

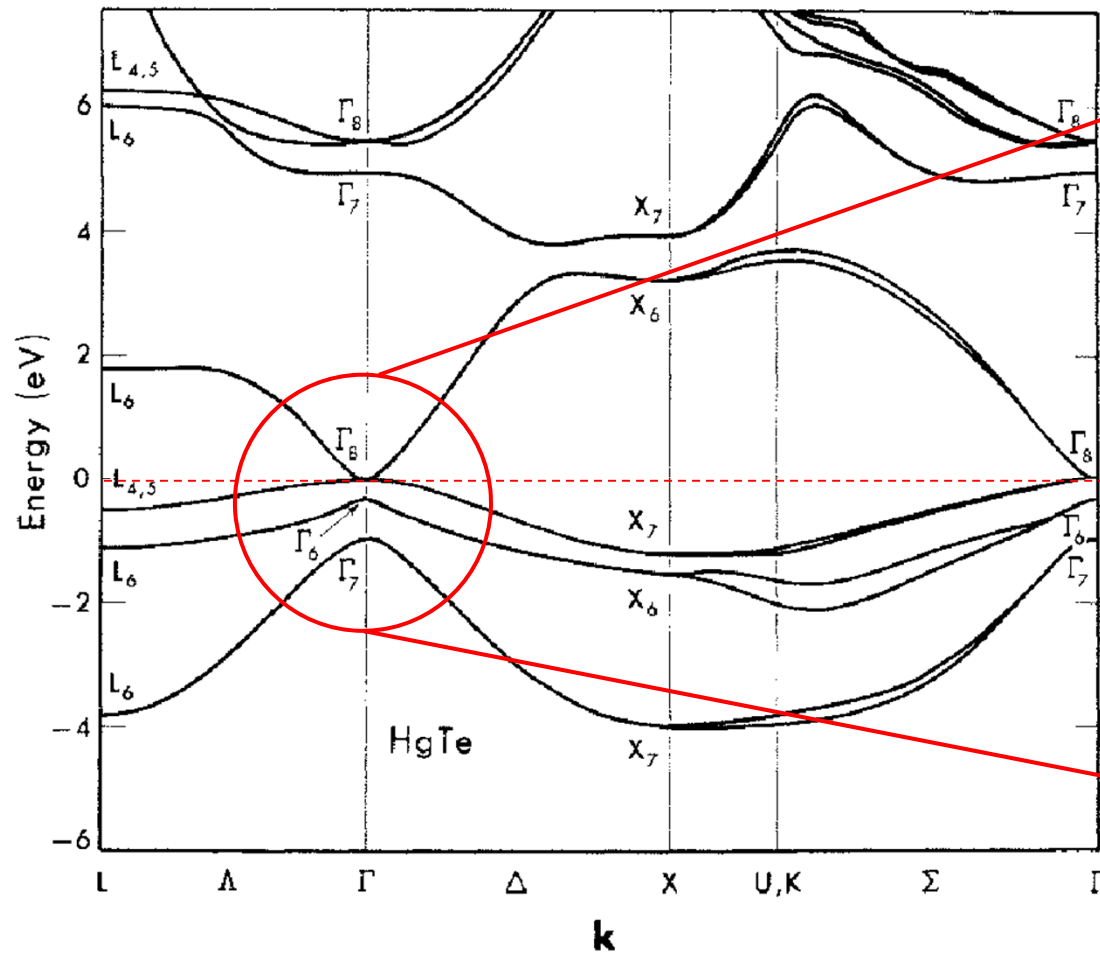
# Dirac Fermions in HgTe

**Laurens W. Molenkamp**

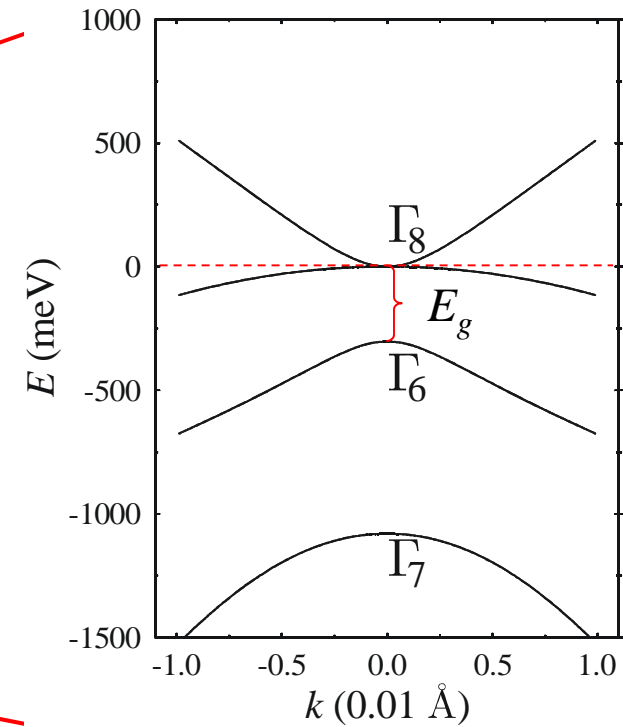
Physikalisches Institut, EP3  
Universität Würzburg

- HgTe/CdTe bandstructure, quantum spin Hall effect
- HgTe as a Dirac system
- Dirac surface states of strained bulk HgTe

**band structure**



**semi-metal or semiconductor**



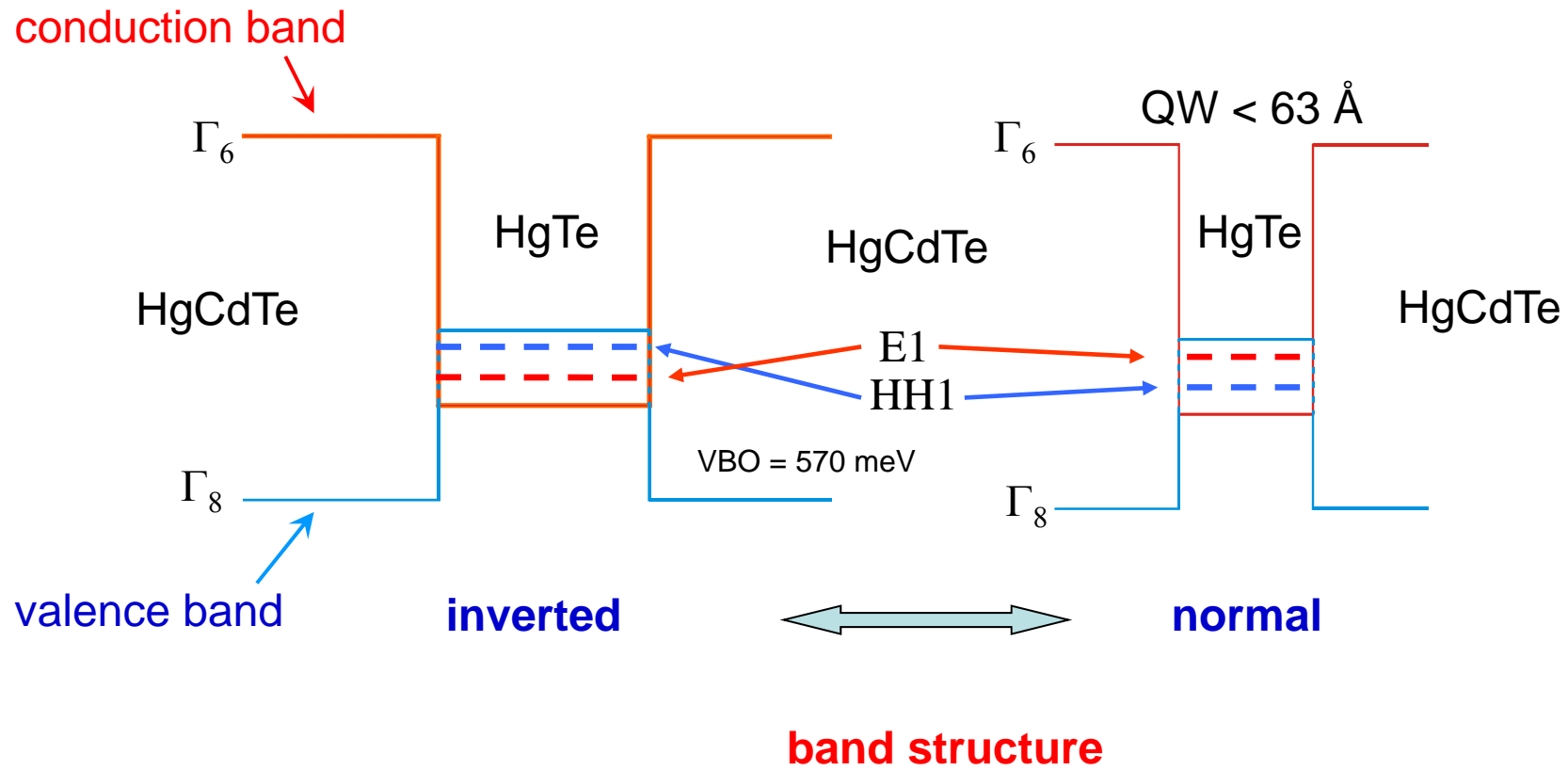
**fundamental energy gap**

$$E^{\Gamma_6} - E^{\Gamma_8} \approx -300 \text{ meV}$$

# HgTe-Quantum Wells



Type-III QW

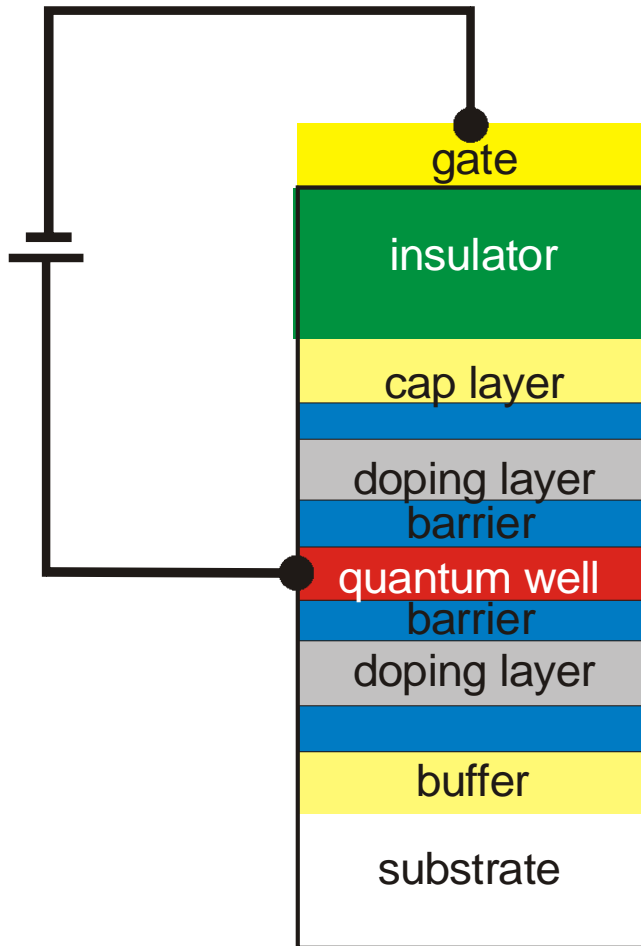


# Layer Structure

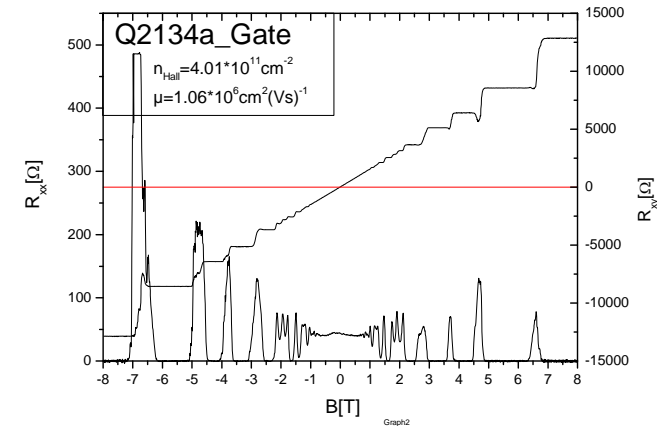


Carrier densities:  $n_s = 1 \times 10^{11} \dots 2 \times 10^{12} \text{ cm}^{-2}$

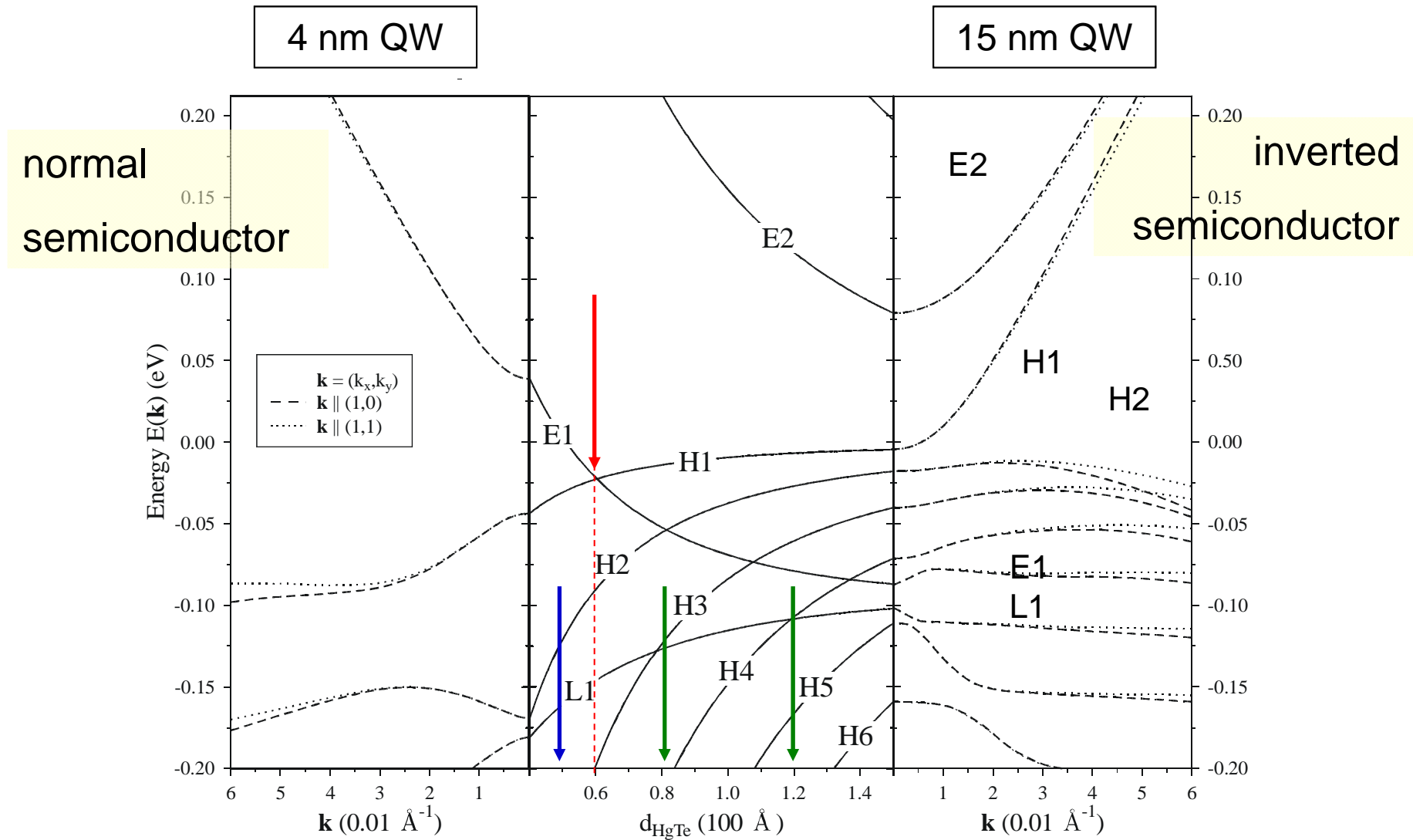
Carrier mobilities:  $\mu = 1 \times 10^5 \dots 1.5 \times 10^6 \text{ cm}^2/\text{Vs}$



