

DE SITTER Holography AND THE CMB

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MOTIVATION

WANT TH'Y OF VERY EARLY UNIVERSE
WHICH RESOLVES / CIRCUMVENTS
INITIAL SINGULARITY

STRATEGY HERE: DS / CFT MAY
MAP EARLY UNIVERSE TO SYSTEM
WITH NO GRAVITY
→ Holography

STROMINGER
⋮

VISION:

QUASI-DE SITTER IN 4D SPACETIME
(INFLATION)



QFT IN 3 EUCLIDEAN DIM'S

THE CMB

STANDARD STORY

INFLATIONARY EPOCH IS NEARLY DE SITTER

→ CMB IS NEARLY SCALE INVARIANT

INTERPRETATION

3) SCALE INVARIANT AT FIXED PT AND CMB IS DESCRIBED BY PERTURBATION THEORY AROUND IT

STRONG PROPOSAL (HOLOGRAPHY)

INFLATIONARY FIXED PT IS IN UNIVERSALITY CLASS OF 3D EUKLIDEAN QFT NEAR RG FIXED PT.

WEAKER PROPOSAL

GRAVITY IN IR IS GOVERNED BY UNIV. CLASSES. THESE ARE USEFUL FOR DESCRIBING CMB (& DARK ENERGY)

OUTLINE OF TALK

DEFINE HOLOGRAPHIC MAP IN ANALOGY WITH ADS / CFT

- COMMENTS:
- 1) NO MICROSCOPICS ("EFFECTIVE FIELD THEORY")
 - 2) COMPUTATIONAL SCHEME DIFFERENT FROM "DS/CFT" MALDACENA
 - 3) CLOSE RELATION TO STANDARD INFLATIONARY COSMOLOGY

APPLICATIONS TO COSMOLOGY

- 1) RECOVER STANDARD PREDICTIONS FOR CMB
- 2) DISCUSS POSSIBLE HOLOGRAPHIC INTERPRETATION

H-J THEORY

NATURAL SET-UP FOR SEMI-CLASSICAL TREATMENT (STANDARD IN COSMOLOGY)

EXAMPLE: CLASSICAL MECHANICS

$$S(q, t) = \int^t dt' L(q, \dot{q}, t')$$

FCT OF $t, q(t)$ EVALUATED ON-SHELL
→ ENFORCE E.O.M

$$\delta_q S = \int dt' \left(\frac{\partial L}{\partial q} - \partial_t \frac{\partial L}{\partial \dot{q}} \right) \delta q + \frac{\delta L}{\delta \dot{q}} \delta \dot{q} \Big|_t$$

$$\Rightarrow \frac{\partial S}{\partial q} = \frac{\partial L}{\partial \dot{q}} = p.$$

SEMI-CLASSICAL WAVEFCT

$$\Psi \sim e^{iS(q, t)}$$

$$\Rightarrow \hat{p} = -i \frac{\partial}{\partial q} \text{ ACTS AS } \frac{\partial S}{\partial q}$$

HAMILTONIAN

$$\frac{\partial S}{\partial t} + H(q, \frac{\partial S}{\partial q}) \Leftrightarrow \text{SCHRÖDINGER EQUATION}$$

INFLATION

ϕ IS SPECTATOR SCALAR FIELD IN DE SITTER

$$ds^2 = a(z)^2 (-dz^2 + d\vec{x}^2)$$

$$a(z) = -\frac{1}{H\tau} \quad \text{NOTE } \tau < 0$$

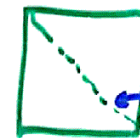
$\tau \rightarrow 0_-$ IS FUTURE

$$S(\phi, \tau) = \int d^3x dz = \int d^4x \sqrt{g} \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$

BODY CONDITION @ $\tau = -\infty$

ACTION FINITE (INTEGRAL CONVERGENT)

PHYSICS: EUKLIDEAN CONTINUATION REGULAR



REGULAR @ "HORIZON"

ANALOGY W. ADS/CFT

PAST HORIZON ~ ORIGIN OF EAdS

$$\Psi(\phi, t) \sim e^{iS} \leftrightarrow Z(\phi)$$

INFRARED DIVERGENCESEXAMPLE: S_{HJ} FOR MASSIVE SCALAR

$$S(\phi, \tau) = \int_{-\infty}^{\tau} d^3x dz \sqrt{-g} \left(\frac{1}{2} \left(\frac{\phi'}{a} \right)^2 - m^2 \phi^2 \right)$$

$$= \int d^3x \frac{1}{2} a^2 \phi \phi' \Big|_{-\infty}^{\tau} + \underbrace{\text{e.o.m.}}_0$$

