

# Topological Quantum Chemistry

Maia G. Vergniory

**ikerbasque**  
Basque Foundation for Science

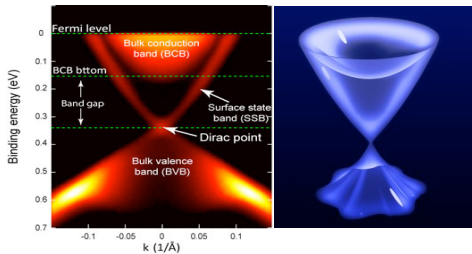


**dipc**  
Donostia International Physics Center

KITP

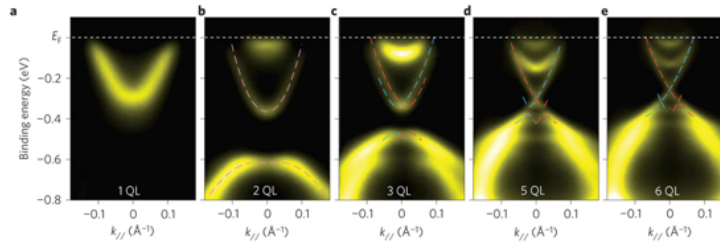
Topological Quantum Matter: Concepts and Realizations

# Discovery of topological materials



Hsieh Nature (2008)  
Zhang Nat Phys (2009)  
Xia Nat Phys (2008)

**3D TIs** theory + exp  
protected by **TRS** :  $\text{Bi}_2\text{Se}_3$



Chang Science (2013)

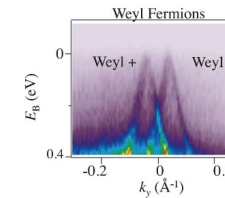
QAHE semimetals

Weng PRX (2015)  
Xu Science (2015)

**Weyl**

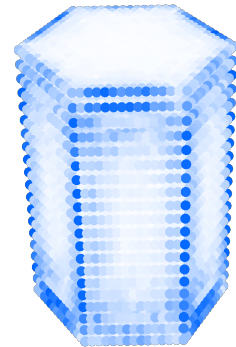
## Discovery of a Weyl fermion semimetal and topological Fermi a

Su Yang Xu,<sup>1,2a</sup> Ilya Belopolski,<sup>1a</sup> Nasser Aldoust,<sup>1,2a</sup> Madhab Neupane,<sup>1,2a</sup> Guang Bian,<sup>4</sup> Chenglong Zhang,<sup>4</sup> Raman Sankar,<sup>2</sup> Guoqing Chang,<sup>6,7</sup> Zhi-Jun Yu,<sup>1</sup> Chi-Cheng Lee,<sup>6,7</sup> Shin-Ming Huang,<sup>6,7</sup> Hao Zheng,<sup>7</sup> Jie Ma,<sup>8</sup> Daniel S. Sanchez,<sup>1</sup> BaoKai Wang,<sup>6,7,8</sup> Arun Bansil,<sup>9</sup> Fangcheng Zhou,<sup>3</sup> Pavel P. Shihayev,<sup>1,10</sup> Hsin Li Shuang Jia,<sup>1,11</sup> M. Zahid Hasan<sup>1,2†</sup>



Bradlyn (2016)

New Nodal Lines  
Fermions

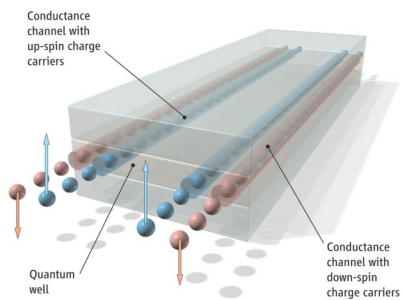


Bzdušek (2016)

2006	2007	2008	2009	2010	2011	2012	2013	2014	2015	2016	2017	2018
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## Prediction HgTe 2D TI

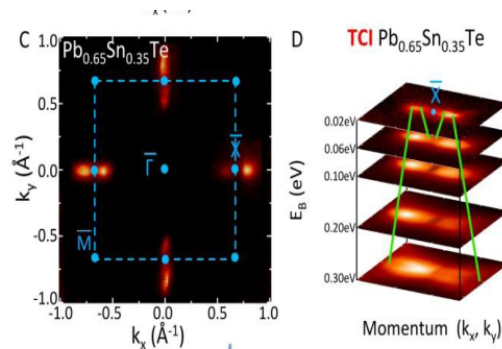
Bernevig Science (2006)  
König Science (2007)



Kane & Mele, PRL (2015)

## Mirror Chern insulators

Hsieh Nat Comm (2012)  
Tanaka Nat Phys (2012)



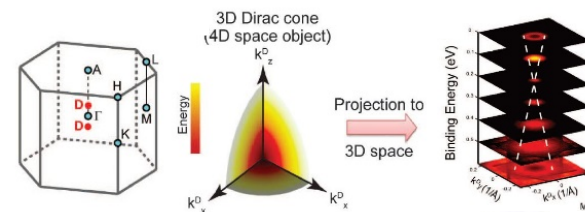
## Type II Weyls

Soluyanov (2015)

## Dirac semimetals

## Discovery of a Three-Dimensional Topological Dirac Semimetal, Na<sub>3</sub>Bi

Z. K. Liu,<sup>1\*</sup> B. Zhou,<sup>2,3\*</sup> Y. Zhang,<sup>3</sup> Z. J. Wang,<sup>4</sup> H. M. Weng,<sup>4,5</sup> D. Prabhakaran,<sup>2</sup> S.-K. Mo,<sup>3</sup> Z. X. Shen,<sup>1</sup> Z. Fang,<sup>4,5</sup> X. Dai,<sup>4,5</sup> Z. Hussain,<sup>3</sup> Y. L. Chen<sup>2,6†</sup>

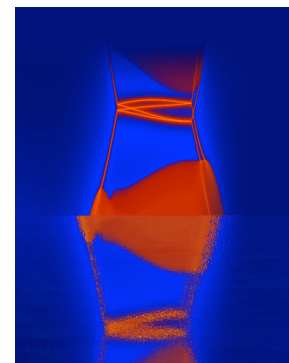


## High Order TIs

Schindler (2018)

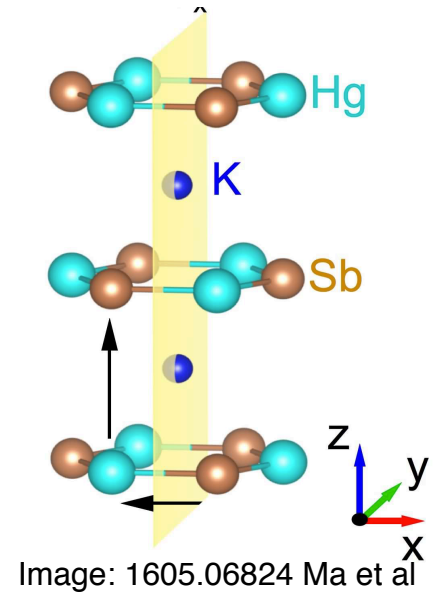
## Non-symmorphic TIs

Alexandrinata (2016)



➔ Recall: a space group is a set of symmetries that defines a crystal structure in 3D

- 230 Space-Groups
- Ingredients:
- unit lattice translations ( $\mathbb{Z}^3$ )
  - point group operations (rotations, reflections)
  - non-symmorphic (screw, glide)
  - orbitals
  - atoms in some lattice positions



How do we go from real space orbitals sitting on lattice sites to electronic bands (without a Hamiltonian)?



### ELEMENTARY BAND REPRESENTATIONS

# Band as Representations

**Band Representation (BR):** set of **bands linked to a localized orbital**  
*(respecting all the crystal symmetries and TRS)*

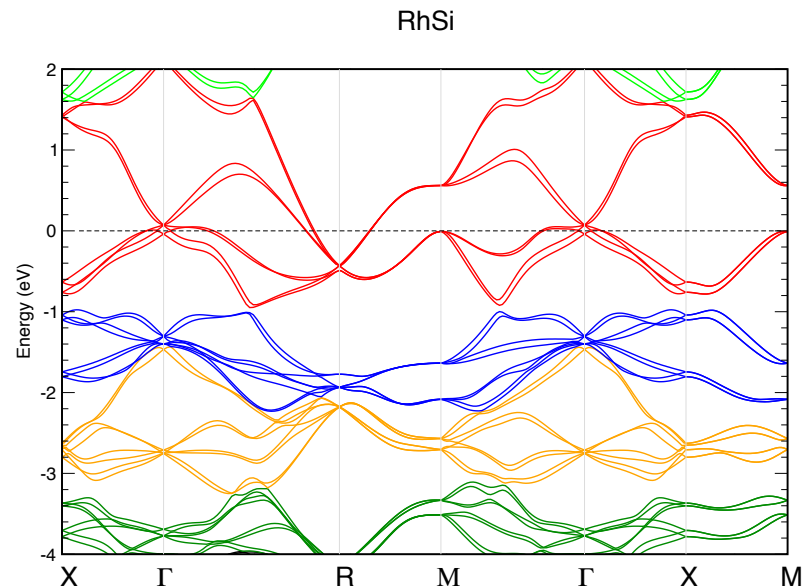


**Elementary BR:** smallest set of bands cannot be decomposed in elementary bands

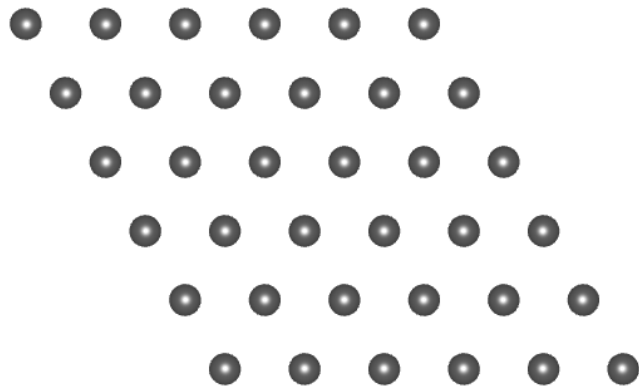
**Physical Elementary BR:** when EBR also respects TR symmetry

**Composite BR:** A BR which is not elementary is a “composite”

(P)EBRs are connected along the BZ

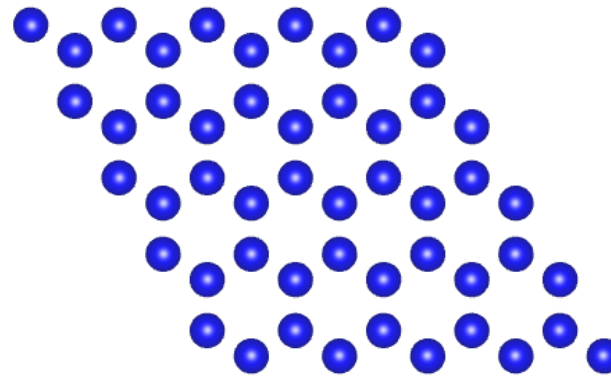


# Induction of a (P)EBR



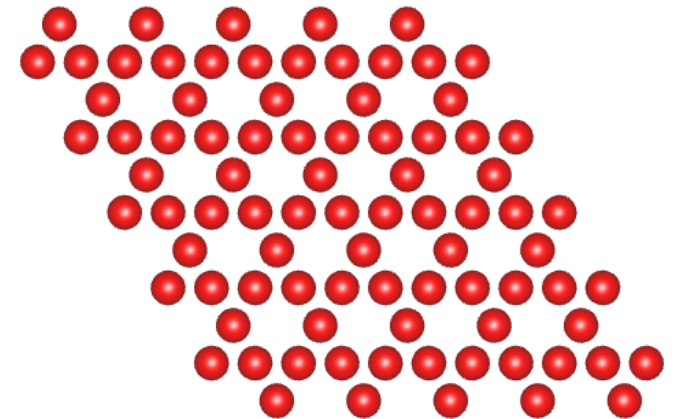
1 atom/unit cell  
(triangular)

1a



2 atoms/unit cell  
(honeycomb)

2b



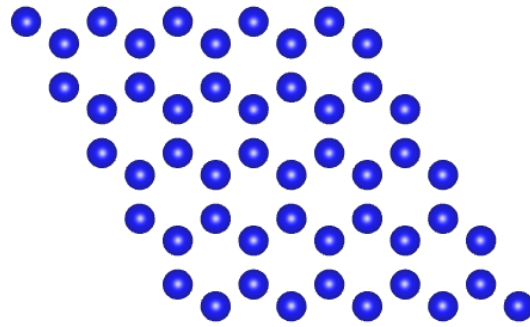
3 atoms/unit cell  
(kagome)

3c

- Within the same SG many ways to arrange atoms
- Each arrangement determines different representations

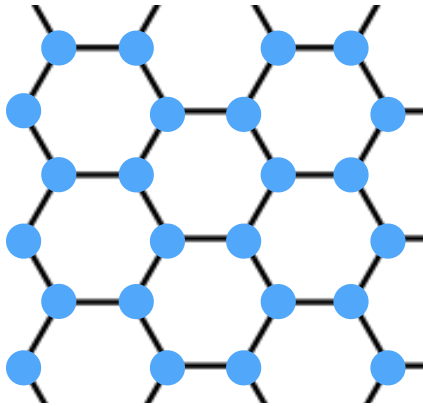
# Induction of a (P)EBR

- Within the same lattice, different orbitals

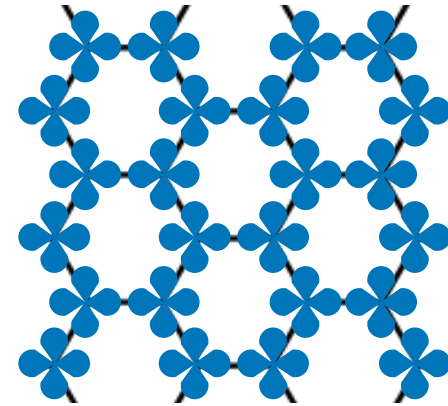


2 atoms/unit cell

**s (or  $p_z$ ) orbitals**

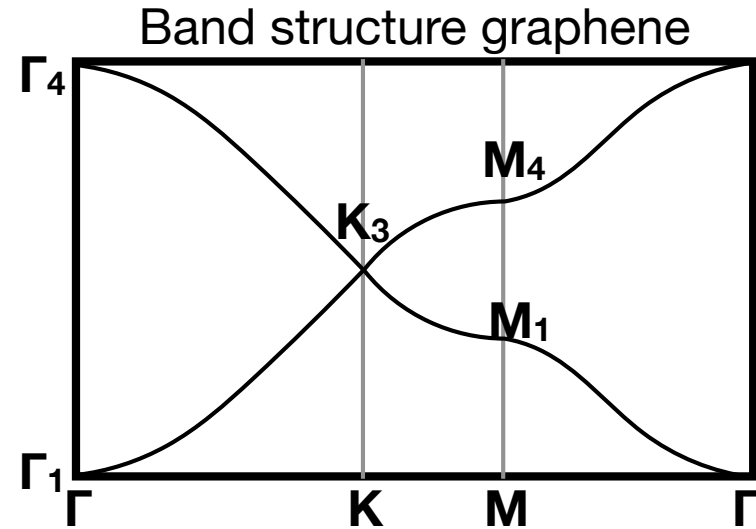
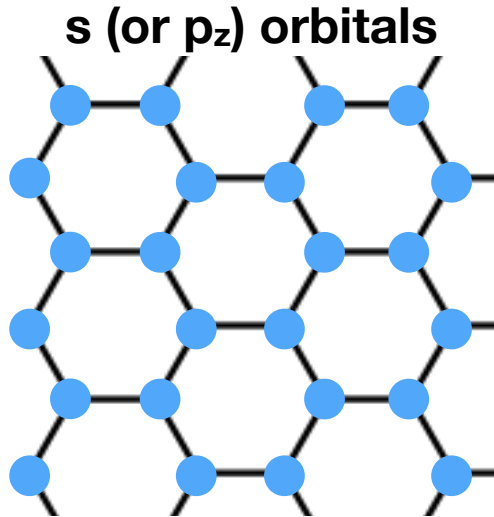


