

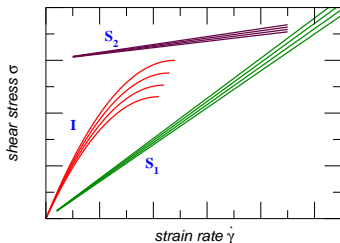
# Heterogeneous Flows and Constitutive Behavior

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Department of Physics, Georgetown University

KITP, 18 January 2018

Homogeneous constitutive relations for a hypothetical fluid:



- I Newtonian:
- $S_1$  flow-induced phase
- $S_2$  flow-induced “gel”

- Identify dynamical variables (flow, composition, and structural information).
- “Phases” = steady state homogeneous solutions to equations of motion:

$$\begin{aligned}
 D_t \phi &= -\nabla \cdot M \nabla \mu(\phi, \nabla \mathbf{v}, \mathbf{Q}) &= 0 & \text{(composition)} \\
 \rho D_t \mathbf{v} &= \nabla \cdot \boldsymbol{\sigma}(\phi, \nabla \mathbf{v}, \mathbf{Q}) &= 0 & \text{(flow)} \\
 D_t \mathbf{Q} &= \mathcal{L}(\phi, \nabla \mathbf{v}, \mathbf{Q}) &= 0 & \text{(microstructure)}
 \end{aligned}$$

- Calculate steady state flow curves as a function of concentration.

# Stress-concentration couplings

- Two-fluid models [Helfand/Fredrickson 1989, Milner 1992], flow-induced migration [Leighton & Acrivos 1987, Schmidt/Marques/Lequeux PRE 1995].

$$\begin{aligned}\rho D_t \mathbf{v} &= \nabla \cdot G(\phi) \mathbf{W} - \phi \nabla \frac{\delta F(\phi)}{\delta \phi} + 2 \nabla \cdot \eta(\phi) \mathbf{D} - \nabla p_0, \\ D_t \phi &= -\nabla \cdot \zeta^{-1}(\phi) \left[ \nabla \cdot G(\phi) \mathbf{W} - \phi \nabla \frac{\delta F}{\delta \phi} + 2 \nabla \cdot \eta(\phi) \mathbf{D} \right] \\ (\partial_t + \mathbf{v}_m \cdot \nabla) \mathbf{W} + \dots &= +2 \mathbf{D}_m - \frac{\mathbf{W}}{\tau(\phi)} + \frac{\ell^2}{\tau(\phi)} \nabla^2 \mathbf{W}\end{aligned}$$

- Concentration builds up in more viscous regions.
- Polymer deformation in stress gradient.
- Transverse diffusion due to gradient in shear rate and hence collision rate; . . . .
- Migration in response to stress gradients imposed by geometry (e.g. Poiseuille flow).

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**Careful:**  $\left\{ \begin{array}{l} \text{Total stress tensor} \quad \mathbf{T} \neq \mathbf{T}_{\text{micro}} - p \mathbf{I} !! \\ \text{The pressure is given by} \quad p = p_0 - \text{Tr}(\mathbf{T}_{\text{micro}}). \end{array} \right.$



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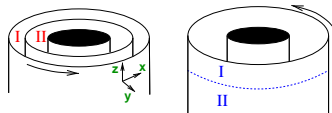
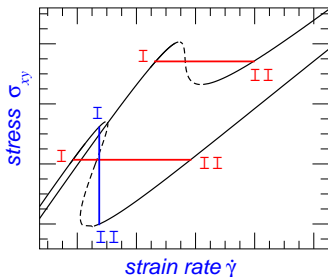
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# Geometry determines “field variables”

Definition of **field and density variables** depends on coexistence **geometry**.



Phase separate according to  $\bar{\phi} = \alpha \phi_I + (1 - \alpha) \phi_S$ .

## Common Stress

Gradient Banding

$$\bar{\dot{\gamma}} = \alpha \dot{\gamma}_I + (1 - \alpha) \dot{\gamma}_{II}$$

$\sigma$  = “field variable”

$\dot{\gamma}$  = “density variable”

## Common Strain Rate

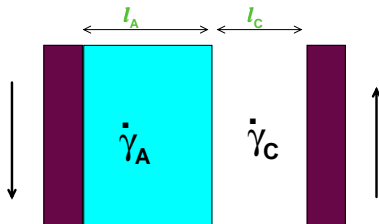
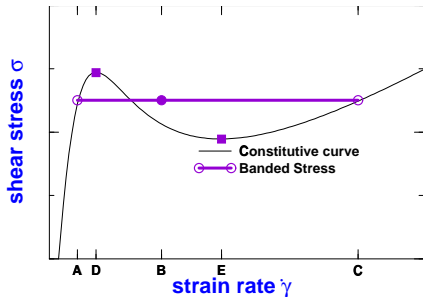
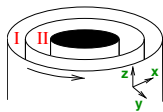
Vorticity Banding

$$\bar{\sigma} = \alpha \sigma_I + (1 - \alpha) \sigma_{II}$$

$\dot{\gamma}$  = “field variable”

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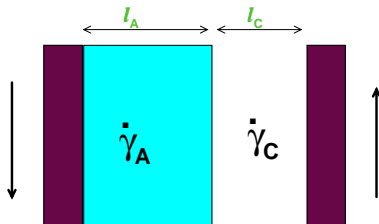
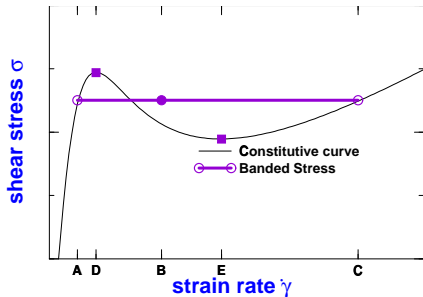
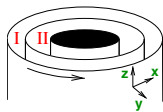
1. Stationary states
2. Stress Balance
3. No  $\phi$  flux:



# Coexistence Conditions: gradient banding [PDO Rheo Acta 2008]

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2. Stress Balance
3. No  $\phi$  flux:

$$\partial_t \mathbf{Q} = 0 \Rightarrow \sigma(\dot{\gamma})$$



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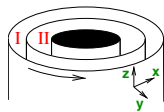
2. Stress Balance

$$\nabla \cdot \boldsymbol{\sigma} = 0 \Rightarrow$$

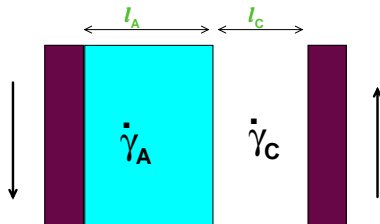
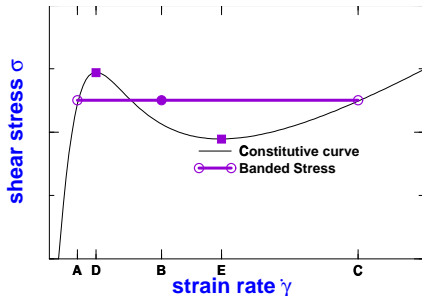
$$\sigma_{xy}^I = \sigma_{xy}^{II}$$

(uniform  $\sigma_{xy}, \sigma_{yy}$ )

$$\sigma_{yy}^I = \sigma_{yy}^{II} \Rightarrow p(y)$$



3. No  $\phi$  flux:



The pressure in the two bands may differ because of normal stress continuity.

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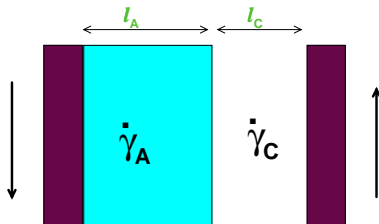
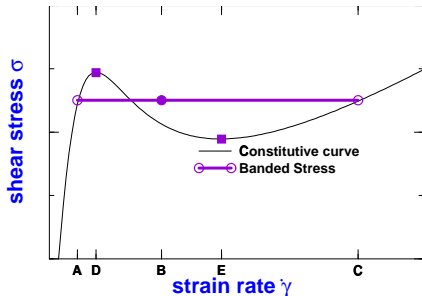
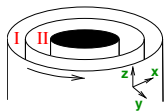
$$\nabla \cdot \mathbf{J} = 0 \Rightarrow$$

$$\partial_t \mathbf{Q} = 0 \Rightarrow \boldsymbol{\sigma}(\dot{\gamma})$$

$$\sigma_{xy}^I = \sigma_{xy}^{II}$$

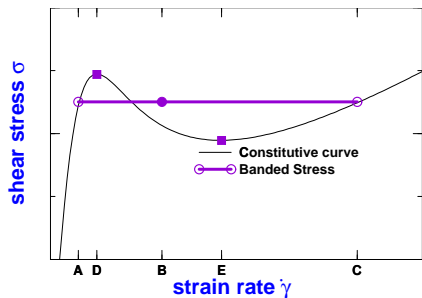
$$\sigma_{yy}^I = \sigma_{yy}^{II} \Rightarrow p(y)$$

$$\mu = \frac{\delta F}{\delta \phi} = \text{uniform}$$

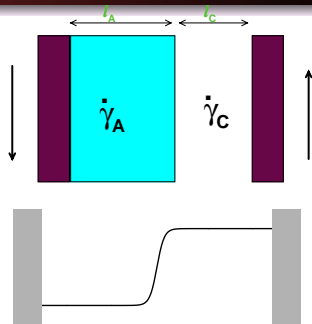


The pressure in the two bands may differ because of normal stress continuity.

# Stress selection and 'non-equilibrium phase coexistence'

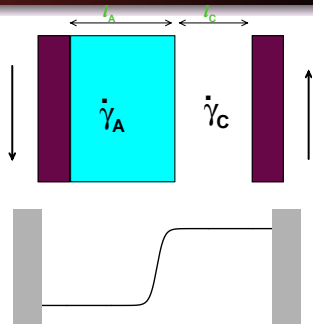
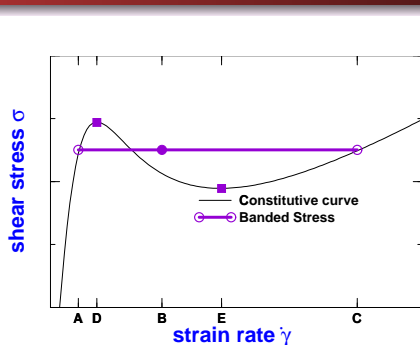


$$\sigma_{tot} = \sigma_p + \eta \dot{\gamma}$$



$$\partial_t \sigma_p = f(\sigma_p, \dot{\gamma})$$

# Stress selection and 'non-equilibrium phase coexistence'



$$\sigma_{tot} = \sigma_p + \eta \dot{\gamma}$$

$$\partial_t \sigma_p = f(\sigma_p, \dot{\gamma}) + \mathcal{D} \nabla^2 \sigma_p$$

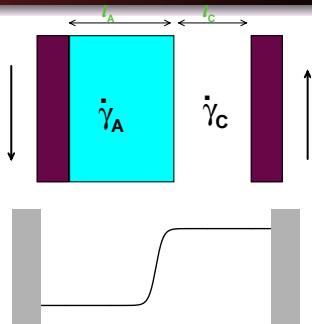
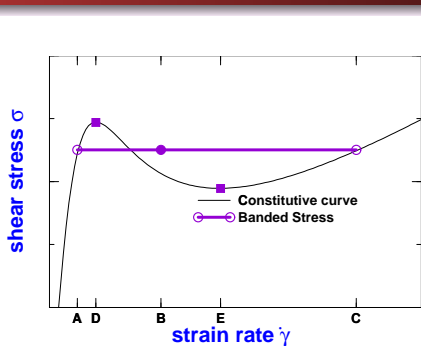
- **Unique  $\mathcal{D}$ -independent** stress determined by **inhomogeneous terms**.

