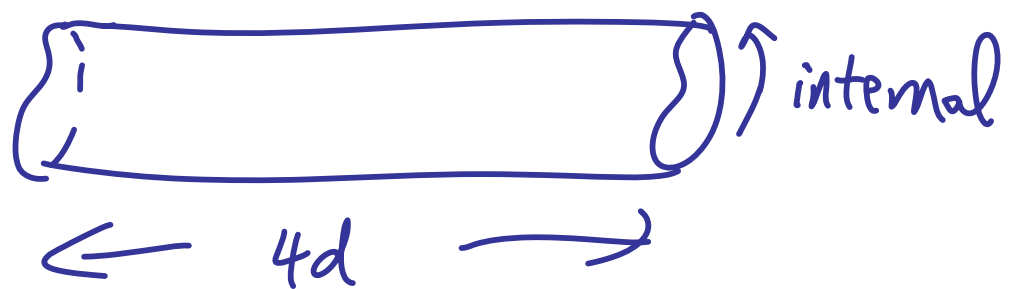




Toward 4d quantum gravity  
in string theory

Work in Progress with Joe Polchinski

We'd like a non-perturbative formulation of 4d physics:



with a hierarchy of energy scales

$$m_{\text{string}} \gg \frac{1}{L_{4d}}$$


AdS/CFT (and previously, BFSS matrix theory) formulate some backgrounds non-perturbatively, but did not (yet) get down to 4d :

BFSS : 11d  $\leftrightarrow$  N D0-brane Q.M.

4d (max susy)  $\leftrightarrow$  D7-branes on  $T^7$   
 $\hookrightarrow$  codim 2  $\rightarrow$  log potential  $\rightarrow C_{N \leq 24}$

AdS/CFT: ①  $AdS_2 \times S^2 \times CY$   
 want  $L_{AdS} \rightarrow \infty$  small  $\checkmark$   
 $\hookrightarrow$  IR divergences in  $AdS_2$

AdS/CFT (2)  $AdS_4 \times \left\{ \begin{array}{l} S^7 \quad (M) \\ S^2/\mathbb{Z} \\ CP^3 \quad (IIA) \end{array} \right.$

  
 $L_{\text{internal}} \sim L_{\text{AdS}}$

No hierarchy of scales in Freund-Rubin compactifications.

Basic reason: In 11/10d Einstein

equations  $\underbrace{R_{MN} - \frac{1}{2} R G_{MN}}_{\text{Internal + 4d}} = 8\pi G \underbrace{T_{MN}}_{\text{flux}}$

all three contributions are of the same order in the solution

For future reference, let us reproduce this in the language of the 4d effective potential energy:

$$S = \int d^{10}x \frac{\sqrt{G}}{g_s^4} \left( \frac{R}{g_s^2} + F_p^2 + \dots \right)$$

→ 4d potential energy

$$U_4 = \frac{1}{g_s^2} \left\{ - \int d^6x \frac{\sqrt{G_6}}{g_s^2} R^{(6)} + \int d^6x \frac{\sqrt{G_6}}{g_s^2} F_p^2 \right\}$$

$$\sim \frac{R^6}{g_s^2 g_s^2} \left( - \frac{1}{R^2} + g_s^2 \frac{Q^2}{R^{2p}} + \dots \right)$$

where  $R \equiv$  size in string units.

$$U_4 \sim M_p^4 \left( \frac{g_s^2}{R^6} \right) \left( - \frac{1}{R^2} + g_s^2 \frac{Q^2}{R^{2p}} + \dots \right)$$

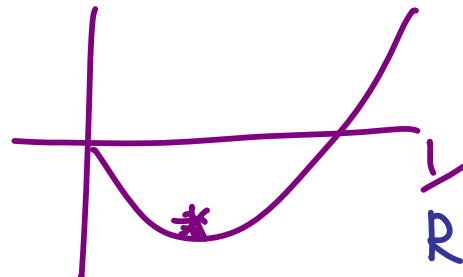
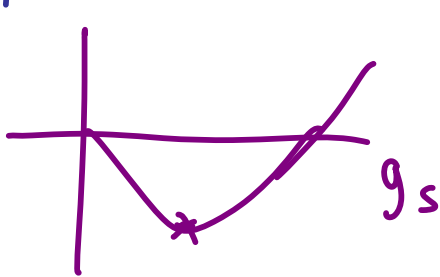
in Einstein frame

