Solving quantum light propagation through atomic ensembles with matrix product states

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Designer Quantum Systems out of Equilibrium November 18, 2016

# Acknowledgments





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## Motivation

- Many emerging experimental systems that can achieve strong interactions between photons
  - Quantum information processing or strongly correlated states of light





Atoms / Fabry- Atoms and solid-state emitters / Perot cavities nanophotonic cavities

**Circuit QED** 

Beyond cavity QED

**Cavity QED** 



a=371nm g=250nm g=250nm g=250nmg=250nm



Atomic Rydberg gases Atoms coupled to nanofibers and PhC waveguides

Waveguide QED

## Theoretical approaches

- Theoretical complexity of cavity QED vs. non-cavity QED systems differs dramatically
- Cavity QED: "easy" to solve



Input-output equation: relates observable fields outside the cavity to dynamics inside

$$a_{out}(t) = a_{in}(t) + \sqrt{\kappa}a(t)$$

## Theoretical approaches

### Non-cavity QED systems



- Common feature: coupling to *propagating* fields
- One photon already represents an infinite Hilbert space!

$$\int d\omega \ f(\omega) a_{\omega}^{\dagger} | \text{vac} \rangle$$

# Many-body physics with light?

• Possible to prepare strongly correlated states of light?

#### **Crystal of photons?**

#### Many interesting proposals:

Crystallization of strongly interacting photons in a nonlinear optical fibre

D. E. CHANG<sup>1</sup>, V. GRITSEV<sup>1</sup>, G. MORIGI<sup>2</sup>, V. VULETIĆ<sup>3</sup>, M. D. LUKIN<sup>1</sup> AND E. A. DEMLER<sup>1\*</sup>

#### Dissipation induced Tonks-Girardeau gas of photons

M. Kiffner<sup>1</sup> and M. J. Hartmann<sup>1</sup> <sup>1</sup>Technische Universität München, Physik-Department I, James-Franck-Straße, 85748 Garching, Germany

#### Wigner Crystallization of Single Photons in Cold Rydberg Ensembles

Johannes Otterbach\* Physics Department. Harvard University. Cambridge. 02138 Massachusetts. USA

Matthias Moos, Dominik Muth, and Michael Fleischhauer Fachbereich Physik und Forschungszentrum OPTIMAS, Technische Universität Kaiserslautern, 67663 Kaiserslautern, Germany (Received 30 April 2013; revised manuscript received 16 July 2013; published 9 September 2013)

### Highly nonlocal optical nonlinearities in atoms trapped near a waveguide

EPHRAIM SHAHMOON,<sup>1,2,3,\*</sup> PJOTRS GRIŠINS,<sup>4</sup> HANS PETER STIMMING,<sup>5,6</sup> IGOR MAZETS,<sup>4,6</sup> AND GERSHON KURIZKI<sup>1</sup>

#### Correlated photon dynamics in dissipative Rydberg media

Emil Zeuthen,<sup>1, 2, \*</sup> Michael J. Gullans,<sup>3</sup> Mohammad F. Maghrebi,<sup>3</sup> and Alexey V. Gorshkov<sup>3</sup>

### • Challenge to analytically solve or numerically verify!

• Continuum, driven, open, out-of-equilibrium, ...

## Theoretical approaches

• Work in regimes of classical optics or Gaussian physics



Photon storage (EIT), EPR entanglement of atomic ensembles and light, spin squeezing, ...

 Computational black box for quantum case: discretize space and solve quantum discretized wave equations



- Hilbert space explodes, even for empty space!
- In practice, limited to 2-3 total excitations in system, excluding quantum jumps
- Goal: approach where limit of empty space becomes trivial

### Outline of new approach

• Real degrees of freedom are just the atoms, encoded in a discrete "spin" Hamiltonian

$$- |e\rangle - |e\rangle - |e\rangle - |e\rangle - |e\rangle$$

$$H_{eff} \sim G(r_i, r_j) \sigma_{eg}^{(i)} \sigma_{ge}^{(j)}$$

$$- |g\rangle - |g\rangle - |g\rangle$$

Re-construct quantum field from generalized input-output equation

 $\rho_{\text{atoms}}(t)$ Photodetectors  $\langle E^{\dagger}(t)E^{\dagger}(t')E(t')E(t) \rangle$ 