**$p_x$  and  $p_y$  orbitals**

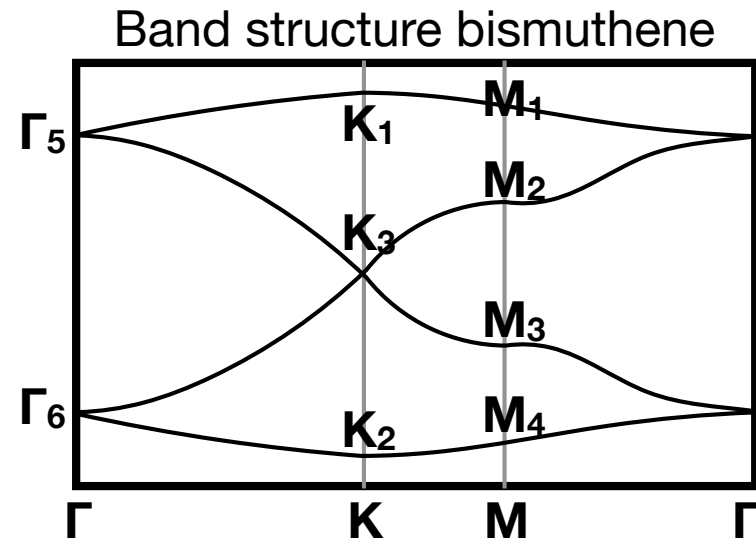
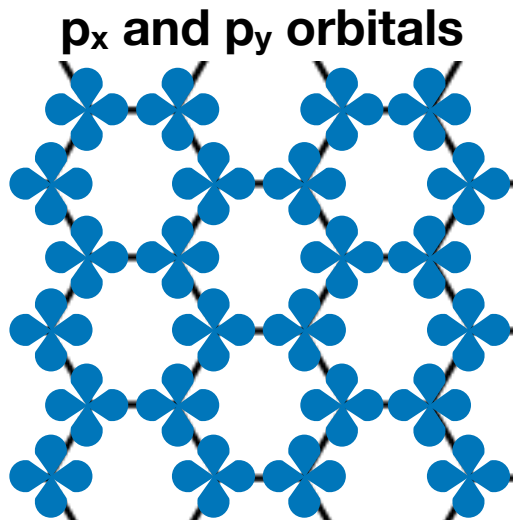


# Induction of a (P)EBR

Site-symmetry group,  $G_q$ , leaves  $q$  invariant

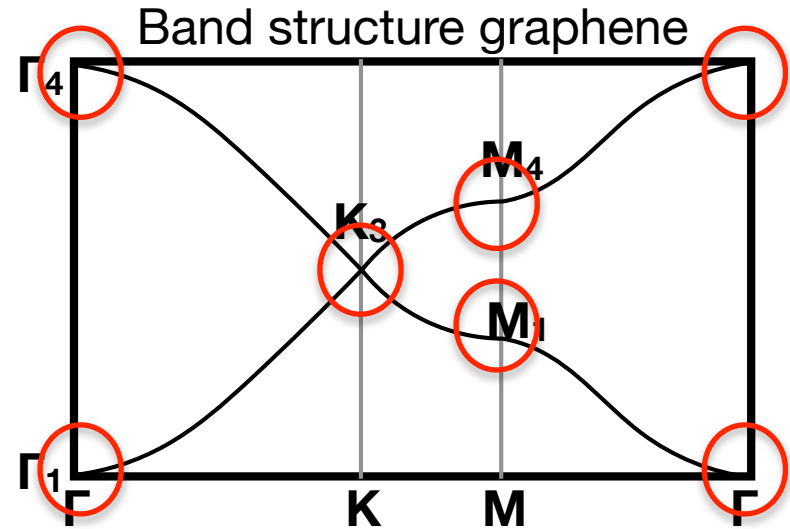
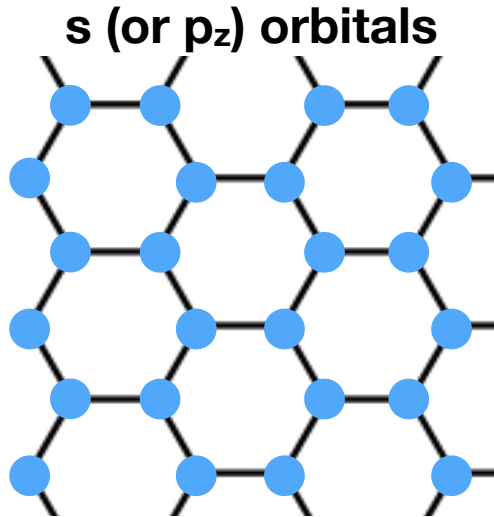


**Real space vs momentum space**

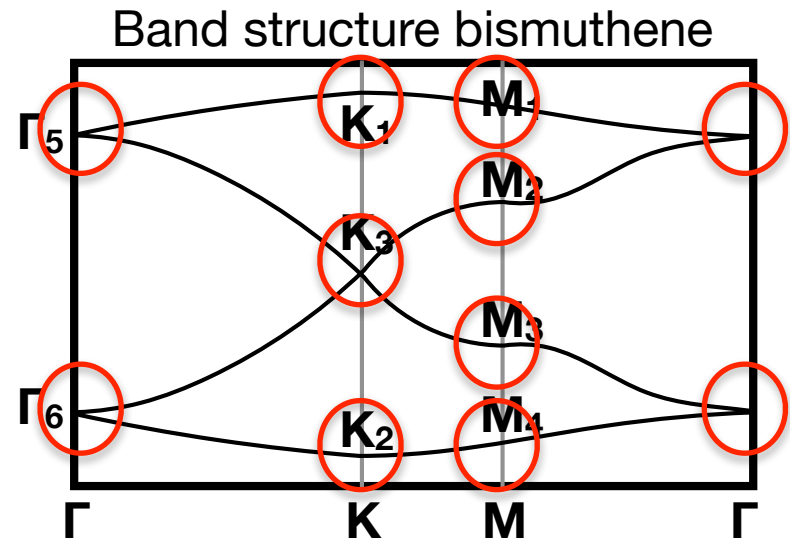
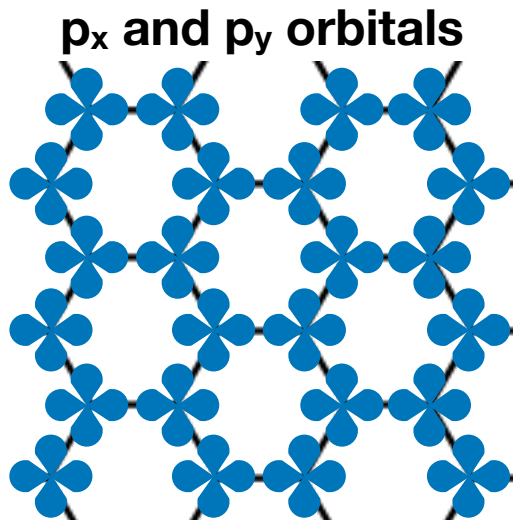


# Induction of a (P)EBR

Each arrangement/orbital determines symmetry representations in the Brillouin zone



**Real space vs momentum space**

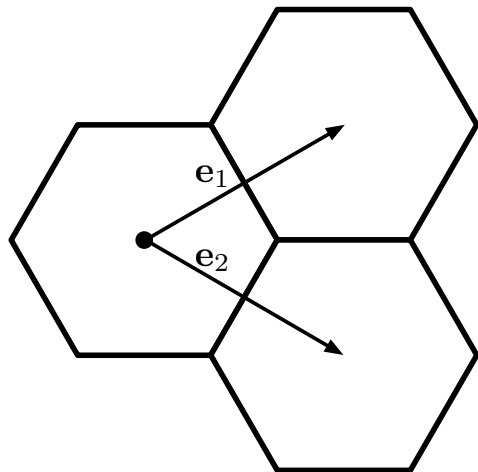




# Induction of a (P)EBR: Example of the honeycomb lattice

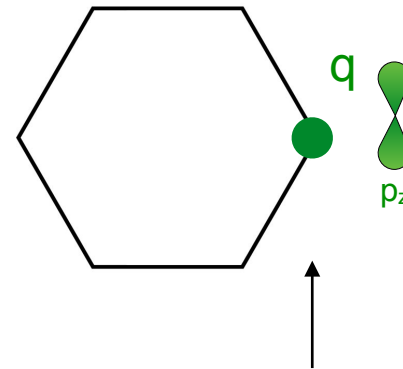
Lets consider the generators of 2D P6mm:  $\{C_2, C_3, m_{1\bar{1}}\}$

Lattice vectors:



$$\begin{cases} e_1 = \sqrt{3}/2x + 1/2y \\ e_2 = \sqrt{3}/2x - 1/2y \end{cases}$$

Lattice site: Wyckoff 2b, spinfull  $p_z$



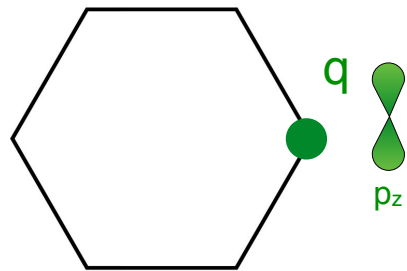
Site-symmetry group,  $G_q$ , leaves  $q$  invariant

Coset decomposition of a Space Group :

$$G = \bigcup_{\alpha=1}^n (g_{\alpha}) (G_q \times \mathbf{Z}^3), \quad g_{\alpha} \notin G_q$$

# Induction of a (P)EBR: Example of the honeycomb lattice

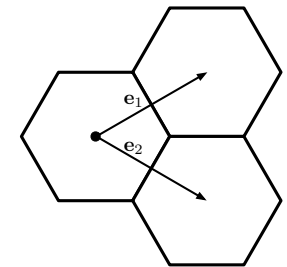
Consider one lattice site:



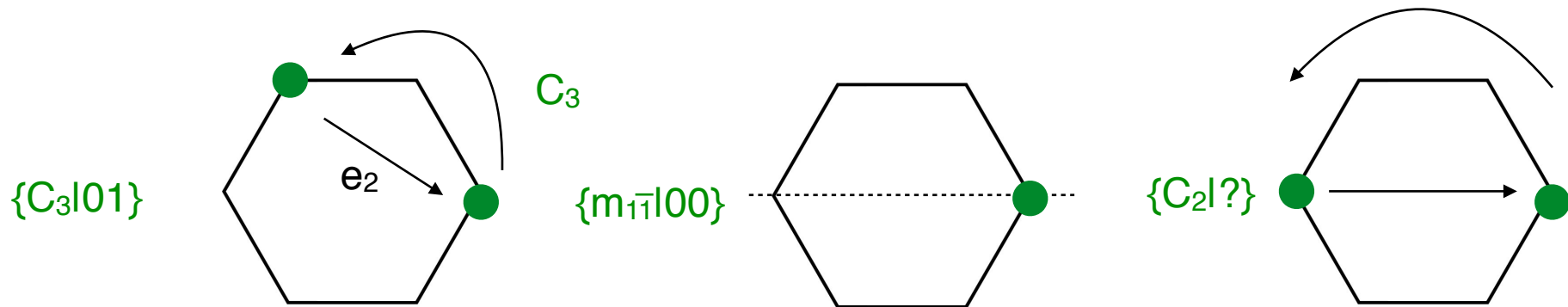
$$G = \bigcup_{\alpha} (g_{\alpha}) \quad (G_q \rtimes \mathbf{Z}^3)$$

(2)                      (1)

$\{C_2, C_3, m_{1\bar{1}}\}$

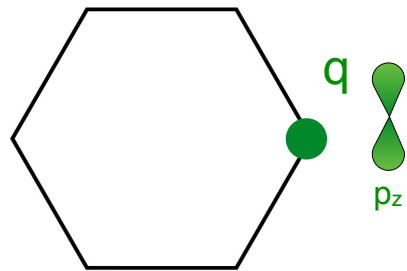


- (1) Site-symmetry group,  $G_q$ , leaves  $q$  invariant  $\{C_3|01\}, \{m_{1\bar{1}}|00\} \approx C_{3v}$   
 → Orbitals at  $q$  transform under a rep,  $\rho$ , of  $G_q$



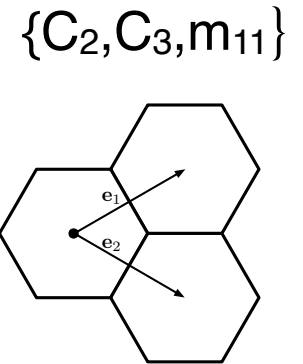
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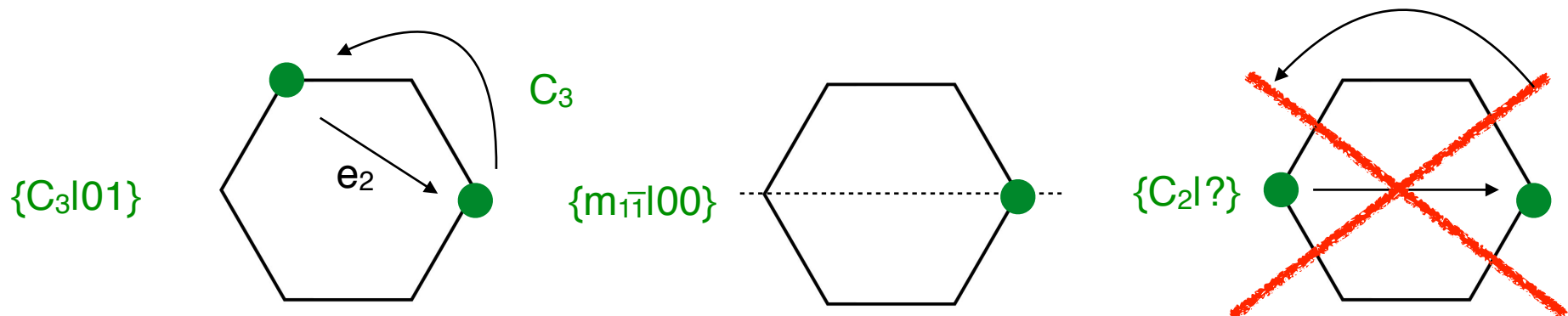


$$G = \bigcup_{\alpha} (g_{\alpha}) \quad (G_q \rtimes \mathbf{Z}^3)$$

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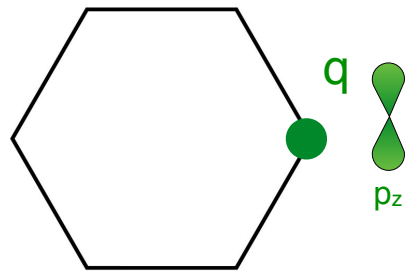


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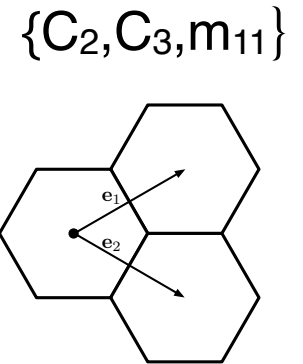
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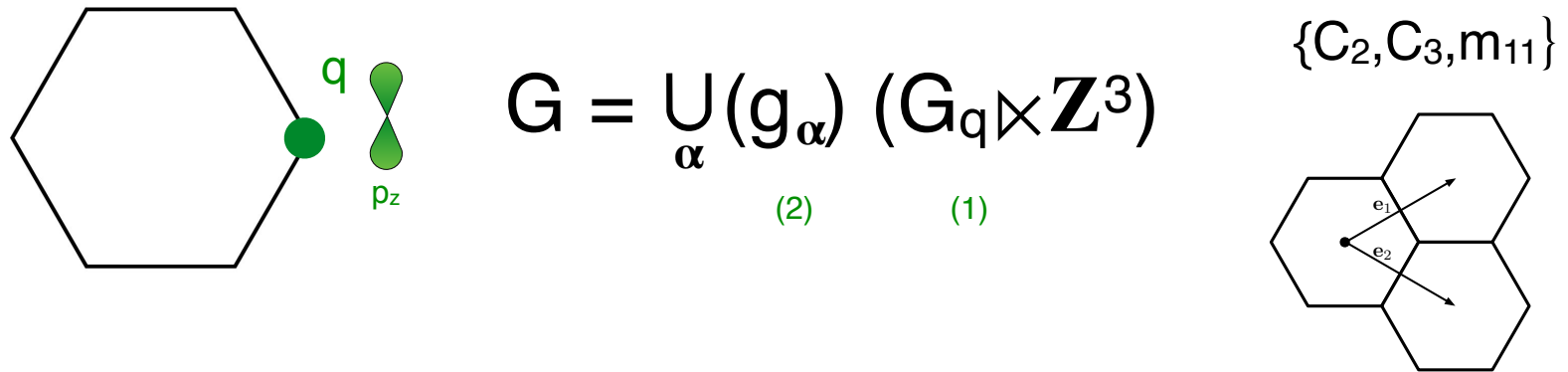
- (1) Site-symmetry group,  $G_q$ , leaves  $q$  invariant  $\{C_3|01\}, \{m_{11}|\bar{1}00\} \approx C_{3v}$   
 → Orbitals at  $q$  transform under a rep,  $\rho$ , of  $G_q$