↑  
 $U_R$

↑  
 $U_{F_p}$

e.g. IIA on  $\mathbb{C}P^3$  w/  $F_6$  &  $F_2$  :

$$\frac{U_4}{M_p^4} \sim -\frac{g_s^2}{R^8} + \frac{g_s^4 Q_2^2}{R^{10}} + \frac{g_s^4 Q_6^2}{R^{18}}$$



Equivalently

$$U_4 \sim M_p^2 \left( -\frac{1}{R^2 g'} + \frac{g_s^2 Q^2}{R^{2p} g'} + \dots \right)$$

$\sim \Lambda^2 \frac{1}{R_{\text{AdS}}^2 g'}$

i.e.  $\Lambda_{\text{min}} \sim \Lambda_R$  in Freund-Rubin

- n.b.  $\text{CFT}_3$  dual to IIA/ $\mathbb{C}P^3$  is  $\left\{ \begin{array}{l} \cdot \text{IR limit of } D2, D6, \text{ KK} \\ \cdot \text{strongly coupled CS} \end{array} \right.$   
 or M thy /  $S^2/\mathbb{P}^1$   
 Schwarz, BL, ABJM, et seq.

In fact there is a large class of 3d CFTs obtained via RG flow from gauge theory

with flavors

Appelquist/Heinz HET  
Sachdev... CNT

$$\frac{1}{g^2} = \underbrace{\frac{1}{g_0^2}}_{\text{classical}} + \underbrace{\frac{\Delta N}{E}}_{\text{1-loop screening}} + \dots$$

(In 3d  $N=4$ ,  $\Delta N = N_f - 2N_c$   
and this is exact)

→ dimensionless coupling has fixed point

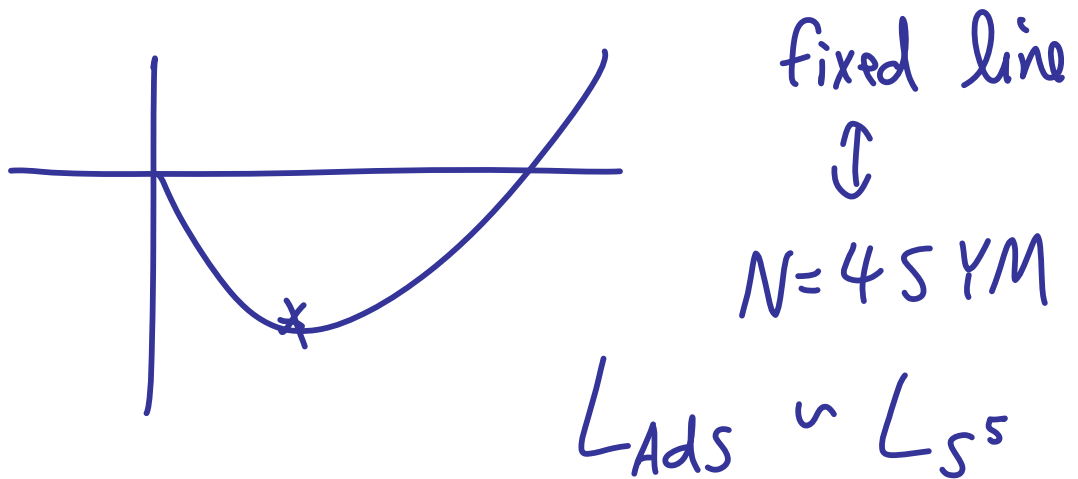
$$\frac{1}{\bar{g}^2} = \left. \frac{E}{g^2} \right|_{E \rightarrow 0} = \Delta N + \dots$$

controlled at large  $N_f$  independently of SUSY.

e.g. D2-D6, D2-D6-orbifold; Hanany-Witten

... examples: Still Freund-Rubin (or stringy)

Similarly, compactification  
on  $S^5$  w/  $F_5 \rightarrow AdS_5 \times S^5$



---

In general, the negative  
term(s) in the potential are key

e.g. {  
positive curvature  
0-planes  
 $-3|W|^2$



Coming from the other direction,  
we can construct

$$(A)dS_4 \times X_{\text{small}}$$

in an apparently large number  
of ways  $B^2$  DRS BP MSS GKP + KKLT...

suggesting a rich set of dual  
CFT<sub>3</sub>s.