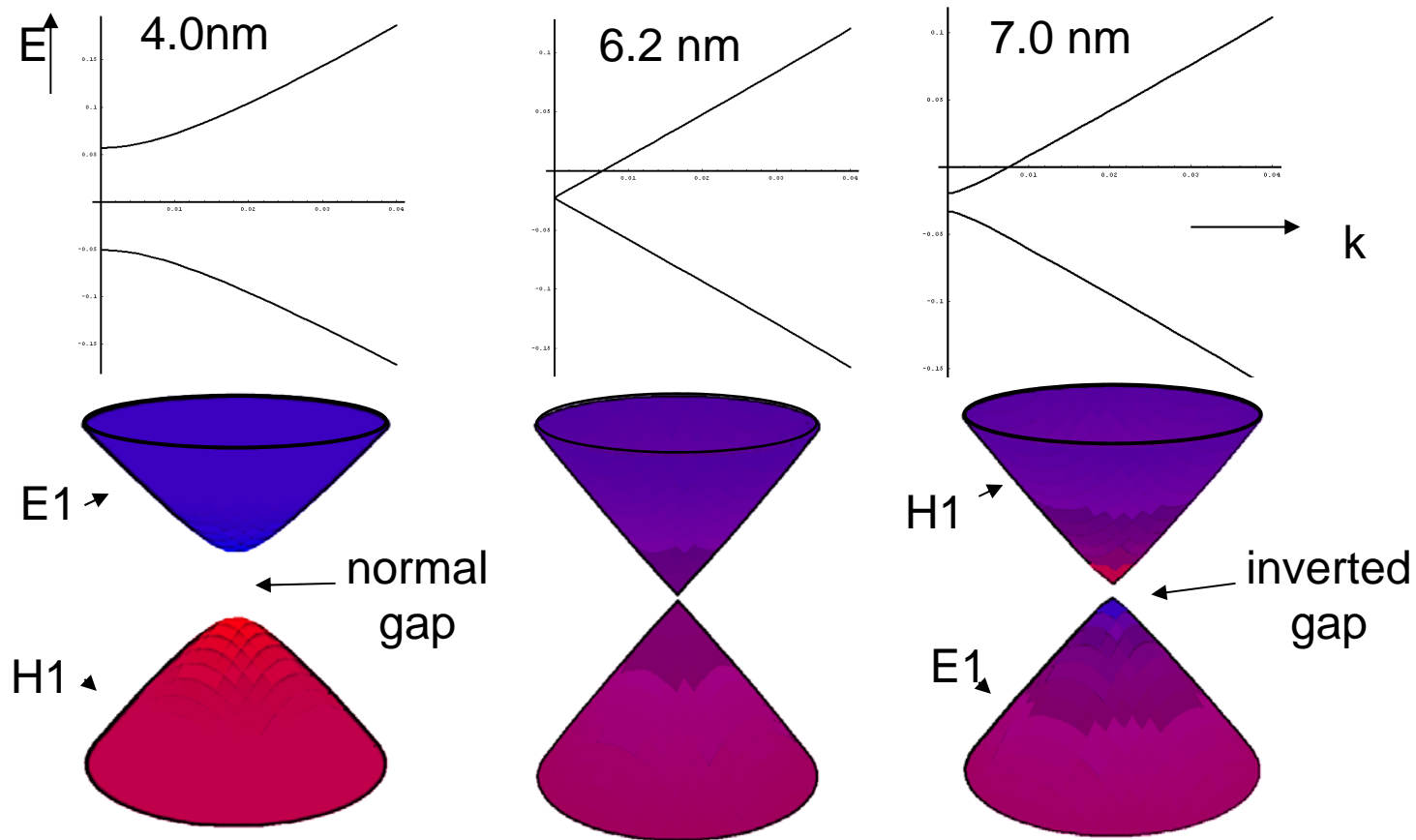
- Au
- 100 nm  $\text{Si}_3\text{N}_4/\text{SiO}_2$
- 25 nm CdTe
- 10 nm HgCdTe  $x = 0.7$
- 9 nm HgCdTe with I
- 10 nm HgCdTe  $x = 0.7$
- 4 - 12 nm HgTe
- 10 nm HgCdTe  $x = 0.7$
- 9 nm HgCdTe with I
- 10 nm HgCdTe  $x = 0.7$
- 25 nm CdTe
- CdZnTe(001)



symmetric or asymmetric doping



B.A Bernevig, T.L. Hughes, S.C. Zhang, Science **314**, 1757 (2006)



# **QSH**

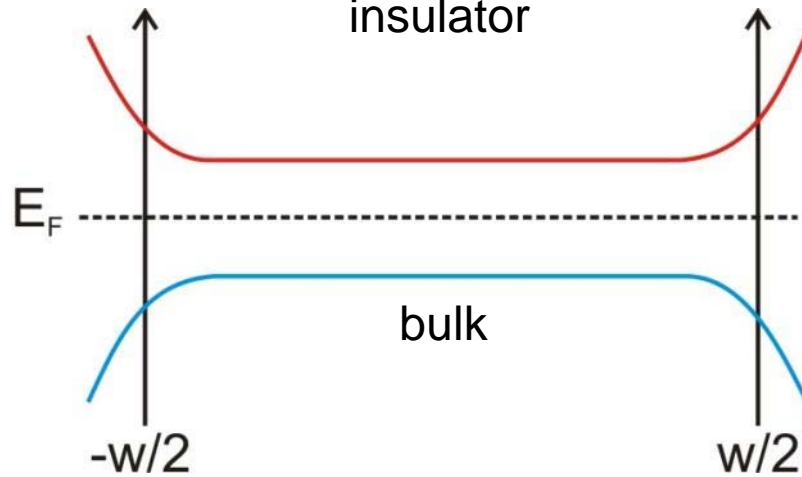
# **insulator**



# QSHE, Simplified Picture

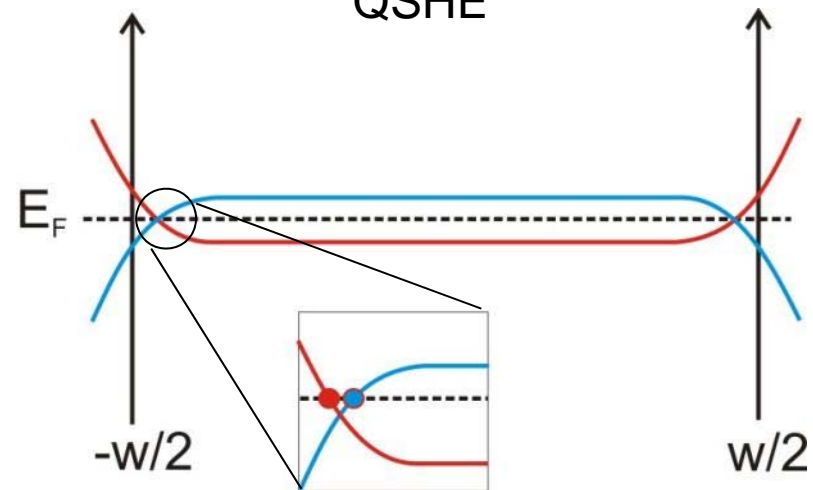
$m > 0$

normal  
insulator

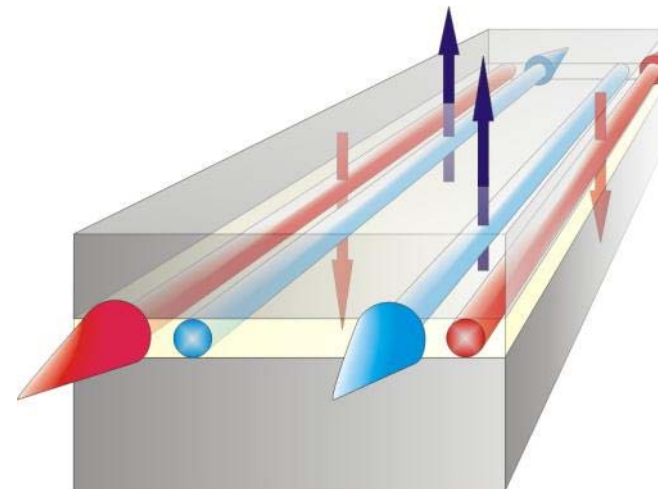


$m < 0$

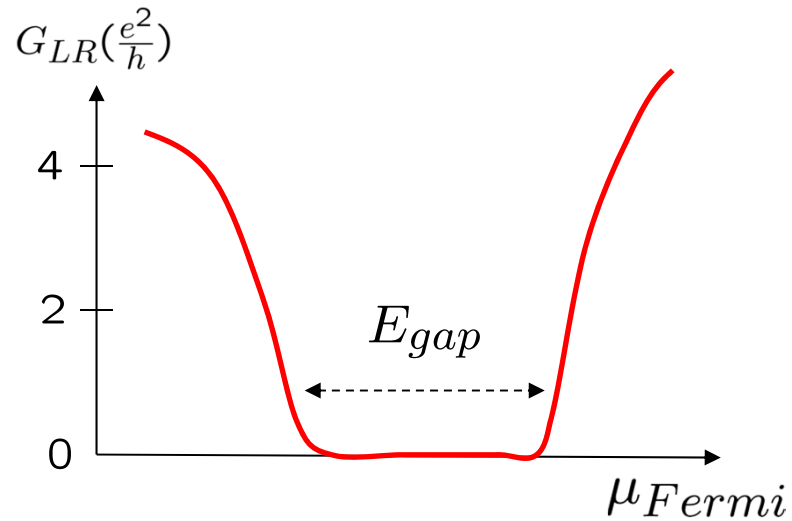
QSHE



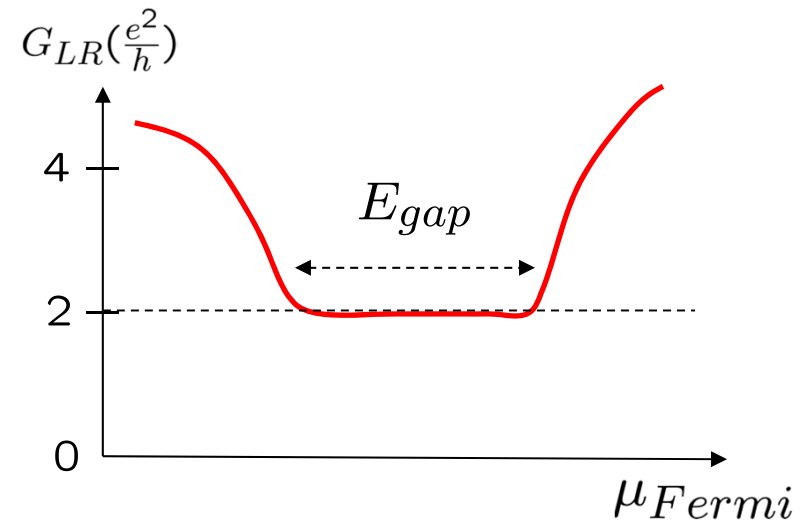
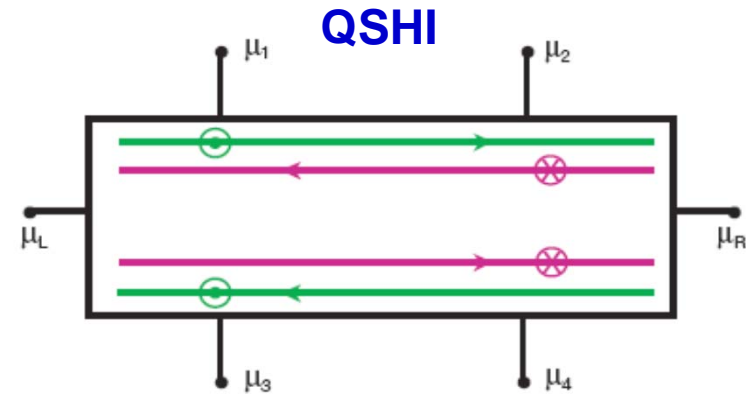
entire sample  
insulating



