

# TOPOLOGICAL MEMBRANES

with

Lilia Angelova

Paul de Medeiros

hep-th/0507089

## TOPOLOGICAL M-THEORY \*

- o topological theory on 7-manifolds with  $G_2$ -structure
- o suggested as unification of A/B model topological strings
- o unification of lower dim form theories of gravity

\* hep-th/0411073

R. Dijkgraaf, A. Neitzke, S. Gukov, C. Vafa

hep-th/0412021

N. Nekrasov

## HITCHIN'S THEORY

$M$  with  $G_2$ -holonomy  $\leftrightarrow$   
associative 3-form  $\phi$

$$d\phi = 0 \quad d^* \phi = 0$$

classical action which  
extremizes on these conditions

$$V(\phi) = \int_M \phi \wedge * \phi$$

[+ in terms of  $\phi$ ]

$$\text{variation } \phi = \phi_0 + dB$$

$$d\phi_0 = 0$$

Reduction to 6d:

$$M_7 = M_6 \times S^1$$

Hitchin's ansatz

$$\phi = S + \alpha \wedge dt$$

$$*\phi = \bar{S} \wedge dt + \frac{1}{2} \alpha \wedge \alpha$$

$$V(d) = \int \phi \wedge *\phi =$$

$$= \frac{1}{2} \int \alpha \wedge \alpha \wedge dt - \frac{i}{4} \int \alpha \wedge \bar{\alpha} \wedge dt$$

$$S = S + i\bar{S}$$

A and  $B + \bar{B}$  model actions.

- topological F-theory  
on  $\text{Spin}(7)$  - manifolds

Spin(7) holonomy metric

$\leftrightarrow$  Cayley 4-form  $\Psi$

$$d\Psi = 0 \quad * \Psi = \Psi$$

- action is more difficult to define

For  $CY \times T^2 = M_8$

$$V(\Psi) = \frac{1}{2} \int_{M_8} * \Psi \wedge \Psi$$

Torus reduction of Cayley-form

$$V(\Psi) = \int_{M_6 \times T^2} dx dy \left( \frac{1}{2} \alpha_1 \alpha_2 \alpha_3 - \frac{i}{4} \beta_1 \bar{\beta}_2 \right)$$

modular transformations of extra terms

$$\begin{pmatrix} dx \\ dy \end{pmatrix} \rightarrow \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} dx \\ dy \end{pmatrix}$$

$$ad - bc = 1$$

$$dX = \begin{pmatrix} dx \\ dy \end{pmatrix} \quad \tilde{\zeta} = \begin{pmatrix} \tilde{s} \\ s \end{pmatrix}$$

$$\exists = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \quad SL(2, \mathbb{Z}) \text{ invariant}$$

Cayley form

$$\Psi = dX^T \exists \tilde{\zeta} + \frac{1}{2} dX^T \exists dX + \frac{1}{2} \alpha \wedge \alpha$$

modular invariant if

$$\begin{pmatrix} \tilde{s} \\ s \end{pmatrix} \rightarrow \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \tilde{s} \\ s \end{pmatrix} \quad \tilde{s} \rightarrow s$$

topological S-duality?

Construct topological theory  
of membranes wrapping  
associative cycles.

Strategy:

Similar to topological A-model.

Consider space of maps from  
3-manifold to  $G_2$ .

Construct topologically invariant  
q.l. that localize on associative  
3-cycles.

USE MATHAI - GUILLEN  
FORMALISM

FUNDAMENTAL FORMULATION  
OF TOPOLOGICAL M-THEORY?

IS THERE A FORMULATION  
IN TERMS OF TOPOLOGICAL  
MEMBRANES?

- membranes wrapping  
associative cycles contribute  
 $N=1$  superpotential \*

- in 7d membrane dual  
to string  
 $\rightarrow$  topological  $G_2$ -string\*

\* hep-th/9808060 Harvey and Moore  
10304115 Beasley and Witten  
\* hep-th/0506211 de Boer, Naggi, Shomer

## MATHAI - GUILLEN FORMALISM

- o integral representation of Euler class of vector bundles
- o constructs topological invariants of infinite dim vector bundles

manifold  $M$ , vector bundle  $E \rightarrow M$

Euler number  $\chi(M)$

- a. topological  
choose generic section  $s$

$$\chi(M) = \sum_{x, s(x)=0} \pm 1$$

Count # of zeros

MG representative of Euler class

$$e_{D_1 S}(E) = \frac{1}{(2\pi)^n} \int dX e^{-\frac{1}{2} g_{ab} s^a s^b + i ds^a X_a + \frac{1}{2} \partial_a^{ab} X_a X_b}$$

- o Euler number obtained for any choice of  $D_1 S$

- o interpolates between two expression of Euler number

$$s \rightarrow y s$$

- o infinite dimensional vector bundles:  $e_D(E)$  has no meaning  
MG expression is used to define regularized Euler number

