Star Formation at Small Scales

Formation of Circumstellar Disks and Outflows

(or what happens in sub and sub-sub grid models)

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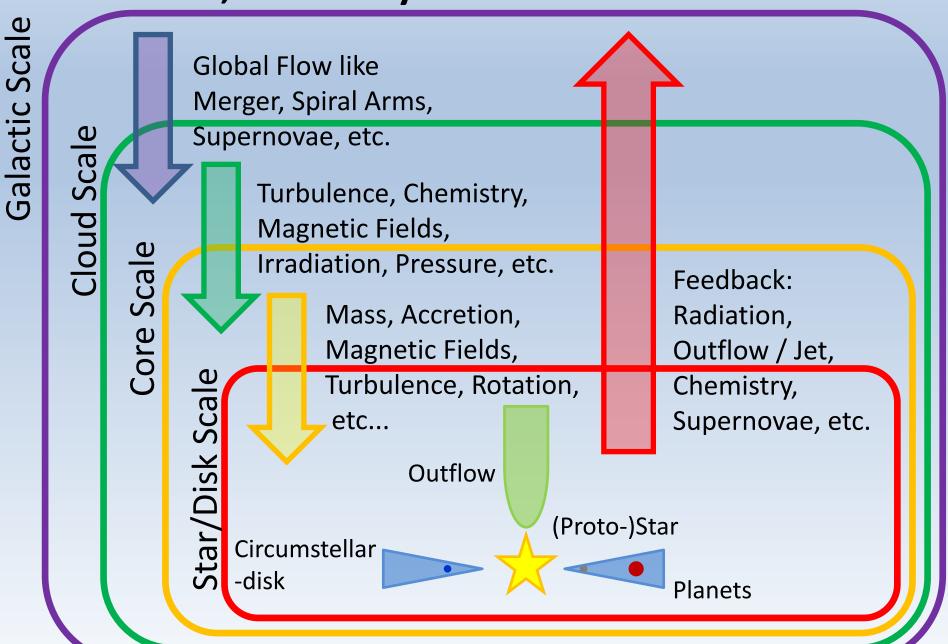
References: Tomida et al., 2013, ApJ, 763, 6 Tomida, Okuzumi & Machida in prep.

Topics

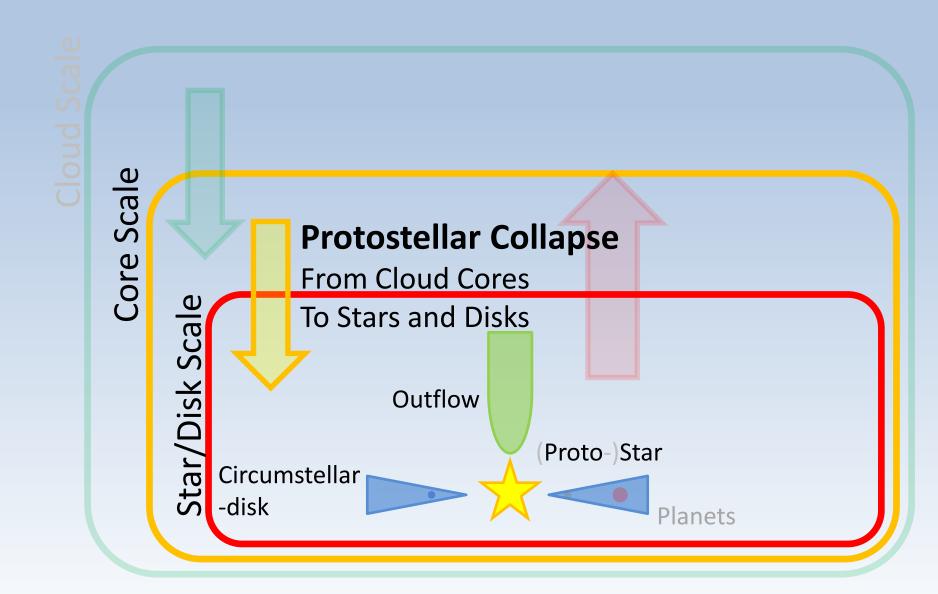
- Introduction
 - Protostellar Collapse
 - Angular Momentum Problem
 - Magnetic Braking Catastrophe
- RMHD simulations of Protostellar Collapse
 - Ideal RMHD
 - Resistive RMHD
 - Ambipolar Diffusion (very preliminary!)
- Observations
 - Young Circumstellar Disks
 - First Core Candidates
- Summary

Introduction

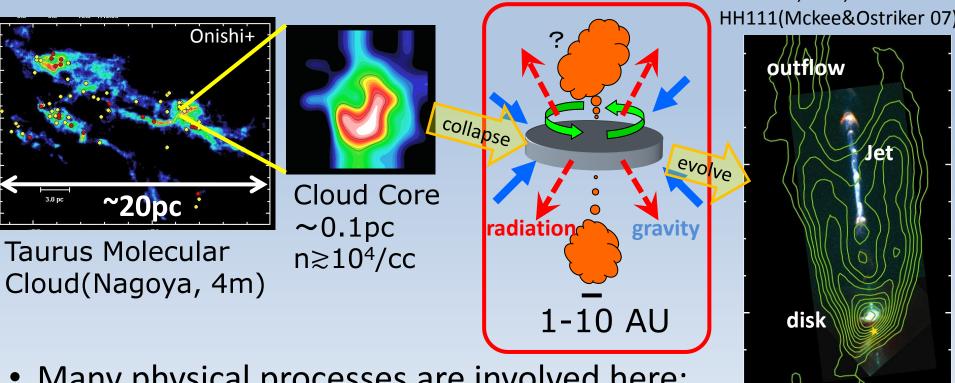
Multi-Scale, Multi-Physics nature of Star Formation



What this talk covers



Protostellar Collapse



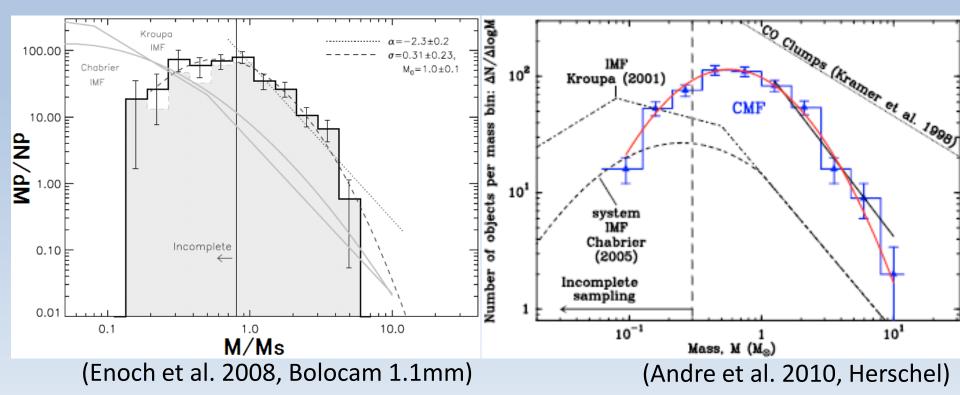
Protostar, Disk, Outflow

- Many physical processes are involved here:
 self-gravity, magnetic fields, radiation transfer, turbulence, chemistry, non-ideal MHD effects, etc...
- Huge dynamic range: 0.1 pc / 1 Rs ~4.5 x **10**⁶
- ⇒Sophisticated numerical simulations are required

Why do we care such a small scale?

- Ultimate goal: the origin of the initial mass function
 - > Relation between CMF and IMF (?)
 - Star Formation Efficiency (at core scale)
 - Binary / Multiple Formation

Core and Stellar Mass Functions



Mass Function of Dense Cores looks like the IMF -- with some shift.

