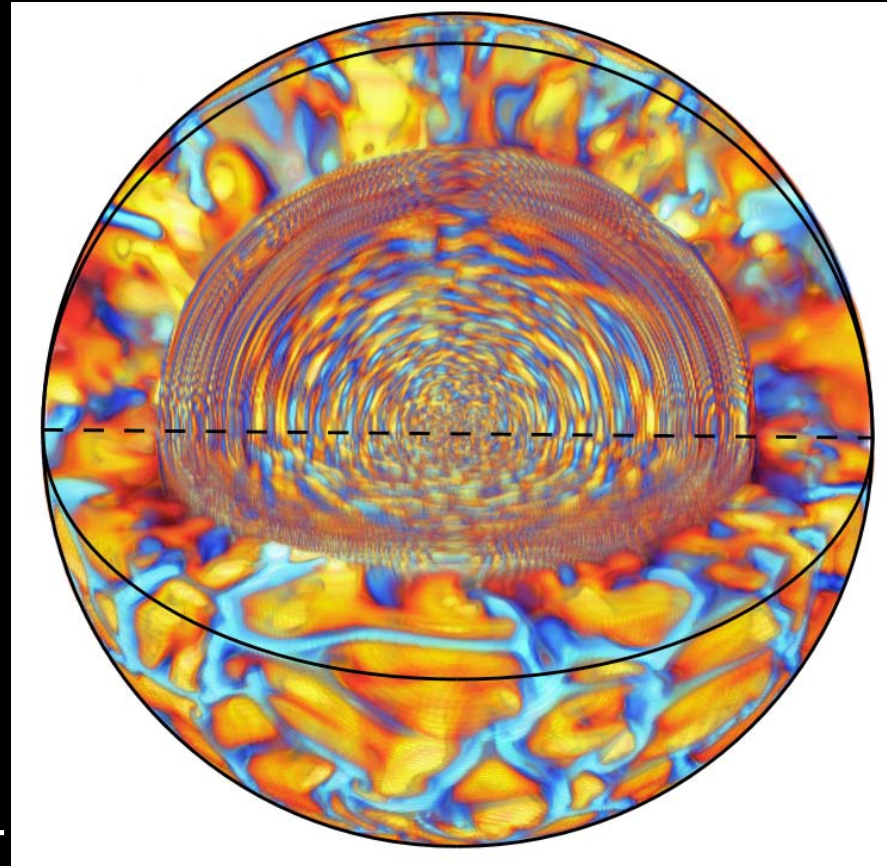


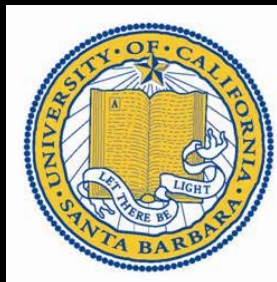
Wave-Driven Pre-Supernova Outbursts

Jim Fuller

KITP/Caltech



Alvan et al. 2014



JIM FULLER'S DIGNITY MEMORIAL
BURRITO ENDOWMENT



1025

DATE 4/18/2017

PAY TO THE
ORDER OF

Rich Townsend

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Security Features
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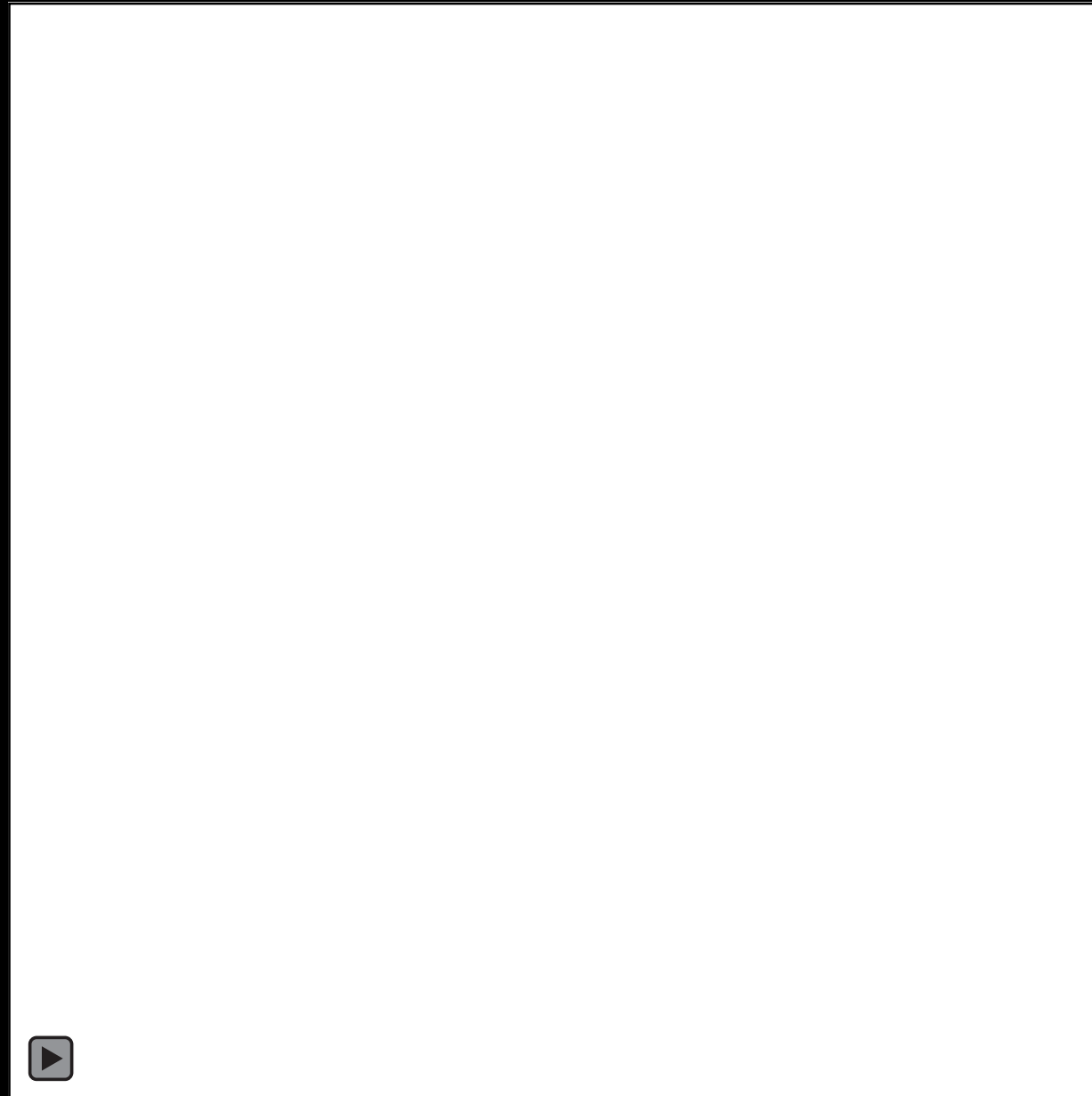
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Motivation

- Pre-SN outbursts common in interacting SNe
 - Occur in last \sim years of star's life
 - Mass loss rates enhanced by factors of $\sim 10^3$, can be $\gg 10^{-3} M_{\text{sun}}/\text{yr}$
- Mild outbursts may be common even in normal SNe
- Unexplained by standard stellar evolution
- Wave-driven outbursts possible solution

Quataert & Shiode (2012)
Shiode & Quataert (2014)

Convection excites gravity waves



Movie made by
Andrea Cristini



Wave Power

- Convection puts energy into waves at a rate

$$L_{\text{wave}} \sim \mathcal{M}_{\text{con}} L_{\text{con}}$$

Lecoanet et al. 2013
Rogers et al. 2013

- Where the convective Mach number is

$$\mathcal{M}_{\text{con}} = \frac{v_{\text{con}}}{c_{\text{sound}}} \ll 1$$

- The convective velocity can be estimated from mixing length theory

Late Stage Massive Stellar Evolution

- During late burning stages, neutrinos carry away burning energy causing burning timescales to be short

$$t_{\text{dyn}} \ll t_{\text{nuc}} \ll t_{\text{therm}}$$

$$L_{\text{nuc}} \gg L_*$$

- Consequently, there are situations where

$$L_{\text{wave}} \gg L_*$$

- Waves can transport energy to surface on timescale of $\sim t_{\text{dyn}}$

Wave Power in Massive Stars

- Huge energy fluxes during late burning phases

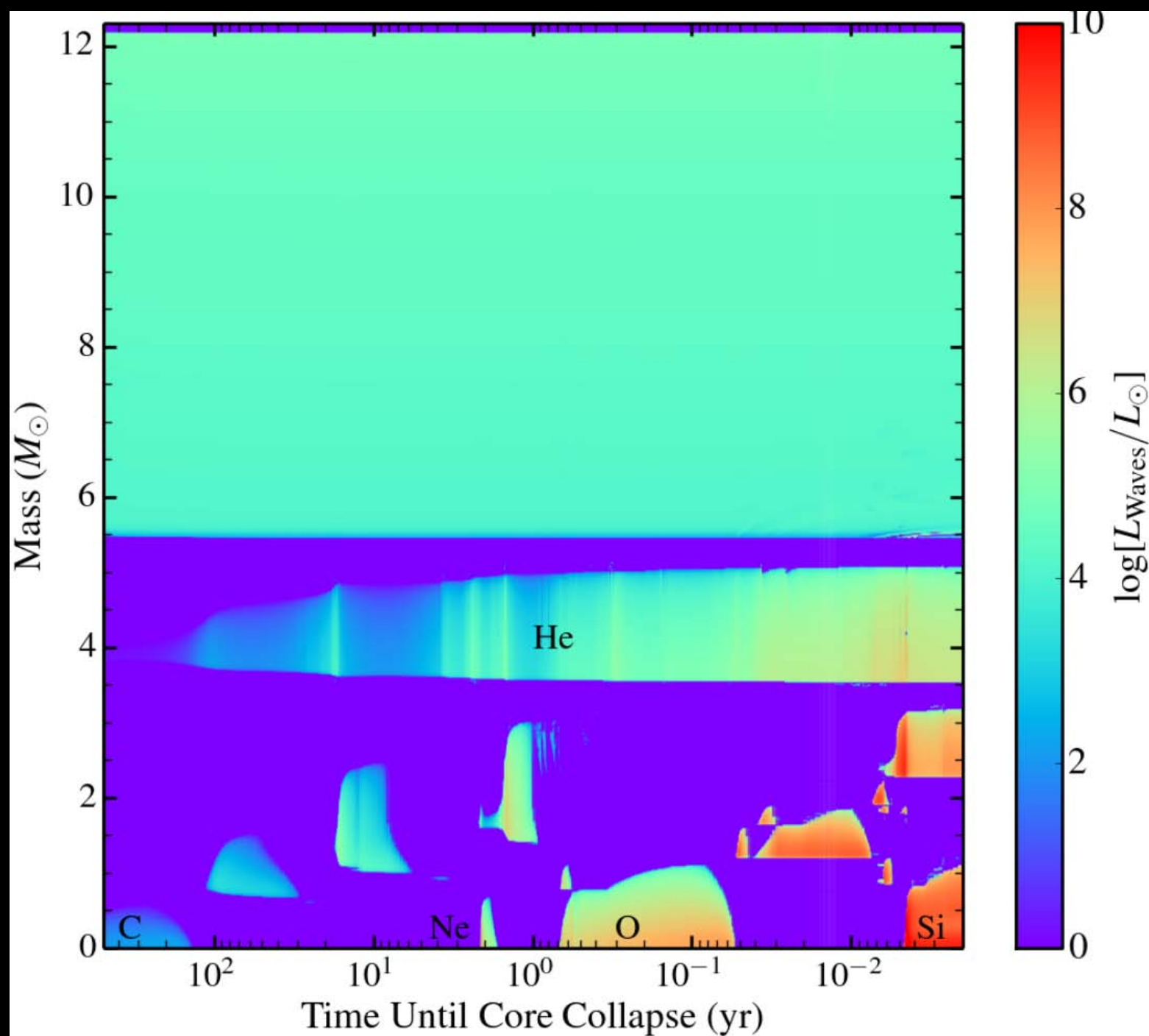
$$L_{\text{wave}} \sim \mathcal{M}_{\text{conv}} L_{\text{conv}} \sim 10^8 \left(\frac{L_{\text{conv}}}{10^{10} L_{\odot}} \right) \left(\frac{\mathcal{M}_{\text{conv}}}{0.01} \right) L_{\odot}$$

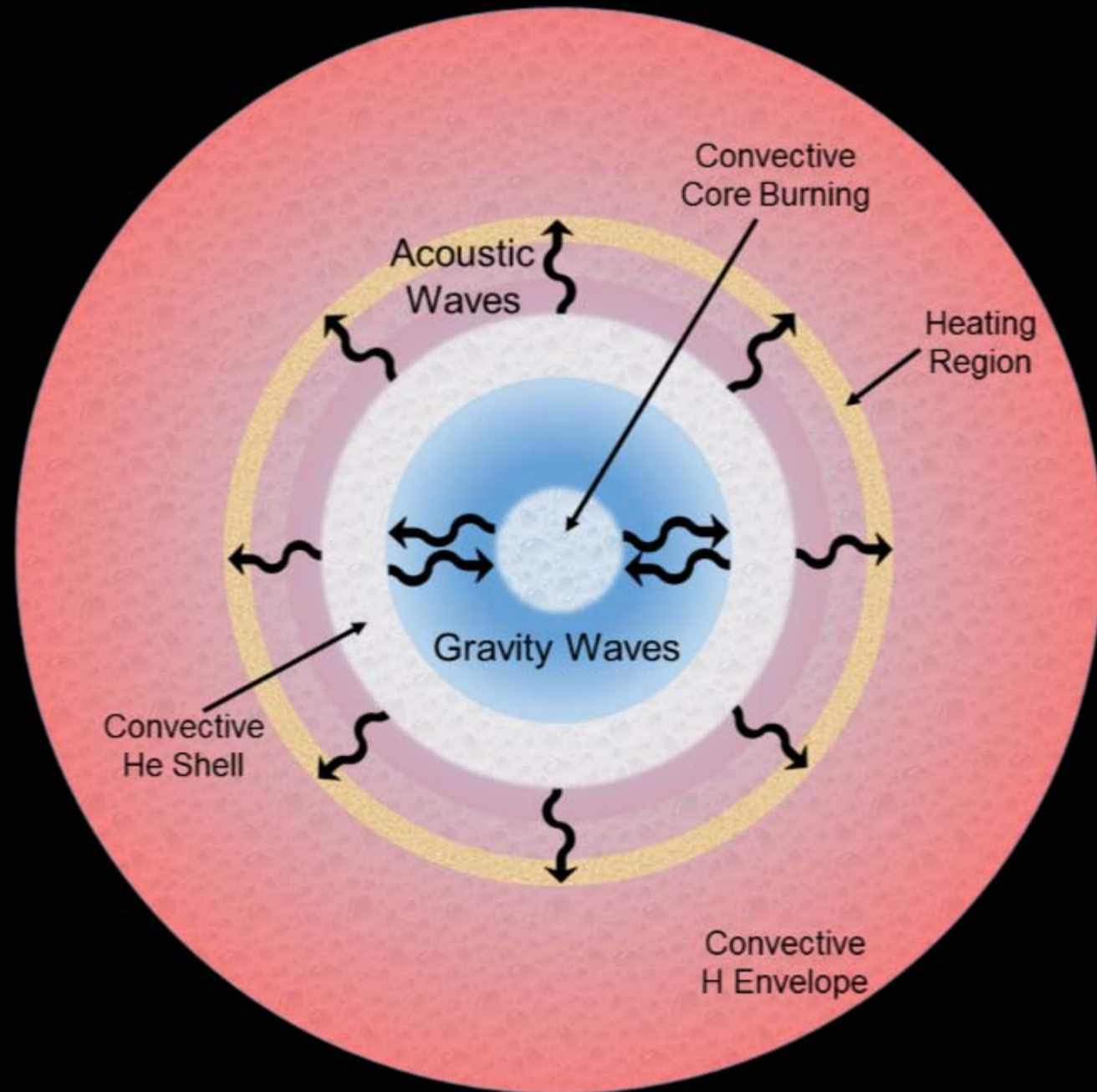
Table 1. Late stages of massive stellar evolution.

Stage	Duration (t_{nuc})	$L_{\text{fusion}} (L_{\odot})$	Mach ($\mathcal{M}_{\text{conv}}$)	τ_c (s)
Carbon	$\sim 10^3$ yr	$\sim 10^6$	~ 0.003	$\sim 10^{4.5}$
Neon	~ 1 yr	$\sim 10^9$	~ 0.01	$\sim 10^3$
Oxygen	~ 1 yr	$\sim 10^{10}$	~ 0.02	$\sim 10^3$
Silicon	~ 1 d	$\sim 10^{12}$	~ 0.05	$\sim 10^2$

$$E_{\text{waves}} \sim 10^{47-48} \text{ erg}$$

Quataert
& Shiode (2012)



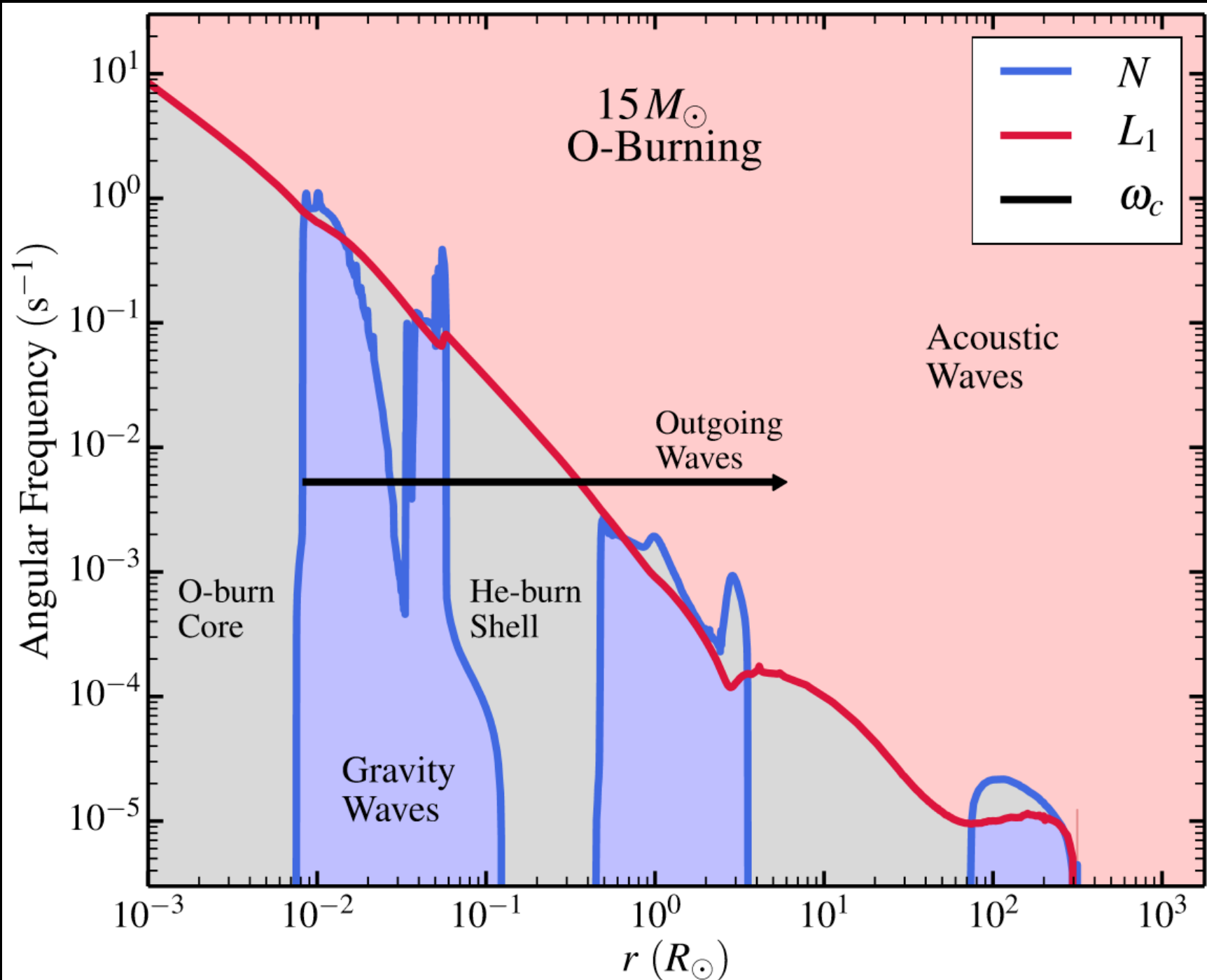


Wave Propagation

- Gravity excited in core must tunnel into stellar envelope as acoustic waves
- Acoustic waves damp in envelope, converting wave energy to thermal energy

$$v_{\text{con}} = \left[L_{\text{con}} / (4\pi\rho r^2) \right]^{1/3}$$

$$\omega_{\text{con}} = 2\pi \frac{v_{\text{con}}}{2\alpha_{\text{MLT}} H}$$



Methods

- Run MESA models including the effects of wave energy transport

$$\gamma_\nu = \frac{\delta \epsilon_\nu}{\epsilon} \simeq \frac{\Gamma_1^2 \nabla_{\text{ad}}^2 g^2}{N^2 c_s^4} \left(\frac{\partial \ln \epsilon_\nu}{\partial \ln T} \right)_\rho \epsilon_\nu$$