Rep	E	$C_3$	$M$	$\bar{E}$
→ $\bar{\Gamma}_6$	2	1	0	-2

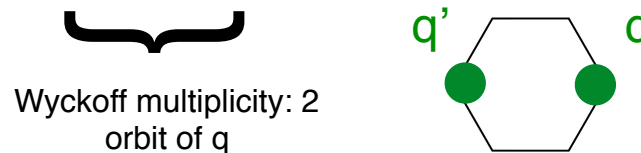
Character table for the double-valued representation of  $C_{3v}$

# Induction of a (P)EBR: Example of the honeycomb lattice

Consider one lattice site:



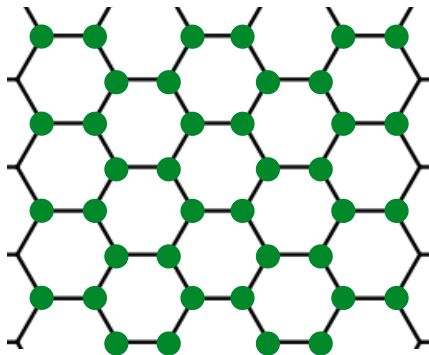
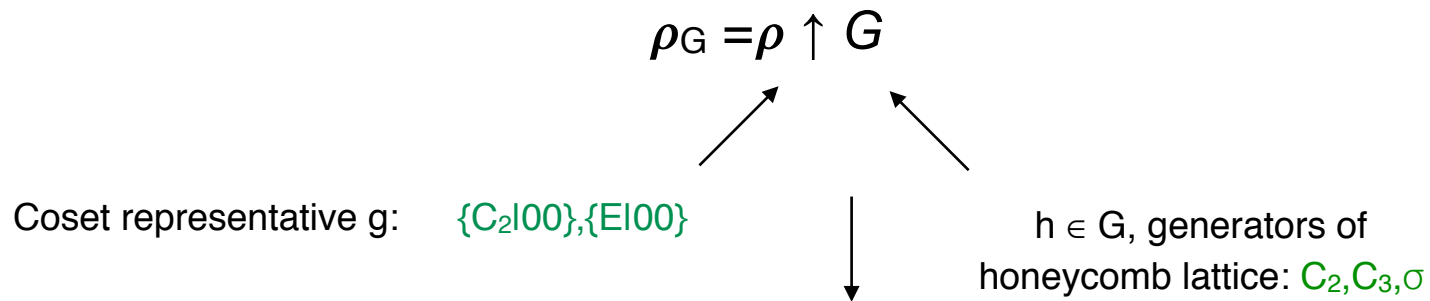
- (1) Site-symmetry group,  $G_q$ , leaves  $q$  invariant  $\{C_3|01\}, \{m_{1\bar{1}}|00\} \approx C_{3v}$   
 → Orbitals at  $q$  transform under a rep,  $\rho$ , of  $G_q$
- (2) Elements of space group  $g \notin G_q$  (coset representatives) move sites in an orbit “Wyckoff position”  $\{C_2|00\}, \{E|00\}$



# Induction of a (P)EBR: Example of the honeycomb lattice

$\bar{\Gamma}_6$  induced in  $C_{6v}$

electron bands sitting at p<sub>z</sub> orbitals in  
Wyckoff 2b in wall paper group 17



$$\rho_{i_{\alpha}, j_{\beta}}(\mathbf{h}) = \rho_{ij}(\mathbf{g}_{\alpha\beta})$$

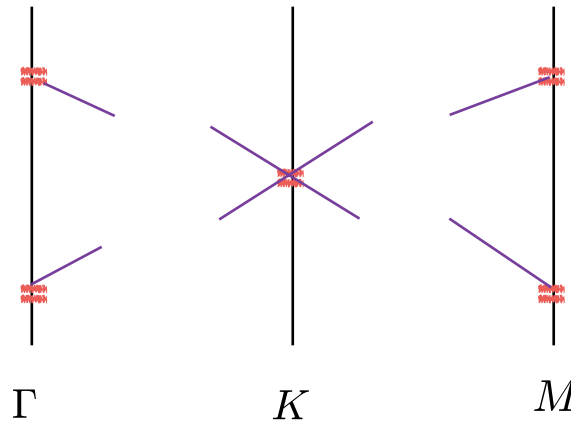
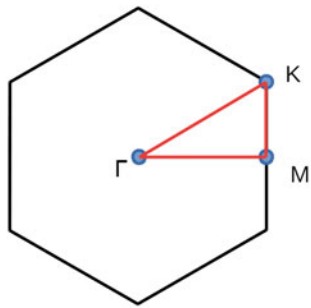
$$\mathbf{g}_{\alpha\beta} = \mathbf{g}_{\alpha}^{-1} \{E|t_{\alpha\beta}\} \mathbf{h} \mathbf{g}_{\beta}$$

$$\rho_G^k(\mathbf{h}) = e^{-i(\mathbf{k} \cdot \mathbf{t}_{\alpha\beta})} \rho_{ij}(\mathbf{g}_{\alpha\beta})$$

dimension of this band representations = connectivity in the Brillouin zone

# Subduction in momentum space

$$(\rho \uparrow G) \downarrow G_k$$



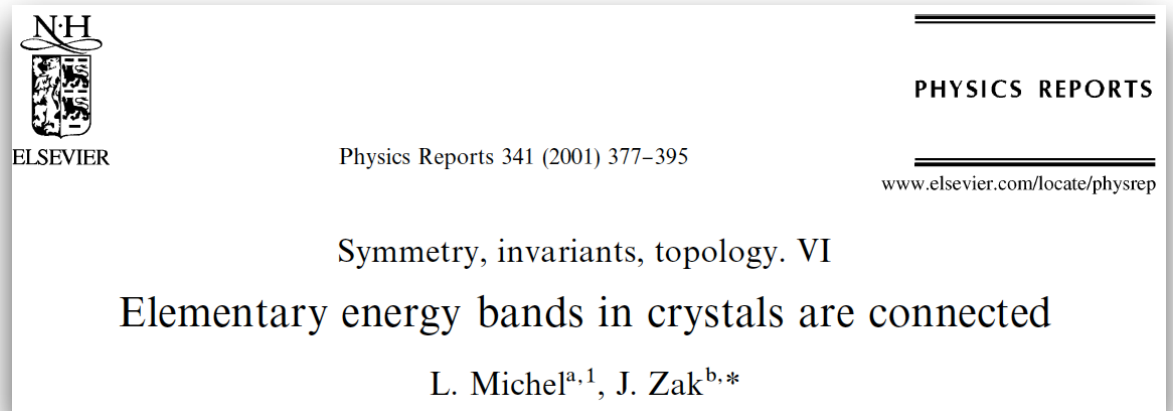
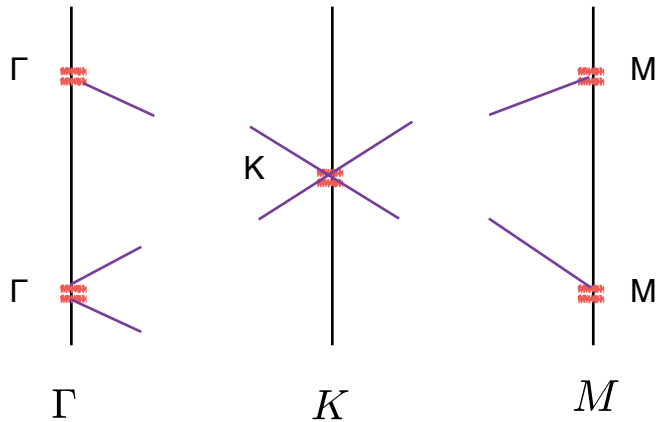
$$\rho_G^\Gamma = \bar{\Gamma}_7 \oplus \bar{\Gamma}_8$$

**Table 1.5** Table of characters of the group  $C_{6v}$

$C_{6v}$	$E$	$C_3^\pm$	$C_2, \bar{C}_2$	$C_6^\pm$	$m_{11}$	$m_{1\bar{1}}$	$\bar{E}$	$\bar{C}_3^\pm$	$\bar{C}_6^\pm$
$\rho_G^\Gamma$	4	2	0	0	0	0	-4	-2	0
$\bar{\Gamma}_7$	2	1	0	$-\sqrt{3}$	0	0	-2	-1	$\sqrt{3}$
$\bar{\Gamma}_8$	2	1	0	$\sqrt{3}$	0	0	-2	-1	$-\sqrt{3}$
$\bar{\Gamma}_9$	2	-2	0	0	0	0	-2	2	0

# ATOMIC LIMIT

- Restricting to the little group at  $k$  to find irreps at each  $k$  point (subduction) -> **all bands connected**
- EBR is defined by a **maximal Wyckoff position and the irreps in real space**



## ○ Topology?

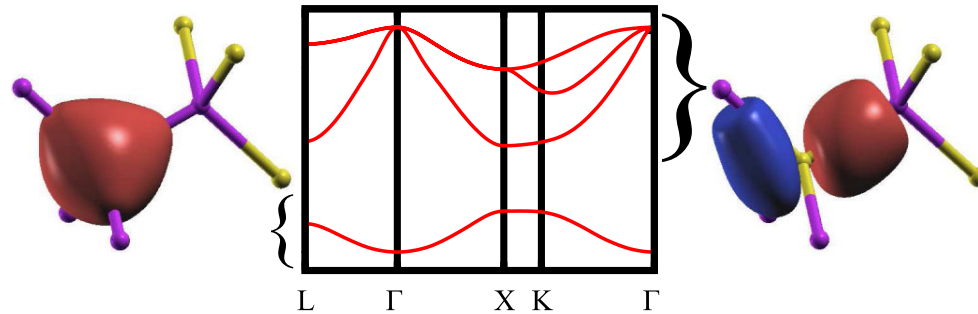
By construction, a **band representation has an atomic limit**,  
and all atomic limits yield a band representation



Recall: **Topological bands CANNOT Have Maximally Localized** Wannier Function



# ATOMIC LIMIT

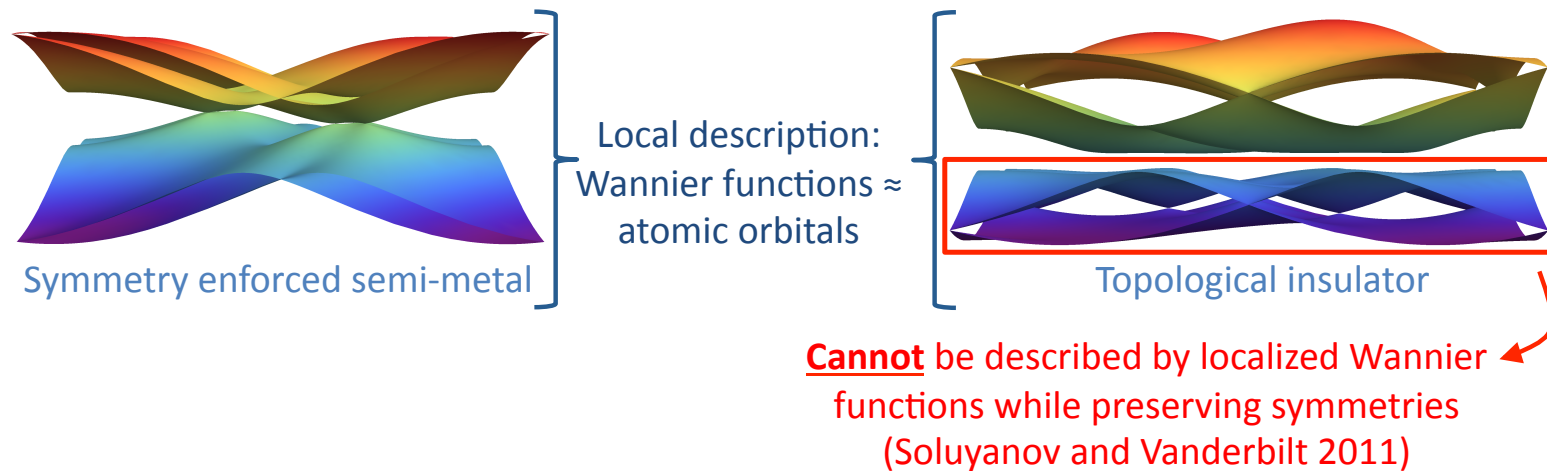


- 1) Bands in  $\rho_G$  are **connected** (this phase can always be realized) in the Brillouin zone
- 2) Bands in  $\rho_G$  are **not connected**: at least one **topological band**  
(Disconnected (P)EBR = set of disconnected bands that connected form an (P)EBR)

# GRAPHENE

What makes the disconnected bands topological?

All four bands come from a single set of localized orbitals ( $p_z$ , spin up/down)



Disconnected bands are topological because they lack localized Wannier functions that obey TR

## TQC Statement

All sets of bands induced from symmetric, localized orbitals, are topologically trivial by design.

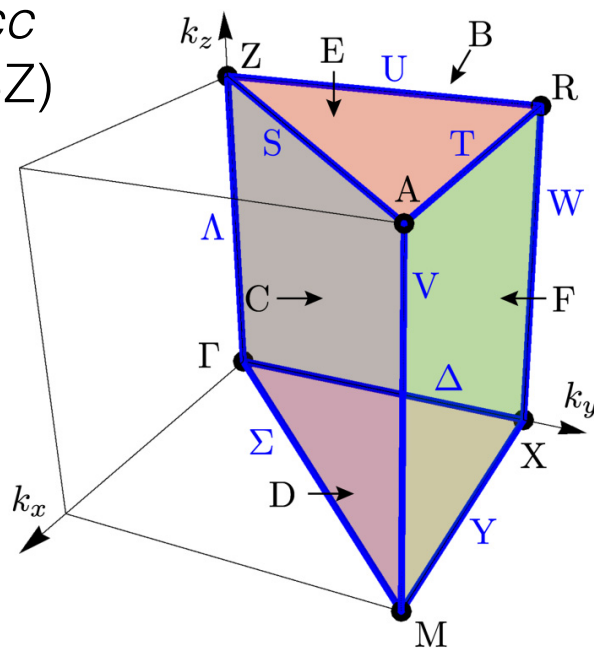
# Maximal k-vectors and paths

For all the 230 SG:

**maximal k-vectors** + minimal set non-redundant **connections**

**k** vector in a manifold is **maximal** if its little co-group  
it's not a subgroup of another manifold of vectors **k'**  
(in general coincides with high-symmetry k-vector)

*P4/ncc*  
(first BZ)



k-vec	mult.	Coordinates	Little	Maximal co-group	TR
$\Gamma$	1	(0,0,0)	$4/m\bar{m}m(D_{4h})$	yes	yes
Z	1	(0,0,1/2)	$4/m\bar{m}m(D_{4h})$	yes	yes
M	1	(1/2,1/2,0)	$4/m\bar{m}m(D_{4h})$	yes	yes
A	1	(1/2,1/2,1/2)	$4/m\bar{m}m(D_{4h})$	yes	yes
R	2	(0,1/2,1/2)	$m\bar{m}m(D_{2h})$	yes	yes
X	2	(0,1/2,0)	$m\bar{m}m(D_{2h})$	yes	yes
$\Lambda$	2	(0,0,w), $0 < w < 1/2$	$4mm(C_{4v})$	no	no
V	2	(1/2,1/2,w), $0 < w < 1/2$	$4mm(C_{4v})$	no	no
W	4	(0,1/2,w), $0 < w < 1/2$	$mm2(C_{2v})$	no	no
$\Sigma$	4	(u,u,0), $0 < u < 1/2$	$mm2(C_{2v})$	no	no
S	4	(u,u,1/2), $0 < u < 1/2$	$mm2(C_{2v})$	no	no
$\Delta$	4	(0,v,0), $0 < v < 1/2$	$mm2(C_{2v})$	no	no
U	4	(0,v,1/2), $0 < v < 1/2$	$mm2(C_{2v})$	no	no
Y	4	(u,1/2,0), $0 < u < 1/2$	$mm2(C_{2v})$	no	no
T	4	(u,1/2,1/2), $0 < u < 1/2$	$mm2(C_{2v})$	no	no
D	8	(u,v,0), $0 < u < v < 1/2$	$m(C_s)$	no	no
E	8	(u,v,1/2), $0 < u < v < 1/2$	$m(C_s)$	no	no
C	8	(u,u,w), $0 < u < w < 1/2$	$m(C_s)$	no	no
B	8	(0,v,w), $0 < v < w < 1/2$	$m(C_s)$	no	no
F	8	(u,1/2,w), $0 < u < w < 1/2$	$m(C_s)$	no	no
GP	16	(u,v,w), $0 < u < v < w < 1/2$	1(1)	no	no

