupled to spin

# Topological spin electromagnetic field (Adiabatic limit)

- Effective U(1) gauge field Adiabatic limit

$$A_{\mathbf{s},\mu}^z = \frac{1}{2}(1 - \cos \theta) \partial_\mu \phi$$



- Spin electromagnetic fields

$$E_{\mathbf{s},i} = -\nabla_i A_{\mathbf{s},0}^z + \partial_t A_{\mathbf{s},i}^z = -\frac{1}{2} \mathbf{n} \cdot (\dot{\mathbf{n}} \times \nabla_i \mathbf{n})$$

$$B_{\mathbf{s},i} = (\nabla \times A_{\mathbf{s}}^z)_i = \frac{1}{4} \sum_{j,k} \epsilon_{ijk} \mathbf{n} \cdot (\nabla_j \mathbf{n} \times \nabla_k \mathbf{n})$$



$B_{\mathbf{s}}$ : Spin Berry's phase

$E_{\mathbf{s}}$ : Spin motive force

$E_{\mathbf{s}}$  and  $B_{\mathbf{s}}$  couple spin structure and electron transport

# Effects of spin-orbit interaction

- Inverse (and direct) spin Hall effects
  - Spin-orbit  $\Rightarrow$  No longer in the adiabatic limit
  - Spin relaxation effects
- Spin relaxation effects on  $E_s$

- Spin relaxation  $\beta$

*Duine PRB'08*

$$\mathbf{n} \cdot (\dot{\mathbf{n}} \times \nabla_i \mathbf{n}) \Rightarrow \beta \mathbf{n} \cdot (\dot{\mathbf{n}} \times (\mathbf{n} \times \nabla_i \mathbf{n})) = \beta (\dot{\mathbf{n}} \nabla_i \mathbf{n})$$

- Onsager relation

*Saslow PRB'07, Tserkovnyak PRB'08*

- Rashba effects

- Weak  $sd$  Transport (Diagrams) *Takeuchi&GT J.Phys.Soc.Jpn'12*

- Strong  $sd$  "Chiral derivative" *Kim PRL12*

- Strong  $sd$  Transport (Diagrams) *GT PRB'13, Nakabayashi&GT'13*

# Rashba-induced spin electromagnetic field

- Strong sd coupling  $\Delta_{\text{sd}}$  + Rashba interaction  $\alpha_{\mathbf{R}}$

$$H = \int d^3r c^\dagger \left[ \left( \frac{\hbar^2}{2m} \nabla^2 + \epsilon_F \right) + \Delta_{\text{sd}} (\mathbf{n} \cdot \boldsymbol{\sigma}) - \frac{i}{2} \alpha_{\mathbf{R}} \cdot \left( \overleftrightarrow{\nabla} \times \boldsymbol{\sigma} \right) \right] c$$

- Diagram calculation

*Nakabayashi&GT,arXiv:1308.0152*

- Result

- Electromagnetic field description works

*Linear order of  $\alpha_{\mathbf{R}}$*

$$\begin{aligned} \nabla \times \mathbf{E}_{\mathbf{R}} + \dot{\mathbf{B}}_{\mathbf{R}} &= 0 \\ \nabla \times \mathbf{B}_{\mathbf{R}} - \epsilon_{\mathbf{R}} \mu_{\mathbf{R}} \dot{\mathbf{E}}_{\mathbf{R}} &= \mu_{\mathbf{R}} \mathbf{j} \\ \mathbf{j}_{\text{Hall}} &= \sigma_{\mathbf{R}} (\mathbf{E} \times \mathbf{B}_{\mathbf{R}}) \end{aligned}$$

$$\mathbf{E}_{\mathbf{R}} = -\frac{m}{e\hbar} (\alpha_{\mathbf{R}} \times \dot{\mathbf{n}})$$

$$\mathbf{B}_{\mathbf{R}} = \frac{m}{e\hbar} [\nabla \times (\alpha_{\mathbf{R}} \times \mathbf{n})]$$

- Spin vector potential *Linear order*

$$\mathbf{A}_{\mathbf{R}} = -\frac{m}{e\hbar} (\alpha_{\mathbf{R}} \times \mathbf{n})$$