- At each time step, compute
 - Wave generation by nuclear burning convective zones in core
 - Wave propagation and fraction of energy tunneling into the envelope
 - Wave damping energy deposition rate per unit mass at each shell within star

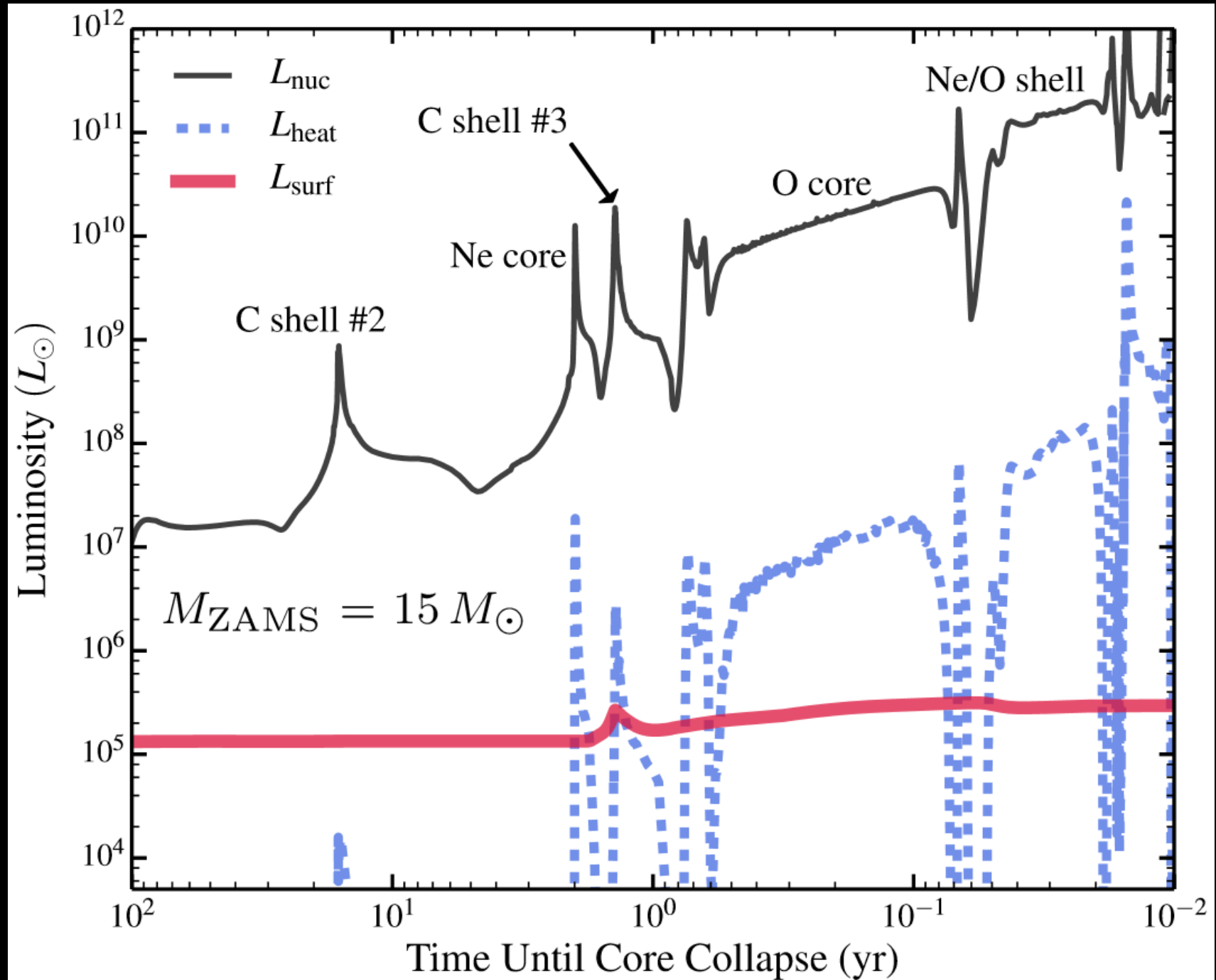
$$L_{\text{heat}} = f_{\text{esc}} L_{\text{wave}} = \left[1 + \frac{T_{\text{shell}}^2 + x_\nu}{T_{\text{min}}^2} \right]^{-1} L_{\text{wave}}$$

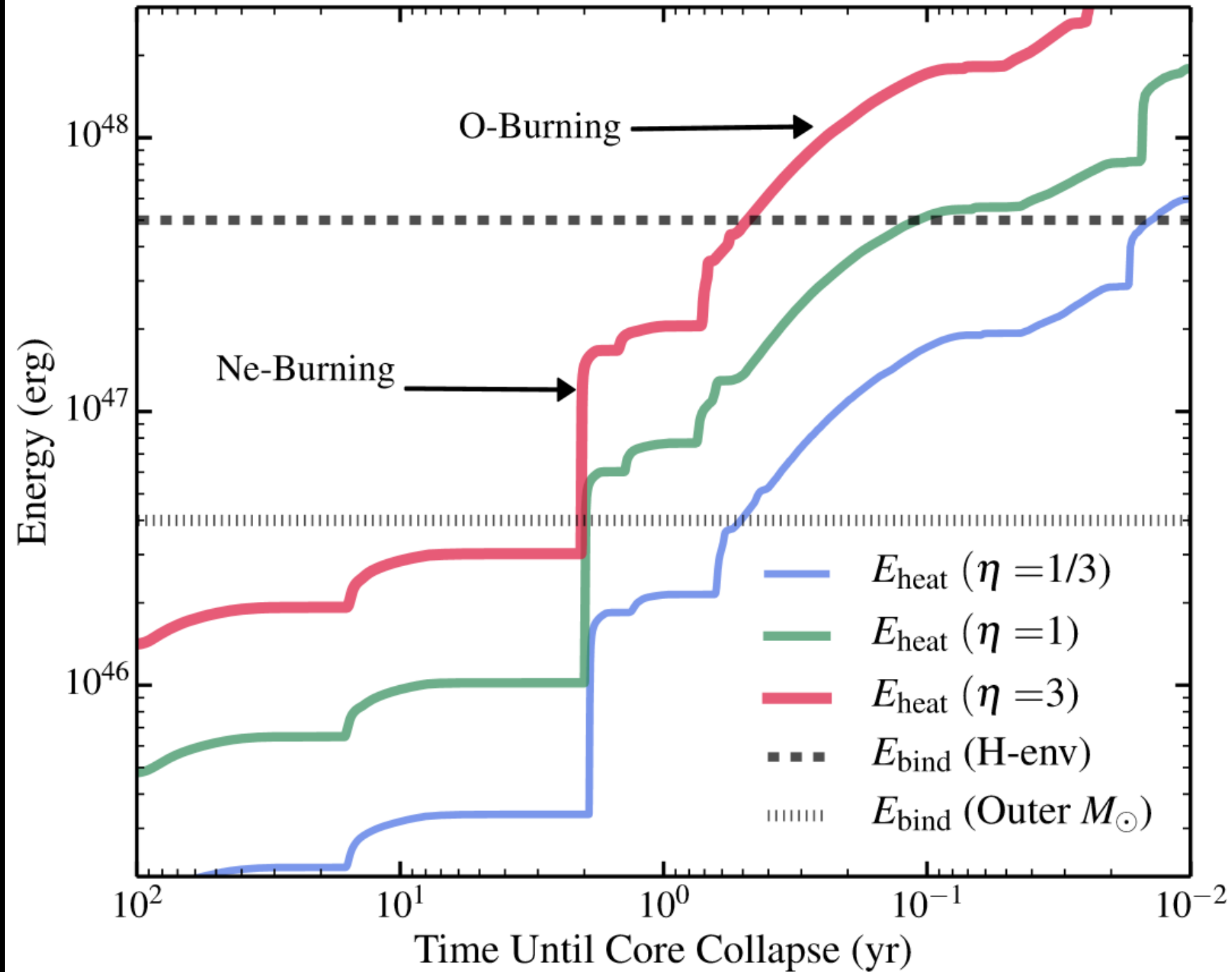
$$T_{1,2}^2 = \exp \left(-2 \int_{r_1}^{r_2} |k_r| dr \right)$$

- Use hydrodynamic version of MESA to allow for high velocity flows

$$k_r^2 = \frac{(N^2 - \omega^2)(L_l^2 - \omega^2)}{\omega^2 c_s^2}$$

Wave Power



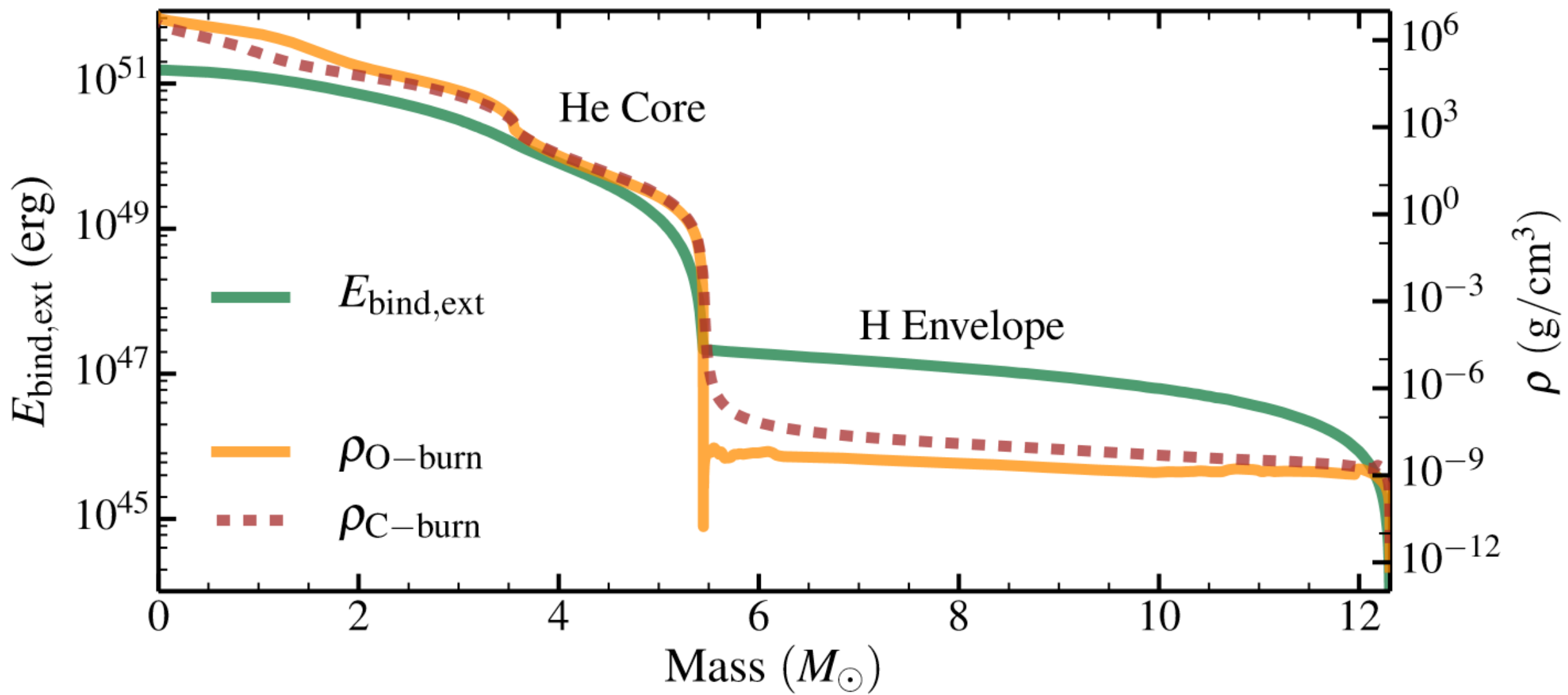


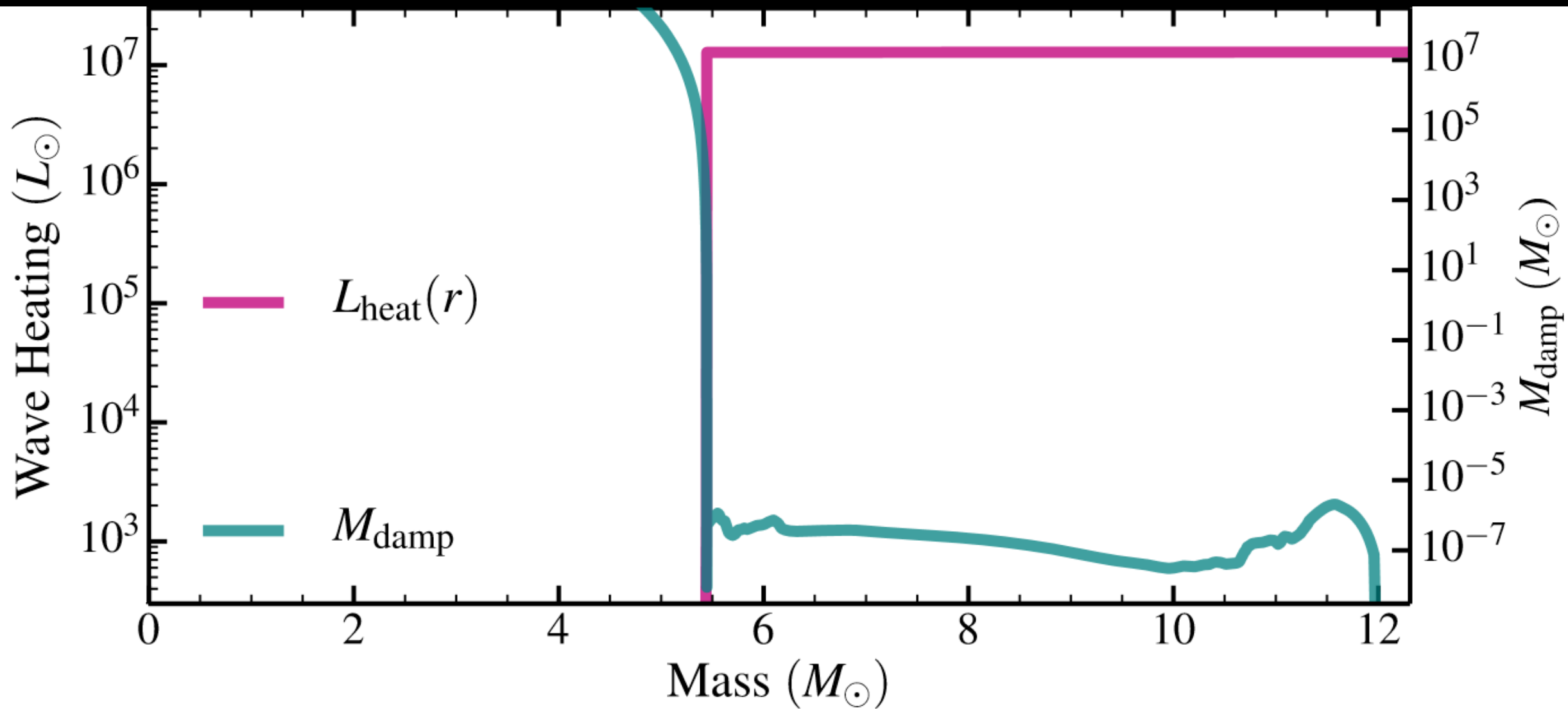
More Methods

- Run MESA models including the effects of wave energy transport
- At each time step, compute
 - Wave generation by nuclear burning convective zones in core
 - Wave propagation and fraction of energy tunneling into the envelope
 - Wave damping energy deposition rate per unit mass at each shell within star
- Use hydrodynamic version of MESA to allow for high velocity flows

$$\frac{dL_{\text{wave}}}{dM} = -\frac{L_{\text{wave}}}{M_{\text{damp}}}$$

$$M_{\text{damp}} = \frac{4\pi\rho r^2 c_s^3}{\omega^2 K} = \frac{3\pi\rho^3 r^2 c_s^3 c_p \kappa}{4\sigma_{\text{SB}}\omega^2 T^3}$$

