#### Telling the time: Using the clockwork of turbulence to answer open questions in fluid dynamics

#### Colm-cille P. Caulfield

#### BP Institute & DAMTP, University of Cambridge



Tom S. Eaves (DAMTP now UBC) Rich R. Kerswell (Bristol) Dan Lucas (DAMTP soon Keele) Igor Mezic (UC Santa Barbara)



Engineering and Physical Sciences Research Council









•Motivation:

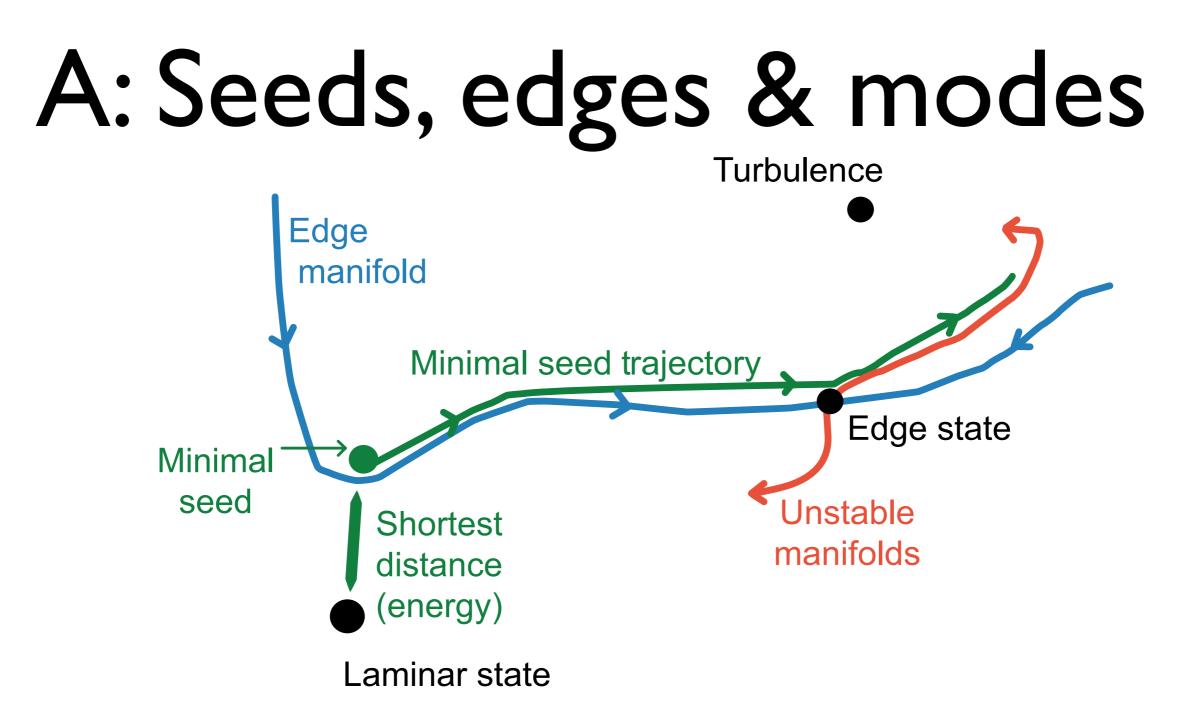
- Turbulence "obviously" important...
- "Structures not statistics" appealing viewpoint
- But can these structures actually give insight to real problems?

2.Approach: (HOW?) Two examples:

A.Can Koopman modes give insight into edge state/transition?

**B.**Can UPOs give insight into layering/mixing in stratified flows? (Recurrence, self-organization and the dynamics of turbulence)

**3**.Conclusions...time to be quantitative...

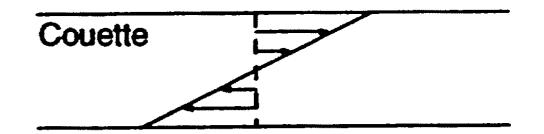


- Turbulent and laminar state separated by an "edge"
- Can we find the lowest energy state to get across the edge?
- Does "seed" follow special route to the "edge state" (I unstable)?

# A: Seeds, edges & modes

- Linear Optimal Perturbations are "easy": Linear Algebra available
- Variational formulation also possible (Schmid 07)
- Using "adjoint" operators (Hill 95, Corbett & Bottaro 00)
- Linear adjoints are "nice": completely decoupled...
- But nonlinear adjoints can be defined...and calculated
- Pringle & Kerswell 10, Cherubini et al 11, Monokrousos et al 11
- Rabin et al 2012 Duguet et al 2013 Kerswell et al 2014...
- Perturbation can feed back: what does the fluid want to do?

## Formulation



- Fix ideas: consider plane Couette flow:
- Butler & Farrell geometry:

$$L_x = 13.66, \ L_y = 2, \ L_z = 3.11, \ U = ye_x, \ Re = \frac{1}{\nu} = 1000$$

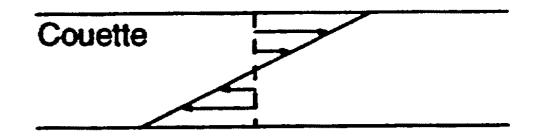
• Hypothesis: Maximize gain of (nonlinear) perturbation  $\hat{u} = u - U$ 

$$G(T) = \frac{E(T)}{E(0)}; E(T) = \frac{1}{2} \langle \hat{u}^2 + \hat{v}^2 + \hat{w}^2 \rangle$$

- Across all time horizons & amplitudes E(0) (note 1/2...!)
- Define spatial and temporal averages (dagger is c.c. transpose):

$$\langle \mathbf{b}, \mathbf{a} \rangle = \frac{\mathbf{I}}{\mathbf{V}} \int_{\mathcal{D}} \mathbf{b}^{\dagger} \mathbf{a} \mathrm{d} \mathbf{V}, \ [\mathbf{b}, \mathbf{a}] = \frac{\mathbf{I}}{\mathbf{VT}} \int_{\mathbf{0}}^{\mathbf{T}} \int_{\mathcal{D}} \mathbf{b}^{\dagger} \mathbf{a} \mathrm{d} \mathbf{V} \mathrm{d} \mathbf{t}$$

## Lagrangian

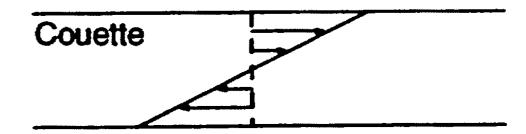


• Constrained variational problem:

$$\begin{split} \mathcal{L} &= \langle \mathbf{q}\left(\mathbf{T}\right), \mathbf{q}\left(\mathbf{T}\right) \rangle - \left[\partial_{t}\mathbf{q} + \mathbf{N}(\mathbf{q}) + \nabla \mathbf{p}, \mathbf{a}\right] + \left[\nabla . \mathbf{u}, b\right] \\ &- \left(\langle \mathbf{q}_{0}, \mathbf{q}_{0} \rangle - \mathbf{E}(\mathbf{0})\right)\mathbf{c} + \langle \mathbf{q}_{0} - \mathbf{q}(\mathbf{0}), \mathbf{a}_{0} \rangle \\ \mathbf{q} &= \hat{u}, \qquad \mathbf{N}(\mathbf{q}_{i}) = \mathbf{U}_{i}\partial_{j}\mathbf{q}_{i} + \mathbf{q}_{i}\partial_{i}\mathbf{U}_{j} + \mathbf{q}_{i}\partial_{j}\mathbf{q}_{i} - \nu\partial_{j}\partial_{j}\mathbf{q}_{i} \end{split}$$

- Final (perturbation) energy to be maximized
- Vector L-multiplier a imposes (nonlinear) Navier-Stokes for all t
- b imposes incompressibility (pressure?)
- c imposes initial energy perturbation
- a0 imposes initial perturbation (structure)  $\langle \mathbf{b}, \mathbf{a} \rangle = \frac{I}{V} \int_{\mathcal{D}} \mathbf{b}^{\dagger} \mathbf{a} \mathrm{d} V, \ [\mathbf{b}, \mathbf{a}] = \frac{I}{VT} \int_{0}^{T} \int_{\mathcal{D}} \mathbf{b}^{\dagger} \mathbf{a} \mathrm{d} V \mathrm{d} t$

