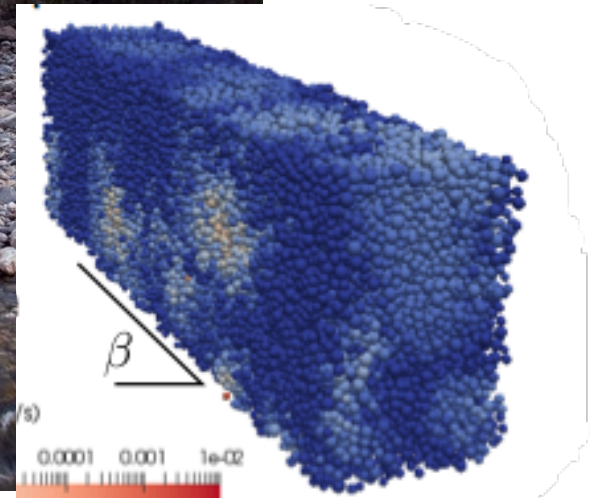
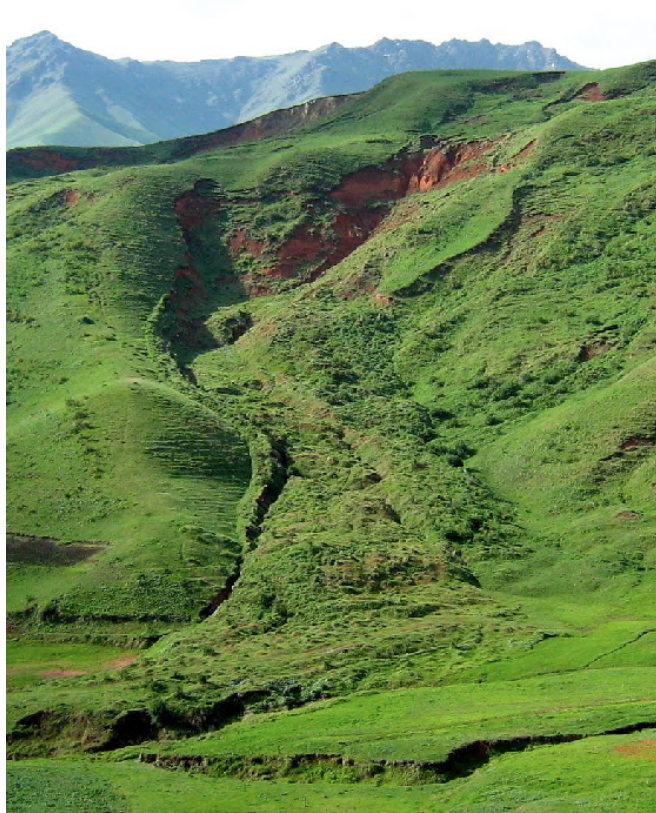


# Glassy dynamics of landscapes

KITP, 2 March 2018

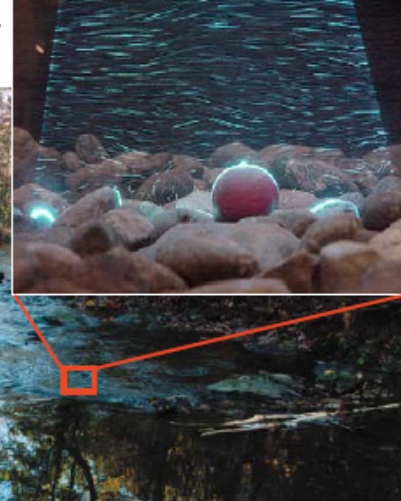
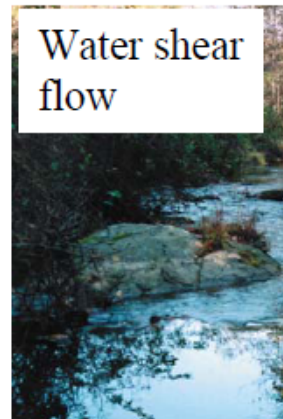


**Douglas Jerolmack**<sup>1,2</sup>, **Behrooz Ferdowsi**<sup>3</sup>, **Carlos Ortiz**<sup>4</sup>, **Morgane Houssais**<sup>5</sup>, **Colin Phillips**<sup>6</sup> and **Douglas Durian**<sup>7</sup>

1. PennSeD, Earth & Environmental Science, Univ. Pennsylvania
2. Mechanical Engineering & Applied Mechanics, Univ. Pennsylvania
3. Geosciences, Princeton Univ.
4. Deloitte
5. Levich Institute, City University of New York
6. Civil & Environmental Engineering, Northwestern Univ.
7. Physics & Astronomy, Univ. Pennsylvania

# Earth's surface is a granular-fluid interface

Shape of landscapes → balance of forces

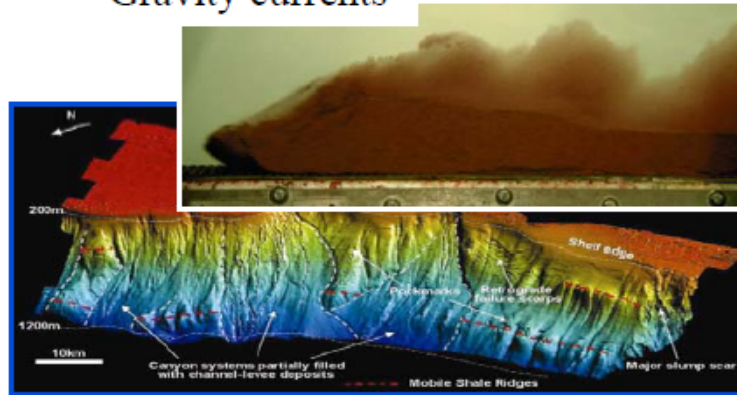
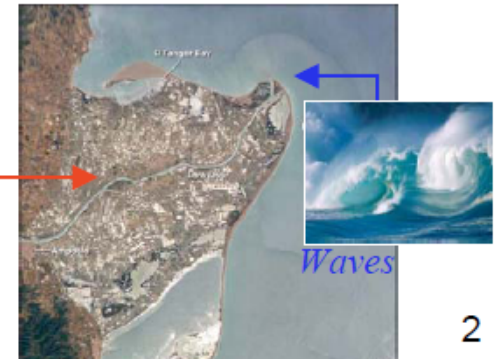


"Hillslope" processes

"Fluvial" processes

Wind shear flow

Gravity currents



Experiments on a gravity-free dispersion of large solid spheres in a Newtonian fluid under shear

THE FLOW OF COHESIONLESS GRAINS IN FLUIDS

By R. A. BAGNOLD, F.R.S.

By R. A. BAGNOLD, F.R.S.

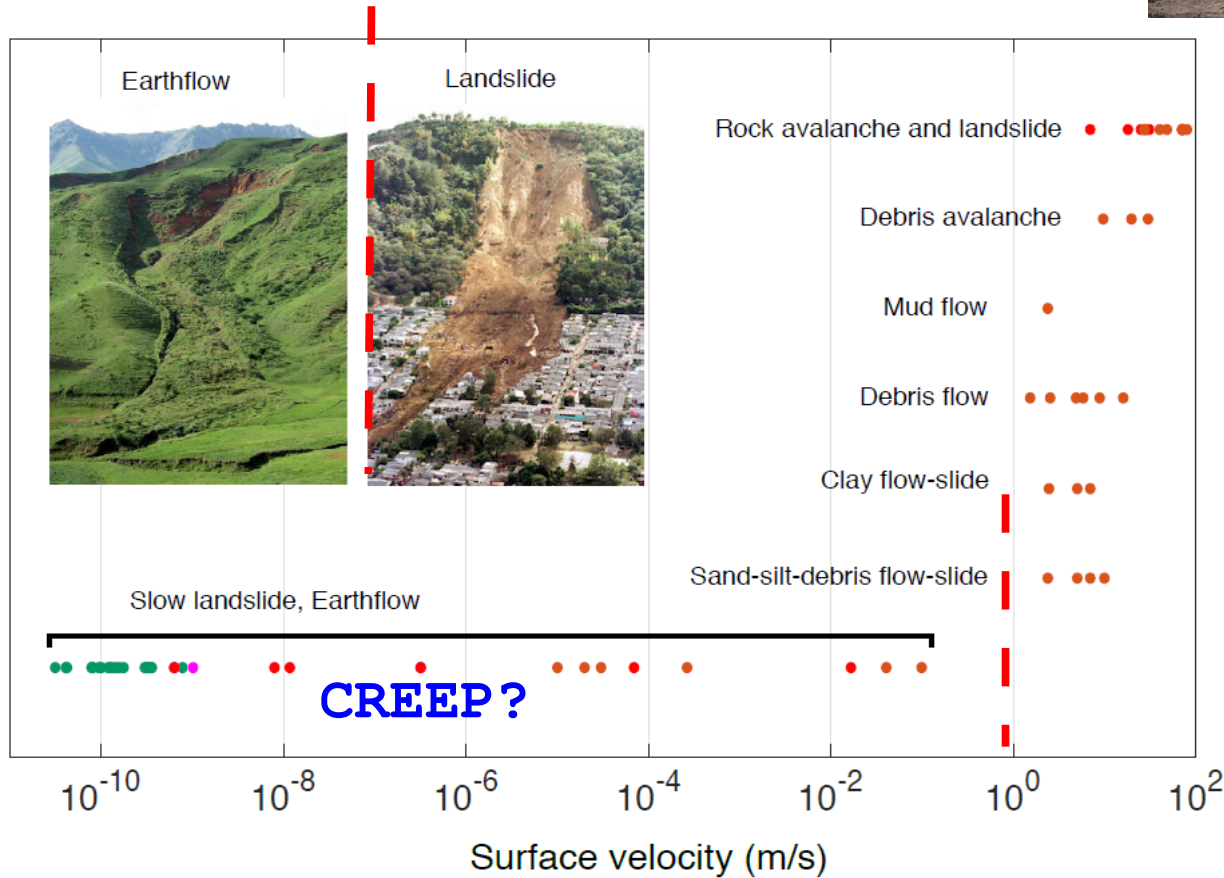
(Received 20 February 1954)

(Received 31 January 1955—Revised 14 April 1956)

# Landscapes are near critical: creeping ↔ flowing



Santa Barbara, 9 January 2018



# Creep in amorphous solids → DISORDER

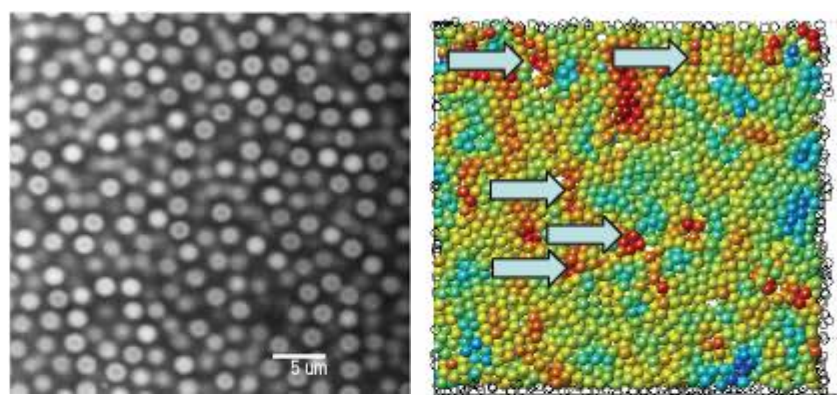
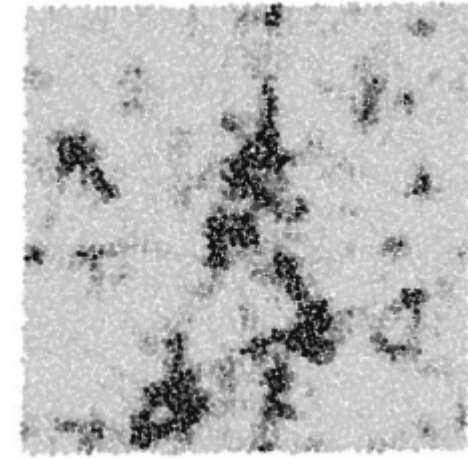
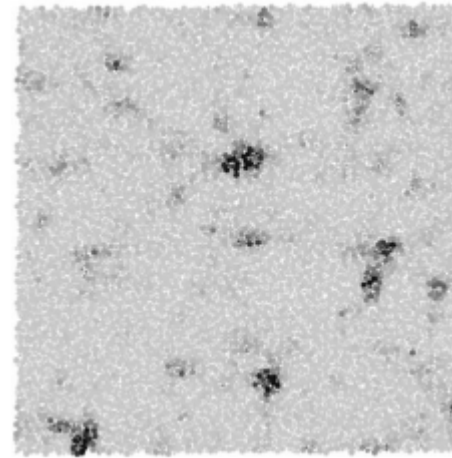
PHYSICAL REVIEW E