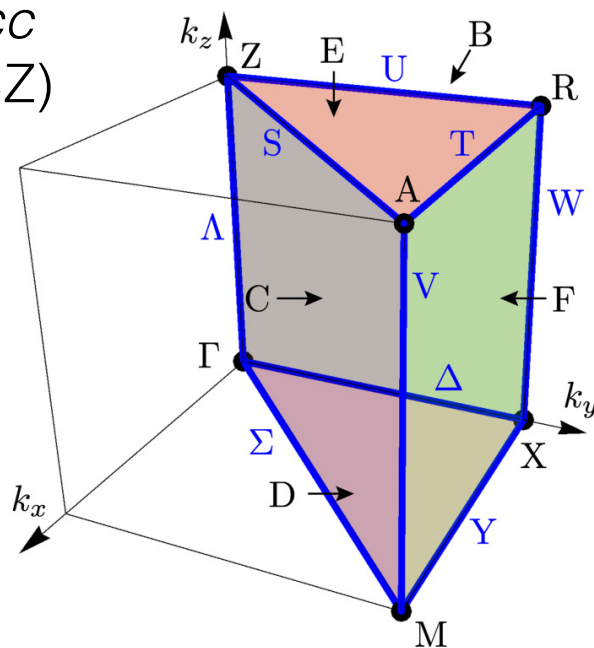
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k-vec	mult.	Coordinates	Little	Maximal co-group	TK
$\Gamma$	1	(0,0,0)	$4/m\bar{m}m(D_{4h})$	yes	yes
Z	1	(0,0,1/2)	$4/m\bar{m}m(D_{4h})$	yes	yes
M	1	(1/2,1/2,0)	$4/m\bar{m}m(D_{4h})$	yes	yes
A	1	(1/2,1/2,1/2)	$4/m\bar{m}m(D_{4h})$	yes	yes
R	2	(0,1/2,1/2)	$m\bar{m}m(D_{2h})$	yes	yes
X	2	(0,1/2,0)	$m\bar{m}m(D_{2h})$	yes	yes
$\Lambda$	2	(0,0,w), 0 < w < 1/2	$4mm(C_{4v})$	no	no
V	2	(1/2,1/2,w), 0 < w < 1/2	$4mm(C_{4v})$	no	no
W	4	(0,1/2,w), 0 < w < 1/2	$m\bar{m}2(C_{2v})$	no	no
$\Sigma$	4	(u,u,0), 0 < u < 1/2	$m\bar{m}2(C_{2v})$	no	no
S	4	(u,u,1/2), 0 < u < 1/2	$m\bar{m}2(C_{2v})$	no	no
$\Delta$	4	(0,v,0), 0 < v < 1/2	$m\bar{m}2(C_{2v})$	no	no
U	4	(0,v,1/2), 0 < v < 1/2	$m\bar{m}2(C_{2v})$	no	no
Y	4	(u,1/2,0), 0 < u < 1/2	$m\bar{m}2(C_{2v})$	no	no
T	4	(u,1/2,1/2), 0 < u < 1/2	$m\bar{m}2(C_{2v})$	no	no
D	8	(u,v,0), 0 < u < v < 1/2	$m(C_s)$	no	no
E	8	(u,v,1/2), 0 < u < v < 1/2	$m(C_s)$	no	no
C	8	(u,u,w), 0 < u < w < 1/2	$m(C_s)$	no	no
B	8	(0,v,w), 0 < v < w < 1/2	$m(C_s)$	no	no
F	8	(u,1/2,w), 0 < u < w < 1/2	$m(C_s)$	no	no
GP	16	(u,v,w), 0 < u < v < w < 1/2	1(1)	no	no

# Maximal k-vectors and paths

All possible connection between maximal and non-maximal  $\mathbf{k}$ -vectors

→ 2 manifolds are connected if:

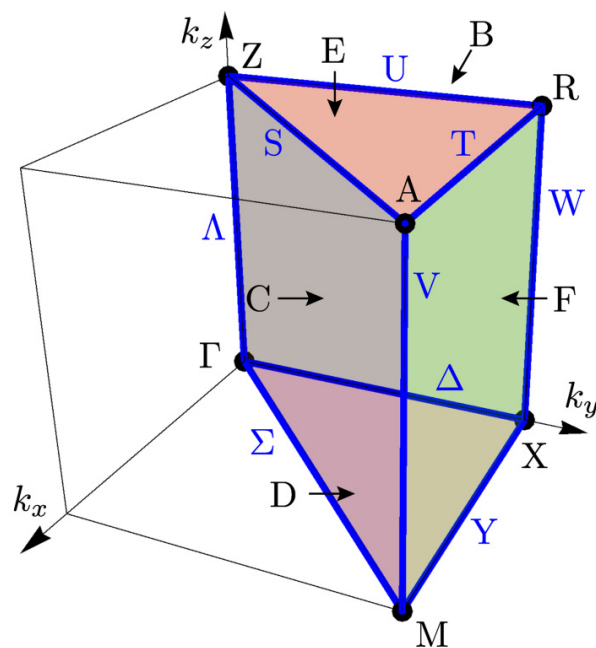
$$\mathbf{k}_i(\mathbf{u}_1) = \mathbf{k}_1$$

$$\mathbf{k}_i(\mathbf{u}_2) = \mathbf{k}_2$$



for each max.  $\mathbf{k}$  in  $^*\mathbf{k}$  and  $\mathbf{k}_i$  non-maximal

*P4/ncc*  
(first BZ)



Maximal $\mathbf{k}$ -vec	Connected $\mathbf{k}$ -vecs	Specific coordinates	Connections with the star
$\Gamma: (0,0,0)$	$\Lambda: (0,0,w)$	$w = 0$	2
	$\Delta: (0,v,0)$	$v = 0$	4
	$\Sigma: (u,u,0)$	$u = 0$	4
	$B: (0,v,w)$	$v = w = 0$	8
	$C: (u,u,w)$	$u = w = 0$	8
	$D: (u,v,0)$	$u = v = 0$	8

$\Gamma$ : 3 lines and 3 planes

# Compatibility relations

All possible connection between maximal and non-maximal  $\mathbf{k}$ -vectors

→ 2 manifolds are connected if:

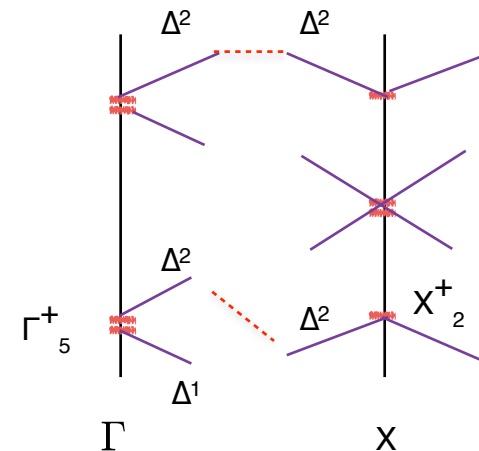
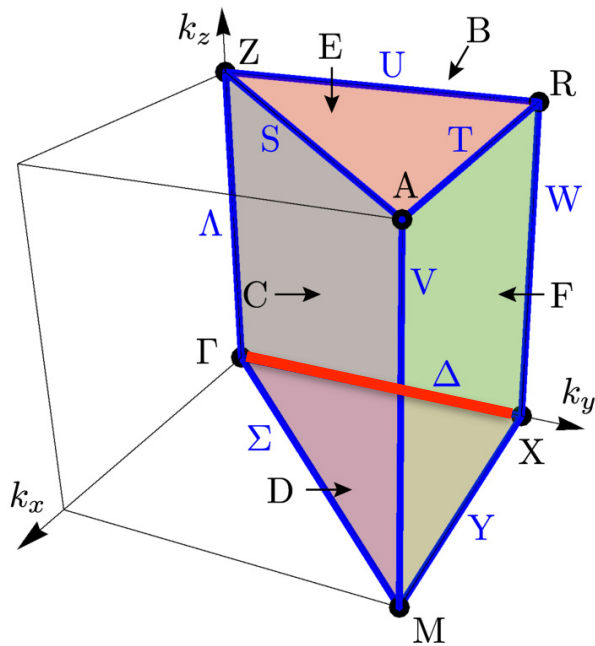
$$\mathbf{k}_i(\mathbf{u}_1) = \mathbf{k}_1$$

$$\mathbf{k}_i(\mathbf{u}_2) = \mathbf{k}_2$$



for each max.  $\mathbf{k}$  in  $^*\mathbf{k}$  and  $\mathbf{k}_i$  non-maximal

*P4/ncc*  
(first BZ)



# Compatibility relations

All possible connection between maximal and non-maximal  $\mathbf{k}$ -vectors

→ 2 manifolds are connected if:

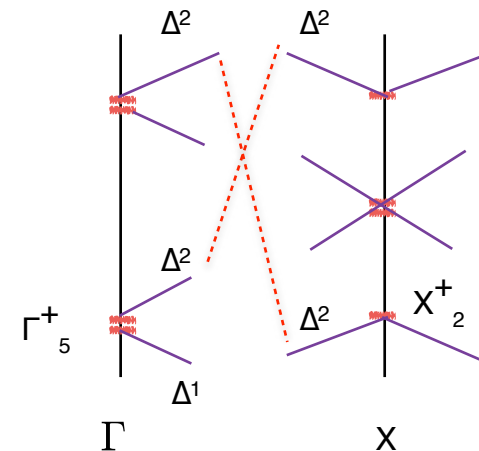
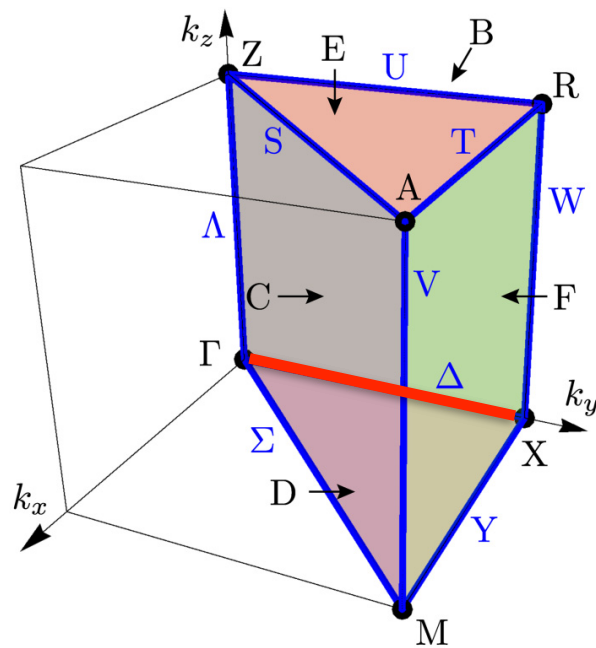
$$\mathbf{k}_i(\mathbf{u}_1) = \mathbf{k}_1$$

$$\mathbf{k}_i(\mathbf{u}_2) = \mathbf{k}_2$$



for each max.  $\mathbf{k}$  in  $^*\mathbf{k}$  and  $\mathbf{k}_i$  non-maximal

*P4/ncc*  
(first BZ)





# Compatibility relations

Reducing the number of paths

**(i) Paths are subspace of other paths**

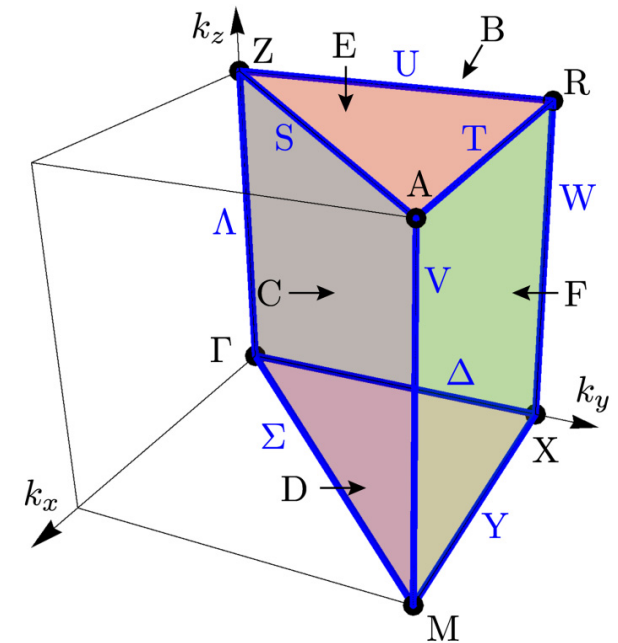
$k^1_M$  and  $k^2_M$  connect through  $k_p$  and  $k_l$ ,  $k_p$  is redundant

**(ii) Paths related by symmetry operations**

A single line or plane of the  $*k$  gives all independent restrictions

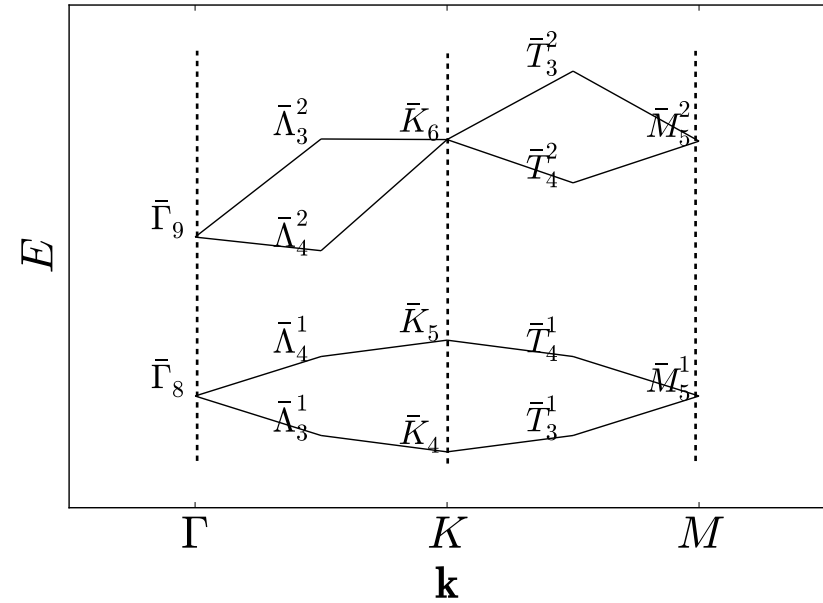
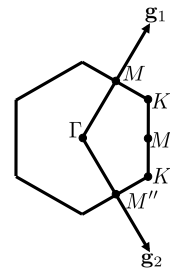
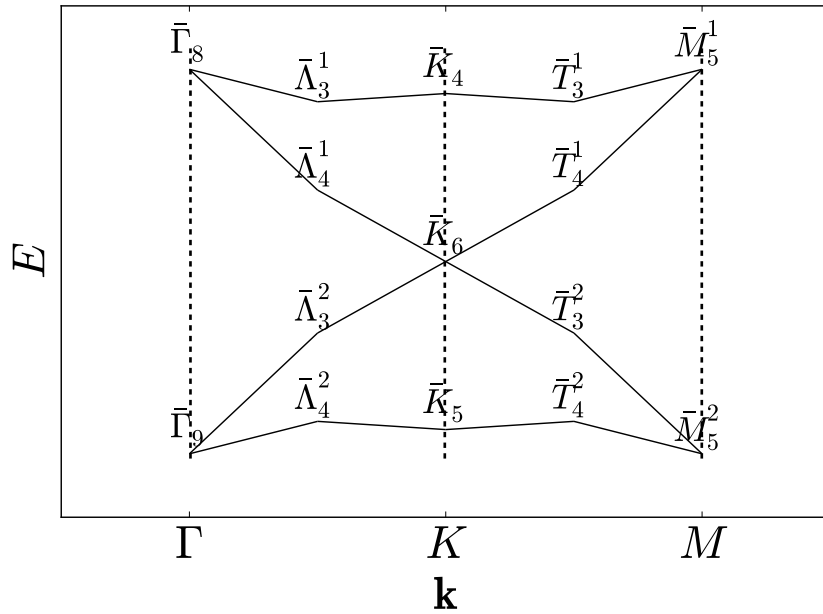
**(iii) Paths that are combinations of other paths**