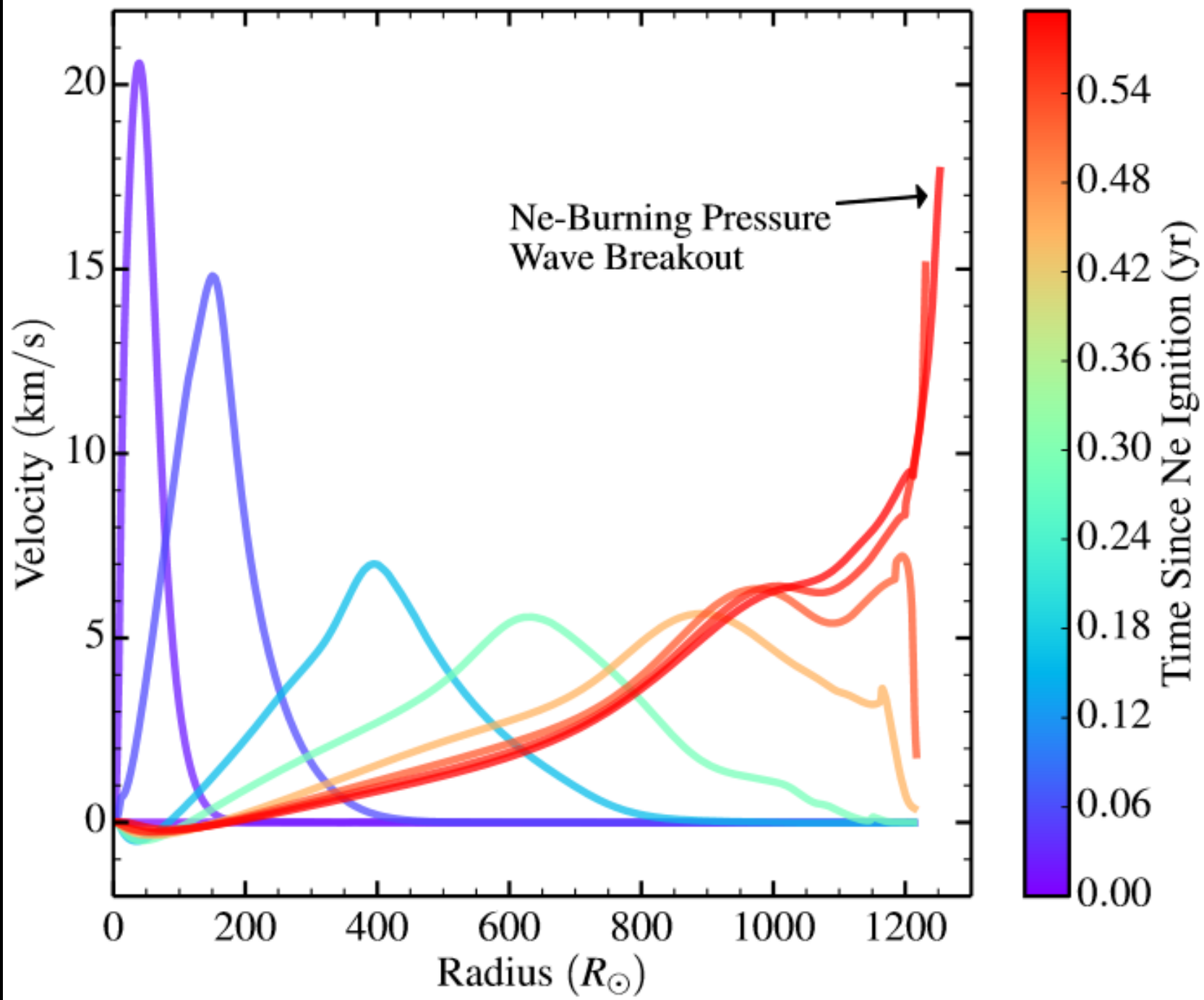


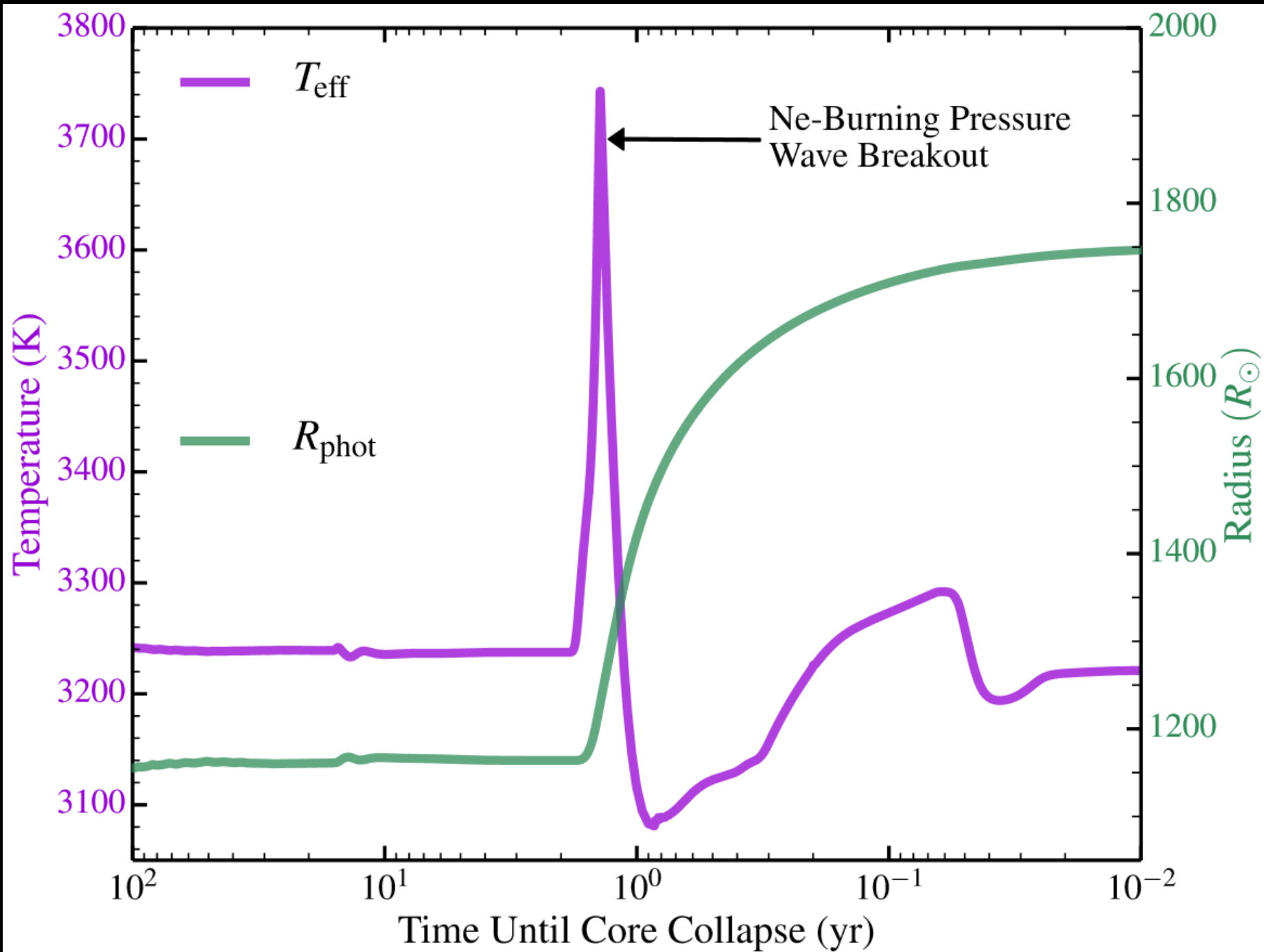


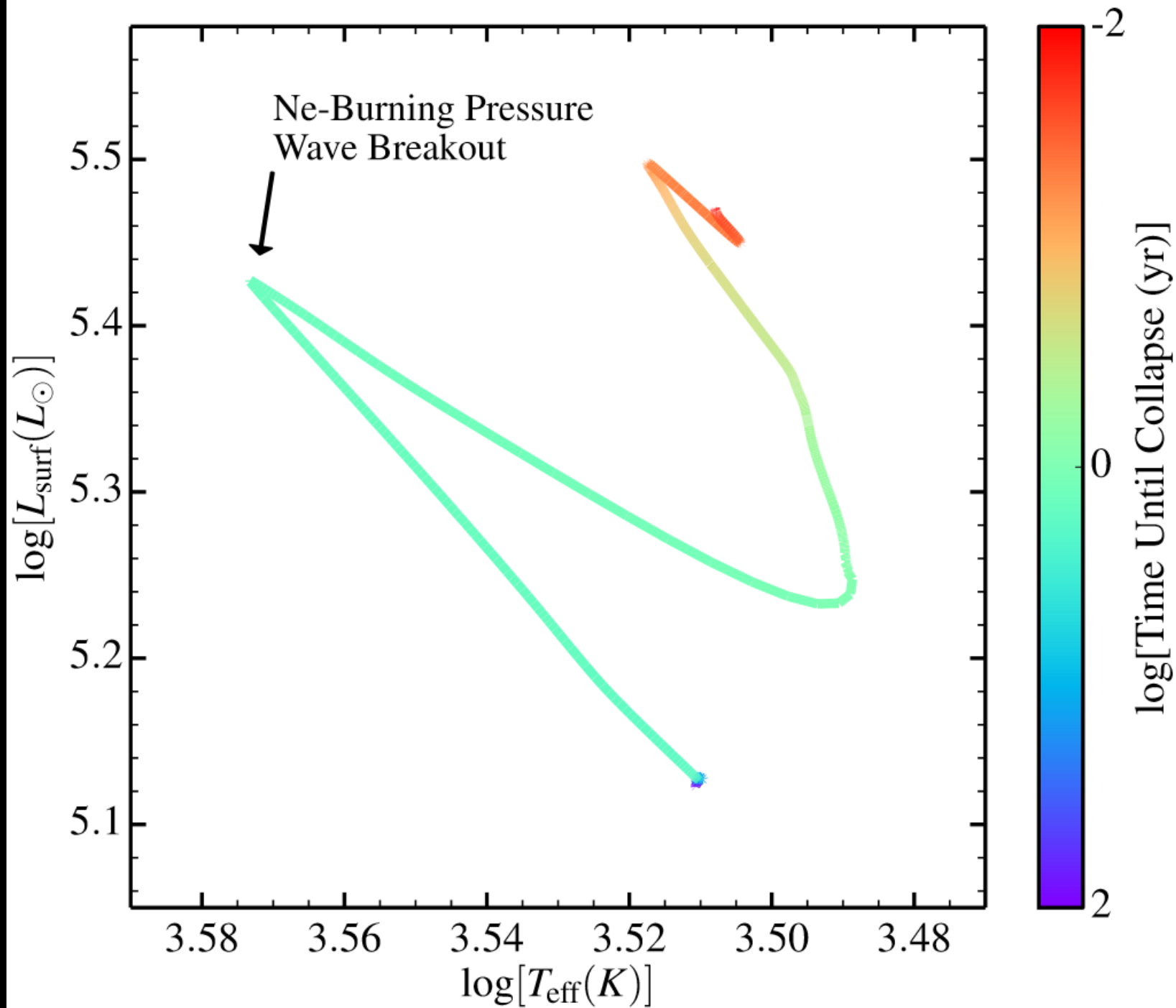
15 M_{sun}
Red
Supergiant

4/25/2017

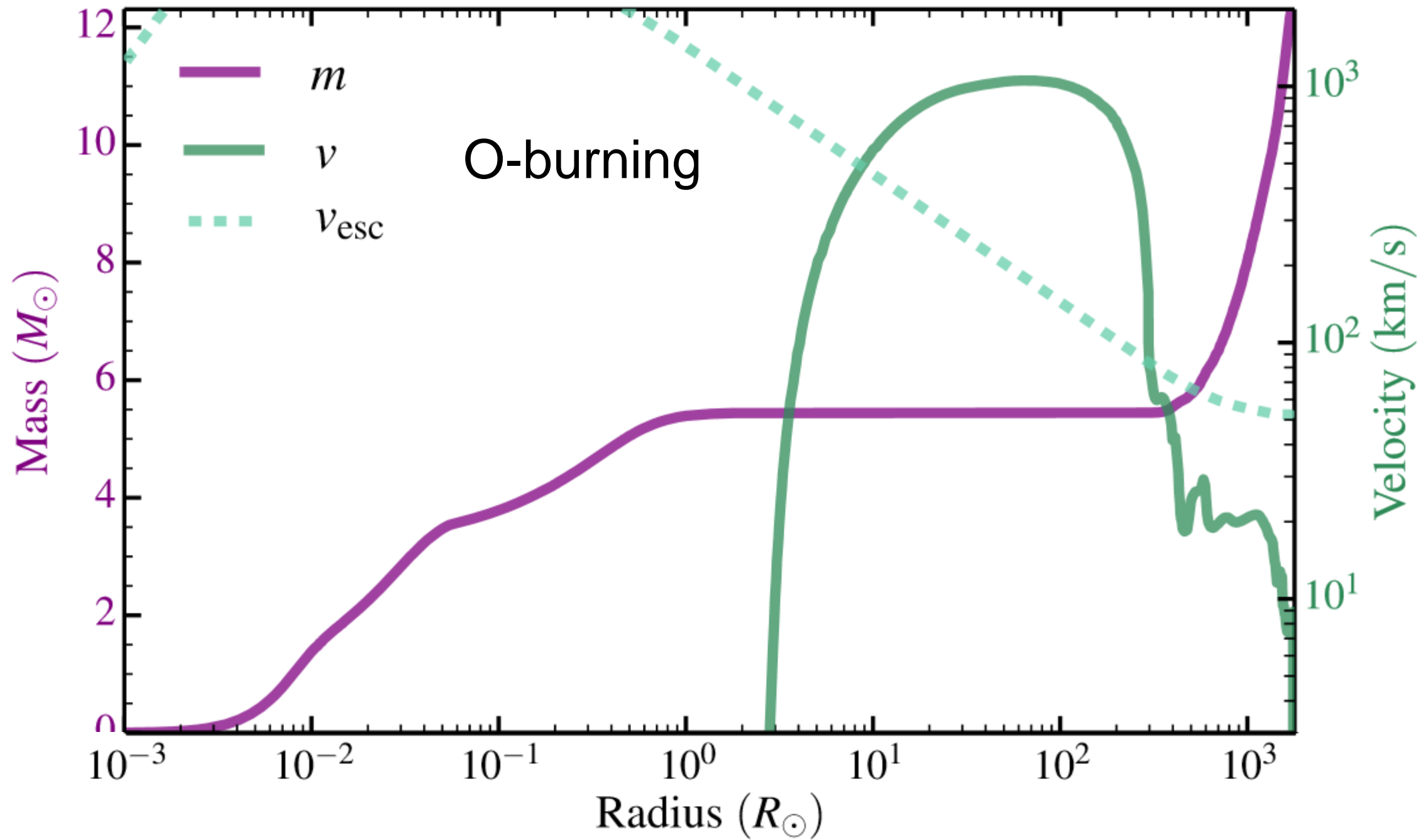




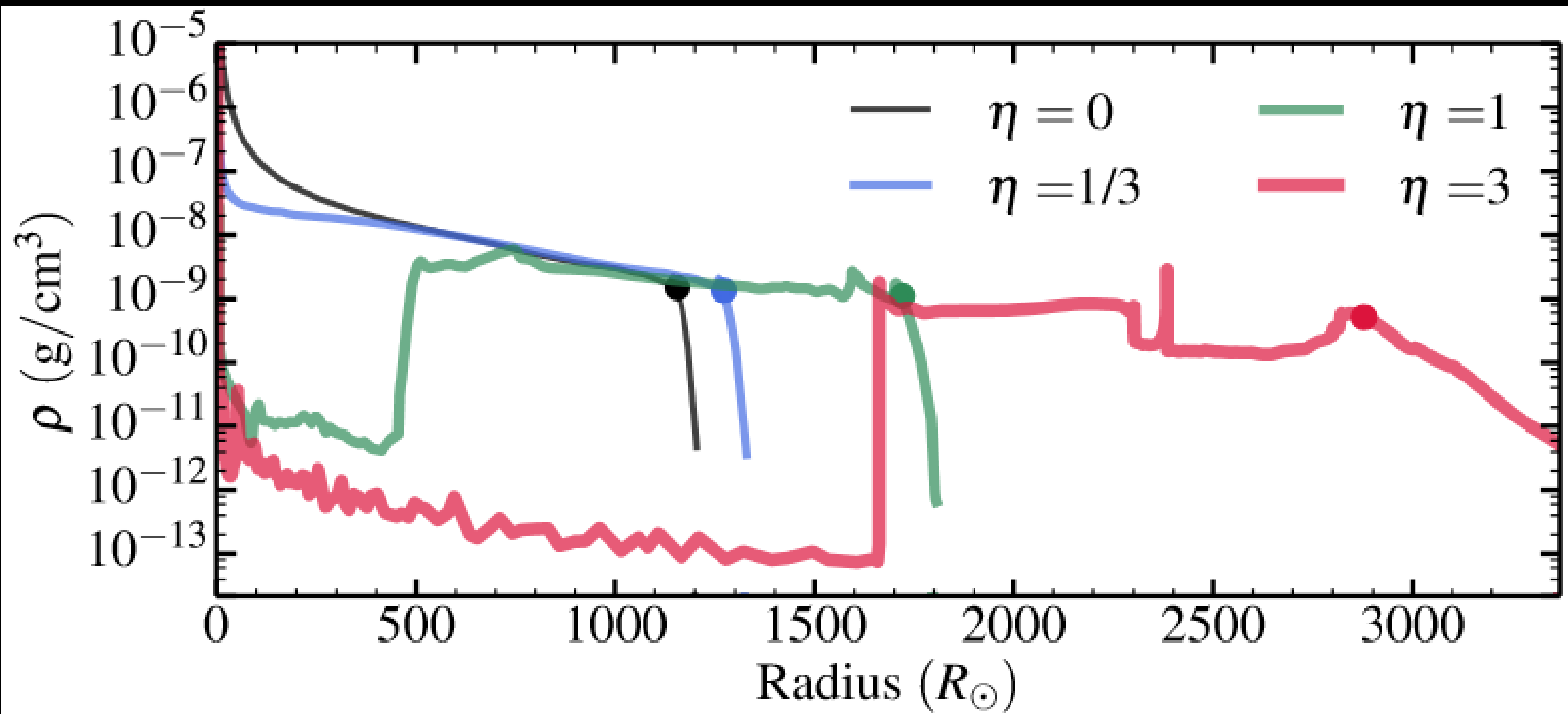




Brightening/dimming
by factor of a few in
last couple years

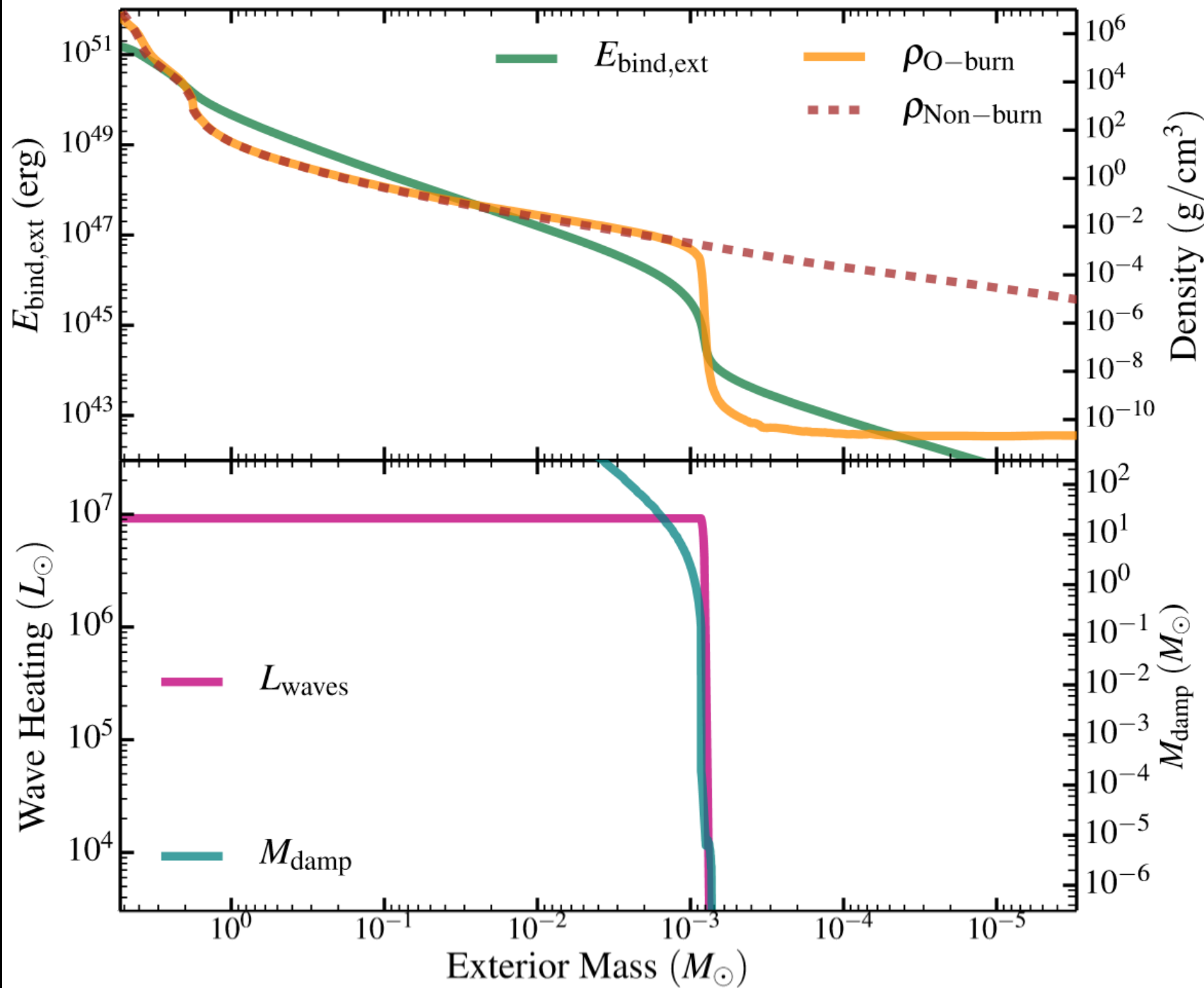


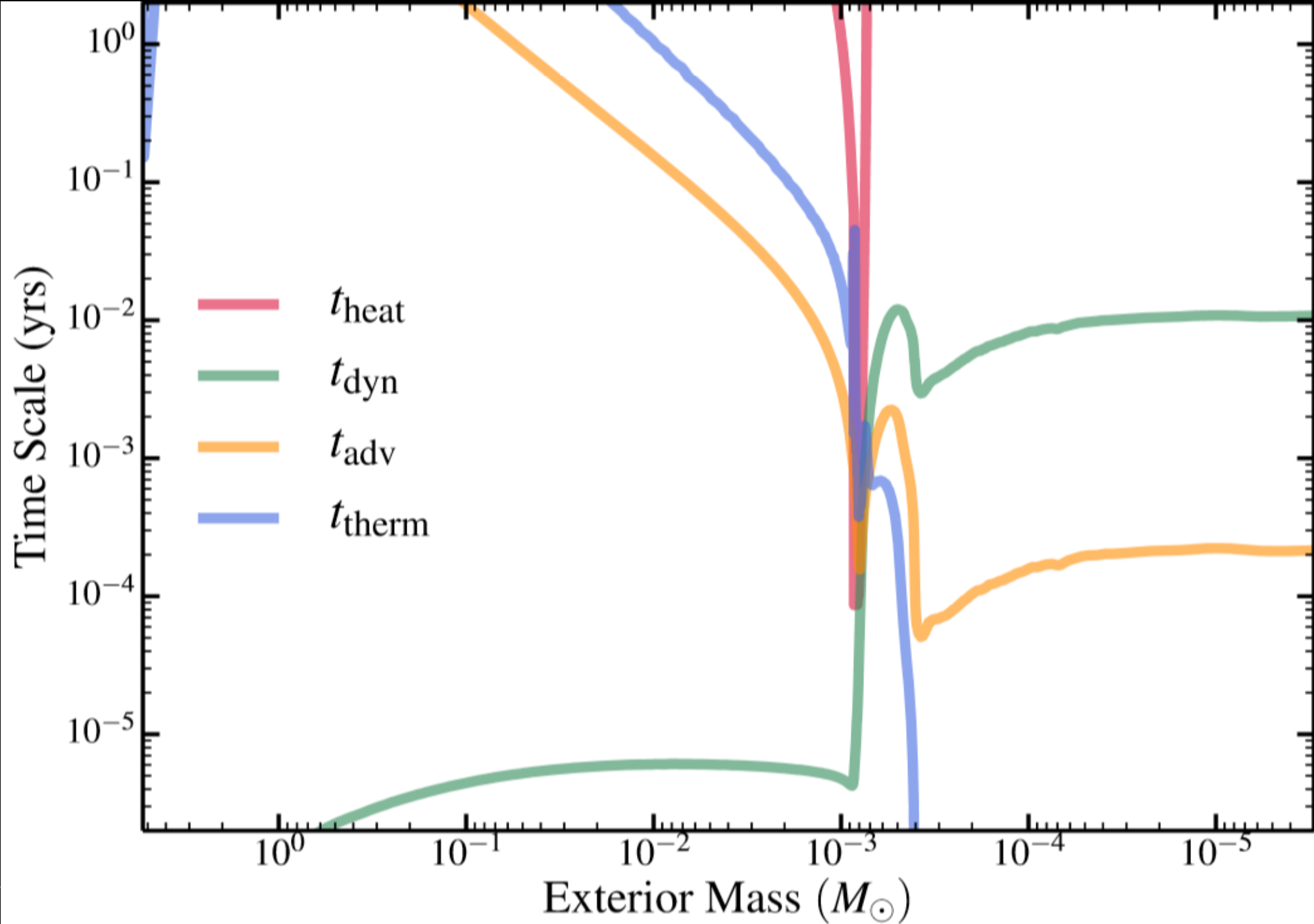




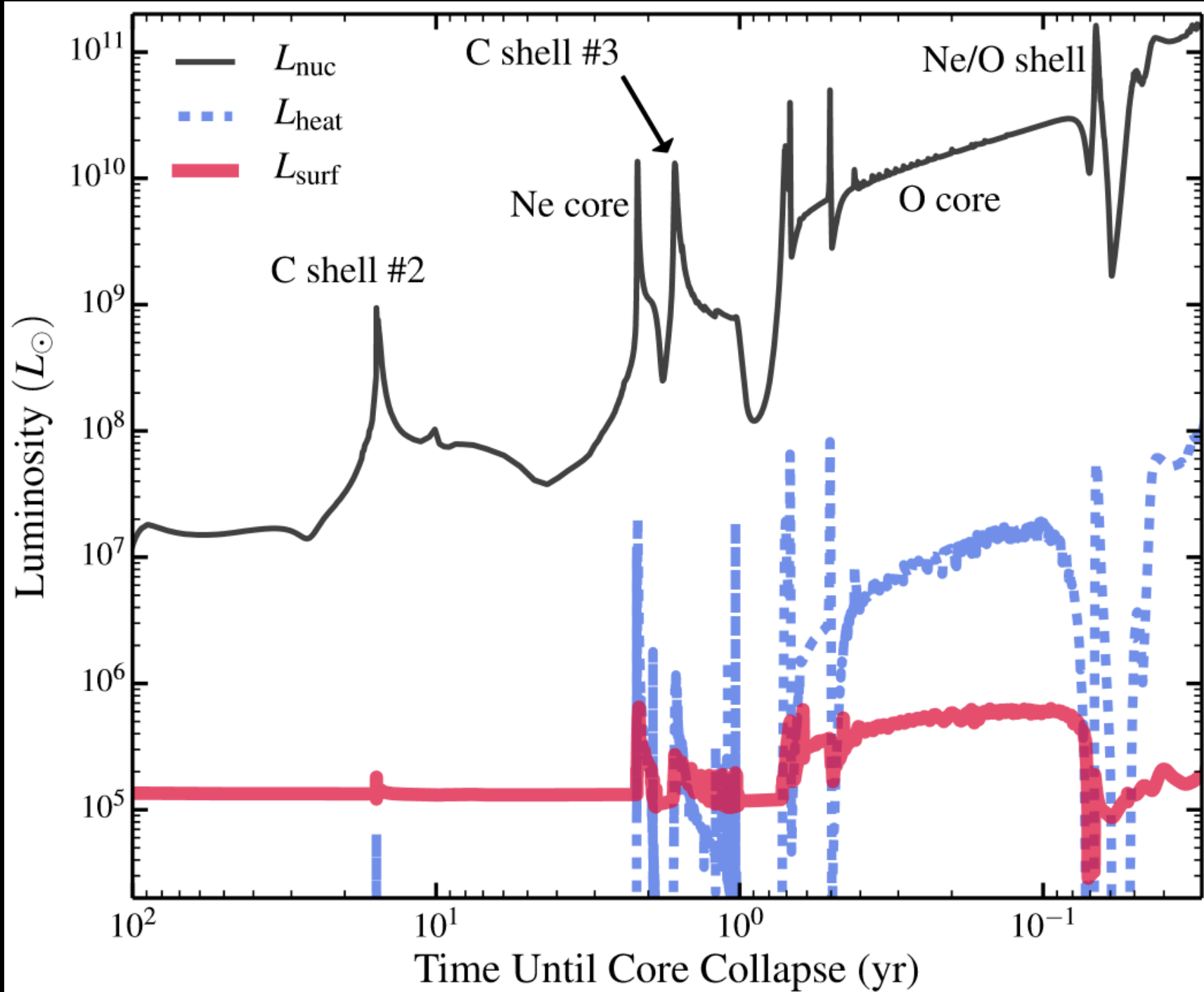
Stripped Stars

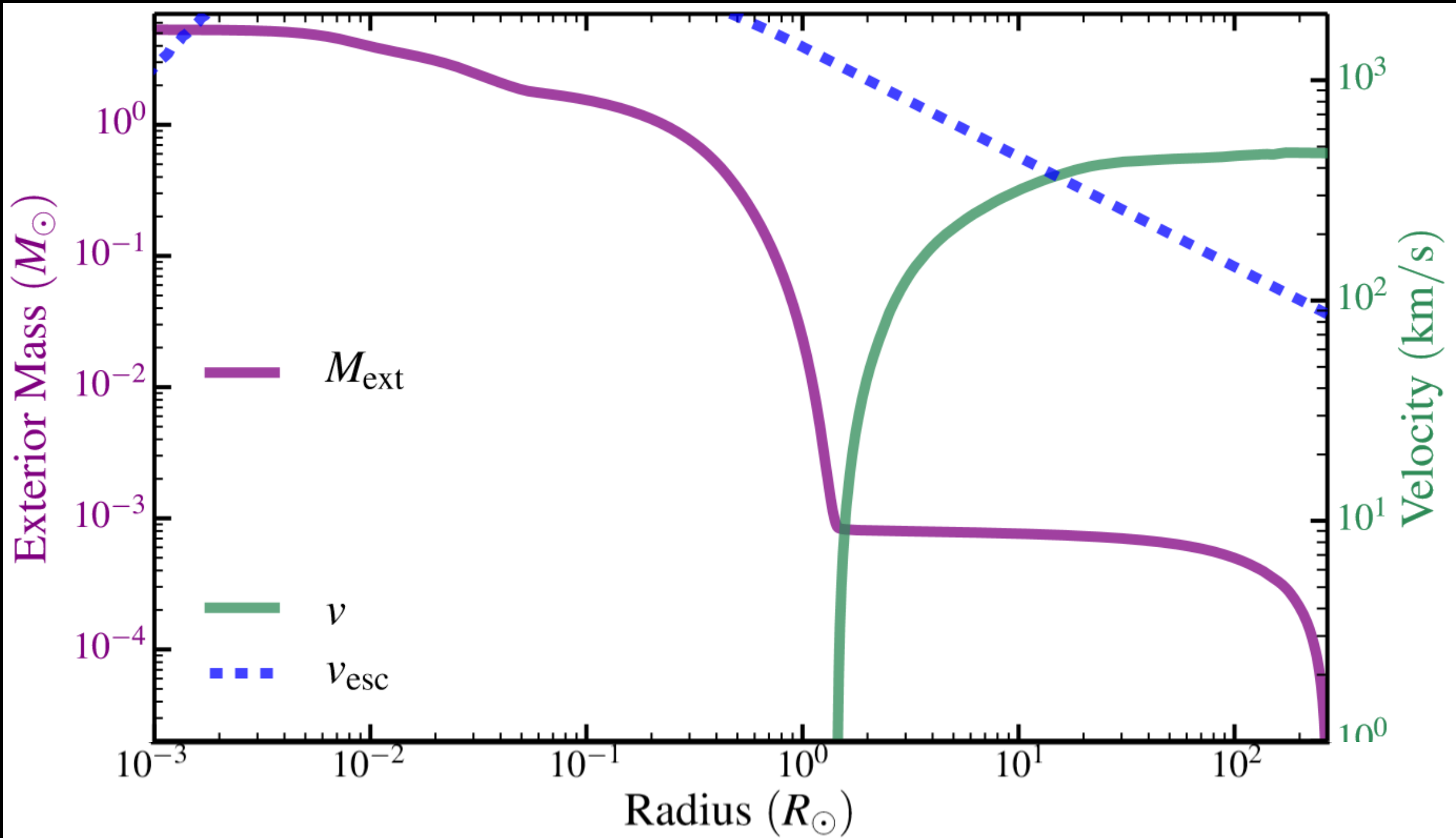
- No thick hydrogen envelope
- Waves still damp at edge of helium core, but now there is very little mass above wave heating region
- Wave heat unbinds atmosphere and launches continuum-driven super-Eddington wind