VOLUME 57, NUMBER 6

Dynamics of viscoplastic deformation in amorphous solids

M. L. Falk and J. S. Langer

Shear transformation zones  
in metallic glass simulations

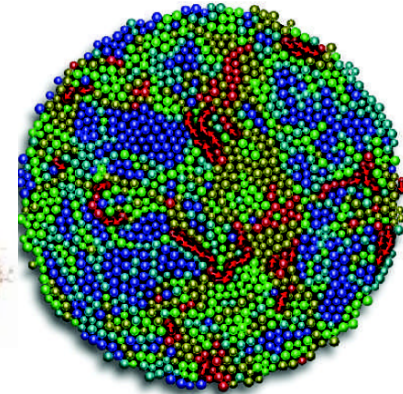


[Weitz group, Harvard]

Confirmed in:

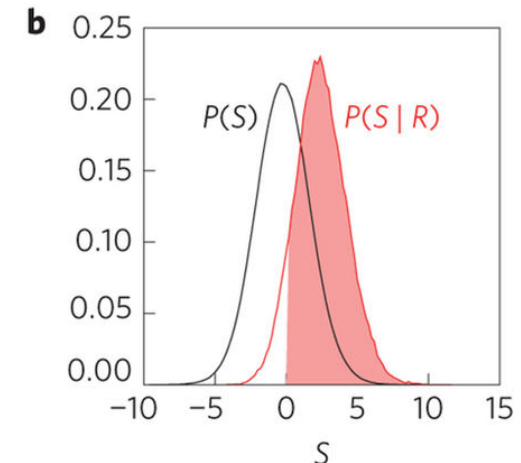
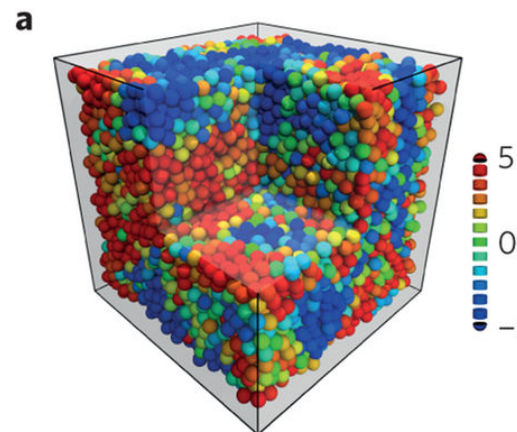
-Suspensions

-Granular materials



[Keys et al., *Nature Physics* 2007]

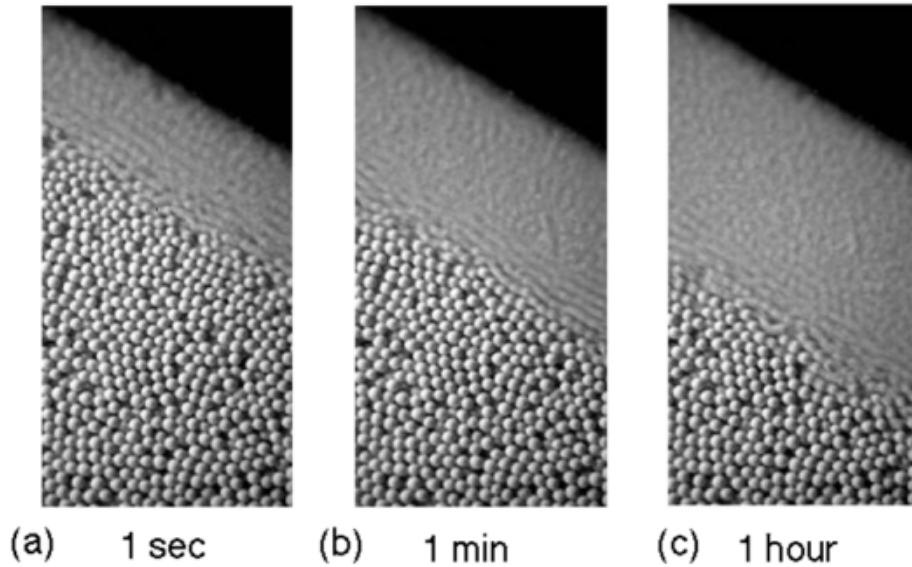
Structural defects  
(disorder) → “soft spots” →  
dislocation/flow



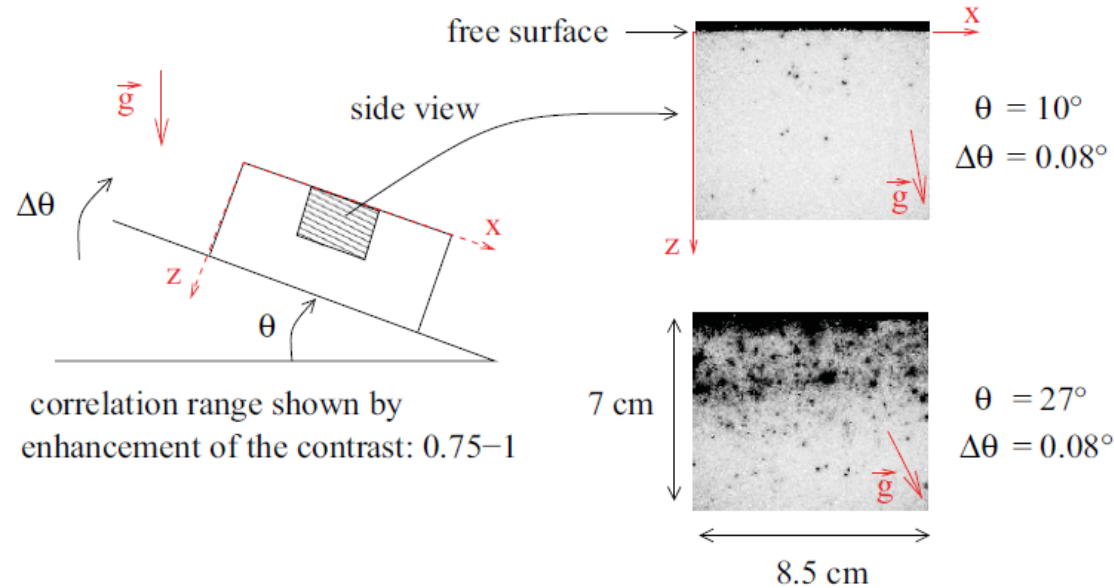
[Liu group, University of Pennsylvania]

# Creep to flowing transition: glassy dynamics

[Komatsu et al., PRL 2001]



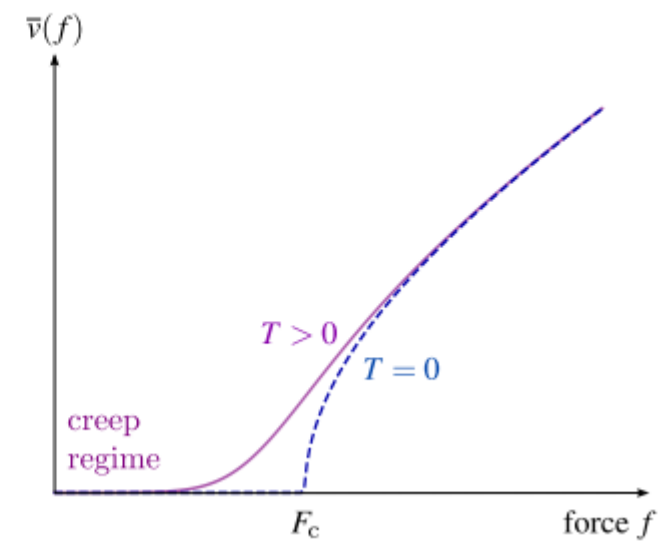
[Amos et al., Phys. Rev.E 2013]



[Agoritsas et al., J. Stat. Phys. 2017]

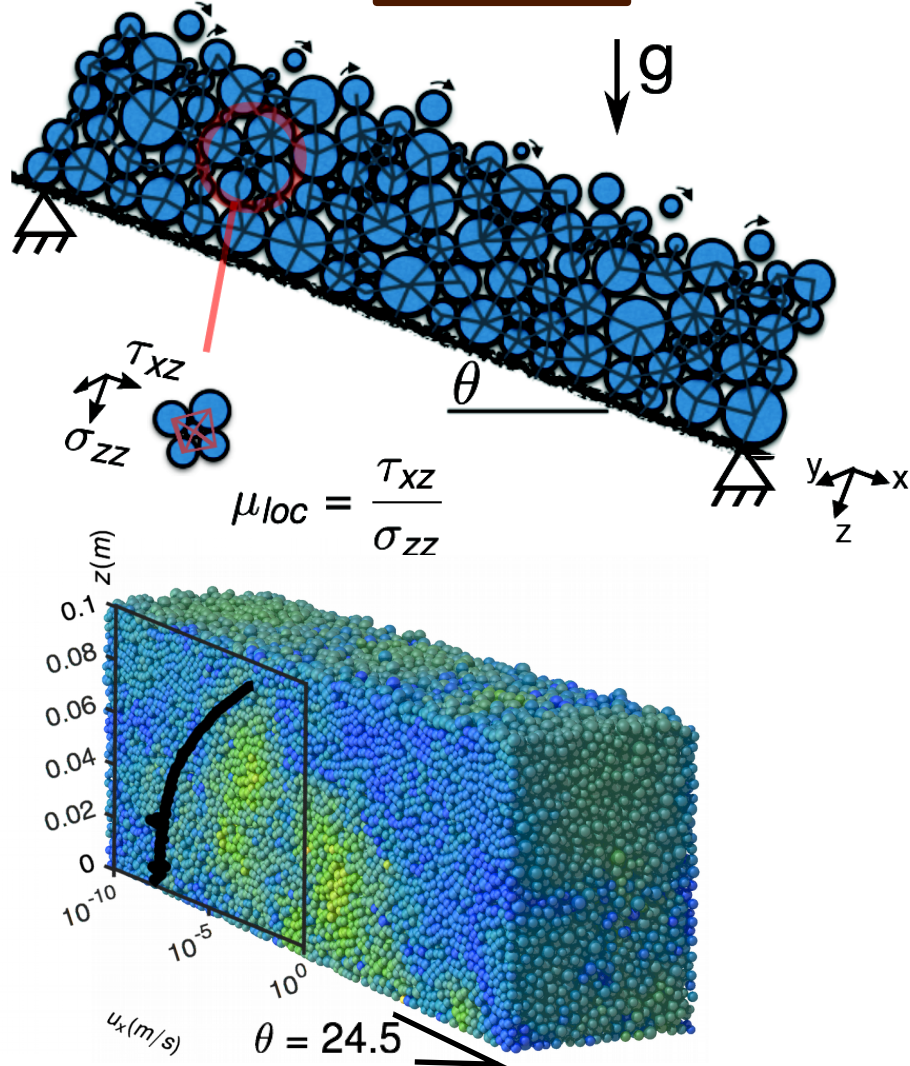
## Glassy failure model

- Avalanches of plasticity.
- At yield: avalanches merge across sample.
- Exponential (CREEP) to power-law (FLOW).
- Temperature - enhances creep.