[PDO/PMG PRA 1992; Lu/PDO/Ball PRL 2000] -

Concentration, finite stiffness, liquid crystallinity



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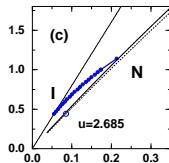
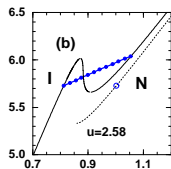
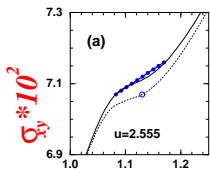
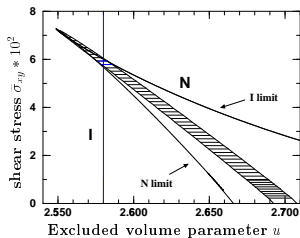
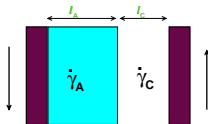
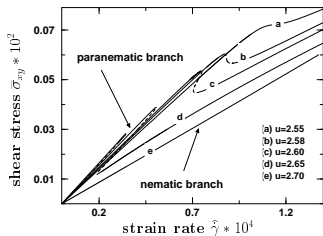
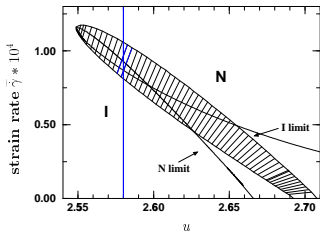


$$\sigma_{tot} = \sigma_p + \eta \dot{\gamma}$$

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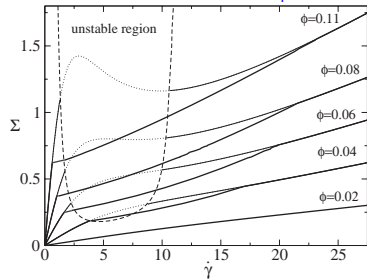
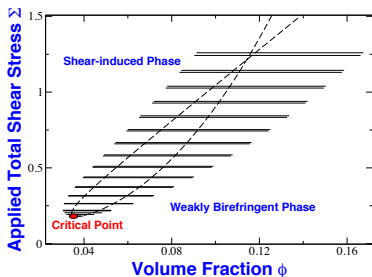
- **Unique  $\mathcal{D}$ -independent** stress determined by **inhomogeneous terms**.  
[PDO/PMG PRA 1992; Lu/PDO/Ball PRL 2000] - Concentration, finite stiffness, liquid crystallinity
- The **pressure difference** can lead to an  $N_2$ -driven instability at a free meniscus and sample ejection [Skorski & PDO, 2011].

# Liquid Crystals: Phase Diagram for Doi Model

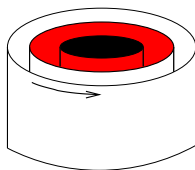


Strain Rate  $\dot{\gamma} * 10^4$

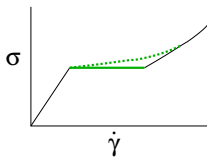
## Coupling to concentration fluctuations [Fielding/PDO, PRE 2003; EPJE (2003)]



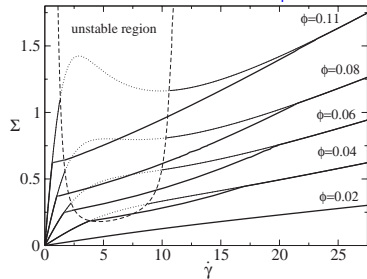
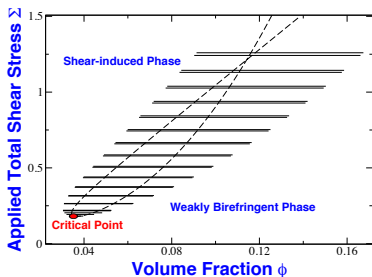
Common stress



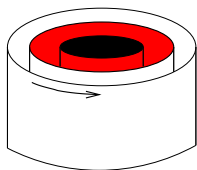
(shear thinning)



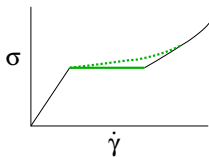
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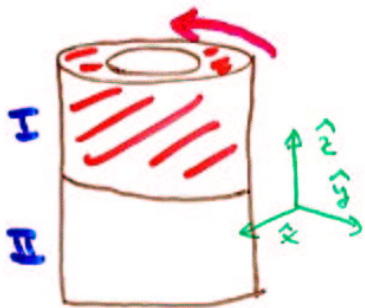
Common stress



(shear thinning)



Note instability even with positive slope  $d\Sigma/d\dot{\gamma} > 0$ .



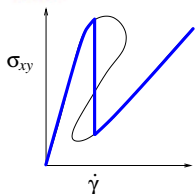
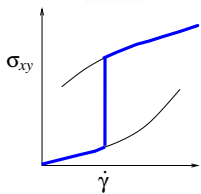
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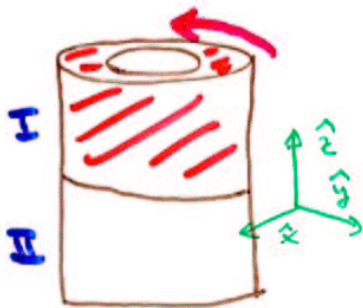
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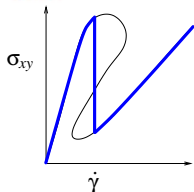
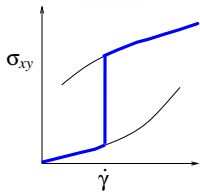

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$$\dot{\gamma}_I = \dot{\gamma}_H$$





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$$D_t \mathbf{Q} = \mathcal{L}(\phi, \nabla \mathbf{v}, \mathbf{Q})$$

$$\nabla \cdot \boldsymbol{\sigma} = 0$$

$$\nabla \parallel \hat{\mathbf{z}}$$

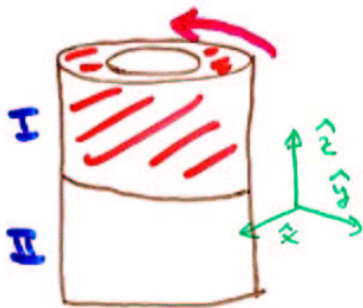
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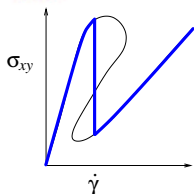
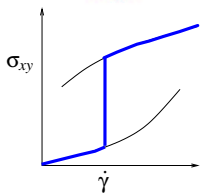
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$$D_t \phi = 0 \Rightarrow \mu^I = \mu^{II}$$