- Not a priori realized as near-horizon limit of brane system
- Can read off interesting properties:

$$N_{\text{d.o.f.}} \sim L_{(A)dS}^2 M_p^2 \leq N \quad \begin{array}{l} b \leftarrow \text{betti \#} \\ \leftarrow \text{flux \#} \end{array} \quad \begin{array}{l} ES \\ AAB \end{array}$$

Plan: Start from known,  
Freund-Rubin dual pair:

$AdS \times S \leftrightarrow QFT$   
(brane construction)

add ingredients  
which cancel or  
nearly cancel

$\leftrightarrow$

additional  
field content,  
couplings of  
QFT

$U_q$

stabilize the  
moduli  $\rightarrow AdS_4$

$\leftrightarrow$

$CFT_3$

7-branes contribute to  $U$  naively as

$$U_7 \sim M_p^4 \left( \frac{g_s^2}{R^6} \right)^2 \cdot \left( \underbrace{\tau_7}_{\substack{\text{tension in} \\ \text{string units}}} \cdot R^4 \right)$$

compare to curvature energy

$$U_{\text{curv}} \sim M_p^4 \left( \frac{g_s^2}{R^6} \right)^2 \cdot R^6 \cdot \left( -\frac{1}{R^2 g_s^2} \right)$$

$\Rightarrow$  for  $\tau_7 \sim \frac{1}{g_s^2}$ , i.e.  $(p, q) 7Bs$ ,

they compete.

cf Aharony  
Fayazuddin  
Maldacena

Of course  $7Bs$ , being codimension 2,  
have large IR back reaction...

The interplay between curvature  
 & 7-brane energy is accurately  
 captured using the techniques  
 of F-theory :

Vafa '96

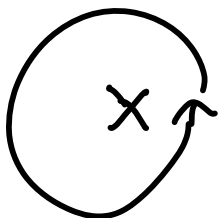
$$T^2 \rightarrow X$$

$$\downarrow$$

$$B$$



$$\tau_{T^2} = C_0 + \frac{i}{g_s} \text{ in } \mathbb{H}^2$$



D7-brane:  $\tau(u) \sim \frac{1}{2\pi i} \log u$

$\tau \rightarrow \tau+1$  monodromy

$(p, q)$  7-brane  $\tau \rightarrow \frac{(1+pq)\tau - p^2}{q^2\tau + (1-pq)}$

Plan: Start from known,

Freund-Rubin dual pair:

$AdS \times S \leftrightarrow QFT$   
(brane construction)

add  $7$ -branes  
which cancel or  
nearly cancel

$U_q$

$\leftrightarrow$  additional  
 $flavors^*$   
couplings of  
 $QFT$

stabilize the  
moduli  $\rightarrow AdS_4$

$\leftrightarrow$   $CFT_3$

\* in general, both electric  
& magnetic cf Douglas/Shenker,  
Argyres/Douglas, Argyres-Plesser-Seiberg-Witten

# D3, D7 and Electric/Magnetic Matter

- 4d  $N=2$   $SU(2)$  SYM w/  $N_f$  hypermultiplets

Seiberg -  
Witten  
solution

monopole  
⊗

dyon  
⊗

$u$  (Coulomb branch)

⊗ ← quark

AD/APSW : can change mass matrix

$M$  such that

mutually nonlocal  
matter is light.

monopole  
⊗

dyon + quark  
⊗

⊗

- In brane constructions (Sen, Banks, Douglas, Seiberg, ...)

$u \leftrightarrow$  D3 position

⊗  $\leftrightarrow$  7B position

The  $T^2$  varying over  $B$

can be described as

Vafa  
Morrison-Vafa  
Kachru Intriligator  
Morrison Vafa

$$y^2 - x^3 - x f(u) z^4 - g(u) z^6 = 0$$

↑                      ↗  
coordinates on  $B$

i.e. as a degree 6 hypersurface  
in  $WP^2(2,3,1)$ .

