Modeling suspected spin liquids on the pyrochlore lattice NaCaNi2F7 and MgCr2O4

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11/20/2019 KITP

Spin liquids

- Exactly solvable models
 - Topological degeneracy of the ground state
 - Long-range quantum entanglement
 - Fractionalized excitations
- Material candidates and experiments
 - Transport
 - Spectroscopy
- This talk
 - Spin-1 NaCaNi2F7 and spin-3/2 MgCr2O4
 - Antiferromangetic Heisenberg model on a pyrochlore lattice
 - Neutron scattering probe of spin correlations
 - Modeling the spin dynamics

References

Article | Published: 29 October 2018

Continuum of quantum fluctuations in a threedimensional S = 1 Heisenberg magnet

K. W. Plumb 🔀, Hitesh J. Changlani, A. Scheie, Shu Zhang, J. W. Krizan, J. A. Rodriguez-Rivera, Yiming Qiu, B. Winn,

R. J. Cava & <u>C. L. Broholm</u>

Nature Physics 15, 54–59 (2019) | Download Citation ⊻

arXiv.org > cond-mat > arXiv:1810.09481

Condensed Matter > Strongly Correlated Electrons

Dynamical structure factor of the three-dimensional quantum spin liquid candidate NaCaNi₂F₇

Shu Zhang, Hitesh J. Changlani, Kemp W. Plumb, Oleg Tchernyshyov, Roderich Moessner (Submitted on 22 Oct 2018)

arXiv.org > cond-mat > arXiv:1810.11869

Condensed Matter > Strongly Correlated Electrons

Magnetic excitations of the classical spin liquid MgCr2O4

X. Bai, J. A. M. Paddison, E. Kapit, S. M. Koohpayeh, J.-J. Wen, S. E. Dutton, A. T. Savici, A. I. Kolesnikov, G. E. Granroth, C. L. Broholm, J. T. Chalker, M. Mourigal (Submitted on 28 Oct 2018)

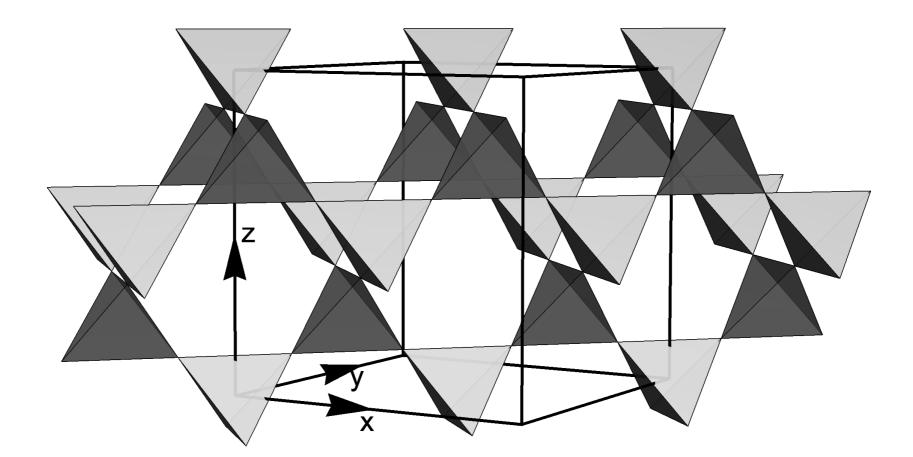
Outline

- Heisenberg pyrochlore antiferromagnet:
 - Classical model: A classical spin liquid
 - Quantum model: A quantum spin liquid?
 - Material realizations
 - NaA'B2F7 (A' = Ca, B = Ni spin-1; A' = Sr, B = Mn spin-5/2)
 - ACr2O4 (A = Cd; Zn; Mg spin-3/2; ...)
 - •
- A general picture for spin dynamics
 - Theoretical modeling of NaCaNi2F7 and MgCr2O4
- Discussion

Heisenberg pyrochlore antiferromagnet

Pyrochlore: corner-sharing tetrahedra

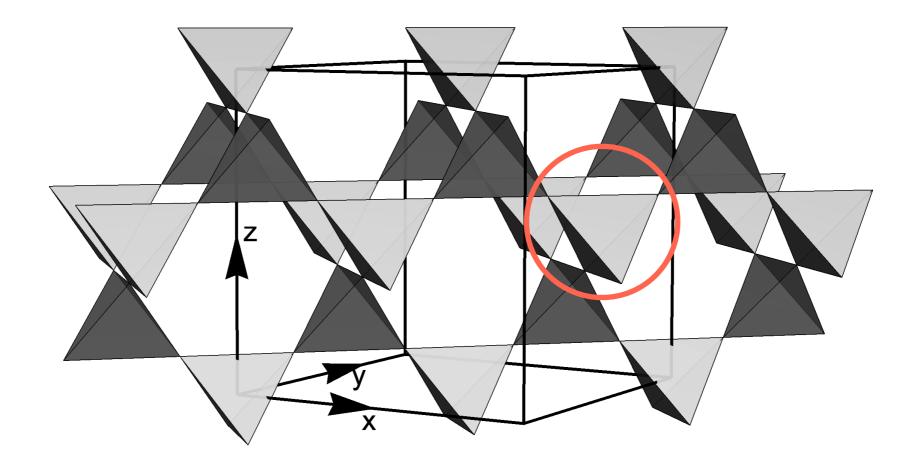
Antiferromagnetic Heisenberg interaction $H = J \sum_{\langle ij \rangle} \mathbf{s}_i \cdot \mathbf{s}_j, \quad J > 0$



Heisenberg pyrochlore antiferromagnet

Pyrochlore: corner-sharing tetrahedra

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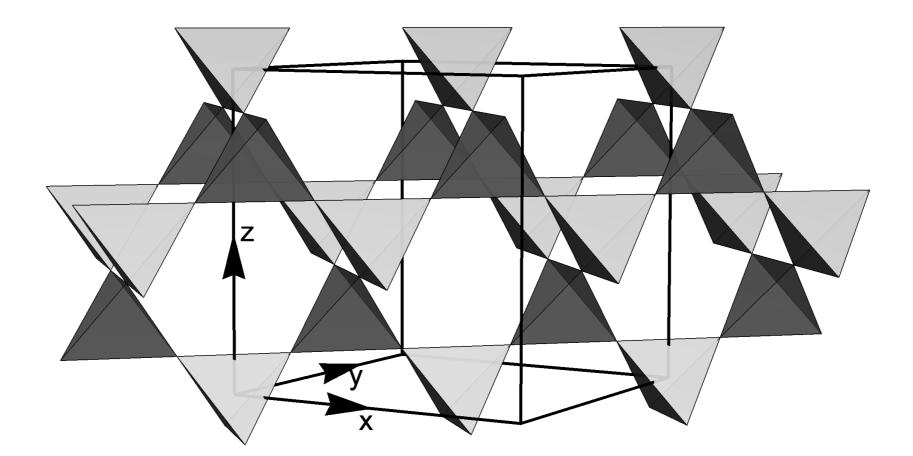


 $H_{\boxtimes} = \frac{J}{2} (\sum_{i \in \boxtimes} \mathbf{s}_i)^2 + \text{const.} \implies \sum_{i \in \boxtimes} \mathbf{s}_i = 0 \quad \text{Magnetic frustration}$ Moessner and Chalker 1998

Heisenberg pyrochlore antiferromagnet

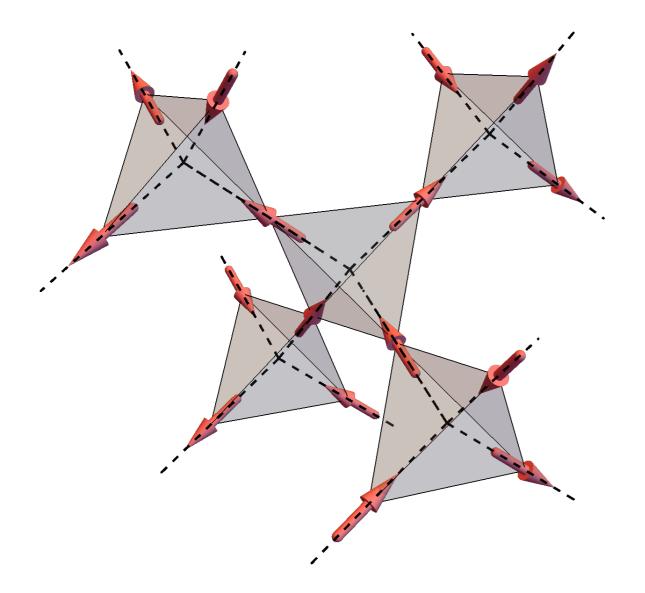
Pyrochlore: corner-sharing tetrahedra

Antiferromagnetic Heisenberg interaction $H = J \sum_{\langle ij \rangle} \mathbf{s}_i \cdot \mathbf{s}_j, \quad J > 0$



Searching for the Coulomb phase...

