

$$\frac{\partial t}{\partial r} \propto \frac{\partial \Omega}{\partial r}$$

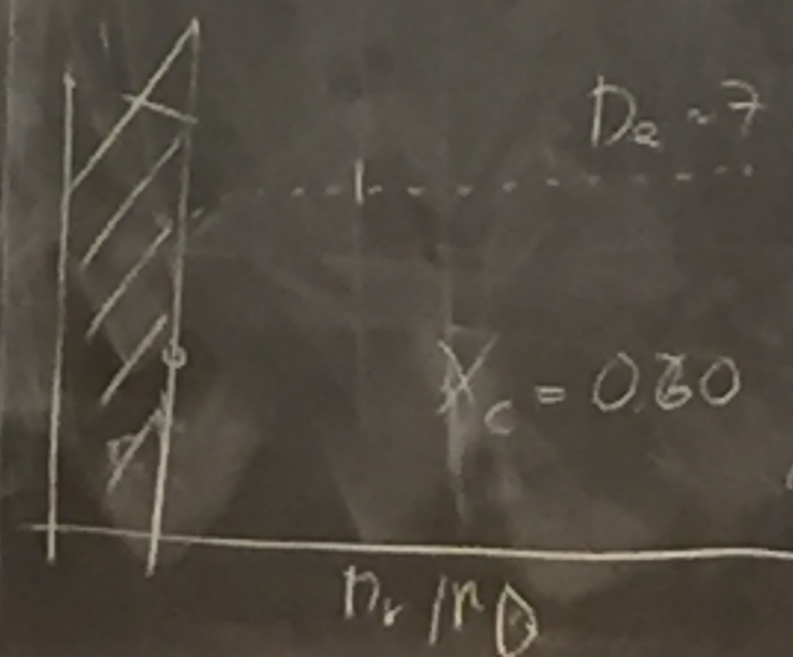
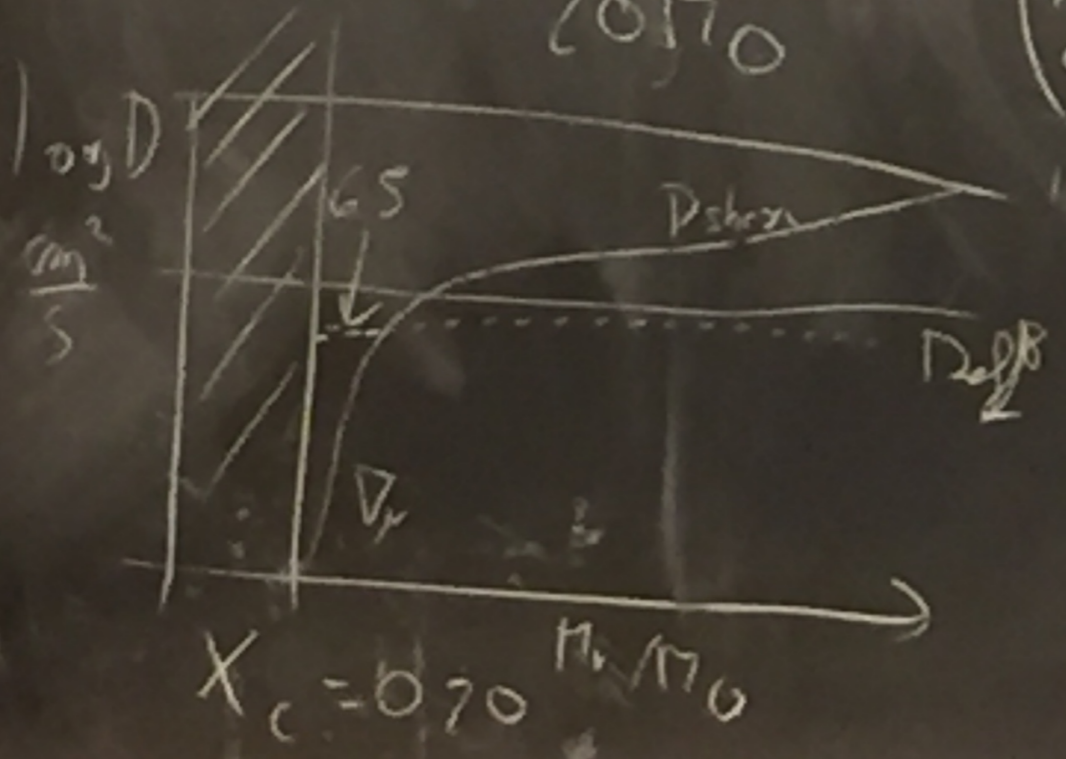
shear diff

effective diffusion coeff

weak coupling

mixing of ch $\rightarrow D_s \propto \left(\frac{\partial \Omega}{\partial r} \right)^2$

strong coupling (solid bond)
mixing $\propto D_e \propto U \propto S$



Where we stopped yesterday
(blackboard courtesy Georges Meynet)

