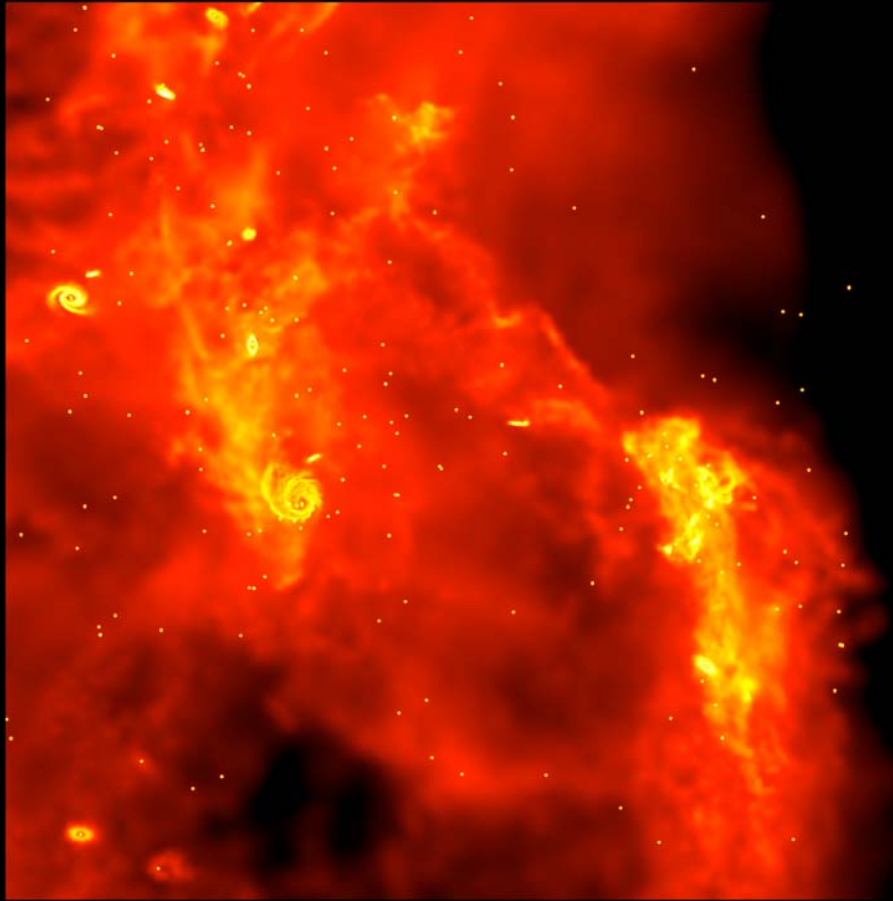


The Competitive Accretion Debate



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What is CA and how does it work?

- Theory to help explain the form of the IMF, motivated by two properties of star formation:
 - Stars form in groups, (associations, clusters).
 - Clusters tend to be mass-segregated, perhaps from birth.
- Works by considering how much mass a newly formed star can accrete, over and above the mass which went into collapse to produce it.

Accretion and the IMF...

Accretion rate:

$$\dot{M}_* = \pi \rho V_{\text{rel}} R_{\text{acc}}^2$$

Zinnecker (1982)

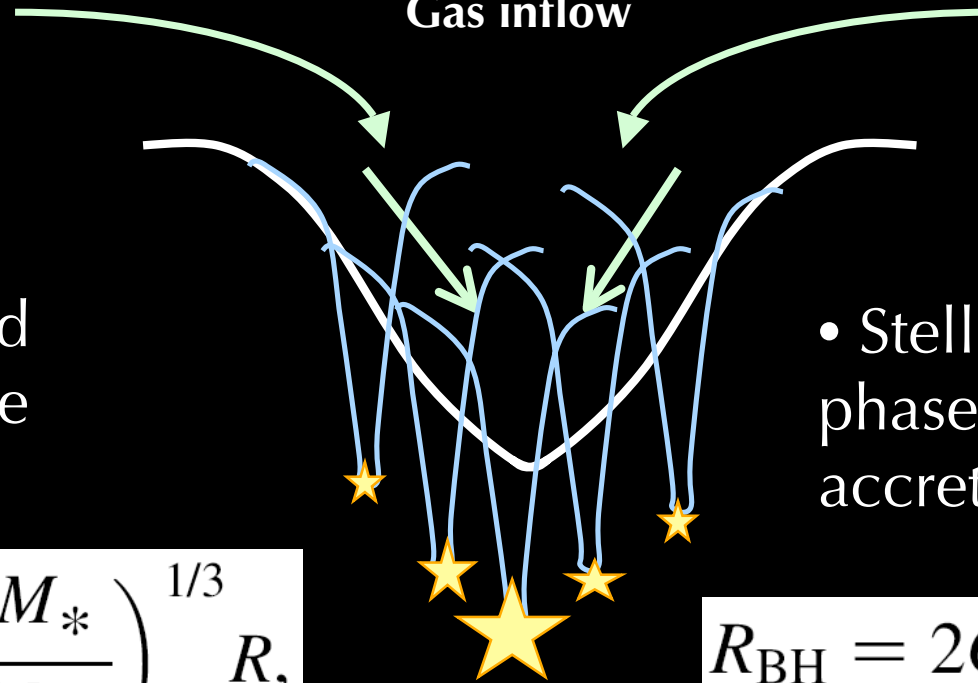
Bonnell et al (2001a,b)

- Gas dominated phase: tidal-lobe accretion,

$$R_{\text{tidal}} \approx 0.5 \left(\frac{M_*}{M_{\text{enc}}} \right)^{1/3} R,$$

$$\frac{dn}{dm} \propto m^{-1.5}$$

Gas inflow



- Stellar dominated phase: Bondi-Hoyle accretion,

$$R_{\text{BH}} = 2GM_*/(V_{\text{rel}}^2 + c_s^2)$$

$$\frac{dn}{dm} \propto m^{-2.5}$$

Hierarchical process...

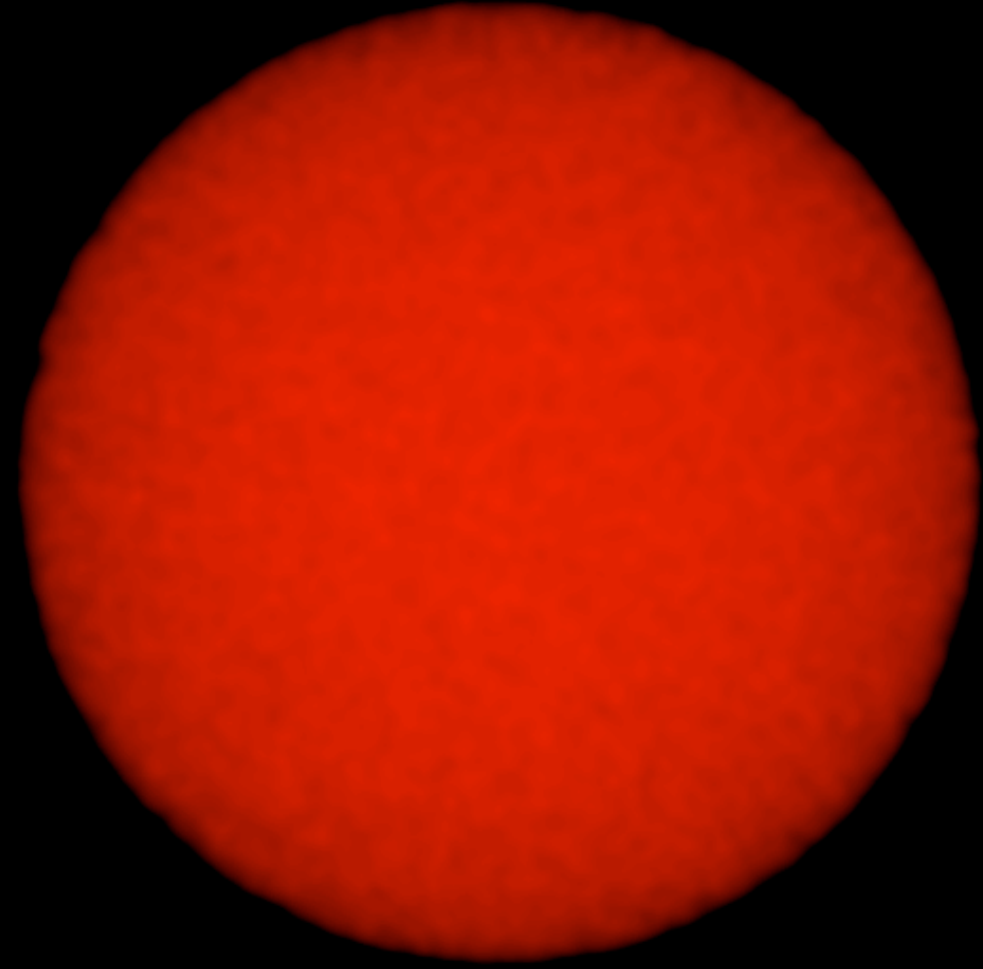
- Hierarchical dissipation of turbulence
- Small scales loose support first:

$$t_{\text{disp}} \sim t_{\text{cross}} \sim L/\sigma(L)$$

$$t_{\text{cross}} \sim L^{0.5}$$

- Followed by collapse of progressively larger regions

The Formation of a Stellar Cluster

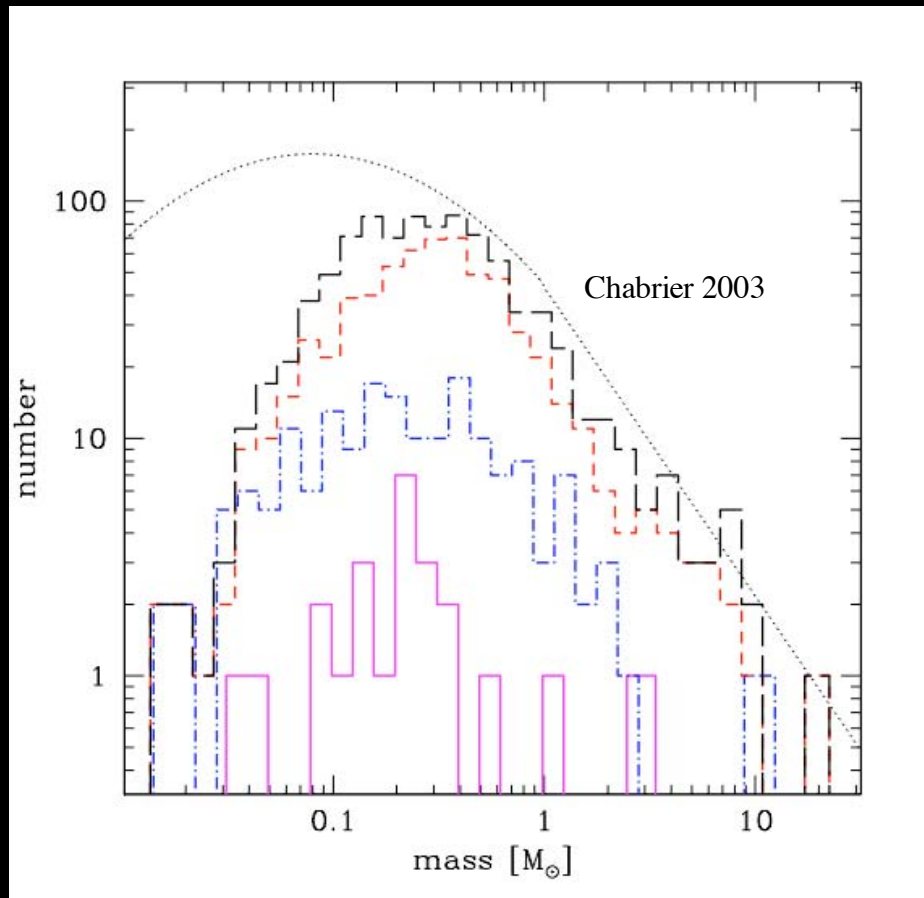


0.25pc

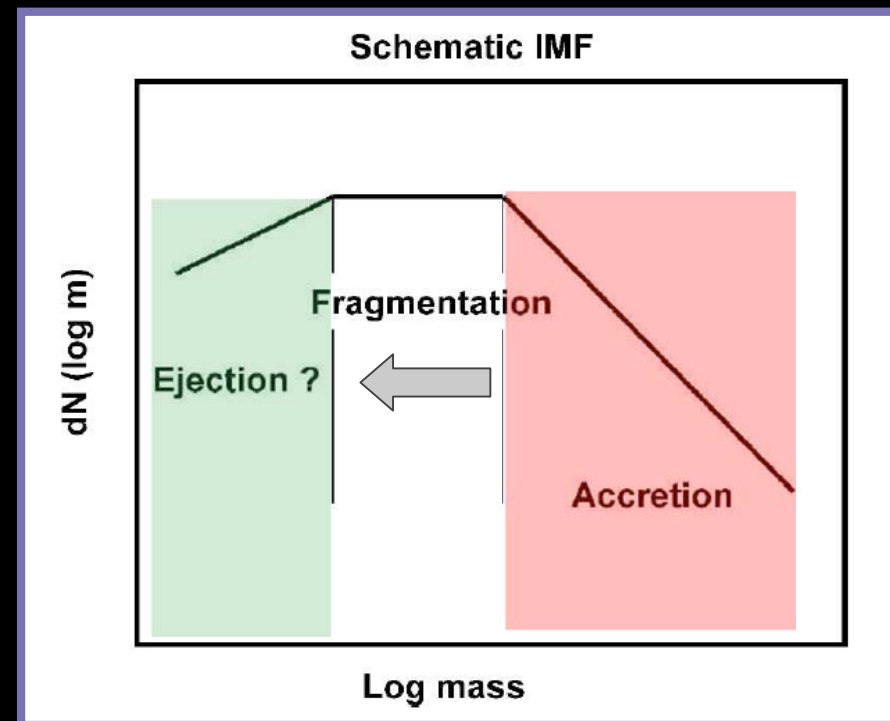
Bonnell, Bate & Vine (2003)

Features of the CA mass function

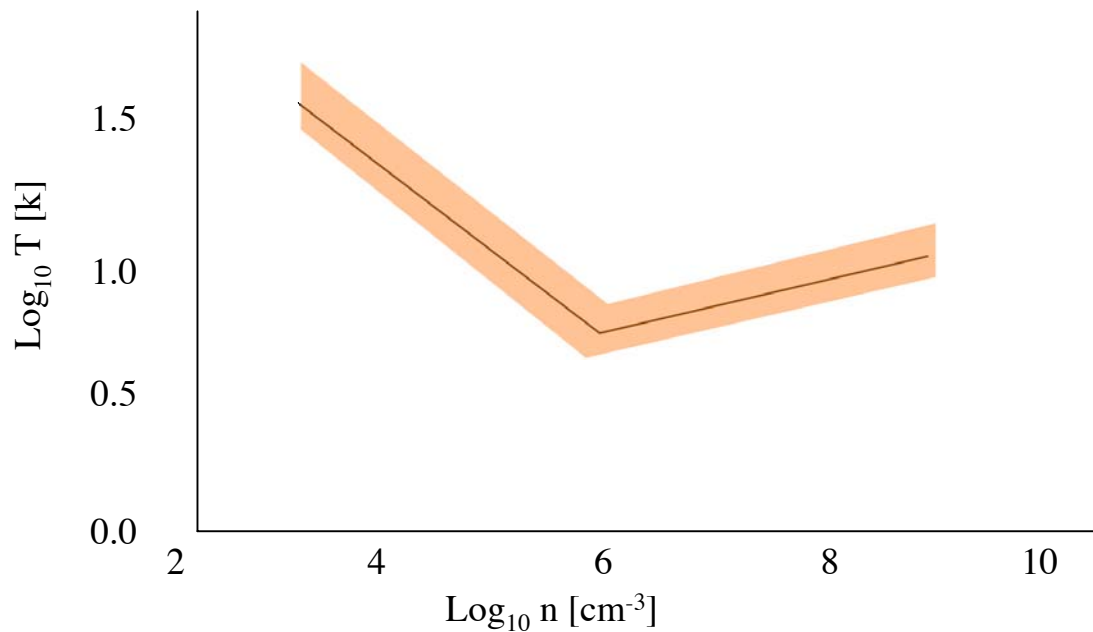
Grows in time



All mass bins are related



Sensitivity to cloud conditions....



Found that changing the initial Jeans mass in the set-up, alters the position of the 'knee' in the IMF.

Does competitive accretion really need such fine tuning?

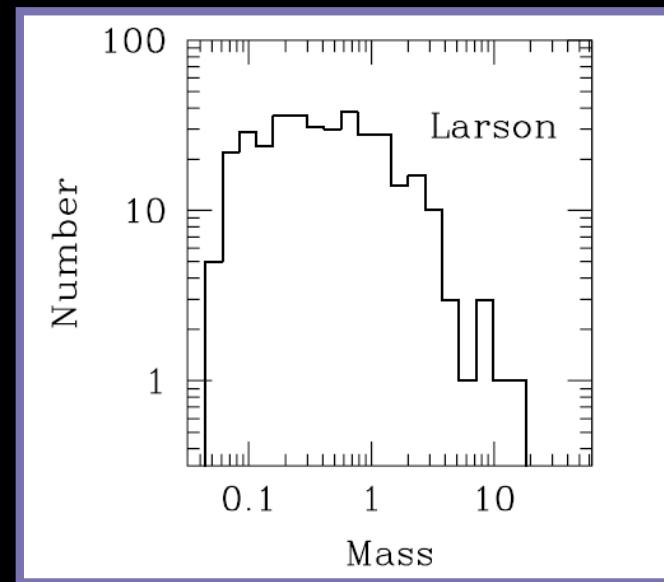
Larson (1985, 2005):