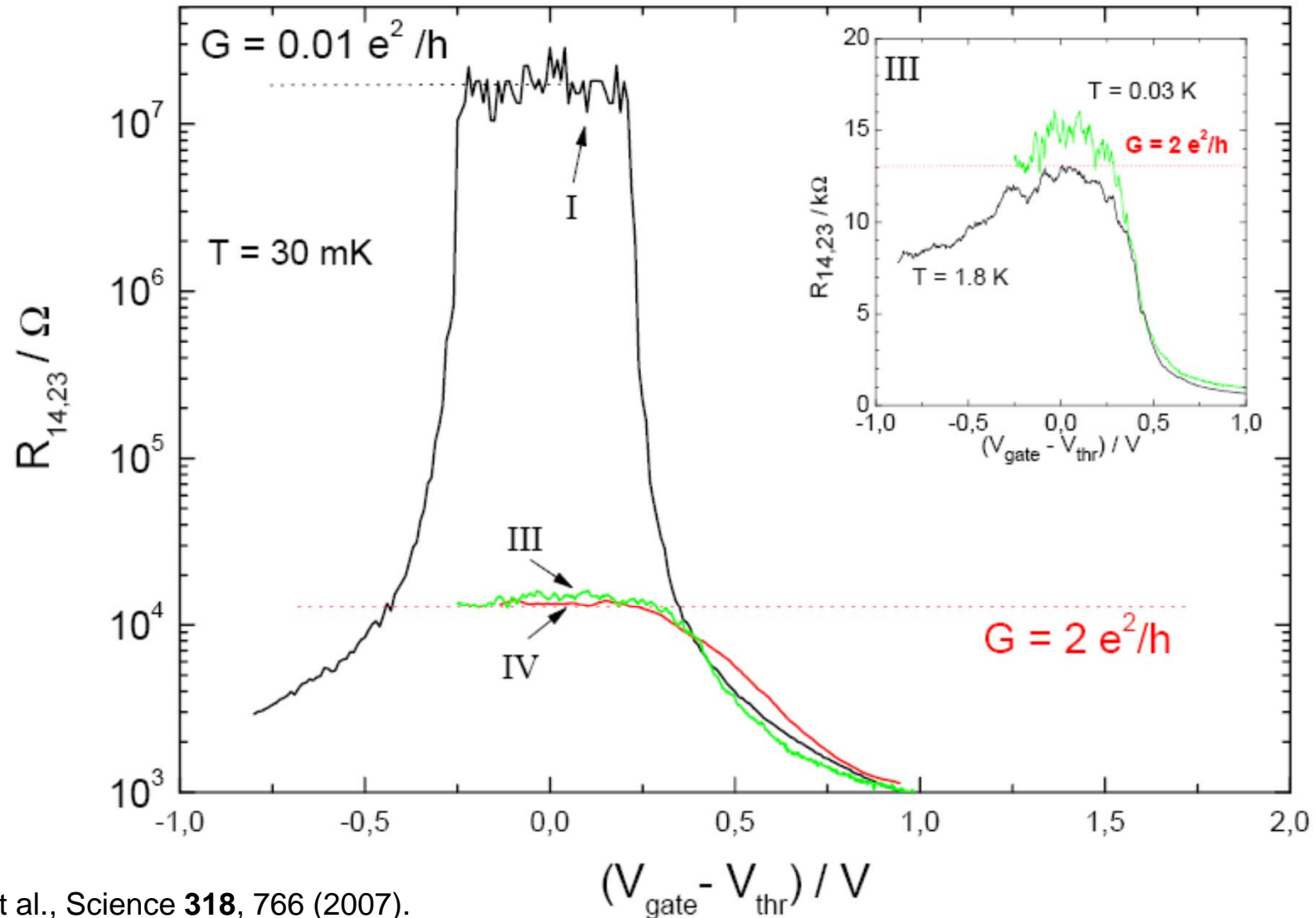
normal insulator state

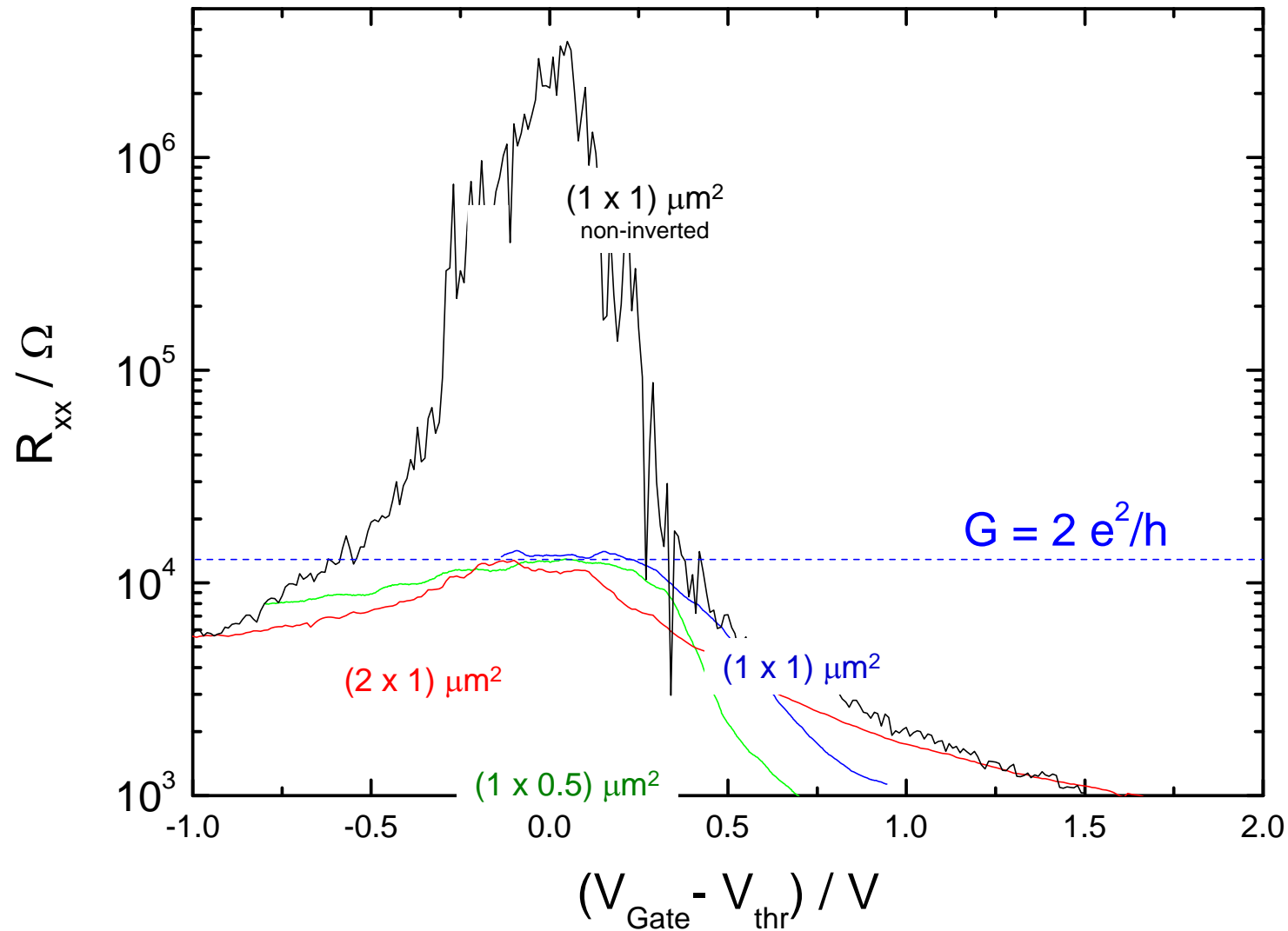


$d < d_c$ , normal regime



$d > d_c$ , inverted regime

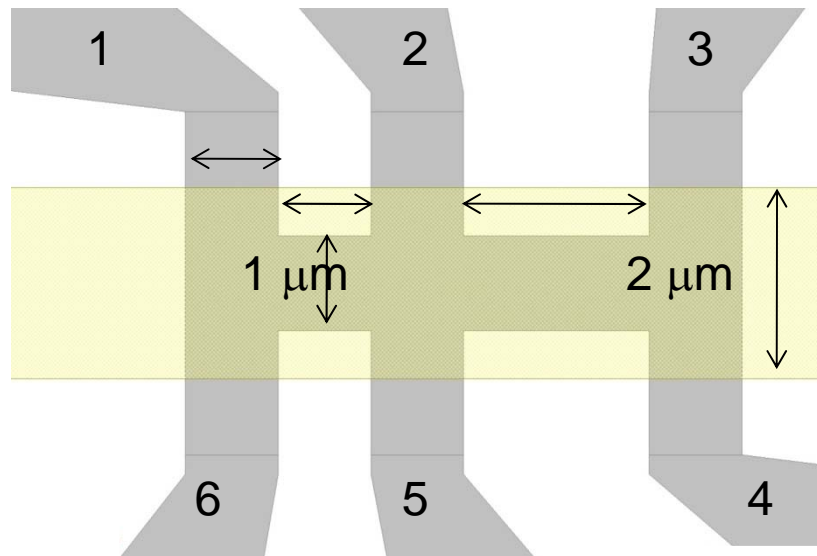




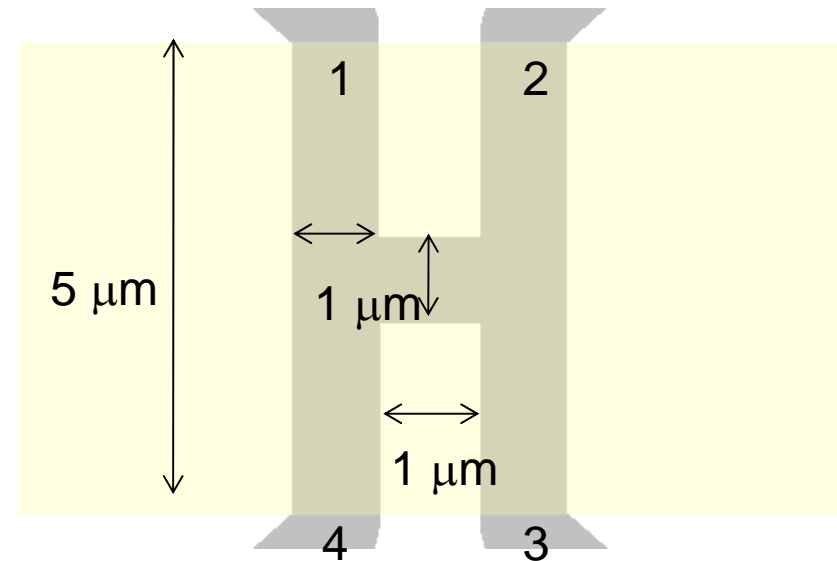
# Verify helical edge state transport



(a)

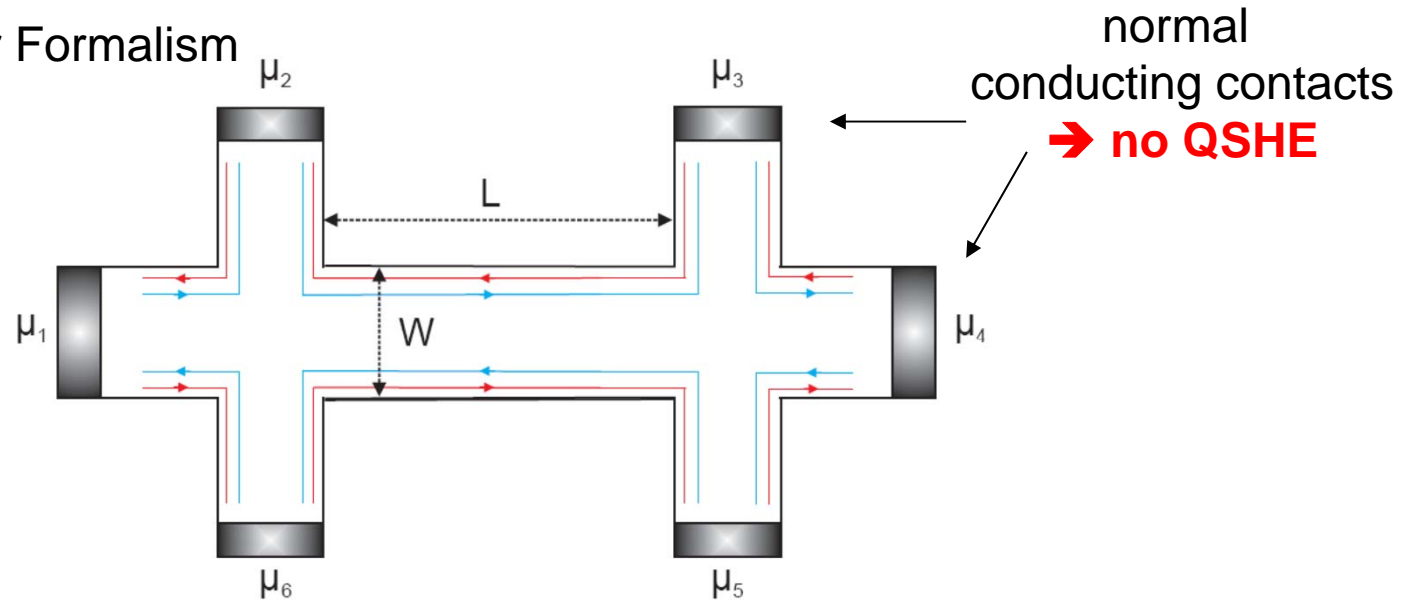


(b)



Multiterminal /Non-local transport samples

Landauer-Büttiker Formalism



$$T = \begin{pmatrix} -2 & 1 & 0 & 0 & 0 & 1 \\ 1 & -2 & 1 & 0 & 0 & 0 \\ 0 & 1 & -2 & 1 & 0 & 0 \\ 0 & 0 & 1 & -2 & 1 & 0 \\ 0 & 0 & 0 & 1 & -2 & 1 \\ 1 & 0 & 0 & 0 & 1 & -2 \end{pmatrix}$$

