



Dynamic topological orders in periodically driven systems

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Ashvin Vishwanath
(Harvard)

Plan

New non-equilibrium dynamical phases from periodic driving

$$H(t + T) = H(t)$$

Floquet symmetry protected topological phases (SPTs)

- aka driven + interacting topological insulators

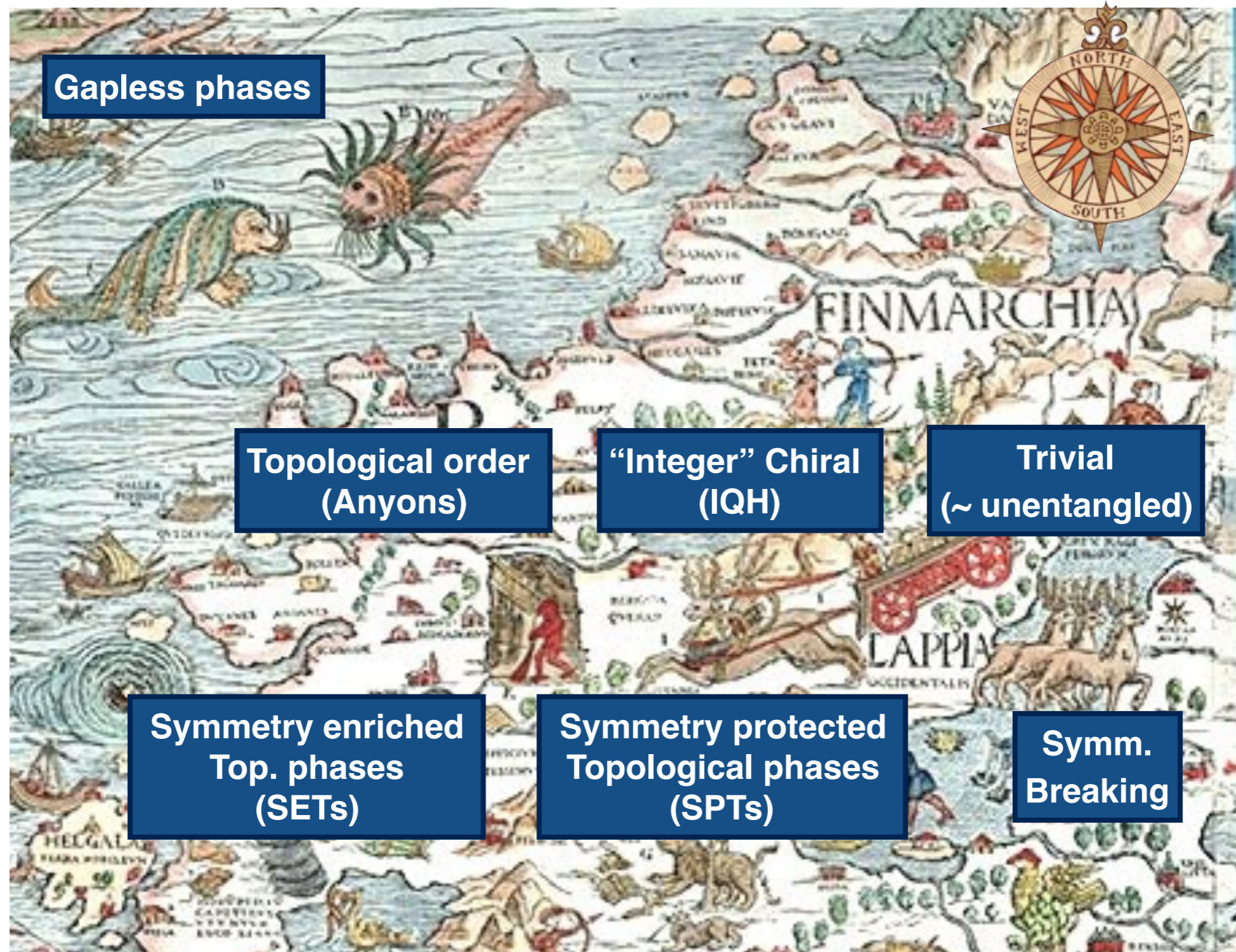
1. *ACP, T. Morimoto, A Vishwanath Phys. Rev. X* **6**, 041001 (2016)
2. *ACP. T Morimoto, arXiv:1610.03485*

Chiral Floquet phases

- Edge: one-way pumping of quantum information

3. *Po, Fidkowski, Morimoto, ACP, Vishwanath arXiv:1609.00006*

“World Map” of quantum matter (Equilibrium, $T=0$)



← (Ground-state) Entanglement

Isolated (“closed”) quantum many-body systems



Ultracold atoms

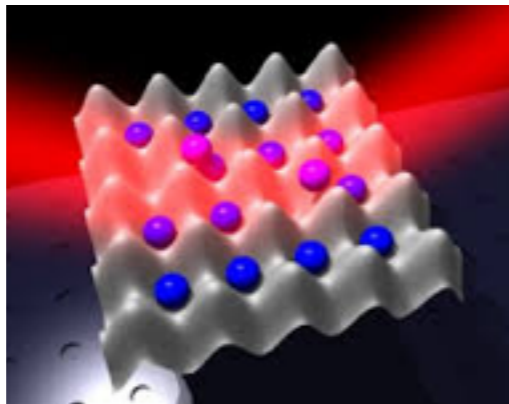


image: <http://www.lens.unifi.it>

Trapped Ions

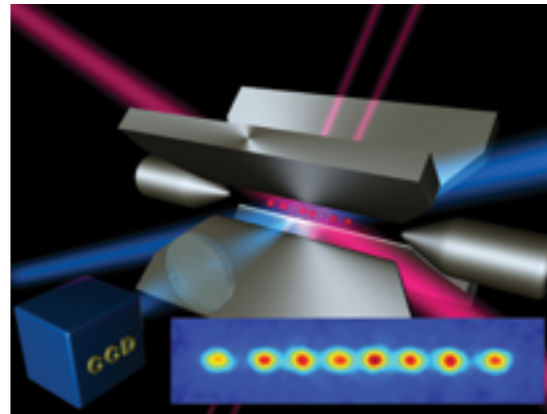


image: www.laserfocusworld.com

NV Centers

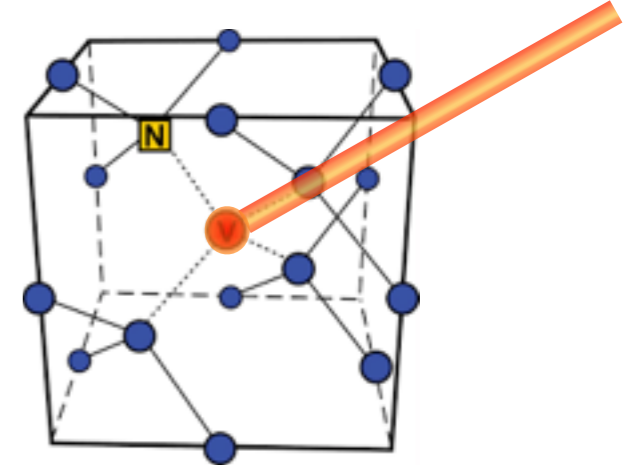
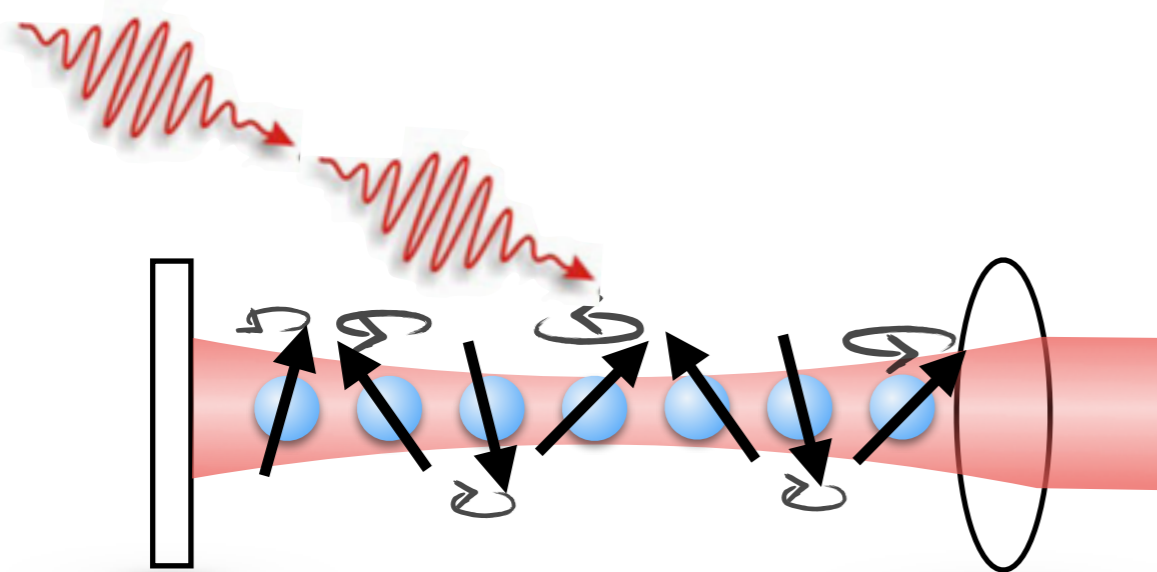


image: www.labnews.co.uk



$$|\Psi(0)\rangle = |\psi_1\rangle \otimes |\psi_2\rangle \otimes \cdots \otimes |\psi_L\rangle = \sum_n c_n |n\rangle$$

$$|\Psi(t)\rangle = e^{-iHt} |\Psi(0)\rangle$$

$$H(t+T) = H(t)$$

New universal quantum phenomena in dynamics?

Periodic Driving 1: Engineer new interactions

$$H(t + T) = H(t) \quad \rightarrow \quad H_{\text{eff}} \approx H_0 + \frac{iT}{2} \int_0^T dt_1 \int_0^{t_1} dt_2 [H(t_1), H(t_2)] + \dots$$

$$U(T) = e^{-iH_{\text{eff}}T}$$

Lindner, Refael, Galitski Nat. Phys. '11
Many others...

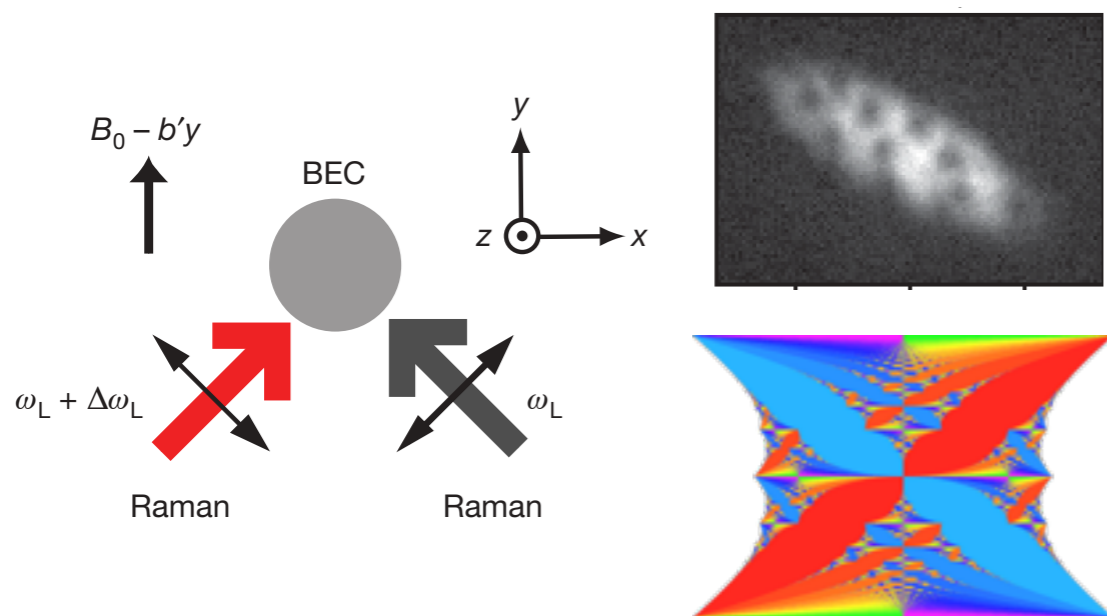
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Cold atoms



Synthetic Gauge Fields

Lin et al. (Speilman group) Nature '09;
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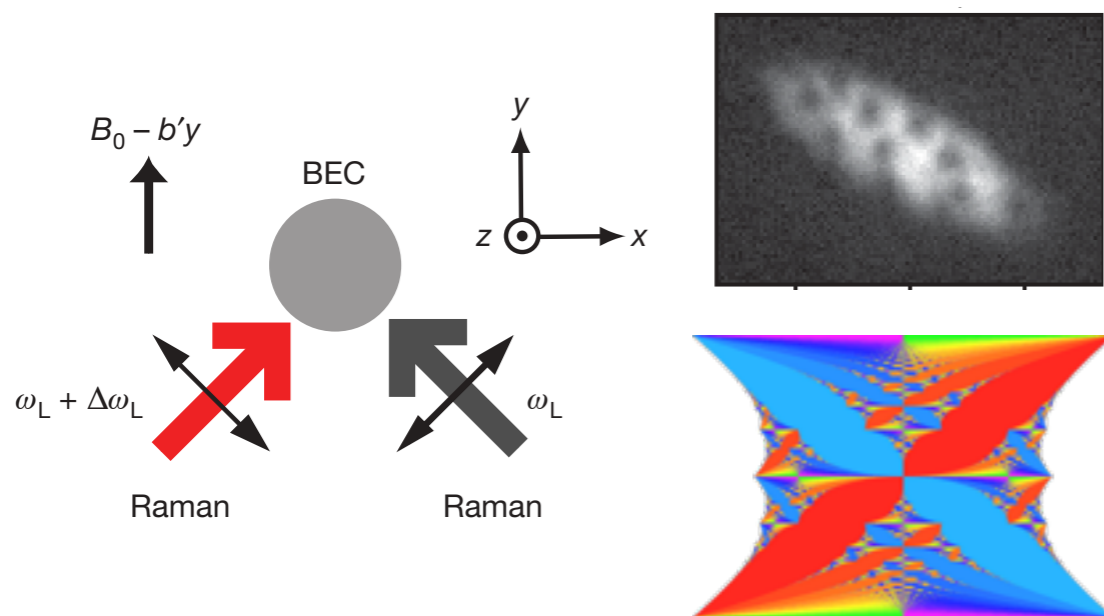
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Solids

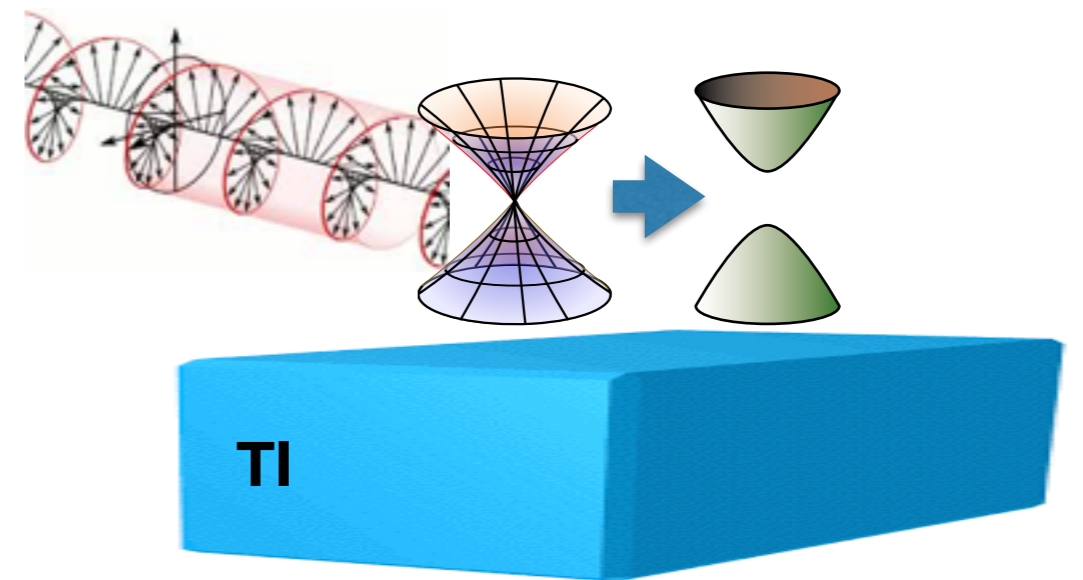
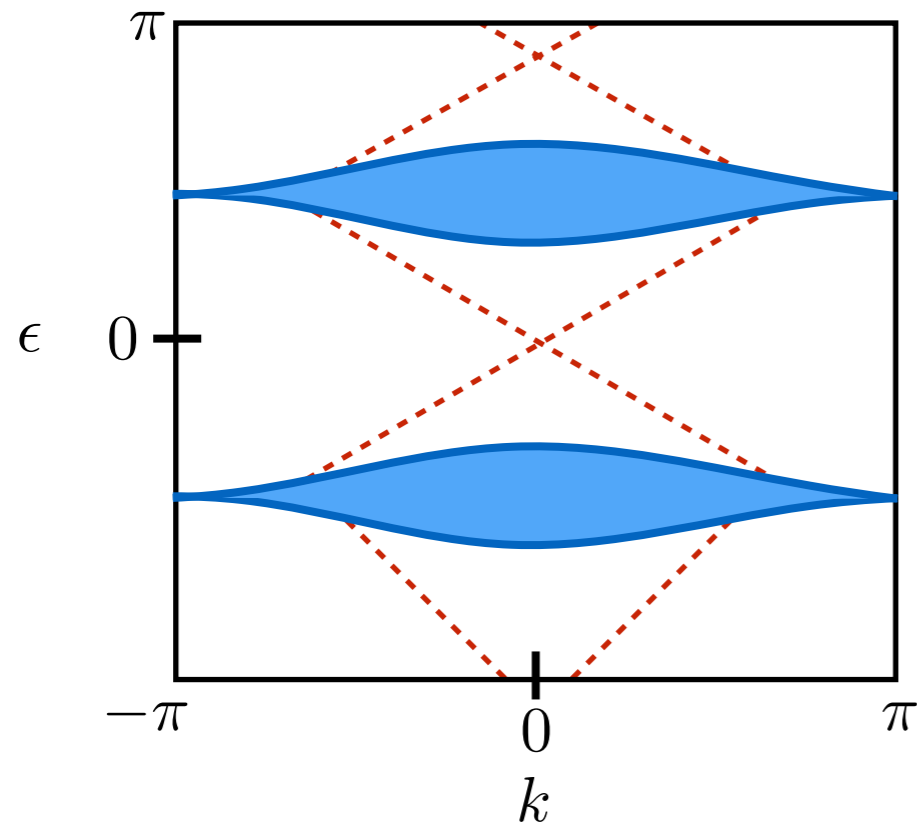


Photo-induced Anomalous Quantum Hall state in TI surface

Wang et al. (Gedik Group) Science '13

Periodic Driving 2: Fundamentally new dynamical phases



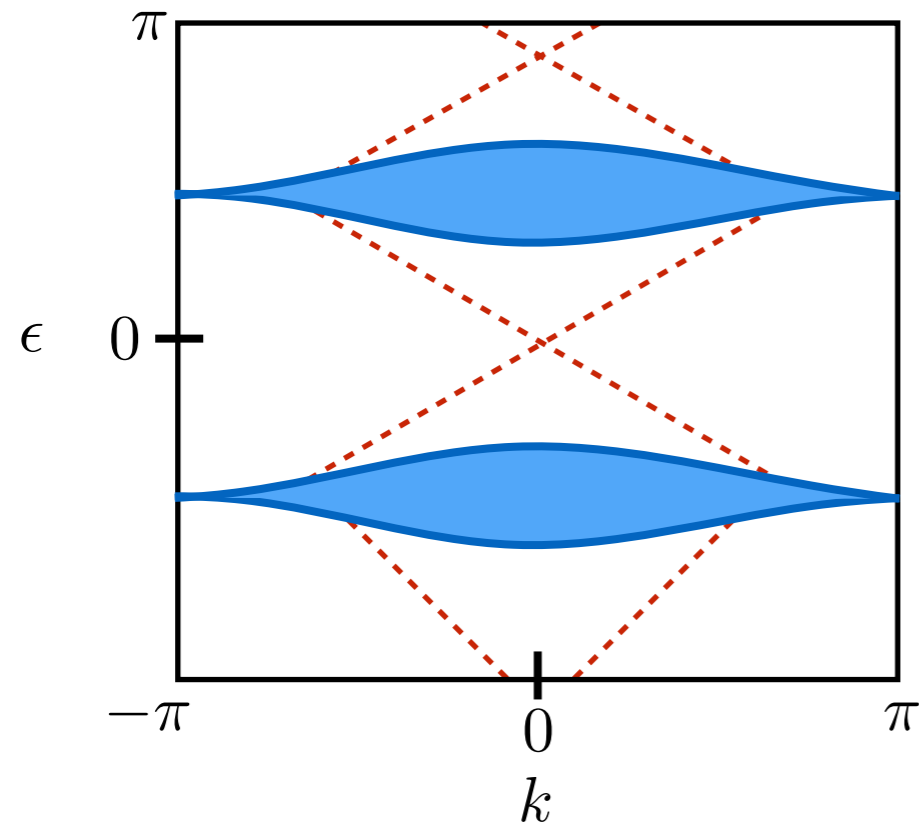
New types of band topology due to periodicity of quasi-energy

$$U(T) = \mathcal{T}\{e^{-\int_0^T dt H(t)}\} \quad \epsilon_n \simeq \epsilon_n + \frac{2\pi}{T}\mathbb{Z}$$

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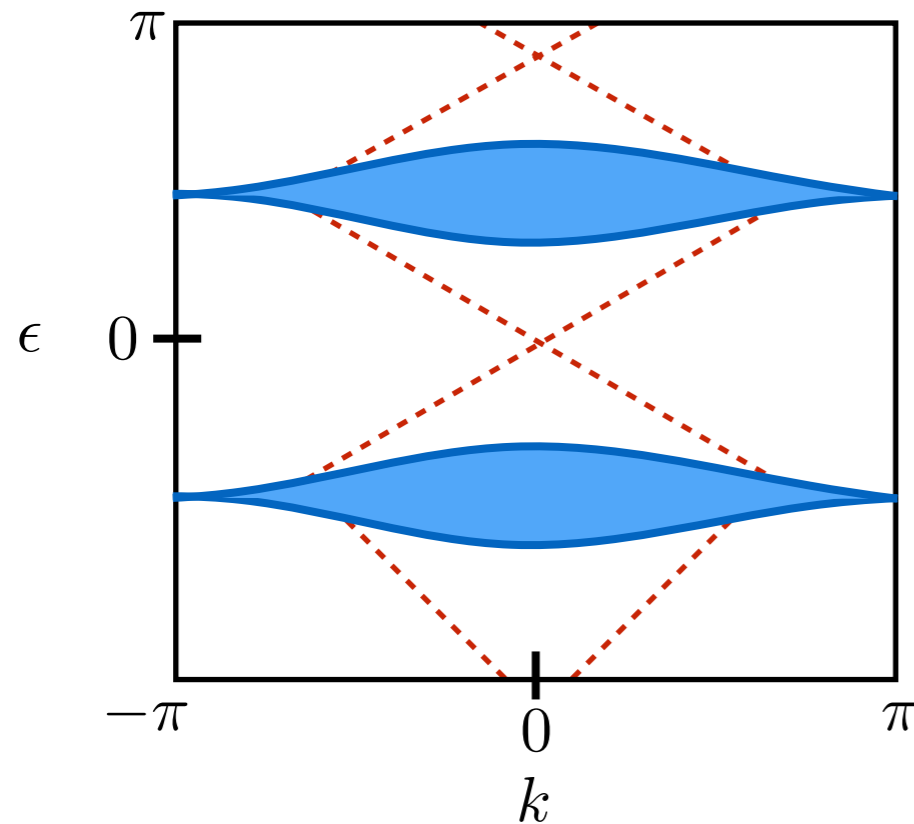
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Why do we pay attention to non-interacting band topology?