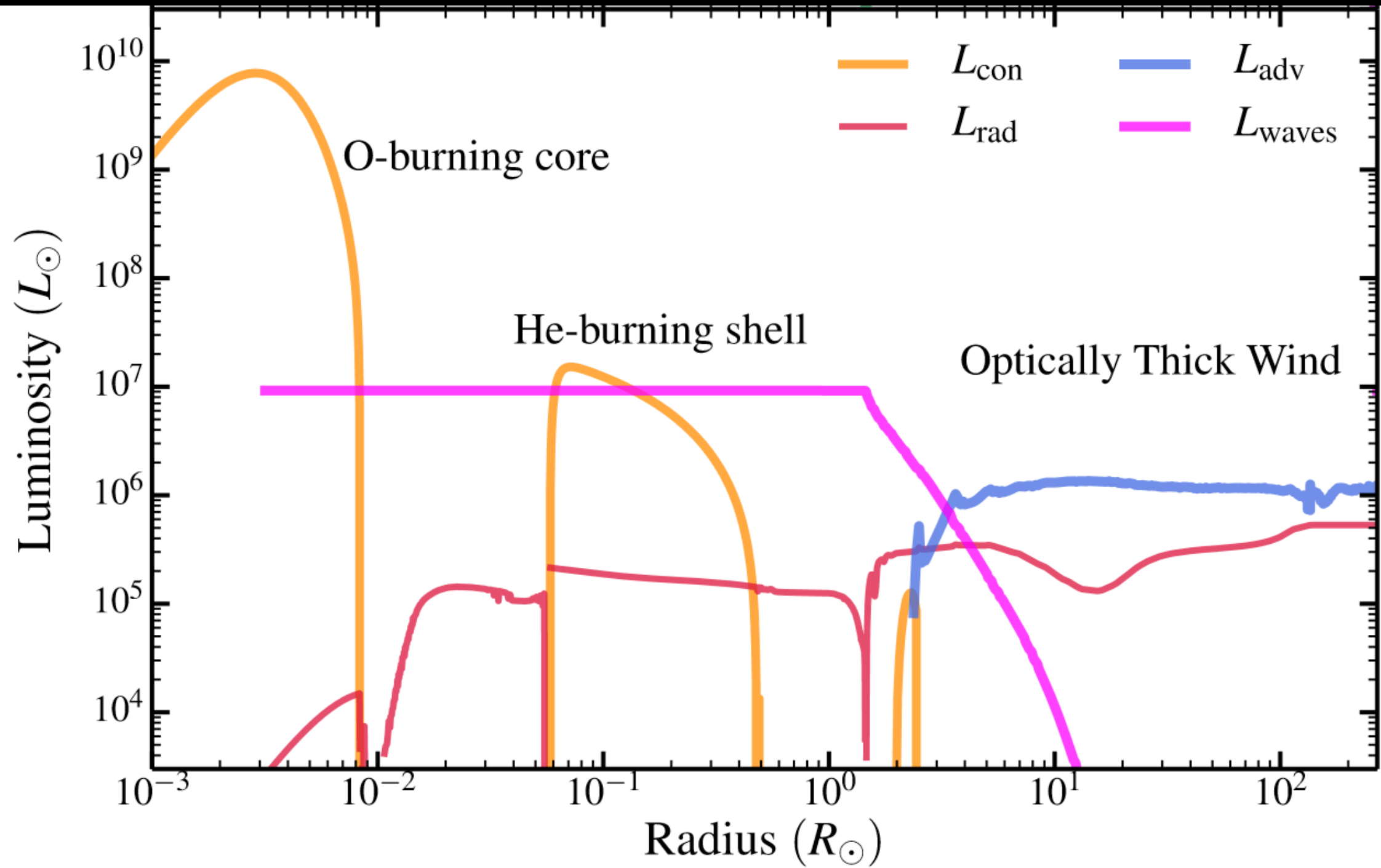


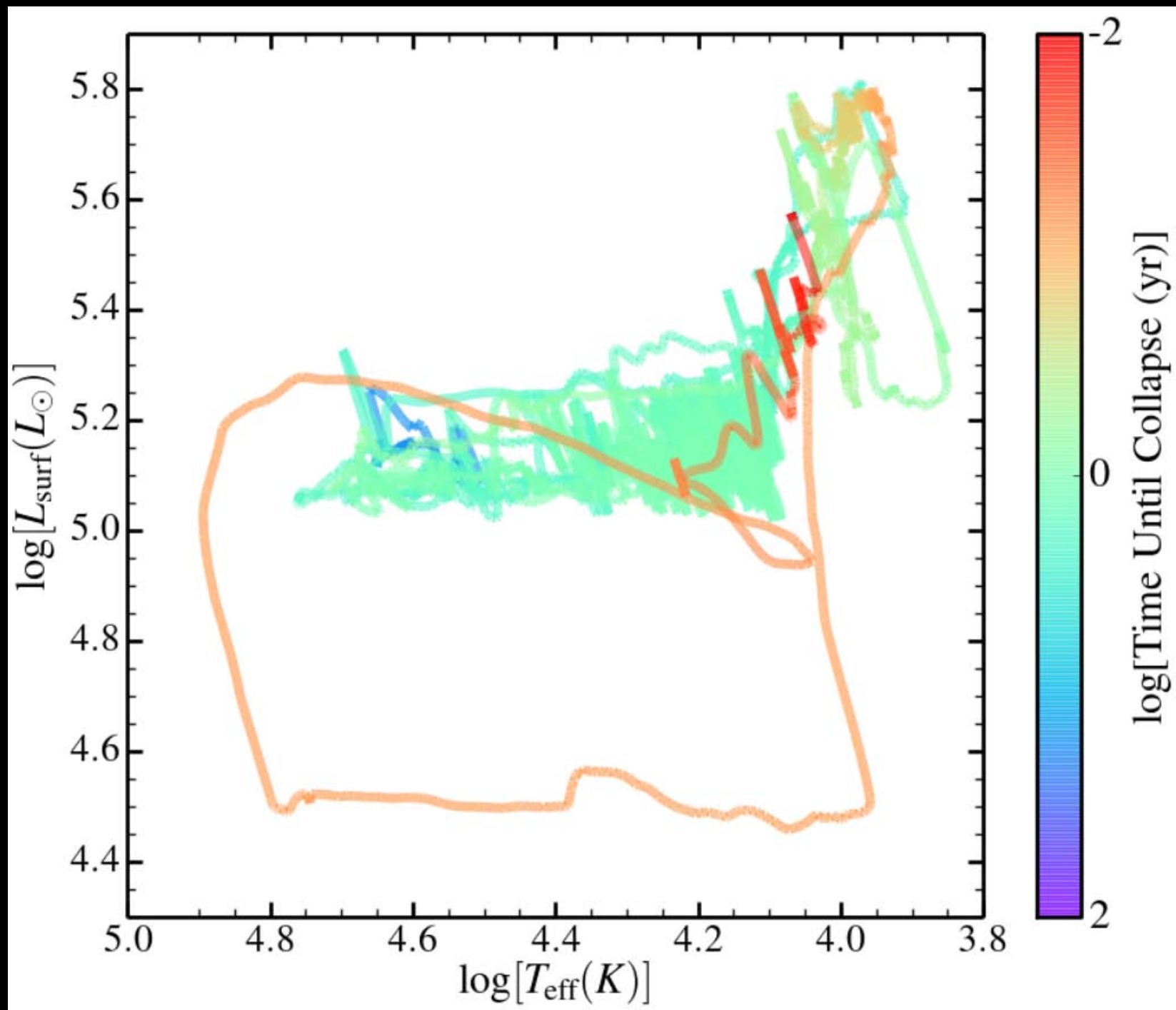


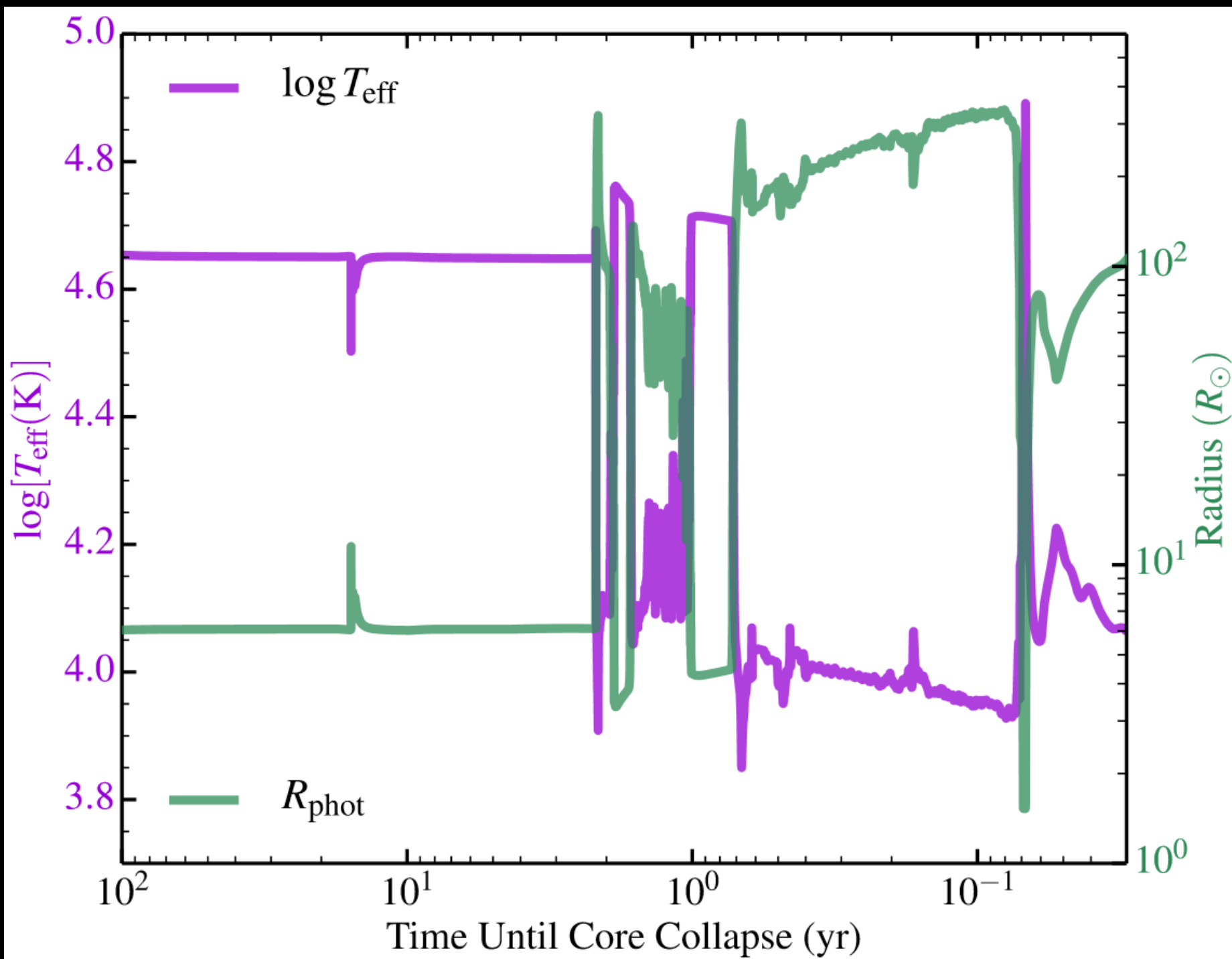


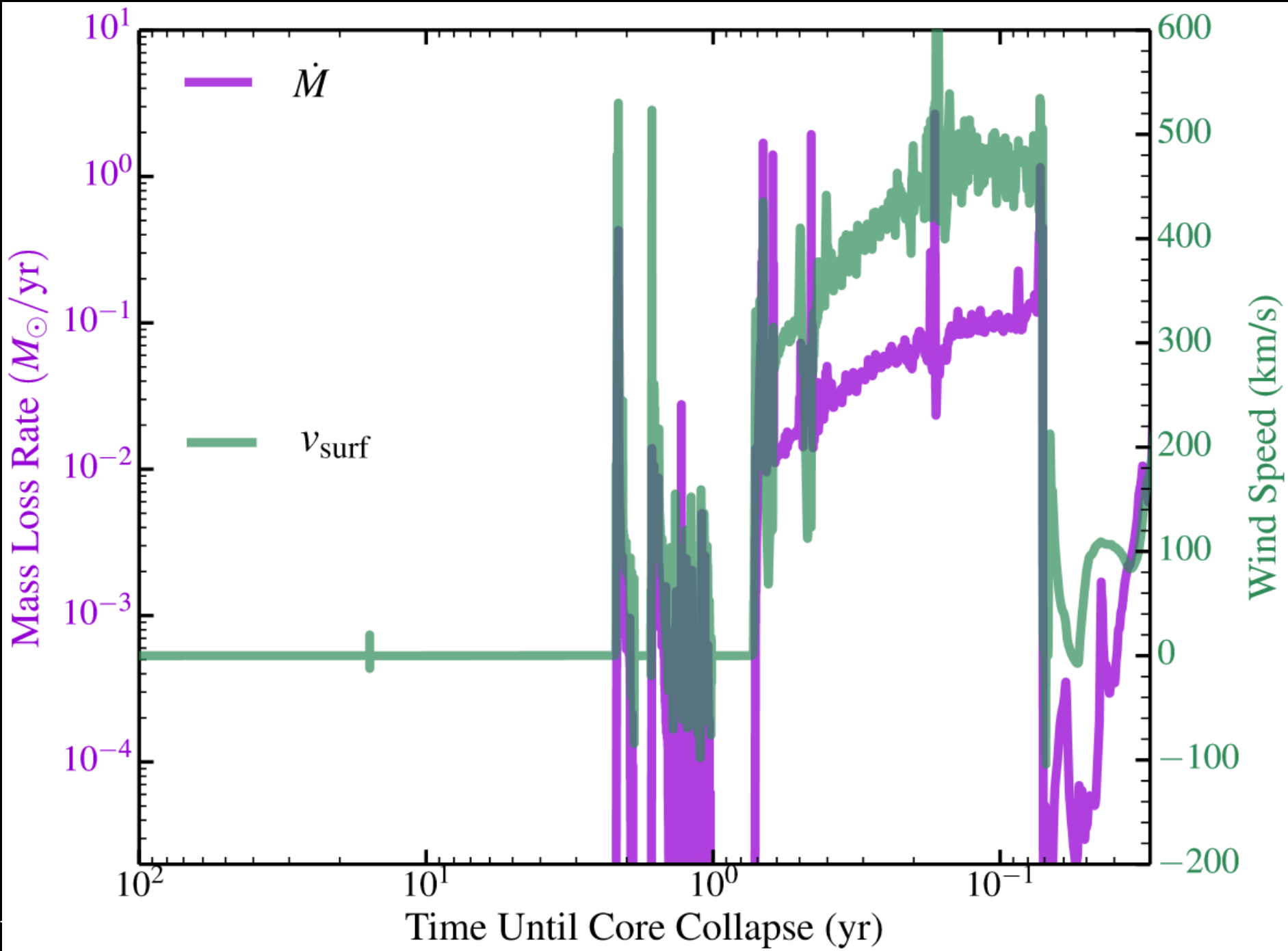












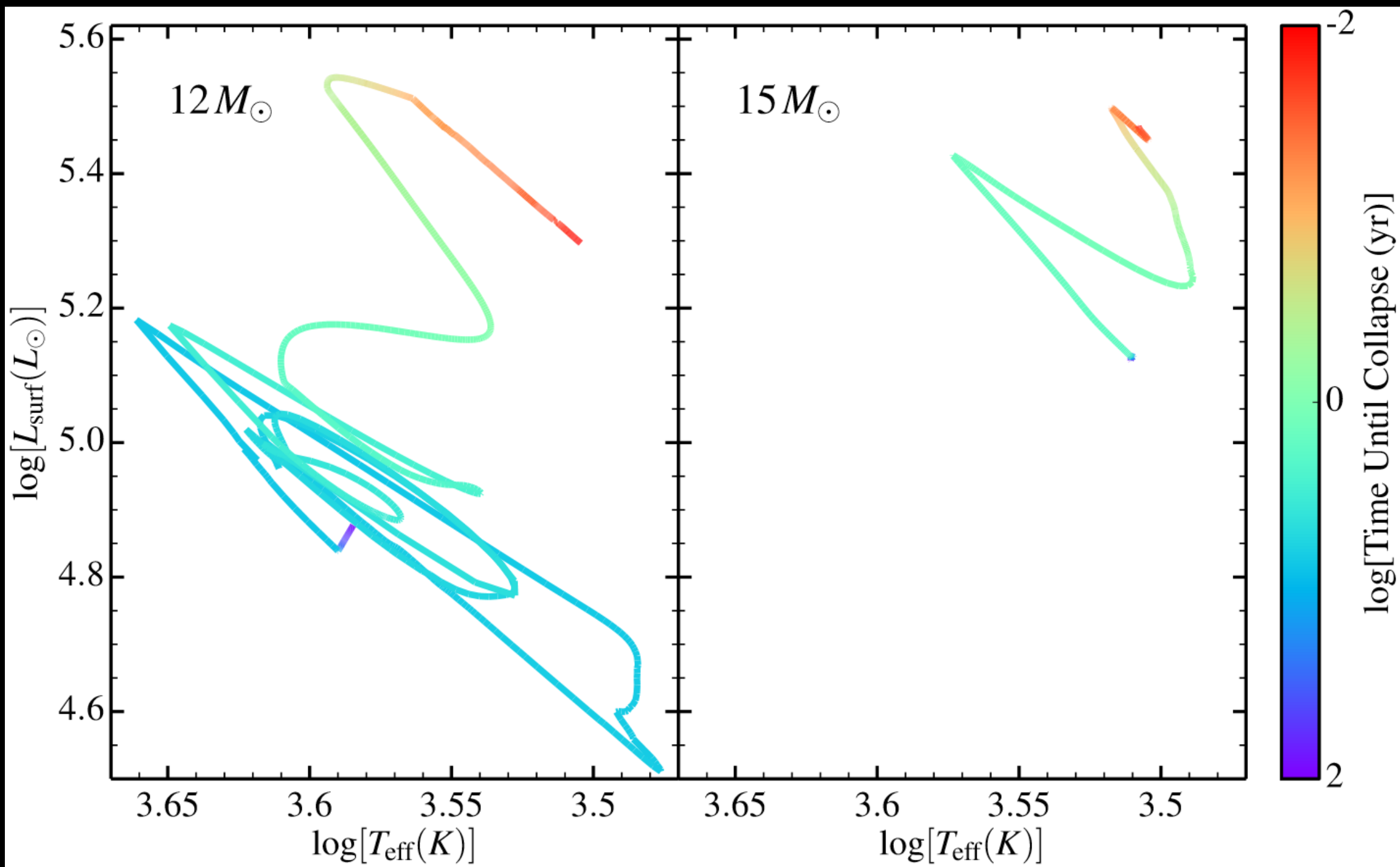
Conclusions and Discussion

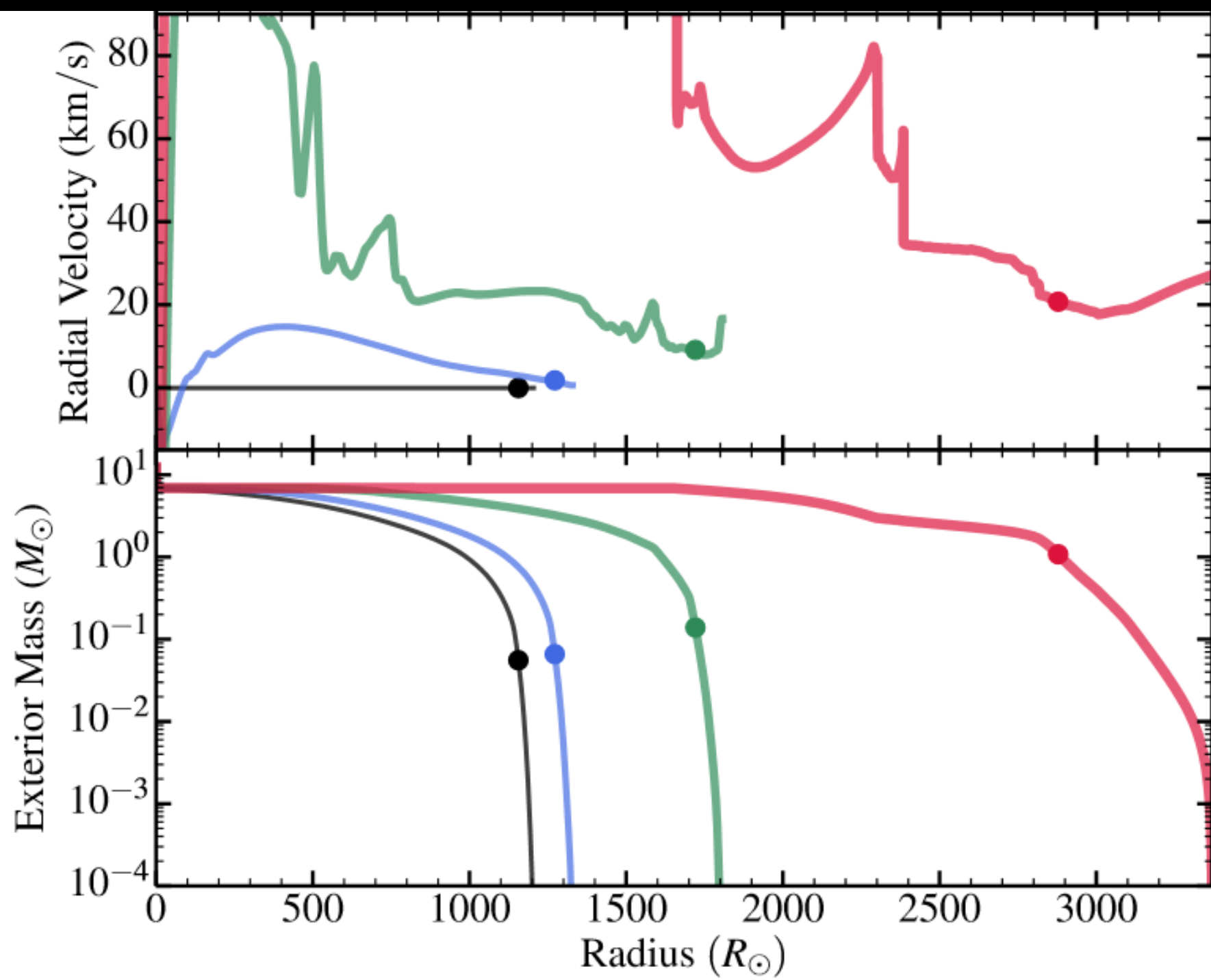
- Wave heating in RSGs creating Type II SNe can cause mild outbursts, enhanced mass loss (up to $\sim 1 M_{\text{sun}}$) in final years before SNe
 - Wave heating unlikely to lead to very luminous Type IIn SNe
- Wave heating is very good candidate to create flash-ionized SNe
- Heating affects envelope density structure, likely contributes to type II SNe diversity
- Unlikely to substantially alter core structure (binding energy $\sim 10^{51}$ erg)