$$T = 4.4 (\rho / 10^{-18})^{-0.27} \text{ K},$$

$$(\rho < 10^{-18} \text{ gcm}^{-3})$$

$$T = 4.4 (\rho / 10^{-18})^{+0.07} \text{ K},$$

$$(\rho > 10^{-18} \text{ gcm}^{-3})$$

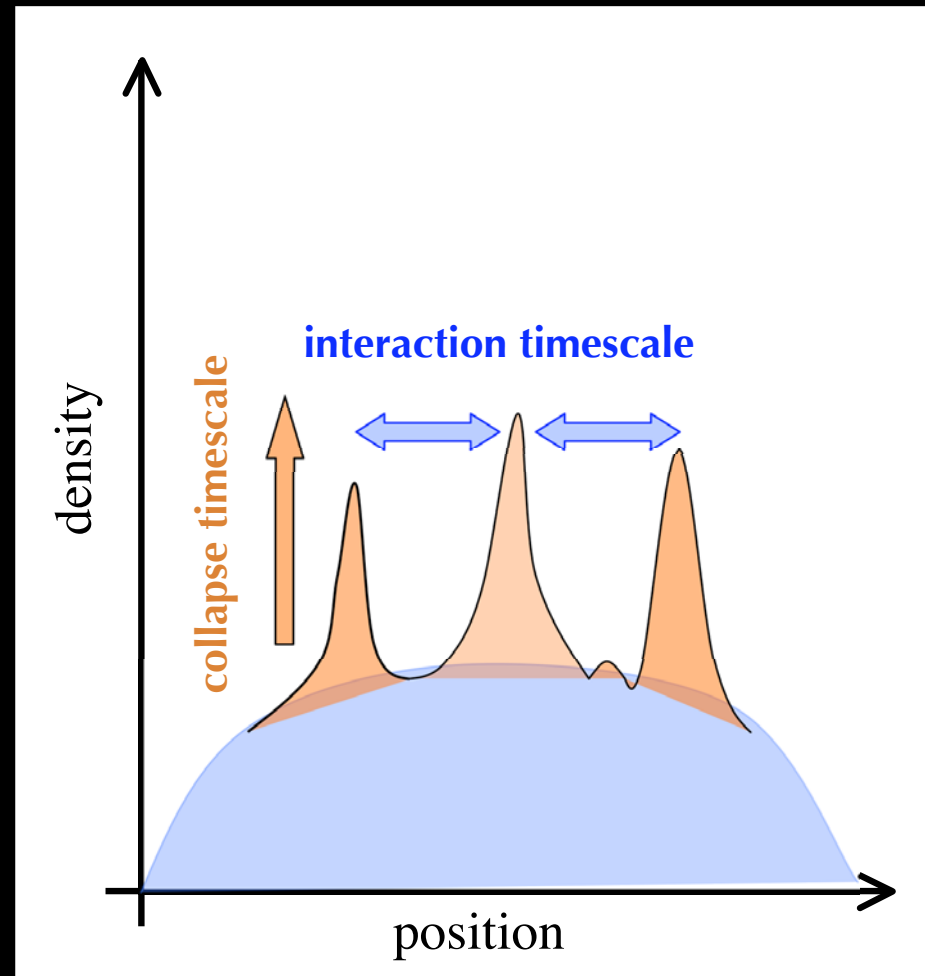


Conditions for CA (1)

- Competitive accretion requires a region in which the collapse timescale and interaction timescale are similar.
- If the clump densities and cloud density are roughly equal, then:

$$t_{\text{inter}} \sim t_{\text{ff}}$$

- Any region with multiple Jeans masses automatically satisfies this requirement.

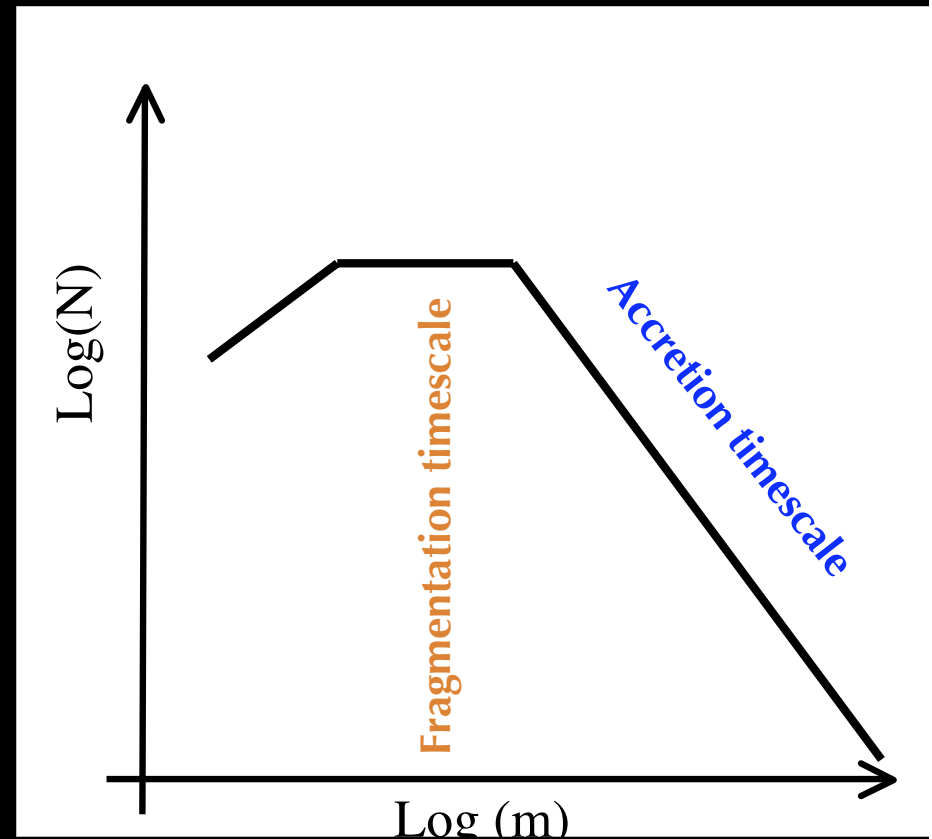


Conditions for CA (2)

- If the ratio of the mass above and below the Salpeter break is to remain the same, then:

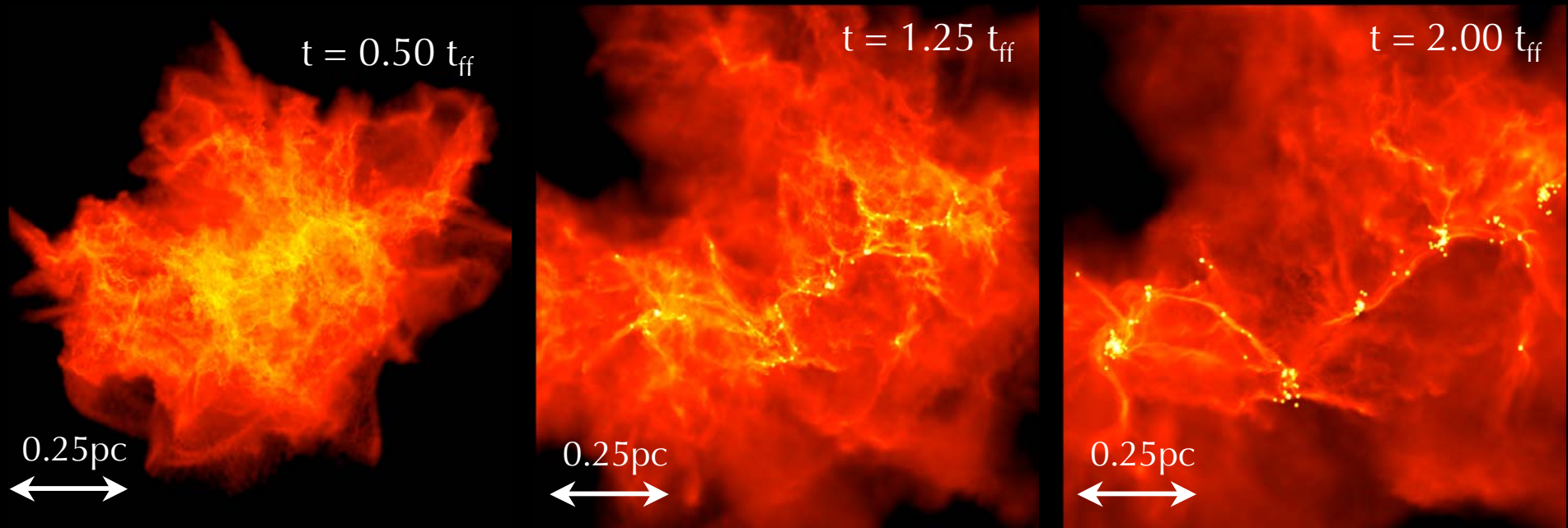
$$t_{\text{frag}} \sim t_{\text{acc}}$$

- Any region characterised by a common density, by which both fragmentation and accretion are dictated, satisfies this requirement.



Unbound clouds

KE = 2 × PE (initially), 1000 solar masses, 0.5pc



No global collapse:

local $t_{\text{ff}} <$ global interaction time-scale

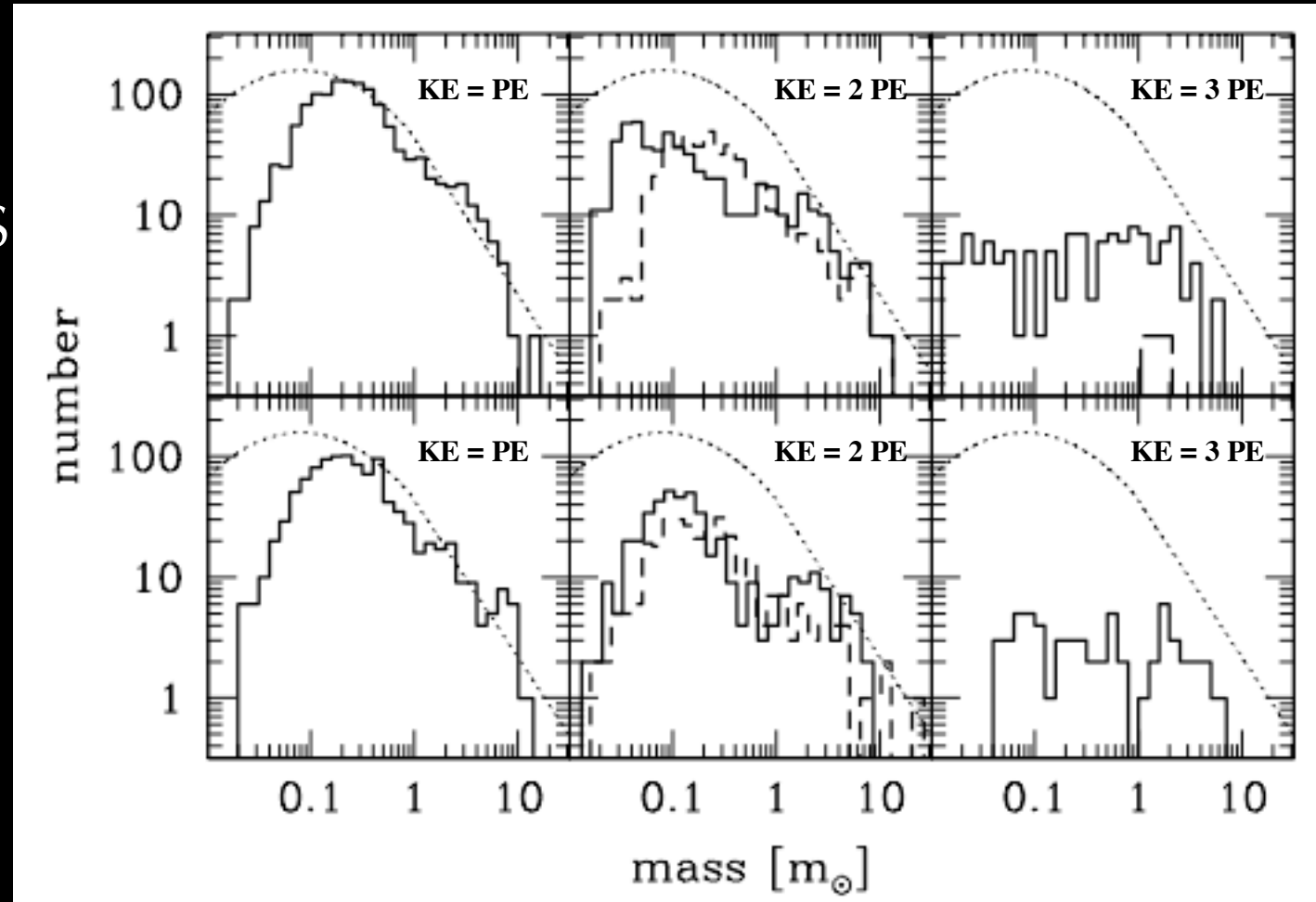
$t_{\text{ff}} \sim 2 \times 10^5$ years

Clark, Bonnell & Klessen (2007)

Mass functions?

Isothermal EOS

Barotropic,
Larson (2005),
Style EOS



RT! (without actually doing RT)

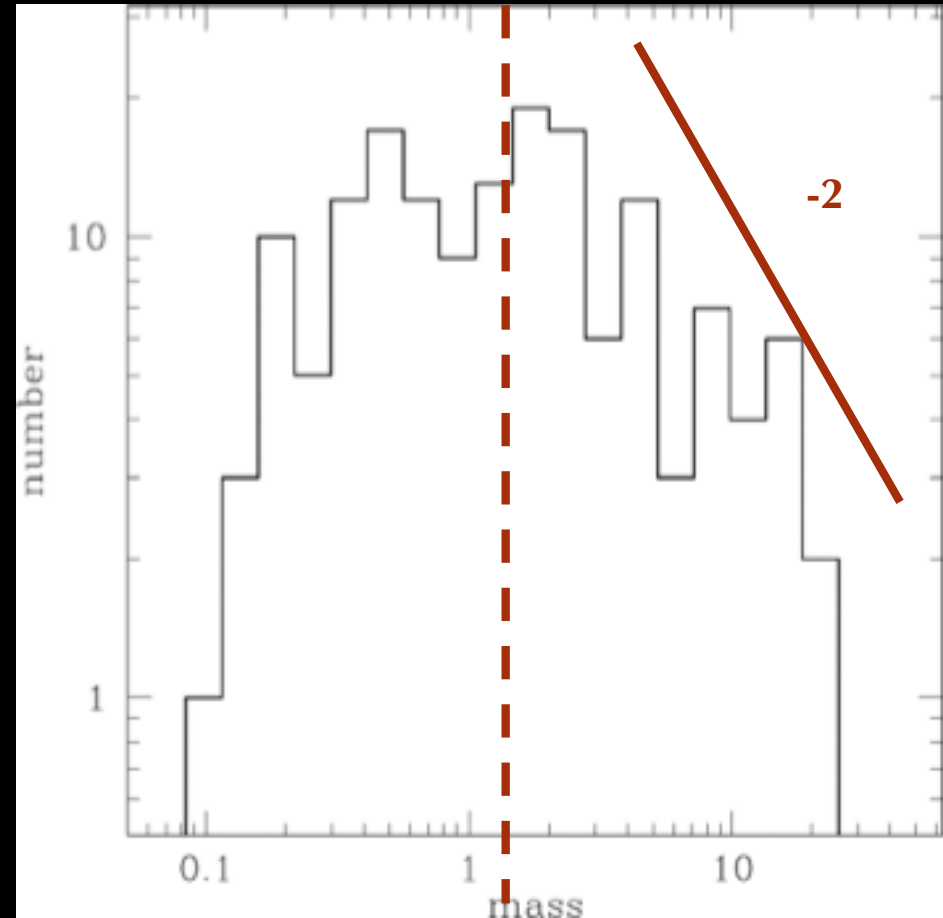
Fit (fudge) to MC models:

$$T = 100\text{K} \left(\frac{M}{10 M_{\odot}} \right)^a \left(\frac{R}{1000 \text{ AU}} \right)^q$$