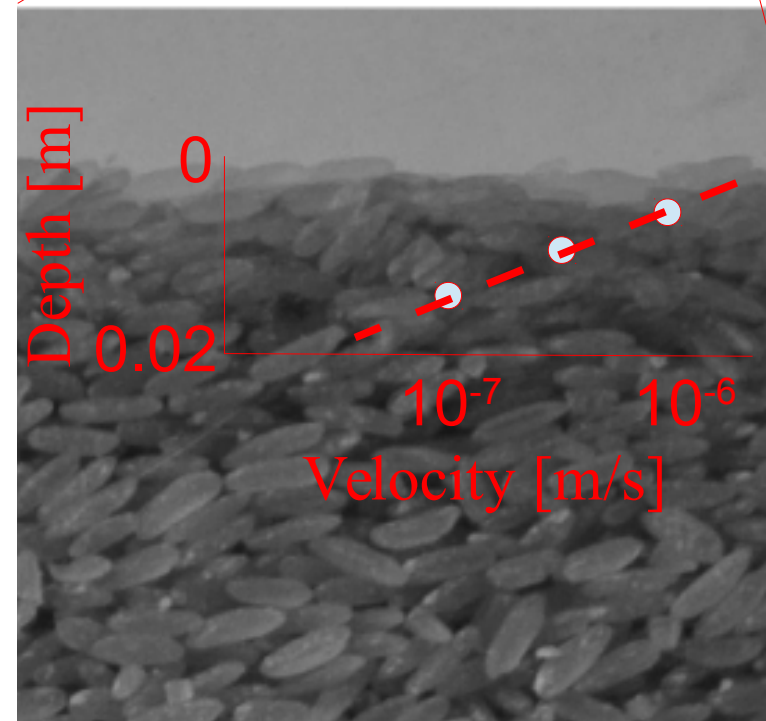
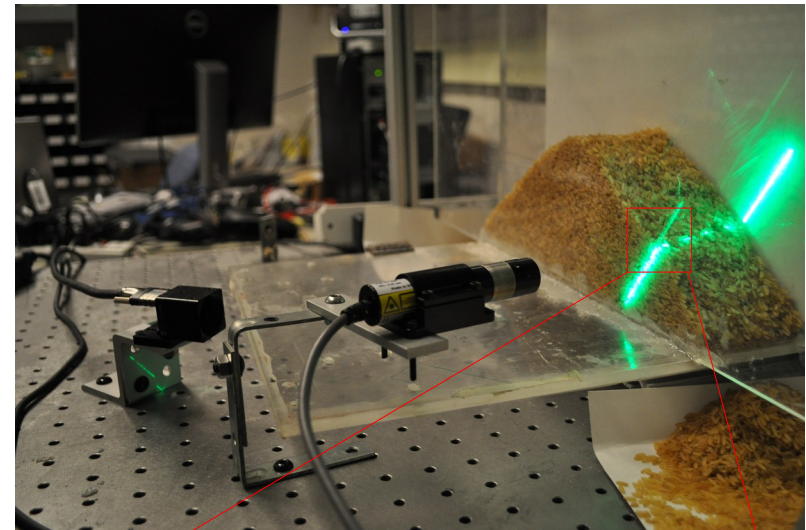


# Creep on hillslopes without disturbance

Creep: DEM simulations.

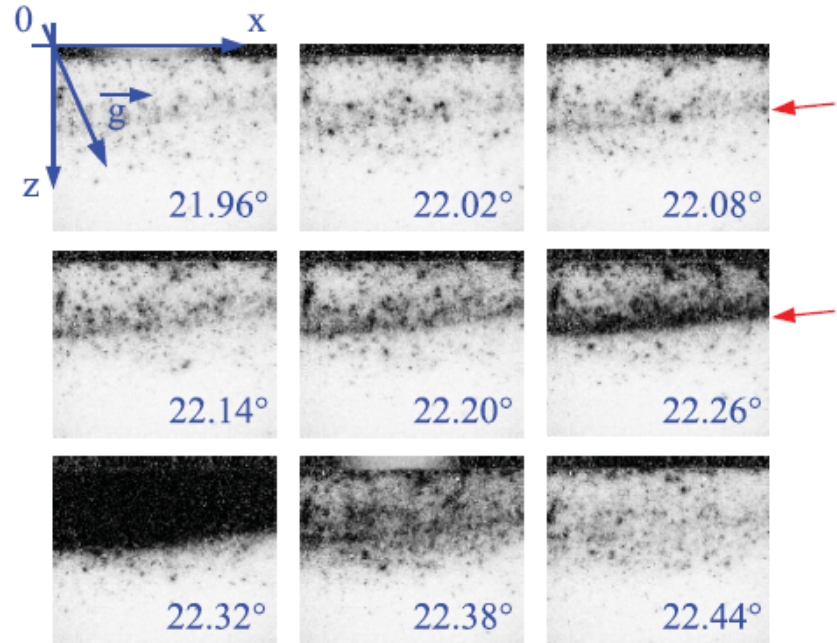
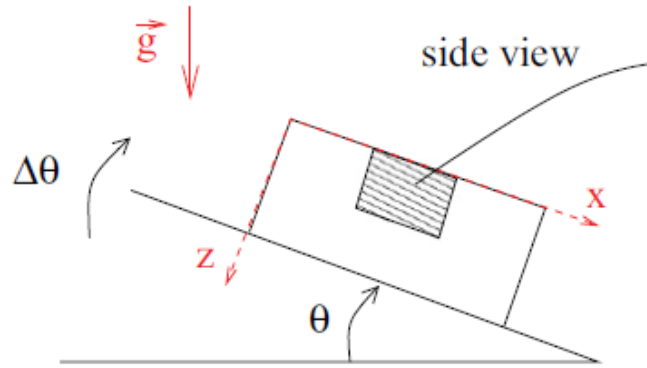


Creep: rice pile experiments.

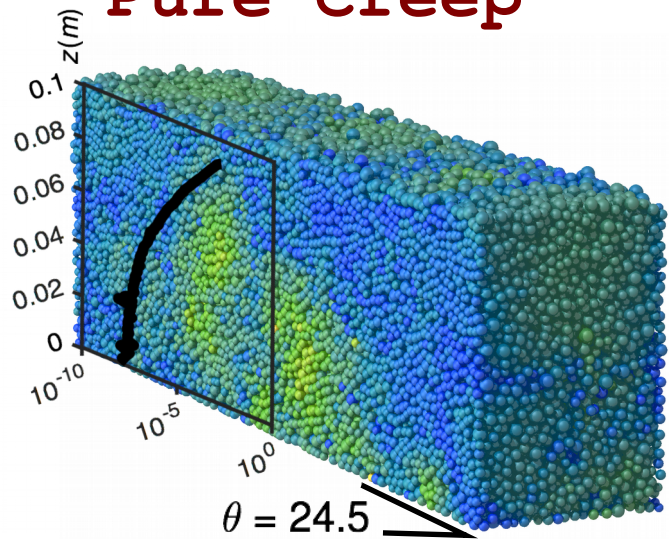


# Glassy dynamics in granular hillslopes

[Amos et al., *Phys. Rev.E* 2013]

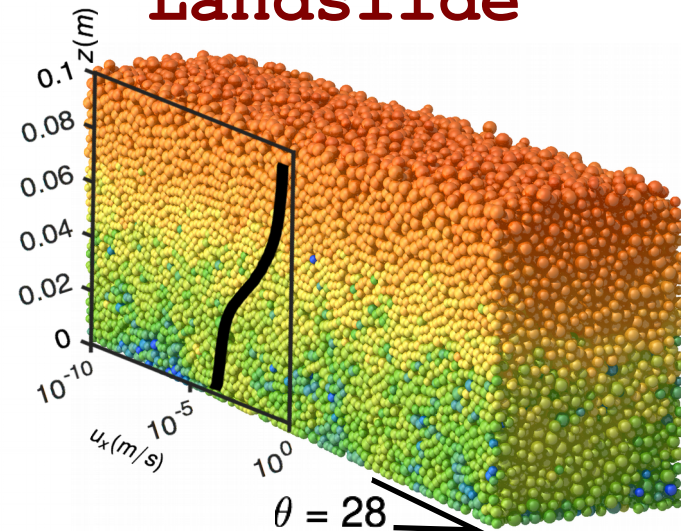


## Pure Creep



Creep occurs at sub-critical slopes.

## "Landslide"



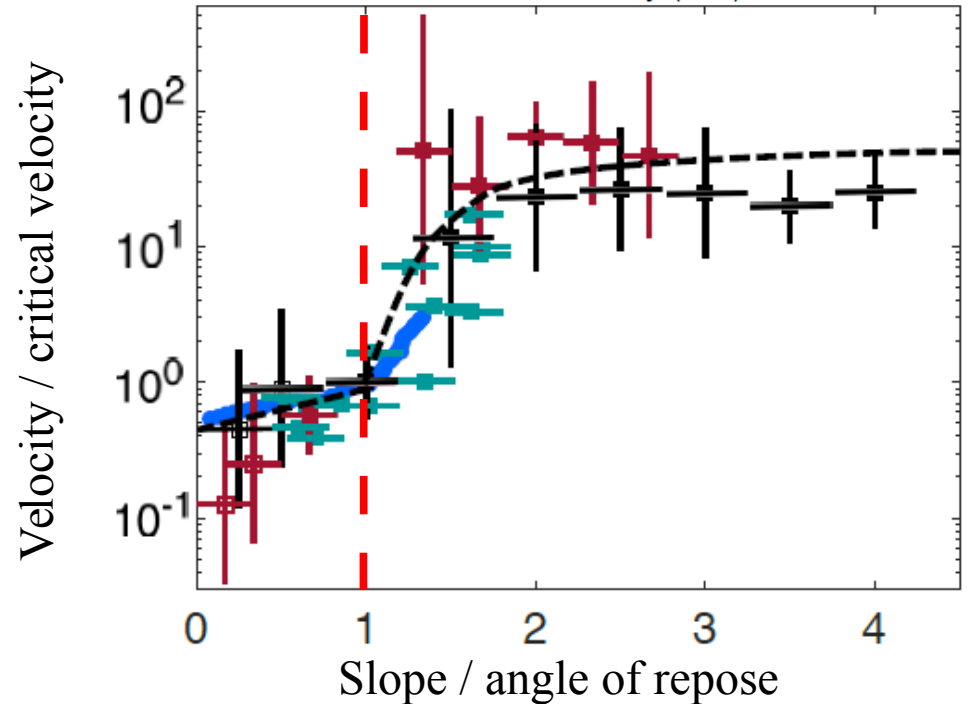
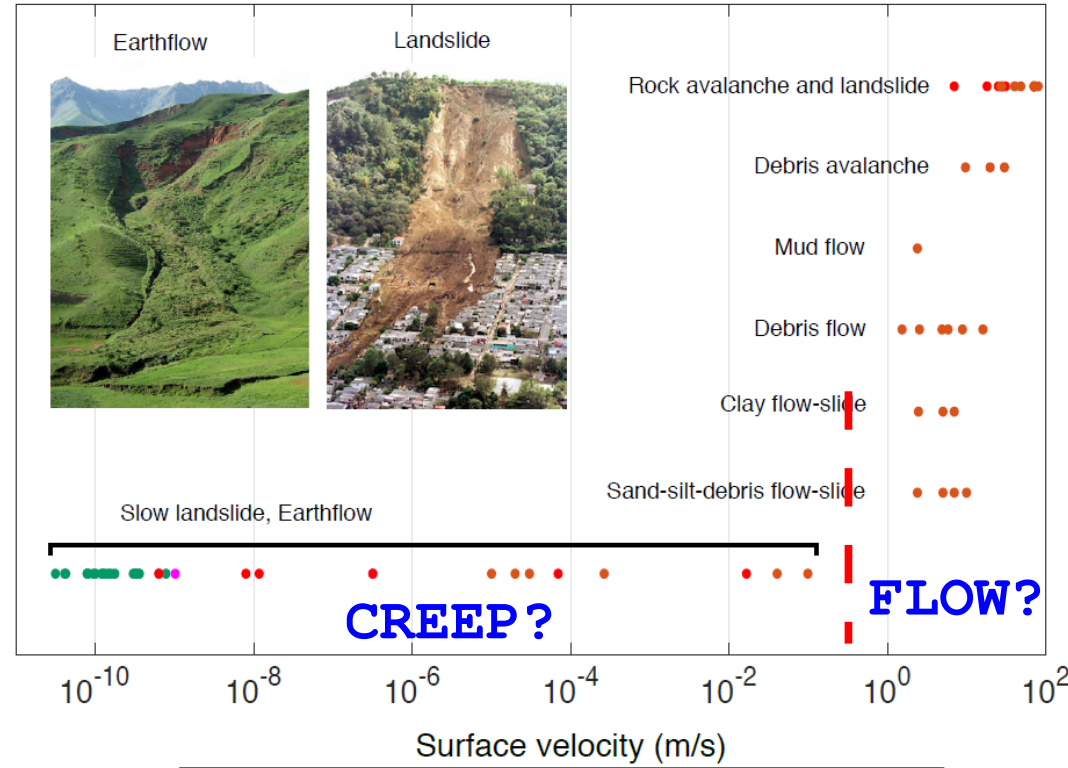
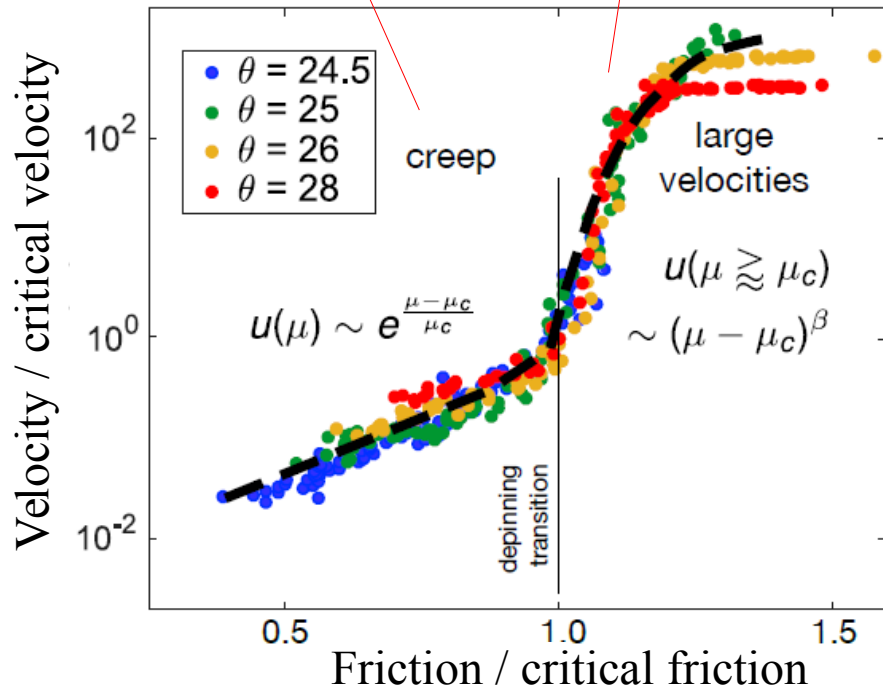
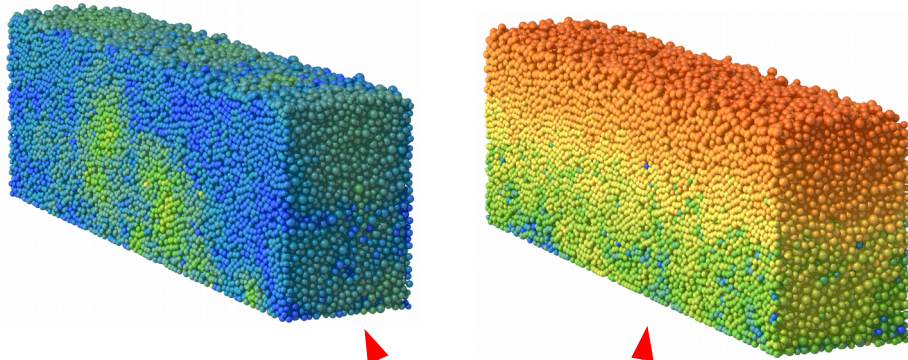
Creep  $\rightarrow$  Landslide is glassy