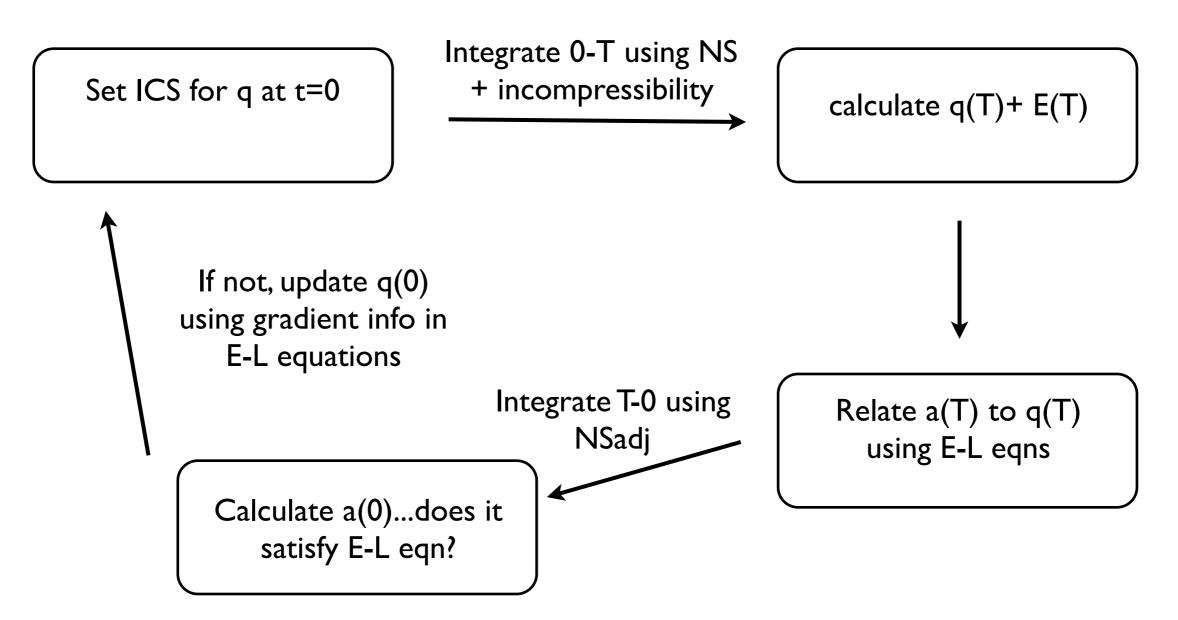
## Adjoint



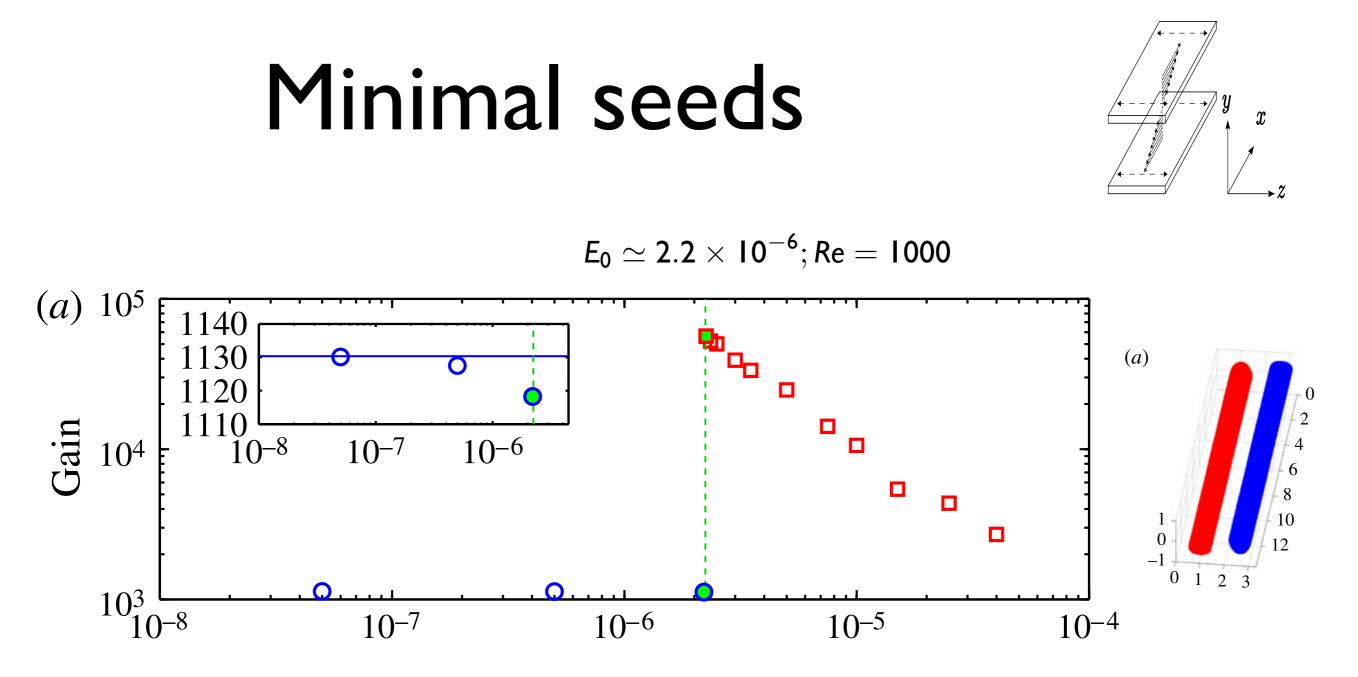
- "Obviously" variations with respect to a, b,c, a0
  - $\mathcal{L} = \langle \mathbf{q}(\mathbf{T}), \mathbf{q}(\mathbf{T}) \rangle [\partial_t \mathbf{q} + \mathbf{N}(\mathbf{q}) + \nabla \mathbf{p}, \mathbf{a}] + [\nabla . \mathbf{u}, \mathbf{b}] \\ (\langle \mathbf{q}_0, \mathbf{q}_0 \rangle \mathbf{E}(\mathbf{0}))\mathbf{c} + \langle \mathbf{q}_0 \mathbf{q}(\mathbf{0}), \mathbf{a}_0 \rangle \\ \mathbf{q} = \hat{\mathbf{u}} \qquad \mathbf{N}(\mathbf{q}) \mathbf{U}\partial_t \mathbf{q} + \mathbf{q}\partial_t \mathbf{q} + \mathbf{q}\partial_t \mathbf{q} \mathbf{u}\partial_t \partial_t \mathbf{q}$
  - $\mathbf{q} = \hat{\mathbf{u}}, \qquad \mathbf{N}(\mathbf{q}_i) = \mathbf{U}_j \partial_j \mathbf{q}_i + \mathbf{q}_i \partial_i \mathbf{U}_j + \mathbf{q}_j \partial_j \mathbf{q}_i \nu \partial_j \partial_j \mathbf{q}_i$
- Recover NS/incompressibility/ICs...but what about:  $\frac{\delta \mathcal{L}}{\delta \mathbf{q}} = \partial_t \mathbf{a} + N^{\dagger}(\mathbf{a}, \mathbf{q}) + \nabla b - (\mathbf{a} + 2\mathbf{q})|_{t=T} + (\mathbf{a} - \mathbf{a}_0)|_{t=0} = \mathbf{0}$   $\frac{\delta \mathcal{L}}{\delta \mathbf{p}} = \nabla \cdot \mathbf{a} = \mathbf{0}$   $\frac{\delta \mathcal{L}}{\delta \mathbf{q}_0} = \mathbf{a}_0 - 2\mathbf{q}_0\mathbf{c} = \mathbf{0} \quad \frac{\delta \mathcal{L}}{\delta T} = \frac{\delta \mathbf{E}}{\delta T} = \mathbf{0}$   $N^{\dagger}(\mathbf{a}_i, \mathbf{q}) = \partial_j \left(\mathbf{q}_j \mathbf{a}_i\right) - \mathbf{a}_j \partial_i \mathbf{q}_j + \partial_j \left(\mathbf{U}_j \mathbf{a}_i\right) - \mathbf{a}_j \partial_i \mathbf{U}_j + \nu \partial_j \partial_j \mathbf{a}_i$ 
  - Adjoint equation inevitably coupled to q as well as a...

# Algorithm

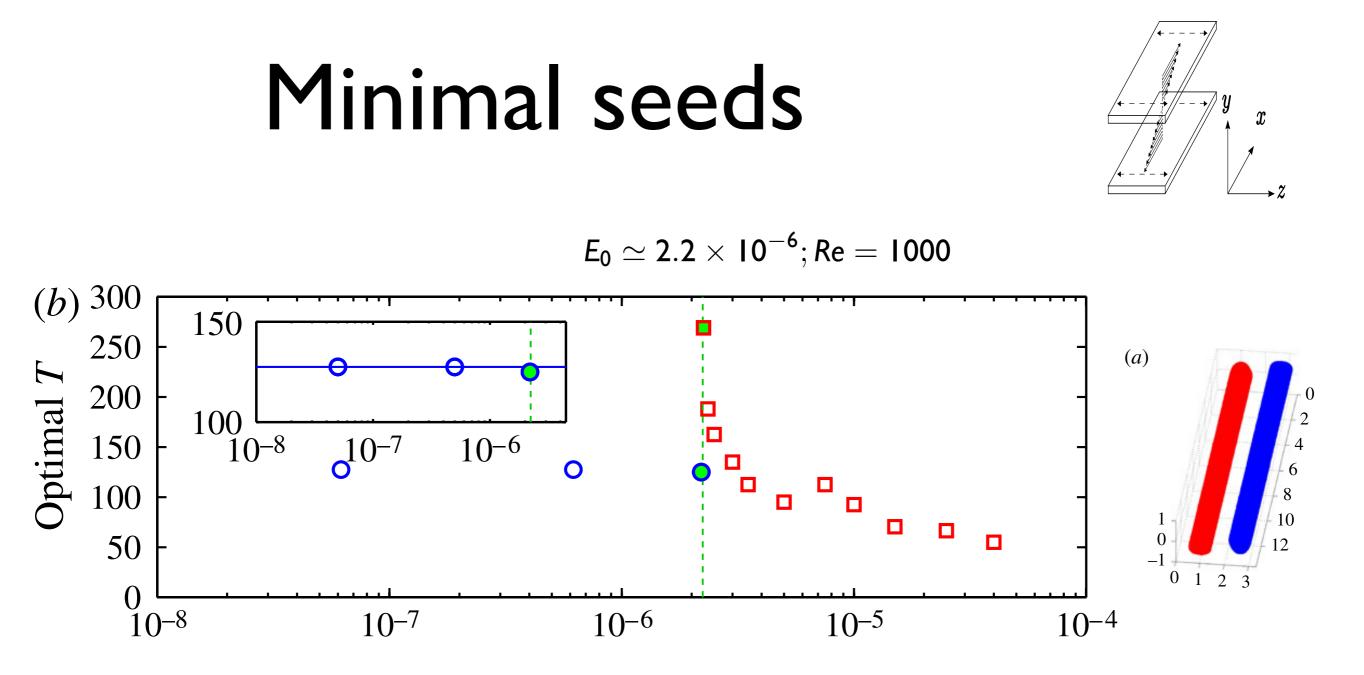
• And now the algorithm is obvious (diffusion has opposite sign!)



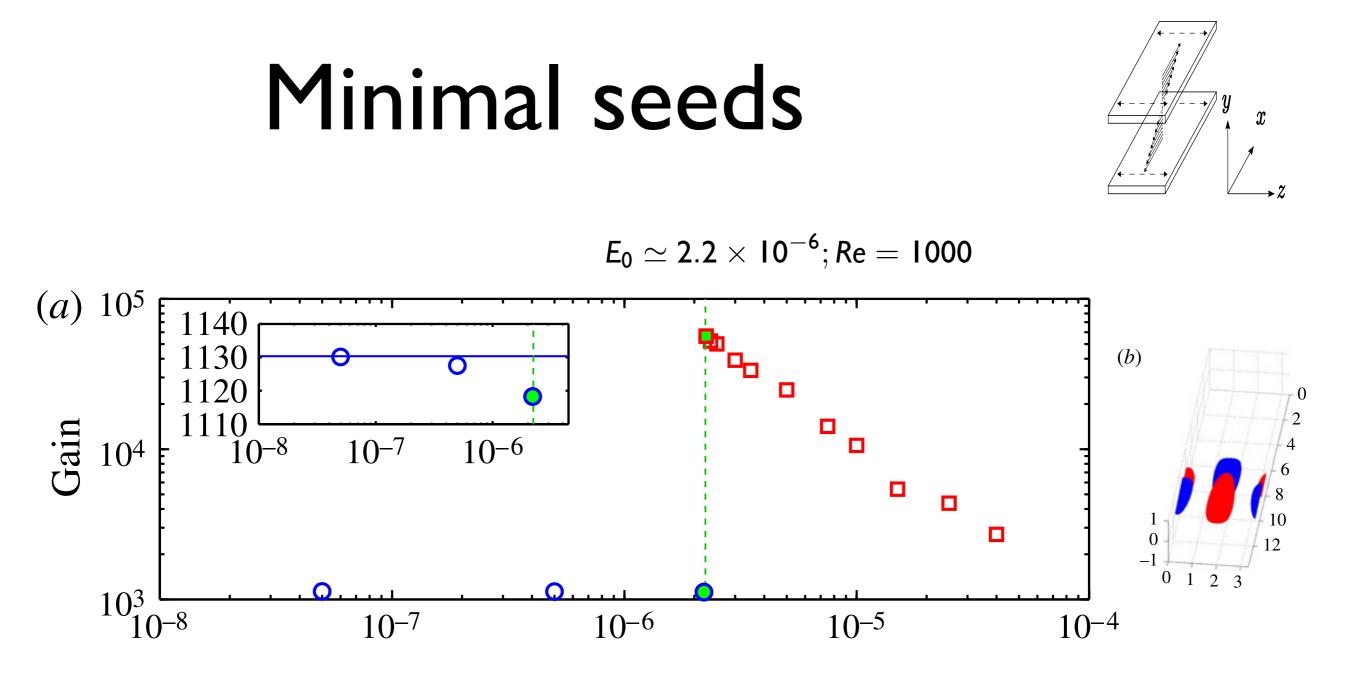
- For small amplitudes: recover Quasi-Linear-Optimal Perturbations
- What happens as the amplitude increases...Re is high enough?



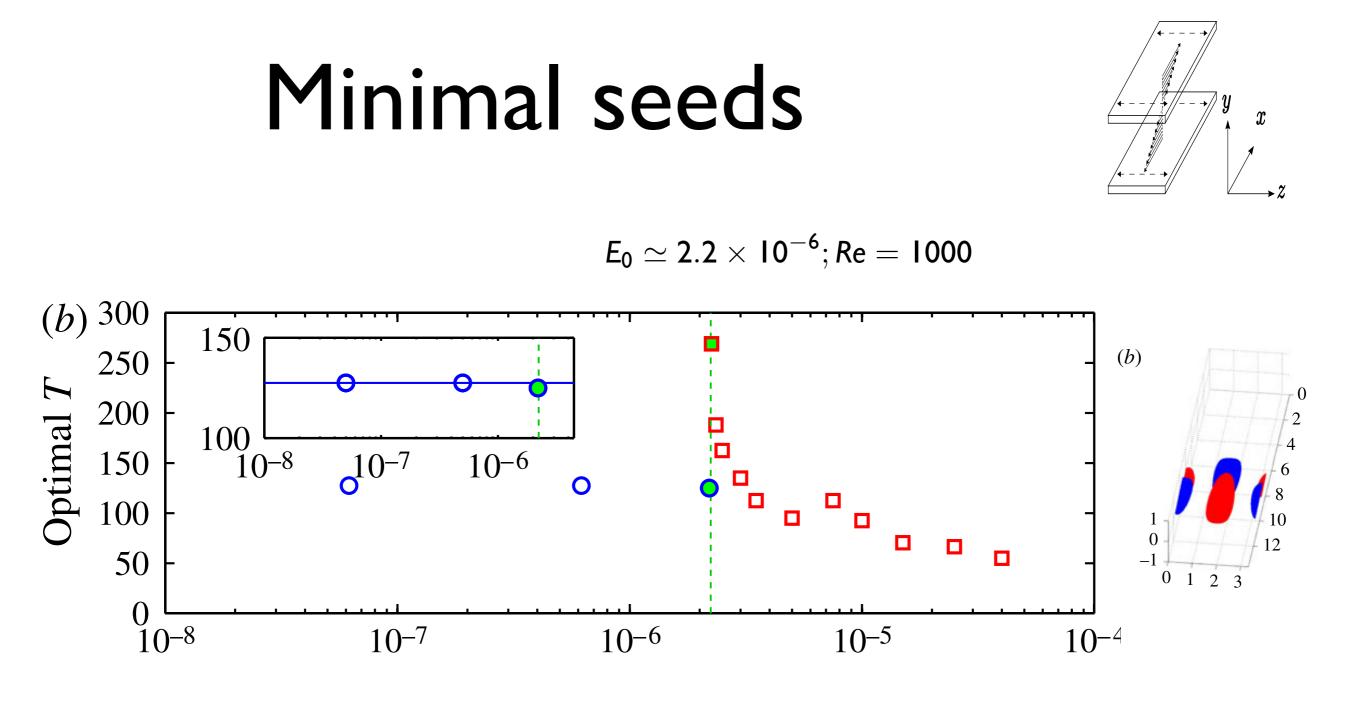
For small initial energy...agrees very closely with linear optimals in gain