(all systems are interacting, not perfectly clean, etc...)

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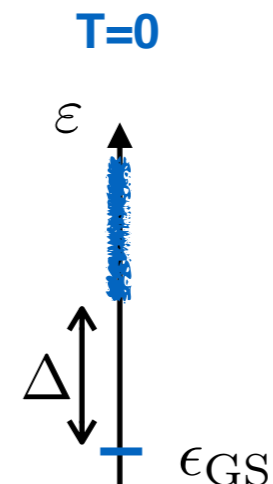
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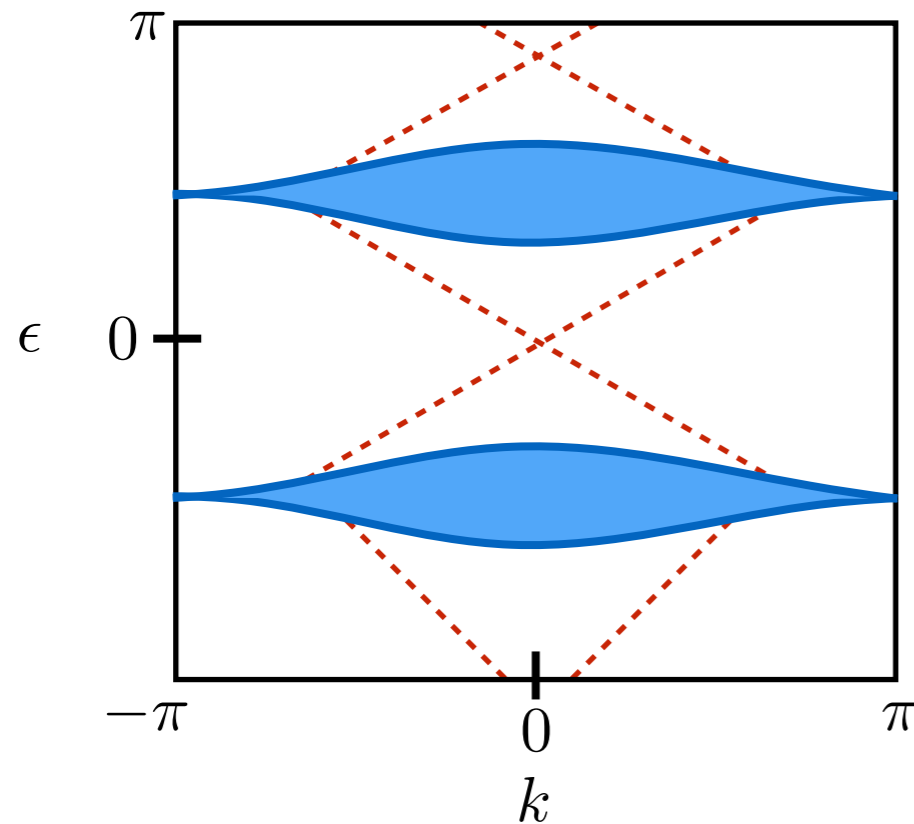
Usual (equilibrium, $T \rightarrow 0$):

- ground-state + gap \Rightarrow physics dominated by filled bands
- weak interactions/disorder don't close the gap (band topology inherited by many-body system)



Many-body spectrum

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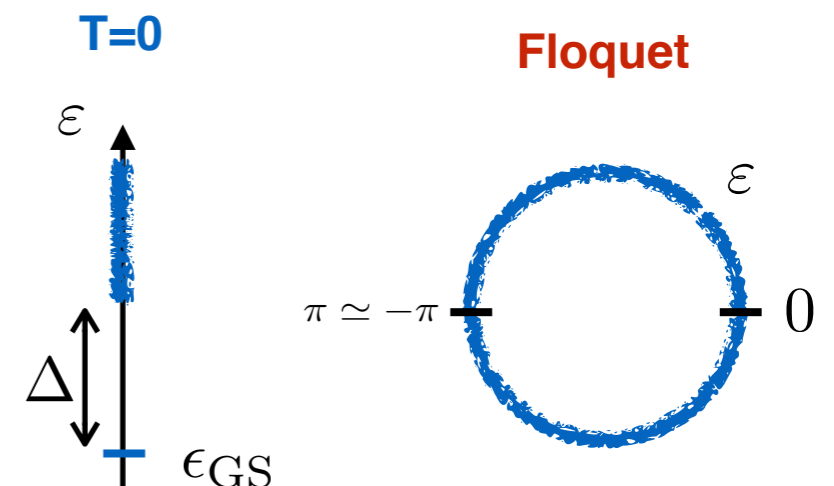
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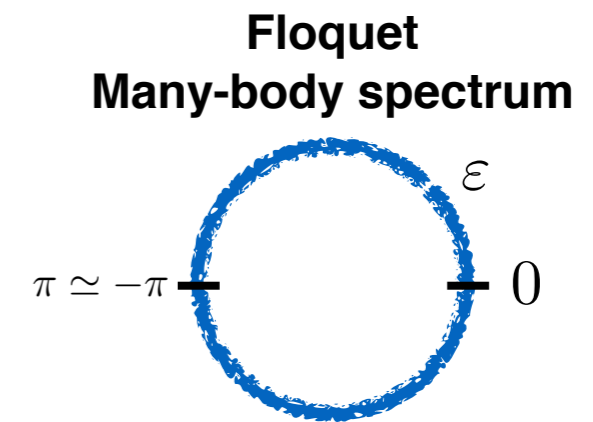
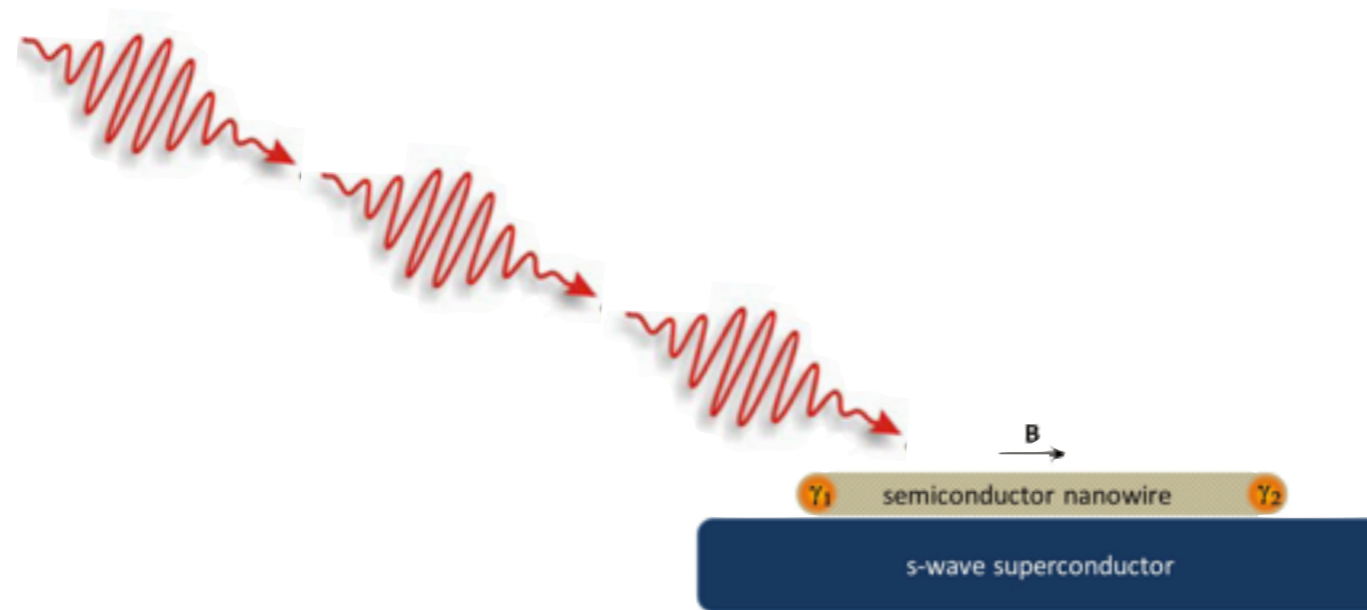
Floquet:

- no “lowest” energy state,
- What single-particle levels are “filled”?
(is there a sense of equilibrium?)

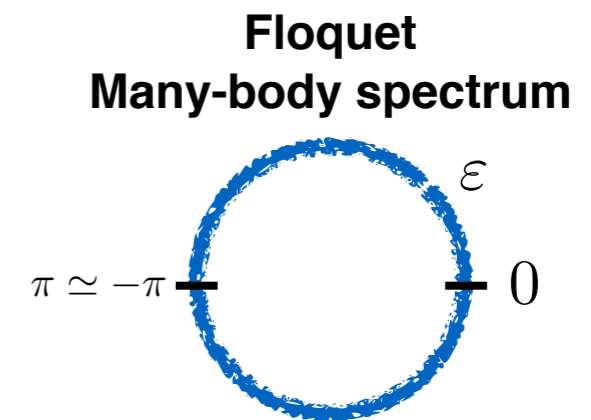
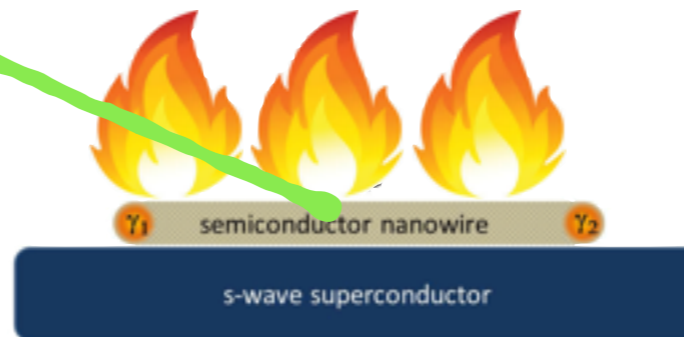
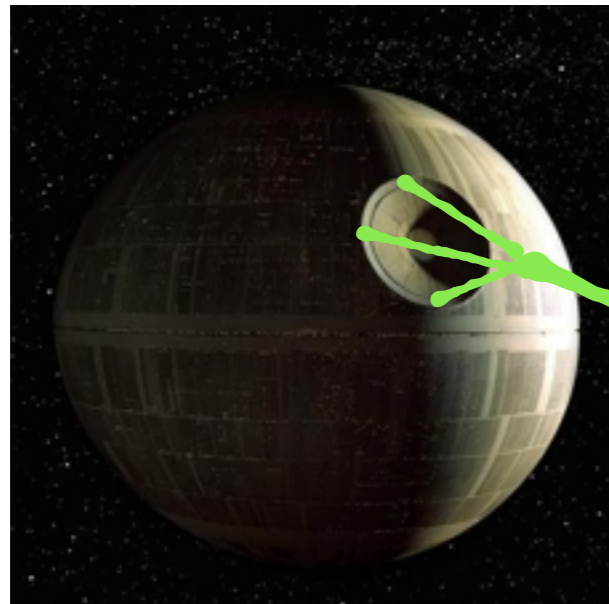


Many-body spectrum

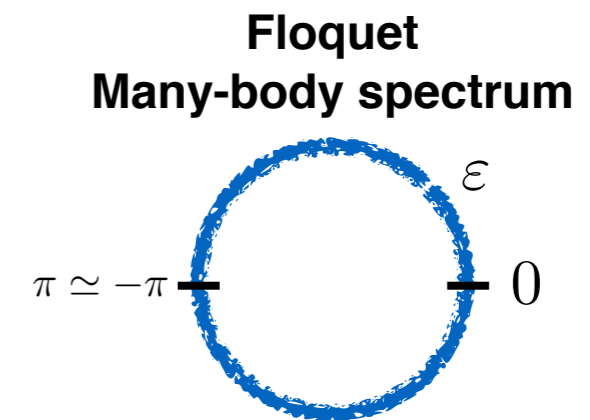
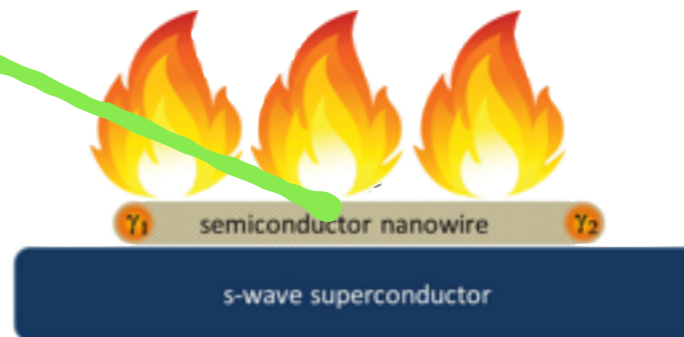
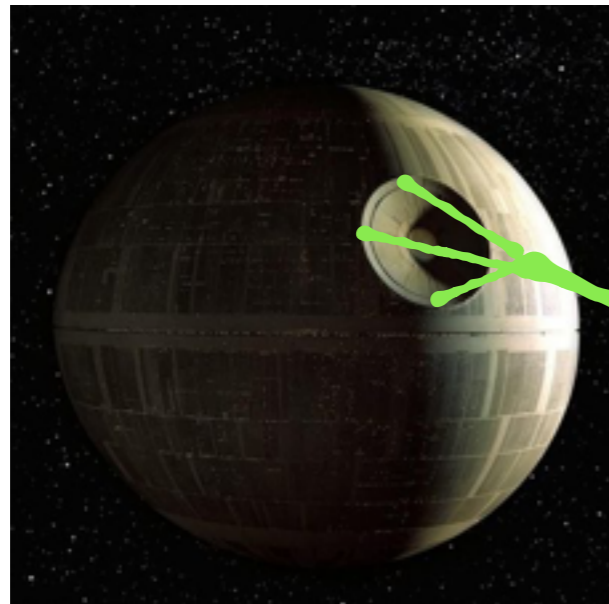
Fate of generic driven/interacting system?



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Fate of generic driven/interacting system?



- runaway heating to effective infinite temperature steady state
- no sense of distinct dynamical “phases”
(all infinite temperature states are equivalent without a phase transition)
- bulk band topology does not govern physics system
(even for arbitrarily weak interactions)

Solution: Many body localization (MBL)

Strong disorder + Isolated system \Rightarrow No thermal equilibrium

Basko, Aleiner, Altshuler; Pal, Huse; many others

All eigenstates have local (area-law) entanglement
(like gapped ground-states)

Bauer, Nayak

“T=0” Quantum order possible in all eigenstates

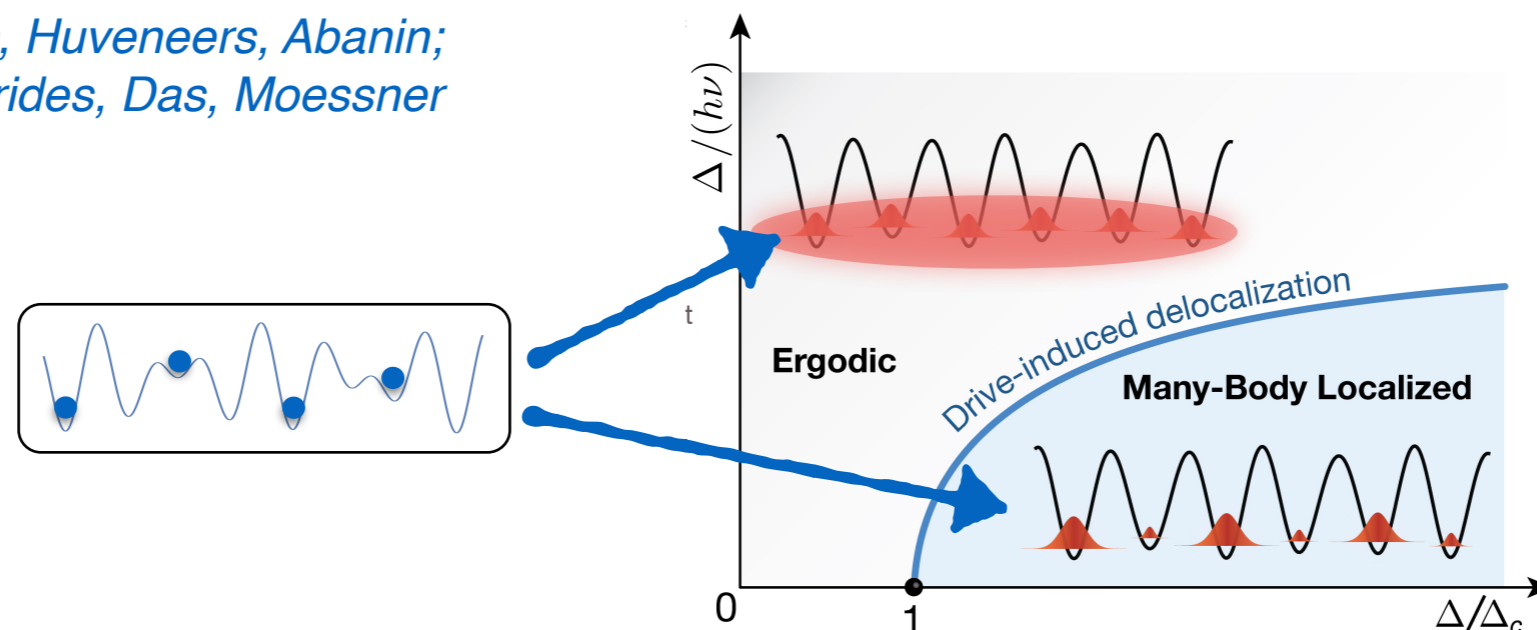
Huse et al.; Bauer, Nayak

Experimental realizations: cold atoms, trapped ions,
(likely many more to come)

Bloch & Monroe Groups

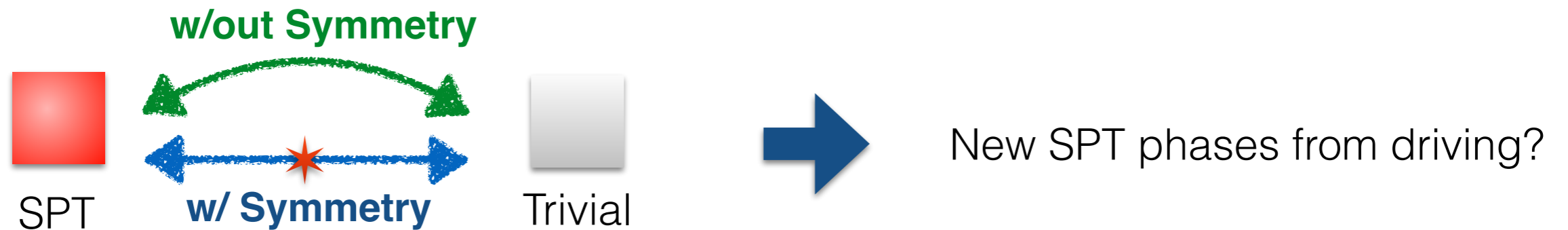
Can (periodically) drive without heating

*Ponte, Papic, Huveneers, Abanin;
Lazarides, Das, Moessner*



Bordia et al. (Bloch group) - arXiv:1607.07868

Driven SPT phases



Haldane Spin-Chain, Topological Insulators,...

1D: ACP, T. Morimoto, A Vishwanath PRX '16
[see also: Keyserlingk, Sondhi; Else, Nayak; Roy, Harper]

2D: ACP, T Morimoto arXiv '16



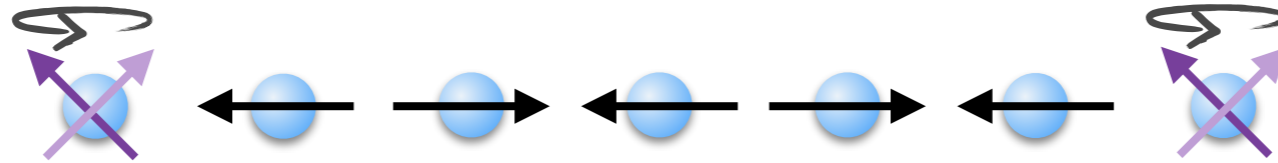
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Ex: Periodically driven Ising chain

(dual of pi-SG model in Khemani, Lazarides, Moessner, Sondhi PRL '16)



$$H(t) = \begin{cases} \sum_i h_i \sigma_i^x & 0 \leq t < T_1 \\ \sum_i J \sigma_i^z \sigma_{i+1}^z & T_1 \leq t < T \end{cases}$$

$$g = \prod_i \sigma_i^x$$

\mathbf{Z}_2 Symmetry

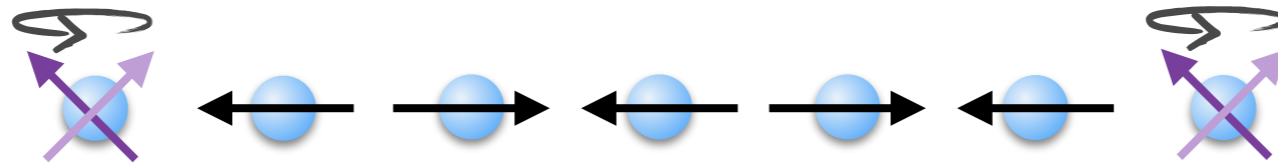
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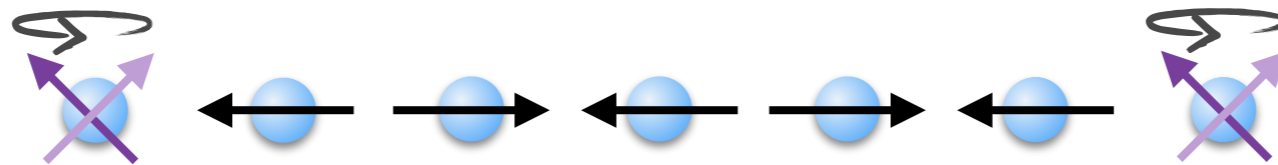
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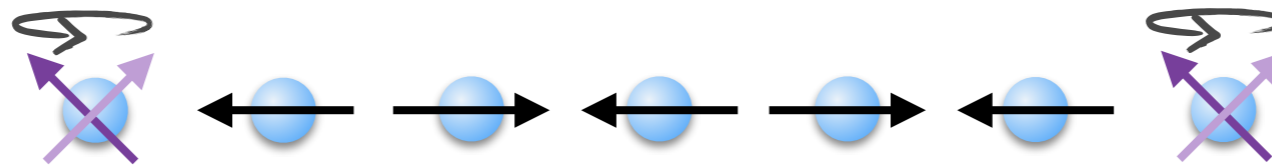
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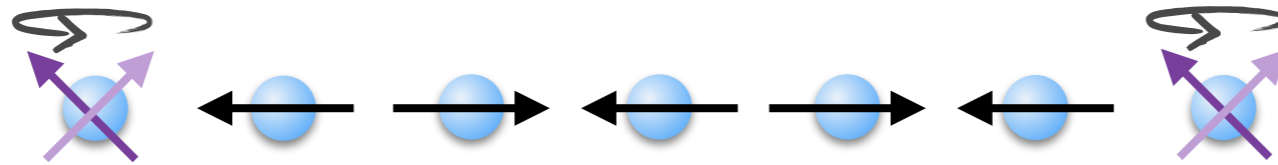
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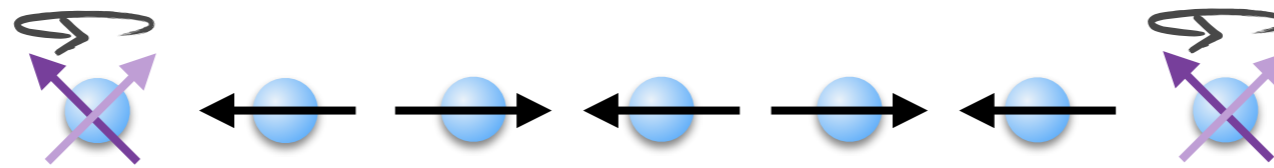
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- Edge spins flip around z each period,
- “Spin-echo” away any symmetry preserving edge field
- Spin-echo usually requires fine-tuning (pi-pulse), **but this is stable to errors!**

ACP, T. Morimoto, A Vishwanath PRX '16

see also: CW von Keyserlingk, S Sondhi PRB '16; DV Else, C Nayak PRB '16, Harper, Roy PRB '16

What protects the edges?