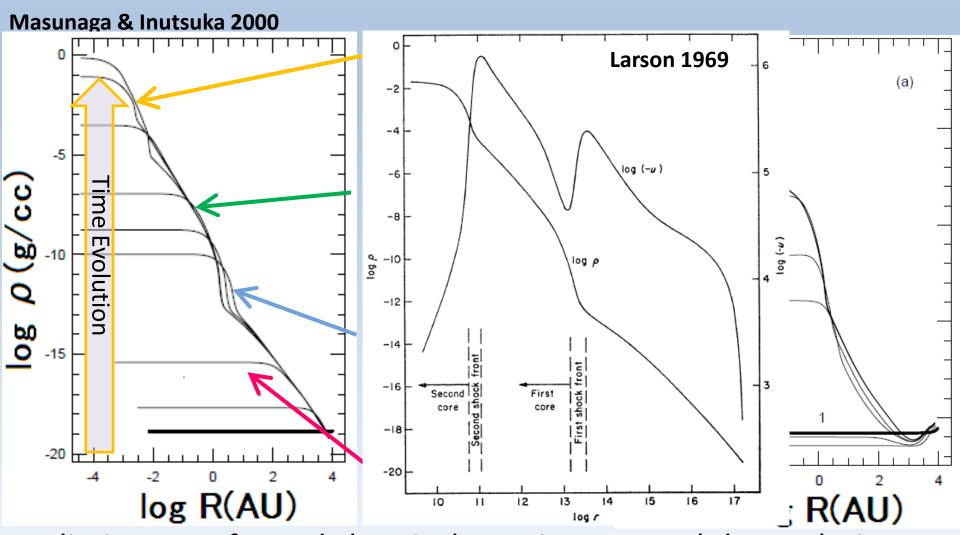
- > Is there such a simple relation? Is **IMF** imprinted in **CMF**? If so ...
- ➤ What is the origin of the CMF?
- ➤ What is the origin of the shift, or **efficiency**? How about **binaries**?

To understand this relation, we have to study protostellar collapse.

Why do we care such a small scale?

- Ultimate goal: the origin of the initial mass function
 - > Relation between CMF and IMF (?)
 - Star Formation Efficiency (at core scale)
 - Binary / Multiple Formation
- Goal of this talk: the origin of circumstellar disks
 - > Angular momentum redistribution
 - ➤ Binary / Multiple / Planet Formation
 - Outflow driving
 - > Feedback: anisotropic radiation, outflows
 - "Flash light effect" (Yorke & Bodenheimer 1999)
 - Disks can be optically thick without dust (Vaidya+ 2009)
- Significant progress in observations with ALMA
- I have a bad allergy to small particles, like pollen & sink particles

Protostellar Collapse: 1D RHD



Radiation transfer and chemical reactions control the evolution.

This scenario is well established based on 1D RHD simulations.

"Problems" in Protostellar Collapse

Angular Momentum Problem

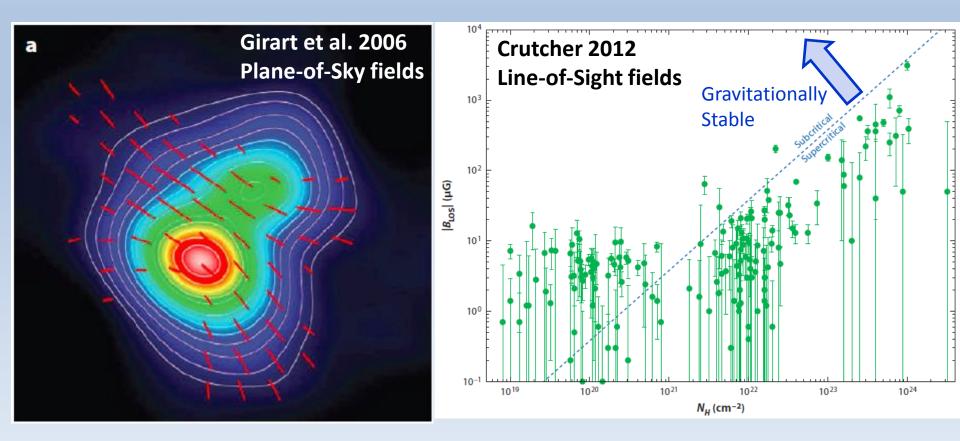
Cloud Cores
$$j_{cl} \approx 5 \times 10^{21} \left(\frac{R}{0.1 \text{pc}}\right)^2 \left(\frac{\Omega}{4 \text{km s}^{-1} \text{pc}^{-1}}\right) \text{cm}^2 \text{s}^{-1} >> j_{\bullet} \approx 6 \times 10^{16} \left(\frac{R_{\bullet}}{2R}\right)^2 \left(\frac{P}{10 \text{day}}\right)^{-1} \text{cm}^2 \text{s}^{-1}$$
 Stars

- →Efficient angular momentum transport during protostellar collapse⇒Gravitational torque, magnetic braking, outflows
- Magnetic Flux Problem
 - Similarly, magnetic flux in cloud cores >> stellar magnetic flux
- → Magnetic fields must dissipate during the collapse
 - ⇒Ohmic dissipation, ambipolar diffusion, (Hall effect), turbulence
- "Magnetic Braking Catastrophe" (Mellon & Li 2008,09, Li+ 2011, etc.)

 Magnetic barking is too efficient; no circumstellar disk is formed

 ⇒Long-term accretion, non-ideal MHD effects, turbulence
- ⇒Realistic **3D simulations with many physical processes**

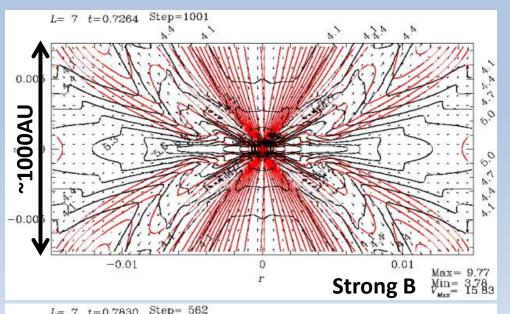
Magnetic Fields

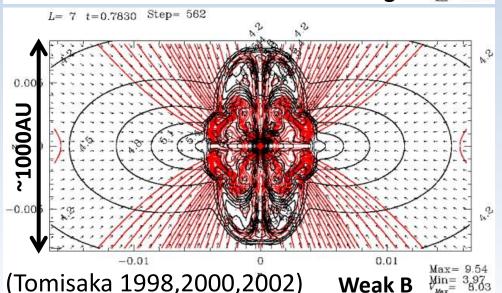


Observations suggest that cloud cores are considerably (supercritical to marginally subcritical) magnetized ($\mu \sim 2-10$). Therefore magnetic fields must have significant effects, actually even in the supercritical regime.

NOTE: these observations are difficult and can have large uncertainties.

Magnetic Braking and Outflows





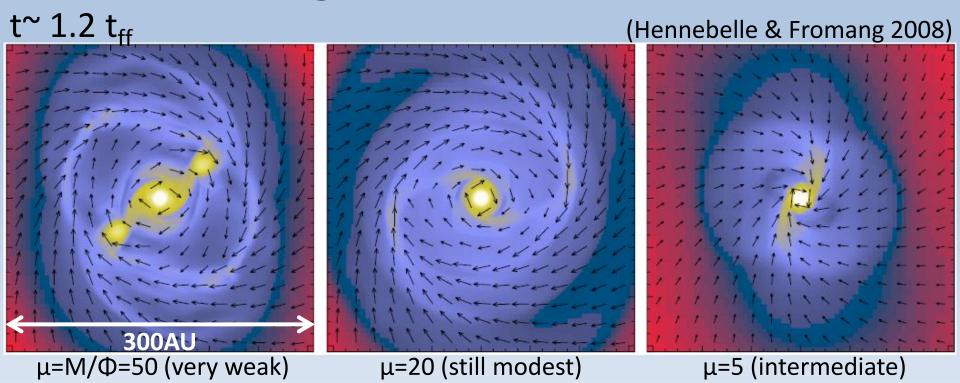
As a result of interaction between magnetic fields and rotation, bipolar outflows are launched from the collapsing cloud. Those outflows and magnetic braking transport angular momentum very efficiently.