 For small initial energy...agrees very closely with linear optimals and T



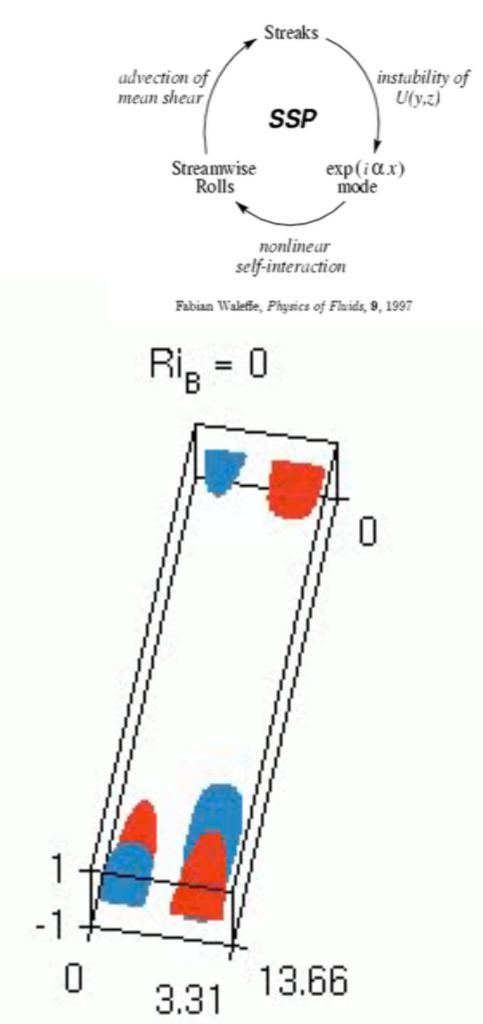
- For small initial energy...agrees very closely with linear optimals
- But sudden change at high enough energy....
- Inherently localised structure



- For small initial energy...agrees very closely with linear optimals
- But sudden change at high enough energy....
- Inherently localised structure...with turbulence at late time

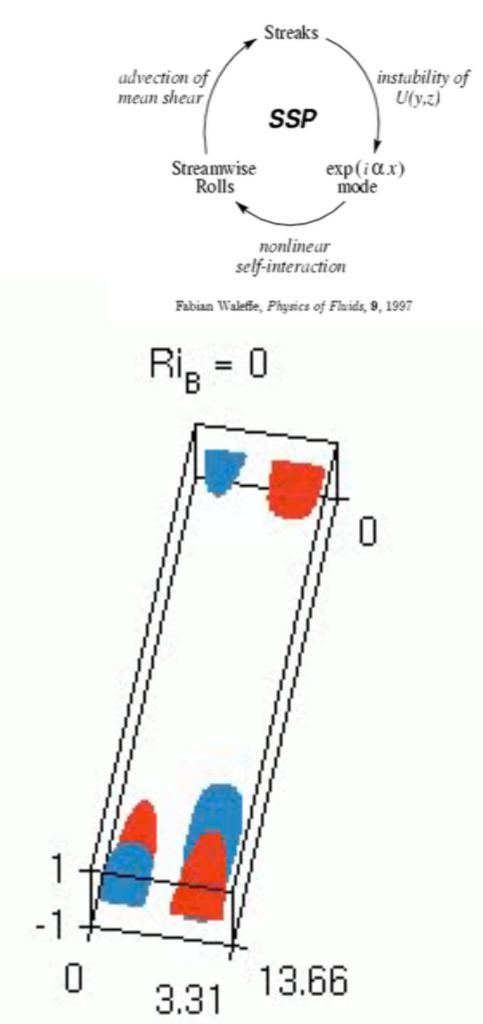
### Minimal seeds

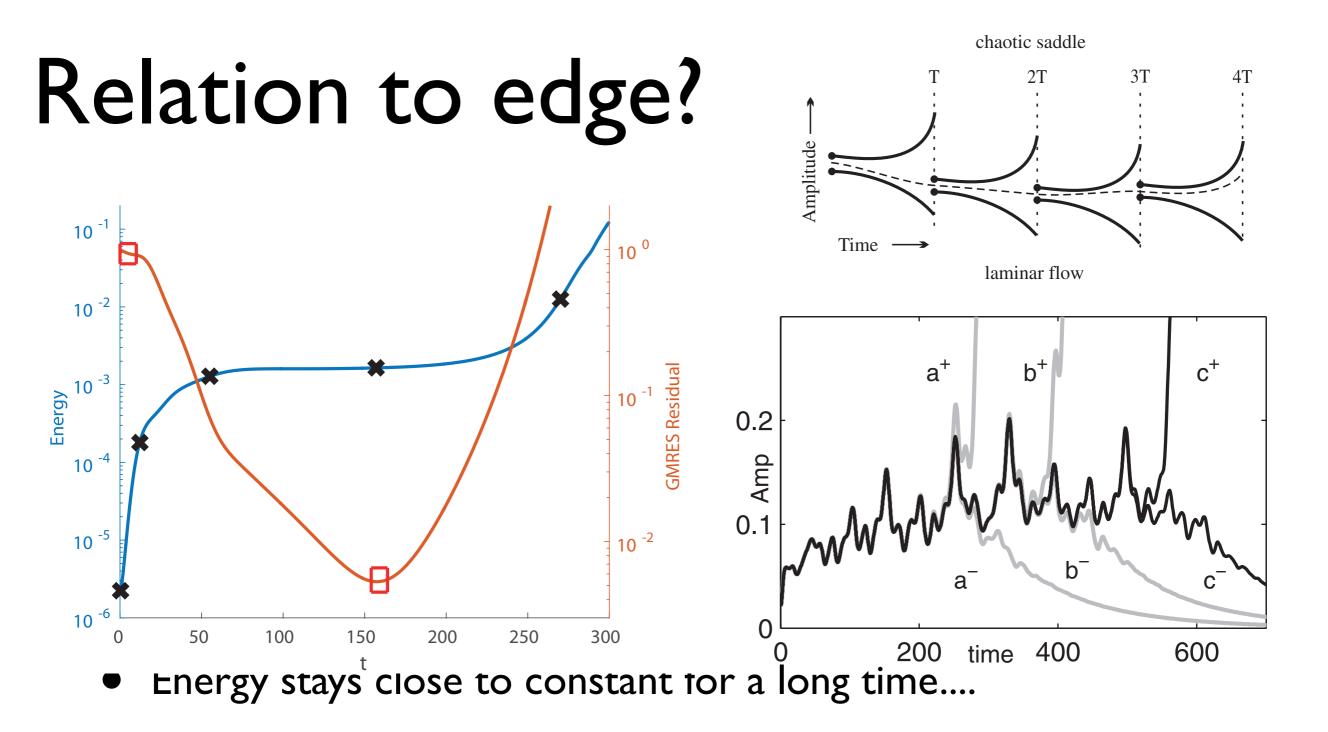
- SSP: Waleffe/VWI: Hall
- Hierarchy of growth mechanisms
- Bootstrapping inherently nonlinearly
- Duguet et al 2013
- Kerswell et al 2014
- Unpack to streaks to bending waves
- to breakdown...



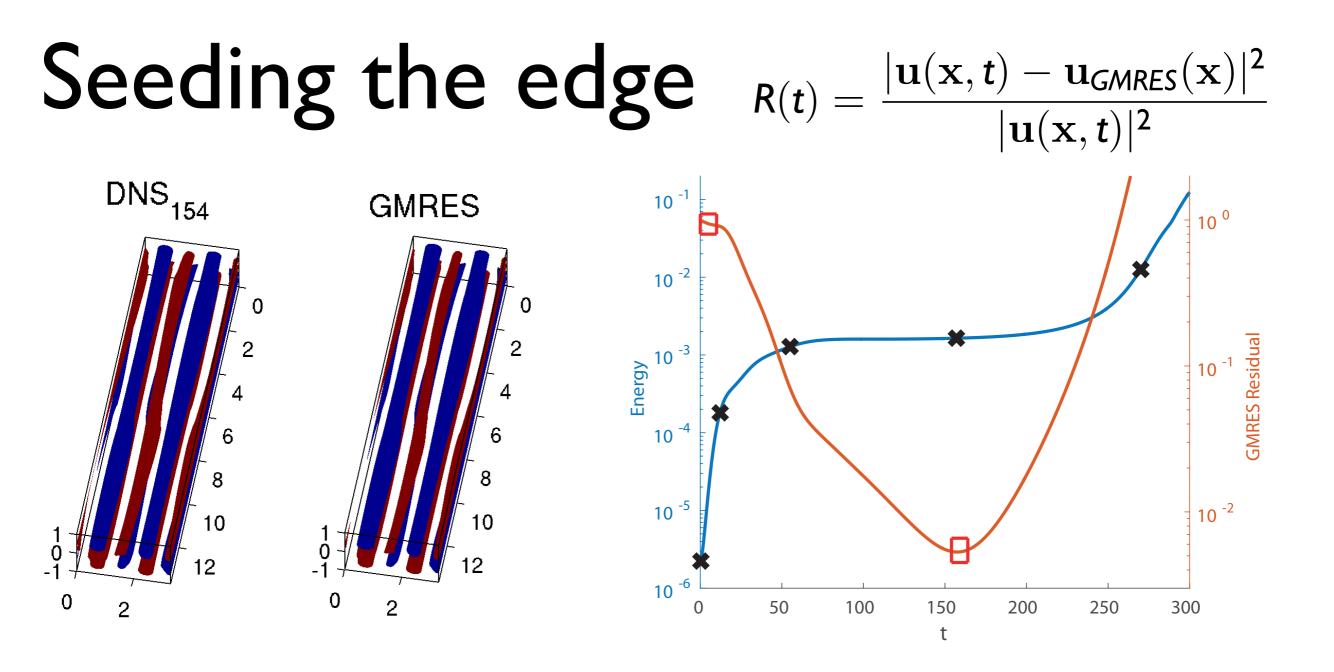
### Minimal seeds

- SSP: Waleffe/VWI: Hall
- Hierarchy of growth mechanisms
- Bootstrapping inherently nonlinearly
- Duguet et al 2013
- Kerswell et al 2014
- Unpack to streaks to bending waves
- to breakdown...





- Suggestive of being a specific way to tiptoe along the edge
- Suggesting approach on stable manifold of edge state Skufca et al 06 etc



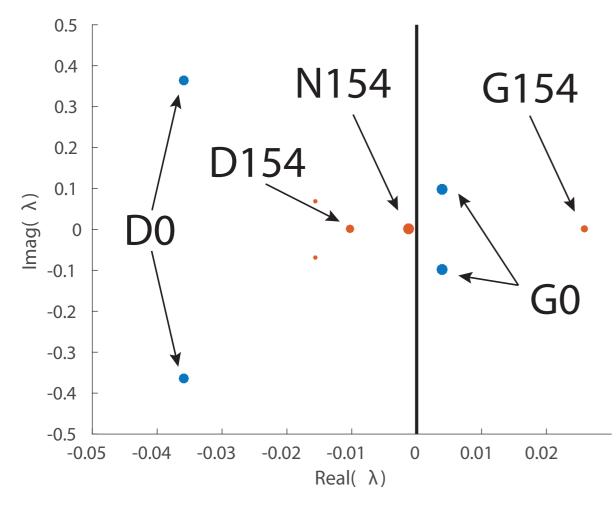
- Use DNS results as guess for Generalized Minimal Residual method
- Finds a steady edge state with very clear structure
- Is it possible to identify reduced description of approach/departure?

### Koopman/DMD Modes

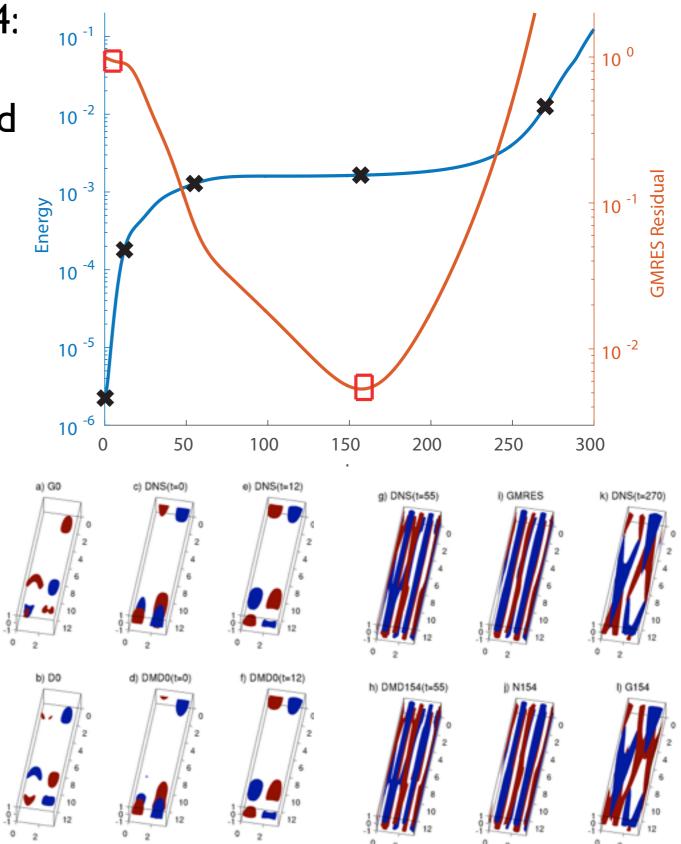
- Dynamical information about flow: take snapshous  $= \mathbf{q}(\mathbf{x}, t_n)$
- Assume (for every n...system not changing much... $\mathbf{\hat{q}}_{n+1} = \mathbf{N}_0 \mathbf{q}$
- Use DMD (Schmid 2010) to find spectrum  $\Delta I_0$
- Decompose (amplitudes "fit" snapshotq)( $\mathbf{x}, \mathbf{t}$ ) =  $\sum a_n \mathbf{m}_n(\mathbf{x}) \exp(\lambda_n \mathbf{t})$
- Finite-dimensional/fixed in time approximation of Koopman operator
- Mezic (2005) linear in the observable  $\mathbf{g}_t \mathbf{q}(\mathbf{x}, \mathbf{0}) = \mathbf{q}(\mathbf{x}, t)$  : effinished equations of the experimental equation equation of the experimental equation eq
- On attractq(x, t)  $\overline{\mathbf{q}}(\mathbf{x}) = \sum a_n \mathbf{m}_n(\mathbf{x}) \exp(\lambda_n t) + \int e^{2\pi i \alpha t} [dE(\alpha)\mathbf{q}]$
- Does this yield a low-dimensional description of the flow?