*Kim'12, Nakabayashi&GT,arXiv*



## Total effective electromagnetic field

- Ferromagnetic metals
- Charge + Spin electromagnetic fields

$$\mathbf{E}_{\text{eff}} = \mathbf{E} + \mathbf{E}_{\text{S}},$$

$$\mathbf{E}_{\text{S}} = \mathbf{E}_{\text{S,top}} + \mathbf{E}_{\text{R}}$$

$$\mathbf{B}_{\text{eff}} = \mathbf{B} + \mathbf{B}_{\text{S}},$$

$$\mathbf{B}_{\text{S}} = \mathbf{B}_{\text{S,top}} + \mathbf{B}_{\text{R}}$$

- Ampère's law

- $\nabla \times (\mathbf{B} + \mathbf{B}_{\text{S}}) - \frac{\partial}{\partial t} (\epsilon\mu\mathbf{E} + \epsilon_{\text{S}}\mu_{\text{S}}\mathbf{E}_{\text{S}}) = \mu\mathbf{j}$

## Total effective electromagnetic field

- Ferromagnetic metals
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$$\begin{aligned} E_{\text{eff}} &= E + E_{\text{S}}, & E_{\text{S}} &= E_{\text{S,top}} + E_{\text{R}} \\ B_{\text{eff}} &= B + B_{\text{S}}, & B_{\text{S}} &= B_{\text{S,top}} + B_{\text{R}} \end{aligned}$$

- Ampère's law

- $\nabla \times (B + B_{\text{S}}) - \frac{\partial}{\partial t} (\epsilon\mu E + \epsilon_{\text{S}}\mu_{\text{S}} E_{\text{S}}) = \mu j$

- $\nabla \times B - \epsilon\mu \frac{\partial E}{\partial t} = \mu(j + j^{(\text{SEMF})})$

$$\begin{aligned} j^{(\text{SEMF})} &\equiv -\frac{1}{\mu} \left( \nabla \times B_{\text{S}} - \epsilon_{\text{S}}\mu_{\text{S}} \frac{\partial E_{\text{S}}}{\partial t} \right) \\ &= -\frac{1}{\mu} (\nabla \times B_{\text{S}}) + \sigma_{\text{S}} E_{\text{S}} \quad (\omega\tau \ll 1) \end{aligned}$$

Electric current induced by spin electromagnetic field

# Rashba-induced spin electromagnetic field

- Rashba-induced current

$$\begin{aligned}j^{(\text{SEMF})} &= -\frac{1}{\mu}(\nabla \times B_{\mathbf{R}}) + \sigma_{\mathbf{s}} E_{\mathbf{R}} \\ &= -\frac{1}{\mu}(\nabla \times (\nabla \times (\alpha_{\mathbf{R}} \times \mathbf{n}))) + \frac{m}{e\hbar} \sigma_{\mathbf{s}} (\alpha_{\mathbf{R}} \times \dot{\mathbf{n}})\end{aligned}$$

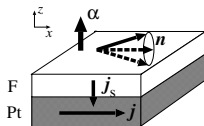
- Dynamic part

- Spin pumping + Inverse spin Hall by Rashba (?)

$$j_x = \alpha_{\mathbf{R}} j_{\mathbf{s},z}^y, j_{\mathbf{s},z} \propto \dot{\mathbf{n}}$$

- Present model is not directly applicable

- Rashba interaction
- Uniform system (Rashba and  $sd$  coexist)



Junction is argued by including electron diffusion

## Rashba-induced spin electromagnetic field

- Where is essential spin pumping contribution ?
  - Include spin relaxation

$$\mathbf{E}_{\mathbf{R}} = \frac{m}{e\hbar} [(\boldsymbol{\alpha}_{\mathbf{R}} \times \dot{\mathbf{n}}) + \beta_{\mathbf{R}}[\boldsymbol{\alpha}_{\mathbf{R}} \times (\mathbf{n} \times \dot{\mathbf{n}})]]$$

$\beta_{\mathbf{R}}$  Spin relaxation rate



$$j^{(\text{SEMF})} = -\frac{m}{\mu e \hbar} \sigma_{\mathbf{s}} \boldsymbol{\alpha}_{\mathbf{R}} \times [\dot{\mathbf{n}} + \beta_{\mathbf{R}}(\mathbf{n} \times \dot{\mathbf{n}})]$$