Two modes of outflows:

Strong fields result in Magnetocentrifugal mode (Blandford & Payne 1982), while weak fields drive magnetic-pressure mode.

(see also, Mouschovias, & Paleologou 1979, 80, Kudoh et al. 1998, etc.)

Magnetic Braking Catastrophe and/or Fragmentation Crisis



Magnetic fields actually transport angular momentum "too efficiently".

Circumstellar disks are not formed, fragmentation is strongly suppressed.

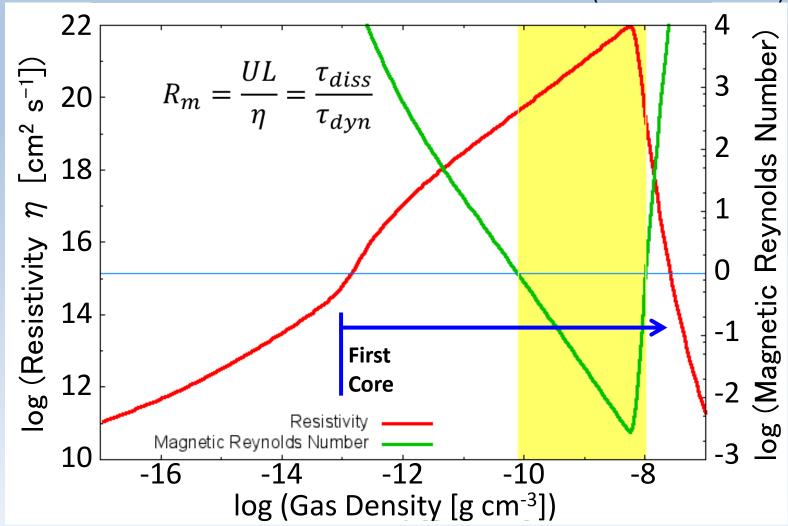
This is a serious problem: Binary rate is known to be high (M: >30% G:

>50%, A: ~80%), and we know lots of circumstellar disks and planets exist.

(see also, Mestel & Spitzer 1956, Mellon & Li 08, 09, Li et al. 11, Hennebelle & Ciardi 09, etc.)

Non-Ideal MHD effects

(Tomida et al. 2013)

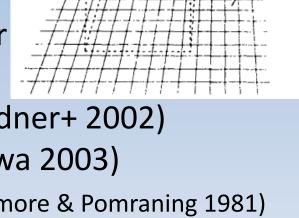


Ohmic resistivity with $\xi=10^{-17}$ s⁻¹ (neglecting Cosmic Ray shielding) Significant flux loss occurs in the first core. (cf. Kunz & Mouschovias 09, 10)₅

RMHD Simulations of Protostellar Collapse

ngr³mhd code

- Huge dynamic range: 3D nested-grids
- MHD → HLLD (Miyoshi & Kusano 2005)
 - (+ Carbuncle care→shock detection + HLLD-)
 - ✓ Fast, robust and as accurate as Roe's solver
 - ✓ Independent from the details of EOS
- div B=0 constraint→Hyperbolic cleaning (Dedner+ 2002)
- Self-gravity→Multigrid (Matsumoto & Hanawa 2003)
- Radiation→Gray Flux Limited Diffusion (Levermore & Pomraning 1981) +Implicit (BiCGStab + ILU decomposition (0) preconditioner)
- EOS including chemical reactions (H₂, H, H⁺, He, He⁺, He²⁺ and e⁻)
- Ohmic dissipation→Super Time Stepping (Alexiades+ 1996)
- (preliminary!) Ambipolar Diffusion with STS
- OD & AD tables are derived using a chemical network with dusts



Ziegler &

Yorke

1997

Basic Equations (w/o div B cleaning)

$$\begin{split} \frac{\partial \rho \mathbf{v}}{\partial t} + \nabla \cdot \left[\rho \mathbf{v} \otimes \mathbf{v} + \left(p + \frac{1}{2} |\mathbf{B}|^2 \right) \mathbb{I} - \mathbf{B} \otimes \mathbf{B} \right] &= -\rho \nabla \Phi + \frac{\sigma_R}{c} \mathbf{F}_r, \text{ Eq. of motion} \\ \frac{\partial \mathbf{B}}{\partial t} - \nabla \times \left(\mathbf{v} \times \mathbf{B} - \underline{\eta} \nabla \times \mathbf{B} \right) &= 0, \text{ Induction eq.} \\ \nabla \cdot \mathbf{B} &= 0, \text{ div B=0} \\ \frac{\partial e}{\partial t} + \nabla \cdot \left[\left(e + p + \frac{1}{2} |\mathbf{B}|^2 \right) \mathbf{v} - \mathbf{B} (\mathbf{v} \cdot \mathbf{B}) - \underline{\eta} \mathbf{B} \times (\nabla \times \mathbf{B}) \right] &= \mathbf{Gas \; Energy \; Eq.} \\ -\rho \mathbf{v} \cdot \nabla \Phi - c \sigma_P (a T^4 - E_r) + \frac{\sigma_R}{c} \mathbf{F}_r \cdot \mathbf{u}, \\ \nabla^2 \Phi &= 4 \pi G \rho, \end{split}$$
 Poisson's Eq.

 $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0,$ Mass Conservation

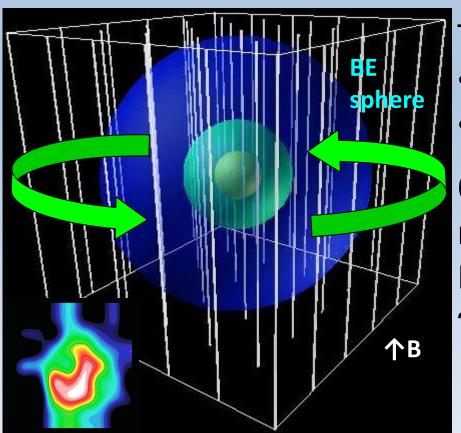
 $\nabla \cdot \mathbf{B} = 0$, div B=0

 $abla^2\Phi=4\pi G
ho$, Poisson's Eq.