2. differential geometric

if  $E$  equipped with  
connection  $\nabla$

Euler class

$$e_0(E) = \frac{1}{(2\pi)^n} \text{Pf}(\Omega_D) =$$

$$= \frac{1}{(2\pi)^n} \int dX e^{\frac{1}{2} \Omega_D^{ab} X_a X_b}$$

$\Omega_D$  : curvature of connection

$$\chi(M) = \sum_E e_0(E)$$

MATHAI - GUILLEN

$$e_{0,s}(E)$$

Regularized Euler number

choose appropriate section so  
that zero locus  $M_s \subset M$  is  
finite dimensional

$$\chi_s(M) = \chi(M_s) = \sum_E e_{s,0}(E)$$

$\circ$  depends on section chosen  
But natural choice of sections  
may carry interesting topological  
information.

Example SGM

$M = \text{Map}(S^1, X)$  space of loops

$$E = TM \quad \{x^i(t)\} \quad x^i(1) = x^i(0)$$

section associated to vector field

$x^i$  generating constant shifts

$$\dot{x}^i = dx^i / dt = \omega_i^a x^a \quad x^a = e^a_i \bar{x}_i$$

MG representation of Euler number  
is a path integral with action

$$I_{\text{Sanc}} = \int d\tau \left( \frac{1}{2} g_{ij} \dot{x}^i \dot{x}^j + i \bar{\psi}_i D_\tau \psi^i - \frac{1}{4} R_{ijk\ell} \psi^i \psi^j \bar{\psi}^\ell \bar{\psi}^k \right)$$

$\bar{\psi}^i = \psi^i + \Gamma_{jk}^i \psi^j \psi^k$   
invariant under BRST

$$\delta x^i = \dot{\psi}^i \quad \delta \psi^i = 0$$

$$\delta \bar{\psi}_i = i g_{ij} \dot{x}^j + \Gamma_{ij}^k \psi^j \bar{\psi}^k$$

path integral computes  
Euler number of manifold

- topological A-model

$$M = \text{Map}(\Sigma_3, X)$$

$M_S$  = span of holom maps

## TOPOLOGICAL MEMBRANES

$$M = \text{Map}(\Sigma_3, X)$$

$a=1,2,3$

$\Sigma_3$  : membrane worldvolume  $\{\sigma^a\}$

$X$  : manifold of  $G_2$ -holonomy

$$\{x^I(\sigma)\} \quad I = 1..7$$

$$TX = T_x(\Sigma_3) \oplus N_x(\Sigma_3)$$

$$\{x^a(\sigma), y^i(\sigma)\} \quad a=1..3, i=1..4$$

static gauge  $x^a(\sigma) = \sigma^a$

$M_S$  = maps corresponding to  
associative cycles embedded  
in  $X$

localize P.I. on  $M_S$

associative 3-cycles calibrated

$$\phi|_{x_3} \leq \text{vol } x_3$$

[equality for assoc. 3-cycle  $\tilde{x}_3$ ]

equivalently

$$\star \phi_{IJKL} dx^I dx^K dx^L |_{\tilde{x}_3} = 0$$

$$x_3 = x(I_3)$$

section

$$\Theta_I = \frac{1}{6} \star \phi_{IJKL} \partial_a x^I \partial_b x^K \partial_c x^L \varepsilon^{abc}$$

in static gauge

$$G_i = \delta_{ij}{}^a \partial_a y^j + \frac{1}{6} \phi^{abc} \star \phi_{IJKL} \partial_a y^I \partial_b y^K \partial_c y^L$$

- on circle reduces to A-model holomorphic curves

note  
 linearized equations can be  
 derived equivalent  
 $\partial_a y^i = \star \phi_{IJKL} \partial_a y^I$   
 $\rightarrow$  deformations theory  
 of associative 3-cycles  
 moduli space of associative  
 3-cycles is generically  
 obstructed  
 colourful of twisted Dirac  
 operator

## Mathai - Quillen action

$$I_n = \int_{\Sigma_2} d^3\sigma \left( \frac{1}{2} g^{IJ} \partial_I \theta_J + i \chi^I (\delta \theta_I - \Gamma_{IJ}^K \psi^J \theta_K) - \frac{1}{4} R_{IJKL} \psi^I \psi^J \chi^K \chi^L \right)$$

$$\delta \theta_I - \Gamma_{IJ}^K \psi^J \theta_K = \frac{1}{2} * \phi_{IJKL} \partial_A \psi^J \partial_B \chi^K \partial_C \chi^L \varepsilon^{abc}$$

invariant under BRST

$$\delta \chi^I = \psi^I \quad \delta \psi^I = 0$$

$$\delta \chi^I = i g^{IJ} \theta_J - \Gamma_{JK}^I \psi^J \chi^K$$

## Topological term

$$\int_{\Sigma_2} x^*(\phi) = \int_{\Sigma_2} d^3\sigma \frac{1}{6} \phi_{IJKL} \partial_A x^I \partial_B x^J \partial_C x^K \varepsilon^{abc}$$