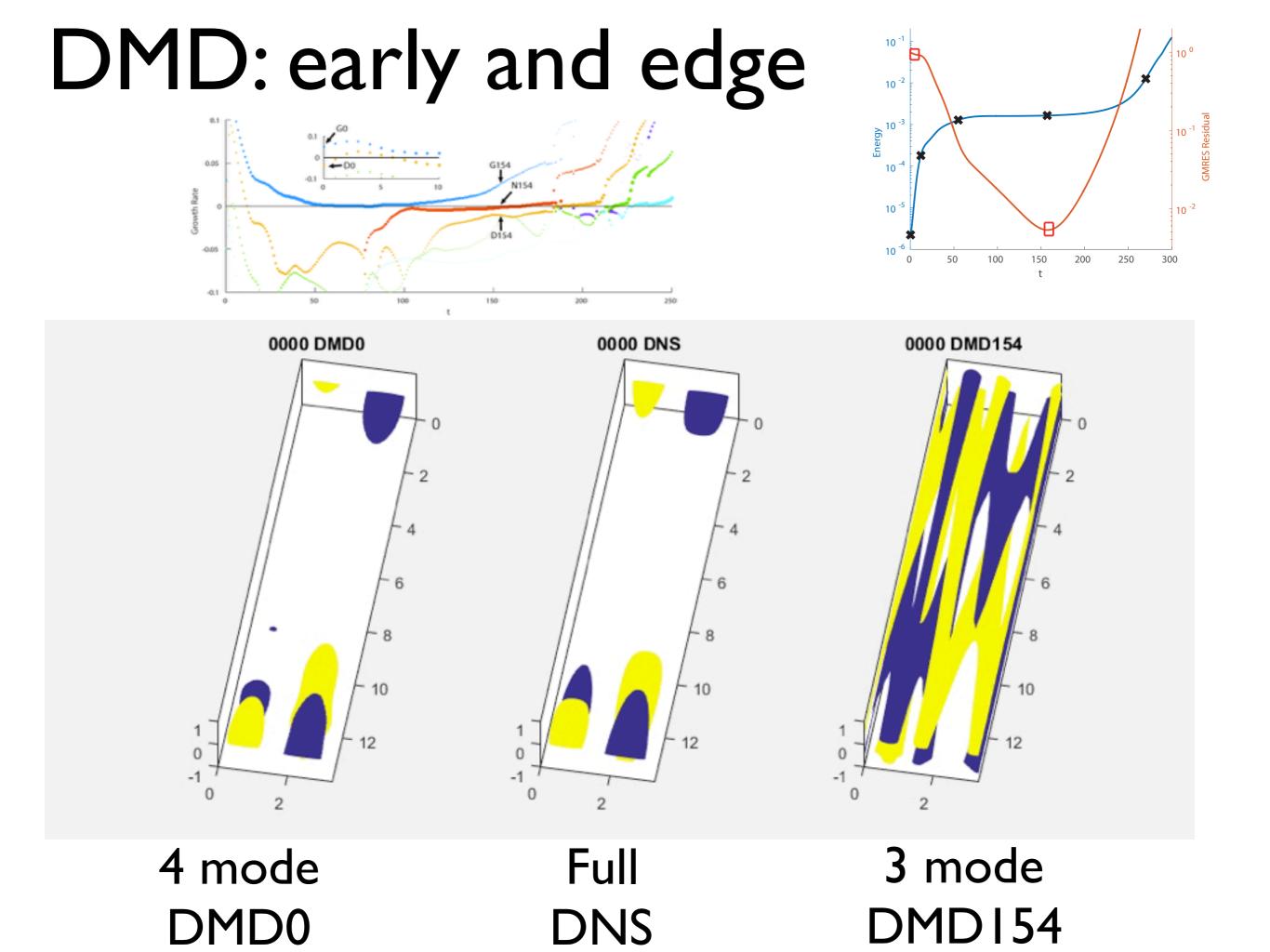
### DMD: early and edge

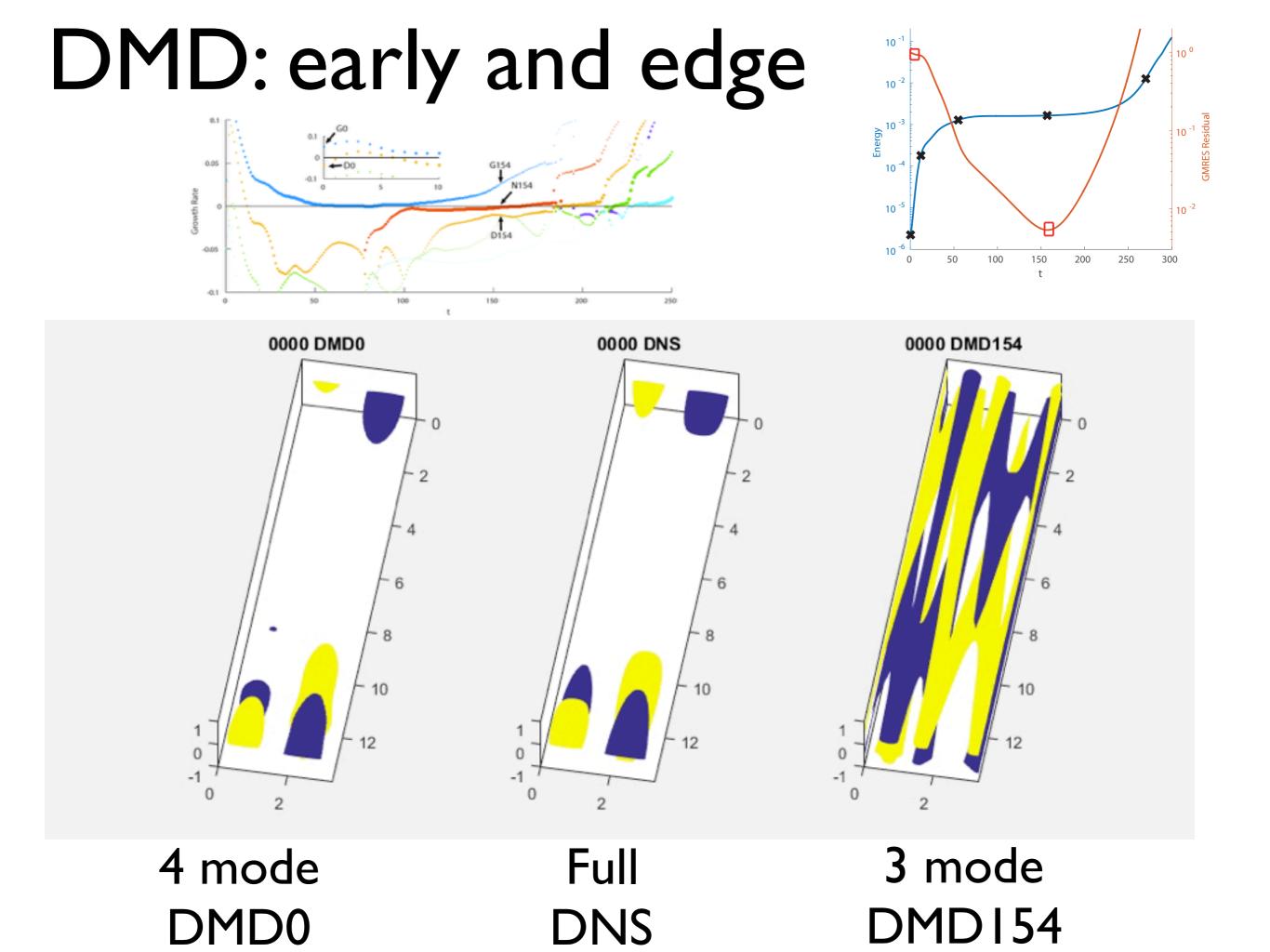
- II snapshots: 0-10 and 154-164:
- Spectrum: 4 (early) 3 (late) mod



- Captures early & late evoluti
- Using constant eigenvalues

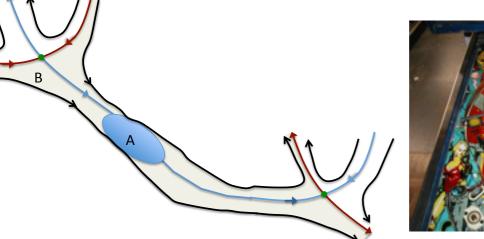




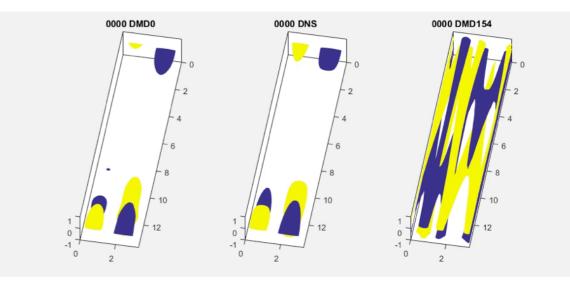


## A: Conclusions

- Direct-Adjoint Looping method rides along edge to edge state
- DAL method gives specific route on stable manifold
- DAL method leaves edge state along unstable manifold
- Koopman modes appear to be the natural description
- Can we use this process to play pinball turbulence?



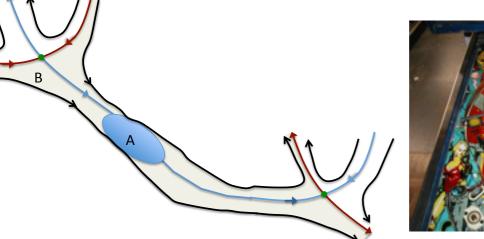




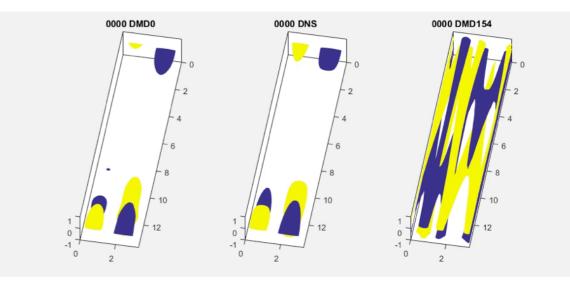
arcadegamessuperstore.com

## A: Conclusions

- Direct-Adjoint Looping method rides along edge to edge state
- DAL method gives specific route on stable manifold
- DAL method leaves edge state along unstable manifold
- Koopman modes appear to be the natural description
- Can we use this process to play pinball turbulence?