 $\frac{\partial E_r}{\partial t} + \nabla \cdot [\mathbf{u}E_r] + \nabla \cdot \mathbf{F}_r + \mathbb{P}_r : \nabla \mathbf{u} = c\sigma_P(a_rT_g^4 - E_r),$ Radiation Transfer

 $\mathbf{F}_r = \frac{c\lambda}{\sigma_R} \nabla E_r, \quad \lambda(R) = \frac{2+R}{6+2R+R^2}, \quad R = \frac{|\nabla E_r|}{\sigma_R E}, \text{ + Eq. of state}$

Simulation Setup



Two rotating models:

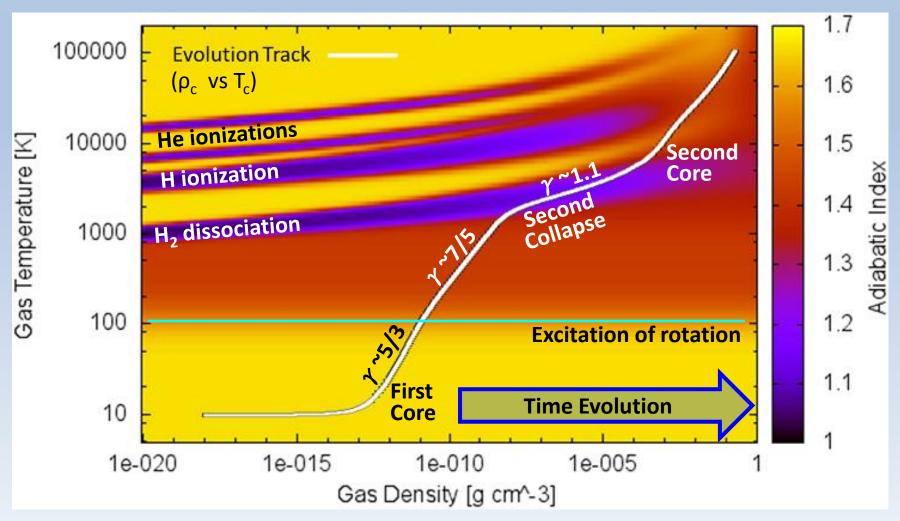
- Ideal MHD model
- Resistive MHD model

64³ x 23 levels, 16 cells / λ_{Jeans} min(Δx) ~ 6.6 x 10⁻⁵AU ~ 0.014Rs End of simulations. Tc ~ 10⁵ K, ~1 yr after 2nd core formation

Huge dynamic range > 10⁸!

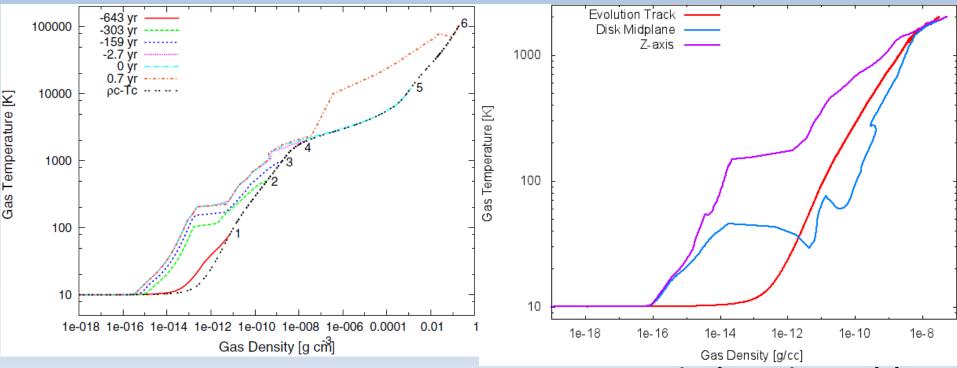
- 1Ms unstabilized BE sphere (ρ_c =1.2 x 10⁻¹⁸ g/cc, T=10K, **R=8800AU**)
- Bz=20 μ G (μ ~3.8), Ω =0.046/t_{ff} ~2.4 x 10⁻¹⁴ s⁻¹, aligned rotator
- 10% m=2 density perturbation
- Opacity: Semenov+ 2003 (dust), Ferguson+ 2005, Seaton+ 1994 (OP)

Thermal Evolution



The central gas element evolves following EOS in $\rho > 10^{-12}$ g /cc. The evolution is consistent with MI2000, except for details of EOS.

Importance of Radiation (M)HD



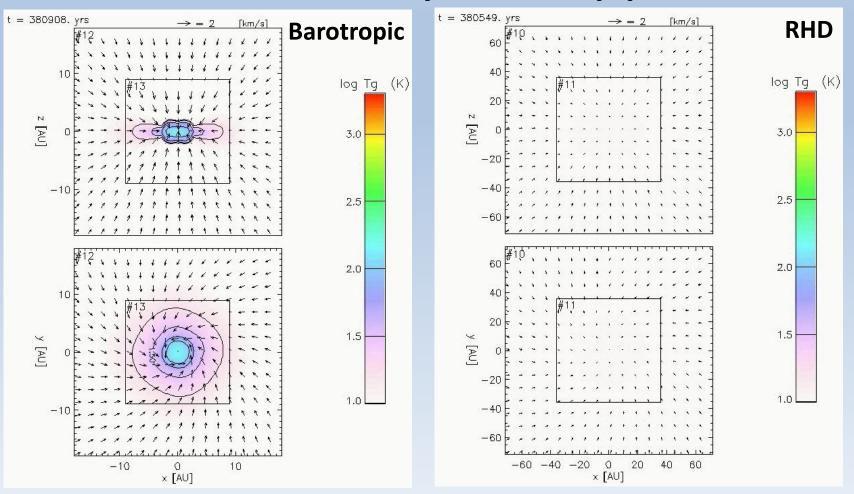
Spherical Model

Non-magnetized Rotating Model

Barotropic approximation can under/overestimate the temperature.

- Non/weakly rotating cases: the temperature tend to be underestimated because shock and radiation heating are neglected.
- More complicated in rotationally-supported disks
- \sim x5 difference in T \Rightarrow x10 difference in Jeans Mass, and stability

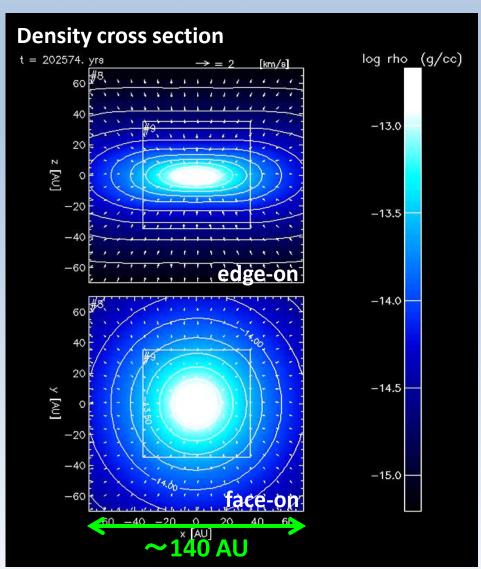
RHD Sims with Gray FLD Approx.

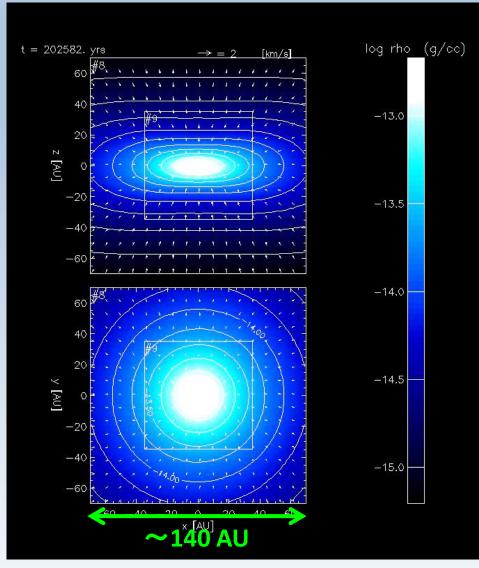


Thicker disks, spiral arms, larger first core mean that RHD is more stable. This also has a significant impact on prediction for observations.

(see, Offner et al. 09, Commercon et al. 10, 11, Tomida et al. 10a,b, Bate 10, 11, etc.)

