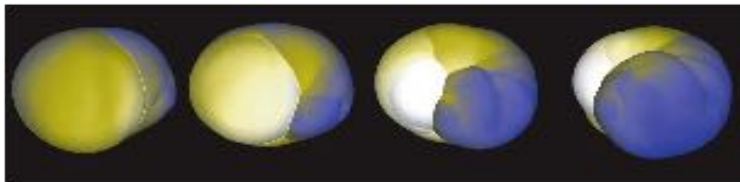
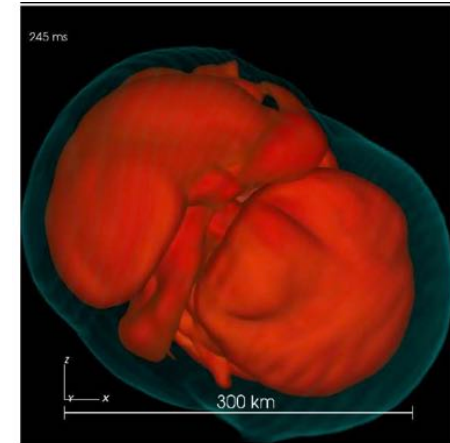
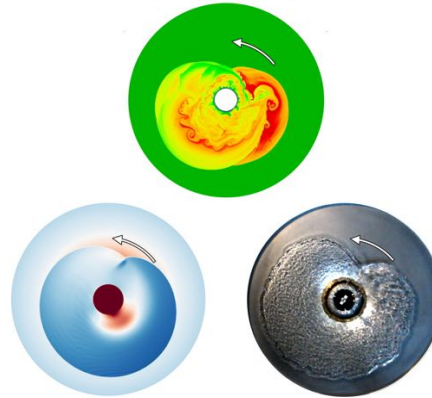


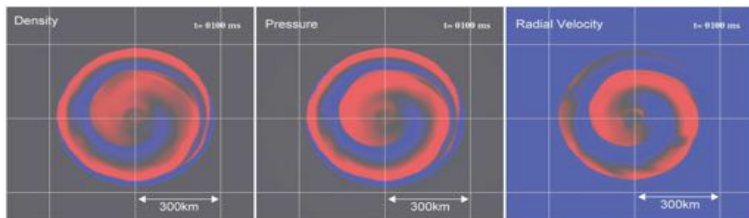
The core-collapse angular momentum budget



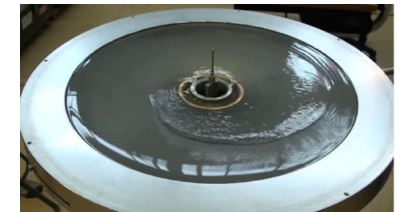
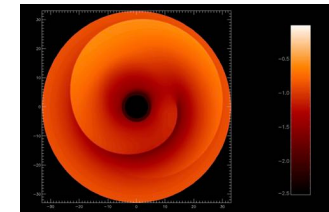
Blondin & Mezzacappa 07



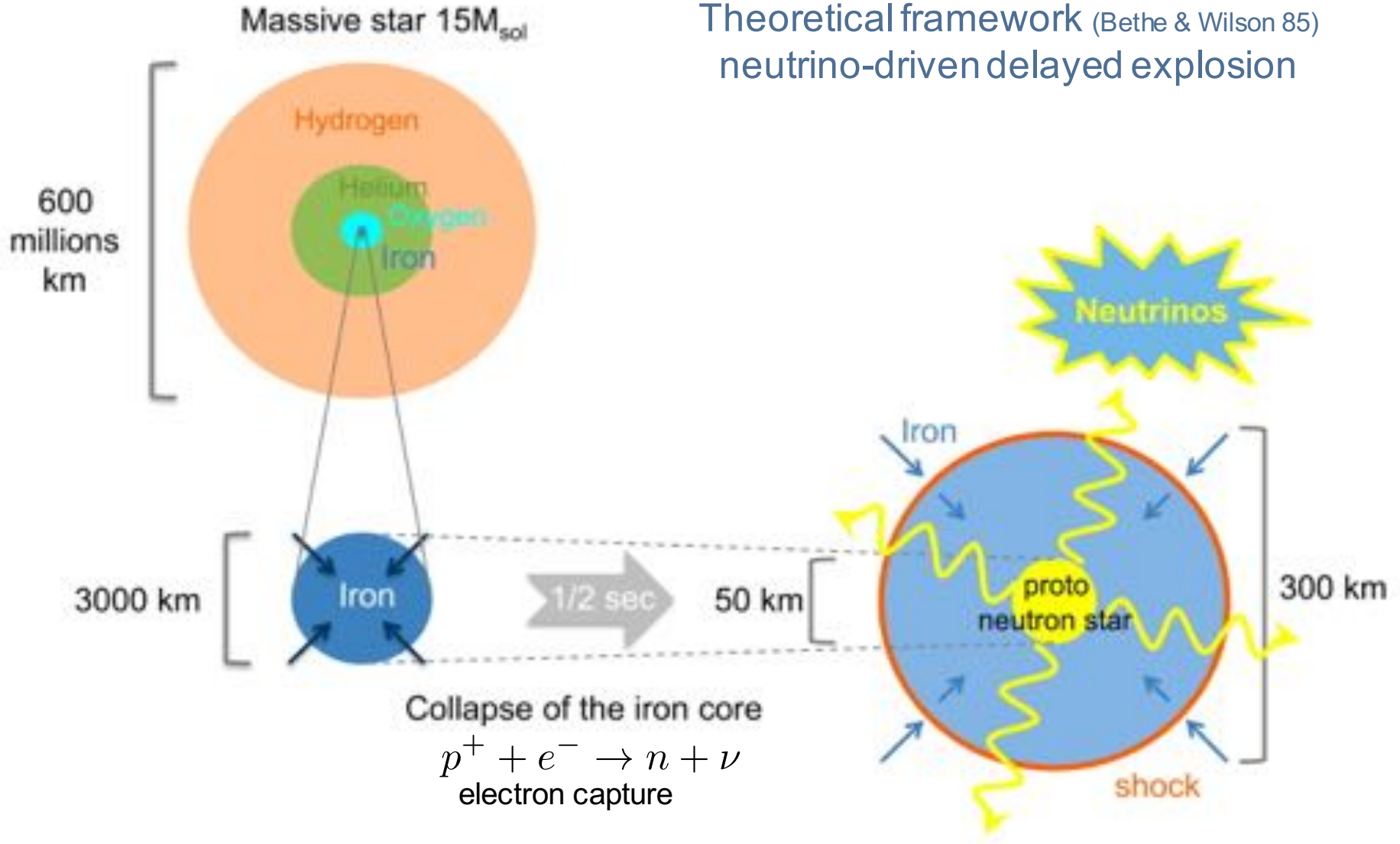
Hanke+13



Takiwaki+16



Theoretical framework (Bethe & Wilson 85)
neutrino-driven delayed explosion

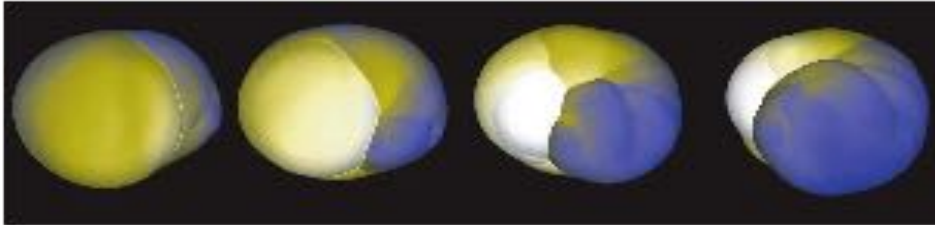


$$\frac{GM_{\text{ns}}^2}{R_{\text{ns}}} \sim 2 \times 10^{53} \text{erg} \left(\frac{30\text{km}}{R_{\text{ns}}} \right) \left(\frac{M_{\text{ns}}}{1.5M_{\text{sol}}} \right)^2$$

modest energy in differential rotation: $E_{\text{diff}} < E_{\text{rot}} \sim 2.4 \times 10^{50} \text{erg} \left(\frac{M_{\text{ns}}}{1.5M_{\text{sol}}} \right) \left(\frac{R_{\text{ns}}}{10\text{km}} \right)^2 \left(\frac{10\text{ms}}{P_{\text{ns}}} \right)^2$