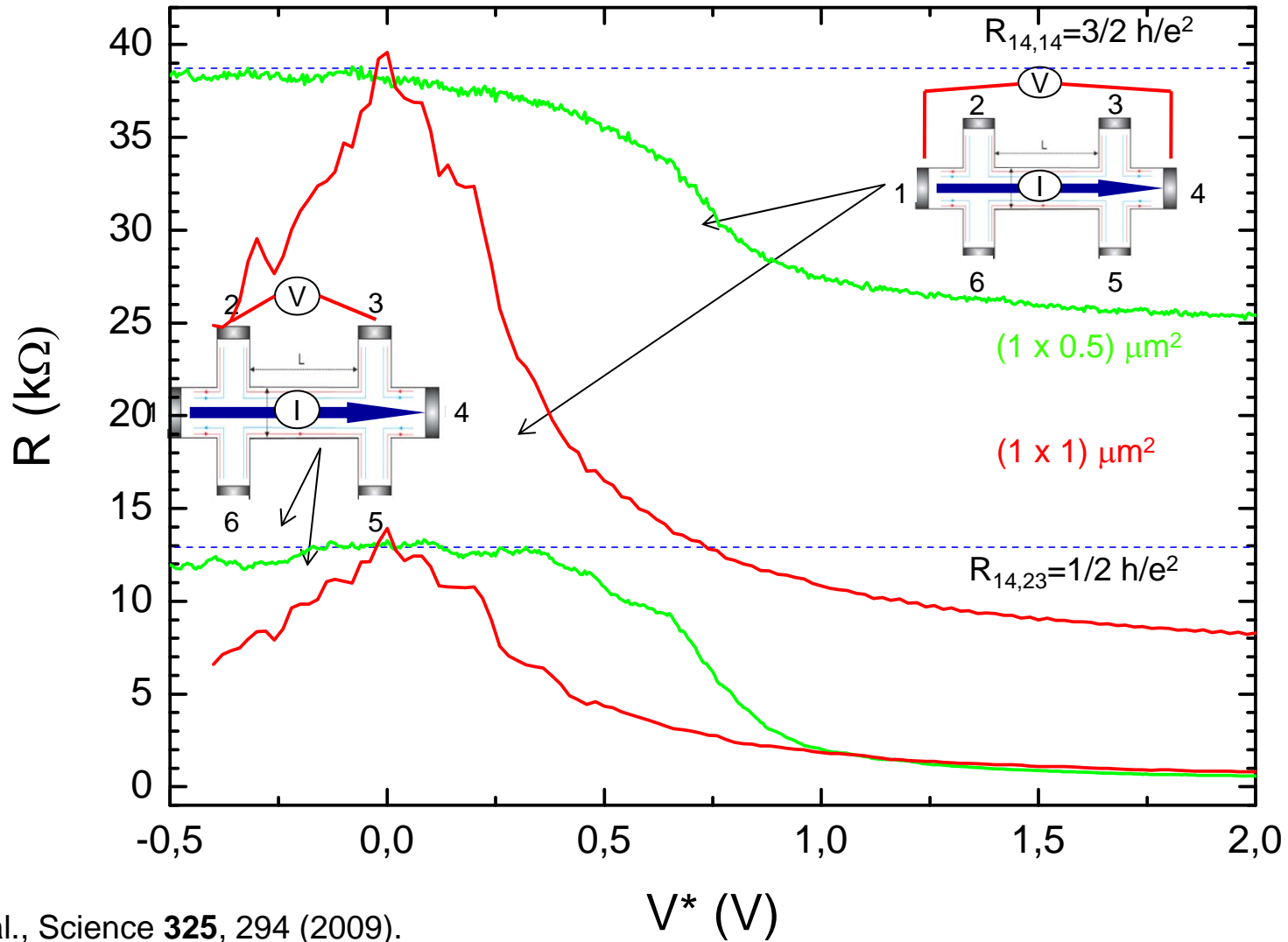
$$\Rightarrow \begin{cases} G_{2t} = \frac{I_{14}}{\mu_4 - \mu_1} = \frac{2e^2}{3h} \\ G_{4t} = \frac{I_{14}}{\mu_3 - \mu_2} = \frac{2e^2}{h} \end{cases}$$