E.O.M. $\Rightarrow \phi(\tau) = c_- \tau^{\lambda_-} + c_+ \tau^{\lambda_+}$

$$\lambda_{\pm} = \frac{3}{2} \left(1 \pm \sqrt{1 - \left(\frac{2m}{3H} \right)^2} \right)$$

S REGULARITY @ $\tau_i = -\infty \Rightarrow \lambda_- \text{ SOL ONLY}$

$$S_{HJ} = \int d^3x \frac{1}{2H^2} c_-^2 \lambda_- \tau^{-3+2\lambda_-} \Big|_{\tau \rightarrow 0}^{\infty}$$

\Rightarrow H-J FUNCTIONAL DIVERGES
AT LATE TIMES \sim LONG WAVELENGTH

 \Rightarrow INFRARED DIVERGENCELOCAL DIVERGENCES: GENERAL

WANT TO FIND GENERAL EXPRESSION FOR DIVERGENCES.

HAMILTONIAN CONSTRAINT \Rightarrow H-J EQUATION

$$\left(\sqrt{\frac{1}{g}} \frac{1}{g^{ij}} \frac{\delta S}{\delta g_{ij}} \right)^2 - 2 \sqrt{\frac{1}{g}} \frac{\delta S}{\delta g_{ij}} \frac{1}{g^{ij}} \frac{\delta S}{\delta g_{ij}} - \frac{1}{2} \left(\frac{\delta S}{\delta \phi} \right)^2$$

$$= V - \frac{1}{2} R + \frac{1}{2} \vec{\partial} \phi \vec{\partial} \phi$$

GENERAL CASE (INCLUDING GRAVITY)

ANSATZ: DERIVATIVE EXPANSION

$$S = \int d^3x \left[V(\phi) + M(\phi) \vec{\partial} \phi \vec{\partial} \phi + \mathcal{E}(\phi) R + \dots \right]$$

$$\Rightarrow V(\phi) + \frac{1}{2} (\partial_{\phi} V)^2 - \frac{3}{4} V(\phi)^2 = 0$$

+ EQ'S FOR $M(\phi), \mathcal{E}(\phi)$

$$\Rightarrow V(\phi) = -2H(\phi)$$

\uparrow
SLOWLY EVOLVING ϕ
 \rightarrow EVOLVING H
 \rightarrow DIV DEPENDS ON
POSITION OF STUDY

COMMENTS

- 1) $V(\varphi)$, $M(\varphi)$, $\Phi(\varphi)$ ARE ALL DIVERGENCES (HIGHER ORDERS IN DERIVATIVE EXPANSION FINITE AS $z \rightarrow 0$)
- 2) FUNCTIONAL DEPENDENCE ON φ IS GENERAL - NOT LIMITED TO HOMOGENEOUS BACKGROUNDS.

RENORMALIZED H-J FUNCTIONAL

GENERAL DIVERGENCES

$$S_{\text{loc}} = \int_{\tau_0}^3 d^3x [V(\varphi) + M(\varphi) \delta\varphi \delta\varphi + \Phi(\varphi) R]$$

$$\Rightarrow S_{\text{ren}} = S - S_{\text{loc}}$$

IS FINITE. INTERPRETATIONS:

1) S_{ren} IS ACTION RELATIVE TO (DIVERGENT) BACKGROUND ACTION

→ H-J FUNCTIONAL OF FLUCTUATIONS

2) IN LOCAL FIELD THEORY WE CAN (MUST) SUBTRACT LOCAL COUNTERTERMS TO CANCEL UV DIVERGENCES

IF INFLATION DUAL TO LOCAL 3D QFT
 THE IR DIV'S OF GRAVITY \sim UV DIV'S OF QFT.

→ HOLOGRAPHY.

