

Background-1

Gromov-Witten Invariants

X = compact (almost) Kähler manifold

$X_{g,n,d}$ = space of stable curves in X

\uparrow genus
 \uparrow degree $\in H_2(X)$
 \uparrow number of marked pts p_1, \dots, p_n

The primary GW invariants are

$$\int_{X_{g,n,d}} \wedge \text{ev}_i^*(\alpha_i)$$

where $\alpha_i \in H^*(X)$, $\text{ev}_i: X_{g,n,d} \rightarrow X$
evaluation at p_i

Descendants. At p_i have tangent line to curve \rightarrow Line bundle L_i on $X_{g,n,d}$ with Euler class ψ_i

$$\int_{X_{g,n,d}} \wedge_i [\text{ev}_i^*(\alpha_i) \cdot \psi_i^{k_i}] = \text{descendants}$$

Background-2

Descendent Potential. Complete invariant is encoded in a generating function

$$\mathcal{D}_X = \exp \left\{ \sum_{g \geq 0} \hbar^{g-1} \mathcal{F}_g(X) \right\}$$

$$\mathcal{F}_g = \sum_{n,d} \frac{\mathcal{Q}^d}{n!} \int_{X_{g,n,d}} \wedge_{i=1}^n \left(\sum \text{ev}_i^*(a_k) \psi_i^k \right)$$

a function of the formal power series

$$\sum_{k \geq 0} a_k \psi^k \in H^*(X) [[\psi]]$$

This is conjectured (and sometimes known) to satisfy miraculous identities, such as a family of differential equations forming $\frac{1}{2}$ Virasoro algebra.

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Ancestors. A standard topologist's approach is to factor integration on $X_{g,n,d}$ via the push-forward to Deligne-Mumford space \overline{M}_g^n . **Problems:**

- No DM spaces for low g, n
- ψ on $X_{g,n,d} \neq \psi$ on \overline{M}_g^n .

One then defines the ancestor potential similarly to the descendent one (slight changes..) but using the DM ψ 's and existing DM spaces only.

Kontsevich-Manin solved the problem of relating descendents to ancestors in principle, in terms of genus 0 data.

Givental gave an explicit formula using his Fock space formalism.

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Givental's Fock space constructions

On $H^*(X)((\psi)) = H((\psi))$ have skew form

$$\Omega(f, g) = \oint \langle f(-\psi), g(\psi) \rangle d\psi$$

$H[[\psi]]$ is Lagrangian for that \Rightarrow can interpret \mathcal{D}_X as vector in Fock space.

That carries an action of the loop group

$$L_\psi GL(H) \cap Sp(H((\psi)))$$

Givental observed that many known relations about \mathcal{D}_X are expressible in terms of quantised quadratic Hamiltonians:

- Virasoro equations
- Genus 0 "Reconstruction Theorem" (Dubrovin; Dijkgraaf-Witten)
- Ancestor / Descendent relation

Backgrd 5

Flavour of formulae. Involve the "2-point gravitational descendent in genus 0"

$S_a(1/z)$ an $a \in H$ -dependent formal series in $\text{End}(H)[[1/z]]$ satisfying

$$S_a^*(-1/z) S_a(1/z) = \text{Id}$$

and defined by

$$(d, S_a \beta) = (d, \beta) + \sum_{n,d} \frac{Q^d}{n!} \int_{X_{0,n+2,d}} \text{ev}_1^*(\alpha) \cdot \prod_{i=2}^{n+1} \text{ev}_i^*(a) \cdot \frac{\text{ev}_{n+2}^* \beta}{z - \Psi_{n+2}}$$

Then, roughly

$$\mathcal{D}_X = \hat{S}_a^{-1} \cdot \mathcal{A}_{X,a}$$

where \mathcal{A} is the ancestor potential and \hat{S}_a = action of S_a on Fock space.

NB: There is a correction term from genus 1.

Backgrd 6

Higher genus Reconstruction.

Givental also conjectured (and proved in some cases) a formula for \mathcal{D}_X in terms of S_a in the case when the quantum cohomology of X is

generically semi-simple; essentially

$$\mathcal{D}_X = \hat{S}_a^{-1} \circ \hat{\Psi} \circ \hat{R}_a \quad (\mathcal{D}_{N \text{ points}})$$

where the constant matrix $\Psi \in \text{End}(H)$ and the "upper-triangular" loop group element $R_a \in \text{End}(H)[[\Psi]] \cap \text{Sp}$ are defined from S_a , using Dubrovin's theory of Frobenius structures.

Note: Givental Conjecture implies the Virasoro relations.