From a Kepler LC to core overshoot, near-core rotation, and envelope mixing



Starring
May Gade Pedersen
Timothy Van Reeth



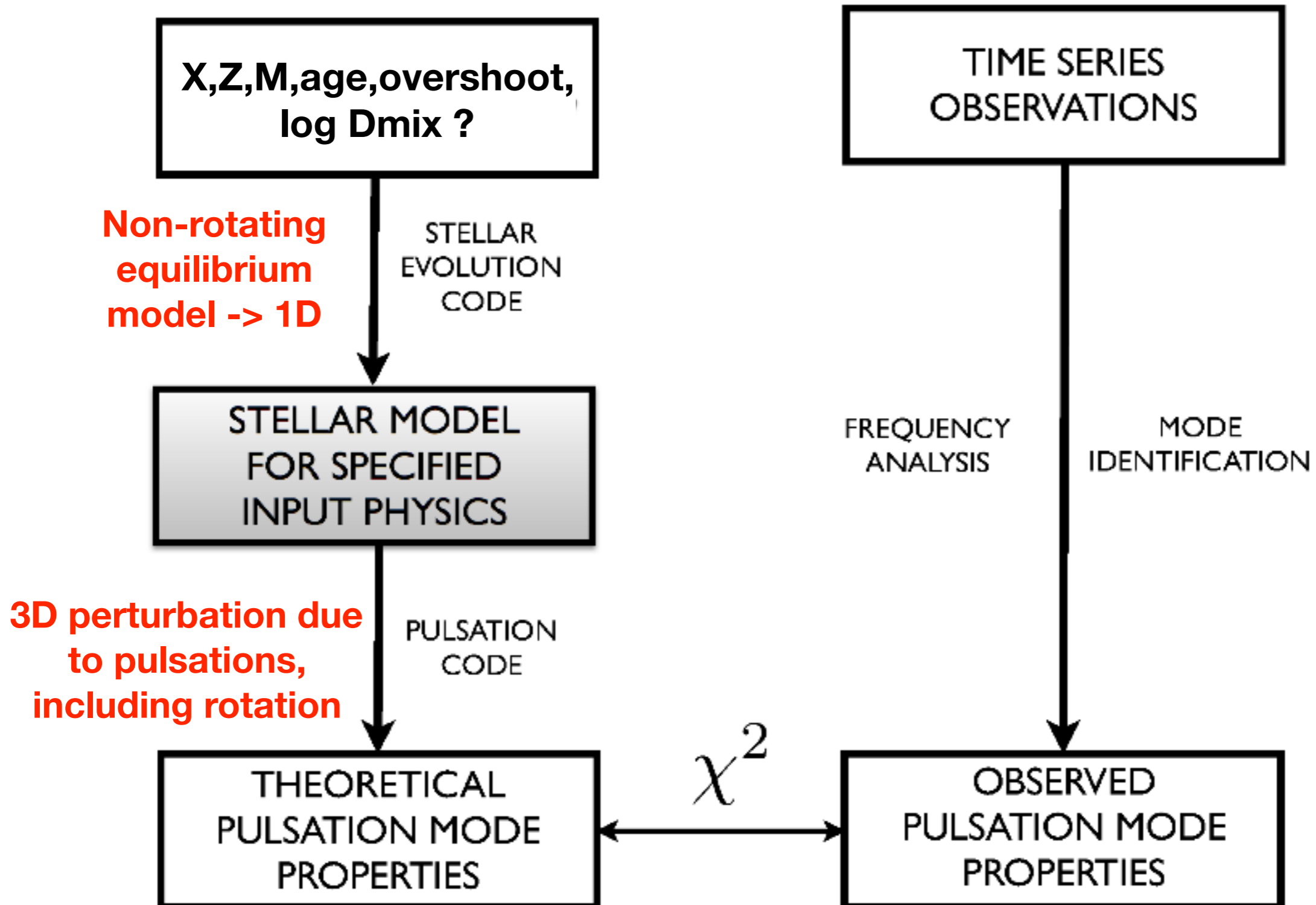
KITP, 7 April 2017

MAMSIE
WAMZIE



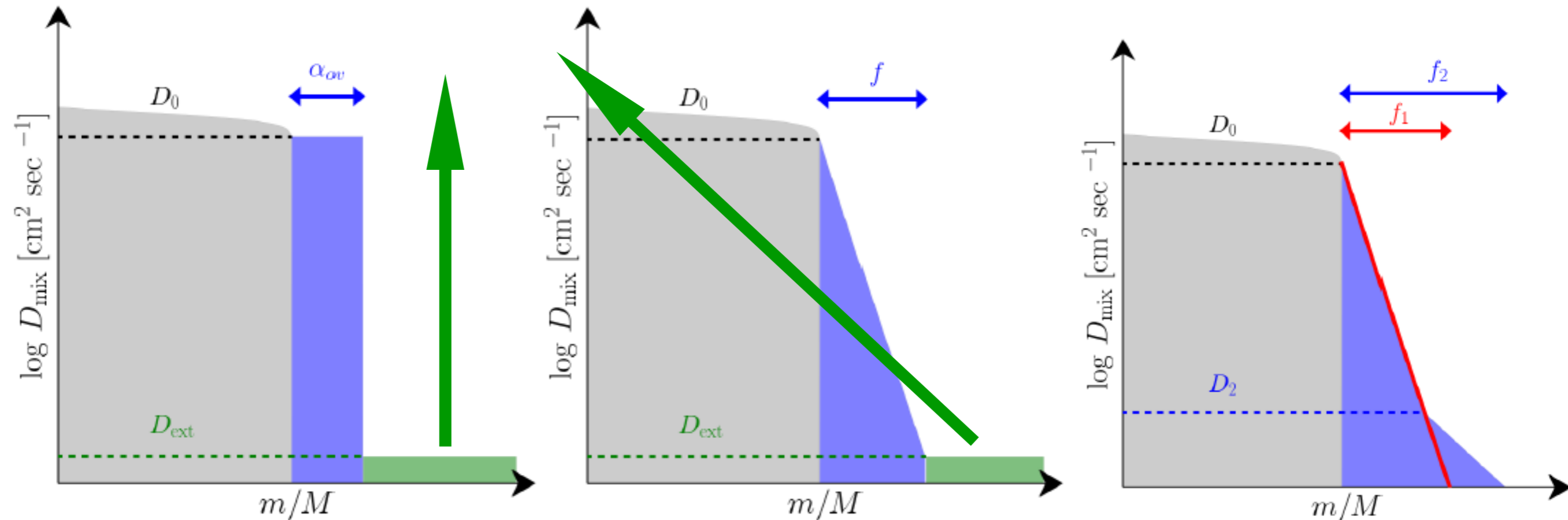


Forward seismic modelling: simplest case



Core overshoot: value & shape

Core overshoot + envelope mixing



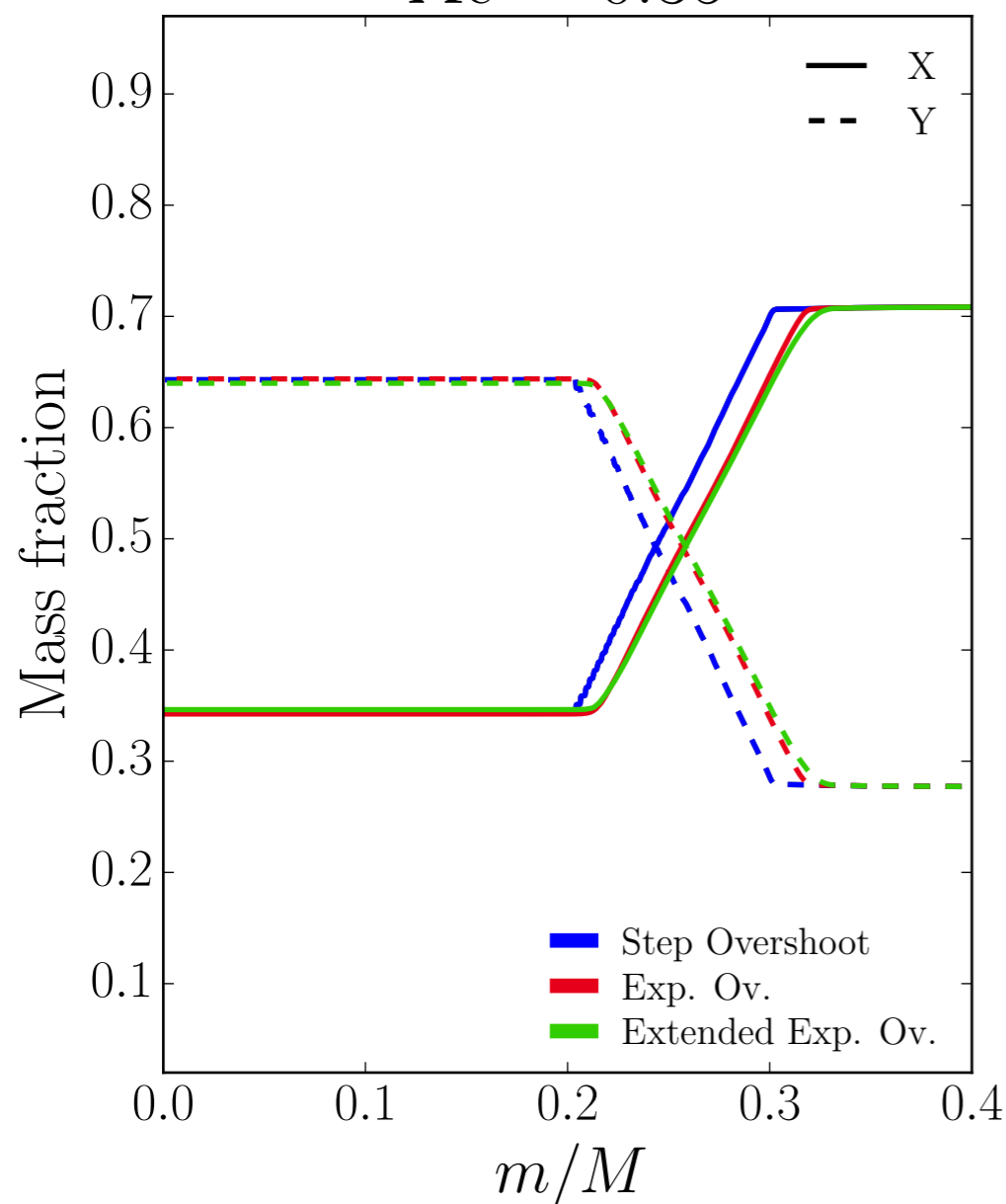
**PhD Thesis, May Gade Pedersen:
plug in results from 2D/3D hydrodynamical simulations
(cf. tutorials Tami, Daniel next Thursday)**

Core overshoot: value & shape

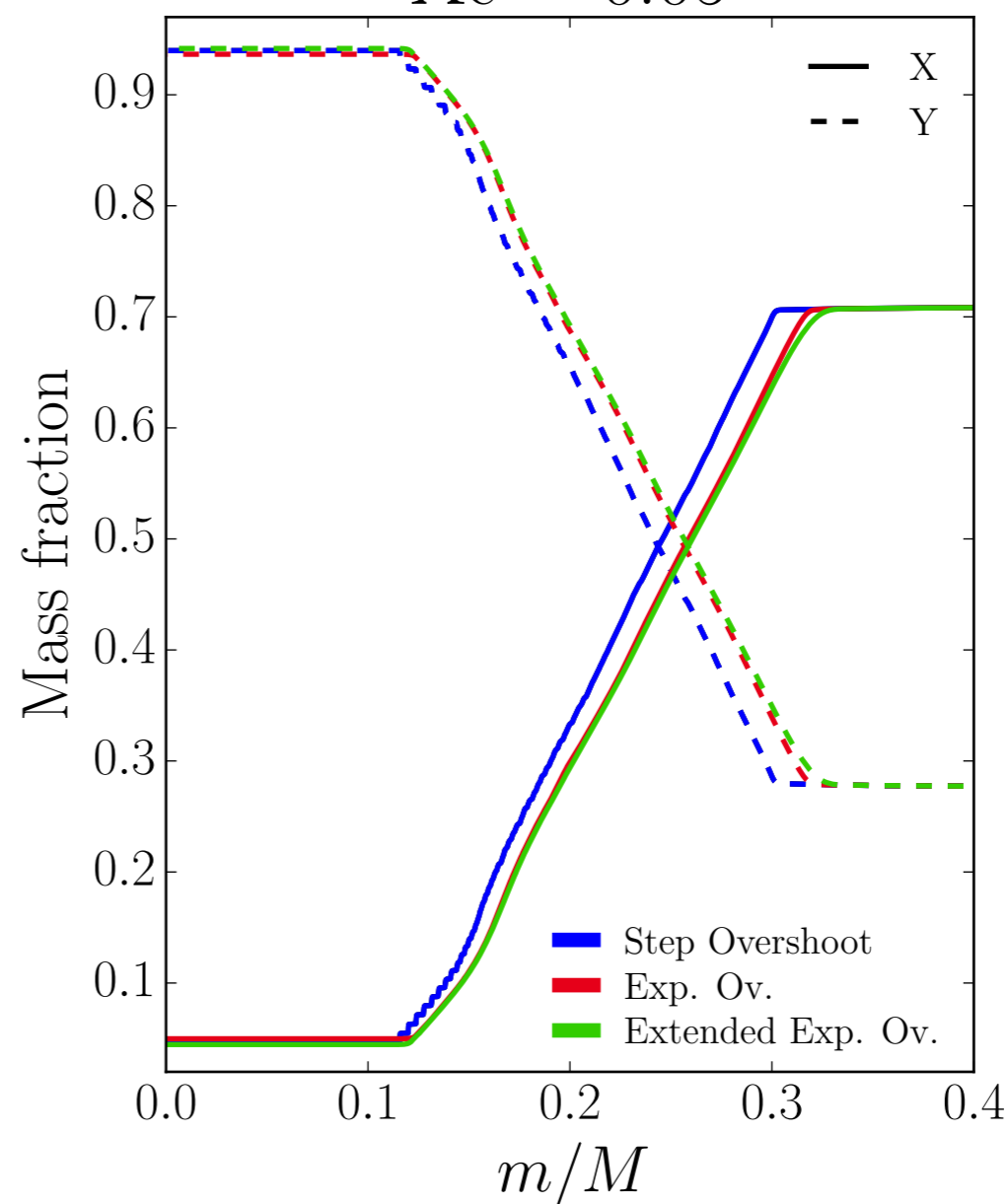
H and He profiles

PhD Thesis, May Gade Pedersen

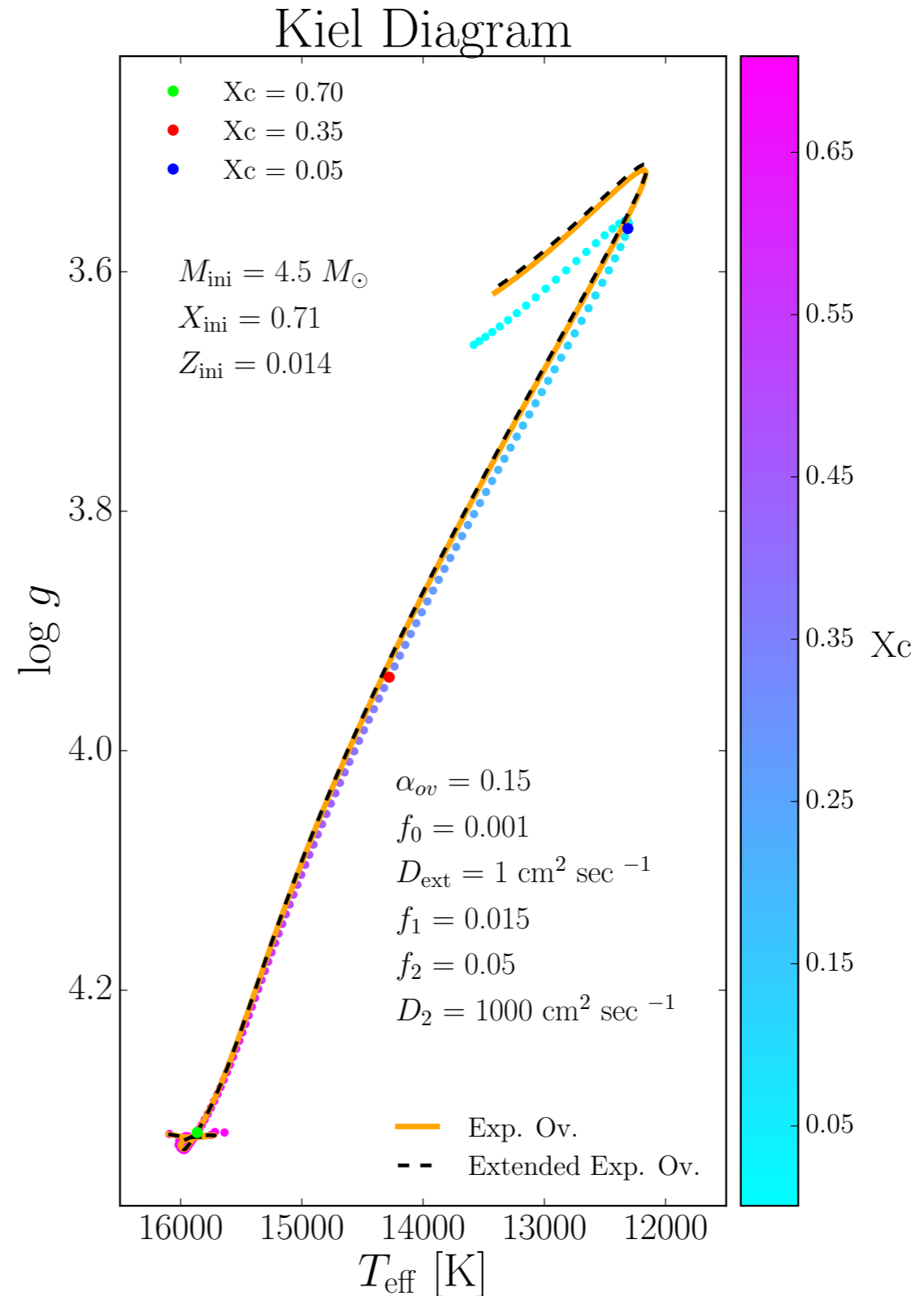
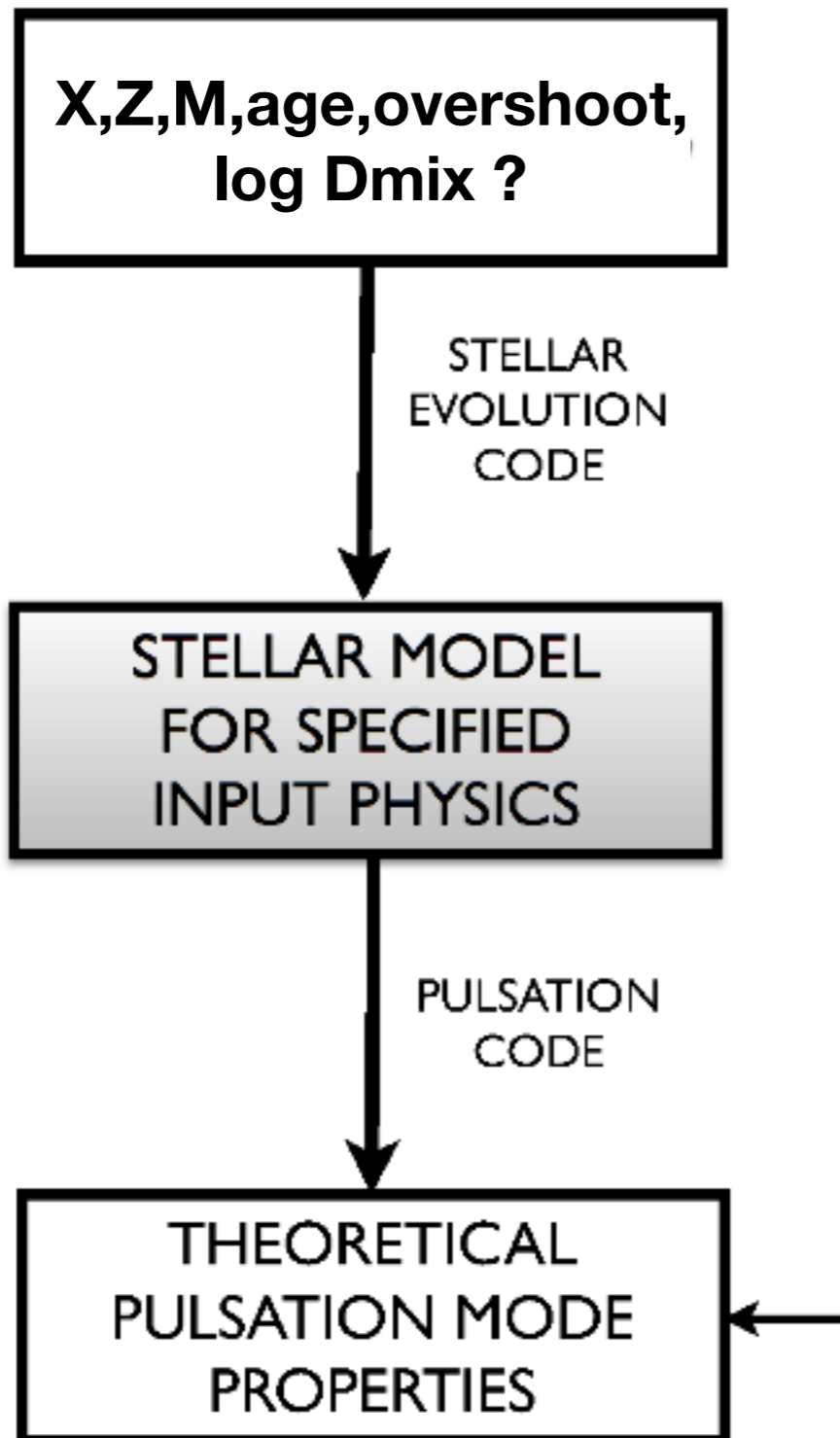
$X_c = 0.35$