Spin ice



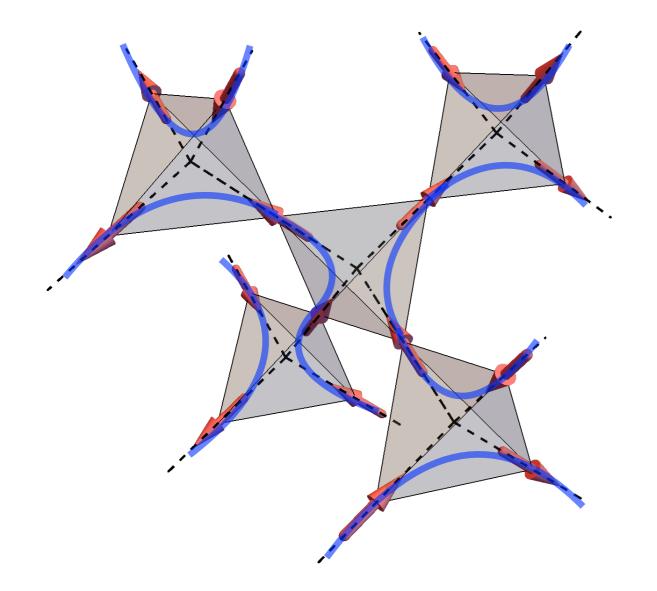
$$H = J \sum_{\langle ij \rangle} \sigma_i \sigma_j, \quad J > 0$$
$$\sigma = \pm 1$$

 $\sum_{i\in\boxtimes}\sigma_i=0$

Ice rule: two in two out

Gingras and McClarty 2014

Spin ice



Divergence free magnetic field

 $\nabla \cdot \mathbf{B} = 0$

Magnetostatics

$$\mathcal{Z}[\mathbf{B}(\mathbf{r})] = \int \mathcal{D}(\mathbf{B}) \exp\left[-\frac{\beta}{8\pi} \int d^3 \mathbf{r} \ \mathbf{B}^2\right] \qquad \beta = 1/k_B T$$

 $\nabla \cdot \mathbf{B} = 0$ $\mathbf{B} = \nabla \times \mathbf{A}$ Vector potential

Momentum space

$$\langle B_{\mu}(\mathbf{q})B_{\nu}(-\mathbf{q})\rangle \propto \delta_{\mu\nu} - \frac{q_{\mu}q_{\nu}}{q^2}$$

Magnetostatics

$$\mathcal{Z}[\mathbf{B}(\mathbf{r})] = \int \mathcal{D}(\mathbf{B}) \exp\left[-\frac{\beta}{8\pi} \int d^3 \mathbf{r} \ \mathbf{B}^2\right] \qquad \beta = 1/k_B T$$

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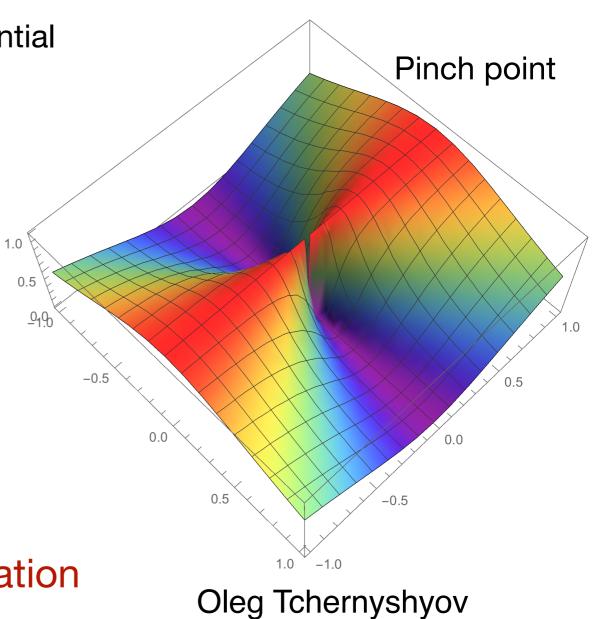
Momentum space

$$\left. \left\langle B_{\mu}(\mathbf{q}) B_{\nu}(-\mathbf{q}) \right\rangle \propto \delta_{\mu\nu} - \frac{q_{\mu}q_{\nu}}{q^2} \right. \\ \left. \left\langle B_z(\mathbf{q}) B_z(-\mathbf{q}) \right\rangle \right|_{q_x=0} \propto \frac{q_y^2}{q_z^2 + q_y^2} \right.$$

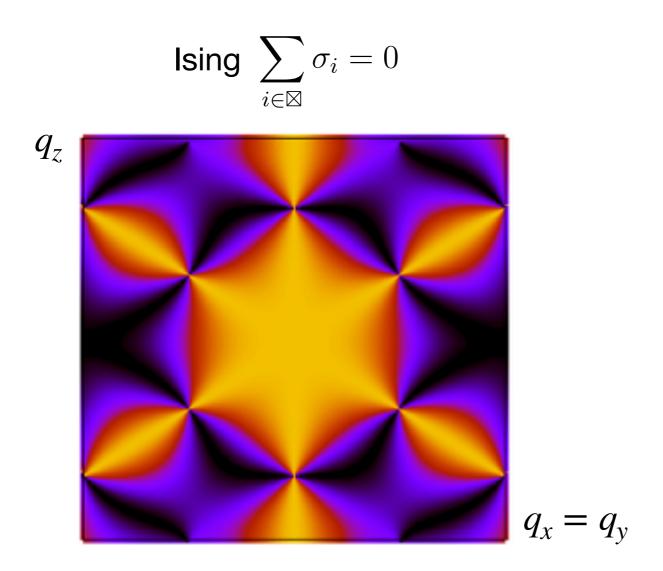
Real space

$$\langle B_{\mu}(\mathbf{r})B_{\nu}(0)\rangle \propto \frac{3r_{\mu}r_{\nu}-\delta_{\mu\nu}r^2}{r^5}$$

Dipolar correlation



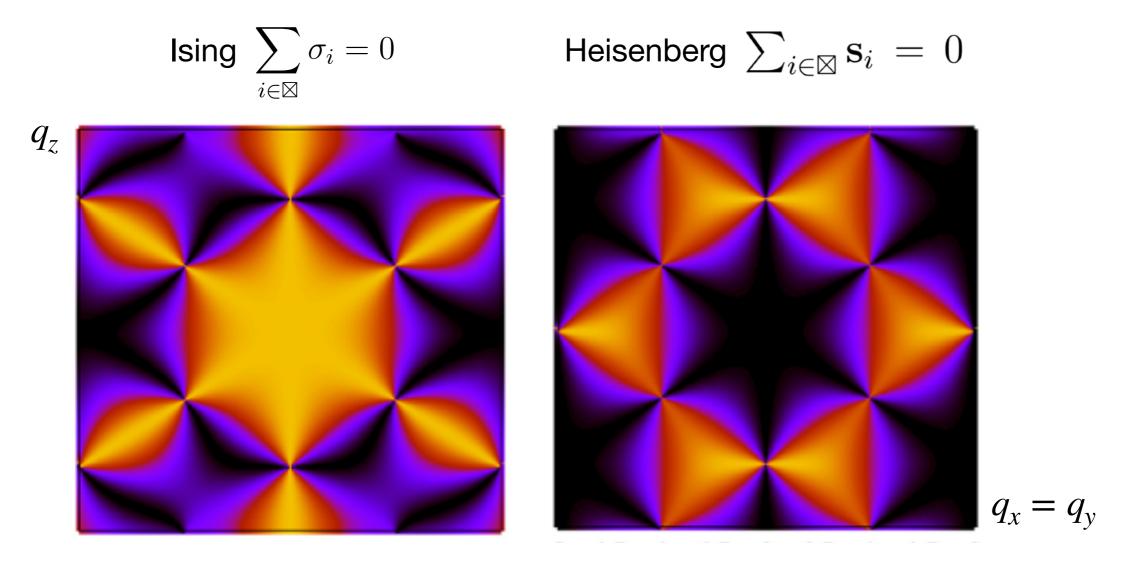
Coulomb phase



Spin correlation with power-law decay: Disordered but strongly correlated

Canals and Lacroix 1998 Isakov et. al. 2004 Henley 2005

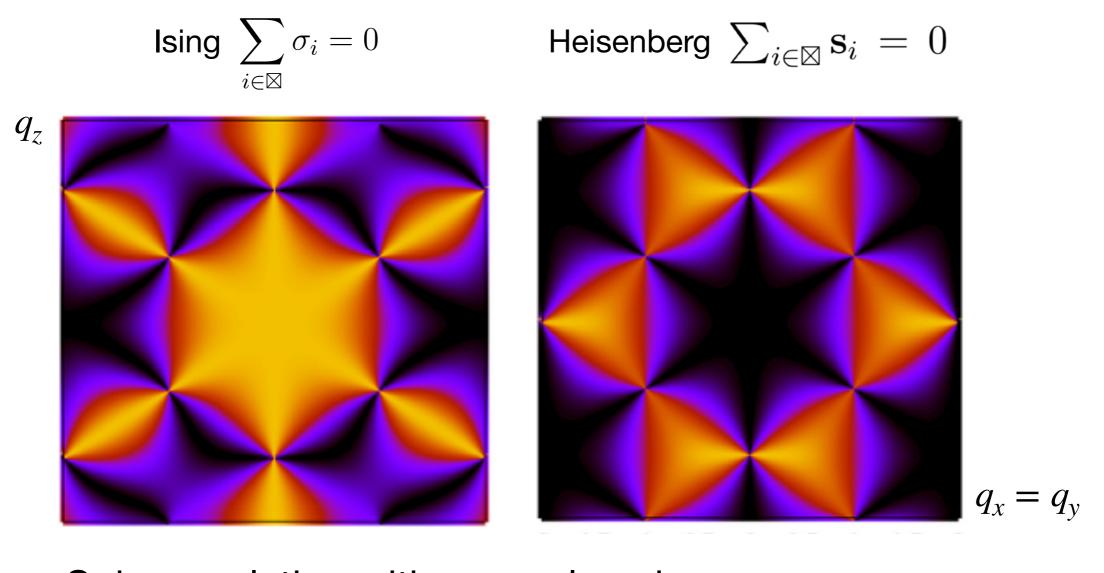
Coulomb phase



Spin correlation with power-law decay: Disordered but strongly correlated

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Spin correlation with power-law decay: Disordered but strongly correlated