\* additional restrictions in non-symmorphic groups (monodromy)



# Compatibility relations

2 possible connectivities for  $\bar{E}_1 \uparrow G(4)$

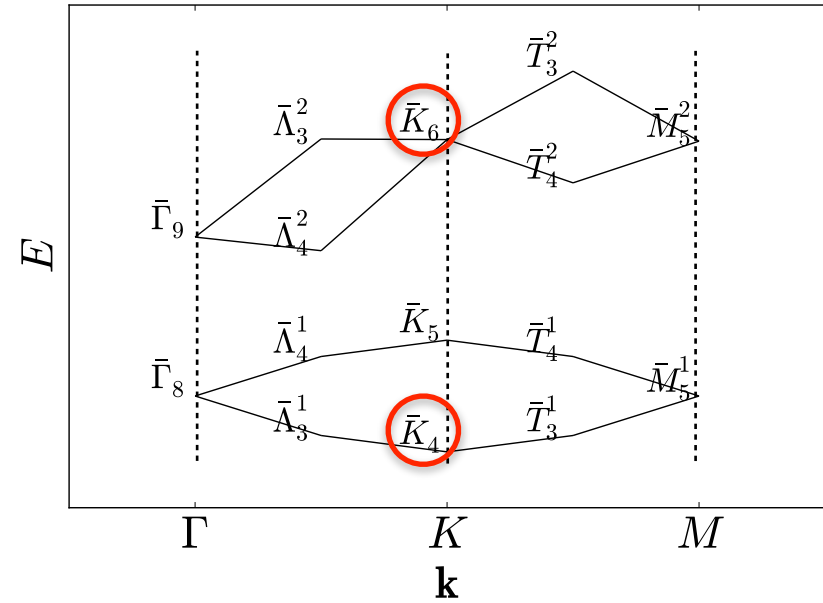
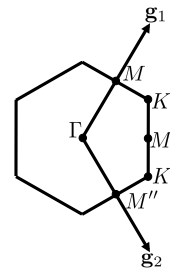
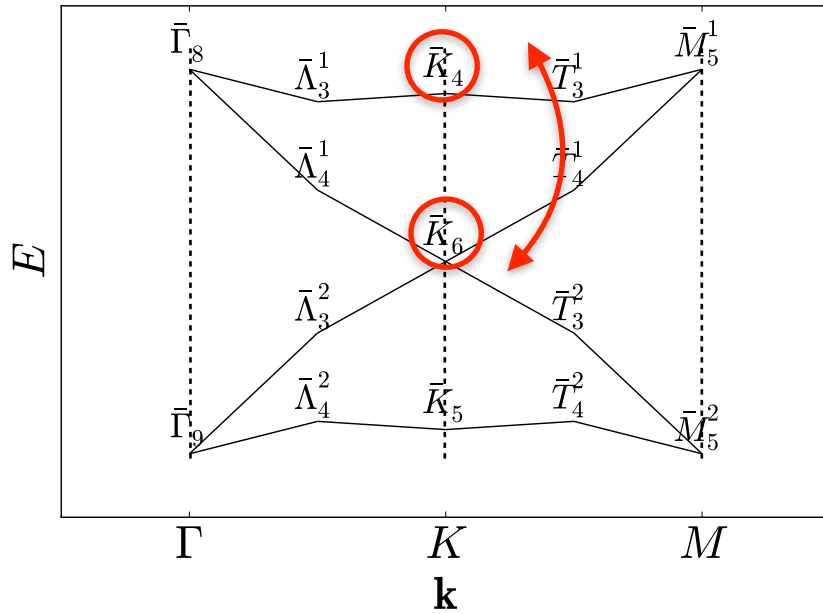


- Single connected component
- Fully connected

- Splitting of EBR
- Topological bands

# Compatibility relations

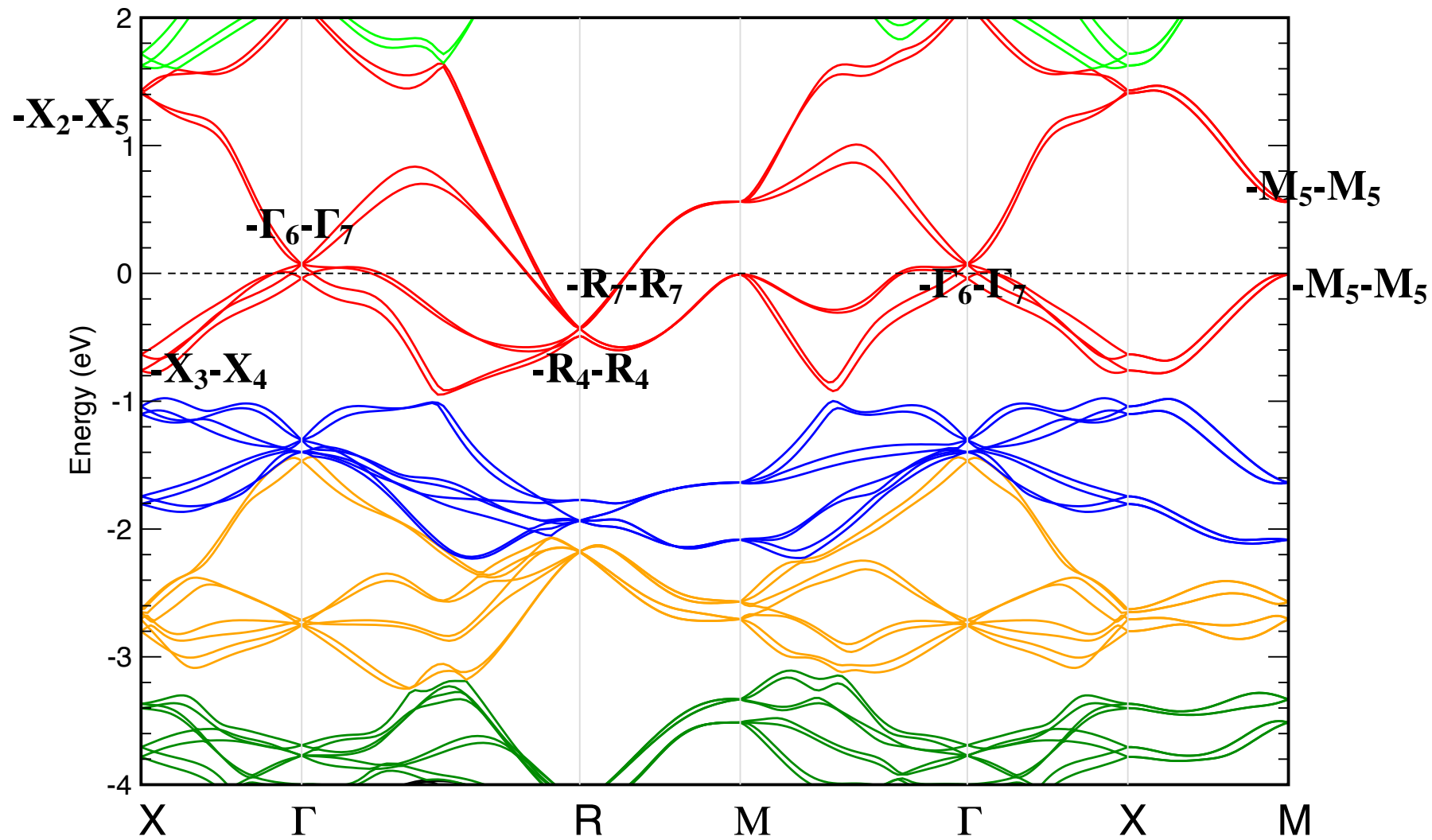
2 possible connectivities for  $\bar{E}_1 \uparrow G(4)$



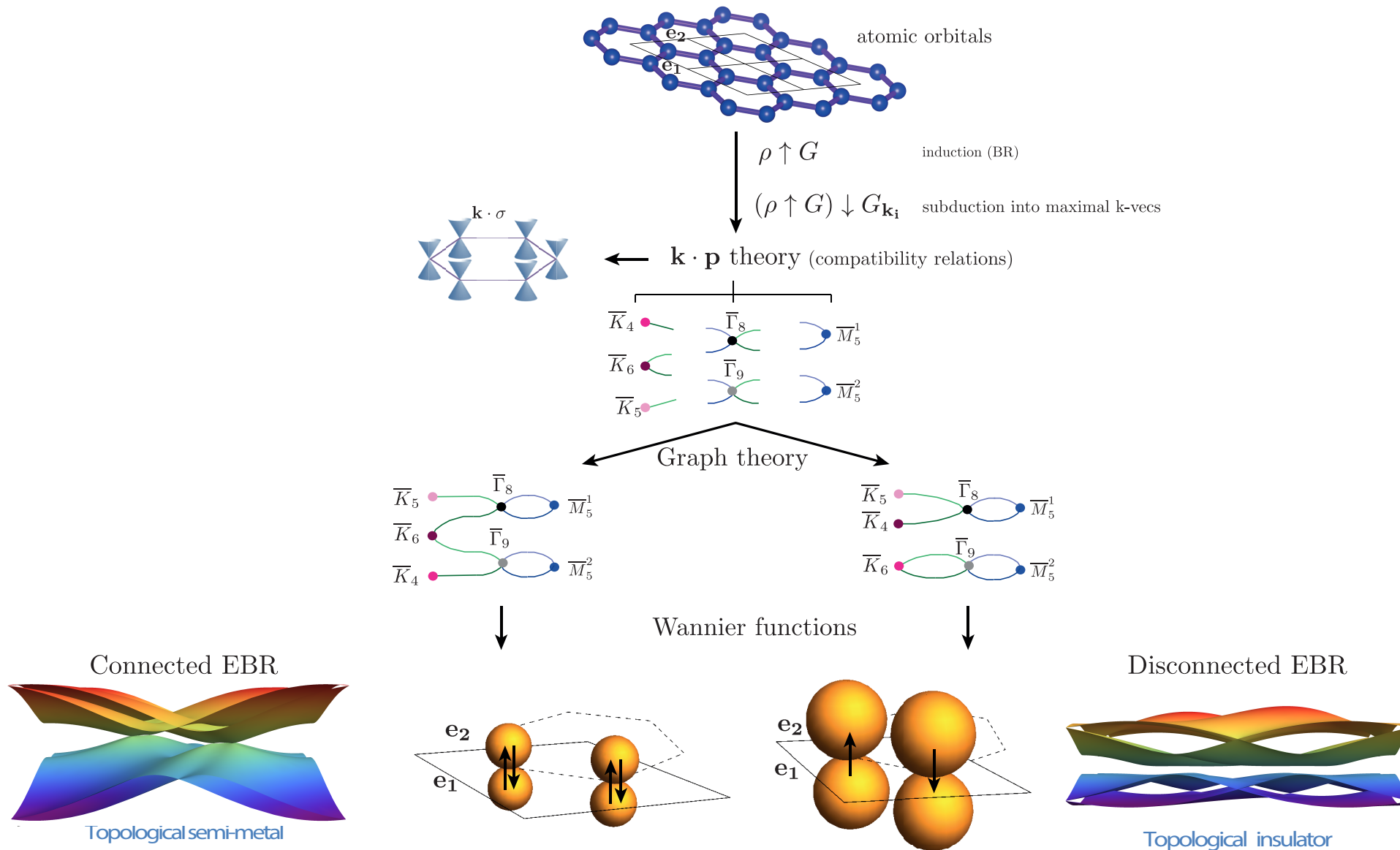
- Single connected component
- Fully connected

- Splitting of EBR
- Topological bands

# RhSi



# TOPOLOGICAL QUANTUM CHEMISTRY



# Classification of crystalline atomic limits

**10398** real-space **atomic limits** of materials

SG	MWP	WM	PG	Irrep	Dim	KR	Bands	Re	E	PE	SG	MWP	WM	PG	Irrep	Dim	KR	Bands	Re	E	PE
1	1a	1	1	$\Gamma_1$	1	1	1	1	e	e	131	2d	2	8	$\Gamma_2^-$	1	1	2	1	e	e
1	1a	1	1	$\bar{\Gamma}_2$	1	2	2	2	e	e	131	2d	2	8	$\Gamma_4^+$	1	1	2	1	e	e
2	1a	1	2	$\Gamma_1^+$	1	1	1	1	e	e	131	2d	2	8	$\Gamma_4^-$	1	1	2	1	e	e
2	1a	1	2	$\Gamma_1^-$	1	1	1	1	e	e	131	2d	2	8	$\Gamma_3^+$	1	1	2	1	e	e
2	1a	1	2	$\bar{\Gamma}_3$	1	2	2	2	e	e	131	2d	2	8	$\Gamma_3^-$	1	1	2	1	e	e
2	1a	1	2	$\bar{\Gamma}_2$	1	2	2	2	e	e	131	2d	2	8	$\bar{\Gamma}_5$	2	1	4	1	e	e
2	1b	1	2	$\Gamma_1^+$	1	1	1	1	e	e	131	2d	2	8	$\bar{\Gamma}_6$	2	1	4	1	e	e
2	1b	1	2	$\Gamma_1^-$	1	1	1	1	e	e	131	2e	2	14	$\Gamma_1$	1	1	2	1	e	e
2	1b	1	2	$\bar{\Gamma}_3$	1	2	2	2	e	e	131	2e	2	14	$\Gamma_4$	1	1	2	1	e	e

**SG:** Space Group

**MWP:** Maximal Wyckoff Position

**WM:** Wyckoff multiplicity in the primitive cell

**PG:** Point group number of the site-symmetry

**Irrep:** Name of the Irrep of the site-symmetry for each BR

**KR:** 1 for PEER, 2 for EBR (f and s)

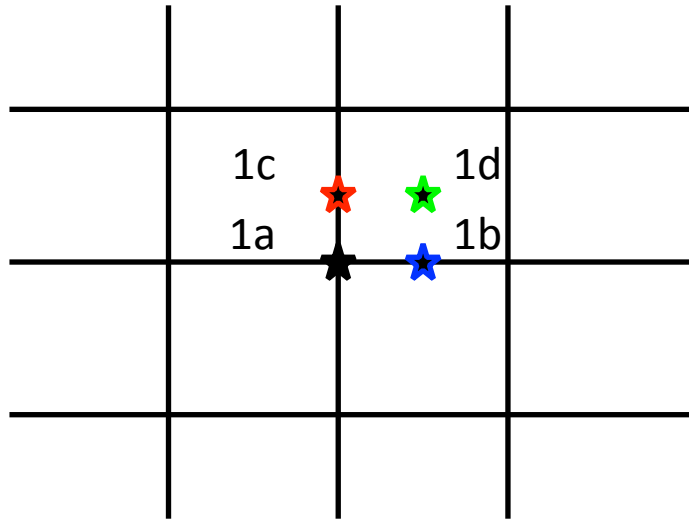
**Bands:** Total number of bands

**Re:** 1 for TRS at each k, 2 for connection with its conjugate

**E:** e for elementary, c for composite

**PE:** e for elementary, c for composite

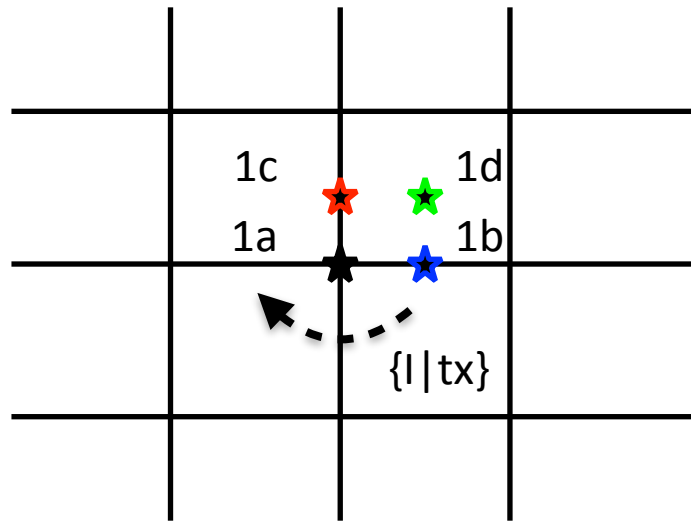
Spinless 2D material,  $\mathcal{I}$ , s and p orbitals