[Ferdowsi, Ortiz and Jerolmack, *REJECTED* 2017]

# Creep to landslide transition is glassy!

[Ferdowsi, Ortiz and Jerolmack, *REJECTED 2017*]

Exponential (CREEP) to power-law (LANDSLIDE)





# Are mountains glassy? Oregon Coast Range test

**Mass conservation:**

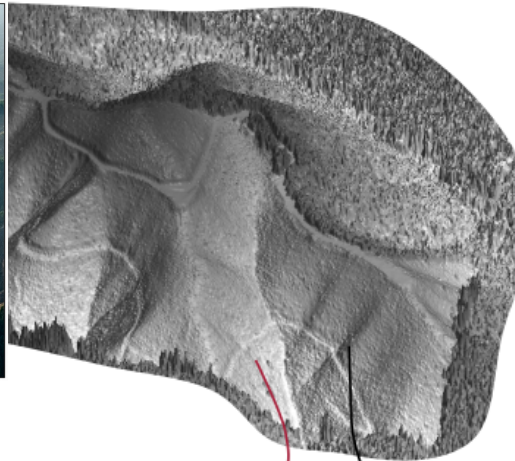
Steady state: Erosion = uplift

$$-\rho_s \frac{\partial z}{\partial t} = \rho_s \nabla \cdot \mathbf{q}_s + \rho_r C_o$$

**Glassy flux model:**

$$\frac{q_s}{q_{sc}} = e^{\frac{s-s_c}{s_c}} \quad |s < s_c$$

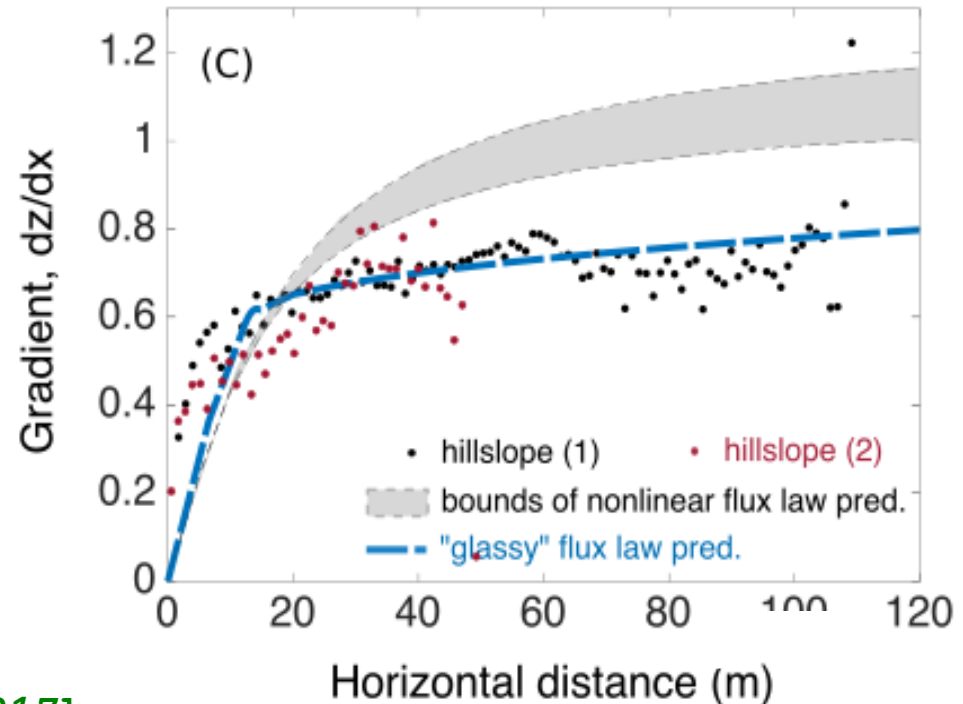
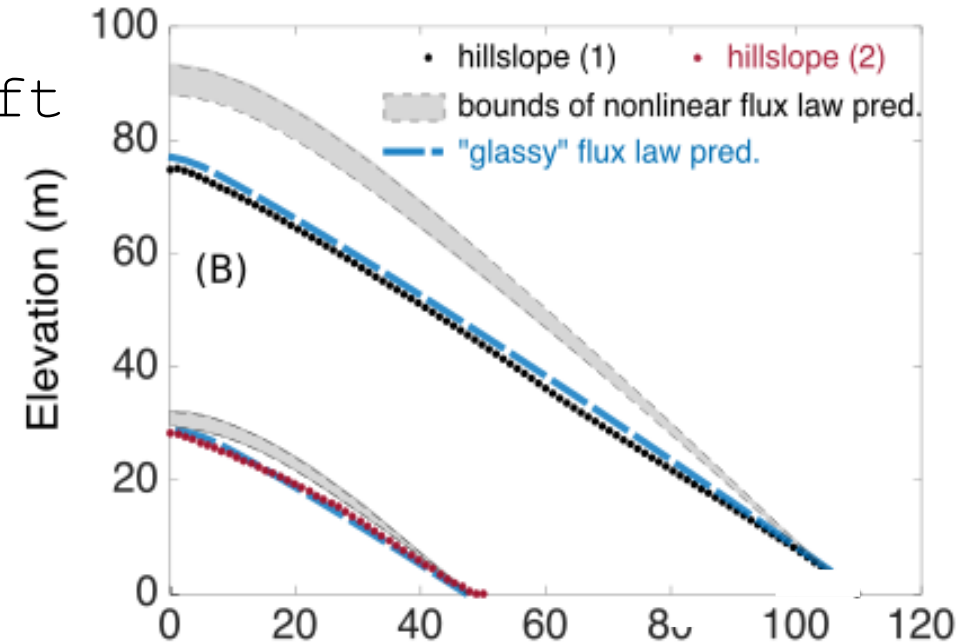
$$A(S - S_c)^\beta \quad |s > s_c$$



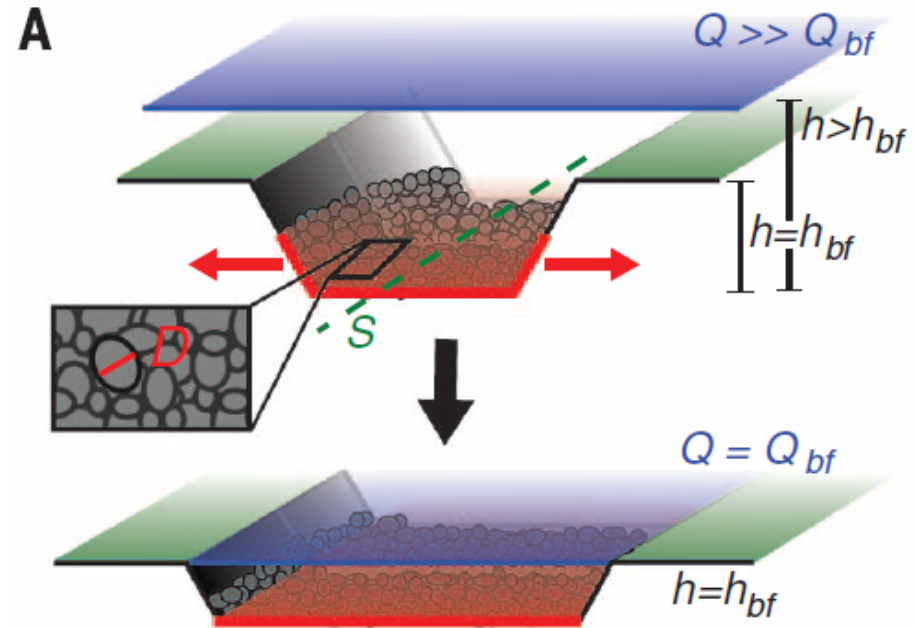
hillslope (2) hillslope (1)

**Constant uplift.**

**Best-fit determination of  $s_c$ .**



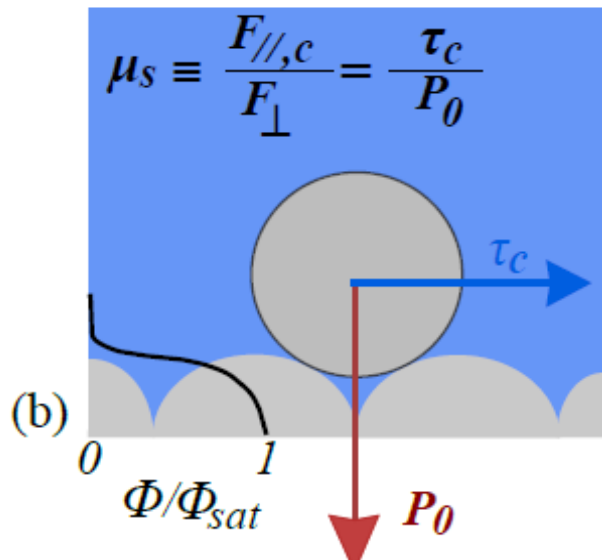
# Rivers are near also near critical



[Phillips & Jerolmack, Science 2016]

$$(1 + \epsilon) \tau_{*c} = \tau_{*bf} \approx 1.2 \tau_{*c}$$

[Parker, J. Fluid Mech. 1978]