$X_c = 0.05$

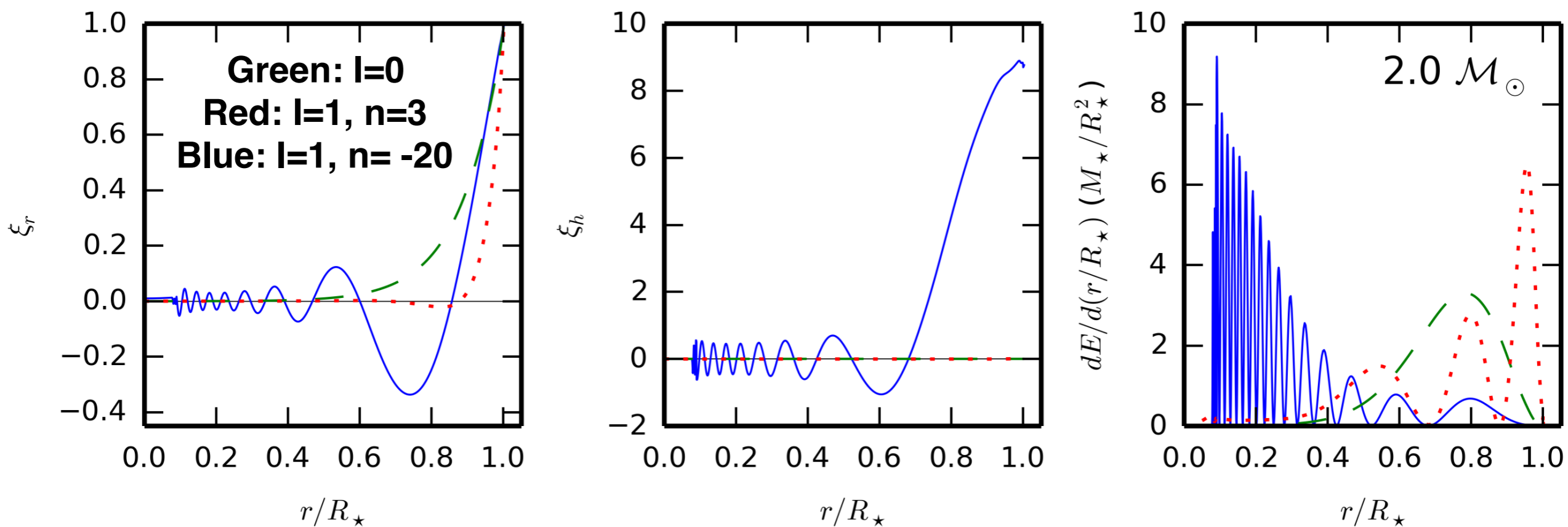


Seismic modelling: interior, not HRD or Kiel

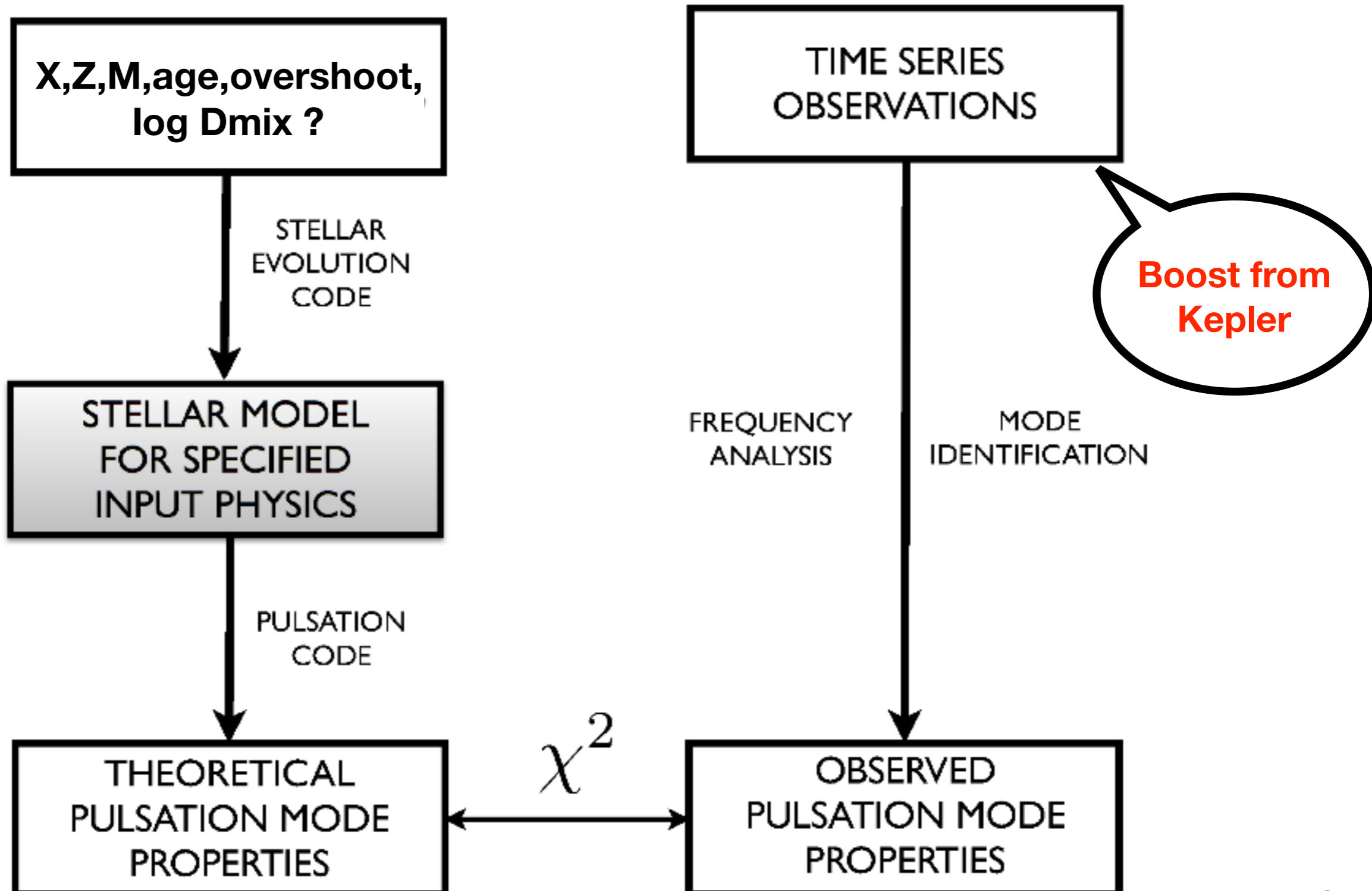


Beauty of gravity modes (only since Kepler)

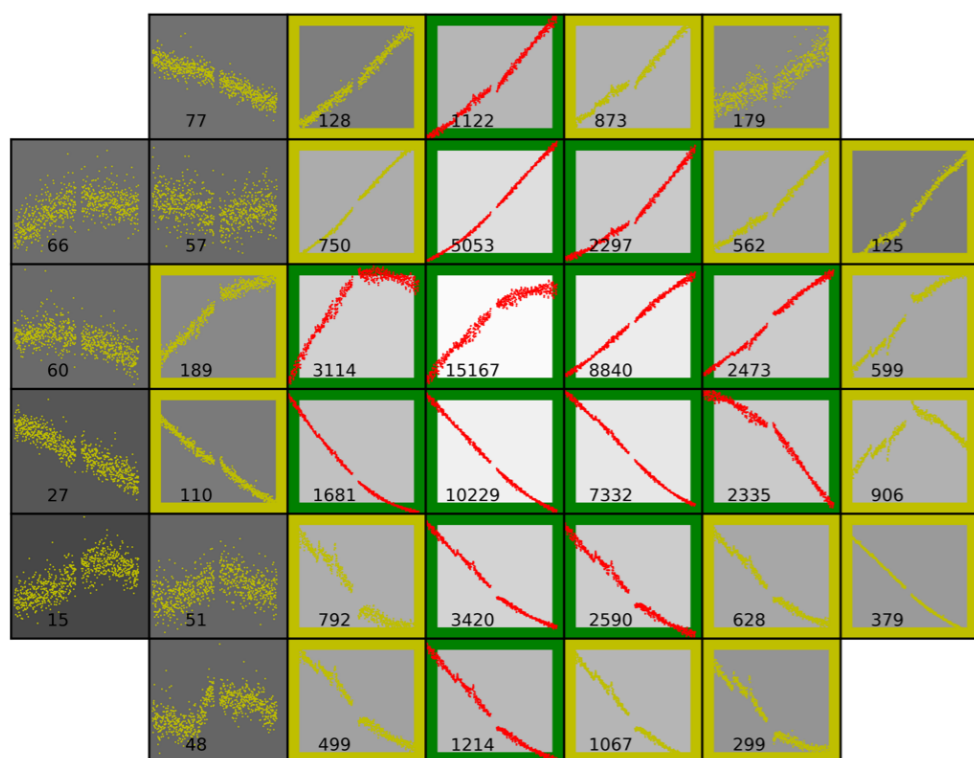
Allow probing of near-core regions & determine X, Z, M, age, overshoot, $\log D_{\text{mix}}$ in B & F-type stars