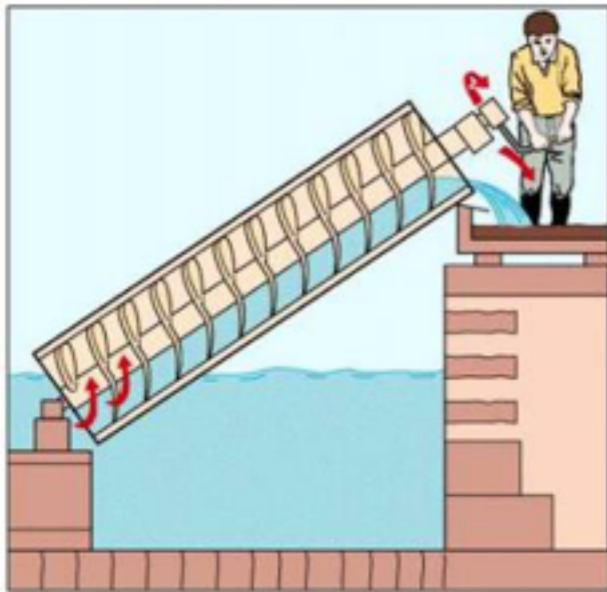


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$$\sigma_{\text{edge}}^x = \pm 1 \rightarrow \mp 1$$

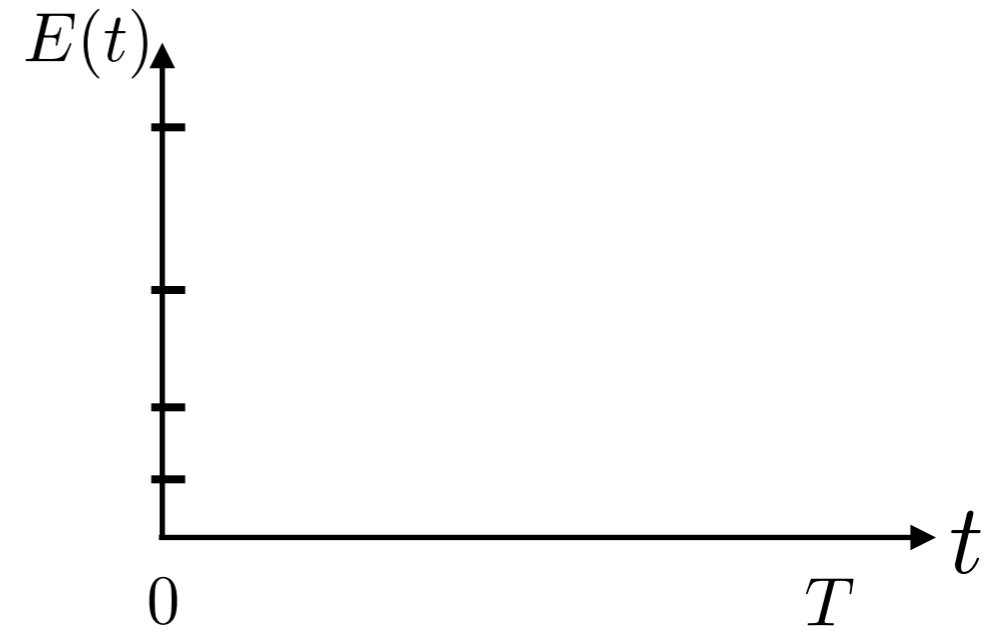
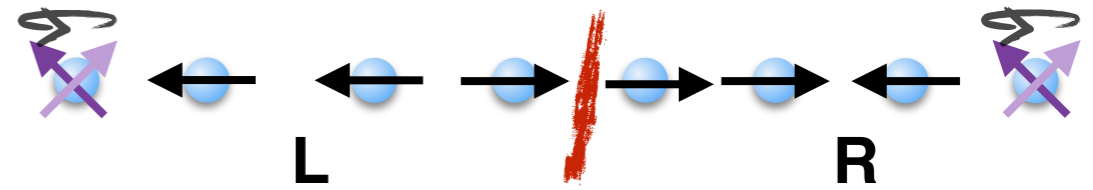


“Quantum archimedes screw”

- Symmetry “charge” (irreducible representation) pumped onto edge during each period
- Possible charges are discrete, “quantized” can’t be continuously altered by small perturbations
- Pumping different symmetry charges (irreducible representations) \Leftrightarrow different FSPT phases

Formal classification

Entanglement spectrum (micromotion)



ACP, T. Morimoto, A Vishwanath PRX '16

ACP, T. Morimoto, A Vishwanath PRX '16 (see also Else & Nayak)

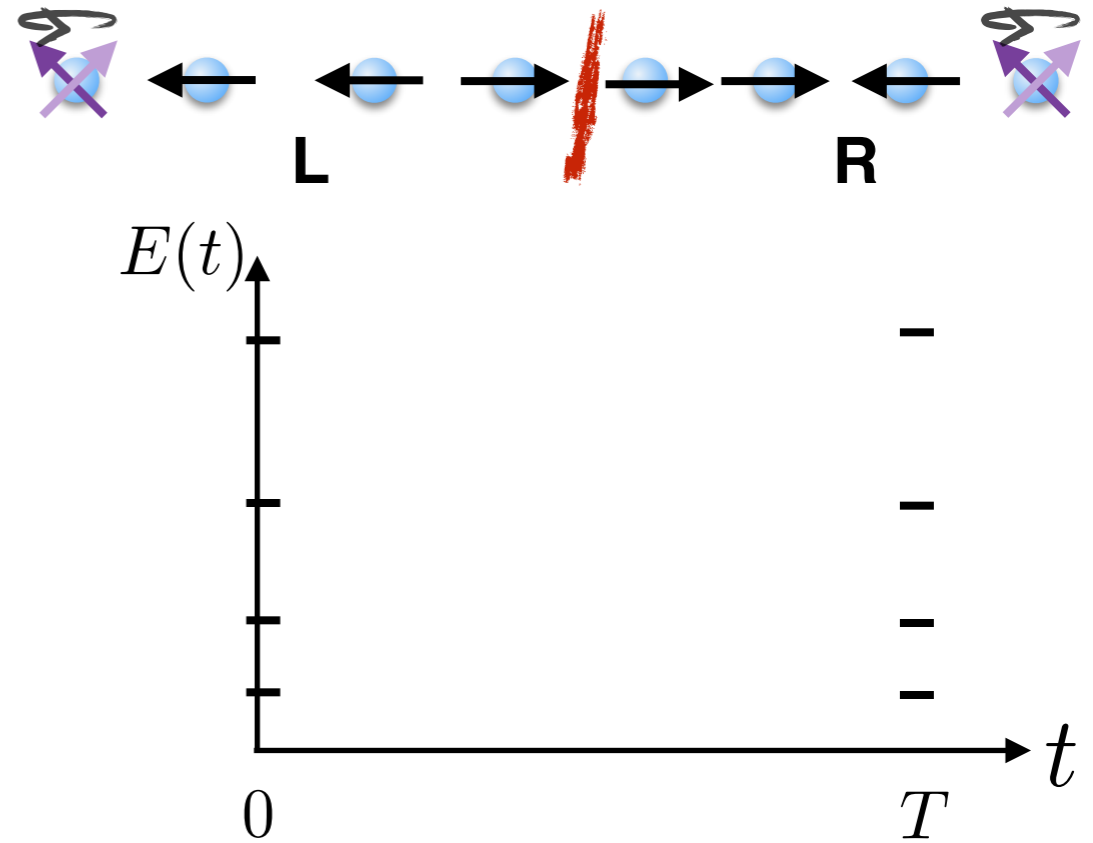
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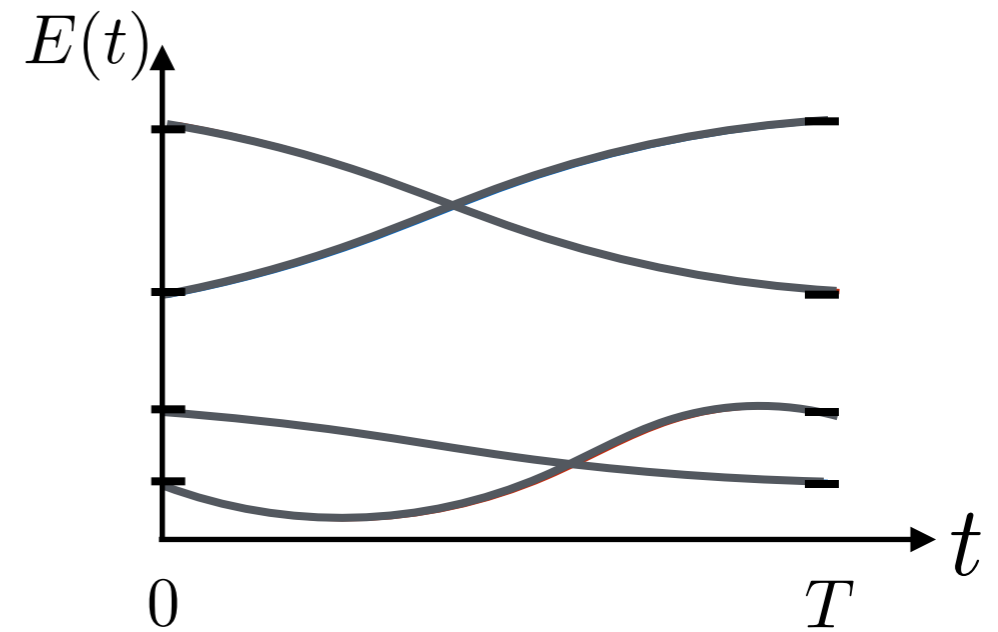
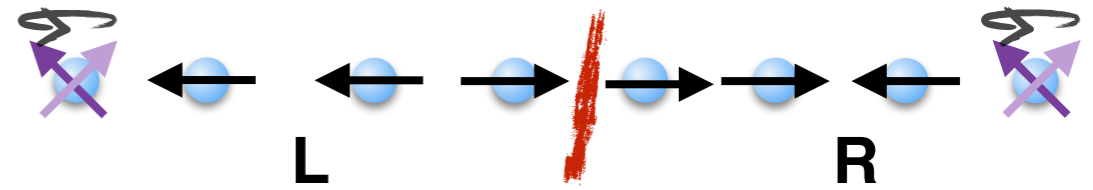
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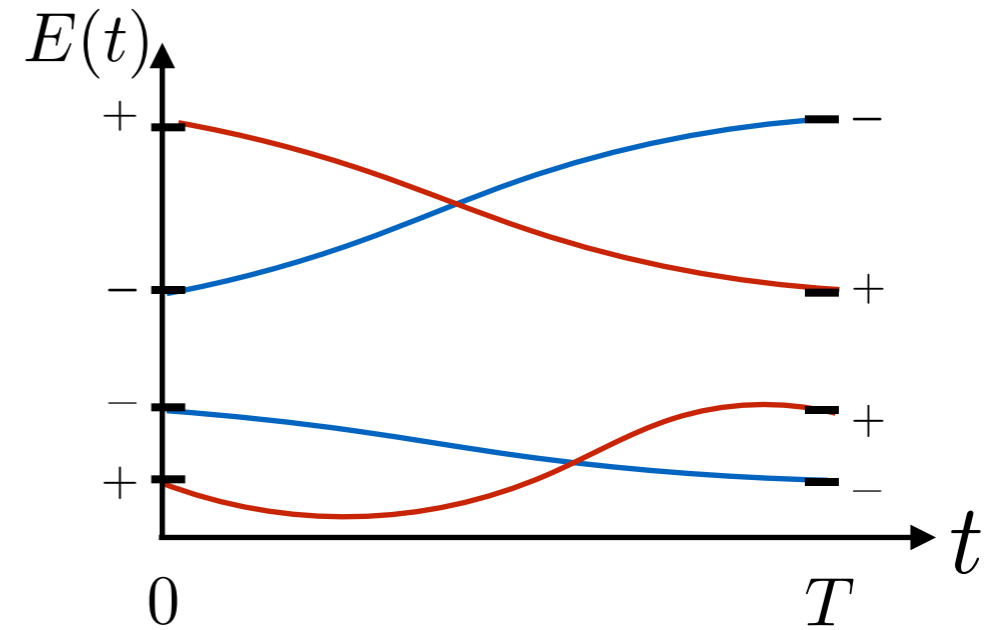
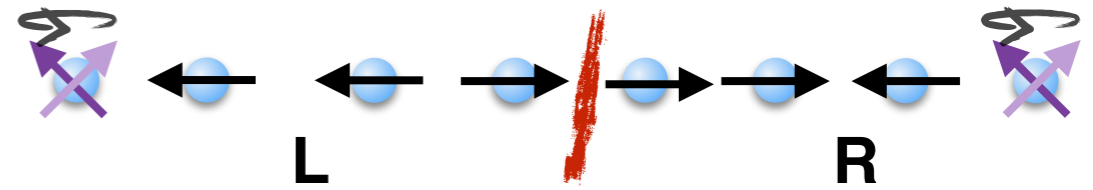
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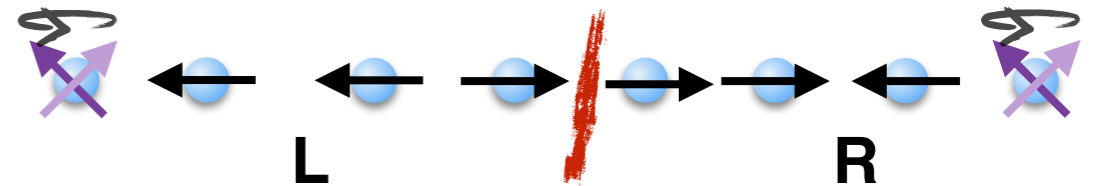
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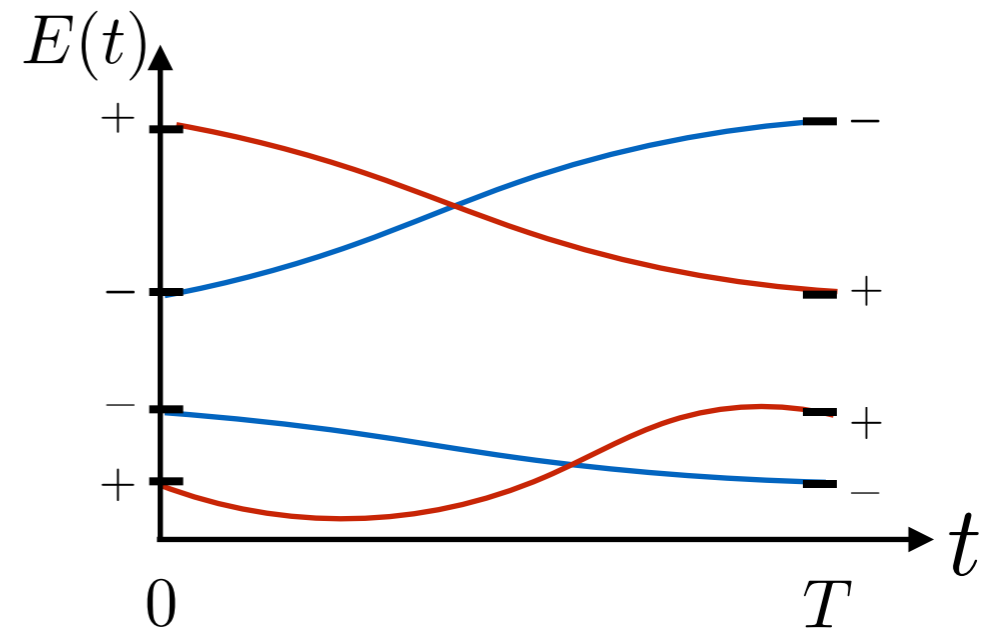
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Cohomology: Extra time-translation symmetry

$$G \longrightarrow G \times \mathbb{Z}$$

Kunneth Formula: $\mathcal{H}^2(G \times \mathbb{Z}, U(1)) = \underbrace{\mathcal{H}^2(G, U(1))}_{\text{1D Equilibrium}} \times \underbrace{\mathcal{H}^1(G, U(1))}_{\text{New Floquet Phases}}$
 (boson phases)

1D Equilibrium New Floquet Phases

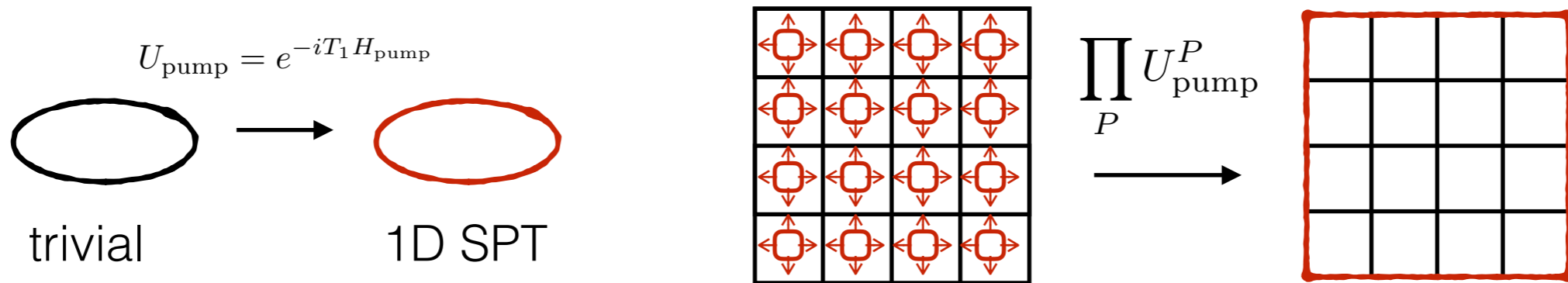
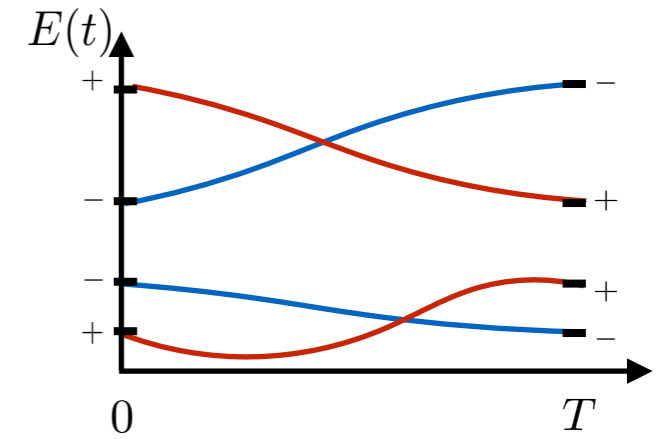
Possible charges = $H^1(G, U(1))$

ACP, T. Morimoto, A Vishwanath PRX '16 (see also Else & Nayak)

Generalizations to 2D

2D Floquet SPTs:

- 0D charge = 0D SPT => 1D SPT



*ACP, T. Morimoto ACP, T. Morimoto arXiv 1610.03485
(see also Else & Nayak PRB '16)*

Floquet enriched topological phases

- Example: gauged Floquet SPT
- More general: pumping 1D topological chains of emergent anyons
- Anyons get permuted each pumping cycle

ACP, T. Morimoto arXiv 1610.03485

Chiral Floquet phases

So far: discrete time-translation = extra symmetry of dynamics

Is this all?