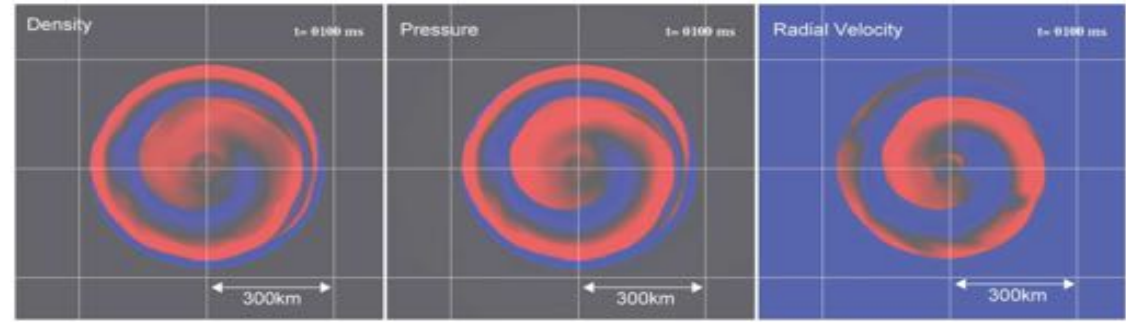
Few 3D simulations including rotation

slow rotation ($j = 10^{15} \text{ cm}^2/\text{s}$): spiral SASI

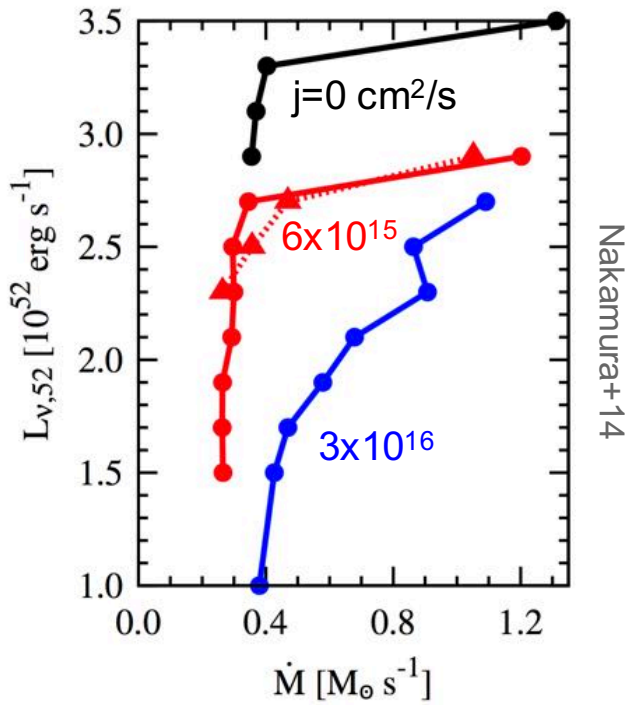


Blondin & Mezzacappa 07

Very fast rotation ($j = 4 \times 10^{16} \text{ cm}^2/\text{s}$): low- $T/|W|$ instability

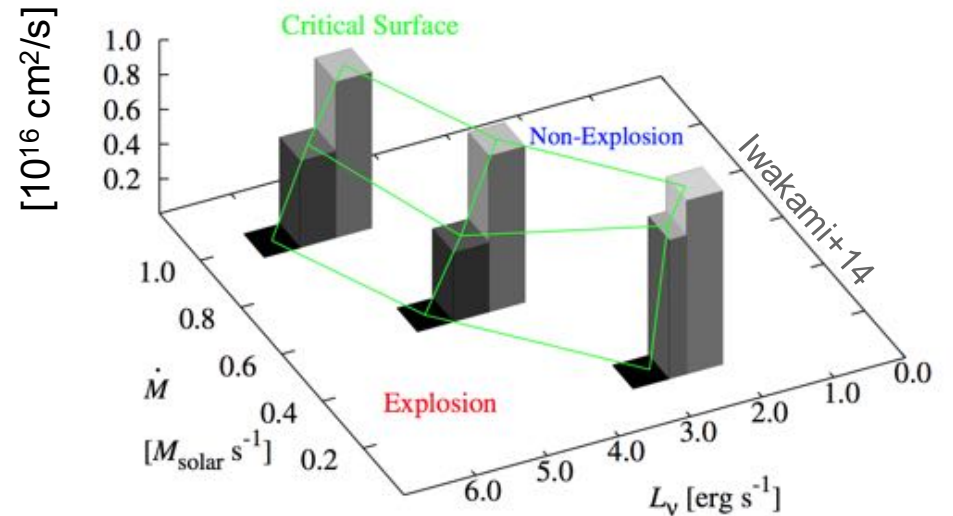


Takiwaki+16



$\Delta L_v / L_v \sim 10\%$

for $j = 5 \times 10^{15} \text{ cm}^2/\text{s}$
 (~ms pulsar)
 modest effect compared
 to the rotational kinetic
 energy involved



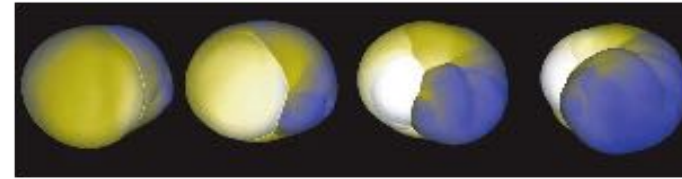
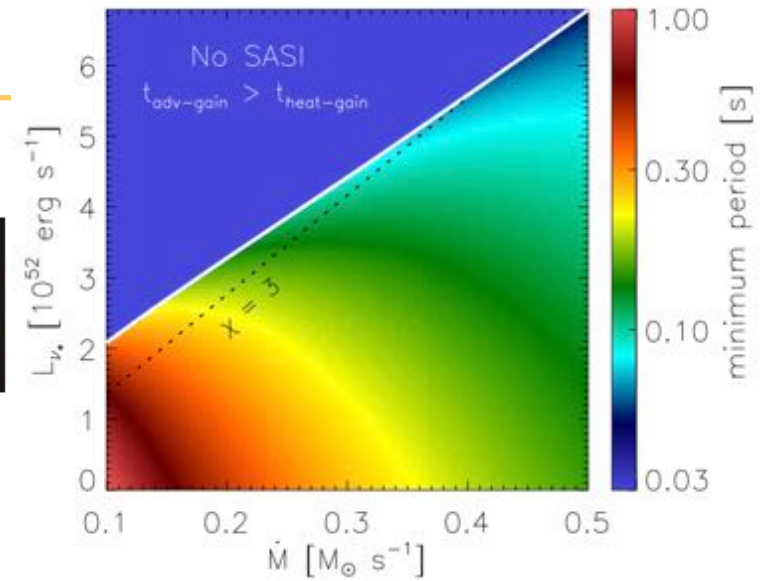
$$E_{\text{rot}} \sim 1.5 \times 10^{52} \text{ erg} \left(\frac{M_{\text{ns}}}{1.5 M_{\text{sol}}} \right) \left(\frac{10 \text{ km}}{R_{\text{ns}}} \right)^2 \left(\frac{j}{5 \times 10^{15} \text{ cm}^2 \text{ s}^{-1}} \right)^2$$

Angular momentum in the final stages of stellar evolution

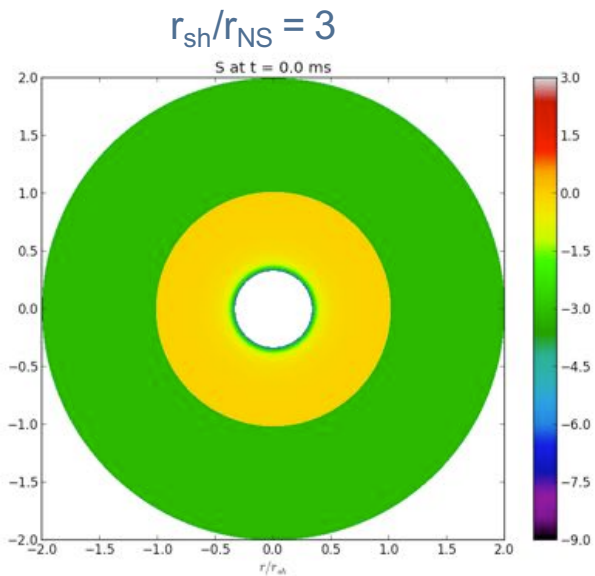
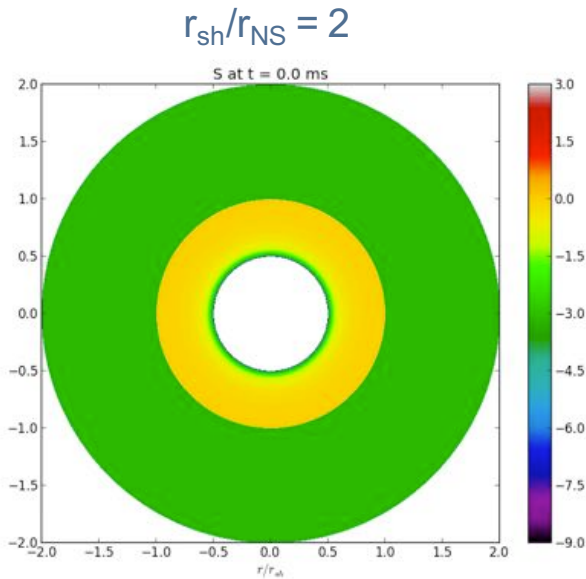
RSG/BSG	O → Si burning	Si → Fe burning	core-collapse dynamics	NS, BH
$R_{\text{He}} \sim 10^7 \text{ km}$ millions yrs	$R_{\text{O}} \sim 3-8 \times 10^3 \text{ km}$ months	$R_{\text{Fe}} \sim 1700 \text{ km}$ hours	$R_{\text{sh}} \sim 150 \text{ km}$ $R_{\text{PNS}} \sim 50 \text{ km}$ $< \text{sec}$	$R_{\text{NS}} \sim 10 \text{ km}$ $> 10^3 \text{ yrs}$
hints of slow core rotation from RG asterosismology (Cantiello+14) possible IGW-magnetized core coupling? (Jim's talk)	convective inhomogeneities (Müller+16)	convectively excited internal gravity waves (Fuller+15)	ν -driven convection, SASI , low $T/ W $ (Kazeroni+17)	electromagnetic spindow shorter timescales? r-mode braking, ν -powered magnetic wind, fallback+propeller braking strong $B > 10^{13} \text{ G}$ field?
from "slow" to $\sim 4 \times 10^{14} \text{ cm}^2/\text{s}$ with Tayler-Spruit dynamo to $\sim 5 \times 10^{16} \text{ cm}^2/\text{s}$ without (Heger+05)		$\sim 10^{13-14} \text{ cm}^2/\text{s}$ stochastic	$\sim 10^{13-14} \text{ cm}^2/\text{s}$ stochastic	uniform distribution $P_{\text{NS}} \sim 10-100 \text{ ms}$ at birth $j_{\text{NS}} \sim 6 \times 10^{13-14} \text{ cm}^2/\text{s}$ (Popov & Turolla 12) ms breakup threshold $\sim 6 \times 10^{15} \text{ cm}^2/\text{s}$