- Looks more like Spin pumping + Inverse spin Hall

$$j_{\mathbf{s},z} = g_{\uparrow\downarrow}(\mathbf{n} \times \dot{\mathbf{n}}) + g'_{\uparrow\downarrow} \dot{\mathbf{n}} \quad g_{\uparrow\downarrow} \simeq \frac{m}{\mu e \hbar} \sigma_{\mathbf{s}} \beta_{\mathbf{R}}$$

- No undefined quantity

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↓

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## Rashba-induced spin electromagnetic field

- Rashba-induced spin electric field

$$E_{\mathbf{R}} = \frac{m}{e\hbar} [(\alpha_{\mathbf{R}} \times \dot{\mathbf{n}}) + \beta_{\mathbf{R}}[\alpha_{\mathbf{R}} \times (\mathbf{n} \times \dot{\mathbf{n}})]]$$

↓

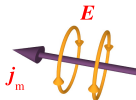
$$\nabla \times E_{\mathbf{S}} + \dot{B}_{\mathbf{S}} = j_{\mathbf{m}} \neq 0$$

- Emergence of spin damping monopole from spin relaxation

*Takeuchi, J. Phys.Soc.Japan '12*

$$j_{\mathbf{m}} = \beta_{\mathbf{R}} \nabla \times (\alpha_{\mathbf{R}} \times N)$$

$N \equiv \mathbf{n} \times \dot{\mathbf{n}}$ : Spin damping vector



- Non-conservation of spin (dissipation)  $\Rightarrow$  Monopole (?)
- No topological meaning (?)
- But is physical

$$\mathbf{n} \times \dot{\mathbf{n}} \Rightarrow j_{\mathbf{m}} \Rightarrow \nabla \times E_{\mathbf{R}} = j_{\mathbf{m}} \Rightarrow \boxed{E}$$

Monopole current converts spin damping to electric voltage

# Summary

- Strong  $sd$  + Rashba spin-orbit system
- Effective electromagnetic field description
  - Beautiful    Physicists' job is to seek beauty ..?
  - No undefined parameters
  - More works necessary (junction, diffusion)  
    Post docs, good students

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# RIKEN positions

- Postdoc/researcher positions available in Spin Physics Theory Group  
[gen.tatara@riken.jp](mailto:gen.tatara@riken.jp)
- Positions offered by RIKEN <http://www.riken.jp/en/careers/programs/>
- Postdocs, Graduate school students (PhD candidates)

## For Doctoral Candidates International Program Associate



### Qualifications

IPA applicants must be (or soon to be) holders of master's degrees and be enrolled as doctoral candidates at universities that have signed (or are willing to sign) an agreement with RIKEN to participate in the Joint Graduate School Program. They should not have more than 6 years of research experience after earning their master's degree. The RIKEN researchers in charge of supervising successful applicants must also hold concurrent positions as visiting faculty members at the collaborating Japanese or overseas universities where the applicants are enrolled.



### Duration

In principle, IPAs can participate in the program for a maximum of 3 years.



### Support from RIKEN

A daily living allowance of 5,200 yen will be paid. Residence in RIKEN campus housing for the duration of the IPA's stay at RIKEN is free of charge. When campus housing is unavailable, RIKEN will pay a housing allowance of up to 70,000 yen per month for off-campus housing. (The IPA must pay his or her own utility costs.) The roundtrip transportation between RIKEN and the IPA's home country will be paid by RIKEN.

## For Postdoctoral Researchers Foreign Postdoctoral Researcher Program



### Qualifications

In principle, applicants should have no more than 5 years of post doctoral research experience.



### Duration

In principle, the contract period is for a maximum of 3 years.



### Remuneration

The base salary is 487,000 yen per month before taxes and the social insurance premium deduction. Commuting, housing, and relocation allowances will also be provided based on individual circumstances. An annual research budget of approximately 1 million yen is allocated to the laboratories hosting the FPRs.



### Application and selection process

Applications are publicly solicited each year and reviewed by a screening committee of scientists in and outside RIKEN working in fields relevant to applications.



For more details please refer to the URL below or send an Email message to: [fpr@riken.jp](mailto:fpr@riken.jp)