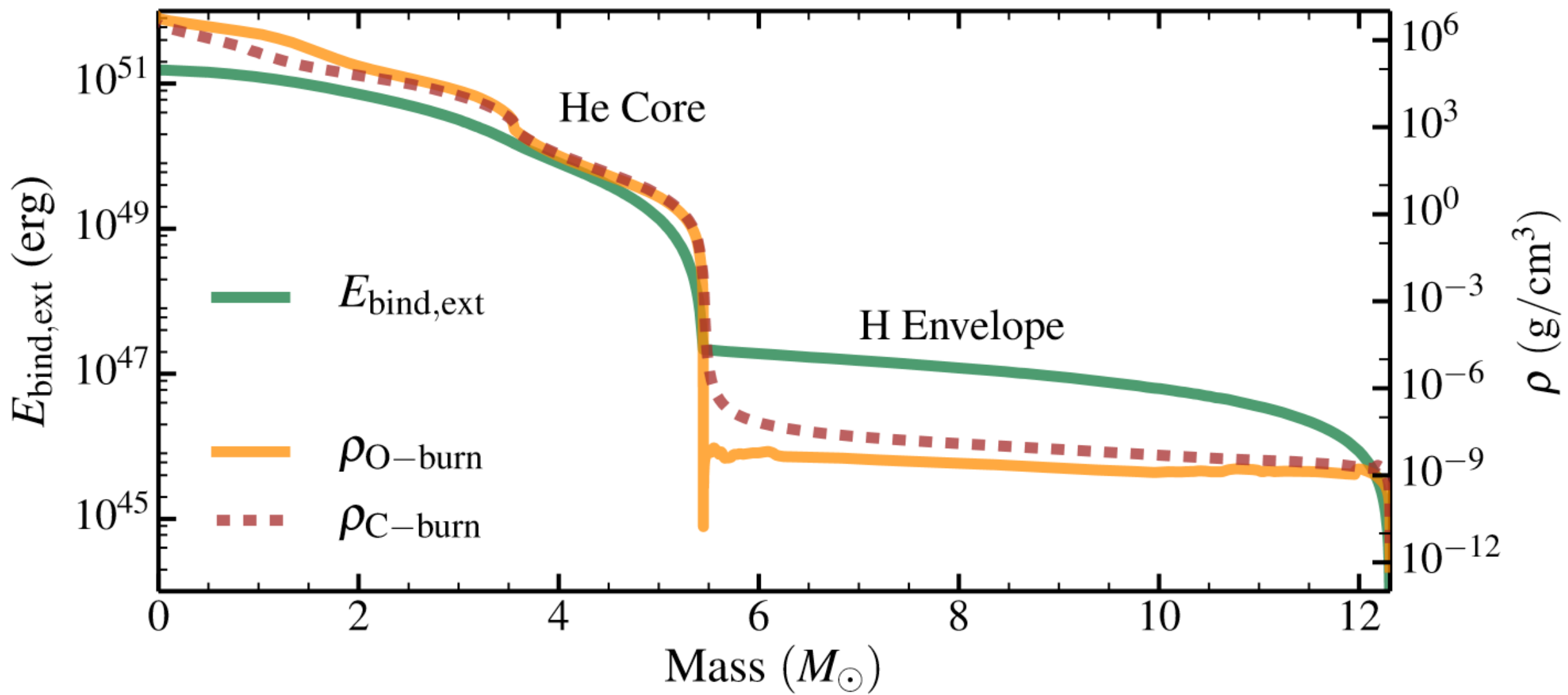
Bonus Material!

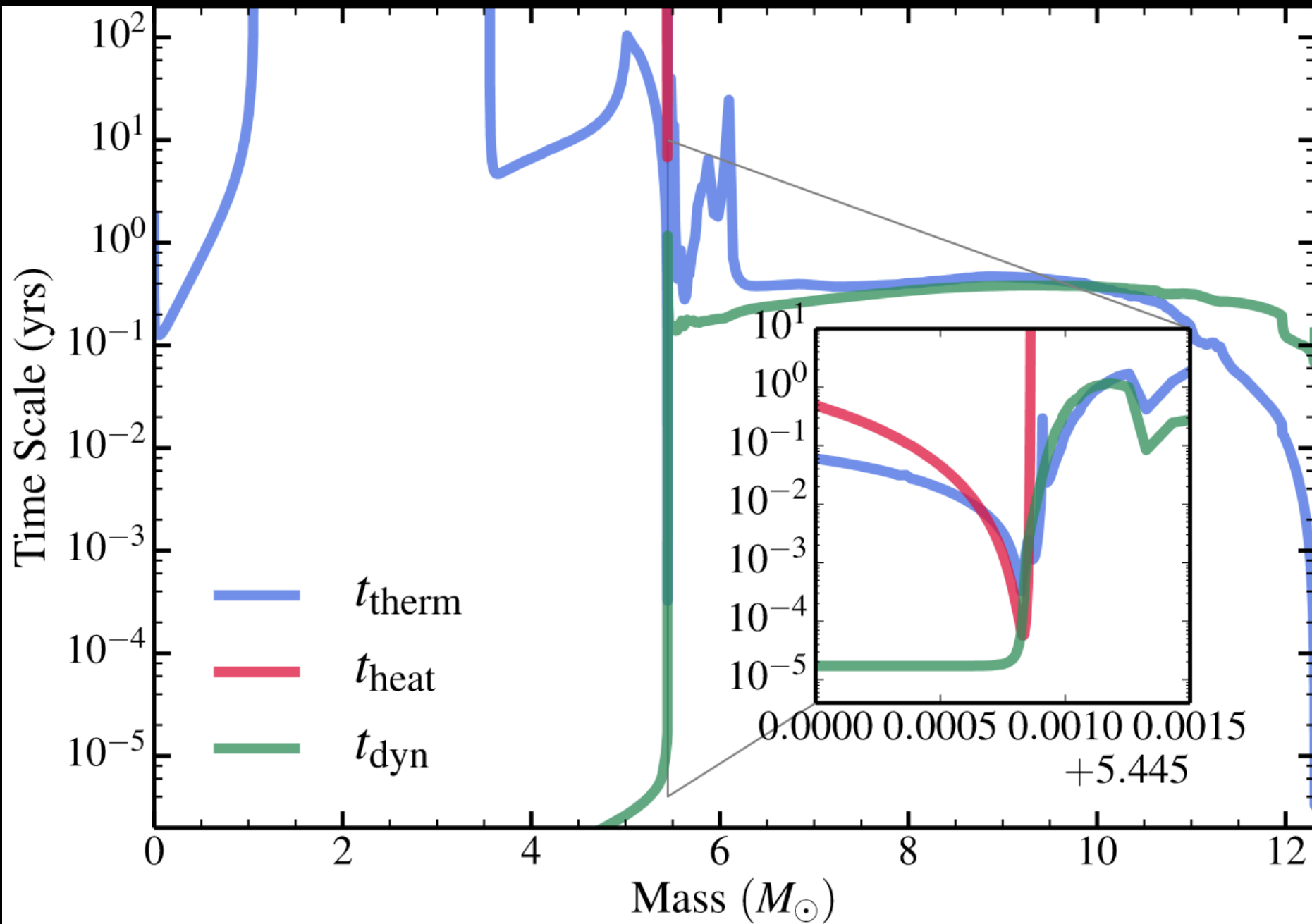
Caveats

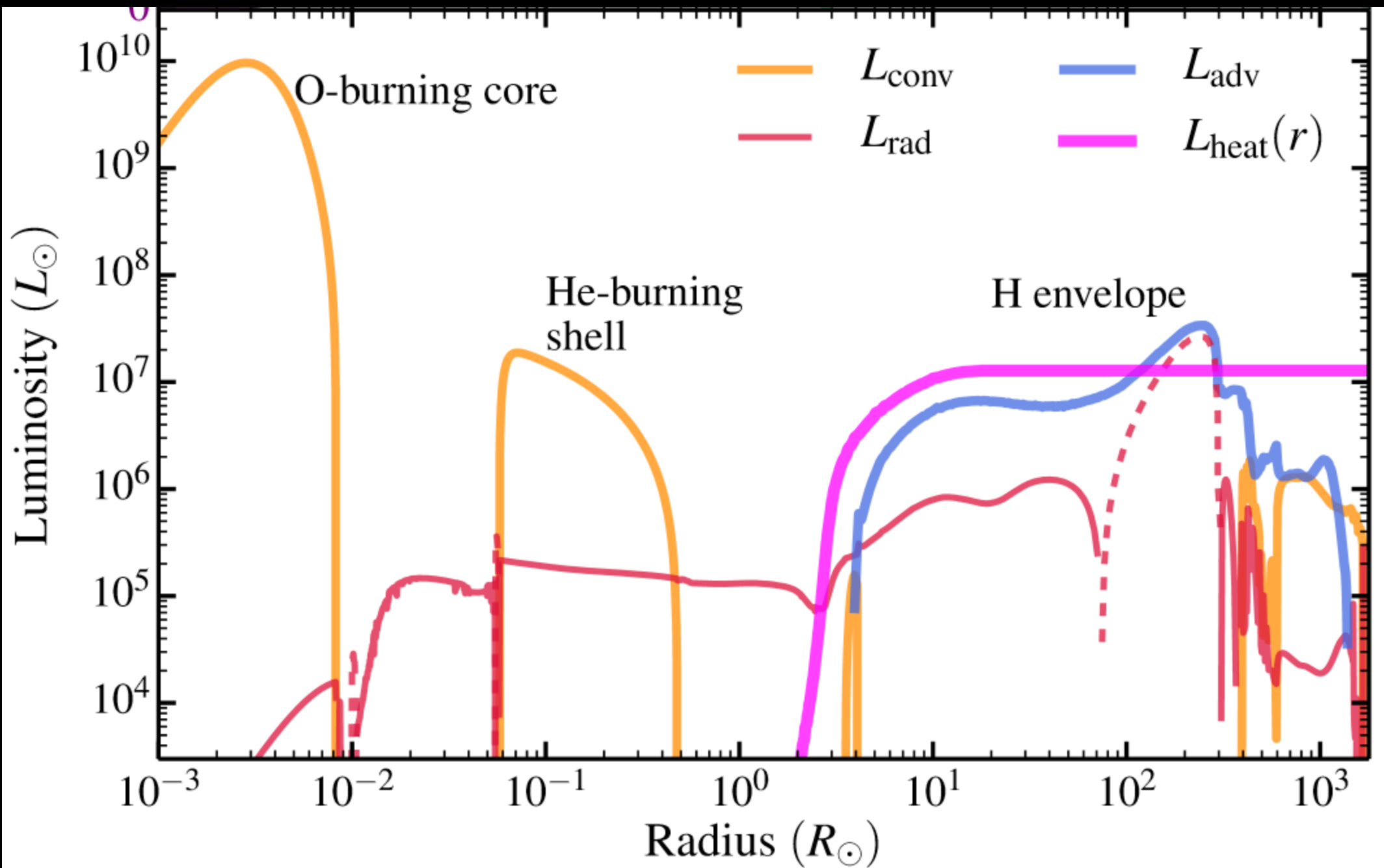
- Binaries, rotation, and magnetic fields
 - Radial expansion may trigger binary interaction → II in supernova (McCleary & Soker 2014)
 - Core magnetic fields can inhibit gravity wave propagation (talk by Cantiello)
- Hard to explain IIbs, small variability of 2011dh (Szczygieł et al. 2012, Hosseinzadeh et al. 2015), small radius/mass loss of iPTF2013bvn (Cao et al. 2013)
- Progenitor landscape is complex
- Burning instabilities also possible (e.g., Meakin & Arnett 2006, 2007, 2016, Smith & Arnett 2014, Woosley & Heger 2015, Soker & Gilkis 2016,)