Spin up of the neutron star induced by the spiral mode of SASI

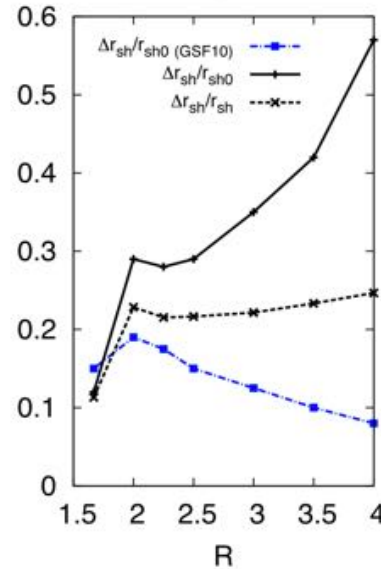
Guilet & Fernandez 14



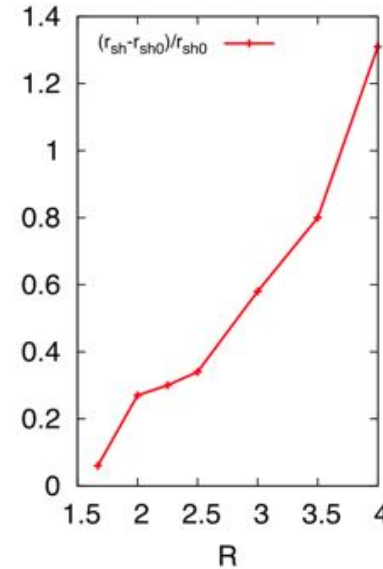
Blondin & Mezacappa 07



Saturation amplitude of SASI

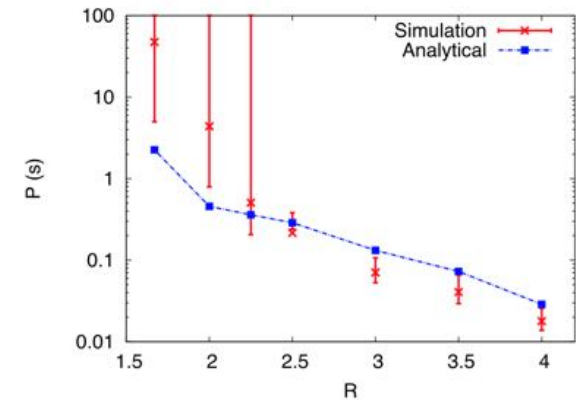


Shock expansion



Kazeroni+16

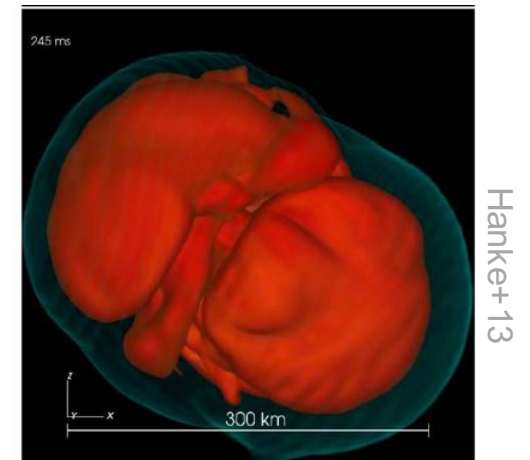
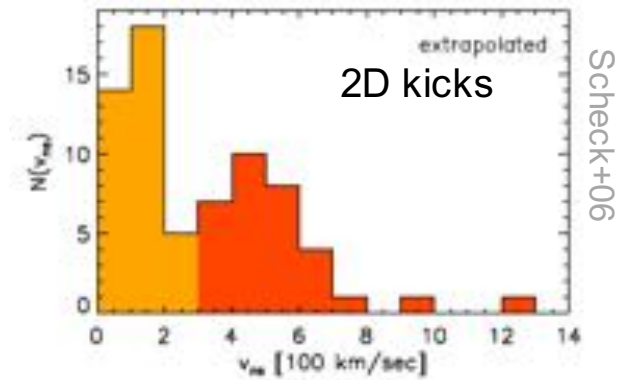
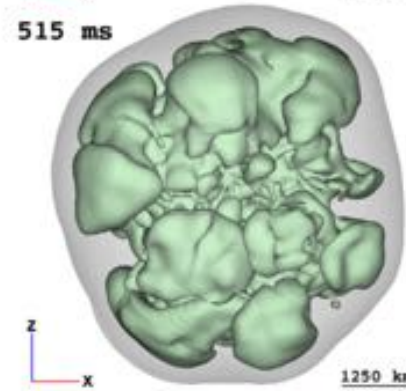
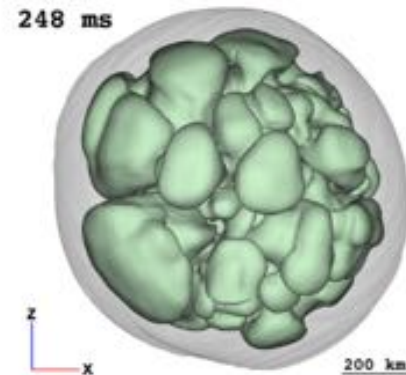
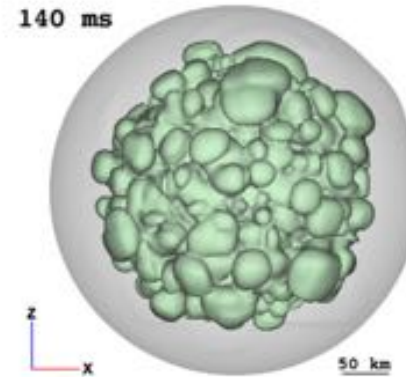
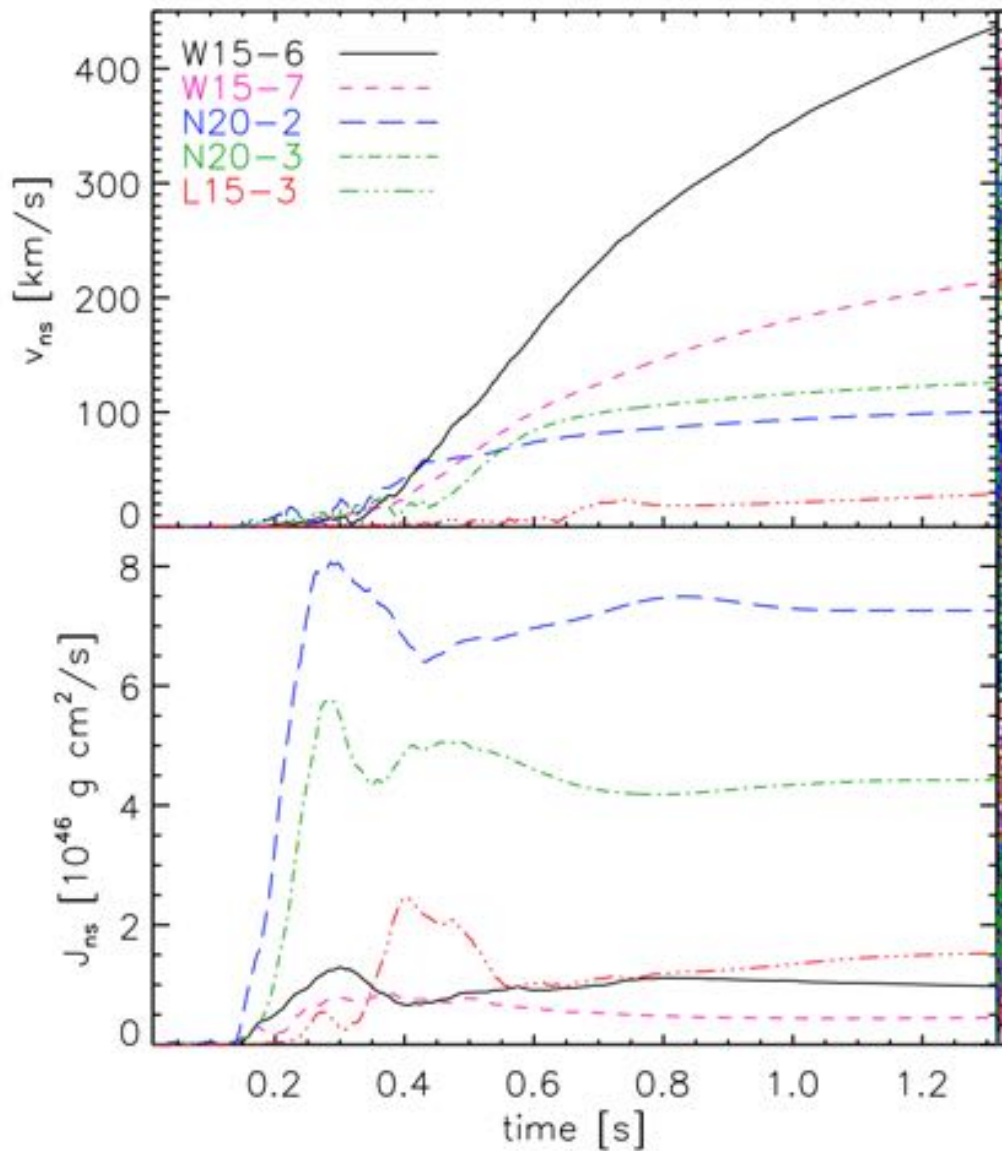
2D cylindrical
adiabatic shock
perfect gas $\gamma=4/3$
 ν -cooling function
no ν -heating