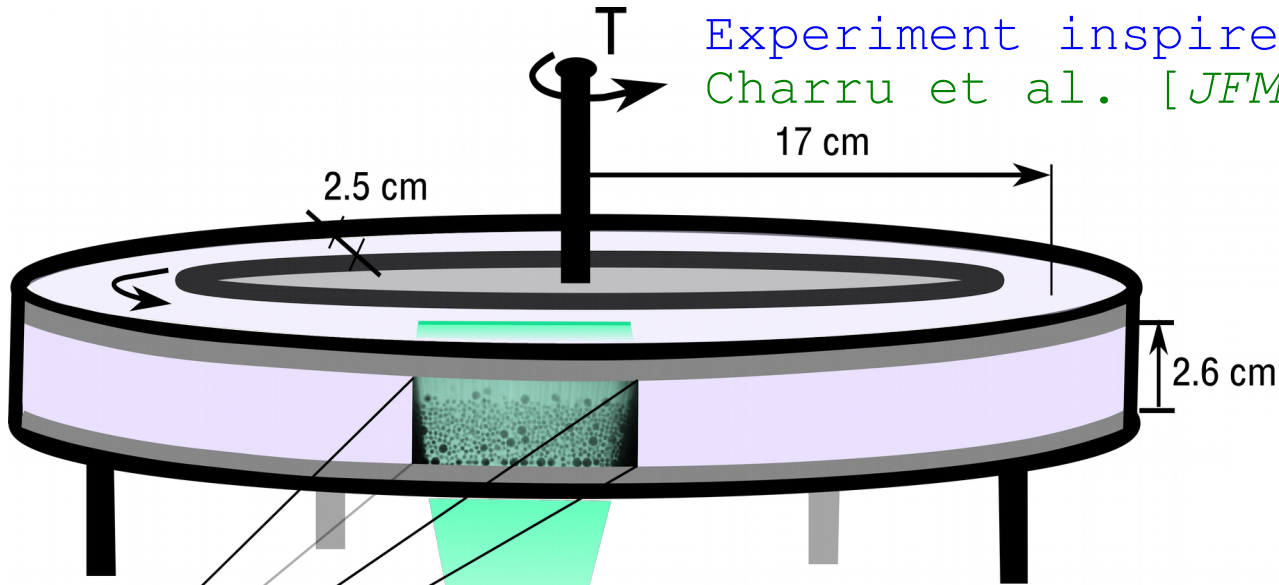
$$\tau^* = \frac{\tau}{(\rho_p - \rho)gd}$$



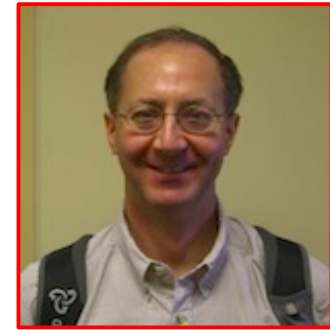
Ethan Mora:

<https://vimeo.com/22140684>

# The onset of erosion: glassy dynamics?



Experiment inspired by  
Charru et al. [*JFM* 2004]



[Houssais, Ortiz, Durian and  
Jerolmack, *Nature Comm.* 2015]

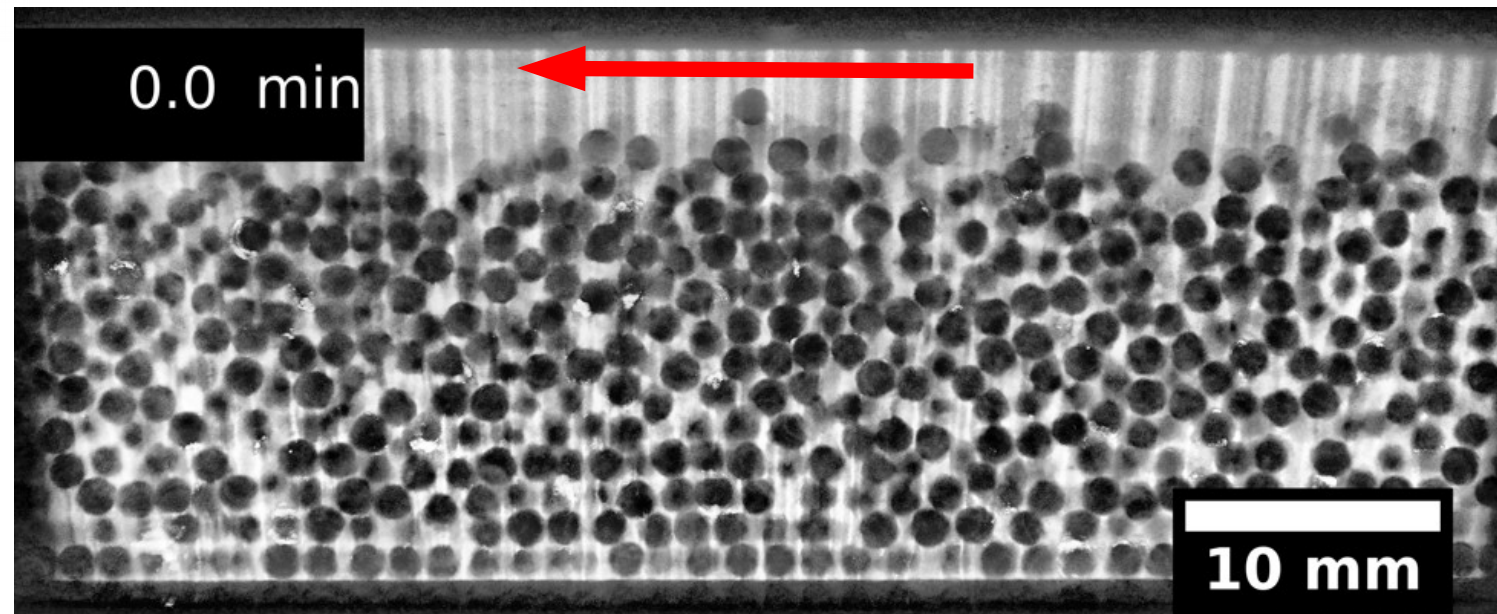
Camera  
lense

Laser  
sheet

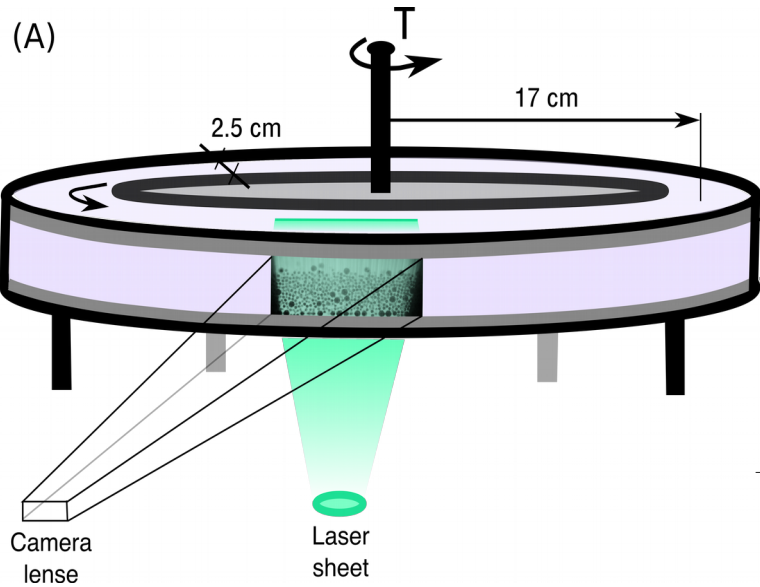
1.5 mm spheres

$Re < 5$  :  
Laminar flow

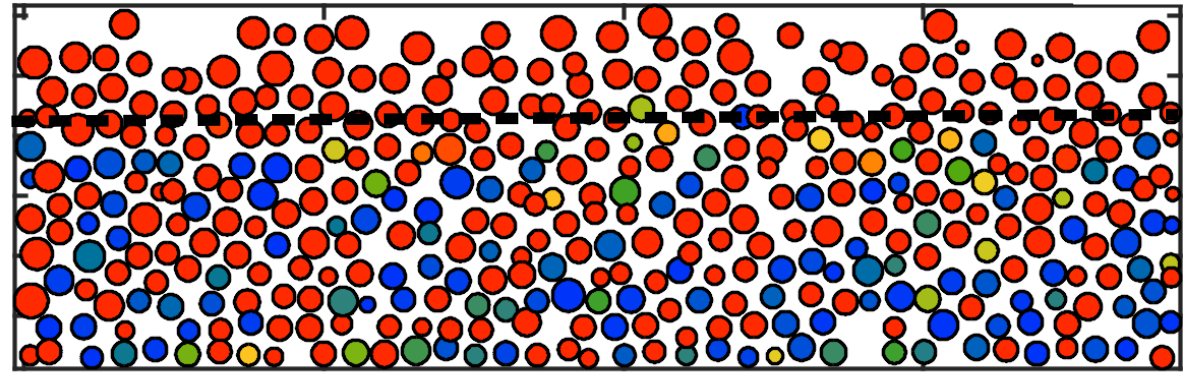
Refractive-Index  
Matched Scanning



# "Glassy" dynamics near threshold

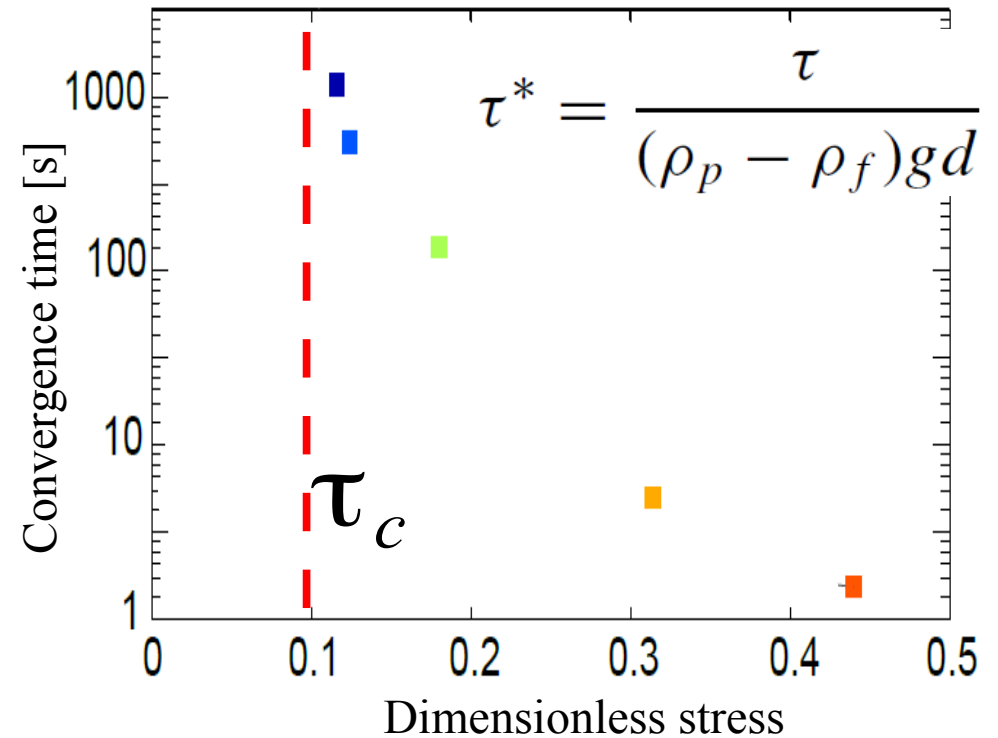
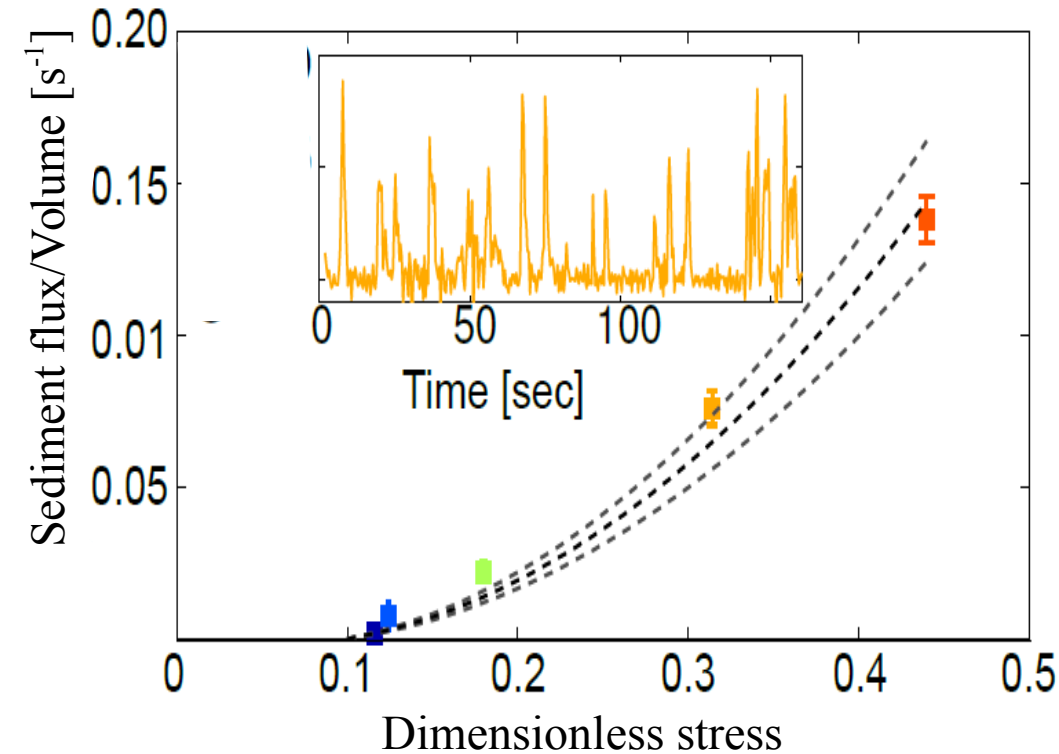