Tautological classes on \overline{M}_g^n

\overline{C}_g^n universal curve
 ↓
 \overline{M}_g^n

- has tangent line L_i at i th marked point
- Euler class $\psi_i \in H^2(\overline{M}_g^n)$
- has fibrewise canonical bundle K , $e(K) = -\psi$,
- $\kappa_i = \int_{\text{fibre}} \psi_i^{i+1} \in H^{2i}(\overline{M}_g^n)$.

The rational Mumford conjecture

(Madsen-Weiss) + Harer stability \Rightarrow

$$H^*(M_g) = \mathbb{Q}[\kappa_i] \text{ if } * < g/2$$

Older result of Loijenga (Bödigheimer - Tillman: over \mathbb{Z})

$$H^*(M_g^n) = H^*(M_g)[\psi_1, \dots, \psi_n]$$

in stable range

Structure of semi-simple FTFT's

- Smooth surfaces, parametrised boundaries:

$\widetilde{M}_g^{p,2}$ = moduli orbifold of such

$$\text{Need } \widetilde{Z}_g^{p,2} : A^{\otimes p} \rightarrow H^*(\widetilde{M}_g^{p,2}) \otimes A^{\otimes 2}$$

satisfying all gluing laws

Datum: $\widetilde{Z}_\infty \in H^*(M_\infty; A)$
 group-like primitive for the monoid structure of M_g

$$\widetilde{Z}_\infty = \exp \left\{ \sum_{i \geq 0} a_i \kappa_i \right\} \quad a_i \in A$$

and with $a_0 = \frac{1}{2} \sum_j \log \theta_j \cdot P_j$.

Note: a fixed $\Sigma_g^{p,1}$ gives $\sum_j \theta_j^{-g} P_j$
 $\alpha \rightarrow$

Formula for $\widetilde{Z}_g^{p,2}$:

$$A^{\otimes p} \xrightarrow{\text{mult.}} A \xrightarrow{\times \widetilde{Z}_\infty} A \xrightarrow{\text{comult}} A^{\otimes 2}$$

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Structure of FTFT's

Smooth surfaces, unparametrised ∂ :

Need $Z_g^{p,2}: A^{\otimes p} \rightarrow H^*(M_g^{p,2}) \otimes A^{\otimes 2}$

Data: \tilde{Z}_∞ as before

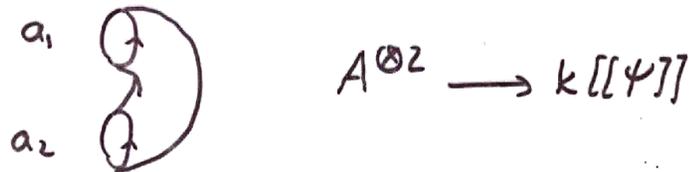
$E(\Psi): A \rightarrow A[[\Psi]]$ (= Id mod Ψ)

(Think of E as implementing the "bordism" from a point to a circle $\textcircled{1}$)

Apply $E(\Psi)$ to incoming states and $E(\Psi)^{-1}$ to outgoing ones

$$Z_g^{p,2}: A^{\otimes p} \xrightarrow{\prod_{i=1}^p E(\Psi_i)} A^{\otimes p} \xrightarrow{\tilde{Z}_g^{p,2}} A^{\otimes 2} \xrightarrow{\prod_{j=1}^2 E^{-1}(\Psi_j)} A^{\otimes 2}$$

Example: $Z_0^{2,0}(a_1, a_2) = \Theta(E(\Psi)a_1 \cdot E^{-1}(\Psi)a_2)$



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Structure of FTFT's

• Deligne-Mumford theories:

Need $\bar{Z}_g^{p,2}: A^{\otimes p} \rightarrow H^*(\bar{M}_g^{p,2}) \otimes A^{\otimes 2}$

Data: \tilde{Z}_∞, E as before

• $F: A \rightarrow A[[\Psi]]$ (= Id mod Ψ)

F describes the extension of Z -classes to the Deligne-Mumford boundary strata.



$Z = \text{Id}$

$(FE)^{-1}(\Psi_2) \circ FE(\Psi_1)$

A boundary stratum in \bar{M}_g has the form $M_{g_1}^1 \times M_{g_2}^1$. Inserting the above formula between the propagators $Z_{g_{1,2}}$ gives the restriction of \bar{Z} to that stratum.

