Particle acceleration by relativistic shock fronts

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- First order Fermi processes—basics
- Diffusive shock acceleration vs. acceleration at relativistic shocks
- Shock physics
- Transport properties
- Injection

DSA (1977/8)

Scattering by (self-excited) turbulence around shock front



and escape probability downstream = 4u'/v

Combining acceleration and escape gives

$$N(p) \propto p^{2-s}$$
 with $s = \frac{3r}{r-1}$ and $r = \frac{u}{u'}$

- Power-law index a function of compression ratio r alone
- For a strong shock front $r = 4 \Rightarrow s = 4$

Nonrelativistic (DSA)

pitch-angle diffusion \Rightarrow near-isotropy \Rightarrow spatial diffusion

solution of PDE in x, p required

power-law of index s = 3r/(r-1), independent of scattering law small escape probability, small $\langle \Delta p \rangle / p$ per cycle

Relativistic

pitch-angle diffusion, particles in narrow, forward directed cone

solution of PDE in μ, x, p required

power-law, index depends on scattering law

escape probability ~ 0.5 , $\langle \Delta p \rangle / p \sim \Gamma^2$ for first cycle, then ~ 2 Eigenfunction expansion \Rightarrow angular dependence:

$$\frac{\exp\left(-\frac{1+\mu_{\rm s}}{1-\mu_{\rm s}u/c}\right)}{\left(1-\mu_{\rm s}u/c\right)^s}$$

Power-law index and compression ratio:

As $\Gamma \rightarrow \infty$, $s \rightarrow 4.23$ Universal index? Kirk et al ApJ 542, 235 (2000)



Loss cone



Comparison of MC/analytic angular distributions

Achterberg et al MNRAS 328, 393 (2001)



Ultra rel. MC



Monte-Carlo simulations in limit $\Gamma \to \infty$

Application to Blazar



and to GRB

Observations of GRB 970508 (Galama et al 1998)



• Equation of state: for relativistic gas ($e \gg \rho$):

$$\frac{v_-}{v_+} \to 3 \qquad \Gamma_{\rm rel} \to \frac{\Gamma_{\rm shock}}{\sqrt{2}}$$

 But, inclusion of ambient magnetic field reduces compression, parameterised by

$$\sigma = \frac{B_{-}^2}{4\pi w} \quad \text{or} \quad M_{\text{fast}} = v_{-} \sqrt{\frac{1+\sigma}{\sigma+c_{\text{s}}^2}}$$

$$(w = 4e/3 = 4 \langle \gamma \rangle n/3$$
, and $c_s = 1/\sqrt{3}$.)

Magnetic field

Effect of ambient magnetic field



Field amplification (e.g., Weibel instability)

$$T^{\mu\nu} = \left(w + \frac{B^2}{4\pi}\right)u^{\mu}u^{\nu} + \left(p + \frac{B^2}{8\pi}\right)g^{\mu\nu} - \frac{B^{\mu}B^{\nu}}{4\pi}$$

If B in shock plane and $\langle \mathbf{B} \rangle = 0$, $\gamma_{\text{eff}} = 4(1 + \sigma_+)/(3 + 2\sigma_+)$

$$\frac{v_{-}}{v_{+}} \to (\gamma_{\rm eff} - 1)^{-1}$$
 and $\Gamma_{\rm rel} \to \Gamma_{\rm shock} \sqrt{\frac{2 - \gamma_{\rm eff}}{\gamma_{\rm eff}}}$

Effect of shock-generated magnetic field





- No scattering ⇒ no stochastic acceleration
- Regular vs. stochastic deflection upstream—minor effect on spectrum (Monte-Carlo: Achterberg et al 2001)
- Presence/absence of downstream scattering—also a minor effect (Monte-Carlo: Ostrowski & Bednarz 2000)
- Anisotropic diffusion?

Relativistic Weibel instability Yang et al (1993, 1994) Medvedev & Loeb (1999)

Wave length: $\lambda_{\perp} \approx \lambda_{\max} \approx c/\omega_{p}$ Growth-length: $\lambda_{\parallel} \approx v_{+}/\omega_{p}$



Effect of anisotropic scattering



$$D_{\mu\mu} \propto \frac{1-\mu^2}{\sqrt{\mu^2 + \left(\lambda_{\parallel}/\lambda_{\perp}\right)^2}}$$

For large $\Gamma_{\rm shock}$ and $\lambda_{\parallel}/\lambda_{\perp}=1/10$

 $\Delta s \approx 0.02$

- Thermalised vs. accelerated particles: $\lambda_{gyro} > \lambda_{\parallel}$ e^-e^+ -plasma $\Rightarrow \gamma > \text{few} \times \Gamma_{\text{shock}}$ e^-p -plasma $\Rightarrow \gamma > (M/m) \times \Gamma_{\text{shock}}$
- Pre-acceleration? Synchrotron resonance for positron-ion coupling Gallant et al (1994)
- Soft spectrum $s > 4 \Rightarrow$ *linear* effect (in nonrelativistic theory highly nonlinear)

- Robust prediction of spectral slope (s = 4.2-4.3 for strong ultrarelativistic shock)
- Nonlinearity weak (compared to nonrelativistic case)
- Injection mechanism needed for e⁻p plasmas (synchrotron maser?)