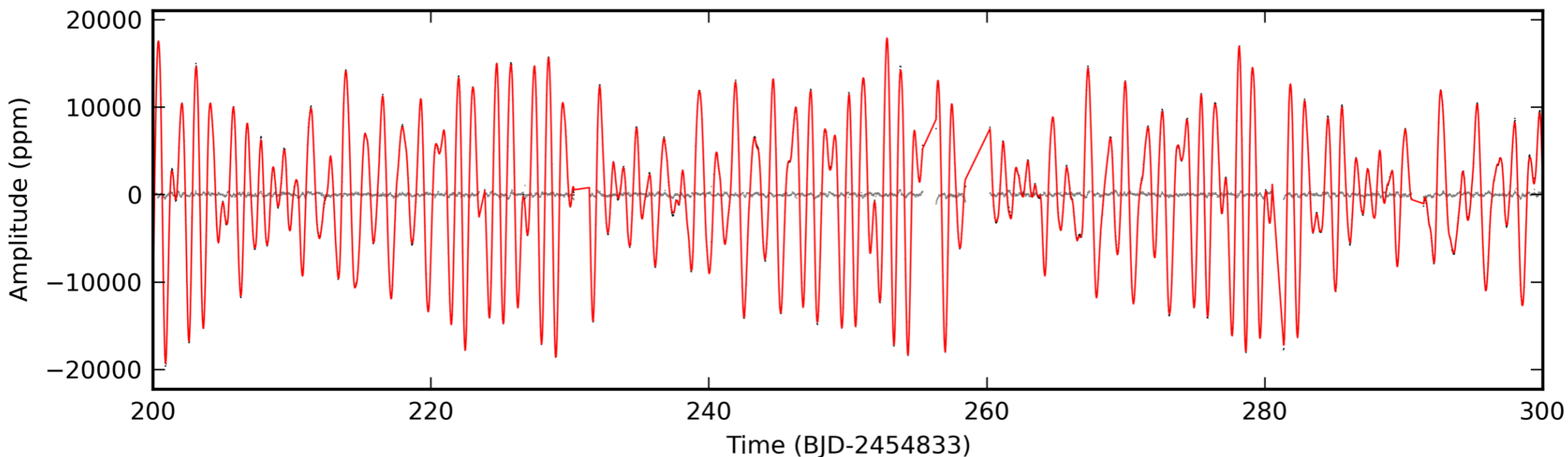
Forward seismic modelling: simplest case



Kepler LCs from raw pixel data (MAST)



Pápics et al. (2014):
LC: 29.42 min (sum of 270 exposures)
Nyquist frequency: 24.47/d



Fourier analysis

Fourier transform of $x(t)$:

$$F(f) \equiv \int_{-\infty}^{+\infty} x(t) \exp(2\pi i f t) dt$$

Fourier transform $F(f)$ of sum of harmonic functions with frequencies f_1, \dots, f_n and amplitudes A_1, \dots, A_n :

$$x(t) = \sum_{k=1}^n A_k \exp(2\pi i f_k t) : F(f) = \sum_{k=1}^n A_k \delta(f - f_k)$$

For $x(t) =$ sine with frequency f_1 , $F(f) \neq 0$ for $f = \pm f_1$

For $x(t) =$ sum of n harmonic functions with frequencies f_1, \dots, f_n ,

$F(f) =$ sum of δ -functions $\neq 0$ for $\pm f_1, \dots, \pm f_n$

Real data set: $x(t)$ known for a discrete number of times t_j ,
 $j=1, \dots, N$

Discrete Fourier transform

$$F_N(f) \equiv \sum_{j=1}^N x(t_j) \exp(2\pi i f t_j)$$

$F_N \neq F$! but connected through window function: $w_N(t) \equiv \frac{1}{N} \sum_{j=1}^N \delta(t - t_j)$

Hence:
$$\frac{F_N}{N} = \int_{-\infty}^{+\infty} x(t) w_N(t) \exp(2\pi i f t) dt$$

Discrete Fourier transform of window function = spectral window $W_N(f)$:

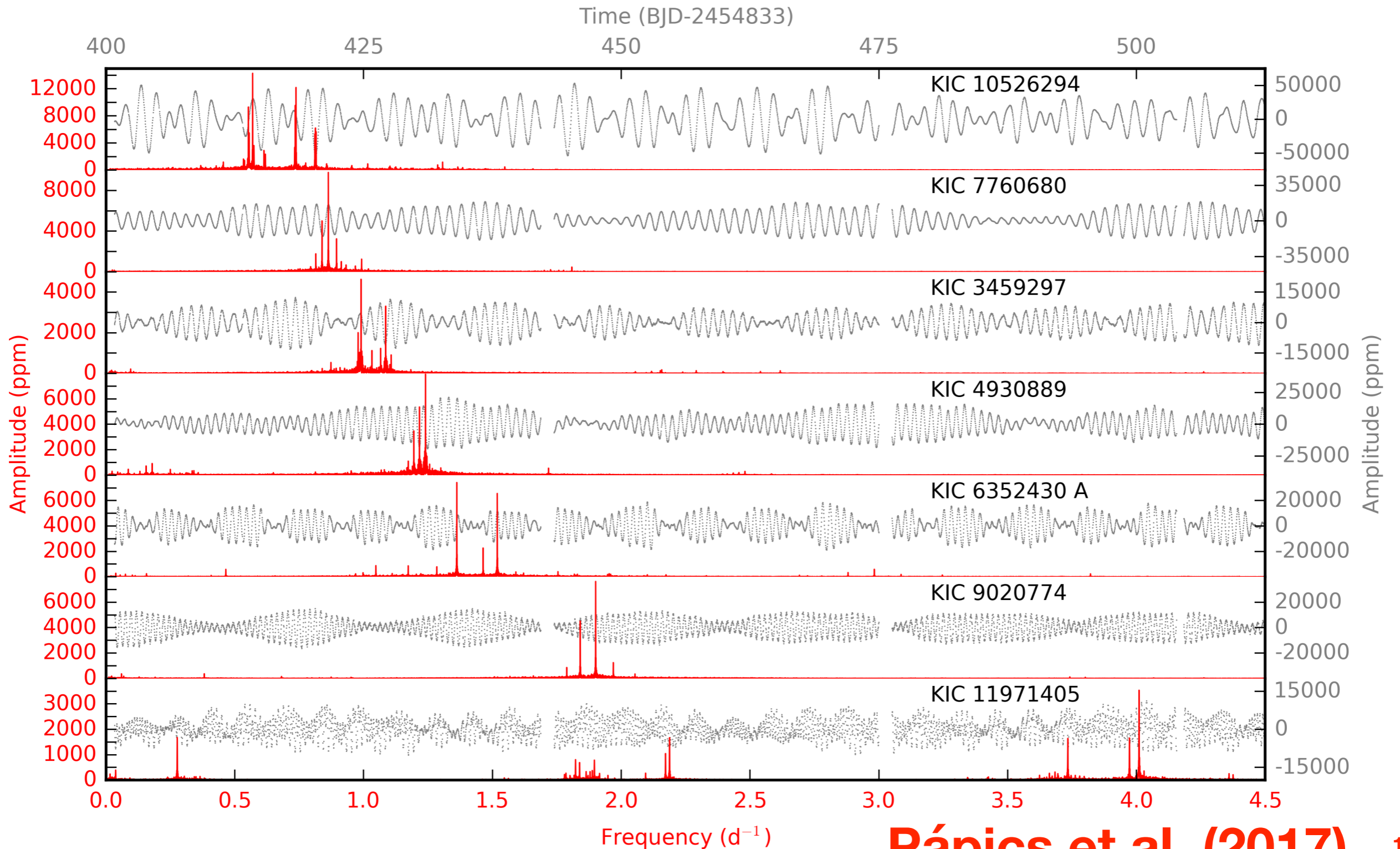
$$W_N(f) = \frac{1}{N} \sum_{j=1}^N \exp(2\pi i f t_j)$$

Discrete Fourier transform = convolution of spectral window and

Fourier transform:
$$F_N(f)/N = F(f) * W_N(f)$$



Detecting gravity-mode oscillations



Pápics et al. (2017)

Prewhitening & residuals

Least-squares fitting with f fixed: $x_i(t_i) = A \sin [2\pi(ft_i + \psi)] + C$

Variance reduction in $\in [0,1]$:

$$1 - \frac{\sum_{i=1}^N \{x_i - [A \sin (2\pi(ft_i + \psi)) + C]\}^2}{\sum_{i=1}^N (x_i - \bar{x})^2}$$

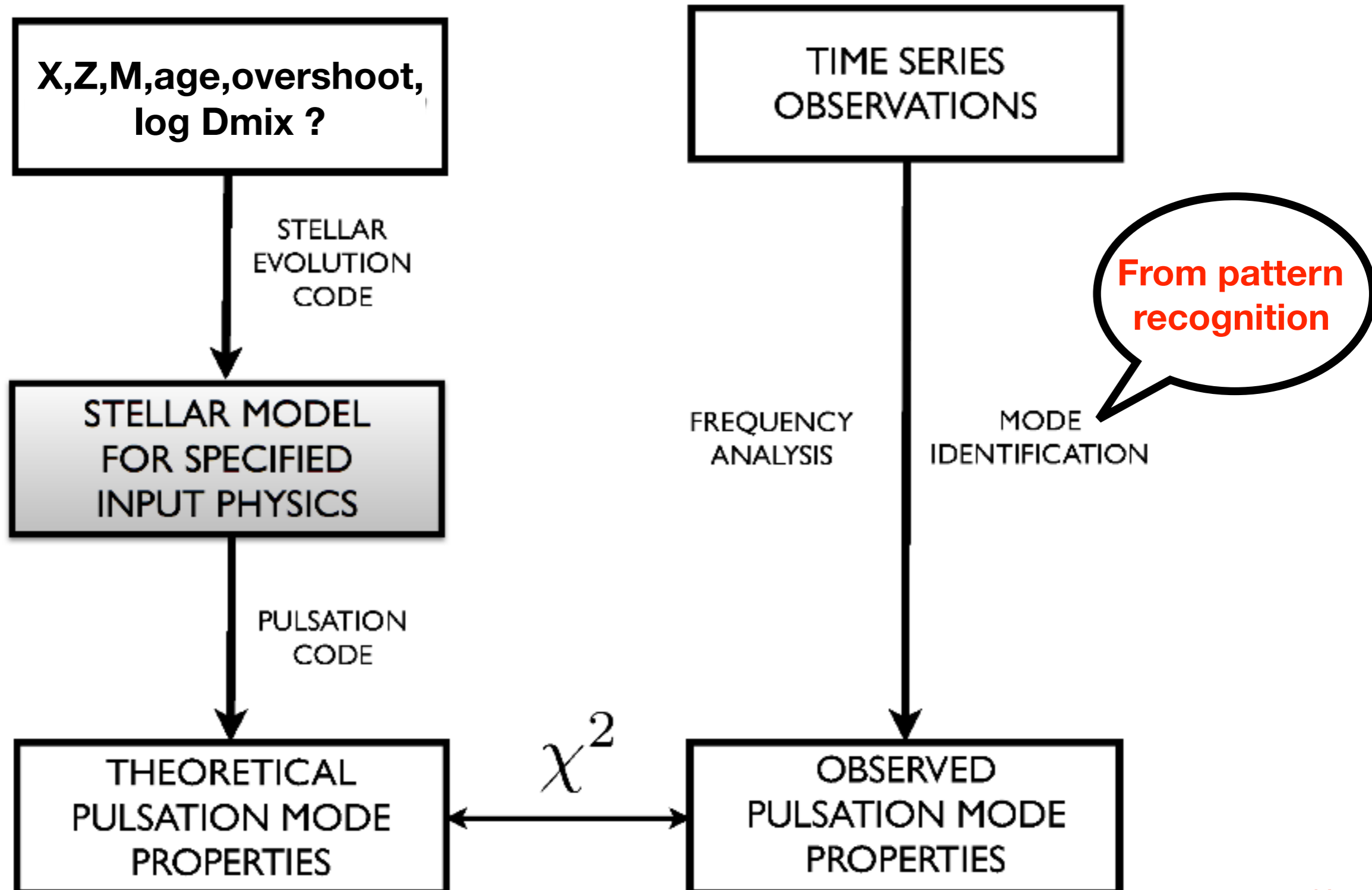
Search for new frequencies in residuals $R_i(f) \equiv x_i - x_i^c(f)$ with

$$x_i^c(f) \equiv A \sin [2\pi(ft_i + \psi)] + C$$

and so on. BUT: frequency is only known up to certain precision:

optimising f within uncertainty interval is necessary: $\sigma_f = \frac{\sqrt{6}\sigma_R}{\pi\sqrt{N}AT}$
do NLLS fitting + prewhitening