For a Kähler base  $B$ , one can  
formulate the  $T^2$  fibration  $T^2 \rightarrow X$

as a hypersurface in  $B \times WP^2(2,3,1)$ ,

and as the target space of a

(2,2) gauged linear  $\sigma$ -model written

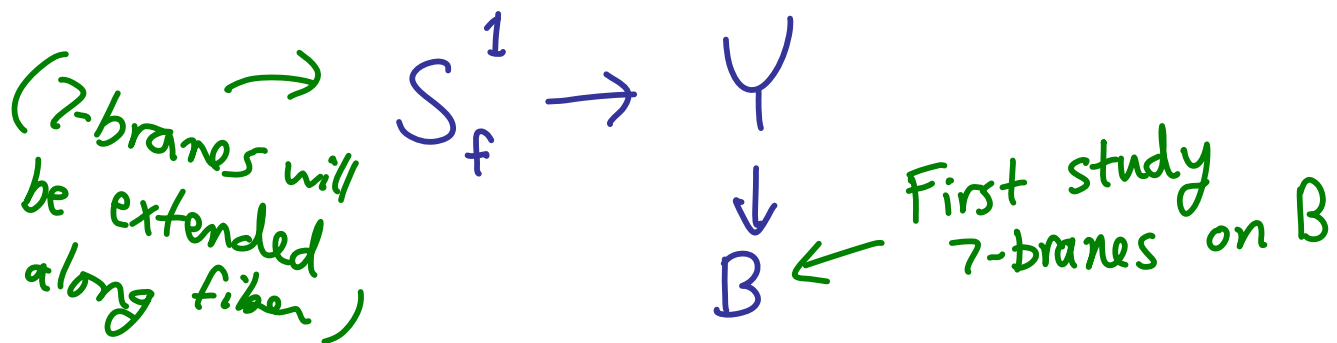
$7$ -branes live at the locus

$$\Delta = 27g^2 + 4f^3 = 0$$

Let us start with  $\mathbb{I}B$  on

$$Y_5 \times S^1 \quad \text{with} \quad \int_Y F_5 = N_c$$

where  $Y_5$  is an  $S^1$  (Hopf) fibration over a Kähler base



Examples:  $Y = S^5, B = \mathbb{C}P^2 \text{ or } \mathbb{W}P^2$

Topologically  $S^2 \times S^3$

}	$Y = T^{1,1}$	$B = \mathbb{C}P^1 \times \mathbb{C}P^1$ Klebanov / Isevtin / Strass
	$Y = Y_{pq}$	$B = \dots$ Gauntlett Martelli Sparks Waldram
	$Y = L_{abc}$	$B = \dots$ Cretic Liu Pope



$$S_f^1 \xrightarrow{\text{size } R_f} Y \xrightarrow{\text{size } R} B \cong B \text{ with "metric flux" } F_{\text{met}} = J_B$$

(useful description for small fiber)

Two classes

of candidate examples:  $Y \times S^1_{\perp} \times (4d)$

① 7-branes nearly cancel  $U_R^{(B)}$

$$U \sim M_p^4 \left( \frac{1}{R^4 R_6 R_f} \right) \left( \frac{R_f^2}{R^4} - \frac{\epsilon}{R^2} + \frac{N_c^2}{R^8 R_f^2} + \frac{Q_1^2}{R_{\perp}^2} + \frac{Q_3^2}{R_{\perp}^2 R^4} \right)$$

$g_s \sim 1$

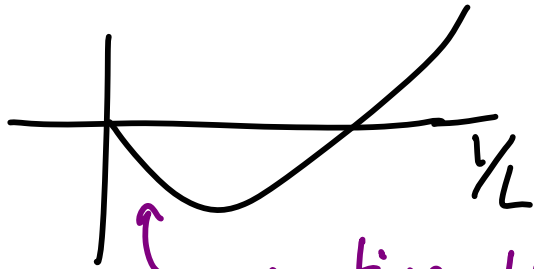
enforce with e.g.

$E_n$  7-branes

$\hookrightarrow$  stable minimum with

$$R_f \ll R \ll R_{\text{Ads}}$$

②  $U_R + U_7 = 0$  (F-theory on CY)



negative term from 0-planes

→ again  $R_{\text{Ads}} \gg R \gg R_f$

\* Are 7-brane moduli tachyonic?

On  $S^5$



allowed tachyon for  $R \sim R_{\text{Ads}}$   
but what about  $R \ll R_{\text{Ads}}$ ?