 Most systems mappable to effective 1D model with N < 100 atoms



## **Electromagnetic Green's function**

 Formalism: quantum atom-light interactions encoded in (classical) E&M Green's function

 $\left[ (\nabla \times \nabla \times) - \omega^2 \epsilon(r, \omega) / c^2 \right] G_{\alpha\beta}(r, r', \omega) = \delta(r - r') \otimes I$ 

 G describes electric field at point r, of a normalized oscillating dipole at r'



 $\alpha, \beta$  encode vector nature of dipole source and field

## The fields: classical analogy

Classical scattering from polarizable particles



- Know the radiation pattern for a dipole
- Can calculate the total field

$$\hat{E}(r,\omega) = \hat{E}_{in}(r,\omega) + \alpha \sum G(r,r_i,\omega)p_i(\omega)$$

**Becomes convolution in time domain** 

- Classical and quantum fields propagate the same way
- Generalized "input-output" equation in time for atoms

$$\hat{E}(r,t) = \hat{E}_{in}(r,t) + \mu_0 \omega_{eg}^2 d_0^2 \sum_i G(r,r_i,\omega_{eg}) \sigma_{ge}^i(t)$$
 Field encoded in atoms!

### Substitute field into equations of motion for atoms

- Master equation for atoms alone
- Coherent evolution (emission and re-absorption)

$$H_{\rm eff} = -\mu_0 d_0^2 \omega_{eg}^2 \sum_{i,j} \left( \operatorname{Re} G(r_j, r_i, \omega_{eg}) \right) \sigma_{eg}^i \sigma_{ge}^j$$

**Atomic dynamics** 

• Dissipation (spontaneous emission)

$$\dot{\rho} = L[\rho] = -\mu_0 d_0^2 \omega_{eg}^2 \sum_{i,j} \left( \operatorname{Im} G(r_j, r_i, \omega_{eg}) \right) \left( \sigma_{eg}^i \sigma_{ge}^j \rho + \rho \sigma_{eg}^i \sigma_{ge}^j - 2\sigma_{ge}^j \rho \sigma_{eg}^i \right)$$

• In short (non-Hermitian Hamiltonian):

$$H_{\rm eff} = -\mu_0 d_0^2 \omega_{eg}^2 \sum_{i,j} G(r_j, r_i, \omega_{eg}) \sigma_{eg}^i \sigma_{ge}^j + \text{quantum jumps}$$

# Mapping to 1D



- Write down effective 1D equation:  $(\partial_t + \partial_z)E(z,t) \sim iP_{ge}(z,t)$
- Approximate: independent emission into other modes
- Idea: invent an equivalent 1D "waveguide" with huge interaction probability and N  $\sim 10^2$  atoms

Ein





• Spin model in 1D

$$H_{\rm eff} = -\mu_0 d_0^2 \omega_{eg}^2 \sum_{i,j} G(r_j, r_i, \omega_{eg}) \sigma_{eg}^i \sigma_{ge}^j \longrightarrow -\frac{i\Gamma_{1D}}{2} \sum_{i,j} e^{ik|z_i - z_j|} \sigma_{eg}^i \sigma_{ge}^j$$

Plane wave propagation in 1D

• Single atom:

$$H_{\rm eff} = -\frac{i\Gamma_{1D}}{2}\sigma_{ee}$$
 •  $\Gamma_{1D}$  is single-atom spontaneous emission rate

• Not very physical – photon always emitted into 1D channel

$$- |e\rangle \Gamma_{1D}$$

$$- |g\rangle$$

## 1D spin model

• Add phenomenological decay rate into other channels

$$H_{\text{eff}} = -\frac{i\Gamma_{1\text{D}}}{2} \sum_{i,j} \exp(ik|z_i - z_j|) \sigma_{eg}^i \sigma_{ge}^j - \frac{i\Gamma'}{2} \sum_i \sigma_{ee}^i$$