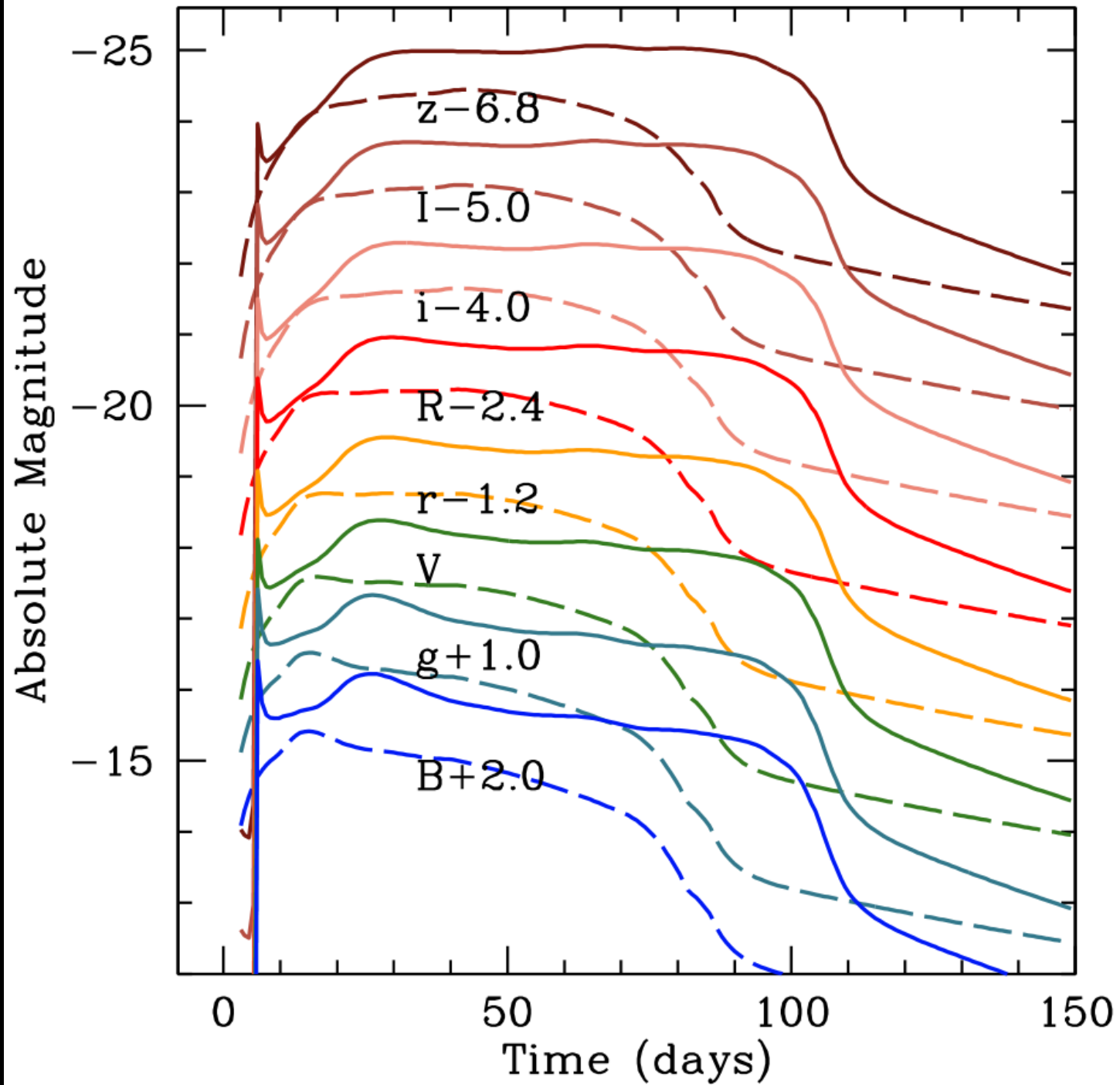


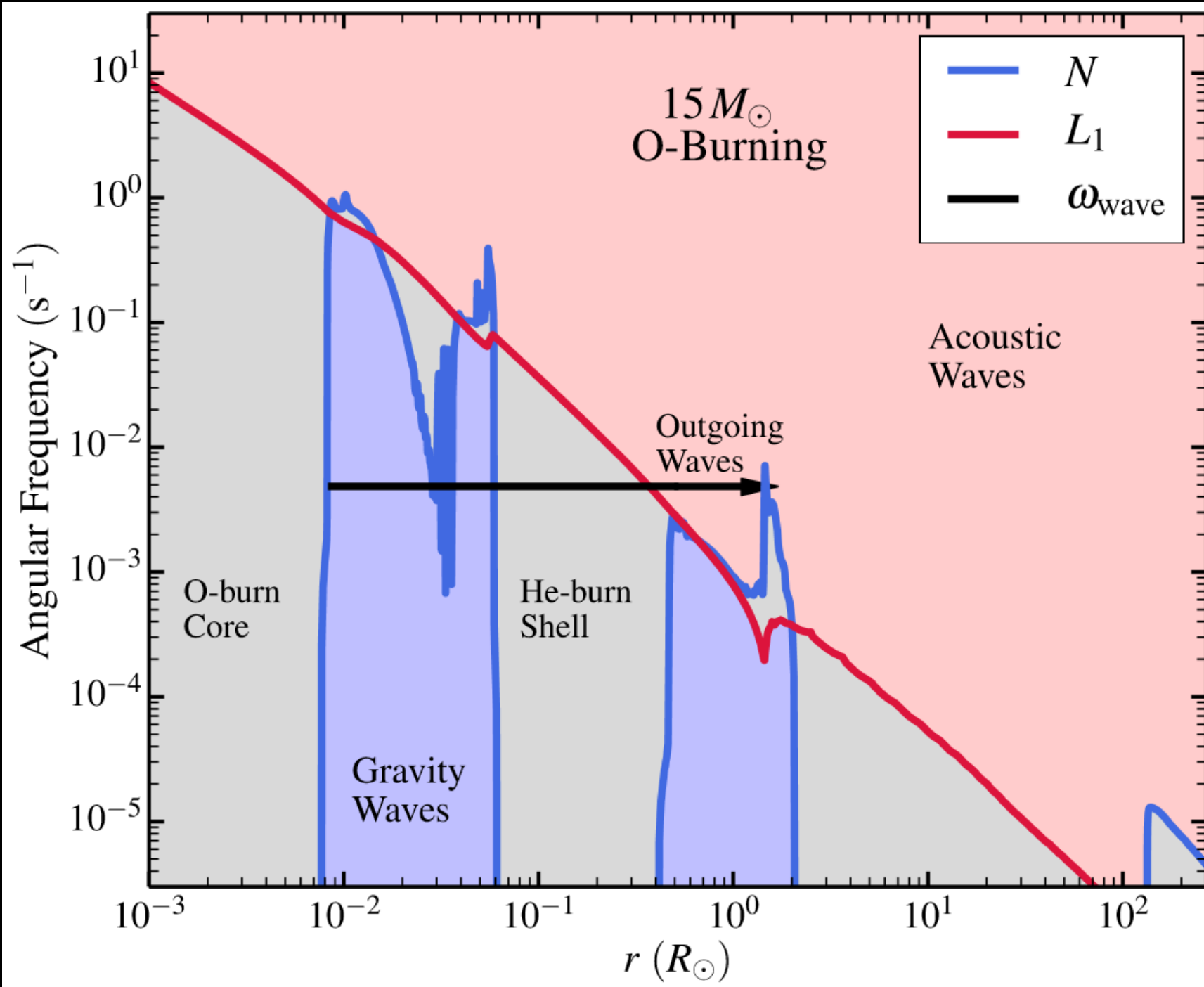


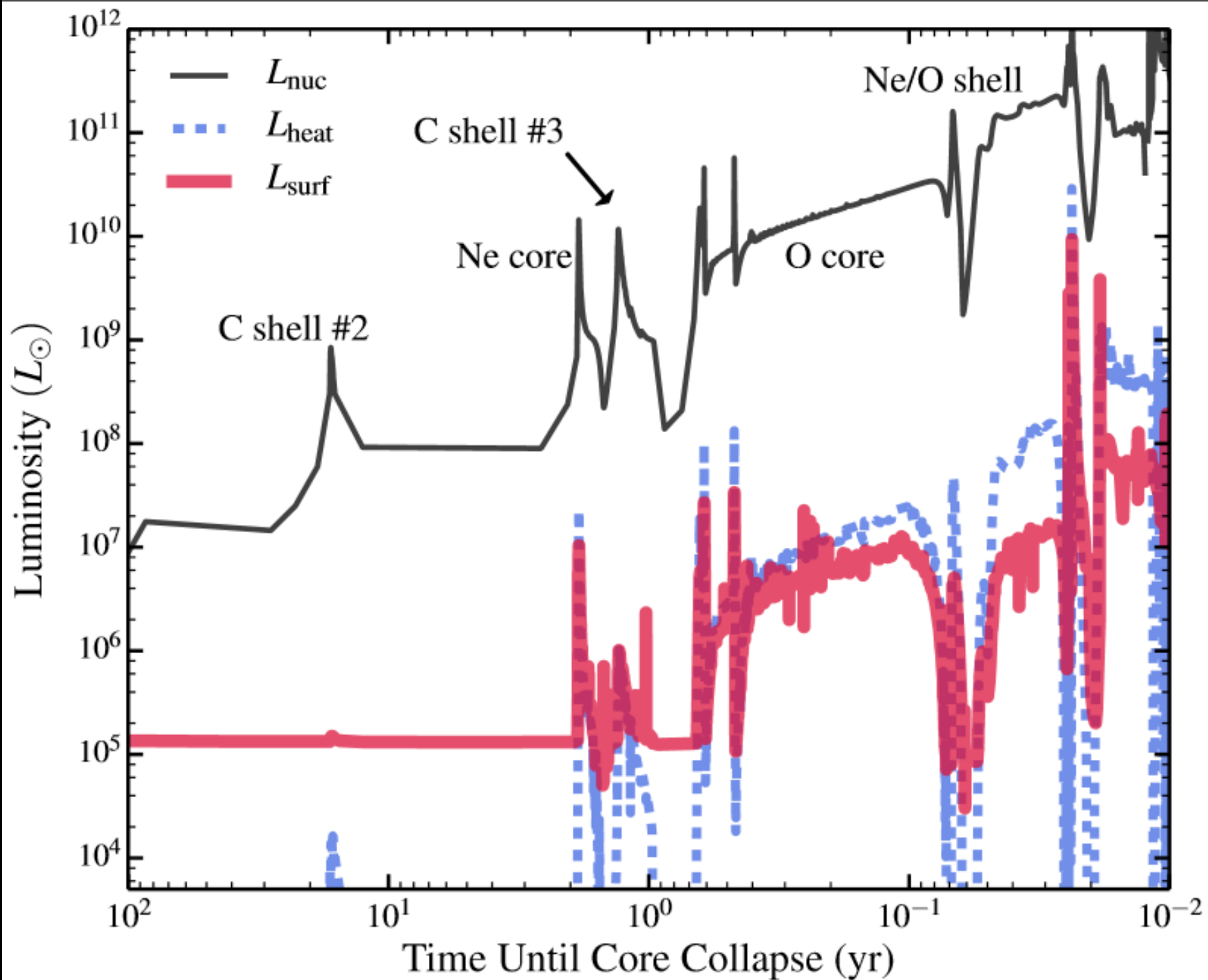




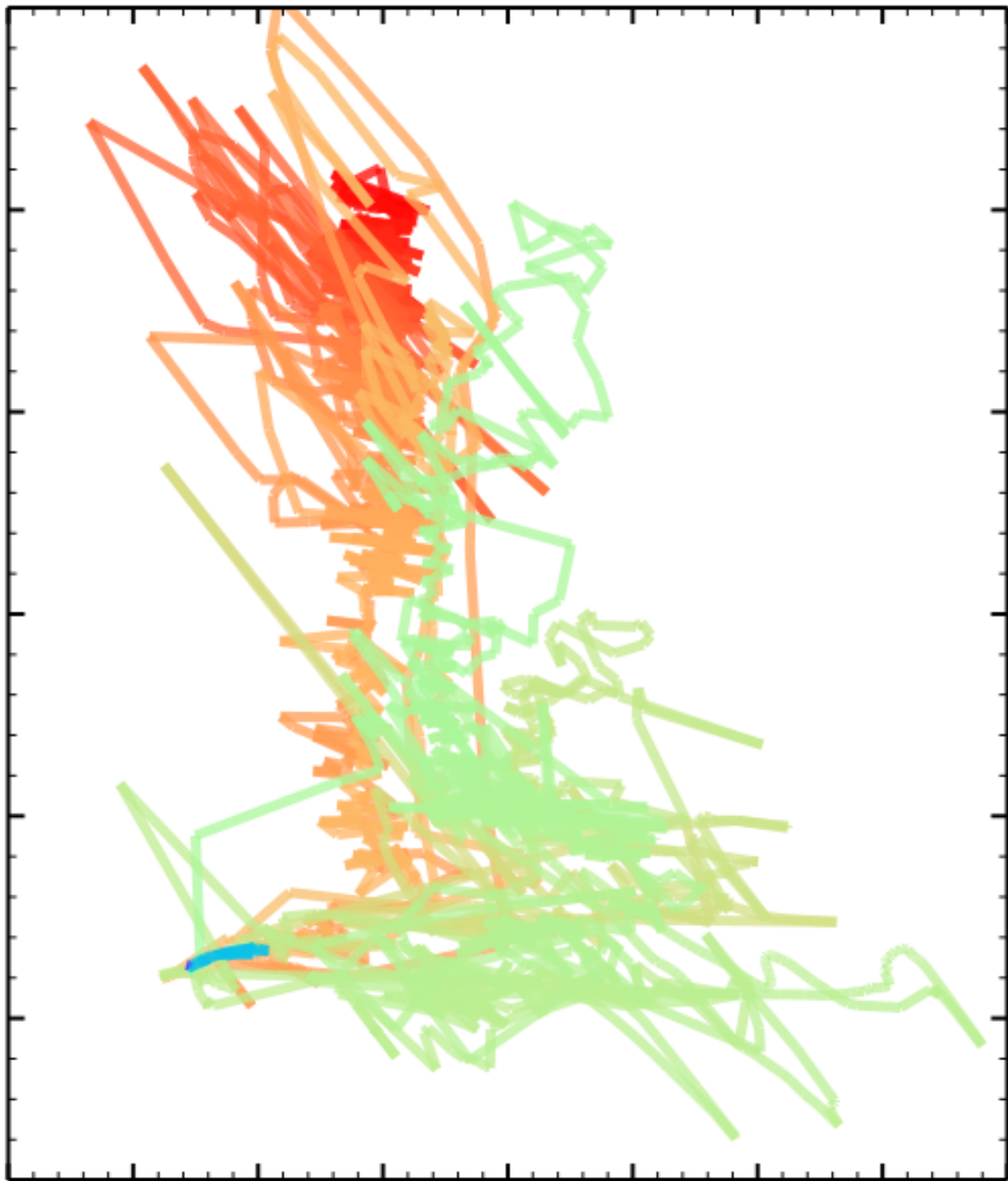




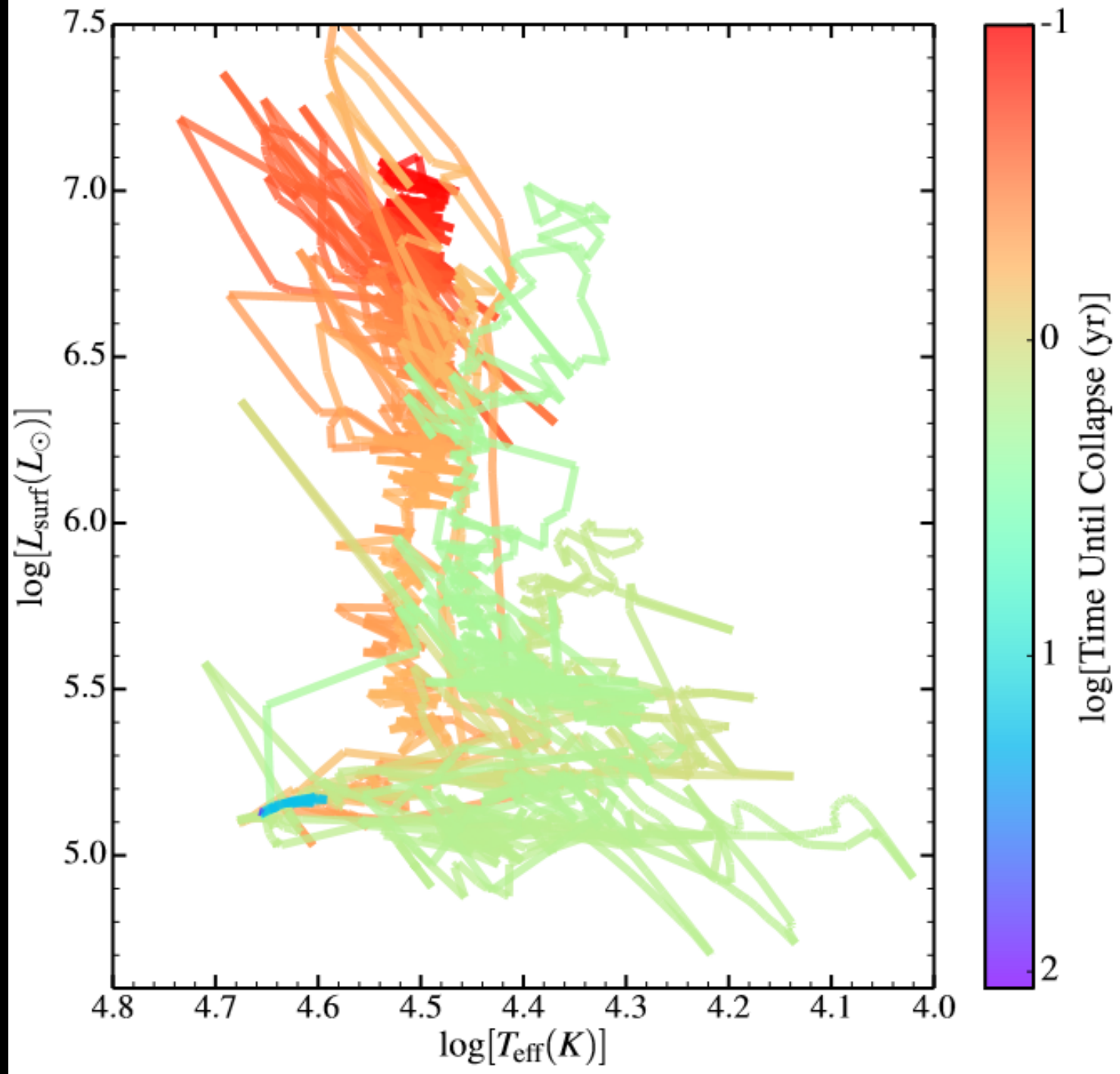




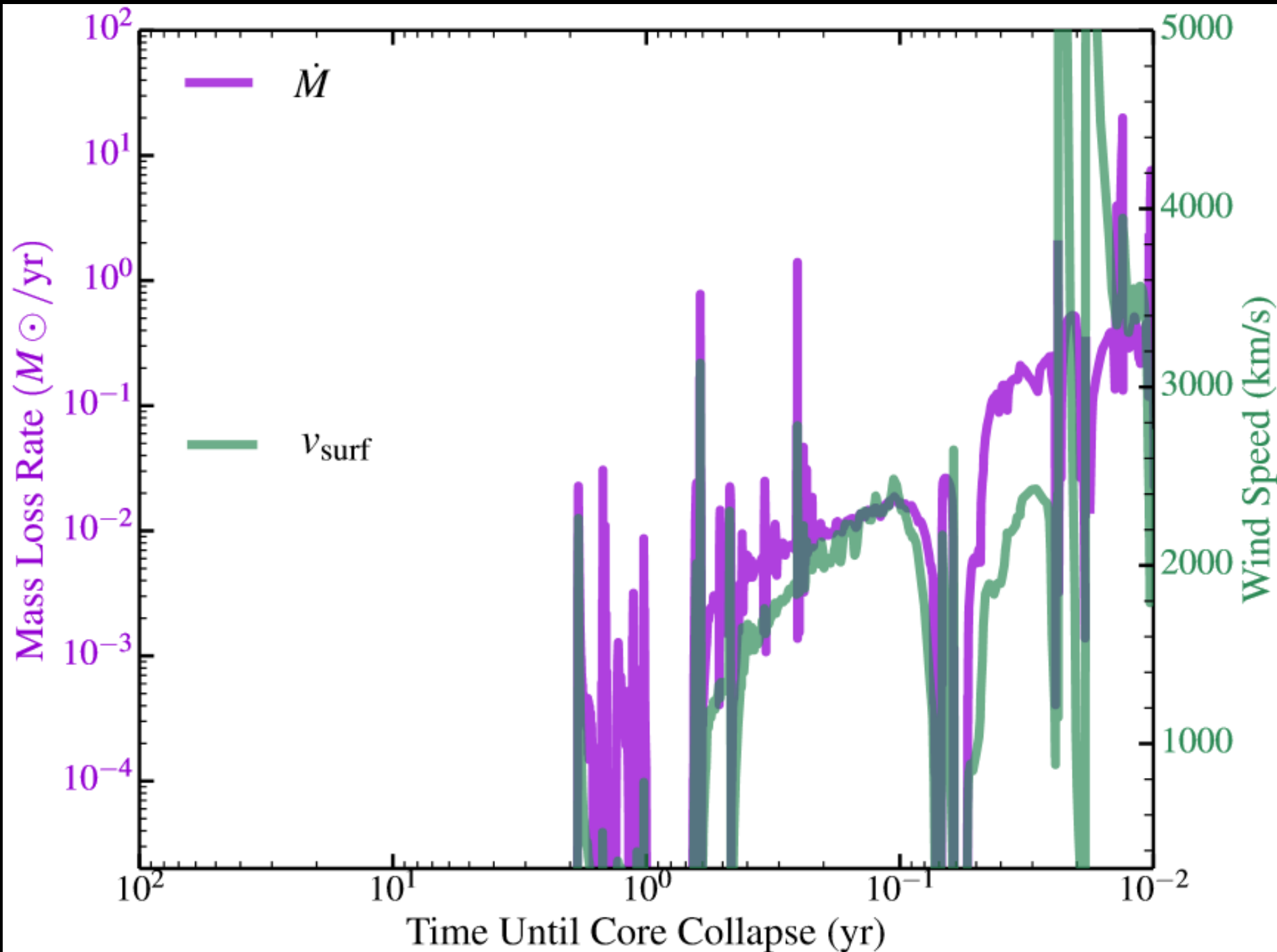
Short thermal time at wave heating region allows for bright outbursts at surface