Outflows

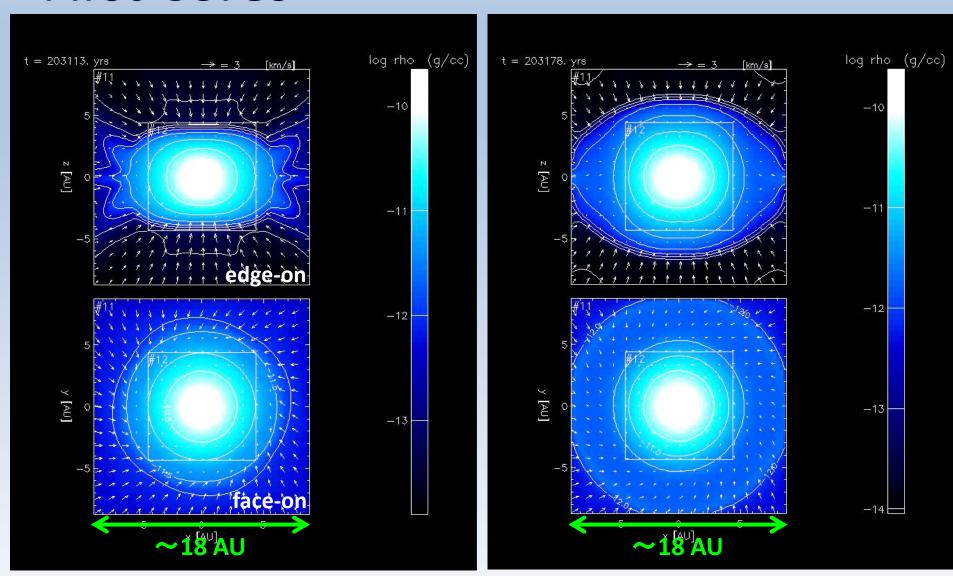




Ideal MHD

Resistive MHD

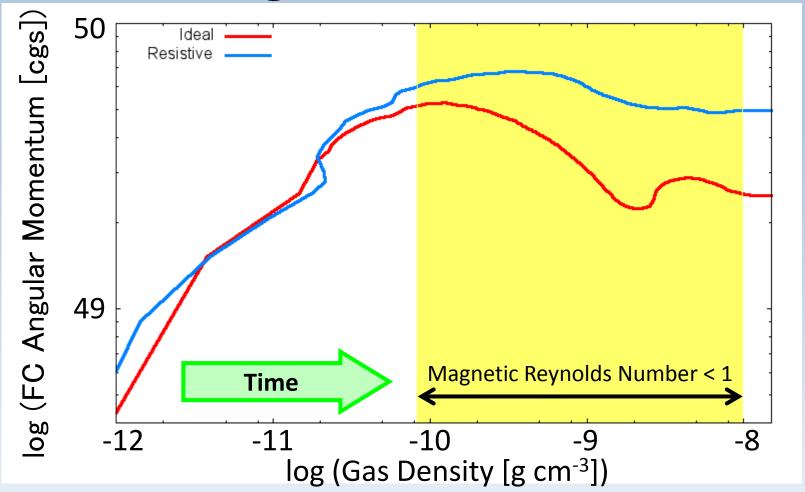
First Cores



Ideal MHD Resistive MHD

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First Core Angular Momentum

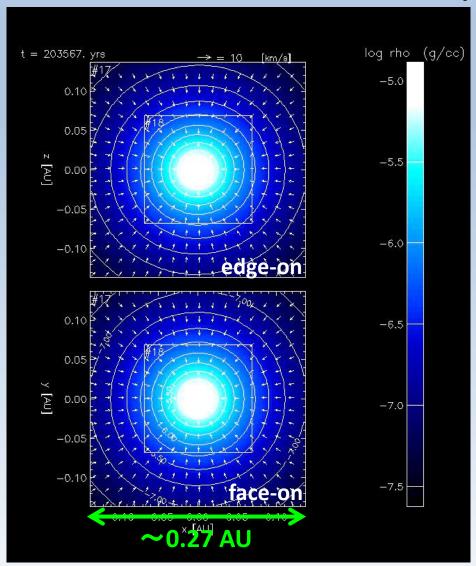


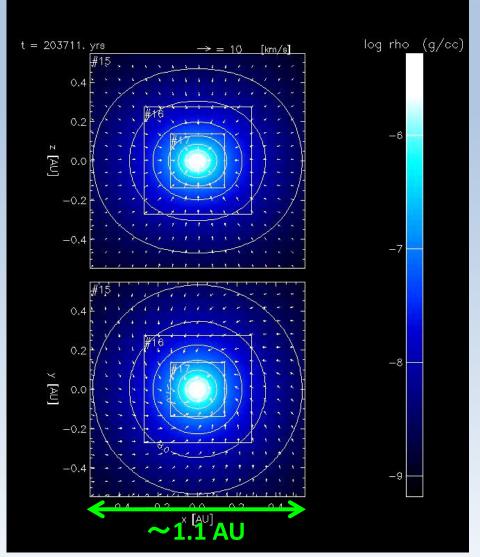
Angular momentum transport is suppressed when resistivity works.

Twice larger angular momentum in FC (more significant in PC)

Longer lifetime due to rotational support:Ideal 800yrs→Resistive 950yrs₅

Protostellar Cores, Disk and Jet





Ideal MHD

Resistive MHD

Protostellar Cores

Radii, Masses, Angular momenta⇒

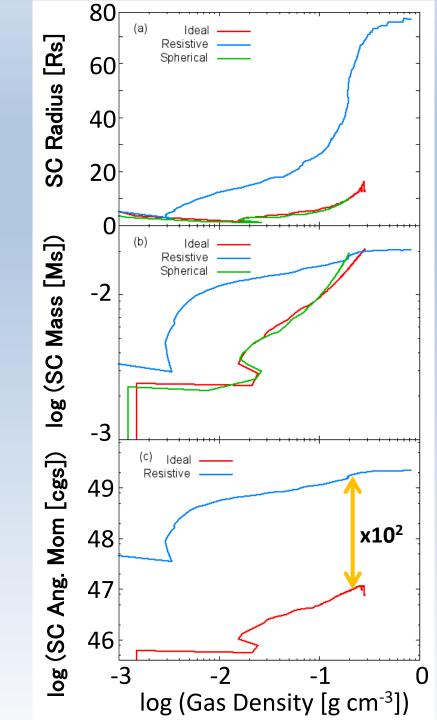
PCs acquire ~0.02 Ms in ~ 1yr

Ideal MHD model = virtually spherical
←very low angular momentum
Circumstellar disk is not formed
"Magnetic Braking Catastrophe"

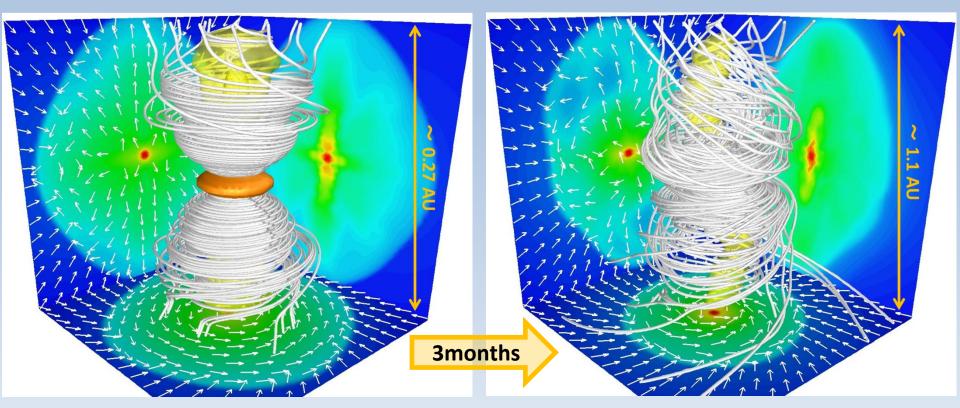
Resistive MHD: large ang. momentum

→rotationally supported disk is formed
Rdisk ~ 0.3 AU at the end of simulation
It will continuously grow via accretion

