

Friction law and hysteresis in granular materials

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Granular materials (T=0, hard)



Andreotti, Forterre, Pouliquen



Poppy seeds, Chicago group

- Solid and Liquid phases separated by a jamming transition
- Mesoscopic scale: collective effects important
- Local scale: interaction law between grains complicated
- Rich interplay between these two scales (e.g. shear thickening)

Hysteresis at the jamming transition



Hysteresis and Velocity weakening



Kuwano, Hatano 2013,

Microscopic theory?

Importance of hysteresis in particulate materials

- Hysteresis: once it starts, flow is fast
- Landslides
- Earthquakes



Granular matter Zooming in a fault

Fault that velocity weaken prone to earthquake

 $\frac{\mu_{start} - \mu_c}{\mu_c} \sim 10$

Observations to explain in granular matter

- Inertia favors hysteresis
- Friction appears necessary *Peyneau roux 08*
- Acoustic vibrations can eliminate hysteresis



Dupont et al, 2003



Lastakowski et al 2015, Johnson and Jia 05

Dimensional analysis and constitutive relations

GDR Midi, Roux, Radjai, Lemaitre, Pouliquen, Forterre,...

- Δ = deformation at contacts / diameter ~ 10⁻⁵
- Rigid limit: if limit $\Delta \rightarrow 0$ not singular, dimensional analysis implies one dimensionless number controls flow $\dot{\epsilon}: strain \ rate$

$$\mathcal{I} = \dot{\epsilon} D \sqrt{\rho/p}$$

 ϵ : strain rate D: particle diameter ρ : mass density

- $\mathcal{I} < 0.1\,$ dense flows. Most numerical studies report $\mathcal{I} > 10^{-3}\,$

(show data)

Velocity strengthening:

$$\mu \approx \mu_c + c_0 \mathcal{I}$$

Simulations

- Static and dynamic friction coefficient identical $\,\mu_f$ Restitution coefficient e Small N to avoid shear bands 15 <u>×10</u>-3 $\times 10^{-3}$ $- - \Delta = 10^{-3.8}$ $N = 10^{2.2}$ (c) 10 $\Delta = 10^{-2.8}$ $N = 10^{2.6}$ 10 $- \Delta = 10^{-1.8}$ $= 10^{3}$ $\mu - \mu(I_0)$ $\mu - \mu(I_0)$ 5 5 0 0 (b) -5 -5 10⁻³ 10⁻² 10^{-4} 10⁻³ 10^{-4} 10^{-2} Ι
- Effects disappear as Δ increases (or friction removed)

Two ideas

Near jamming, dense
 Network of contacts



1/ Collisions induce elastic vibration of the network. These vibrations make some contacts slide. *Melosh 80's*



2/ soft spots in granular flows (~ dislocation in metals)? Contact near the coulomb cone.

Macroscopic friction and fraction of sliding contacts

• Gedanken experiment: Add vibrations to sample.



- Fraction χ of contacts near Coulomb Cone will slide in average
- Flow threshold will decrease $~~ \tilde{\mu}_c = \mu_c (1 g(\chi))$ With ~g(0) = 0
- Same if noise is endogenous, leading to:

$$\mu(\mathcal{I}) = \mu_c(1 - g(\chi)) + c_0 \mathcal{I}$$

- Below: use and test model with $\ g(\chi) = b\chi$ (b interaction dependent)

Macroscopic friction and fraction of sliding contacts



- Phenomenological Model agrees well with data
- Limit $\Delta \rightarrow 0$ singular for χ !
- $\chi(\mathcal{I},\Delta)\sim \chi(\mathcal{I}\Delta^{-0.25})$??

Microscopic observables in granular flows

Simple model of frictionless suspensions



Lerner et al, PNAS 2012

- Velocity fluctuations
- Strain scale ϵ_v

$$\mathcal{L} = \frac{V_r}{\dot{\epsilon}D}$$

Extracting Key microscopic properties

DeGiuli et al, PRE 2015

$$C(\epsilon) = \langle \vec{V}_r(0) \cdot \vec{V}_r(\epsilon) \rangle$$

• Captures velocity fluctuations:

 $C(0) = \langle \vec{V}_r^2(0) \rangle \sim \mathcal{L}^2$

appears to diverge near jamming

• Captures ϵ_v (where memory is lost, function drops off)





Menon Durian 97 DeGiuli et al, PRE 2016

Idea behind $\epsilon_v \sim \mathcal{I}$

DeGiuli et al, PRE 2016

• Network almost force balance



F grows in between collisions

- Newton equation: velocity increases exponentially. Characteristic strain ${\mathcal I}$
- Collision must occur on that strain scale in a stationary state.

Energy balance

DeGiuli et al, PRE 2016

 $\Omega\sigma\dot{\epsilon}$: power injected = power dissipated \mathcal{D}_{tot}

$$\mathcal{D}_{tot} = \mathcal{D}_{col} + \mathcal{D}_{sliding}$$

$$\mathcal{D}_{col}$$
 = N * (collisional rate) * (kinetic energy of particles)
 $\sim N \times (\dot{\epsilon}/\epsilon_v) \times (m\mathcal{L}^2\dot{\epsilon}^2)$

 $\mathcal{D}_{sliding}$ = (number of sliding particles)* (transverse force) *(velocity) $\sim (N\chi) imes (\mu_f p) imes (\mathcal{L}\dot{\epsilon})$

 χ : Fraction of sliding contact μ_f : friction coefficient

Sliding dissipation dominates in dense flows



$$\mathcal{D}_{tot} \approx \mathcal{D}_{col} \quad \rightarrow \quad \mu_p \mathcal{L} \chi \sim 1$$

All quantities $\, {\cal L}, \epsilon_v, {\cal D}_{col}\,$ can be expressed in terms of ${\cal I}, \chi$

Estimating the mechanical noise δf

- vibrational energy $\, \sim au \, D_{col}$
- As to be compared to potential energy $\sim \Delta^2$
- -> explains why rigid limit is singular
- Finally one finds

$$\tilde{\delta}f \equiv \frac{\delta f}{pD^{d-1}} \sim (\mathcal{I}\Delta^{-1/4})/\chi$$



Estimating $\,\chi$ from the mechanical noise



$$\chi \sim \int_0^{\delta f} P(x) dx$$

Close the problem

$$\chi \sim (I/\Delta^{1/4})^{\alpha},$$

$$\mathcal{L} \sim (I/\Delta^{1/4})^{-\alpha},$$

$$\mathcal{R} \sim (I/\Delta^{1/4})^{\gamma},$$

$$\delta f \sim (I/\Delta^{1/4})^{\beta}$$

Further tests



Conclusion

- Hysteresis emerges as collective effect, even if not present at contact level
- Mechanical noise that lubricate contacts, explains main observations
- Role of Delta testable
- Key role of contacts
- close to the coulomb cone
- Suspensions?
- Earthquakes?