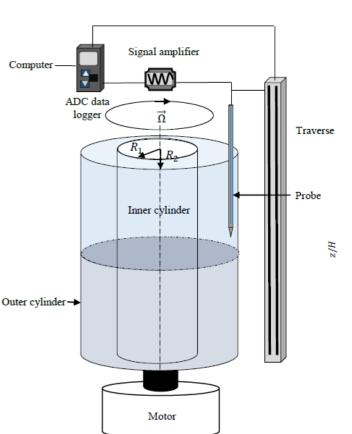


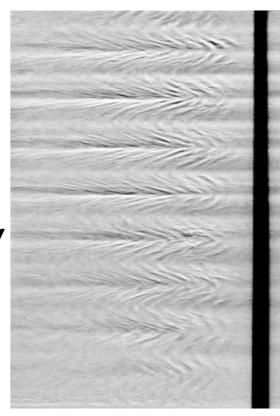


arcadegamessuperstore.com

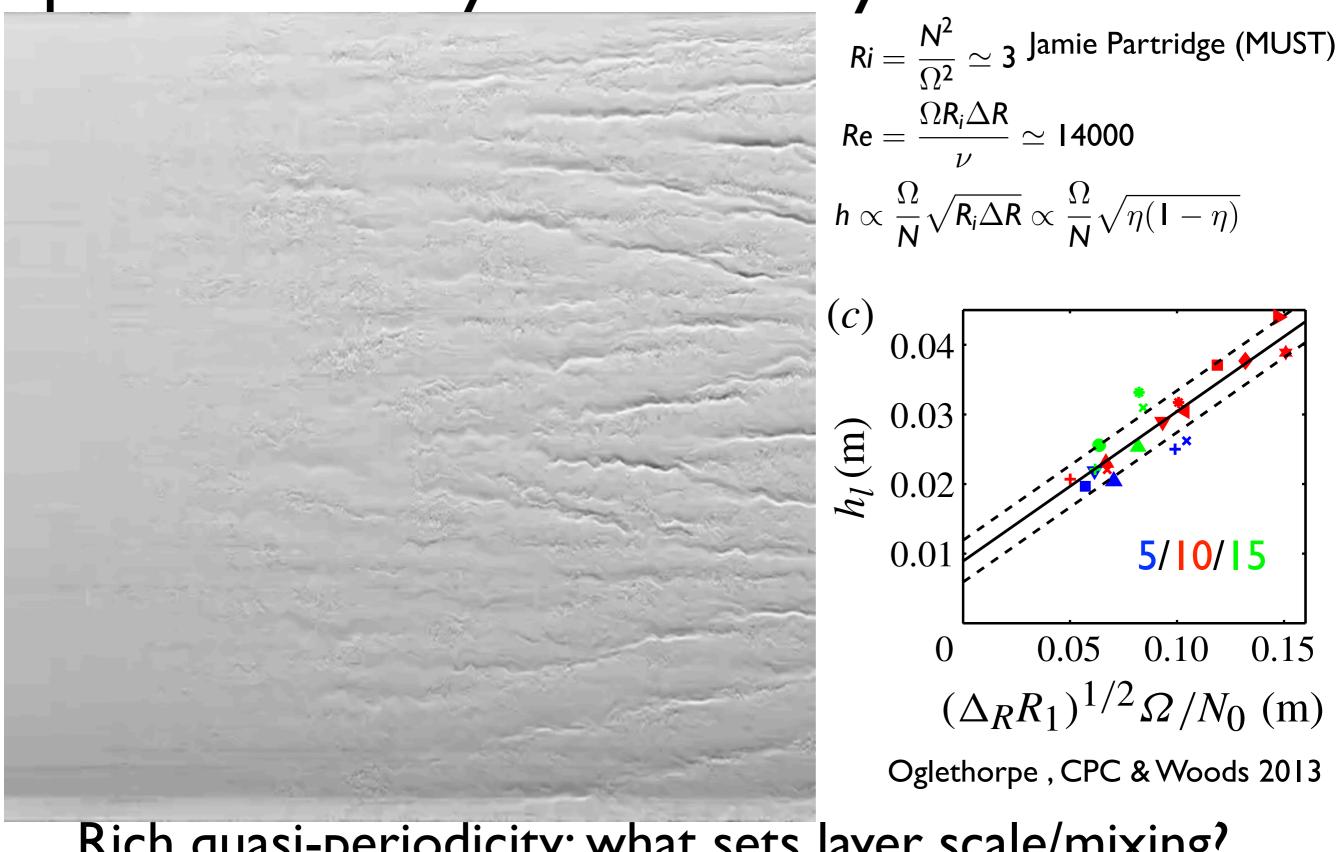
# B: Layering & UPOs

- Initially linearly stratified flows layer spontaneously
- "anti-diffusion" Phillips(1972) Park et al (1994)
- Holford & Linden (1998) vertical rod spontaneously layers:
- Horizontal shear can also be important: stratified T-C flow?
- Outer cylinder 24.7 cm: inner 5/10/15 cm
- Radius ratio:  $\eta = \frac{R_i}{R_0} = 0.208, 0.417, 0.625$
- Re > 10000, initially linearly stratified
- Lots of horizontal shear: visualize with shadow



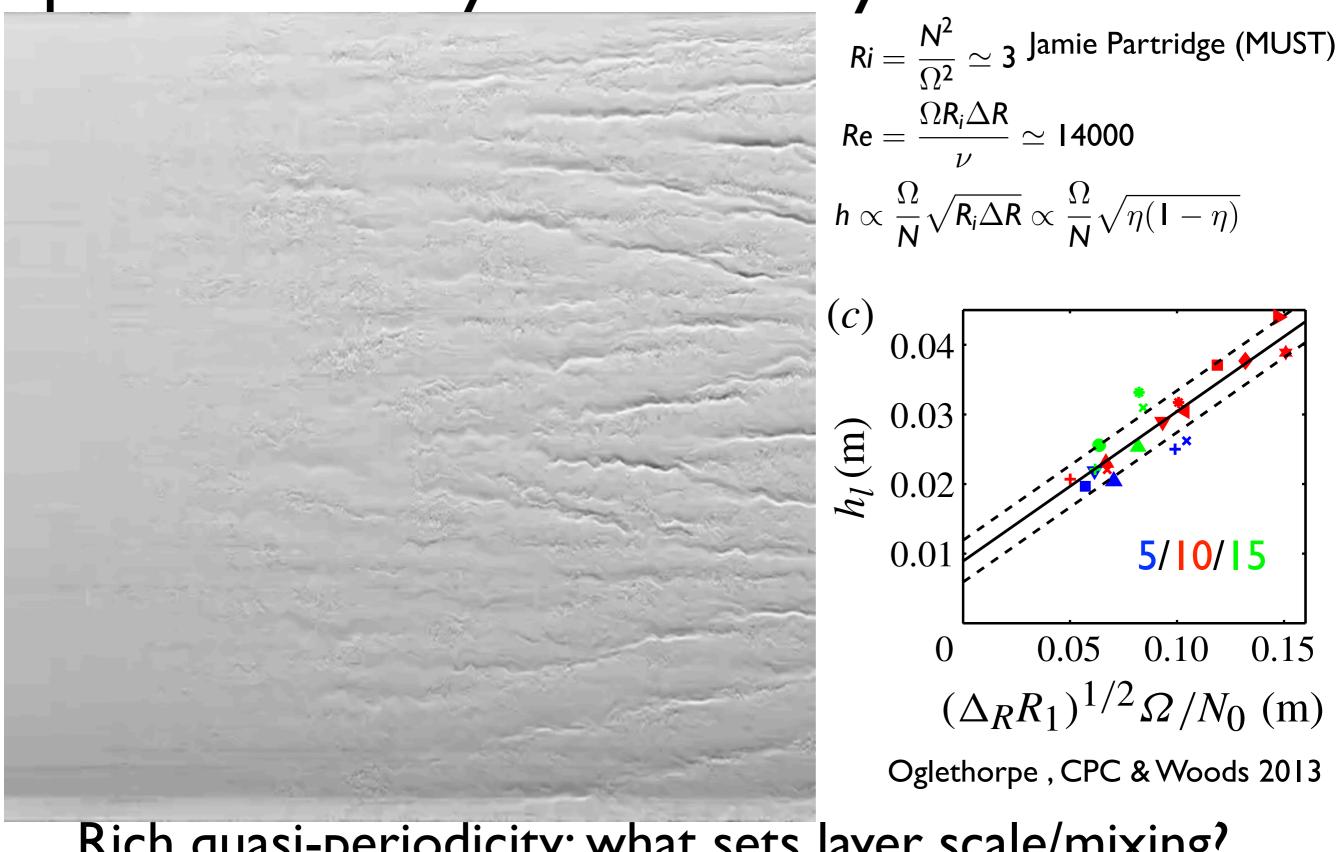


#### Spontaneous layers: NOT Taylor vortices



Rich quasi-periodicity: what sets layer scale/mixing?

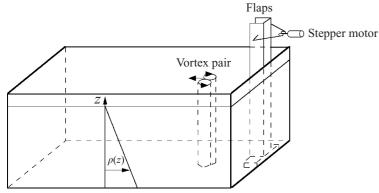
#### Spontaneous layers: NOT Taylor vortices



Rich quasi-periodicity: what sets layer scale/mixing?

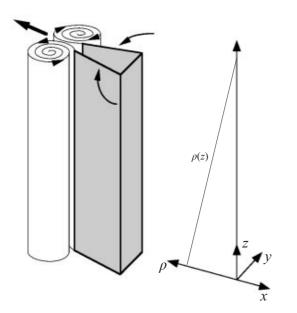
# Layers, zig-zags & connections

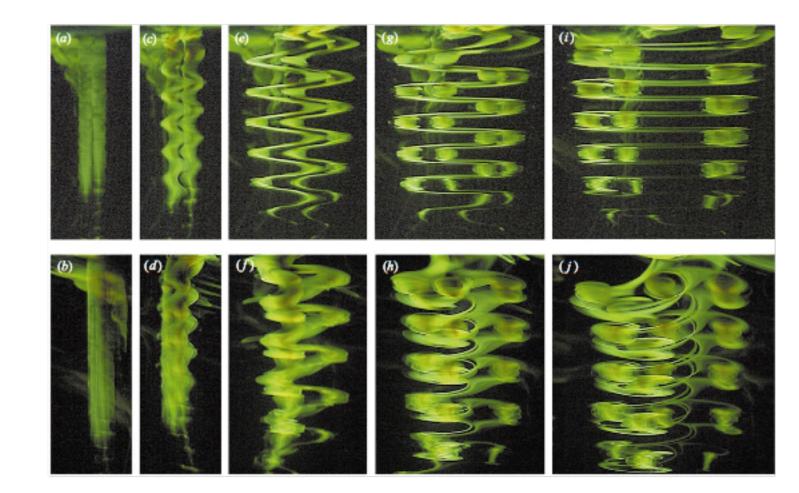
- Layer depth scales like U/N for characteristic velocity U...
- Reminiscent of zig-zag instability of Billant & Chomaz (2000a,b)



Linear stratification

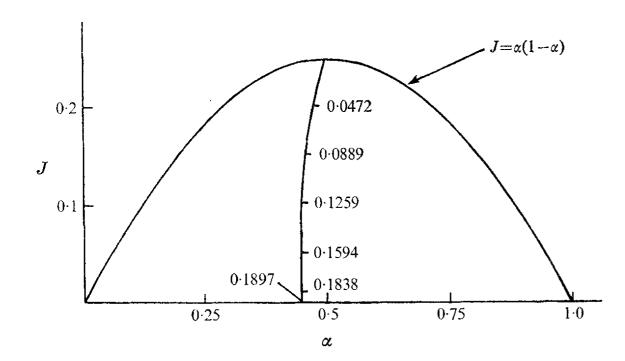
FIGURE 1. Sketch of experimental set-up.

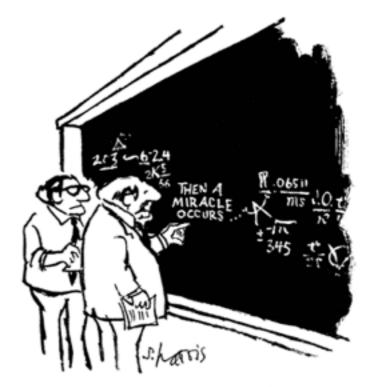




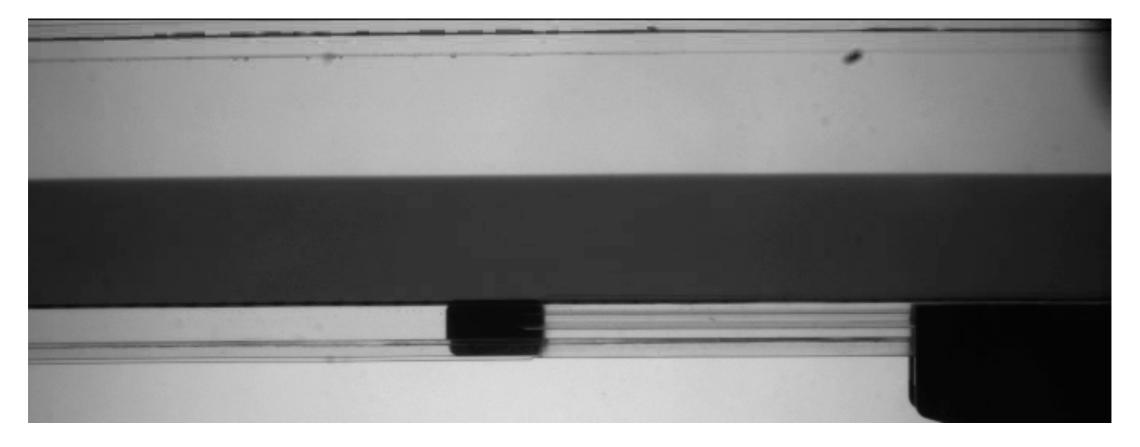
- Experimental miracle: linear mode to structure?