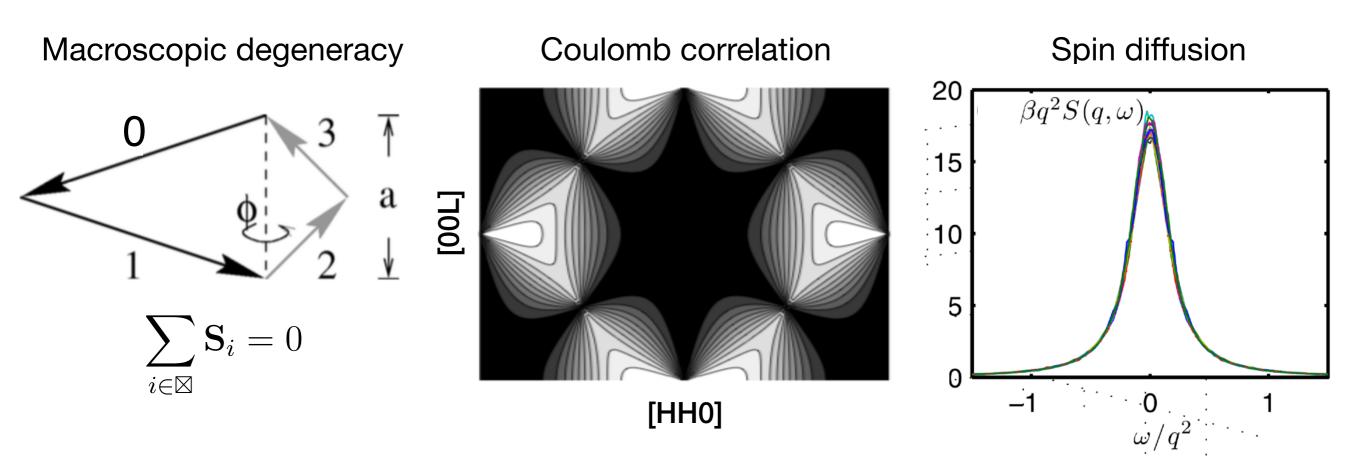
Measured by neutron scattering

Canals and Lacroix 1998 Isakov et. al. 2004 Henley 2005

A classical spin liquid

Heisenberg pyrochlore antiferromagnet

$$H = J \sum_{\langle ij \rangle} \mathbf{s}_i \cdot \mathbf{s}_j, \quad J > 0$$

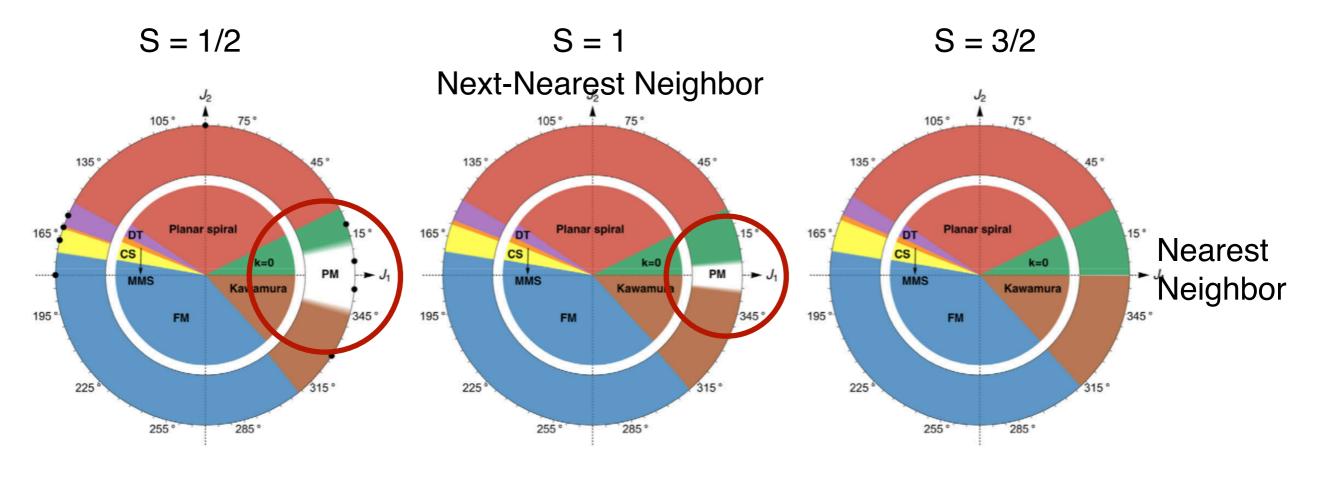


Moessner Chalker 1999 Isakov Greger Moessner Sondhi 2004 Conlon Chalker 2009

A quantum spin liquid?

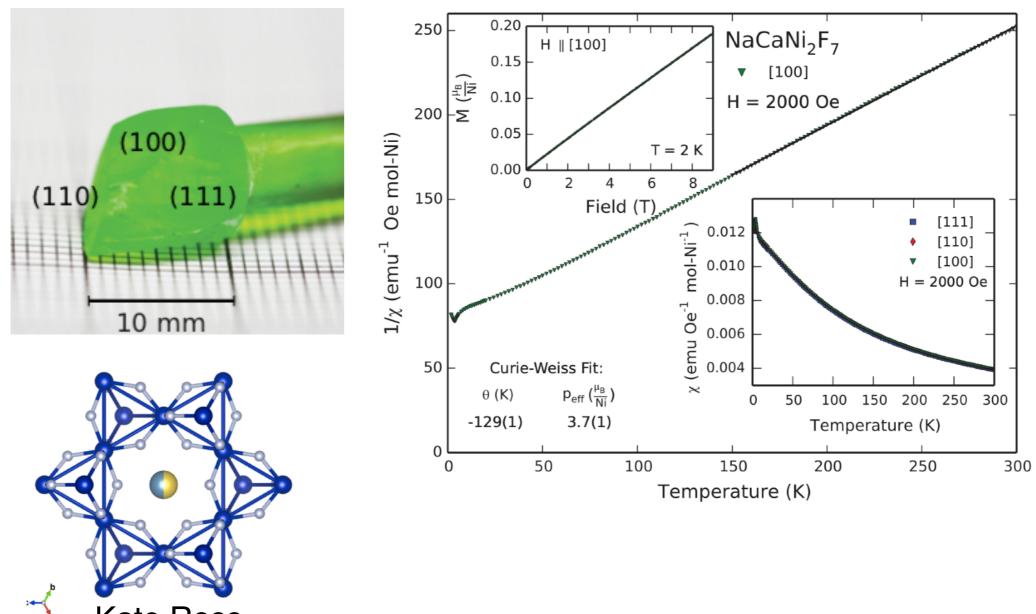
Heisenberg pyrochlore antiferromagnet (quantum version)

$$H = J \sum_{\langle ij \rangle} \mathbf{s}_i \cdot \mathbf{s}_j, \quad J > 0$$



lqbal et.al. 2019 PFRG

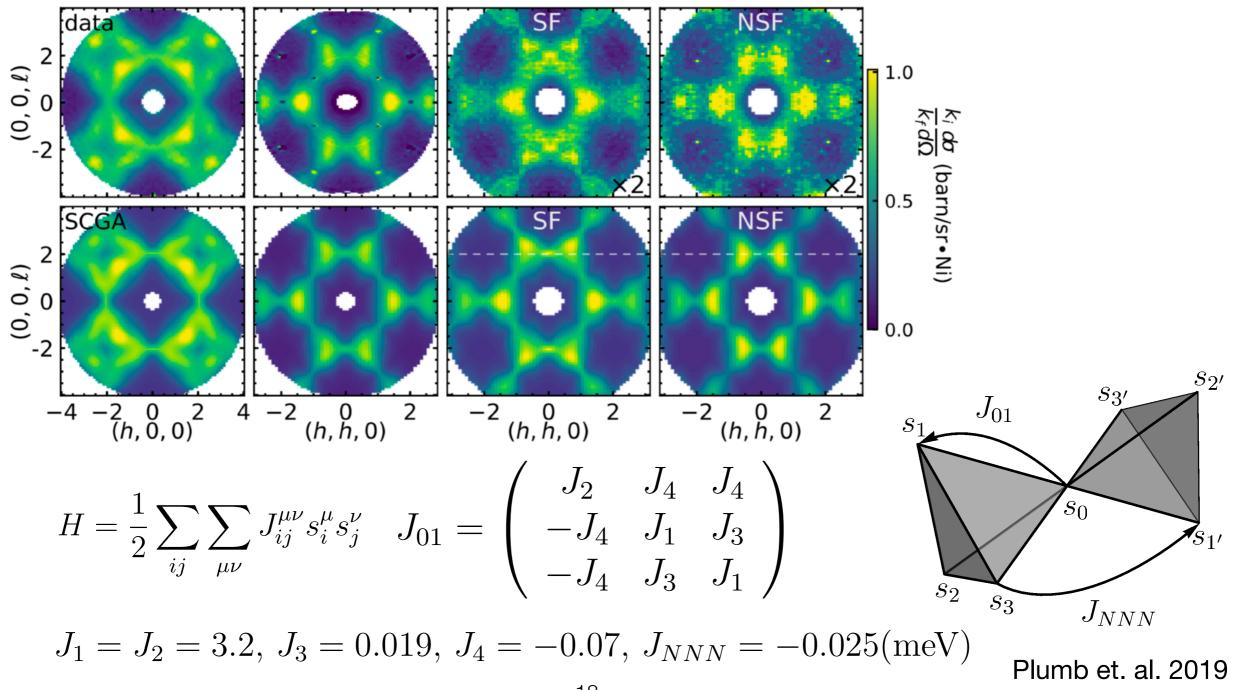
• Ni²⁺ on one pyrochlore lattice and Na¹⁺/Ca²⁺ on another