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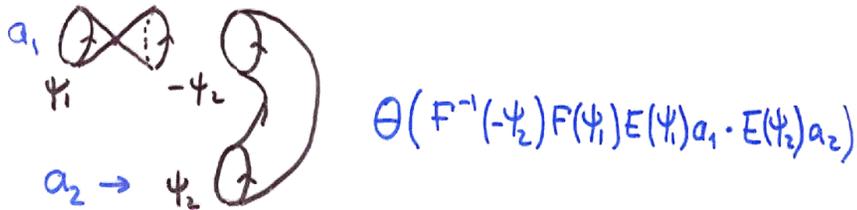
The Symplectic Condition

Theorem. On the data for DM theories, we have the single constraint

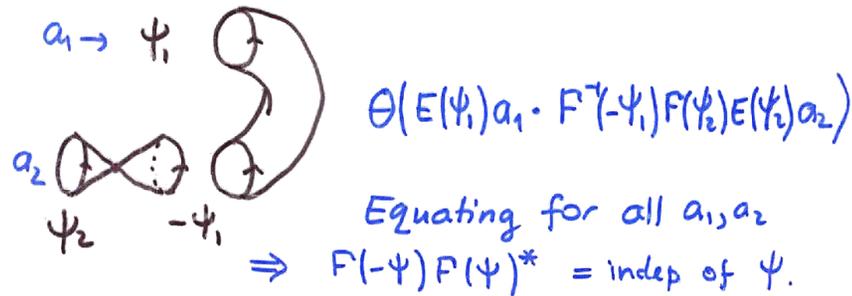
$$F(-\Psi) F(\Psi)^* = Id.$$

(So $F \in LGL(A) \cap Sp$)

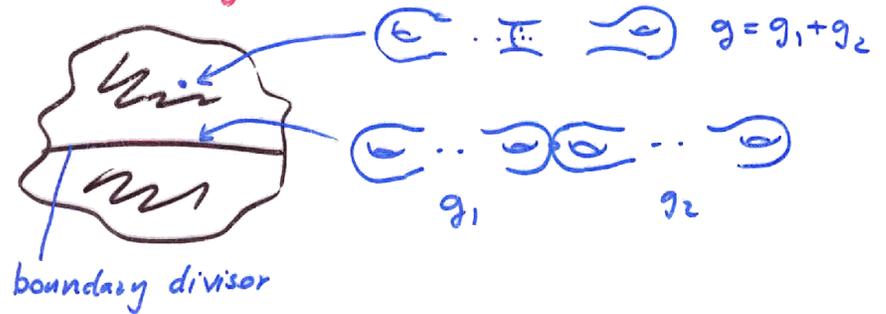
Proof: By pictures



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Extending Z over boundary DM strata



From \tilde{Z} , E and F we know the restriction of \bar{Z} to the "bulk" and to the boundary divisor.

The Euler class of the normal bundle is $\psi_1 + \psi_2$.

If either g_1 or g_2 is large, this is **not** a zero-divisor through a range of degrees \Rightarrow Patched class is **Unique** (In the range)

Stabilisation trick allows us to increase genus w/o loss of information

Large genus stabilisation

- When the surface has a parametrised boundary, this is obvious:

Attach a large genus surface to embed M_g to M_{g+c} .

Propagator α^G is invertible in A
 \Rightarrow no loss of information.

(Note: α invertible $\Leftrightarrow A$ semi-simple)

- When the surface has only a marked point (= unparametrised boundary): attach a "moving" surface of large genus G .

Resulting moduli space: circle bundle over $M_g' \times M_G'$, with Euler class

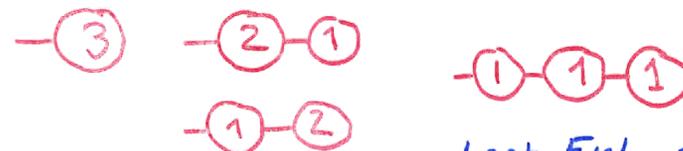
$\psi_1 + \psi_2$. But ψ_2 is a free generator!
 \Rightarrow No loss of information.

Inductive extension of Z to \overline{M}

The Euler class observation is the basis of an inductive procedure showing uniqueness of \overline{Z} from our data.

We single out a marked point for stabilisation and take care that each Euler class involves a ψ on that irreducible component.

Bad example: (4 strata in \overline{M}_3')



Last Euler class is $\boxed{0}$ -divisor even after stabilisation!!

Correct: First glue $\text{---} \textcircled{3} + \text{---} \textcircled{2} \text{---} \textcircled{1}$
 Then glue $\text{---} \textcircled{1} \text{---} \textcircled{2} + \text{---} \textcircled{1} \text{---} \textcircled{1} \text{---} \textcircled{1}$. Then glue together.

Stratification of \overline{M}_2^1

Attaching order	d	graph	substrata
0	4		—
1	3		—
2	2		
2	2		—
4	1		
4	1		—
6	0		

Stratification of \overline{M}_2^2

0	1	2
2	4	4
6	6 (d=1)	6
9	9 (d=0)	9