Note that 7-brane moduli

$$y^2 - X^3 - x \overbrace{f(u)} z^4 - g(u) z^6 = 0$$

are flat directions in the CY  
and nearly flat for  $U_R \propto \frac{\Sigma}{R^2}$

To be specific, consider

$$Y = S^5 \quad (\text{topologically})$$

Start from the  $\mathbb{C}P^2$  model:

(2,2) chiral multiplets  $U_1, U_2, U_3$

$U(1)$  Gauge symmetry

$$(U_1, U_2, U_3) \cong e^{2\pi i \varphi} (u_1, u_2, u_3)$$

$$\Rightarrow D^2 = \left( |u_1|^2 + |u_2|^2 + |u_3|^2 - R^2 \right)^2$$

$\hookrightarrow D=0$  alone gives  $S^5$

$\varphi$  parameterizes  $S^1$  fiber

$$S^1_f \rightarrow \begin{matrix} S^5 \\ \downarrow \\ \mathbb{C}P^2 \end{matrix}$$

$$ds^2_{S^5} = d\sigma_{\mathbb{C}P^2}^2 + R_f^2 (d\alpha + A)^2$$

$$dA = J$$

Gibbons, Pope

To add the 7-branes, want a  $T^2$  fibration over  $B = \mathbb{C}P^2$

Gauged Linear  $\sigma$ -model becomes

		$u_1$	$u_2$	$u_3$	$X$	$Y$	$Z$	$P$
$T^2$	$\{ U(1) \}$	0	0	0	2	3	1	-6
$\mathbb{C}P^2$	$\{ U(1) \}$	1	1	1	$g_x$	$\frac{3}{2}g_x$	0	$-3g_x$

$$S_w = \int d^2\sigma d^2\theta P \left( y^2 - x^3 - f(u) x z^4 - g(u) z^6 \right)$$

Now,  $\sum_{\text{fields } I} g_I = 0$  is the Calabi-Yau condition (ensuring anomaly-free  $U(1) \times U(1)$  R-symmetries appropriate to (2,2) SCFT) written

The running of  $R^2$  in

$$D^2 = \left( |u_1|^2 + |u_2|^2 + |u_3|^2 + g_x |x|^2 + \frac{3}{2} g_x |y|^2 - |p|^2 - R^2 \right)^2$$

is  $M \frac{\partial R^2}{\partial M} \sim \sum_I g_I$

Now  $\sum g_I = 3 - \frac{1}{2} g_x$

and the degree of  $G = y^2 - x^3 - fxz^4 - gz^6$

is  $\deg_G = \deg_g = 3g_x = 18 - 6 \sum g_I$

Fully canceling curvature energy

means  $\sum g_I = 0 \Rightarrow \deg_g = 18$

$\Rightarrow \deg \Delta = 36 \Rightarrow 36$  7-branes

(This agrees with naive result from  $U_7 \sim \dots (r_7 \times \text{vol}) r$ )

- The 7Bs are extended along  $U(1)$  fiber
  - Similarly, on  $\mathbb{C}P^1 \times \mathbb{C}P^1$ , 48 7Bs cancel  $U(1)$ .
- 

To stabilize, need negative term(s) in the potential.

→ introduce orientifold (mutually susy)

Altogether, take  $S_f^1 \rightarrow S^5 \times S^1_{\perp}$   
 $\downarrow$   
 $\mathbb{C}P^2$

with also worldvolume flux  $g$  on  $S_f^1 \times S^1_{\perp}$   
 (+ branes which compensate tadpole...)

$$U \sim \frac{M_p^4}{R_f R_{\perp} R^4} \left( \frac{R_f^2}{R^4} - \frac{1}{R_f R_{\perp} R^2} + \frac{N_c^2}{R^8 R_f^2} + \frac{g^2}{R_{\perp}^2 R_f^2 R^2} \right)$$

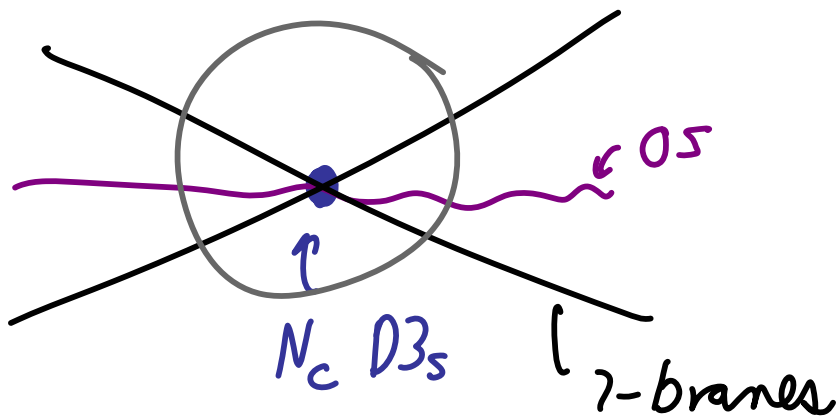
→ minimum, with hierarchy

$$R \sim (g^2 N_c)^{1/4} \quad R_F \sim \frac{R}{g}$$

$$R_{\perp} \sim \frac{g^3}{R} \quad R_{\text{Ads}} \sim g R$$

---

QFT side?