Kate Ross

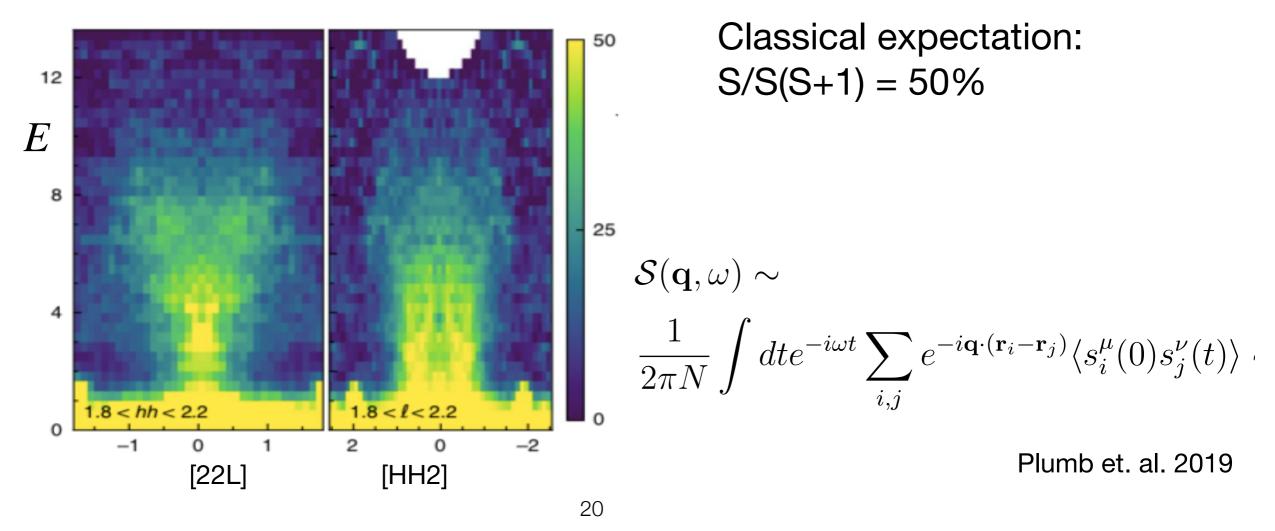
Krizan and Cava 2015

- Ni²⁺ on one pyrochlore lattice and Na⁺/Ca²⁺ on another
- Heisenberg interaction

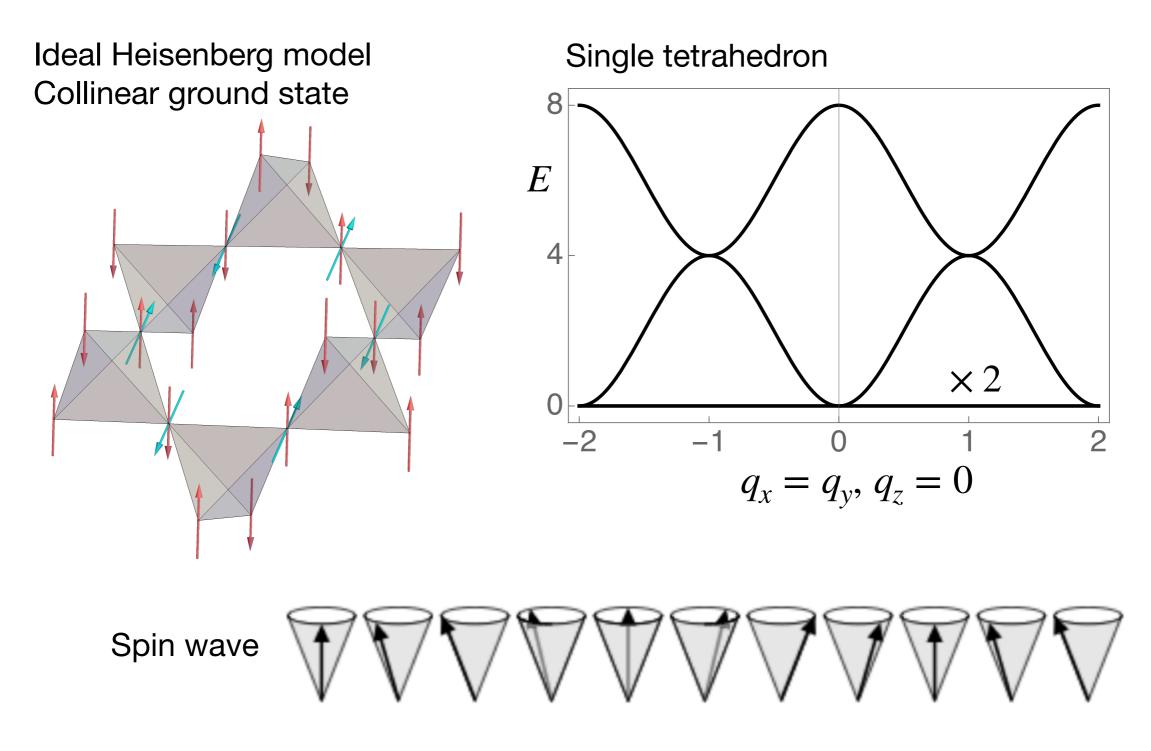


- Ni²⁺ on one pyrochlore lattice and Na⁺/Ca²⁺ on another
- Heisenberg interaction
- Spin-1 system

- Ni²⁺ on one pyrochlore lattice and Na⁺/Ca²⁺ on another
- Heisenberg interaction
- Spin-1 system
- No magnetic ordering down to 0.35K $J_1 = J_2 = 3.2 \text{ (meV)}$
- ~90% spectral weight from inelastic scattering at 1.8K



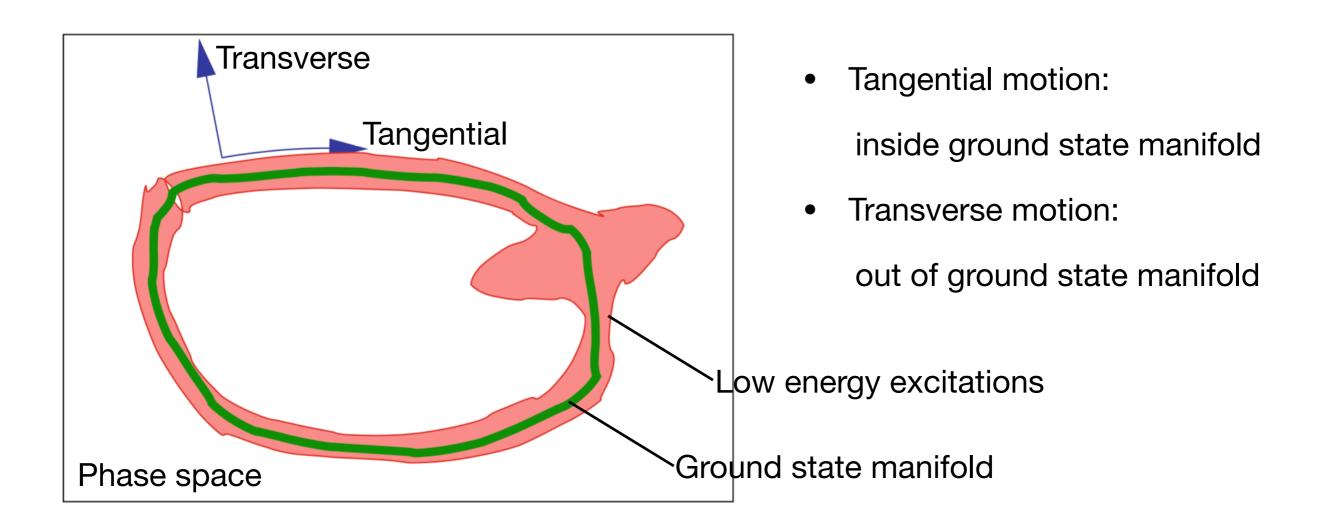
An example of spin dynamics



Canals and Lacroix 1998 Tchernyshyov, Moessner, and Sondhi 2002

Spin dynamics of NaCaNi₂F₇

- Ground states: many degenerate/nearly-degenerate disordered states
- Fluctuations: finite-frequency spin waves
- Dynamical picture: spin waves drive the motion between ground states



- Stochastic Large-n model
 - Tangential motion: slow
 - Transverse motion: fast

- Brownian motion
$$\frac{d\mathbf{v}}{dt} = \frac{1}{m}\mathbf{F}(t) + \boldsymbol{\xi}(t)$$

White noise

- Stochastic Large-n model
 - Tangential motion: slow
 - Transverse motion: fast

• Brownian motion
$$\frac{d\mathbf{v}}{dt} = \frac{1}{m}\mathbf{F}(t) + \boldsymbol{\xi}(t)$$

• To conserve spins (angular momentum)

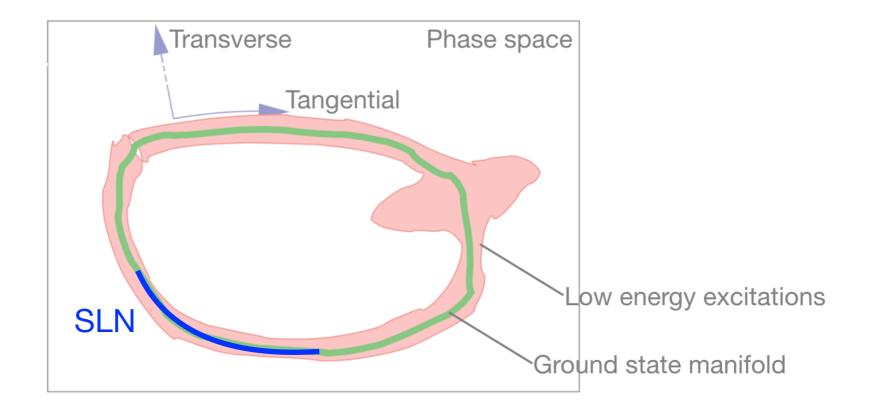
$$\frac{d\mathbf{s}}{dt} = -\gamma \nabla \cdot \mathbf{j}(t) + \boldsymbol{\xi}(t)$$
$$\mathbf{j}(t) = -\nabla \frac{\partial E}{\partial \mathbf{s}}$$