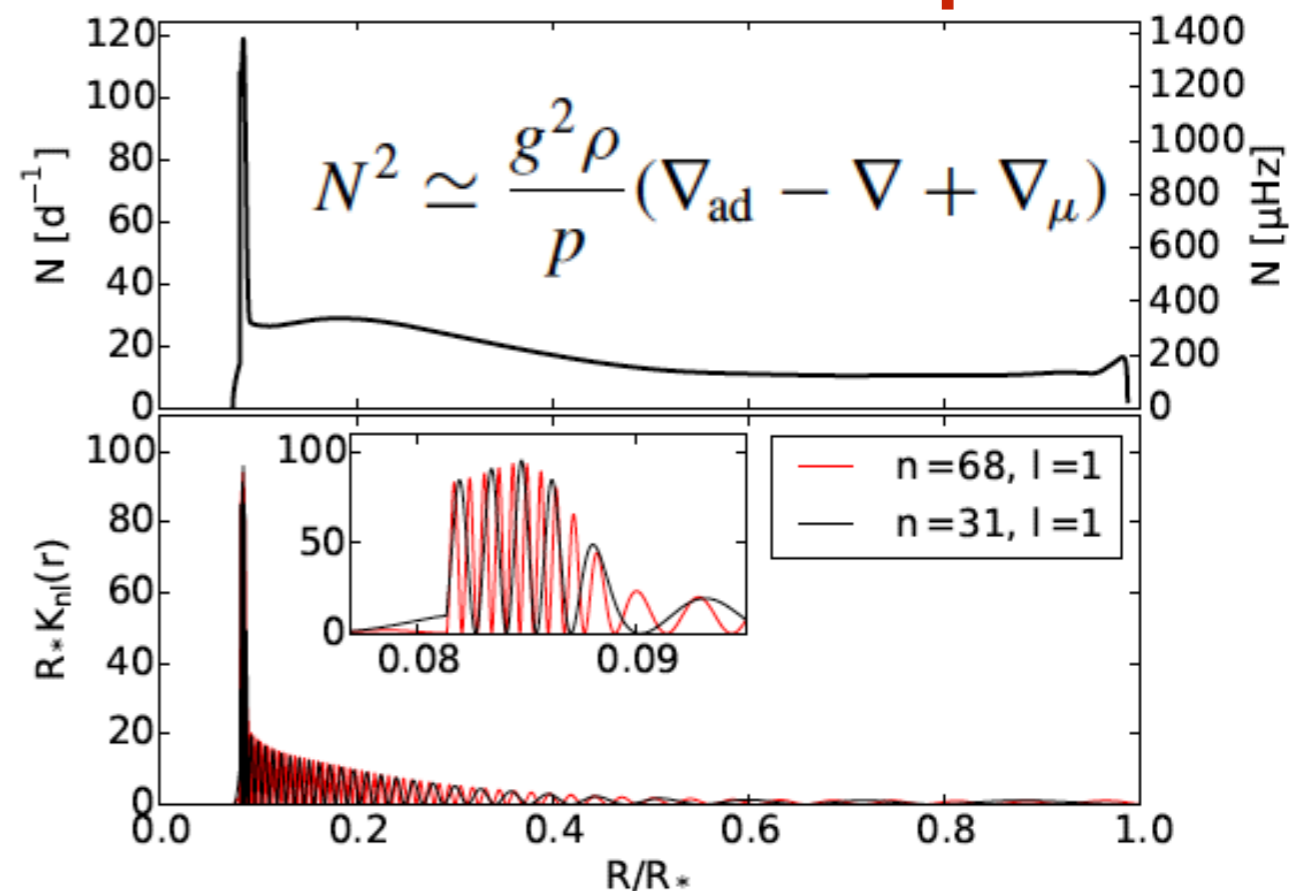
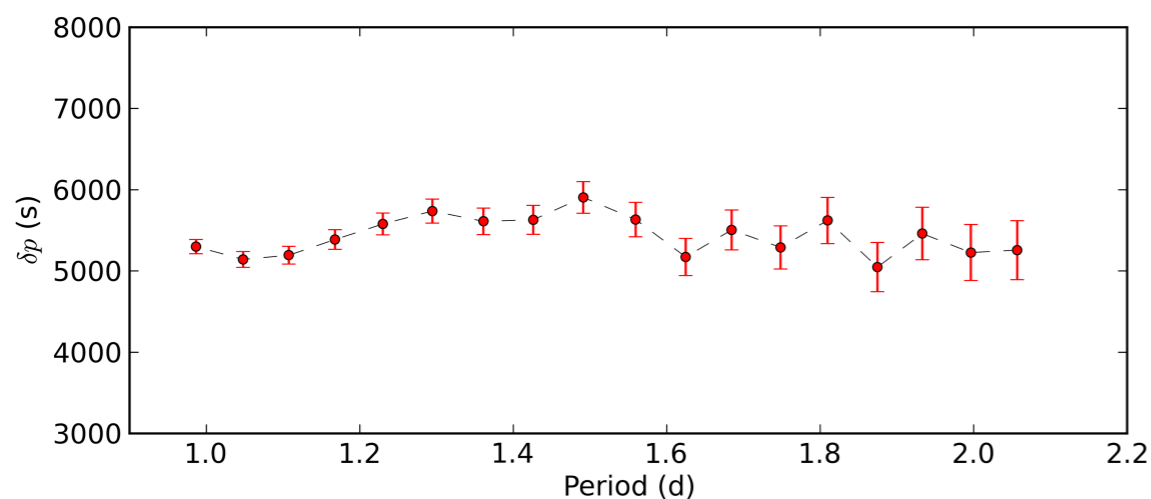
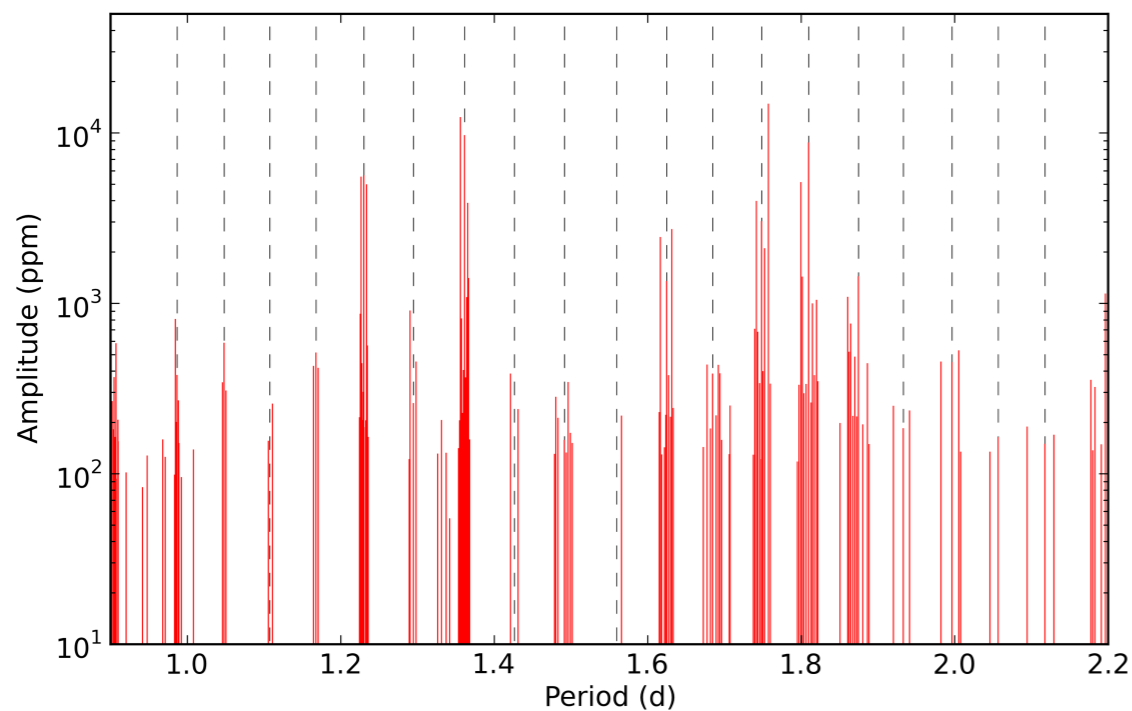
Forward seismic modelling: simplest case



MESA summerschool 2016

$$P_{nl} = \frac{\Pi_0}{\sqrt{l(l+1)}}(n + \alpha_g), \quad \Delta\Pi_l = \frac{\Pi_0}{\sqrt{l(l+1)}}, \quad \Pi_0 = 2\pi^2 \left(\int_{r_1}^{r_2} N \frac{dr}{r} \right)^{-1}$$

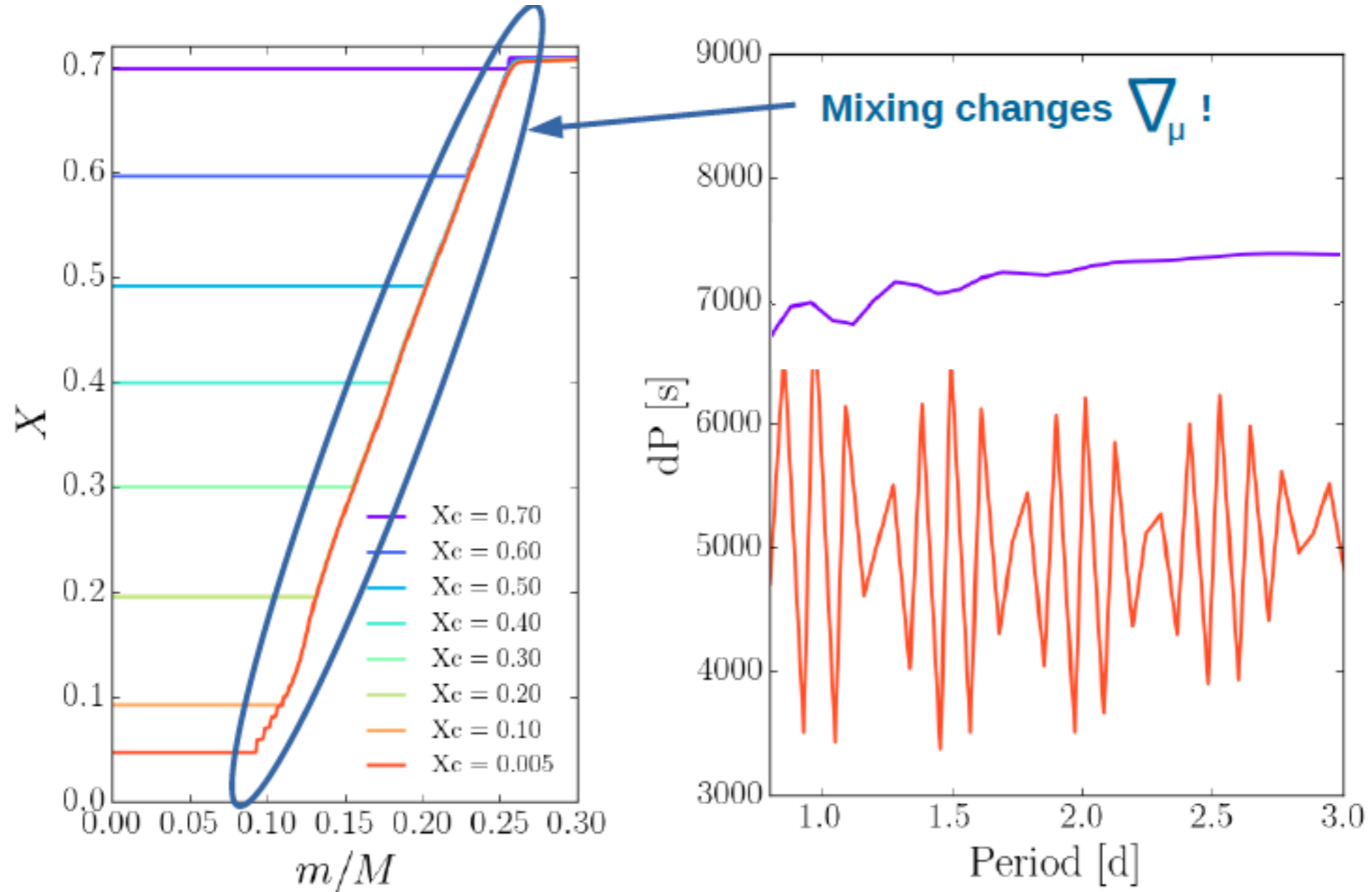
No rotation, no magnetic field, convective core, radiative envelope