pulsar spin periods down to ~ 50 ms

- the strength of SASI increases with the radius ratio $R = r_{\text{sh}}/r_{\text{NS}}$
- unexpected stochasticity

Kick and spin from multi-D simulations: different timescales



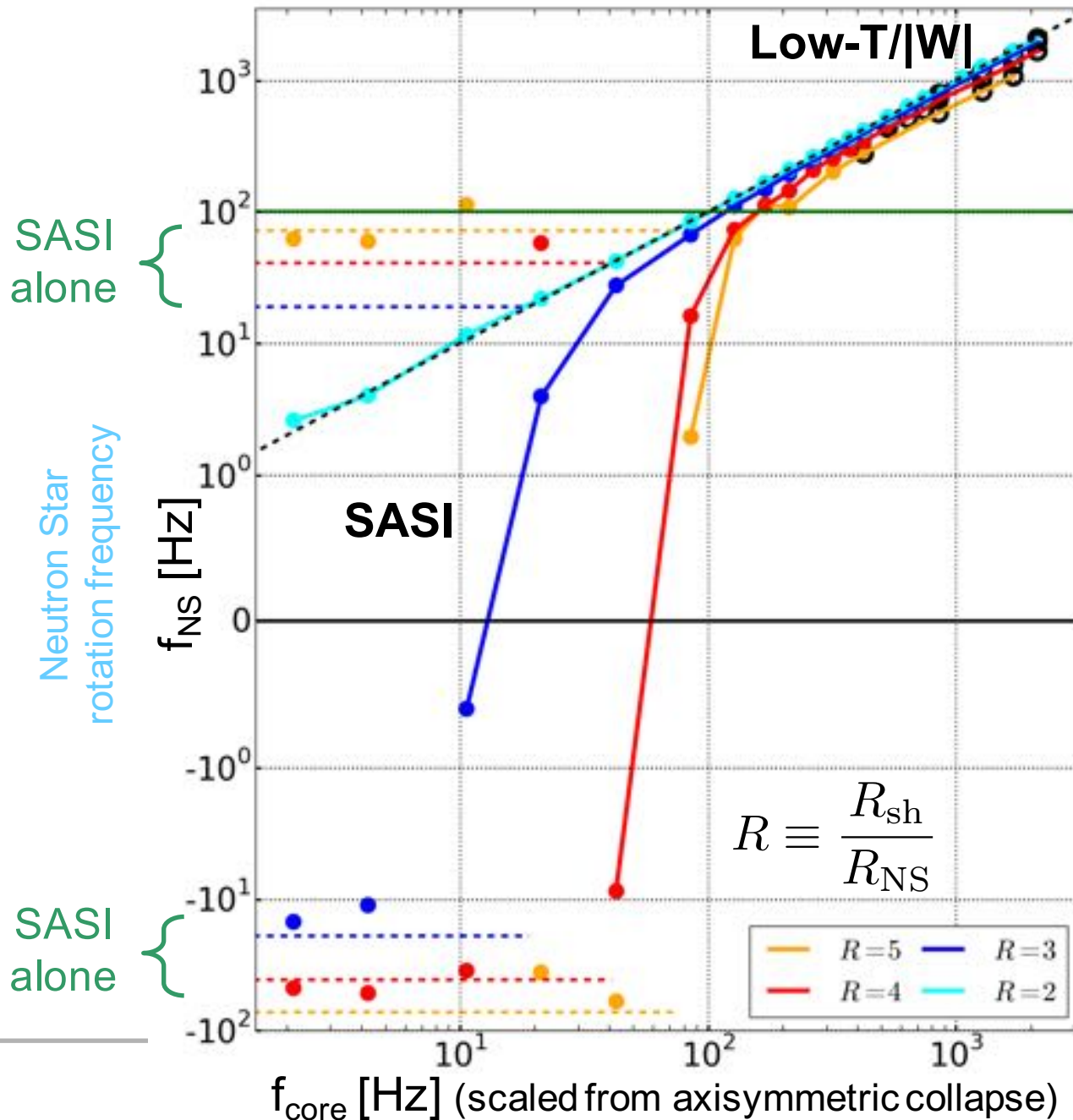
$$\chi \equiv \int_{\text{sh}}^{\text{gain}} \omega_{\text{BV}} \frac{dr}{v_r} < 3$$

$P_{ns} \sim 100\text{ms to } 8\text{s}$ (Wongwathanarat+13)

SASI vs convection?

Spin-up or spin-down of the neutron star?

(Kazeroni+17)

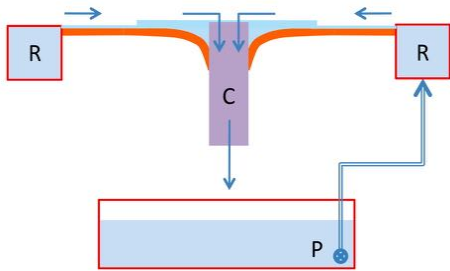


range of NS spin at birth

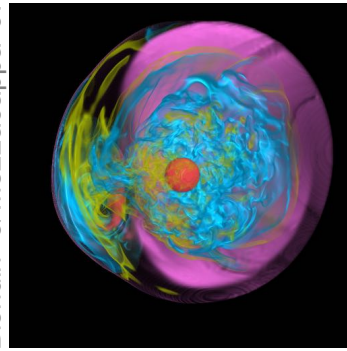
For a strong rotation rate, the corotation instability decelerates the neutron star by less than 30%.

SWASI: an experimental analogue of SASI

Shallow Water Analogue of a Shock Instability



Blondin & Mezzacappa 07



adiabatic gas

$$c_s^2 \equiv \frac{\gamma P}{\rho}$$

$$\Phi \equiv -\frac{GM_{\text{ns}}}{r}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho v) = 0$$

$$\frac{\partial v}{\partial t} + (\nabla \times v) \times v + \nabla \left(\frac{v^2}{2} + \frac{c_s^2}{\gamma - 1} + \Phi \right) = \frac{c_s^2}{\gamma} \nabla S$$

Inviscid shallow water is analogue to an isentropic gas $\gamma=2$

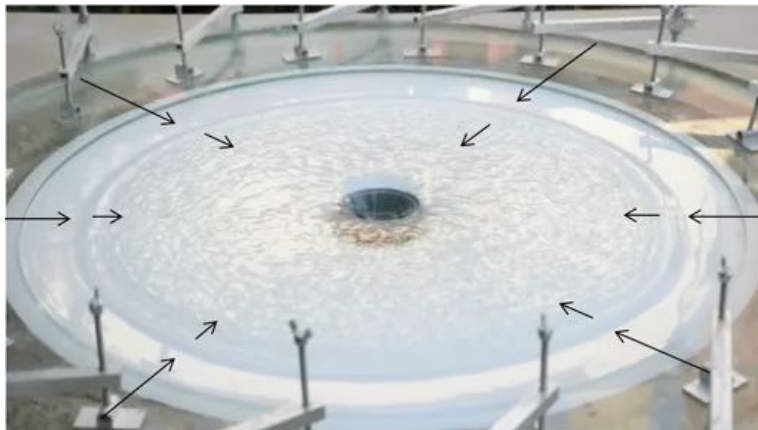
St Venant

$$c_{\text{sw}}^2 \equiv gH$$

$$\Phi \equiv gH_\Phi$$

$$\frac{\partial H}{\partial t} + \nabla \cdot (Hv) = 0$$

$$\frac{\partial v}{\partial t} + (\nabla \times v) \times v + \nabla \left(\frac{v^2}{2} + c_{\text{sw}}^2 + \Phi \right) = 0$$



acoustic waves
shock wave
pressure

surface waves
hydraulic jump
depth

expected scaling

$$\frac{t_{\text{ff}}^{\text{sh}}}{t_{\text{ff}}^{\text{jp}}} \equiv \left(\frac{r_{\text{sh}}}{r_{\text{jp}}} \right) \left(\frac{r_{\text{sh}} g H_{\text{jp}}}{GM_{\text{NS}}} \right)^{\frac{1}{2}} \sim 10^{-2}$$