Particle motion relative to neighbors

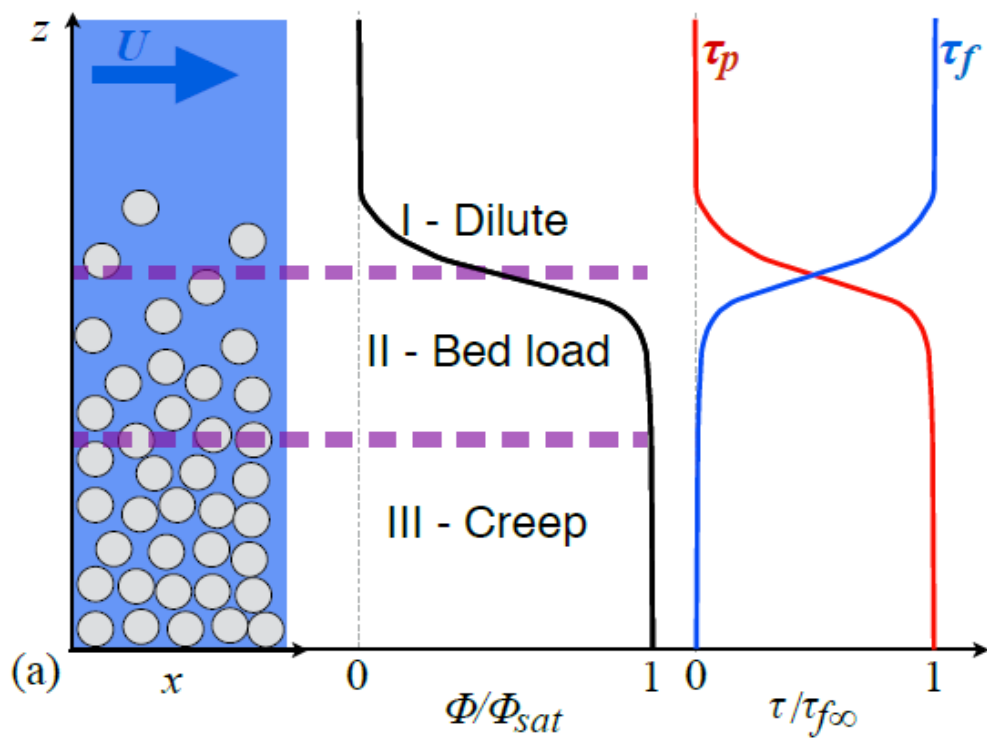
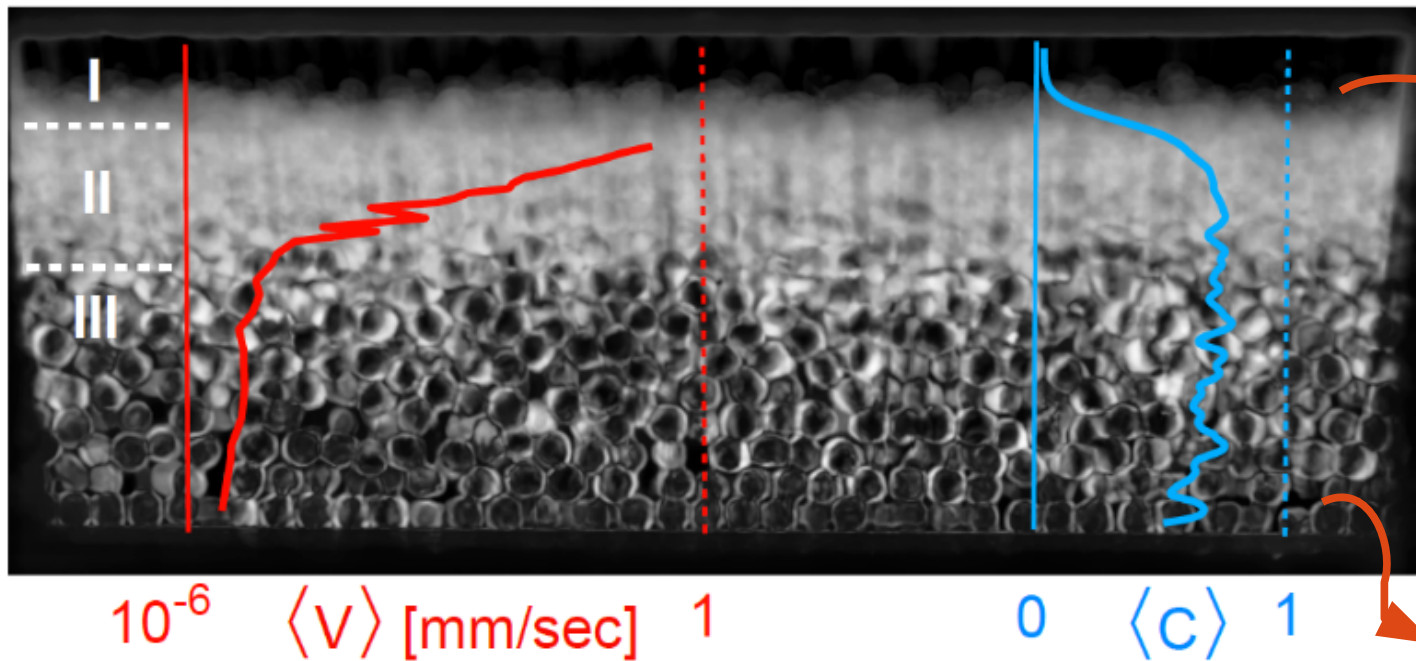


"Spatially heterogeneous dynamics":  
Collective particle rearrangements.

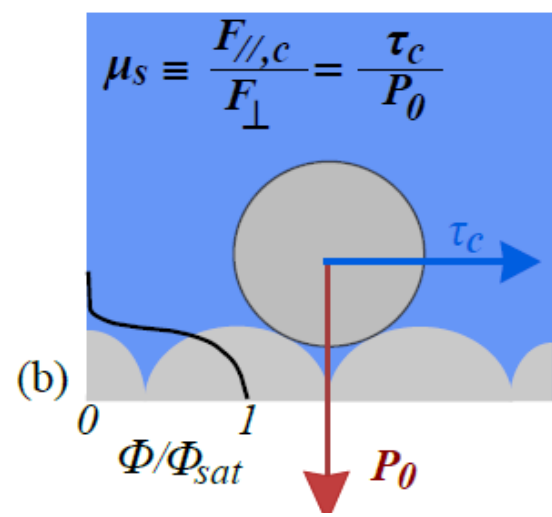
[Houssais, Ortiz, Durian and Jerolmack, Nature Comm. 2015]



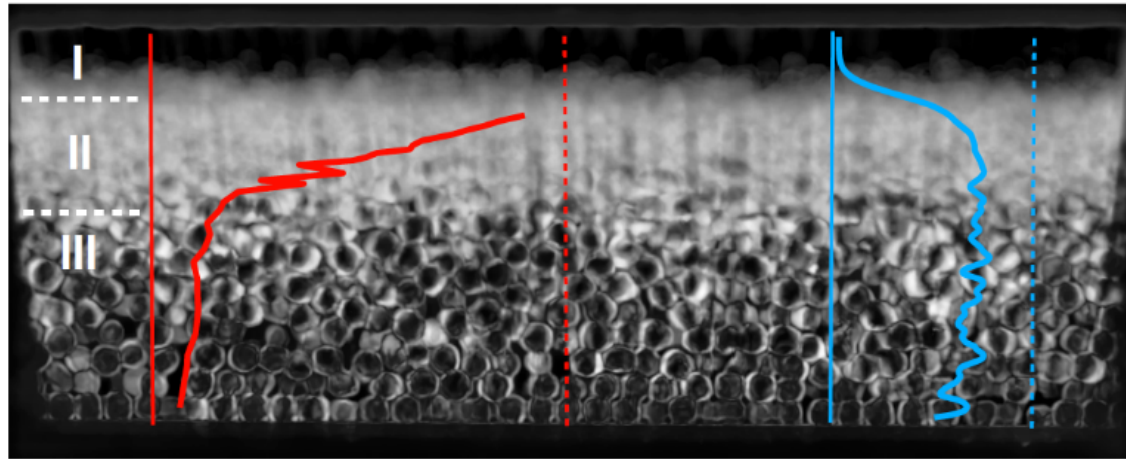
# Averaged vertical profiles



**Three transport regimes.**  
**What is the (ir) rheology?**



# Averaged vertical profiles



$10^{-6}$   $\langle V \rangle$  [mm/sec] 1      0  $\langle C \rangle$  1

Viscous # f (depth):

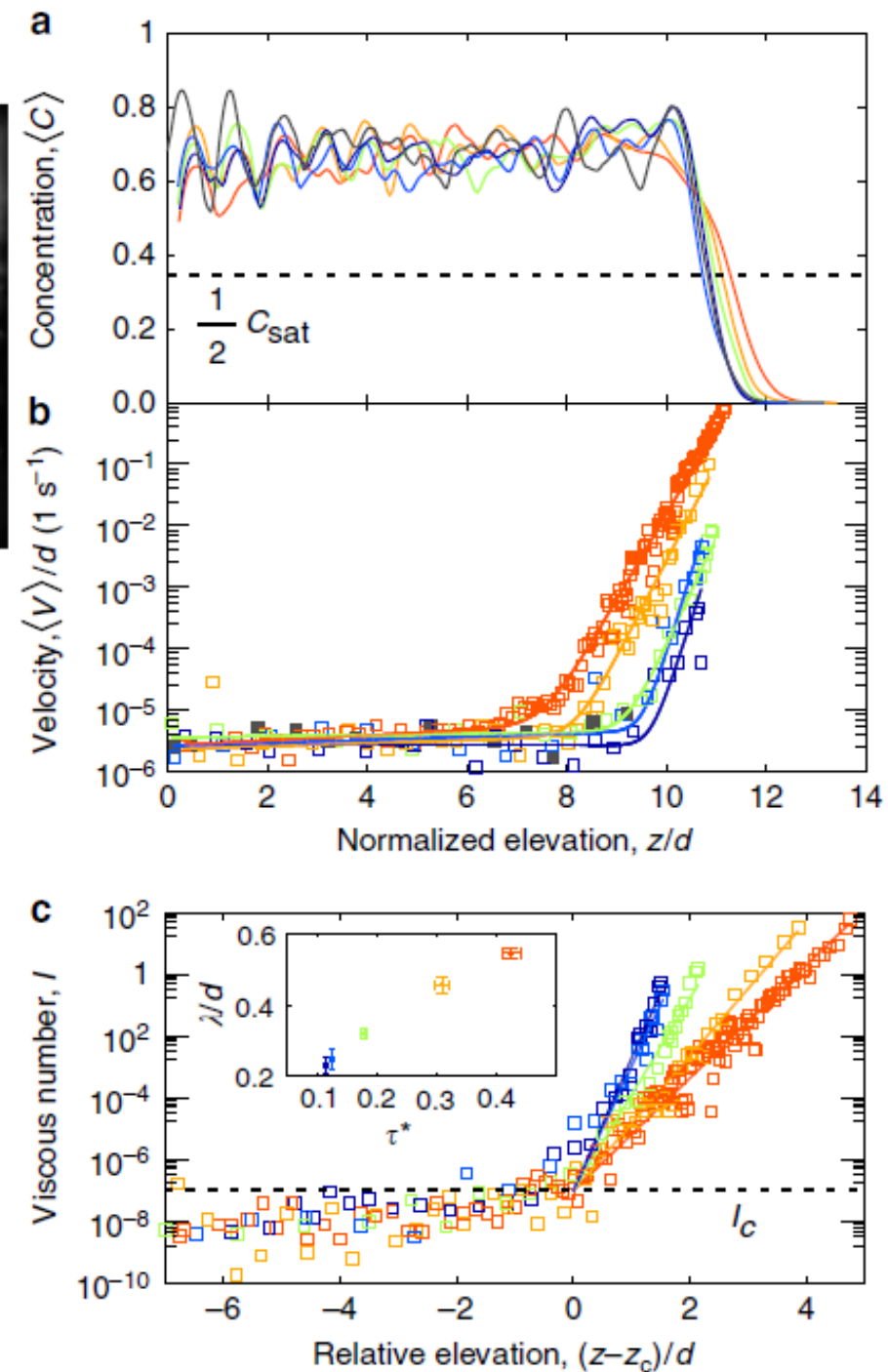
$$I_v = \frac{\eta_f |\dot{\gamma}|}{P_p}$$