Receding core & μ -gradient zone

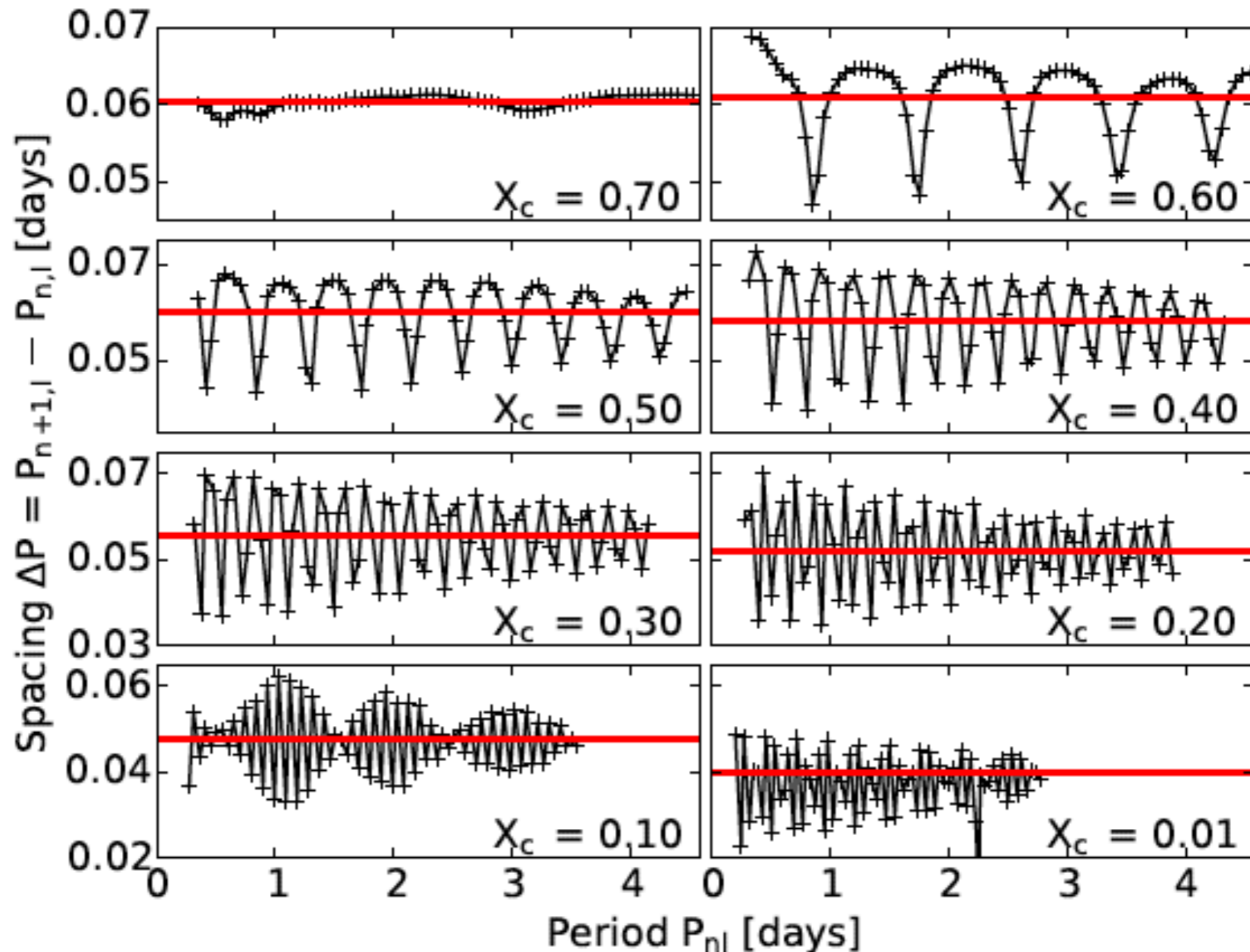
PhD Thesis, May Gade Pedersen



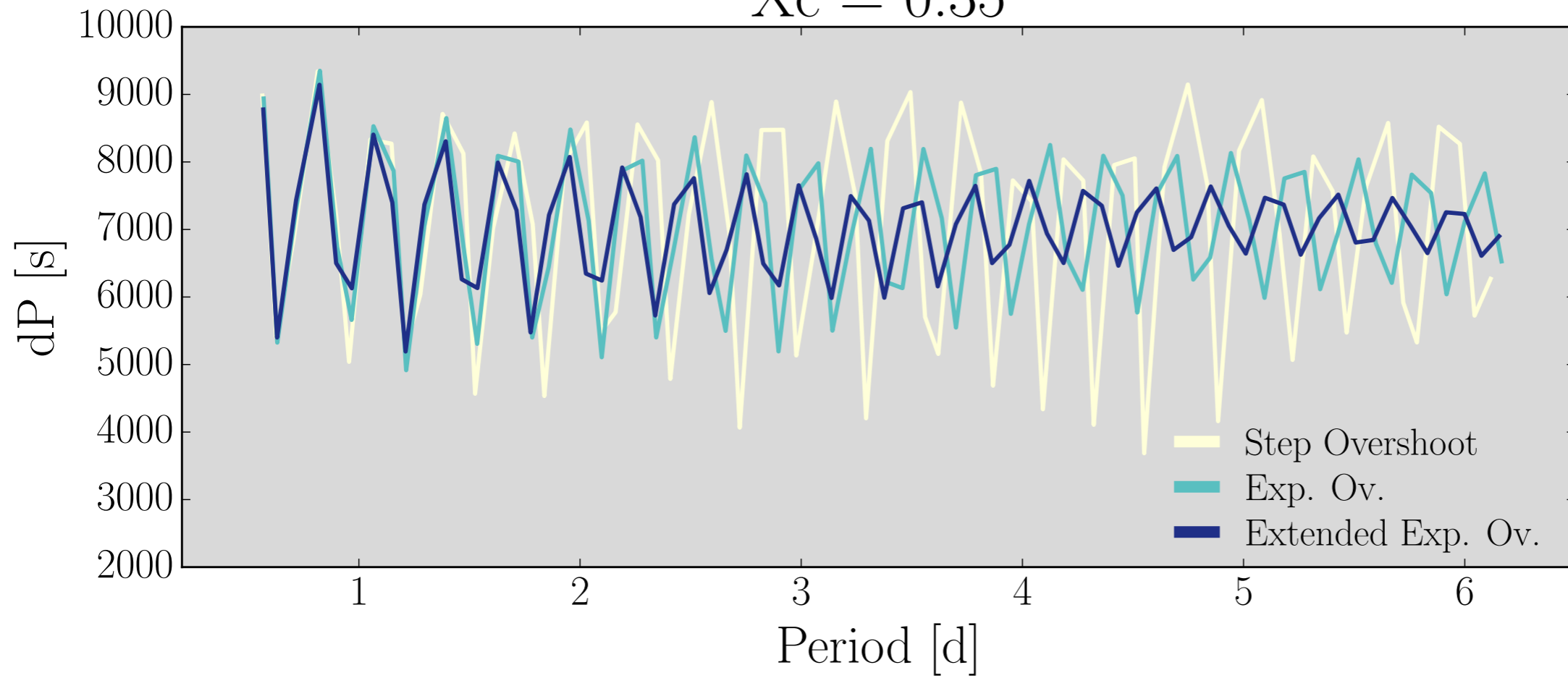
Minilab 2: stellar pulsations



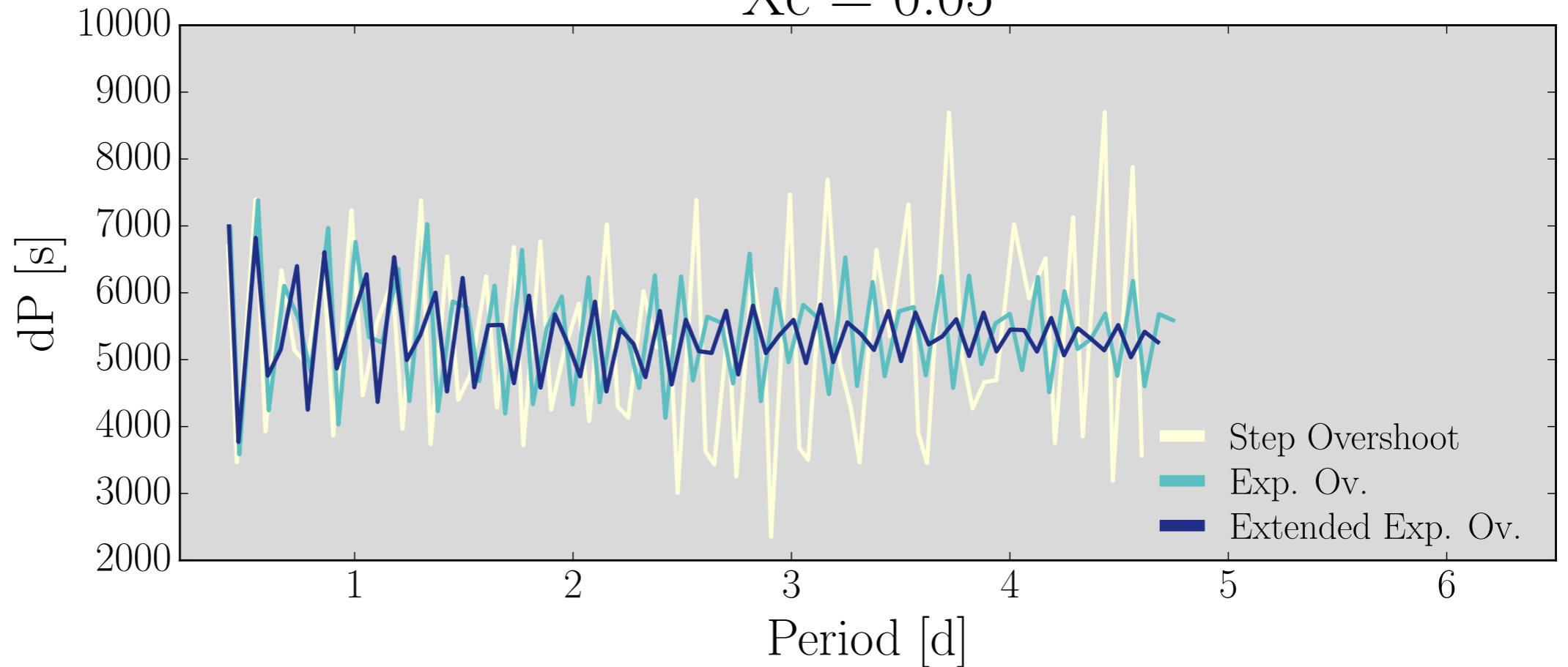
Results for the $3M_{\odot}$ model



$X_c = 0.35$



$X_c = 0.05$

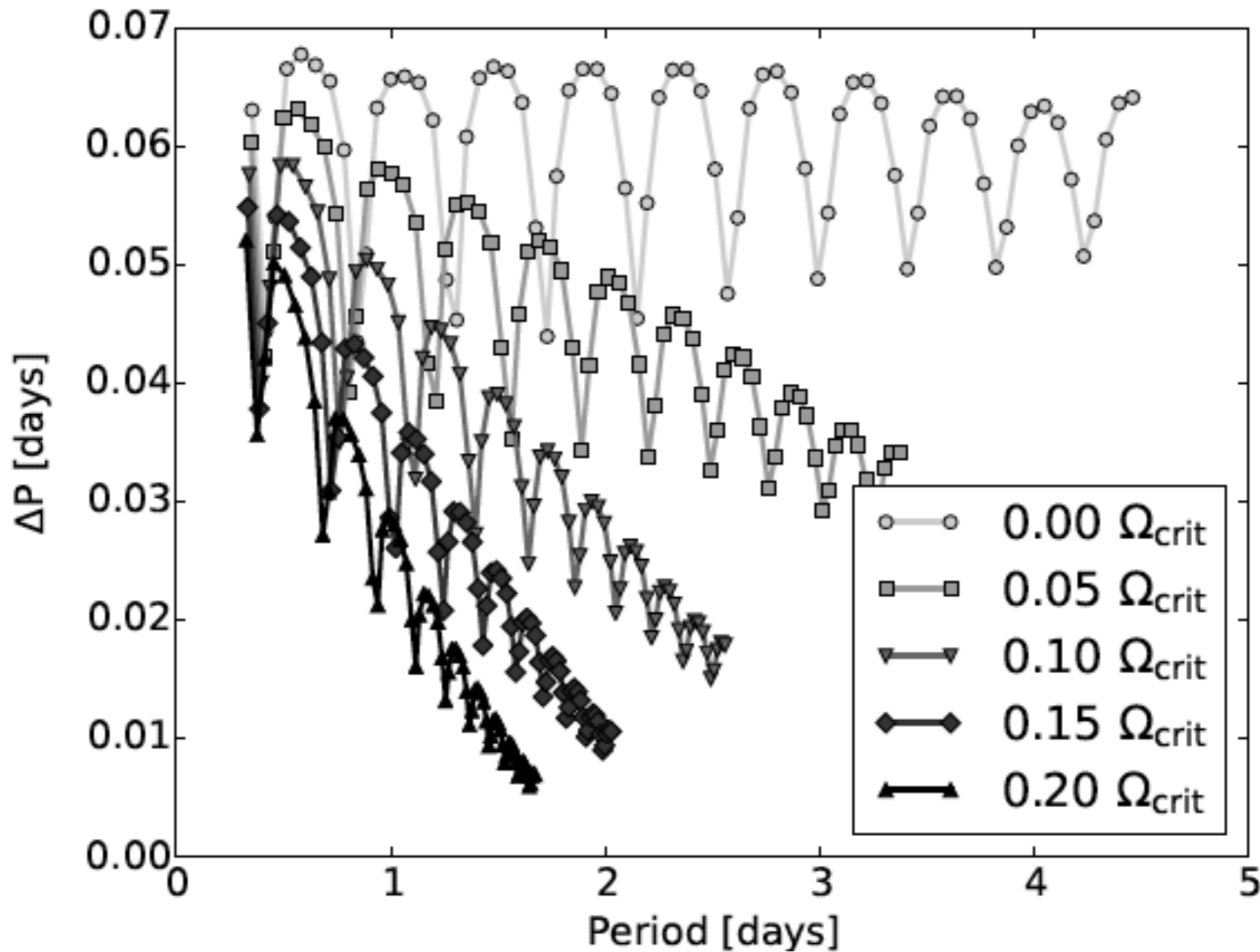


**PhD
Thesis,
May
Gade
Pedersen**

Minilab 3: rigid stellar rotation



Results for the $3M_{\odot}$ model at $X_c = 0.50$



Dips & chemical mixing

Various types of mixing in MESA