shock radius $\times 10^{-6}$

200 km \rightarrow 20 cm

oscillation period $\times 10^2$

30 ms \rightarrow 3 s

SWASI: simple as a garden experiment

November 2010



October 2010



June 2010



May 2010

February 2012



November 2013



June 2014

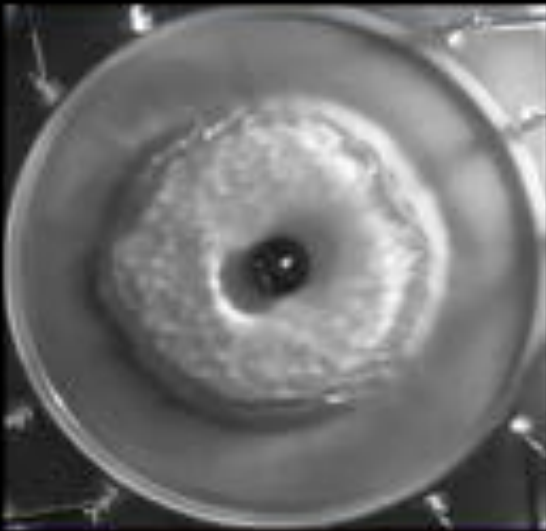


February 2017

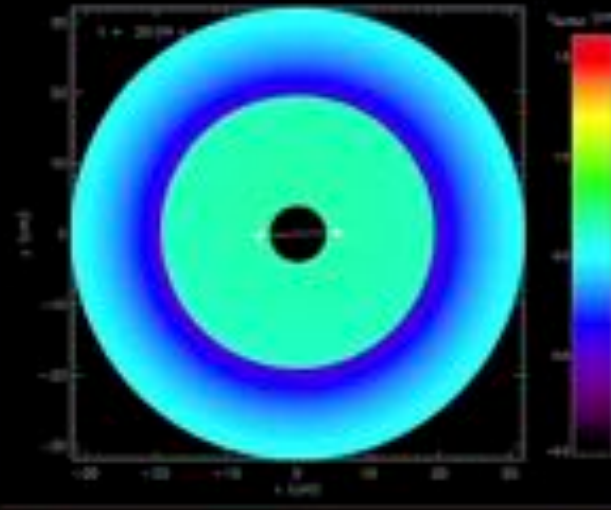
Dynamics of water in the fountain

Dynamics of the gas in the supernova core

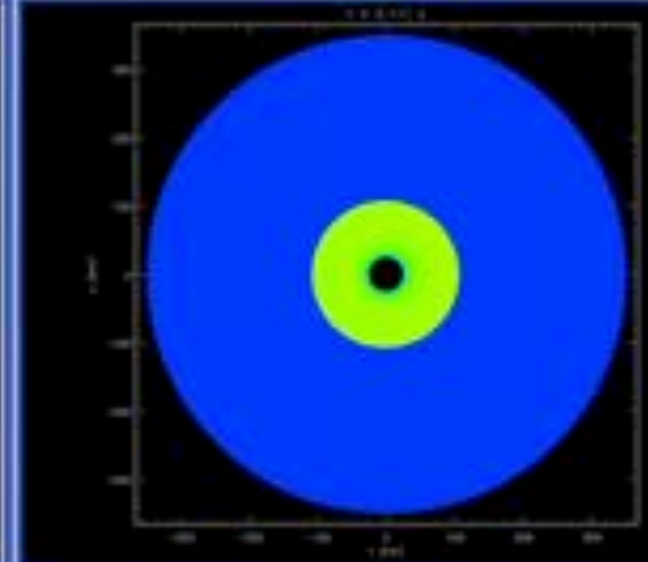
diameter 40cm ← 1 000 000 x bigger → diameter 400km
3s/oscillation ← 100 x faster → 0.03s/oscillation



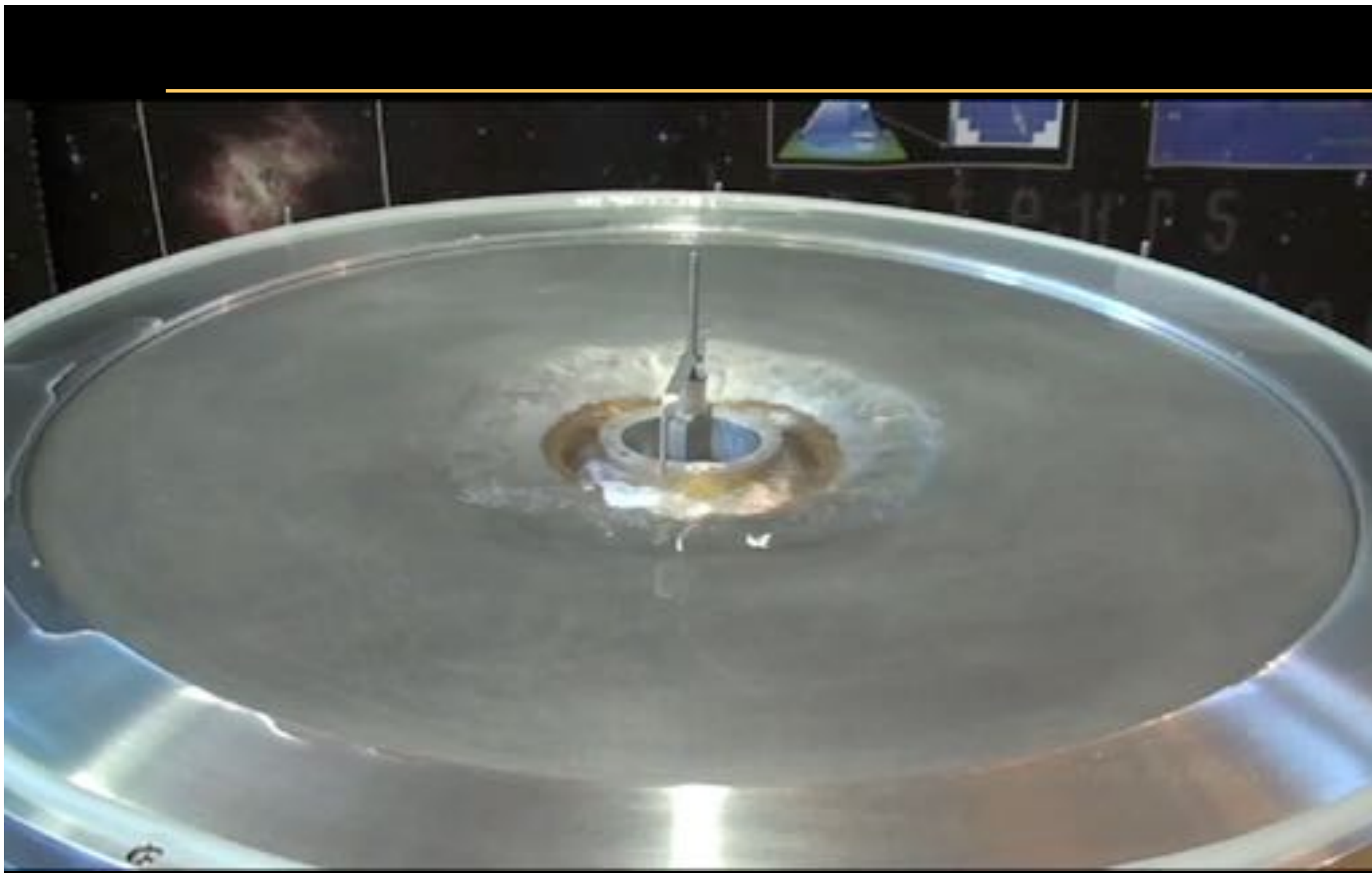
Expérience hydraulique



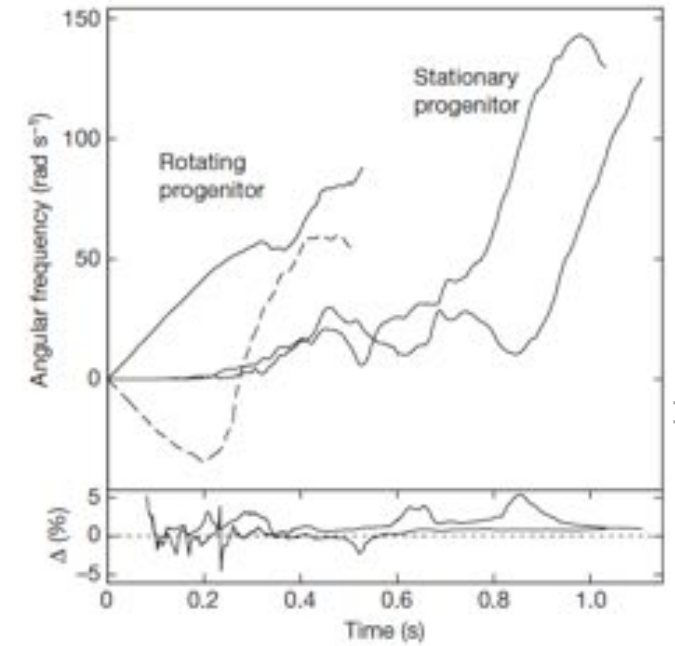
Simulation numérique de l'expérience hydraulique