$$D_t \mathbf{Q} = 0 \quad \dot{\gamma} \text{ branches}$$



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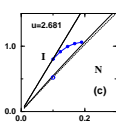
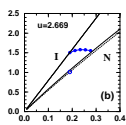
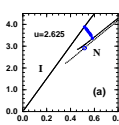
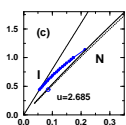
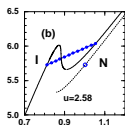
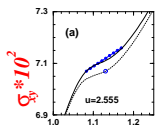
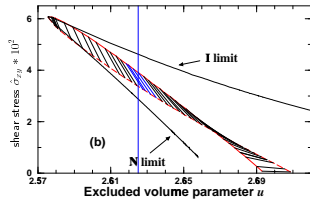
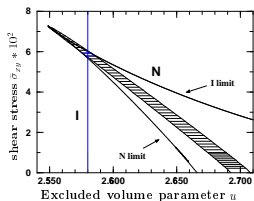
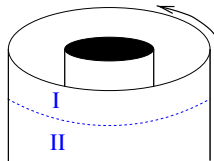
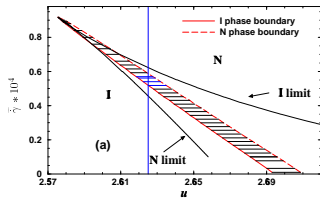
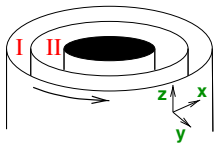
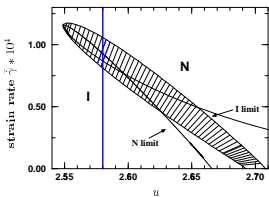
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Inhomogeneous pressure  $p(z)$ ;  
measure  $\sigma_{yy}(z)$ ; stable meniscus?

## Gradient Banding

## Vorticity Banding

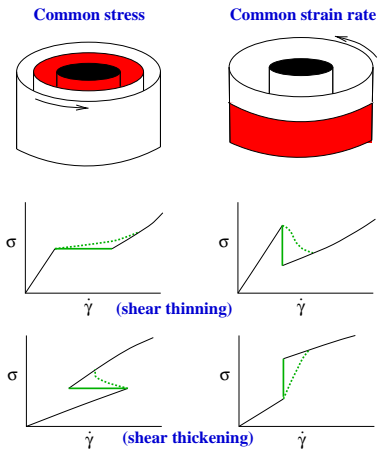


Strain Rate  $\dot{\gamma} * 10^4$

$\dot{\gamma} * 10^4$



# Typical Rheological Signatures of Banding [PDO, *Europhys. Lett.* (1999)]



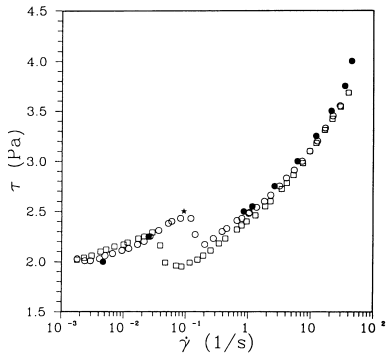
	<b>Gradient</b>	<b>Vorticity</b>
<b>Thin</b>	worms, liq. crystals	colloids, onions, sph. micelles
<b>Thick</b>	worms	onions DST?

- **Flat/Vertical plateaus:** same concentration.
- **Sloped plateaus:** different coexisting concentrations [Schmitt *et al.* 1996].
- **Vorticity banding:** few examples, not well studied.

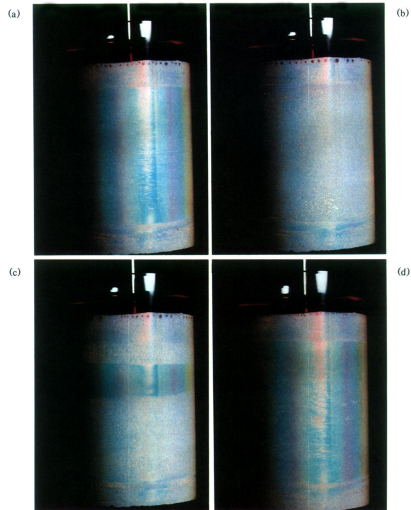
# Thinning Vorticity Banding

## Colloidal Crystalline Suspensions [Chen, Zukowski *et al.* PRL 1992, Langmuir 1994].

$$\phi = 0.45, 0.53; D = 230 \text{ nm.}$$



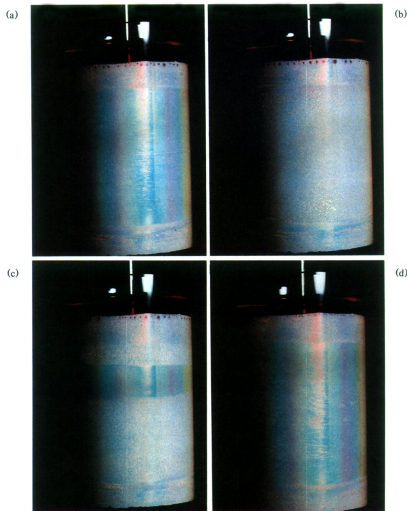
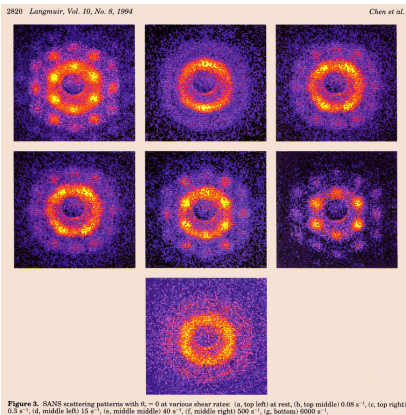
Polycrystalline +  
ordered layer sliding.



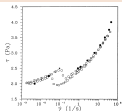
# Thinning Vorticity Banding

## Colloidal Crystalline Suspensions [Chen, Zukowski *et al.* PRL 1992, Langmuir 1994].

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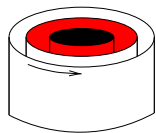


Polycrystalline +  
ordered layer sliding.

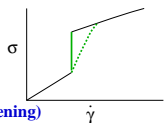
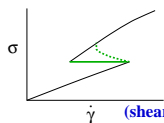
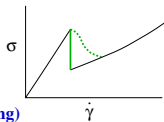
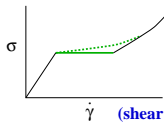
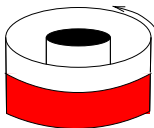


# Kinetics during Vorticity Banding?

Common stress

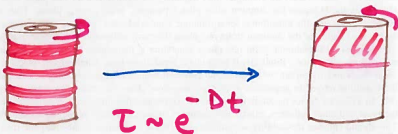


Common strain rate



- No symmetry-breaking field.
- 1D coarsening ?