4/25/2017

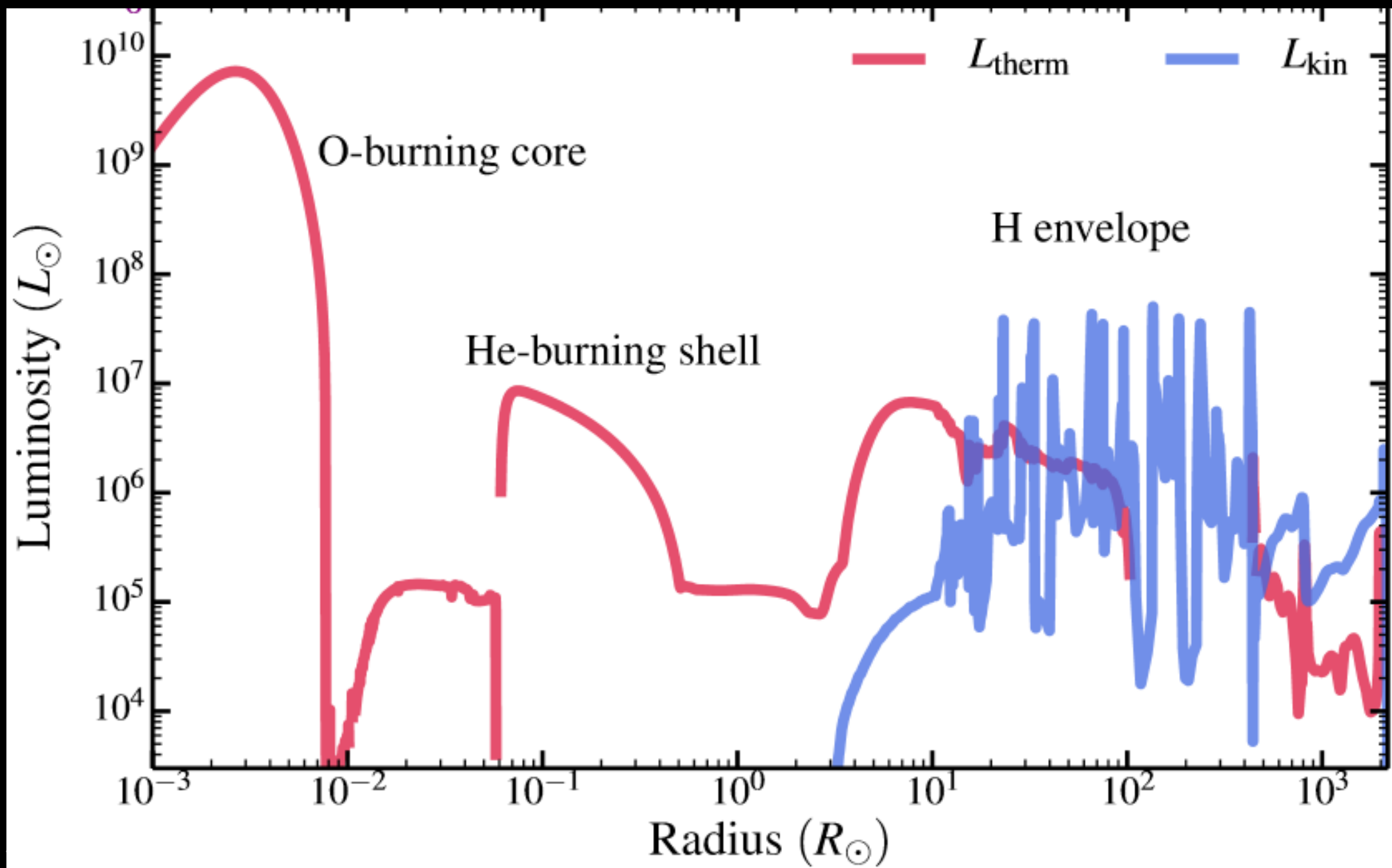


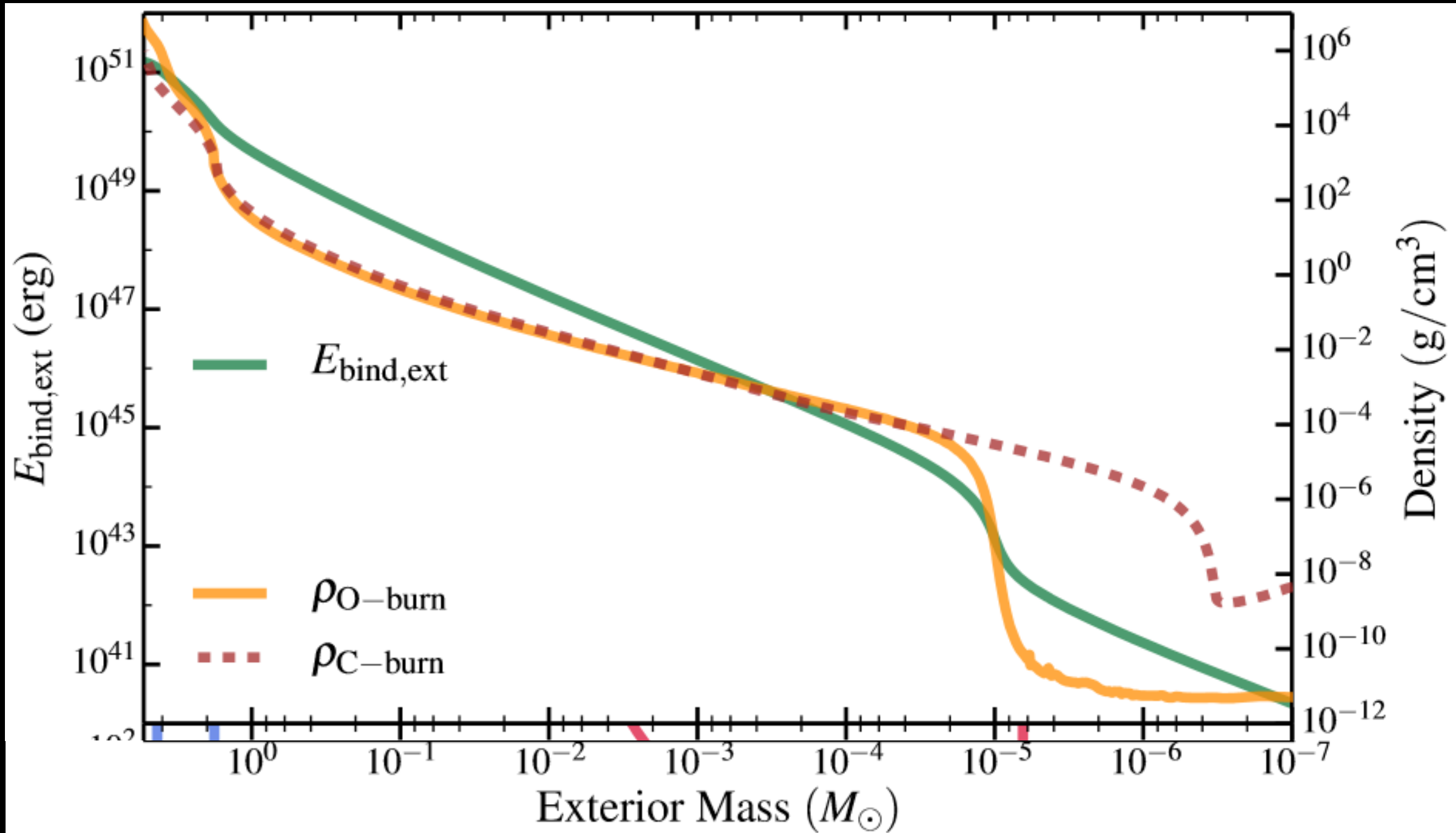
Extreme photometric variability limited to final ~3 years

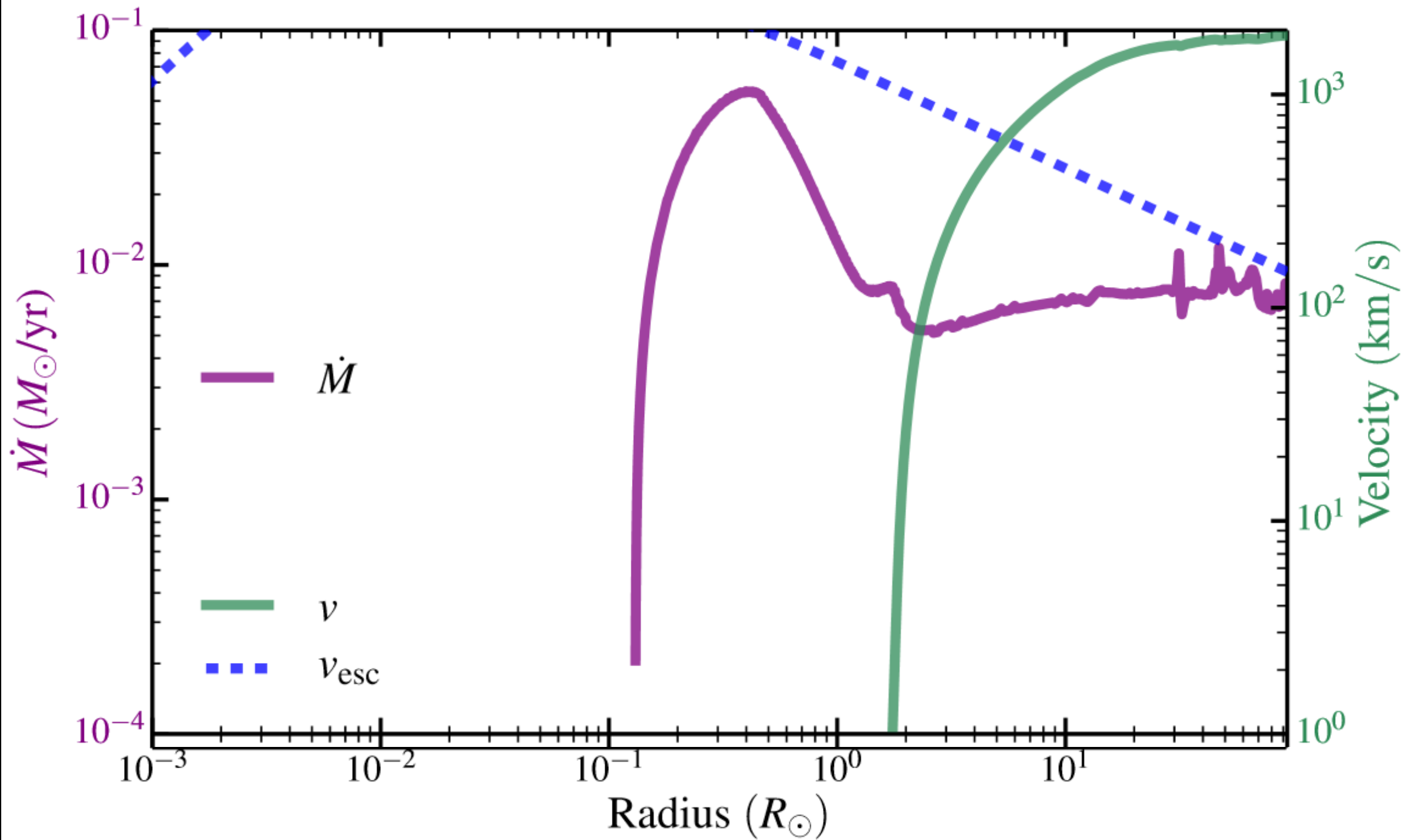


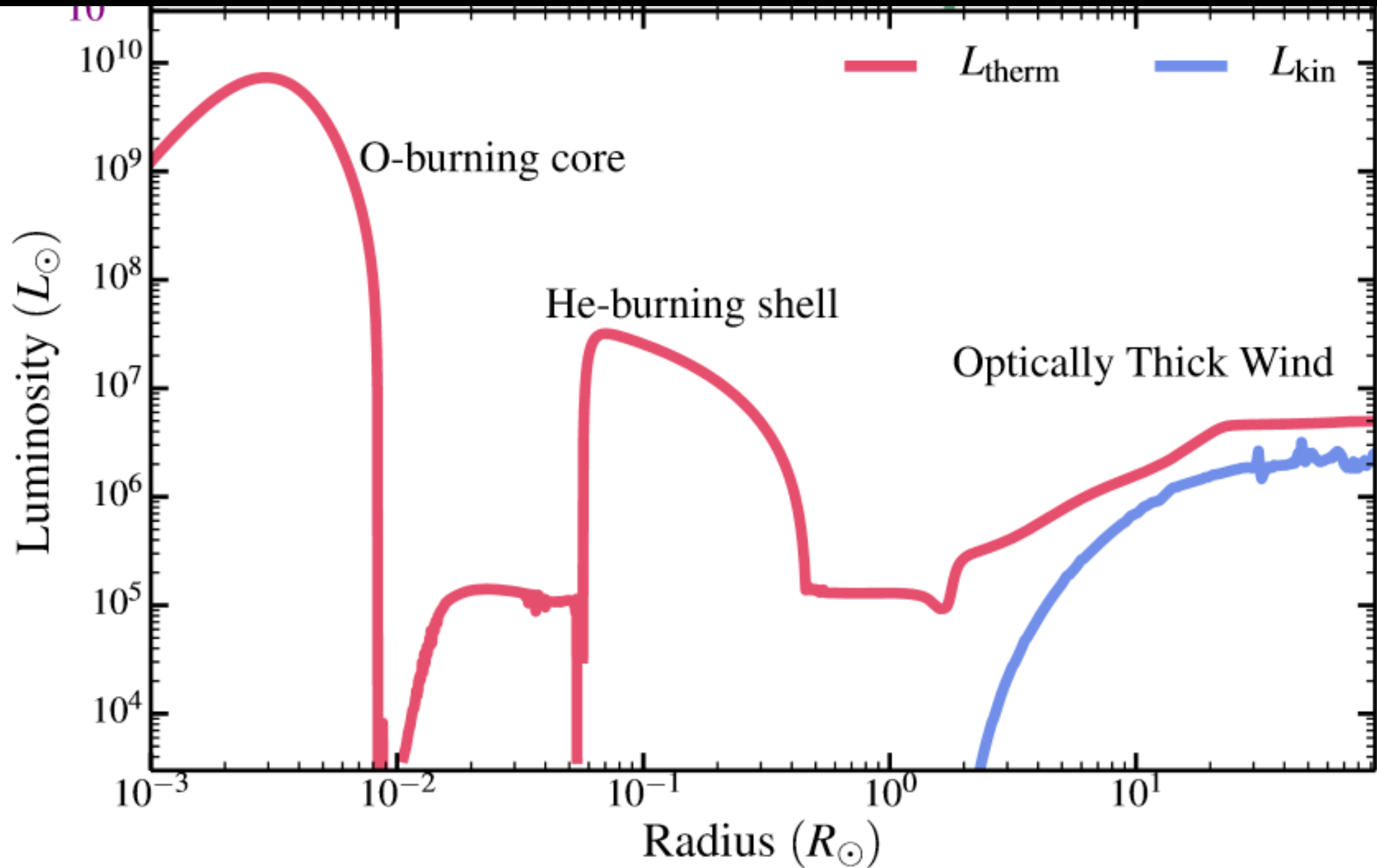
Typical wind speeds of ~ 2000 km/s, mass loss rates of $10^{-2} M_{\text{sun}}/\text{year}$