#### Connecting linear modes to self-organization?

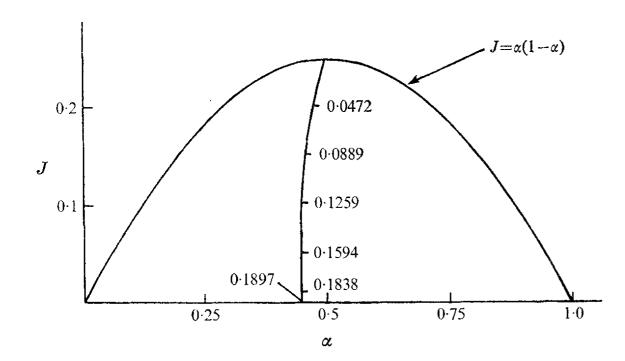


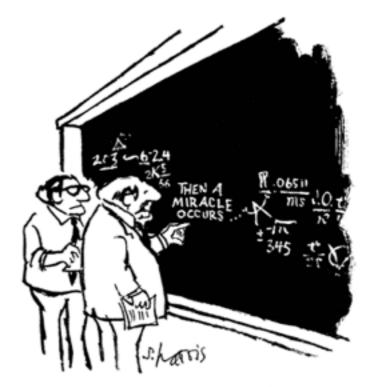


"I think you should be more explicit here in step two."

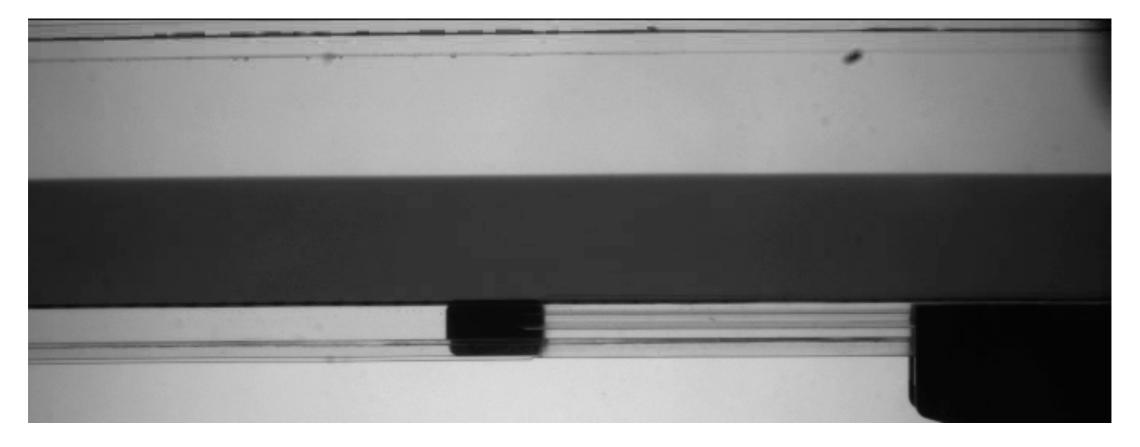


#### Connecting linear modes to self-organization?





"I think you should be more explicit here in step two."



# HS Kolmogorov Flow

- Horizontal Kolmogorov flow  $U \equiv u_{lam} = Re \sin(y) \hat{\mathbf{x}}, \ \rho_B = -z$
- Forced back to laminar flow with three control parameters

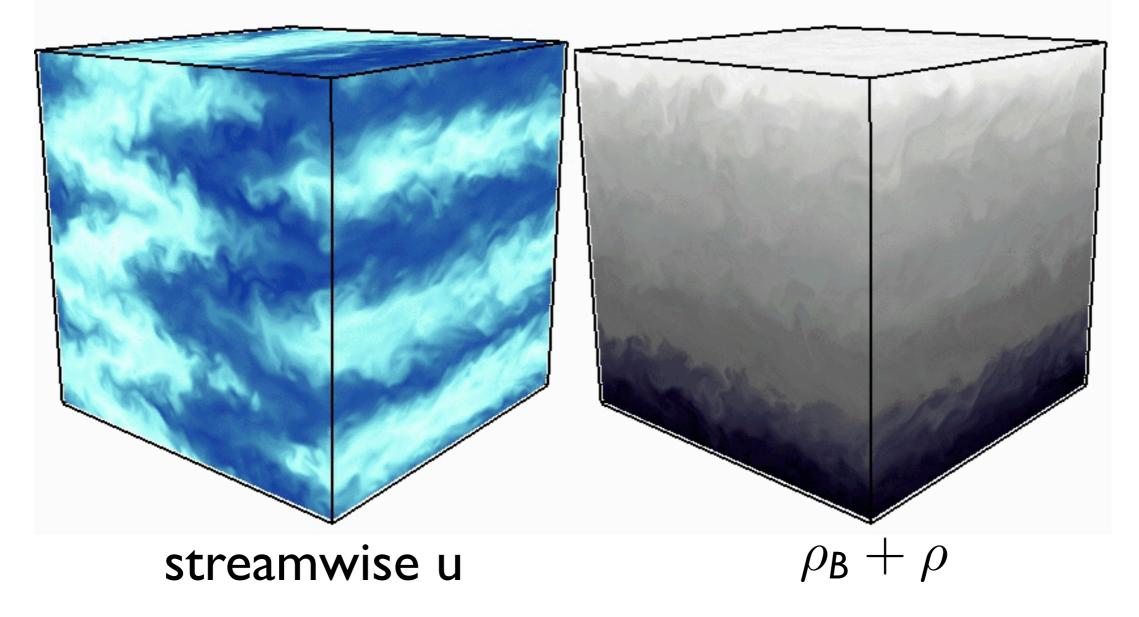
  <sup>∂</sup>u / ∂t + u · ∇u + ∇p = <sup>1</sup>/<sub>Re</sub> Δu + sin(y) x̂ Bρ ẑ
  <sup>2</sup>/<sub>x</sub>
  <sup>j</sup>/<sub>x</sub>
  <sup>j</sup>/<sub>y</sub>
  <sup>j</sup>/<sub>x</sub>
  <sup>j</sup>/<sub>x</sub>
  <sup>j</sup>/<sub>y</sub>
  <sup>j</sup>/<sub>x</sub>
  <sup>j</sup>/<sub>y</sub>
  <sup>j</sup>/<sub>x</sub>
  <sup>j</sup>/<sub>y</sub>
  <sup>j</sup>/<sub></sub>

U(y)

- Buoyancy parameter B
- Aspect ratio  $\alpha = L_y/L_x = L_z/L_x$
- Horizontal shear so vertical vorticity...

# HS Kolmogorov Flow

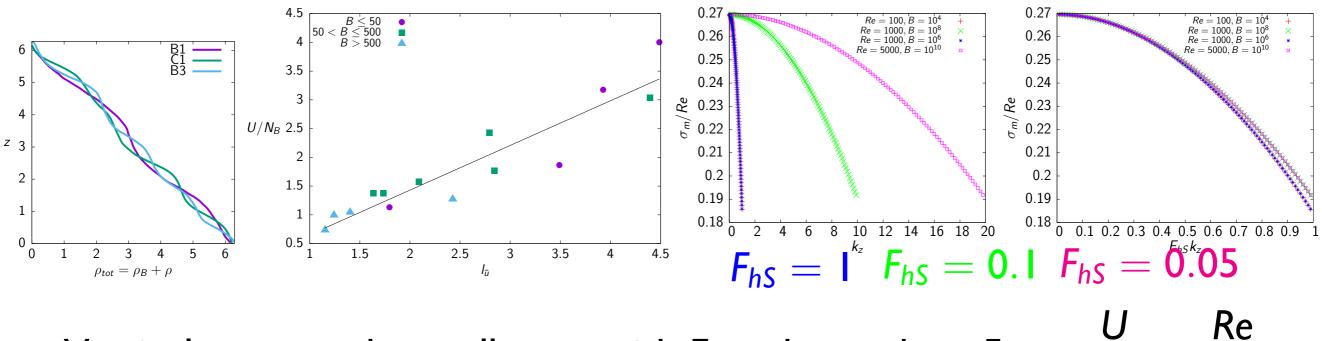
• Horizontal stratified Kolmogorov flow self-organizes into layers!



• Any connection between linear stability and nonlinear dynamics?

## Layer scaling

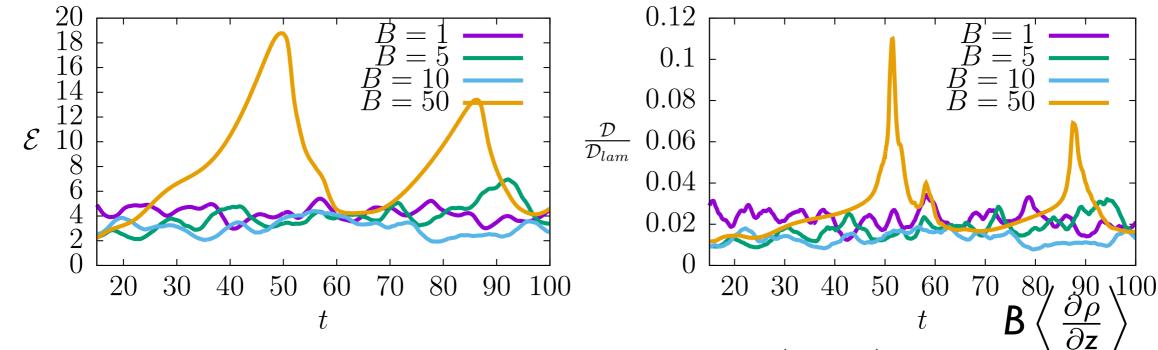
• Layer scaling consistent with U/N...and zig-zag like linear instability



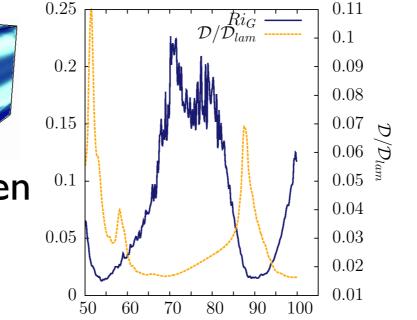
- Vertical wavenumber collapses with Froude number:  $F_{hS} = \frac{U}{LN} = \frac{Re}{\sqrt{R}}$
- Scaling for  $F_{hS} \ll I$ ;  $Re \gg I$
- Stratified turbulence scaling of Billant & Chomaz/Lindborg etc
- Stability properties just like zigzag/Deloncle et al 2007
- Nonlinear properties, particularly mixing?

# Mixing: UPOs?

Particularly in strongly stratified flow: mixing very intermittent



- Complex spatio-temporal structure of  $Ri_G(y, z, t) =$
- t=70,88,90,100
- <**Ri**> drops..
- Then turbulence...then
- Ri up again: UPO??

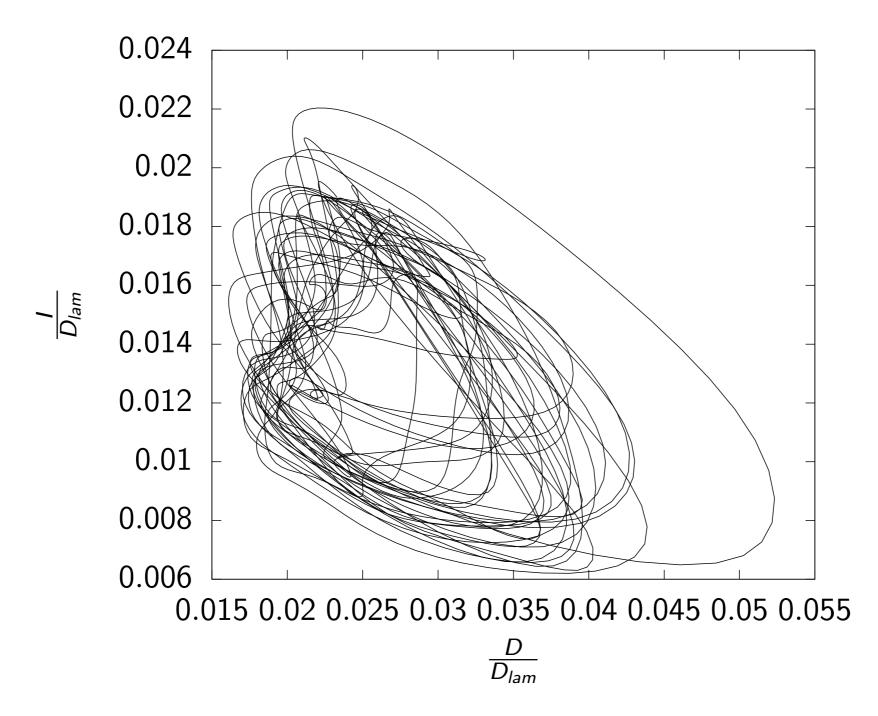


t

1e + 0254 N 3 × 3 1e + 012 1e + 001e-01 54 1e-02 × 3 × 3 1e-03  $2 \ 3$ 4 5 6 3 524 6 0 1 yy