⇒NO Magnetic Braking Catastrophe



Fast outflow from protostellar core



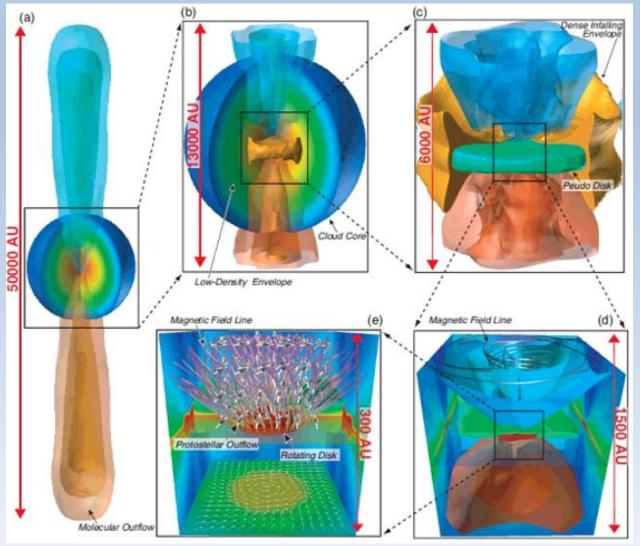
Toroidal fields are rapidly amplified by rotation in resistive case.

→Fast outflow(≥15km/s) is driven by magnetic pressure

Consistent w/ previous MHD sims (Machida et al. 08 etc.)

The magnetic tower is disturbed by the kink instability.

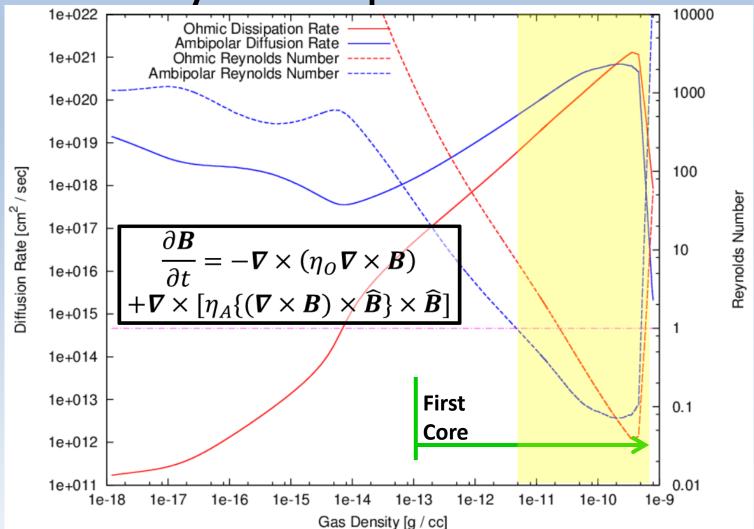
Fate of the disk and outflows



Machida & Hosokawa 13 (they are not allergic.)

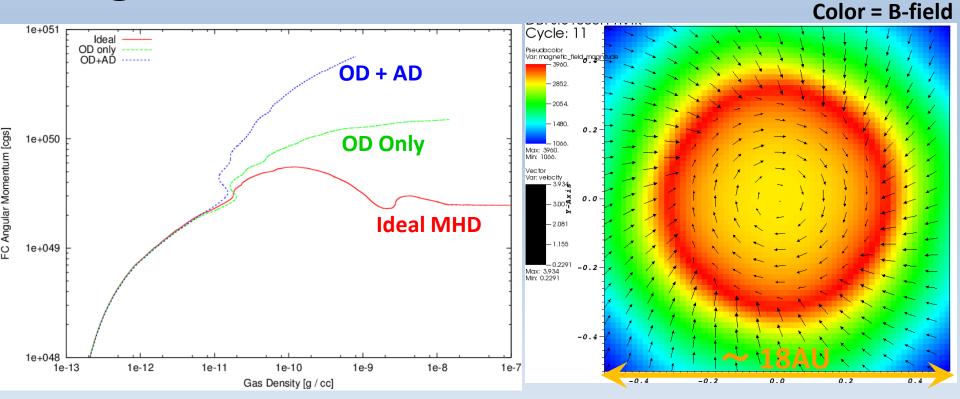
Long-term (till class-I phase) MHD simulation using a sink particle. Outflows and disks grow continuously, $R_{disk} \sim 100 \text{ AU}$

Preliminary: Ambipolar Diffusion



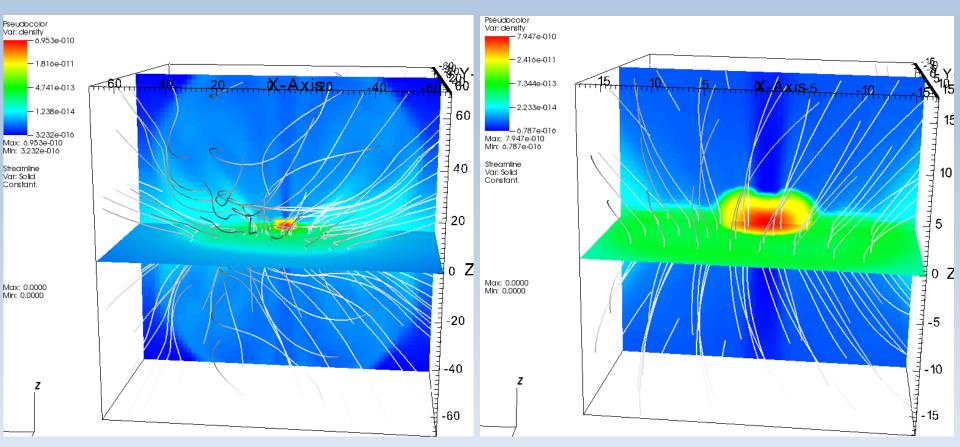
Ambipolar diffusion rate is very high and works in a relatively low density region ⇒ more magnetic flux loss will occur.

Angular Momentum in FC



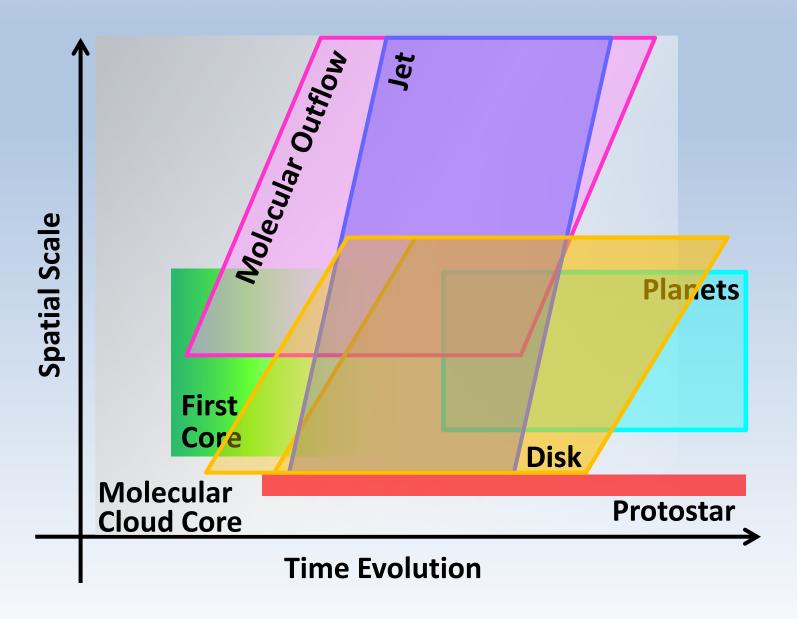
(in the middle of the first core phase, $Tc \sim 800K$)
Large angular momentum remains in the $FC \Rightarrow$ earlier disk formation.
The rotating disk is resistive: magnetic flux is removed from the dense central region and piles up outside the disk.