generally

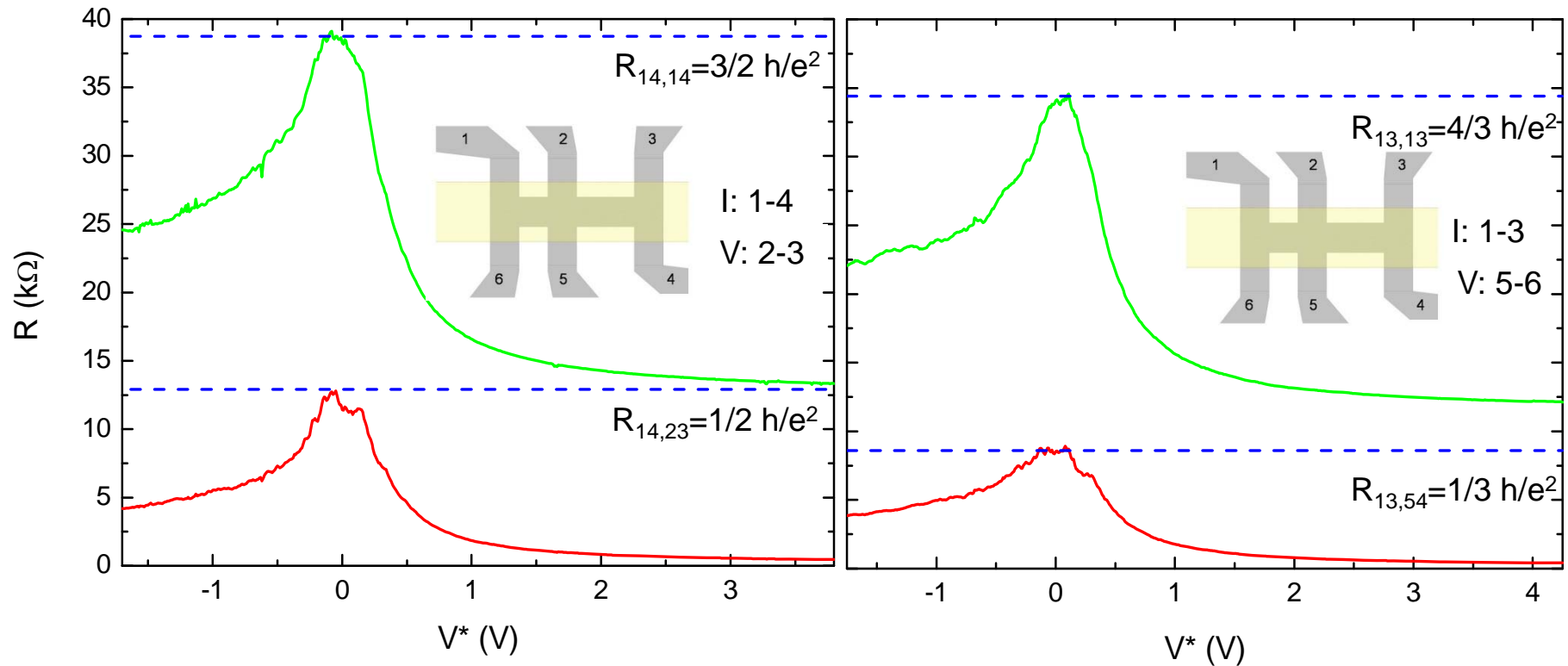
$$R_{2t} = \frac{(n+1)h}{2e^2}$$

$$G_{4t, \text{exp}} \approx 2 \frac{e^2}{h}$$

$$\left. \frac{R_{2t}}{R_{4t}} \right|_{\text{exp}} \approx 3$$

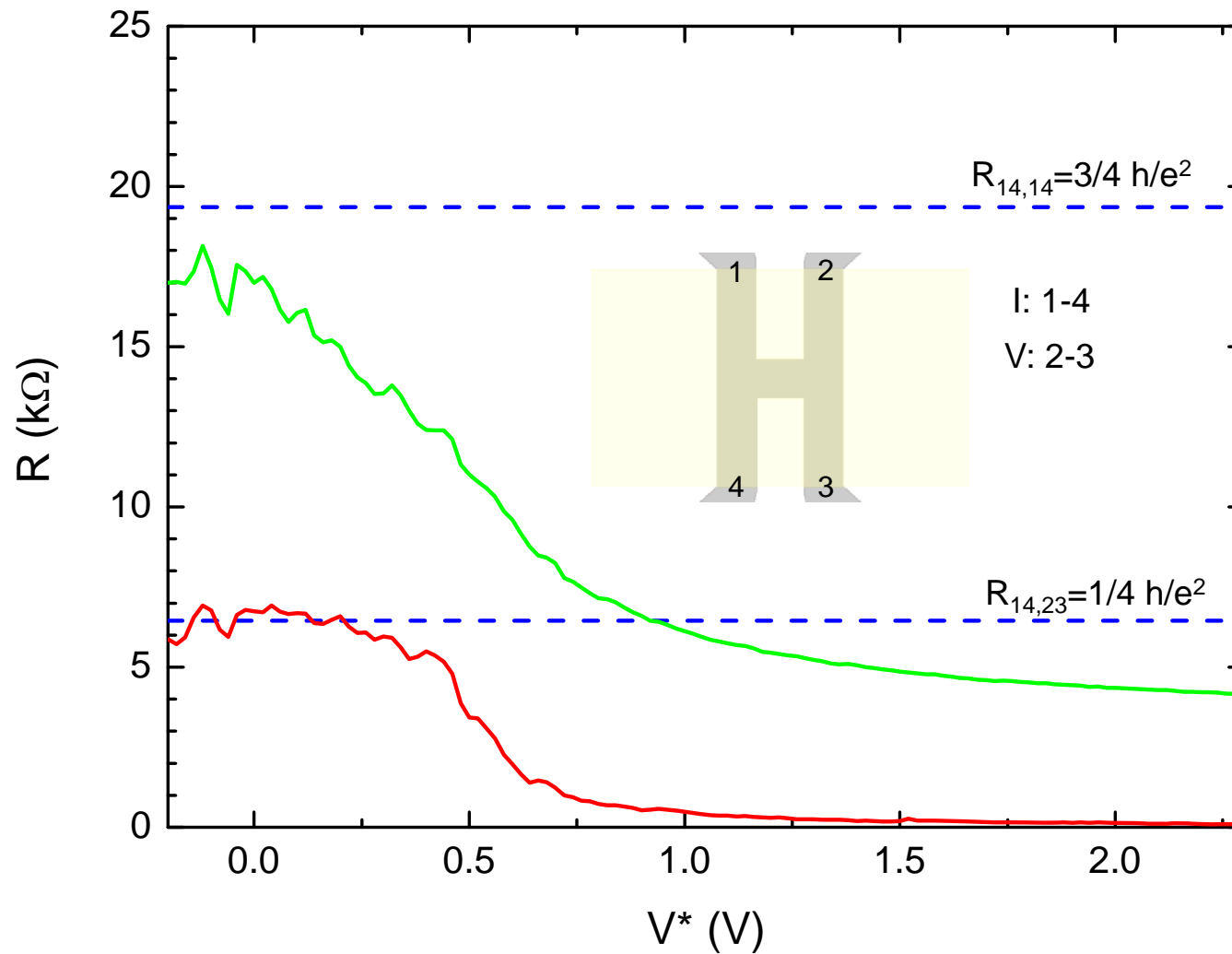


A. Roth et al., Science **325**, 294 (2009).



Configurations would be equivalent in quantum adiabatic regime

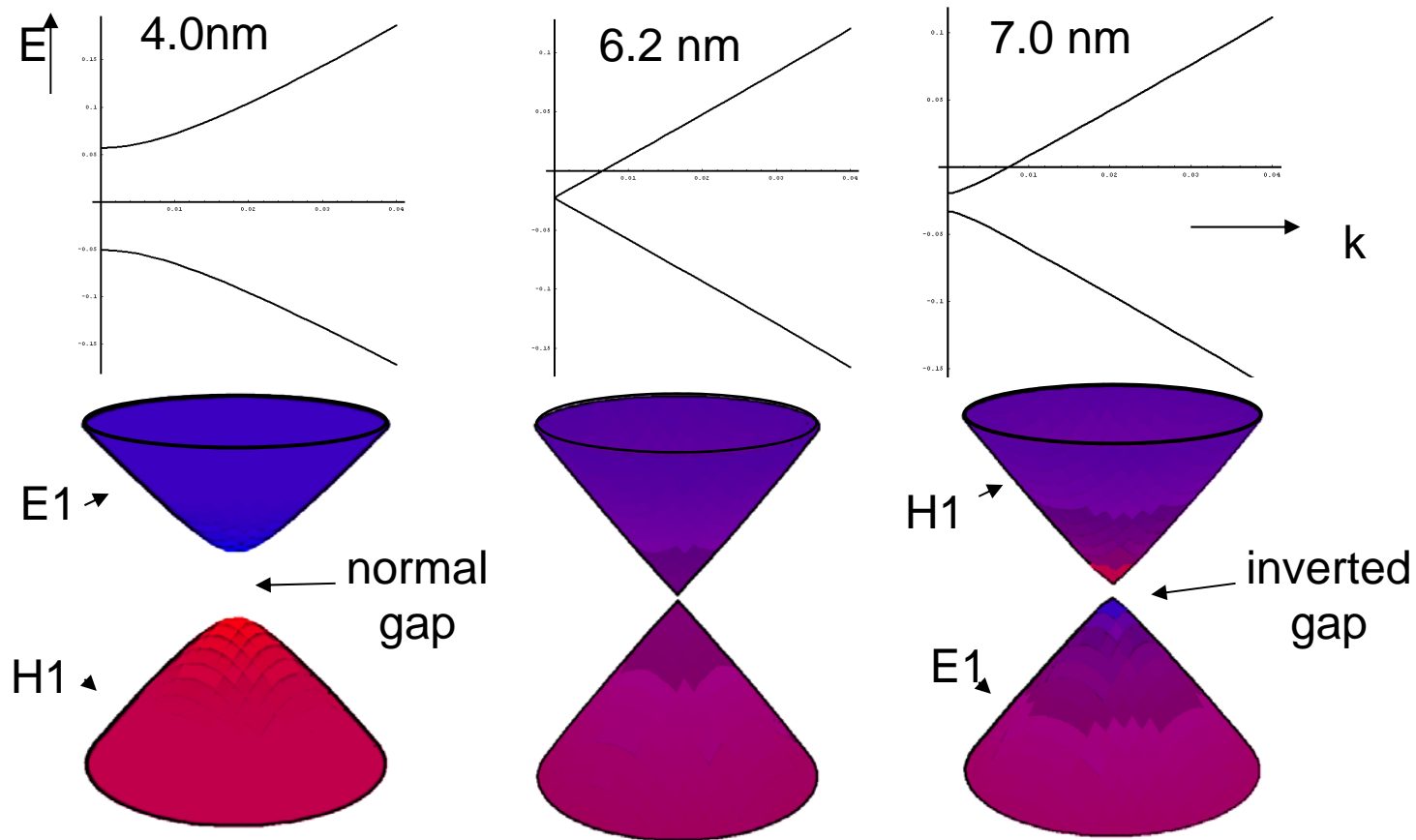


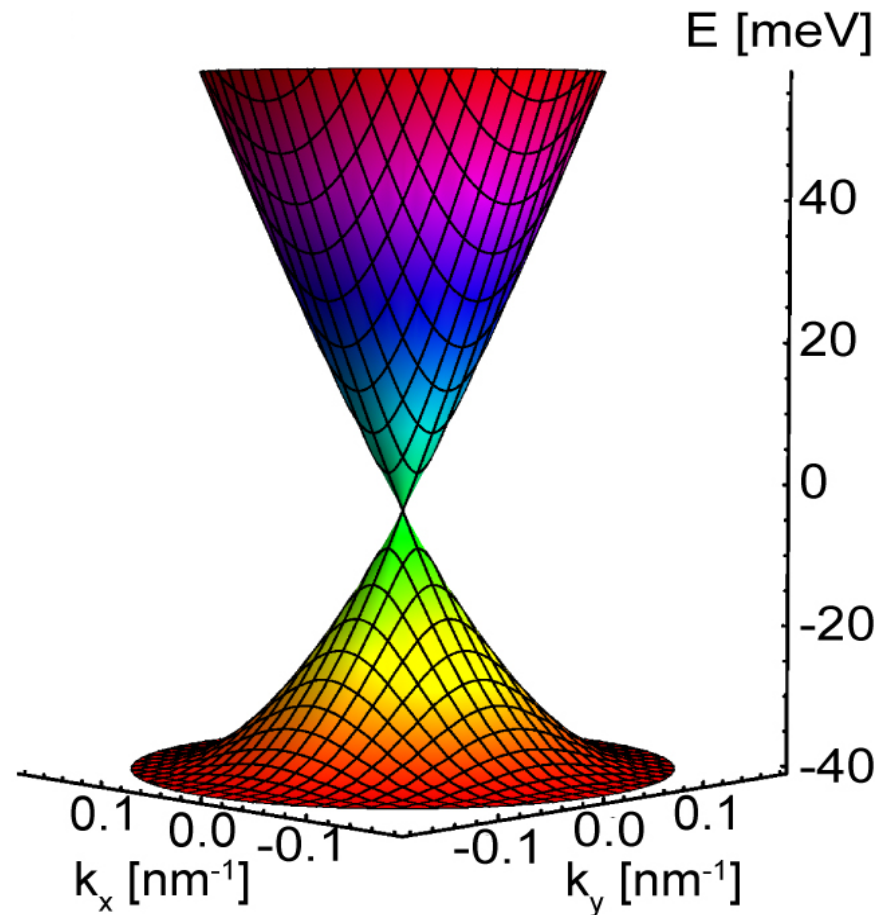
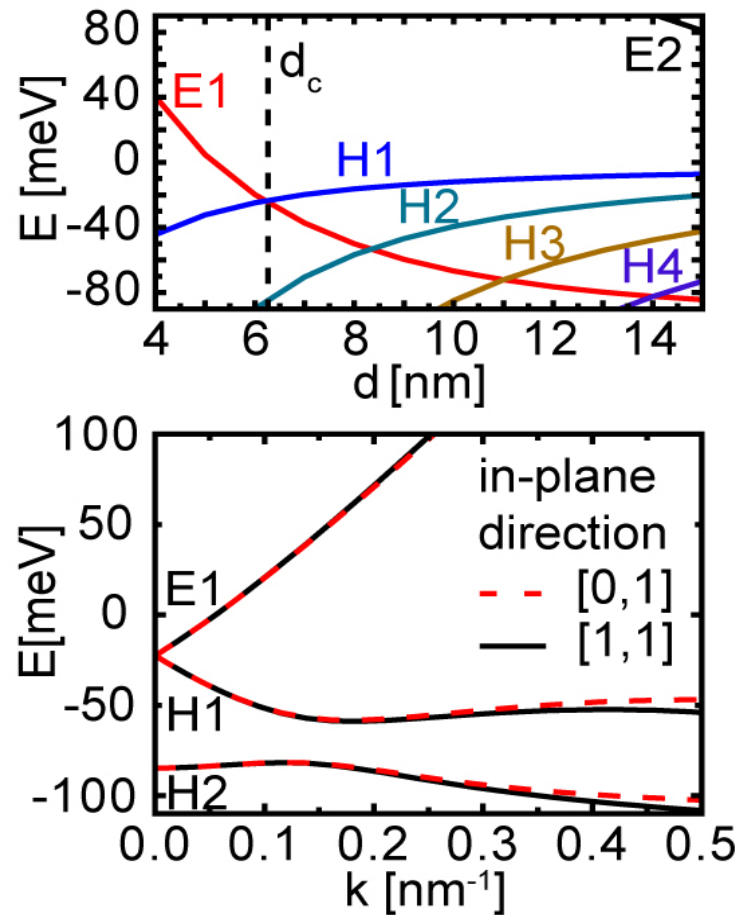


A. Roth et al., Science **325**, 294 (2009).

# Zero gap HgTe well as a Dirac system

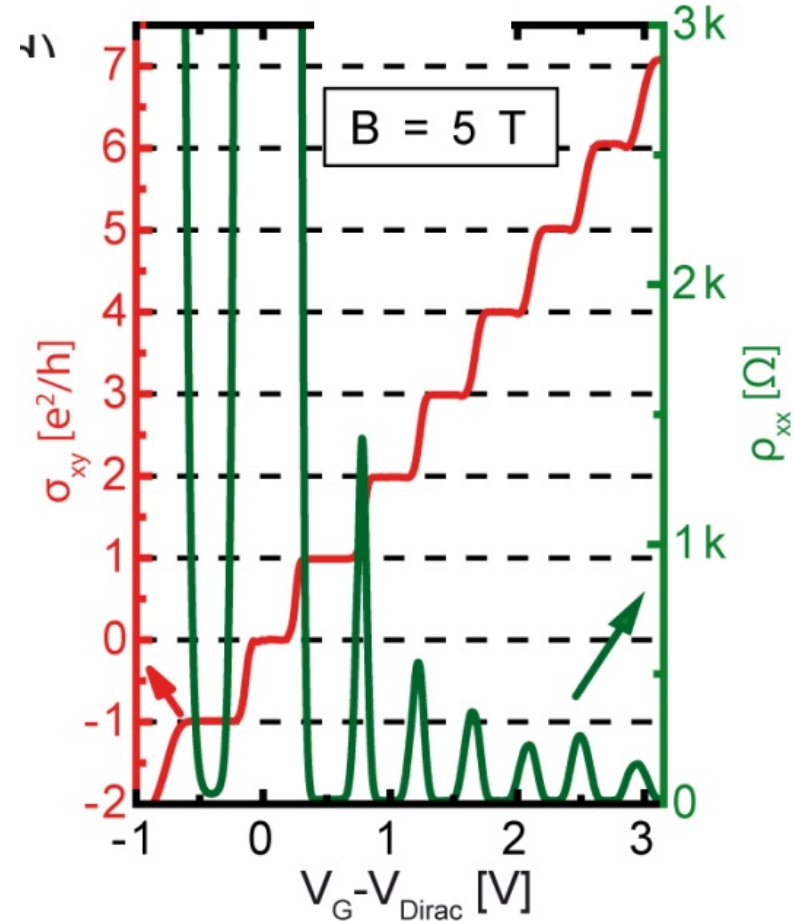
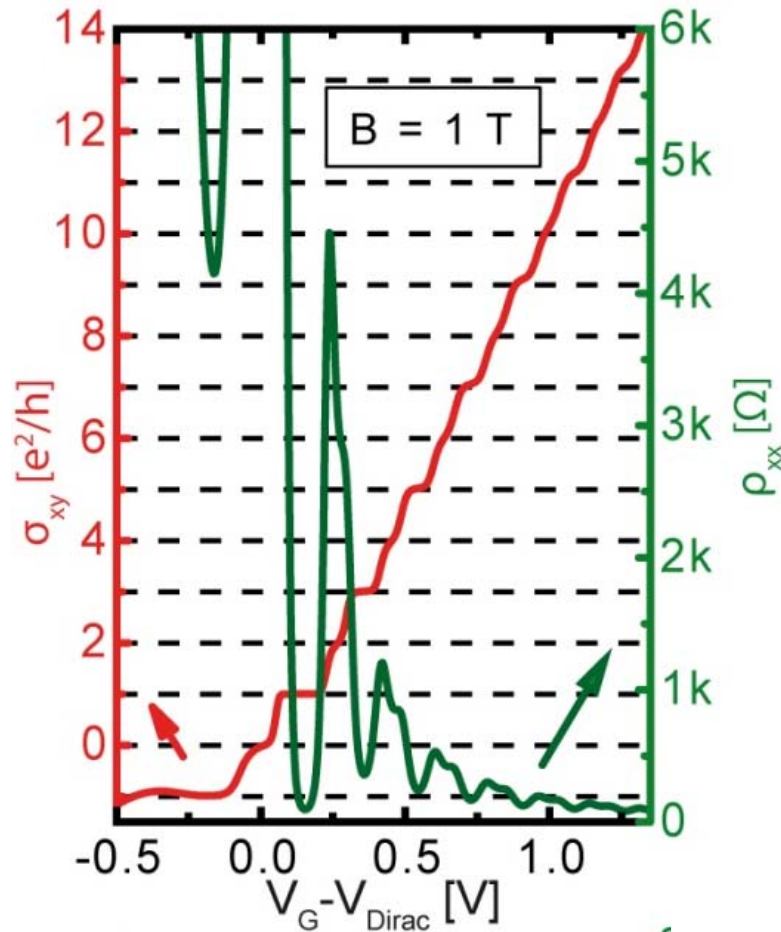
B.A Bernevig, T.L. Hughes, S.C. Zhang, Science **314**, 1757 (2006)





For well thickness  $d=6.3$  nm, the gap closes,  
especially the conductionband shows a linear dispersion: **single** Dirac cone

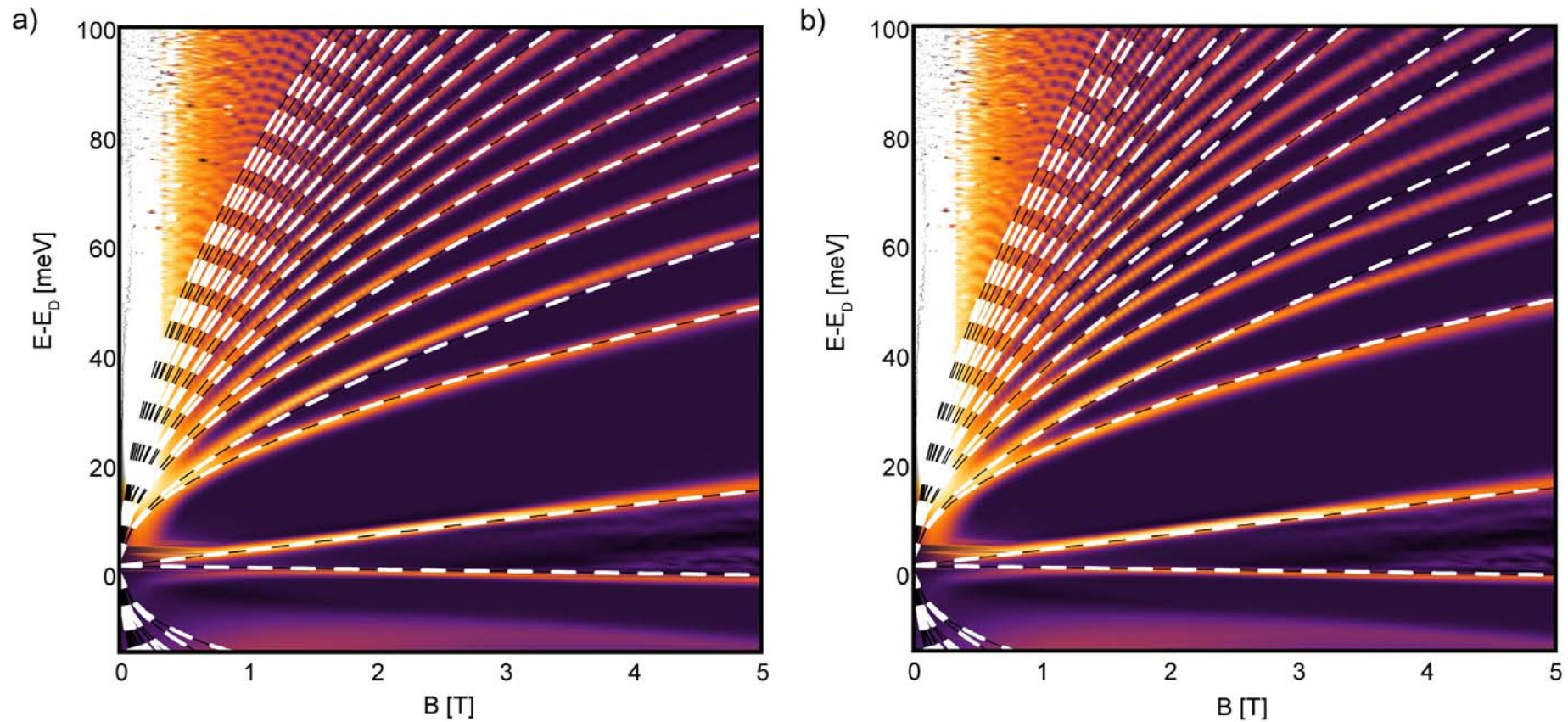
# Quantum Hall effect shows Berry phase



B. Büttner et al., Nature Physics 7, 418 (2011).

Large g-factor ( $g=55$ ) responsible for spin splitting already at low fields. Hall quantization reflects single valley character of the band structure: a HgTe quantum well at  $d=6.3$  nm is half-graphene.