$$a: 0.33 \text{ } M_{\odot} < 10$$

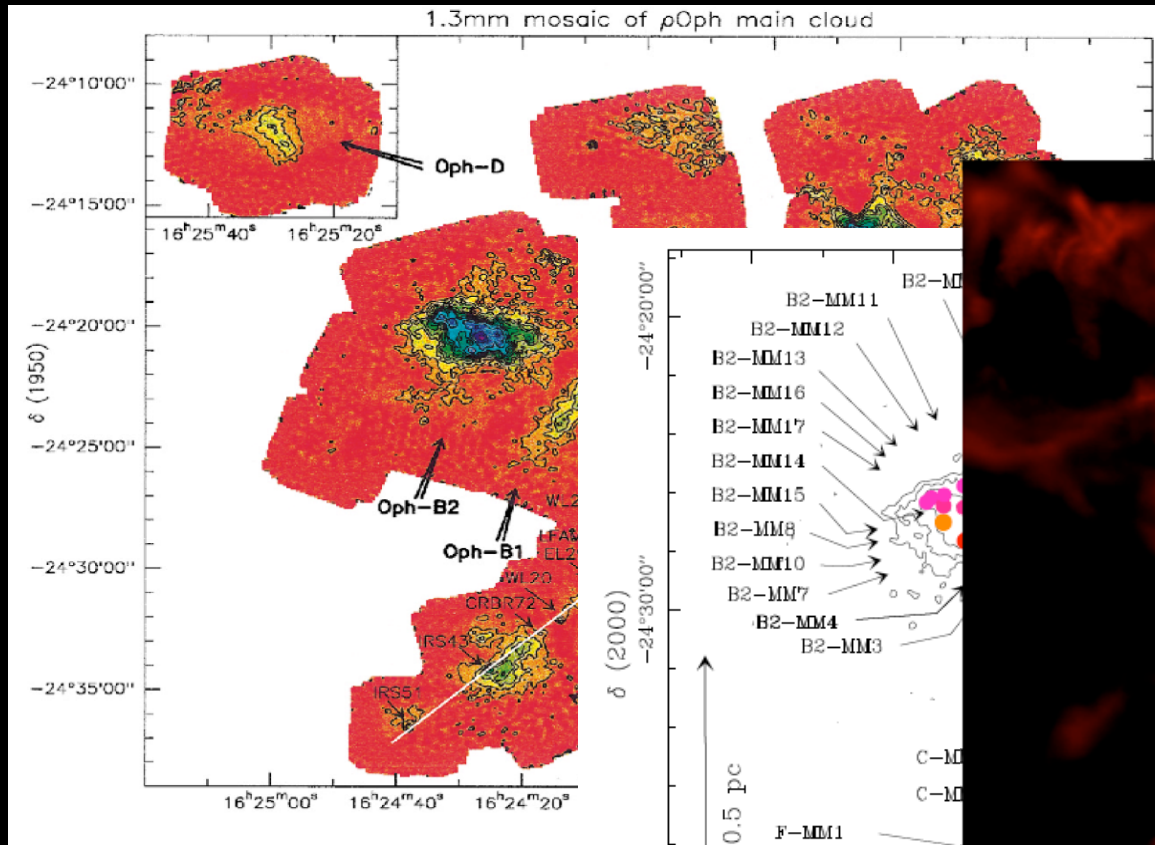
$$a: 1.1 \text{ } M_{\odot} > 10$$

$$q: -0.4 \text{ to } -0.5$$

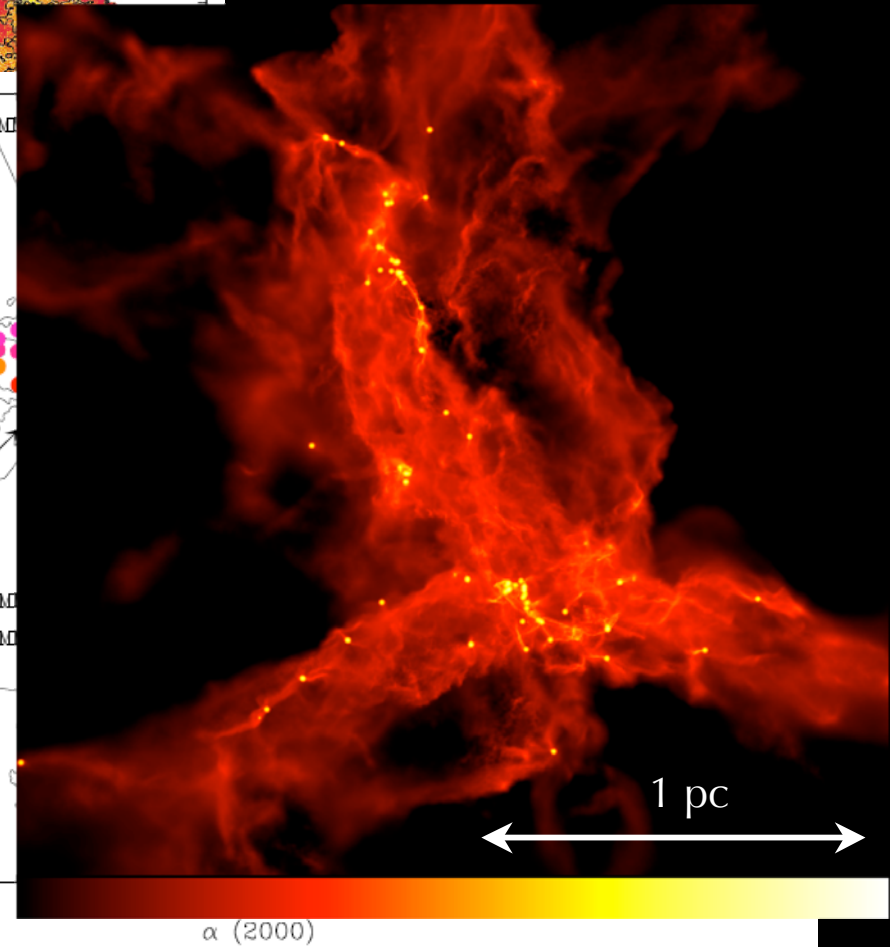


Robitaille et. al. 2006

Observational tests?



Bonnell, Clarke & Bate (2006)



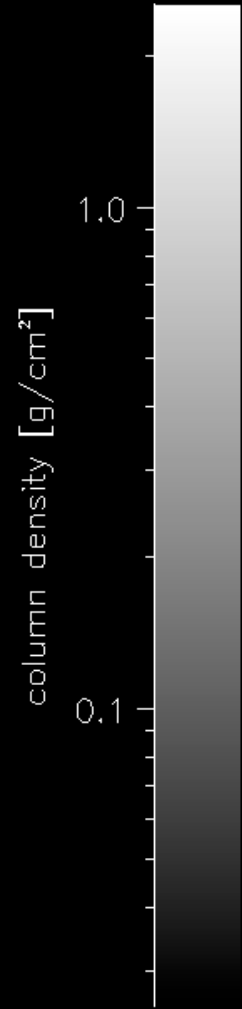
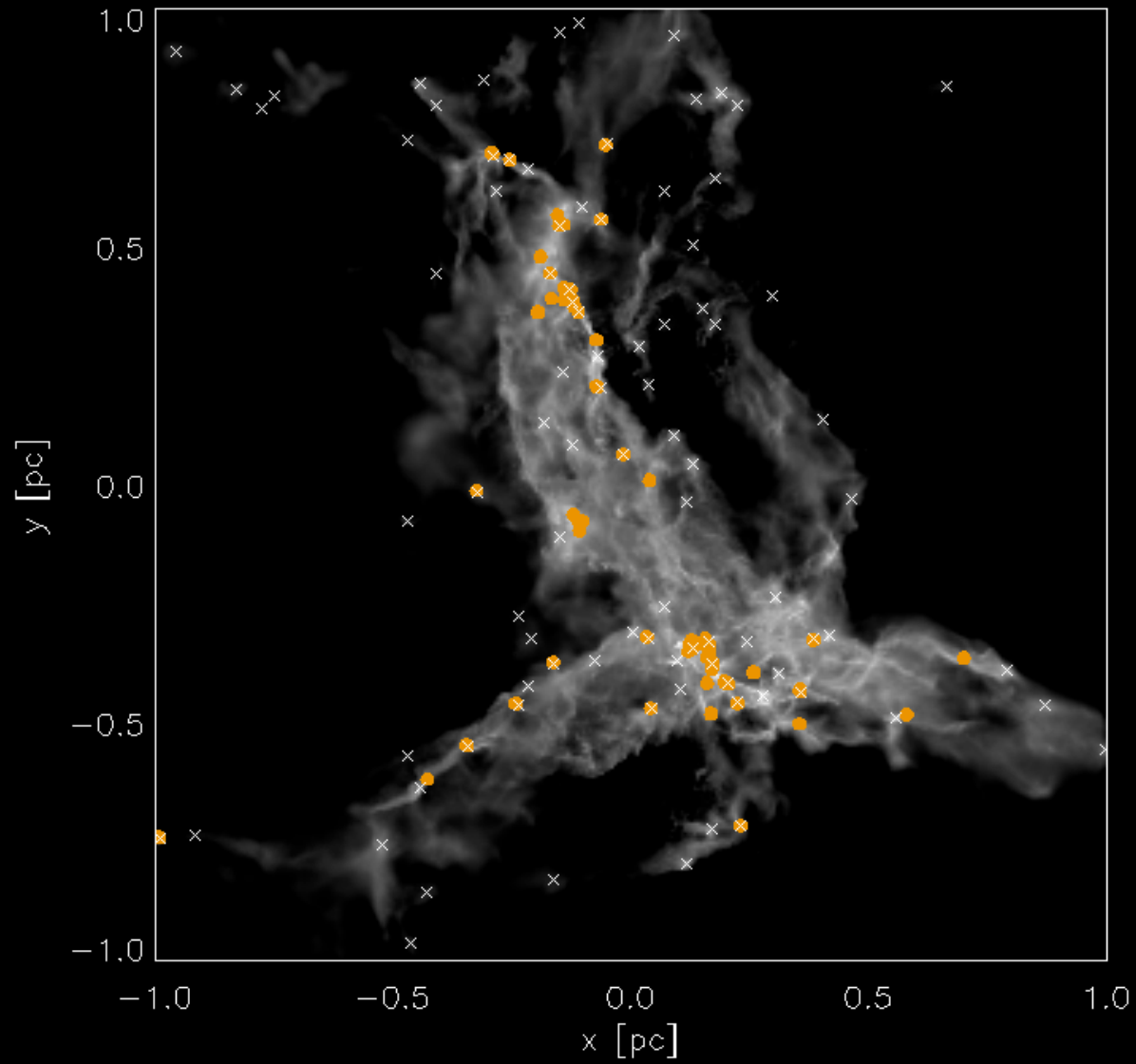
Motte, André & Neri (1998)

André et al. (2007)