Effective friction & viscosity:

$$\tau = \mu P_p = \eta_{\text{eff}} \dot{\gamma}$$

Confining pressure f (depth)

$$P_p(z) = P_0 + (\rho_p - \rho_f)g \int_z^{+\infty} \langle \phi \rangle dz$$



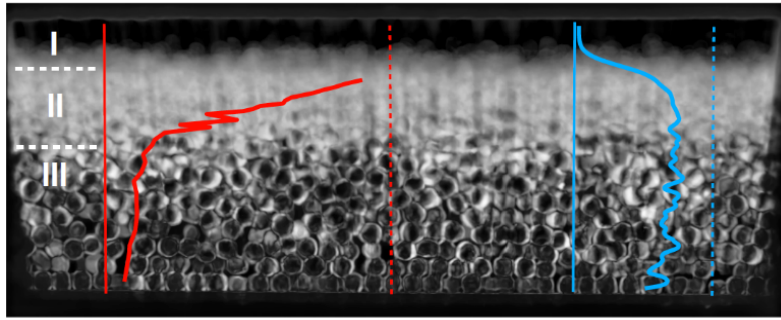
[Houssais, Ortiz, Durian & Jerolmack, *Nature Comm.* 2015; *PRE* 2017]

# Testing $\mu(I)$ rheology of Boyer et al. [PRL 2011]

Remarkable agreement for  $10^{-5} < I_v < 1$

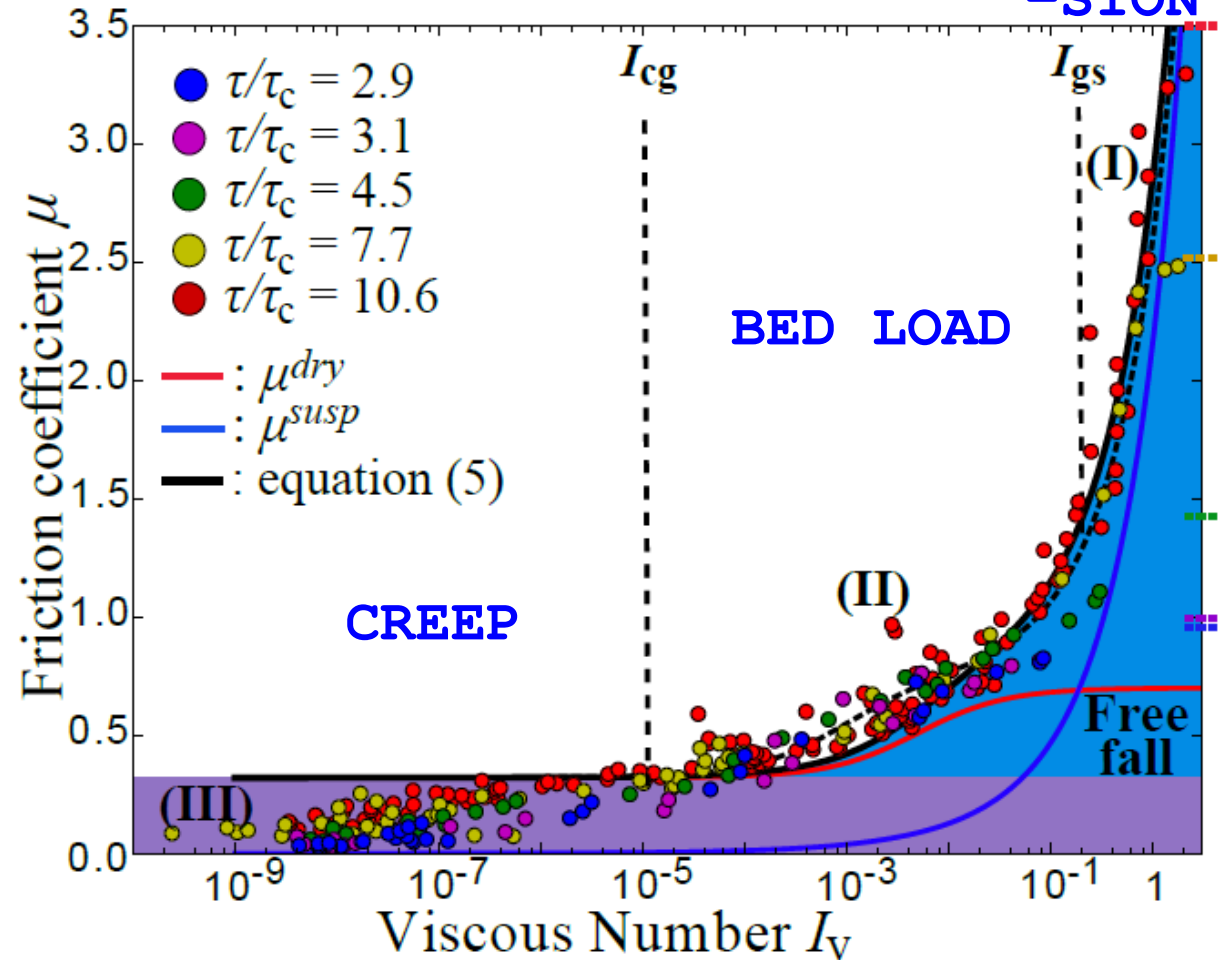
Sediment transport: Dense granular flow to suspension.

SUSPENSION



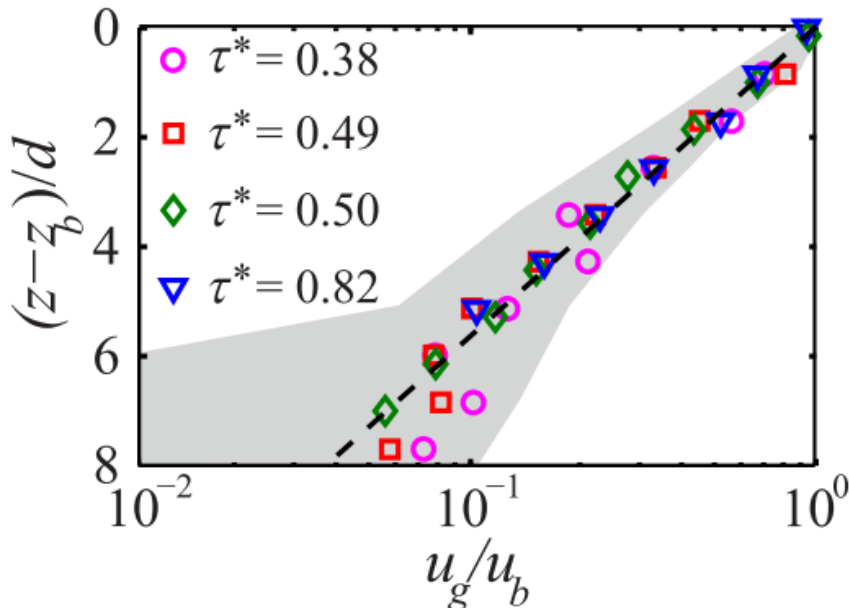
$10^{-6} \langle V \rangle$  [mm/sec] 1      0  $\langle c \rangle$  1

[Houssais, Ortiz, Durian & Jerolmack, *Phys. Rev. E* 2017]



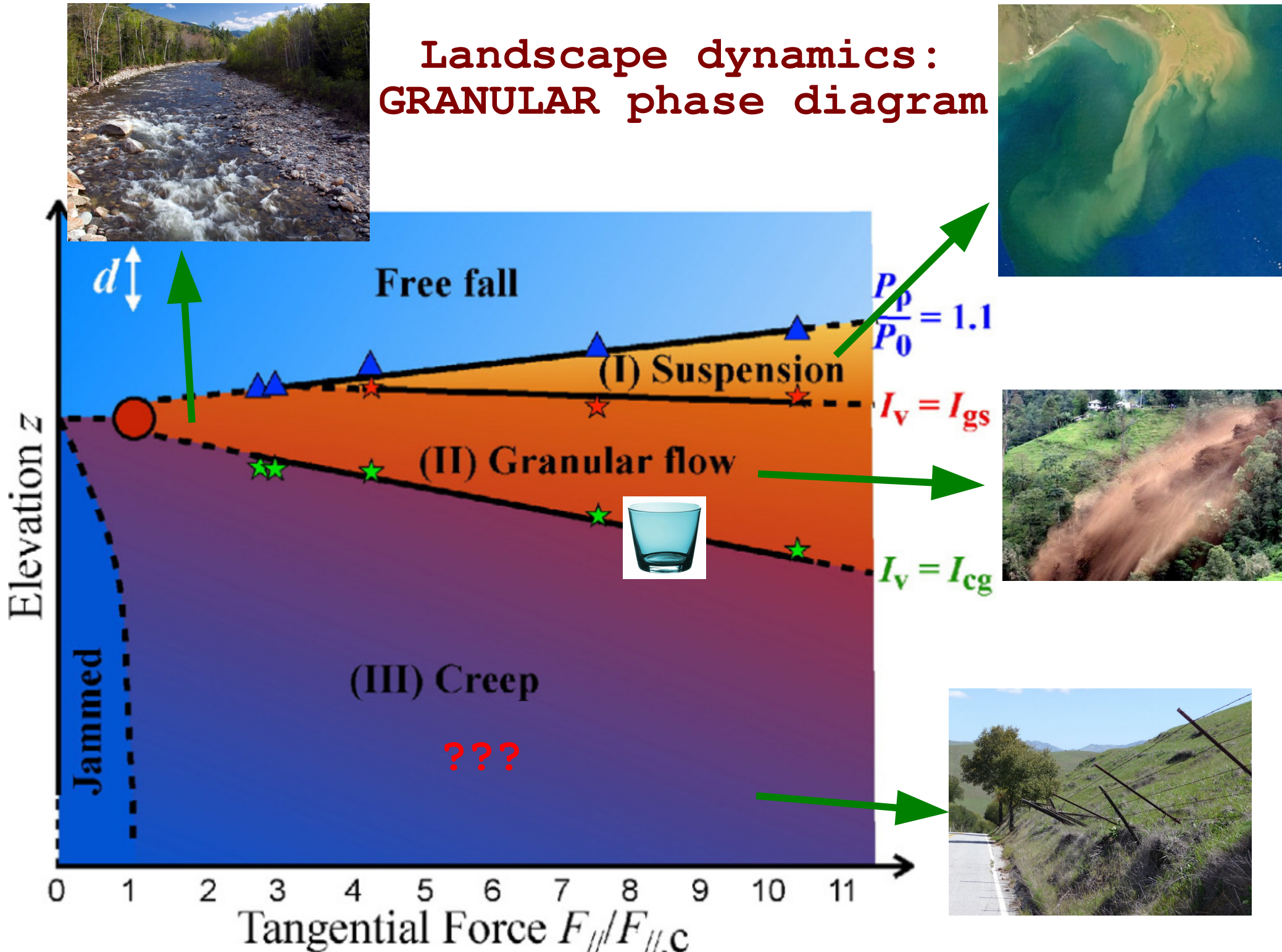
Creep breaks local rheology!

Nonlocal rheology models (e.g., Kamrin) appear to work.