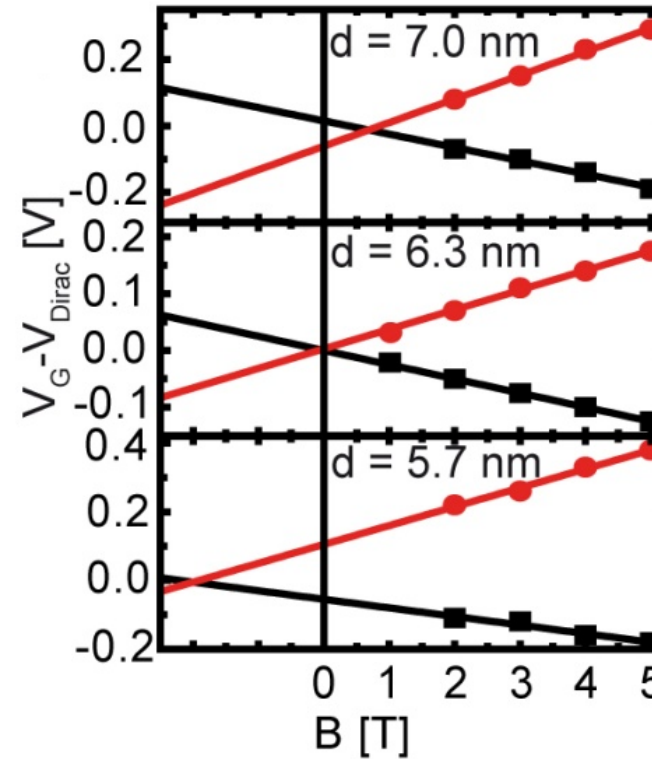
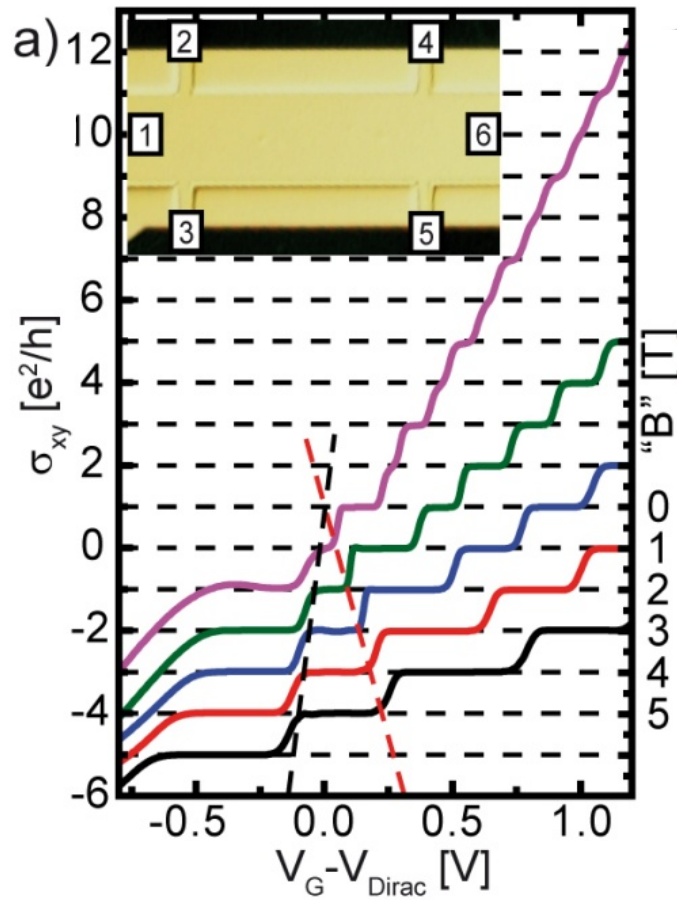




B. Büttner et al., Nature Physics **7**, 418 (2011).

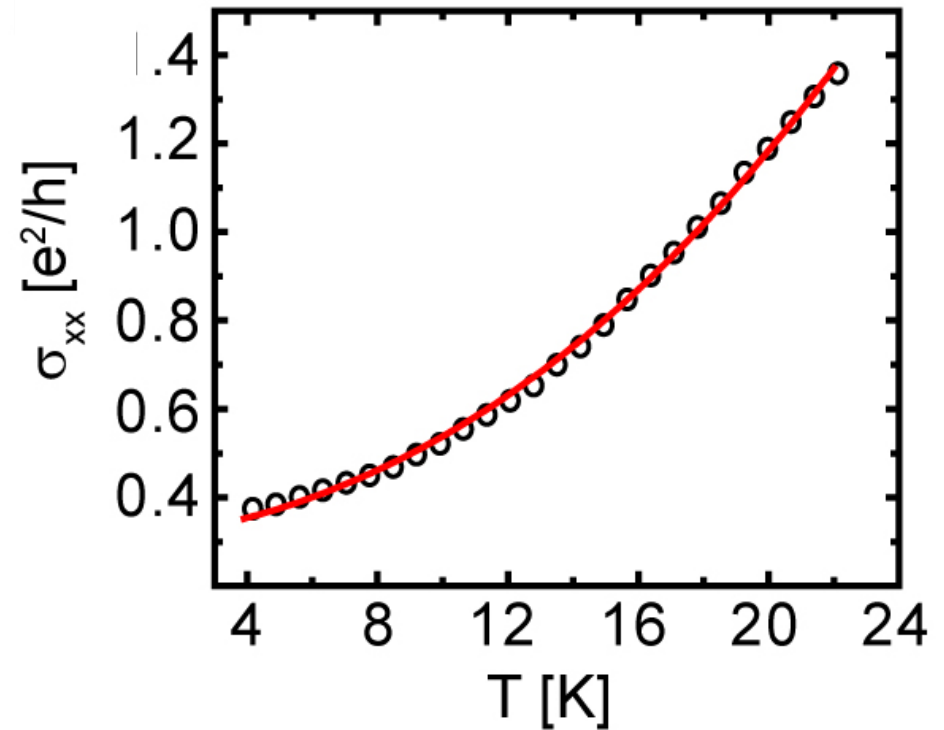
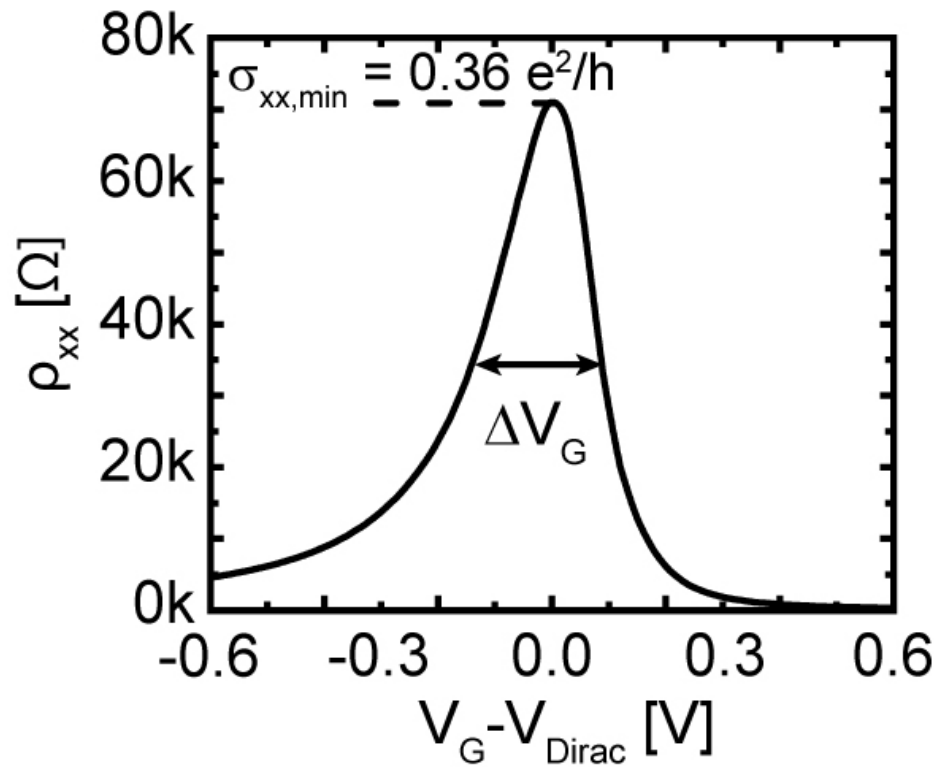
Color coded: gate voltage derivative of longitudinal resistivity.  
Fits: left – 8-band Kane model, right – Dirac Hamiltonian

# Zero mode dispersion



B. Büttner et al., Nature Physics 7, 418 (2011).

Zero mode spin splitting allows to select sample at  $d_c$ .