Stars



Clumps



Clump mass functions

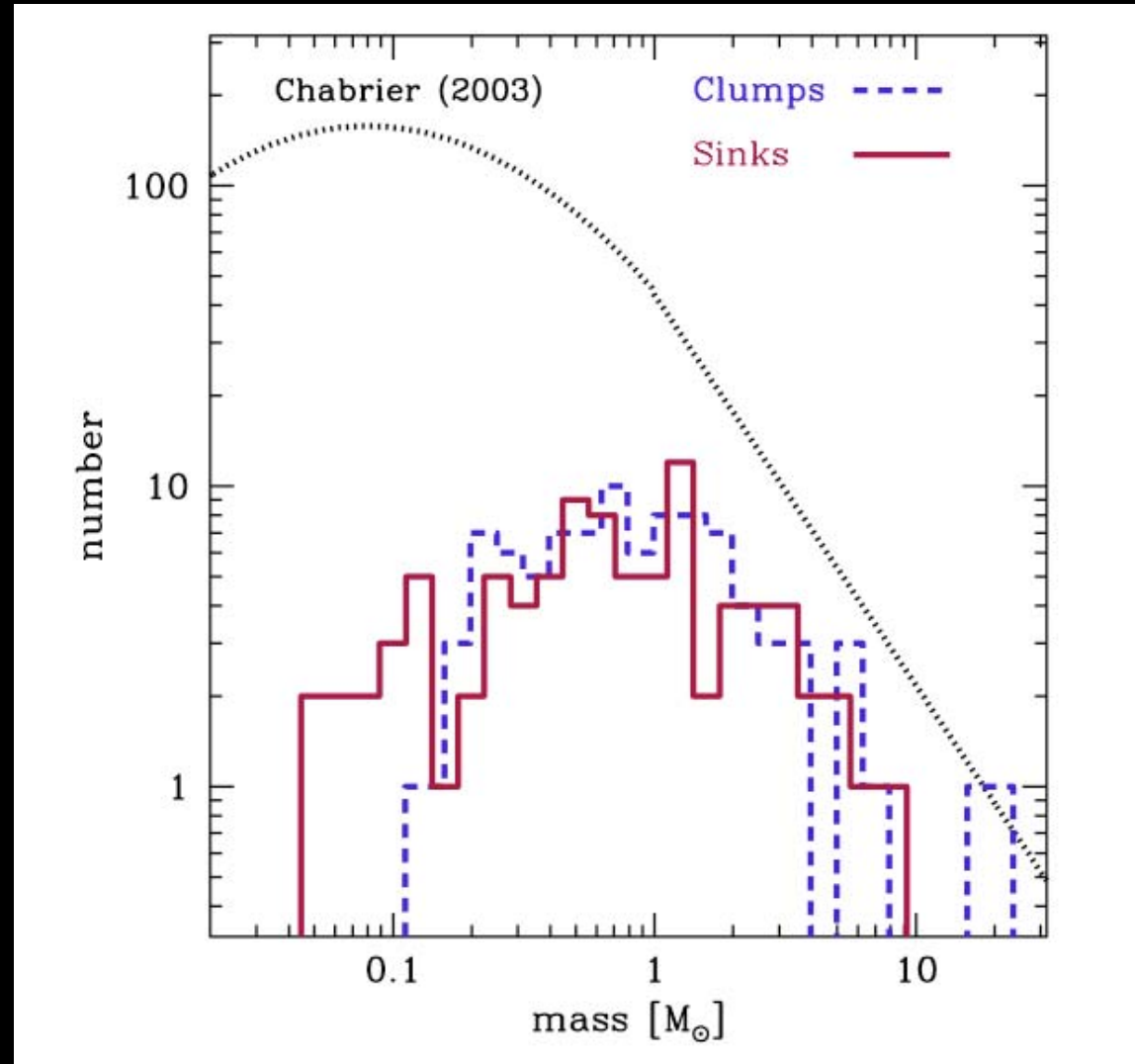
88 sink particles

SPH data mapped to a
2D grid with resolution
 $\sim 1000 \times 1000$ au

Column densities
limited to range
 $0.02 - 2.00 \text{ cm}^{-2}$

Clumps required to
have a density contrast
of a factor 2 in column
density

91 "sink-less" clumps



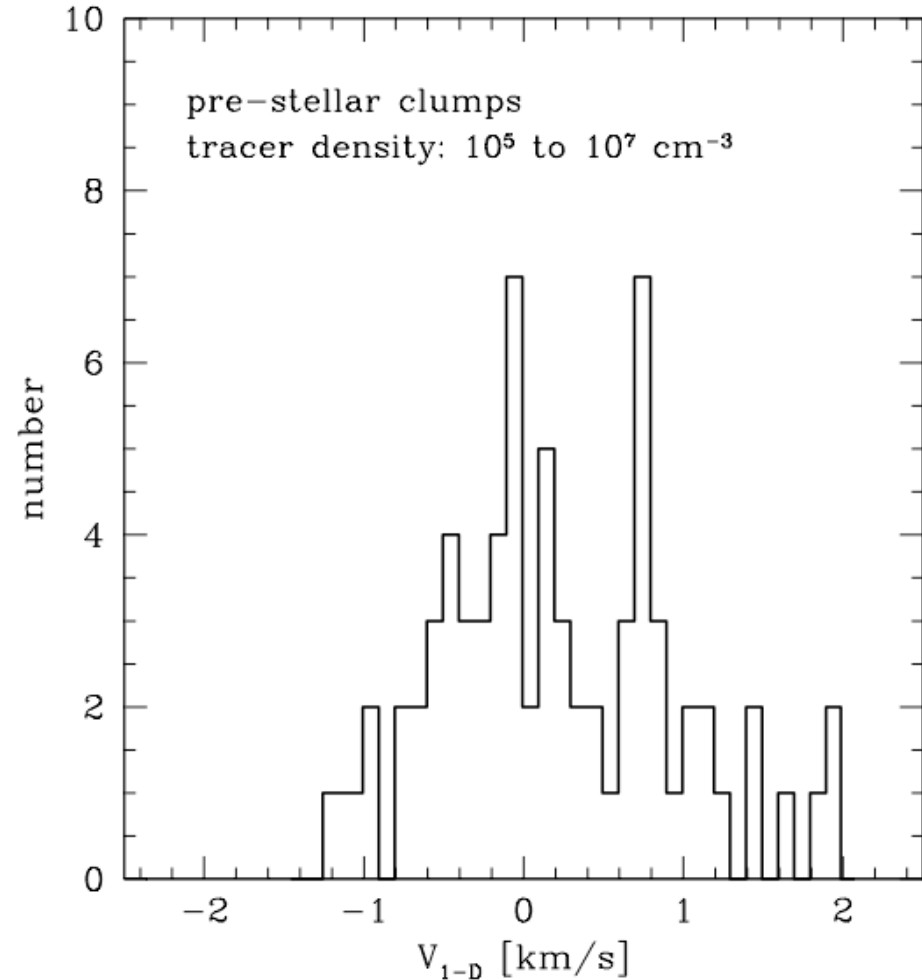
Clump velocity dispersions

Each cluster has its own central velocity

Distribution around this velocity is $\sim 0.25\text{km/s}$, and the mean is only $\sim 0.7\text{ km/s}$ for the whole region.

Typical of turbulent stagnation points (e.g Padoan et al 2001)

Similar velocities to *André et al (2007)*



Can you see competitive accretion?

André et al (2007):

- Used the *clump velocity* to estimate interaction
- From Binney (1987),

$$\frac{t_{coll}}{t_{cross}} = \frac{1}{2} \sqrt{\frac{\pi}{3}} \times \frac{R^2}{N_{cond} r_{cond}^2} \times \frac{1}{1 + \Theta}$$

where,

$$1 + \Theta \equiv 1 + GM_{cond}/(\sigma_{1D}^2 r_{cond})$$

- Using (for L1688):

For the Bonnell et al (2006) cloud:

$$t_{cross} \sim 1.7 \text{ Myr}$$

$$t_{coll}/t_{cross} \sim 13.5$$

$$0.55 \text{ pc}$$

$$\text{km/s} = 57$$

$$6 \text{ km/s}$$

$$R_{cond} \sim 2500 \text{ AU}$$

$$M_{cond} \sim 0.4 M_{\odot}$$

$$t_{cross} \sim 1.78 \text{ Myr}$$

- Get time-scale ratio:

$$t_{coll}/t_{cross} \sim 9$$

Threshold for massive star formation?

- McKee and Tan (2003): $\Sigma_{\text{crit}} \sim 1 \text{ g/cm}^2$

CHARACTERISTIC SURFACE DENSITIES OF REGIONS OF HIGH-MASS STAR FORMATION

Object	M (M_{\odot})	$R_{1/2}$ (pc)	Σ (g cm^{-2})	\bar{P}_{cl}/k (K cm^{-3})	References
Galactic star-forming clumps.....	3800 ^{a,b}	0.5 ^b	1.0	4×10^8	1
Orion Nebula Cluster.....	4600 ^a	0.8	0.24	2×10^7	2
Arches cluster.....	2×10^4	0.4	4	7×10^9	3, 4
Galactic globular clusters.....	2×10^{5a}	3.4	0.8	3×10^8	5, 6
NGC 1569 A1, A2.....	4×10^{5a}	2.2	2.7	3×10^9	7, 8
NGC 5253.....	$(0.6-1.5) \times 10^{6c}$	1.0	20-50	$(2-11) \times 10^{11}$	9

^a Virial mass estimates.

^b The half-mass radius is not well defined for the Plume et al. 1997 clouds, since the mass distribution on larger scales is not known. We therefore evaluate $\Sigma = M/\pi R^2$ using the typical radius and virial mass that they observe.

^c Extrapolation from inferred Lyman continuum luminosity of H II region based on Salpeter IMF with a lower mass limit $m_{\ell} = 1, 0.1 M_{\odot}$.

REFERENCES.—(1) Plume et al. 1997. (2) Hillenbrand & Hartmann 1998. (3) Figer et al. 1999. (4) Kim et al. 2000. (5) Binney & Merrifield 1998. (6) van den Bergh et al. 1991. (7) Gilbert & Graham 2001. (8) de Marchi et al. 1997. (9) Turner et al. 2000.

Conditions required by CA

- In CA calculations, roughly $500 M_{\odot}$ gas accreted before massive star ($8-10 M_{\odot}$) is formed.

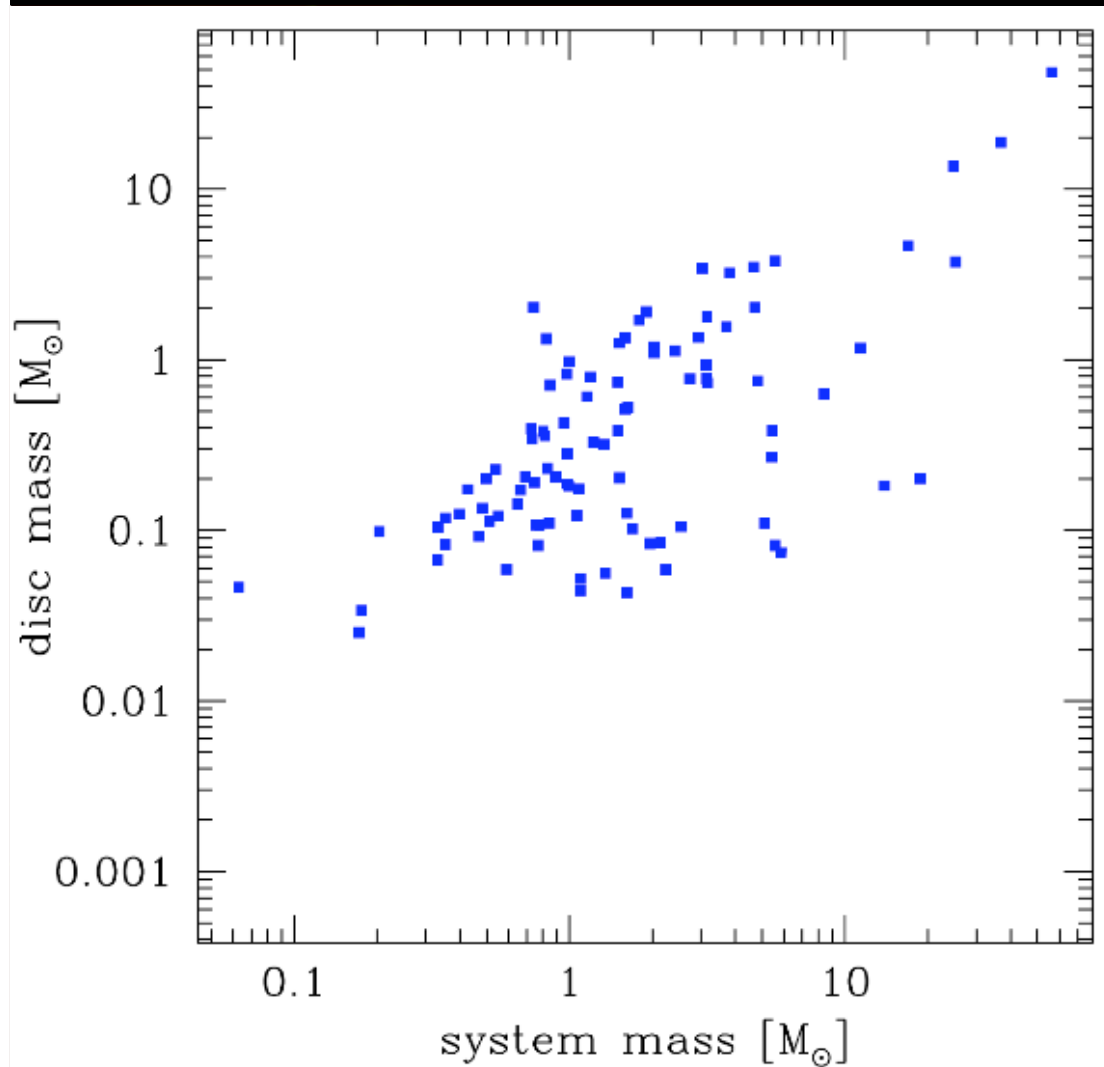
$$f_{\text{sfe}} = \frac{M_{*}}{M_{*} + M_{\text{gas}}} = \frac{M_{*}}{M_{\text{clump}}} \quad M_{\text{clump}} = M_{*} / f_{\text{sfe}}$$