Adrian Po
(Berkeley)



Lukasz Fidkowski
(Stony Brook)



Takahiro Morimoto
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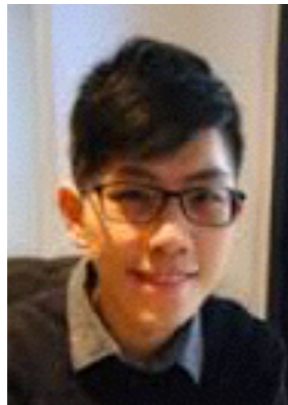


Ashvin Vishwanath
(Harvard)

Chiral Floquet phases

So far: discrete time-translation = extra symmetry of dynamics

Is this all? **No!**



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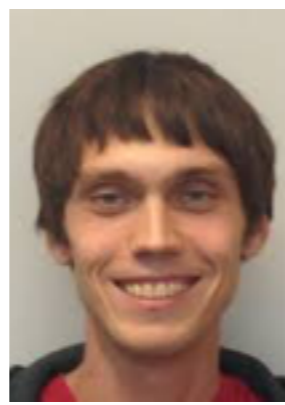
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Intrinsically topological dynamics (no symmetry)



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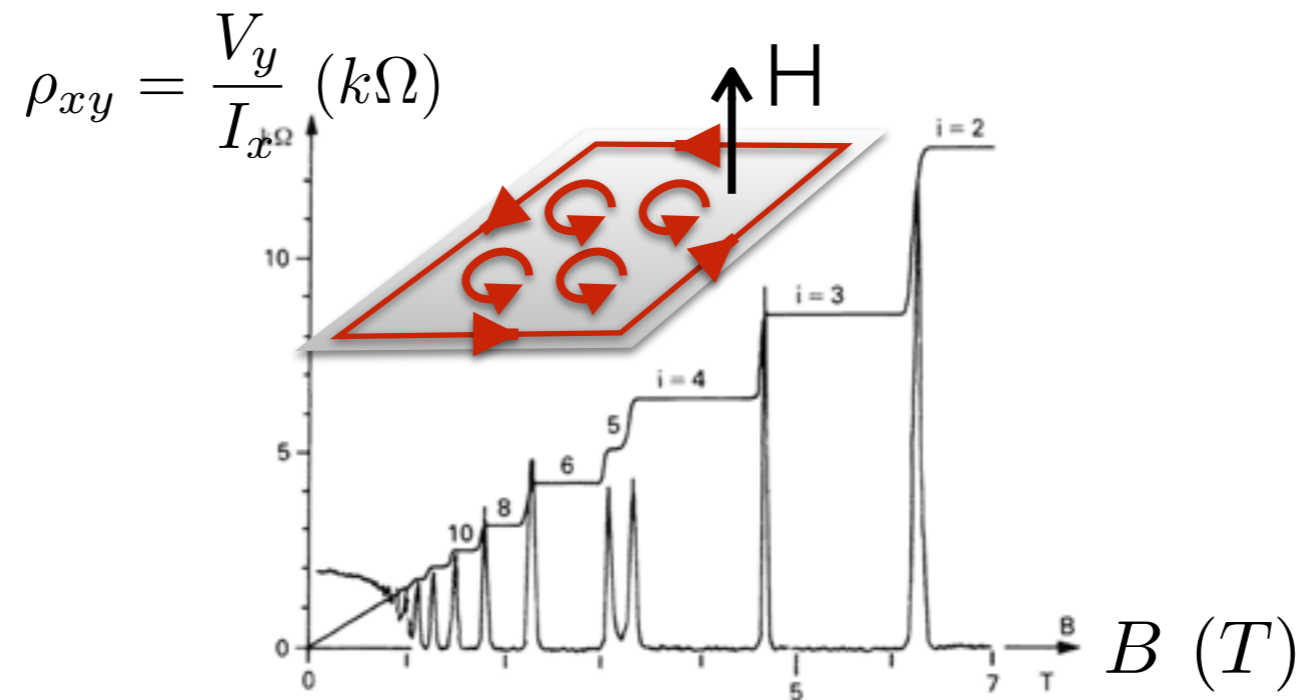


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Non-Equilibrium chiral matter?



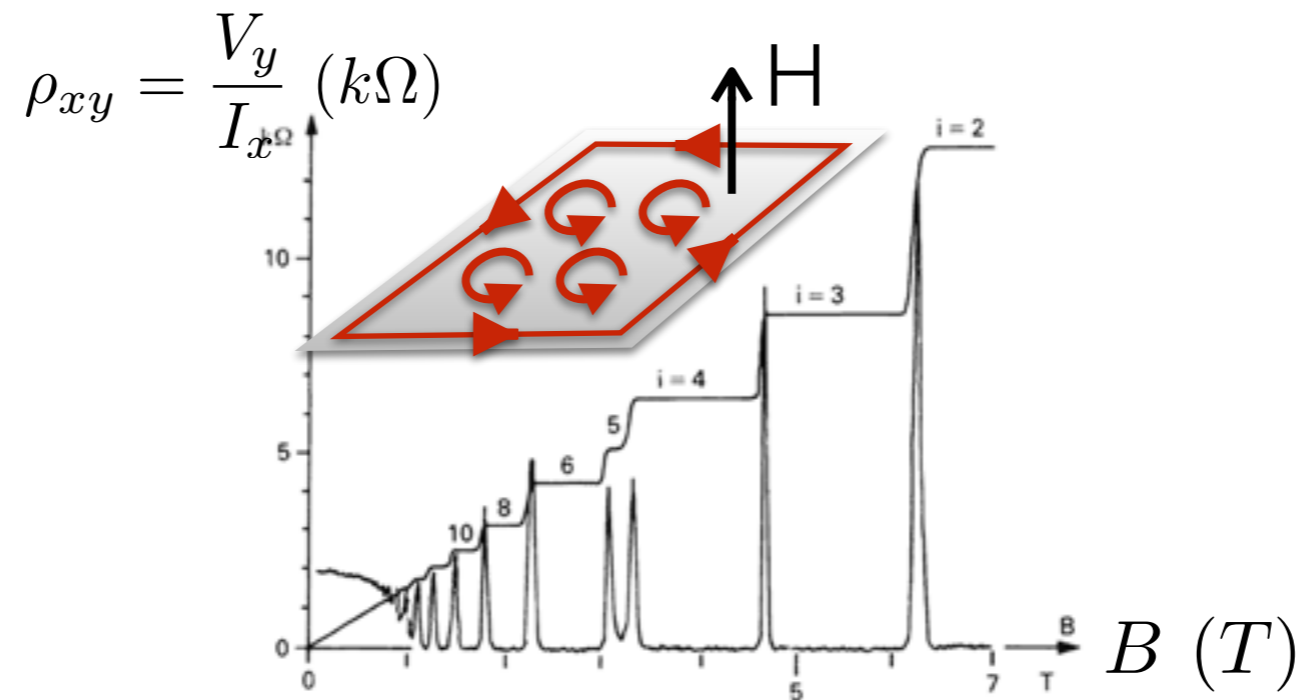
No chiral edges in energy conserving MBL systems

- Chern number = obstacle to localization
- Can prove even in the presence of arbitrary interactions
 - any MBL system will not have “gravitational anomaly” (no thermal quantum Hall effect)

*Halperin, '82
Nandkishore ACP, PRB '14*

*Kitaev '06
ACP, Vishwanath arXiv '15*

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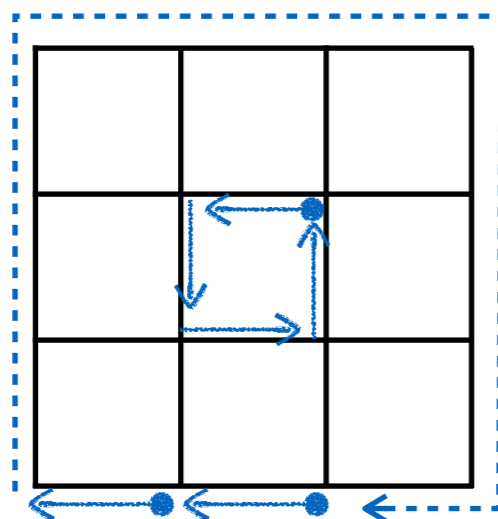
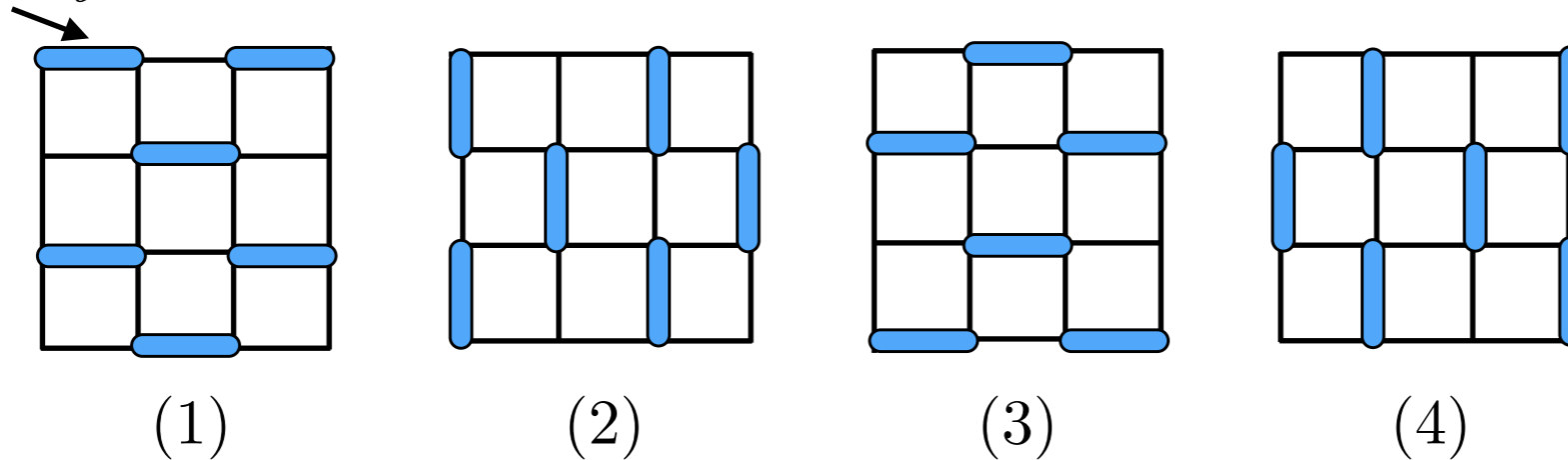
What if energy is not conserved?

SWAP model

(Direct bosonic analog of free-fermion version
by Rudner, Berg, Levin, Titum, Lindner, Refael PRX '13, '16)

$$U(T) = e^{-iH_5} e^{-iH_4} e^{-iH_3} e^{-iH_2} e^{-iH_1}$$

$$\text{SWAP}_{ij} = e^{i\pi \vec{S}_i \cdot \vec{S}_j}$$



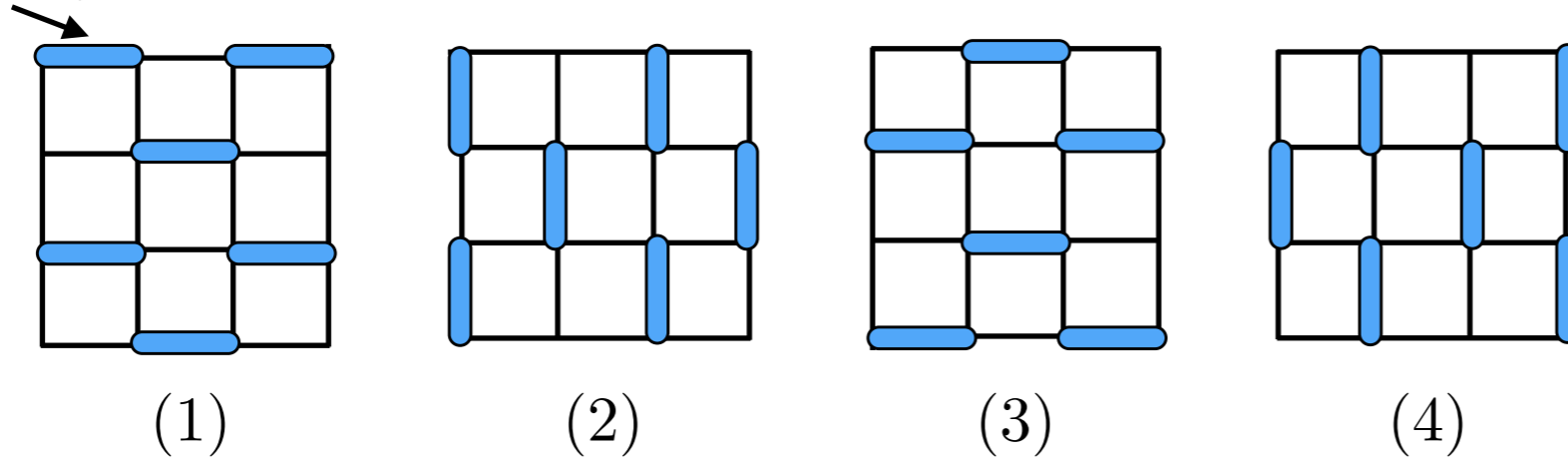
$$U_{1\dots 4}(T) = \begin{cases} 1 & \text{bulk} \\ \hat{T}_1 & \text{edge} \end{cases}$$

SWAP model

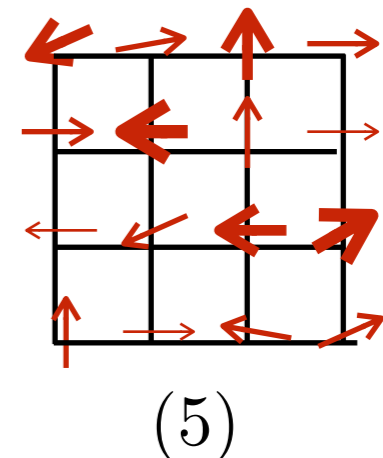
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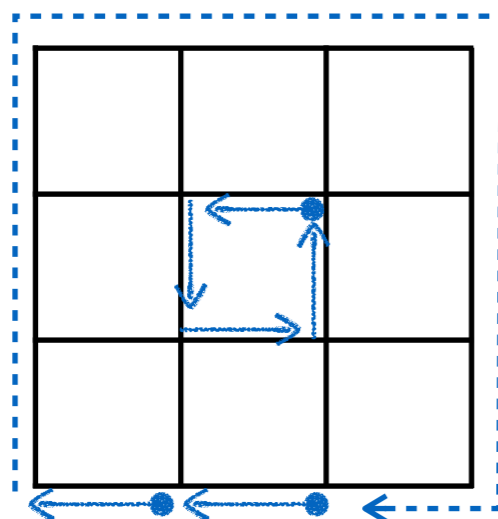


Disorder step



$$H_5 = \sum_i \vec{h}_i \cdot \vec{S}_i$$

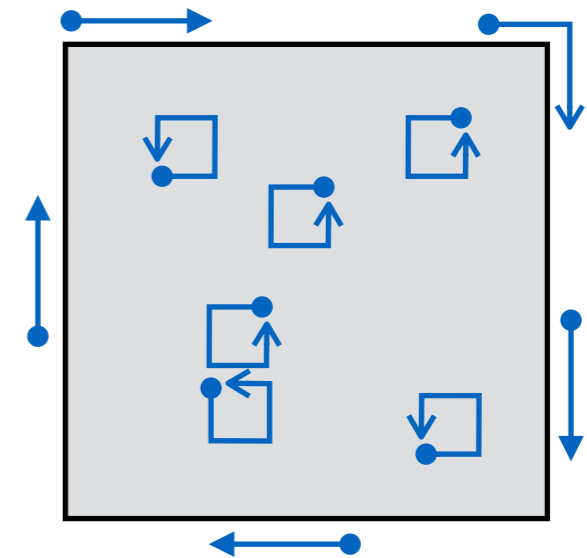
Localized bulk:
edge dynamics
occurs separate
from bulk



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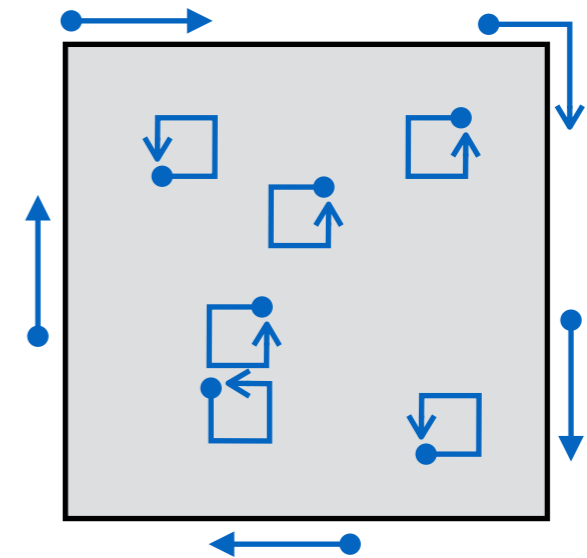
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Rudner, Berg, Levin, PRX '2013



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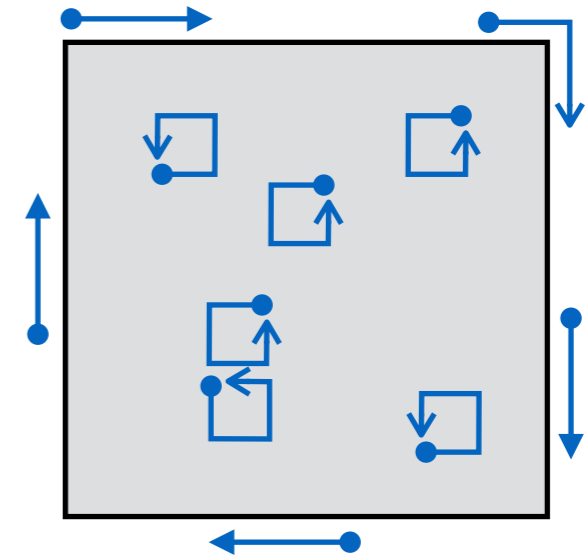
Single particle winding invariant (fermion version) *Rudner, Berg, Levin, PRX '2013*



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Many-particle invariant?

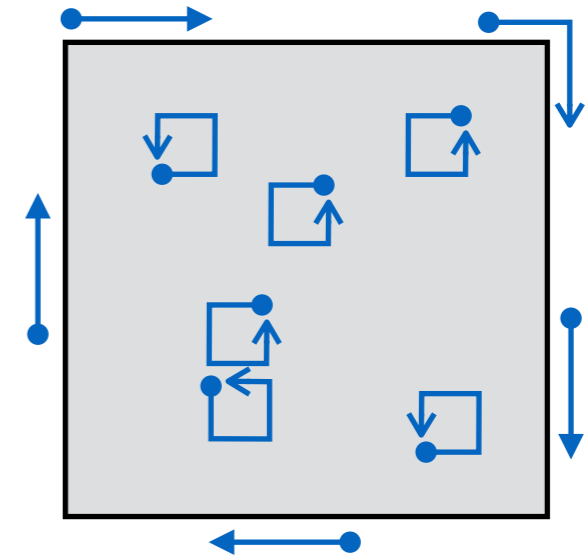


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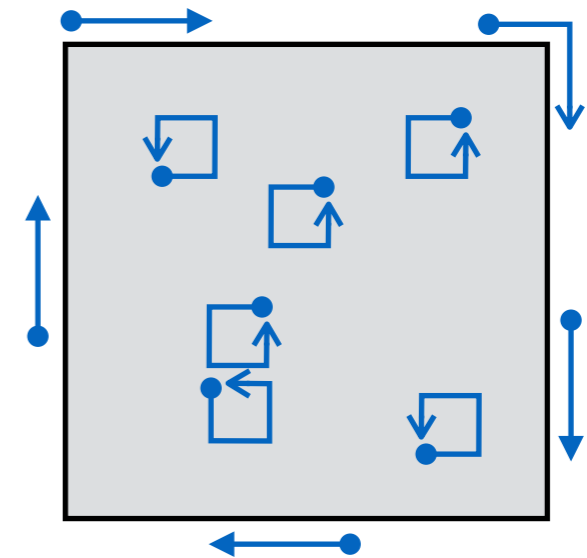
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Nathan et al. arXiv:1610.03590