Zeroth Order: How To Use The Eigenvalues of EBRs to Find Topological Materials:

1. Find the Atomic Limit Real Space Orbitals

Spinless 2D material,  $\mathcal{I}$ , s and p orbitals

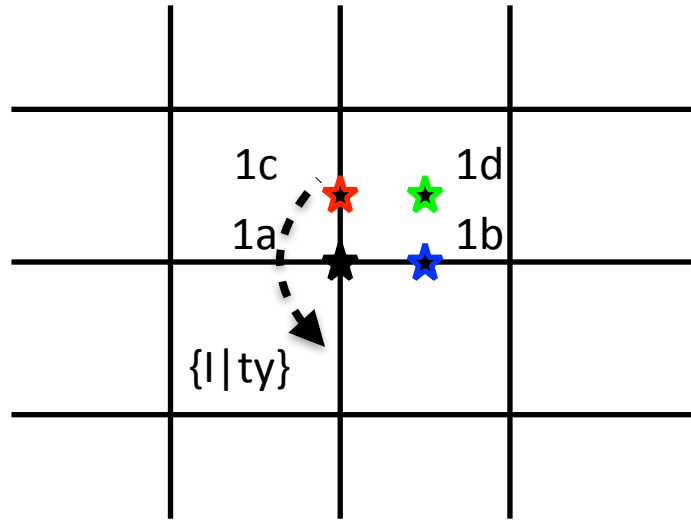


Zeroth Order: How To Use The Eigenvalues of EBRs to Find Topological Materials:

1. Find the Atomic Limit Real Space Orbitals



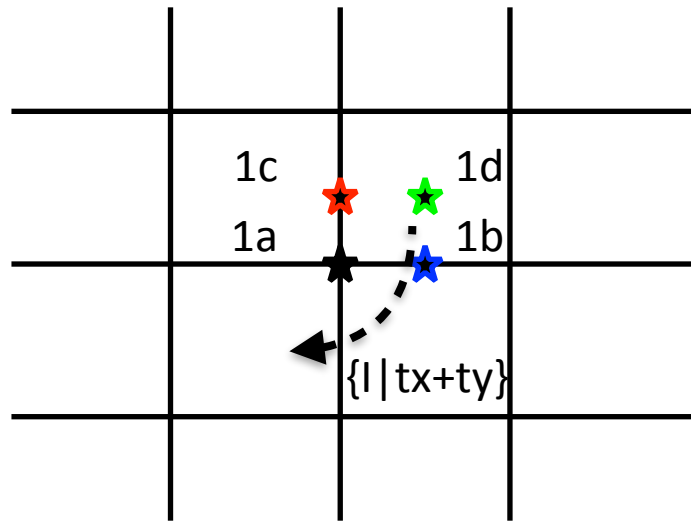
Spinless 2D material,  $\mathcal{I}$ , s and p orbitals



Zeroth Order: How To Use The Eigenvalues of EBRs to Find Topological Materials:

1. Find the Atomic Limit Real Space Orbitals

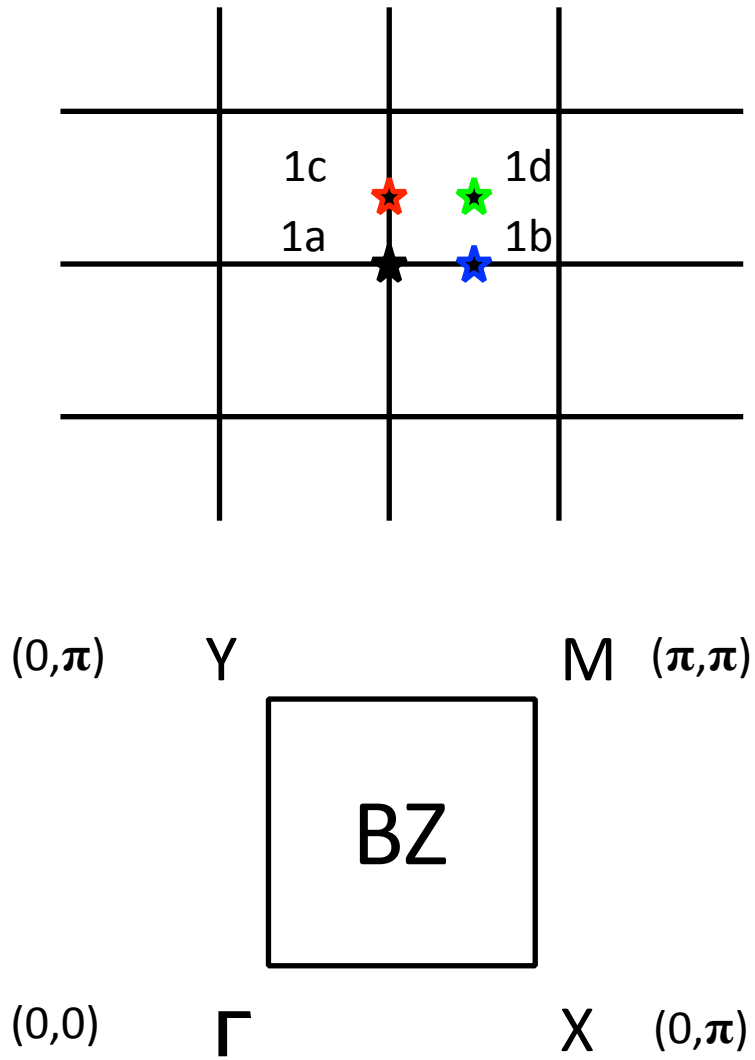
Spinless 2D material,  $\mathcal{I}$ , s and p orbitals



Zeroth Order: How To Use The Eigenvalues of EBRs to Find Topological Materials:

2. Obtain the Induced Representation in  $k$ (momentum) space

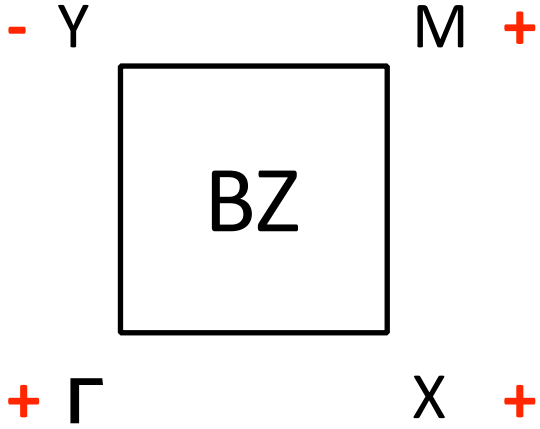
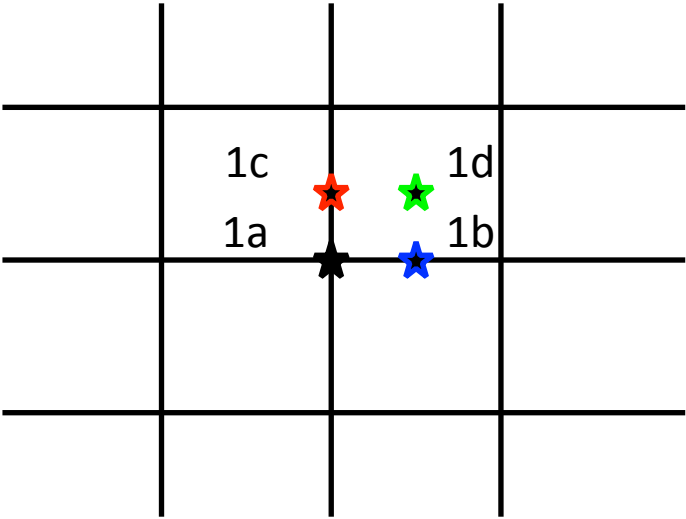
Spinless 2D material,  $\mathcal{I}$ ,  $\mathbf{s}$  and  $\mathbf{p}$  orbitals



	$\Gamma$	$X$	$Y$	$M$
s 1a	+	+	+	+
p 1a	-	-	-	-
s 1b	+	-	+	-
p 1b	-	+	-	+
s 1c	+	+	-	-
p 1c	-	-	+	+
s 1d	+	-	+	-
p 1d	-	+	+	-

↑  
eik  
↓

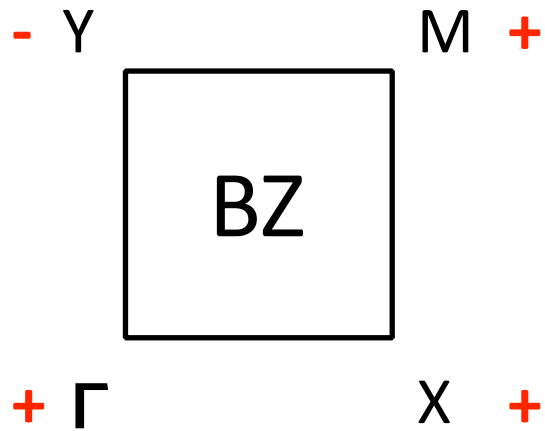
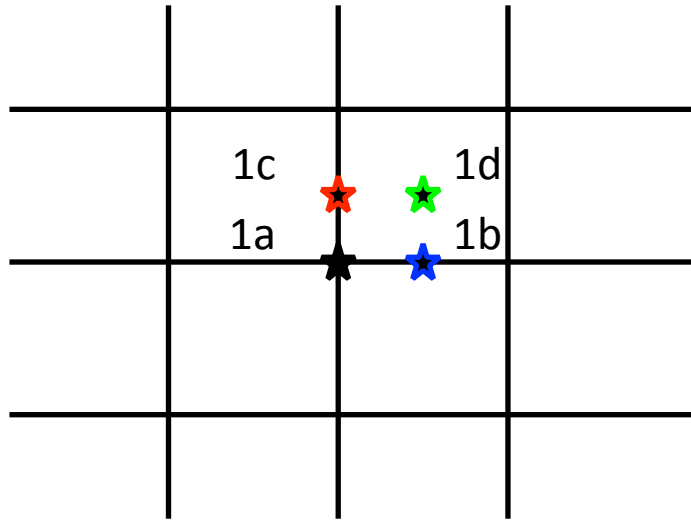
Spinless 2D material,  $\mathcal{I}$ , s and p orbitals



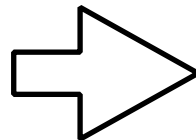
odd # in  $\mathcal{I}$  eigenvalues

	$\Gamma$	X	Y	M
s 1a	+	+	+	+
p 1a	-	-	-	-
s 1b	+	-	+	-
p 1b	-	+	-	+
s 1c	+	+	-	-
p 1c	-	-	+	+
s 1d	+	-	+	-
p 1d	-	+	+	-

Spinless 2D material,  $\mathcal{I}$ , s and p orbitals



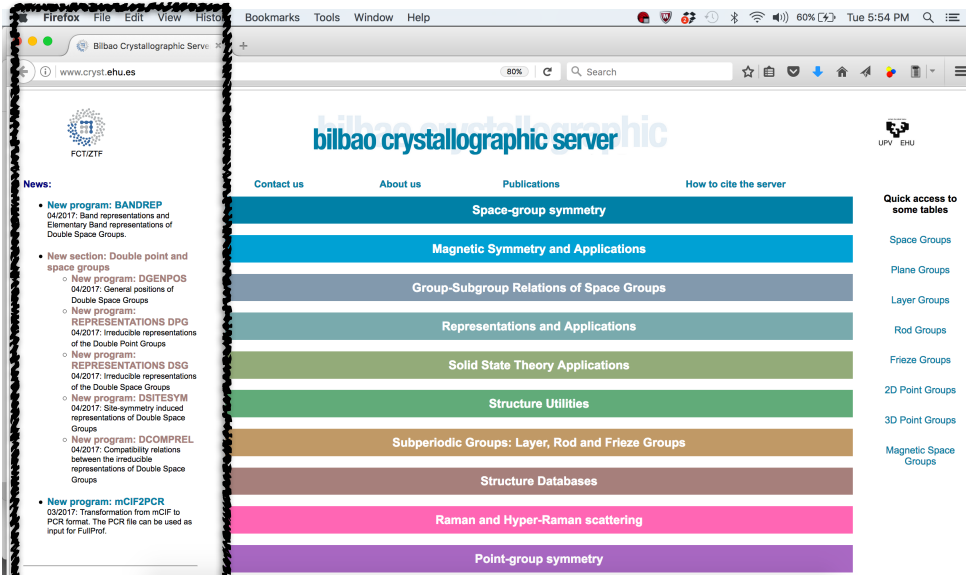
odd # in  $\mathcal{I}$  eigenvalues



	$\Gamma$	X	Y	M
s 1a	+	+	+	+
p 1a	-	-	-	-
s 1b	+	-	+	-
p 1b	-	+	-	+
s 1c	+	+	-	-
p 1c	-	-	+	+
s 1d	+	-	+	-
p 1d	-	+	+	-

TOPOLOGICAL

<http://www.cryst.ehu.es/cryst/bandrep>



All 230 double groups

For each group, reps/wyckoff/k point reps

Compatibility relations in the Brillouin Zone

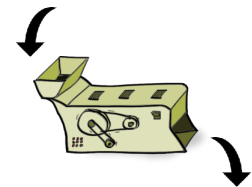
All 10398 Atomic limits

Disconnected EBR's/Graphs/Topological bands

Wyckoff pos.	1a(6mm)	1a(6mm)	1a(6mm)	1a(6mm)	1a(6mm)	1a(6mm)	2b(3m)
Band-Rep.	$A_1 \uparrow G(1)$	$A_2 \uparrow G(1)$	$B_1 \uparrow G(1)$	$B_2 \uparrow G(1)$	$E_1 \uparrow G(2)$	$E_2 \uparrow G(2)$	$A_1 \uparrow G(2)$
Decomposable/Indecomposable	Indecomposable	Indecomposable	Indecomposable	Indecomposable	Indecomposable	Indecomposable	Indecomposable
$\Gamma: (0,0,0)$	$\Gamma_1(1)$	$\Gamma_2(1)$	$\Gamma_4(1)$	$\Gamma_3(1)$	$\Gamma_6(2)$	$\Gamma_5(2)$	$\Gamma_1(1) \oplus \Gamma_4(1)$
A: (0,0,1/2)	$A_1(1)$	$A_2(1)$	$A_4(1)$	$A_3(1)$	$A_6(2)$	$A_5(2)$	$A_1(1) \oplus A_4(1)$
K: (1/3,1/3,0)	$K_1(1)$	$K_2(1)$	$K_2(1)$	$K_1(1)$	$K_3(2)$	$K_3(2)$	$K_3(2)$
H: (1/3,1/3,1/2)	$H_1(1)$	$H_2(1)$	$H_2(1)$	$H_1(1)$	$H_3(2)$	$H_3(2)$	$H_3(2)$
M: (1/2,0,0)	$M_1(1)$	$M_2(1)$	$M_4(1)$	$M_3(1)$	$M_3(1) \oplus M_4(1)$	$M_1(1) \oplus M_2(1)$	$M_1(1) \oplus M_4(1)$
L: (1/2,0,1/2)	$L_1(1)$	$L_2(1)$	$L_4(1)$	$L_3(1)$	$L_3(1) \oplus L_4(1)$	$L_1(1) \oplus L_2(1)$	$L_1(1) \oplus L_4(1)$

Atom arrangement

Orbital



High-symmetry points

# Band Representations in the Bilbao Crystallographic Server

<http://www.cryst.ehu.es/cryst/bandrep>

Bilbao Crystallographic Server → BANDREP

[Help](#)

## Band representations of the Double Space Groups

### Band Representations

This program calculates the band representations (BR) induced from the irreps of the site-symmetry group of a given Wyckoff position.

Alternatively, it gives the set of elementary BRs of a Double Space Group.

In both cases, it can be chosen to get the BRs with or without time-reversal symmetry.

The program also indicates if the elementary BRs are decomposable or indecomposable. If it is decomposable, the program gives all the possible ways to decompose it.