Outflows and First Core with AD



The bipolar outflows are very similar to the previous cases. First core becomes disk-like because of rotational support. However, the disk radius is still small, about 5 AU.

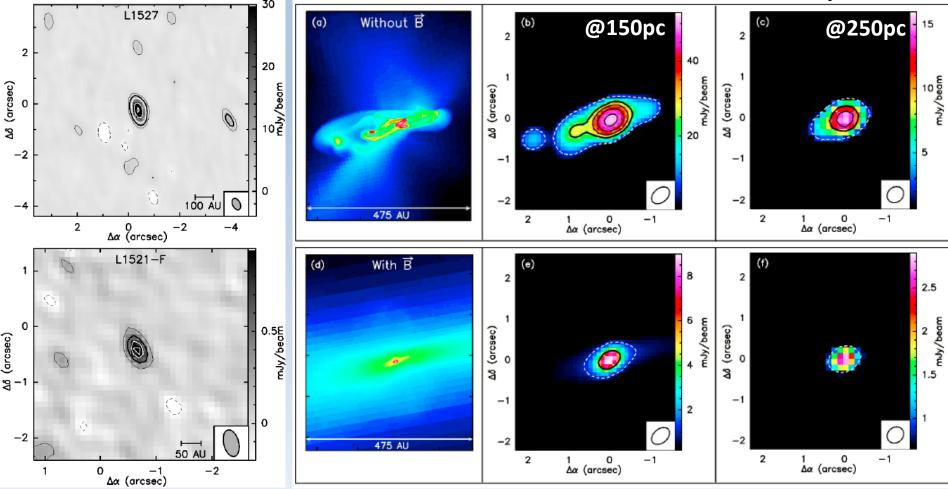
To Summarize: A Schematic Picture



Implications from / for Observations

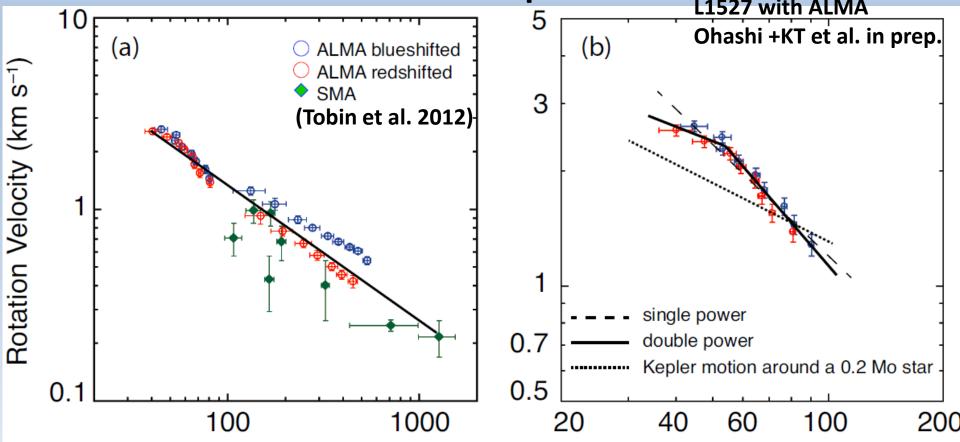
Observations of Young Disks

Maury et al. 2010



1.3mm Dust continuum observations of Class-0 sources with PdBI. The observed disks are small and more consistent with the MHD models.

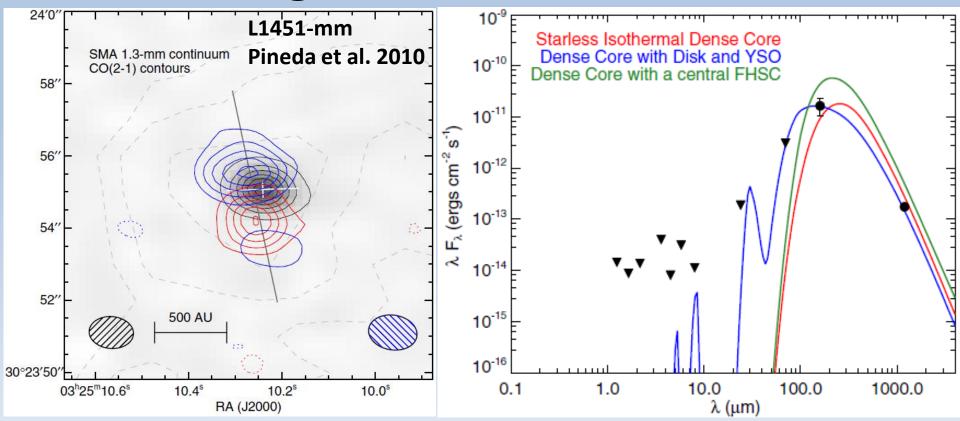
A well-studied example: L1527 IRS



Tobin+ 2012 (SMA & CARMA): R~120 AU disk around 0.2 Ms protostar Ohashi+ in prep. (ALMA Cycle-0): R < 60AU disk around 0.3 Ms protostar ⇒Disks can be formed early, but should be small in the early phase

Rotation Radius (AU)

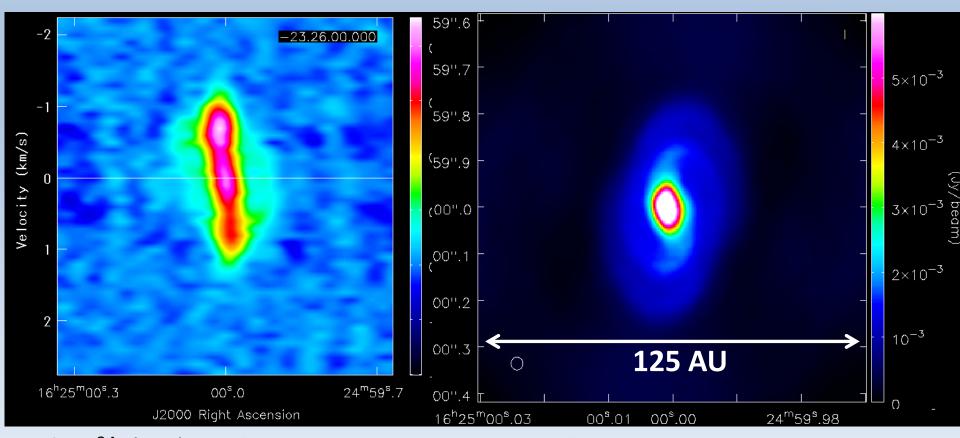
Even Younger: First Core Candidates



Recent first core candidates: L1451-mm, Barnard 1-bN, Per-Bolo 58 etc.

- Faint compact molecular cores without stellar NIR emission
- Associated with compact, slow outflows without fast jet
- However: it must be rare: ~1 FC in 100-1000 molecular cloud cores
- Predicted in Larson 1969 but not confirmed observationally yet

Synthetic Observations



Left: C³⁴S(5-4), Cycle-0, extended, 0.5arcsec, 4h, P-V diagram Right: 345GHz continuum, Full ALMA, 0.02arcsec, 4h (Tomida, PhD Thesis)

Motivated by ALMA, theoretical predictions based on the results of R(M)HD simulations are actively performed. (especially for first cores.)