$$\Sigma_{\text{crit}} \sim M_{\text{clump}}^{1/3} \rho_{\text{frag}}^{2/3} \phi_{\text{geom}}^{2/3} \quad \phi_{\text{geom}} = (4\pi/3)$$

Larson (2005) suggests this is set by the transition from line to dust cooling: $\rho_{\text{frag}} 10^{-20} - 10^{-17} \text{ g/cm}^3$

$$\Sigma_{\text{crit}} \sim 1 \text{ g cm}^{-2} \left[\frac{M_{*}}{500 M_{\odot}} \right]^{1/3} \left[\frac{f_{\text{sfe}}}{0.5} \right]^{1/3} \left[\frac{\rho_{\text{frag}}}{2 \times 10^{-19} \text{ g cm}^{-3}} \right]^{2/3} \left[\frac{\phi_{\text{geom}}}{4.2} \right]^{2/3}$$

Discs?



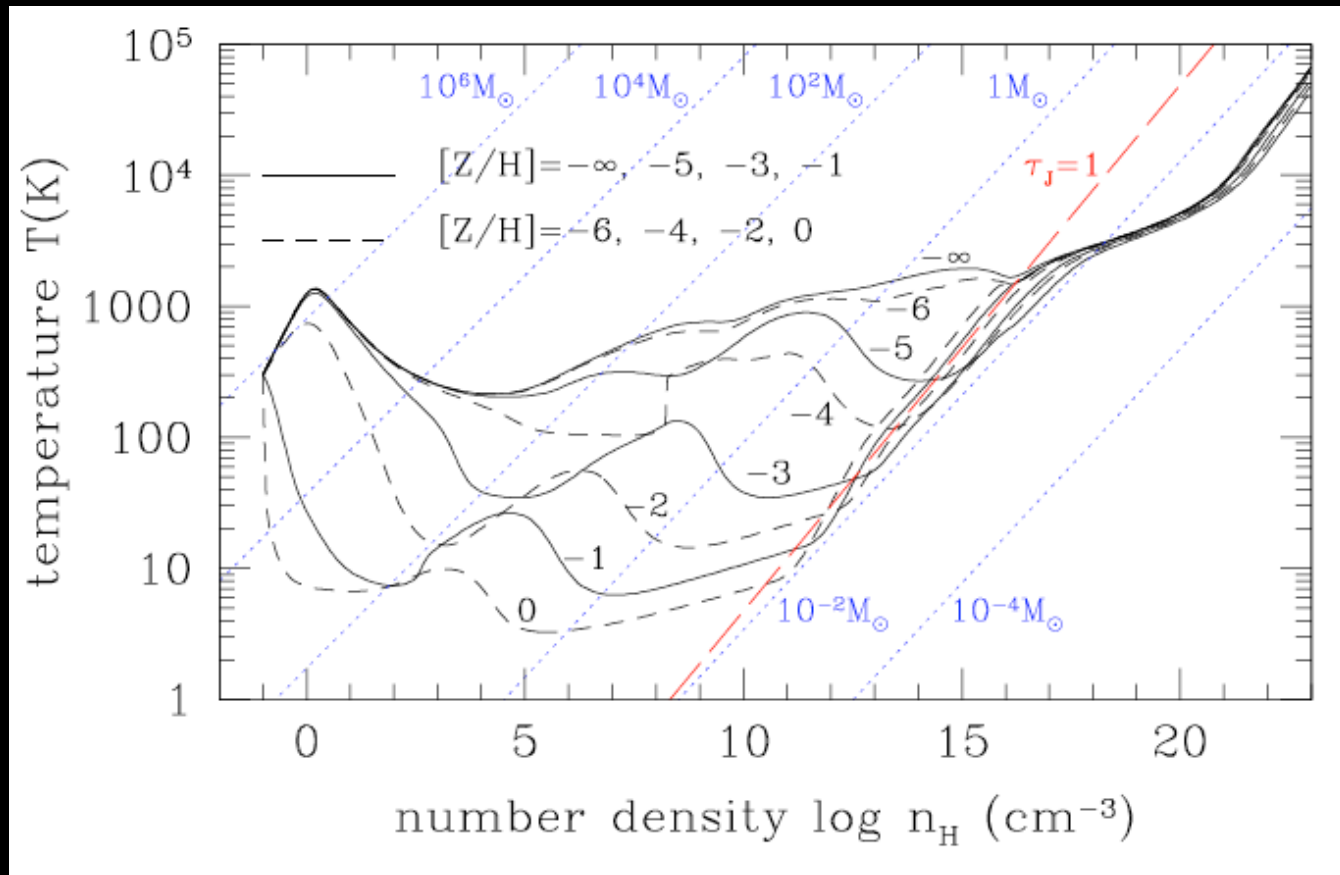
- Relationship between disc mass and protostellar system mass:

$$m_{\text{disc}} \propto m_{\text{sys}}^{1.5 - 2}$$

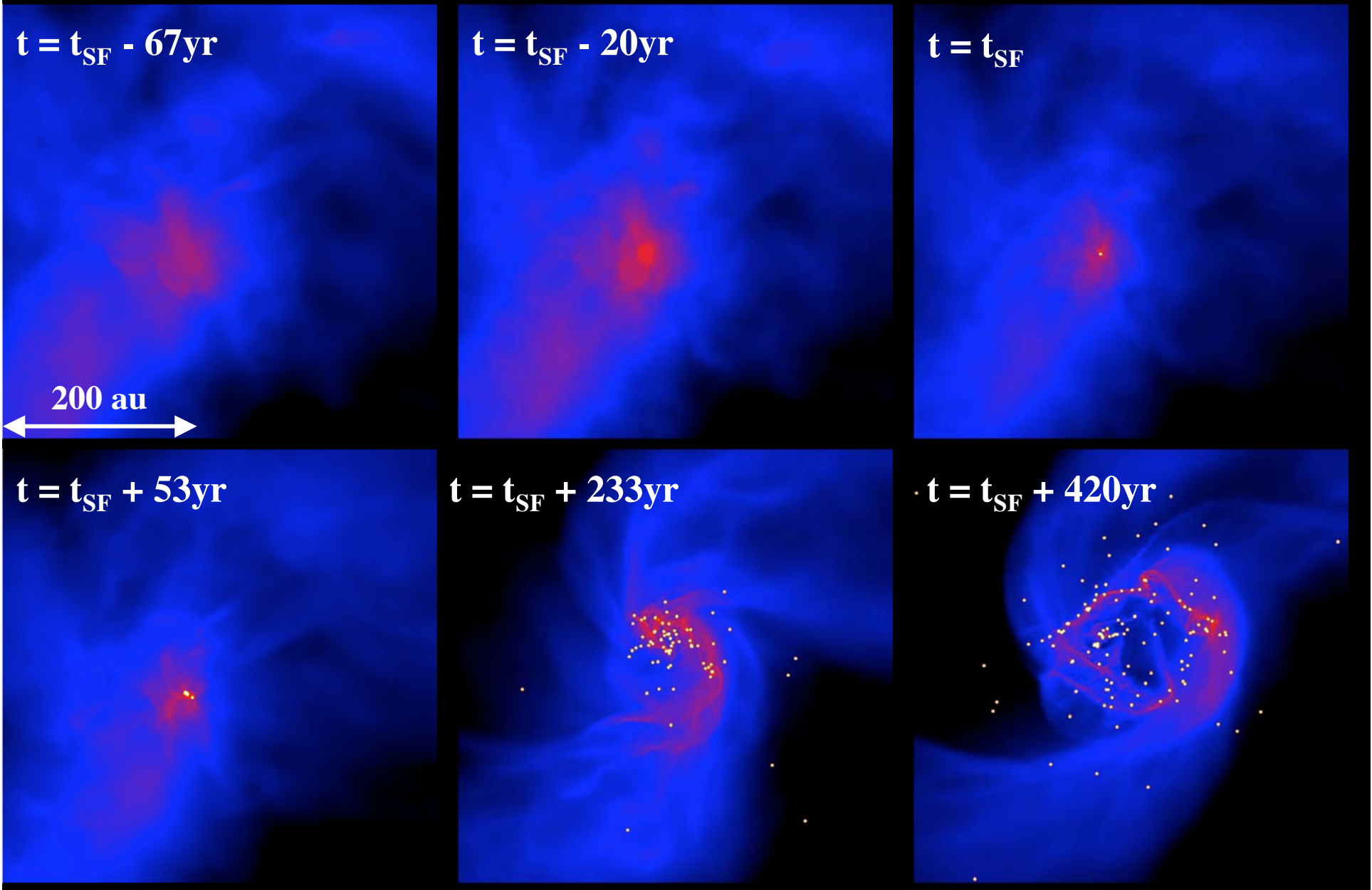
- Note that discs come and go!
- Angular momentum vector can change!