Please, enter the sequential number of group as given in the *International Tables for Crystallography*, Vol. A

choose it

1. Get the elementary BRs without time-reversal symmetry
2. Get the elementary BRs with time-reversal symmetry
3. Get the BRs without time-reversal symmetry from a Wyckoff position
4. Get the BRs with time-reversal symmetry from a Wyckoff position

[Elementary](#)

[Elementary TR](#)

[Wyckoff](#)

[Wyckoff TR](#)

Bilbao Crystallographic Server  
<http://www.cryst.ehu.es>

For comments, please mail to  
[administrador.bsc@ehu.es](mailto:administrador.bsc@ehu.es)

<http://www.cryst.ehu.es/cryst/bandrep>

Output

12b(2)	8a(3)	8a(3)	12b(2)
$B \uparrow G(6)$	$1 \bar{E}^2 \bar{E} \uparrow G(8)$	$\bar{E} \bar{E} \uparrow G(8)$	$1 \bar{E}^2 \bar{E} \uparrow G(12)$
Indecomposable	Indecomposable	Decomposable	Decomposable
$2 \Gamma_4(3)$	$2 \bar{\Gamma}_5(2) \oplus \bar{\Gamma}_6 \bar{\Gamma}_7(4)$	$2 \bar{\Gamma}_6 \bar{\Gamma}_7(4)$	$2 \bar{\Gamma}_5(2) \oplus 2 \bar{\Gamma}_6 \bar{\Gamma}_7(4)$
$H_1(1) \oplus H_2 H_3(2) \oplus H_4(3)$	$2 \bar{H}_5(2) \oplus \bar{H}_6 \bar{H}_7(4)$	$2 \bar{H}_6 \bar{H}_7(4)$	$2 \bar{H}_5(2) \oplus 2 \bar{H}_6 \bar{H}_7(4)$
$P_1(2) \oplus P_2(2) \oplus P_3(2)$	$\bar{P}_5(1) \oplus \bar{P}_6(1) \oplus 2 \bar{P}_7(3)$	$2 \bar{P}_4(1) \oplus 2 \bar{P}_7(3)$	$\bar{P}_4(1) \oplus \bar{P}_5(1) \oplus \bar{P}_6(1) \oplus 3 \bar{P}_7(3)$
$PA_1(2) \oplus PA_2(2) \oplus PA_3(2)$	$\bar{PA}_5(1) \oplus \bar{PA}_6(1) \oplus 2 \bar{PA}_7(3)$	$2 \bar{PA}_4(1) \oplus 2 \bar{PA}_7(3)$	$\bar{PA}_4(1) \oplus \bar{PA}_5(1) \oplus \bar{PA}_6(1) \oplus 3 \bar{PA}_7(3)$
$3 N_1(1) \oplus 3 N_2(1)$	$4 \bar{N}_3 \bar{N}_4(2)$	$4 \bar{N}_3 \bar{N}_4(2)$	$6 \bar{N}_3 \bar{N}_4(2)$



<http://www.cryst.ehu.es/cryst/bandrep>

Output

	branch 1	branch 2
1	$\bar{H}_5, \bar{\Gamma}_5, \bar{P}_5, \bar{P}_6, \bar{N}_4, \bar{N}_4$	$\bar{H}_6, \bar{H}_7, \bar{\Gamma}_6, \bar{\Gamma}_7, \bar{P}_4, \bar{P}_7, \bar{N}_3, \bar{N}_3, \bar{N}_4, \bar{N}_4$
2	$\bar{H}_6, \bar{\Gamma}_6, \bar{P}_4, \bar{P}_6, \bar{N}_4, \bar{N}_4$	$\bar{H}_7, \bar{H}_5, \bar{\Gamma}_5, \bar{\Gamma}_7, \bar{P}_5, \bar{P}_7, \bar{N}_3, \bar{N}_3, \bar{N}_4, \bar{N}_4$
3	$\bar{H}_7, \bar{\Gamma}_7, \bar{P}_4, \bar{P}_5, \bar{N}_4, \bar{N}_4$	$\bar{H}_5, \bar{H}_6, \bar{\Gamma}_5, \bar{\Gamma}_6, \bar{P}_6, \bar{P}_7, \bar{N}_3, \bar{N}_3, \bar{N}_4, \bar{N}_4$

<http://www.cryst.ehu.es/cryst/bandrep>

Output

## Elementary band-representations without time-reversal symmetry of the Double Space Group $I2_13$ (No. 199)

The first row shows the Wyckoff position from which the band representation is induced.  
In parentheses, the symbol of the point group isomorphic to the site-symmetry group.

The second row gives the symbol  $\rho \uparrow G$ , where  $\rho$  is the irrep of the site-symmetry group.  
In parentheses, the dimension of the representation.

The output shows the decomposition of the band representations into irreps of the little groups  
of the given k-vectors in the first column.  
In parentheses, the dimensions of the representations.

Minimal set of paths and compatibility relations to analyse the connectivity

Show all types of k-vectors

<http://www.cryst.ehu.es/cryst/bandrep>

Output

Maximal k-vec	Compatibility relations	Intermediate path	Compatibility relations	Maximal k-vec
$\Gamma:(0,0,0)$	$\Gamma_1(1) \rightarrow \Delta_1(1)$ $\Gamma_2(1) \rightarrow \Delta_1(1)$ $\Gamma_3(1) \rightarrow \Delta_1(1)$ $\Gamma_4(3) \rightarrow \Delta_1(1) \oplus 2 \Delta_2(1)$ $\bar{\Gamma}_5(2) \rightarrow \bar{\Delta}_3(1) \oplus \bar{\Delta}_4(1)$ $\bar{\Gamma}_6(2) \rightarrow \bar{\Delta}_3(1) \oplus \bar{\Delta}_4(1)$ $\bar{\Gamma}_7(2) \rightarrow \bar{\Delta}_3(1) \oplus \bar{\Delta}_4(1)$	$\Delta:(0,v,0)$	$H_1(1) \rightarrow \Delta_2(1)$ $H_2(1) \rightarrow \Delta_2(1)$ $H_3(1) \rightarrow \Delta_2(1)$ $H_4(3) \rightarrow 2 \Delta_1(1) \oplus \Delta_2(1)$ $\bar{H}_5(2) \rightarrow \bar{\Delta}_3(1) \oplus \bar{\Delta}_4(1)$ $\bar{H}_6(2) \rightarrow \bar{\Delta}_3(1) \oplus \bar{\Delta}_4(1)$ $\bar{H}_7(2) \rightarrow \bar{\Delta}_3(1) \oplus \bar{\Delta}_4(1)$	$H:(1,1,1)$
$\Gamma:(0,0,0)$	$\Gamma_1(1) \rightarrow \Lambda_1(1)$ $\Gamma_2(1) \rightarrow \Lambda_2(1)$ $\Gamma_3(1) \rightarrow \Lambda_3(1)$ $\Gamma_4(3) \rightarrow \Lambda_1(1) \oplus \Lambda_2(1) \oplus \Lambda_3(1)$ $\bar{\Gamma}_5(2) \rightarrow \bar{\Lambda}_5(1) \oplus \bar{\Lambda}_6(1)$ $\bar{\Gamma}_6(2) \rightarrow \bar{\Lambda}_4(1) \oplus \bar{\Lambda}_6(1)$ $\bar{\Gamma}_7(2) \rightarrow \bar{\Lambda}_4(1) \oplus \bar{\Lambda}_5(1)$	$\Lambda:(-u,u,-u)$	$H_1(1) \rightarrow \Lambda_1(1)$ $H_2(1) \rightarrow \Lambda_2(1)$ $H_3(1) \rightarrow \Lambda_3(1)$ $H_4(3) \rightarrow \Lambda_1(1) \oplus \Lambda_2(1) \oplus \Lambda_3(1)$ $\bar{H}_5(2) \rightarrow \bar{\Lambda}_5(1) \oplus \bar{\Lambda}_6(1)$ $\bar{H}_6(2) \rightarrow \bar{\Lambda}_4(1) \oplus \bar{\Lambda}_6(1)$ $\bar{H}_7(2) \rightarrow \bar{\Lambda}_4(1) \oplus \bar{\Lambda}_5(1)$	$H:(1,1,1)$
$\Gamma:(0,0,0)$	$\Gamma_1(1) \rightarrow \Lambda_1(1)$ $\Gamma_2(1) \rightarrow \Lambda_2(1)$ $\Gamma_3(1) \rightarrow \Lambda_3(1)$ $\Gamma_4(3) \rightarrow \Lambda_1(1) \oplus \Lambda_2(1) \oplus \Lambda_3(1)$ $\bar{\Gamma}_5(2) \rightarrow \bar{\Lambda}_5(1) \oplus \bar{\Lambda}_6(1)$ $\bar{\Gamma}_6(2) \rightarrow \bar{\Lambda}_4(1) \oplus \bar{\Lambda}_6(1)$ $\bar{\Gamma}_7(2) \rightarrow \bar{\Lambda}_4(1) \oplus \bar{\Lambda}_5(1)$	$\Lambda:(-u,u,-u)$	$P_1(2) \rightarrow \Lambda_1(1) \oplus \Lambda_2(1)$ $P_2(2) \rightarrow \Lambda_2(1) \oplus \Lambda_3(1)$ $P_3(2) \rightarrow \Lambda_1(1) \oplus \Lambda_3(1)$ $\bar{P}_4(1) \rightarrow \bar{\Lambda}_4(1)$ $\bar{P}_5(1) \rightarrow \bar{\Lambda}_5(1)$ $\bar{P}_6(1) \rightarrow \bar{\Lambda}_6(1)$ $\bar{P}_7(3) \rightarrow \bar{\Lambda}_4(1) \oplus \bar{\Lambda}_5(1) \oplus \bar{\Lambda}_6(1)$	$P:(1/2,1/2,1/2)$

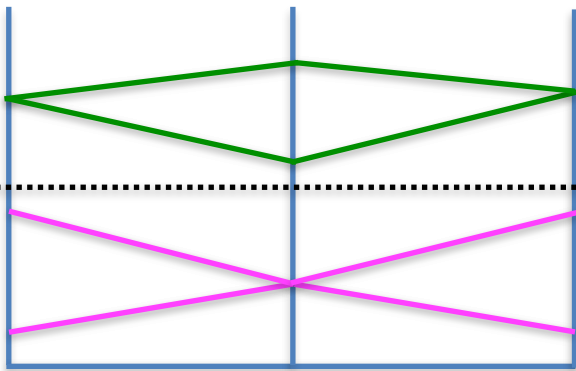
# Materials search

Topological “insulating” classes  $\left\{ \begin{array}{l} \text{EBR1} \\ \text{EBR2} \end{array} \right.$