- Mixing disconnected with rotation
- Rotationally induced mixing

$$D = D_{Non-rot} + f_c ($$

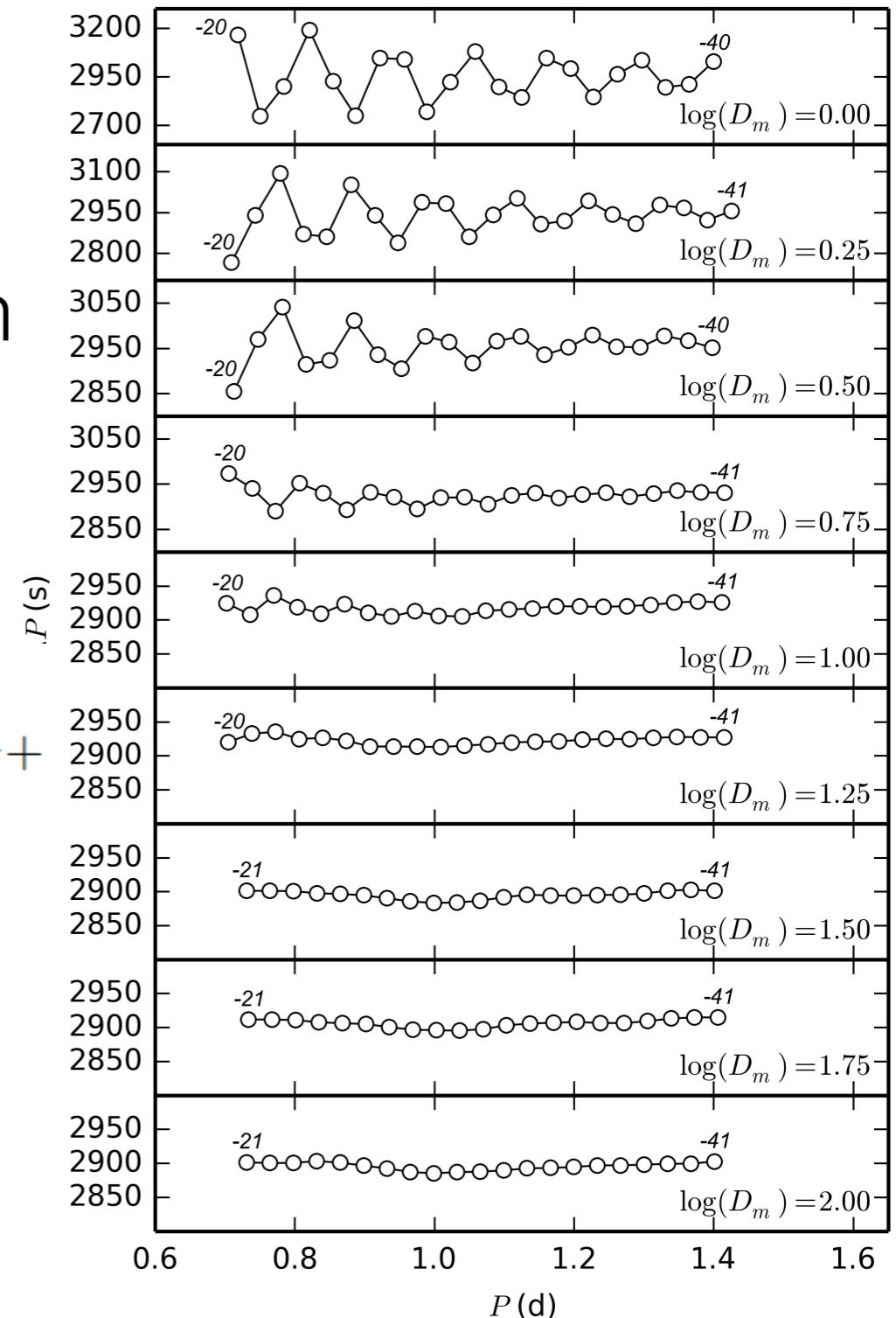
$$\cancel{D_{DSI, factor} \cdot D_{DSI}} + \cancel{D_{SH, factor} \cdot D_{SH}} + \cancel{D_{SSI, factor} \cdot D_{SSI}} +$$

$$\cancel{D_{ES, factor} \cdot D_{ES}} + \cancel{D_{GSF, factor} \cdot D_{GSF}} + \cancel{D_{ST, factor} \cdot D_{ST}})$$

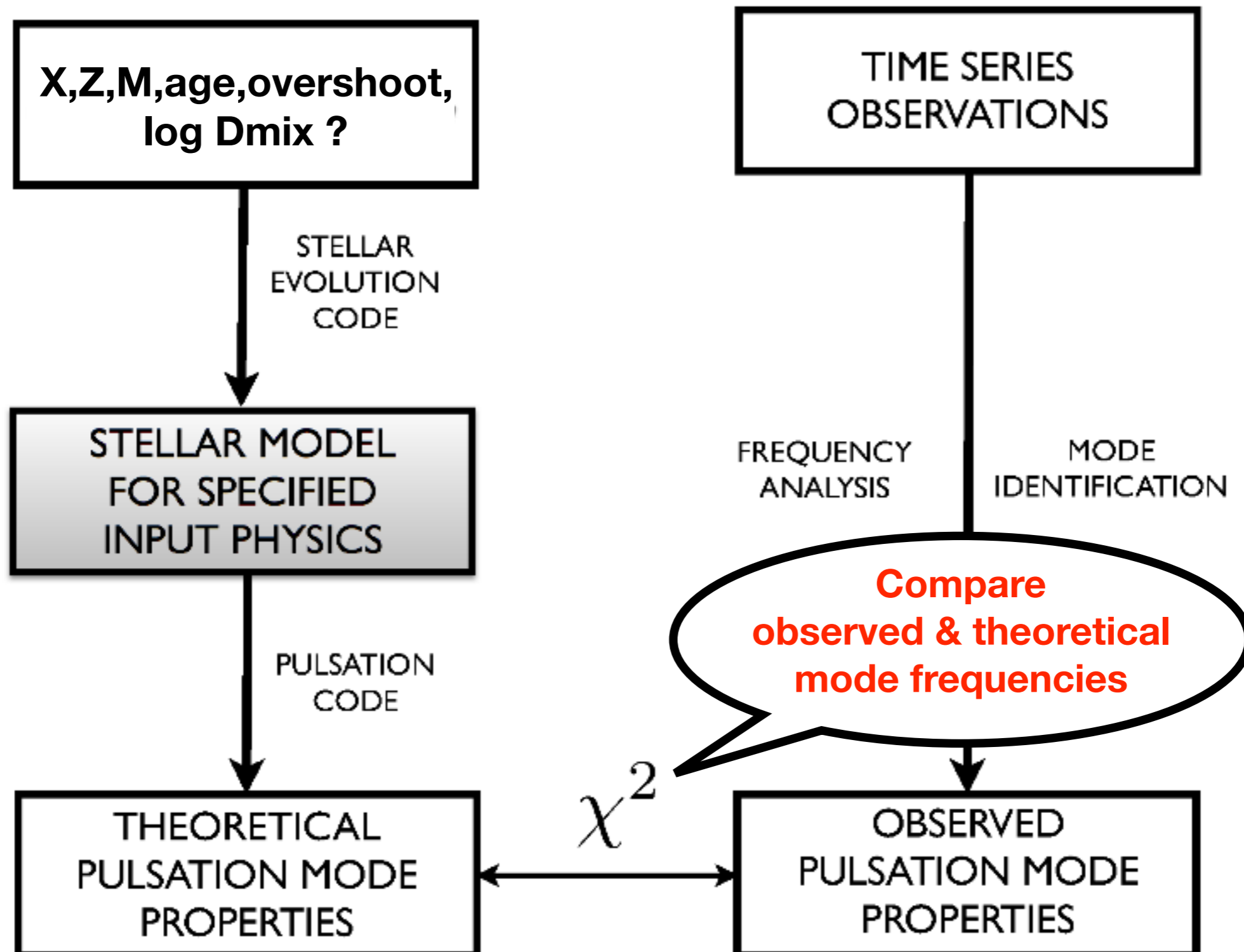
→ **log Dmix**

$$M = 1.76 M_{\odot}, Z = 0.0125, f_{ov} = 0.005, X_c = 0.20$$

Schmid & Aerts (2016)

