

# Expectations from GRB Afterglows

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A discussion based largely on:  
Levinson et al., 2002

# Afterglow Model

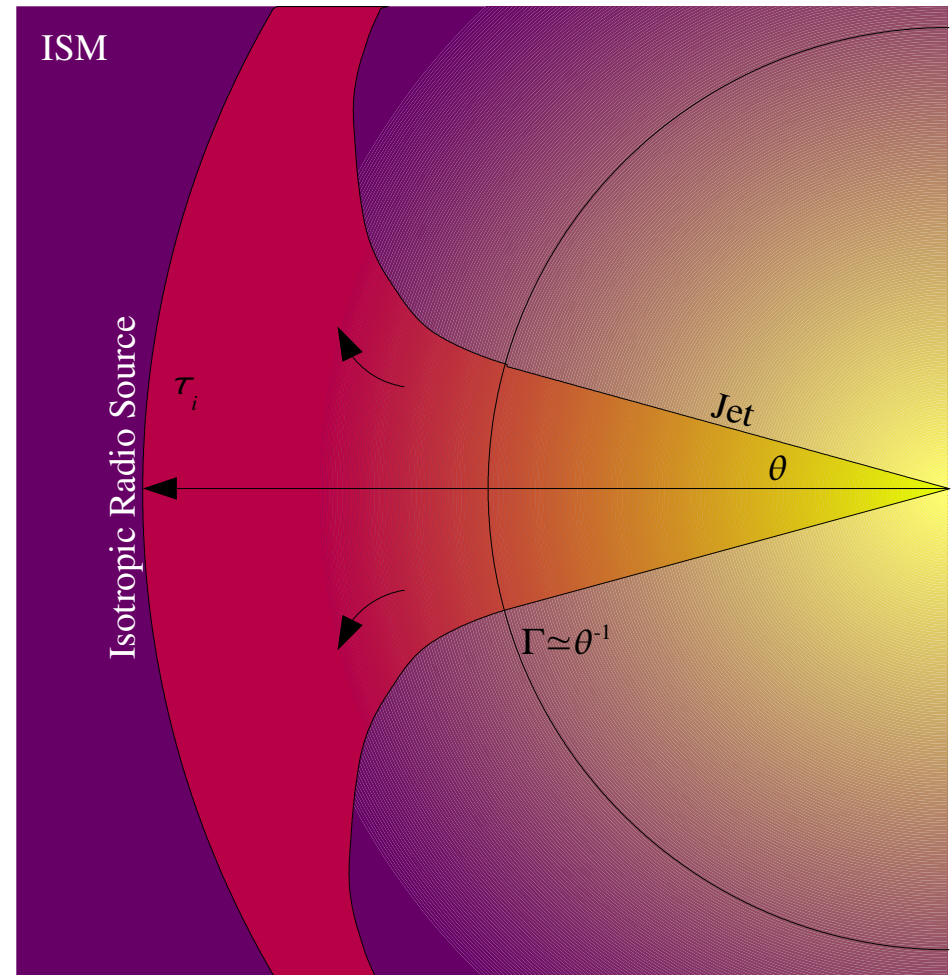
## The Basics

- “Ultrarelativistic, jetlike ejecta propagating into an ambient medium...”
- Non-relativistic once  $\Gamma \simeq \theta^{-1}$
- Source becomes isotropic

$$\tau_i \simeq 3$$

$$t \equiv \tau t_{\text{SNT}}(1 + z)$$

$$t_{\text{SNT}} \simeq 6 \times 10^6 (E_{51}/n_0)^{1/3} \text{S}$$



# Afterglow Model

## Radio Emission

- After the transition to a non-relativistic jet, the power is dominated by synchrotron emission.

- ✧ Peak synchrotron frequency,

$$\nu_{\star} \approx 1 \left( \frac{1+z}{2} \right)^{-1} \left( \frac{\xi_e}{0.3} \right)^2 \left( \frac{\xi_B}{0.3} \right)^{1/2} n_0^{1/2} \text{ GHz}$$

- ✧ Observed flux above  $\nu_{\star}$ ,

$$f_{\nu} \approx 2h^2 \left( \frac{1+z}{2} \right)^{1/2} \left( \frac{\xi_e}{0.3} \right) \left( \frac{\xi_B}{0.3} \right)^{3/4} \left( \frac{d_L}{R_0} \right)^{-2} n_0^{3/4} E_{51} \nu_9^{-1/2} \tau^{-9/10} \text{ mJy}$$

- GRB 970508 (Frail et al. 2000),  $n_0 \approx 1, \xi_e \sim \xi_B \sim \frac{1}{3}$

# Afterglow Model

## Observation Rate

- Age at which flux drops below detection limit,

$$\tau_m = \eta(1 + z)^{5/9} (d_L/R_0)^{-20/9}$$

$$\eta = h^{20/9} \left( \frac{f_{\nu, \min}}{1 \text{ mJy}} \right)^{-10/9} \left( \frac{\xi_e}{0.3} \right)^{10/9} \left( \frac{\xi_B}{0.3} \right)^{10/12} \left( \frac{2}{\nu_9} \right)^{5/9} n_0^{10/12} E_{51}^{10/9}$$

- Number of detectable radio afterglows,

$$\frac{dN_R}{dz} = N_0 d_L^2 \frac{dl}{dz} (\tau_m - \tau_i) \Phi(z) \quad N_0 = f_b^{-1} 4\pi R_0^3 \dot{n} t_{\text{SNT}}$$

- Given necessary constants and a cosmological model, we can calculate  $N_R$  by integrating over  $z$ !

# Afterglow Model

## A Example Case

- As an example, take,

$$f_{\nu,\min} = 5 \text{ mJy}, \tau_i = 3, h = 0.75, \dot{n} = .5$$

- So then,  $d_{L,\max} \simeq 0.2R_0 = 0.8 \text{ Gpc}$ ,  $z_{\max} \simeq 0.2$
- Thus we can neglect cosmological effects and we analytically integrate from  $z=0$  to  $z=z_{\max}$  to find,

$$N_R \simeq 18(500f_b)^{-1} \left(\frac{\dot{n}}{0.5}\right) \left(\frac{f_{\nu,\min}}{5 \text{ mJy}}\right)^{-3/2} \left(\frac{\xi_e}{0.3}\right)^{3/2} \left(\frac{\xi_B}{0.03}\right)^{9/8} \left(\frac{\tau_i}{3}\right)^{-7/20} n_{-1}^{19/24} E_{51}^{11/6} \nu_9^{-3/4}$$

## Afterglow Model

# Comments on this Example

$$N_R \simeq 18(500f_b)^{-1} \left(\frac{\dot{n}}{0.5}\right) \left(\frac{f_{\nu,\min}}{5 \text{ mJy}}\right)^{-3/2} \left(\frac{\xi_e}{0.3}\right)^{3/2} \left(\frac{\xi_B}{0.03}\right)^{9/8} \left(\frac{\tau_i}{3}\right)^{-7/20} n_{-1}^{19/24} E_{51}^{11/6} \nu_9^{-3/4}$$

- Weak dependence on  $\tau_i$
- $N_R$  decrease with increasing beam factor ( $\theta$ )
- Observations show that  $\xi_B=0.03$ ,  $n_0=0.1$ ,  $\dot{n}=0.5$  are conservative estimates
- No need to account for beamed sources

# Afterglow Model

## Numerical Results

- For deeper surveys, one must solve numerically,  
$$h = 0.75, \Phi(z) = (1 + z)^3$$