+ quantum jumps

Model connects quantitatively to experimental 1D systems



## Calculating output fields

- Apply input-output formalism to calculate fields
  - Randomly positioned two-level atoms, many runs



Linear transmission spectrum N=10 atoms,  $\frac{\Gamma_{1D}}{\Gamma'} = 0.3$ Randomly distributed on 100 sites (10 runs)

• On resonance:

$$T \approx \exp(-\text{OD})$$
  
 $\text{OD} = 2N_{\text{atoms}}\Gamma_{1\text{D}}/\Gamma' \quad \text{(for }\Gamma_{1\text{D}} < \Gamma')$ 

 Establishes connection with free-space ensembles with much higher N<sub>atoms</sub>!

# • Exact dynamics of spin Hamiltonian on full Hilbert space: $N_{\rm atoms} < 20$

**Dynamics with MPS** 

• Larger systems: use matrix product states

• MPS in a nutshell:  $|\psi\rangle = \sum_{\{i=\uparrow,\downarrow\}} c_{i_1i_2i_3...i_N} |i_1i_2i_3...i_N\rangle$  Exact

**Exact wavefunction** 

**Rank-N tensor** 

Re-shape tensor

 $c_{i_1 i_2 \dots i_N} = M^{(i_1)} M^{(i_2)} \dots M^{(i_N)}$ 



## **Dynamics with MPS**

- Exact representation: matrix size D grows exponentially with N
- Approximation: truncate at max size D x D
- Interpretation:

sa



Good representation if s<sub>a</sub> decays rapidly (low entanglement entropy)



## **Dynamics with MPS**

• Procedure for dynamics



Initial state  $|\psi(t_0)\rangle$ 

$$\begin{split} |\psi(t_0+\Delta t)\rangle \approx \\ (1-iH_{\rm eff}\Delta t)|\psi(t_0)\rangle \end{split}$$



• Size of MPS grows after evolution step



## Test case: vacuum induced transparency

- Does the formalism work??
- Few cases of known many-body solutions
  - Exception: vacuum induced transparency
- Electromagnetically induced transparency:



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- Does the formalism work??
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  - Exception: vacuum induced transparency
- Electromagnetically induced transparency:

- Photon hybridizes with a spin excitation,  $|g\rangle \rightarrow |s\rangle$
- Reduced group velocity  $v_g \propto \Omega_{\rm cl}^2 \ll c$  and transparency window

## Vacuum induced transparency

Replace classical control field with cavity vacuum



- Pulse propagation acquires nonlinearity of atom-cavity interaction
  - Photon number-dependent group velocity,  $v_g \propto n_{
    m ph}$



• Spatial separation of photon number

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Theory: Fleischhauer (PRL, 2010)
Experiment: Vuletic (Science, 2011)
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### • $N_{\text{atoms}} = 100, \Gamma_{1D} = 2, \Gamma' = 1, g = 2, \kappa = 0.03$ (OD=400)

**MPS** simulations

• Coherent state input:



### • $N_{\text{atoms}} = 100, \Gamma_{1D} = 2, \Gamma' = 1, g = 2, \kappa = 0.03$

**MPS** simulations



### • $N_{\text{atoms}} = 100, \Gamma_{1D} = 2, \Gamma' = 1, g = 2, \kappa = 0.03$

**MPS** simulations

• Effect of partial absorption



### • $N_{\text{atoms}} = 100, \Gamma_{1D} = 2, \Gamma' = 1, g = 2, \kappa = 0.03$

**MPS** simulations

• Effect of partial absorption



## **Two-time correlation functions**

• Distortion of a two-photon wavepacket



# **Two-time correlation functions**

• Distortion of a two-photon wavepacket





• Promising technique to solve for many-photon dynamics



- Why does it work, and when does it fail?
  - Nature of entanglement growth in dissipative systems
- Atom-light interactions as a quantum spin model
  - Other interesting consequences?