• Discrete lattice

Generalized force

Garanin and Canals 1999 Conlon and Chalker 2009

- Stochastic Large-n model
- Long time scale
- Relaxation



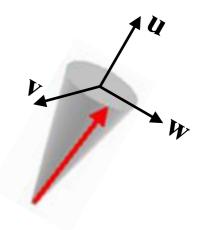
- Stochastic Large-n model
- Long time scale
- Relaxation

• Linear spin wave theory

$$\hbar \frac{d}{dt} \mathbf{s}_i = -\mathbf{s}_i \times \mathbf{H}_{\text{eff}}$$

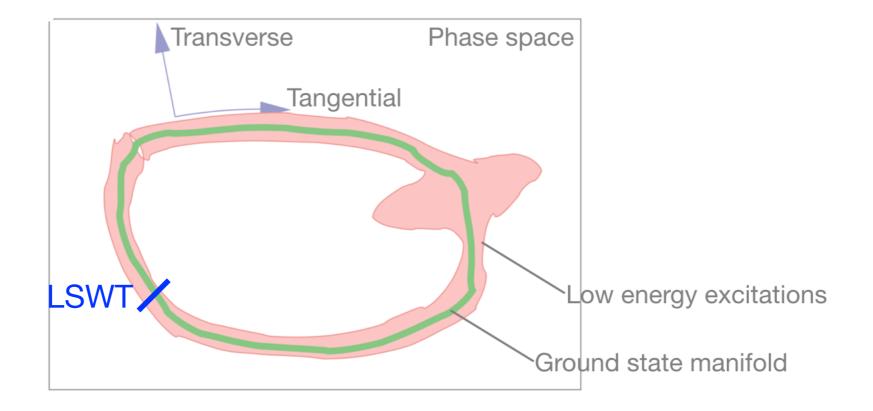
Linearized for small deviations from a ground state

$$\mathbf{s}_{i} = \sqrt{S^{2} - S(x_{i}^{2} + y_{i}^{2})} \mathbf{u}_{i} + \sqrt{S}(x_{i}\mathbf{v}_{i} + y_{i}\mathbf{w}_{i})$$
$$\approx \left(S - \frac{x_{i}^{2} + y_{i}^{2}}{2}\right)\mathbf{u}_{i} + \sqrt{S}(x_{i}\mathbf{v}_{i} + y_{i}\mathbf{w}_{i}).$$

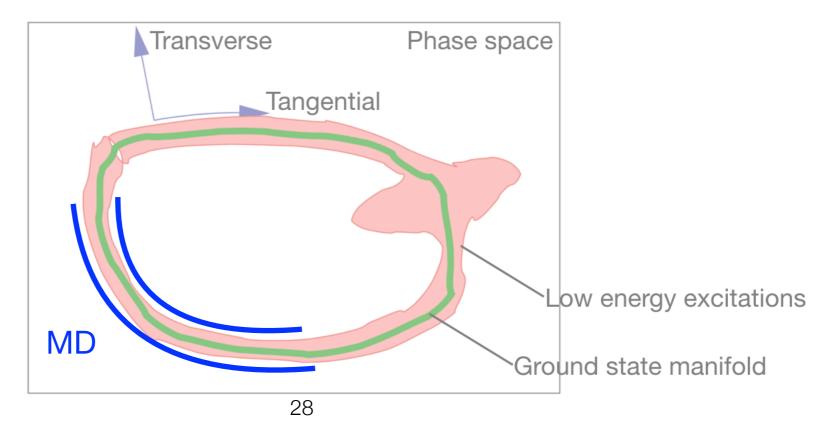


Walker and Walstedt 1980 Savary and Balents 2014

- Stochastic Large-n model
- Long time scale
- Relaxation
- Linear spin wave theory
- Finite-frequency excitations
- Around one ground state



- Stochastic Large-n model
- Long time scale
- Relaxation
- Linear spin wave theory
- Finite-frequency excitations
- Around one ground state
- Molecular dynamics
- Full simulation
- Short + long time scale



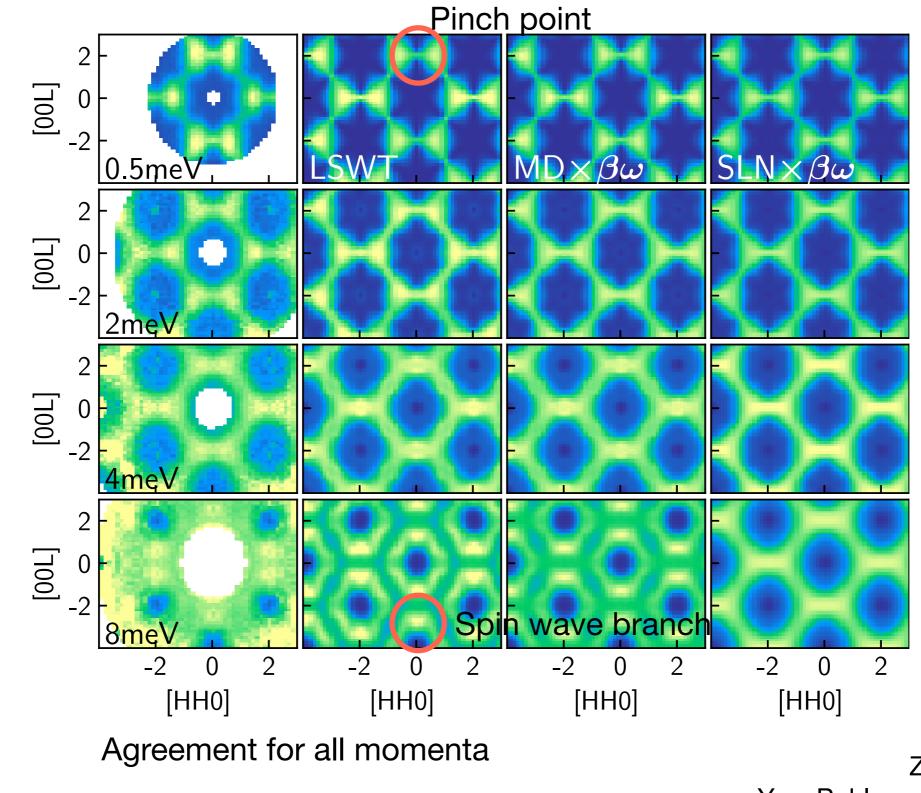
- Stochastic Large-n model
- Long time scale
- Relaxation
- Linear spin wave theory
 Zero temperature
- Finite-frequency excitations
- Around one ground state

- Molecular dynamics
 Thermal excitation
- Full simulation
- Short + long time scale

Considering quantum statistics of spin waves:

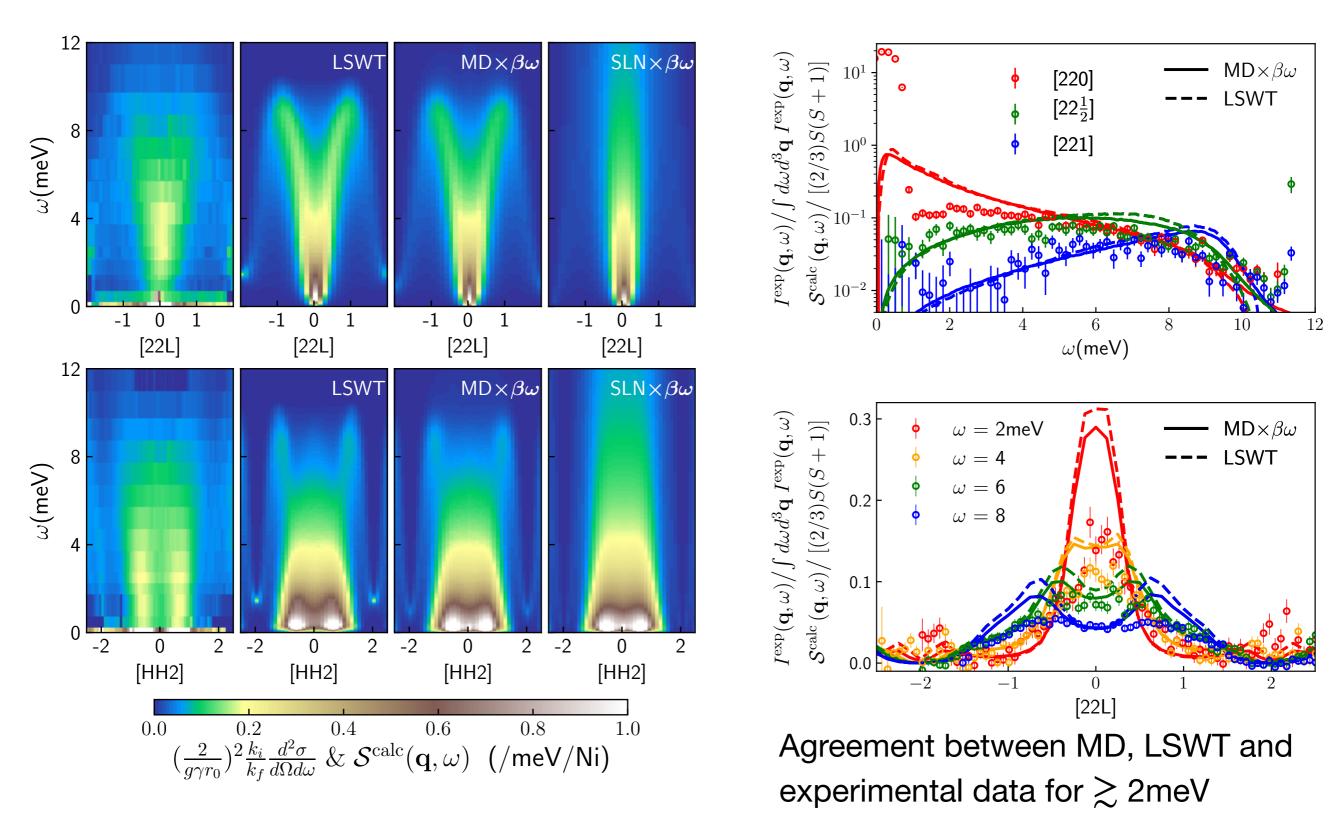
$$\beta \omega S_{\text{MD}}(\mathbf{q}, \omega) \big|_{\text{finite T}} = S_{\text{LSWT}}(\mathbf{q}, \omega) \big|_{\text{T}=0}$$
for $\beta \omega \gg 1$