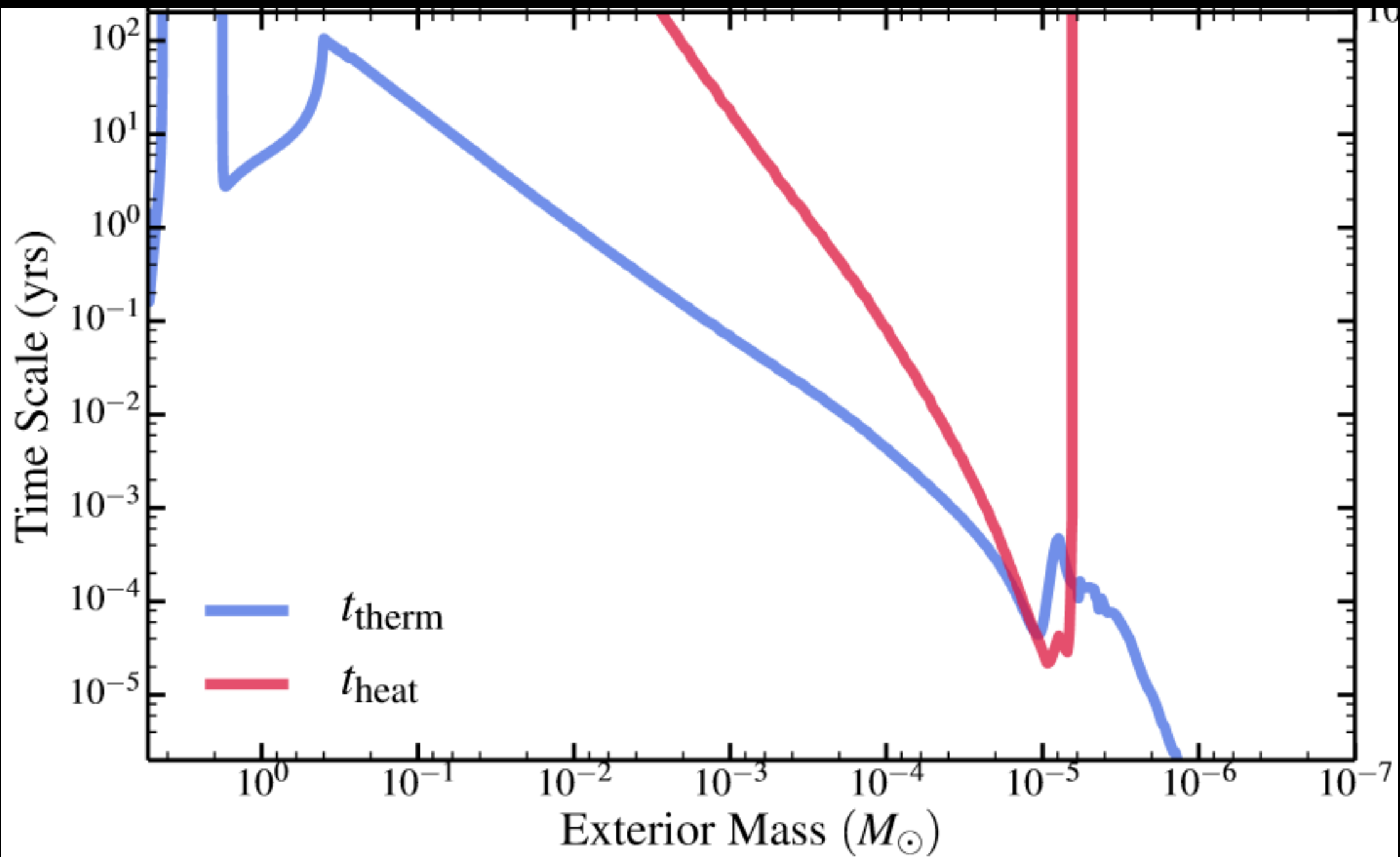
Quataert et al. 2016

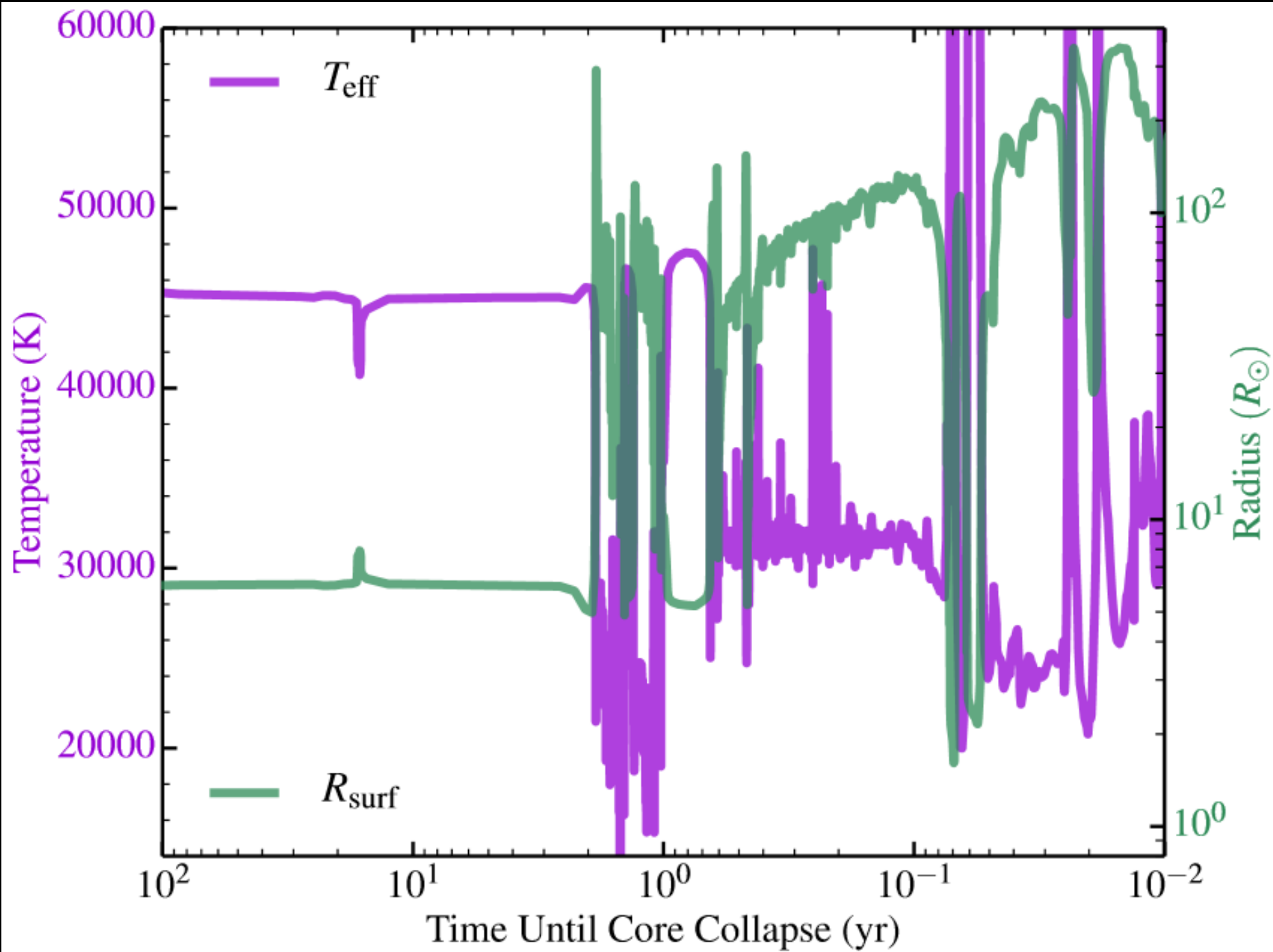


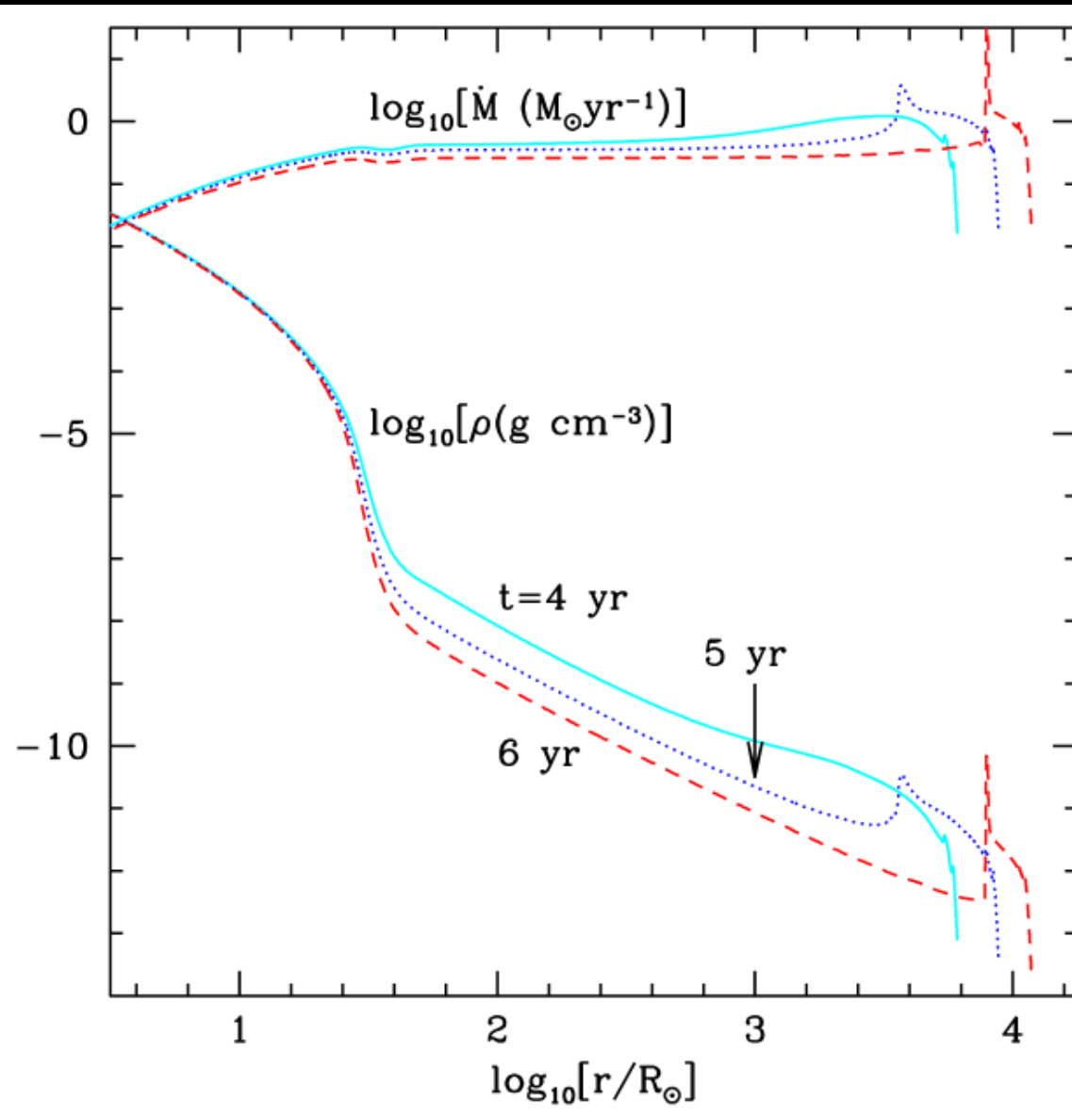
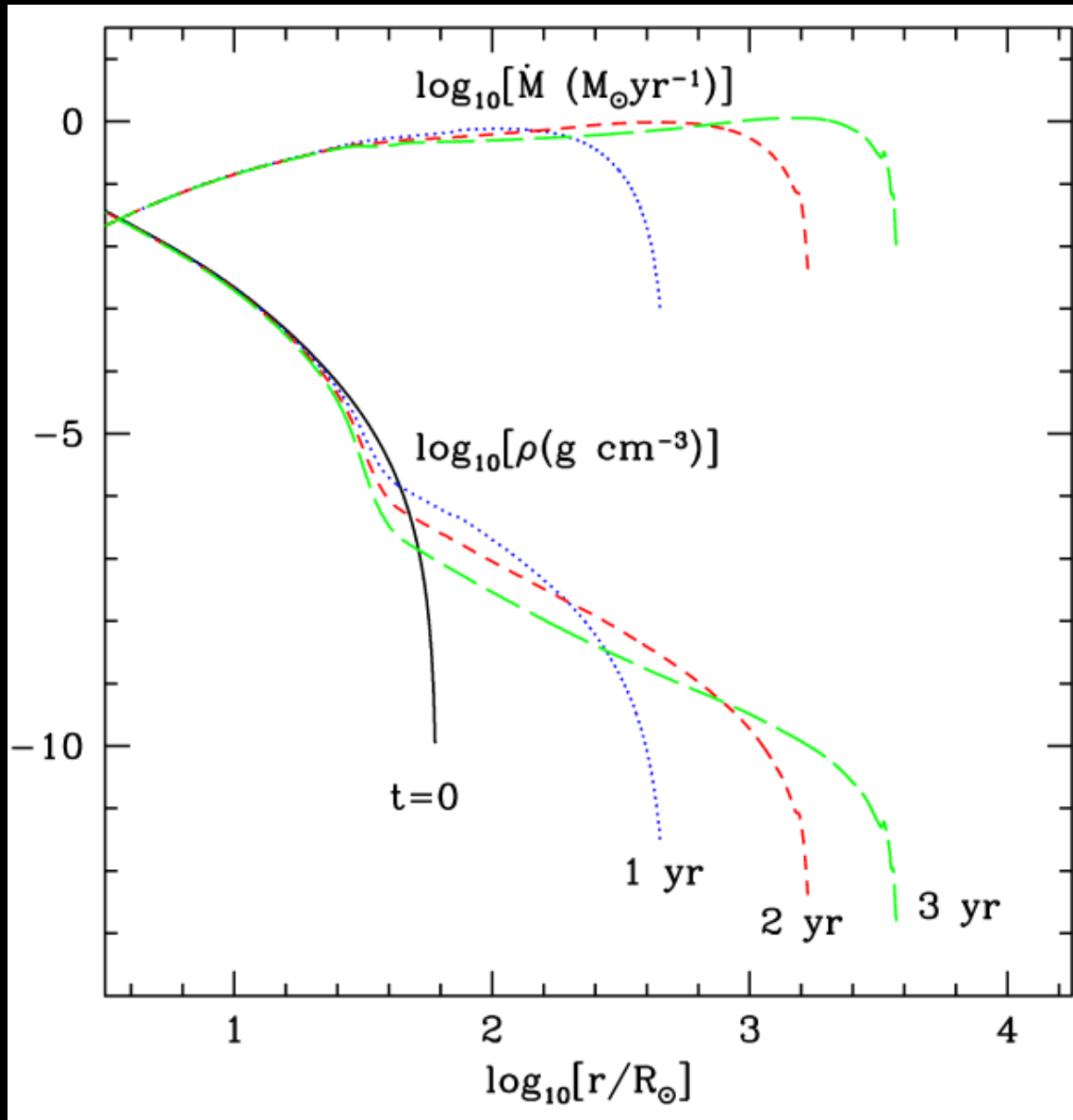


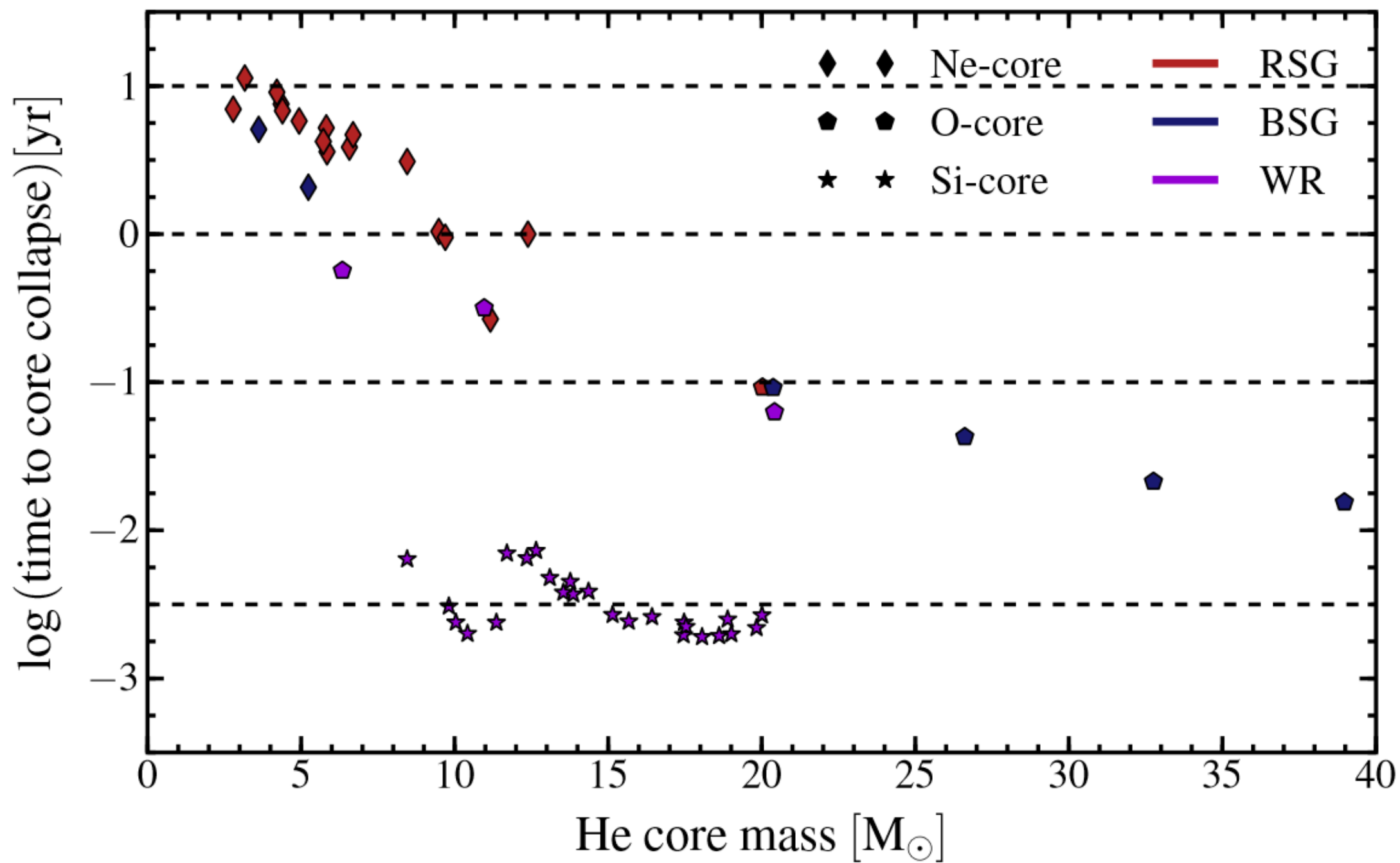












Wave Frequencies

- Wave frequencies determine wave properties
- Convection tends to excite waves with frequency

$$\nu_{\text{wave}} \sim \nu_{\text{con}} \sim \frac{v_{\text{con}}}{2H}$$

- Because convection is slow (sub-sonic), wave frequencies are low, meaning convection most efficiently excites gravity waves

Wave Propagation

- After being excited by convection, waves propagate through the star until they damp out
- Wave propagation determined by local values of the Brunt-Vaisala frequency N , and the Lamb frequency L_l

- Gravity waves where $\omega < N, L_l$

$$L_l^2 = \frac{l(l+1)c_{\text{sound}}^2}{r^2}$$

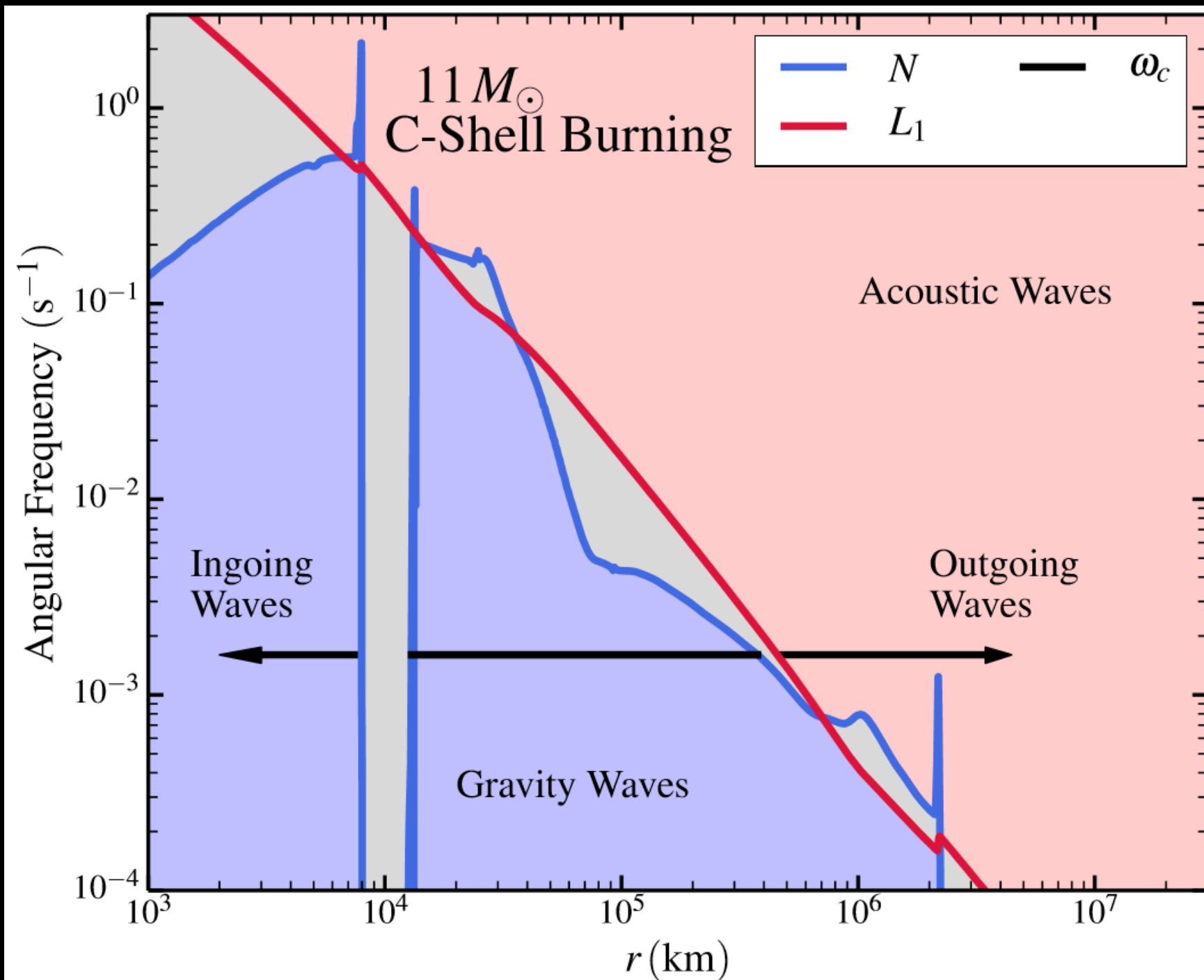
- Sound waves where $\omega > N, L_l$

- Partial wave reflection where $\omega = N, L_l$

Wave Tunneling

- To propagate into envelope of a massive star, waves must switch from gravity waves to pressure waves
- Waves “tunnel” through evanescent region separating gravity wave and sound wave cavities, analogous to particle tunneling in quantum mechanics
- Fraction of wave energy that tunnels through evanescent region is

$$T^2 \sim \left(\frac{r_g}{r_p} \right)^{2l}$$



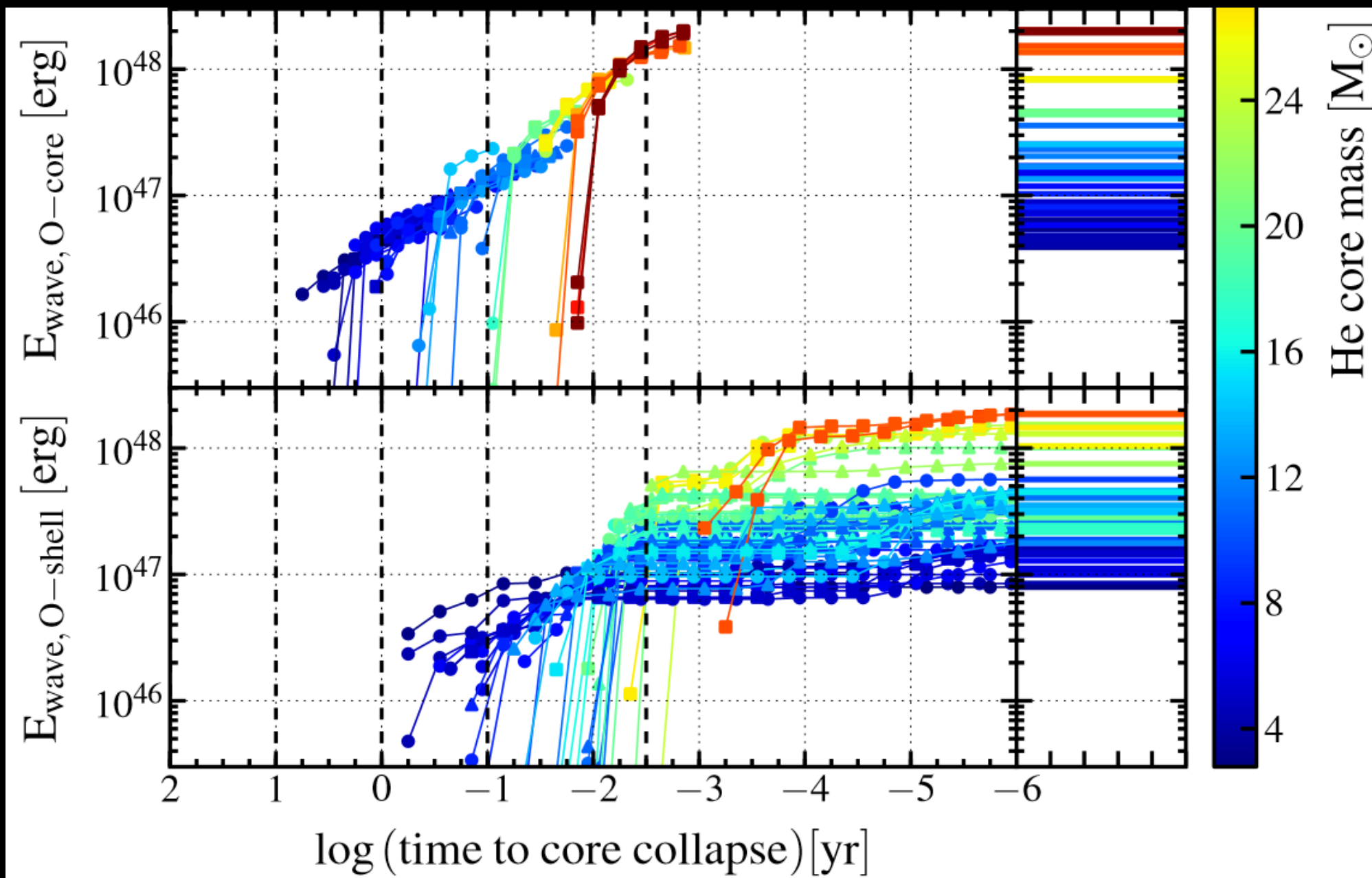
Wave Damping

- Waves can be damped out by various processes such as thermal diffusion, neutrino emission, and non-linear wave breaking

- Wave damping rate due to thermal diffusion is

$$\gamma_{\text{damp}} = k_r^2 K_{\text{therm}}$$

- In our scenario, waves damp out at base of hydrogen envelope, where they deposit energy as heat



Shiode &
Quataert (2014)