- reduces on circle to A-model

$$\int_{\Sigma_2} x^*(\kappa) \quad \kappa : \text{k\"ahler form}$$

measures membrane instantons

- also required for agreement with physical supermembrane action upto quadratic order

$$\int_{\Sigma_2} d^3\sigma \sqrt{\det g_{IJ} \partial_A x^I \partial_B x^J}$$

Observables

- action has global  $U(1)$  ghost number symmetry

$$(x, \psi, \chi) \sim (\sigma_1, +, -1)$$

$$\langle O_{p_1} \dots O_{p_s} \rangle = \int_M O_{p_1} \dots O_{p_s} e^{-tS_M}$$

$$S_n = I_n + \int_{\Sigma_2} x^*(d)$$

restricted to  $H(x) \subset M$  for maps of a given homotopy class  
membrane instanton weight

$$e^{-t \int_{\Sigma_2} x^*(d)}$$

LOCAL OPERATORS

BRST cohomology isomorphic to de Rham cohomology

$$O_{w_p} = \frac{1}{p!} w_{I_1} \dots w_{I_p}(x) + I_1 \dots I_p$$

$$\delta O_{w_p} = \partial O_{w_p}$$

$G_2$ -holonomy decomposition

$$H^0(x, R) = R$$

$$H^1(x, R) = H_2^1(x, R)$$

$$H^2(x, R) = H_2^2(x, R) + H_{14}^2(x, R)$$

$$H^3(x, R) = H_1^3(x, R) + H_2^3(x, R) + H_{23}^3(x, R)$$

[  $H_{12}^3 = 0$  for smooth compact  $G_2$  ]

de Rahm cohomology spanned by

$$\mathcal{O}_\phi = \frac{1}{6} \phi_{IJK} \psi^I \psi^J \psi^K$$

$$\mathcal{O}_\alpha = \frac{1}{2} \phi_{IJK} \delta_{K\alpha} \psi^I \psi^J \psi^K$$

$$\mathcal{O}_p = \frac{1}{2} p_{IJ} \psi^I \psi^J \quad p = p_{IJ}^{12} p^{IJ}$$

### NON-LOCAL OPERATORS

descent equation

$$\delta W_n = d W_{n-1}$$

$$\mathcal{O}^{(n)} = \int_{Y^n} W_n \quad n=1,2,3$$

3 sets of non-local descendants

$$* p_{IJ}^{12} = \delta_I^P \delta_J^Q - \frac{1}{6} \phi_{IJK} \phi^{PQR}$$

$$\mathcal{O}_\phi^{(1)} = \int_{Y^1} \frac{1}{2} \phi_{IJK} \partial_a x^I \partial_b x^J \partial_c x^K d\sigma^a d\sigma^b$$

$$\mathcal{O}_\phi^{(2)} = \int_{Y^2} -\frac{1}{2} \phi_{IJK} \partial_a x^I \partial_b x^J \partial_c x^K d\sigma^a d\sigma^b$$

$$\mathcal{O}_\phi^{(3)} = - \int_{Y^3} x^a (\phi)$$

o local observables same  
as for G2-string

But no non-local operators  
corresponding to ours in G2-string  
[no localization on 3-cycles]

## CONCLUDING REMARKS

- analogy to topological A-model

string worldsheet  $\sim$  membrane world volume  
 holomorphic curves  $\sim$  associative cycles

A-model free energy

$$F = \sum_{g=0}^{\infty} g_s^{2g-2} F_g$$

What is the analogue of genus expansion?

$\rightsquigarrow$  classification of closed orientable 3-manifolds

A-model genus-zero free energy

$$F_0 = \int_X K \wedge K +$$

$$+ \sum_{\text{genus } g} dz \sum_{m=1}^{\infty} \frac{1}{m^g} e^{-m K|z|}$$

Contribution of constant maps?  
anticipate

$$F_{\text{const}}^{G_2} = \int_X d \wedge \star d$$

membrane instanton weights

$$e^{-\int_{Z_3} x^*(d)}$$

## OPEN QUESTIONS

- RELATION TO G2 - STRING?  
S-DUALITY?
  - B-MODEL FROM MEMBRANES?
  - COUPLING TO TOPOLOGICAL  
GRAVITY?  
ONE LOOP?
  - BACKGROUND 3-FORM FIELD?  
→ GENERALIZED CALIBRATIONS
  - DO-BRANE FORMULATION?  
MATRIX THEORY DESCRIPTION?
- CAN WE LEARN ABOUT  
FULL PHYSICAL M-THEORY?