*Simulation numérique de l'onde de choc
dans le cœur de la supernova*



Rotating progenitor: accreted angular momentum changes its sign as SASI grows



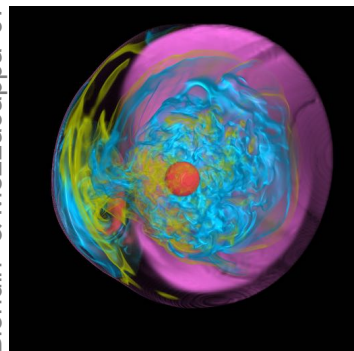
Blondin & Mezzacappa 07

-significant shear even when the centrifugal force $\Omega^2 R$ is weak $\frac{\Omega}{\Omega_{NS}} \propto \left(\frac{R_{NS}}{R}\right)^2$

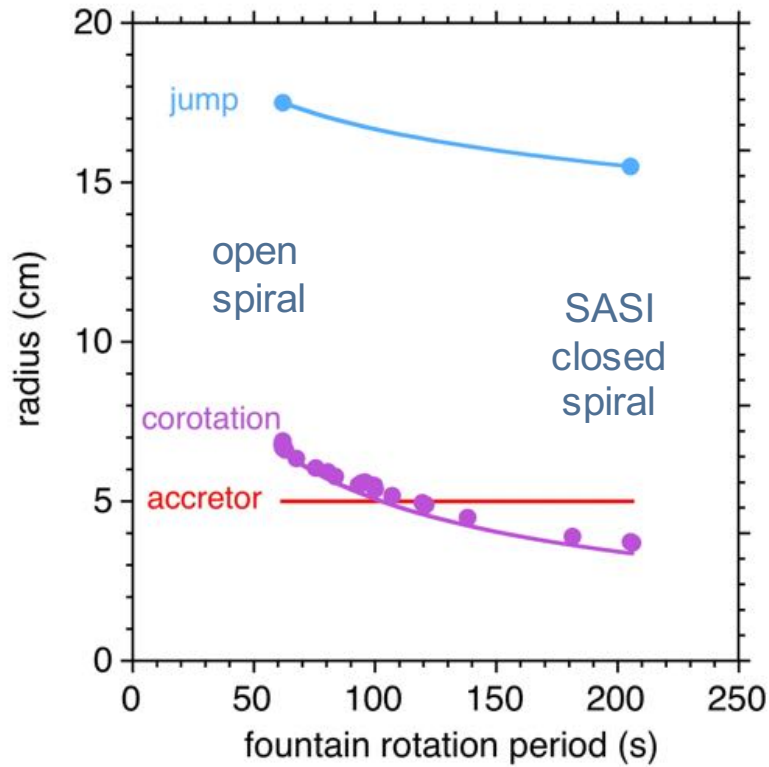
-the prograde mode is favoured by differential rotation as in shocked accretion

Blondin & Mezzacappa 07, Yamasaki & Foglizzo 08, Kazeroni+17

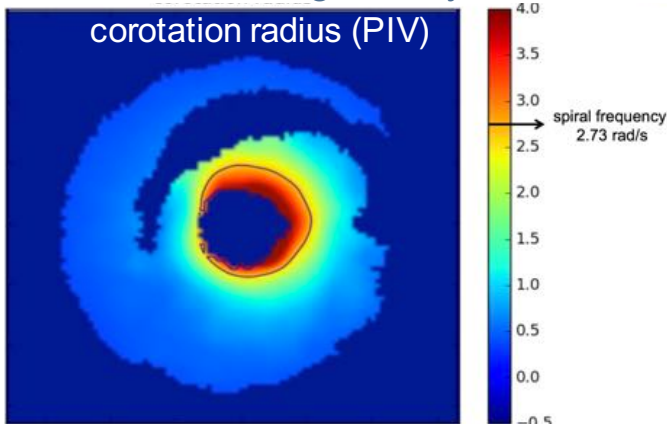
Blondin & Mezzacappa 07



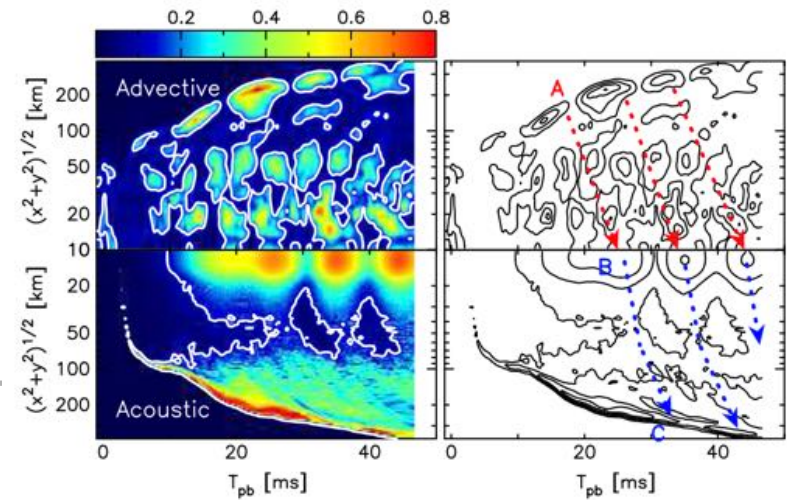
Increasing the rotation rate: continuous transition from SASI to the corotation instability



the rotation period is gradually decreased (205s → 62s)
the flow rate is gradually decreased (1.1 L/s → 0.59 L/s)



→ the gravitational wave signature of the low $T/|W|$ instability may be hard to disentangle from the SASI oscillation (Kuroda+14)



Unexpectedly robust spiral shock driven at the corotation radius when the inner rotation rate reaches 20% Kepler ($\text{low } T/|W|=0.02$)



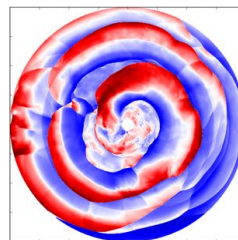
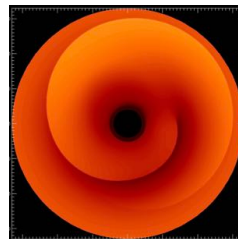
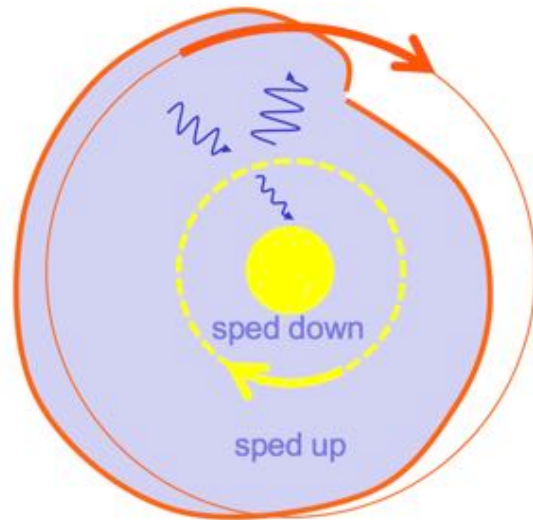
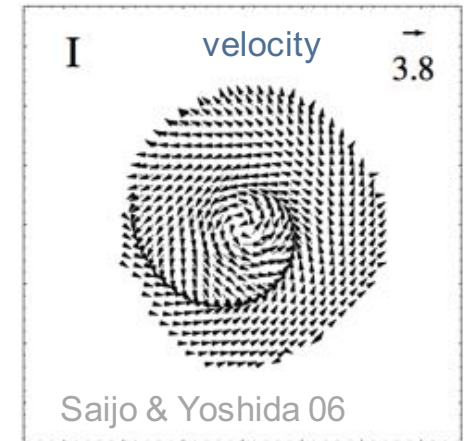
Spiral instability with a weak shock

Radial accretion enforces differential rotation

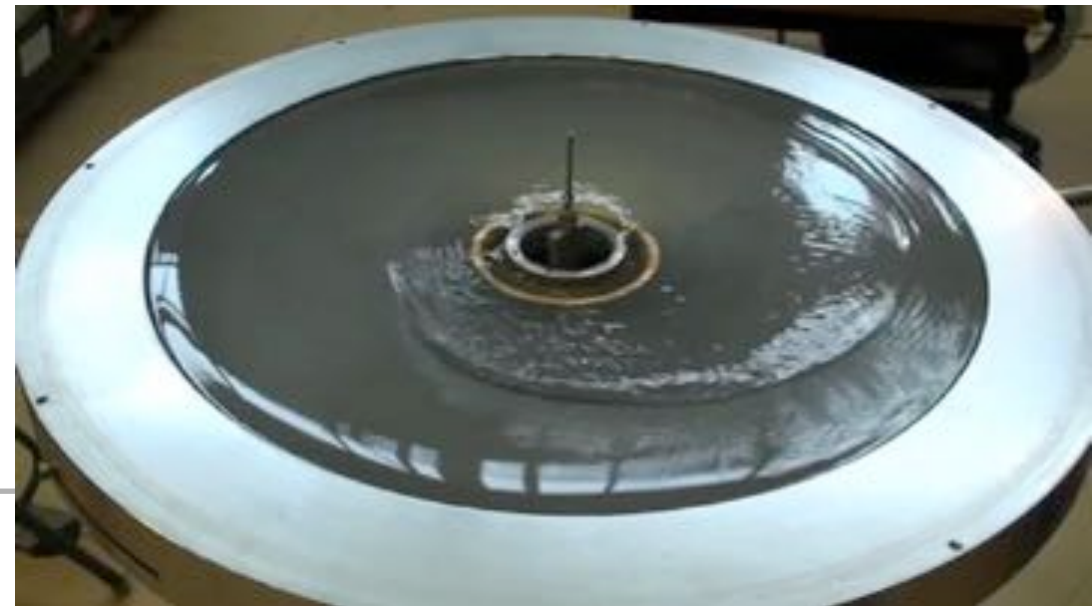
$$\frac{\Omega}{\Omega_{\text{NS}}} \propto \left(\frac{R_{\text{NS}}}{R} \right)^2$$