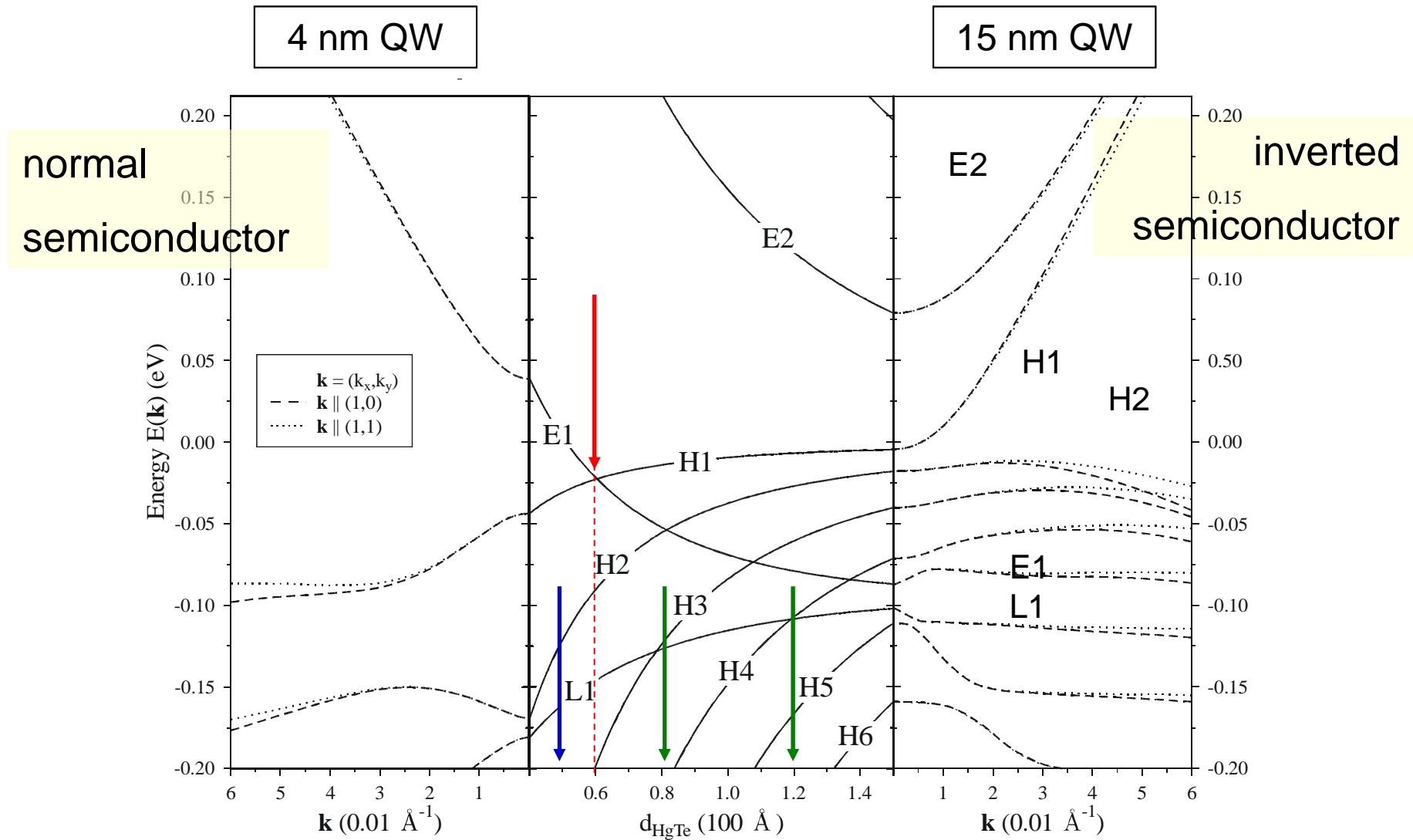


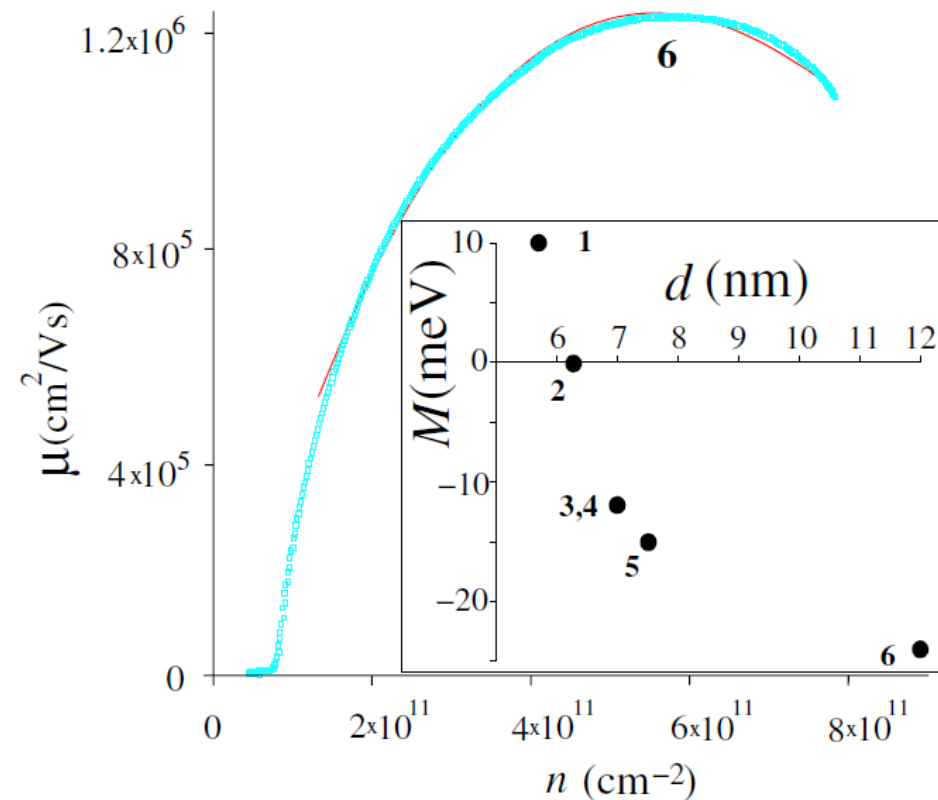
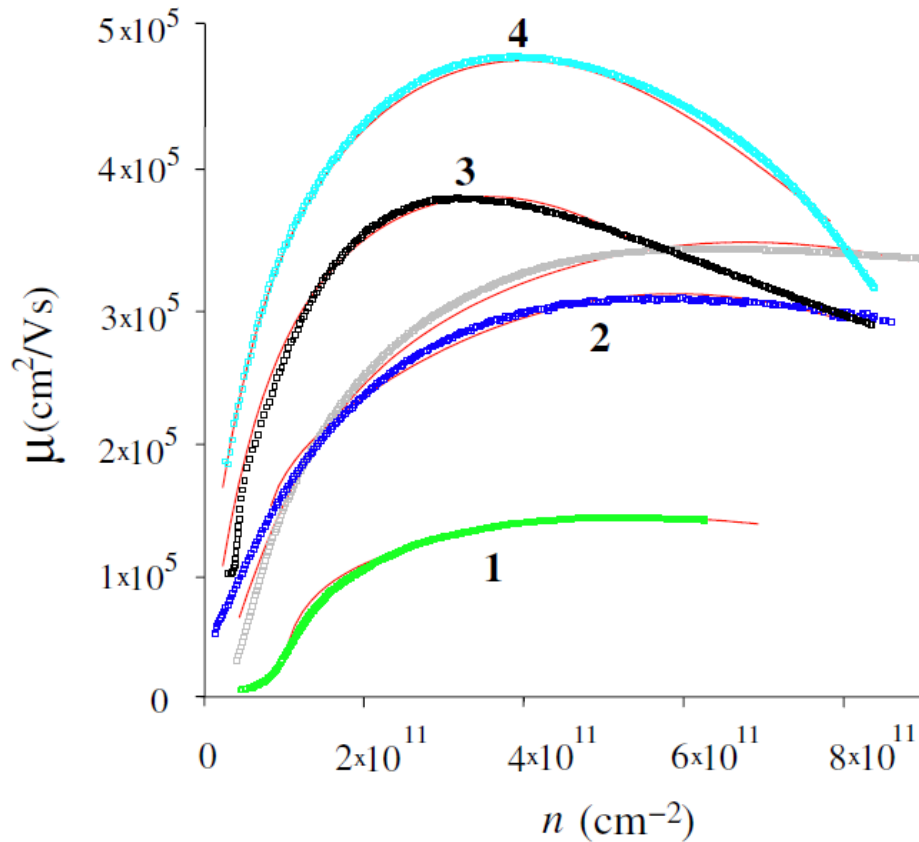
Peak width and mobilities comparable with/better than free standing graphene  
 Scattering mechanisms: probably mass fluctuations + Coulomb (fit is Kubo model)



# Adding a Dirac mass

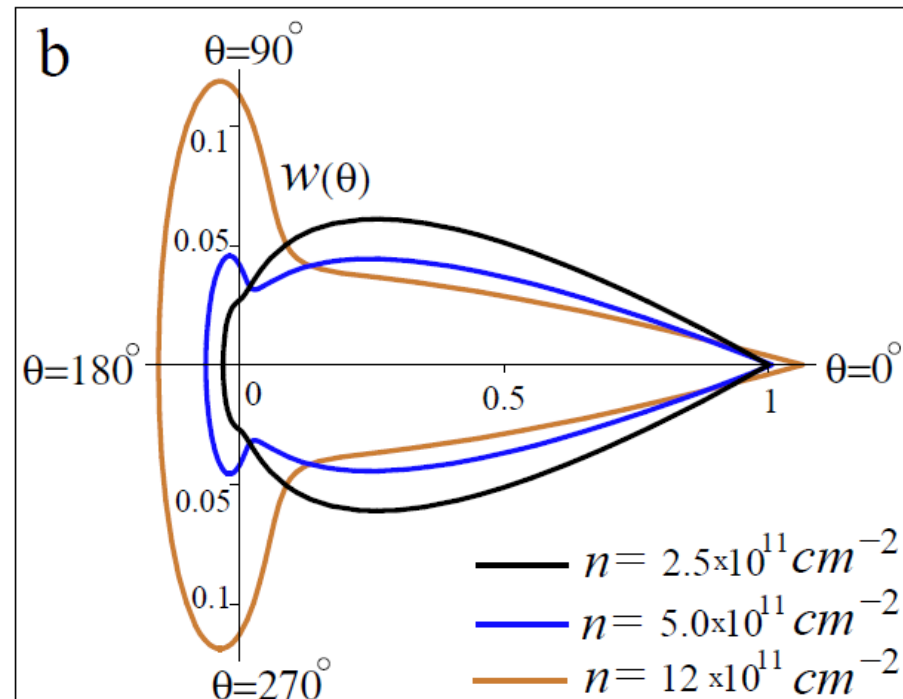
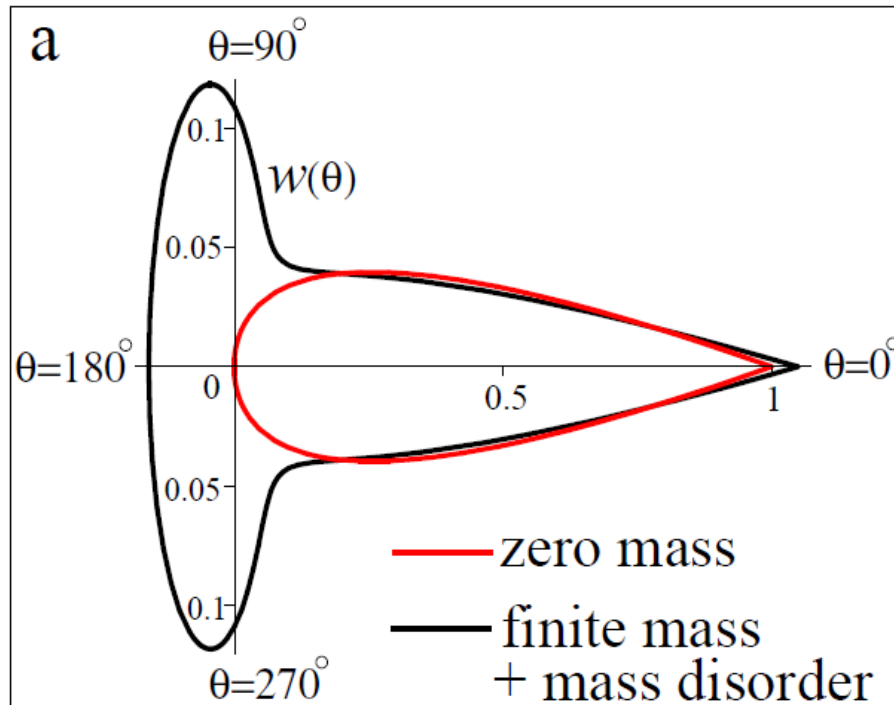
# Changing well width changes Dirac mass





B. Büttner et al., Phys. Rev. Lett. **106**, 076802 (2011).

Originally increase in mobility from reduced impurity scattering,  
then changeover to behavior due to well width (Dirac mass) fluctuations.

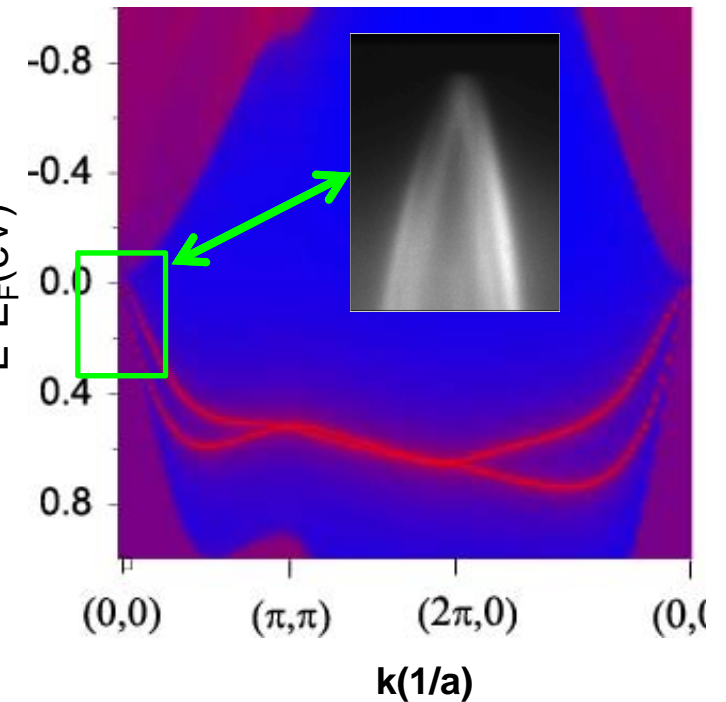
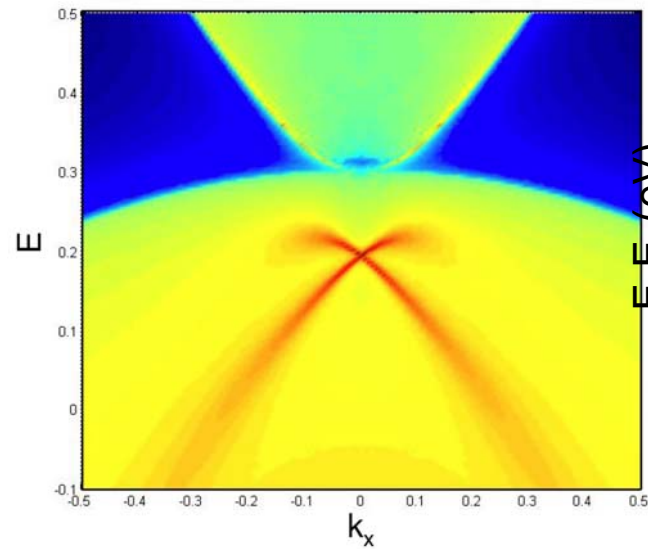
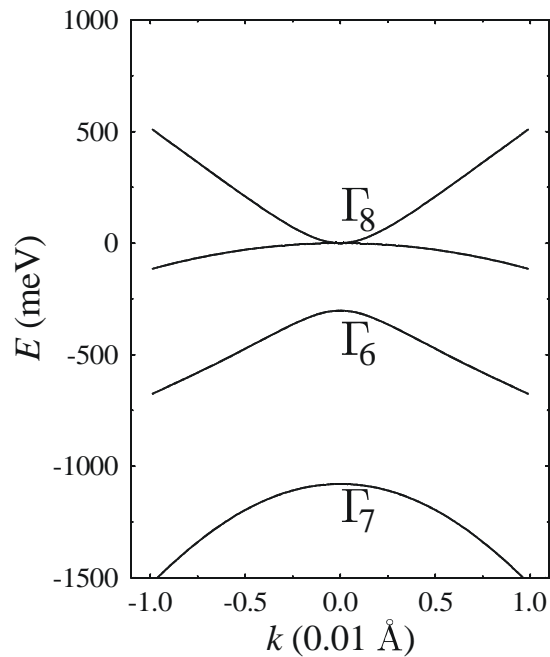


B. Büttner et al., Phys. Rev. Lett. **106**, 076802 (2011).

Modeling by Grigory Tkachov and Ewelina Hankiewicz:  
Mass and disorder induce backscattering of Dirac fermions.

# **Dirac Surface States on strained bulk HgTe**

# Bulk HgTe as a 3-D Topological Insulator'

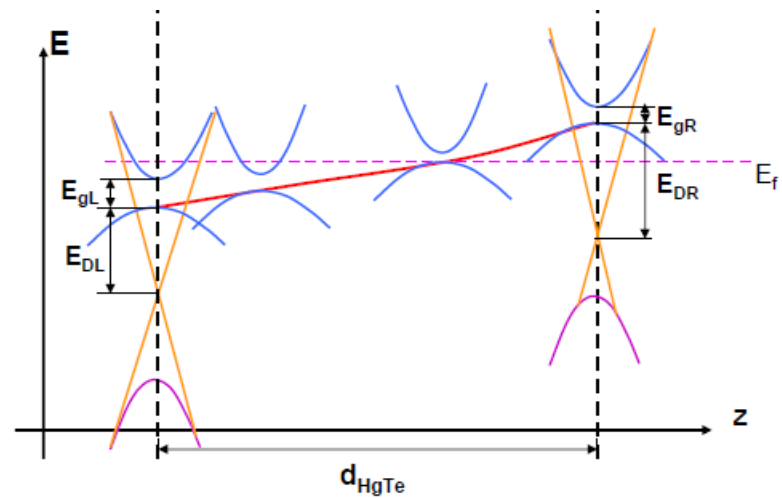
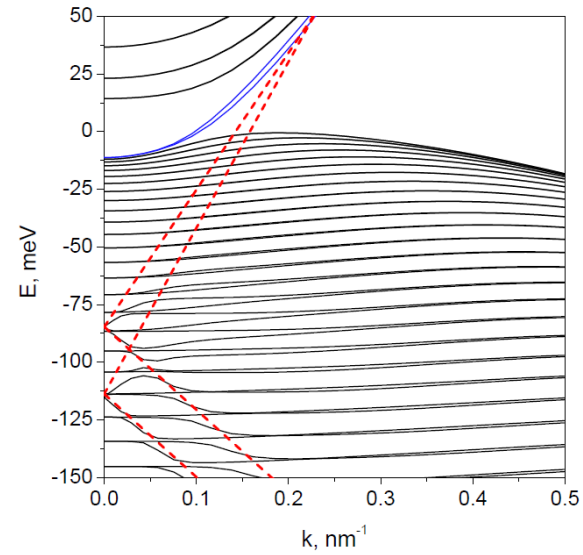
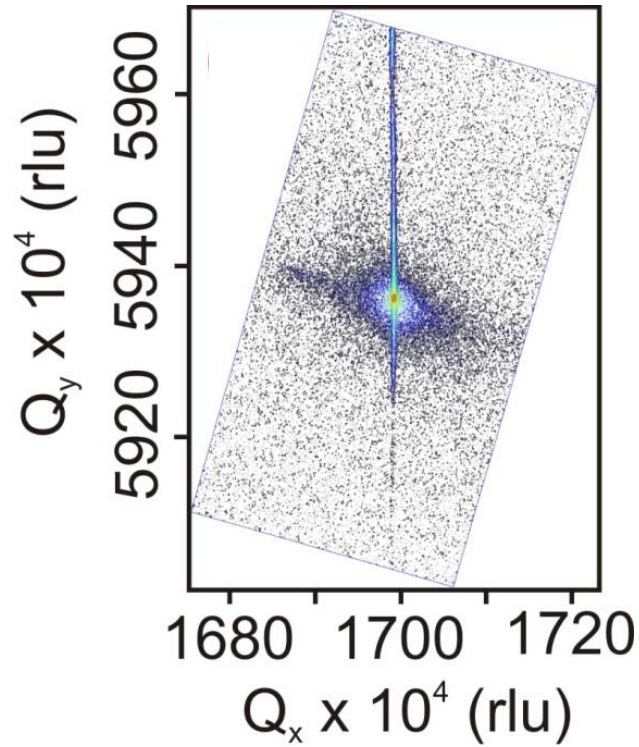


Bulk HgTe is semimetal,  
topological surface state overlaps w/ valenceband.