CA at low metallicity?

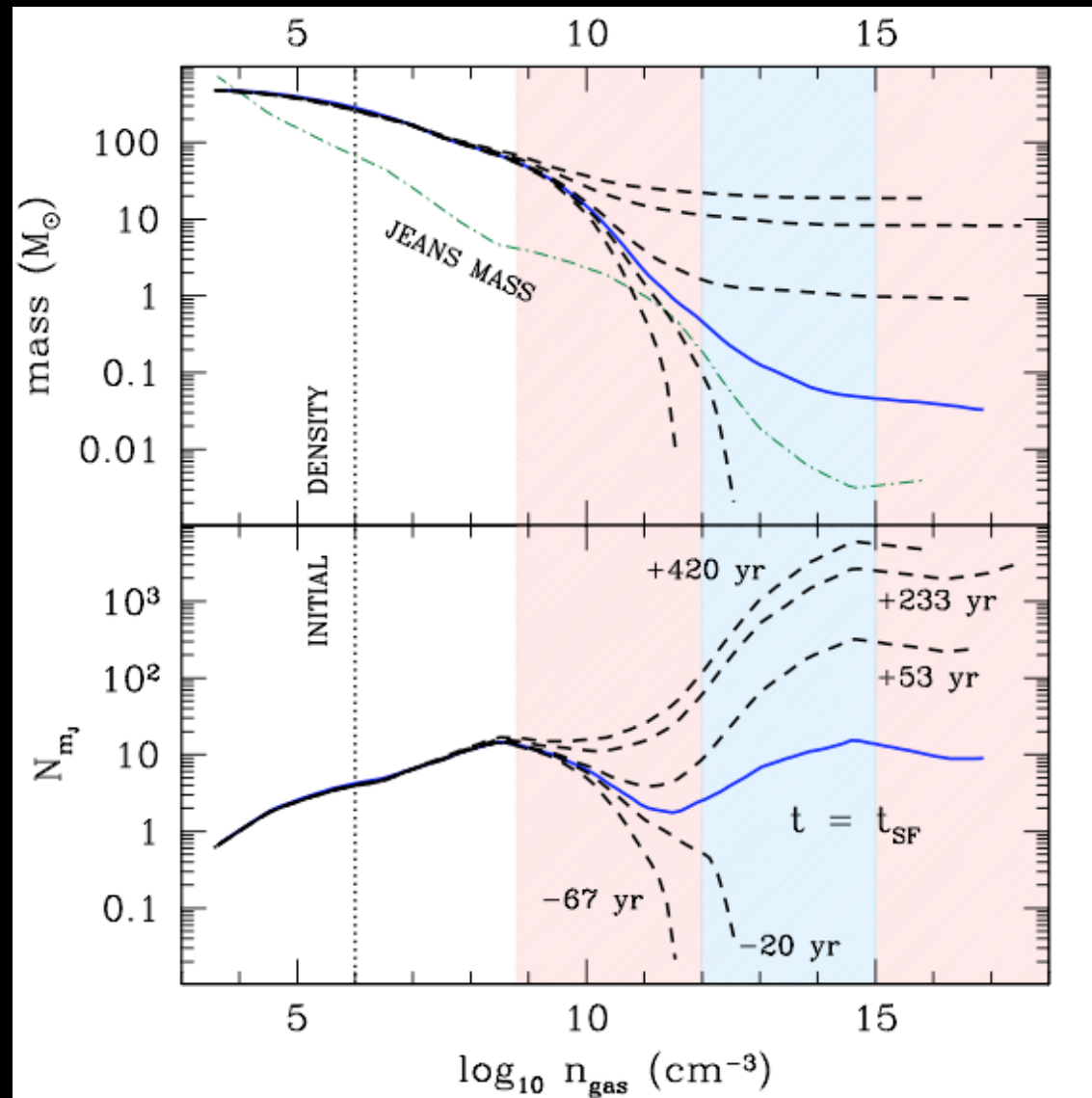
- Omukai et al (2005) suggest that cooling by dust can promote fragmentation, even at very low metallicities:



CA at low metallicity

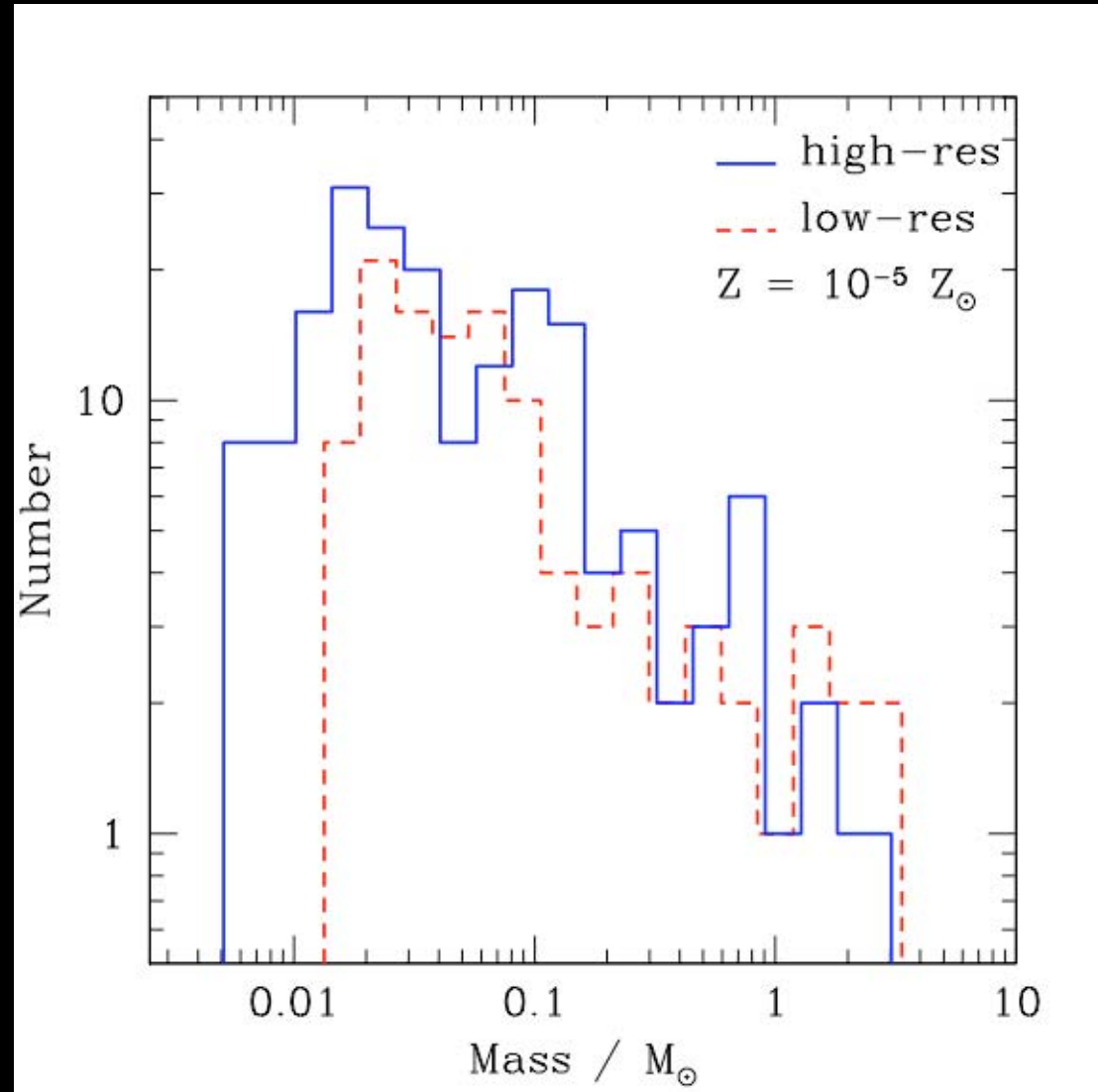


Violent fragmentation...



CA at low metallicity?

- Fragmentation at very low Jeans mass.
- Moves very rapidly into B-H accretion phase.
- Salpeter-type slope extends right down to the fragmentation mass.



Summary

- Competitive accretion requires bound, collapsing regions to produce the 'correct' IMF.
- Difficult to use observed interaction time-scales to estimate the competitive accretion rates: tend to neglect the changing potential which plays a crucial role.
- CA models so far require $\Sigma_{\text{crit}} \sim 1 \text{ g/cm}^2$
- Disc observations may help to determine importance of interactions.