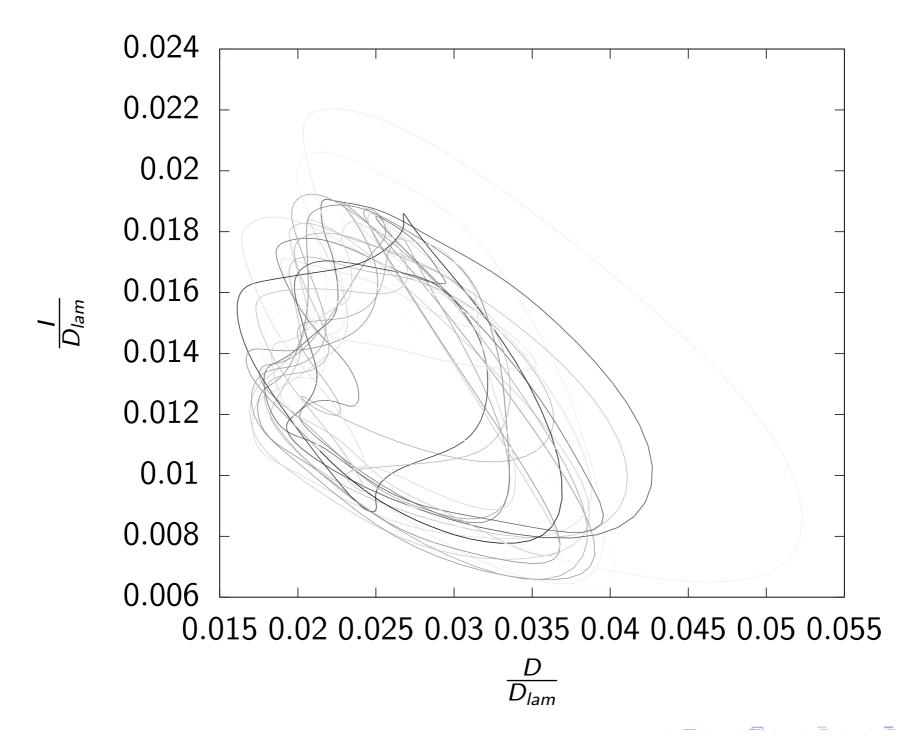
<u>∂u<sub>h</sub></u>

• Project trajectory onto plane:



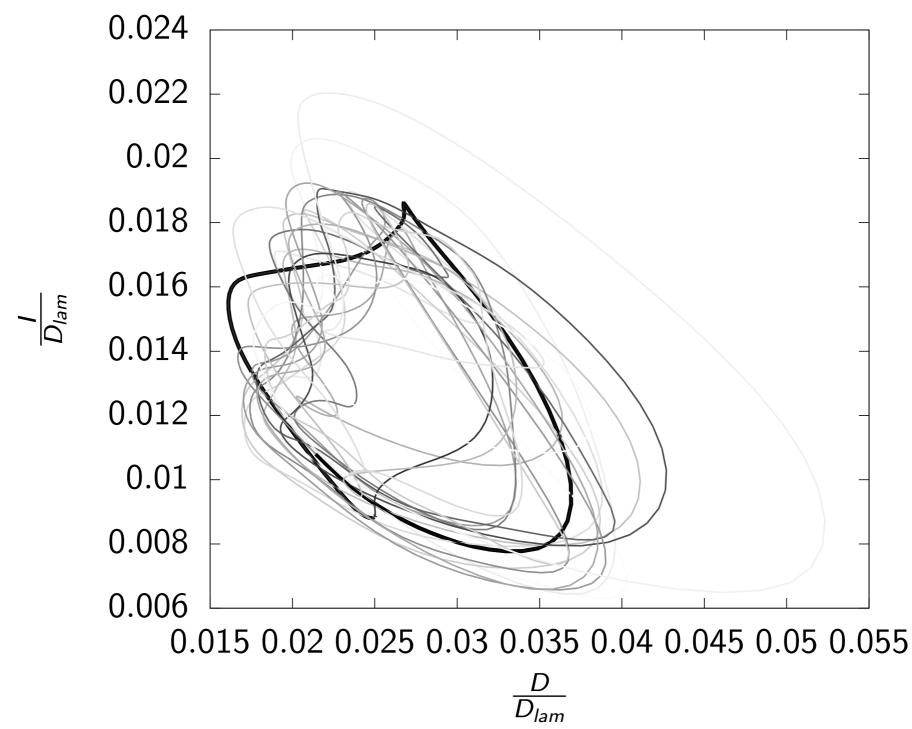
• Looks an unholy mess....but look closer...

• Closer inspection reveals nearly recurrent episodes

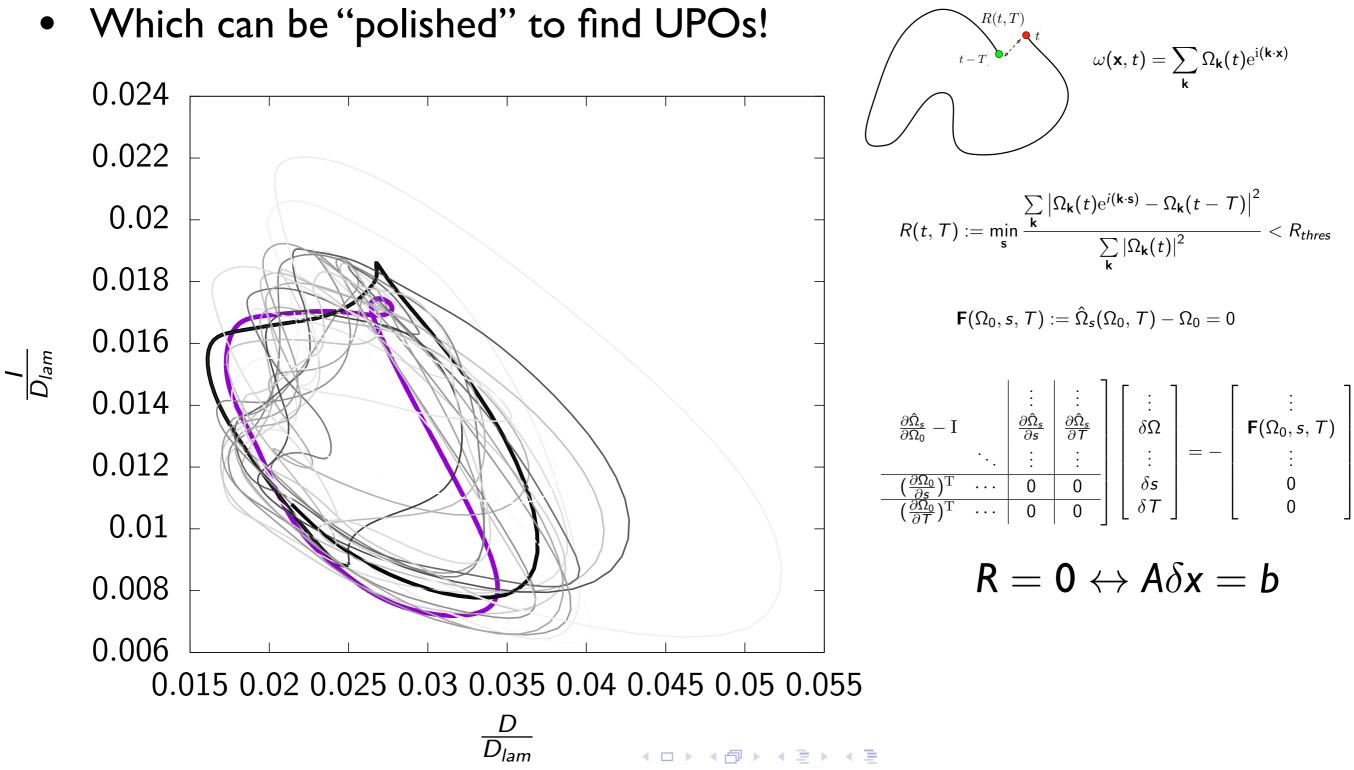


• If you look really closely...

• Closer inspection reveals nearly recurrent episodes

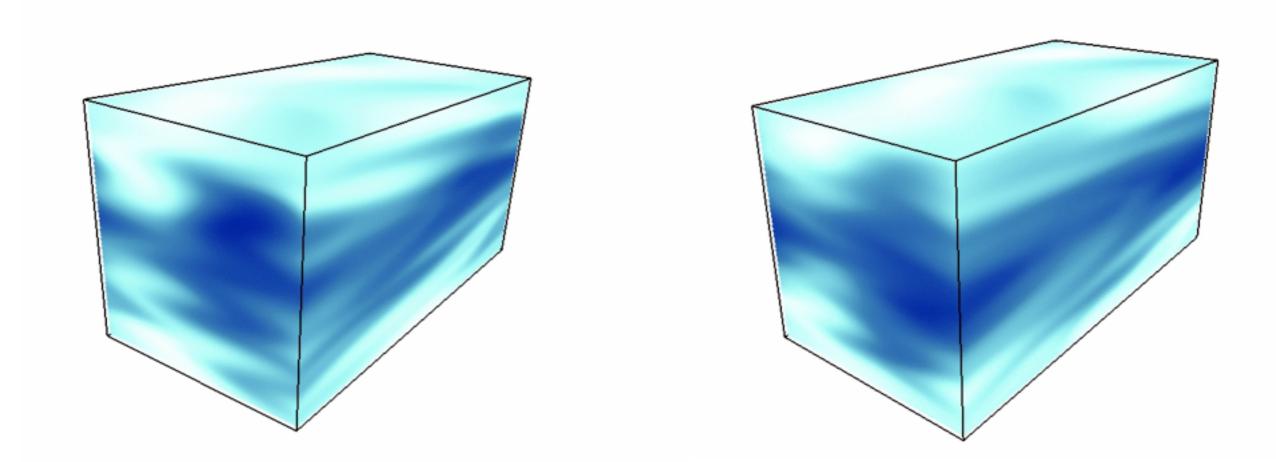


• Which can be used as original guesses for recurrent flow analysis...



See Lucas & Kerswell 2015/16 for Newton Solve/Hookstep/GMRES etc

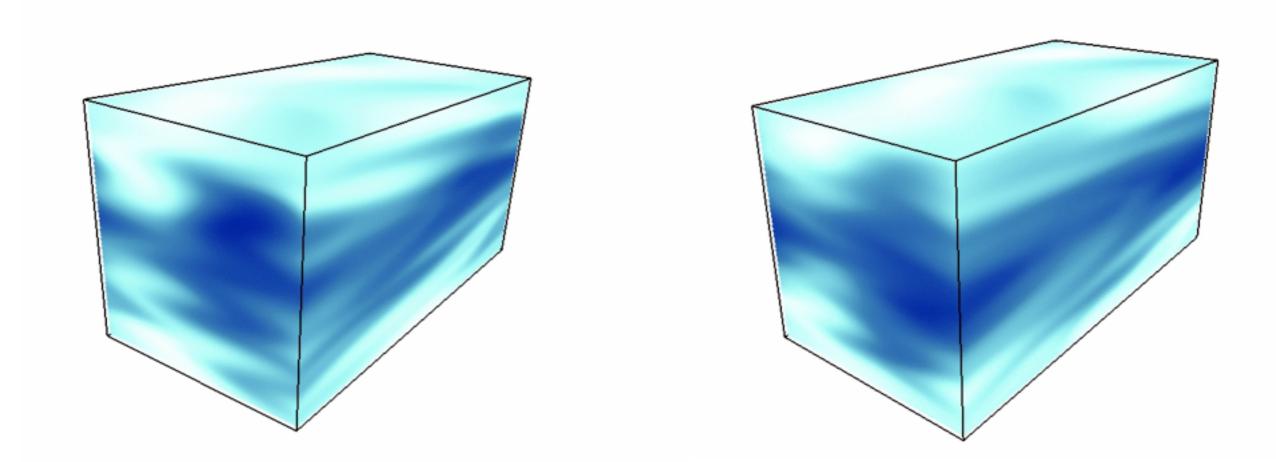
• Can find stratified UPOs, and they really "look" like the flow:



$$Re = 50, B = 50, \alpha = 0.5$$

• But can they tell us something quantitative about the flow behaviour?