Analogue to the "low $T/|W|$ instability" of a neutron star rotating differentially

(Shibata+02,03, Saijo+03,06, Watts+05, Corvino+10, Passamonti & Andersson 15)



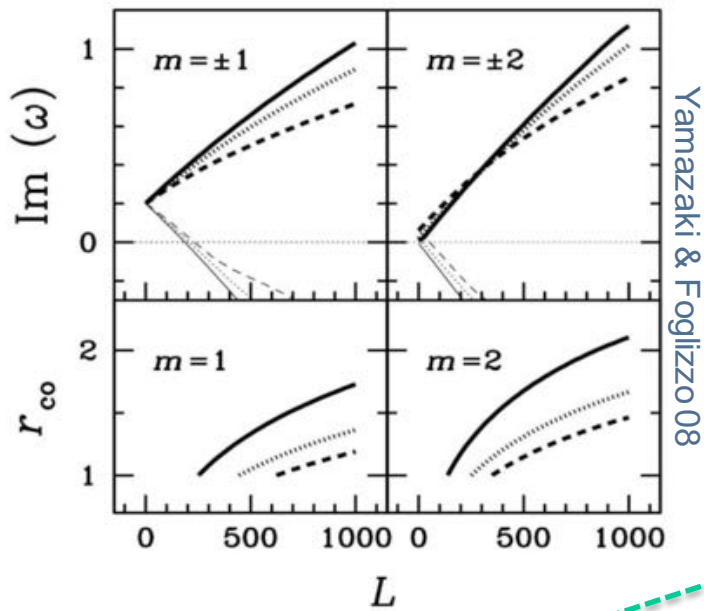
Spiral instability with subsonic accretion



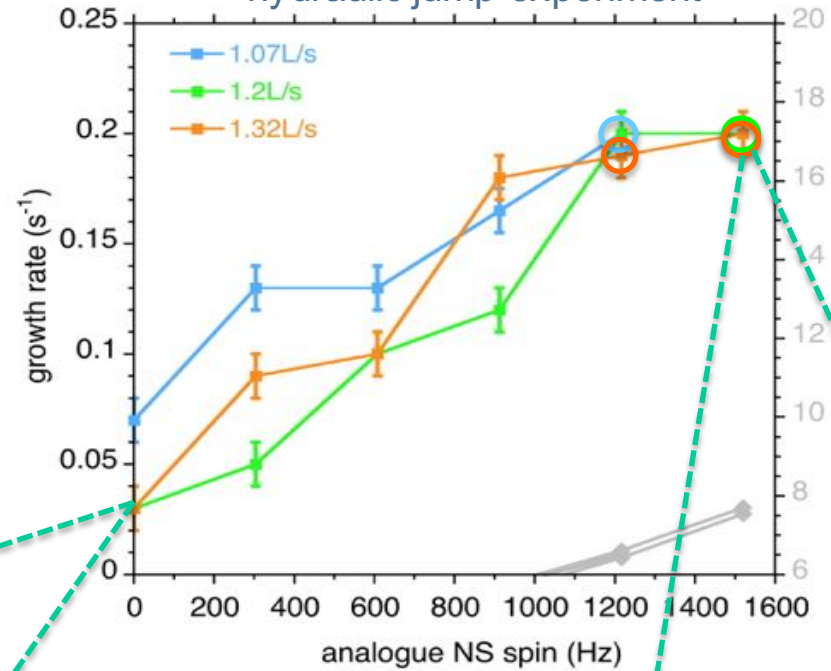
Instability mechanism: interaction of a corotation radius with acoustic waves (Papaloizou & Pringle 84, Goldreich & Narayan 85)

Rotation effects in the experiment compared to gas dynamics

shocked gas dynamics

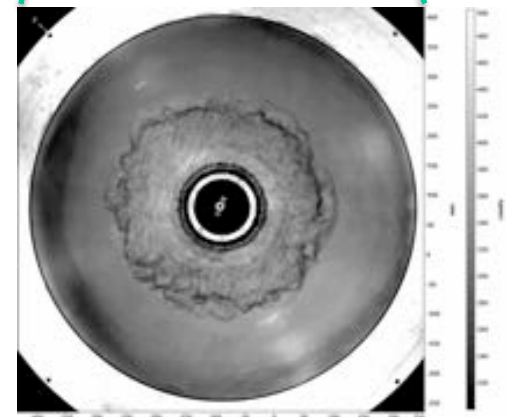
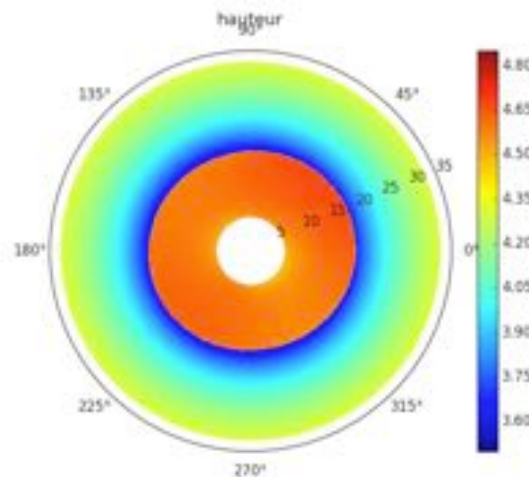
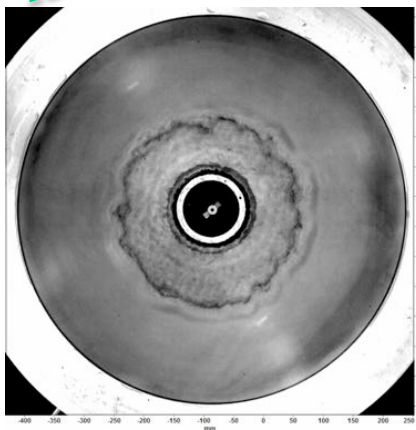


hydraulic jump experiment



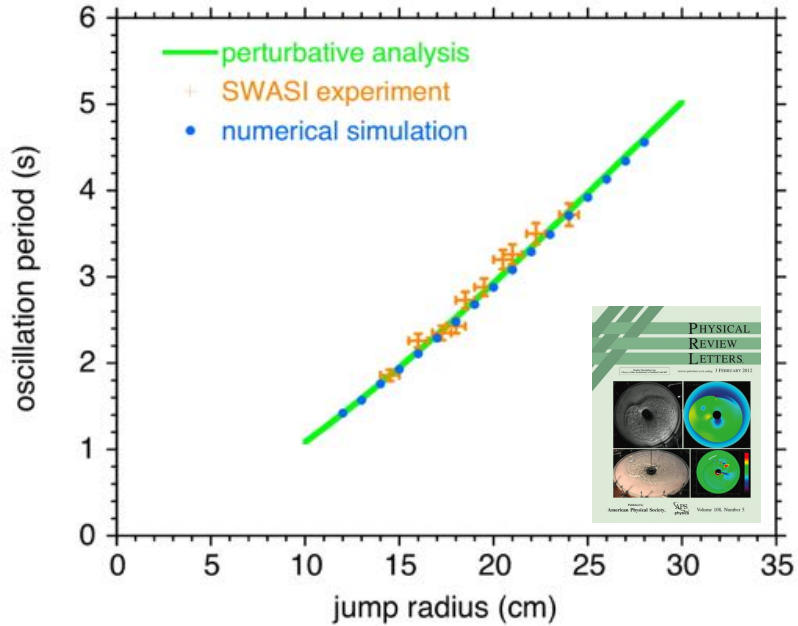
SASI
 $m = \pm 1$
 sloshing \rightarrow spiral

$m = 1, 2$ spiral
 with corotation

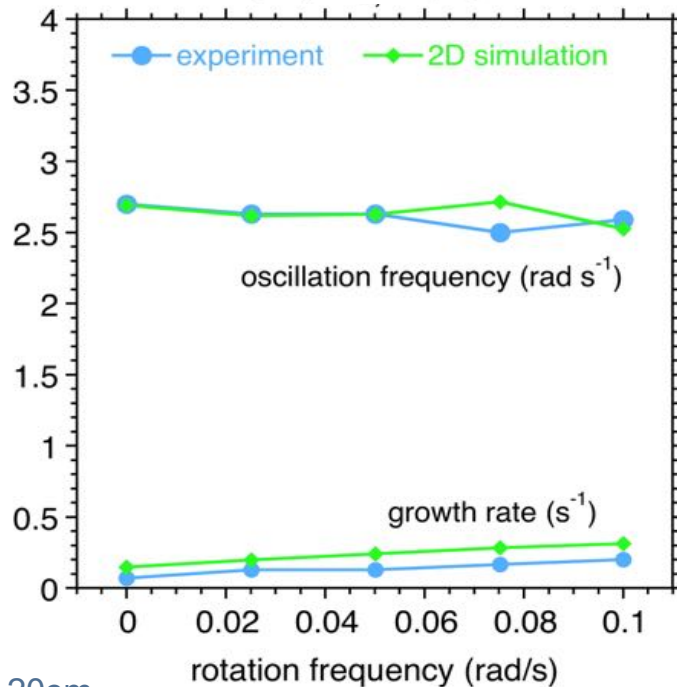


same increase of the growth rate with the core rotation as in gas accretion & same trend towards a $m=2$ linear instability