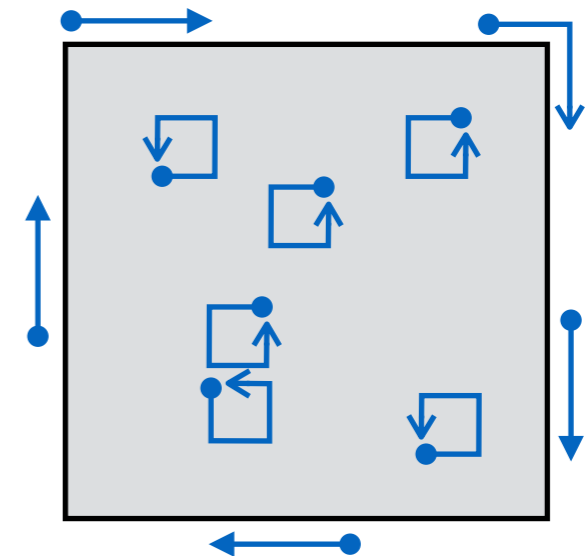
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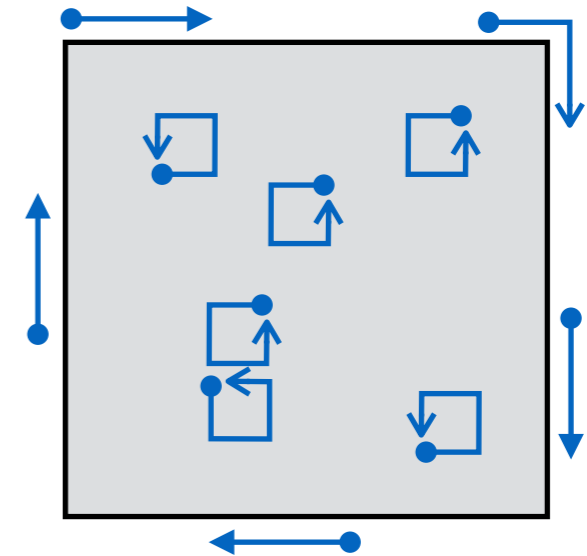
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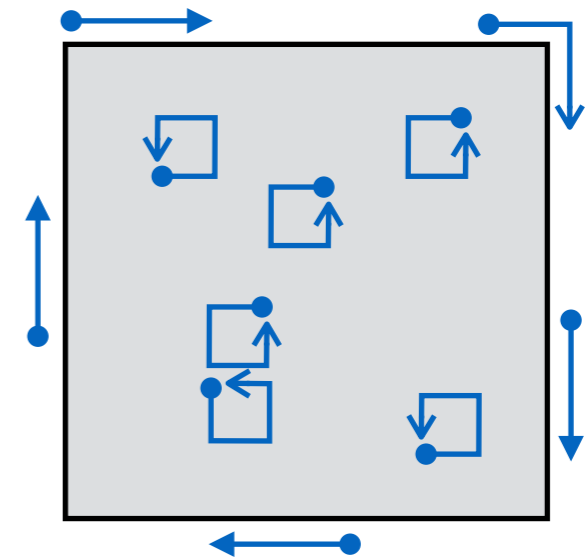
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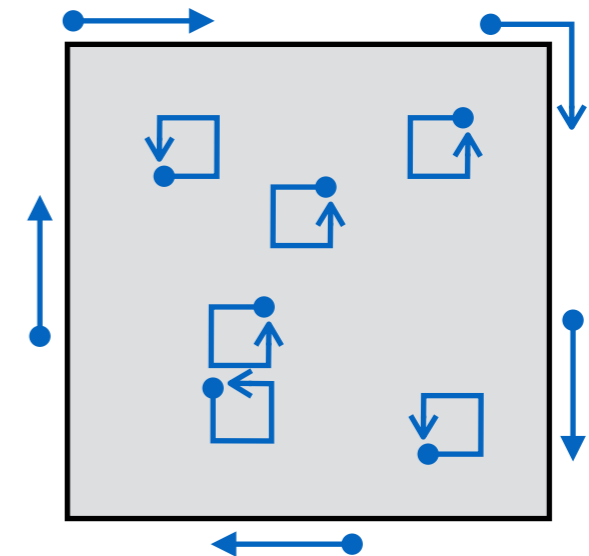
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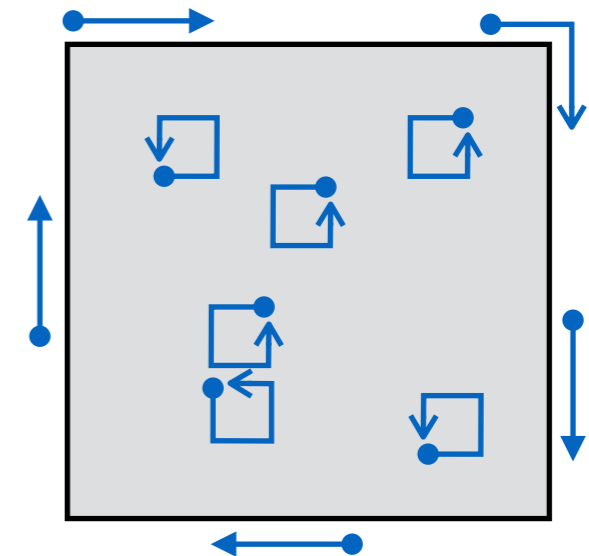
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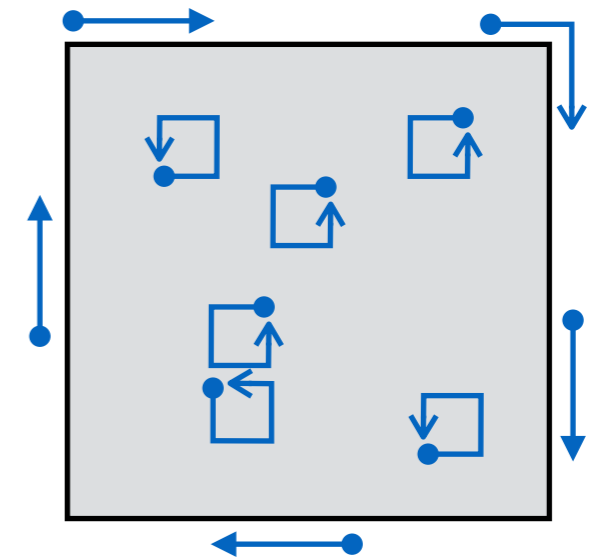
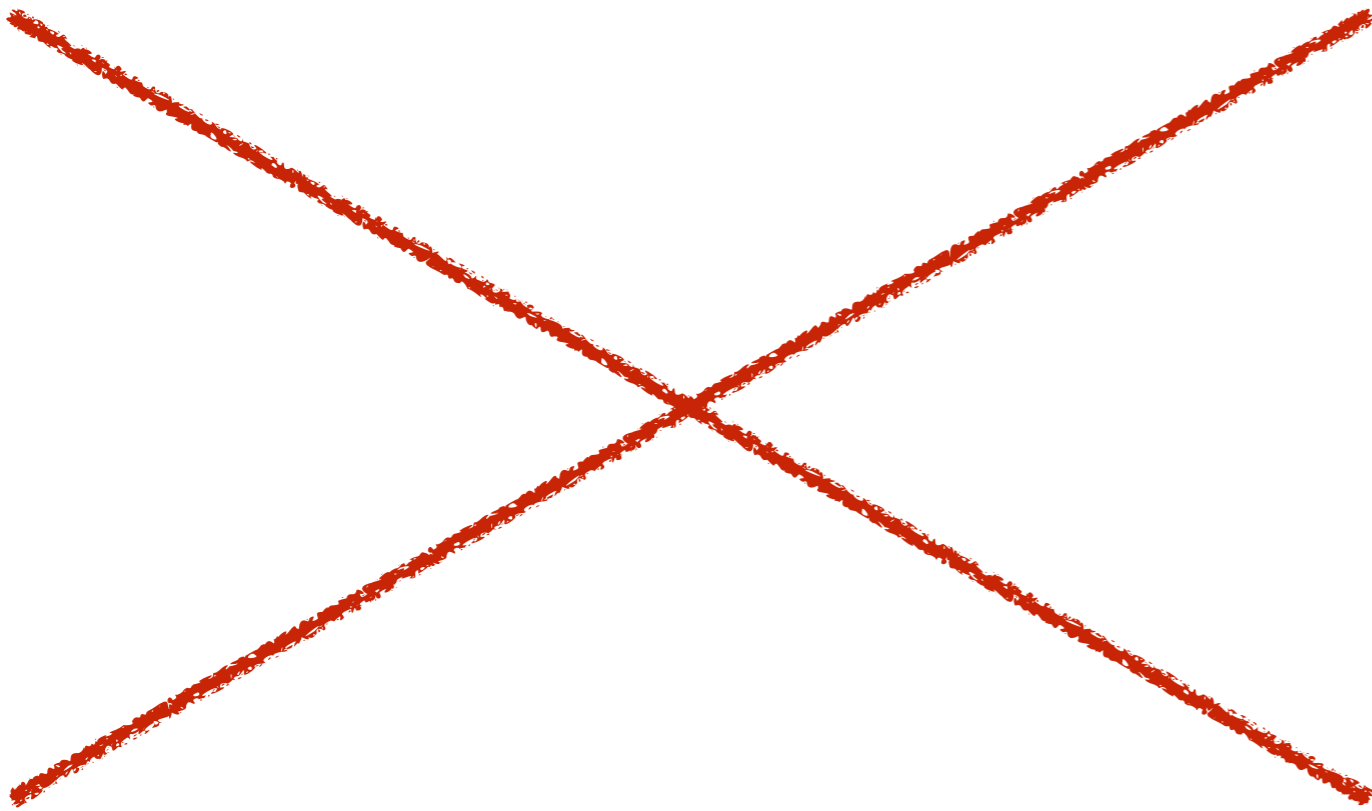
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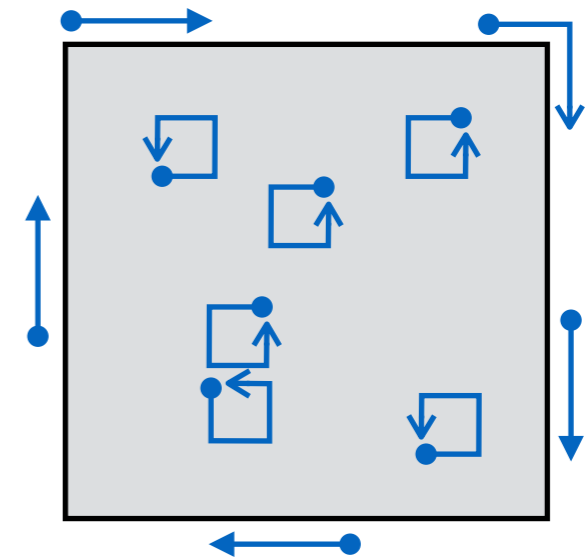
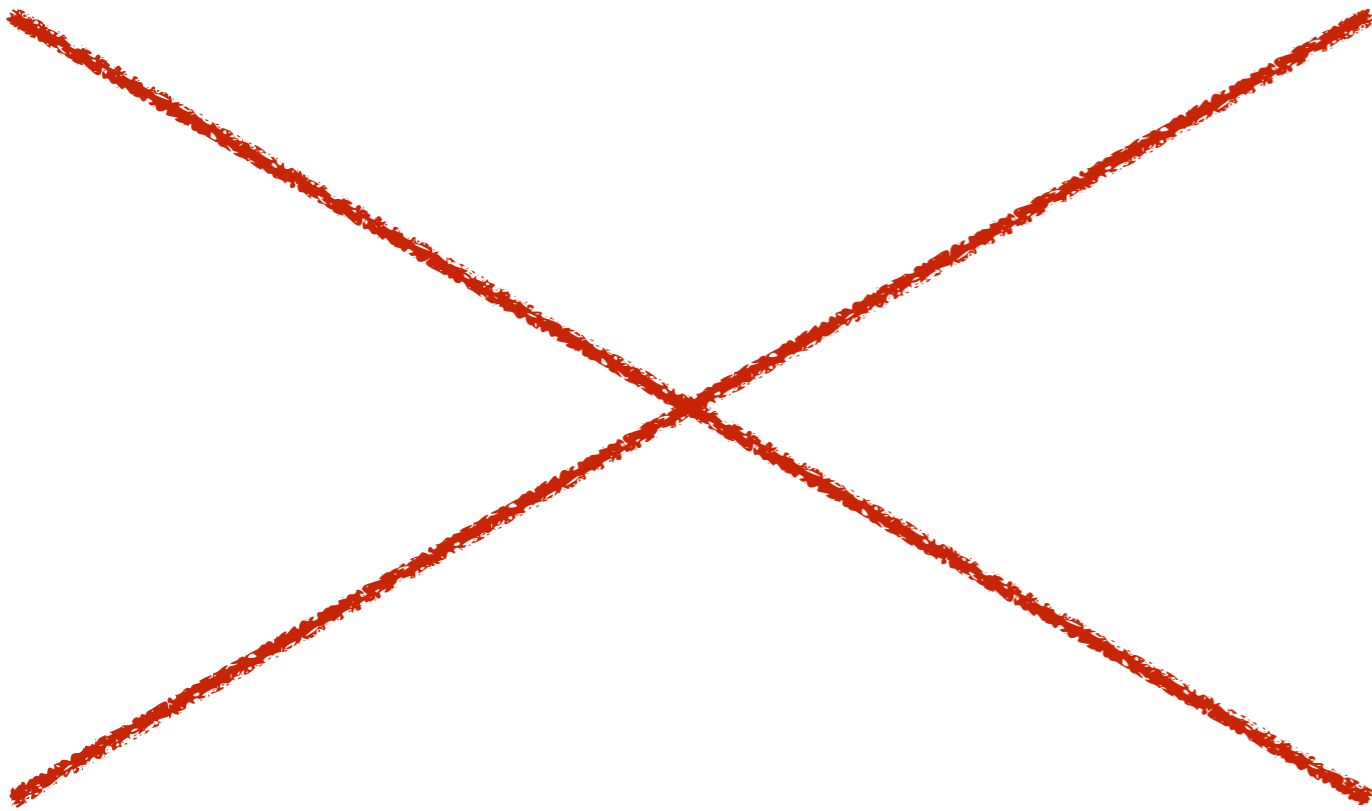
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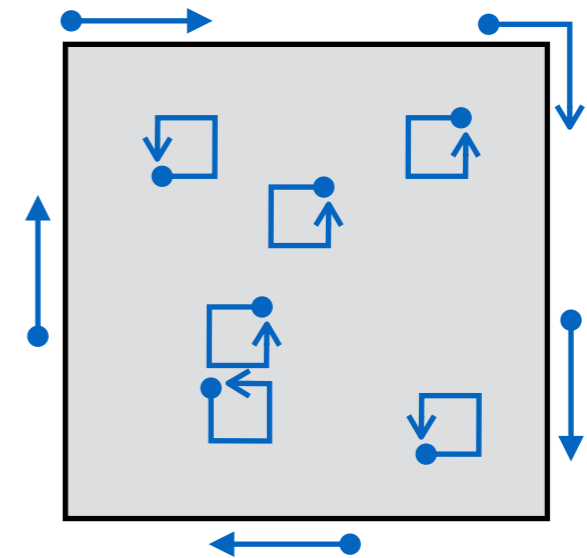
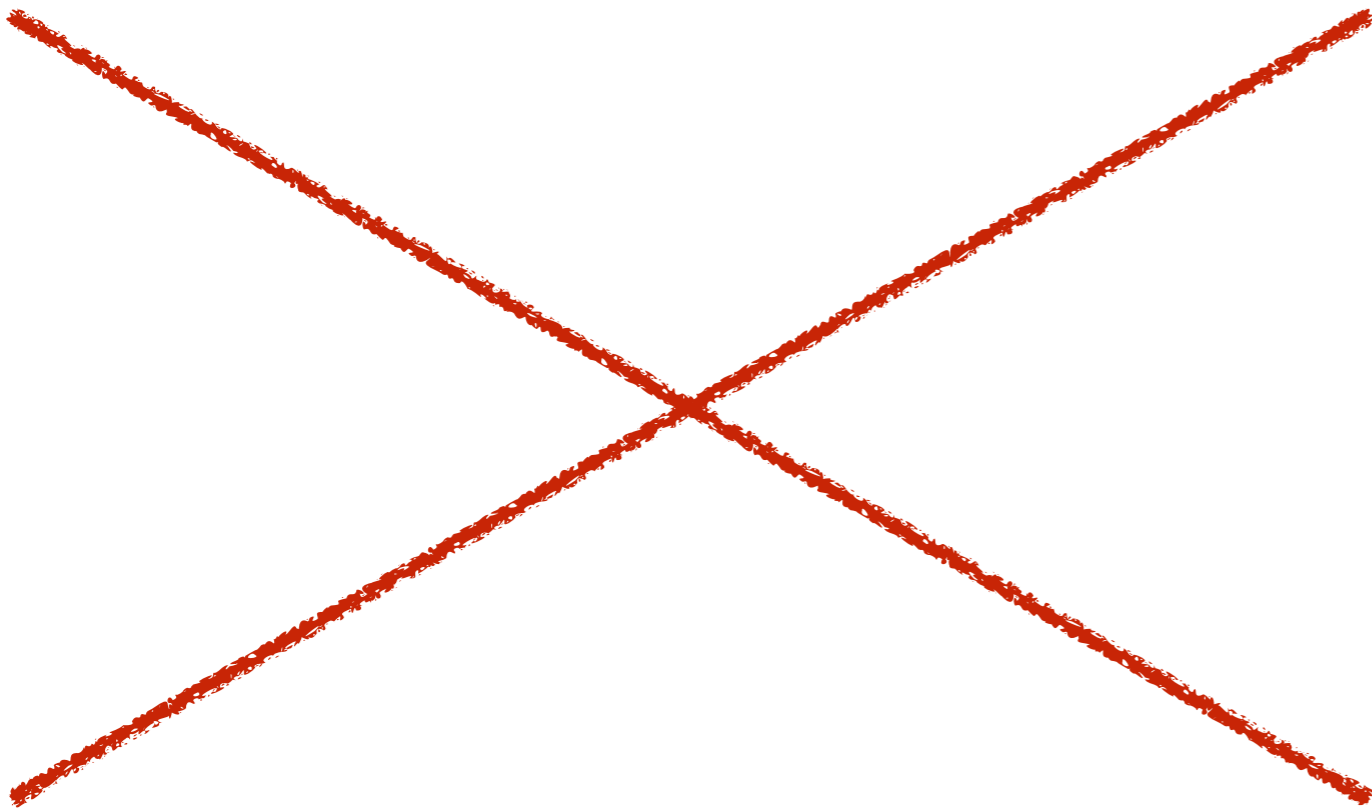
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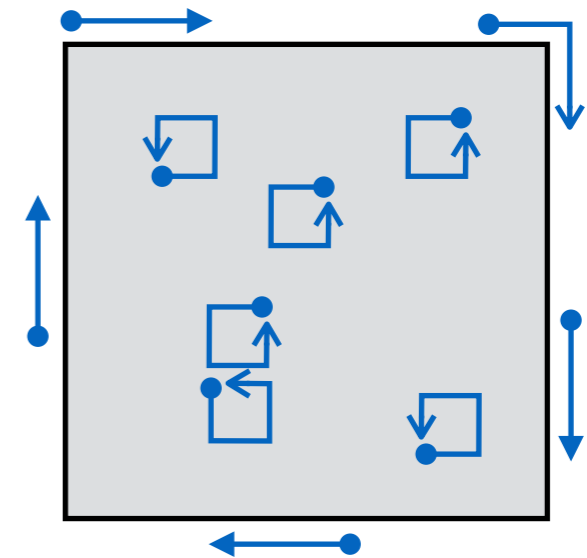
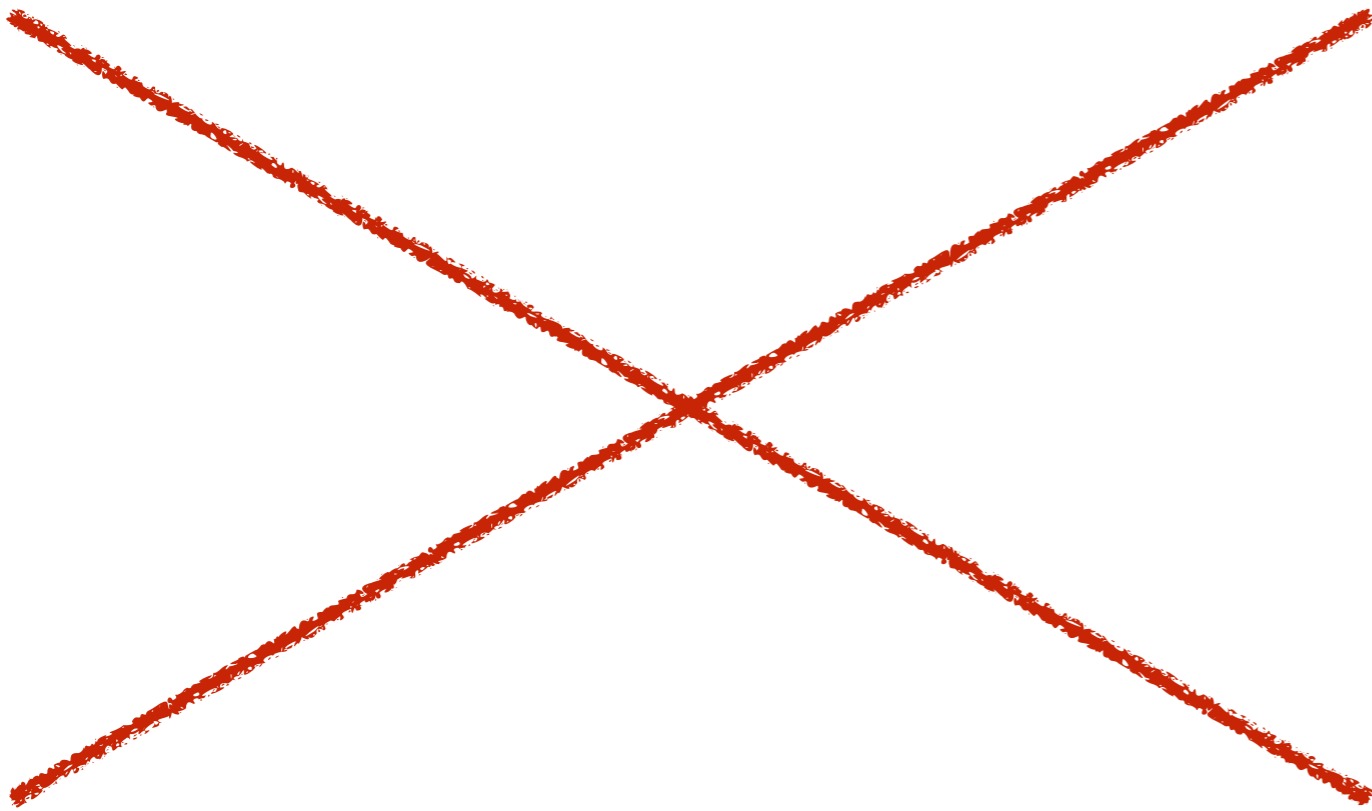
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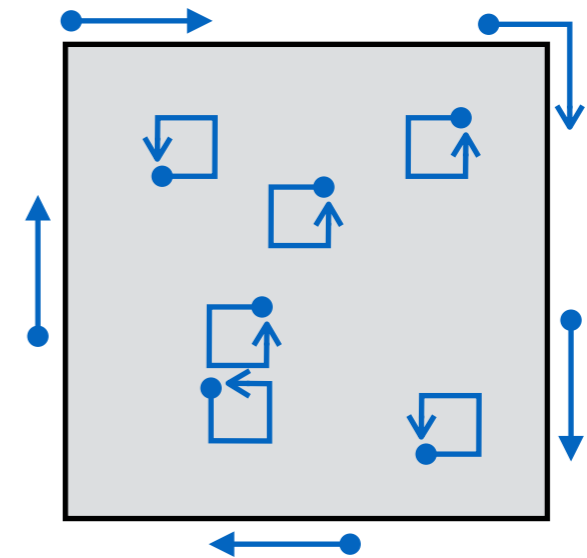