- Expect slow kinetics (in Couette)



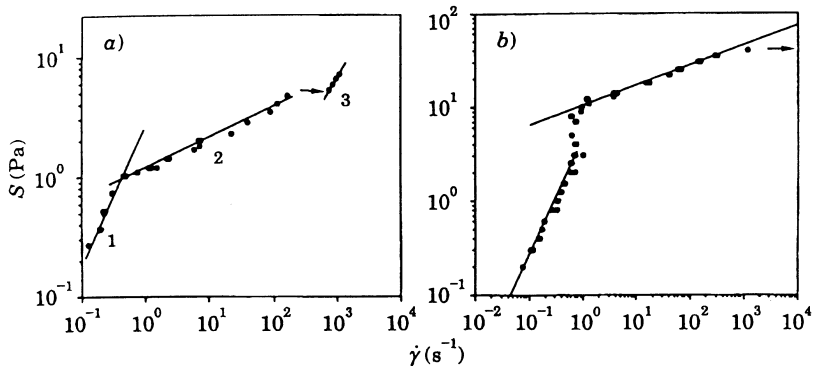
$$D: \partial_t Q = \dots + D \nabla^2 Q$$

1D coarsening

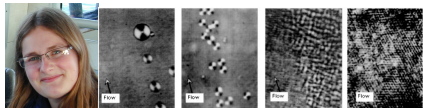
- Curvature of Couette Flow accelerates coarsening of gradient banking

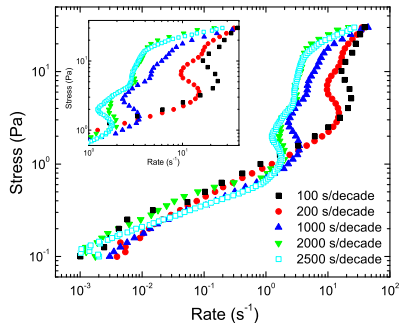
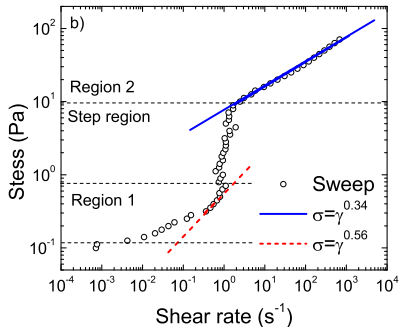


# Lamellar-to-onion transition... [Wilkins & PDO EPJE 2006]



- SDS/dodecane/pentanol/water [Diat & Roux 1993]
- Candidate for vorticity banding on the (“stress cliff”)??

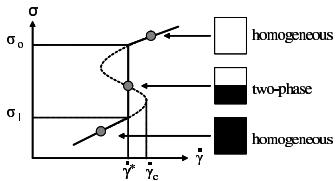




- No “Newtonian regime”
- Yield stress [smectic defect network?] and HB-like.

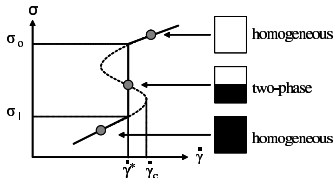
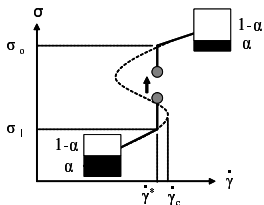
- Several transitions: vorticity bands?
- Hysteresis, slow transition.

# Cliff Region: Response to a stress jump $\sigma \rightarrow \sigma + \delta\sigma$ ?



$$\bar{\sigma}(\dot{\gamma}^*) = \alpha \sigma_0(\dot{\gamma})e + (1 - \alpha) \sigma_L(\dot{\gamma})$$

# Cliff Region: Response to a stress jump $\sigma \rightarrow \sigma + \delta\sigma$ ?

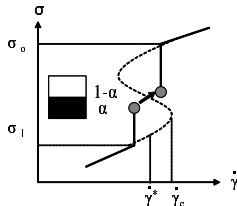
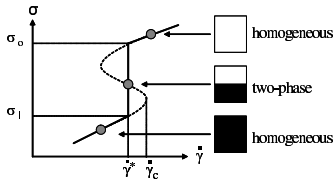
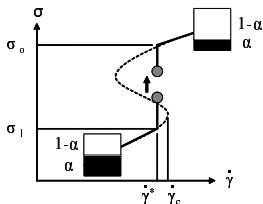


$$\bar{\sigma}(\dot{\gamma}^*) = \alpha \sigma_0(\dot{\gamma})e + (1 - \alpha) \sigma_L(\dot{\gamma})$$

- 1 Nucleate more onions,  $\alpha \rightarrow \alpha + \delta\alpha$  and return to selected strain rate.



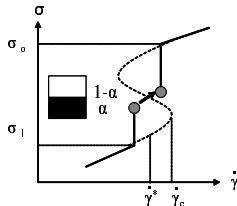
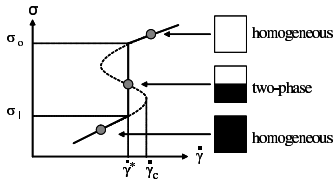
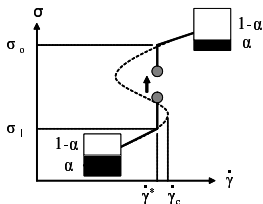
# Cliff Region: Response to a stress jump $\sigma \rightarrow \sigma + \delta\sigma$ ?



$$\bar{\sigma}(\dot{\gamma}^*) = \alpha \sigma_O(\dot{\gamma})_e + (1 - \alpha) \sigma_L(\dot{\gamma})$$

- 1 Nucleate more onions,  $\alpha \rightarrow \alpha + \delta\alpha$  and return to selected strain rate.
- 2 Resist nucleation, evolve strain rate  $\dot{\gamma}^* \rightarrow \dot{\gamma}^* + \delta\dot{\gamma}$ .