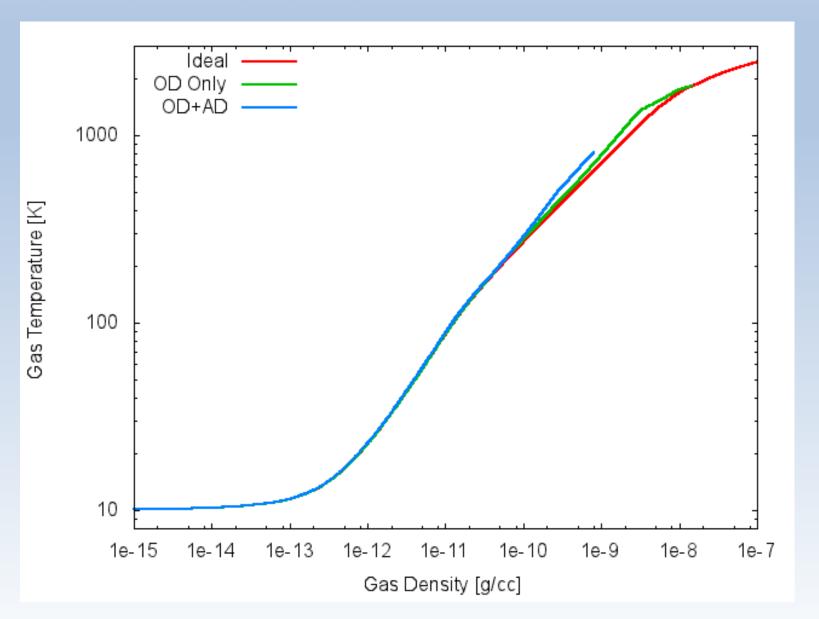
Summary

RMHD simulations of protostellar collapse with non-ideal MHD

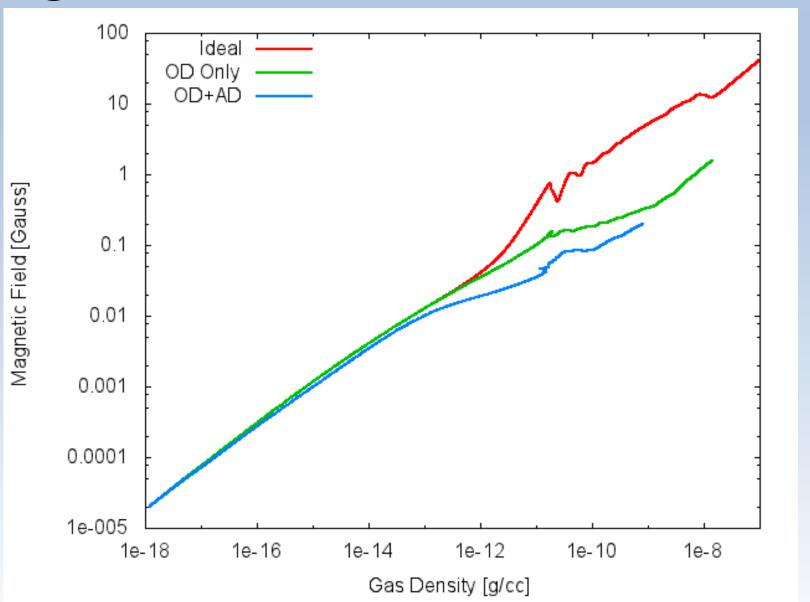
- Radiation transfer and realistic thermodynamics are important
- Magnetic braking is so efficient in the ideal MHD case that no rotationally-supported disks can be formed in the early phase
- Ohmic dissipation enables early formation of disks
- As natural byproducts, two different outflows are launched: slow, loosely collimated outflows from the first core scale and fast, well collimated jets from the protostellar core scale
- (preliminary!) With ambipolar diffusion, disk formation can be possible even before the second collapse (= birth of a star)
- Disks can be formed early, but should be small, will grow later
- Magnetic Braking Catastrophe is not so catastrophic as it sounds, rather a quantitative question: how, when, and how massive?
- ALMA observations of young disks will be crucial

Thank you!

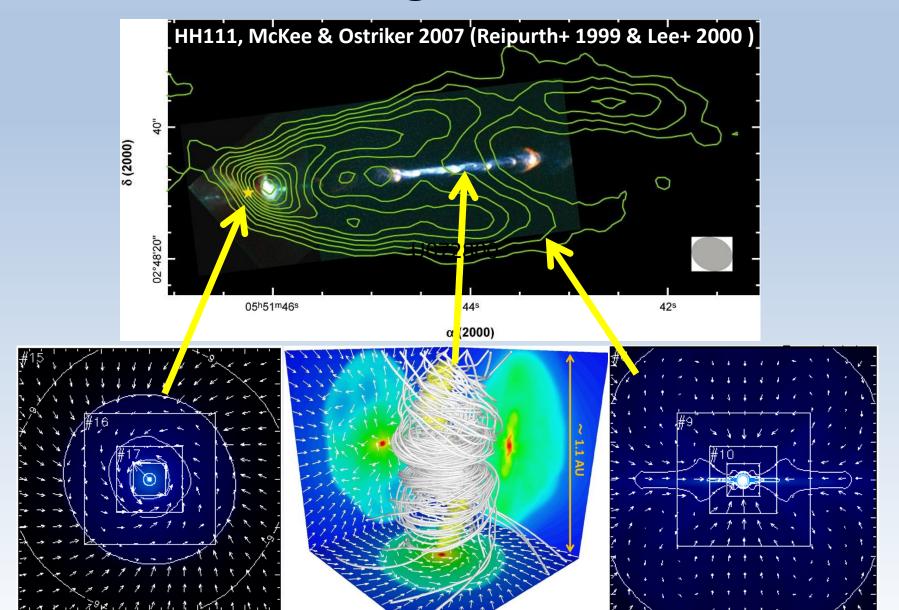
Thermal Evolution



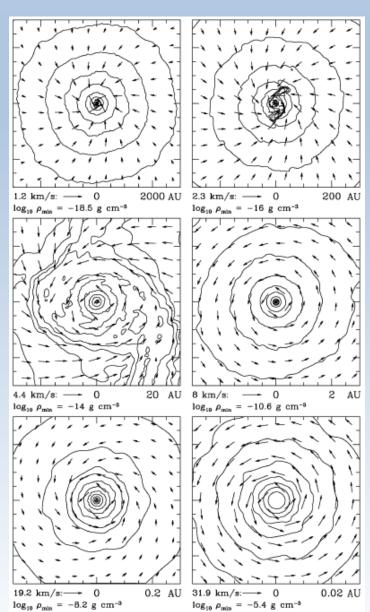
Magnetic Field



With some "imagination"...



Gravitational Torque



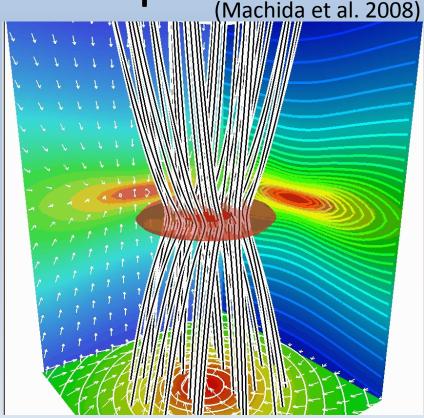
Bate (1998) first performed 3D SPH simulations of protostellar collapse and showed that the rotationally-supported disk becomes unstable and spiral arms are formed.

These non-axis-symmetric structure can transport ang. mom. efficiently and finally a protostar is formed. (see also Matsumoto & Hanawa 03, Saigo et al. 08, Commercon et al. 08, etc.)

Note: Thermodynamics (radiation transfer) is modeled using a fitting formula based on 1D RHD simulations (so-called barotropic approximation)

3D MHD Protostellar Collapse

(Matsumoto 2007, SFUMATO AMR MHD code)



3D simulations show that magnetic fields transport ang. mom. efficiently.

(see also, Matsumoto & Tomisaka 04, Banerjee & Pudritz 06, Hennebelle & Fromang 08, Hennebelle & Teyssier 08, Duffin & Pudritz 09, Hennebelle & Ciardi 09, etc., etc...)