[Allen & Kudrolli, *Phys. Rev. Fluids* 2017]

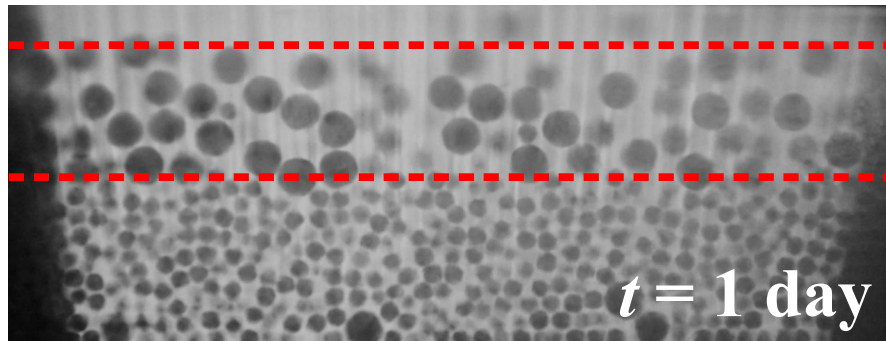
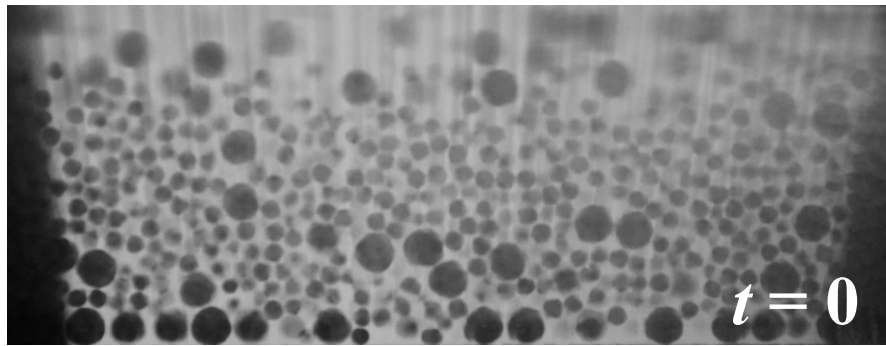
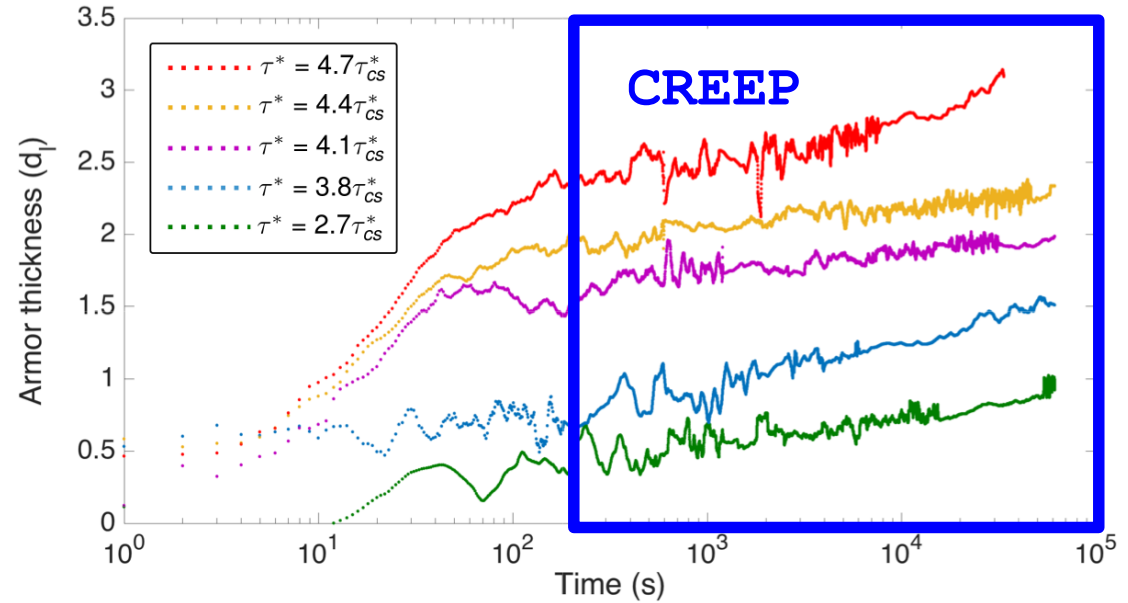
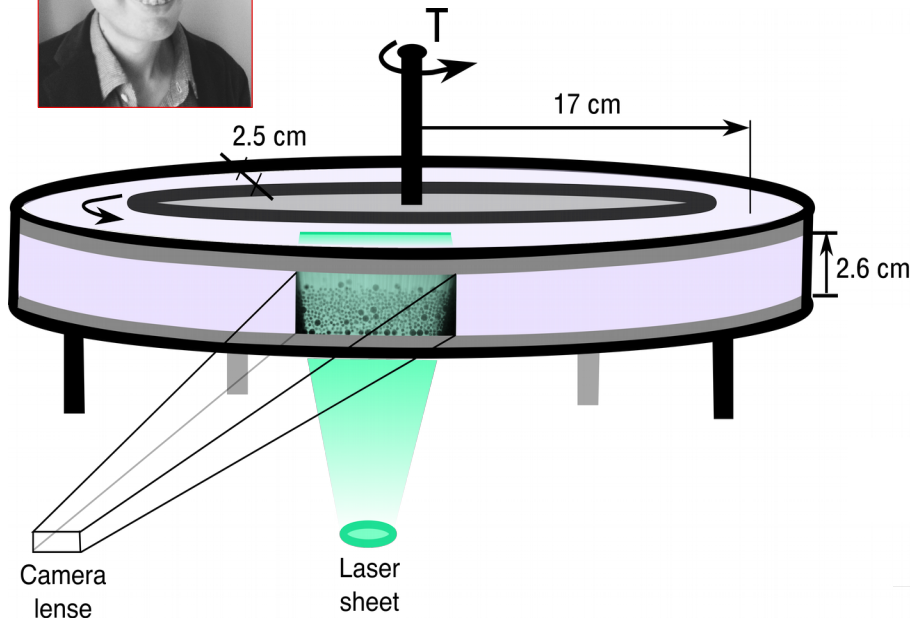
# Landscape dynamics: GRANULAR phase diagram





# Contribution of creep to segregation

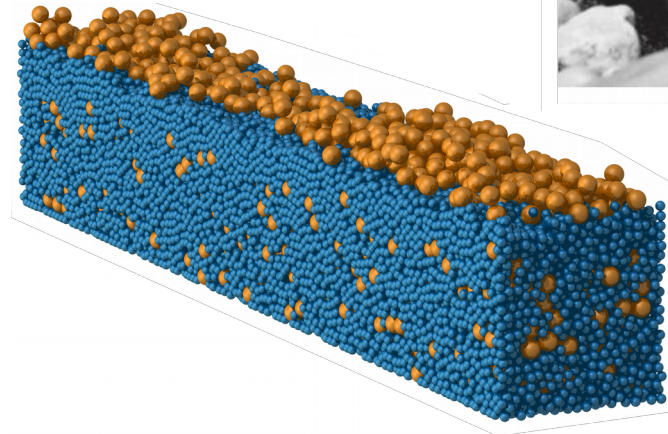
[Ferdowsi, Ortiz, Houssais & Jerolmack, *Nature Comm.* 2017]



## Segregation modes

Bed-load  $\rightarrow$  advection  
Creep  $\rightarrow$  diffusion.

River-bed armoring:  
granular phenomenon



DEM simulations  
reproduce  
experiments,  
WITHOUT FLUID.

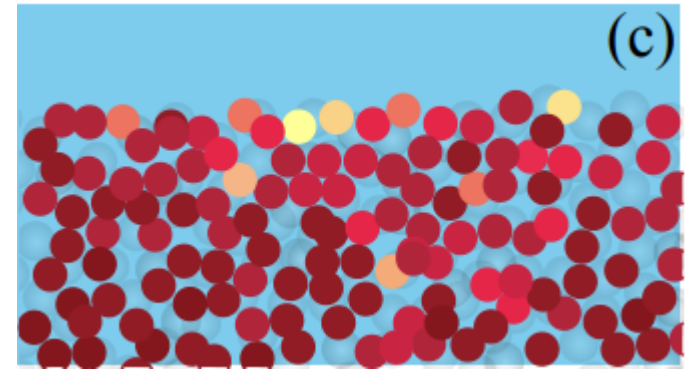
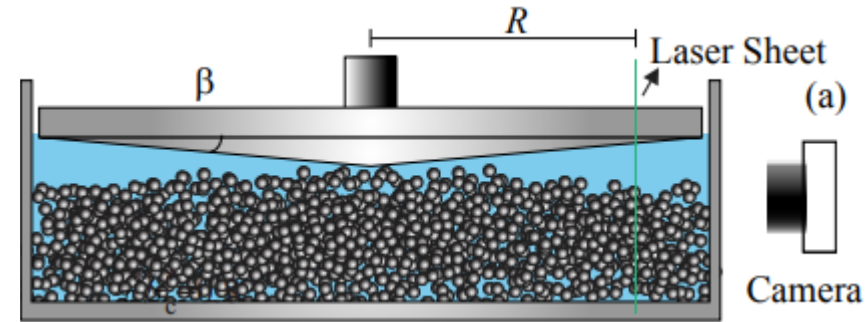
# What does creep do? → “armoring”

Granular bed consolidation, creep and armoring under subcritical fluid flow

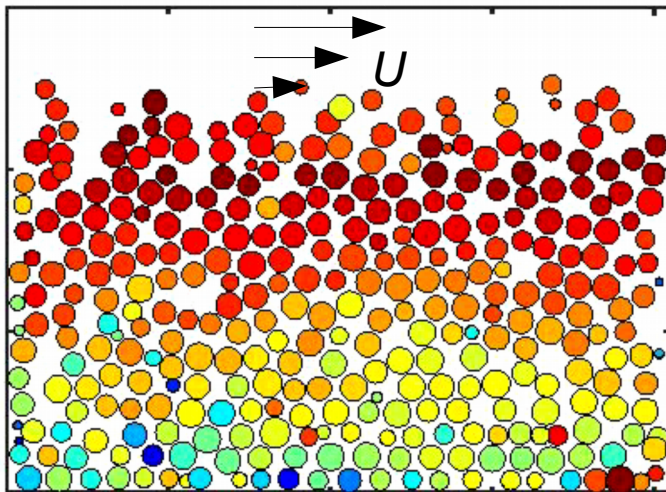
Benjamin Allen and Arshad Kudrolli  
Department of Physics, Clark University, Worcester, MA 01610  
(Dated: December 29, 2017)

We show that a freshly sedimented bed composed of spherical grains settles and creeps forward over extended periods under an applied hydrodynamic shear stress  $\tau^*$ , which is below the critical value  $\tau_c^*$  for bedload transport. The rearrangements are found to last at least over millions of times the sedimentation time scale of a grain in the fluid. Compaction occurs throughout the bed, but creep is observed to decay exponentially with depth, and decreases over time. The granular volume fraction in the bed is found to increase logarithmically, saturating at the random close packing value  $\phi_{rcp} \approx 0.64$ , while the surface roughness is on average essentially unchanged. Thus, we find that bed armoring occurs due to a deep shear-induced relaxation of the bed toward the volume fraction associated with the glass transition.

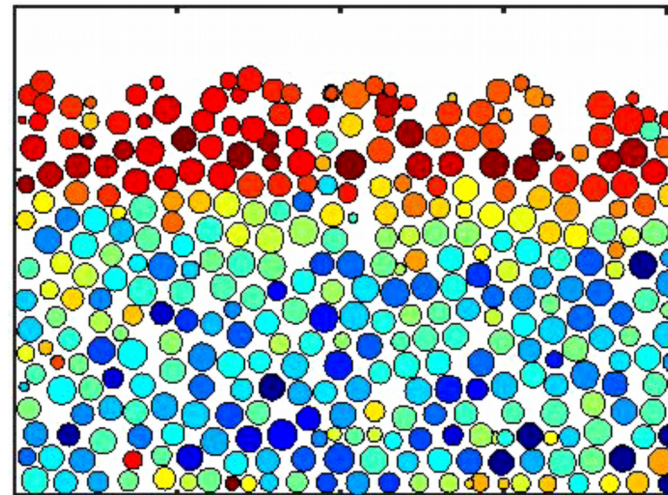
Sub-yield creep in the absence of “flow”.



Grain velocity - start  
→ Loose bed



Grain velocity - 10 min  
→ More rigid



→ Creep drives subsurface development of structure.