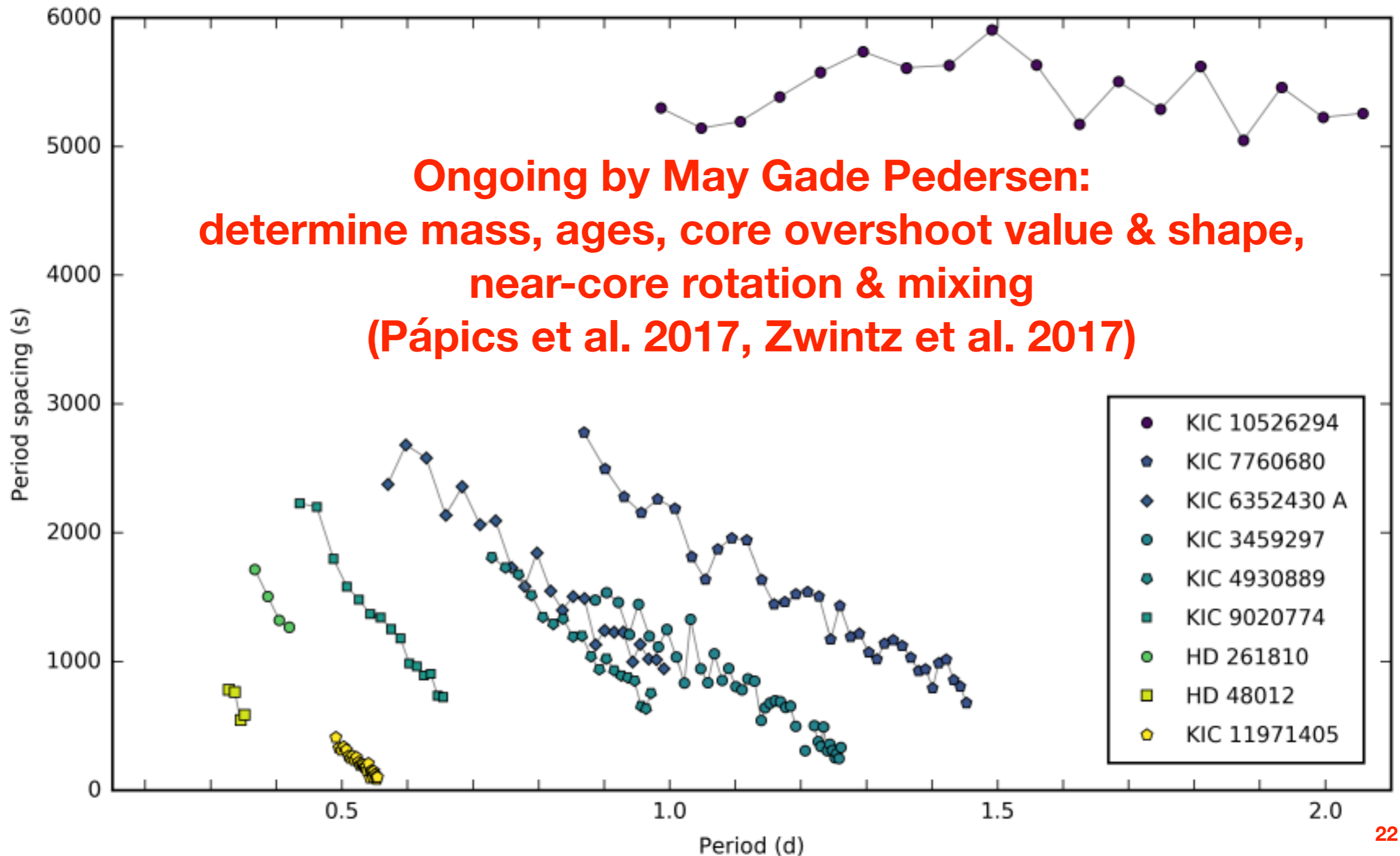


Forward seismic modelling: simplest case

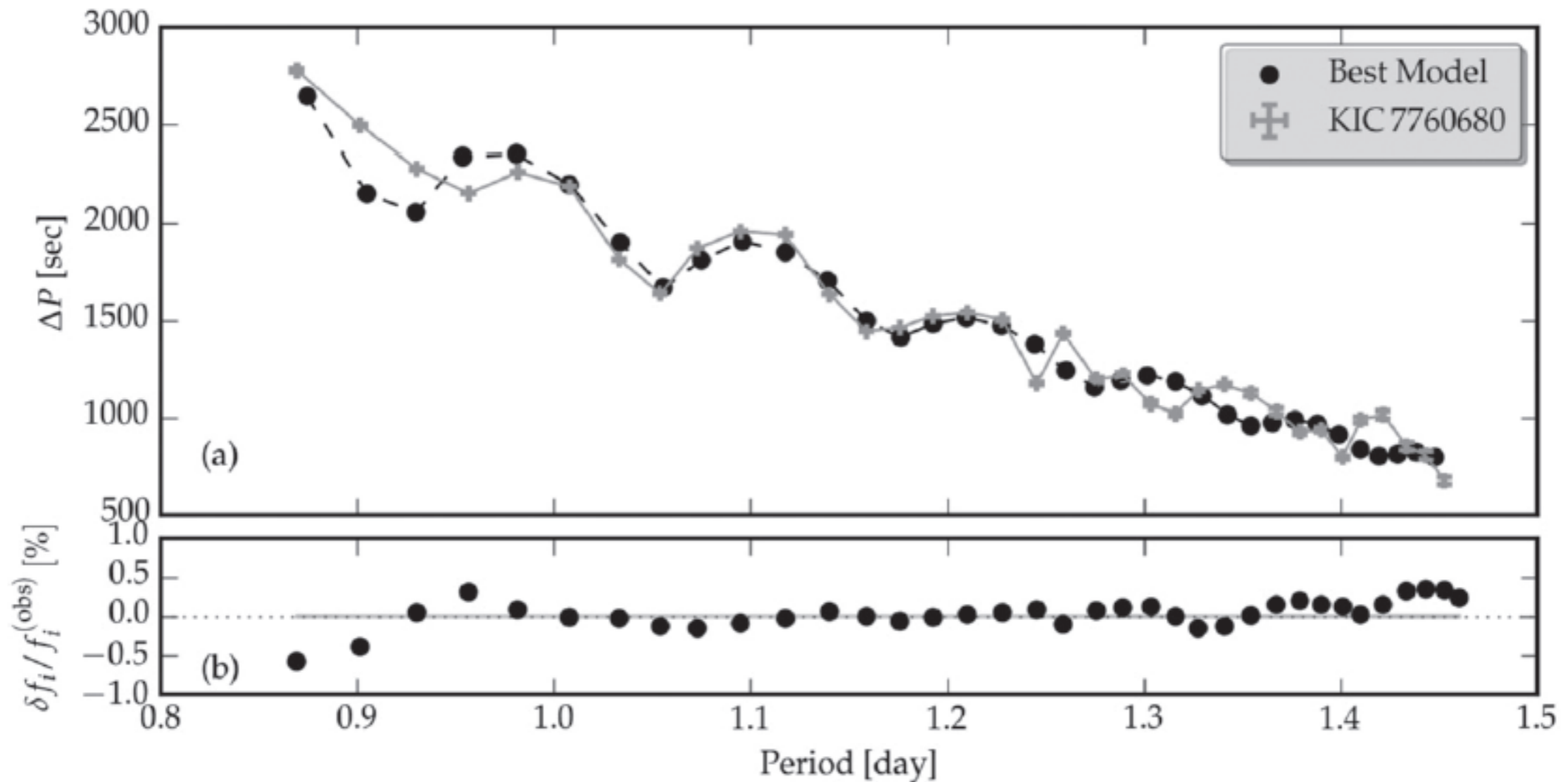


Gravity-mode period spacings massive stars

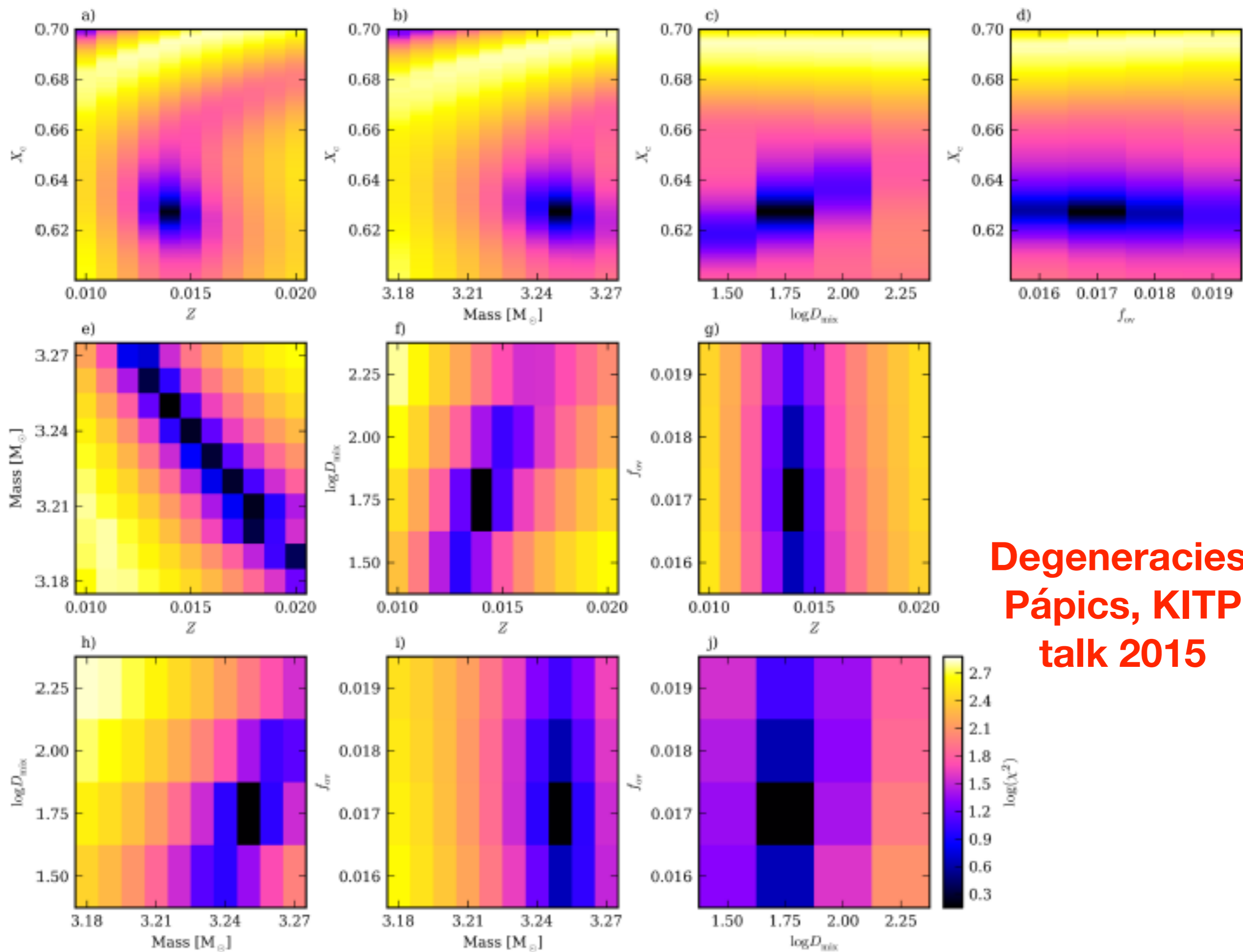
Ongoing by May Gade Pedersen:
 determine mass, ages, core overshoot value & shape,
 near-core rotation & mixing
 (Pápics et al. 2017, Zwintz et al. 2017)



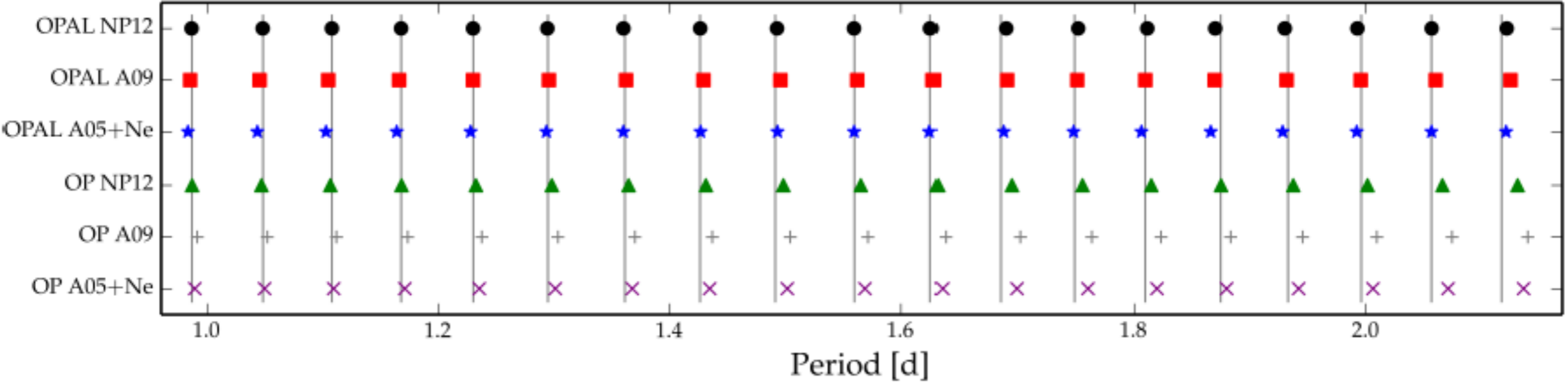
A real star: KIC 7760680



KIC7760680: $M=3.25M_{\odot}$, $X_c=0.50$, $f_{rot}=0.48/d$, $f_{ov}=0.024$ Hp
 $\log D_{mix}=0.75 \pm 0.25$ (Moravveji et al. 2016)



**Degeneracies,
Pápics, KITP
talk 2015**



**Dependence on opacities and chemical mixture
(Moravveji et al. 2015)**