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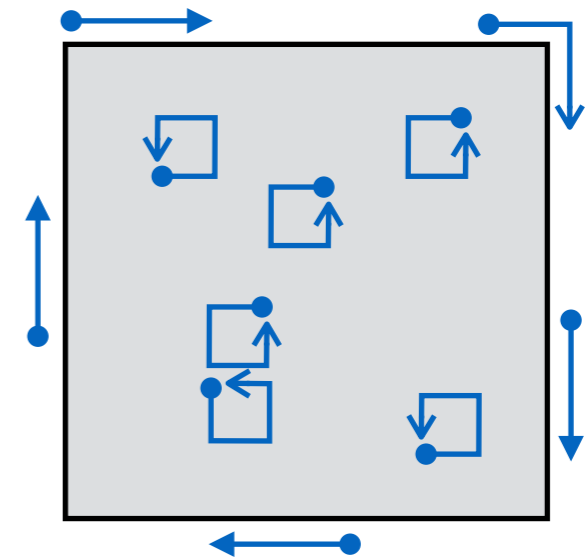
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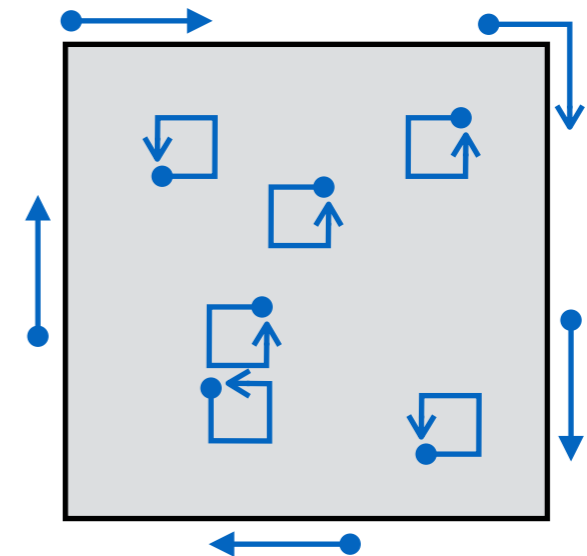
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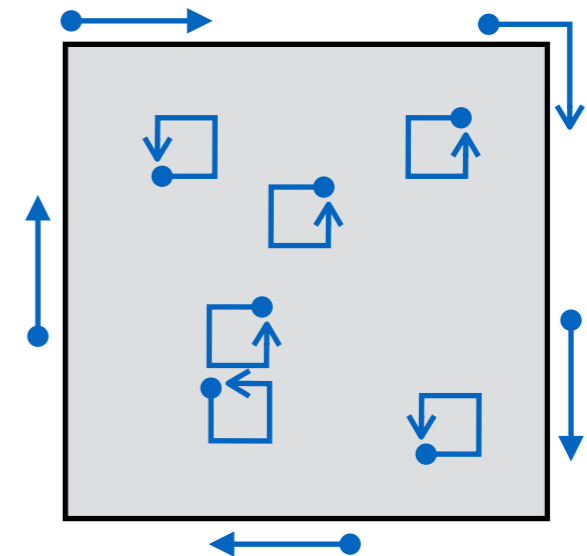
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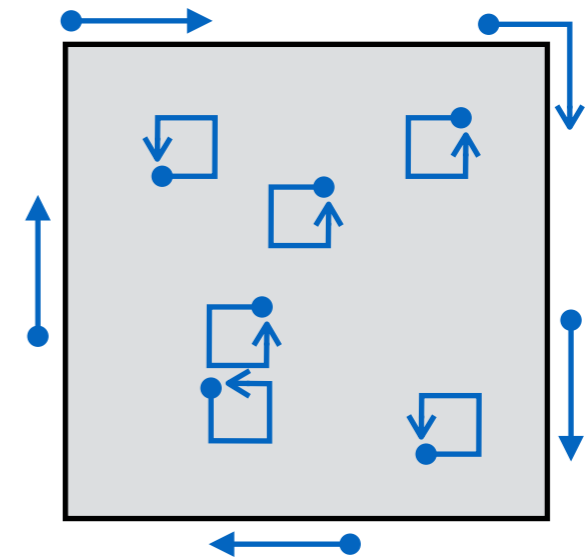
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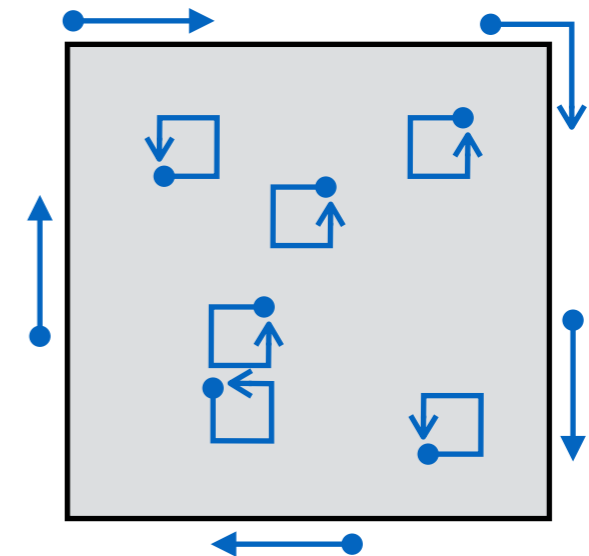
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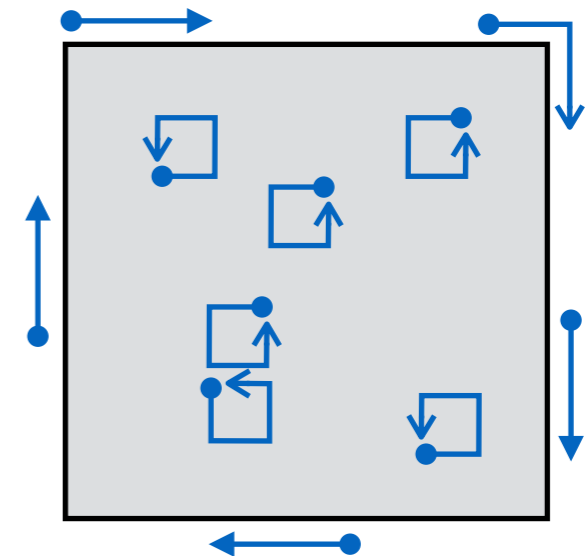
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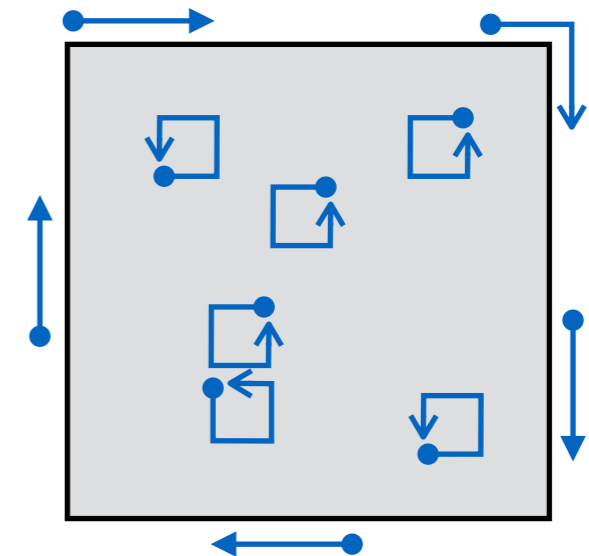
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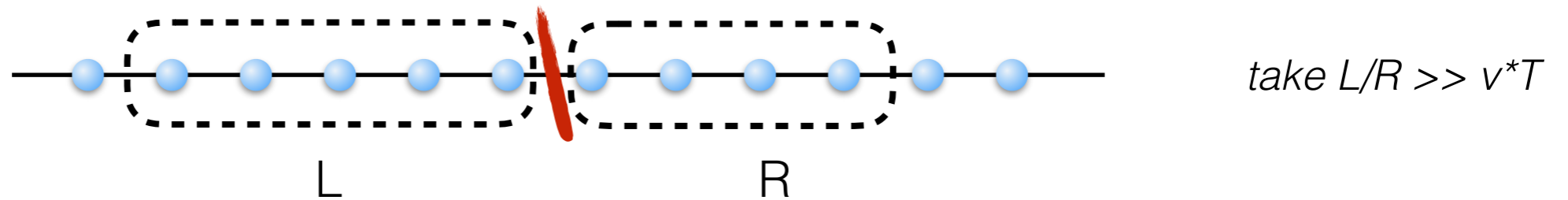
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Quantized pumping of quantum information around edge

A quantum information flow gauge



Algebra of observables: $\mathcal{A} = \left\{ \sum_{i,j=1}^{\mathcal{D}_A} a_{ij} |i\rangle\langle j| ; a_{ij} \in \mathbb{C} \right\}$ $e_{ij} = |i\rangle\langle j|$

Overlap of algebras:

$$\langle \mathcal{A}, \mathcal{B} \rangle = \frac{\sqrt{\mathcal{D}_A \mathcal{D}_B}}{\mathcal{D}_{\text{tot}}} \sqrt{\sum_{i,j=1}^{\mathcal{D}_A} \sum_{l,m=1}^{\mathcal{D}_B} \left| \text{tr} \left(e_{ij}^{a\dagger} e_{lm}^b \right) \right|^2}$$

$$\langle \mathcal{A}, \mathcal{B} \rangle = \begin{cases} 1 & [\mathcal{A}, \mathcal{B}] = 0 \\ \mathcal{D}_A & \mathcal{A} = \mathcal{B} \end{cases}$$

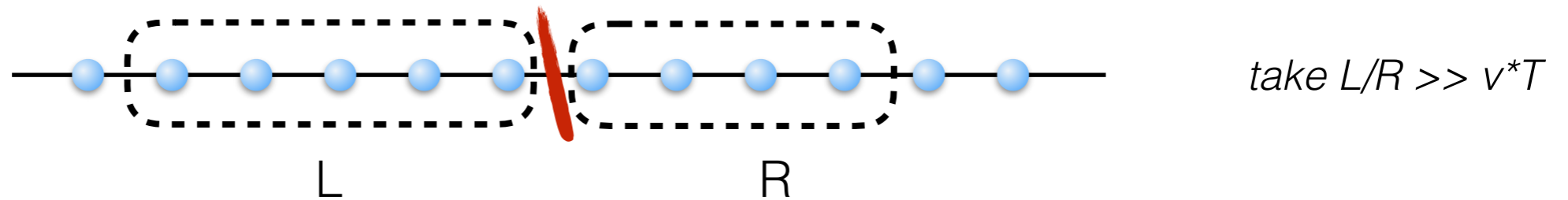
$\log \langle \mathcal{A}, \mathcal{B} \rangle =$ “how much information about \mathcal{A} is contained in \mathcal{B} ”

Topological index (GNVW):

$$\nu = \log \frac{\langle U(\mathcal{A}_L), \mathcal{A}_R \rangle}{\langle \mathcal{A}_L, U(\mathcal{A}_R) \rangle} \in \log \mathbb{Q}_+ \quad \text{imbalance of information flow } (-> \text{left}) - (<- \text{right})$$

A quantum information flow gauge

D. Gross, V. Nesme, H. Vogts, R.F. Werner arXiv:0910.3675



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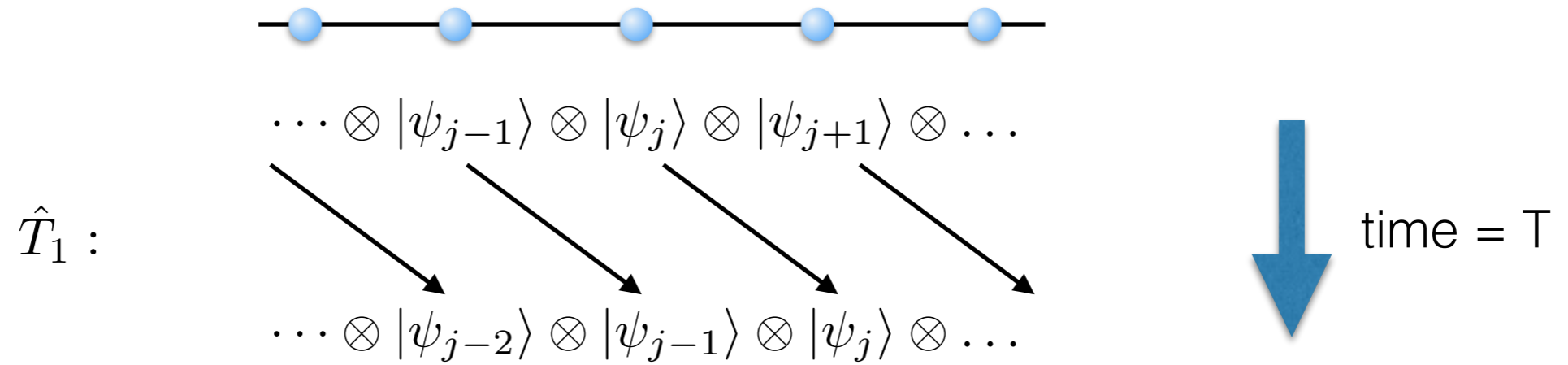
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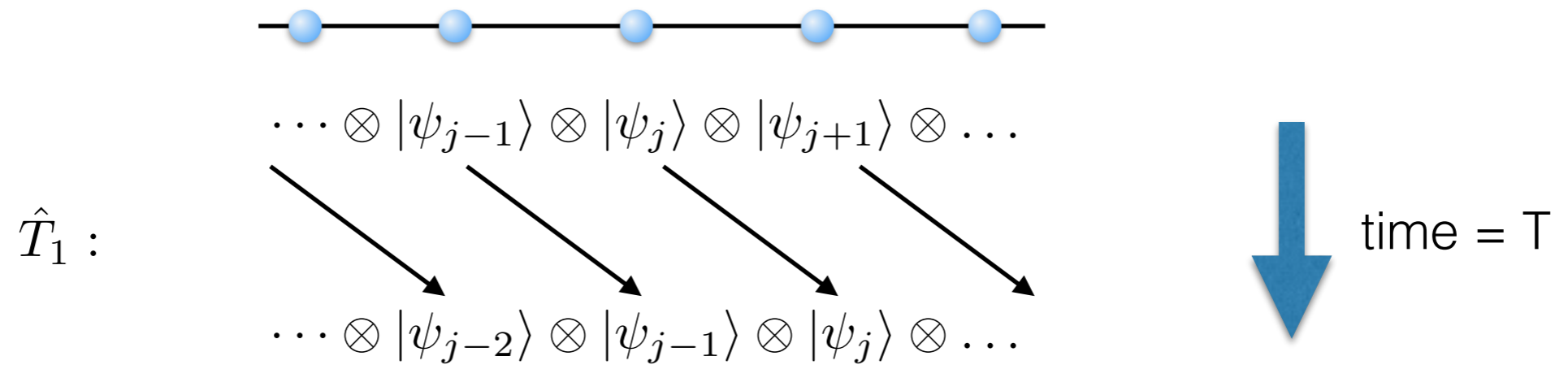
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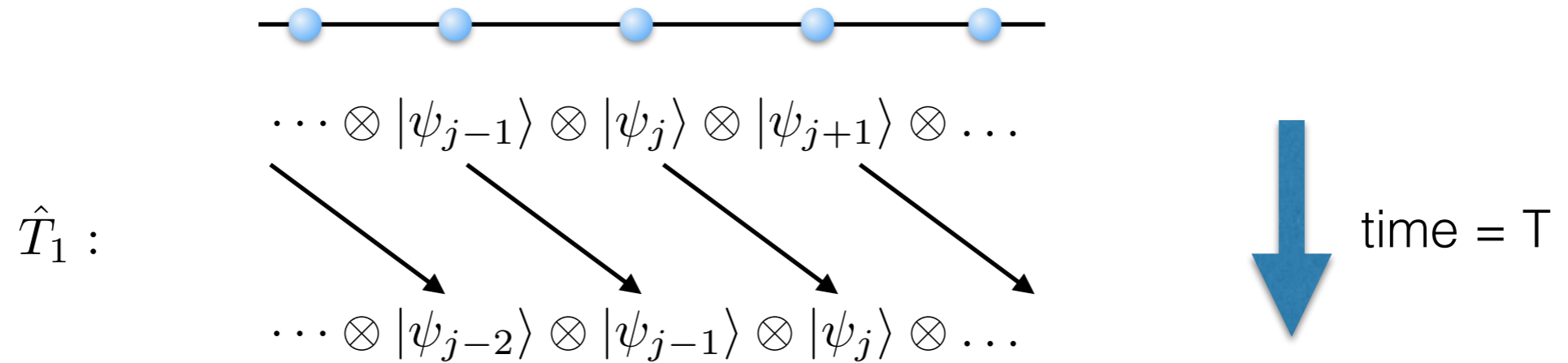


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T_1 is acts locally: (information propagates w/ finite speed)

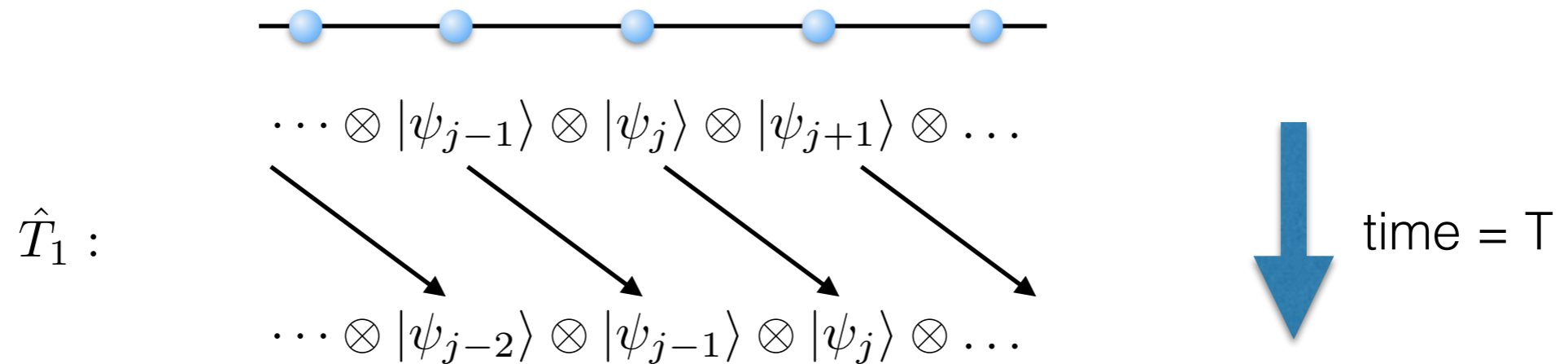
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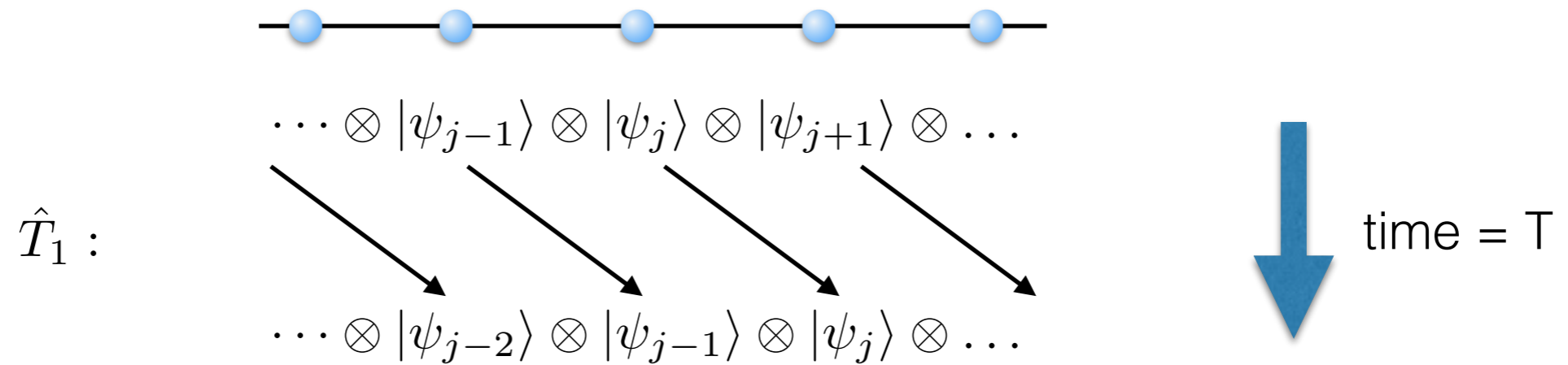
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Cannot be realized with time evolution by a local Hamiltonian

- Intuitive argument:
 - local Hamiltonian can always have edge
 - behavior at edge is sick! ``pile-up'' of states (non-unitary)

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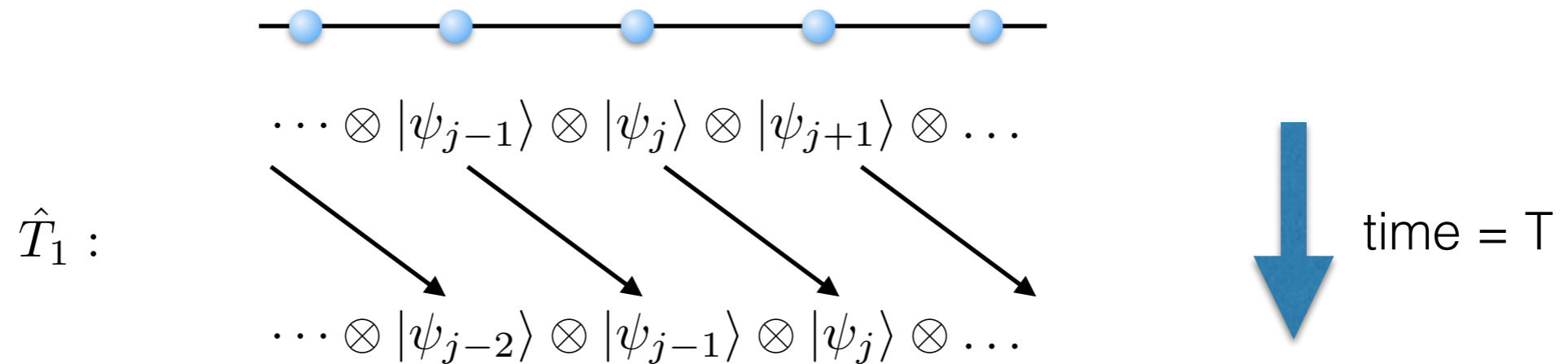
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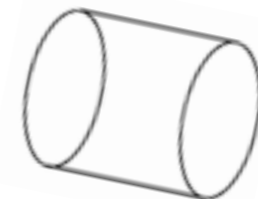
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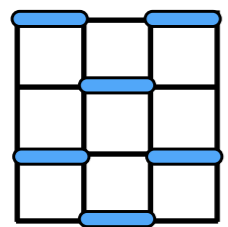
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 - local Hamiltonian can always have edge
 - behavior at edge is sick! “pile-up” of states (non-unitary)
- Proof: [D. Gross, V. Nesme, H. Vogts, R.F. Werner arXiv:0910.3675](#)

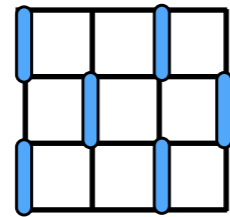
but, OK for boundary of local 2D system (no edge)



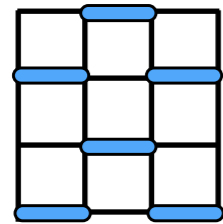
Dynamical chiral index of SWAP model



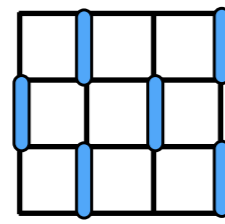
(1)



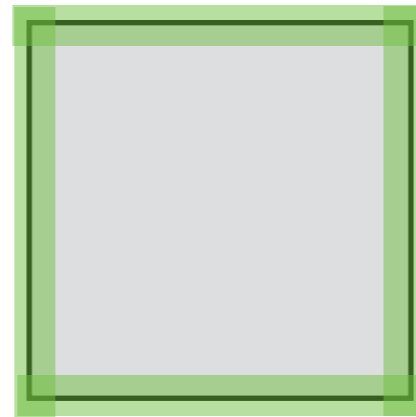
(2)