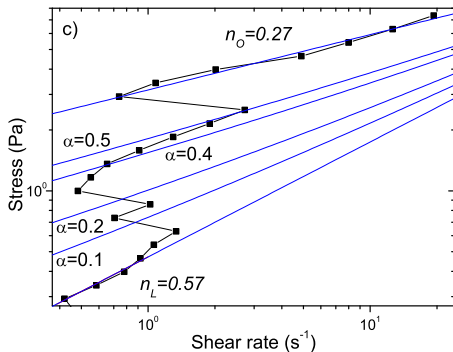
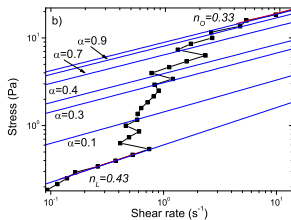
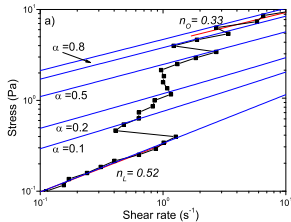
# Cliff Region: Response to a stress jump $\sigma \rightarrow \sigma + \delta\sigma$ ?



$$\bar{\sigma}(\dot{\gamma}^*) = \alpha \sigma_O(\dot{\gamma})e + (1 - \alpha) \sigma_L(\dot{\gamma})$$

- 1 Nucleate more onions,  $\alpha \rightarrow \alpha + \delta\alpha$  and return to selected strain rate.
  - 2 Resist nucleation, evolve strain rate  $\dot{\gamma}^* \rightarrow \dot{\gamma}^* + \delta\dot{\gamma}$ .
- Nucleation/growth harder if vorticity banded:
    - ▶ different interface configuration (not under shear)
    - ▶ No driving force in Couette flow?
    - ▶ Lamellae and onions cannot smoothly evolve to one another?

# Scaling the Stress Cliff? [Wilkins & PDO EPJE 2006]

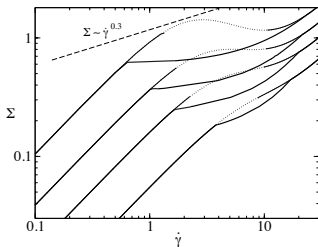


$$\bar{\sigma}(\dot{\gamma}^*) = \alpha \sigma_o(\dot{\gamma})e + (1 - \alpha) \sigma_L(\dot{\gamma})$$

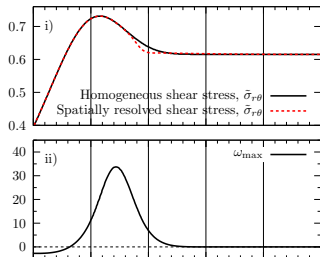
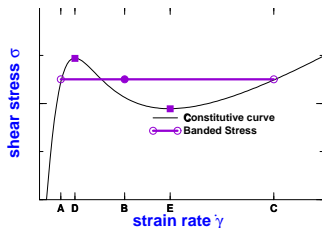
# Instability of shearing flows [Fielding JOR 2016]

**Overshoots** imply instability.....

- $\frac{d\sigma}{d\dot{\gamma}} < 0$  (constitutive)
- $\frac{d\sigma}{d\dot{\gamma}} < 0$  (transient/dynamic)
- concentration/order coupling



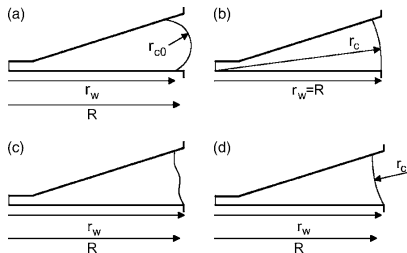
[Fielding & PDO, 2003]



[Adams & PDO, PRL 2008]

# Boundary Effects

T. SCHWEIZER AND M. STÖCKLI



[Schweizer & Stöckli JOR 2008]

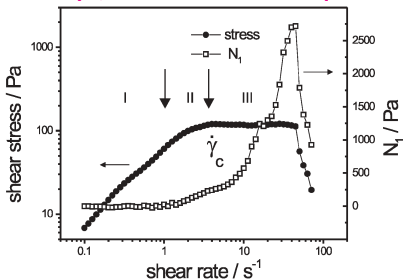
- Free surface balance
- Edge fracture
- Wall slip
- Instability in cone & plate and Couette flows.
- Stress gradient in channel/Poiseuille flows.

# Can shear banding induce edge fracture? [Skorski and PDO, JOR (2011)]

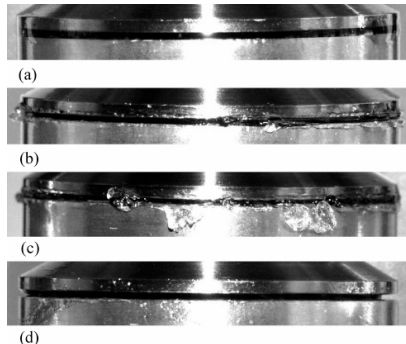
Edge fracture for  $|N_2| > \frac{\gamma}{R}$  [Tanner & Keentok, JOR 1983, Hemingway & Fielding arXiv:1703.05013].

## Wormlike Micelles

[Lopez-Gonzalez, Soft Matter 2006]



## Polymer Solutions



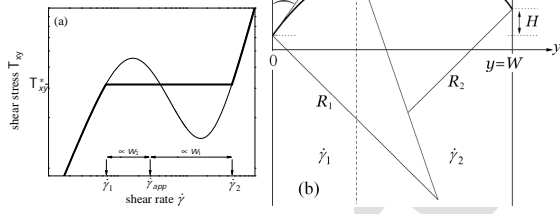
**Figure 2.** Pictures of the meniscus (a) after sample loading, (b) during the initial transient following step-up to 2500 Pa, (c) during the final portion of the transient, and (d) with extruded material removed from the exterior of the fixture. The pictures were taken without temperature control by the Peltier plate for better visualization of the meniscus shape.

[Inn, Wissbrun, and Denn, Macromolecules 2005],  
Heterogeneous flows

# Surface tension $\gamma$ vs. normal stress balance:

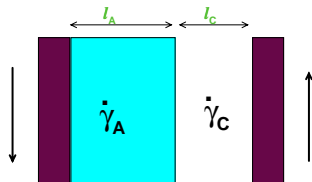
$$N_2 = T_{yy} - T_{zz}$$

side view



- Normal stress balance at meniscus:

$$\Delta N_2 = \Delta \left( \frac{\gamma}{R} \right)$$

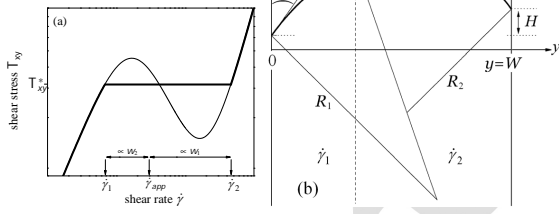


Top view

# Surface tension $\gamma$ vs. normal stress balance:

$$N_2 = T_{yy} - T_{zz}$$

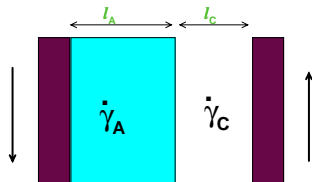
side view



- Normal stress balance at meniscus:

$$\Delta N_2 = \Delta \left( \frac{\gamma}{R} \right)$$

- Fixed contact angle  $\phi$ .