Experimental growth rate and oscillation period compared to shallow water modelling

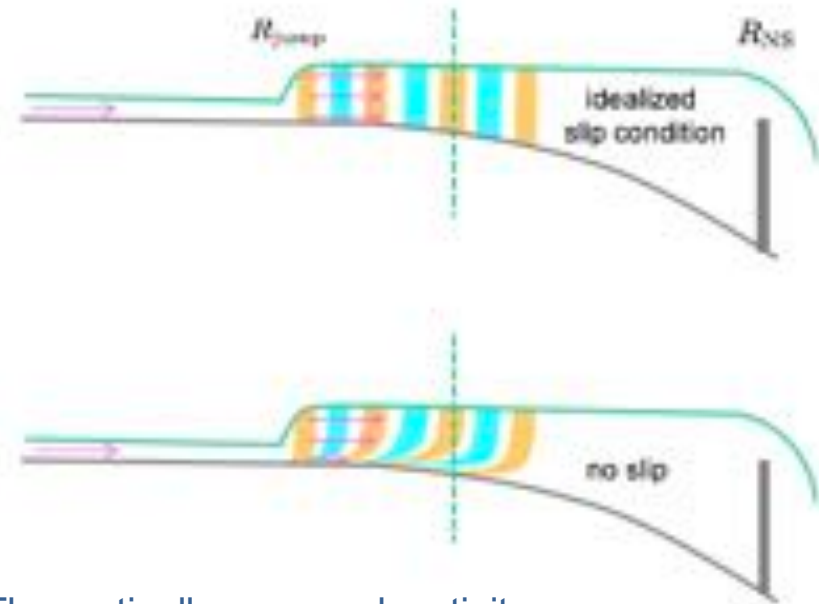


Foglizzo+12



$R_{jp} \sim 20\text{cm}$

- excellent modelling of the oscillation frequency
limited by the measured radial width of the hydraulic jump
- systematic offset of the experimental growth rate
expected phase mixing of the dragged vorticity



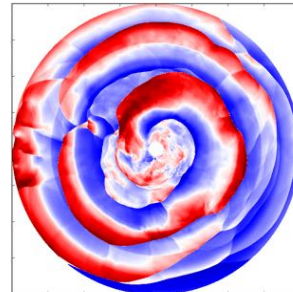
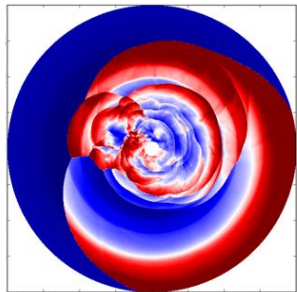
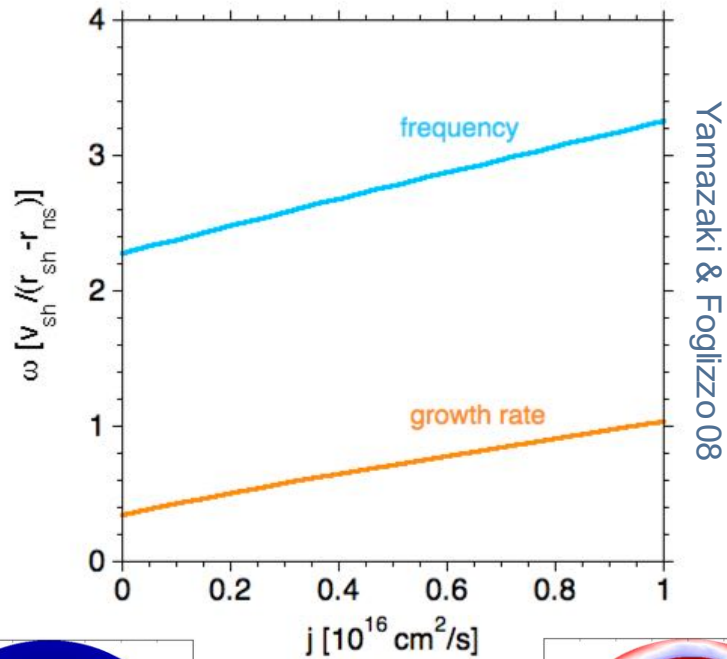
The vertically averaged vorticity is damped by a factor Q

$$Q \sim \int_0^H \frac{dz}{H} \cos \left[\frac{\omega_{\text{SASI}} \Delta R}{v(z)} \right] \sim 0.27 \text{ (laminar)}$$

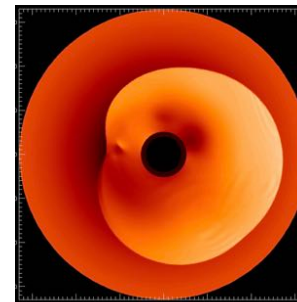
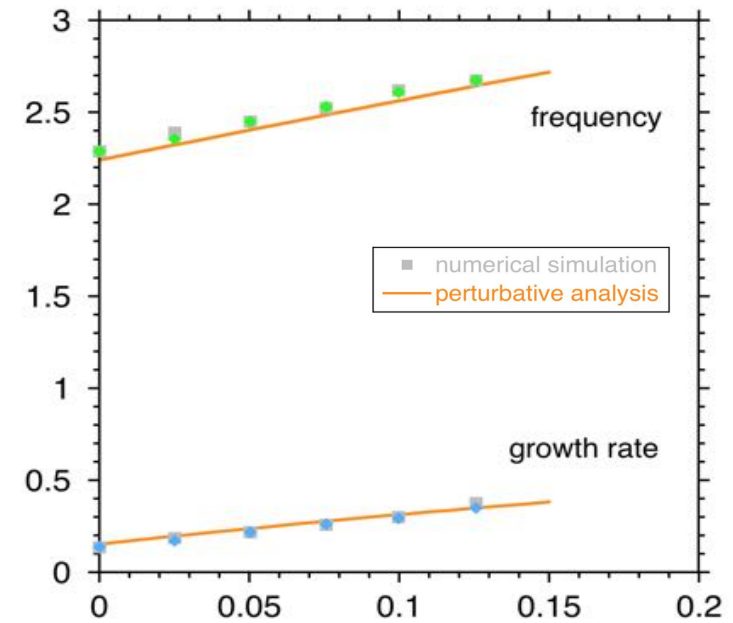
$$\sim 0.52 \text{ (turbulent)}$$

Rotation effect in shallow water equations compared to gas dynamics

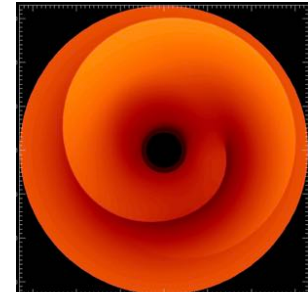
shocked gas dynamics



shallow water equations



fountain rotation



Why is the prograde mode of SASI destabilized by rotation?
 Why is the transition to the corotation instability so smooth?

The St Venant system of shallow water equations is a simpler framework to understand the coupling of SASI with rotation and the transition to the low $T/|W|$ instability:

- adiabatic equations (no neutrino cooling)
- no buoyancy effects

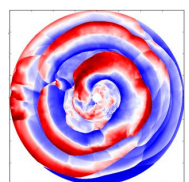
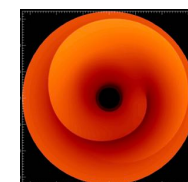
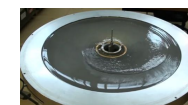
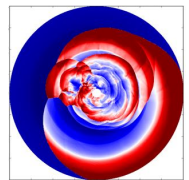
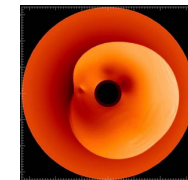
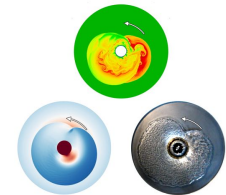
Conclusions

2D Cylindrical gas dynamics suggests that (Kazeroni+17)

- SASI can account for pulsar rotation periods down to ~ 50 ms
- for rotation rates > 100 Hz the corotation instability decreases the pulsar spin by $< 30\%$

Two core collapse instabilities captured in a hydraulic experiment

- an intuitive approach to multi-D processes that produce GW
- experimental results confirmed by a shallow water numerical model
- first** experimental confirmation that spiral SASI can produce a counter-spinning neutron star
- first** experimental demonstration of the 'low $T/|W|$ ' instability
- the corotation instability 'low $T/|W|$ ' connects smoothly to SASI
→GW signatures?



-
- the theory of the corotation instability in a postshock accretion flow is still missing