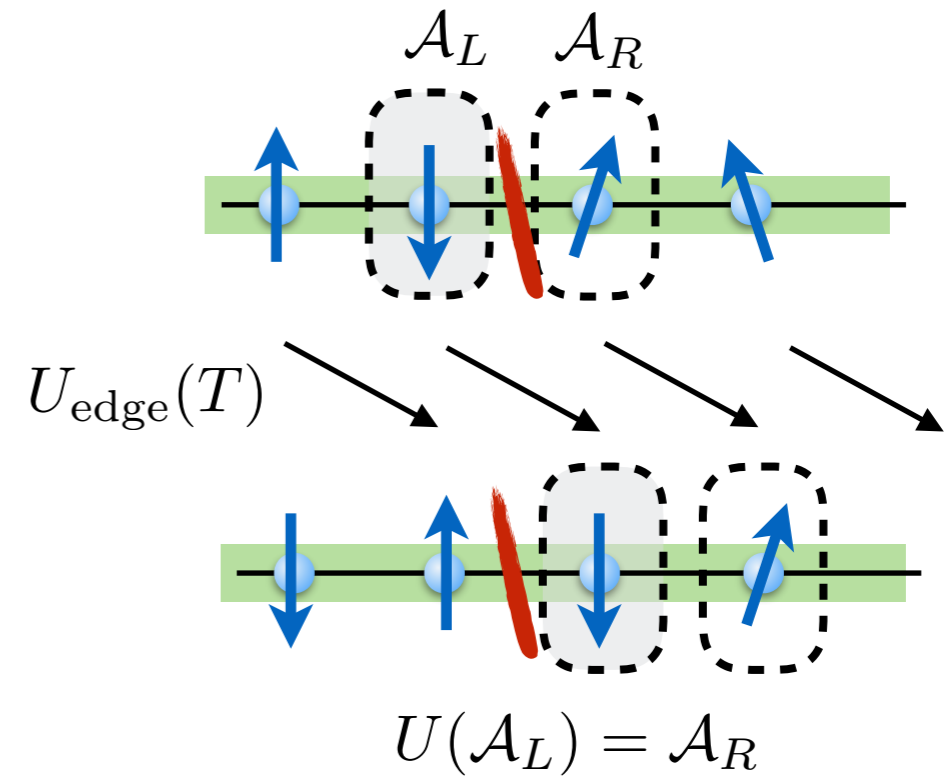
(3)



(4)

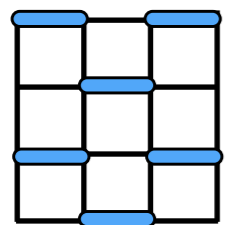


$$U_{(\text{MBL})} = U_{\text{edge}} U_{\text{bulk}}$$

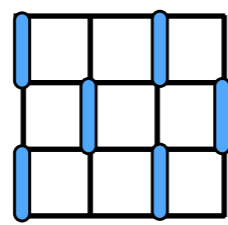


$$\nu_{\text{SWAP}} = \log \frac{\langle U(\mathcal{A}_L), \mathcal{A}_R \rangle}{\langle \mathcal{A}_L, U(\mathcal{A}_R) \rangle} = \log(2)$$

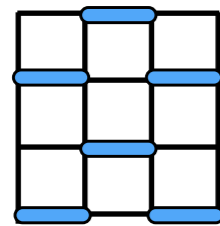
Dynamical chiral index of SWAP model



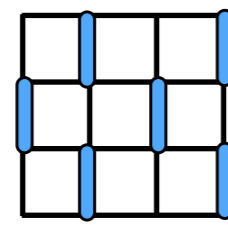
(1)



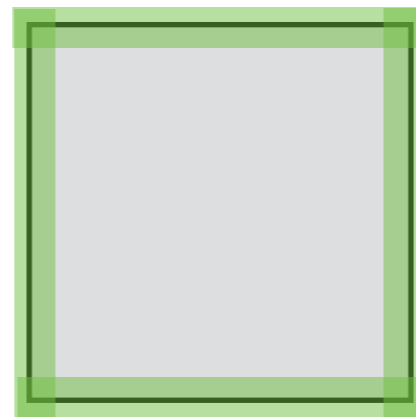
(2)



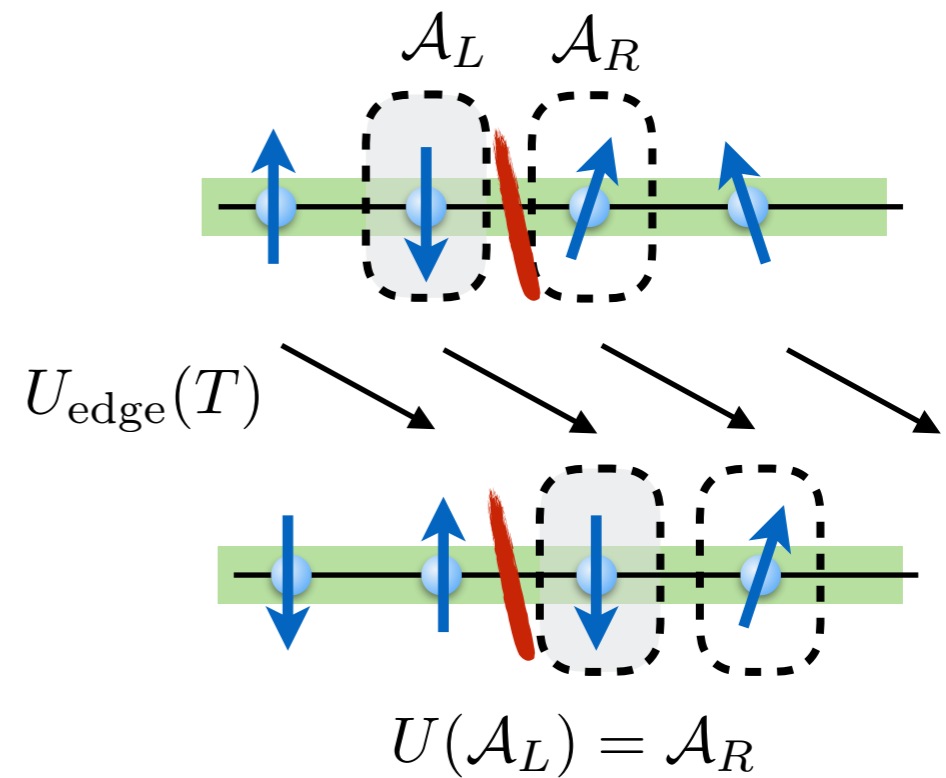
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(4)



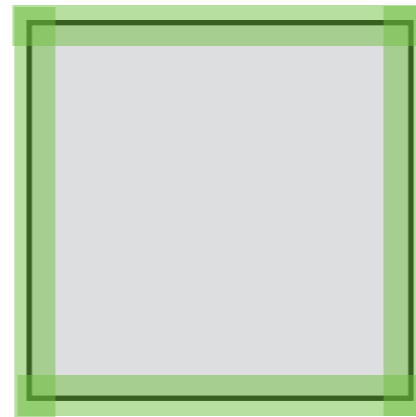
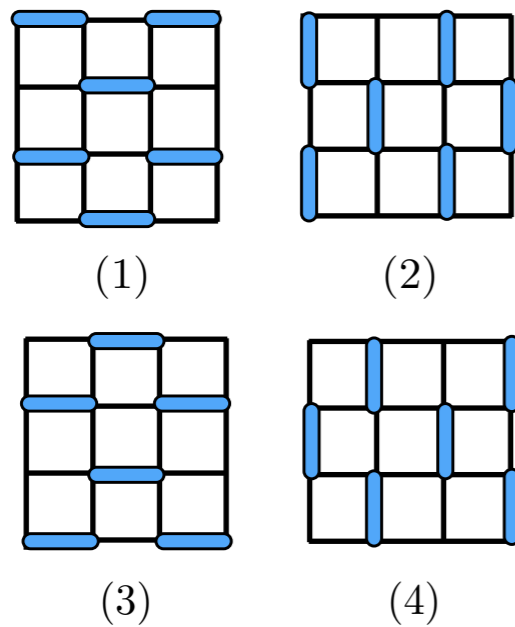
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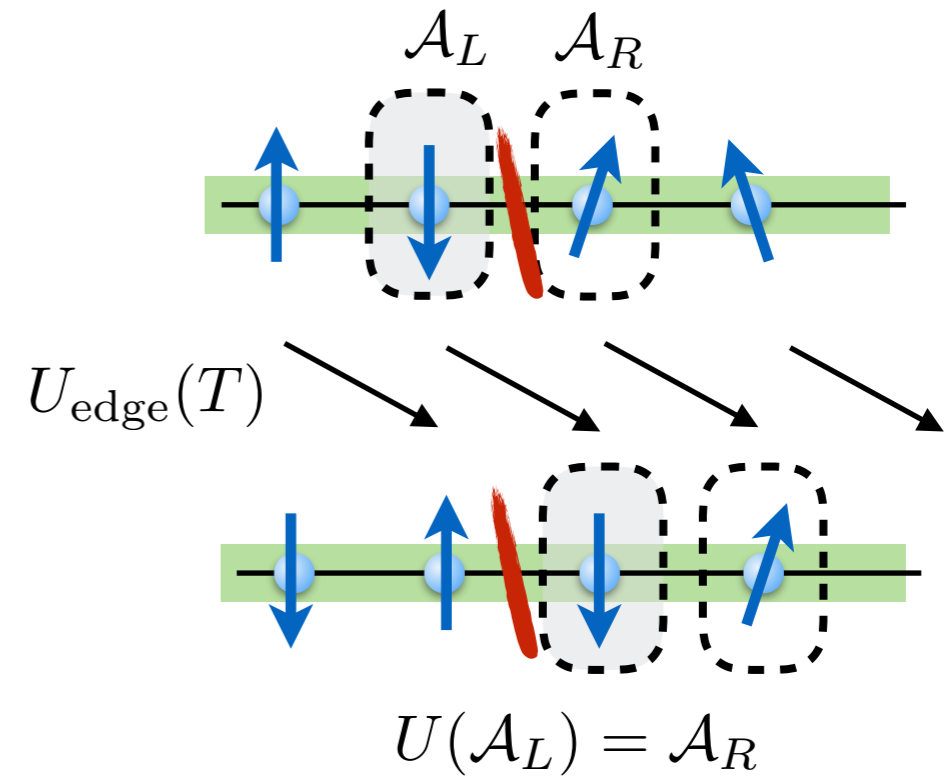
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Contrast to quantum Hall edge states

Dynamical chiral index of SWAP model



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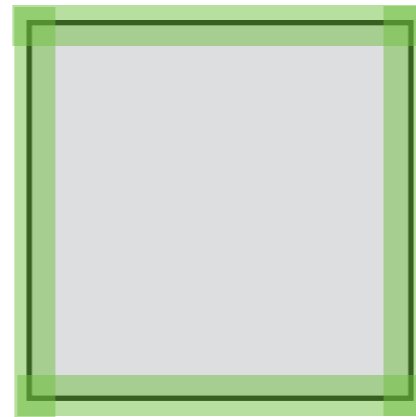
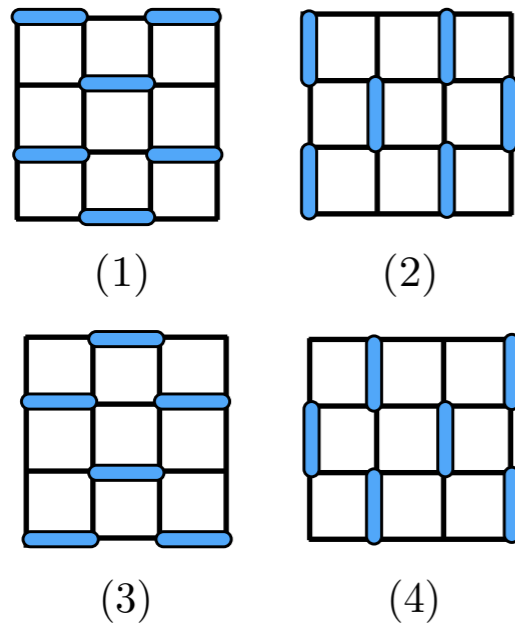


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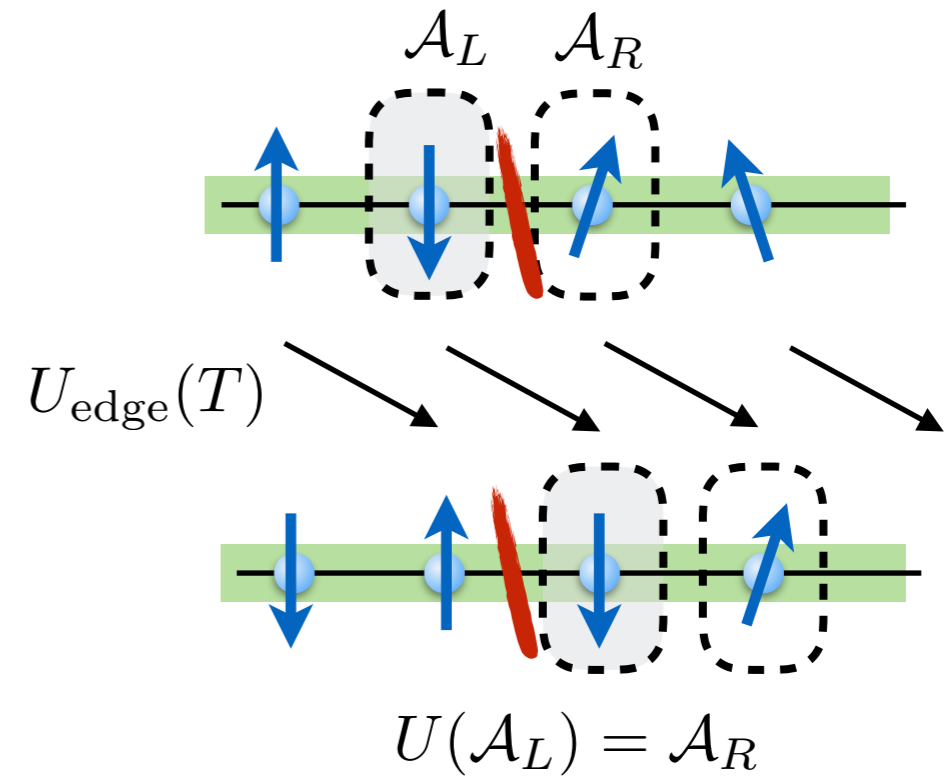
Contrast to quantum Hall edge states

- amount of information sent by QH edge not quantized (depends on v^*T — discrete pumping vs continuous flow)

Dynamical chiral index of SWAP model



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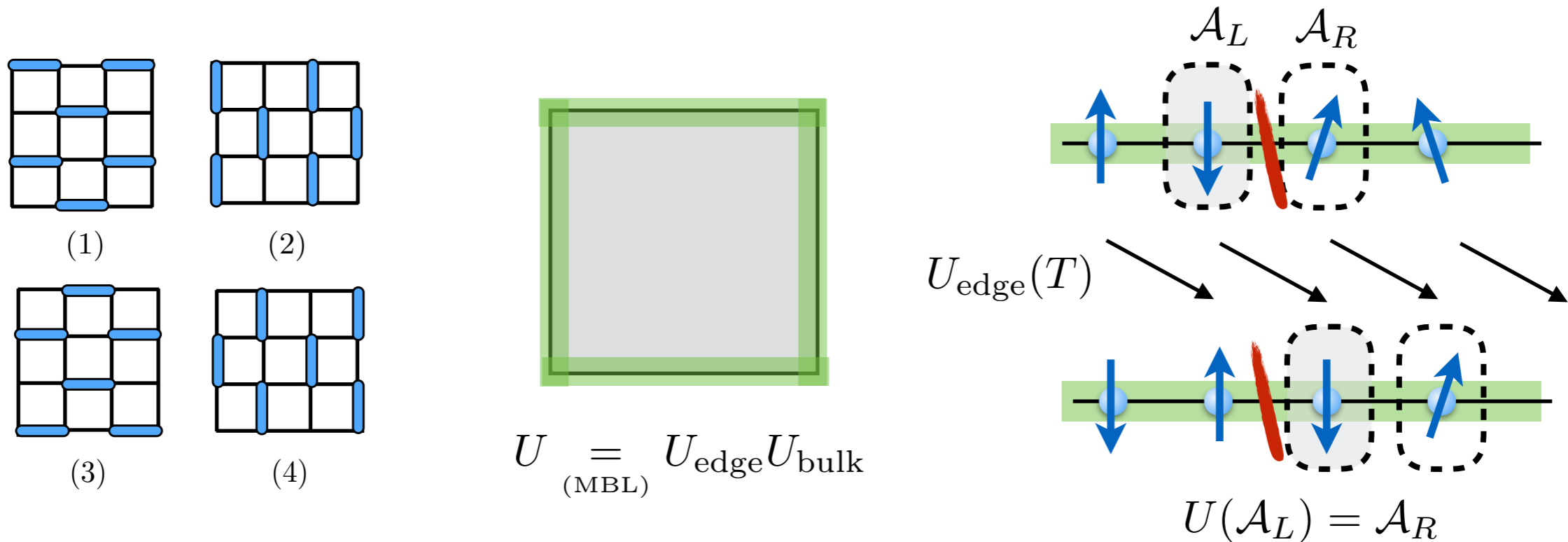


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Contrast to quantum Hall edge states

- amount of information sent by QH edge not quantized (depends on v^*T — discrete pumping vs continuous flow)
- QH edges don't sharply exist out of equilibrium

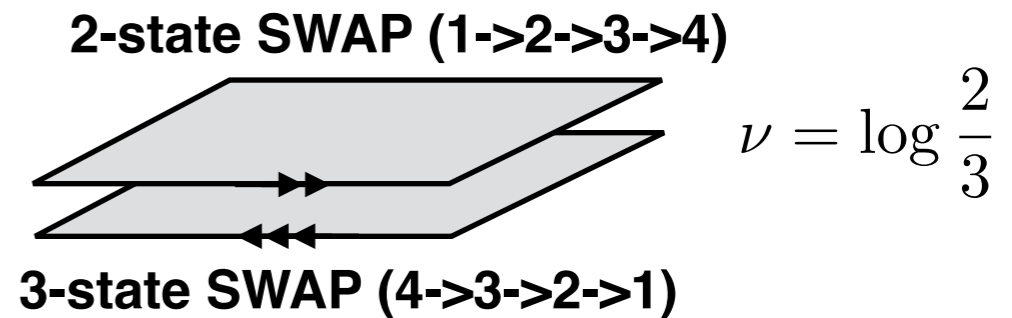
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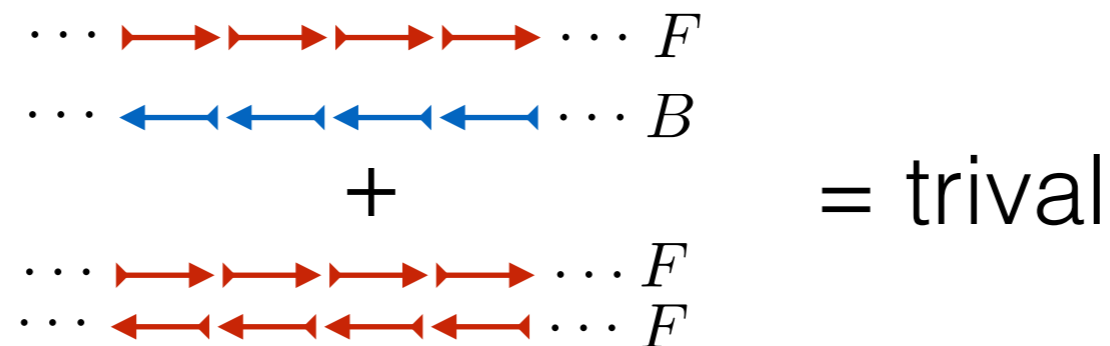
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Contrast to quantum Hall edge states

- amount of information sent by QH edge not quantized (depends on v^*T — discrete pumping vs continuous flow)
- QH edges don't sharply exist out of equilibrium
- Multiplicative rather than additive structure of index



Fermion chiral Floquet phases



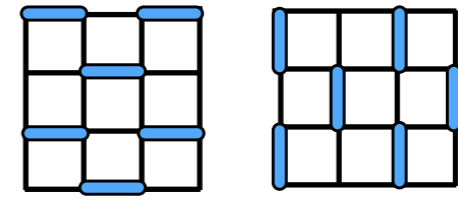
Fermion state equivalent to 2-state boson w/ $\nu_F = \log(2)$

- Can show: counter-propagating fermion and boson edges can be deformed to a trivial state
(requires edge-reconstruction by non-chiral modes — stable topological equivalence)
- Contrast to equilibrium:
8x minimal fermion chiral state = minimal Boson c.s.
- Remaining challenge: generalize formal index to deal with fermionic operator algebras

$$\nu_F = \log 2$$

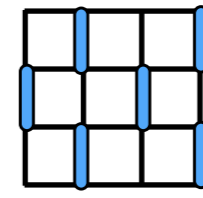
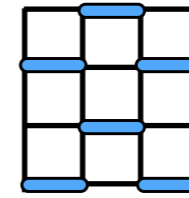
Fractional chiral Floquet phases?