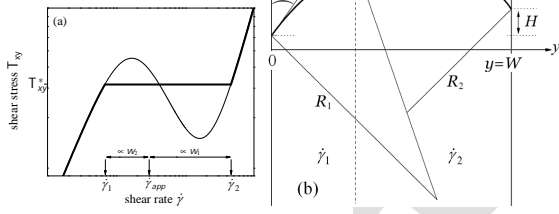
Top view



# Surface tension $\gamma$ vs. normal stress balance:

$$N_2 = T_{yy} - T_{zz}$$

side view

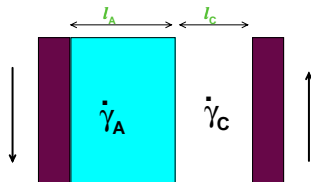


- Normal stress balance at meniscus:

$$\Delta N_2 = \Delta \left( \frac{\gamma}{R} \right)$$

- Fixed contact angle  $\phi$ .
- Control parameter:

$$A = \frac{LG}{\gamma} \frac{\Delta N_2}{G}$$

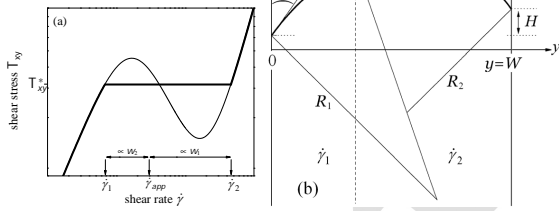


Top view

# Surface tension $\gamma$ vs. normal stress balance:

$$N_2 = T_{yy} - T_{zz}$$

side view

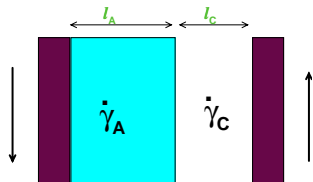


- Normal stress balance at meniscus:

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- Fixed contact angle  $\phi$ .
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$$A = \frac{LG}{\gamma} \frac{\Delta N_2}{G}$$

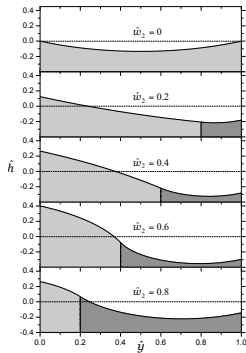


Top view

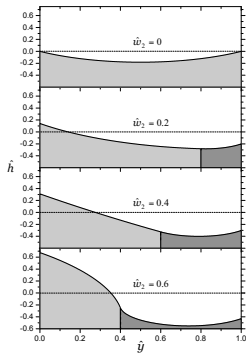
- Elastocapillary length

$$\xi_e = \gamma/G$$

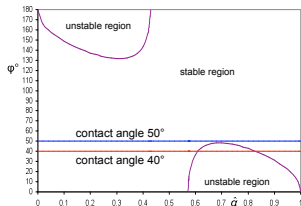
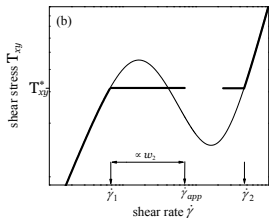
# Meniscus profiles across the plateau – Instability



$$\phi = 60^\circ$$



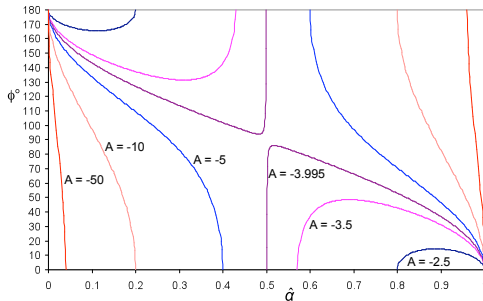
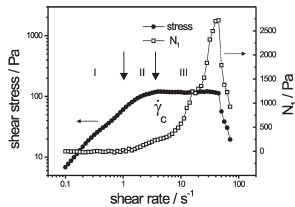
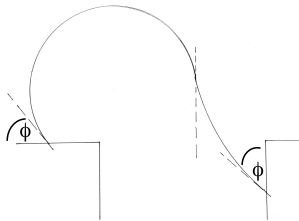
$$\phi = 40^\circ$$



$$A = \frac{LG \Delta N_2}{\gamma G} = -3.5$$

# Instabilities: $A = LG\Delta N_2/\gamma G$

## Expulsion



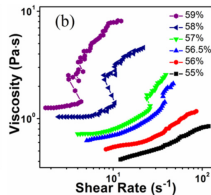
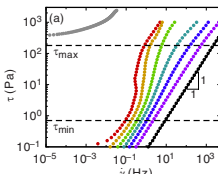
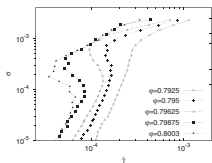
$A < -4$	Unstable
$-4 < A < -2$	Unstable or stable
$-2 < A < 0$	Stable

$$-A \simeq \begin{cases} 10 - 100 & \text{polymer solutions} \\ 0.5 - 4 & \text{wormlike micelles} \end{cases}$$

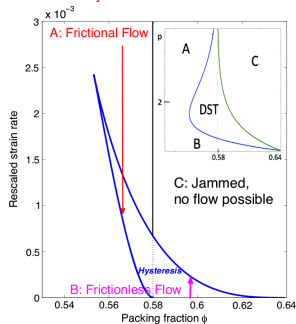
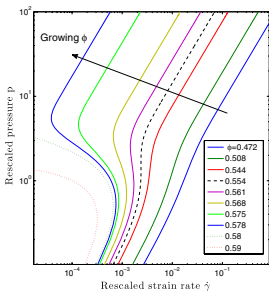
# Vorticity Banding: DST in colloids?