$$\mathbf{b} = \sum_i n_i(\mathbf{k}) = \sum_m a_m \text{EBR}$$

LCEBR

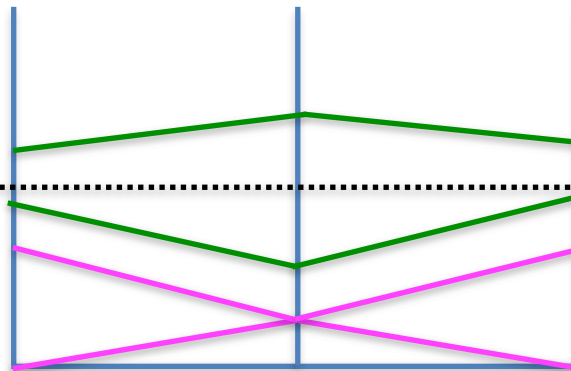
Linear Combination EBR



$$\mathbf{b} = \text{EBR1}$$

SEBR

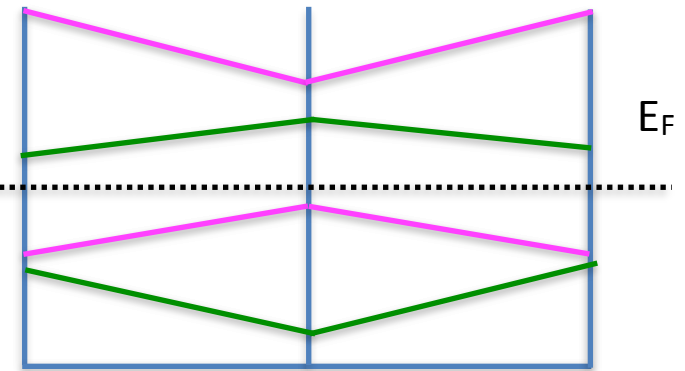
Split EBR



$$\mathbf{b} = \text{EBR1} + 1/2 * \text{EBR2}$$

NLC

Non Linear Combination

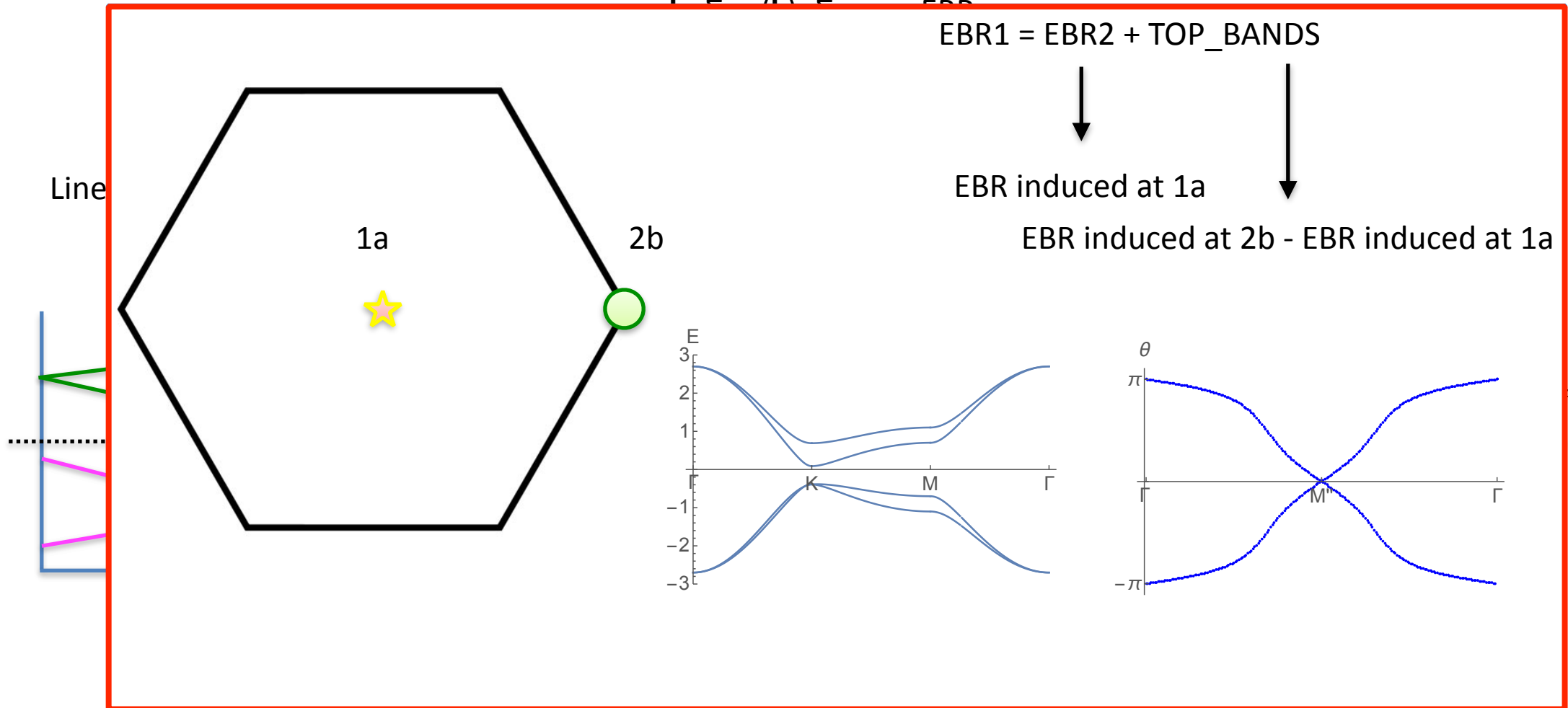


$$\mathbf{b} = ?$$

\* Fragile:  $\text{EBR}_F = \text{EBR1} - \text{EBR2}$

# Materials search

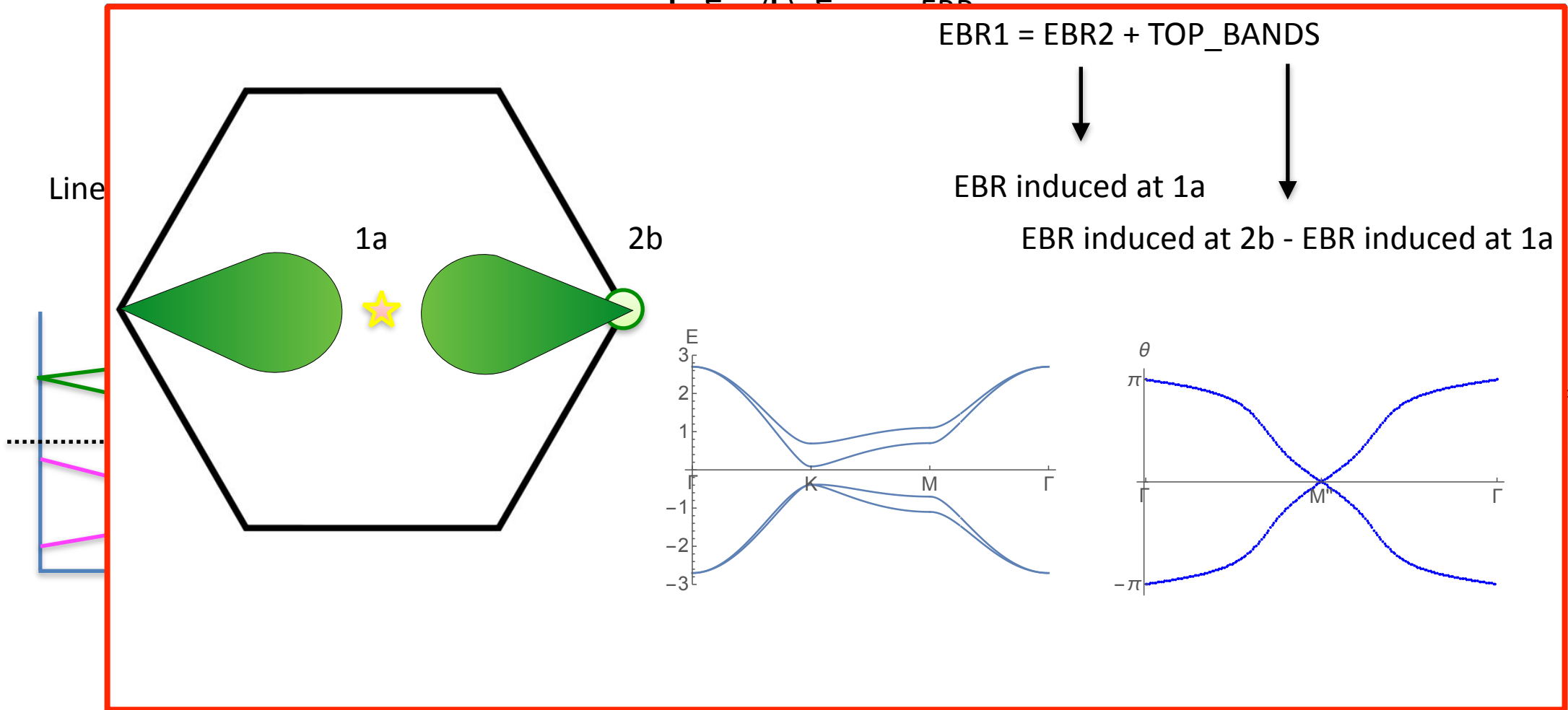
Topological “insulating” classes  $\left\{ \begin{array}{l} \text{EBR1} \\ \text{EBR2} \end{array} \right.$



\* Fragile:  $\text{EBR}_F = \text{EBR1} - \text{EBR2}$

# Materials search

Topological “insulating” classes  $\left\{ \begin{array}{l} \text{EBR1} \\ \text{EBR2} \end{array} \right.$



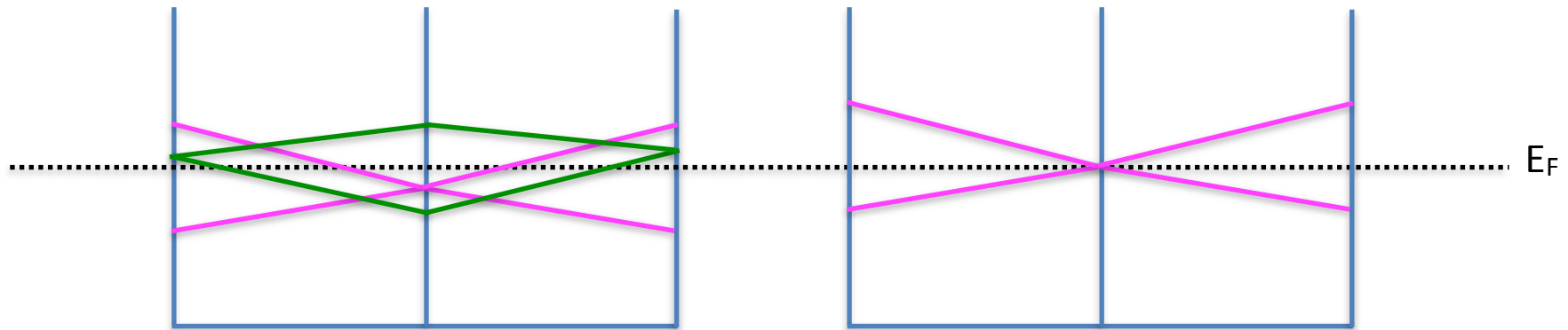
\* Fragile:  $\text{EBR}_F = \text{EBR1} - \text{EBR2}$

# Materials search

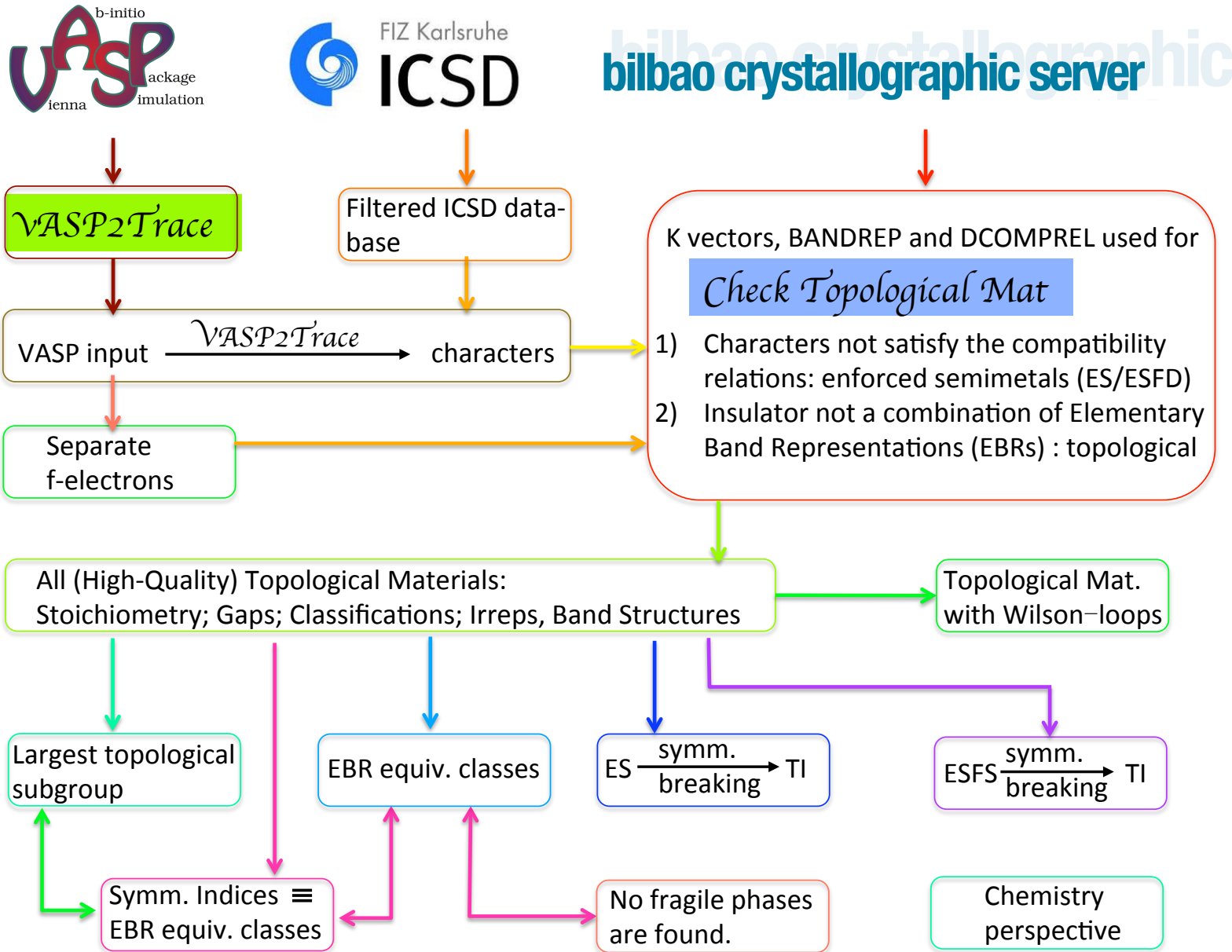
Topological “metallic” classes  $\left\{ \begin{array}{l} \text{EBR1} \\ \text{EBR2} \end{array} \right.$

ES  
Enforced semimetals

ESFD  
ES Fermi Degeneracy

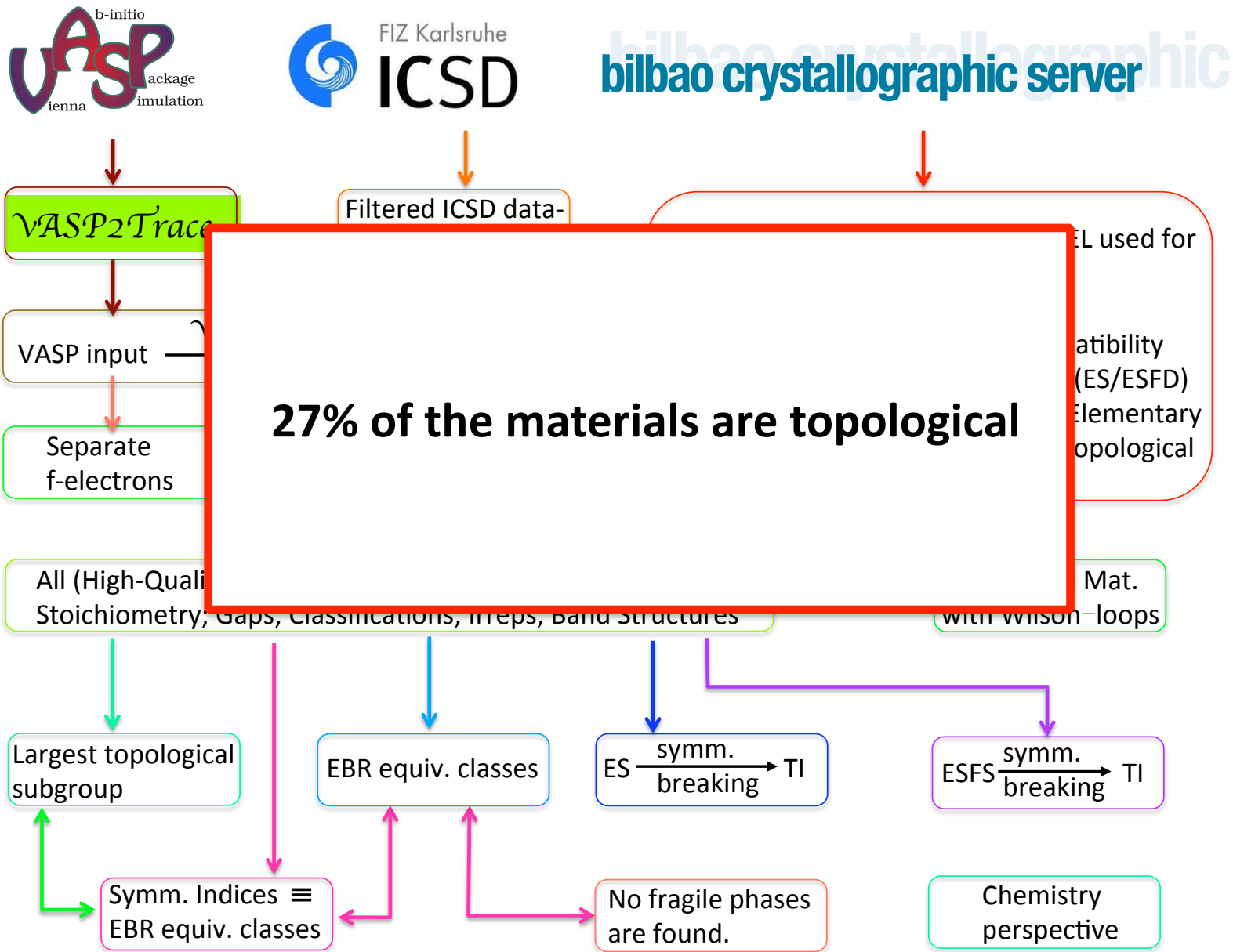


# Materials search





# Materials search







# Topological Materials Database

<https://www.topologicalquantumchemistry.org>

 **Topological Materials Database**

24905 Materials: 4339 Topological Insulators, 10061 Semi-Metals

 Search  About 

Compound Contains

Only these elements  Exclude

ICSD Number

e.g. Bi1 Se2 Ge

eg. O1 N

- or -

eg. 123456

Search

▼ Show Advanced Search

1 H																	2 He
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba	57 La	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn

# Check Topology

[www.cryst.ehu.es/cryst/checktopologicalmat](http://www.cryst.ehu.es/cryst/checktopologicalmat)

## Check Topological Mat

### Check Topological Mat.

Given a file that contains the eigenvalues at each maximal k-vec of a space group, the program gives the set of irreducible representations at each maximal k-vec (time-reversal is assumed). Then, using the compatibility relations and the set of Elementary Band Representations (EBRs), it checks whether the set of bands can be put as linear combinations of EBRs. This (self-explanatory) file shows the format of the file to be uploaded in the menu on the right:

#### [File\\_Description](#)

You can download examples of input files here:

[Example\\_Ag1Ge1Li2](#)      [Example\\_Ag1O2Sc1](#)  
[Example\\_B2Ca3Ni7](#)      [Example\\_Ba3Co10O17](#)  
[Example\\_Ba3Ca1O9Ru2](#)

You can generate the "trace.txt" file in your own computer using VASP and this program (fortran).

#### [vasp2trace](#)

Read the "README.pdf" file for help on the use of vasp2trace.  
If you are using "Check Topological Mat." and/or "vasp2trace" programs in the preparation of an article, please cite this reference:

[M.G. Vergniory, L. Elcoro, C. Felser, B.A. Bernevig, Z. Wang \(2018\) arXiv:1807.10271](#)

Upload your traces.txt file

No file selected.

1. A number of bands that is **not a sum of EBRs is topological**
2. A number of bands that does **not satisfy the compatibility relations** cannot be separated from other bands and describes a **semimetal**

# Check Topology

[www.cryst.ehu.es/cryst/checktopologicalmat](http://www.cryst.ehu.es/cryst/checktopologicalmat)

Bilbao Crystallographic Server → Check Topological Mat.

## Result of the analysis:

- The material is a topological insulator.
- List of topological indices:

$z_{2w,1}=0$

$z_{2w,2}=0$

$z_{2w,3}=0$

$z_4=0$

$z_2=0$

$z_8=4$

```
Number of electrons=10
Number of maximal k-vectors=4
Symbols of the k-vectors, number of bands up to fermi level and set of irreps.
  -GM      -X      -L      -W
    10      10      10      10
-GM6 (2) -X6 (2)  -L9 (2) -W6 (2)
-GM6 (2) -X6 (2)  -L8 (2) -W7 (2)
-GM8 (2) -X8 (2)  -L8 (2) -W7 (2)
-GM11 (4) -X8 (2) -L4-L5 (2) -W6 (2)
              -X9 (2)  -L9 (2) -W7 (2)
```

- The material belongs to the strong topological class: 1
- Clicking on [See the irreps](#) you can see the details about the number of bands and the ider
- The set of bands can be put as linear combination of Elementary Band Representations (l to get some possible linear combinations of EBRs and partial EBRs.
- Click on [Subgroups](#) to check the topological character of the structure in each of its (trans



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