Theoretical modeling



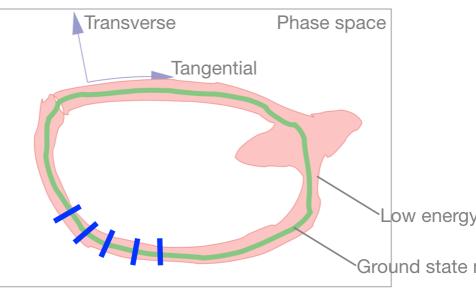
Zhang et. al. 2019 Yan, Pohle, and Shannon 2018

Theoretical modeling

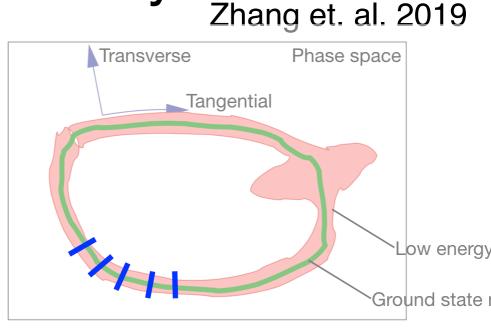


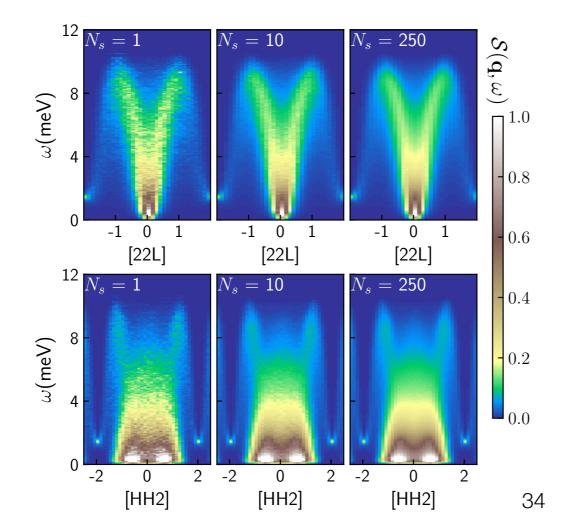
- Not considered
 - Lack of long-range order
 - Quantum renormalization
 - Thermal broadening
 - Multimagnon continuum

- Correct statistics
 - Fluctuations of disordered ground states
 - Averaged among similar spectra

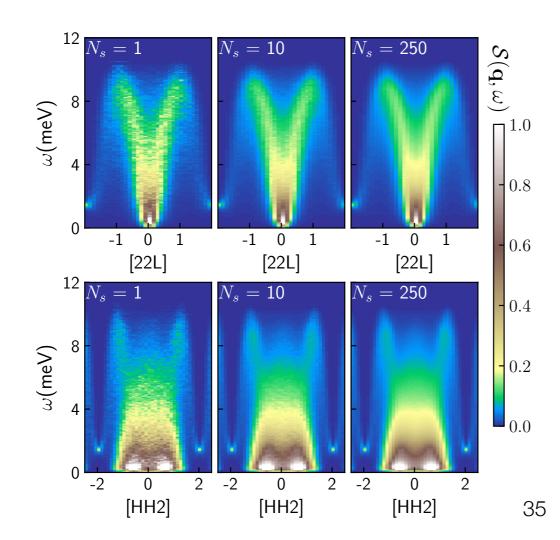


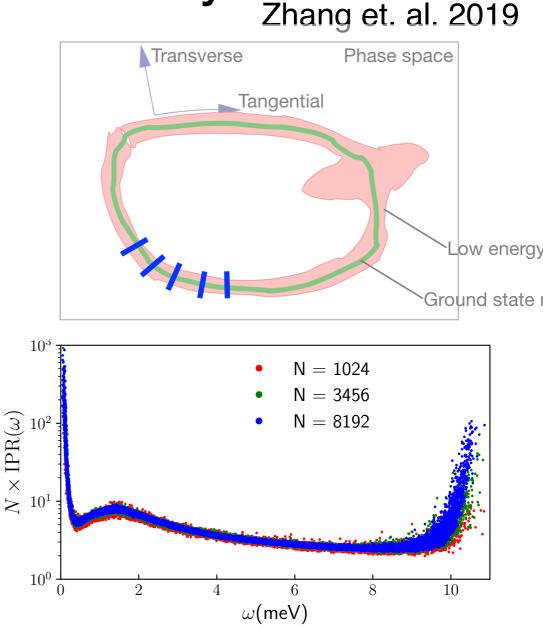
- Correct statistics
 - Fluctuations of disordered ground states
 - Averaged among similar spectra
 - Without well-defined momenta





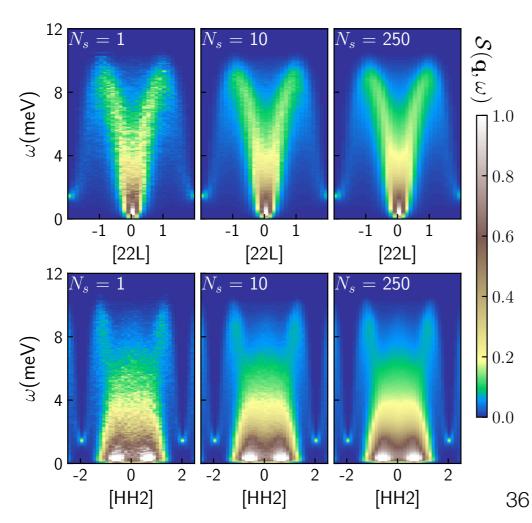
- Correct statistics
 - Fluctuations of disordered ground states
 - Averaged among similar spectra
 - Without well-defined momenta
 - Delocalized

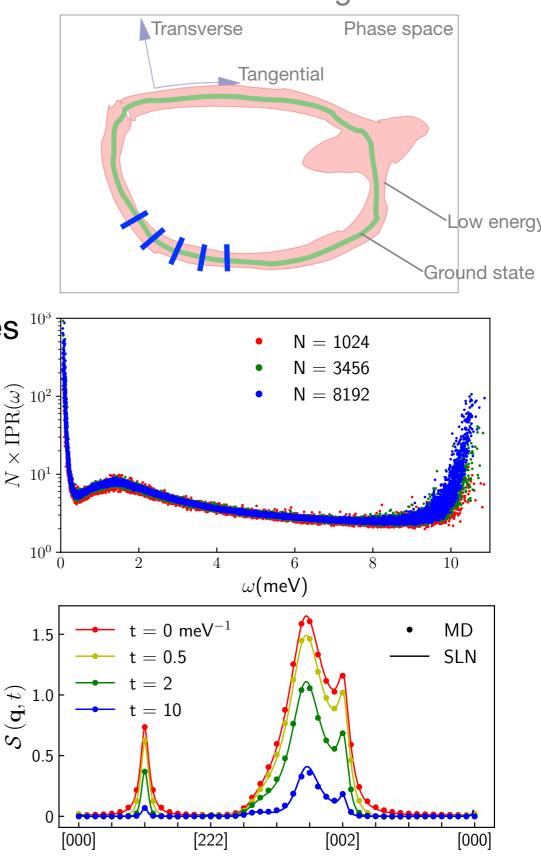




Zhang et. al. 2019

- Correct statistics
 - Fluctuations of disordered ground states
 - Averaged among similar spectra
 - Without well-defined momenta
 - Delocalized
 - Driving the motion between ground states