A complicated generalization of the D3 + D7 brane constructions for Seiberg/Witten, Argyres/Douglas, Plesser, SW

Now generalize to cases where we do not fully cancel the curvature energy. Consider

$T^2$  fibration over  $B = \underline{W}P^2$

Gauged Linear  $\sigma$ -model

		$u_1$	$u_2$	$u_3$	$x$	$y$	$z$	$P$
$T^2$	$\{ U(1) \times$	0	0	0	2	3	1	-6
$\underline{W}P^2$	$\{ U(1)$	<u><math>w_1</math></u>	<u><math>w_2</math></u>	<u><math>w_3</math></u>	$g_x$	$\frac{3}{2}g_x$	0	$-3g_x$

$$S_w = \int d^2\sigma d^2\theta P \left( y^2 - x^3 - f(u) x z^4 - g(u) z^6 \right)$$

Again  $\beta_{R^2, \text{full}} \quad \sum g = \sum w - \frac{1}{2}g_x$   
 in the full system including the 7Bs.



Again  $\beta_{R^2, \text{full}} \sim \Sigma g = \Sigma W - \frac{1}{2} g_x$   
in the full system including the 7Bs.

For  $WP^2$  alone,

$$\beta_{R^2, WP^2} = \Sigma W$$

$\Rightarrow$  IF  $\Sigma W - \frac{1}{2} g_x \ll \Sigma W$

then we almost cancel the curvature.

$$\rightarrow \mathcal{U}_{R, \text{full}} \sim M_p^4 \left( \frac{g_s^2}{\text{Vol}} \right)^2 \cdot \text{Vol.} \left( -\frac{\Sigma}{R^2} \right)$$

with  $\Sigma \sim \frac{\Sigma W - \frac{1}{2} g_x}{\Sigma W}$

(using the NLSM result  $\beta \sim R_{MN}$ )

\* What about the singularities of  $WP^2(w_1, w_2, w_3)$ ?

$$(u_1, u_2, u_3) \cong (\lambda^{w_1} u_1, \lambda^{w_2} u_2, \lambda^{w_3} u_3)$$

$$\rightarrow \text{for } \lambda = e^{\frac{2\pi i}{w_1}} \in U(1)$$

$(u_1, 0, 0)$  is a (generally non-SUSY)  $\mathbb{C}^2 / \mathbb{Z}_{w_3}$  orbifold fixed point.

By itself\*, this has twisted tachyons which condense, smoothing it out

Adams  
Polchinski  
ES ...  
.. Morrison..

★ Don't Panic! • Must include 7-branes  
can restore SUSY

• and  $S^1$  fiber : Removes  $U(1)$  projection entirely

From  $WP^2$  we get candidate examples of ...

① 7-branes nearly cancel  $U_R^{(B)}$



$$U \sim M_p^4 \left( \frac{1}{R^4 R_6 R_f} \right) \left( \frac{R_f^2}{R^4} - \frac{\epsilon}{R^2} + \frac{N_c^2}{R^8 R_f^2} + \frac{Q_i^2}{R_\perp^2} \right)$$

$g_s \sim 1$

enforce

with e.g.

$E_n$  7-branes

$\hookrightarrow$  stable minimum with

$$R_f \ll R \ll R_{\text{Ads}} \sim R_\perp \frac{Q_i^*}{Q_i}$$

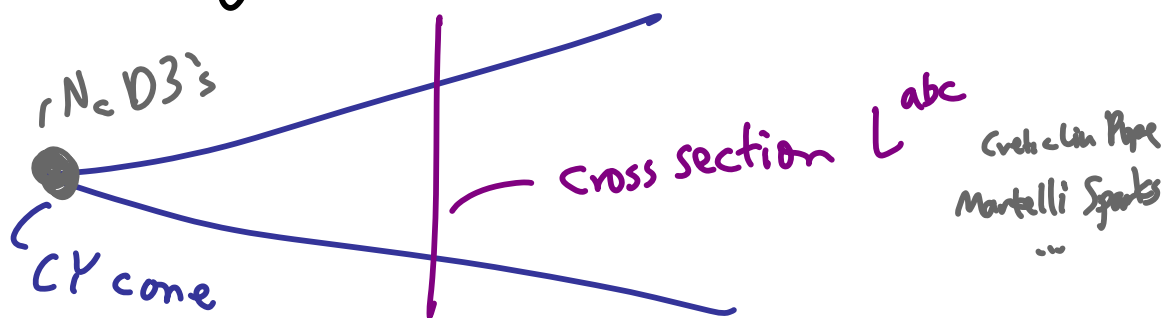
... as long as all other directions are stable (at most allowed tachyons)