GRAVITATIONAL BACKREACTION

PROPER SYSTEM TO CONSIDER: SCALAR + GRAVITY

PERTURBED METRIC (LONGITUDINAL GAUGE)

$$ds^2 = a(z)^2 \left(-(1+2\Phi)dz^2 + (1+2\Psi)d\vec{x}^2 \right)$$

 \Rightarrow 3 SCALAR FIELDS $\Phi, \Psi, \delta\phi$
2 CONSTRAINTS (EOM OF GAUGEFIXED METRIC COMPONENTS $g_{zi}, g_{ij} (i \neq j)$)
 \Rightarrow ONE NET EFFECTIVE SCALAR.

$$\mathcal{S} = \delta\phi - \frac{\phi'}{\mathcal{H}} \Psi$$

(NICE GEOMETRIC INTERPRETATION)
MUKHANOV

THIS SCALAR IS FIELD OF REAL INTEREST

THE WAVE EQUATIONLINEARIZING EINSTEIN EQUATIONS
(HARD WORK) GIVE E.O.M.

$$\mathcal{S}'' + 2\left(\mathcal{H} - \frac{\phi''}{\phi'}\right)\mathcal{S}' - \bar{\nabla}^2 \mathcal{S} + 2\left(\mathcal{H}' - \mathcal{H}\frac{\phi''}{\phi'}\right)\mathcal{S} = 0$$

DIFFICULT TO SOLVE FOR GENERAL $V(\phi)$
(WHICH ENTERS THROUGH BACKGROUND \mathcal{H}, ϕ)SLOW ROLL INFLATION

$$\epsilon = \frac{1}{2} \left(\frac{\partial_\mu V}{V} \right)^2 \ll 1$$

$$\eta = \frac{\partial_\mu^2 V}{V} \ll 1$$

THEN

$$\mathcal{S}_R = |z|^{\frac{1}{2} + \nu - 6} \mathcal{H}_\nu(1kz)$$

$$\nu = \frac{1}{2} - \eta + 3\epsilon$$

 \uparrow
HANKEL FCT.

THE QUADRATIC ACTION

THE H-J FUNCTIONAL IS THE ON-SHELL ACTION LESS COUNTER TERMS

$$S_{\text{ren}} = \int \mathcal{L} - S_{\text{loc}}$$

$$\frac{1}{2} R^{(4)} - \frac{1}{2} \partial_r \phi \partial_r \phi - V(\phi)$$

QUADRATIC FLUCTUATIONS

$$= \int \mathcal{L}_{\text{quad}} - S_{\text{loc}}$$

USE E.O.M. INTEGRATE IMPARTS

$$= -2 \int d^3 k \left(\frac{2}{\phi_1}\right)^2 \frac{ik^3}{H^2} (k\tau_0)^{6\epsilon-2\gamma} \rho_k \rho_{-k}$$

THIS IS RENORMALIZED ACTION

WRITTEN IN TERMS OF GAUGE-INVARIANT VARIABLES

THE FLUCTUATION SPECTRUM

$$\Psi[\rho] = e^{i S_{\text{ren}}[\rho]}$$

$$\Rightarrow \langle \rho_E \rho_{-E} \rangle = \int \mathcal{D}\rho \rho_E \rho_{-E} |\Psi[\rho]|^2$$

$$= \left(\frac{H}{\phi}\right)^2 \frac{H^2}{4k^3} (k\tau_0)^{2\gamma-6\epsilon}$$

POWERSPECTRUM

$$P_\rho = 4\pi \left(\frac{k}{2\pi}\right)^3 \langle \rho_E \rho_{-E} \rangle = \left(\frac{H}{\phi}\right)^2 \left(\frac{H}{2\pi}\right)^2 (k\tau_0)^{2\gamma-6\epsilon}$$

SPECTRAL INDEX

$$P_\rho \sim k^{n_s-1} \Rightarrow n_s-1 = 2\gamma-6\epsilon$$

AMPLITUDE ENHANCEMENT

$$\frac{P_{\rho}}{P_T} = \left(\frac{H}{\phi}\right)^2 = \frac{1}{\epsilon} \gg 1$$

CUT OFF (IN) DEPENDENCE

$$P_Q = \underbrace{\left(\frac{H}{\dot{\phi}}\right)^2 \left(\frac{H}{2\dot{\sigma}}\right)^2}_{\text{BACKGROUND DATA}} (k\tau_0)^{2\gamma-6\epsilon}$$

↑
CUT OFF

BACKGROUND DATA SATISFIES
FRW EQUATION \Rightarrow MILDLY DEPENDENT
ON TIME
↑
SLOWROLL

\Rightarrow ALL τ_0 DEPENDENCE CANCELS

REMOVING CUTOFF GIVES FINITE RESULT.