[Grob et al., PRE 2013, Brown & Jaeger Rep Prog Phys 2014, Pan et al. PRE 2015]

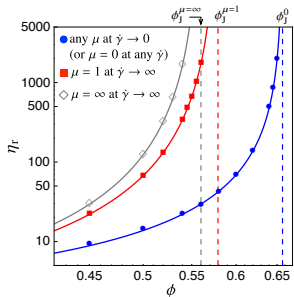
GROB, HEUSSINGER, AND ZIPPILIUS



[Wyart & Cates PRL 2014]



# Wyart/Cates model (non-Brownian suspensions) [PRL 2014]



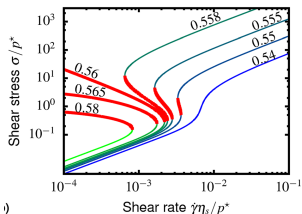
[Seto, Mari, Denn, Morriss, Cates, Wyart, ...]

Viscosity 
$$\eta = \frac{1}{\left(1 - \frac{\phi}{\phi_J(p_p)}\right)^2}$$

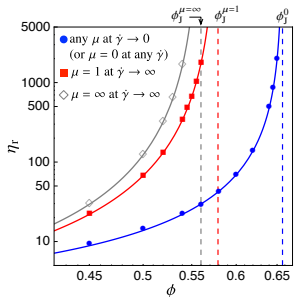
Jamming 
$$\phi_J = \phi_m f(p_p) + \phi_0(1 - f(p_p))$$

Contacts 
$$f(p_p) = (1 - e^{-p_p/p^*})$$

[Seto et al., PRL 2013, JOR 2014]



[Guy et al., JOR 2016]



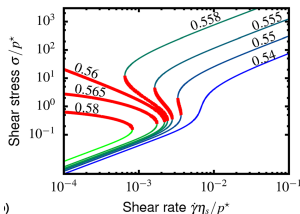
[Seto, Mari, Denn, Morriss, Cates, Wyart, ...]

Viscosity 
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Contacts 
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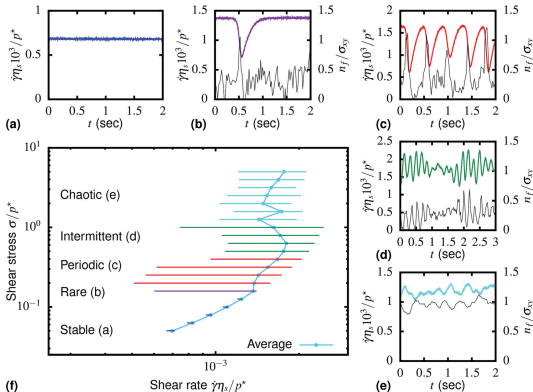
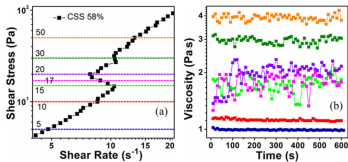


- Correct choice of fabric tensor  $\mathbf{A}$ ?
- Orientational order?
- $\mathbf{T} = \mathbf{T}(\mathbf{A}, f, \mathbf{D}, \dots)$
- Dynamics  $\partial_t f = \dots, \partial_t \mathbf{A} = \dots$  ?
- Frictional models are non-analytic and non-differentiable!

[Guy et al., JOR 2016]

# Unsteady vorticity banding in non-Brownian suspensions?

PHYSICAL REVIEW E **92**, 032202 (2015)



[Pan, de Gagny, Weber and Bonn, PRE 2015]

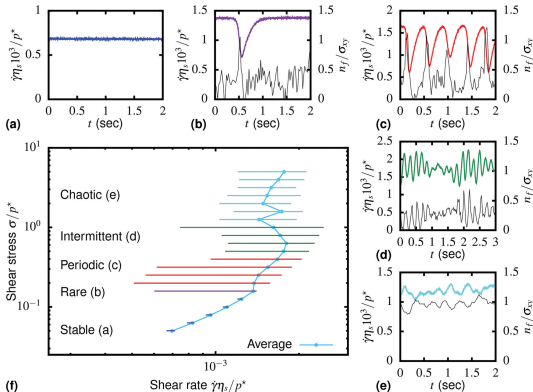
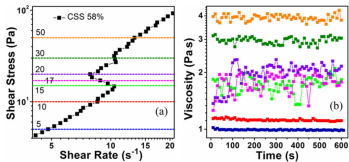
[Hermes et al., JOR 2016]

- **Unsteady bands:** cannot separately balance particle stress  $p_{p,zz}(f, \phi)$  and solvent pressure at the interface? General feature for non-Brownian suspensions?



# Unsteady vorticity banding in non-Brownian suspensions?

PHYSICAL REVIEW E **92**, 032202 (2015)



[Pan, de Gagny, Weber and Bonn, PRE 2015]

[Hermes et al., JOR 2016]

- **Unsteady bands:** cannot separately balance particle stress  $p_{p,zz}(f, \phi)$  and solvent pressure at the interface? General feature for non-Brownian suspensions? **What about the meniscus? [open boundaries in expts]**

## Final thoughts.....

There are many models for yielding materials (fluids and solids); usually these are not treated at a particle level.

- STZ-motivated models [Lemaitre, Langer, Manning, Falk, ...]
- Long-range elastic/viscoplastic models [Picard, Lequeux, Ajdari, Martens, Barrat, ...]
- Fluidity models [Bonn, Mansard/Colin, Ovarlez, Coussot, ...] .

# Work done and funded by...

**G Wilkins** (Patent Attorney)  
**Stan Skorski** (Leeds)  
**Elian Masnada** (Georgetown)

**James Adams** (Surrey)  
**Suzanne Fielding** (Durham)  
**Richard Graham** (Nottingham)

**O Radulescu** (Montpellier)  
**C-Y D Lu** (Taiwan)

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**R Ball** (Warwick)  
**S Manneville** (Lyon)  
**A Colin** (ESPCI)  
**S Lerouge** (Paris)

The logo for EPSCRC, featuring the letters 'EPSCRC' in a bold, purple, sans-serif font, centered between two horizontal green lines.The logo for SoftComp, consisting of a green square with a white stylized 'S' shape inside, followed by the text 'SoftComp' in a bold blue font and 'SOFT MATTER COMPOSITES' in a smaller blue font below it.The logo for dynacop, featuring the word 'dynacop' in a blue, lowercase, sans-serif font with a stylized 'y'. Above the text are several overlapping curved lines in red, yellow, and blue. Above the lines is the text 'Dynamics of Architecturally Complex Polymers' in a small font.

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