\* Ok maybe F is for Five ...

... but if replace  $\frac{Q_i^2}{R_\perp^2}$  with  $\frac{Q_3^2}{R_\perp^2 R^4}$ , full hierarchy

$\rightarrow S^2 \times S^3$  examples

More general candidate examples:



- Start from a Calabi-Yau cone, given by GLSM with charges  $(c, atb-c, -a, -b)$  ( $\sum \text{charges} = 0$ )
- Mod out by an additional  $(2,2) U(1)$  to pick out 4d base in  $S'_f \rightarrow L \downarrow B$
- Add 7-branes as before, imposing  $\beta_{R^2_B}, \text{ with } 7Bs \ll \beta_{R^2_B} \text{ alone}$
- Undo a  $(1,1) U(1)$  to get  $S'_f$  fibers
- On the full cone, this gives our brane construction for the QFT

• On  $Y^{pq}$   $S_f^1 \rightarrow Y^{(pq)}$  units of metric flux  
 $B_4 \cong S^2 \times S^2$

$$ds^2 = ds_{B_4}^2 + R_f^2 (d\chi + A)^2 \quad dA \sim J_4$$

The Naive 7-brane potential energy  
 for 7-branes wrapped on SUSY cycles

$$\Sigma_{1,2} \sim S^3 / \mathbb{Z}_p \quad n_3 < n_4$$

$$\left\{ \begin{matrix} 3 \\ 4 \end{matrix} \right\} \sim S^3 / \mathbb{Z}_{p \pm q} \quad n_1 + n_2 + n_3 + n_4 = 40$$

$$\rightarrow \epsilon \sim \frac{q}{p} \text{ which } \ll 1 \text{ for } p \gg q$$

\* still need to check against F-theory on  $B_4$

## The string coupling

7-branes corresponding to mutually non-local flavors have  $g_s \sim 1$

⊗ e.g.  $T \sim e^{\frac{i\pi}{3}}$   $f(u) = 0$  branch  
Dasgupta/Mukhi

In a (near-) SUSY background,

the (approximate-) moduli are

- $R^2$  (GLSM D-term)
- polynomial coefficients in  $f, g$

(Other modes of the metric + dilaton have KK-scale masses  $\sim \frac{1}{R\sqrt{g_i}}$ .)

In the LSM,  $f + g$  are superpotential couplings  $\Rightarrow$  don't run even at  $\mathcal{O}(\epsilon)$   
 $\Rightarrow$  expect  $|m^2| \leq \mathcal{O}(\epsilon^2)$

## Sen Limit


In F theory,  $\exists$  limit (Sen)

$$\begin{aligned} f &= -3h^2 + \epsilon\eta \\ g &= -2h^3 + \epsilon h\eta - \epsilon^2 \frac{\chi}{12} \end{aligned} \quad \epsilon \rightarrow 0$$

for which  $g_s \rightarrow 0$ . i.e. all the  $(p, q)$  7-branes boil down to  $O7$ -planes +  $D7$ -branes

Such examples, if they can also be stabilized (including  $g_s$ ), would be purely electric on the QFT side.

## Remaining issues

- ~~SUSY~~ scale ( $\ll \frac{1}{R}$ );  (non-) perturbative stability in all directions
- F-theory vs Naive Potential energy in general (agree for  $\mathbb{C}P^1 \times \mathbb{C}P^1, \mathbb{C}P^2$ )
- Analysis of QFT side/brane construction
- Sen limit examples

...



## Future directions

•  $dS_4$  ?



slightly over-cancel  
curvature energy

↳ Now no tachyons are allowed.

• Giddings - Kachru - Polchinski + KKLT

Use 7-branes & F-theory on a positively curved base. Interpret as in above examples?

• Other applications (e.g. AdS/CMT?)

...