C. Brüne et al., Phys. Rev. Lett. **106**, 126803 (2011).

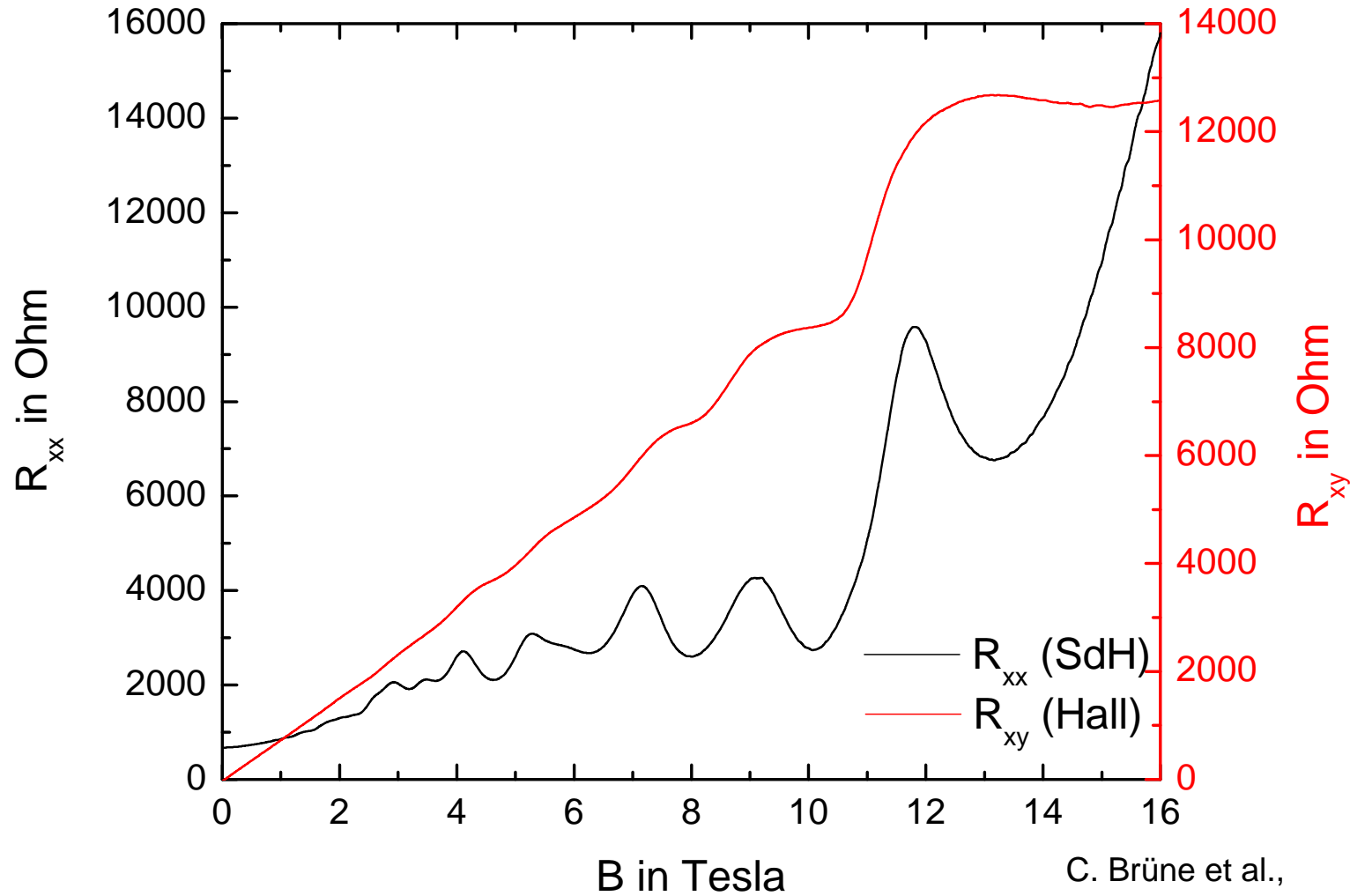
ARPES:  
Yulin Chen, ZX Shen,  
Stanford

# 70 nm layer on CdTe substrate: coherent strain opens gap





# Bulk HgTe as a 3-D Topological 'Insulator'

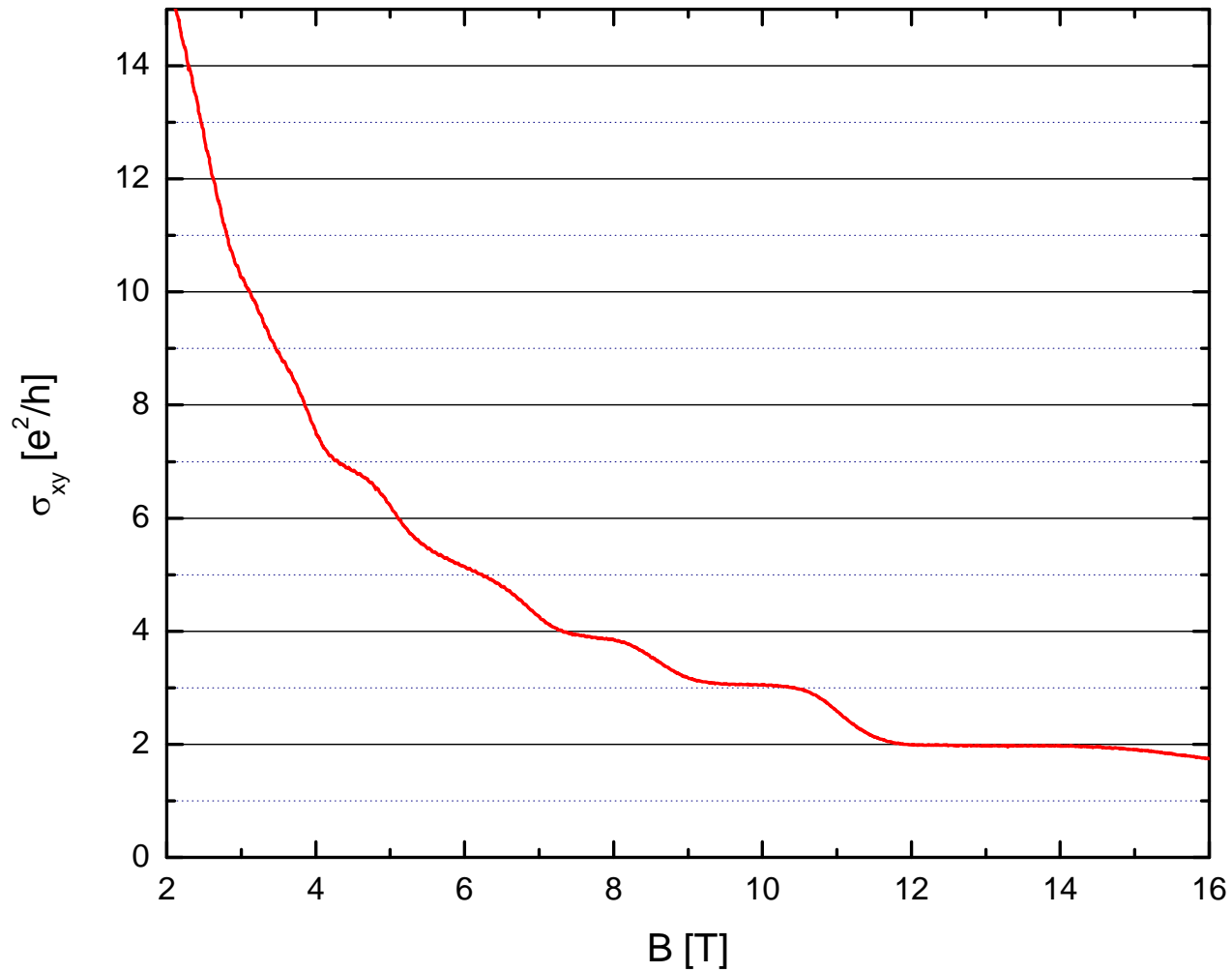


C. Brüne et al.,  
Phys. Rev. Lett. **106**, 126803 (2011).

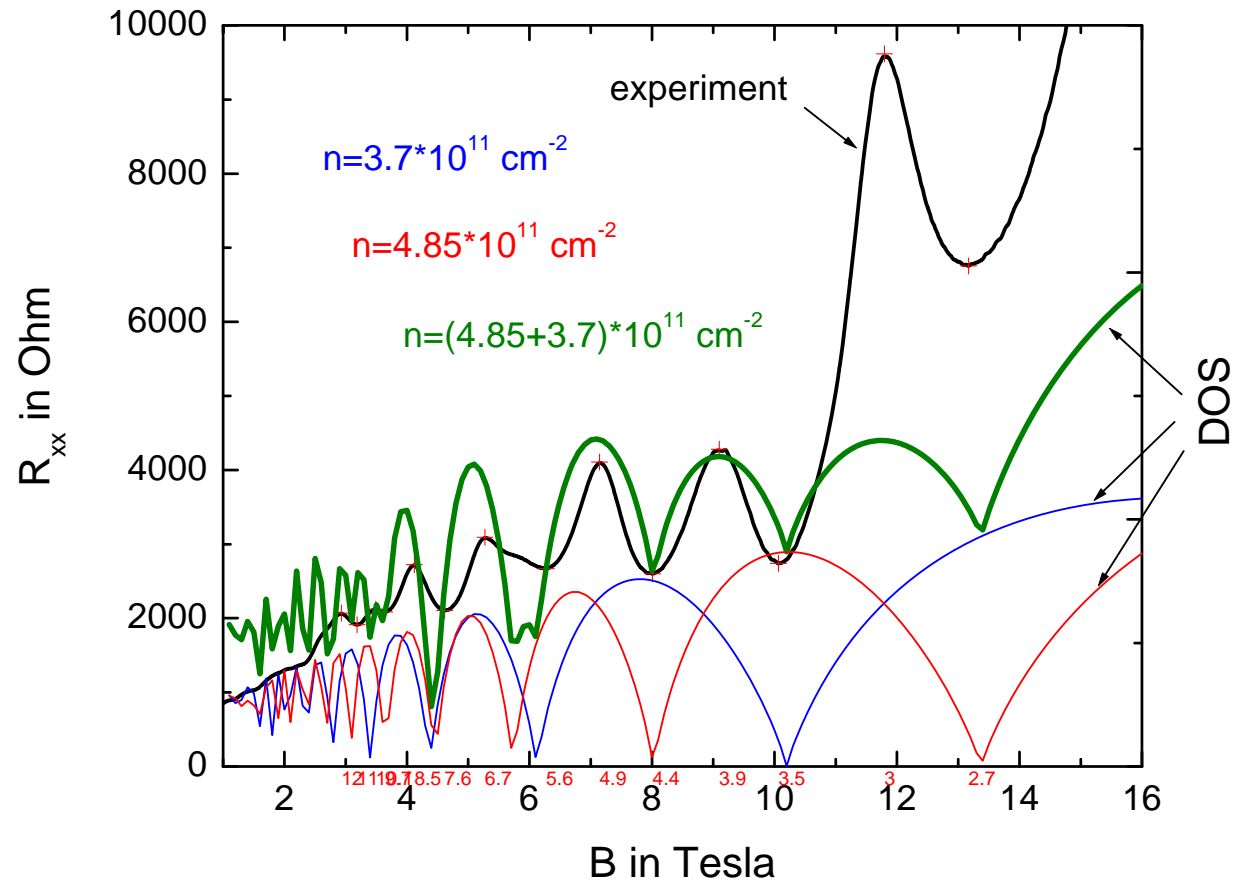
@ 20 mK: bulk conductivity almost frozen out - Surface state mobility ca. 35000 cm<sup>2</sup>/Vs



# Bulk HgTe as a 3-D Topological 'Insulator'



@ 20 mK: same data, plotted as conductivity



C. Brüne et al., Phys. Rev. Lett. **106**, 126803 (2011).

Red and blue lines : DOS for each of the Dirac-cones with the corresponding fixed 2D-density,  
Green line: the sum of the blue and red lines

- HgTe quantum wells: normal and inverted gap, linear (Dirac) dispersion
- First observation of Quantum Spin Hall Effect
- At  $d=dc$ , a HgTe QW is ideal model system for zero mass Dirac fermion physics
- Can conveniently study Dirac fermions w/ finite Dirac mass
- Strained 3D layers show QHE of topological surface states

## Collaborators:

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Theory: Alena Novik, Chaoxing Liu, Ewelina Hankiewicz, Grigory Tkachov, Patrick Recher, Björn Trauzettel (all @ Würzburg), Jairo Sinova (TAMU), Shoucheng Zhang, Xiaoliang Qi (Stanford)

Funding: DFG (SPP Spintronics, DFG-JST FG Topotronics), Humboldt Stiftung, EU-ERC AG “3-TOP”, DARPA