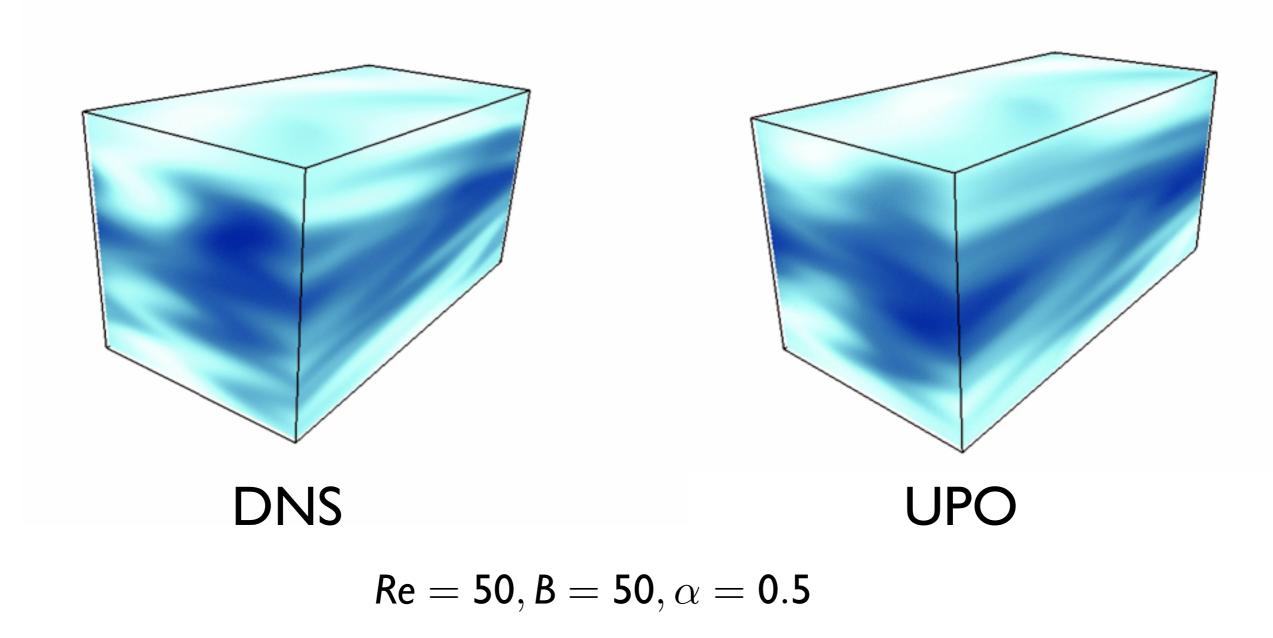
• Can find stratified UPOs, and they really "look" like the flow:



$$Re = 50, B = 50, \alpha = 0.5$$

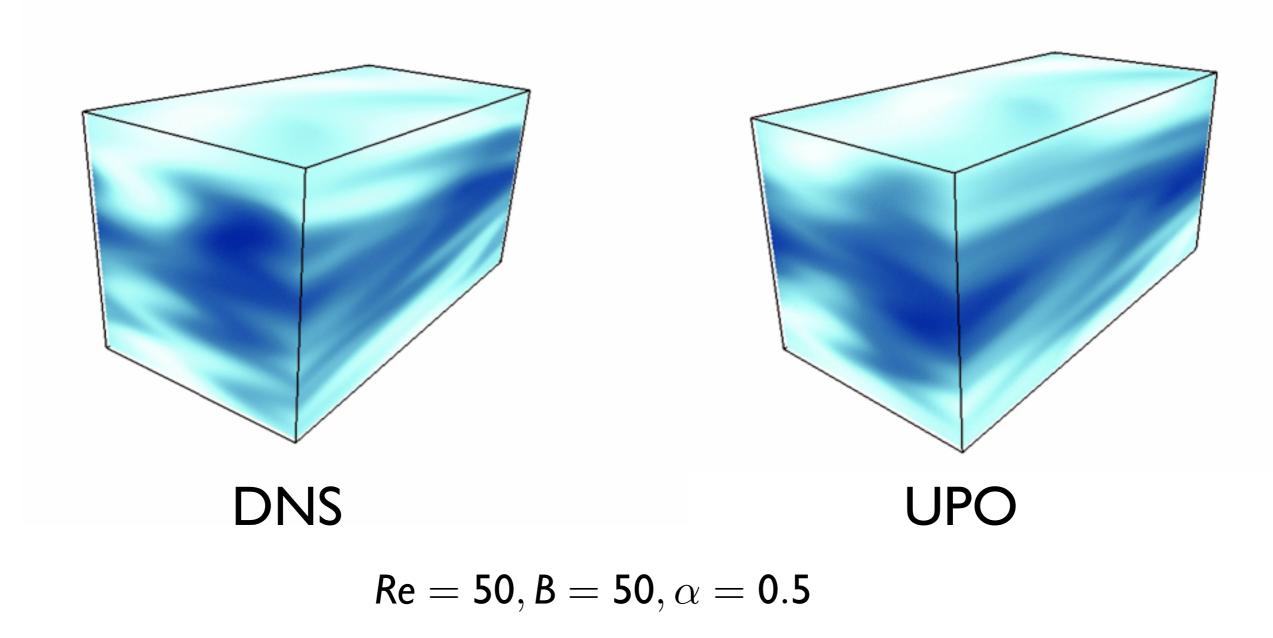
• But can they tell us something quantitative about the flow behaviour?

• Can find stratified UPOs, and they really "look" like the flow:

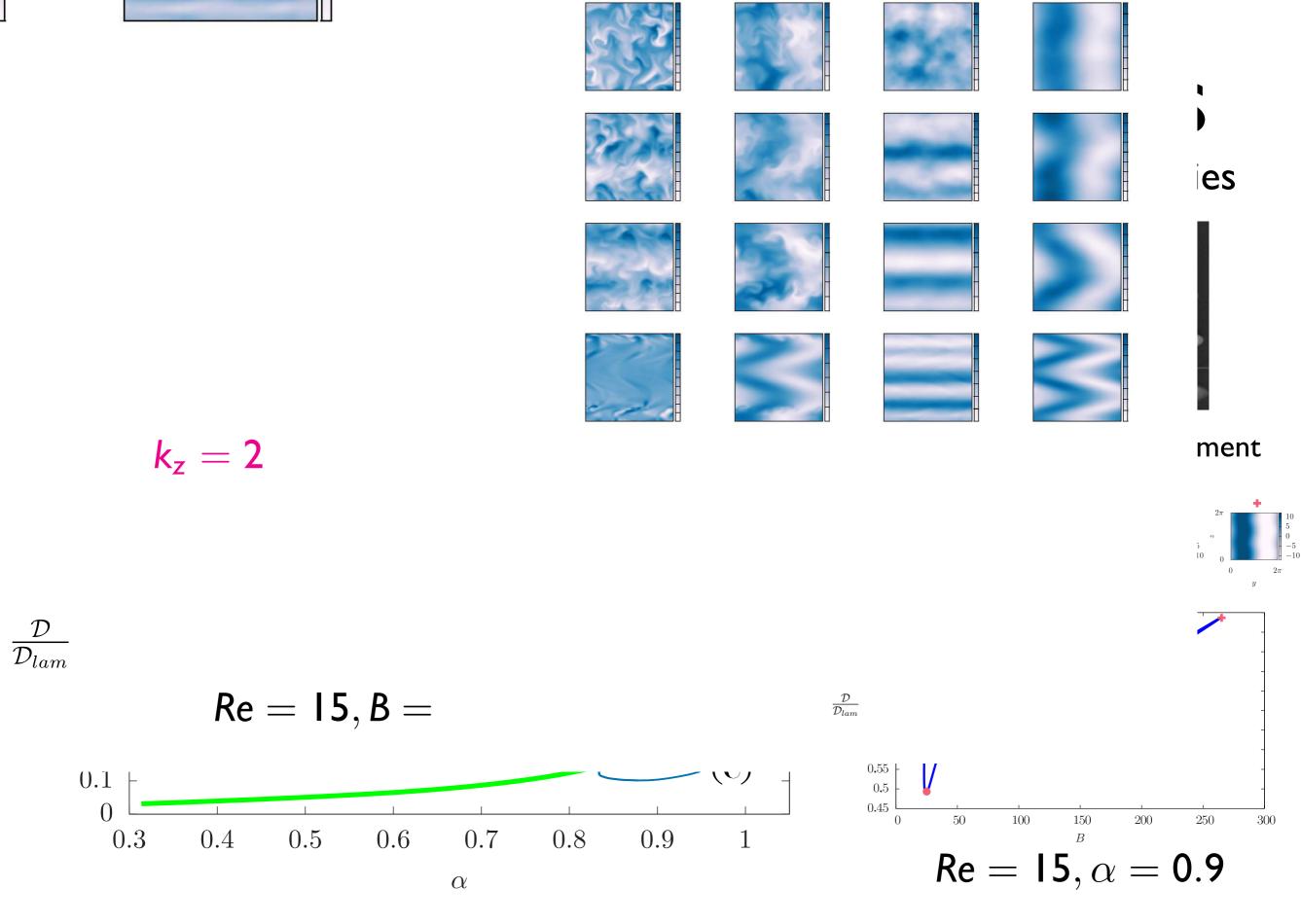


But can they tell us something quantitative about the flow behaviour?

• Can find stratified UPOs, and they really "look" like the flow:



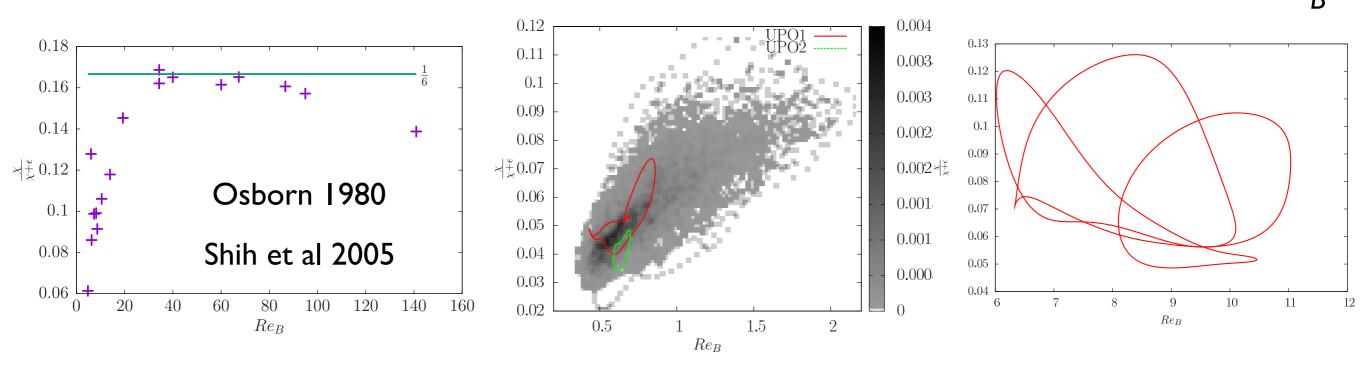
But can they tell us something quantitative about the flow behaviour?



Chevron state is a secondary bifurcation from both primary instabilities!

# Mixing?

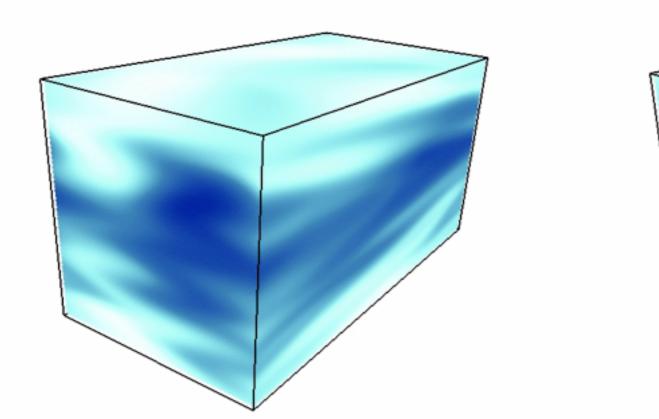
- In stratified turbulence, potential energy sink as well as dissipation
- Question: how much into mixing/how much into dissipation: efficiency
- Buoyancy versus momentum dissipation:  $\chi = B\langle |\nabla \rho|^2 \rangle, \ \epsilon = \frac{1}{Re}\langle |\nabla \mathbf{u}|^2 \rangle$
- Flow is consistent with standard models of other flows with  $Re_B =$

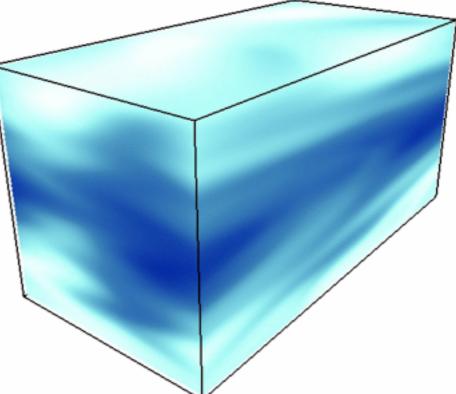


• ECS has the quantitative mixing property of the full flow...higher Re?

## B: Conclusions

- Recurrent flow analysis applicable to stratified shear turbulence
- ECS/UPOs: RECURRENCE & SELF-ORGANIZATION
- Finite amplitude bifurcations from linear instabilities
- Layering structure and mixing **DYNAMICS** set quantitatively
- Can they be used as "modes" for reduced models/explanations?





## B: Conclusions

- Recurrent flow analysis applicable to stratified shear turbulence
- ECS/UPOs: RECURRENCE & SELF-ORGANIZATION
- Finite amplitude bifurcations from linear instabilities
- Layering structure and mixing **DYNAMICS** set quantitatively
- Can they be used as "modes" for reduced models/explanations?