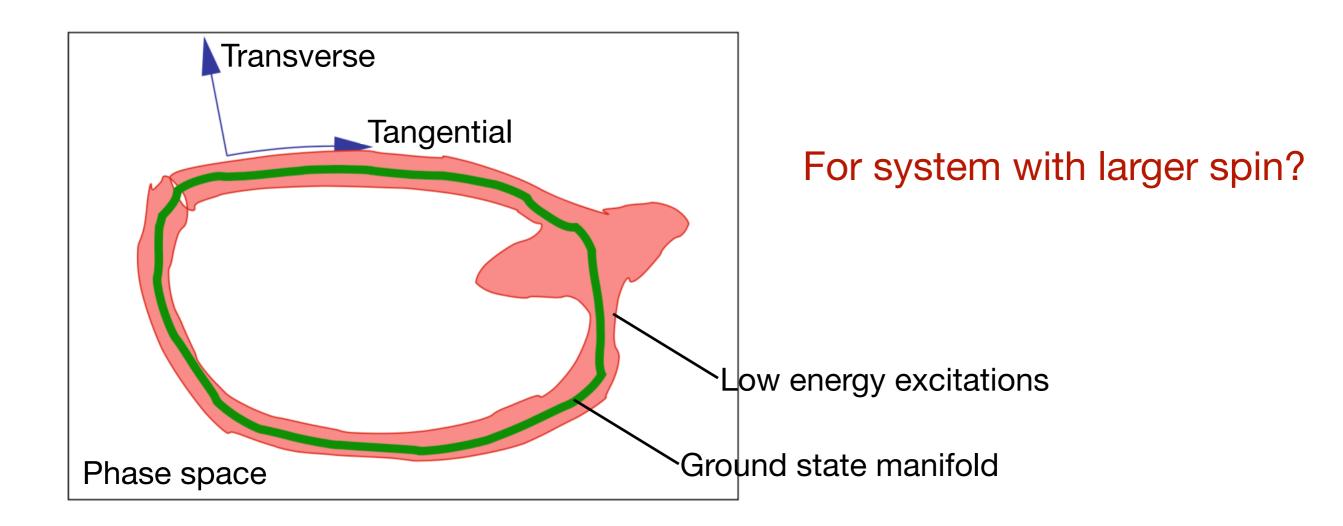


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Spin dynamics

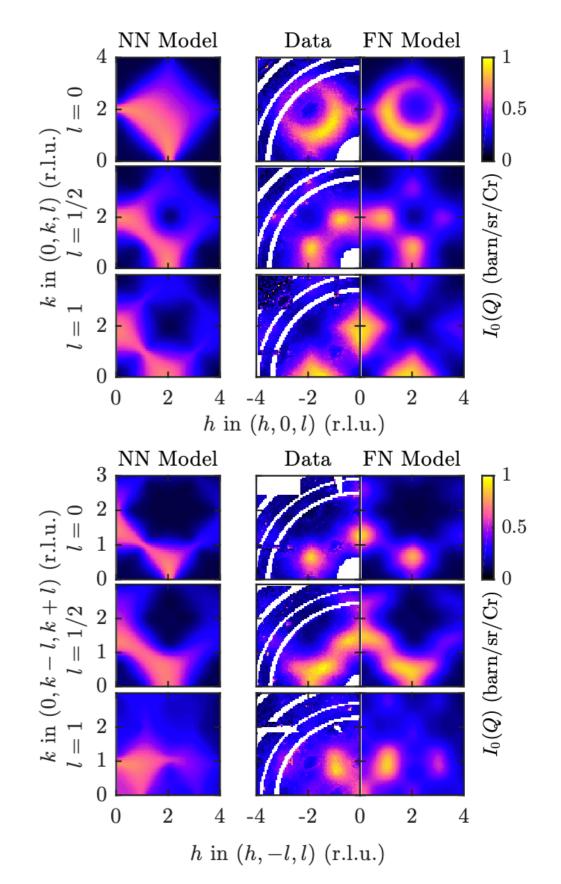
Zhang et. al. 2019

- Ground states: many degenerate/nearly-degenerate disordered states
- Fluctuations: finite-frequency spin waves
- Dynamical picture: spin waves drive the motion between ground states



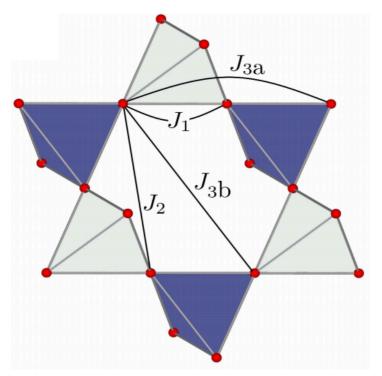
Approximate classical spin liquid MgCr₂O₄

- Cr³⁺ on pyrochlore lattice
- Spin-3/2 system
- Magnetostructural transition at 13K
- Broad scattering patterns at 20K



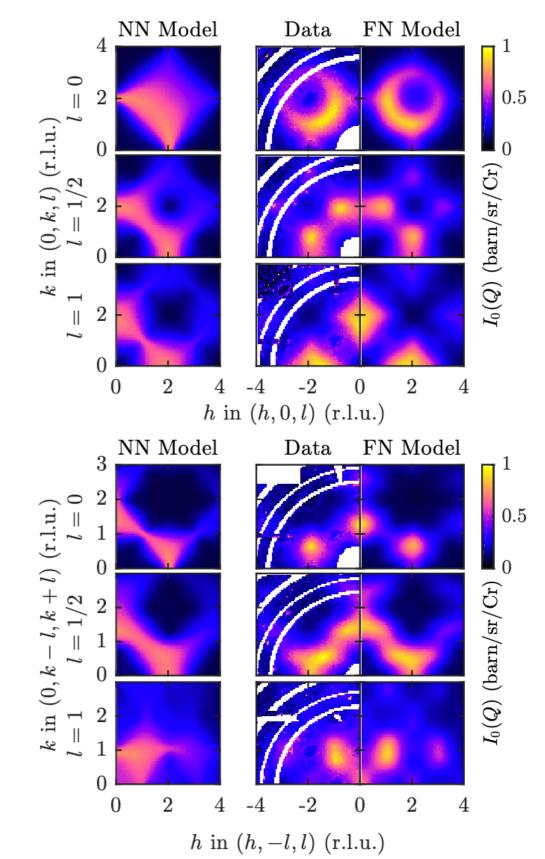
Approximate classical spin liquid MgCr₂O₄

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- Heisenberg interaction

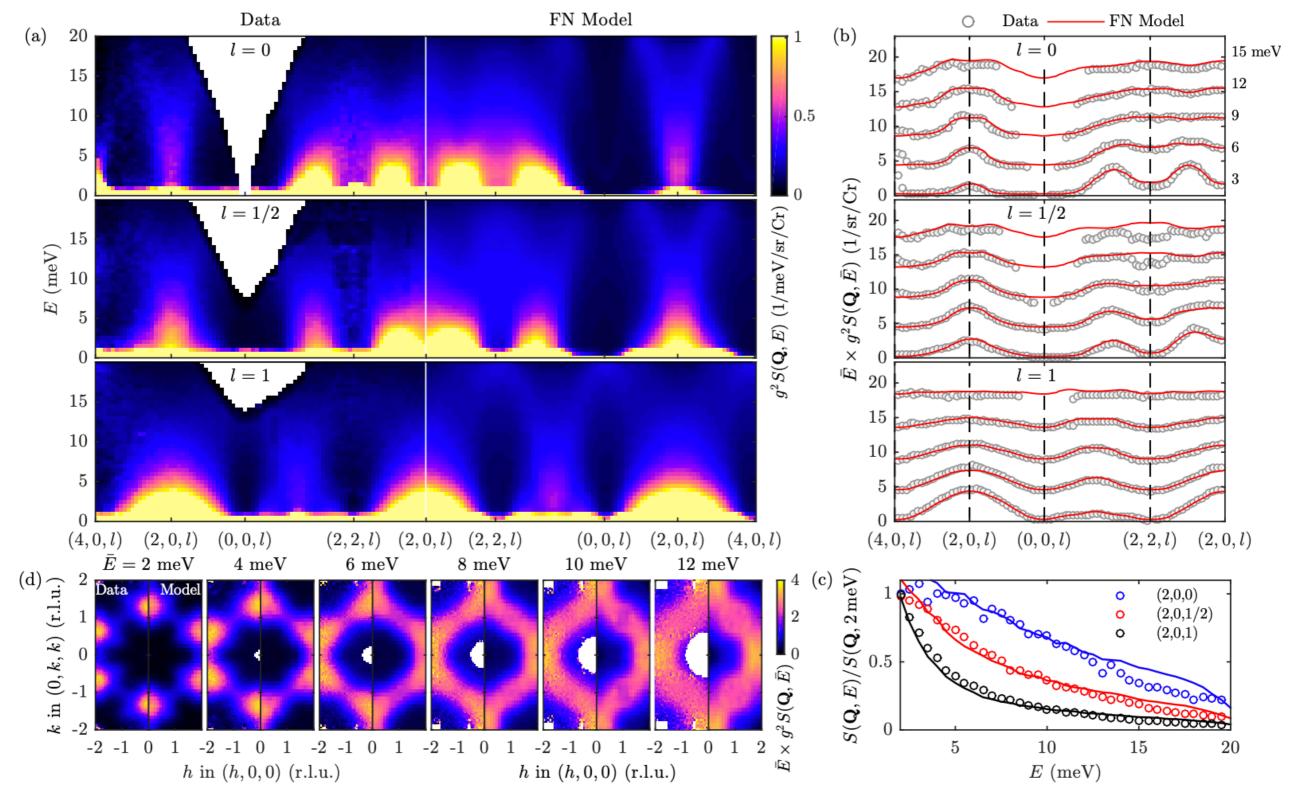


$$J_1 = 3.3, J_2 = 0.27, J_{3a} = 0.34, J_{3b} = 0.03 (\text{meV})$$

Bai et. al. 2019



Approximate classical spin liquid MgCr₂O₄



Bai et. al. 2019

Discussion

Linear spin wave theory as a good description for Heisenberg pyrochlore antiferromagnets



A peek into quantum spin liquids

with classical tools

Summary

- Heisenberg pyrochlore antiferromagnet
- Theoretical modeling
 - Stochastic large-n + linear spin wave theory + molecular dynamics
 - Spin-1 quantum spin liquid NaCaNi₂F₇
 - Spin-3/2 classical spin liquid MgCr₂O₄
- A general picture for spin dynamics
 - Slow & fast motion
- Quantum nature?