Fractional chiral Floquet phases



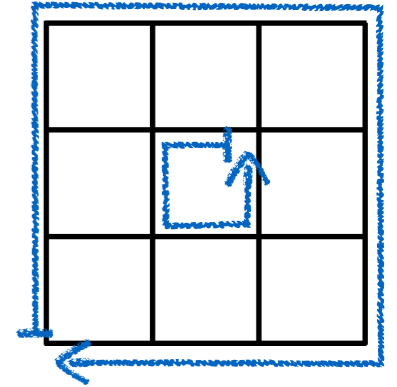
(1)

(2)



(3)

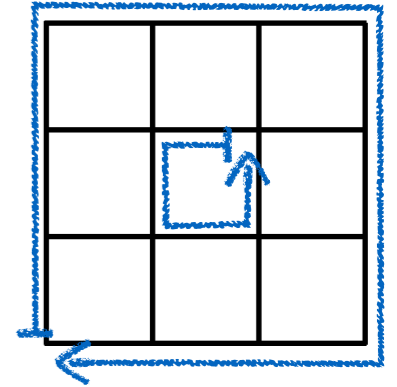
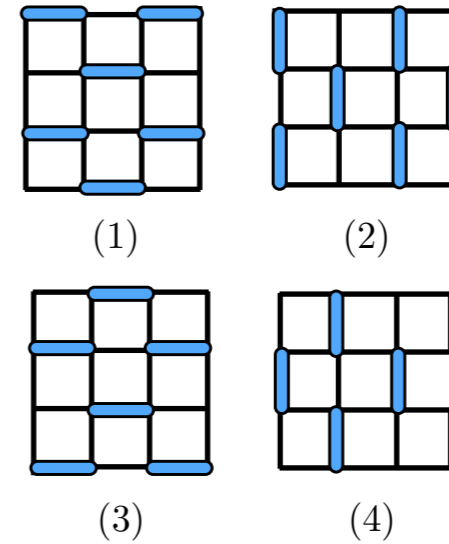
(4)



Fractional chiral Floquet phases

Non-fractional:

- $U(T) \sim 1$ (bulk),
chiral translation (boundary)
- Edge: pumps qudits with integer d



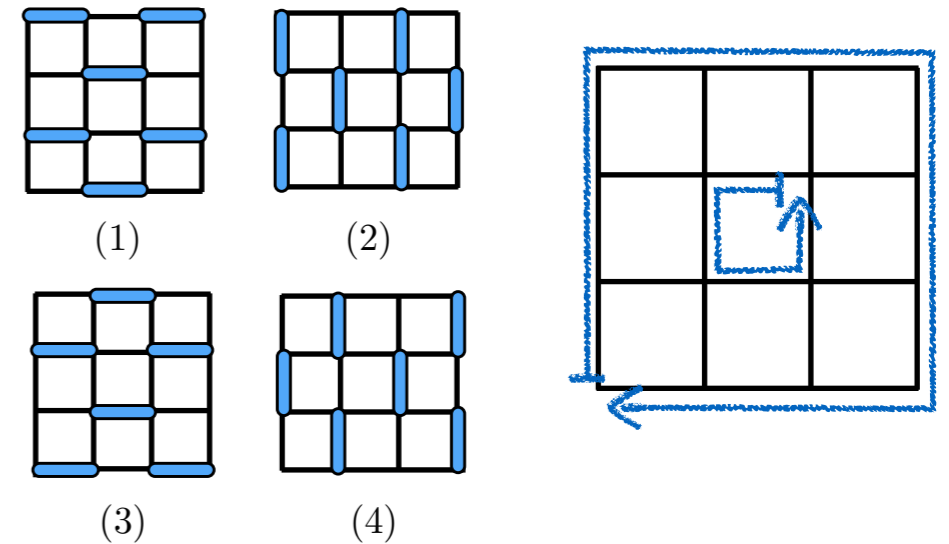
Fractional chiral Floquet phases

Non-fractional:

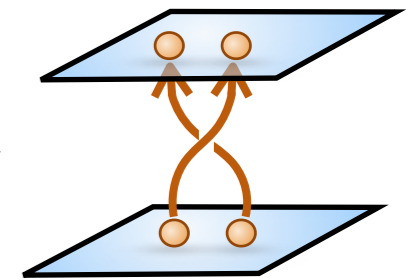
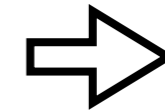
- $U(T) \sim 1$ (bulk),
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- Edge: pumps qudits with integer d

Fractional

- $U(T) \sim e^{-iH_{\text{TO}}}$ (bulk)
chiral translation of anyons (boundary)
- Edge: pumps fractional qudits with irrational d



SWAP

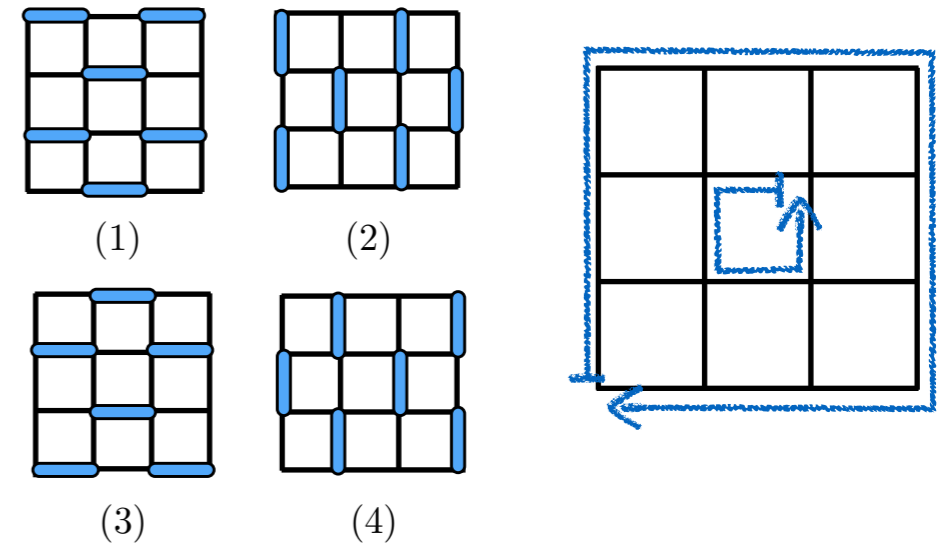


Braid

Fractional chiral Floquet phases

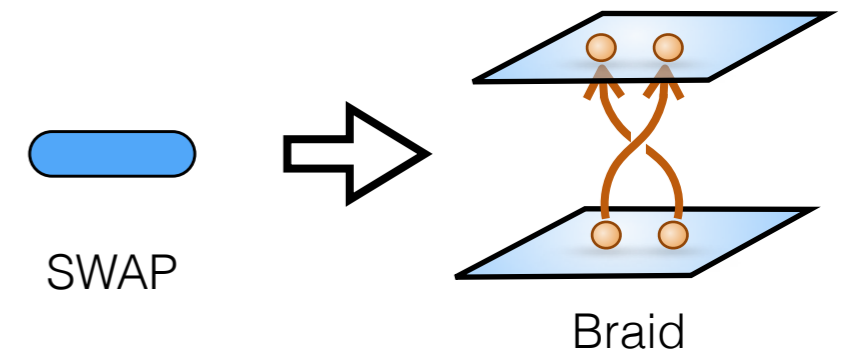
Non-fractional:

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Fractional

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chiral translation of anyons (boundary)
- Edge: pumps fractional qudits with irrational d



Localization: requires Abelian topological order

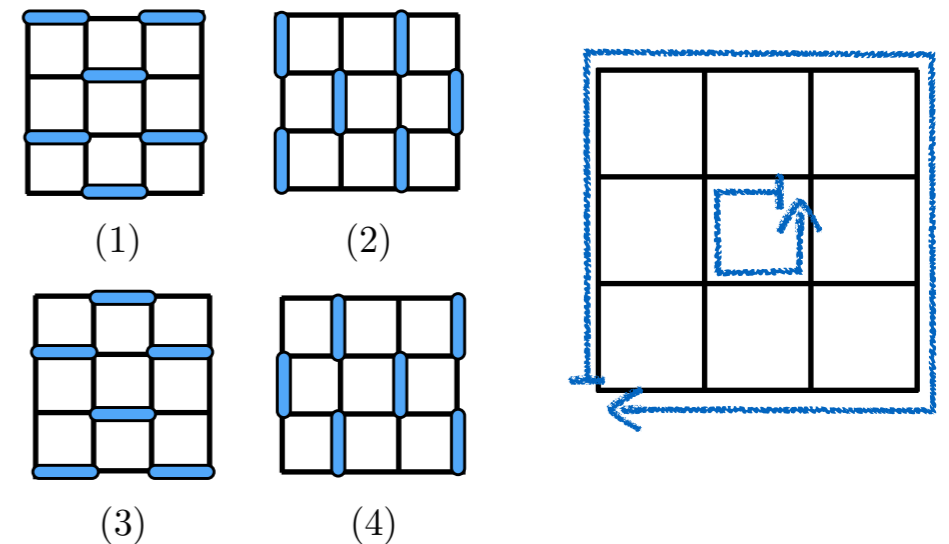
Potter, Vasseur arXiv '16

- Looks like $d=1$ only?
- No!: Can pump “sqrt” of an Abelian anyon

Fractional chiral Floquet phases

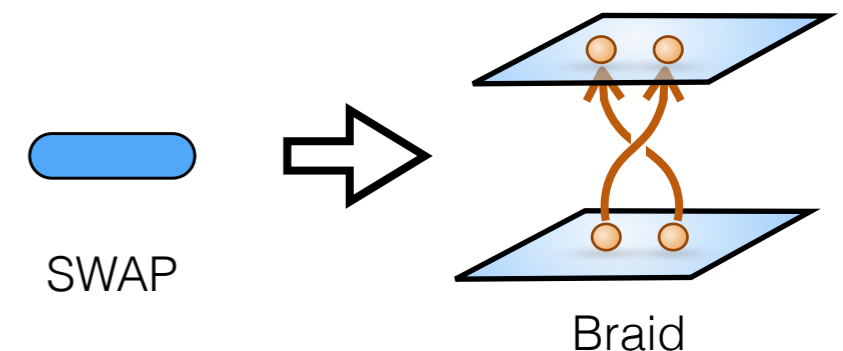
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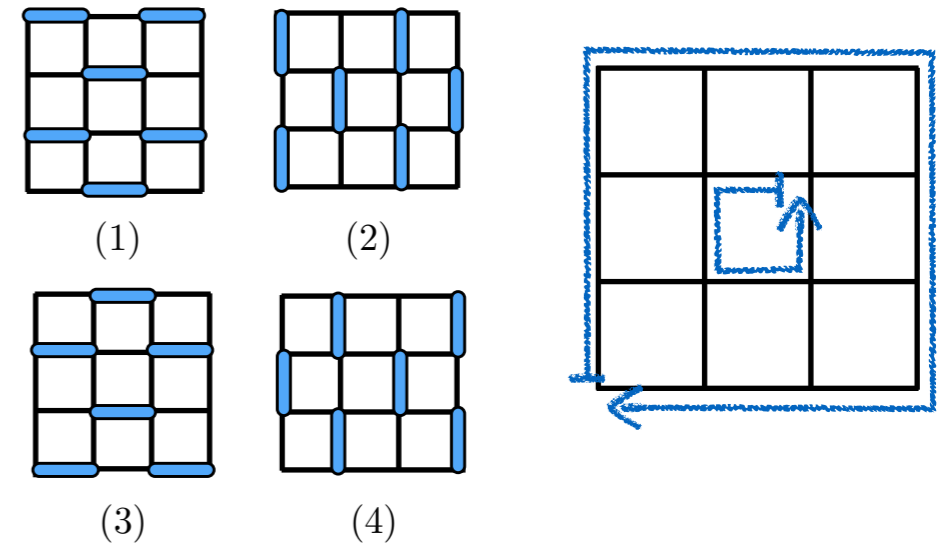
$$\mathbb{Z}_2 \text{ TO: Majorana, } d = \sqrt{2}, \nu = \frac{1}{2} \log 2$$

Po, Fidkowski, Morimoto, Vishwanath, ACP to appear

Fractional chiral Floquet phases

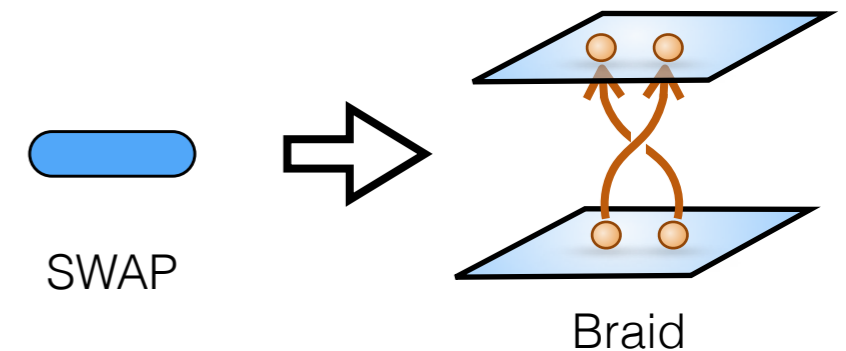
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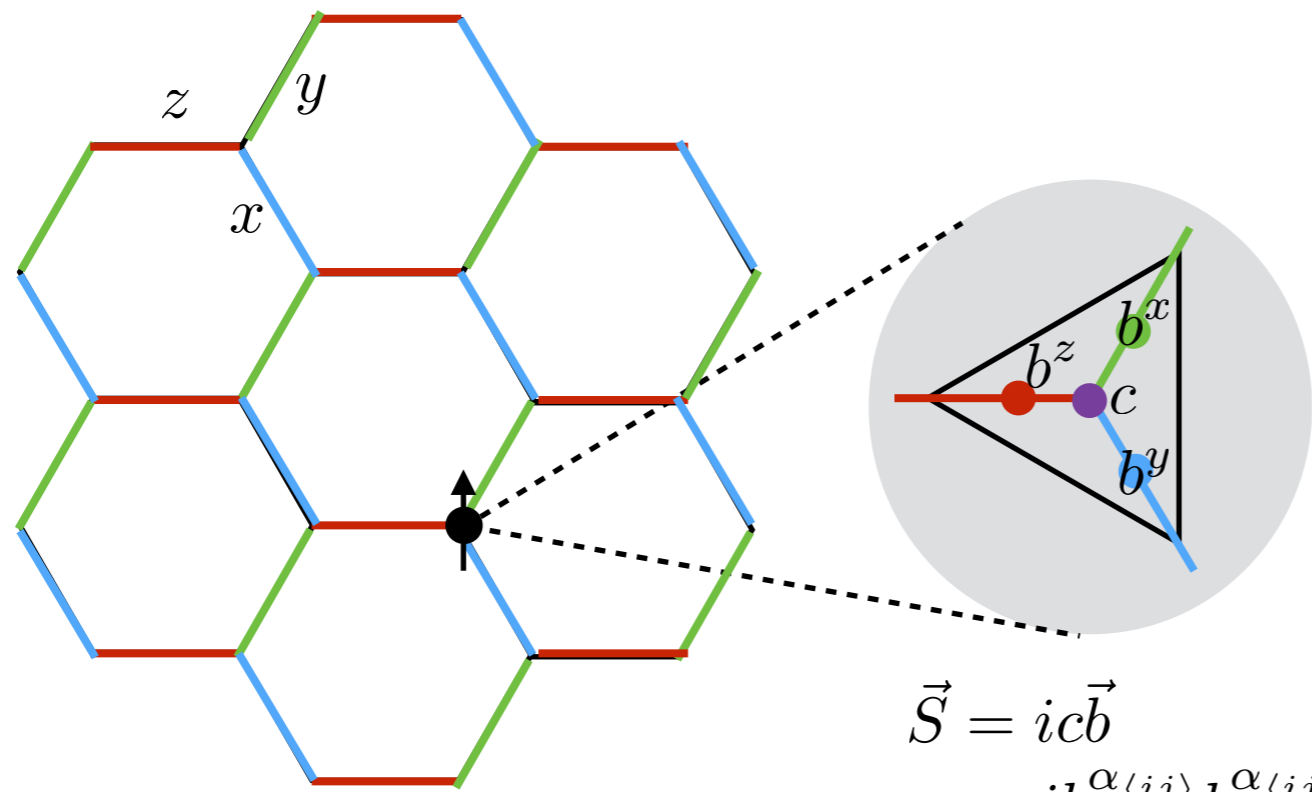
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$$\mathbb{Z}_2 \text{ TO: Majorana, } d = \sqrt{2}, \nu = \frac{1}{2} \log 2 \quad \mathbb{Z}_N \text{ TO: Parafermion, } d = \sqrt{N}, \nu = \frac{1}{2} \log N$$

\mathbf{Z}_2 Example - Driven Kitaev Honeycomb model



Kitaev '05

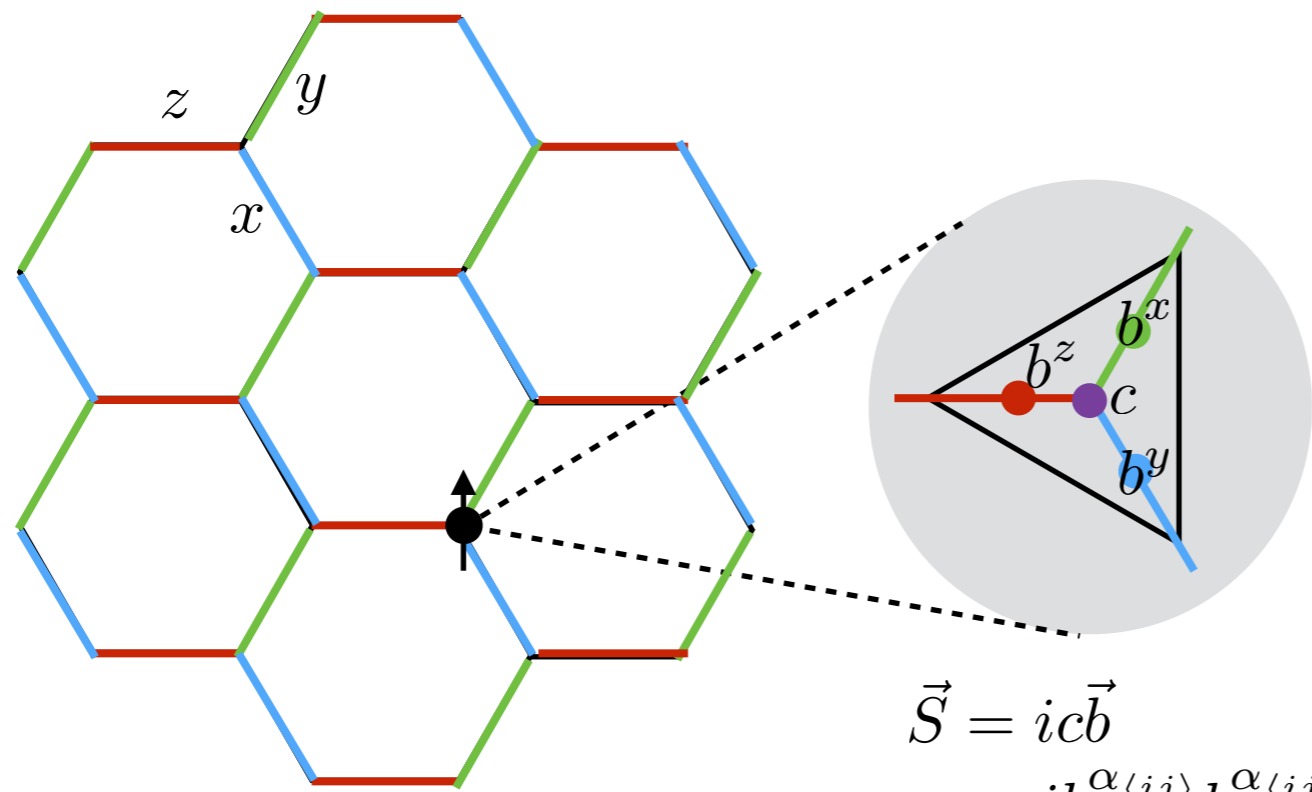
$$\vec{S} = ic\vec{b}$$

$$\sigma_{ij} = ib_i^{\alpha\langle ij\rangle} b_j^{\alpha\langle ij\rangle}$$

$$S_i^\alpha S_j^\alpha = ic_i \sigma_{ij} c_j$$

- $H_1 = \frac{\pi}{2} \sum_{\langle ij \rangle \in X} S_i^x S_j^x$
- $H_2 = \frac{\pi}{2} \sum_{\langle ij \rangle \in X} S_i^y S_j^y$
- $H_3 = \frac{\pi}{2} \sum_{\langle ij \rangle \in X} S_i^z S_j^z$

\mathbf{Z}_2 Example - Driven Kitaev Honeycomb model



Kitaev '05

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- $H_1 = \frac{\pi}{2} \sum_{\langle ij \rangle \in X} S_i^x S_j^x$
- $H_2 = \frac{\pi}{2} \sum_{\langle ij \rangle \in Y} S_i^y S_j^y$
- $H_3 = \frac{\pi}{2} \sum_{\langle ij \rangle \in Z} S_i^z S_j^z$

Bulk:

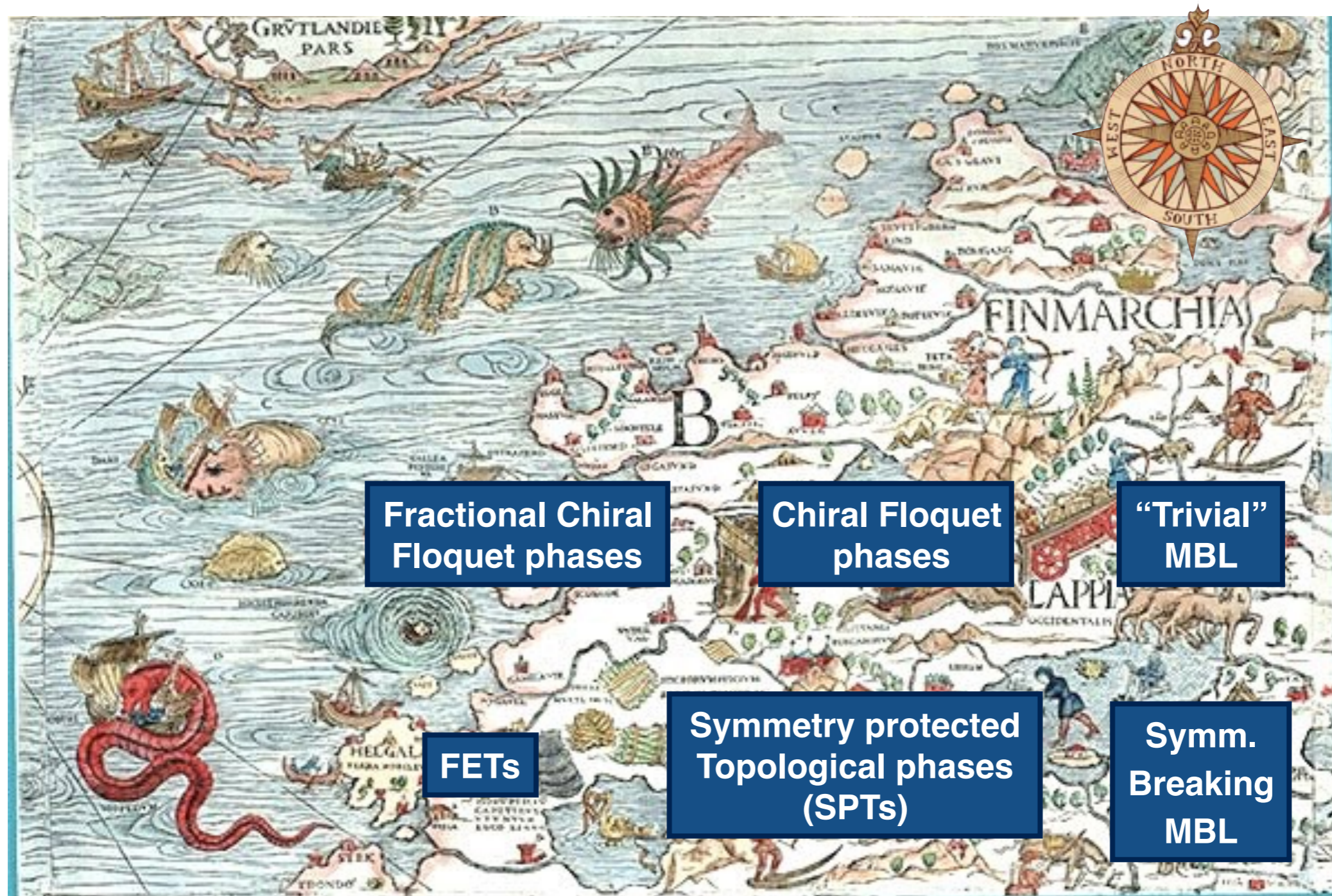
- Bulk evolution = evolution w/r.t. static Hamiltonian w/ \mathbf{Z}_2 TO

Edge:

- c's pumped around edge chirally
- 2x Majorana edge = 1x fermion edge = spin-1/2 boson
- Fractional chiral index: $\nu_M = \frac{1}{2} \log 2$

Po, Fidkowski, Morimoto, Vishwanath, ACP to appear

“World Map” of (Driven) quantum matter



←
(Excited-state) Entanglement