(BUT LINEAR PERT THY AND/OR SLOWROLL
APX MAY NOT BE VALID THROUGHOUT)


RUNNING (PHYSICAL) CUTOFF $\tau_0 = \frac{1}{aH}$

(HORIZON CROSSING) MAY BE MORE APPRIATE

(SUM'S LARGE LOG'S)

PHYSICAL OBSERVABLES

FOUNDATIONAL PROBLEM IN QUANTUM DE SITTER:
WHAT ARE PHYSICAL OBSERVABLES?

E.G.  \Rightarrow X'S SPACE LIKE SEPARATED
 \Rightarrow CORRELATORS META-OBSERVABLES
WITTEN

END OF INFLATION

\Rightarrow MODES REENTER HORIZON
 \Rightarrow BECOME PHYSICALLY OBSERVABLE

"WE ARE METAOBSERVERS"⁴
DANIELSSON

PRECISE CONSTRUCTION (I.E. DETAILS
OF REHEATING ETC) DOES NOT SEEM
IMPORTANT.

THE HOLOGRAPHIC RG EQUATION

S_{HJ} SOLVES H-J EQ ; S_{loc} SOLVES IT TOO

⇒ EQ FOR $S_{REN} = S - S_{loc}$

$$(2g_{ij} \frac{\delta}{\delta g_{ij}} + \beta \frac{\delta}{\delta \phi}) S_{ren}(g_{ij}, \phi) = 0$$

$$\beta \equiv -2 \frac{\partial \phi U}{U} = -\frac{\dot{\phi}}{H}$$

QUADRATIC ORDER

$$(a \frac{\partial}{\partial a} + \beta_0 \frac{\partial}{\partial \phi} + 2\gamma) S^{quad}(a, \phi) = 0$$

$$\gamma = \frac{\partial \beta}{\partial \phi} = 2\epsilon - \eta$$

$$\beta_0^2 = 1 - 2\epsilon$$

SPECTRAL INDEX

$$n_s = 1 - \beta_0^2 - 2\gamma = 1 + 2\eta - 6\epsilon$$

SCALING OF CORRELATORS FROM
CALLAN-SYMANZIK EQUATION !

HOLOGRAPHIC INTERPRETATION

COSMOLOGICAL EVOLUTION ~ RG FLOW

STROMINGER

AS TIME MOVES FORWARD ENTROPY INCREASES

→ MOTION TOWARDS UV
("INTEGRATING IN" D. O. F.)

SO MOTION BACK TOWARDS INFLATIONARY
EPOCH → MOTION TOWARDS IR

SO WE ARE NOT SURPRISED THERE IS
A SCALE INVARIANT FIXED PT IN OUR
PAST !

PREDICTION FIXED PT A STABLE ATTRACTOR

$$\Rightarrow (kz)^{n_s} \rightarrow 0 \quad z \rightarrow \infty$$

⇒ SPECTRUM RED $n_s < 1$
(LOG GROWTH AT LARGE SCALES)

A TOY MODEL

VISION: COMPUTE n_s, n_T, \dots FROM
3D DUAL RATHER THAN $V(\phi)$

TOY MODEL

$$S = \int d^3x \left[\frac{1}{2} \partial_r \phi \partial^r \phi - \frac{1}{6!} g \phi^6 \right]$$

CLASSICALLY SCALE INVARIANT BUT
 ϕ^6 IRRELEVANT IN $SR_p \rightarrow$ JUST WHAT WE WANT
DUE TO QUANTUM

$$\left(\frac{\partial}{\partial \ln r} + \rho \frac{\partial}{\partial g} - 2\lambda \right) \langle \phi^6 \phi^6 \rangle = 0$$

$$\frac{5}{3} \frac{g}{16\pi^2} > 0$$

$$\Rightarrow G_{\phi^6} \sim \left(\frac{\mu}{k} \right)^{2\lambda} \sim (kz)^{-2\lambda}$$

$$\Rightarrow n_s = 1 - 2\lambda$$

SUMMARY

DEVELOPED TH'Y OF CMB IN ANALOGY
WITH ADS/CFT. SUGGESTIVE

FOCUS ON IR DIV'S, FULLY CHARACTERIZED.
INTERPRETATION? UNIVERSALITY?

FUTURE DIRECTIONS

COMPLETE RG ANALYSIS

- HIGHER PT FCTS, BEYOND SLOWROLL

IN PROGRESS

CONNECTION TO STRING TH'Y?