Acknowledgement

• Collaborators:

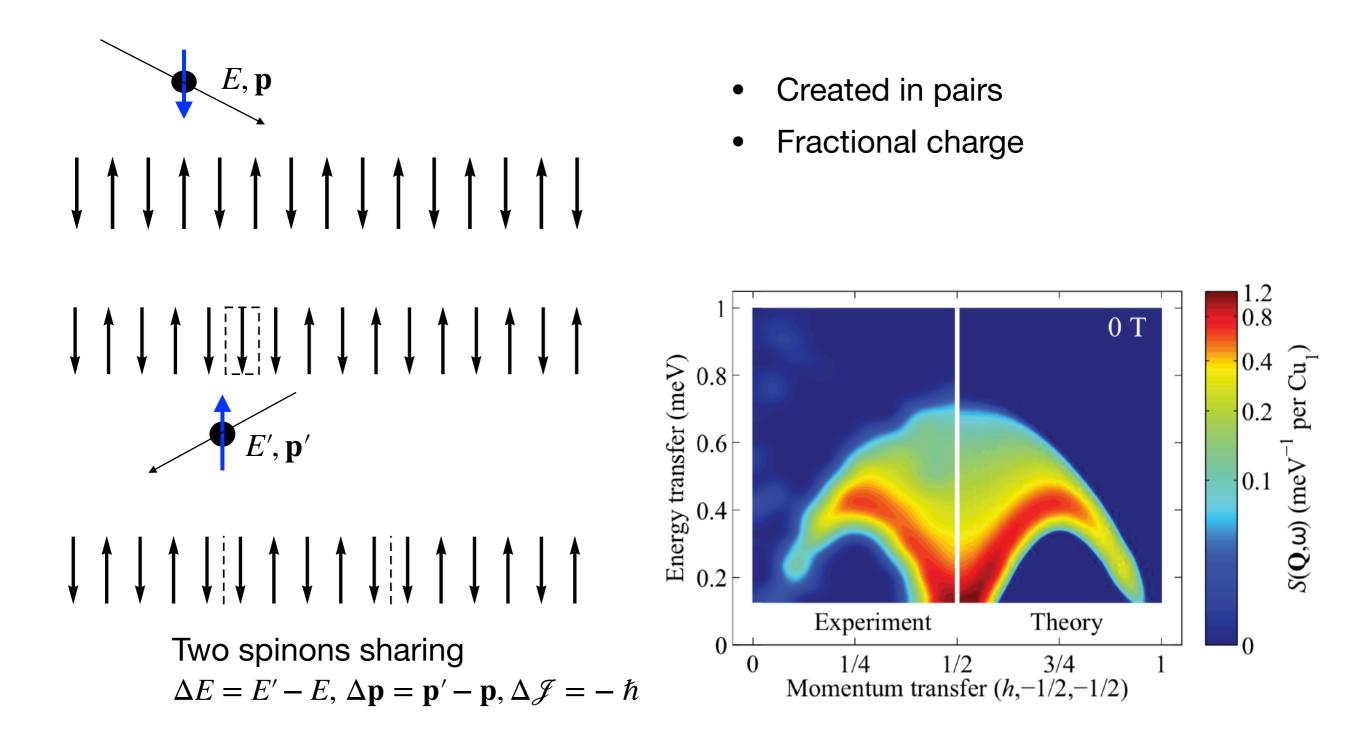
Hitesh J. Changlani (Florida state)
Kemp W. Plumb (Brown)
Oleg Tchernyshyov (JHU)
Roderich Moessner (MPIPKS)

• Funding:

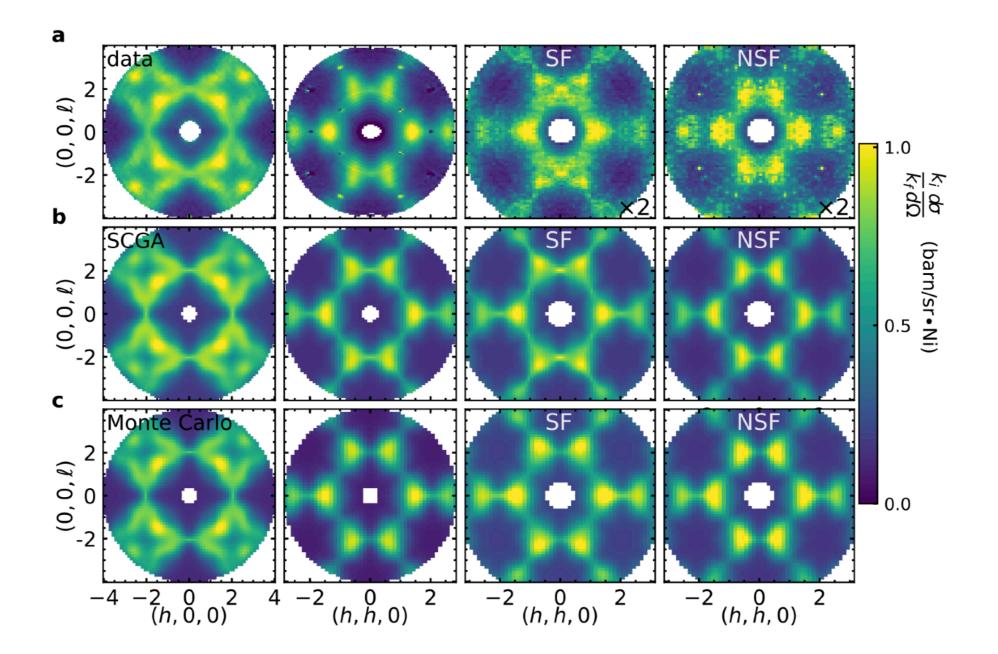
U.S. Department of Energy, Grant DE-FG02-08ER46544 Deutsche Forschungsgemeinschaft via grant SFB 1143

- Computation: HHPC&MARCC, Maryland
- Experiments: NIST&Oak Ridge

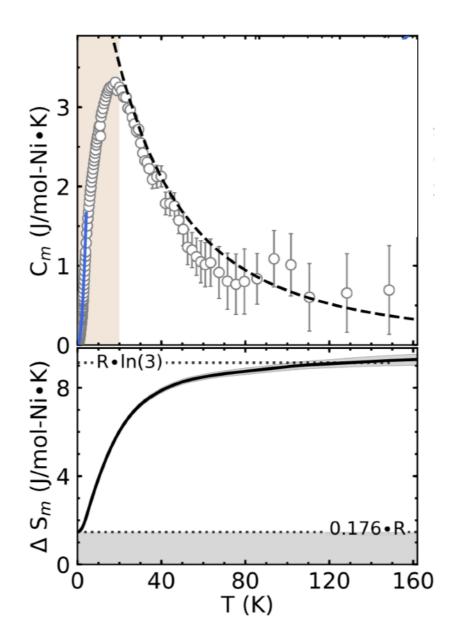
Fractionalized excitations



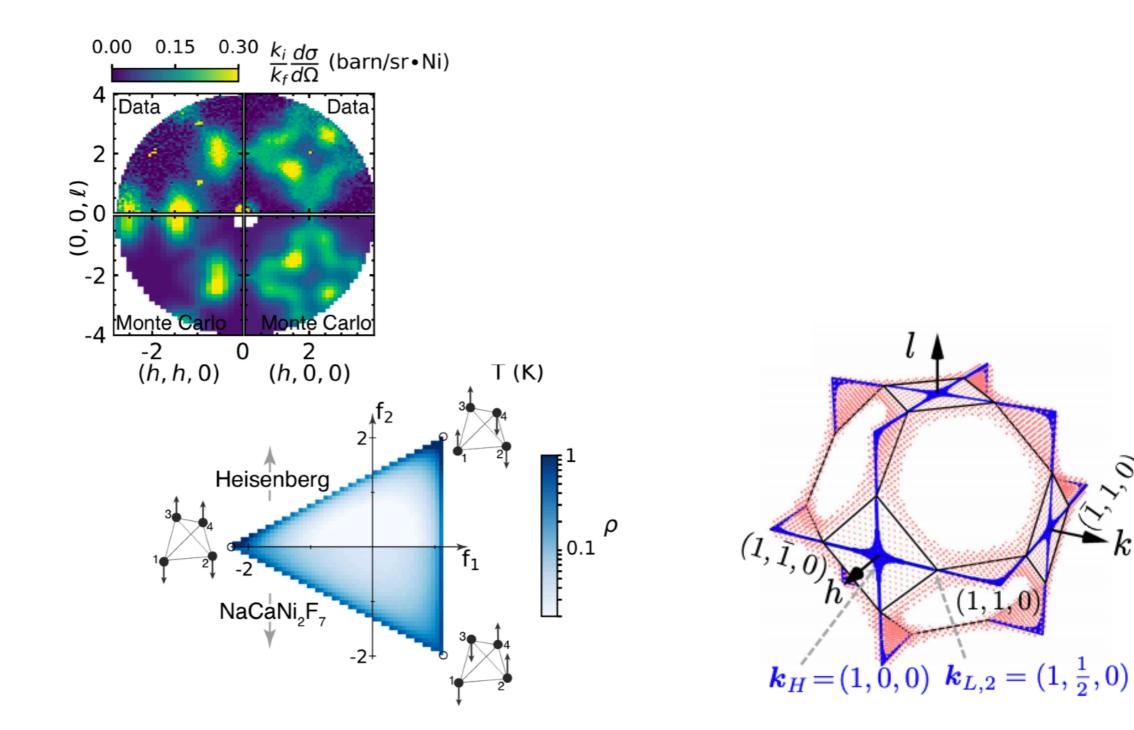
Fit the exchange parameters



Specific heat



Ground states



0

k